Transverse Instability of Megaripples

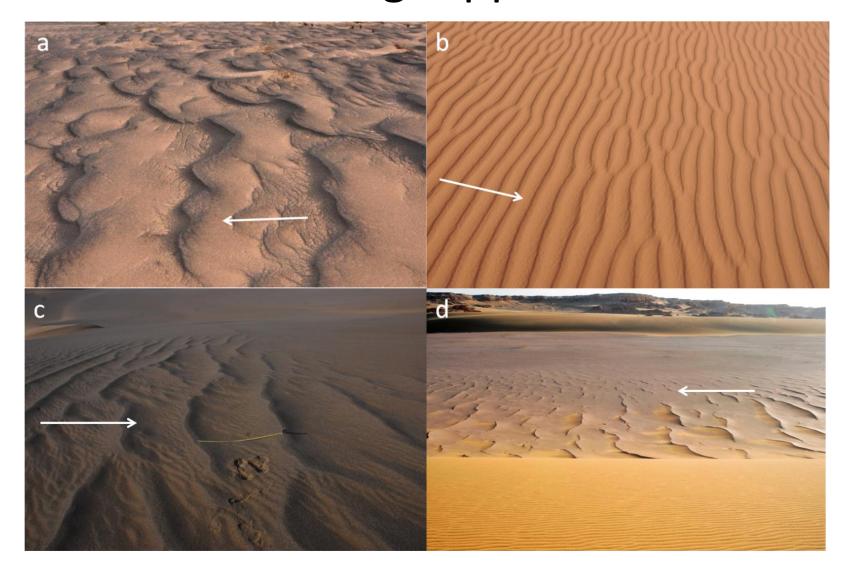
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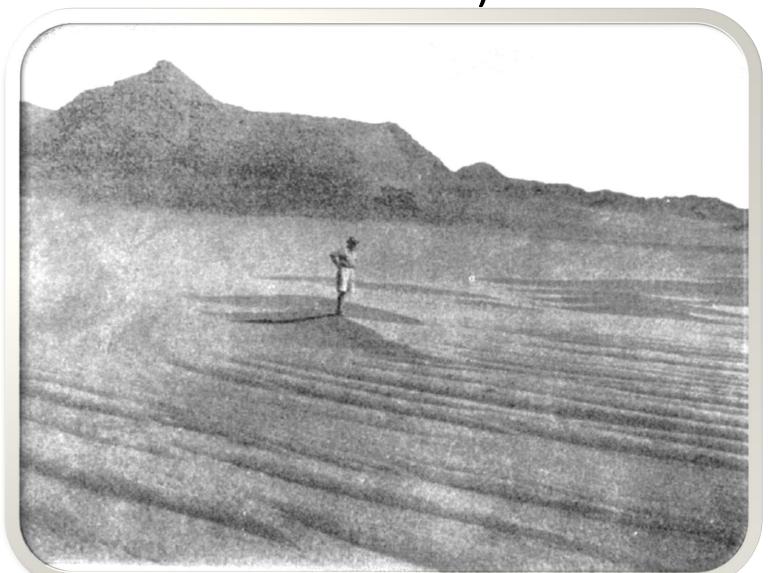
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Kavli Institute-Particle-Laden Flows in Nature

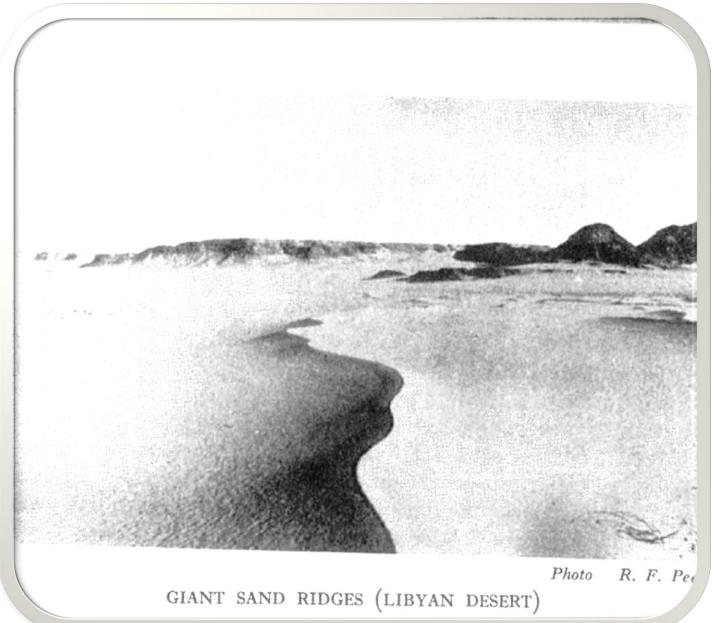
Field Observations of sand ripples and megaripples



Bagnold on a giant sand ridge (Libyan Desert)



Bagnold (1941, Plate 6)



Different types of ripples at Nahal Kasuy







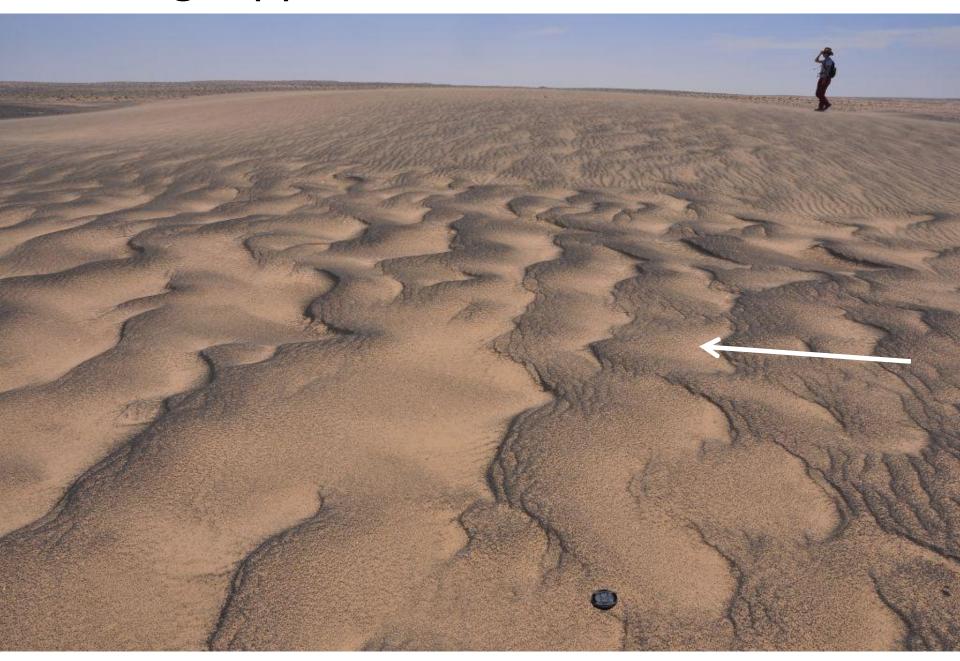




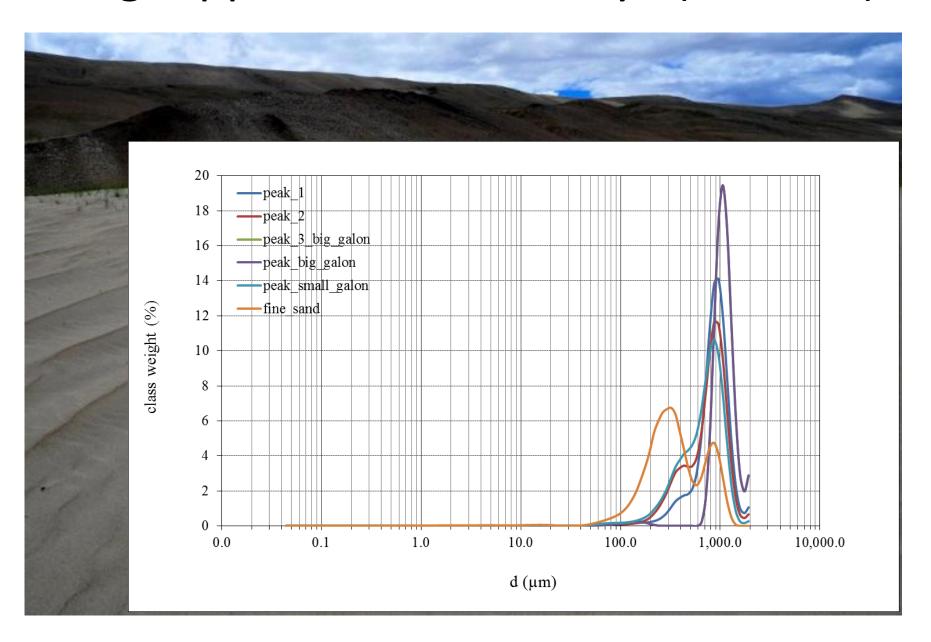
While sand ripples form almost straight lines, megaripples have greater sinuosity due to their transverse instability, a property that causes small megaripple undulations to grow with time.



Megaripples at Grand Falls, Arizona



Megaripples at the Himalaya (4600 m)



The main Questions:

- Why are sand ripples straight whereas megaripples are wavy?
- What is the mechanism behind this transverse instability of megaripples?
- How is the sinuosity of the megaripples depends on the grain size distribution?



Main features of normal ripples and megaripples

	Normal ripples	Megaripples			
Wavelength	Up to	30 cm-20 m ?			
Ripple index*	>15	<15			
Time scale	Minutes	Days and years			
Sorting	Unimodal distribution of	Bimodal distribution of grain			
	grain sizes (typically 0.100-	sizes, with coarse grains			
	0.300 mm in diameter)	0.7-4 mm in diameter			
Basic Process	Saltation and reptation	Saltation and reptation of			
	(creep) of fine grains	fine grains and creep of			
		coarse grains.			
Plan View	Crests form almost straight,	Generally, the crests form			
	continuous lines	wavy and discontinuous			
		lines			
*Ripple index = ratio of wavelength to height					

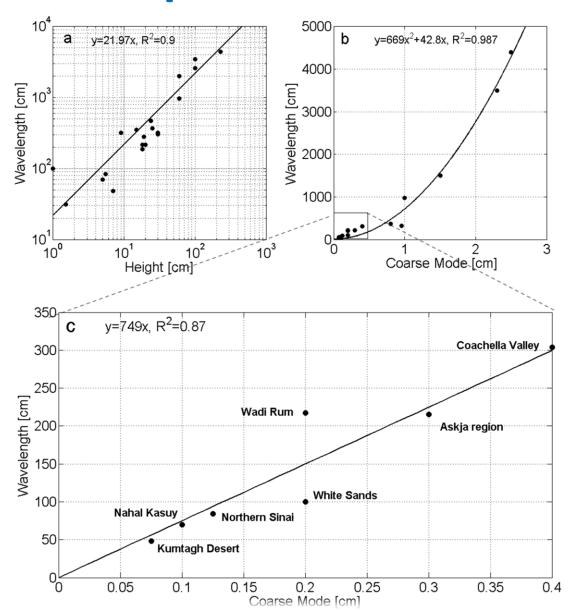
Huge megaripples (or gravel ripples) at the Puna plateau (Argentina)



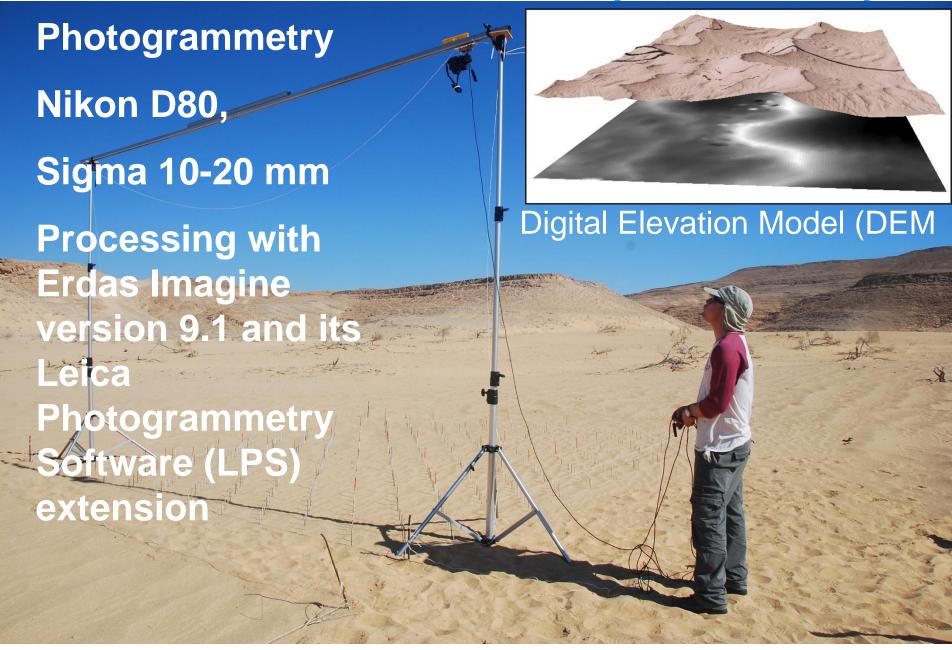
Bagnold (1941, p. 149)

 "the term ripple has been applied to those surface forms whose wavelengths depends on the wind strength, and remains constant as time goes on; and those other forms whose wavelength may increase indefinitely with time are called ridges"

Megaripples morphometry according to previous studies



Evolution of the instability-field study



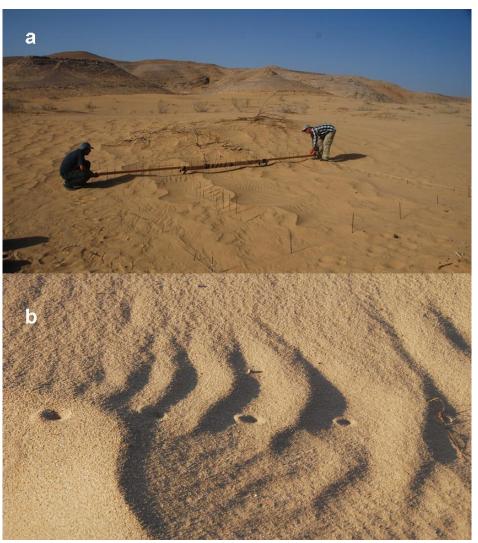
To avoid interfering with plot dynamics, the imaging and the ground control point (GCP) markings had to be made from outside.

Method 1: Iron rods

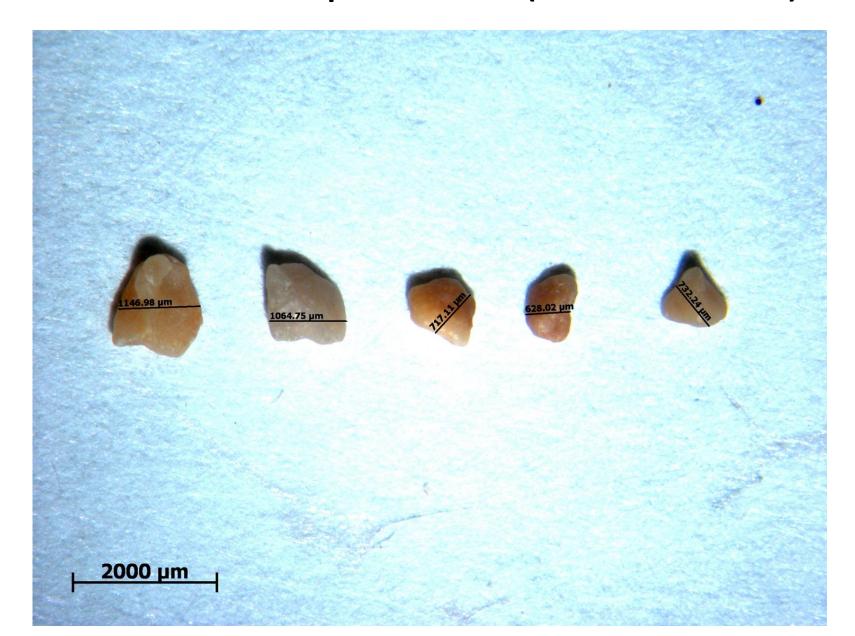


0.6 mm in diameter and 12 cm in height above the surface, 15 cm horizontal spacing

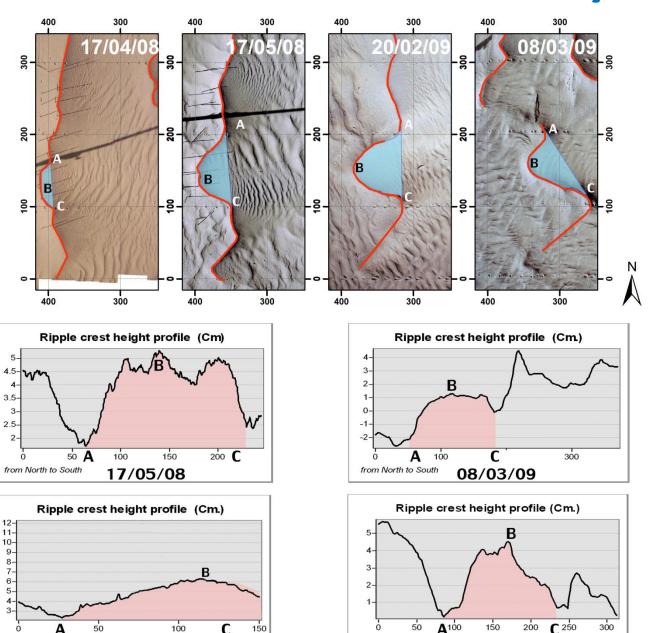
Method 1: Marking holes with "giant comb" (15 cm)



The coarse particles (96% Calcite)



Evolution of the instability-field study



From North to South

17/04/08

From North to South

20/02/09



Size parameters of the megaripple during the period 17/4/08-8/3/09

Date	Bay's	Bay's vertex	Bay's	Bay's
	Width [cm]	Movement	Average	Area
		[cm]	Edges	[cm ²]
			Movement	
			[cm]	
17/4/2008	73			782
_				
18/4-	108	19	44	2798
17/5/2008				
10/5/2000	101	14	24	40FF
18/5/2008	101	14	24	4055
-				
20/2/2009				
21/2-	112	38	55	2831
8/3/2009				

Wind data from Nahal Kasuy

Period	DP [v.u]	RDP [v.u.]	RDD [Deg]	RDP/DP	t[%]	
17/4-	5.34	5.18	279	0.94	7.1	
17/5/2008						
18/5/2008-	15.97	7.28	225	0.46	7.38	
20/2/2008						
21/2-8/3-2009	12.99	10.85	251	0.83	18.2	

Note: DP is the drift potential; RDP is resultant drift potential; RDD is RDP direction; RDP/DP is the wind directionality and t is the time (in percent) that the wind is above the fluid threshold for sand transport (taken as 6 m/s).

$$q_c = (1 - p)\rho c \frac{H}{2}$$

Basic Hypothesis (Yizhaq et al., 2012)

- The origin of the instability is due to variations in megaripple height, which do not diminish over time, and due to the inverse dependence of ripple drift velocity on the height.
- Non-uniform megaripple height may be due to irregularities in the accumulation of coarse particles (both number and sizes)
- The lateral coupling (by lateral reptation flux) in megaripples is smaller than in normal ripples

Growth of the transverse instability

$$\frac{\Delta x}{\Delta t} = c_1 - c_2 = \frac{2q_c(H_2 - H_1)}{(1 - p)\rho H_1 H_2}$$

 c_1, c_2 migration speeds

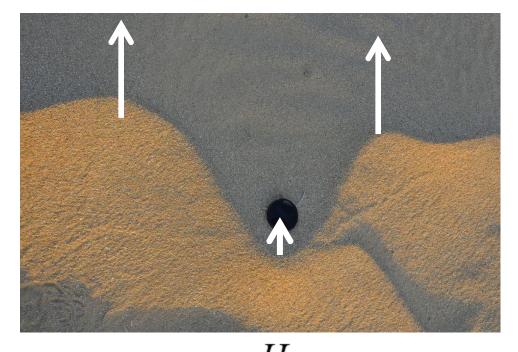
 H_1 height of edges

 H_2 height of vertex

 q_c reptation flux

p porosity

 ρ grain density



$$q_c = (1-p)\rho c \frac{H}{2}$$
 Zimbelman et al., 2009

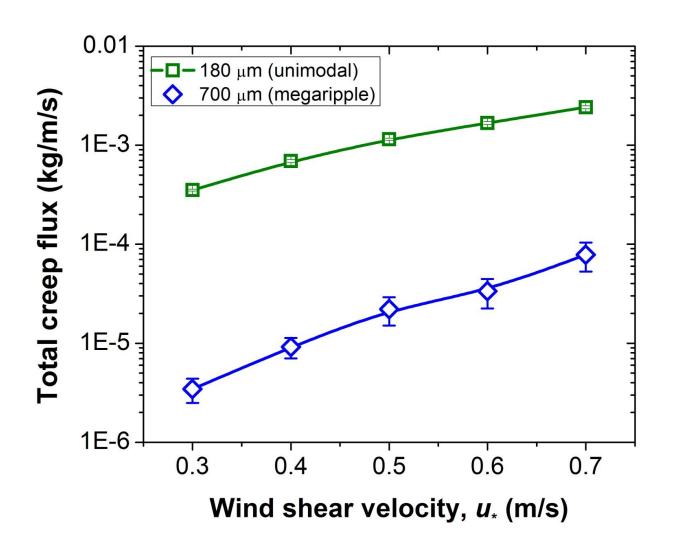
Estimation of Δx for the period 21 February – 8 March

$$q_c = 9.49 \cdot 10^{-6} \text{ kg} \cdot \text{m}^{-2} \cdot \text{s}^{-1}$$
 $p = 0.35$

$$\rho = 2710 \text{ kg/m}^3$$
 $H_1 = 0.04 \text{ m}$
 $H_2 = 0.08 \text{ m}$

 $\Delta t = 18 \text{ days} \Rightarrow \Delta x = 18.62 \text{ cm}$ compare to measured 17 cm

Reptation flux –COMSALT (Kok and Renno, 2009) simulations



A model for sand ripples based on Anderson model (1987)

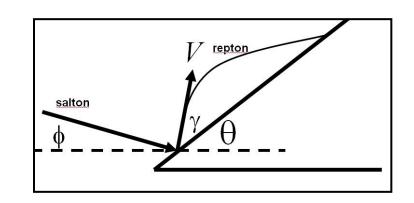
$$(1 - \lambda_p) \rho \frac{\partial \zeta}{\partial t} = -\frac{\partial Q}{\partial x}$$

$$\frac{\partial Q}{\partial x} > 0$$
 Erosion
$$\frac{\partial Q}{\partial x} < 0$$
 Deposition

- λ_p porosity of the bed (about 0.35)
- ρ density of sand
- ζ elevation of the sand surface
- Q flux of sand grains

Number density of impacting grains

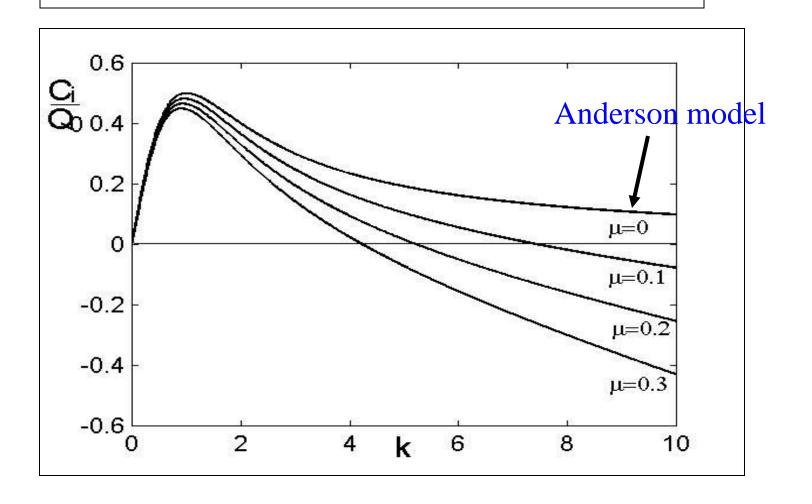
$$N_{im}(x) = N_{im}^{0} \left(1 + \frac{\tan \theta}{\tan \phi} \right) \cos \theta$$
$$= N_{im}^{0} \cot \phi \frac{\tan \phi + \zeta_{x}}{\sqrt{1 + \zeta_{x}^{2}}}$$



 N_{im}^{0} number density of impacting grains on a horizontal surface (about 10^{7} m⁻² s⁻¹) θ inclination of the surface $\tan \theta = \zeta_x$

A flux correction term

•
$$Q_r(x,t) = Q_r^{flat}(x,t) (1 - \mu \zeta_x)$$



The full integral model for sand ripples

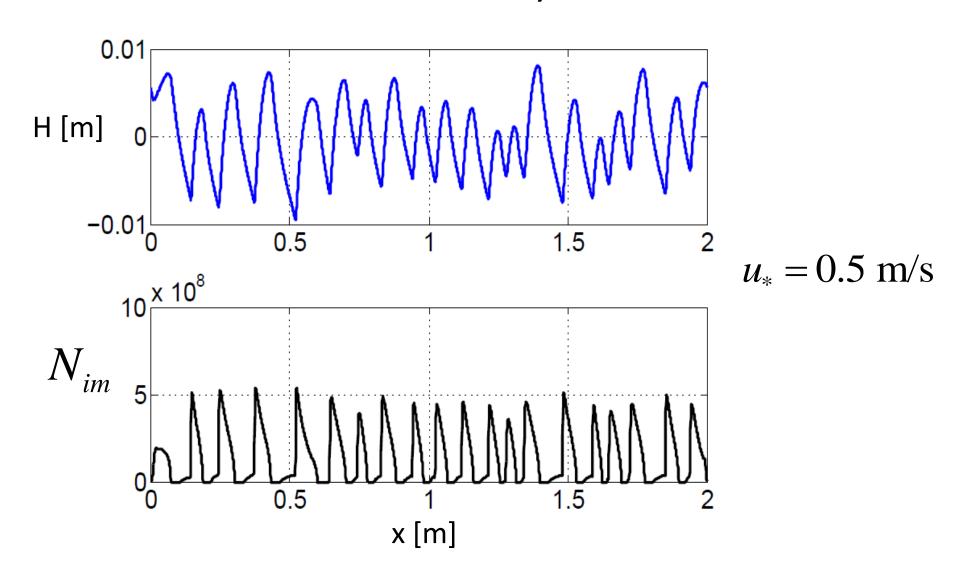
(Yizhaq, H, N.J. Balmforth and A. Provenzale) Physica D, 195, 207-228 (2004), running for martian conditions (Yizhaq et al., 2014)

$$\frac{\partial \zeta}{\partial t} = -Q_0 \frac{\partial}{\partial x} \left[\left(1 - \mu \zeta_x \right) \int_{-\infty}^{\infty} d\alpha \ p(\alpha) \int_{x-\alpha}^{x} F(x') \ dx' \right]$$

$$F(x) = \operatorname{Max} \left\{ \frac{\tan \phi + \zeta_x}{\sqrt{1 + \zeta_x^2}}, 0 \right\}$$

$$Q_0 = \frac{mN_r N_{im}^0 \cot \phi}{\rho (1-n)}$$

2D rippls (parameter values are given by COMSALT)



Long-Wave Approximation

Goal:

Derving a compact or minimal description of the dynamics (from integro-deifferential equation to PDE).

I. Nondimensional Variables:

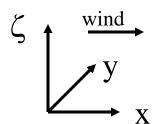
$$\widetilde{x} = \frac{x}{\overline{a}}; \ \widetilde{\alpha} = \frac{\alpha}{\overline{a}}; \ \widetilde{\zeta} = \frac{\zeta}{\overline{a}}; \ \widetilde{t} = \frac{Q_0 t}{\overline{a}}; \ \widetilde{p}(\widetilde{\alpha}) = \overline{a}p(\alpha)$$

II. Near the instability onset: $\zeta(x,t) = \varepsilon \zeta(X,t)$, $X = \varepsilon x$

III. Taylor series expansion of F(X) and $F(X-\epsilon\alpha)$

IV. Assuming
$$\mu \sim \varepsilon$$
 and $T = \varepsilon^2 t$; define: $\overline{a^p} = \int_{-\infty}^{\infty} \alpha^p p(\alpha) d\alpha$

Three Dimensional Ripples mathematical model : long-wave expansion equation



In the lateral direction we assume only rolling.

$$\zeta_{t} = -\partial_{x} Q_{x} - \partial_{y} Q_{y}$$

$$Q_{x} = Q_{r}^{0} (1 - \mu \zeta_{x}) \qquad Q_{y} = -\varepsilon Q_{r}^{0} \zeta_{y}$$

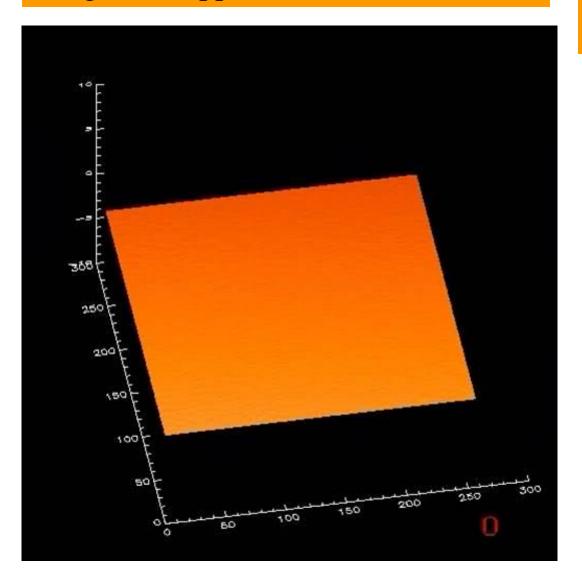
$$\zeta(X,Y,T)_{T} = -(1 - \mu \tan \varphi)\zeta_{XX} + \varepsilon \tan \varphi \zeta_{YY}$$

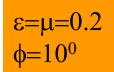
$$+ \frac{\varepsilon}{2} \left[\overline{a^{2}}\zeta_{XXX} + (\tan \varphi + 2\mu)(\zeta_{X}^{2})_{X} \right]$$

$$+ \frac{\varepsilon^{2}}{2} \left[-\frac{\overline{a^{3}}}{3}\zeta_{XXXX} + (\zeta_{X}^{3})_{X} - \frac{\overline{a^{2}}\tan \varphi}{2}(\zeta_{X}^{2})_{XX} \right]$$

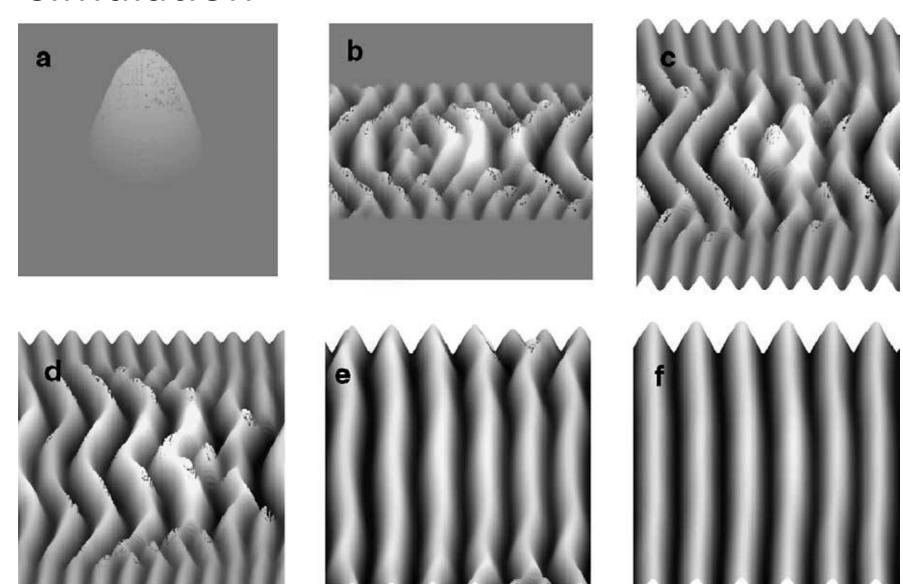
*Yizhaq et al., Physica D, 2004

3D simulation of normal sand ripples (long-wave approximation)





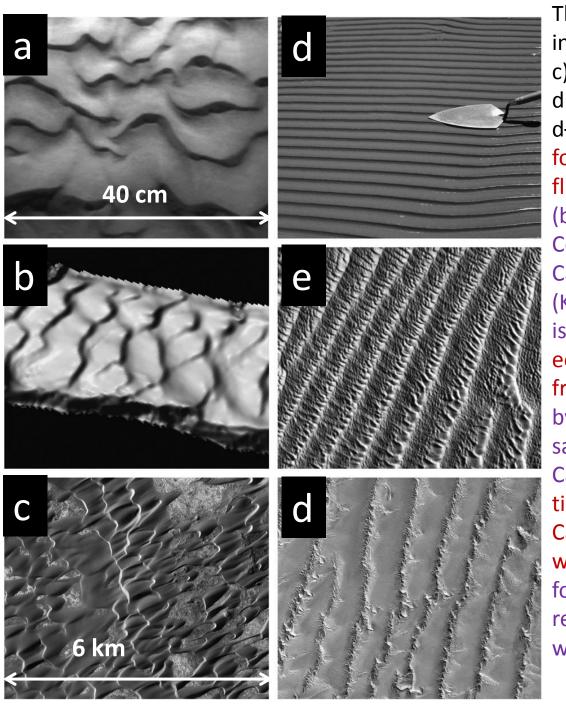
3DNormal ripples numerical simulation



Straightness requires along-crest coupling (Rubin, 2012)

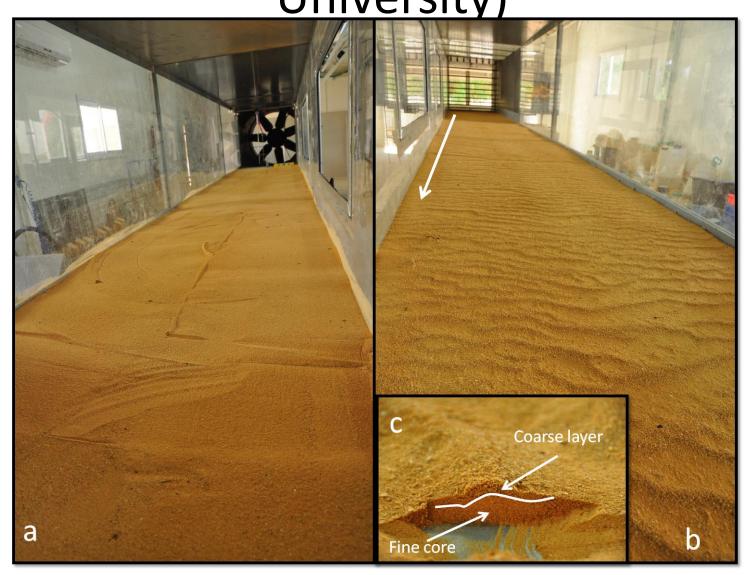
Two-dimensional bedforms and their origins:

- Longitudinal and oblique bedforms in unidirectional flows Flow and transport have an along-crest component.
- Bedforms migrating along a slope (inclined crests) gravity drives along-crest transport.
- Wind ripples along-crest transport caused by ballistic impacts that splash grains laterally.
- Bedforms in reversing flows Flow reversals reduce net across-crest transport, but allow along-crest transport.

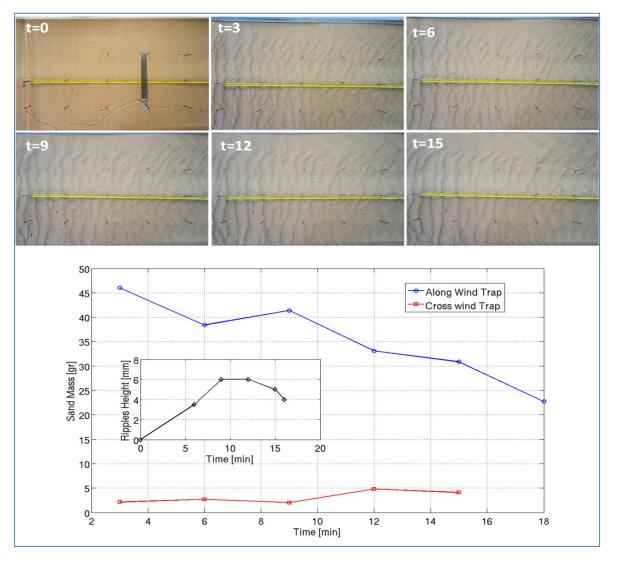


Three-dimensional ripples and dunes in unidirectional flows (left column ac) and two-dimensional ripples and dunes in reversing flows (right column d-f). (a) Ripples formed by unidirectional flow in a lab flume; flow is from top to bottom; (b) Dunes in unidirectional flow in the Colorado River in Grand Canyon viewed by multi beam sonar (Kaplinski et al., 2009); channel width is approximately 70 m. (c) Crescentic eolian dunes on Mars (winds roughly from right to left;); (d) Ripples formed by reversing wave-generated flow on a sand bar in Colorado River in Grand Canyon. (e) Dunes formed by reversing tidal currents, San Francisco Bay, California (Barnard et al., 2011); wavelength is 60 m. (f) Eolian dunes formed by seasonally reversing winds, Namib Desert; dune wavelength is 2 km.

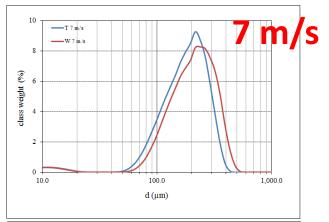
Wind Tunnel Experiments (Ben Gurion University)



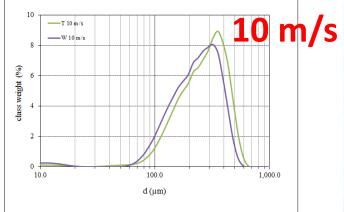
Wind tunnel Experiment with bimodal sand distribution

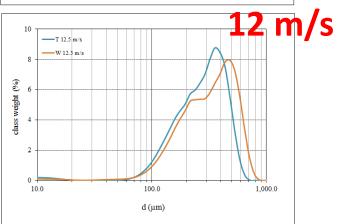


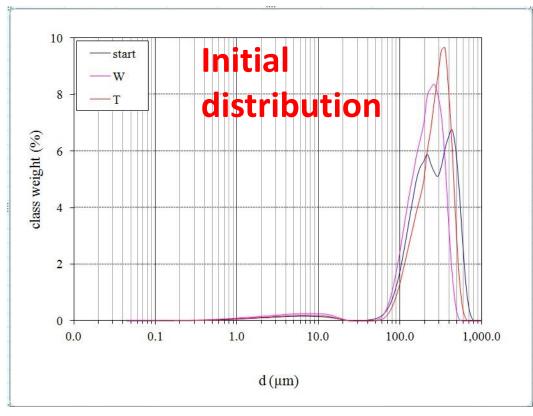
Time evolution of incipient megaripples in the wind tunnel (wind speed 7 m/s) with natural sand collected from Nahal kasuy (time measured in minutes). The bottom graph shows the mass of reptating sand collected in both directions (parallel and transverse to wind direction). The mass in the trap along the wind direction is at least four times larger than that in the cross-wind direction. The inset shows the ripple height (in mm) during the experiment.



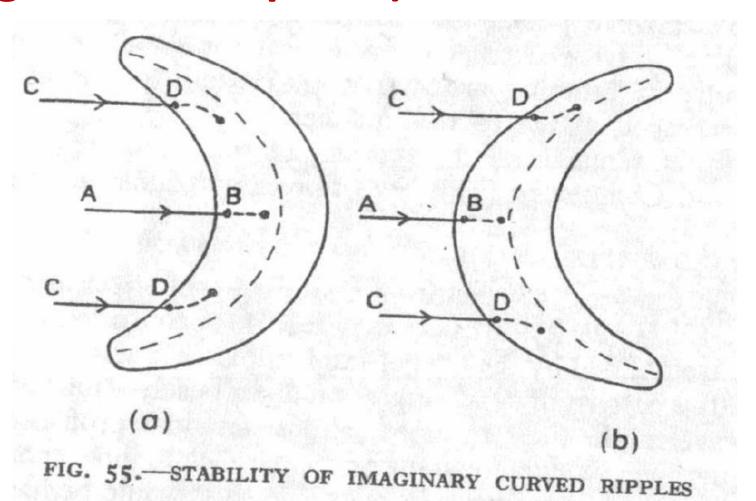
Grain size distribution of reptation for different wind speeds







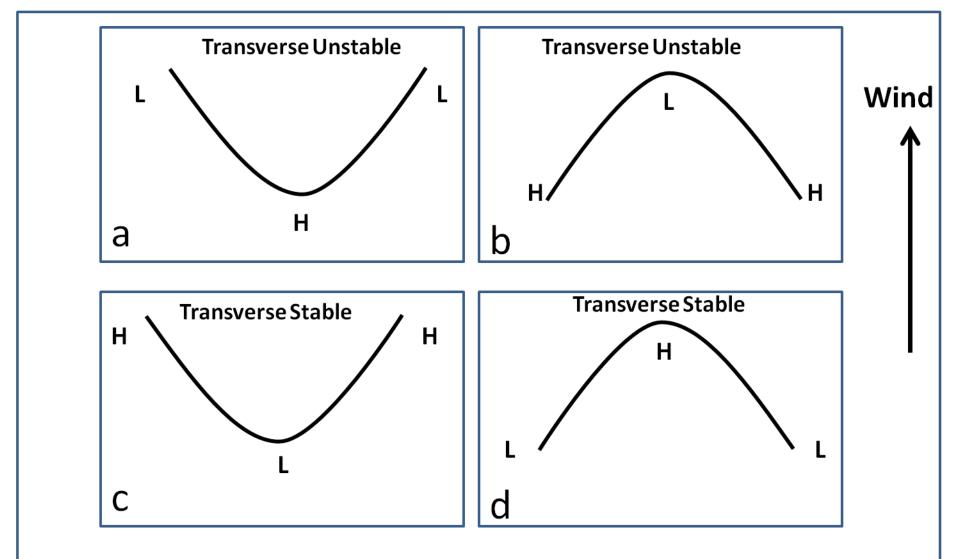
Bagnold's Explanation for the transverse instability of normal ripples (Bagnold 1941, p.161)



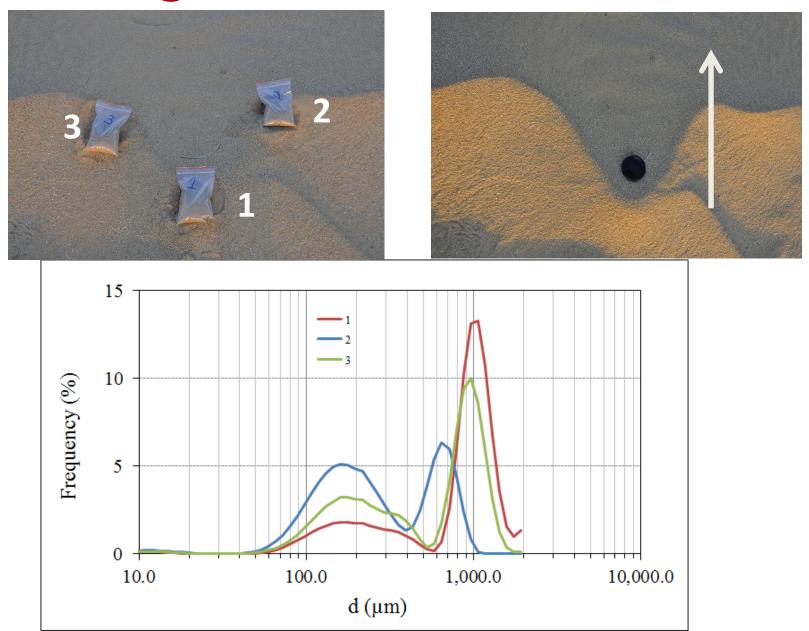
Bagnold's Explanation

- "In ripple (a) the grains roll inwards towards the more transverse portion, and tend to accumulate there. This causes the rate of advance of the transverse portion to become slower than that of the oblique portion."
- The above explanation, though simple, is at present merely tentative".

The four suggested basic dynamics of curved section of the ripple (top view)



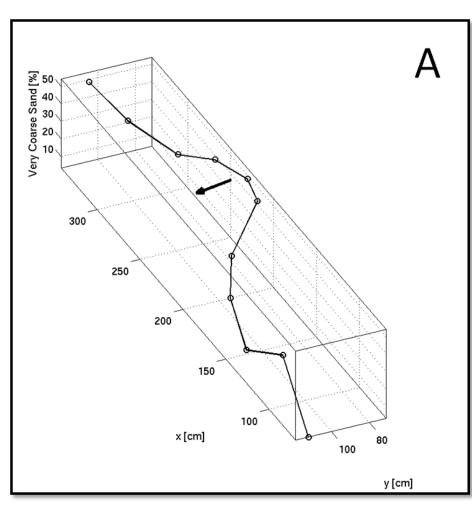
The higher sections are coarser

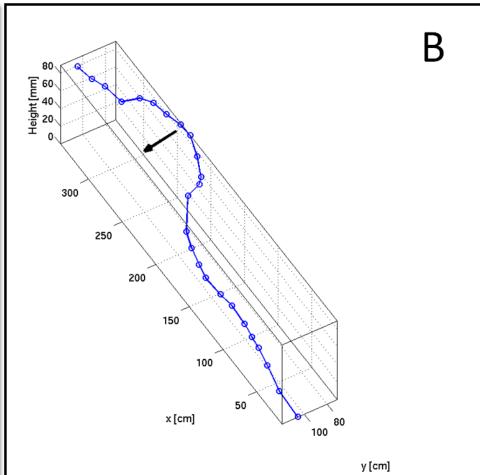


Field measurements



Grain size distribution along the megaripple crest





There are different in the armoring layer along the ripple crest

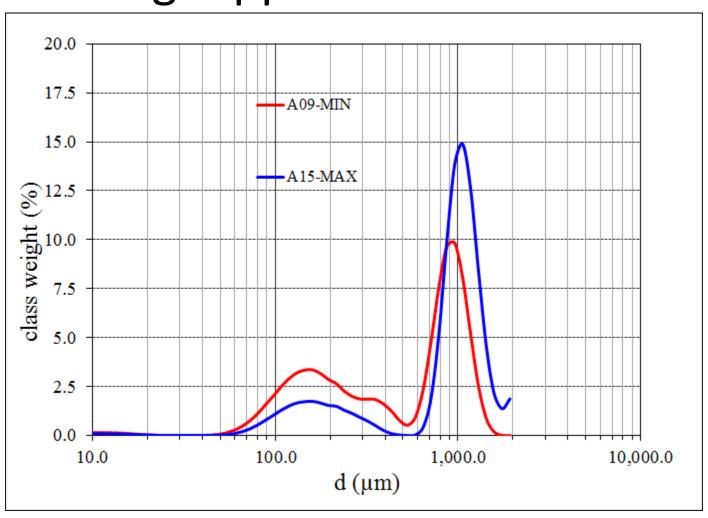




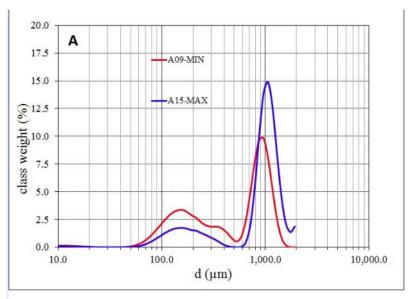
Width=6 mm at the lowest point of the megaripple

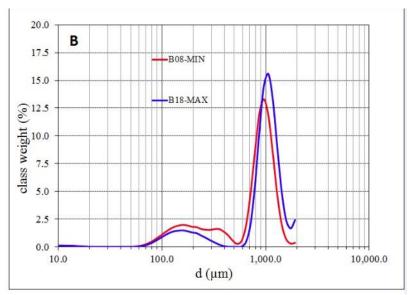
Width=10 mm at the highest point of the megaripple

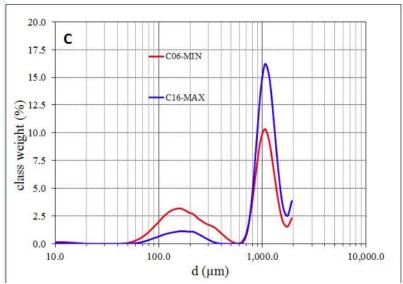
The highest part along the megaripples are coarser

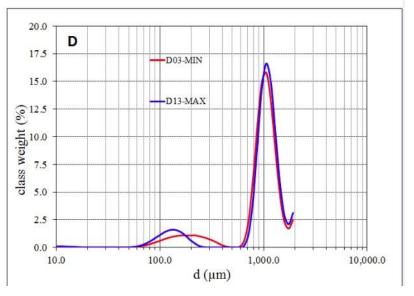


More results from other megaripples









New Indices for bimodal grain size distribution

$$\beta_1 = \frac{D_c - D_f}{D_c}, \quad 0 < \beta_1 < 1$$

Size segregation

$$\beta_2 = \frac{f_c - f_f}{f_c}, -1 < \beta_1 < 1$$

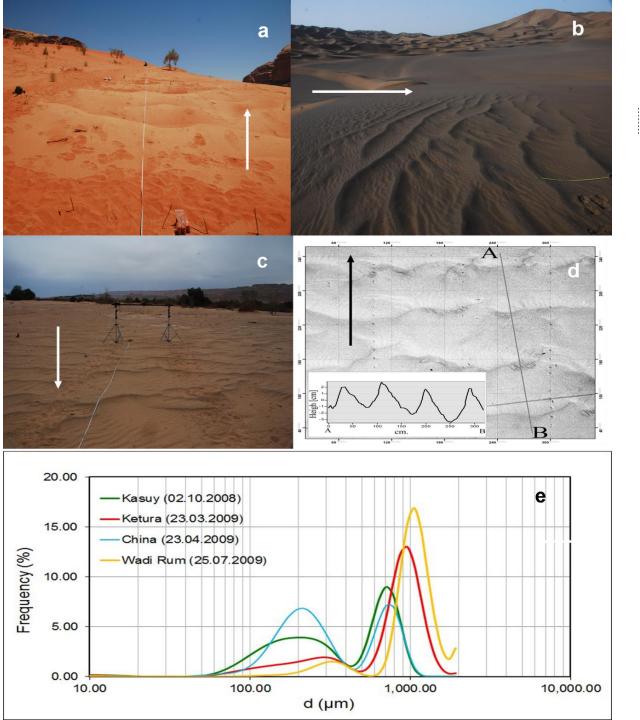
Frequency segregation

c - coarse

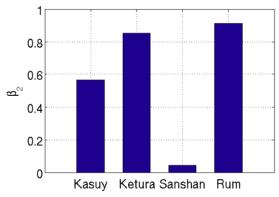
f - fine

D –grain diameter

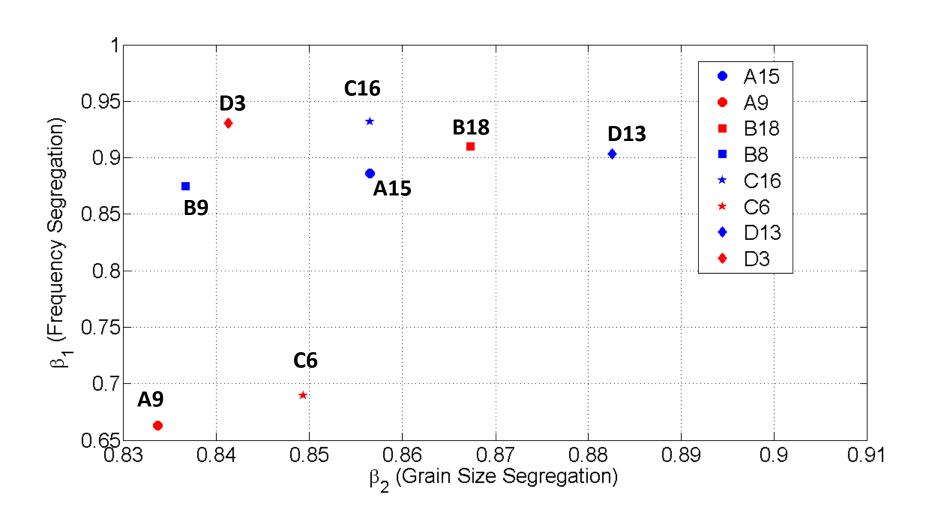
f- frequency



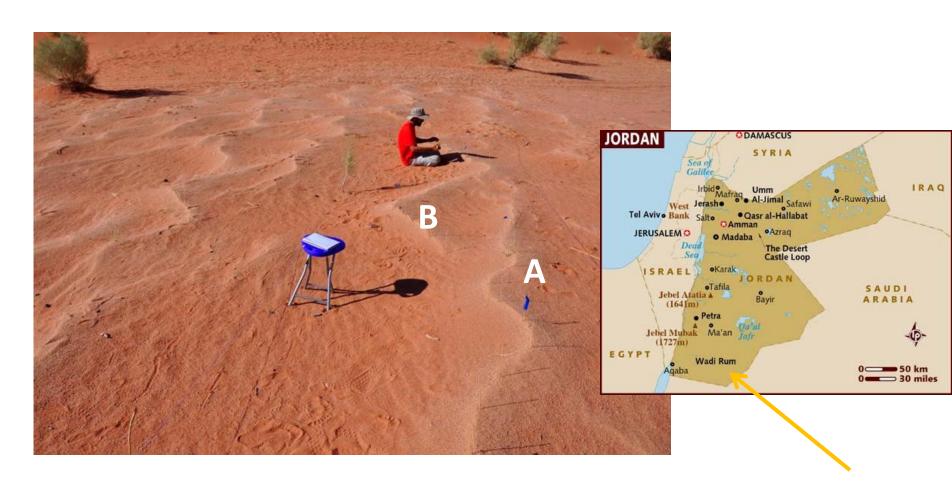




Results for 4 megaripples



Results from Wadi Rum – Jordan



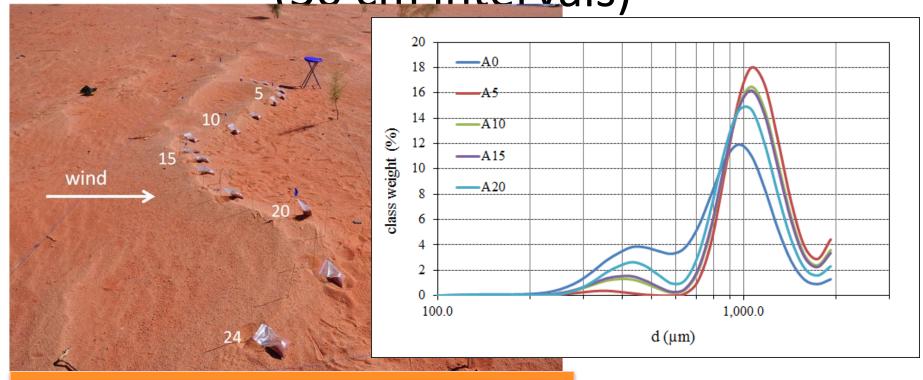
$$SI = \frac{l}{I}$$

Measure of the megaripple sinuosity I actual length along the crest L shortest length

$$SI(A) = 7.2 / 6.3 = 1.19$$

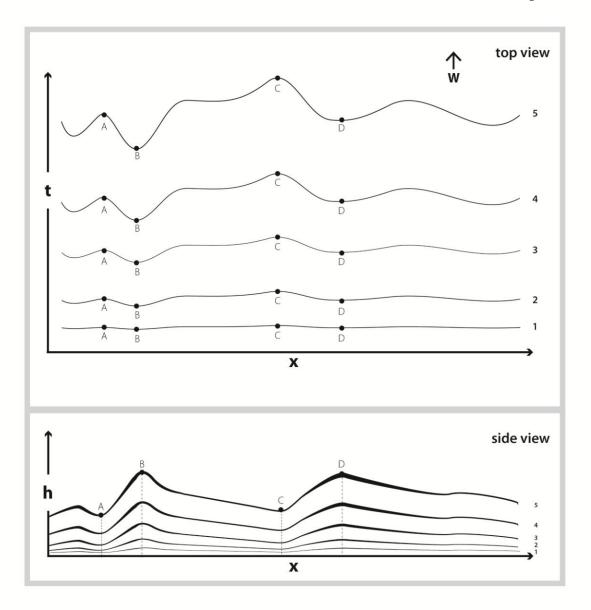
$$SI(B) = 7.2 / 6.6 = 1.09$$

Grain-size distribution along the crest (30 cm intervals)

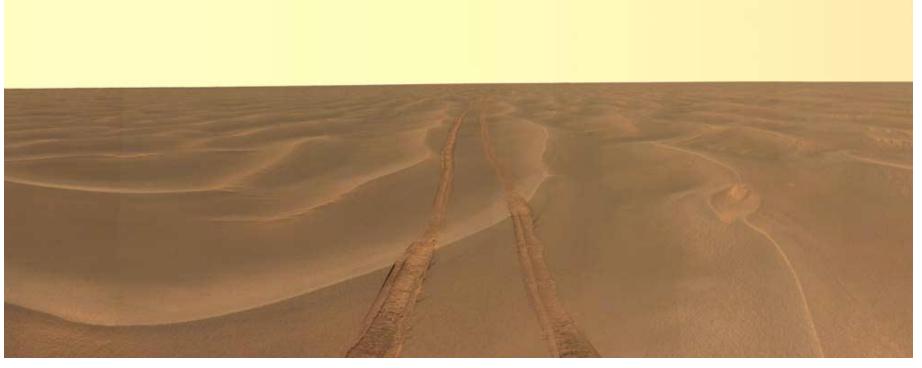


Wadi Ram Jul 2013	A0	A5	A10	A15	A20
Coarse Mode	959	1060	1060	1060	1060
Fine Mode	433	355	392	392	433
Frequency Coarse	11.92	18	16.49	16.09	14.69
Frequency Fine	3.58	0.39	1.31	1.53	2.94
Minimum F	3.31	0	0.24	0.3	0.96
Minimum D	584	584	584	584	584

possible scenario for development of the transverse instability.

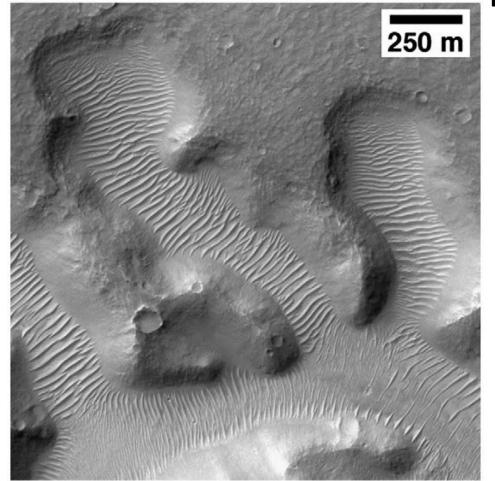


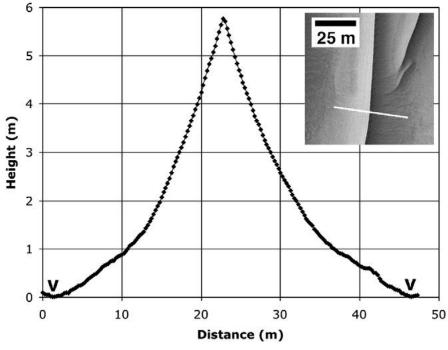
Megaripples on Mars



Approximate true color Pancam panorama of eolian ripples on the Meridiani plains (Squyres et al., 2006)

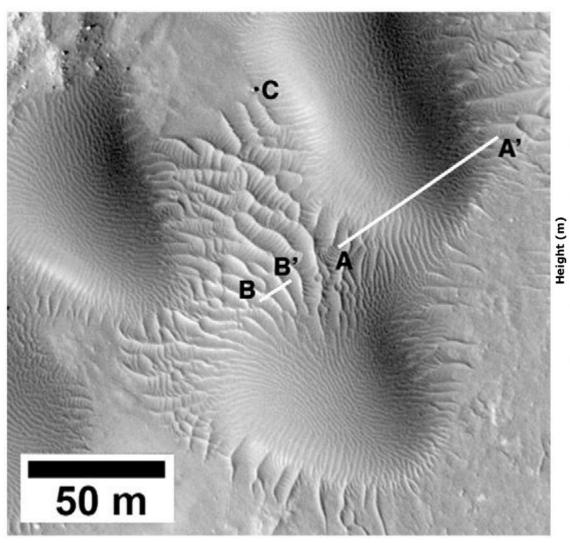
What are the TARs (Transverse Aeolian Ridges)? Small dunes, Large ripples or Megaripples

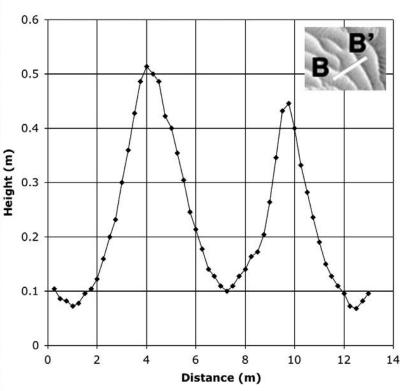




Zimbelman, 2010

Can the sinuosity of the bedforms can help to identify them?





Summary

- Sand ripples are transverse stable whereas megaripples are transverse unstable.
- There is a size segeration along the crest of the megaripples, the highest places are coarser.
- The thickness of the armoring layer along the megaripples crest is non uniform; it is larger at the higher sections of the ripple.
- The lateral coupling in megaripples is smaller than in normal ripples- the reptation flux is smaller in the lateral direction than in the wind direction.

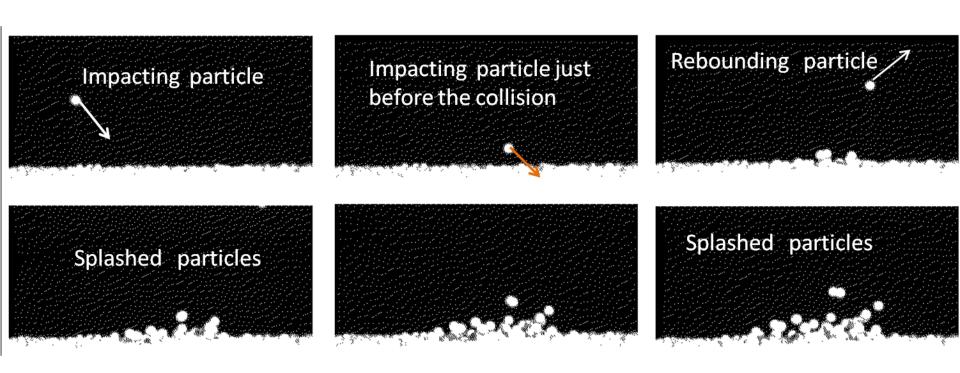
Future Studies *



Studying the spatiotemporal dynamics of a field of megaripples and the relation between the different crest.

- What is the mechanism that prevents segmentation of the crests?
- Mathematical model for megaripples (3D?)

3D study of the splash distribution for different grain size distribution



Beladjine et al. 2007

Relevant Papers

Anderson, R.S., 1987. A theoretical model for impact ripples. Sedimentology, 34, 943-956.

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Gordon, M. and Neuman, C. M., 2011. A study of particle splash on developing ripple forms for two bed materials. *Geomorphology*, 129, 79–91.

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Yizhaq, H., Balmforth, N.J and Provenzale, A. 2004. Blown by wind: Nonlinear dynamics of aeolian sand ripples. *Physica D*, 195, 207-228.

Yizhaq, H., Katra, I., Kok, J. F. and Isenberg, O., 2012. Transverse instability of megaripples. Zimbelman, J.R. 2010. Transverse Aeolian Ridges on Mars: First results from HiRISE images, *Geomorphology*, 121, 22-29, doi: 10.1016/j.geomorph.2009.05.012, 2010.

Geology, 40, 459-462.