

Spin-orbital quantum liquid on the honeycomb lattice

F. Mila

Ecole Polytechnique Fédérale de Lausanne
Switzerland

Collaborators

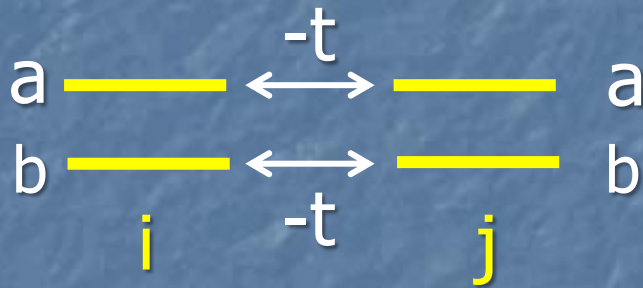
P. Corboz (Zürich), **A. Läuchli** (Innsbruck)

M. Lajko and **K. Penc** (Budapest)

Scope

- Introduction: the spin-orbital symmetric Kugel-Khomskii model as $SU(4)$ quantum permutation
- $SU(4)$ on honeycomb lattice
 - Hartree
 - Flavour-wave theory
 - iPEPS
 - Gutzwiller projected wavefunctions
- Conclusion: no LRO, algebraic correlations

Mott insulator with orbital degeneracy



- 1 electron per site
- on-site repulsion U

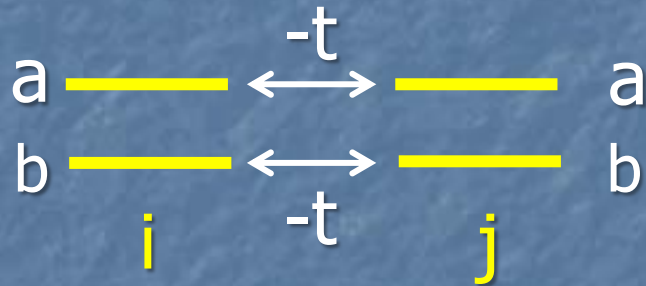
spin

orbital

$$\mathcal{H} = \frac{2t^2}{U} \sum_{\langle i,j \rangle} (2\mathbf{S}_i \cdot \mathbf{S}_j + \frac{1}{2})(2\mathbf{T}_i \cdot \mathbf{T}_j + \frac{1}{2})$$

Symmetric Kugel-Khomskii model

SU(4) formulation



- 1 electron per site
- on-site repulsion U




$$\mathcal{H} = -t \sum_{\langle i,j \rangle, \alpha} (c_{i,\alpha}^\dagger c_{j,\alpha} + \text{H.c.}) + U \sum_{i, \alpha < \beta} n_{i,\alpha} n_{i,\beta}$$

$$\alpha = |a \uparrow\rangle, |a \downarrow\rangle, |b \uparrow\rangle, |b \downarrow\rangle$$

$$\mathcal{H} = \frac{2t^2}{U} \sum_{\langle i,j \rangle} P_{ij}$$

Permutations

Quantum permutations

- Objects with N flavours on a lattice
- Hilbert space = $\{ | \sigma_1 \sigma_2 \dots \sigma_L \rangle \}$
 $\sigma_i = 1, 2, \dots, N$ or $\sigma_i = A, B, C, \dots$ or  ,  ,  ...

$$\mathcal{H} = \sum_{\langle i, j \rangle} P_{ij}$$

$$P_{ij} | \sigma_1 \dots \sigma_i \dots \sigma_j \dots \sigma_N \rangle = | \sigma_1 \dots \sigma_j \dots \sigma_i \dots \sigma_N \rangle$$

SU(N) formulation

$$H = \sum_{\langle i,j \rangle} S_m^n(i) S_n^m(j)$$

$$S_m^n |\mu\rangle = \delta_{n,\mu} |m\rangle$$

$$[S_m^n, S_k^l] = \delta_{n,k} S_m^l - \delta_{m,l} S_k^n$$

S_m^n \rightarrow generators of SU(N)

At each site: fundamental N-dimensional representation

Physical realizations

■ **N=2** → spin-1/2 Heisenberg $P_{ij} = 2\vec{S}_i \cdot \vec{S}_j + 1/2$

■ **N=3** → S=1 biquadratic $P_{ij} = \vec{S}_i \cdot \vec{S}_j + (\vec{S}_i \cdot \vec{S}_j)^2 - 1$

■ **N=4** → symmetric Kugel-Khomskii model

$$H = \sum_{ij} J_{ij} \left(2\vec{s}_i \cdot \vec{s}_j + \frac{1}{2} \right) \left(2\vec{\tau}_i \cdot \vec{\tau}_j + \frac{1}{2} \right)$$

■ **Any N** → N-flavour fermions in optical lattices

General properties

- Soluble in 1D with **Bethe Ansatz**

→ algebraic correlations with **periodicity** $2\pi/N$

Sutherland, 1974

- **Equivalent of SU(2) dimer singlet: N sites**

$$|S\rangle = \frac{1}{\sqrt{N!}} \sum_P \text{sgn}(P) |\sigma_{P(1)} \sigma_{P(2)} \dots \sigma_{P(N)}\rangle$$

with $\{\sigma_1 \sigma_2 \dots \sigma_N\} = \{1 \ 2 \ \dots \ N\}$

Li, Ma, Shi, Zhang, PRL'98

Hartree approximation

$$|\psi\rangle = \prod_i |\varphi_i\rangle$$

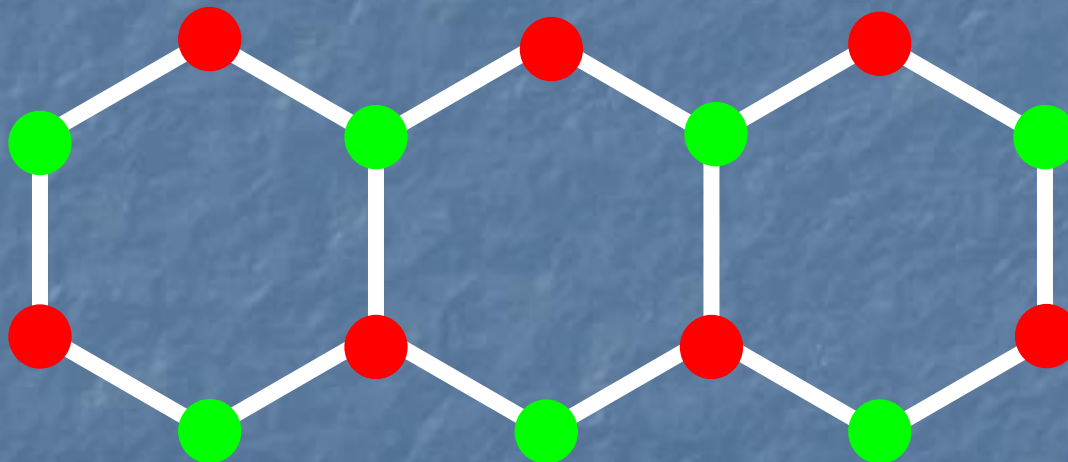
$$\langle\varphi_1\varphi_2|P_{12}|\varphi_1\varphi_2\rangle = \langle\varphi_1\varphi_2|\varphi_2\varphi_1\rangle = |\langle\varphi_1|\varphi_2\rangle|^2$$

→ on 2 sites, **energy minimal** if $\langle\varphi_1|\varphi_2\rangle = 0$

→ on a lattice, Hartree energy certainly **minimal** if **colors on nearest neighbors are different**

SU(4) on honeycomb lattice

Bipartite lattice \rightarrow infinite number of 'Hartree' ground states for more than 2 colors



● or ● \rightarrow ● or ● at any site

Flavour-wave theory

- Schwinger bosons

$$\mathcal{P}_{ij} = \sum_{\mu, \nu \in \{A, B, C, D\}} a_{\mu, i}^\dagger a_{\nu, j}^\dagger a_{\nu, i} a_{\mu, j}$$

$$\sum_{\nu} a_{\nu, i}^\dagger a_{\nu, i} = 1$$

- Local condensation

$$a_{\mu, i}^\dagger, a_{\mu, i} \rightarrow \sqrt{1 - \sum_{\nu \neq \mu} a_{\nu, i}^\dagger \tilde{a}_{\nu, i}}$$

Harmonic theory

$$\mathcal{H} = \sum_{\substack{\alpha, \beta = A, B, C \dots \\ \alpha \neq \beta}} \mathcal{H}_{\alpha\beta}$$

$$\mathcal{H}_{\alpha\beta} = \sum_{\substack{\text{disconnected} \\ \text{clusters } \mathcal{C}}} \sum_{\substack{\langle i, j \rangle \in \mathcal{C} \\ i \in \beta, j \in \alpha}} \mathcal{H}_{\alpha\beta}(i, j)$$

$$\mathcal{H}_{\alpha\beta}(i, j) = (\alpha_i^\dagger + \beta_j)(\alpha_i + \beta_j^\dagger) - 1$$

$$\langle (\alpha_i^\dagger + \beta_j)(\alpha_i + \beta_j^\dagger) \rangle \geq 0 \quad \Rightarrow \quad \langle \mathcal{H}_{\alpha\beta}(i, j) \rangle \geq -1$$

Lower bound saturated for two-site cluster

iPEPS

- iPEPS = infinite Projected Entangled Pair States
- Variational method based on a tensor product wave-function
- Becomes exact if the dimension D of the tensors $\rightarrow \infty$
- Can be seen as a generalization of DMRG

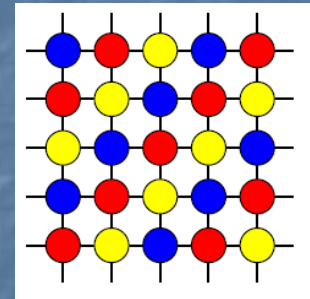
Verstraete and Cirac, 2004

iPEPS and symmetry breaking

- SU(3) on square lattice: **stripe color order**

Toth, Läuchli, FM, Penc, PRL 2010

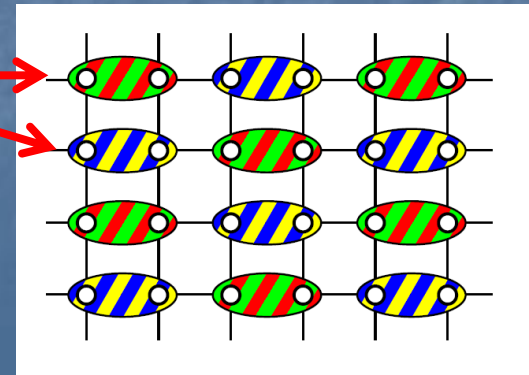
Bauer et al, PRB 2011



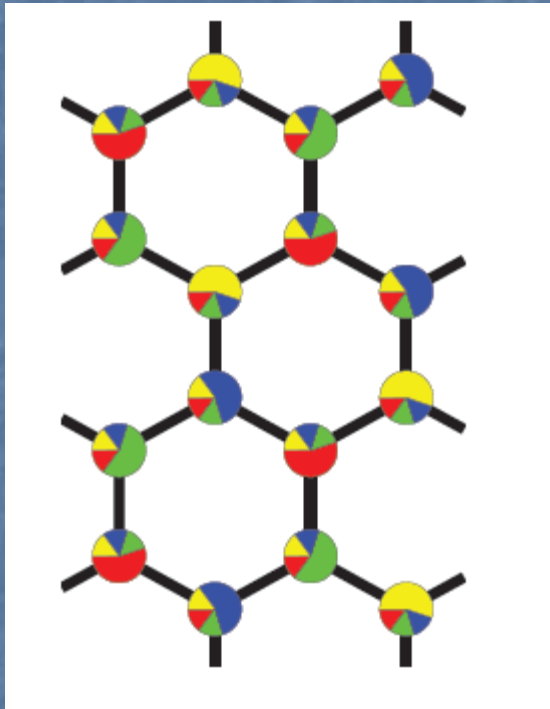
- SU(4) on square lattice: **dimerization**

IRREP dim=6

Corboz, Läuchli, Penc,
Troyer, FM, PRL 2011

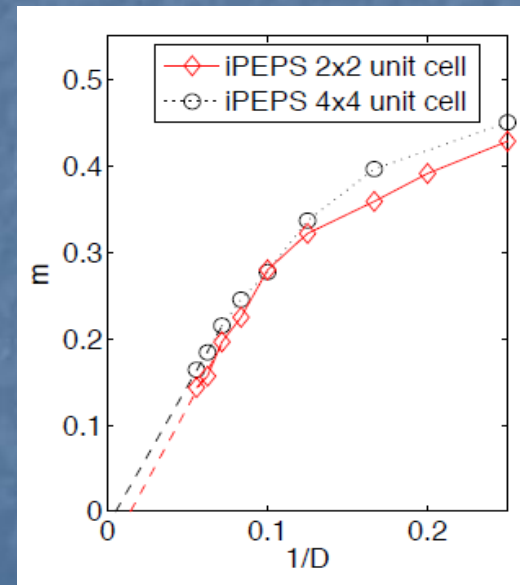


iPEPS for $SU(4)$ on honeycomb

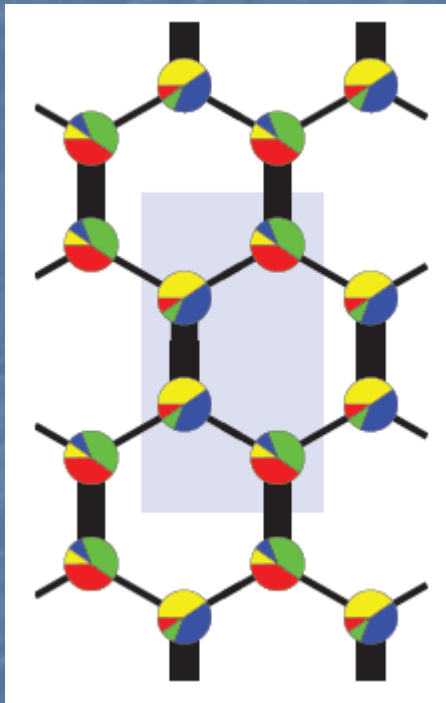


4x4 unit cell

- All bond energies equal
- No color order

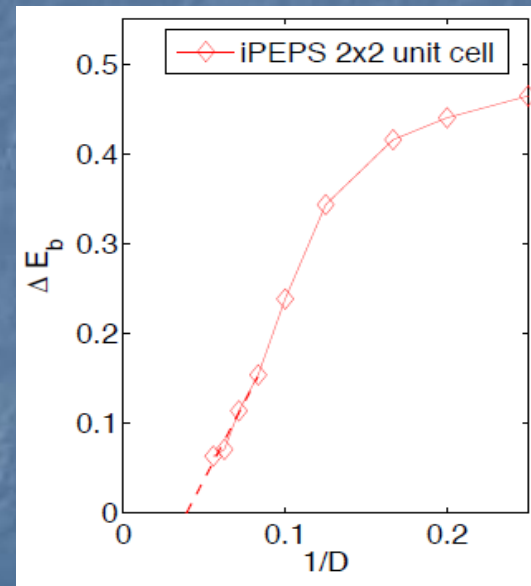


Absence of dimerization



2x2 unit cell

Difference between
bond energies



SU(N) quantum liquids

- RVB? Not very likely (no closed 4-site plaquettes)
- Chiral liquid (N large, N/k atoms per site, $k \geq 5$)

Hermele, Gurarie, Rey, PRL 2009

- Algebraic liquid for SU(4) on square lattice

Wang and Vishwanath, PRB 2009

- Majorana fermion representation
- Half-filled band
- Dirac spectrum for π -flux state

Fermionic approach for SU(4) on honeycomb

- Schwinger fermions

- quarter filling

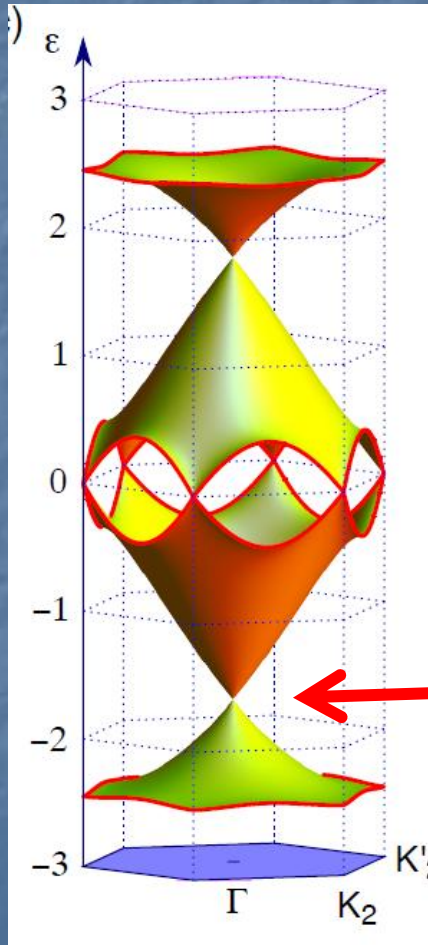
- No flux: Fermi surface

- π -flux: Dirac spectrum

- Majorana fermions

- as for square lattice

π -flux with Schwinger fermions

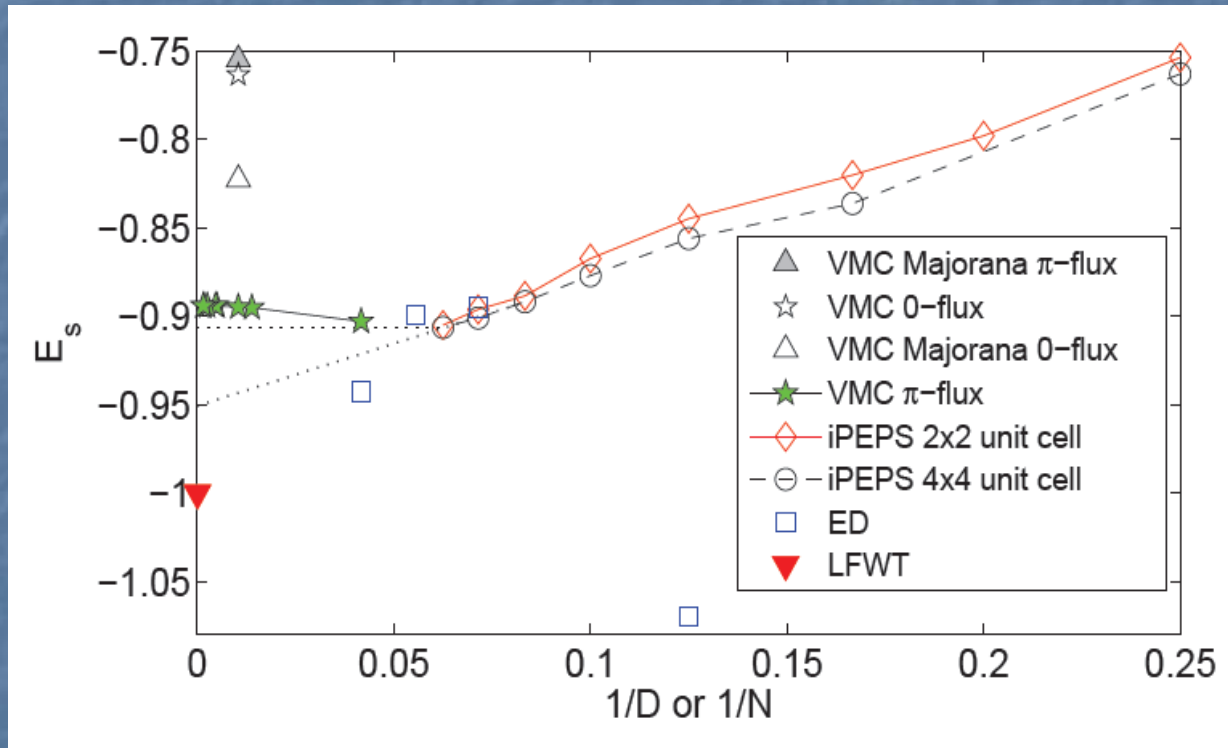


π -flux per hexagonal plaquette



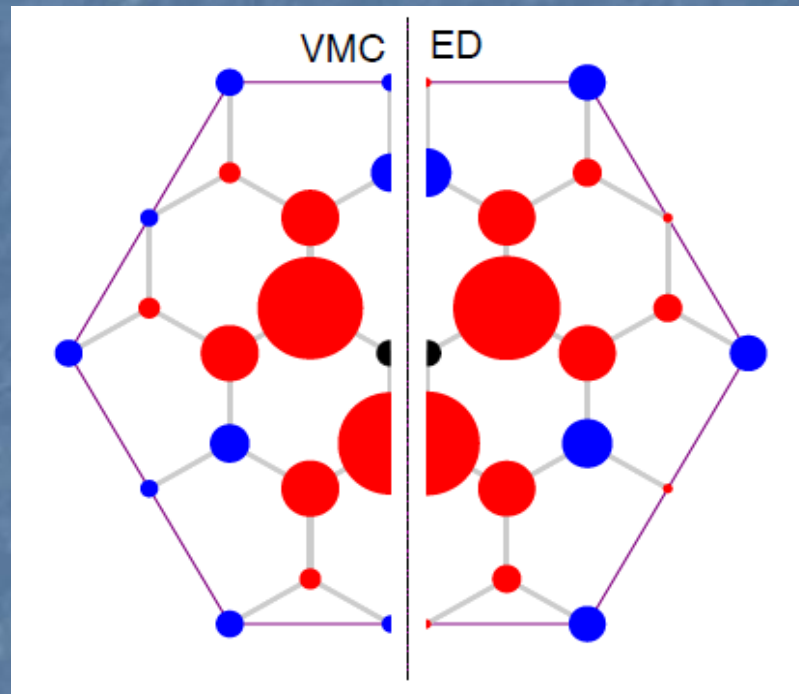
Dirac point at quarter filling

Ground-state energy



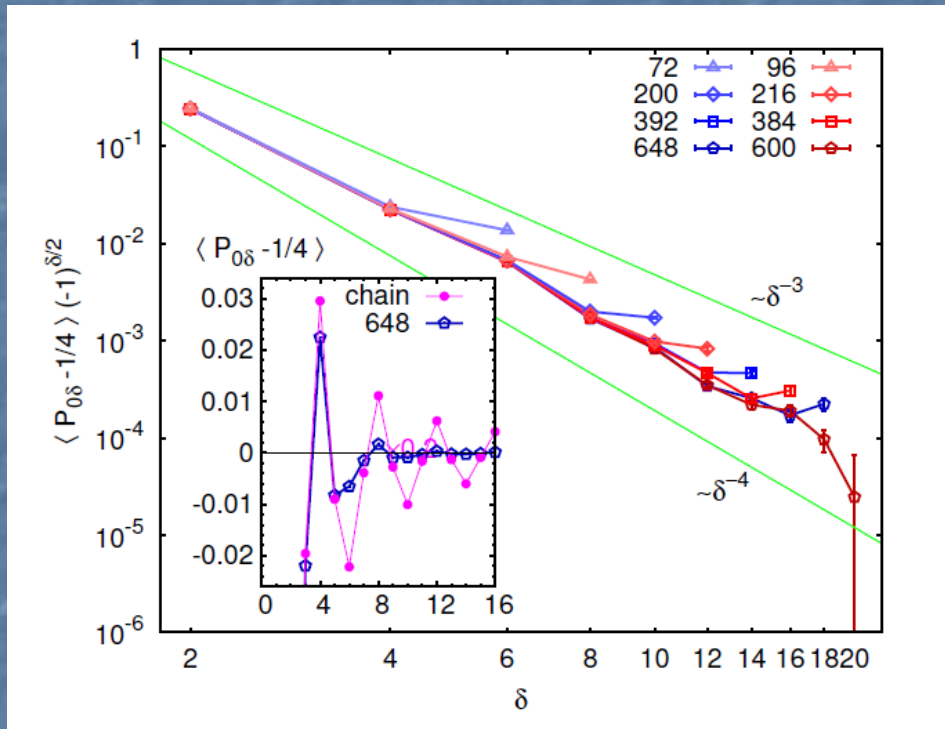
➔ Gutzwiller projected Schwinger fermion π -flux

Comparison ED / π -flux



Short-range color-color correlations: **very good**

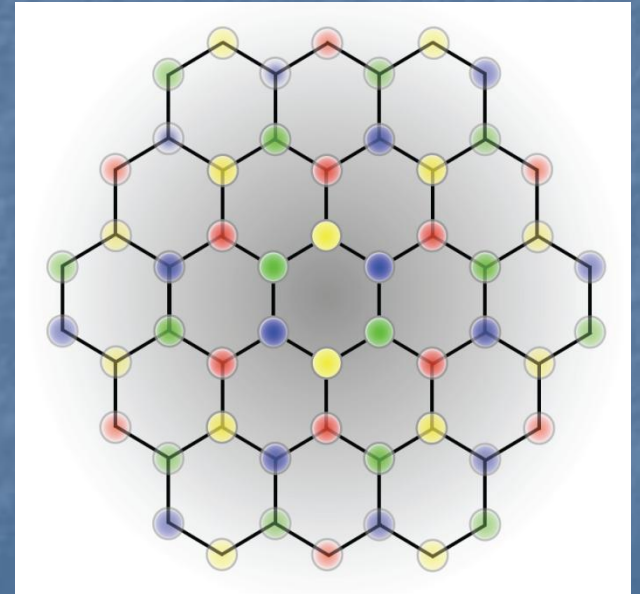
Long-distance correlations



$\langle P_{ij} - 1/4 \rangle \sim |\mathbf{r}_{ij}|^{-\alpha}$, with an exponent α between 3 and 4.

Discussion

- SU(4) in 1D: algebraic correlations with slow decay (exponent 3/2)
- Coupling between chains
 - no LRO
 - faster decay of correlations



Conclusion

- Symmetric Kugel-Khomskii on honeycomb lattice
 - no lattice symmetry breaking
 - no spin or orbital long-range order
 - evidence in favour of an algebraic liquid
- Related to $\text{Ba}_3\text{CuSb}_2\text{O}_9$?