

Cosmological Information from Non-linear Weak Lensing

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Outline

- ✱ **Motivation:** is there significant information in the non-linear regime in WL beyond “usual” LSS statistics ?
 1. Statistical power of dN/dz (cluster counts)
 2. Complementarity of dN/dz and P_l (power spectrum)
- ✱ **Non-linear statistics robust to selection effects:**
 1. One-point function of convergence (analytic)
 2. Peak counts (including non-cluster peaks)
 3. Minkowski Functionals (V_0, V_1, V_2)

Motivation: Clusters

- ✦ **Exquisite statistical error forecast**

from large SZ, XR, WL surveys – but have to know cluster mass M

- ✦ **“Phenomenological” self-calibration**

e.g. $M=M_0(M/M_0)^\alpha(1+z)^\beta$ or more complicated forms

use combination of mass-dependent observables $dN/dz + P(k)$

(e.g. Majumdar & Mohr 2004; Wang et al. 2004)

- ✦ **“Physical” self-calibration**

use parameterized cluster structure model to relate

SZ, XR and WL observables (e.g. Shang, Haiman & Verde 2009; Younger et al. 2006)

- ✦ **WL surveys**

200,000 (?!) clusters in 20,000 sq.deg. – problem: projections/selection

Motivation: WL surveys

- ★ Several large ($\geq 1,000$ sq. deg) WL surveys forthcoming:
(e.g. Pan-STARRS, KIDS, DES, LSST, Euclid)
- ★ Shear power spectrum and related large-scale statistics
(e.g. Kaiser 1992; Jain & Seljak 1997; Hu 1999, 2002;
Huterer 2002; Refregier et al. 2004; Abazajian & Dodelson 2003;
Takada & Jain 2004; Song & Knox 2004;)
E.g. $\sigma(w_0)=0.06$; $\sigma(w_a)=0.1$ from 11-parameter fit to
tomographic shear power spectrum (LSST) + *Planck*
- ★ Comparable statistical errors from cluster number counts
(e.g. Wang et al. 2004, 2005; Fang & Haiman 2007;
Takada & Bridle 2007; Marian & Bernstein 2006, 2008)
E.g. $\sigma(w_0)=0.04$; $\sigma(w_a)=0.09$ from 7-parameter fit to
 $\sim 200,000$ shear-selected cluster counts (LSST) + *Planck*

Motivation: Clusters vs Peaks

- ✱ Cluster counts and shear power spectrum can be considered independent observables – high synergy
Covariance changes parameter-estimates by $< \text{few } \%$
(Fang & Haiman 2007; Takada & Bridle 2007)
- ✱ However, selection effects are (probably) a showstopper in a WL survey alone, due to projection effects
Filter-dependent trade-off between *completeness* and *purity*: “best compromise” values are $\sim 70\%$ for both
(e.g. White et al. 2002; Hamana et al. 2005; Hennawi & Spergel 2005)
- ✱ Why not define observable immune to projection effects?
historical reason: cosmology-dependence of halo mass function calculable from “Press-Schechter”

Peak counts

- ★ A simple statistic: # of shear peaks, regardless of whether or not they correspond to true bound objects as a function of height, redshift and angular size

Kratochvil, Haiman, Hui & May (2010)

- ★ Fundamental questions about “false” (non-cluster) peaks:
 1. How does $N(\text{peak})$ depend on cosmology ?
 2. What is the field-to-field variance $\Delta N(\text{peak})$?
- ★ Requires simulations

N-body Simulation Details

Kratochvil, Haiman, Hui & May (2010)

- pure DM (no baryons, neutrinos, or radiation)
- public code GADGET-2, modified to handle $w_0 \neq -1$
- fiducial Λ CDM cosmology from *WMAP*:
($w_0, \Omega_\Lambda, \Omega_m, H_0, \sigma_8, n$) = (-1.0, 0.74, 0.26, 0.72, 0.79, 1.0)
- fix *primordial* amplitude $\Delta^2_R = 2.41 \times 10^{-9}$ at $k = 0.002 \text{ Mpc}^{-1}$
($\sigma_8=0.79$ vs. 0.75)
- two alternative cosmologies, differ only in $w_0 = -0.8$ (or -1.2)
- 512^3 box, size $200h^{-1} \text{ Mpc}$, $z_{in}=60$, $M_{DM}=4.3 \times 10^9 M_\odot$
- gravitational softening length $\varepsilon_{pl} = 7.5h^{-1} \text{ kpc}$
- output particle positions every $70h^{-1}$ comoving Mpc
- runs at NSF TeraGrid and IBM Blue Gene / Brookhaven

Simulating Weak Lensing Maps

☀ Ray-tracing

- compute 2D potential (2048×2048) in each lens plane
- implement algorithm to follow rays (Hamana & Mellier 2001)
- compute shear (γ), convergence (κ) and reduced shear (μ)

☀ Mock “observational” parameters

- gaussian 1-component shear noise from intrinsic ellipticity:
 $\sigma_\gamma = 0.15 + 0.035z$ (Song & Knox 2004)
- $n_{\text{gal}} = 15 \text{ arcmin}^{-2}$ background galaxies, at $z_s = 1, 1.5, \text{ and } 2$
- smooth κ -map with 2D finite Gaussian 0.25 - 30 arcmin
- use $3 \times 3 \text{ deg}^2$ smoothed convergence maps

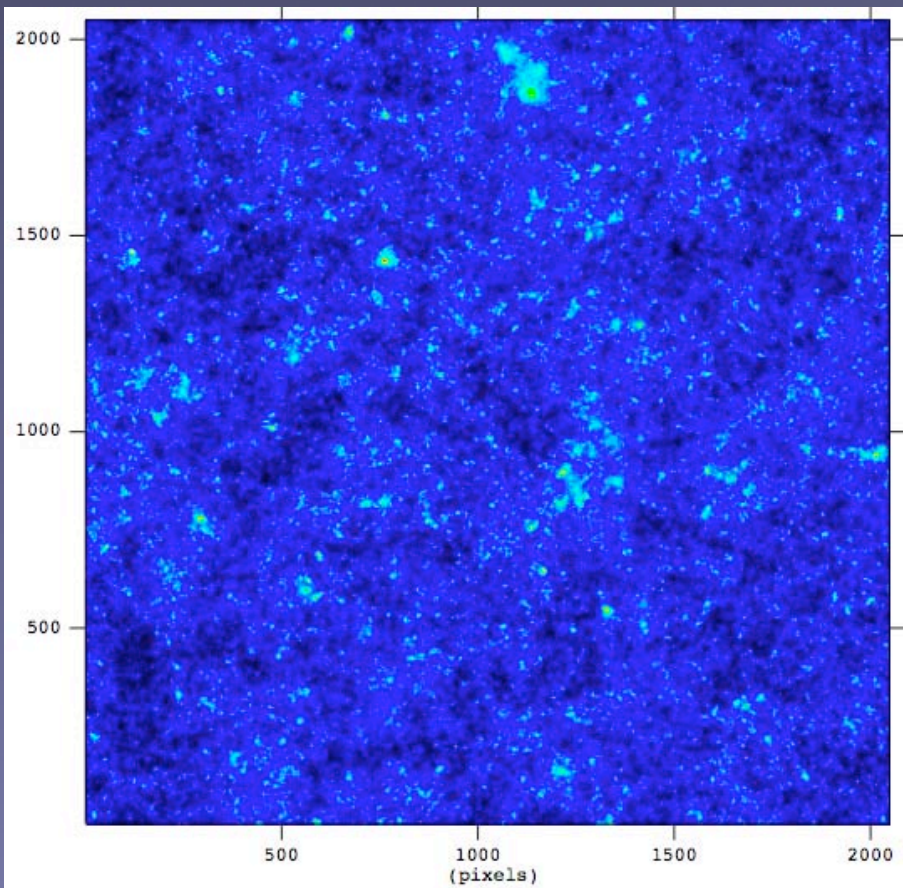
☀ Identifying peaks

- find all local maxima, record their height κ_{peak}

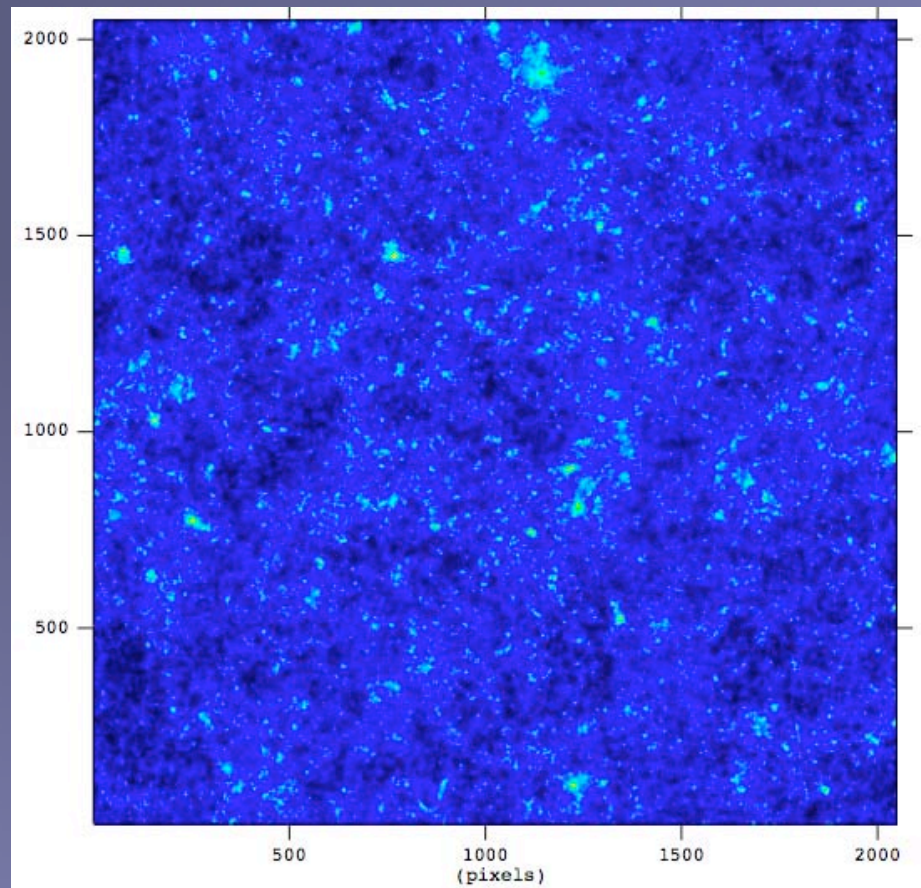
Example

raw convergence map (3×3 deg²; 2048 \times 2048 pixels)

$w=-1$



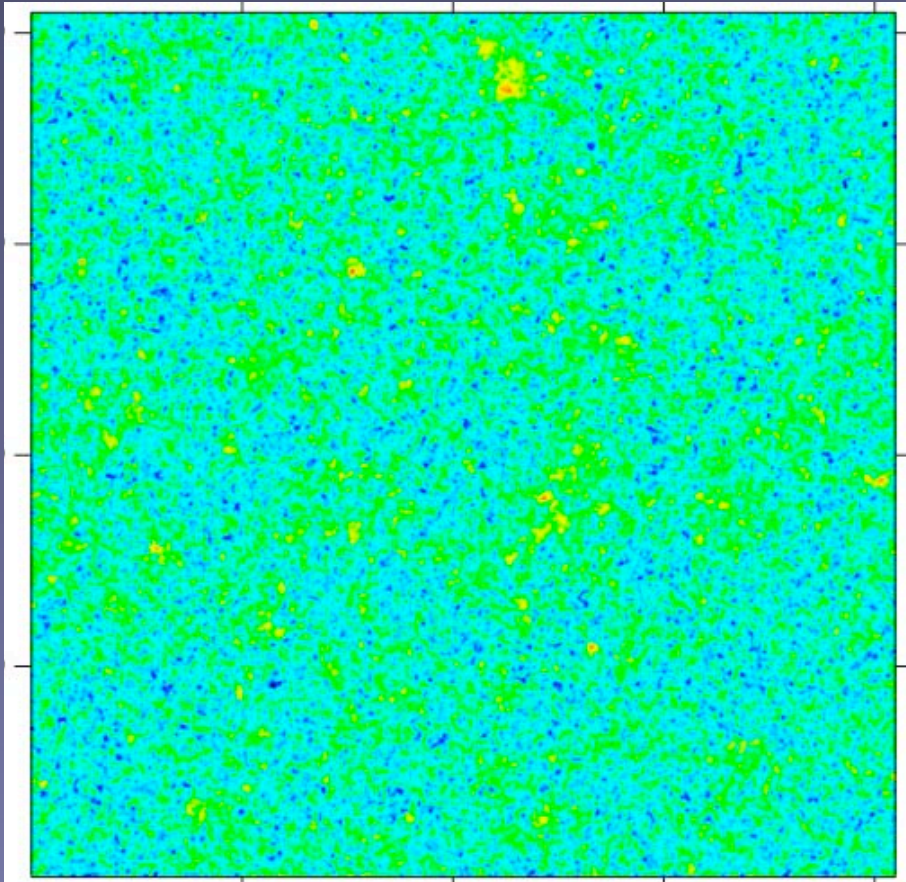
$w=-0.8$



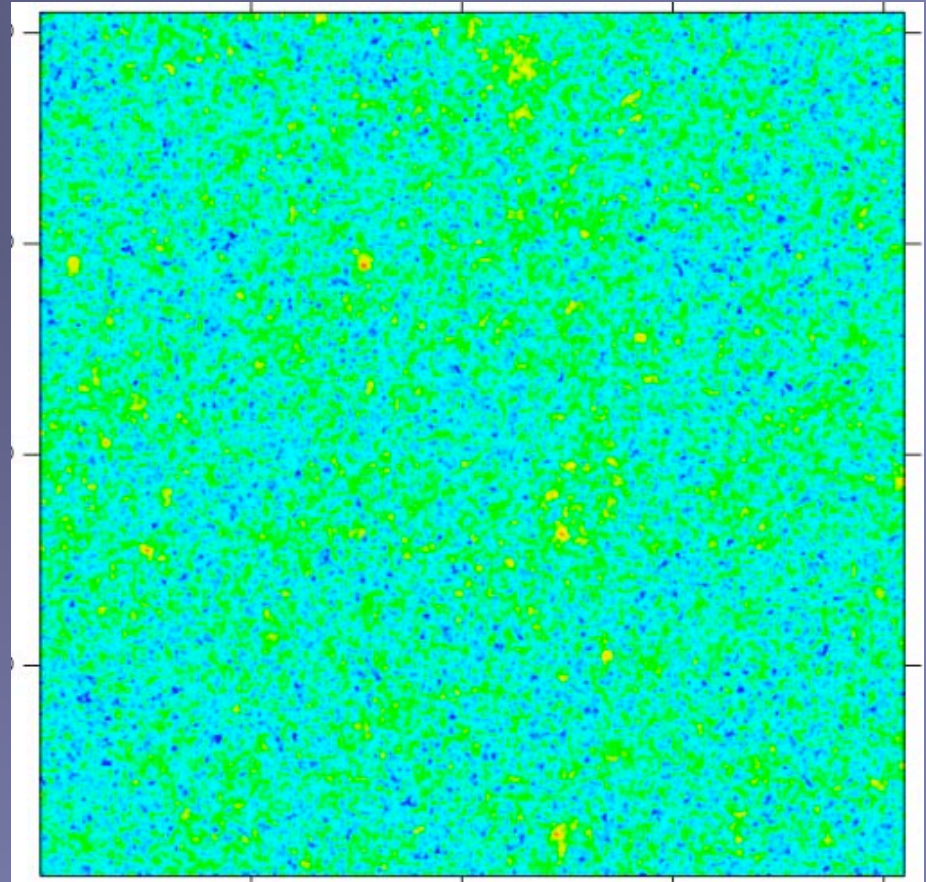
Example

convergence map with noise and 1-arcmin gaussian smoothing

$w=-1$



$w=-0.8$



Basic Results

- 3×3 deg² field, smoothing with 1-arcmin, galaxies at $z_s=2$

- Expectations based on clusters with $\kappa_G \geq 4.5\sigma_\kappa$

(Fang & Haiman 2007)

$$N(\text{clusters}) = 150 \pm 25 \quad \text{for } w=-1$$

$$N(\text{clusters}) = 103 \pm 21 \quad \text{for } w=-0.8$$

→ $S/N \approx 2\sigma$ mostly coming from change in σ_8

- Total peak counts above same threshold [w/no noise]

$$N(\text{peaks}) = 576 \pm 86 \quad [230 \pm 42] \quad \text{for } w=-1$$

$$N(\text{peaks}) = 547 \pm 85 \quad [186 \pm 37] \quad \text{for } w=-0.8$$

→ $S/N \approx 0.3\sigma$: (i) smaller difference, (ii) larger variance

- Total peak counts (all peaks):

$$N(\text{peaks}) = 11,622 \pm 62 \quad \text{for } w=-1$$

$$N(\text{peaks}) = 11,562 \pm 62 \quad \text{for } w=-0.8$$

Statistical Methodology

- ★ Covariance matrix for number of binned, tomographic peaks:

$$\mathbf{C}_{i,j}^m \equiv \frac{1}{R-1} \sum_{r=1}^R \left[N_i^m(r) - \bar{N}_i^m \right] \left[N_j^m(r) - \bar{N}_j^m \right]$$

- $R=500$ realizations in cosmology m (rotate/shift/slice box)
- $i = 15$ (height) \times 3 (source redshift) = 45 bins

- ★ Compute “ χ^2 ” between test (m) and fiducial (n) cosmology:

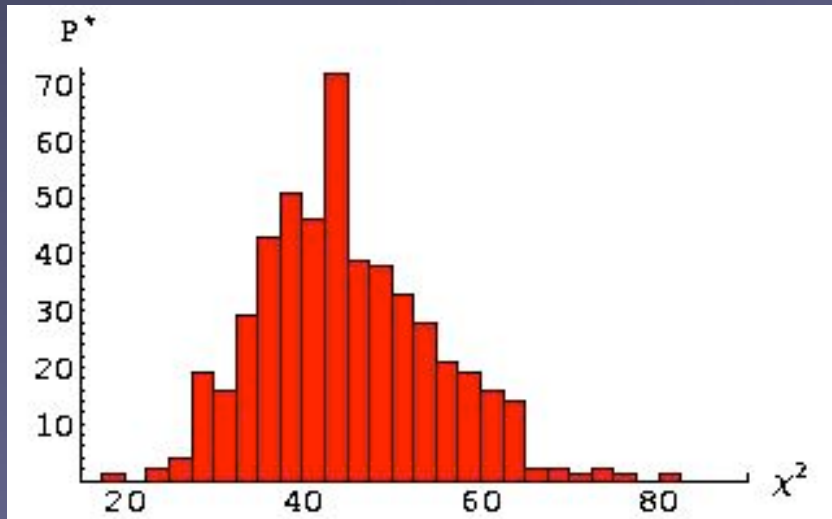
$$\chi^{2;m,n}(r) \equiv \sum_{i,j=1}^{45} \left[N_i^m(r) - \bar{N}_i^n \right] \left[\mathbf{C}_{i,j}^n \right]^{-1} \left[N_j^m(r) - \bar{N}_j^n \right]$$

- ★ Compute likelihood at which cosmology m can be distinguished from cosmology n :

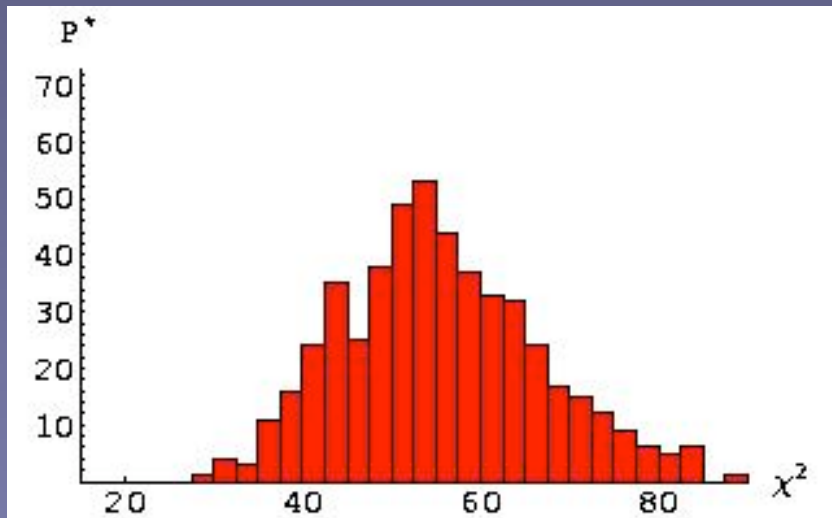
- given by overlap between two distributions $\chi^{2;m,n}$ and $\chi^{2;n,n}$

Chi Square Distributions

3 redshift bins, 15 peak height bins, 0.5 arcmin smoothing



$w = -1$ versus $w = -1$
 $\langle \chi^2 \rangle = 44.89$

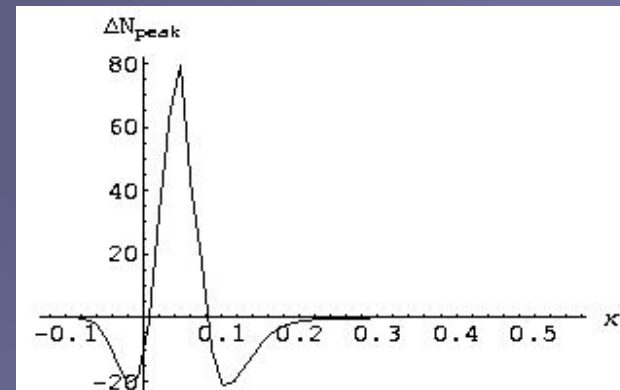
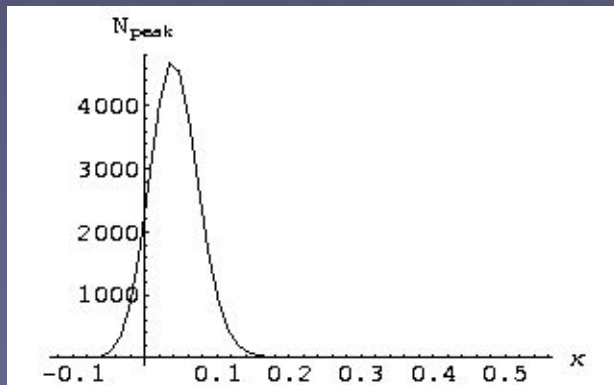


$w = -0.8$ versus $w = -1$
 $\langle \chi^2 \rangle = 55.51$

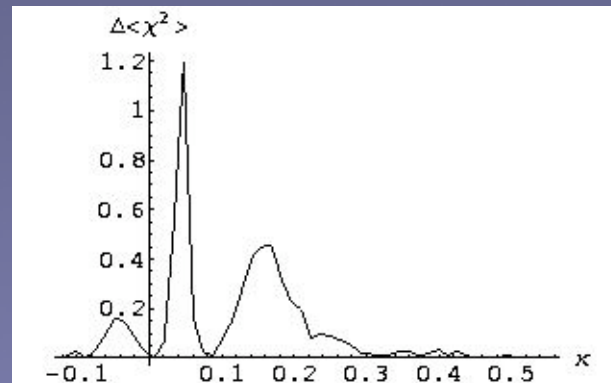
$\rightarrow \Delta \langle \chi^2 \rangle \approx 10$
or 85% confidence

Which peaks dominate constraints?

- smoothing with 0.5 arcmin, galaxies at $z_s=2$
- $w=-1$ more peaks at high+low ends (DE dominates later)
- $w=-0.8$ peaks are more sharply peaked
- medium height ($\kappa \approx 0.04$, or 2σ) peaks dominate the total χ^2



Total # of peaks



Difference in N_{peak}

Contribution to χ^2

Results: statistical significance

- $w=-0.8$ distinguishable at 85% confidence from $w=-1.0$
 - 70% chance for 68% CL
 - 26% chance for 95% CL
- covariance has small effect overall (cuts high- χ^2 tail)
- co-adding several smoothing scales gives only modest ($\sim 10\%$) improvement over single best scale (~ 0.5 arcmin – smaller than best cluster case)
- scaling from $3 \times 3 = 9 \text{ deg}^2$ to $20,000 \text{ deg}^2$:
 - rough guess:
 - significance $\sqrt{(20,000/9)}=50$ times better
 - 1.5σ sensitivity to w_0 is $\Delta w_0 = 0.2/50 = 0.004$

Questions:

- what causes the medium-height peaks?

- (i) one or more individual collapsed halos*
- (ii) mildly over-dense large-scale filaments*
- (iii) unvirialized 'half-collapsed' halos*
- (iv) galaxy shape noise*

- is the information in the medium-height peaks still complementary to the power spectrum?

linear fluctuations in potential, or galaxy shape noise would both produce gaussian random fields

- parameter degeneracies?

*well known in the case of cluster counts;
similar Ω_m - σ_8 degeneracy found for peaks*

(Dietrich & Hartlap 2010)

Suite of N-body Simulations

(Yang, Kratochvil, Wang, Lim, Haiman & May 2011, PRD, submitted)

- 80 new N-body runs; vary w , Ω_m , σ_8
- Higher resolution (4,096 x 4,096) density/potential planes
- 15 gals/amin² at $z_s=1$ and $z_s=2$

	σ_8	w	Ω_m	# of sims
Fiducial	0.798	-1.0	0.26	5
Control	0.798	-1.0	0.26	45
High- σ_8	0.850	-1.0	0.26	5
Low- σ_8	0.750	-1.0	0.26	5
High- w	0.798	-0.8	0.26	5
Low- w	0.798	-1.2	0.26	5
High- Ω_m	0.798	-1.0	0.29	5
Low- Ω_m	0.798	-1.0	0.23	5

TABLE I: *Cosmological parameters varied in each model. The universe is always assumed to be spatially flat ($\Omega_\Lambda + \Omega_m = 1$).*

Origin of Peaks

- **What causes peaks?**

- *identify halos, match them to peaks [use fiducial cosmology]*

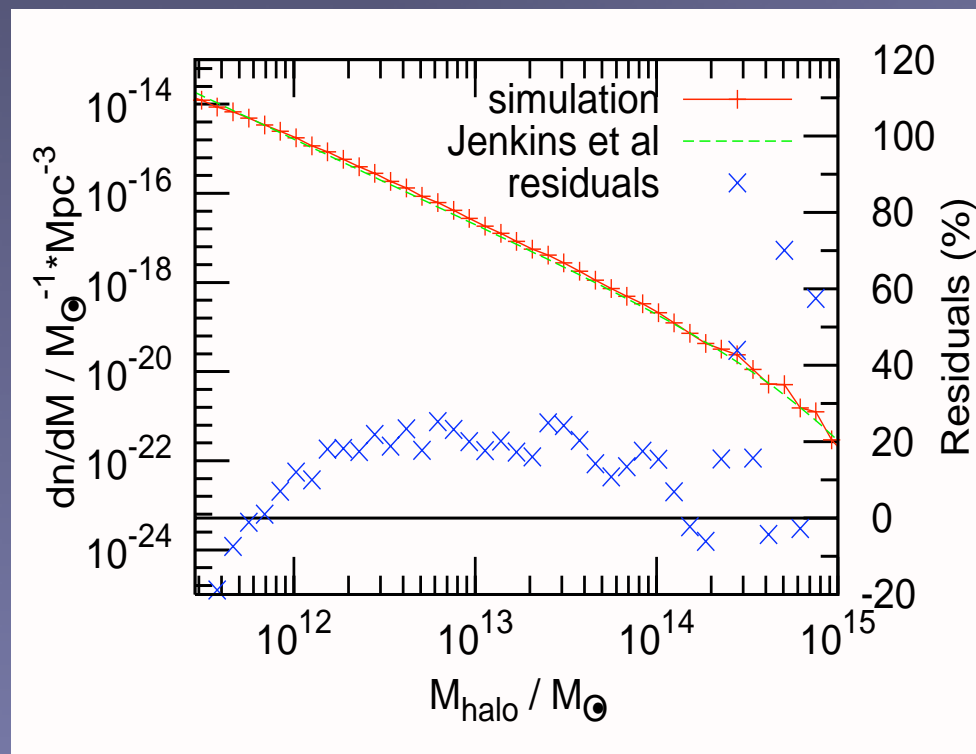
- **What drives the cosmology-dependence of peak counts?**

- *compare two different cosmologies (e.g. vary σ_8) with*

- identical noise realization and (quasi) identical initial condition*

Halo finding

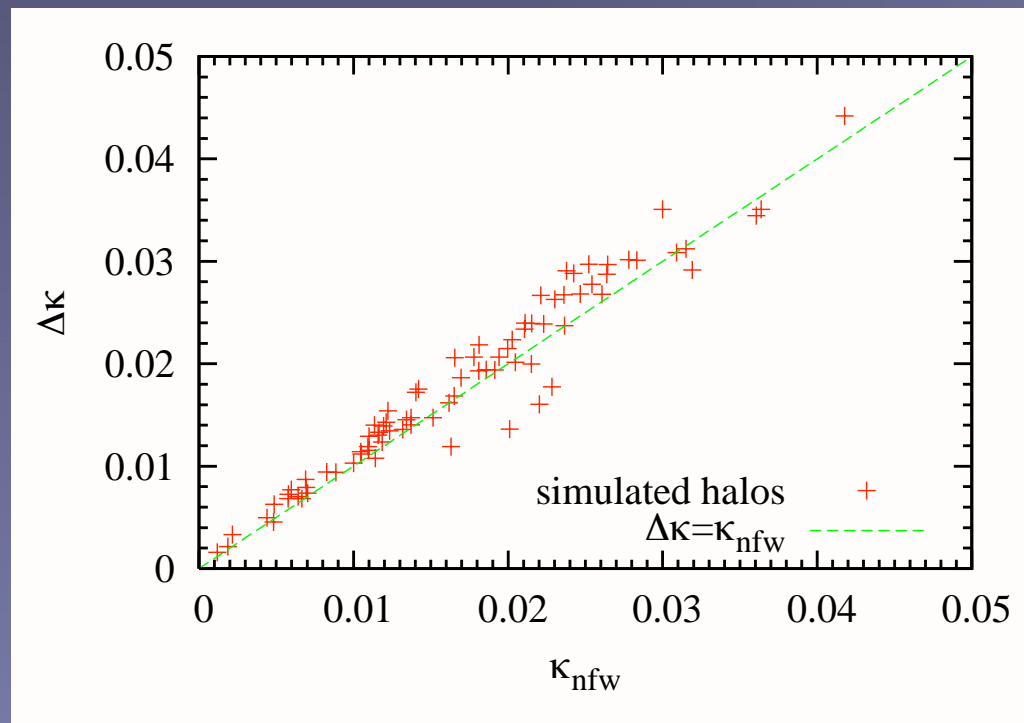
- AMIGA halo finder (Knollmann & Knebe 2009)
- [interesting cropped halos at box edges - make no difference]
- 20% agreement in range $2 \times 10^{11} M_{\odot} < M < 3 \times 10^{14} M_{\odot}$



Halo contributions to convergence

- Assume spherical NFW halo with measured M_{vir}
- Compute $\kappa_{\text{NFW}}(\Phi, M_{\text{vir}})$ as a function of impact parameter Φ
- Compare to $\Delta\kappa$ when ray-tracing is repeated with halo removed

(fractional bias of -0.06; r.m.s. scatter of 0.15)



What causes peaks?

- identify halos, match them to peaks [use fiducial cosmology]
- use 50 noisy maps, 1 arcmin smoothing, 1.8 arcmin cone, $z_s=2$
- Peaks: 2,802 high ($\kappa > 4.8\sigma_{\text{noise}}$) 27,556 medium ($1.1\sigma_{\text{noise}} < \kappa < 1.6\sigma_{\text{noise}}$)
- Halos: 758 22,352

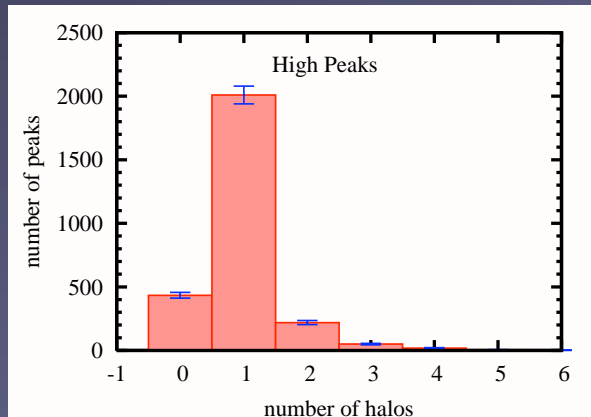
class	matching result	number of matches	
		(high peaks)	(medium peaks)
<i>i</i>	halo \Leftrightarrow peak	526 (0.93)	2653 (4.7)
<i>ii</i>	halo with no paired peak	230 (0.41)	19609 (35)
<i>iii</i>	peak with no paired halo	2264 (4.0)	24709 (44)
<i>iv</i>	halo \Rightarrow peak	2 (0.0035)	90 (0.16)
<i>v</i>	halo \Leftarrow peak	12 (0.021)	194 (0.34)

10% of medium peaks have one-to-one match

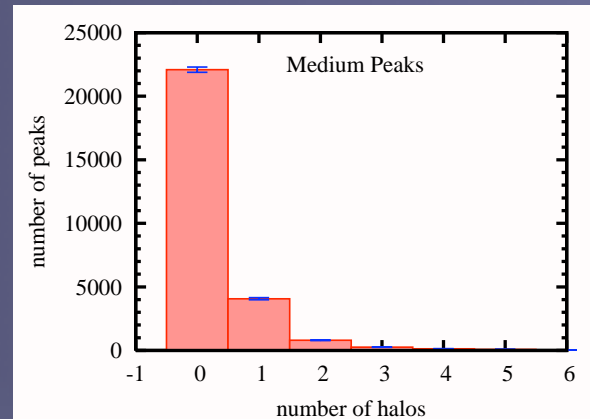
90% of medium peaks/halos have no match

What causes peaks?

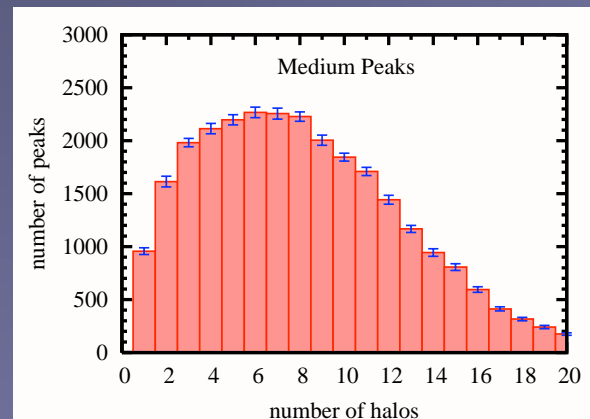
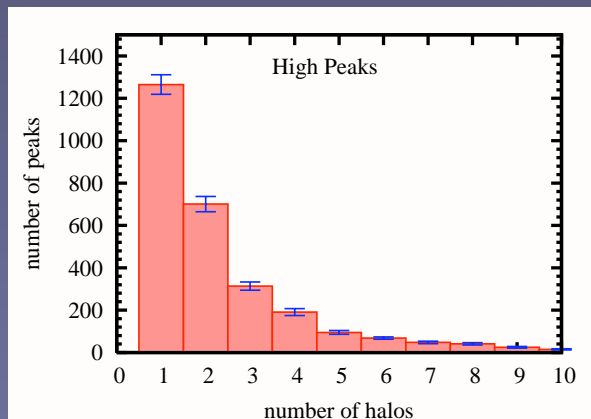
high peaks



medium peaks



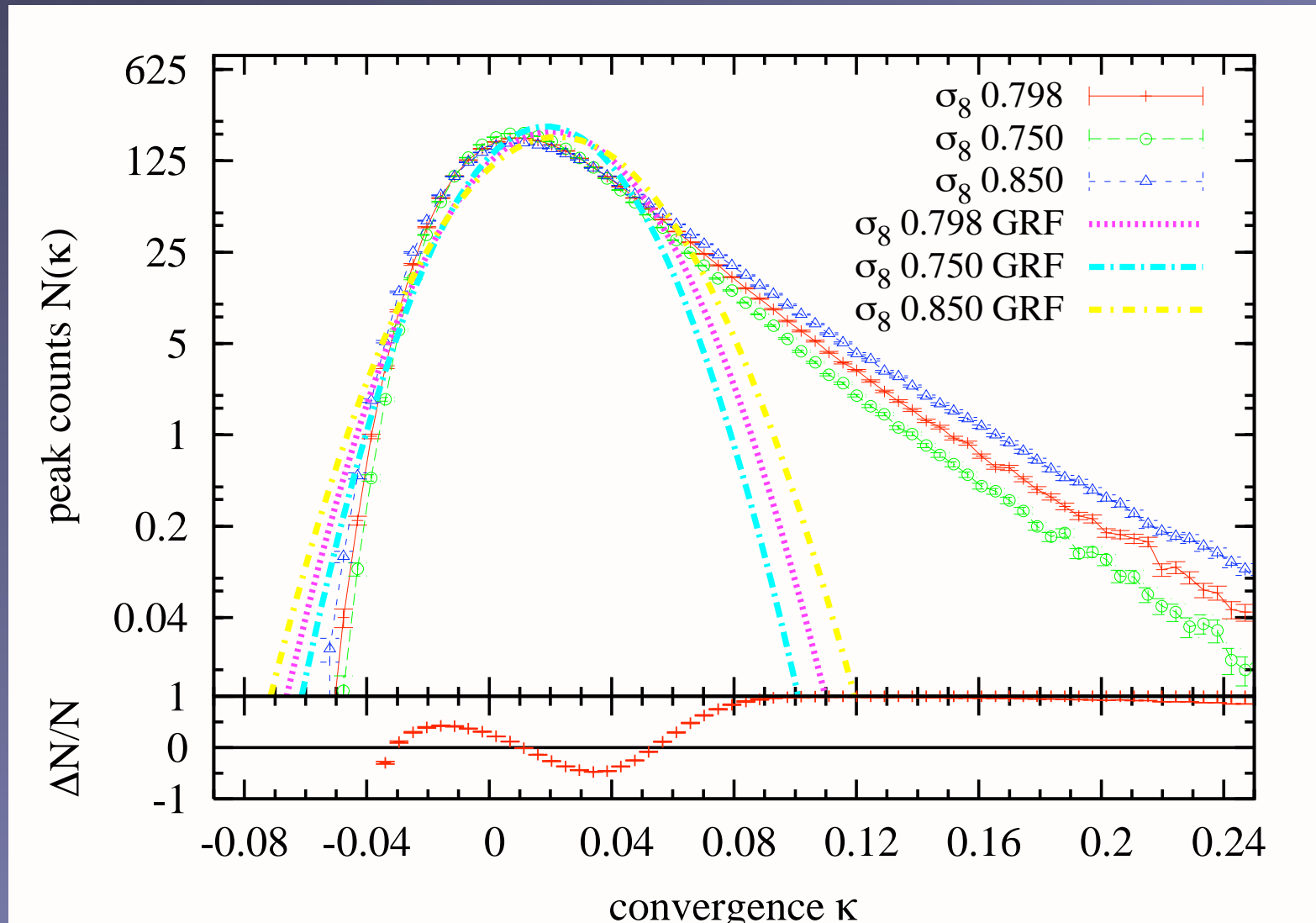
← noise or halo contributions



← halo only contributions

medium peaks are dominated by noise
and the sum of 4-8 halos along the LOS

What drives cosmology-dependence?



Peaks in a Gaussian Random Field:

- Mock maps = analytic prediction (Bond & Efstathiou 1987)
- Can compute covariance matrix numerically in GRF case

$$n_{\max}(\nu)d\nu = \frac{1}{2\pi\theta_*^2} \exp(-\nu^2/2) \frac{d\nu}{(2\pi)^{1/2}} G(\gamma, \gamma\nu) \quad (10)$$

where

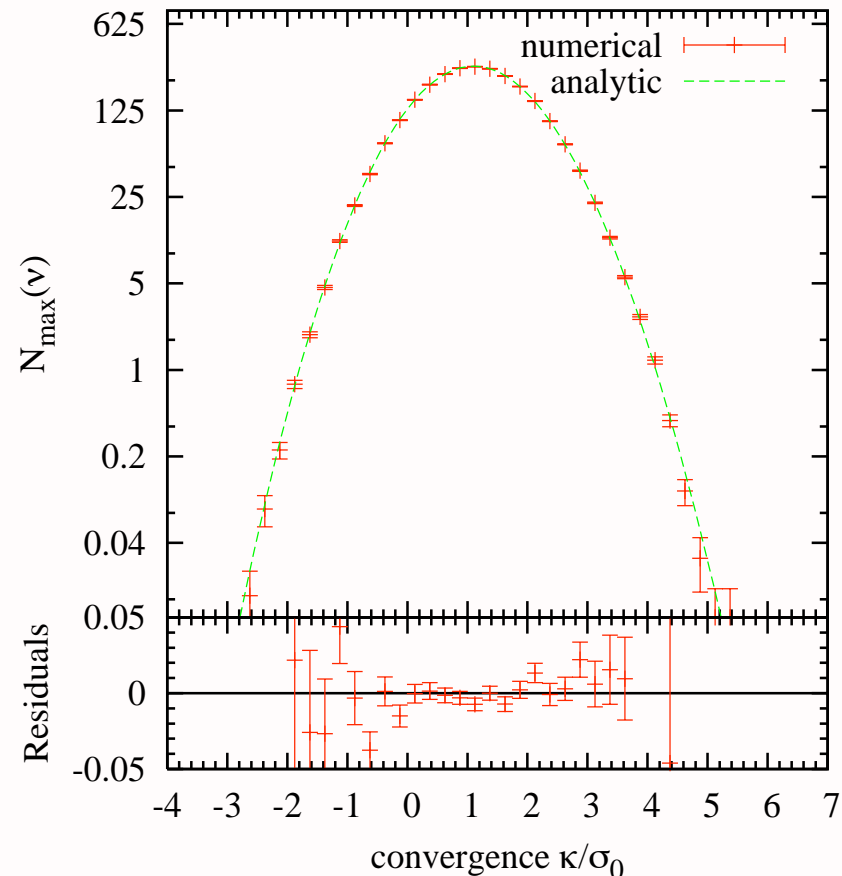
$$\begin{aligned} G(\gamma, x_*) &= (x_*^2 - \gamma^2) \left[1 - \frac{1}{2} \operatorname{erfc} \left\{ \frac{x_*}{[2(1 - \gamma^2)]^{1/2}} \right\} \right] \\ &+ x_* (1 - \gamma^2) \frac{\exp\{-x_*^2/[2(1 - \gamma^2)]\}}{[2\pi(1 - \gamma^2)]^{1/2}} \\ &+ \frac{-x_*^2/(3 - 2\gamma^2)}{(3 - 2\gamma^2)^{1/2}} \\ &\left[1 - \frac{1}{2} \operatorname{erfc} \left\{ \frac{x_*}{[2(1 - \gamma^2)(3 - 2\gamma^2)]^{1/2}} \right\} \right] \end{aligned} \quad (11)$$

$$\gamma = \sigma_1^2 / (\sigma_0 \sigma_2) \quad (12)$$

$$\theta_* = \sqrt{2} \sigma_1 / \sigma_2 \quad (13)$$

$$\begin{aligned} \sigma_p^2 &= \int_0^\infty \frac{\ell d\ell}{2\pi} \ell^{2p} P_\ell \\ &= p! 2^{2p} (-1)^p \frac{d^p \xi}{d(\theta^2)^p} (0), \end{aligned} \quad (14)$$

$$n_{\text{pk}} = (4\pi\sqrt{3})^{-1} \theta_*^{-2}. \quad (15)$$



What drives cosmology-dependence?

- Why does number of medium peaks scale inversely with σ_8 ?

$\sigma_8=0.80$: 2,802 high 27,556 medium

$\sigma_8=0.75$: 1,987 high 29,097 medium

- Match individual peaks in the two cosmologies (via halos)

80-90% of high peaks have 1-1 correspondance (small κ bias)

55% of medium peaks \rightarrow “fragile”: sensitive to halo masses/locations

- Change in κ of individual peaks:

class	low- $\sigma_8 \rightarrow$ fiducial	fiducial \rightarrow low- σ_8
exit to low κ	3303	3420
stay in bin	7683	7515
exit to high κ	5373	4408
total matched	16359	15343
lost (unmatched)	12738	12213

*2/3rd of sensitivity
due to “scatter”
(net \sim 1000 peaks)*

*1/3rd to “destruction”
(net \sim 500 peaks)*

Cosmology Sensitivity

“NEVER show a Table in a talk slide!” - David Spergel,
Princeton, NJ, cca. 2002

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map type	cosmology pair	noiseless $\Delta\chi^2$		noisy $\Delta\chi^2$	
		unscaled	scaled	unscaled	scaled
Sim	F and High- σ_8	5.16	0.46	5.89	4.29
GRF	F and High- σ_8	10.65	0.23	5.87	3.16
Sim	F and Low- σ_8	5.01	0.34	5.09	3.67
GRF	F and Low- σ_8	9.93	0.16	4.98	2.58
Sim	F and High- Ω_m	3.61	0.033	4.02	2.46
GRF	F and High- Ω_m	7.68	0.014	3.77	2.01
Sim	F and Low- Ω_m	4.39	0.053	4.44	2.56
GRF	F and Low- Ω_m	8.79	0.043	4.08	2.11
Sim	F and High- w	0.98	0.47	0.65	0.27
GRF	F and High- w	0.93	0.017	0.46	0.14
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< $\Delta\chi^2$

from peak-counts
between pairs of
models

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Raw sensitivity:

- direct sensitivity
comparable to P_1

- sensitivity to w
is $\sim 10\times$ weaker
(NB no tomography)

Cosmology Sensitivity

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How much of the Information is in σ_{κ} ?:

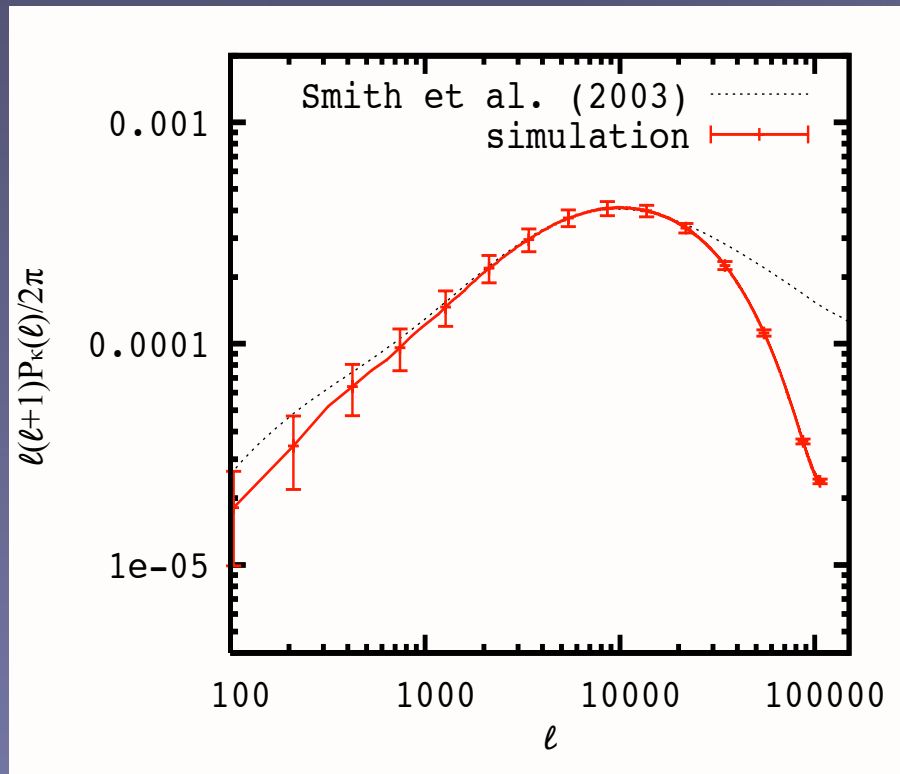
- not most of it!

- note: almost all of the info in the *noiseless* maps is in $\sigma_{\kappa} \rightarrow$ “non-linear interaction” between signal and noise

Convergence power spectrum

- Noise-free, unsmoothed maps, using $z_s=2$
- Theoretical prediction based on NL matter power spectrum

(Nicaea; Kilbinger et al. 2009)



- Agreement in the range $400 < \ell < 30,000$

Is there info beyond power spectrum?

Treat $P_i = l(l+1)P(l)$ in 5 bins $100 < l < 20,000$ as 5 additional observables, compute $C_{ij} = \langle (P_i - \langle P_i \rangle) (P_j - \langle P_j \rangle) \rangle$ and cross-terms $C_{ij} = \langle (P_i - \langle P_i \rangle) (N_j - \langle N_j \rangle) \rangle$

observable type	cosmology	noiseless $\Delta\chi^2$	noisy $\Delta\chi^2$
Peak Counts ($\Delta\chi_N^2$)	Fiducial	5.16	5.89
Power Spectrum ($\Delta\chi_P^2$)	and	17.06	8.12
Combination ($\Delta\chi_{NP}^2$)	High- σ_8	37.07	16.36
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.67	1.17
Peak Counts ($\Delta\chi_N^2$)	Fiducial	5.01	5.09
Power Spectrum ($\Delta\chi_P^2$)	and	13.03	5.76
Combination ($\Delta\chi_{NP}^2$)	Low- σ_8	26.79	11.87
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.49	1.09
Peak Counts ($\Delta\chi_N^2$)	Fiducial	3.61	4.02
Power Spectrum ($\Delta\chi_P^2$)	and	17.69	6.15
Combination ($\Delta\chi_{NP}^2$)	High- Ω_m	32.42	11.65
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.52	1.15
Peak Counts ($\Delta\chi_N^2$)	Fiducial	4.39	4.44
Power Spectrum ($\Delta\chi_P^2$)	and	16.49	5.61
Combination ($\Delta\chi_{NP}^2$)	Low- Ω_m	29.47	10.94
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.41	1.09
Peak Counts ($\Delta\chi_N^2$)	Fiducial	0.98	0.65
Power Spectrum ($\Delta\chi_P^2$)	and	0.92	0.29
Combination ($\Delta\chi_{NP}^2$)	High- w	2.79	0.84
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.46	0.90
Peak Counts ($\Delta\chi_N^2$)	Fiducial	0.44	0.36
Power Spectrum ($\Delta\chi_P^2$)	and	0.64	0.19
Combination ($\Delta\chi_{NP}^2$)	Low- w	1.69	0.51
$\Delta\chi_{NP}^2 / (\Delta\chi_N^2 + \Delta\chi_P^2)$		1.57	0.92

- Raw sensitivities of N and P are comparable
- P(l) degraded more by noise *noise adds no signal δP*
- Synergy from combination: *cosmology-induced changes ($\delta N, \delta P$) do not obey correlations $\langle \Delta N \Delta P \rangle \neq 0$*

Marginalized Errors

- Fisher matrix – use backward/forward finite difference
- Use 15 κ -bins; scale errors by angle $\sqrt{20,000/12} \sim 50$
- **Tomography improves marginalized errors by factor of ~two**
redshift-dependent parameter-sensitivities are non-degenerate
- **Factor of ~2 from adding peak counts to power spectrum**
counts contain information beyond $P(l)$

marginalized error	σ_8	w	Ω_m
z2	0.0065	0.030	0.0057
z1	0.0078	0.036	0.0057
z2+z1	0.0024	0.018	0.0022
Power Spectrum ($z_s = 2$)	0.0047	0.026	0.0028
z2+Power Spectrum	0.0026	0.012	0.0019
z1+Power Spectrum	0.0037	0.020	0.0026
tomography combined	0.0012	0.0096	0.0010
combined/(z2+Power Spectrum)	0.47	0.79	0.52

Minkowski Functionals

- ★ One of several general topological measures of isodensity surfaces (in 3D) or contours (2D)

- an alternative probe of non-Gaussianity

Gott et al. (1982)

- ★ Examples of previous applications

- genus in 3-D galaxy distribution in SDSS

- 2D Minkowski Functionals in CMB: search for primordial non-Gaussianity (f_{NL} limits similar to bispectrum)

- fractional-area $P(>\kappa_{\text{th}})$ in WL – a.k.a. V_0 , one of 3 MFs

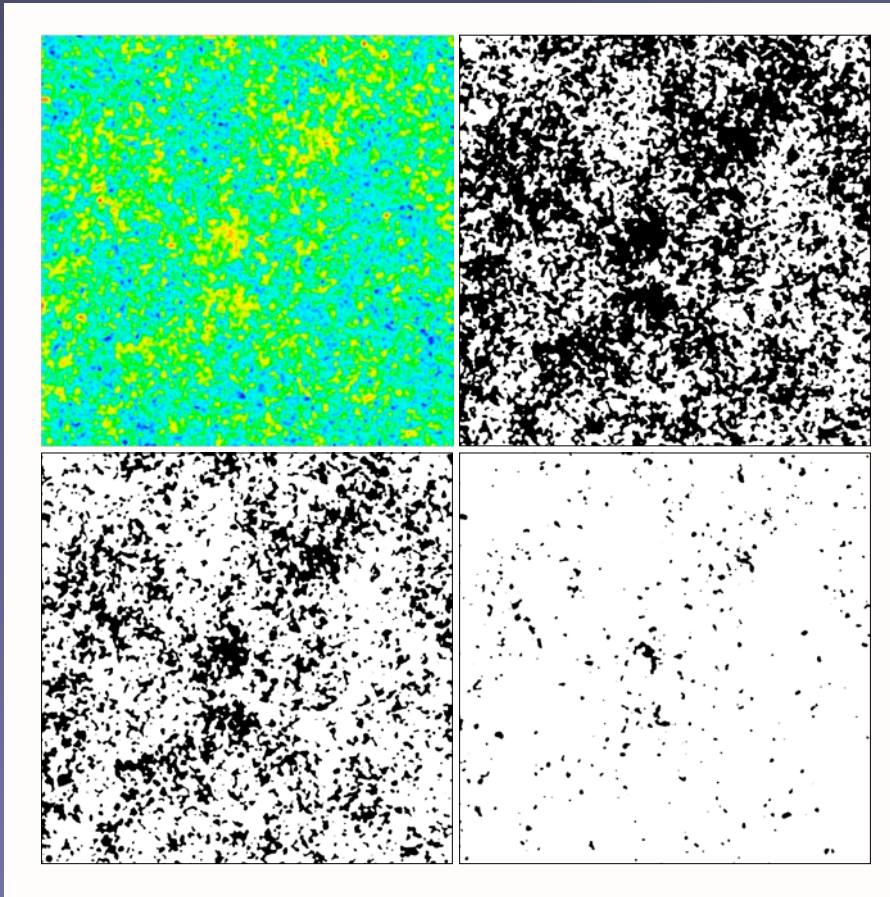
Wang et al. (2009)

- genus in WL maps: V_2 (OCDM vs SCDM; Matsubara & Jain 2001)

- ★ Robustly measurable in data and in our set of simulations

Minkowski Functionals

(Kratochvil, Lim, Wang, Haiman, May & Huffenberger 2011, in prep)



contours as a function of the threshold κ_{th}

- V_0 : *area* above κ_{th}
- V_1 : contour *boundary length*
- V_2 : *Euler characteristic*
#connected regions
above minus #below κ_{th}

Any functional that is additive, translation-invariant, and continuous, is a linear combination of V_1 , V_2 , and V_3 .

Minkowski Functionals

Depend on threshold, and on 1st and 2nd derivatives of $\kappa(\mathbf{x},y)$

$$V_0(\nu) = \int_{\mathbf{R}^2} \Theta(\kappa - \nu) da$$

$$V_1(\nu) = \frac{1}{4} \int_{\mathbf{R}^2} |\nabla\kappa| \delta(\kappa - \nu) da$$

$$V_2(\nu) = \frac{1}{2\pi} \int_{\mathbf{R}^2} |\nabla\kappa| \delta(\kappa - \nu) K da$$

Measure in simulations using finite difference



$$V_0(\nu) = \int_{\mathbf{R}^2} \Theta(\kappa(\mathbf{x}) - \nu) da \quad (8)$$

$$V_1(\nu) = \int_{\mathbf{R}^2} \delta(\kappa(\mathbf{x}) - \nu) \sqrt{\kappa_x^2 + \kappa_y^2} da \quad (9)$$

$$V_2(\nu) = \int_{\mathbf{R}^2} \delta(\kappa(\mathbf{x}) - \nu) \frac{2\kappa_x\kappa_y\kappa_{xy} - \kappa_x^2\kappa_{yy} - \kappa_y^2\kappa_{xx}}{\kappa_x^2 + \kappa_y^2} da \quad (10)$$

Minkowski Functionals in a GRF

Calculable analytically

$$V_0^{\text{GRF}}(\nu) = \frac{1}{2} \left[1 - \text{Erf} \left(\frac{\nu - \mu}{\sqrt{2}\sigma_0} \right) \right]$$

$$V_1^{\text{GRF}}(\nu) = \frac{1}{8\sqrt{2}} \frac{\sigma_1}{\sigma_0} \exp \left(-\frac{(\nu - \mu)^2}{2\sigma_0^2} \right)$$

$$V_2^{\text{GRF}}(\nu) = \frac{\nu - \mu}{4\sqrt{2}} \frac{\sigma_1^2}{\sigma_0^3} \exp \left(-\frac{(\nu - \mu)^2}{2\sigma_0^2} \right)$$

where μ is the mean, and

$$\sigma_0 = \sqrt{\langle \kappa^2 \rangle - \mu^2}$$

is the standard deviation while

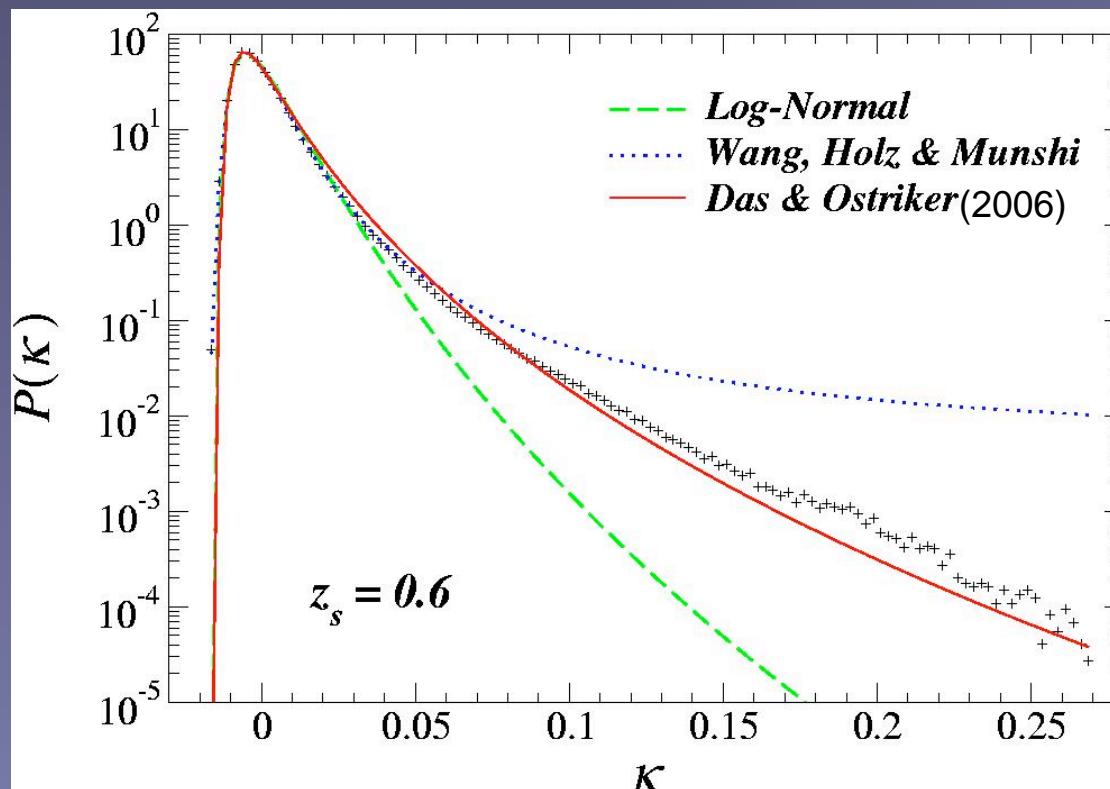
$$\sigma_1 = \sqrt{\langle \kappa_x^2 + \kappa_y^2 \rangle}$$

is its first moment.

Fractional Area of “Hot Spots”

Wang, Haiman & May (2009)

- ★ A simple statistic: one-point function of convergence
i.e. fraction of sky above a fixed threshold $\kappa > \kappa_T = v\sigma_N$
“analytically” calculable, analogous to mass function:

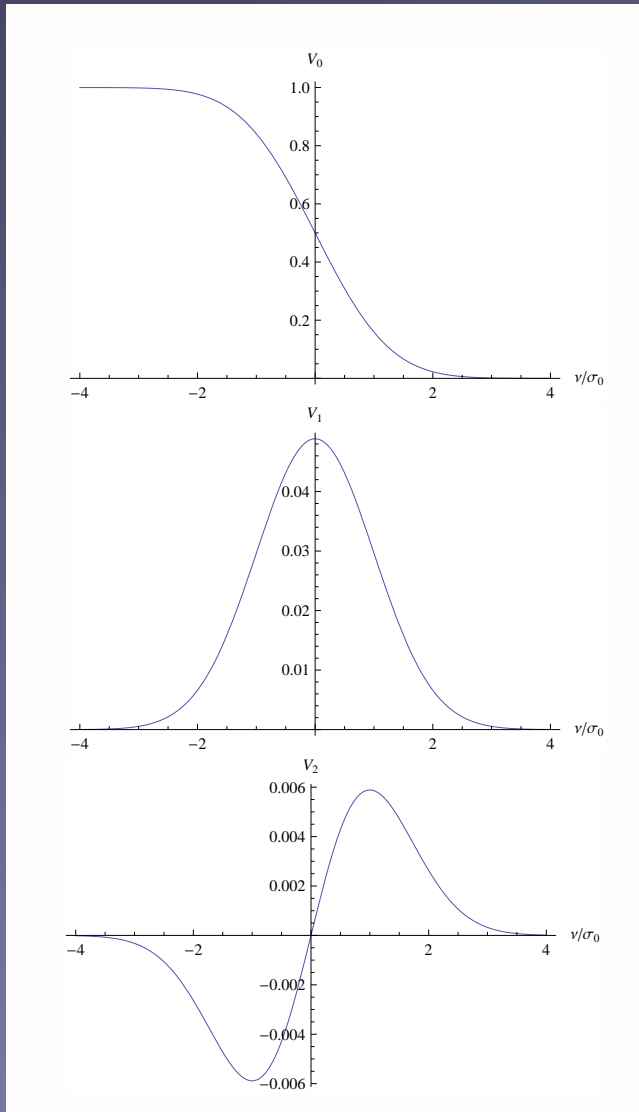


cosmology
dependence
only through
 $\langle \kappa^2 \rangle$ and κ_{\min}

$$F = \int_{v\sigma} P(\kappa) d\kappa$$

Simulations by
M.White (2005)

Minkowski Functionals



- V_0 : *area above κ_{th}*

- V_1 : *contour boundary length*

- V_2 : *Euler characteristic*
genus: #connected regions
above minus #below κ_{th}

Minkowski Functionals

$$V_i(\text{fid})$$

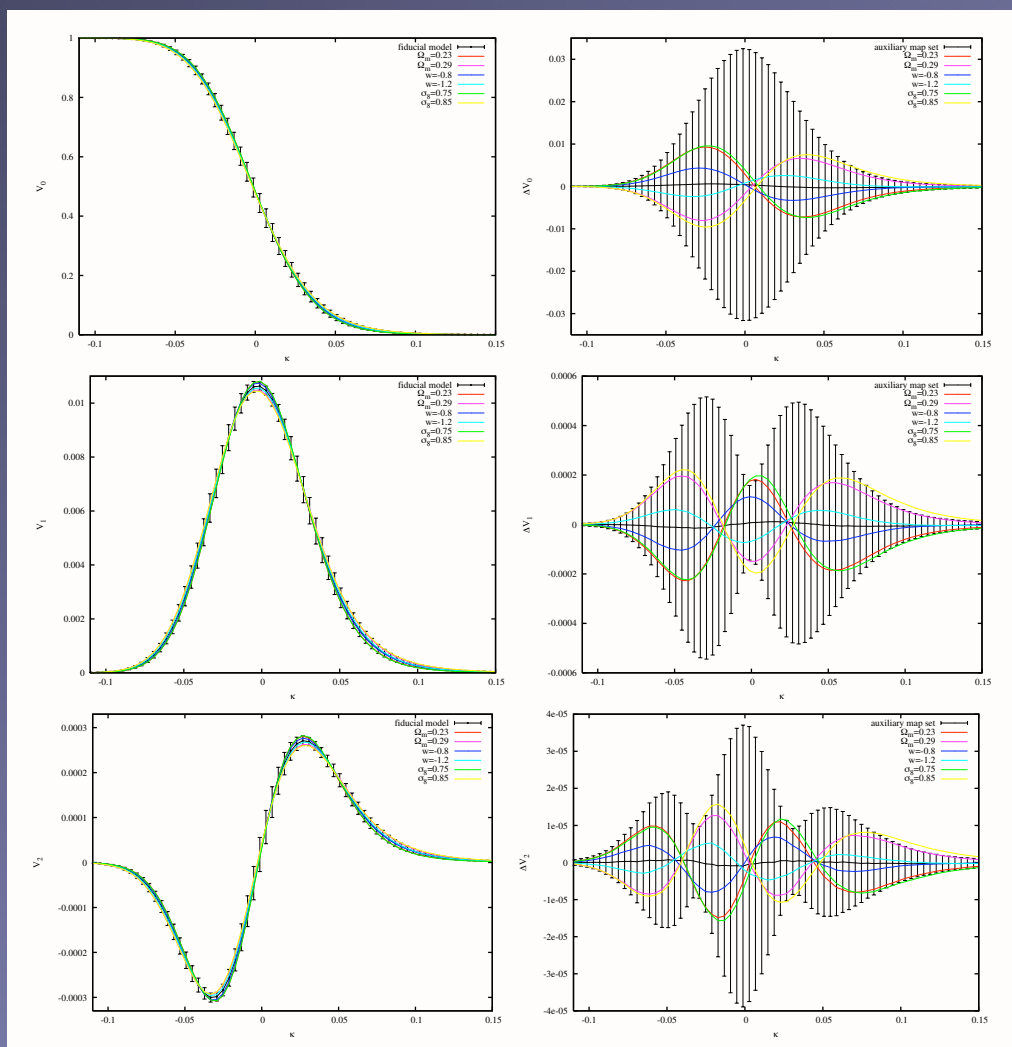
$$\Delta V_i = V_i - V_i(\text{fid})$$

$$z_s = 2$$

$$\theta = 1'$$

$$n = 15/\text{amin}^2$$

1,000 maps



$$V_0$$

$$V_1$$

$$V_2$$

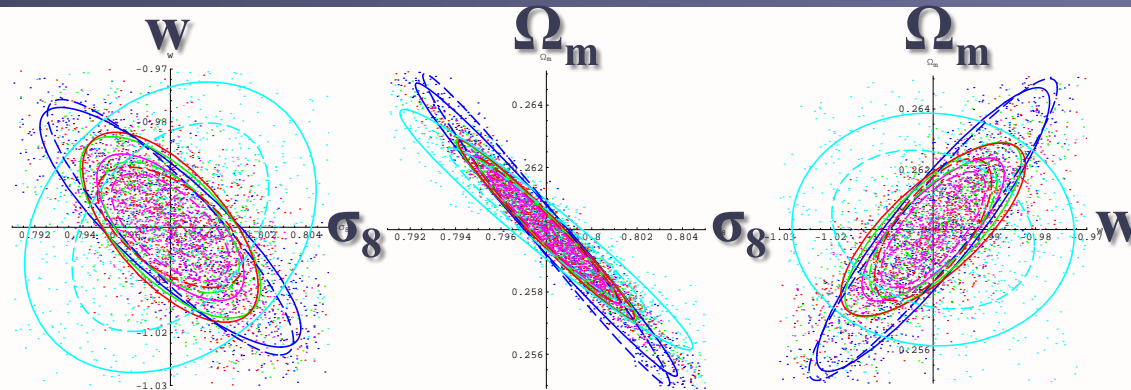
Marginalized Errors

- **MFs deliver factor of 1.5-2 tighter errors than power spectrum**
(V_1 best individual sensitivity; V_1, V_2 have comparable errors)
- **Monte-Carlo errors – good agreement with Fisher matrix**
(slight asymmetry in error bars)
- **Backward vs Forward derivatives**
(small asymmetry)

	$\Delta\Omega_m$	Δw	$\Delta\sigma_8$
$z_s = 2, \theta_G = 1'$			
V_0	-0.00322 (-0.00317) 0.00319 (0.00322)	-0.0151 (-0.0158) 0.0154 (0.016)	-0.00391 (-0.00357) 0.00395 (0.00345)
V_1	-0.0019 (-0.00141) 0.00192 (0.00147)	-0.0113 (-0.00781) 0.0108 (0.00803)	-0.0026 (-0.002) 0.00265 (0.00193)
V_2	-0.00192 (-0.00155) 0.00183 (0.00151)	-0.0123 (-0.00735) 0.0118 (0.00752)	-0.00253 (-0.00206) 0.00257 (0.00202)
PS	-0.00252 (-0.00178) 0.00253 (0.00175)	-0.0183 (-0.0134) 0.0174 (0.0134)	-0.00421 (-0.0029) 0.00427 (0.00294)
MFs	-0.00162 (-0.00129) 0.00162 (0.00126)	-0.00896 (-0.00659) 0.00929 (0.00651)	-0.00225 (-0.00168) 0.00216 (0.00176)

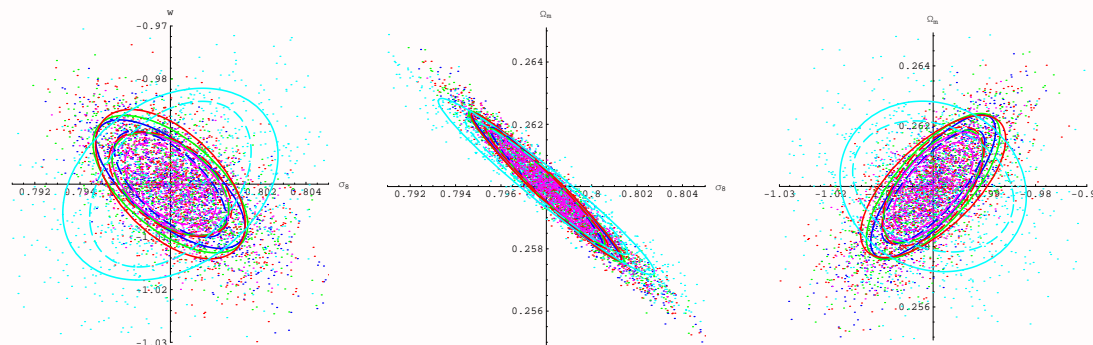
Marginalized Errors (preliminary!)

$z_s=2$
 $\Theta=1'$



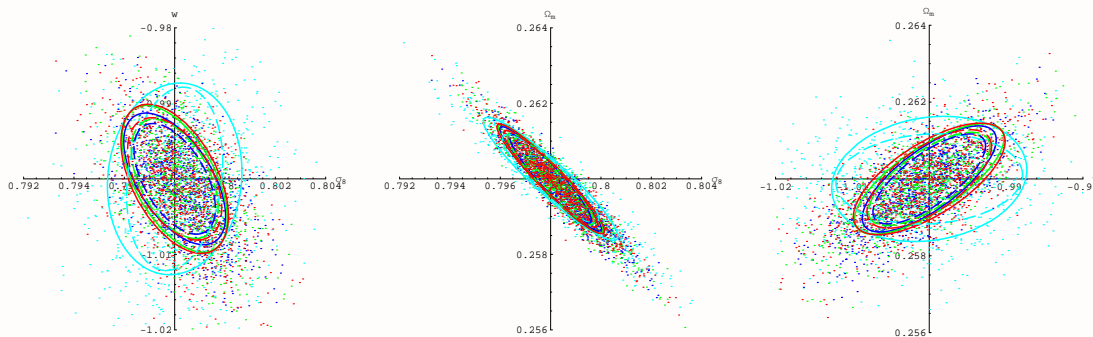
V_0
 V_1
 V_2
 P_1

$z_s=2$
 $\Theta=1,2,3,$
 $5,10'$



note
 V_0
shrinks

$z_s=1,1.5,2$
 $\Theta=1,2,3,$
 $5,10'$



extra
 $\times 1.5-2$

Conclusions

- Peak counts and MFs deliver constraints on Ω_m , w , σ_8
comparable or tighter than the power spectrum
 - This information is new: from non-linear, non-Gaussian regime,
complementary to the power spectrum
 - Peaks: most info is in medium peaks, from halo projections
 - MFs: V_1 most sensitive, V_2 comparable for marginalized
errors; V_0 comparable w/multiple smoothing scales combined
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