

Postselection-free entanglement dynamics via spacetime duality

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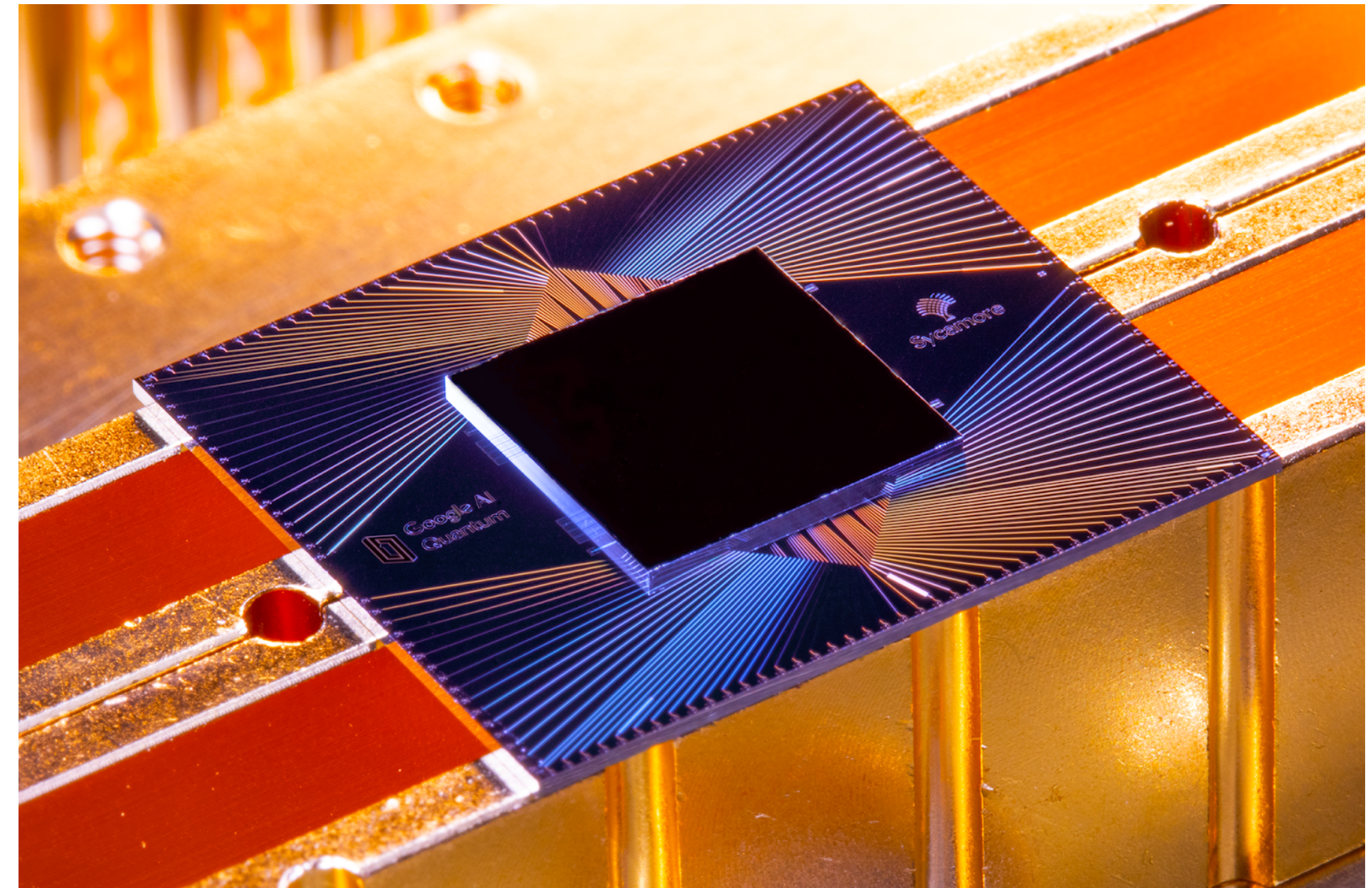
Based on arXiv:2010.xxxxx with V. Khemani

Frontiers of Quantum Computing and Quantum Dynamics, KITP, 10/20/2020

Introduction

Frontiers of quantum dynamics

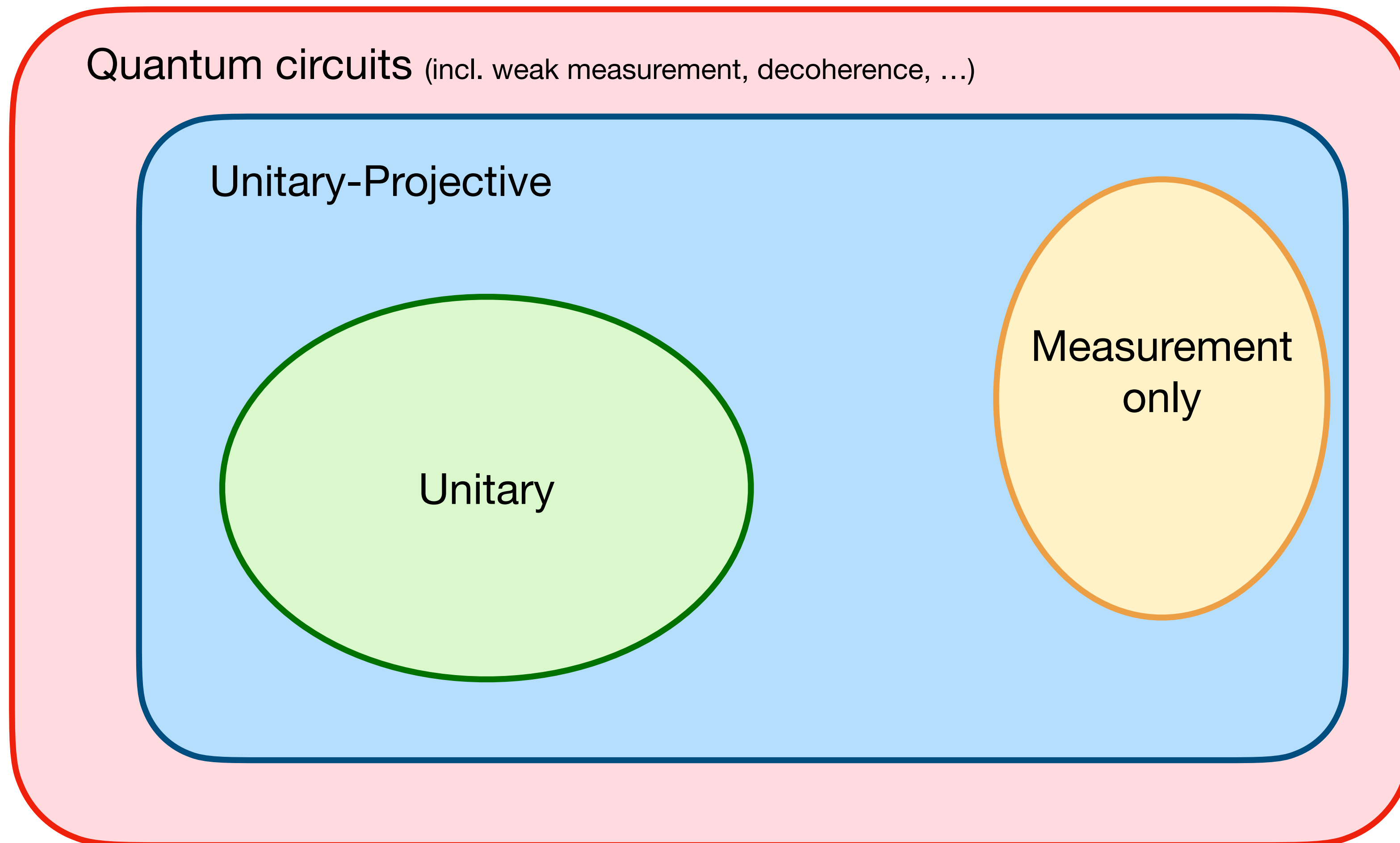
- NISQ devices: unprecedented control, measurement capabilities
- New focus on out-of-equilibrium dynamics, open systems
- Many exciting questions
 - Phases?
 - Universality?
 - Chaos?
 - Classical/quantum simulability?
 - **Entanglement?**



Google's Sycamore chip

Introduction

Entanglement dynamics in quantum circuits



Nahum, Ruhman, Vijay, Haah 16;
von Keyserlingk, Rakovszky, Pollmann, Sondhi 17;
Khemani, Vishwanath, Huse 17;
Zhou, Chen 18;

...

Li, Chen, Fisher 18;
Skinner, Ruhman, Nahum 18;
Choi, Bao, Qi, Altman 19;
Gullans, Huse 19;

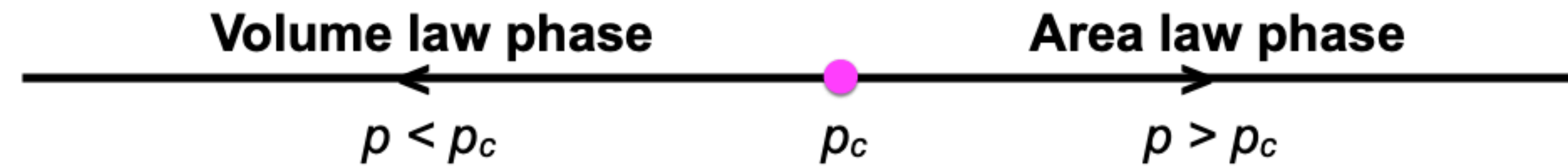
...

Nahum, Skinner 19;
MI, Gullans, Gopalakrishnan, Huse, Khemani 20;
Lavasani, Alavirad, Barkeshli 20;
Sang, Hsieh 20;

Introduction

Entanglement phases in unitary-projective circuits

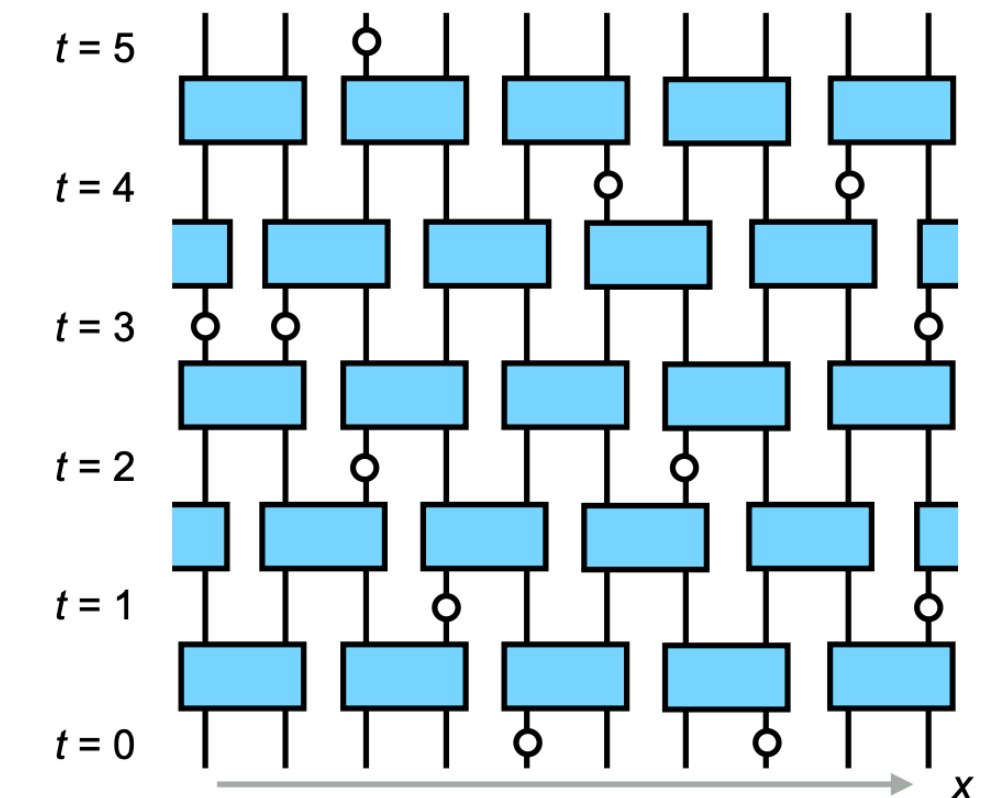
- Random unitary circuits + projective measurements:
entanglement phases [Li, Chen, Fisher 18; Skinner, Ruhman, Nahum 18]
- Many open questions on the nature of phases & transition



Nature of QEC code?
Subleading entropy terms
in Clifford models?

Order?
Phases?

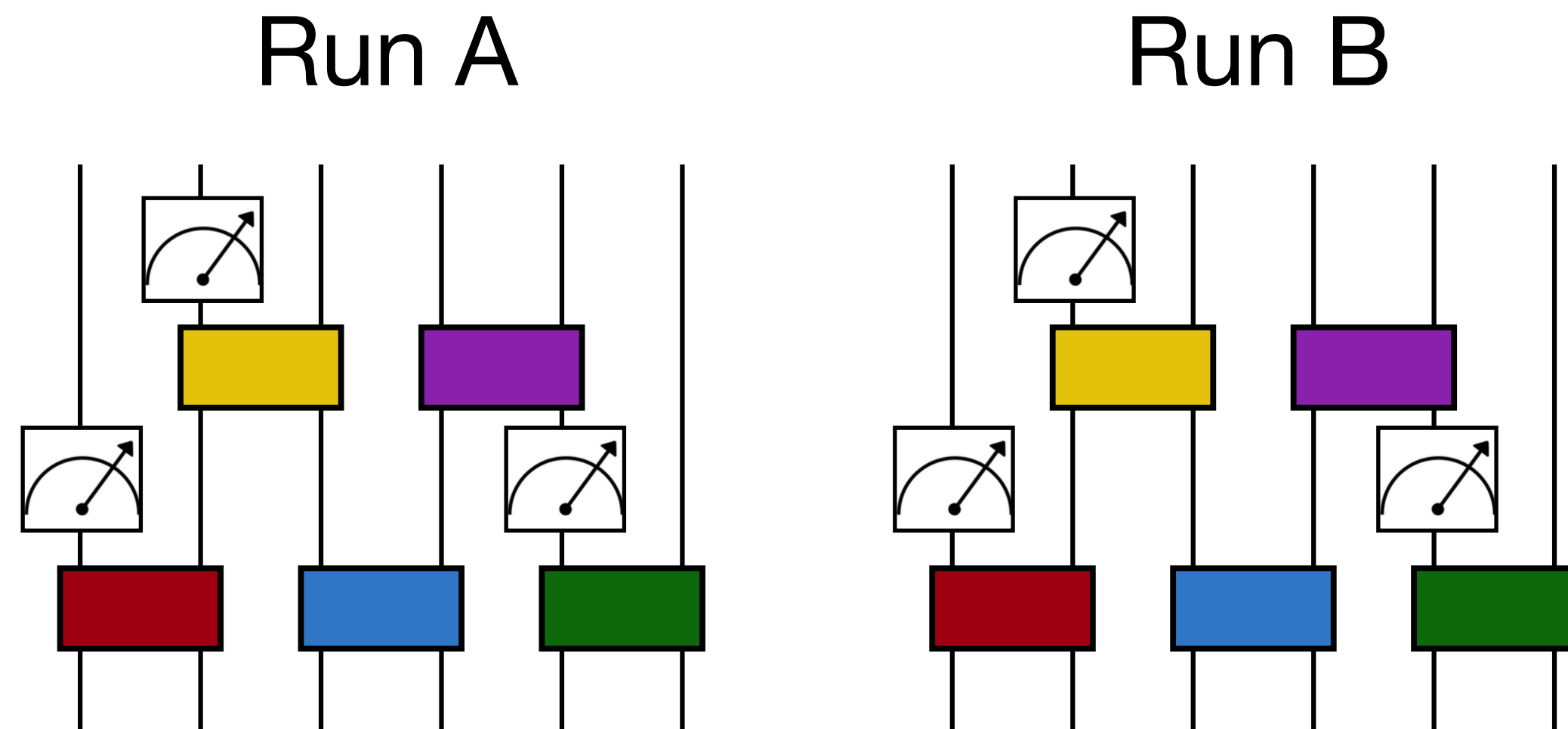
Universality?
CFT?



Li, Chen, Fisher arXiv:1901.08092

The problem of postselection

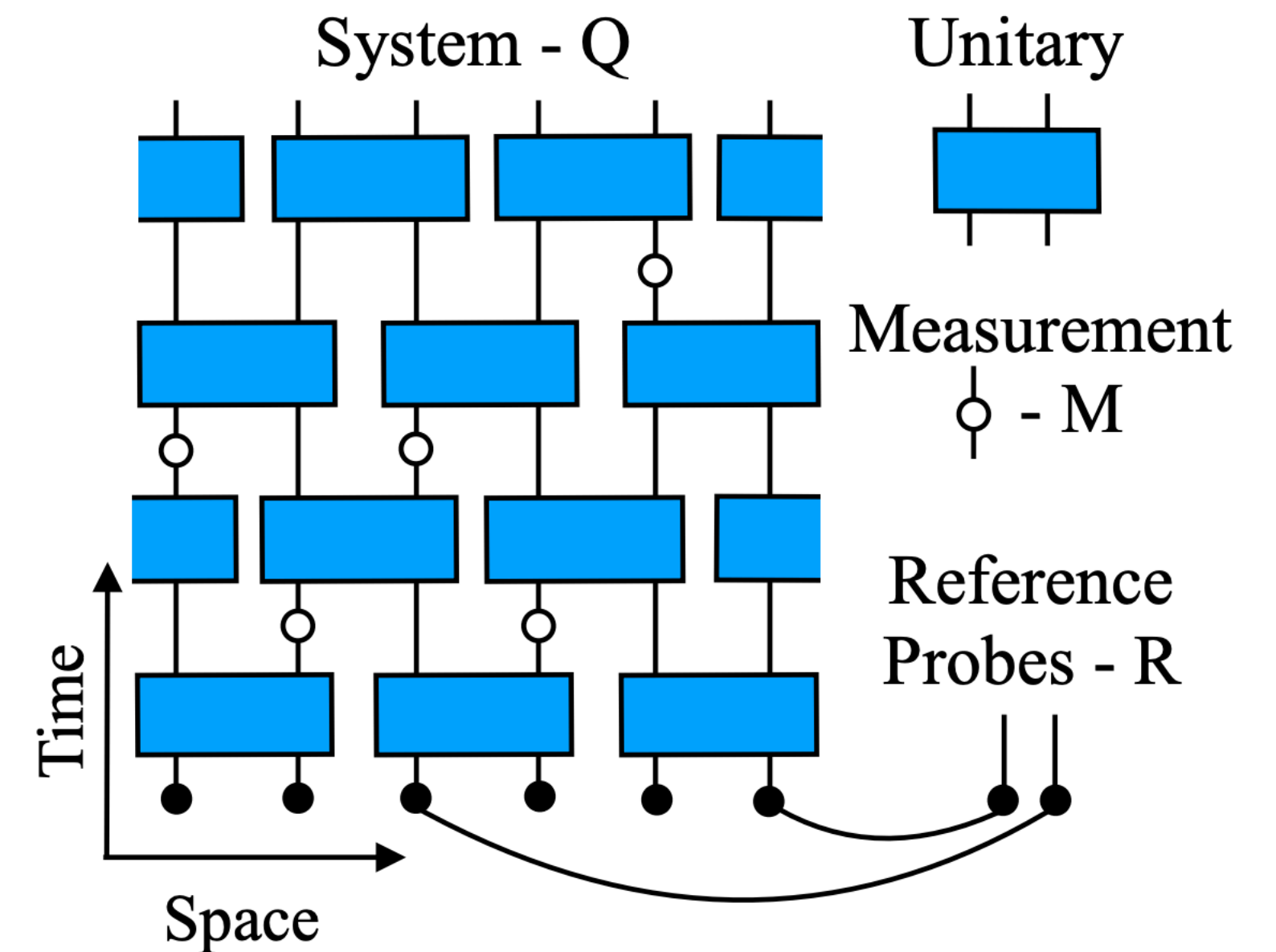
- Measuring entropy: $\text{Tr}(\rho \otimes \rho \cdot \text{SWAP}) = \text{Tr}(\rho^2)$ [Greiner et al 15]
- Producing two identical copies of the state is hard!
 - Quantum randomness of measurement outcomes requires **postselection**
 - Exponential overhead



	Run A	Run B	Same?
M1			
M2			
M3			

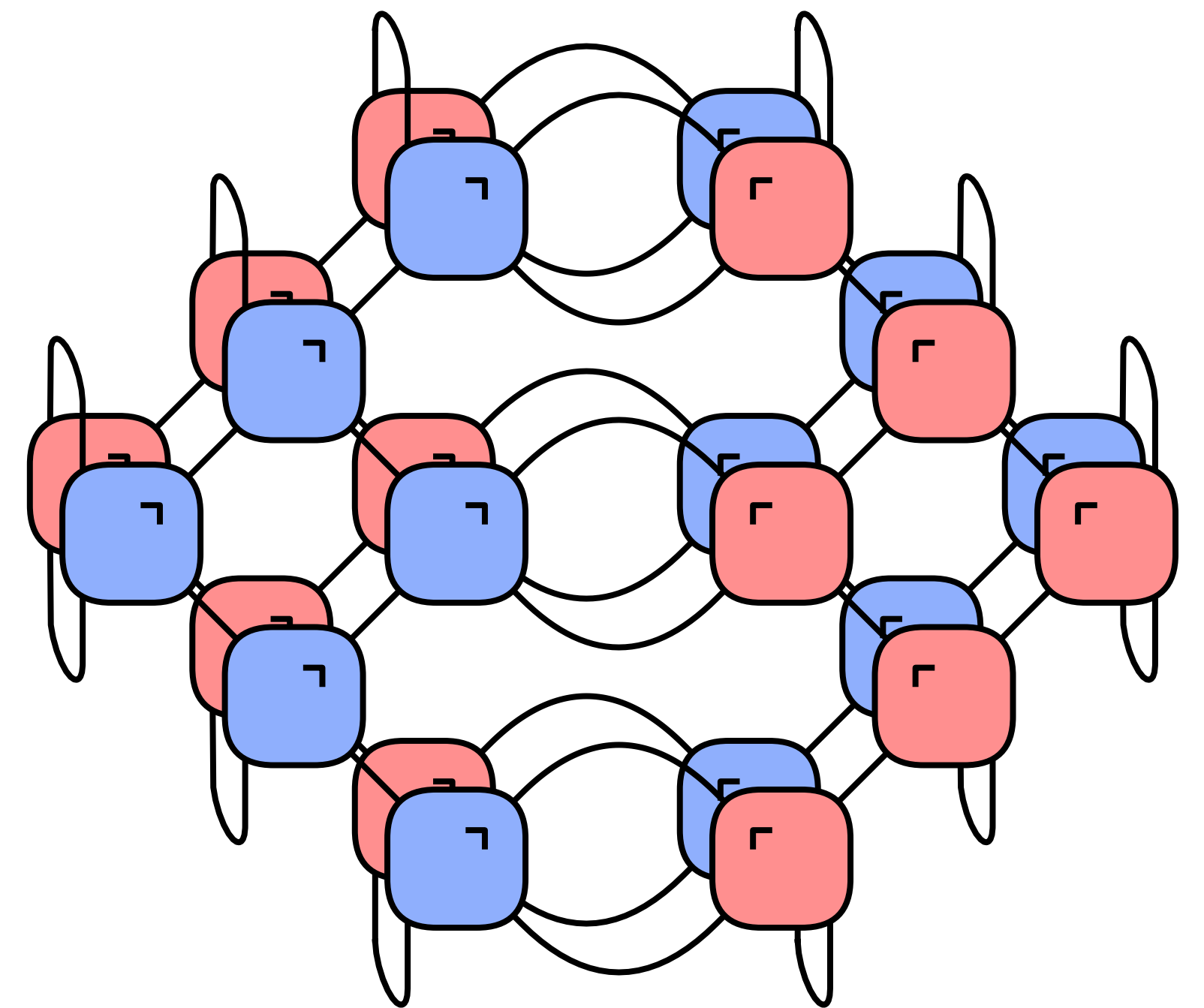
The problem of postselection

- Measuring entropy: $\text{Tr}(\rho \otimes \rho \cdot \text{SWAP}) = \text{Tr}(\rho^2)$ [Greiner et al 15]
- Producing two identical copies of the state is hard!
 - Quantum randomness of measurement outcomes requires **postselection**
 - Exponential overhead
- Ways around this limitation [Gullans, Huse 19]
 - Local order parameter, “decoding lightcone”
 - Classical simulation + feedback (Clifford circuits)
- More direct access to dynamics of entanglement?

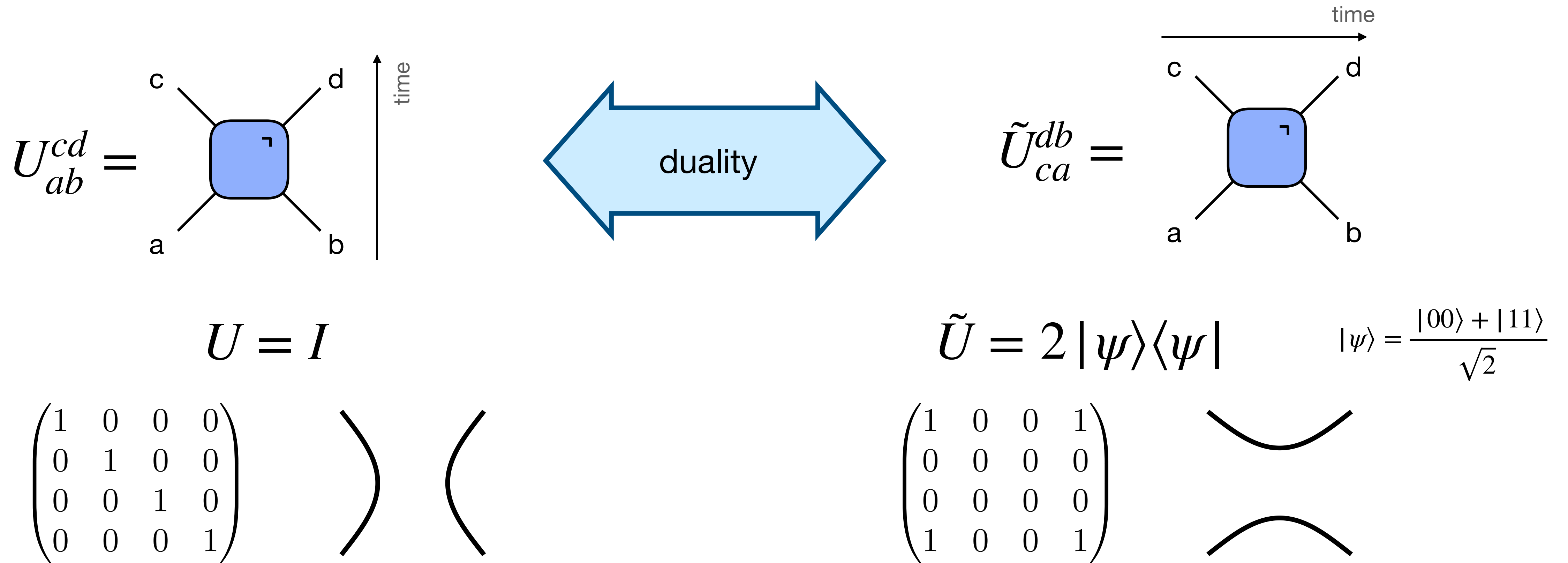


This talk

- A new way of realizing output states of a broad class of **non-unitary circuits**: “*spacetime duals*” of unitary circuits
- Avoids **postselection problem** entirely
- Protocol for direct measurement of the **purity** / **entanglement entropy** of arbitrary subsystems as a correlation function in a unitary evolution



Spacetime duality



- Dualized gates need not be unitary! Generally, **weak measurements**
- **Forced** measurement - always same outcome!

Spacetime duality

Some recent applications

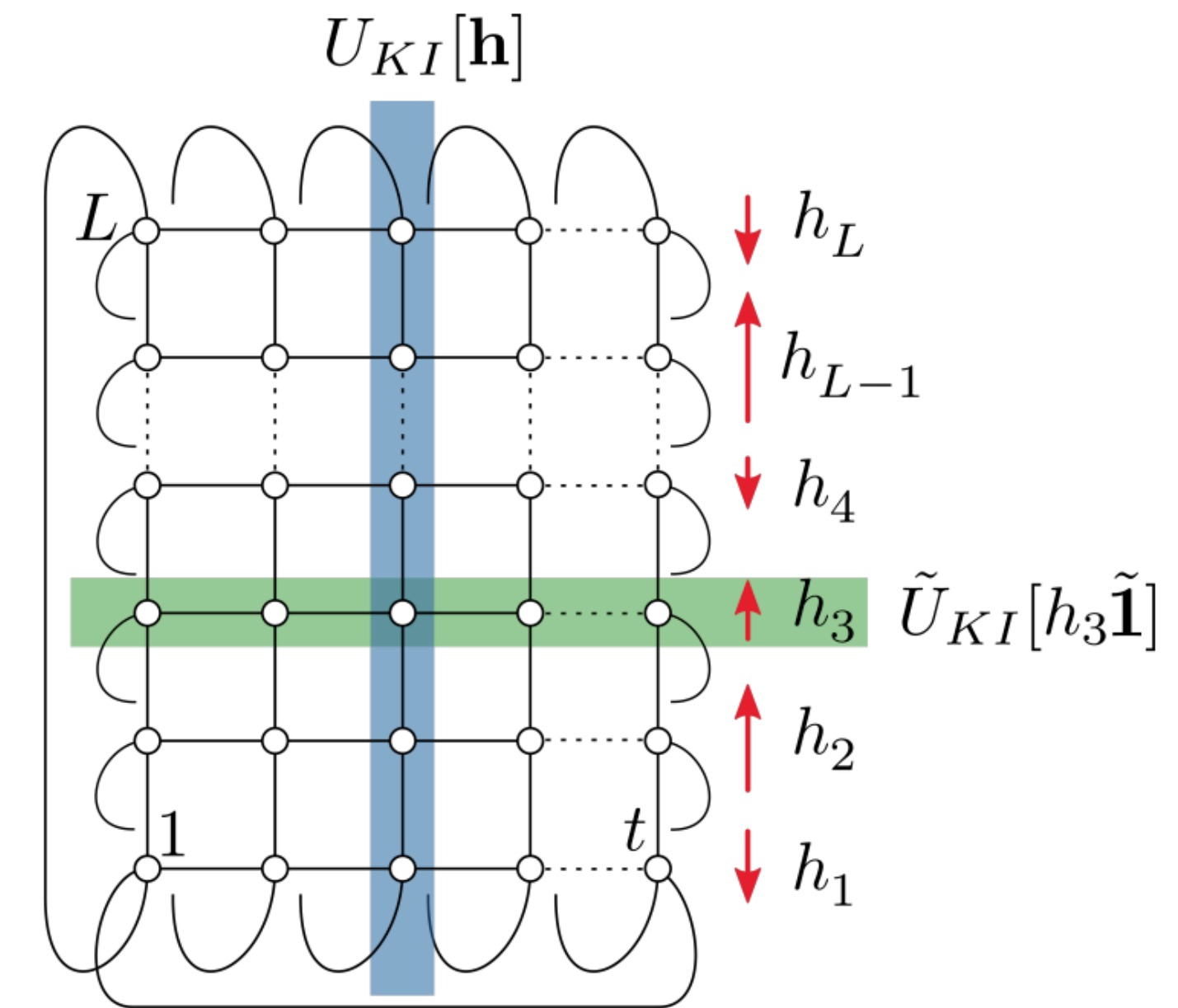
- **Dynamics**

- Dual-unitary circuits: exact connections to random-matrix theory, quantum chaos [Bertini, Kos, Prosen 18]

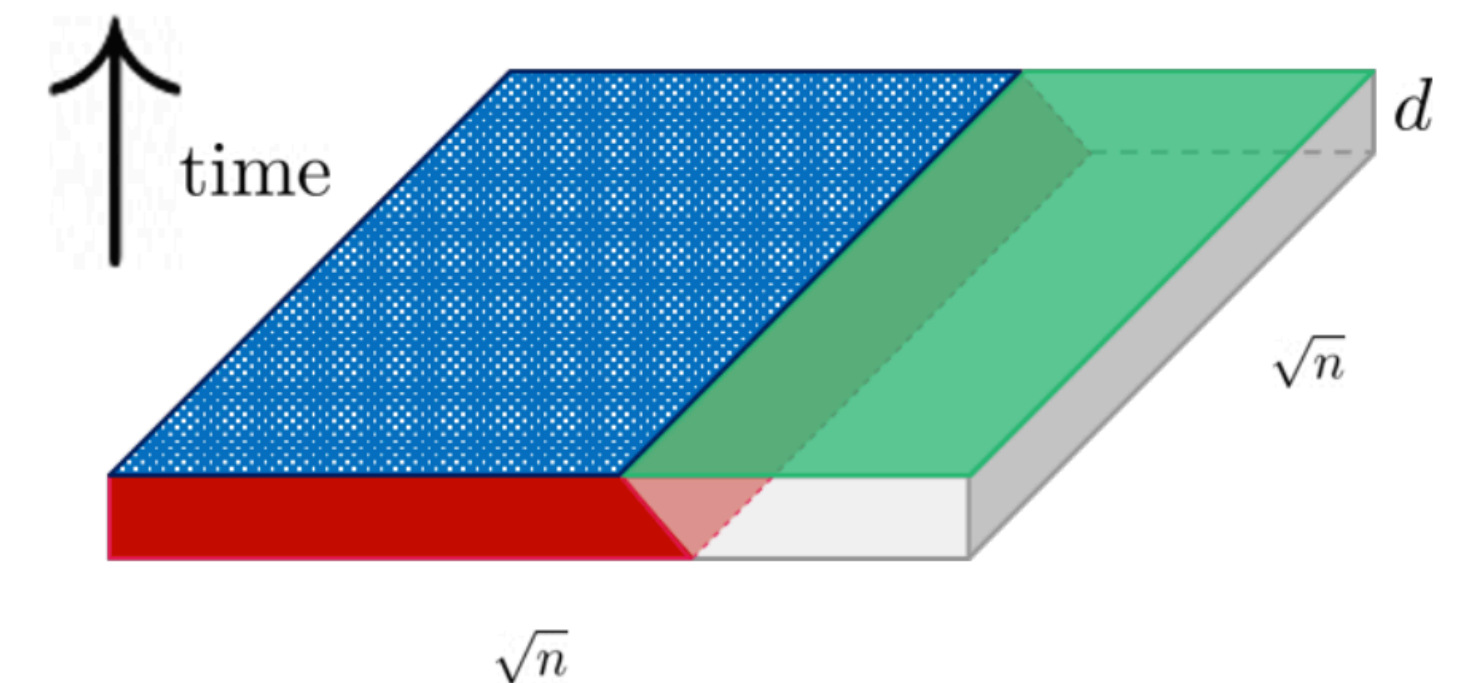
- Analytical control over correlations, OTOCs, entanglement... [Prosen et al 18-20]

- **Complexity**

- Optimized tensor network contractions [Harrow et al 20; Cirac et al 09]



Bertini, Kos, Prosen
arXiv:1812.05090



Napp, La Placa, Dalzell, Brandao, Harrow,
arXiv:2001:00021

Purification dynamics

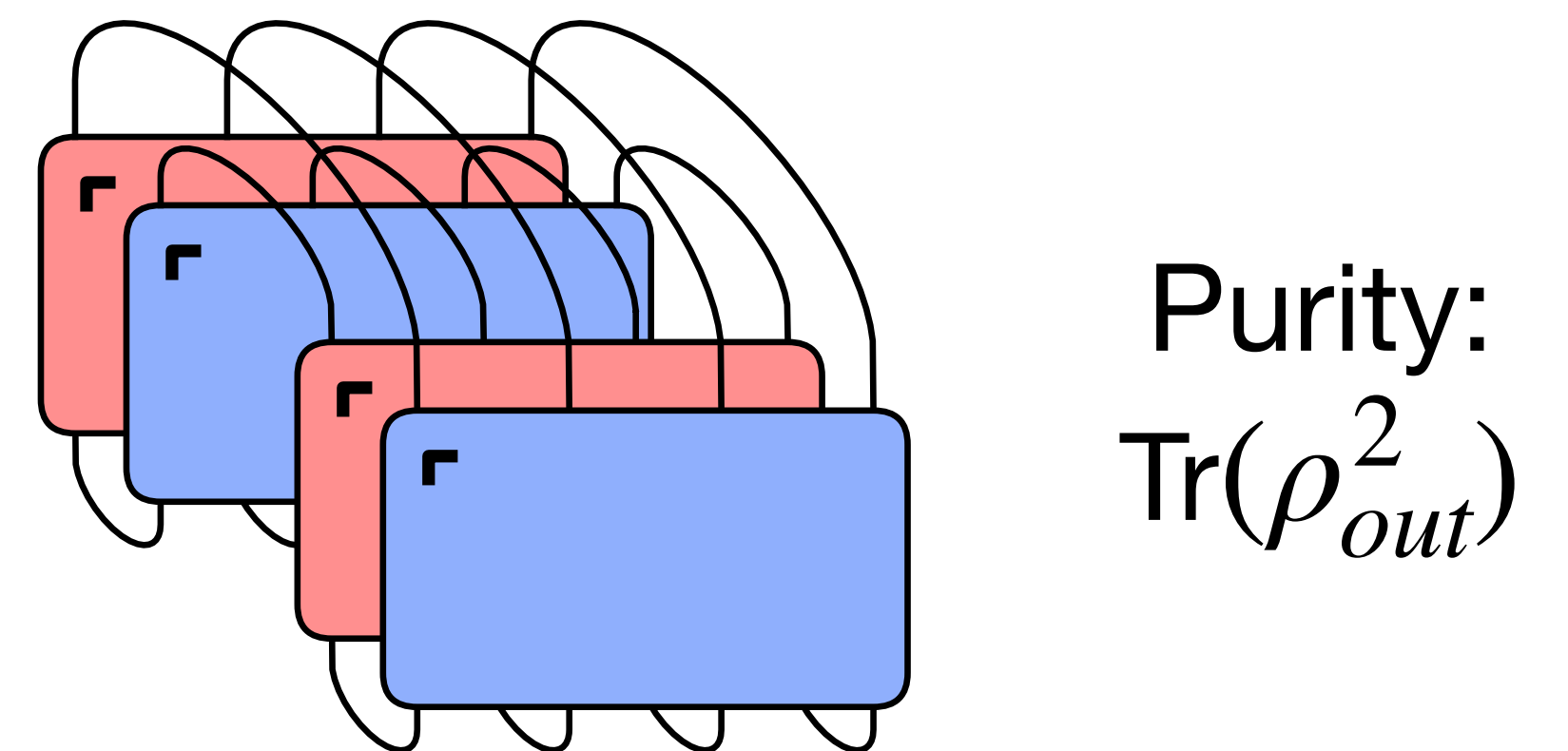
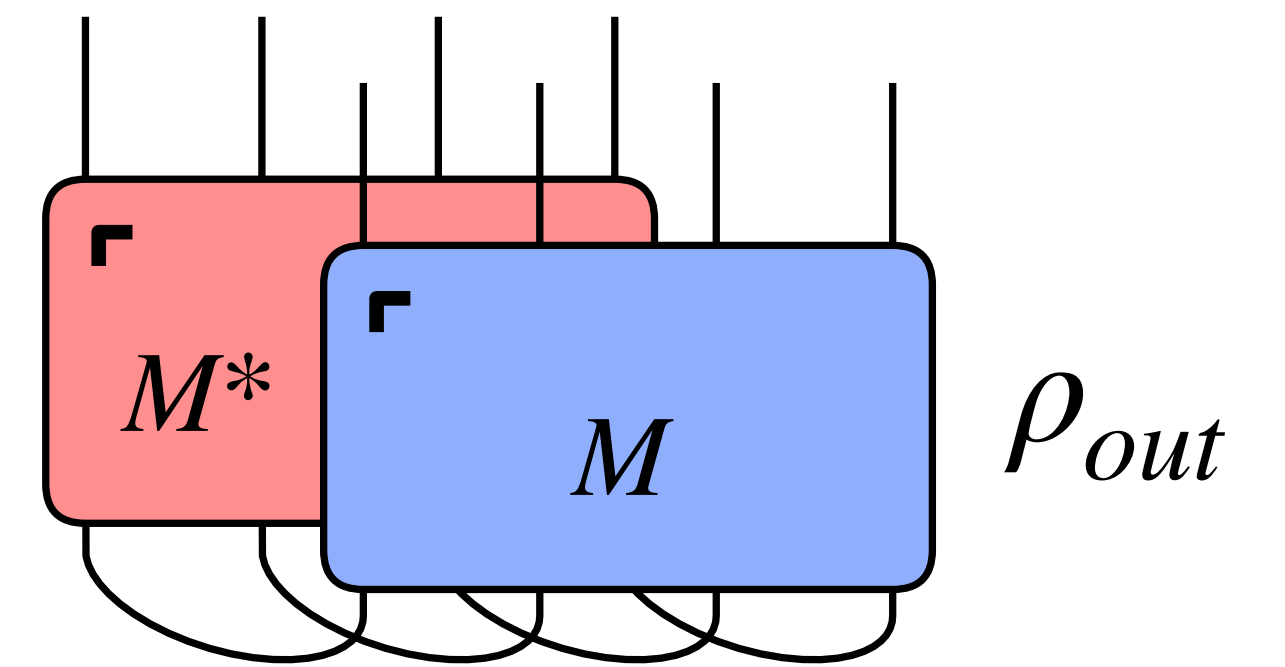
Definition

- Fully mixed state fed into a non-unitary circuit M
- The output state may or may not be purified:

$$\rho_{in} \propto I \mapsto \rho_{out} = M\rho_{in}M^\dagger \propto MM^\dagger$$

- **Purification phases & transitions** [Gullans, Huse 19]
 - Equivalent to the entanglement phases for pure-state inputs

Purification phase	Related Entanglement phase	Purifies in time
Pure	Area-law	$\log(L)$
Critical	Critical	$O(L)$
Mixed	Volume-law	$\exp(L)$

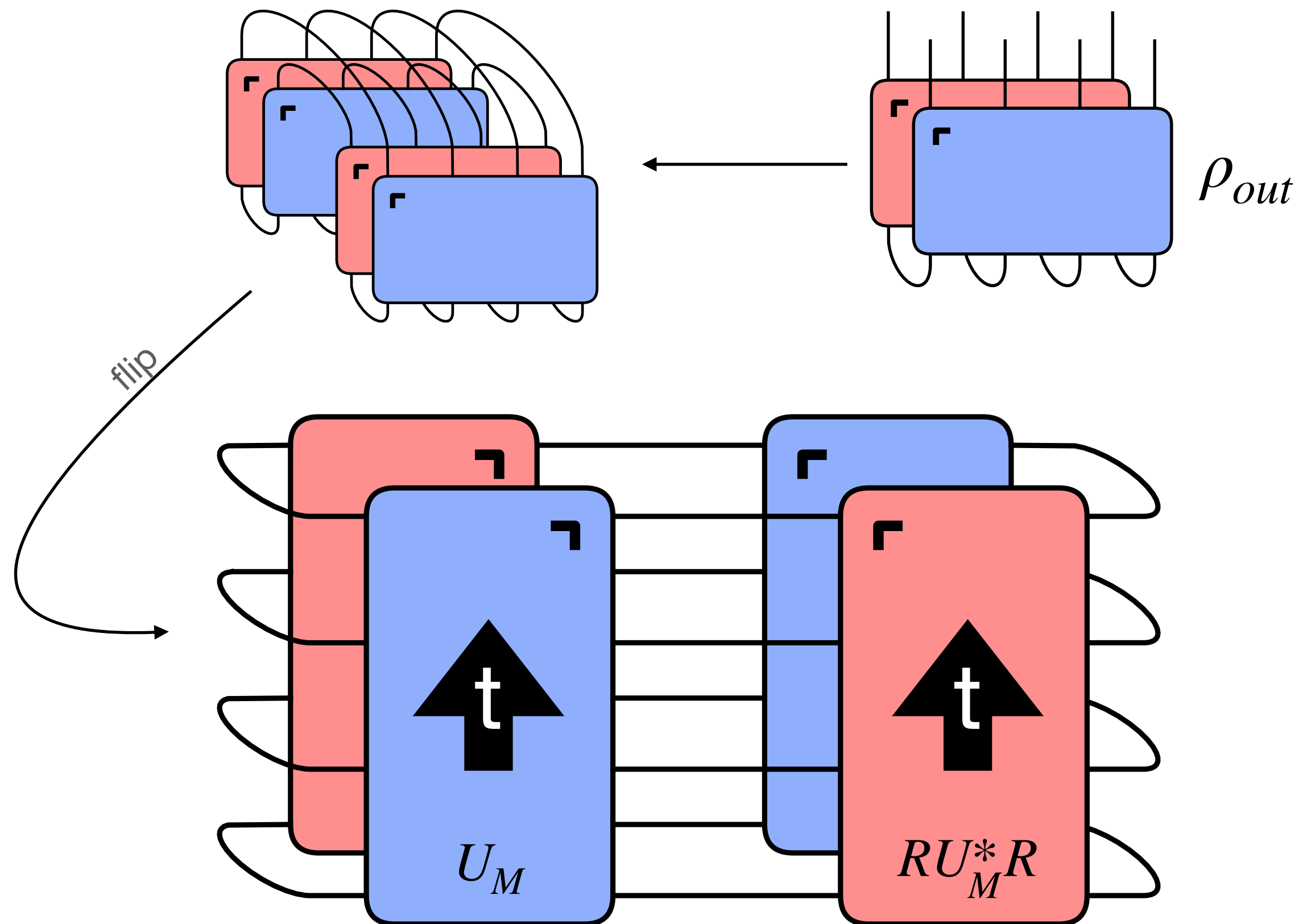


Purification dynamics

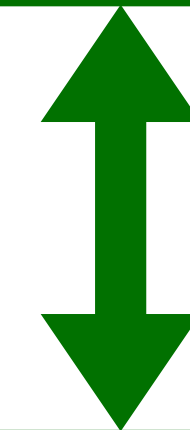
Spacetime duals of unitary circuits

If $M = \tilde{U}_M$, with U_M unitary:

$$\text{Tr}(\rho_{out}^2) = e^{-S_2(\rho_{out})}$$



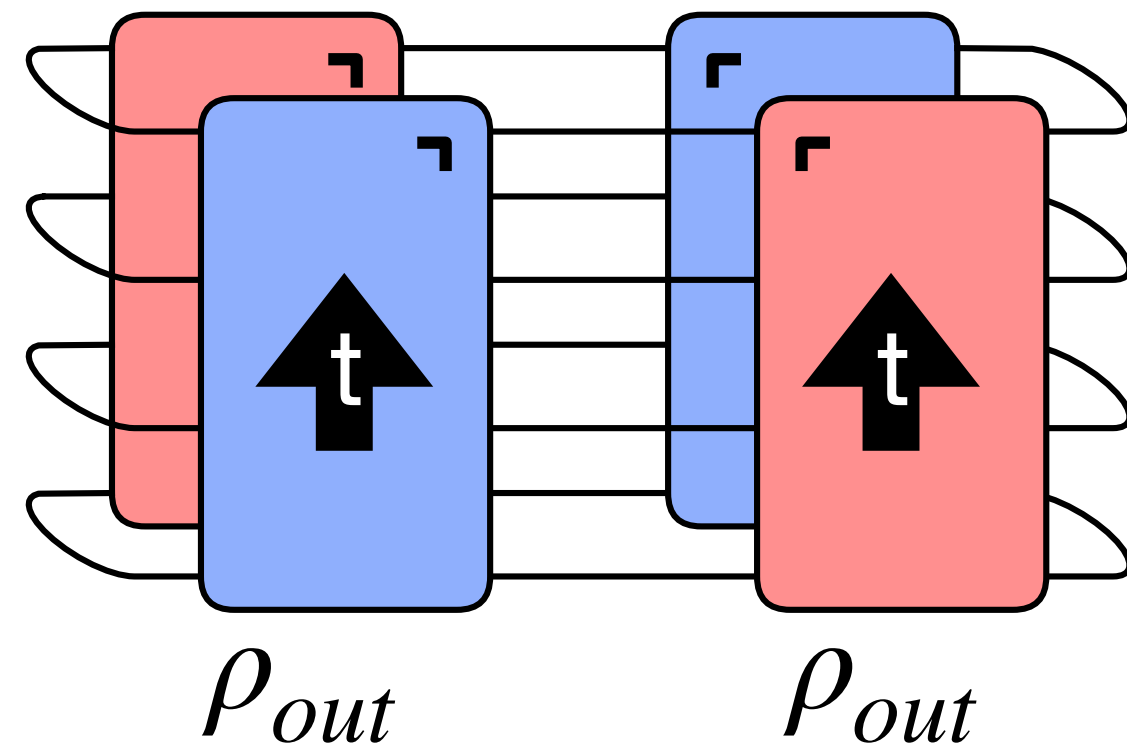
Purity of the output of **hybrid**
circuit M



Correlation function* in a **unitary**
circuit $U_{left} \otimes U_{right}$

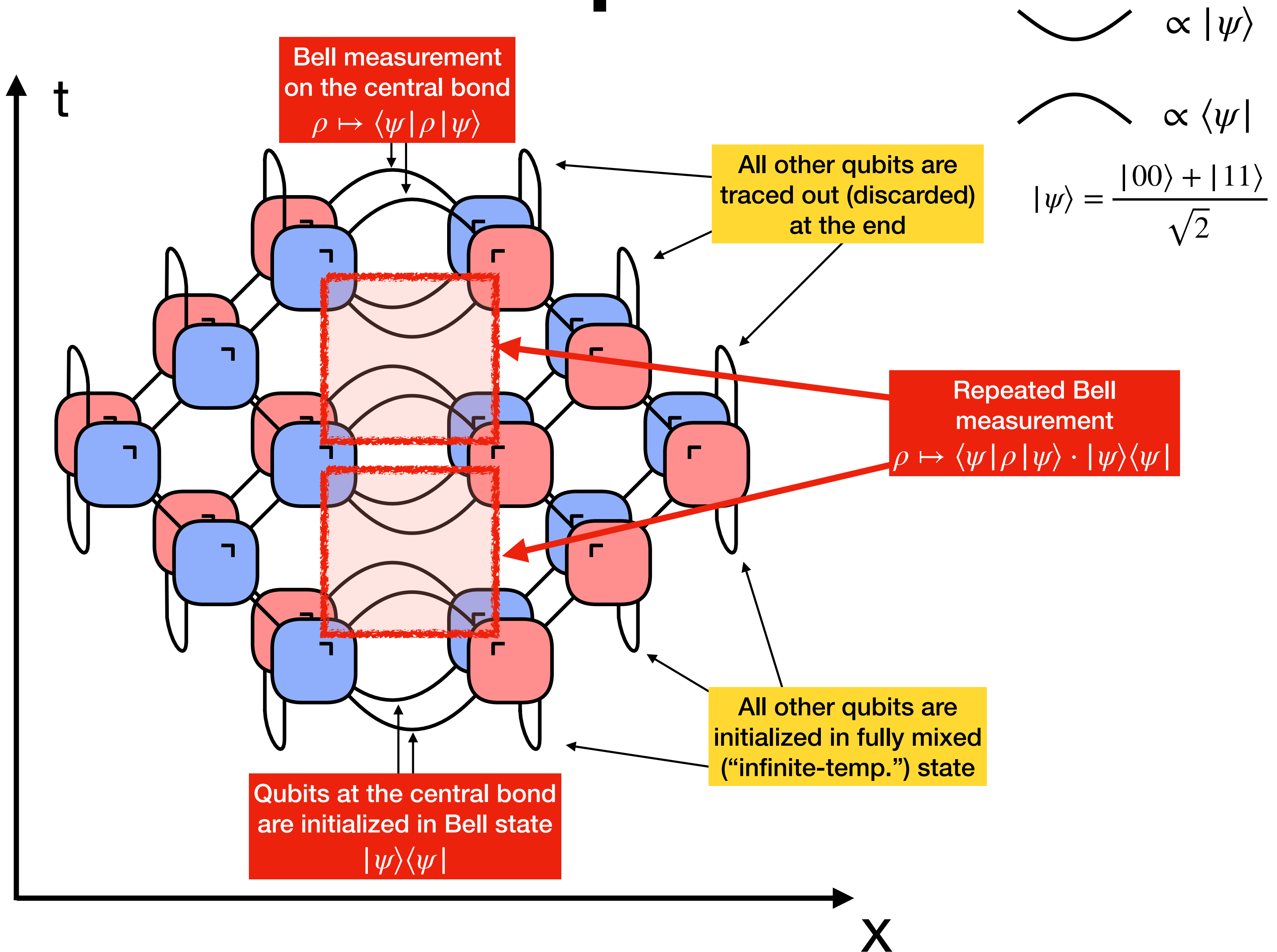
* [weird, multi-point]

Entanglement measurement protocol



Operational protocol:

- Initialize $I \otimes |\psi\rangle\langle\psi| \otimes I$
- Every time a Bell measurement yields “wrong” outcome (not $|\psi\rangle$), **stop** and record a **failure**
- When T consecutive Bell measurements yield “right” outcome, record a **success**

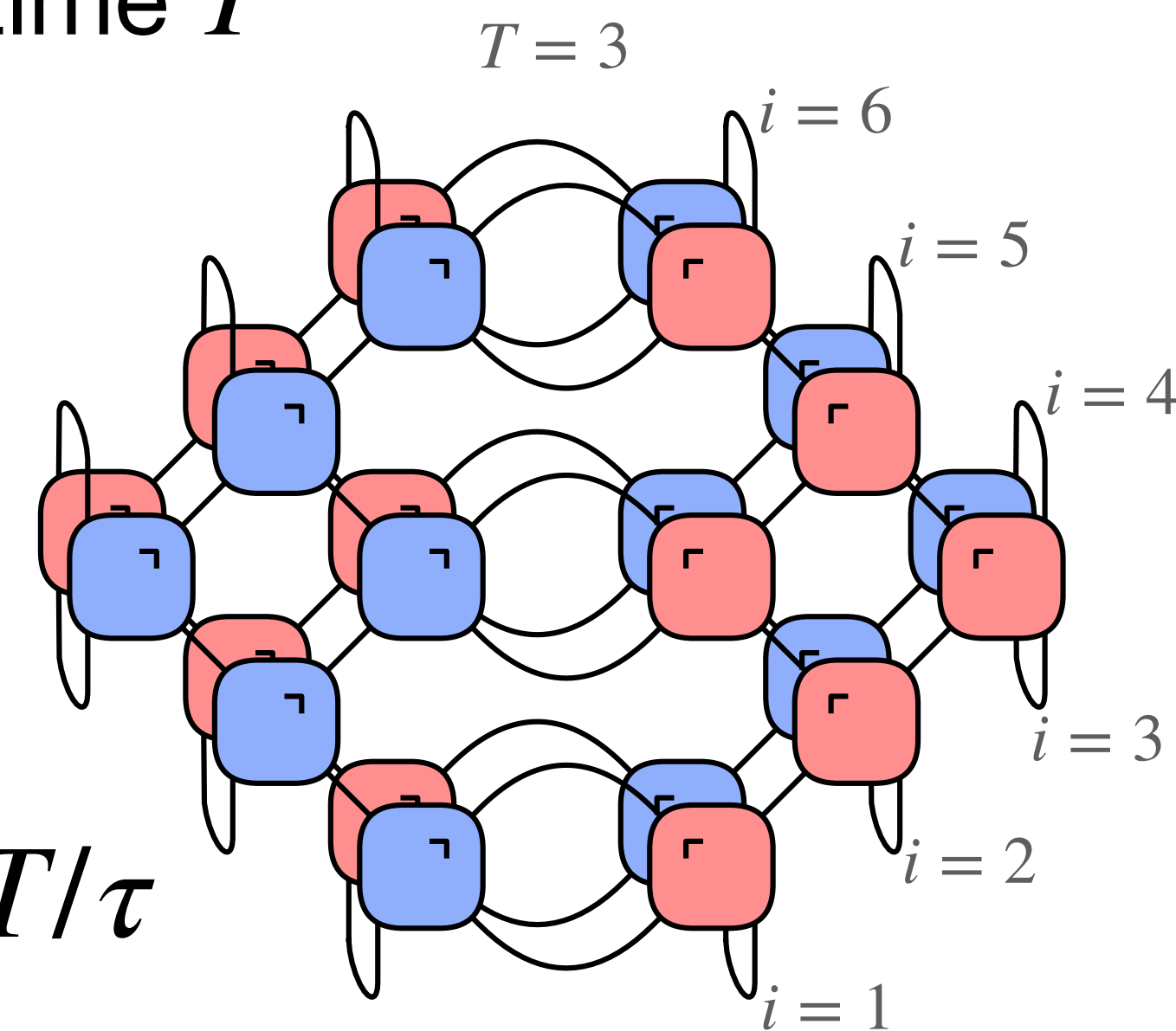


Entanglement measurement protocol

- Run protocol N_{tot} times; $N_+(T)$ counts n. of runs successful up to time at least T
- Measures purification dynamics of $2T$ “dual” qubits evolved for time T

$$\text{Tr}(\rho_{out}^2 |_{2T}) = N_+(T)/N_{tot}$$

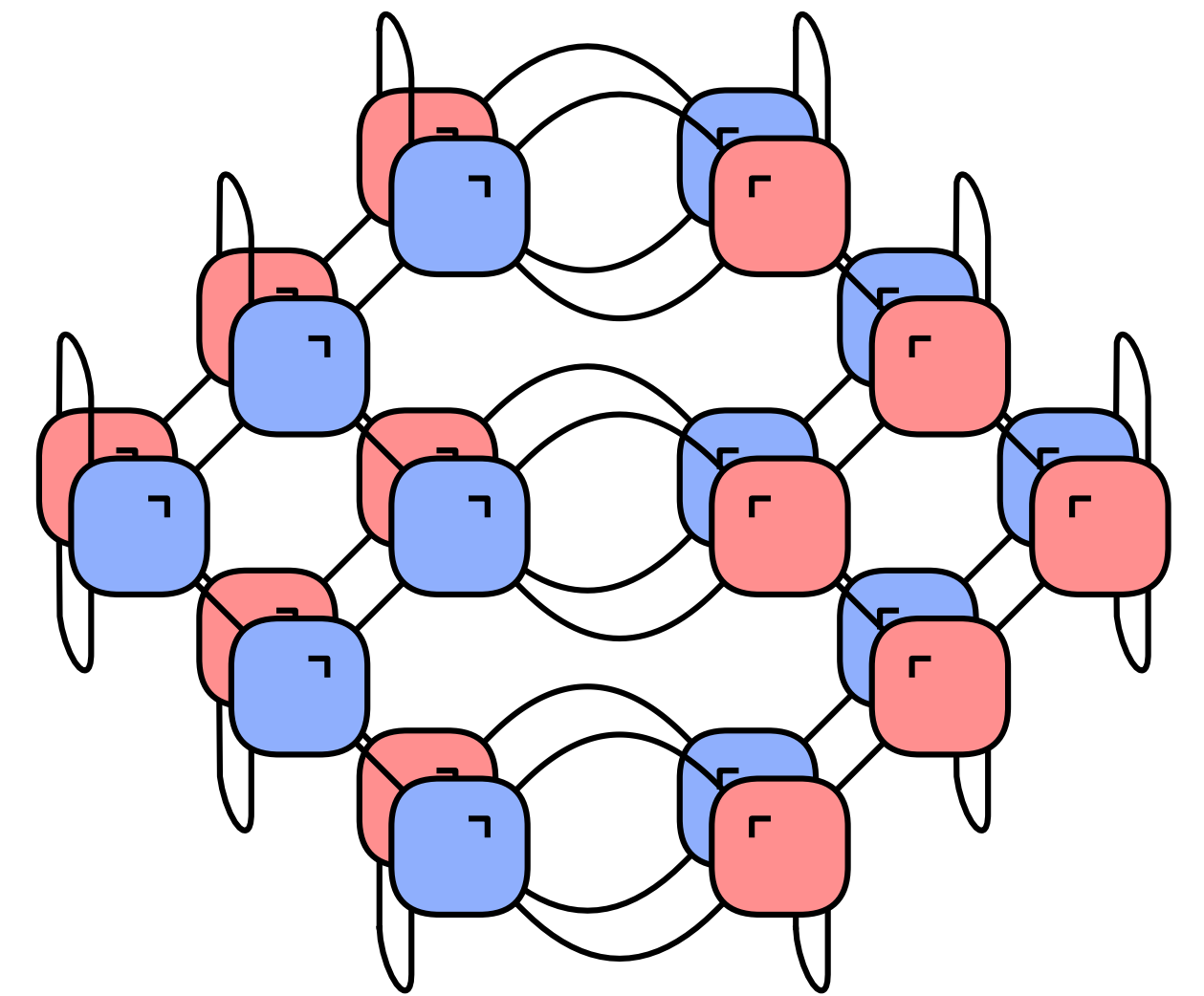
- Reasonable to expect $N_+(T) \sim Ne^{-T/\tau}$ in most cases...
 - ...i.e. a **mixed phase**: $\text{Tr}(\rho_{out}^2 |_{2T}) \sim e^{-T/\tau} \implies S_2 \sim T/\tau$
 - **Entropy density** $s_2 \simeq (2\tau)^{-1}$ directly measurable!



Comments

- Doesn't use **postselection**
 - “Failed” runs contribute to N_{tot} , provide key information
 - Unitary circuits automatically prepare two copies of same state
- Constraint on which models can be implemented
 - **Wide class** - as big as space of unitary circuits
- Constraint on **maximum time** of hybrid dynamics
- The state ρ_{out} is never prepared all at once in space, exists on a **temporal slice** of the circuit
 - But it *can* be prepared in space, too!

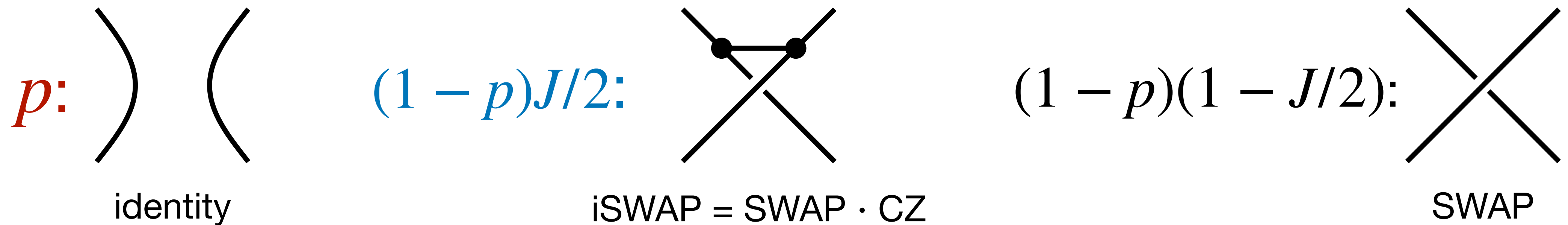
$$\text{Tr}(\rho_{out}^2 |_{2T}) = N_+(T)/N_{tot}$$



Purification phases and their signatures

Model

- Unitary Clifford circuit (efficient simulations); parameters $p, J \in [0,1]$
 - Two-qubit gates & their probabilities:

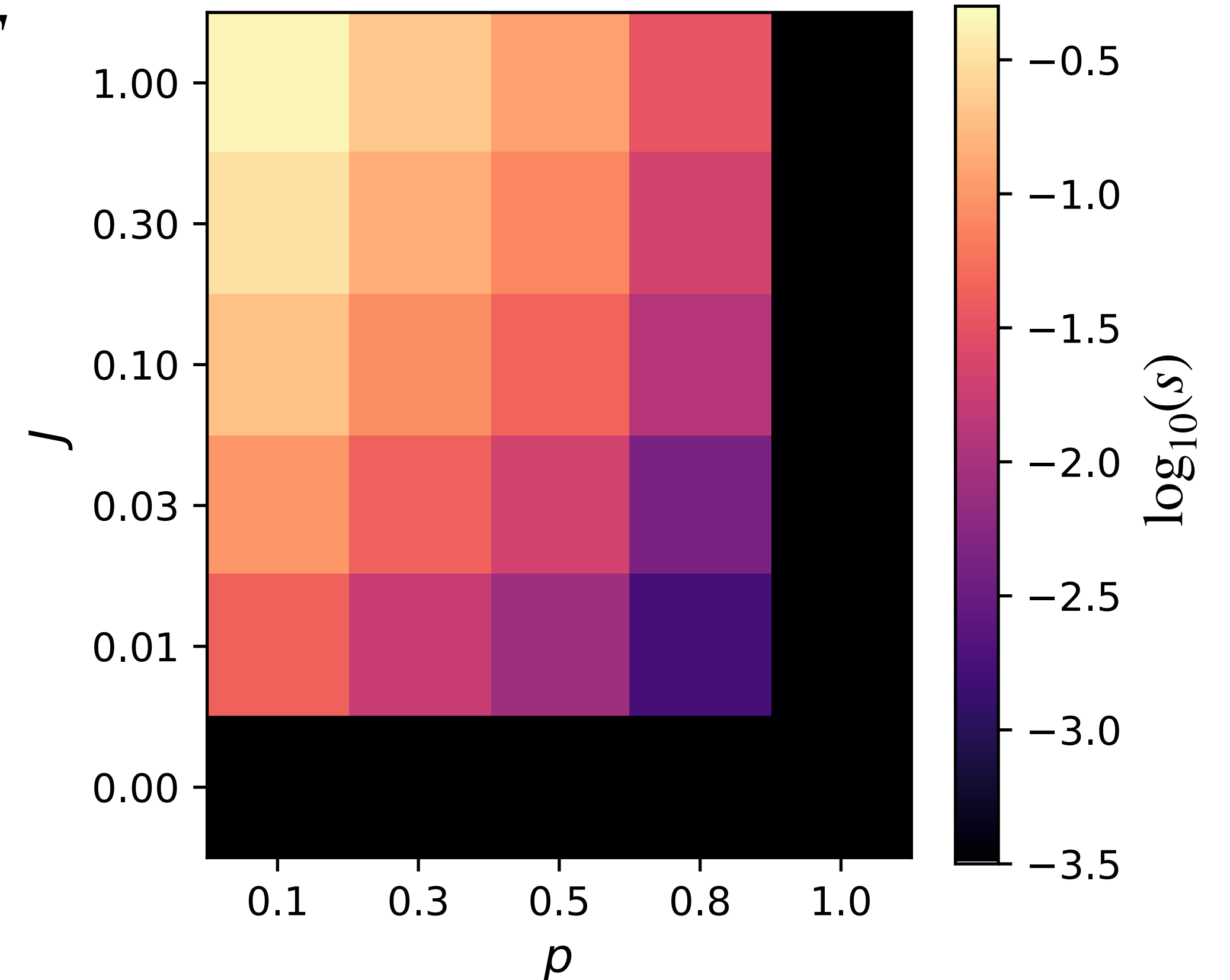
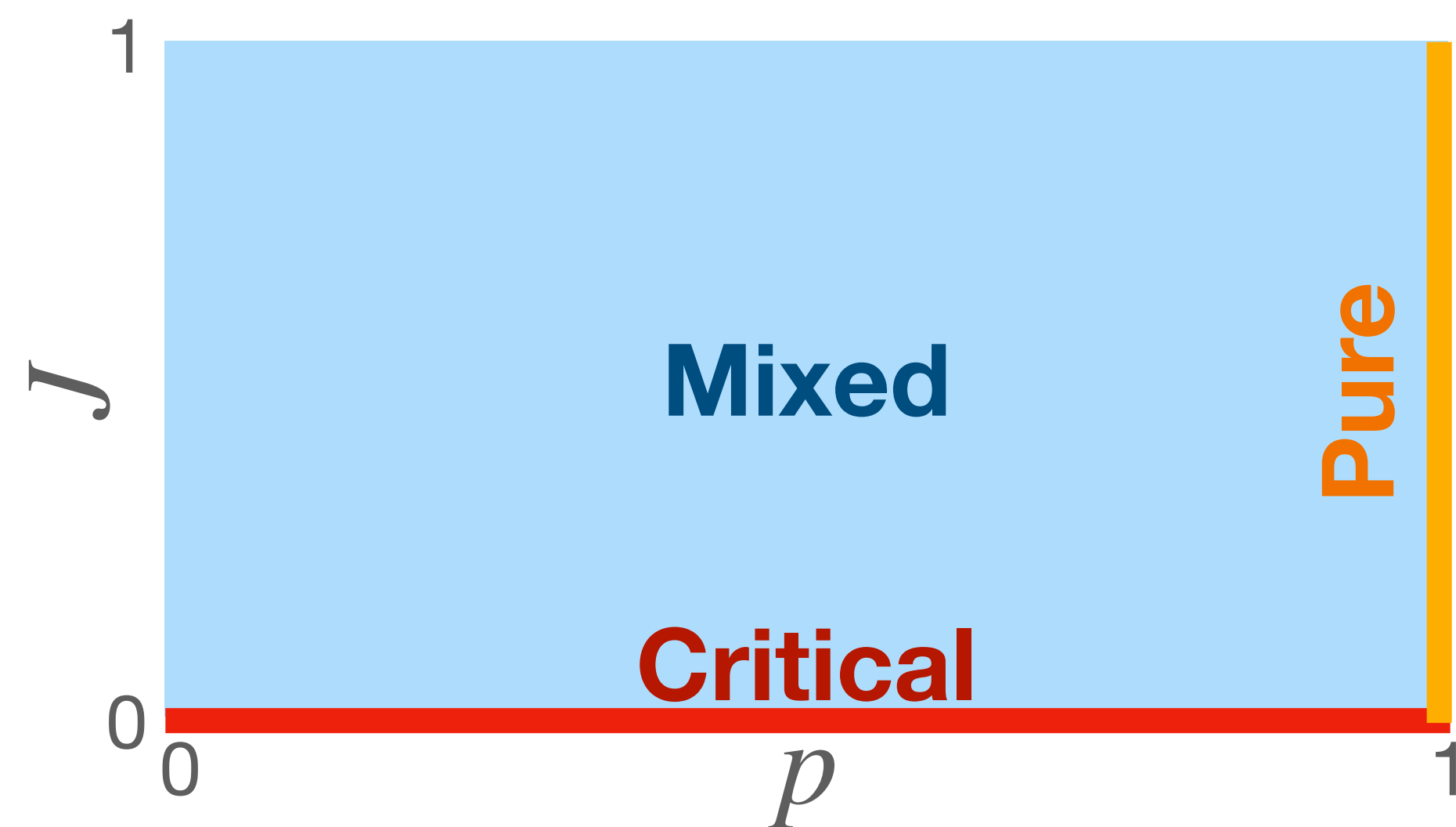


- + random 1-qubit Clifford gates
- p tunes **measurement rate** in dual circuit
- J tunes **interaction rate** ($J = 0$ is noninteracting SWAP circuit)

Purification phases and their signatures

Phase diagram

- Numerical simulation (stabilizer method, $2T \leq 4096$ qubits)
- Entropy density: $s = \lim_{T \rightarrow \infty} S(\rho_{out} |_{2T}) / 2T$
- Purification phase diagram



Purification phases and their signatures

Critical phase

- $J = 0$: SWAP circuit, **random walk** of qubit world lines

- Qubit worldlines can:

- **start & end in ρ_{out}** . Pure Bell pair, $S = 0$

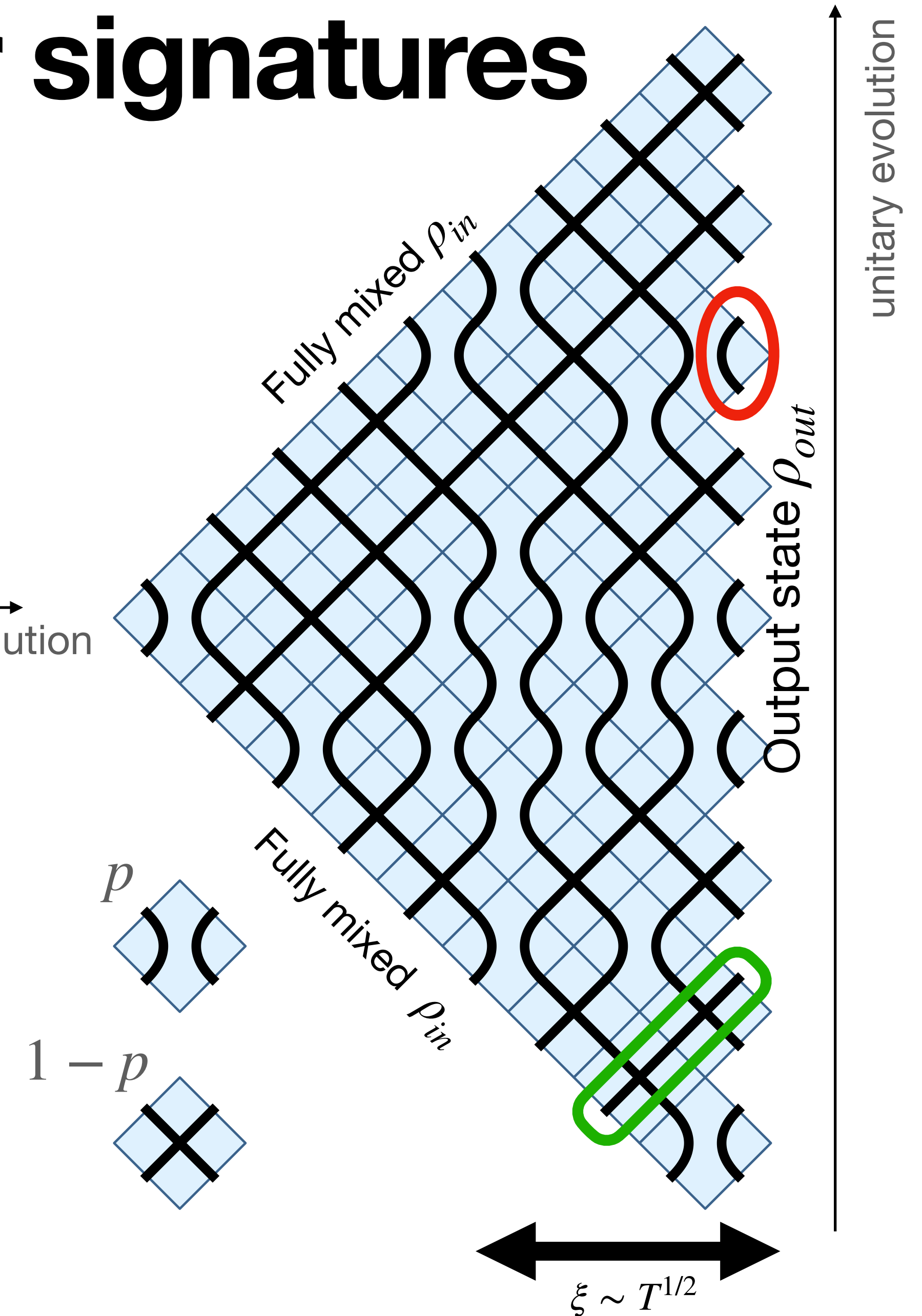
- **start in ρ_{in} , end in ρ_{out}** (or vice versa).
Fully mixed qubit, $S = 1$

- **Diffusion:** $S \sim T^{1/2}$

- **Result:** $\text{Tr}(\rho_{out}^2 |_{2T}) \sim e^{-\#\sqrt{T}}$

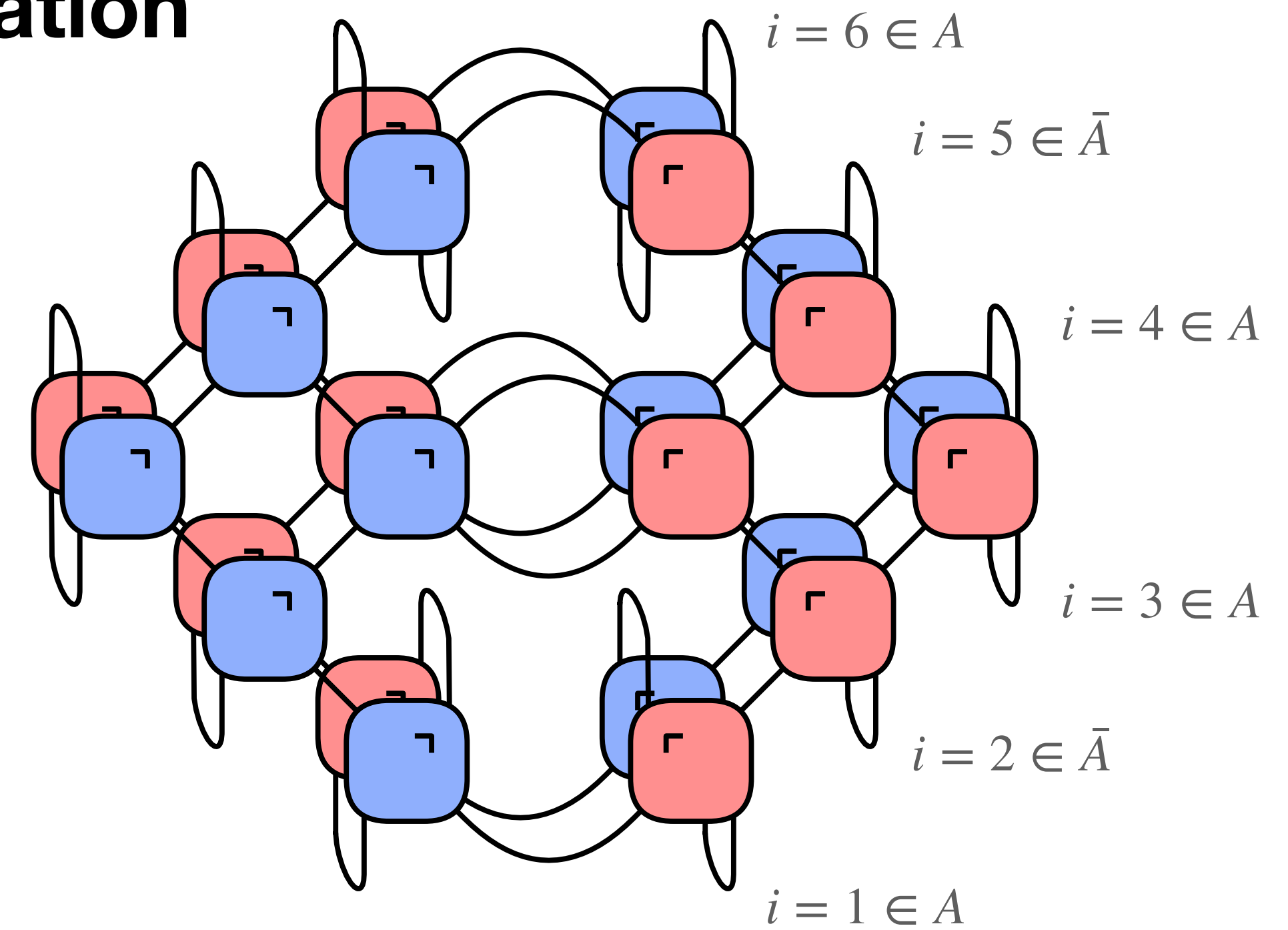
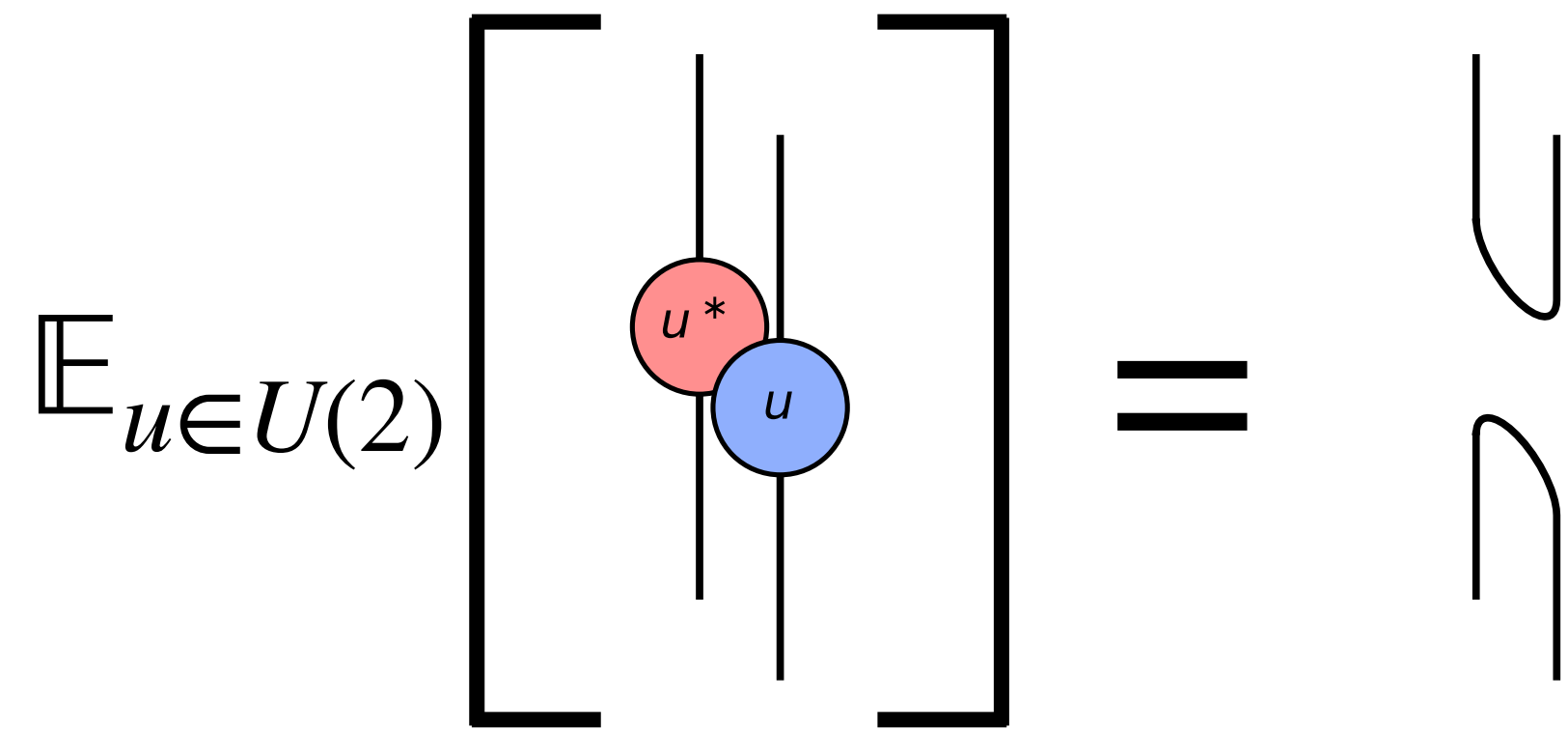
(cf. $e^{-T/\tau}$ for mixed phase)

→ hybrid evolution



Subsystem entropies

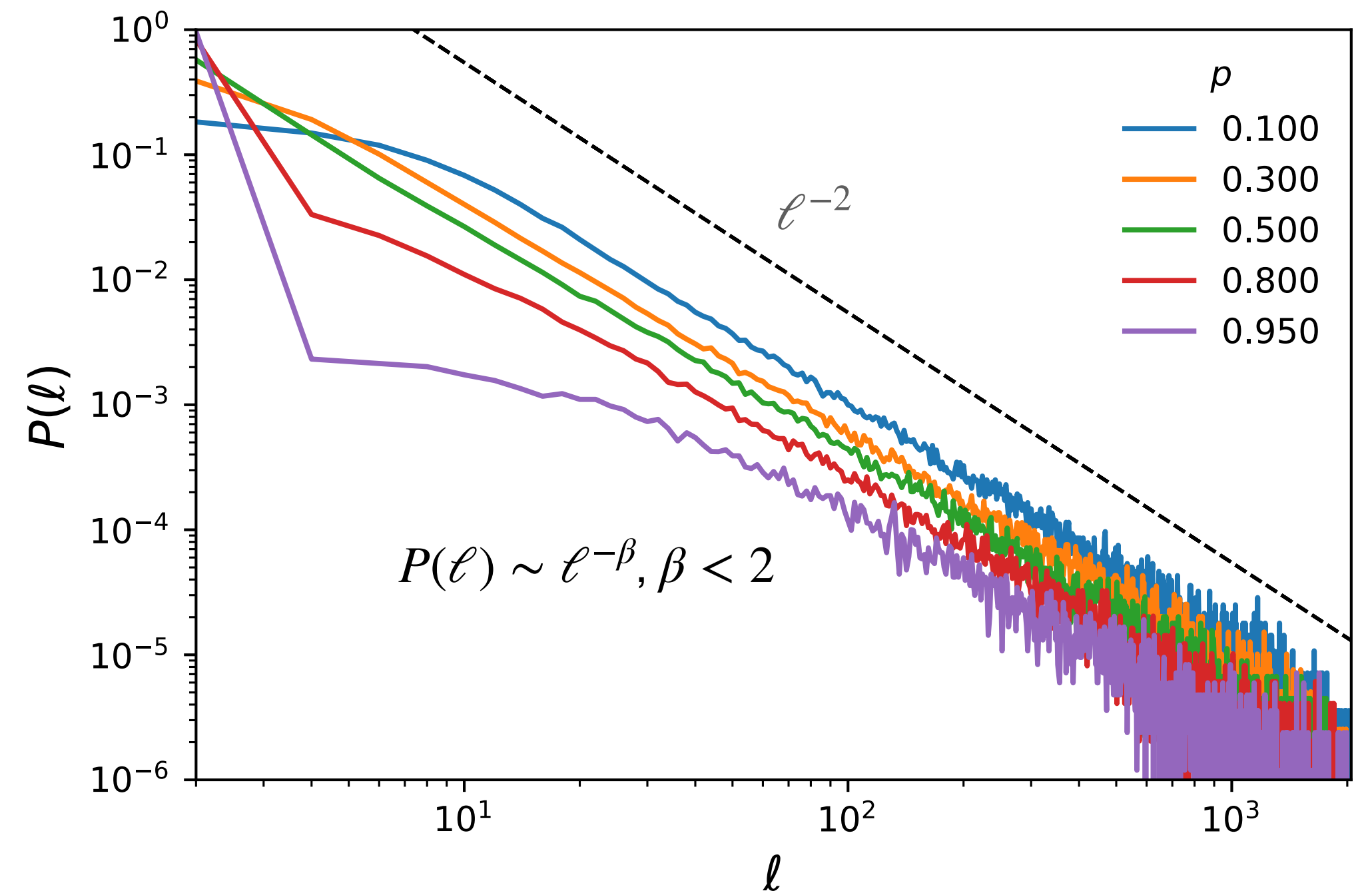
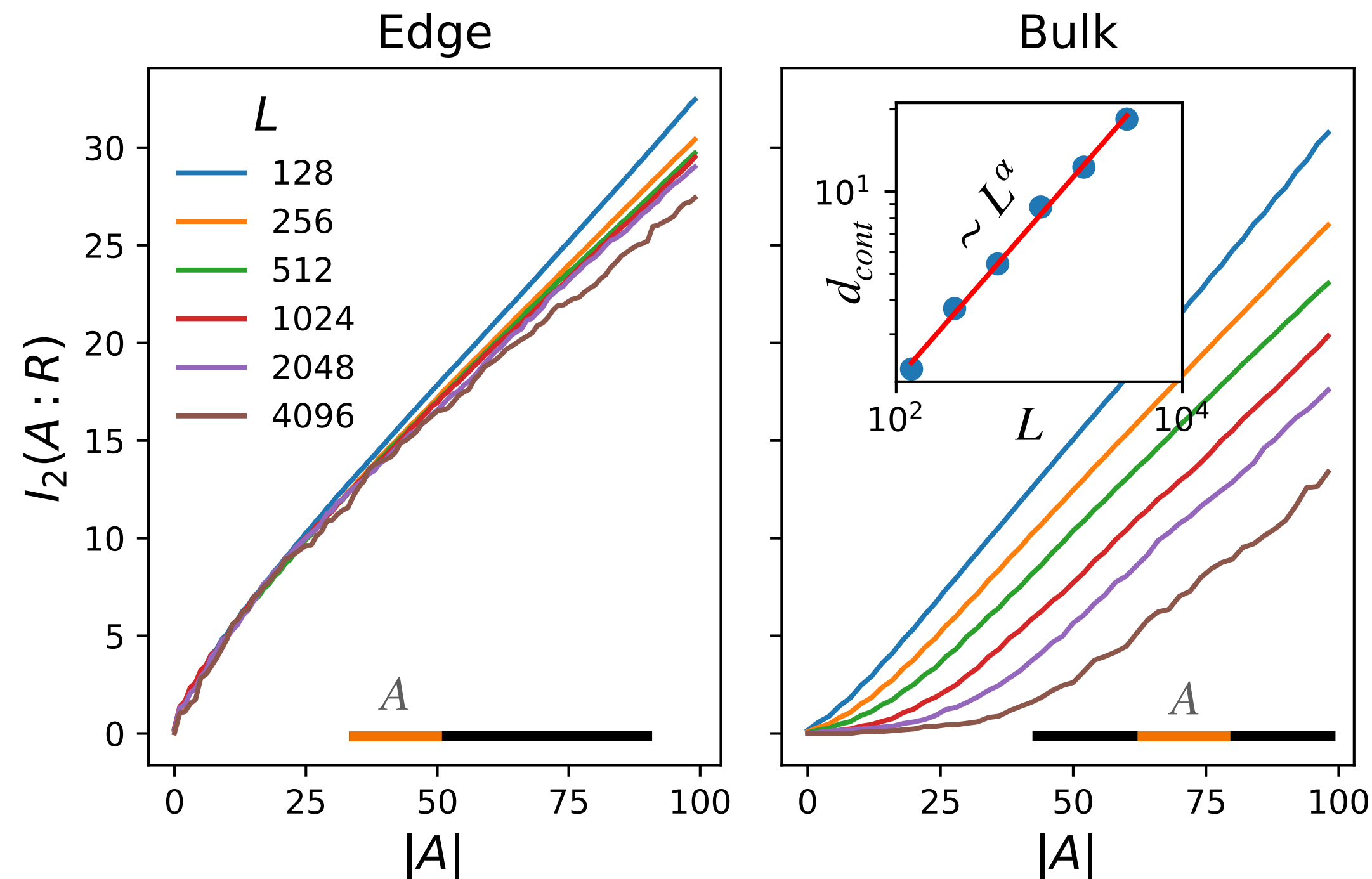
- Method can be generalized to measure $\text{Tr}(\rho_{out,A}^2)$ for **arbitrary subsystem A**
- How to “trace out” a physical qubit? **Depolarization**
 - Random Clifford gate + averaging many runs



- Application: direct access to (Renyi) mutual information, **code properties** in mixed phase

Quantum code properties of the mixed phase

- How special is the mixed phase in this model?
 - Power-law divergence in **mutual information** $I_2(A : \bar{A})$
 - Contiguous **code distance**: $d_{cont} \sim 1$ near edge, $d_{cont} \sim L^\alpha$ in bulk



Conclusion

- Spacetime duality enables forced measurements, sidesteps issue of **postselection**: novel route to investigate broad class of hybrid circuits
- Global & subsystem purity directly measurable: enables access to **quantum code properties** in the mixed phase
- Interesting **constraints** on entanglement/purification dynamics from unitarity in “transverse” direction
- Broader applications for **entanglement dynamics**?

Thank you!

MI, V. Khemani, arXiv:2010.xxxxx

Thanks to M. Gullans, D. Huse and
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State preparation protocol

