





# Andrea Puhm

 $\mathsf{CEA}/\mathsf{Saclay} \to \mathsf{UCSB}$ 

1210.6996

with Steve Avery and Borun Chowdhury

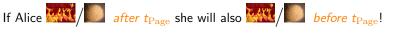
FuzzOrFire KITP, 30 August 2013

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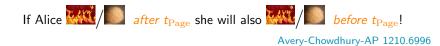
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Avery-Chowdhury-AP 1210.6996

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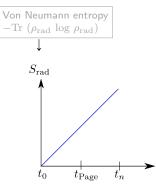
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See also recently: Bousso 1308.2665 w/ similar techniques, Marolf-Polchinski 1307.4706 w/ different techniques

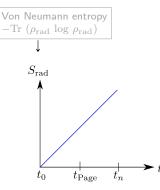
and before firewalls: van Raamsdonk et al 1206.1323 in context of AdS/CFT



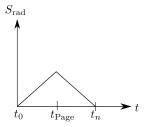


Final radiation state has  $S_{rad} \neq 0$ : Evaporation via max entangled Hawking pairs ( $S_{pair} = 0$ ) is non-unitary!



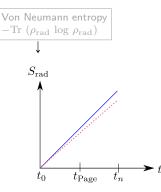


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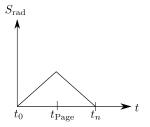


Any normal evaporating body in a pure intial state has a pure final state:  $S_{\rm rad} = 0$ .

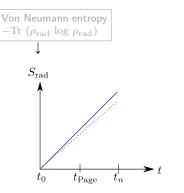




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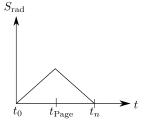
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Purity of final radiation requires  $S_{\text{pair}} \neq 0 \rightarrow \text{hair no later than } t_{\text{Page}}!$ 

Page curve for typical states:



#### I. Outside:

*Purity* of the final radiation state: pair  $\rightarrow$  hair *no later than*  $t_{Page}$ .

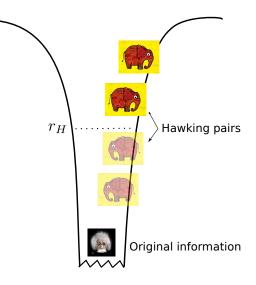
#### II. Infalling:



#### Questions:

Sharper information loss statement from conditions beyond purity? Do young black holes have hair?

## When and how does the original information get out?



- 1 Purity: Pure states evolve into pure states.
- **2** Linearity: The map between initial and final states is linear.
- **3** Preservation of norm: Evolution of states preserves norm.
- **4** Invertibility: The map of initial state to the radiation is invertible.

Do these additional conditions make the statement of information loss in Hawking's process more precise?

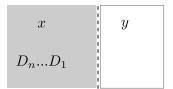
We will model evaporation via qubits and use the following rules:

- The dimension of the physical Hilbert space is constant.
- Fundamental properties of fast scramblers do not depend on the nature of formation.
- The evaporation process can be described within a local framework.
- The general dynamics is state independent.

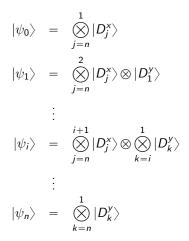
## The 'moving bit' model

#### Unitary model of evaporation:

Move qubits **D** from **x** to **y**.



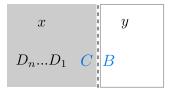
Typical state:  $S_{D_i} \neq 0$ .



## The Hawking model

#### Non-unitary model of evaporation:

Create qubits C, B at interface. Move C to x and B to y.



Maximally entangled:  $S_{B_iC_i} = 0$ .

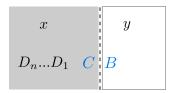
C annihilate with D.

$$\begin{aligned} |\psi_{0}\rangle &= \bigotimes_{j=n}^{1} |D_{j}^{x}\rangle \\ |\psi_{1}\rangle &= \bigotimes_{j=n}^{1} |D_{j}^{x}\rangle \otimes |C_{1}^{x}\rangle \otimes |B_{1}^{y}\rangle \\ &\vdots \\ |\psi_{i}\rangle &= \bigotimes_{j=n}^{1} |D_{j}^{x}\rangle \otimes \bigotimes_{k=i}^{1} \left(|C_{k}^{x}\rangle \otimes |B_{k}^{y}\rangle\right) \\ &\vdots \\ |\psi_{n}\rangle &= \bigotimes_{j=n}^{1} |D_{j}^{x}\rangle \otimes \bigotimes_{k=n}^{1} \left(|C_{k}^{x}\rangle \otimes |B_{k}^{y}\rangle\right) \end{aligned}$$

## General model framework

#### Potentially non-unitary model:

Create qubits **C**, **B** at interface. **D**, **C**, **B**  $\xrightarrow{\mathcal{U}}$   $\widetilde{D}$ ,  $\widetilde{C}$ ,  $\widetilde{B}$ . Move  $\widetilde{C}$  to **x** and  $\widetilde{B}$  to **y**.



Identifying two auxiliary qubits.

Tracing out auxiliary qubits.

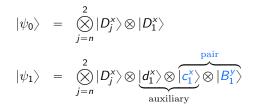
$$\begin{aligned} |\psi_{0}\rangle &= \bigotimes_{j=n}^{1} |D_{j}^{x}\rangle \\ |\psi_{1}\rangle &= \bigotimes_{j=n}^{1} |\tilde{D}_{j}^{x}\rangle \otimes |\tilde{C}_{1}^{x}\rangle \otimes |\tilde{B}_{1}^{y}\rangle \\ &\vdots \\ |\psi_{i}\rangle &= \bigotimes_{j=n}^{1} |\tilde{D}_{j}^{x}\rangle \otimes \bigotimes_{k=i}^{1} \left(|\tilde{C}_{k}^{x}\rangle \otimes |\tilde{B}_{k}^{y}\rangle\right) \\ &\vdots \\ |\psi_{n}\rangle &= \bigotimes_{i=n}^{1} |\tilde{D}_{j}^{x}\rangle \otimes \bigotimes_{k=n}^{1} \left(|\tilde{C}_{k}^{x}\rangle \otimes |\tilde{B}_{k}^{y}\rangle\right) \end{aligned}$$

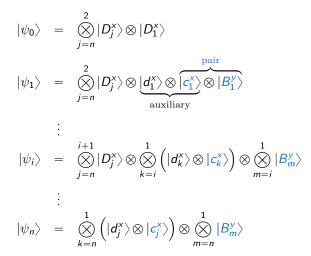
$$|\psi_0\rangle = \bigotimes_{j=n}^2 |D_j^x\rangle \otimes |D_1^x\rangle$$

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$$|\psi_i\rangle = \bigotimes_{j=n}^{i+1} |D_j^x\rangle \otimes \bigotimes_{k=i}^1 \left( |\phi\rangle \otimes |\phi\rangle \right) \otimes \bigotimes_{m=i}^1 |B_m^y\rangle$$
  
:

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Evolution of basis of states:

$$\begin{split} |\psi_0\rangle &= \bigotimes_{j=n}^2 |D_j^x\rangle \otimes |D_1^x\rangle \\ |\psi_1\rangle &= \bigotimes_{j=n}^2 |D_j^x\rangle \otimes \qquad |B_1^y\rangle \end{split}$$

$$\begin{array}{l} \vdots \\ |\psi_i\rangle &= \bigotimes_{j=n}^{i+1} |D_j^x\rangle \otimes \\ \vdots \end{array}$$

$$\bigotimes_{m=i}^{1} |B_{m}^{y}\rangle$$

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 $B_i = D_i$ : rewriting of the moving bit model via pairs  $\Rightarrow$  unitary.

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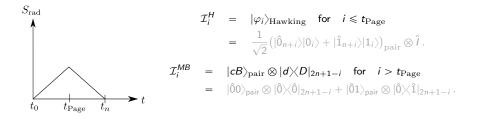
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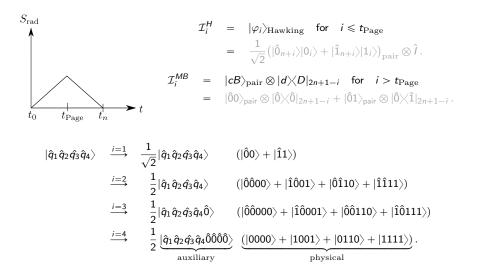
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 $\Rightarrow$  Unitarity requires  $S_{\text{pair}} \neq 0$  at every step for typical states.

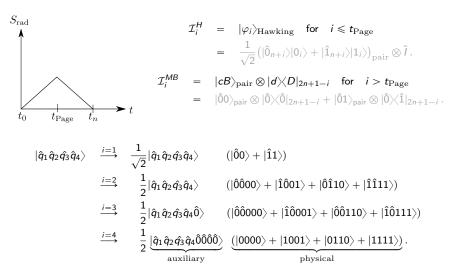
# Hawking + 'moving bit' model



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Final radiation state pure but independent of initial state! Information of the original state never came out!

### Where is the information?

Tracking state:

$$|T\rangle = \frac{1}{2^{\frac{n}{2}}}\sum_{i=1}^{2^{n}}|\widehat{bh}_{i}\rangle|\widetilde{r}_{i}\rangle,$$

 $\{|\hat{w}_i\rangle\}, \{|\tilde{r}_i\rangle\}$  are orthonormal bases for the *n*-qubit initial state and reference system.

- The decrease of the entanglement entropy of the black hole quantifies how much information has left the black hole.
- The entanglement entropy of the external radiation quantifies how much information is now in the radiation.
- Mutual information vanishes in qubit framework. (Moreover conjectured to be zero for black hole evaporation [Giddings,Shi]).
  → Information loss occurs in wrong identification of auxiliary qubits.

## Losing a qubit

How much information is lost if we trace out the wrong qubit?

Initial tracking state:

$$|\mathcal{T}_0\rangle = \frac{1}{2} \Big( |\hat{0}\hat{0}\rangle|\tilde{0}\tilde{0}\rangle + |\hat{0}\hat{1}\rangle|\tilde{0}\tilde{1}\rangle + |\hat{1}\hat{0}\rangle|\tilde{1}\tilde{0}\rangle + |\hat{1}\hat{1}\rangle|\tilde{1}\tilde{1}\rangle \Big) \,,$$

Evolution operator:

$$\mathcal{I}_i = \left( |\hat{0}\hat{0}\rangle|0\rangle\langle\hat{0}|_{(3-i)} + |\hat{0}\hat{0}\rangle|1\rangle\langle\hat{1}|_{(3-i)} \right)_{D_{i-1}} \otimes I_{A_{i-1}} \otimes I_{A_{i-1}} \quad \text{for} \quad i = 1, 2 \,.$$

$$i = 1: \quad |\mathcal{T}_1\rangle = \mathcal{I}_1 |\mathcal{T}_0\rangle = |\hat{0}\hat{0}\rangle \otimes \frac{1}{2} \Big( |\hat{0}0\rangle |\tilde{0}\tilde{0}\rangle + |\hat{0}1\rangle |\tilde{0}\tilde{1}\rangle + |\hat{1}0\rangle |\tilde{1}\tilde{0}\rangle + |\hat{1}1\rangle |\tilde{1}\tilde{1}\rangle \Big) \,.$$

Info loss from wrong identification of auxiliary space with bleaching space:

$$\begin{split} S(\mathcal{A}_1) &= -\mathrm{Tr}(\rho_1^{aux}\mathrm{log}\rho_1^{aux}) \stackrel{}{\underset{\uparrow}{=}} \log 2 \, . \\ \rho_1^{aux} &= |\hat{0}\rangle\langle \hat{0}| \otimes \frac{1}{2} \left( |\hat{0}\rangle\langle \hat{0}| + |\hat{1}\rangle\langle \hat{1}| \right) \end{split}$$

Simple example: maximum amount of information loss on a qubit. General:  $0 < S(A_i) \le \log 2$  for each erroneously traced out physical qubit. Black hole evaporation via Hawking pair production is non-unitary. Need deviation from vacuum at horizon no later than  $t_{Page}$ . ( $\leftarrow$  Purity )

Hawking, Mathur, Avery, AMPS

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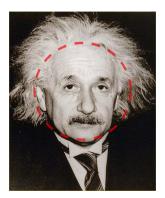
Hawking, Mathur, Avery, AMPS

- Purity does not guarantee that any information can be recovered.
- Information of original state out in every step. (← Invertibility)
- $\Rightarrow$  Need horizon-scale structure/hair before  $t_{Page}$ !

Avery-Chowdhury-AP

# Also young black holes have hair!





## Comments

Unitary, local, state-independent framework: QG kicks in at horizon-scale!

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- The real world is not supersymmetric, so why care? There is a mechanism for the supersymmetric and non-supersymmetric extremal case (Bena-Warner et al) which continues to work when perturbing away from extremality (Bena-AP-Vercnocke)! For near-extremal it is the same mechanism as used in uplifting AdS to dS via antibranes (gazillions of cosmology papers)!

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Unexpected but interesting: learn something deep about QG.