

# Granular flows

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debris flows,  
pyroclastic flows,  
rock avalanches...

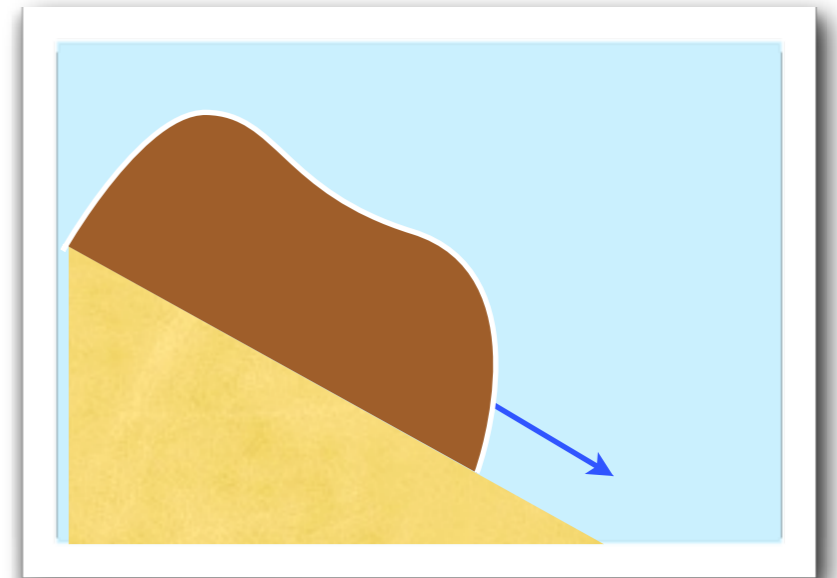


## Questions:

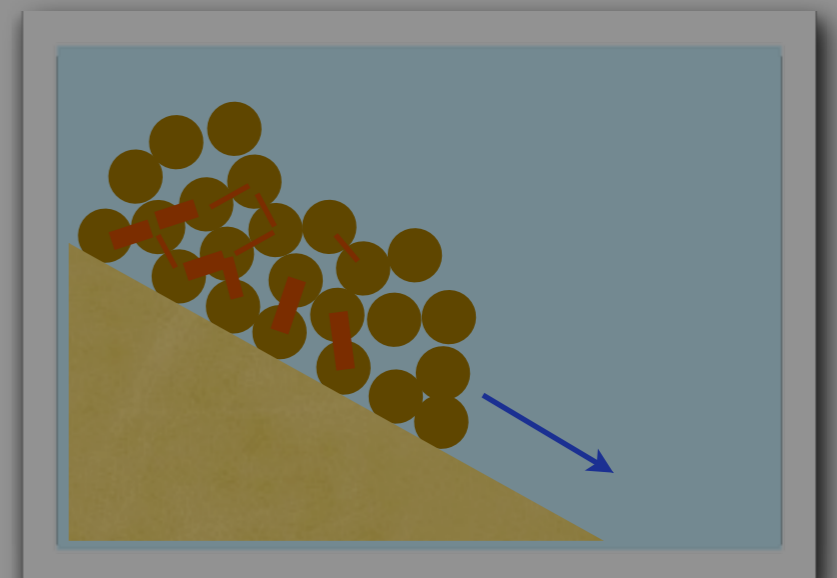
continuum description  
of granular media

==> phenomenology

In this talk:  
a fluid mechanics point of view



understanding the microscopic dynamics.

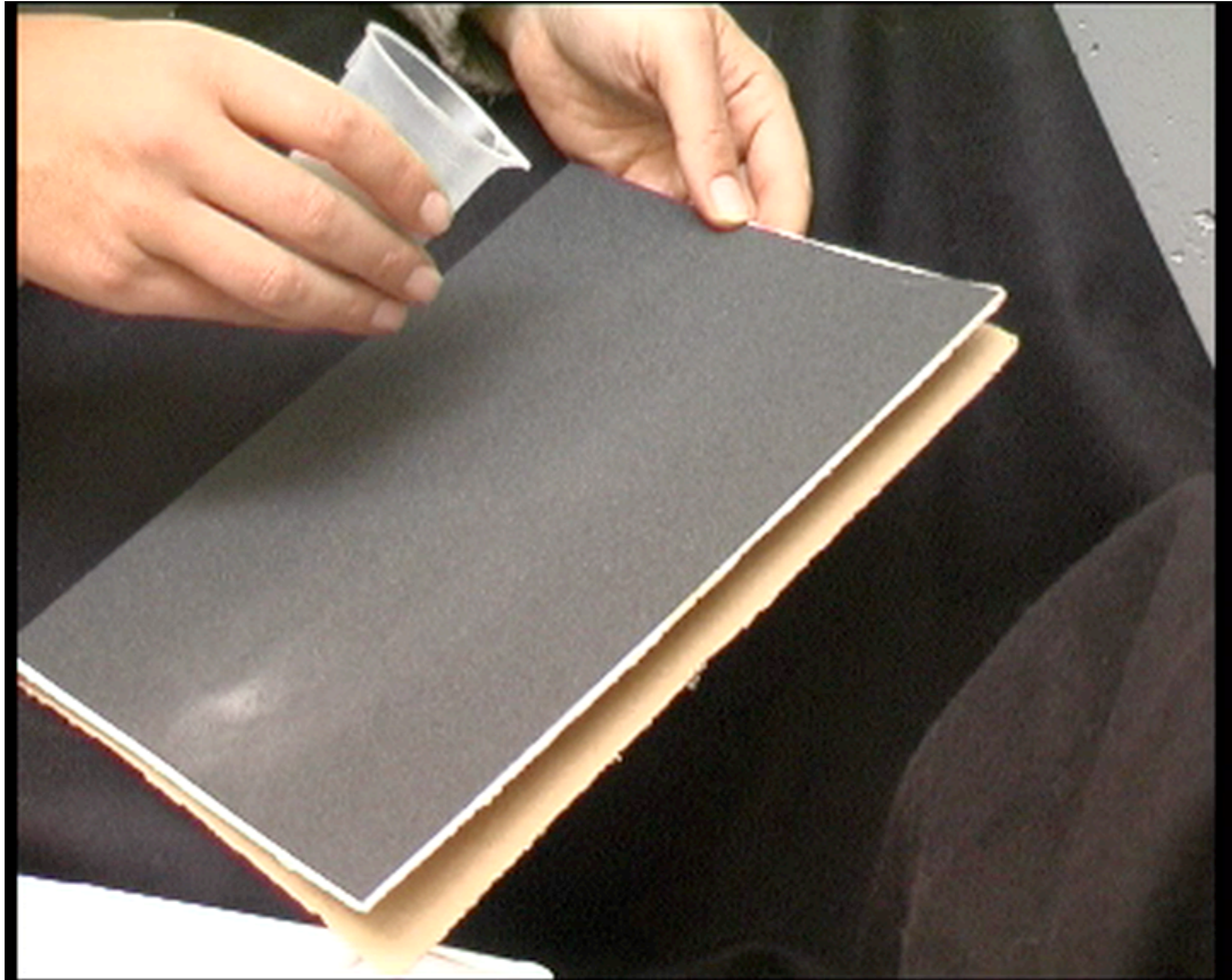


1) Rheology of dry granular flow

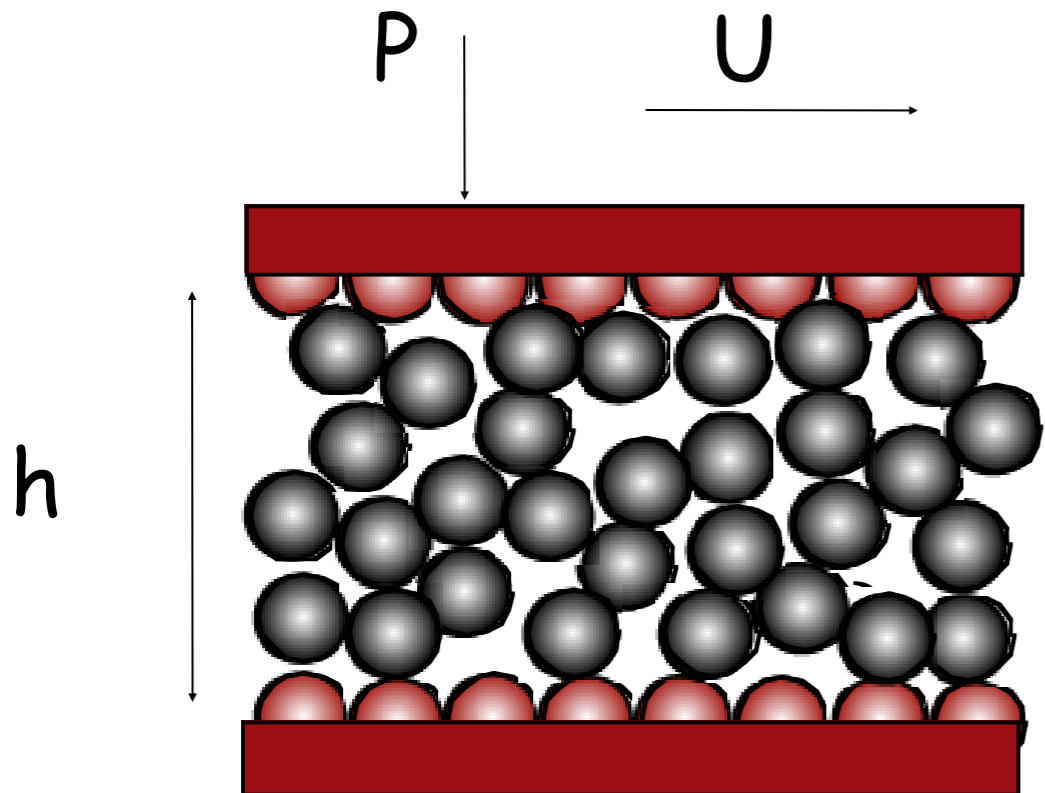
2) Rheology of immersed granular flow

3) dragging objects in a granular medium...

# Dry granular flows



# plane shear under controlled normal stress



Lois et al 2005

Da Cruz et al, PRE 05

GdR Midi, Eur. Phys. J 04

$$\dot{\gamma} = U/h$$

One imposes  $P$  and  $\dot{\gamma}$

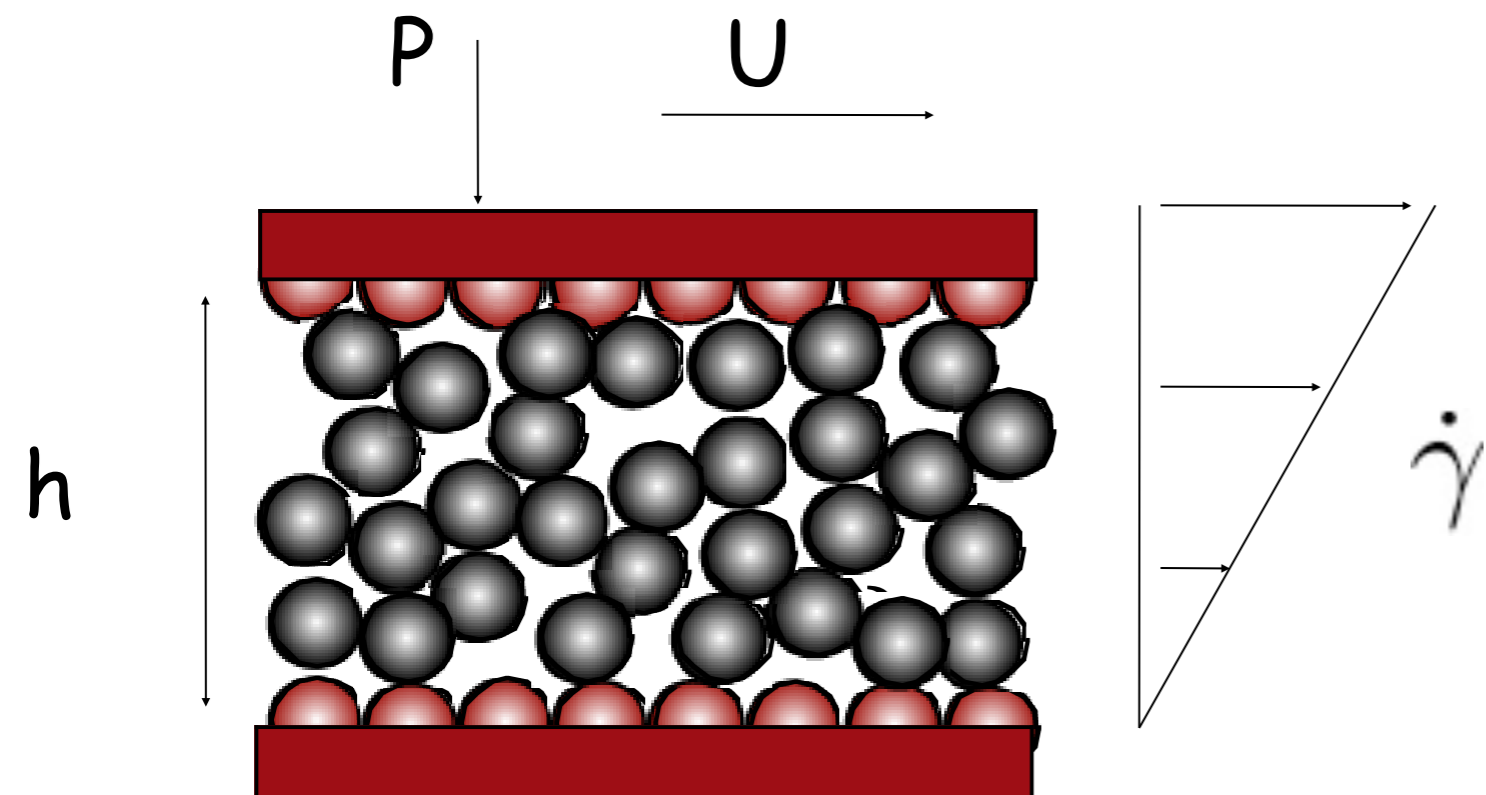
Shear stress  $\tau$ ?

Volume fraction  $\phi$ ?

A single dimensionless number  
(inertial number)

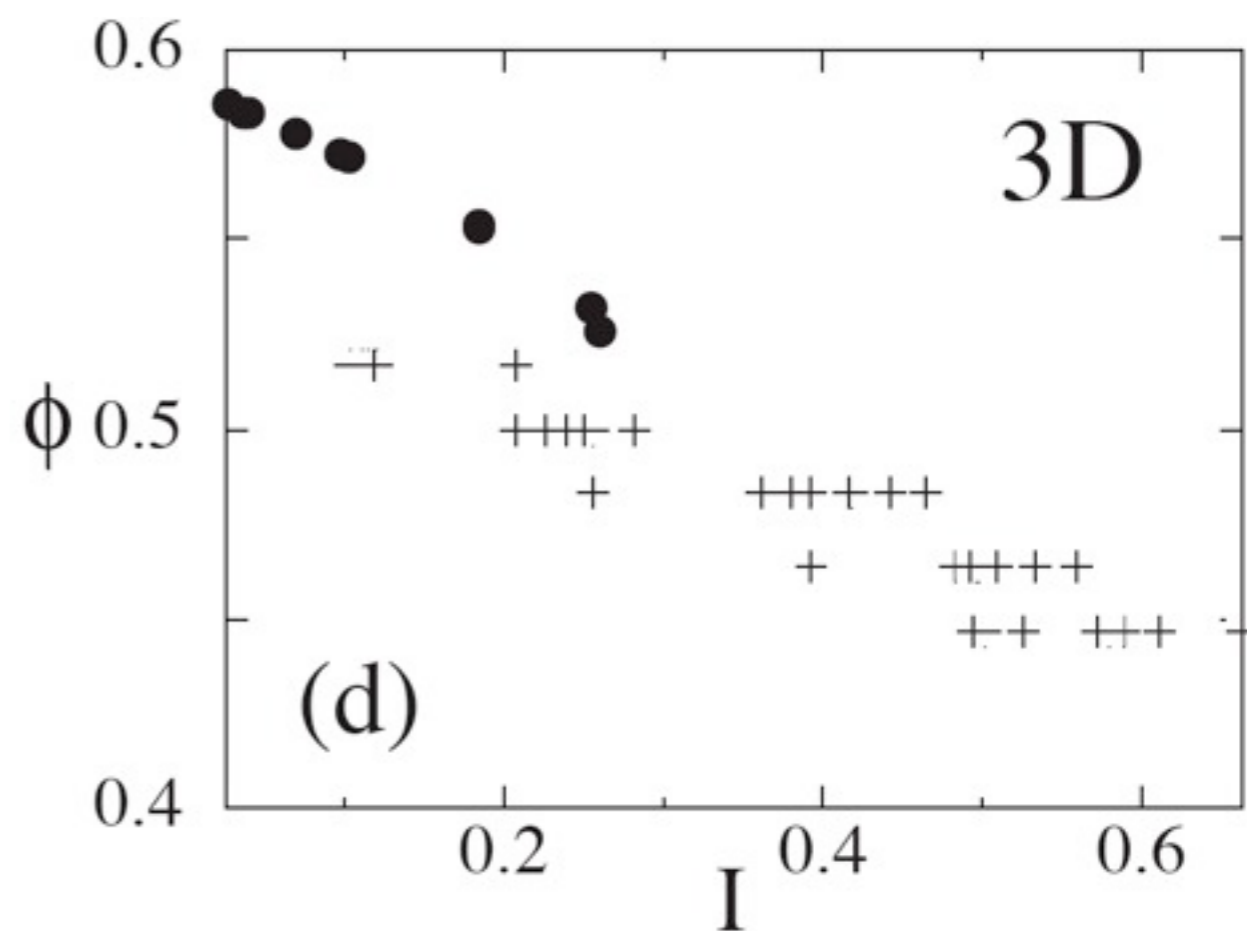
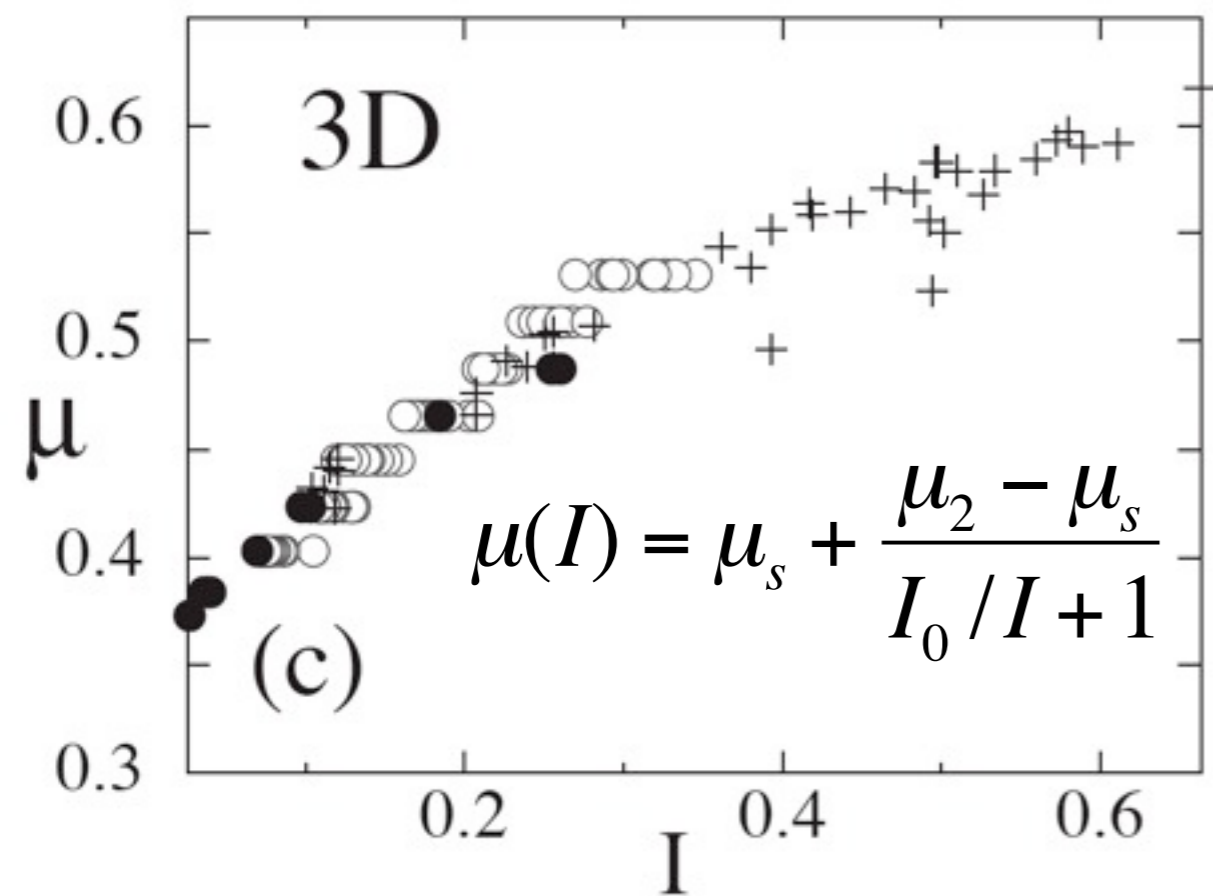
$$I = \frac{\dot{\gamma}d}{\sqrt{P/\rho_s}}$$

(Savage 84,  
Anczyk et al 99)



$$\tau = \mu(I)P$$

$$\Phi = \Phi(I)$$



3D generalisation of the friction law :  
granular flows as a viscoplastic fluid  
(Jop et al Nature 06)

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial z} = 0,$$

$$\rho_s \phi \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial z} \right) = \rho_s \phi g \sin \theta - \frac{\partial P}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xz}}{\partial z},$$

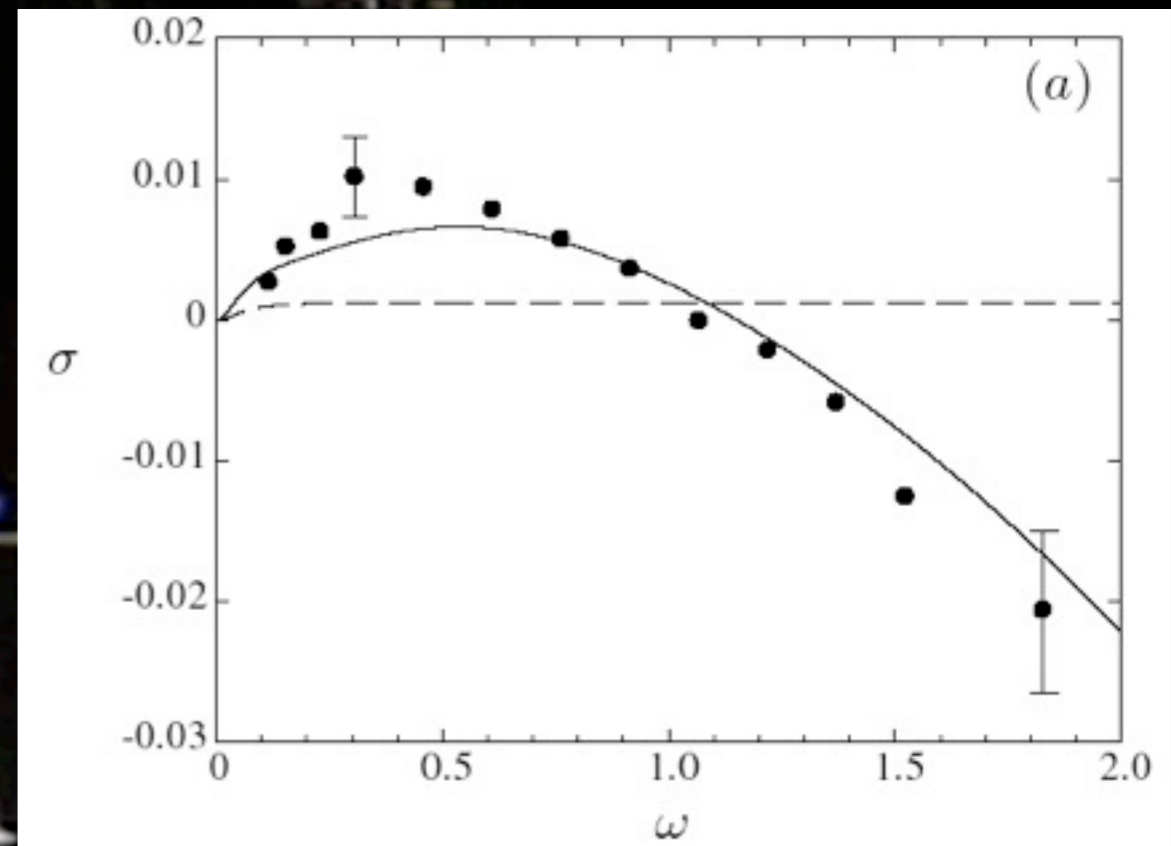
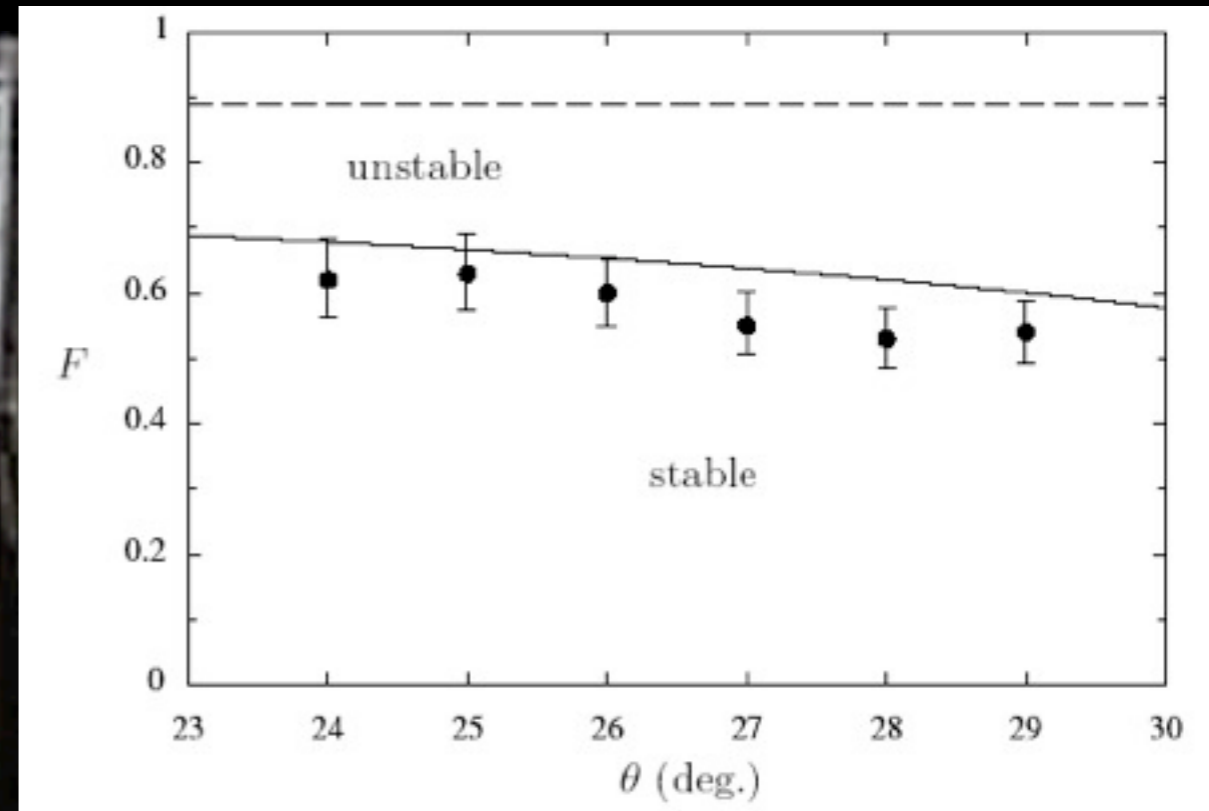
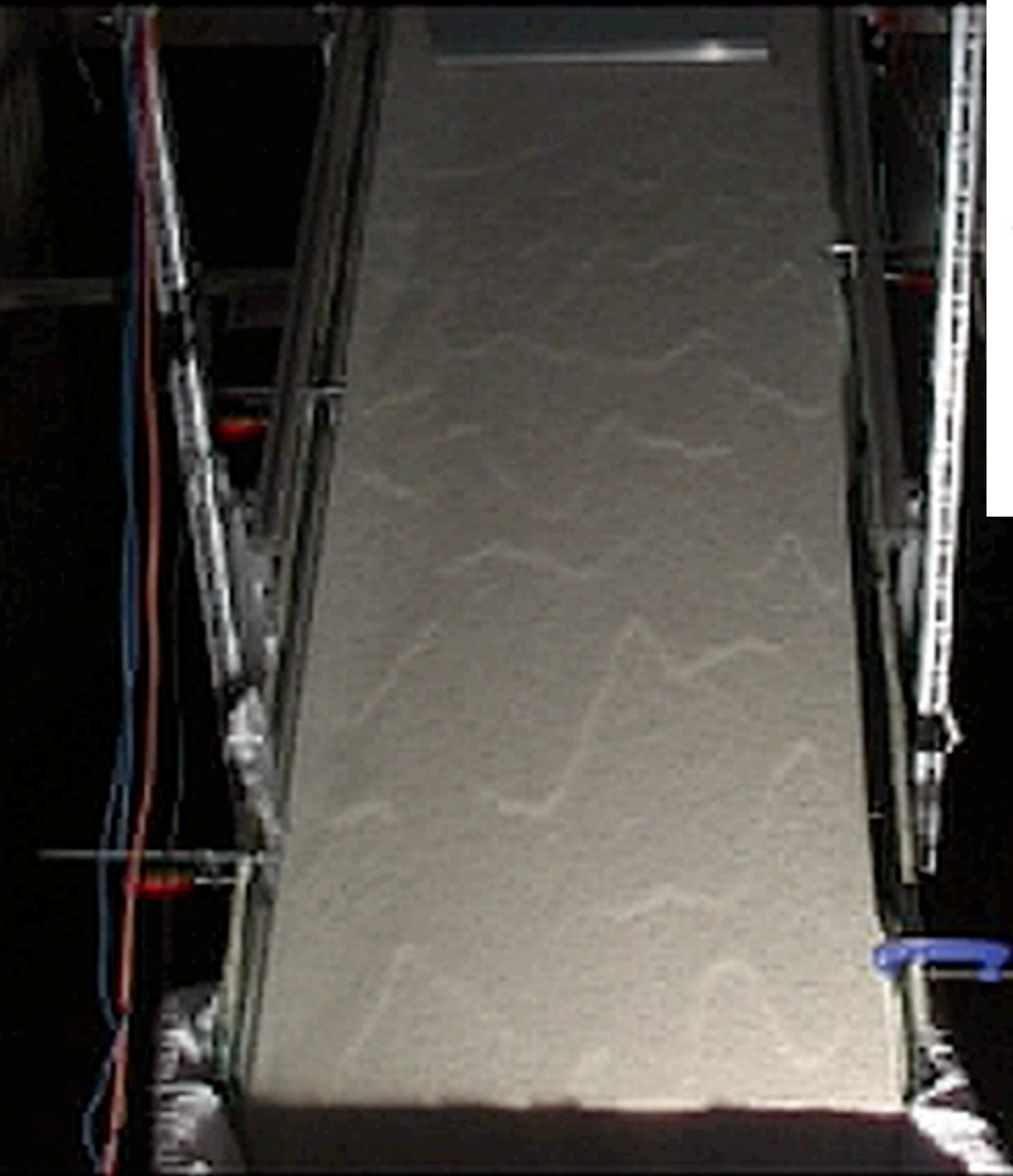
$$\rho_s \phi \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial z} \right) = -\rho_s \phi g \cos \theta - \frac{\partial P}{\partial z} + \frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zz}}{\partial z},$$

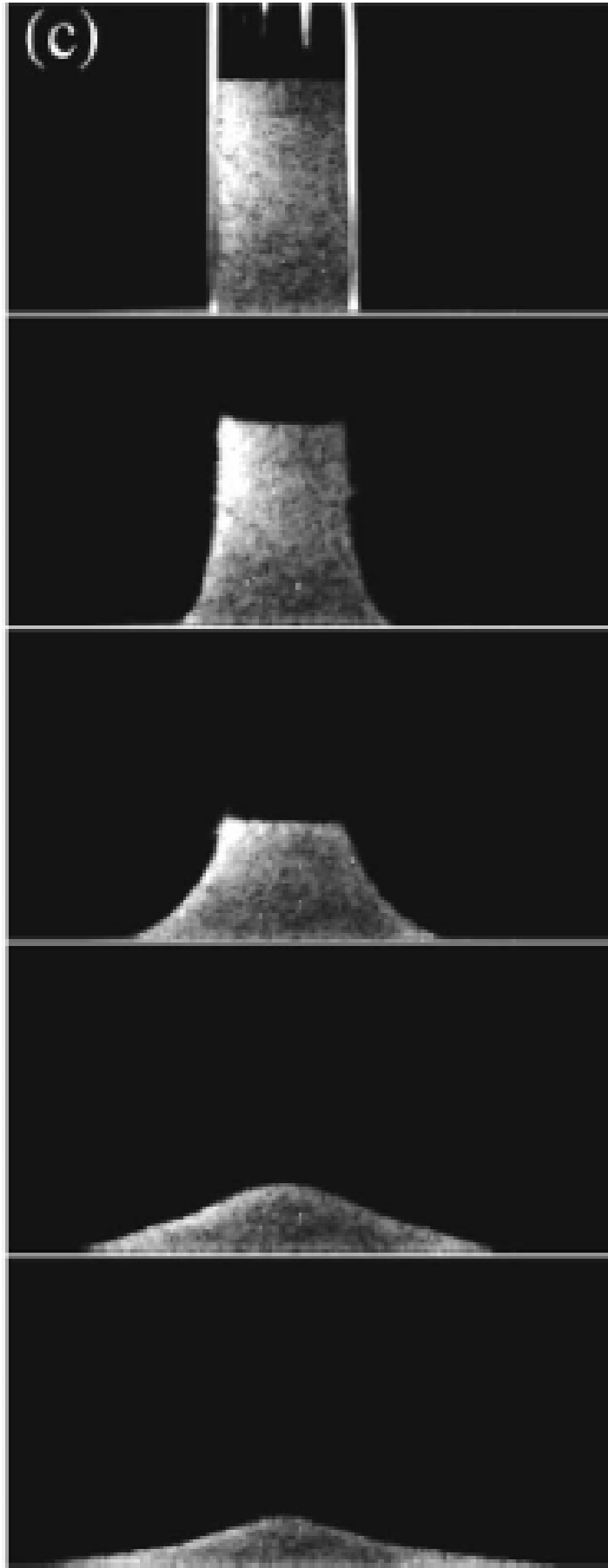
$$\tau_{ij} = \frac{\mu(I)P}{\dot{\gamma}} \dot{\gamma}_{ij} \quad I = \frac{\dot{\gamma}d}{\sqrt{P/\rho_s}}$$

Pressure dependent viscosity

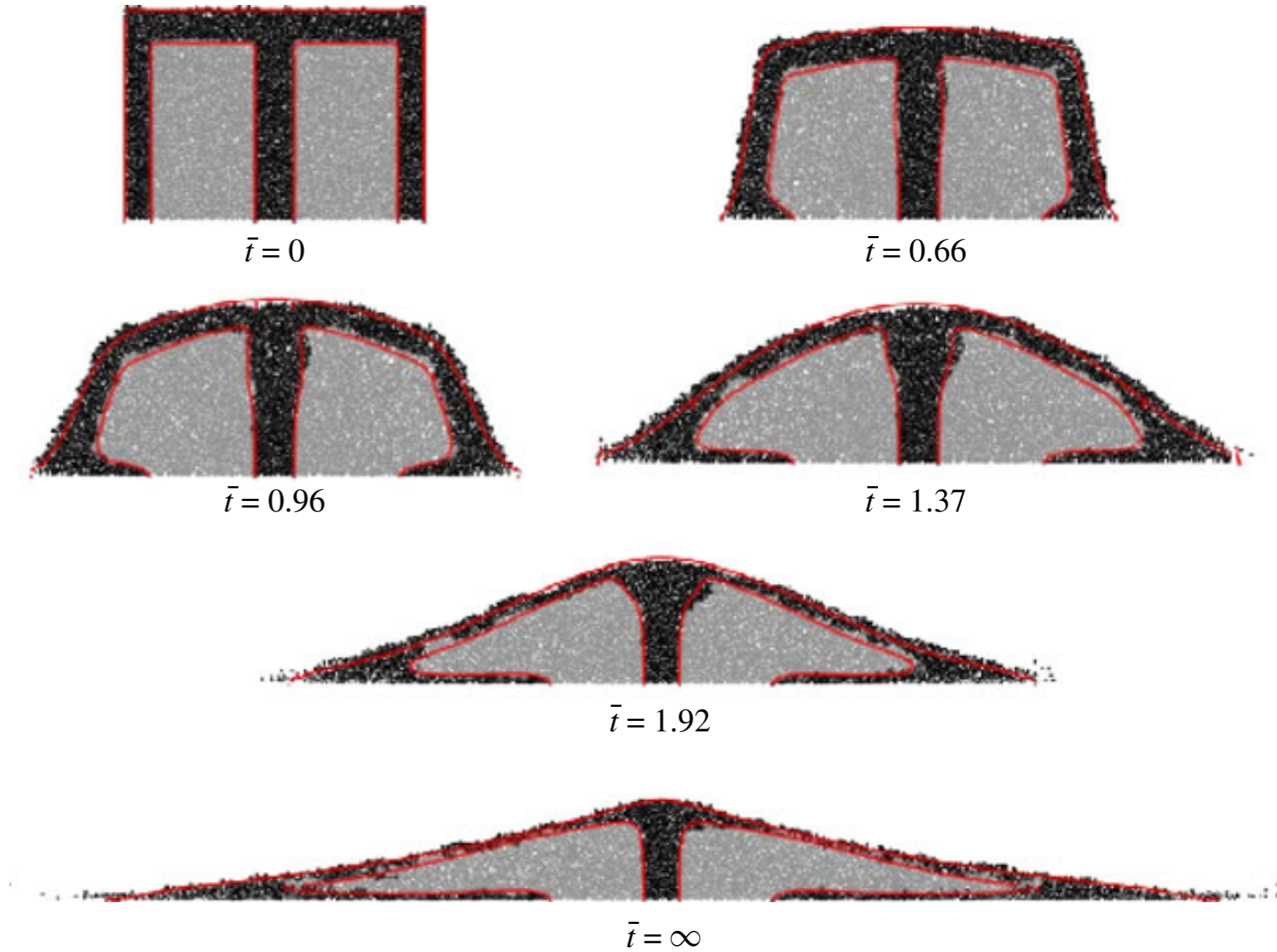


# rolls waves





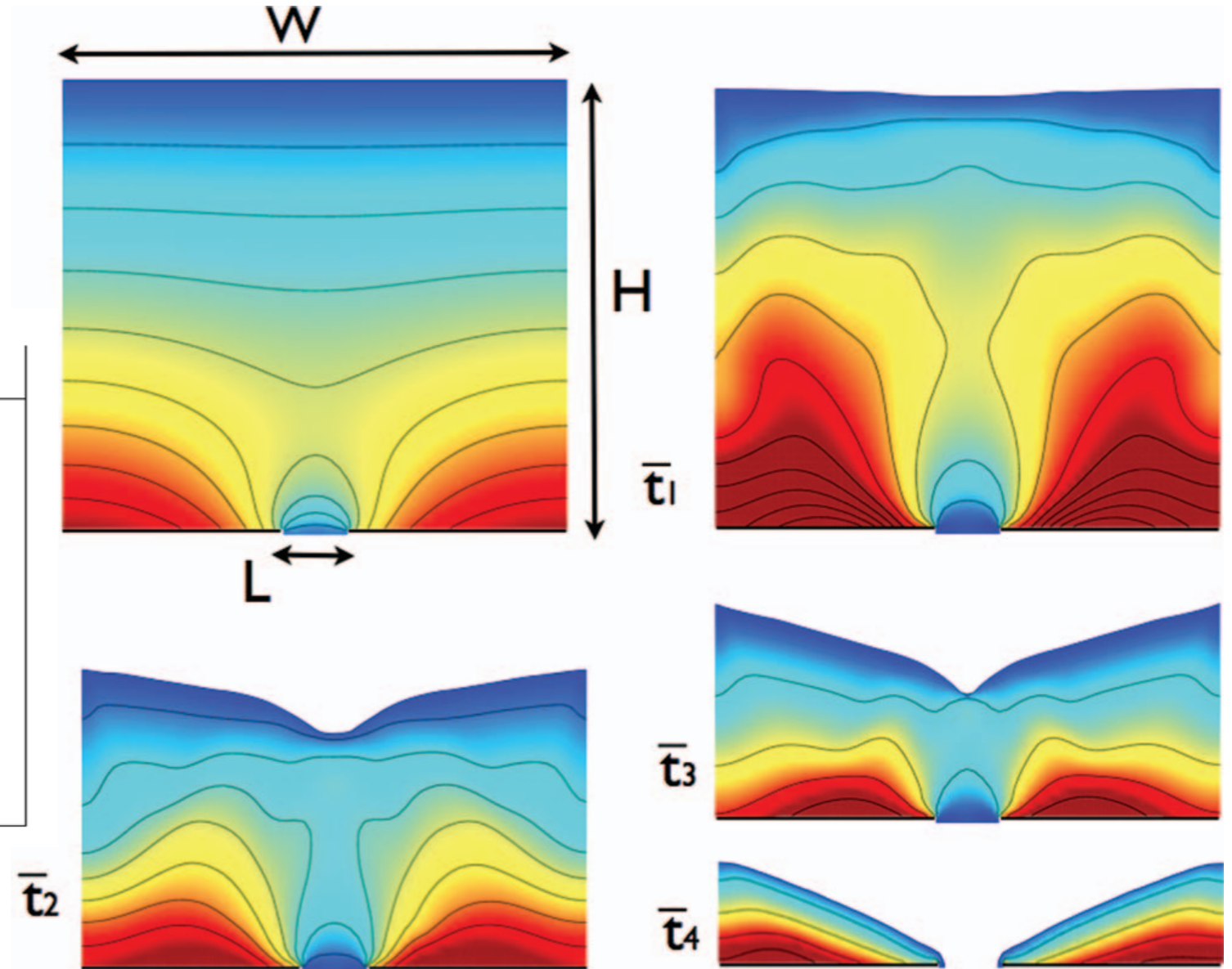
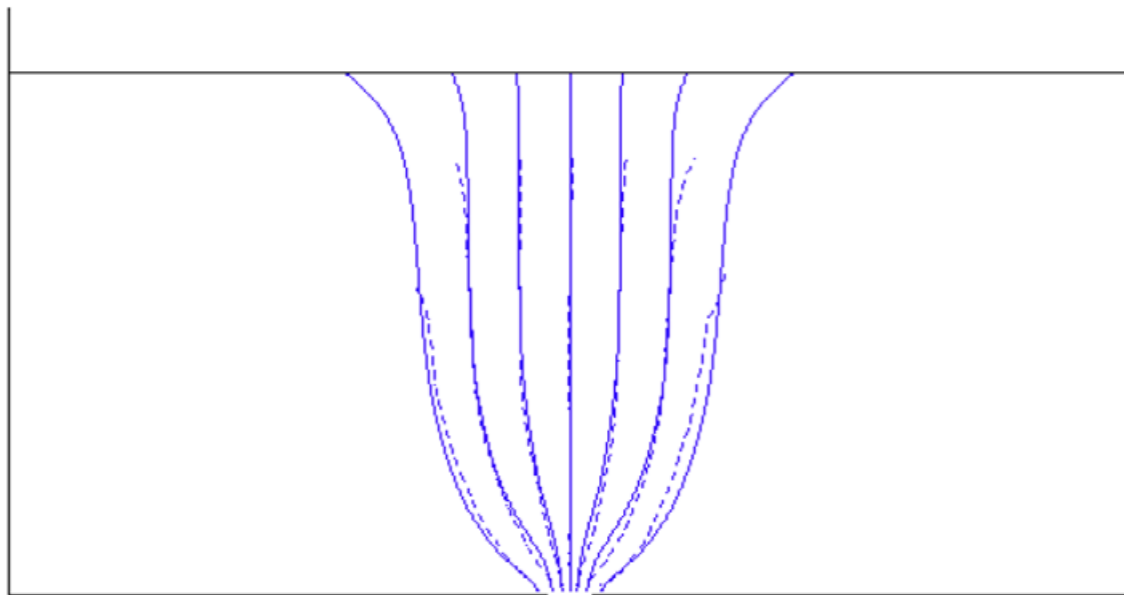
# granular collapse



lajeunesse et al 2005

Lagree et al 2011  
Gerris code..

# flow in a silo



Kamrin 2010

Staron et al 2012



Particles + liquid  
in a dense regime

Difficulties:

multi body problem with :

Hydrodynamic interactions

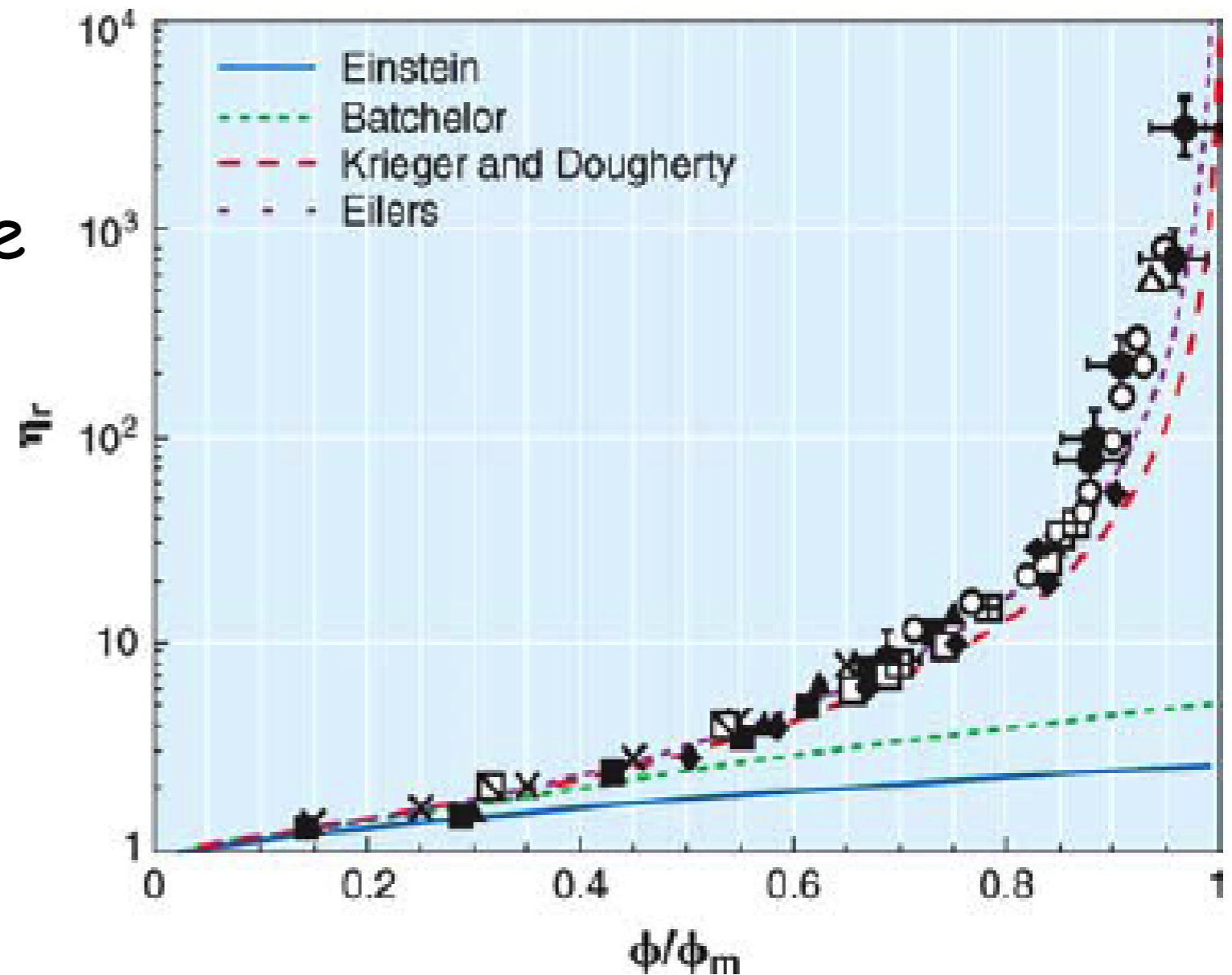
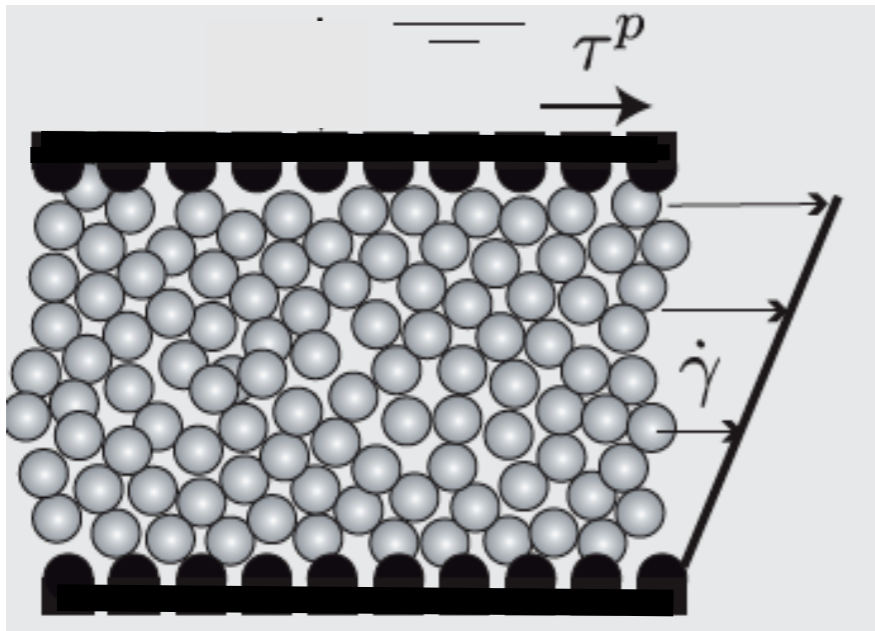
+

Contact between grains



# the big picture of suspensions

- Non brownian rigid spheres
- incompressible and newtonian fluid
- no cohesion, no attractive forces...



$$\tau = \eta_s(\phi)\eta_f\dot{\gamma}$$

Ovarlez et al., *J. Rheol.* 2006

Bonnoit et al., *J. Rheol.* 2009

# A brief history of viscous suspensions

dilute  
suspension

dense  
suspension

jamming

0

$\phi_m$

$\phi$

1906

1970,72

1988 =>...

2007 =>...

time



# A very brief history of viscous suspensions

dilute  
suspension

dense  
suspension

jamming

0

$\phi_m$

$\phi$



1906  
Einstein

1970,72  
Batchelor  
Green

1988 =>...  
Brady, Bossis  
Nott, Morris,..

2007 =>...  
Olson teitel07, Tighe et al 09,  
Heussinger, Andreotti et al 10,12  
Wyart 12, Berthier12,..

Single  
particle

Pair  
interaction

hydrodynamic  
interactions

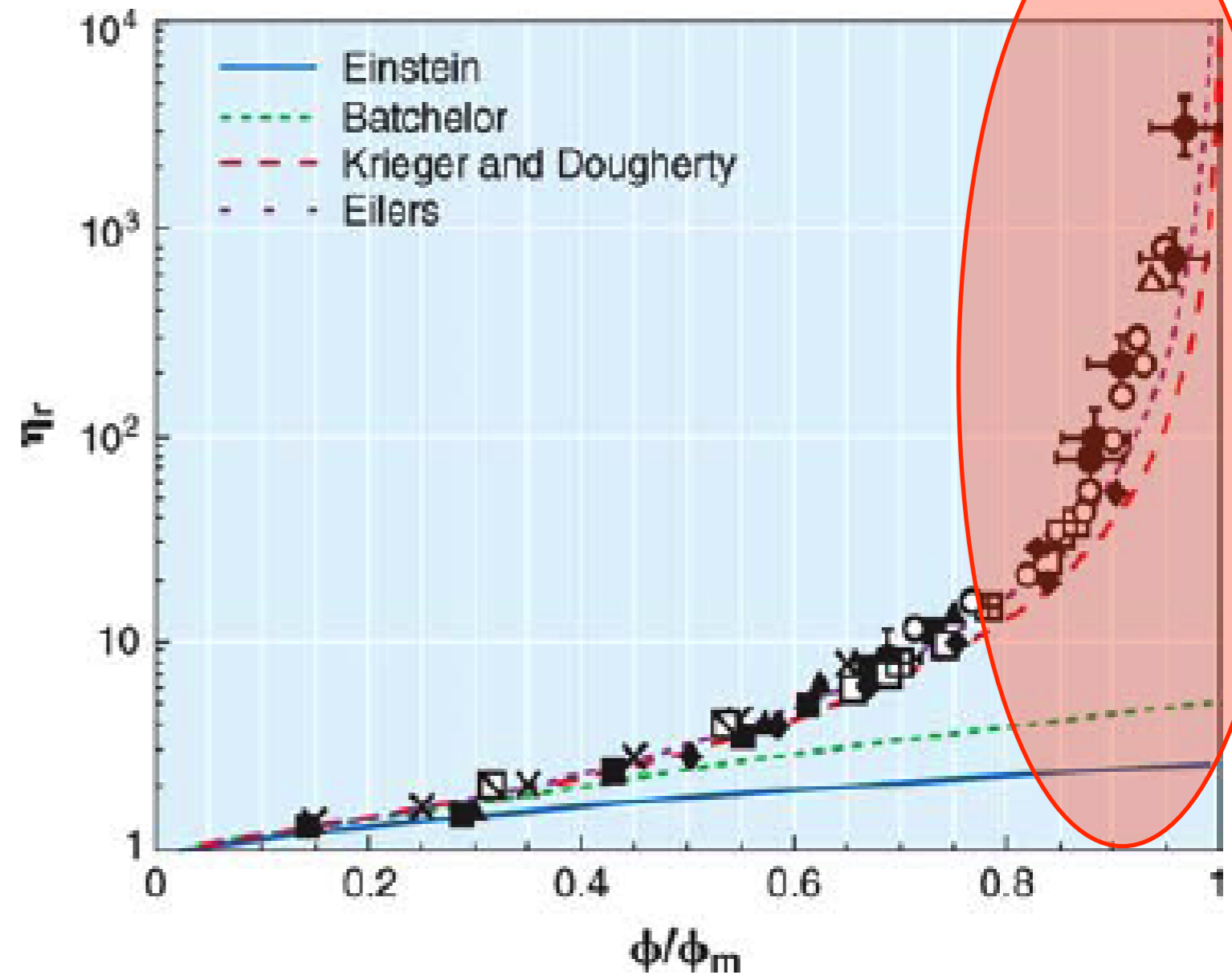
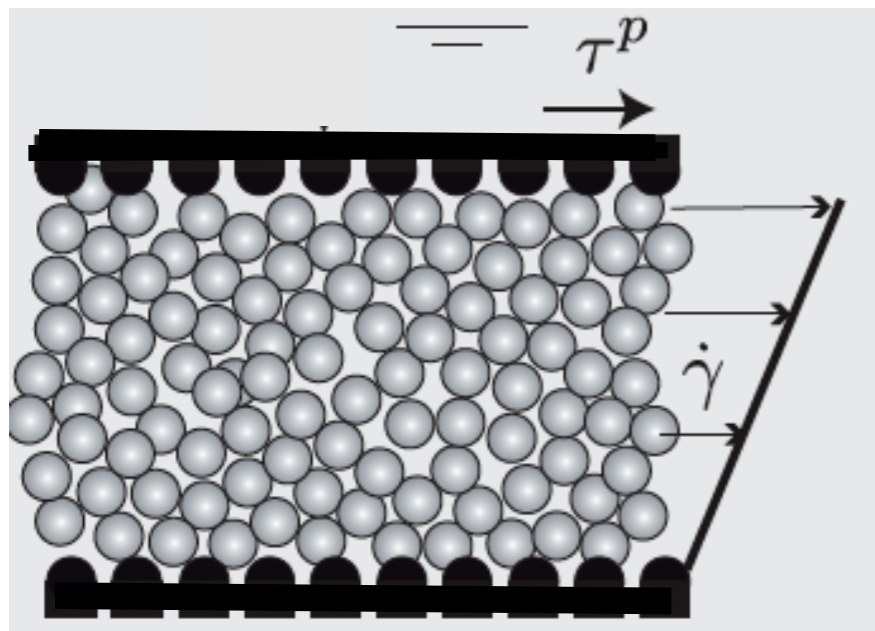
steric/elastic  
interactions

$$\eta = \eta_f \left( 1 + \frac{5}{2} \phi + \alpha \phi^2 .. \right)$$

towards  
Constitutive  
equations

understanding  
the divergence

link with granular media?



$$\tau = \eta_s(\phi)\eta_f\dot{\gamma}$$

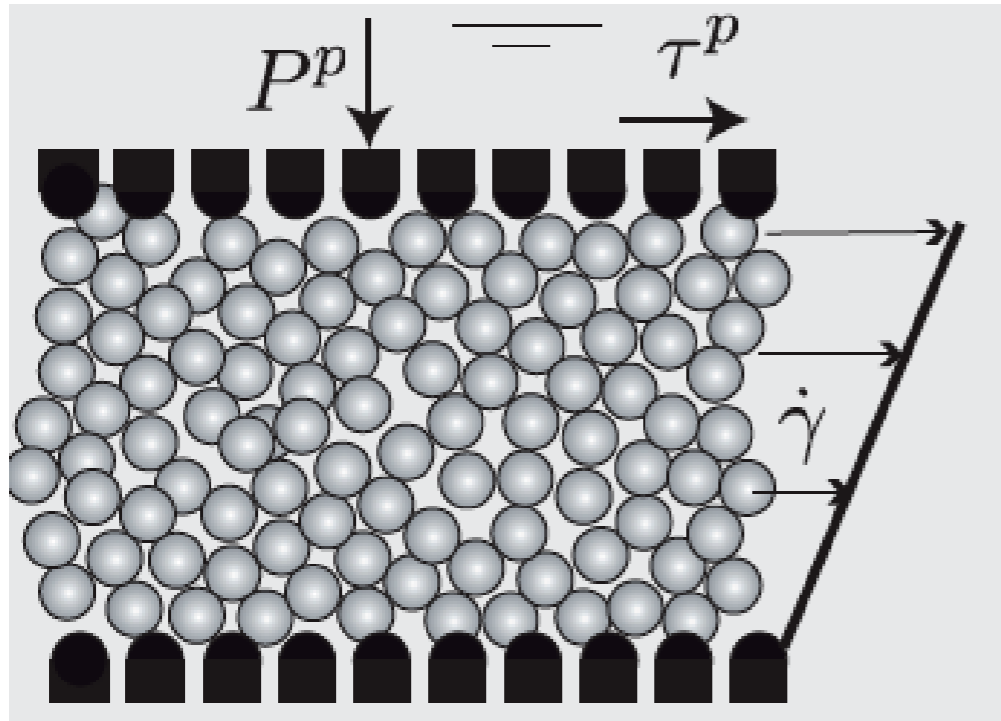
Ovarlez et al., *J. Rheol.* 2006

Bonnoit et al., *J. Rheol.* 2009



# The granular approach of suspensions!!

If inertia is negligible...



One imposes  $P$  and  $\dot{\gamma}$

Shear stress  $\tau$ ?

Volume fraction  $\phi$ ?

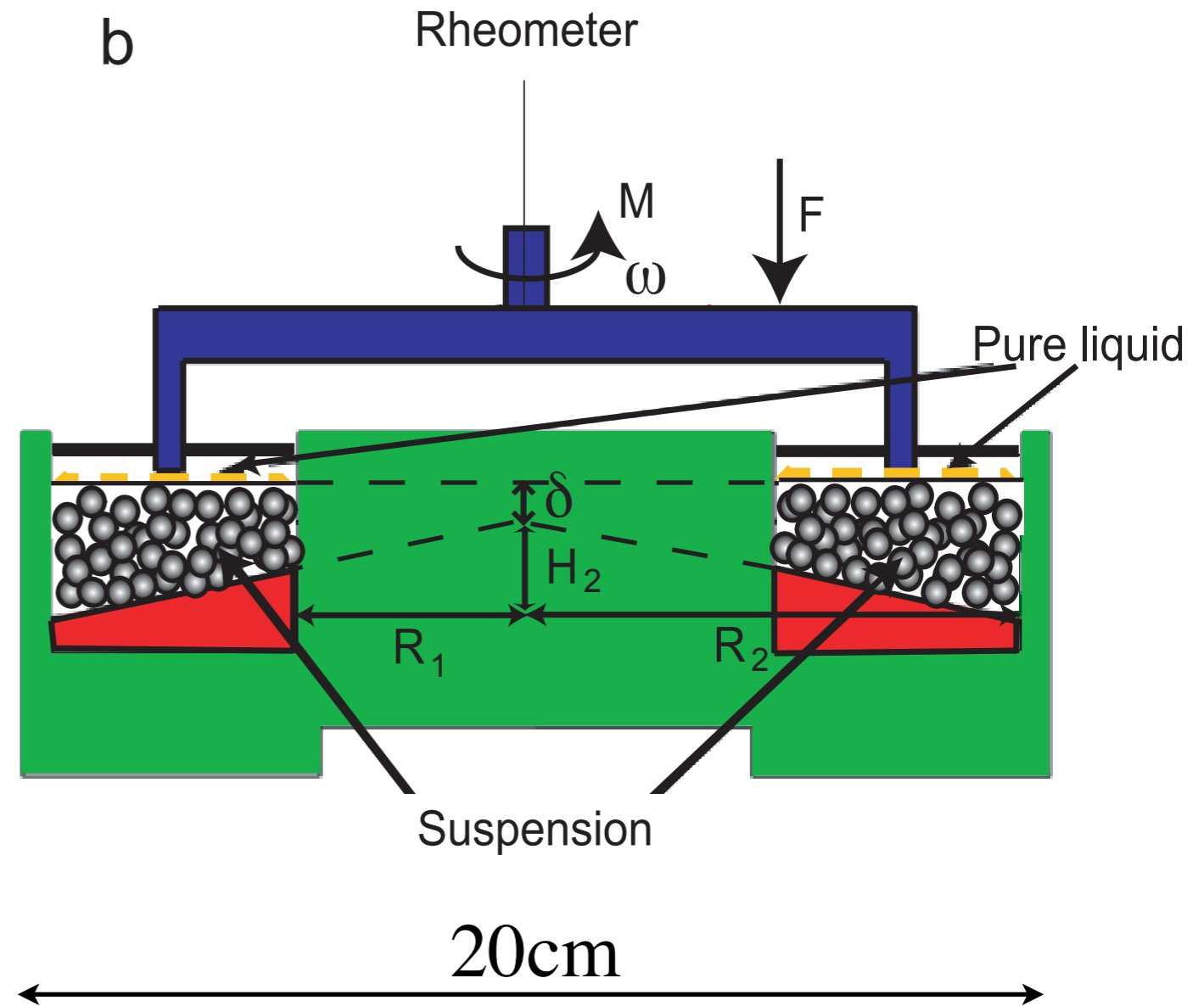
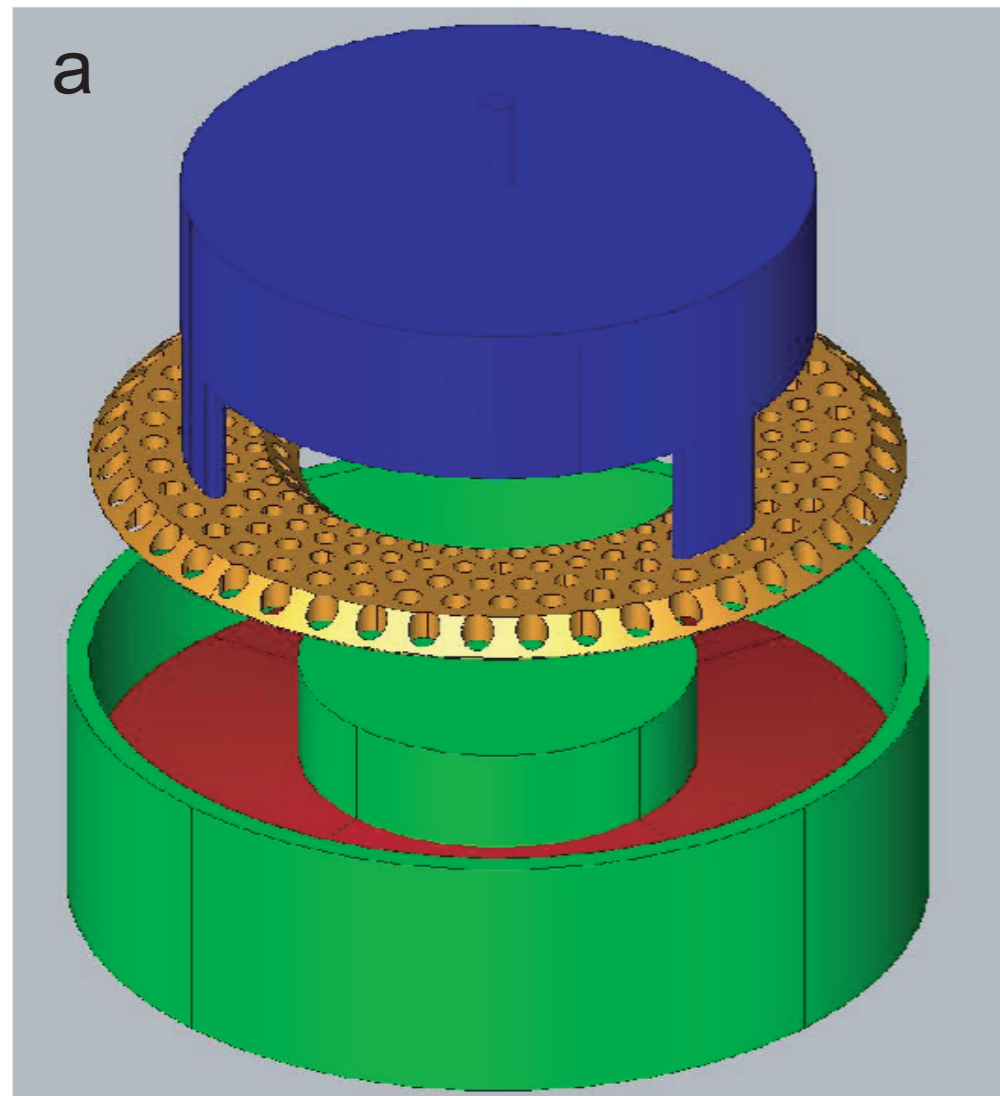
$$t_{micro} = \frac{\eta_f}{Pp}$$

$$I_v = \frac{\eta_f \dot{\gamma}}{P_p}$$

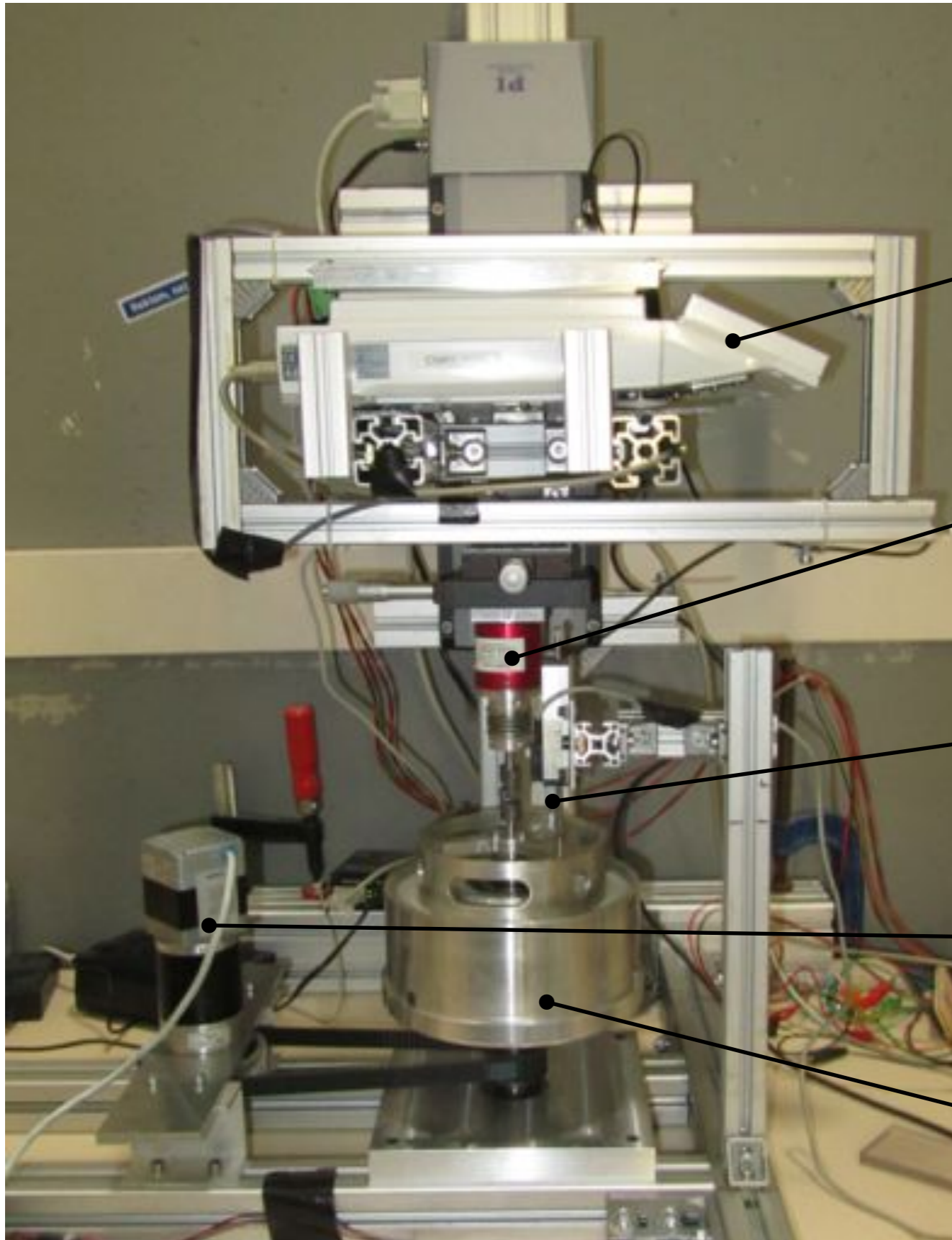
$$\tau = \mu(I_v) P_p$$

$$\phi = \phi(I_v)$$

## François Boyer's Phd...



580  $\mu\text{m}$  and 1mm polystyrene beads in viscous fluid  
(Polyethylene glycol-n) of the same density ( $1.05\text{g}/\text{cm}^3$ )



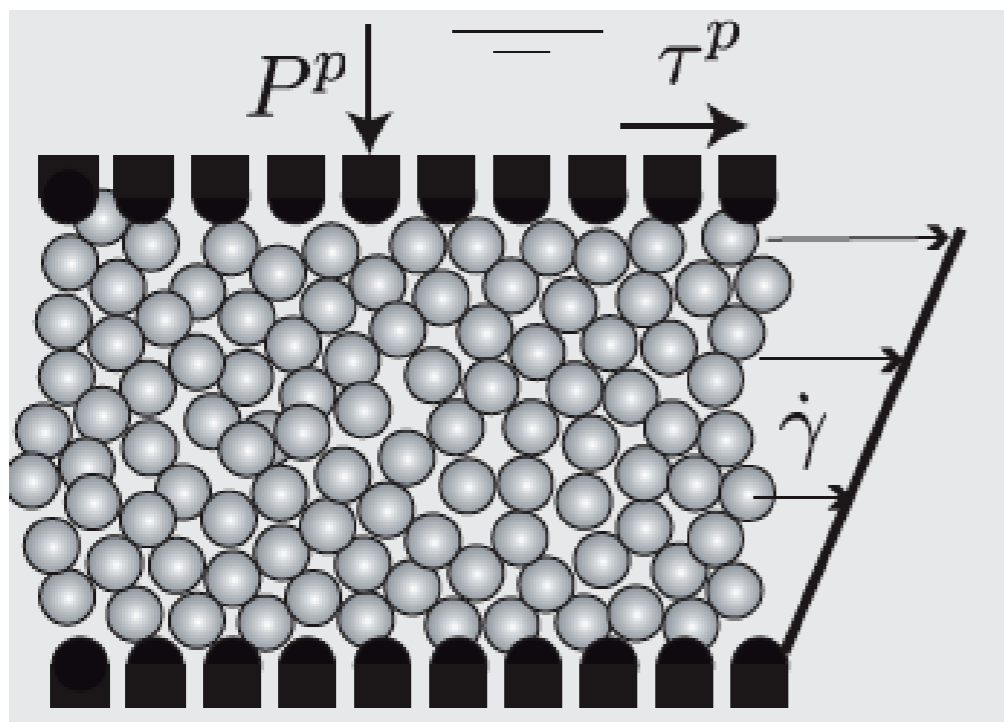
high precision scale

torquemeter

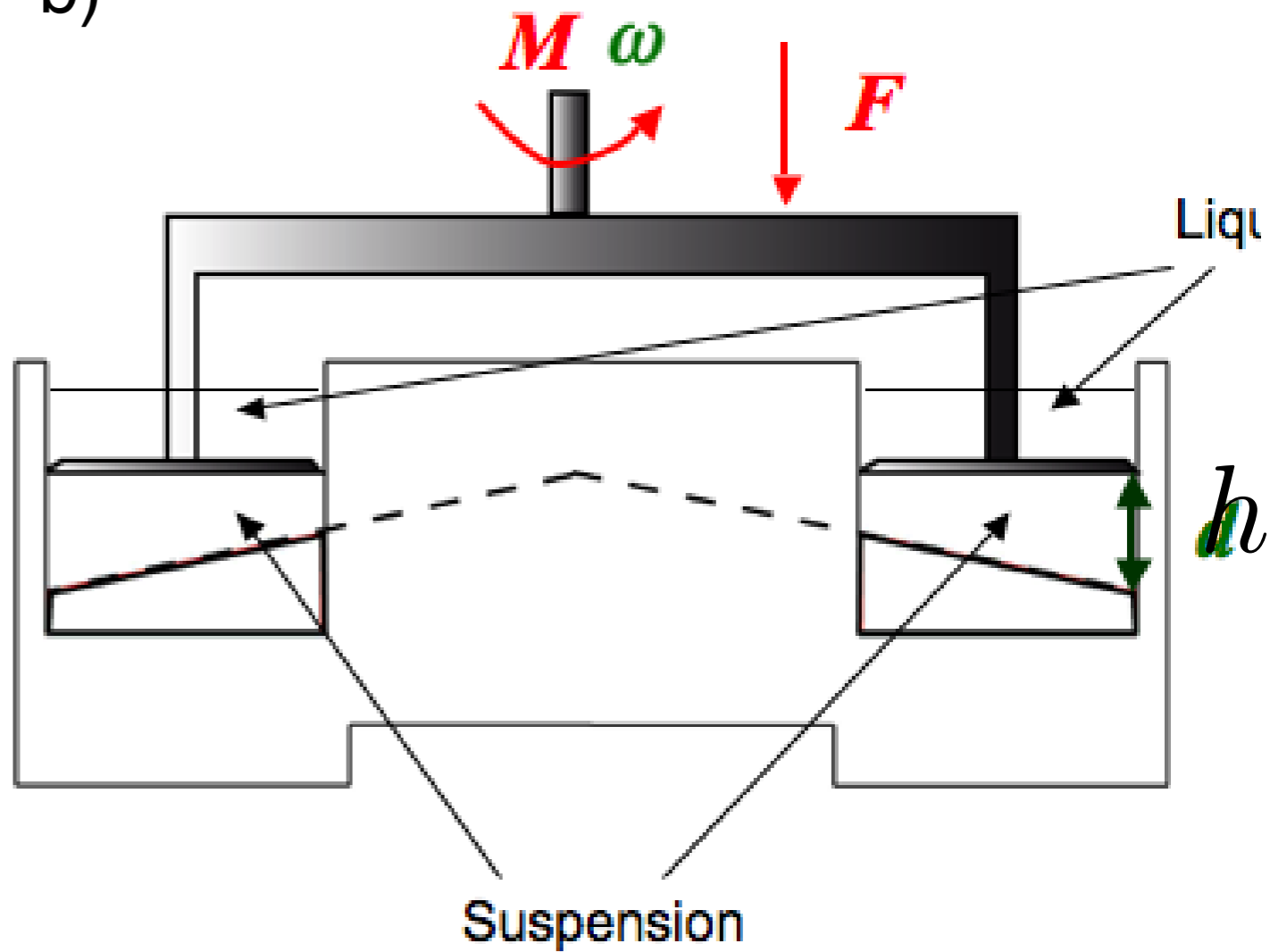
position sensor

motor

shear cell



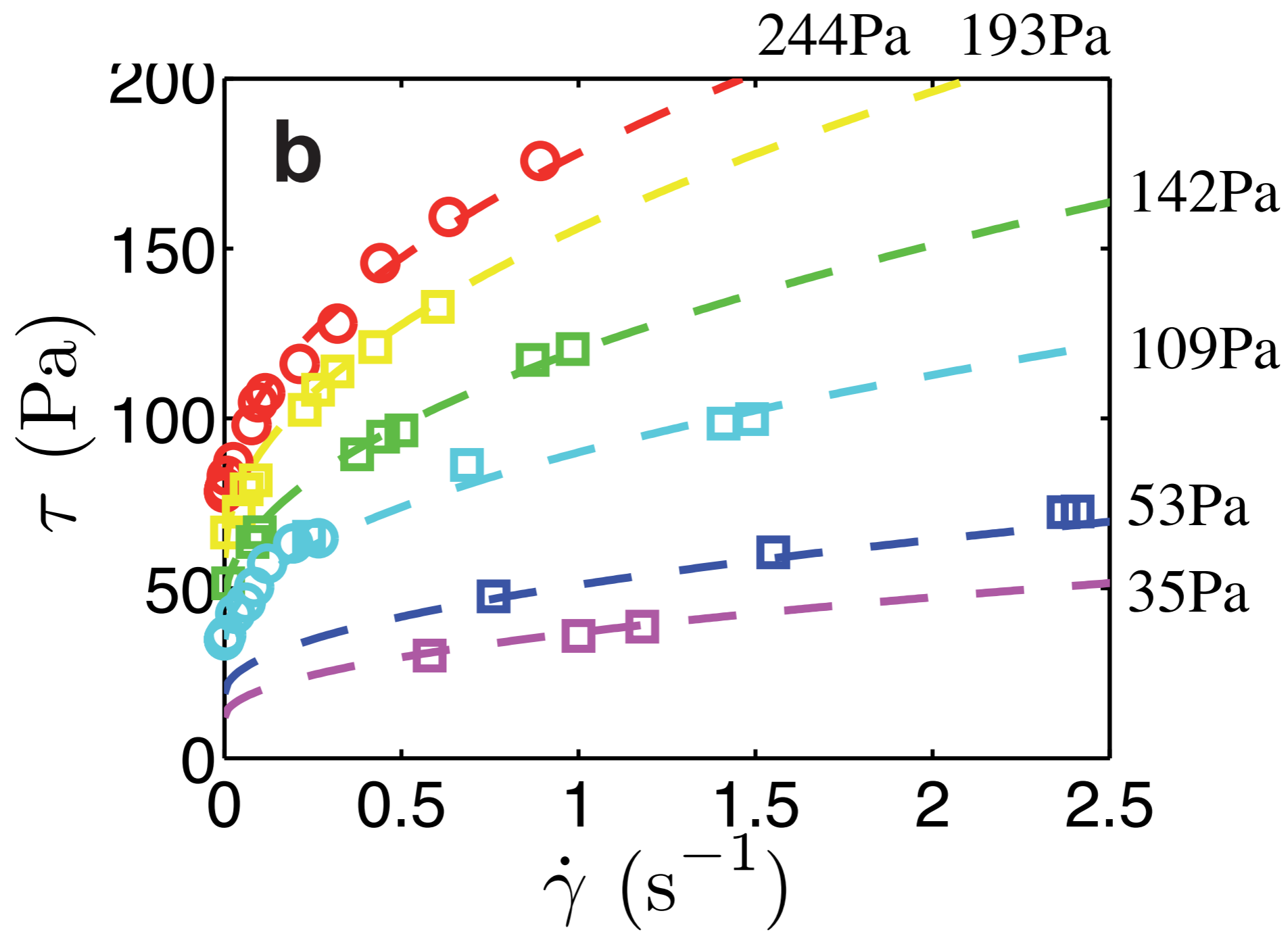
b)



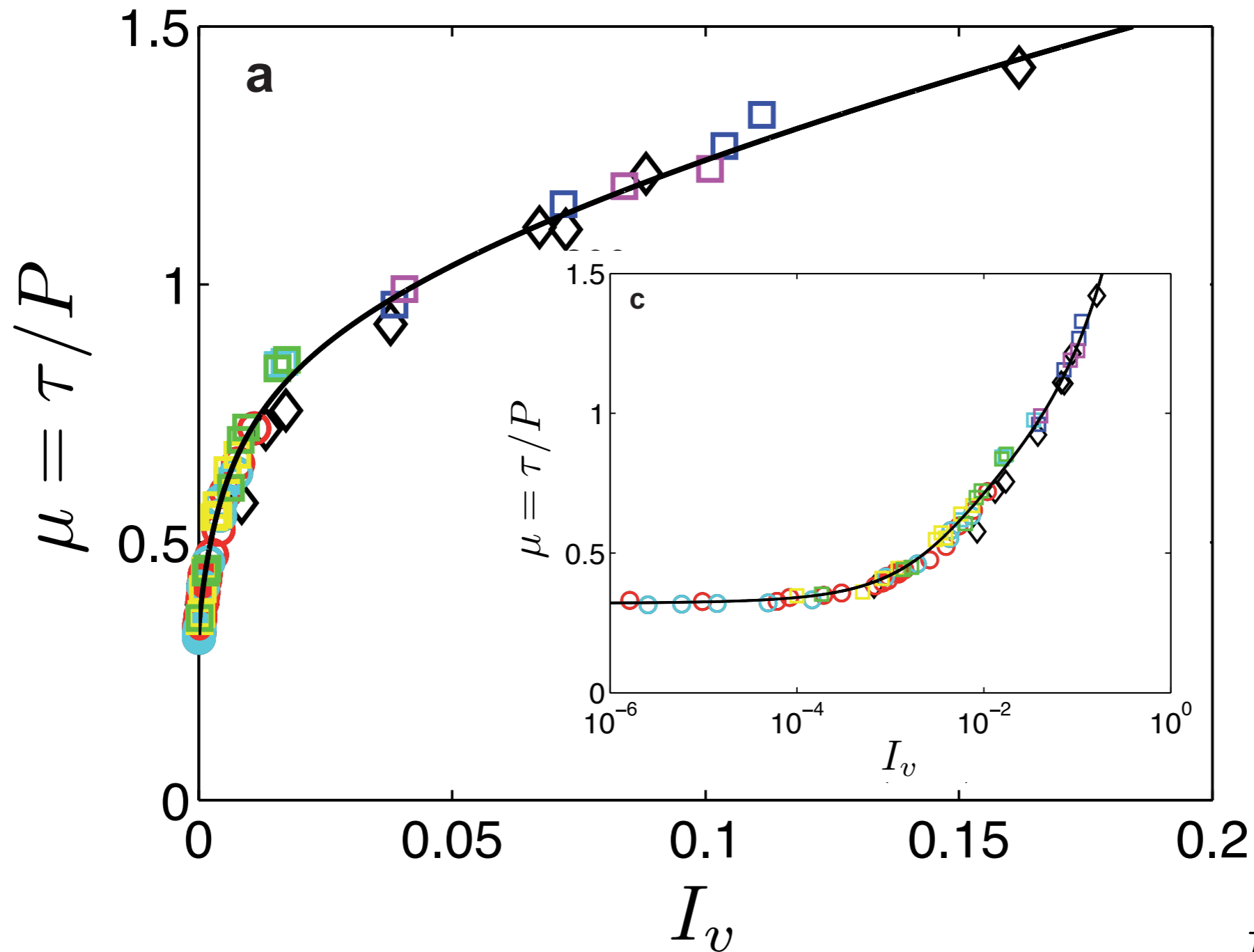
$$\text{Imposed} \left\{ \begin{array}{l} F \implies P_p \\ M \implies \tau \end{array} \right.$$

$$\text{Measured} \left\{ \begin{array}{l} \omega \implies \dot{\gamma} \\ h \implies \phi \end{array} \right.$$

$$\begin{aligned} I_v &= \frac{\eta_f \dot{\gamma}}{P_p} \\ &\implies \mu = \frac{\tau}{P_p} \\ &\phi \end{aligned}$$

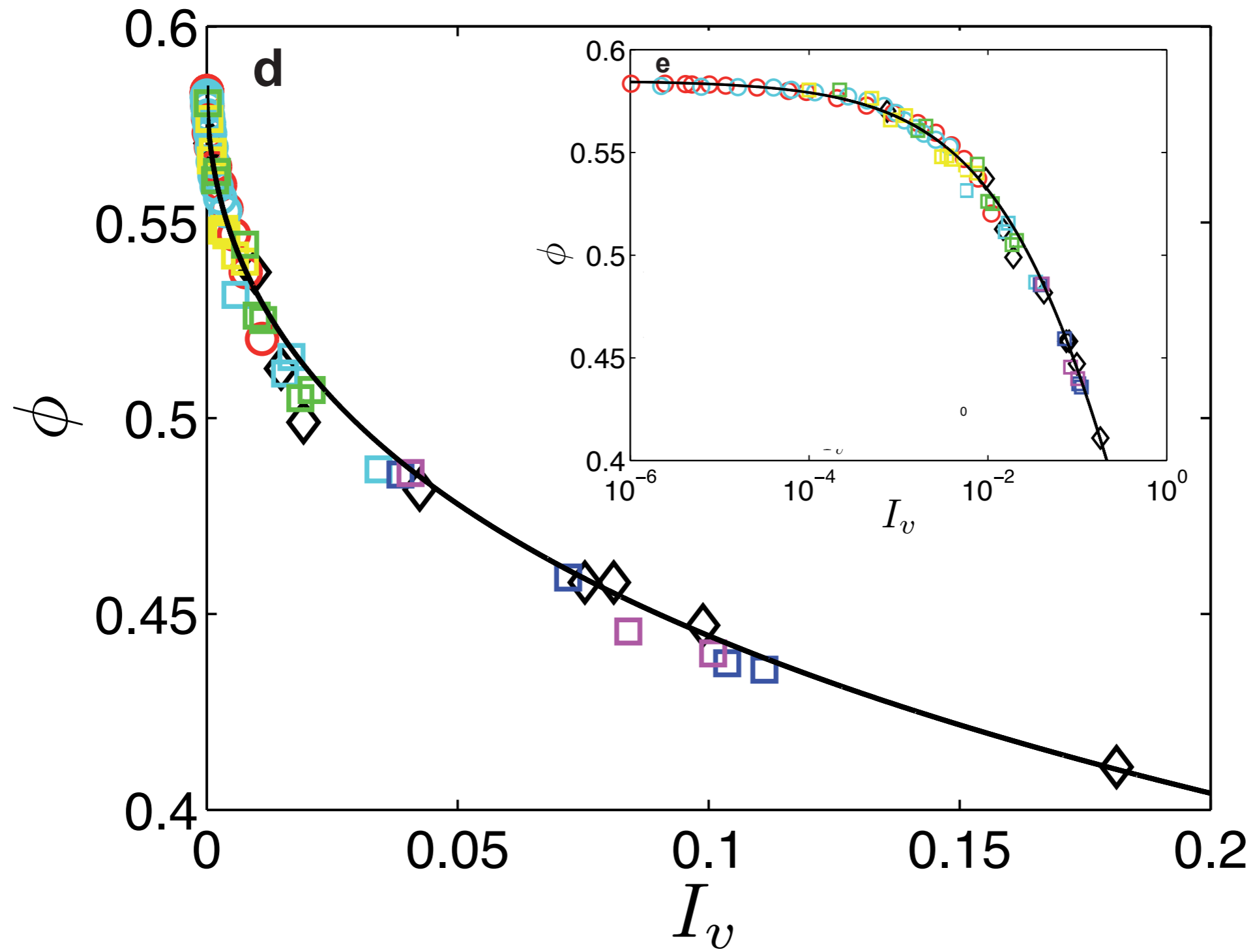


a friction law  $\mu(I_v)$  : different viscosities,  
 pressure, shear rate, particle size,..



$$\mu_c = 0.32$$

$$I_v = \frac{\eta_f \dot{\gamma}}{P_p}$$

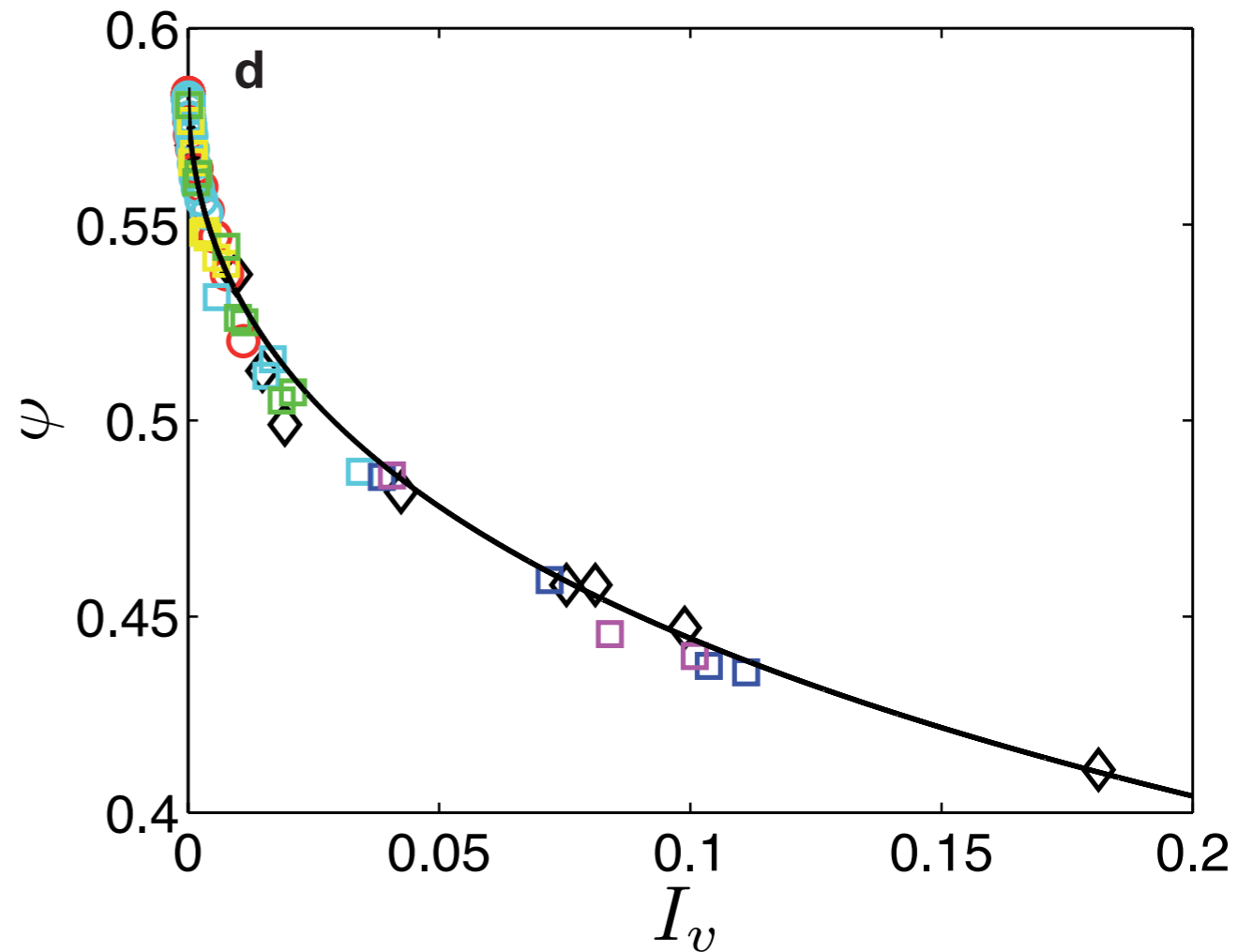
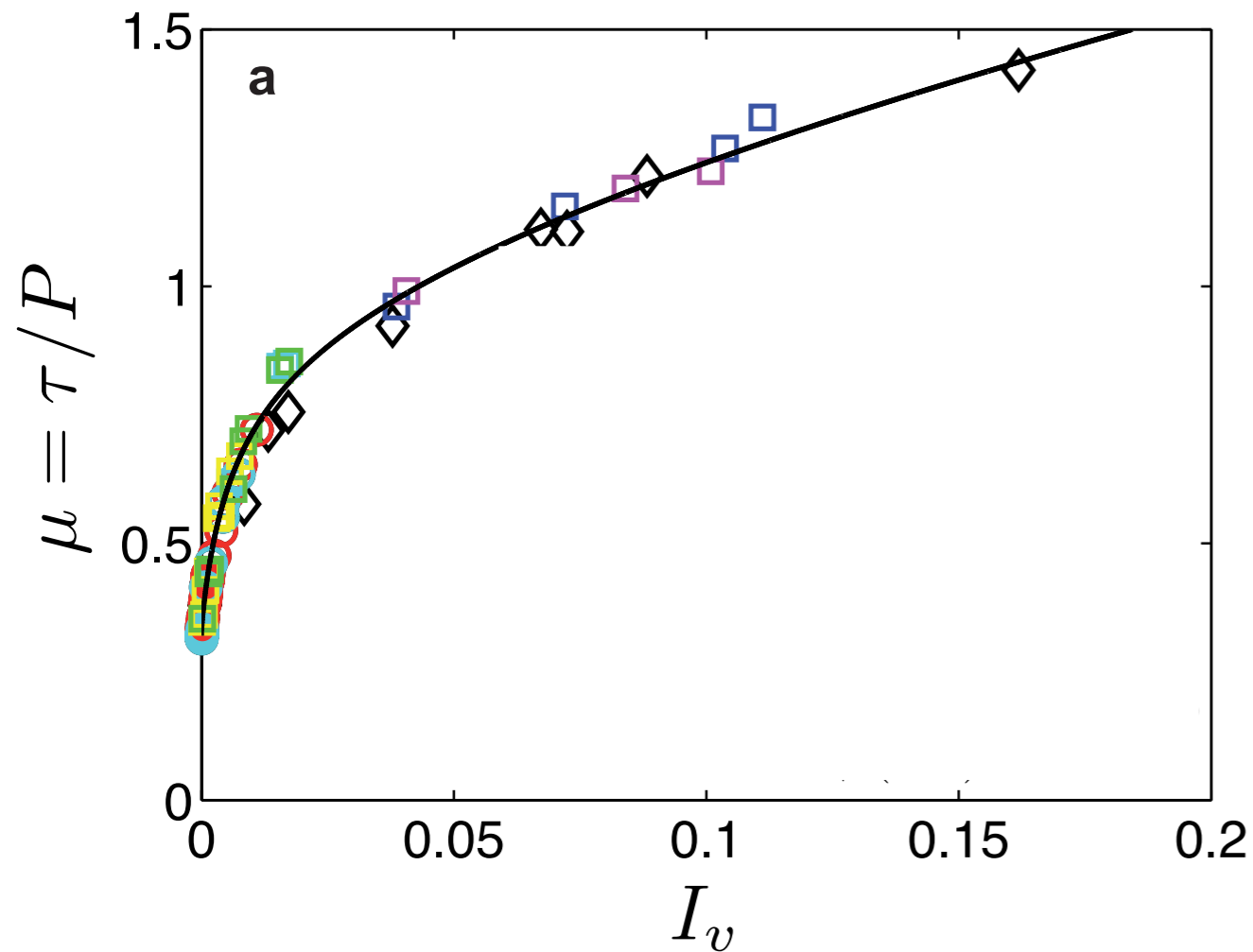


$$\phi_c = 0.585$$

a phenomenological  
rheology:

$$\tau = \mu(I_v) P_p$$

$$\phi = \phi(I_v)$$



$$\mu(I_v) = \mu^c(I_v) + \mu^h(I_v) \text{ with}$$

$$\mu^c(I_v) = \mu_1 + \frac{\mu_2 - \mu_1}{1 + I_0/I_v}$$

$$\mu^h(I_v) = I_v + \frac{5}{2} \phi_m I_v^{1/2}$$

$$\phi(I_v) = \frac{\phi_m}{1 + I_v^{1/2}}$$



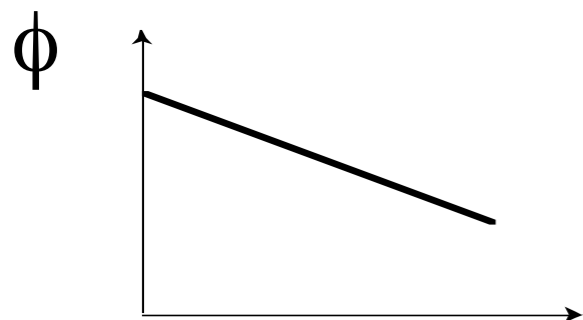
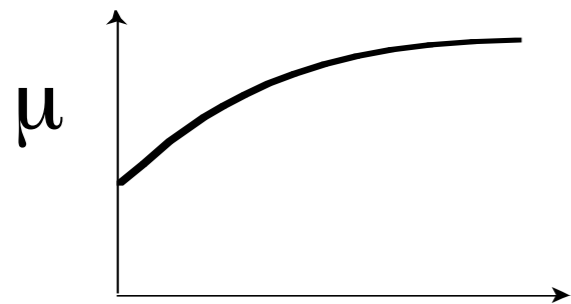
# Link with the rheology of dense suspensions :

imposed pressure

$$\frac{\tau}{P_p} = \mu(I_v)$$

$$\phi = \phi(I_v)$$

$$I_v = \frac{\eta_f \dot{\gamma}}{P_p}$$



Or

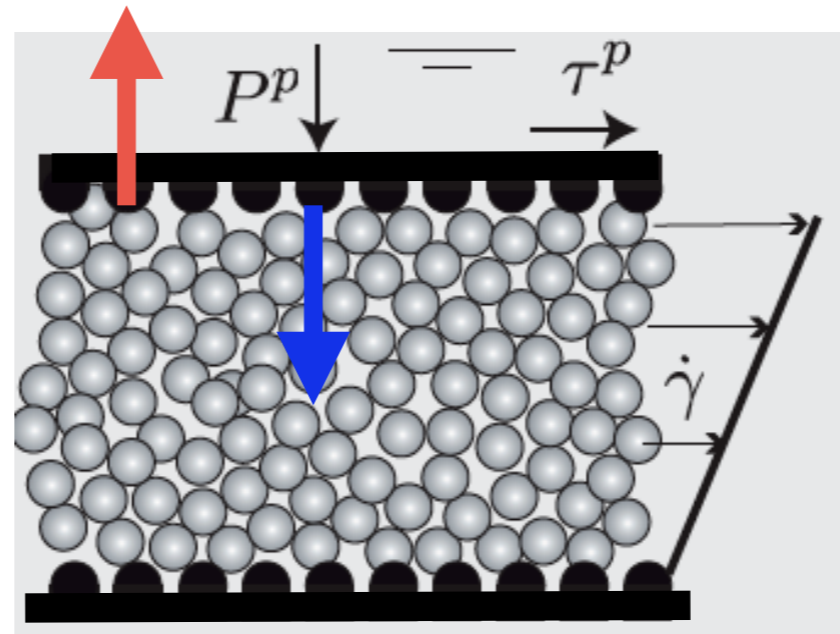
volume imposed

$$\tau = \eta_s(\phi) \eta_f \dot{\gamma}$$

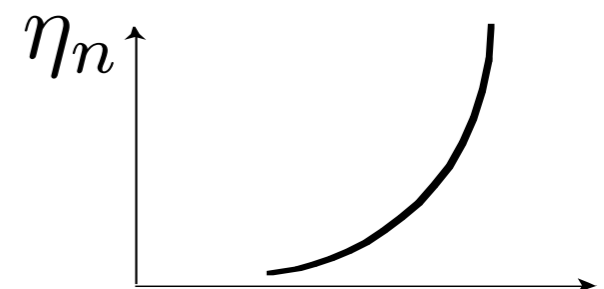
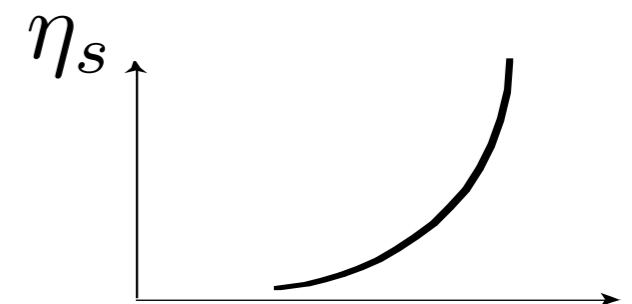
$$P_p = \eta_n(\phi) \eta_f \dot{\gamma}$$

(Morris and Boulay 99)

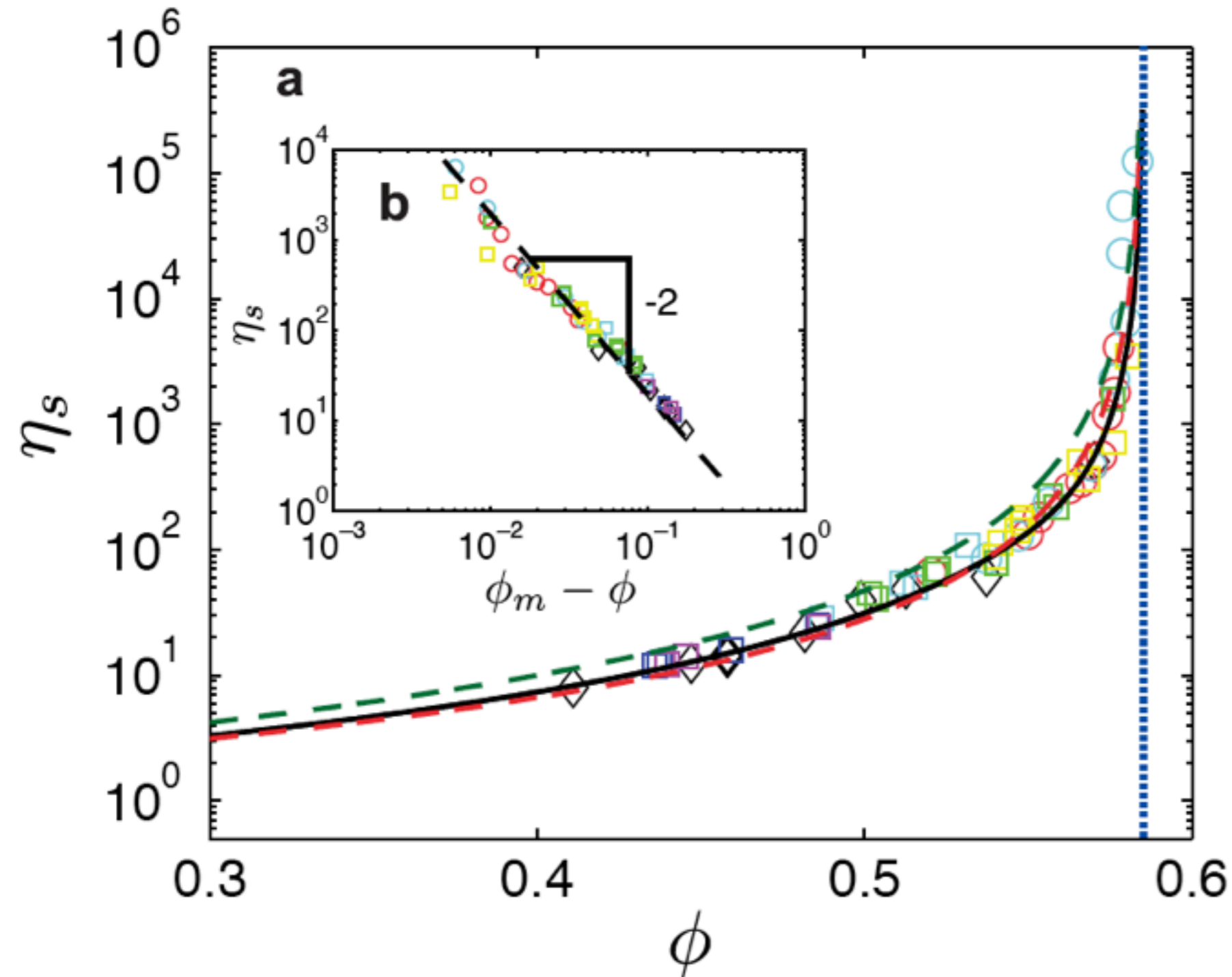
Particles push on the wall



the wall pulls on water



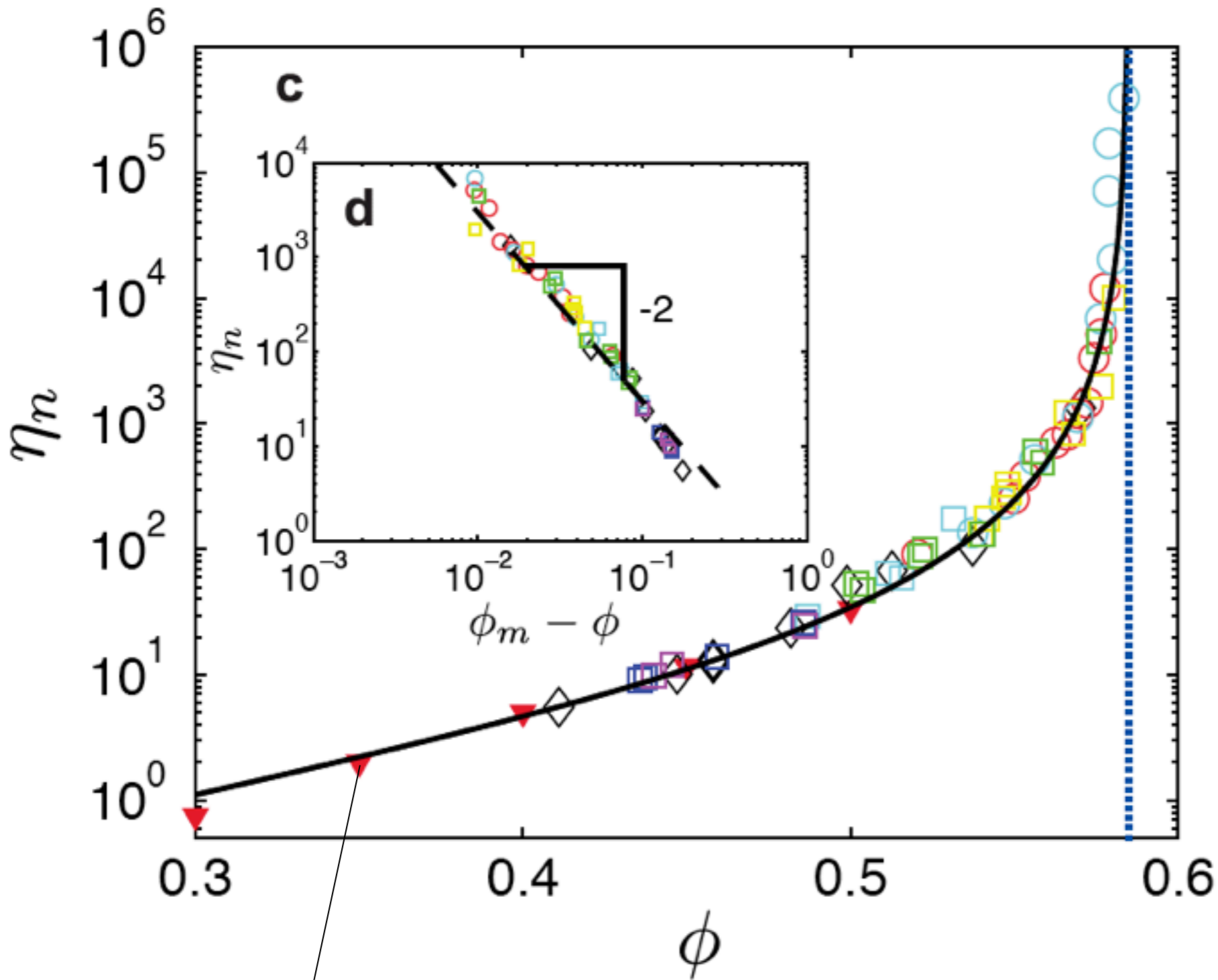
# Link with dense suspensions??



$$\tau = \eta_s(\phi)\eta_f\dot{\gamma}$$

$$\eta_s(\phi)^{-1} \sim \frac{\mu_1}{\phi_m^2} (\phi_m - \phi)^2$$

# Link with dense suspensions??



$$P_p = \eta_n(\phi)\eta_f\dot{\gamma}$$

Deboeuf et al PRL 2010

the pressure imposed rheology is interesting :

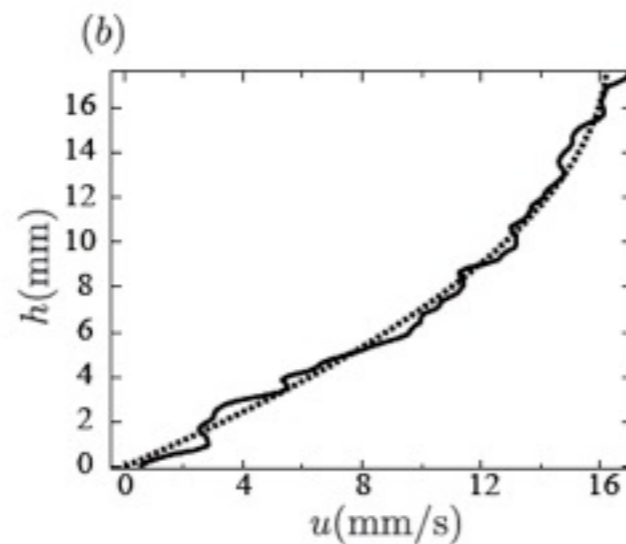
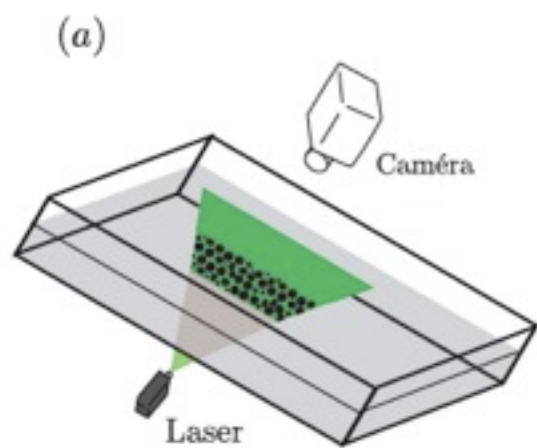
1) From a rheological point of view:

easy to be close to the maximum volume fraction :

precise measure of  $\phi_m$  , study of the divergence...

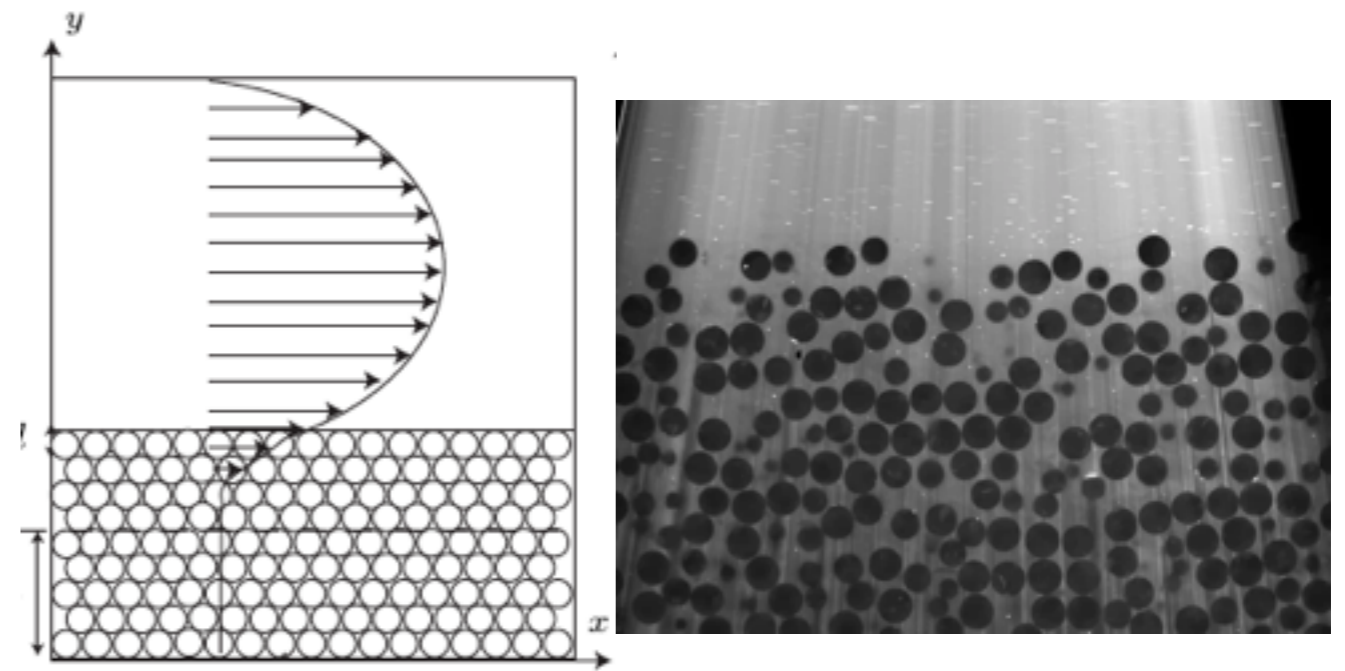
2) the natural description for some configurations:

-submarine avalanches



-sediment transport

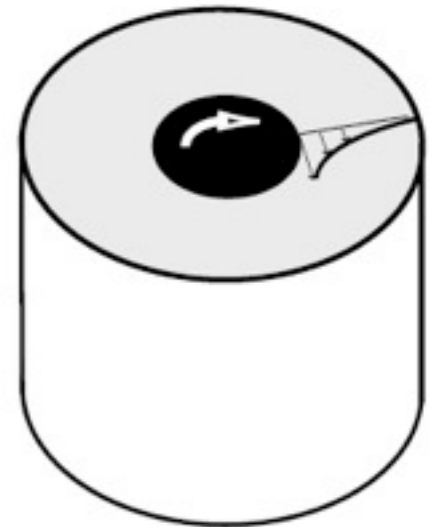
Ouriemi, Aussillous, Guazzelli et al 08,09  
Pailha et al, 12



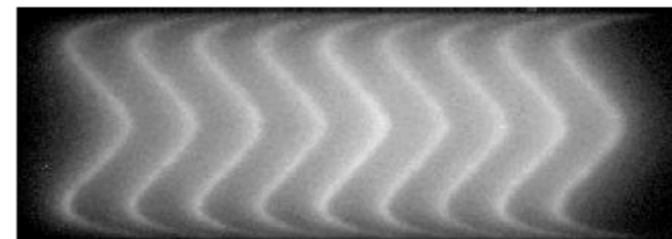
For both dry and immersed granular media, a visco-plastic description is relevant and captures the first order of the viscous nature of the flows.

Beyond  $\mu(I)$  ...

1) Quasistatic flows (shear band, finite size effects)  
A need for non local approach...

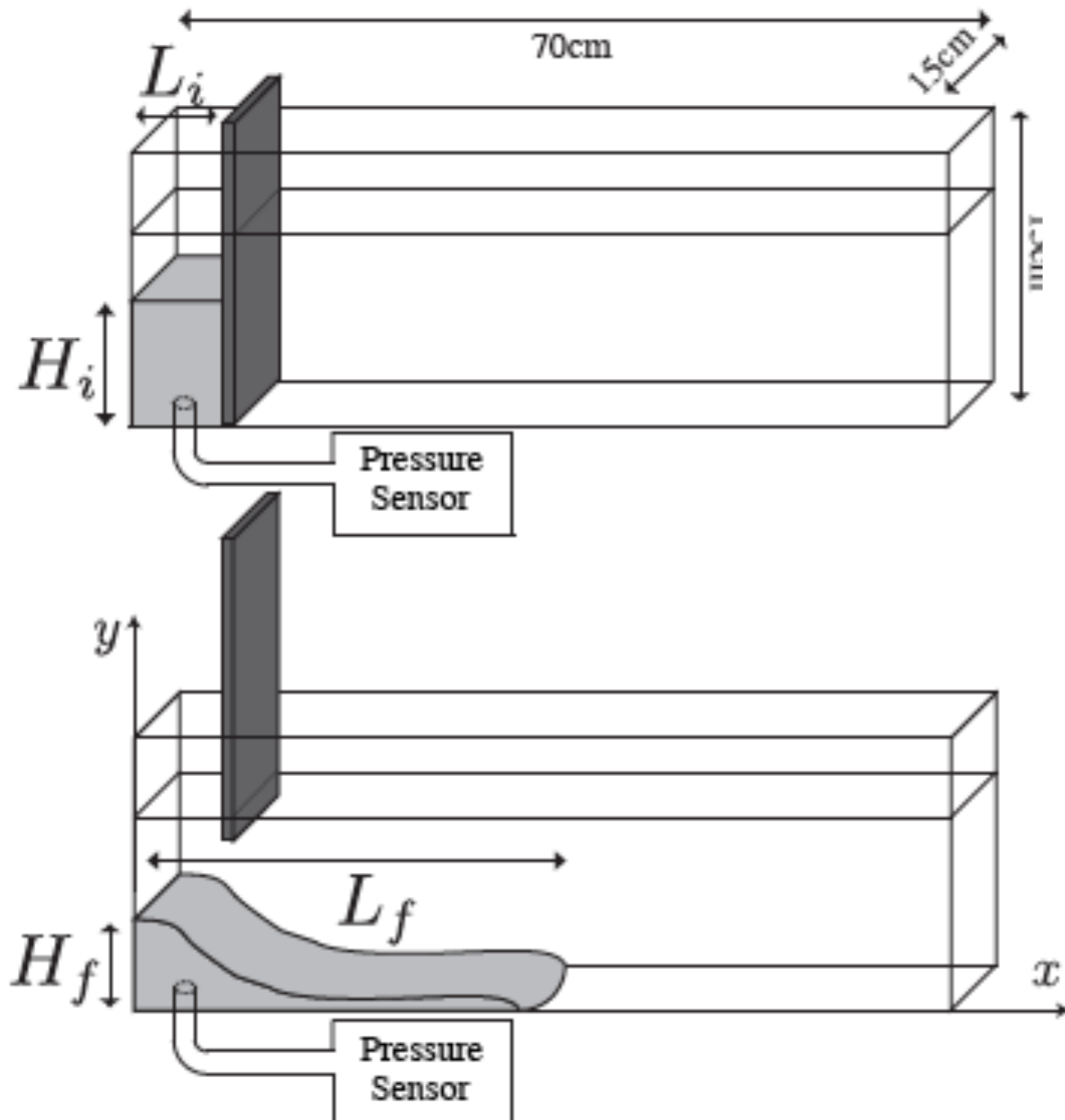


2) Normal stress differences



3) Transient flows when preparation plays a crucial role:

# Role of the preparation on submarine avalanches: Dam break problem (L. Rondon Phd).



- glass beads,  $d = 225 \mu m$  in mean diameter
- Liquid : mix of Ucon-oil and water, 23 times more viscous than water
- Initial aspect ratio  $\mathcal{A} = \frac{H_i}{L_i}$
- $0.55 < \phi_i < 0.62$

Loose  $\phi_i = 0,55$

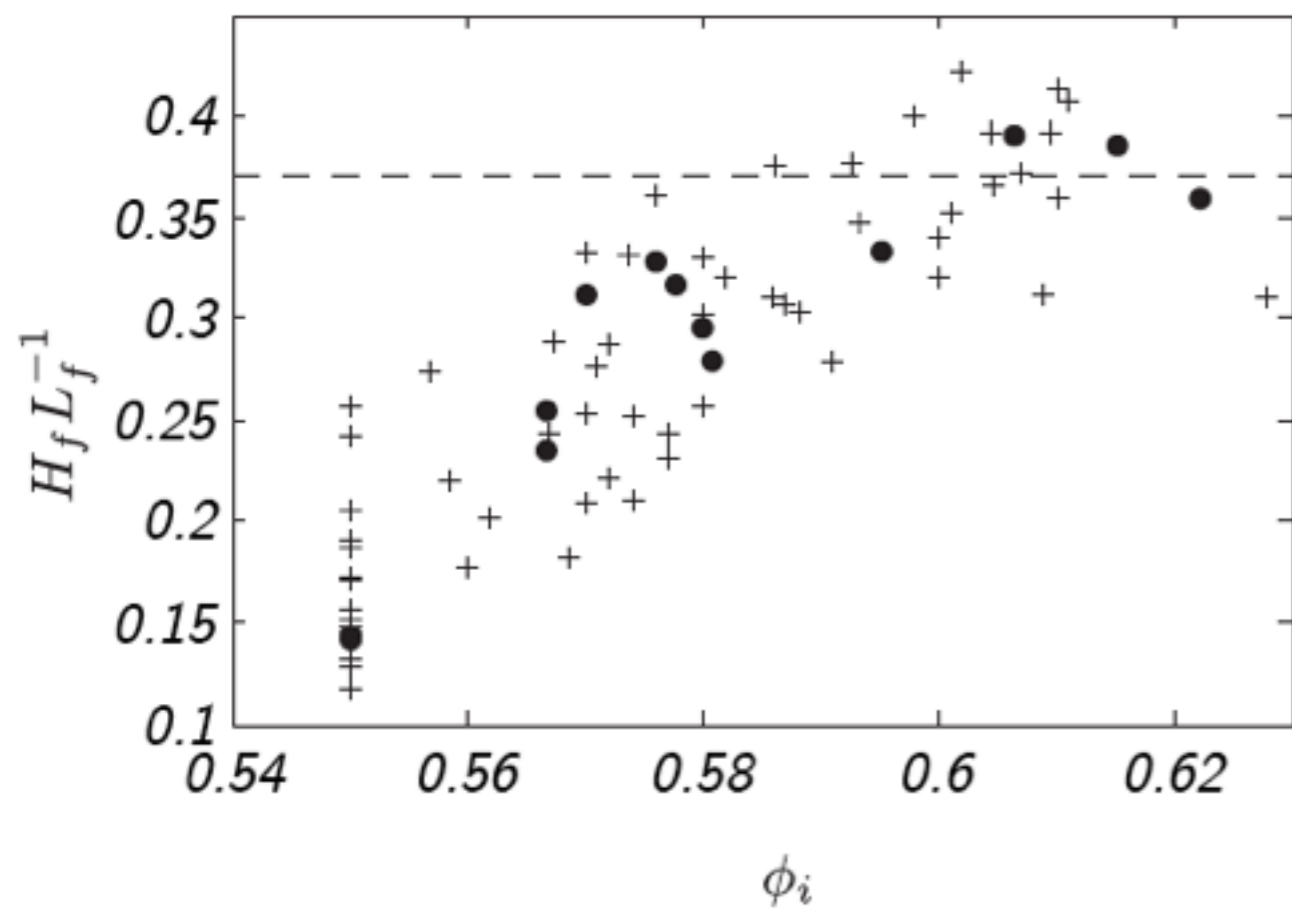
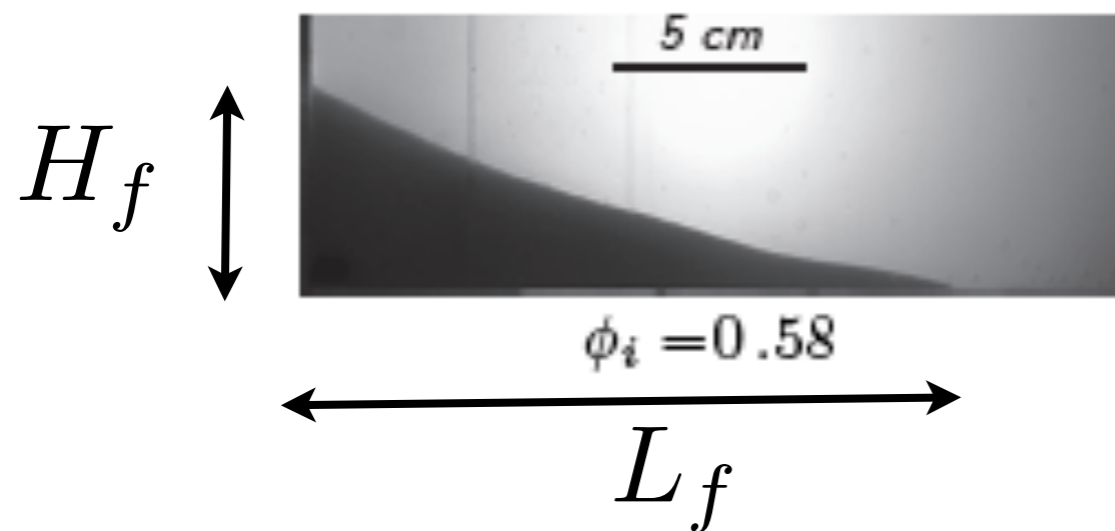


dense  $\phi_i = 0,61$

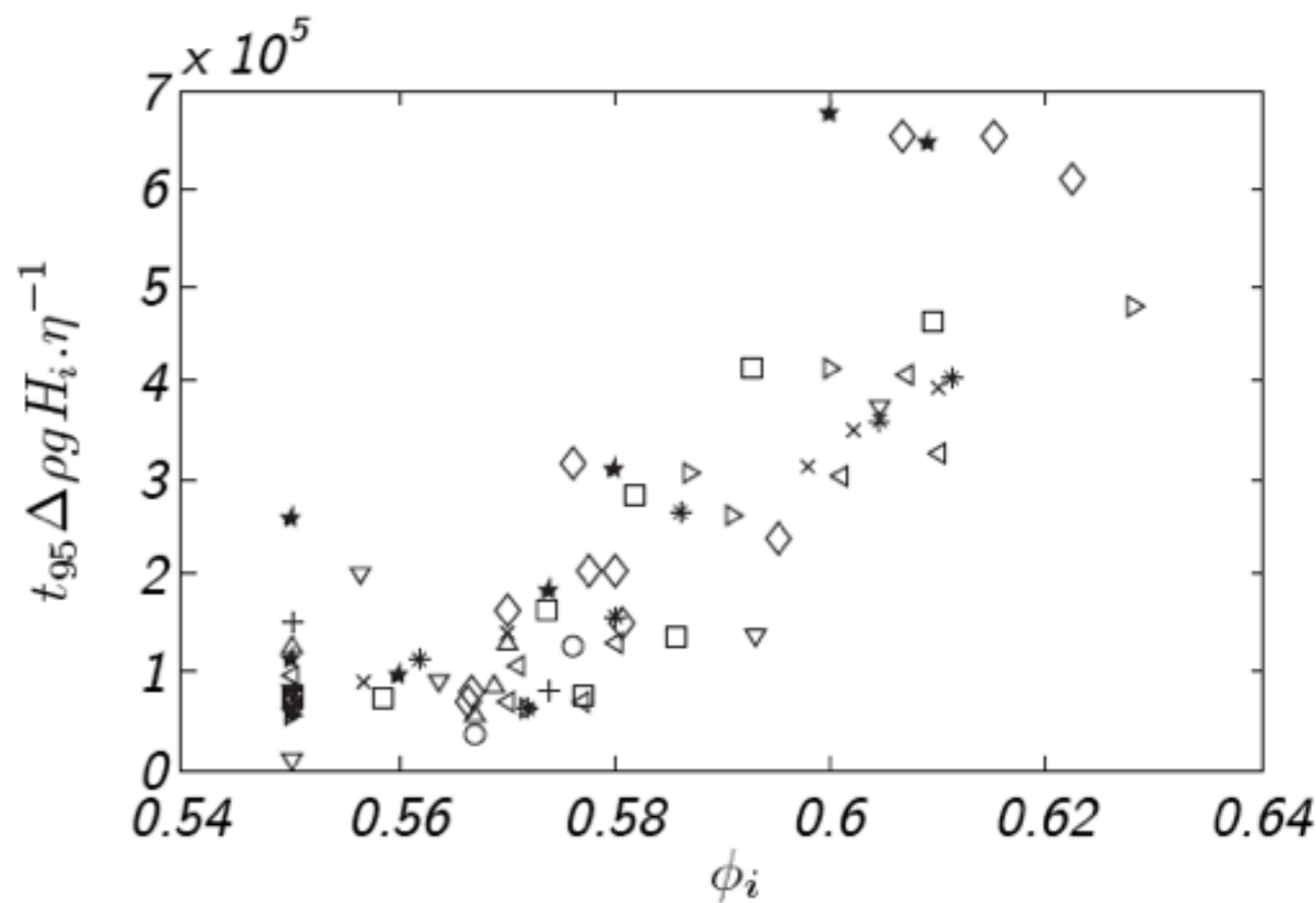




Slope



time

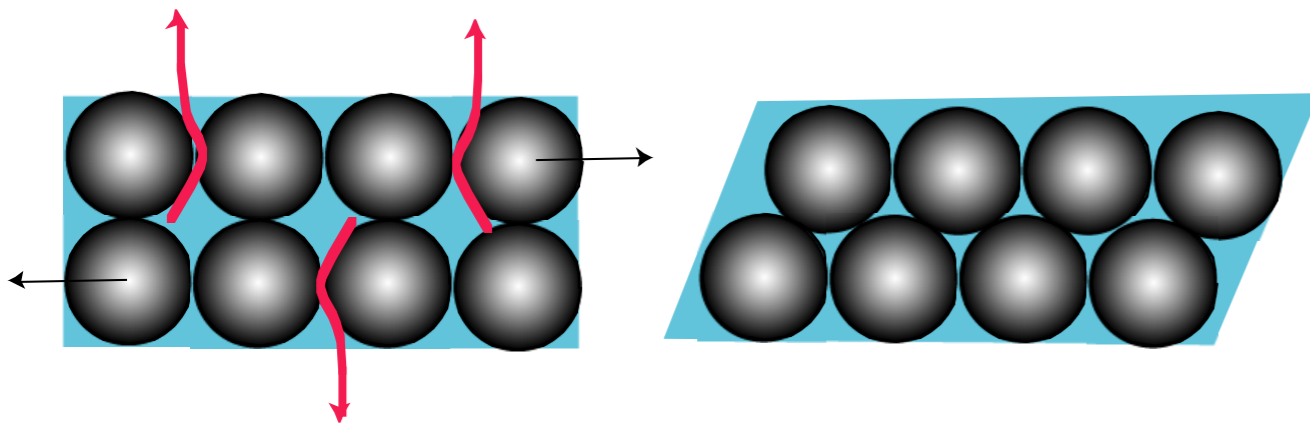


Unsteady flows => variation of  $\phi$   
=> relative motion between grains and fluids  
=> additional stress...

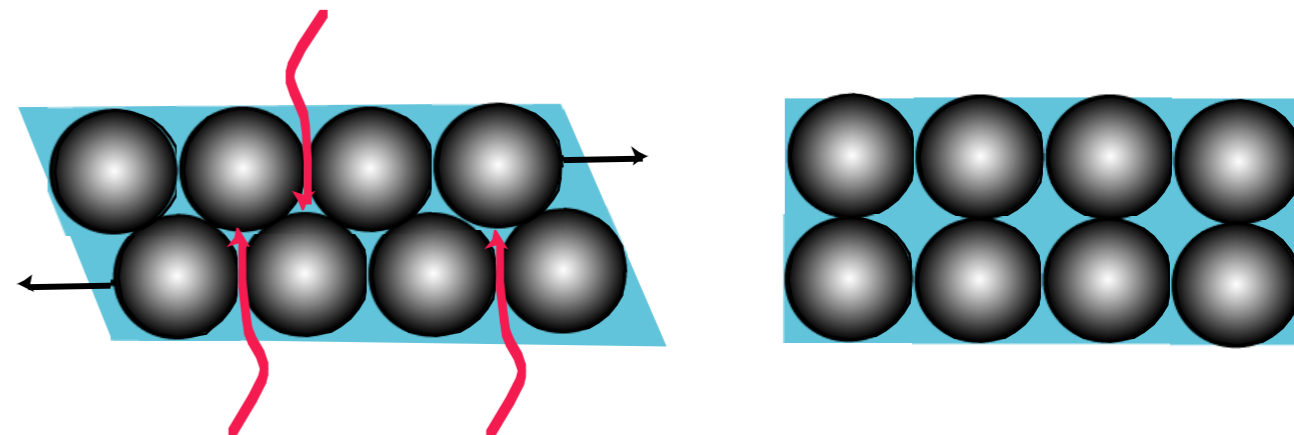
## Pore Pressure feedback argument

(Iverson Rev. Geo. 97, JGR 05)

Loose case



Dense case



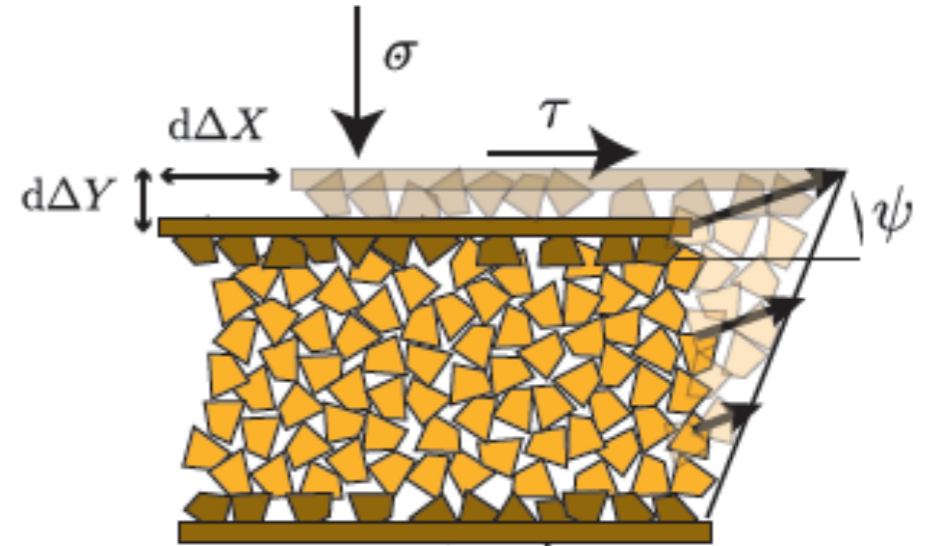
# A shear rate dependent critical state theory

Dilatancy angle:

$$d\phi/dt = \dot{\gamma} \tan \psi$$

$$\tan \psi = K(\phi - \phi(I_v))$$

$$\tau_b = (\mu(I_v) + \tan \psi) P^p$$



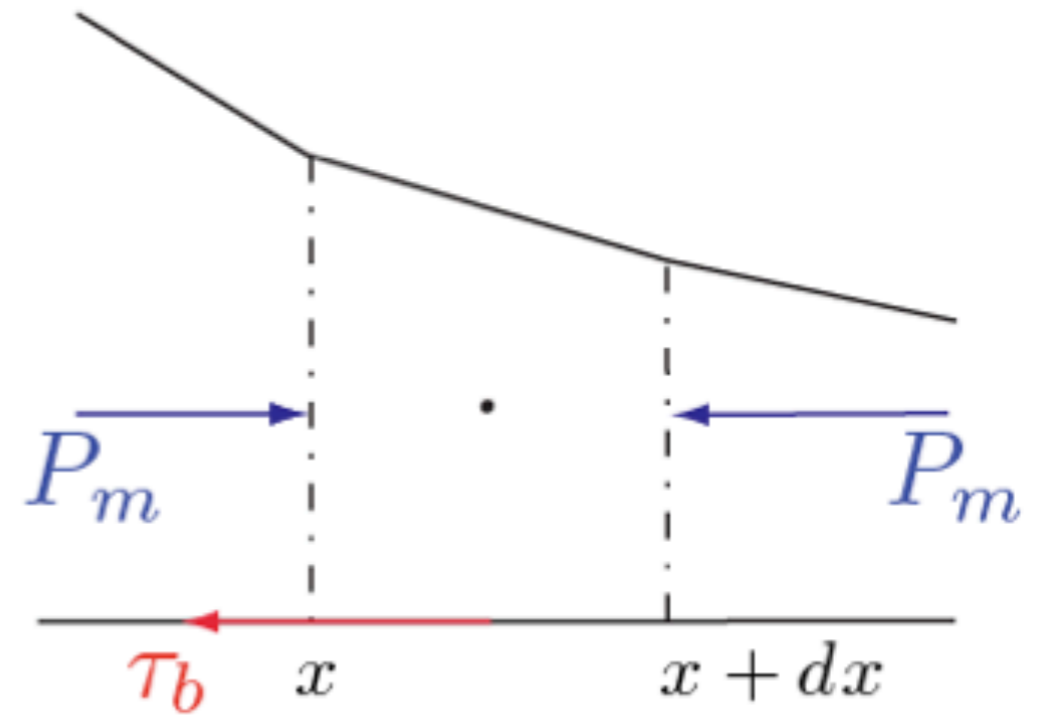
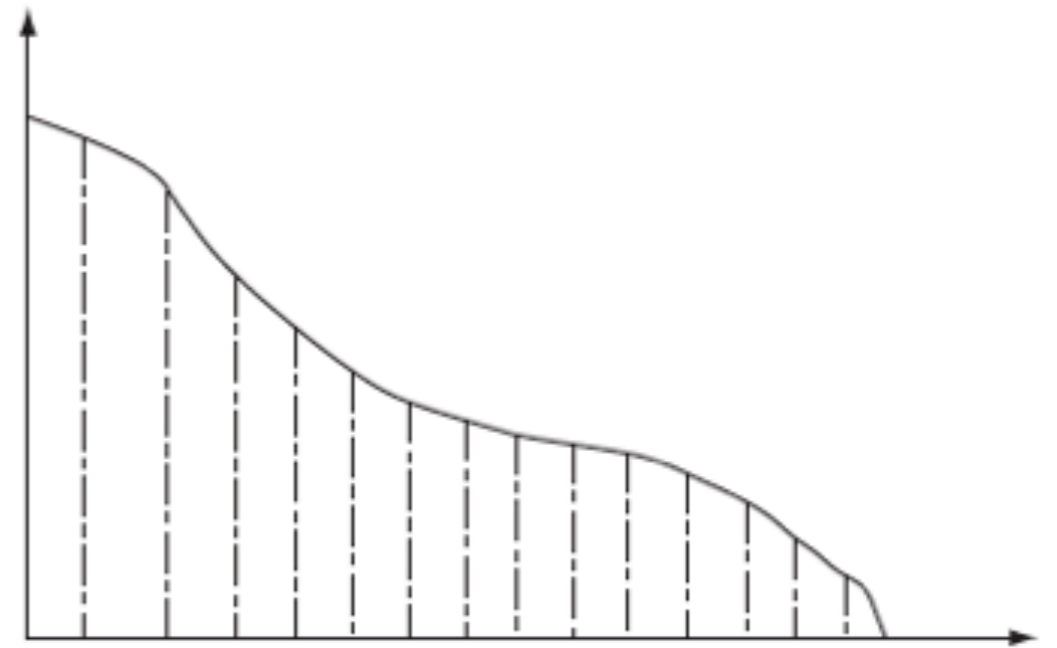
The granular pressure:

$$P^p = \rho g h \cos \theta + \alpha \frac{\eta h}{d^2} u \tan \psi$$

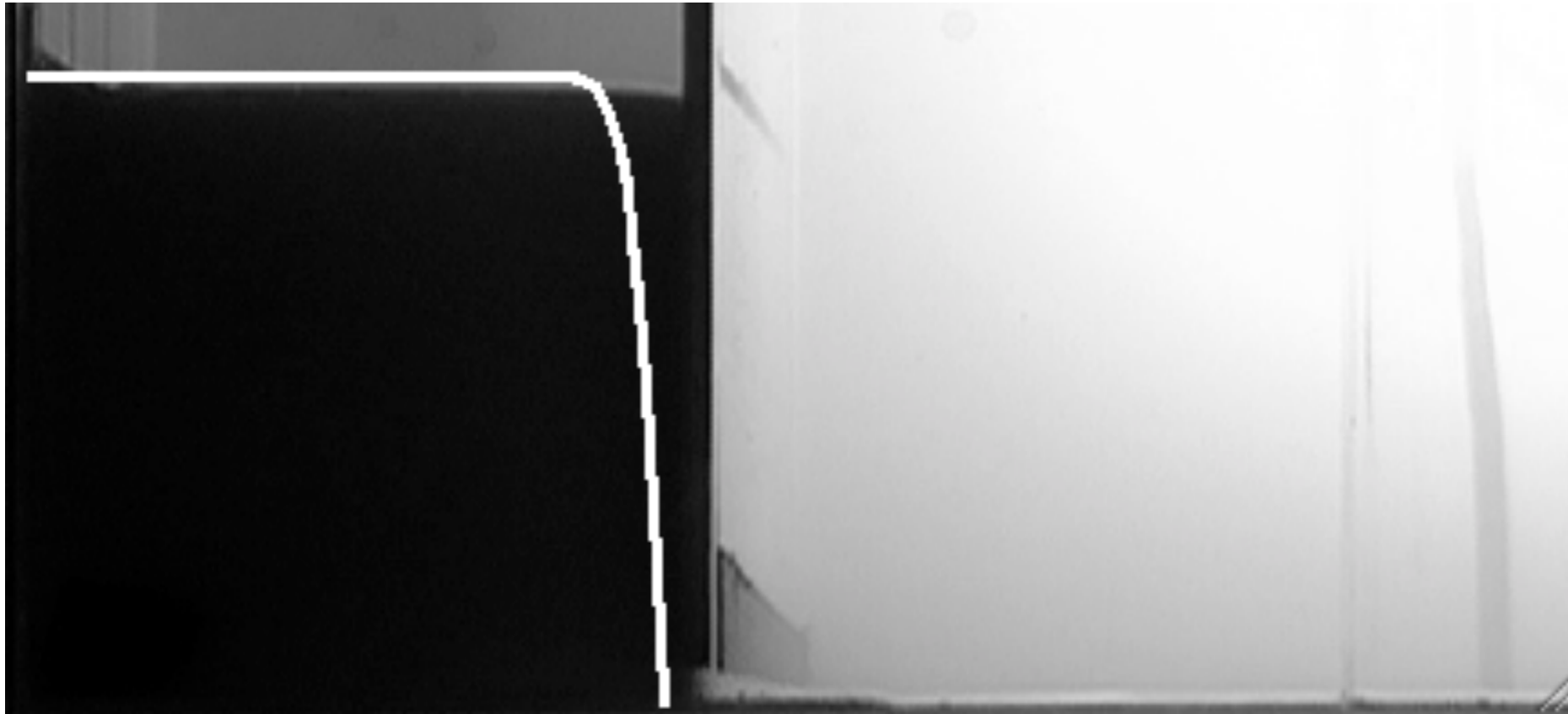
the steady  
granular rheology  
+  
dilatancy  
effects

in a depth averaged model

=> semi-quantitative predictions!



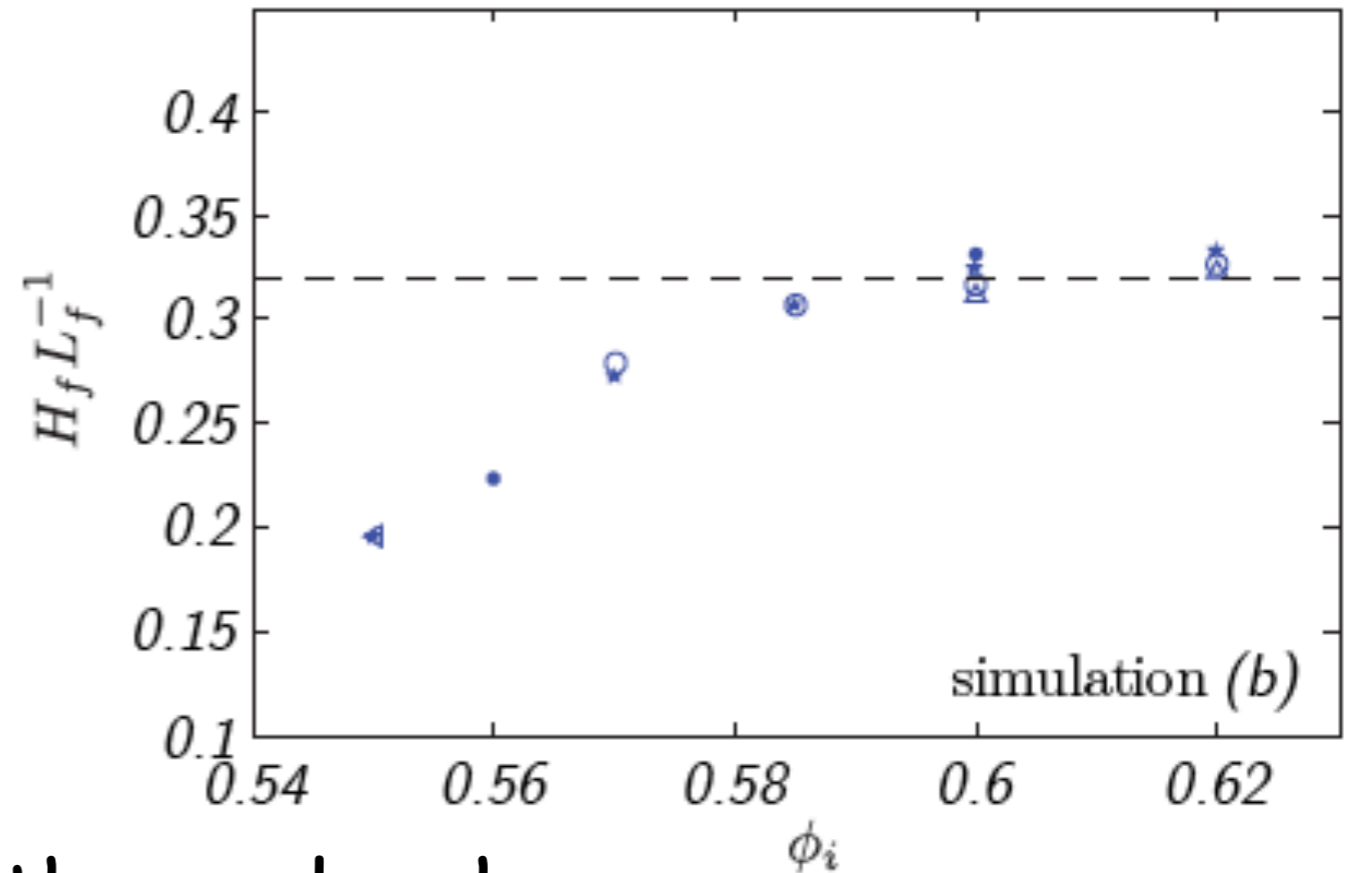
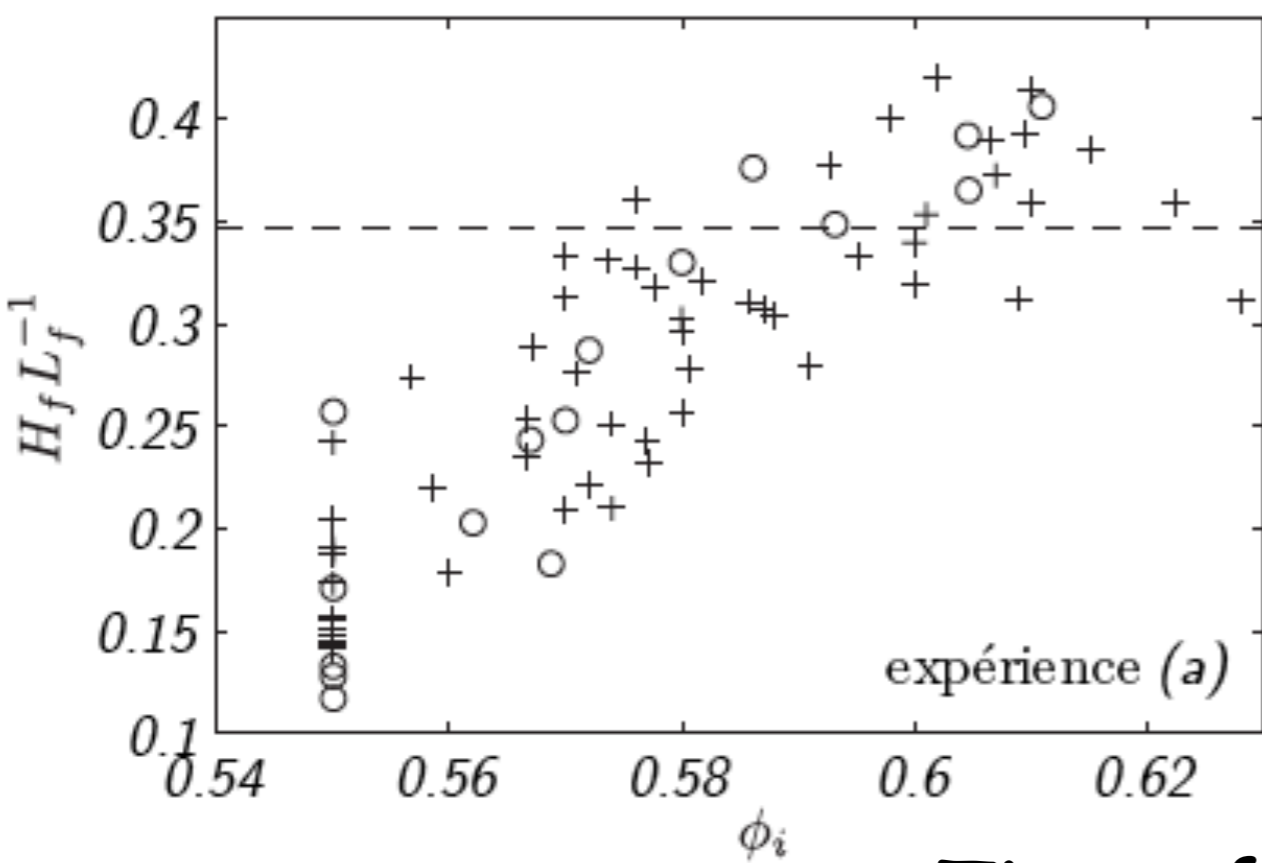
dense initial state



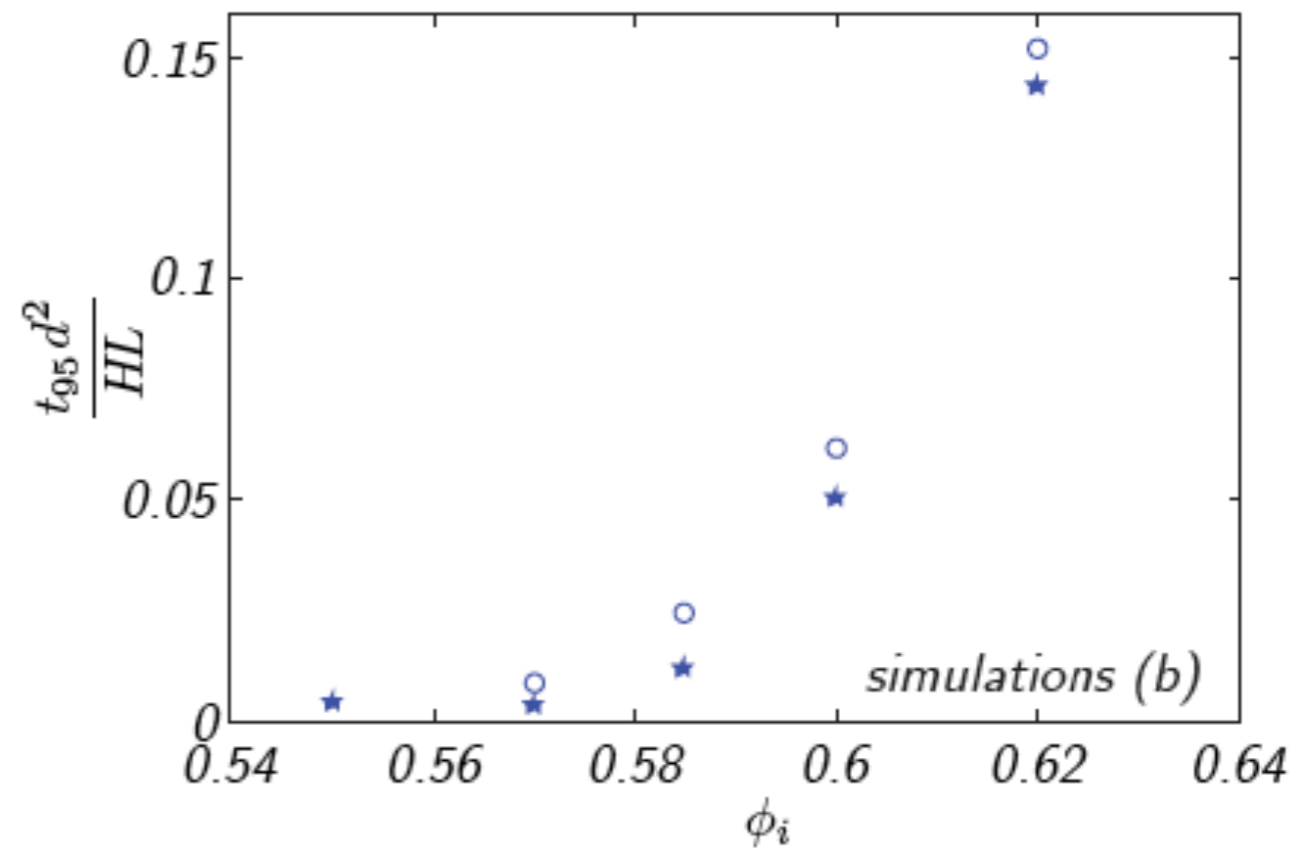
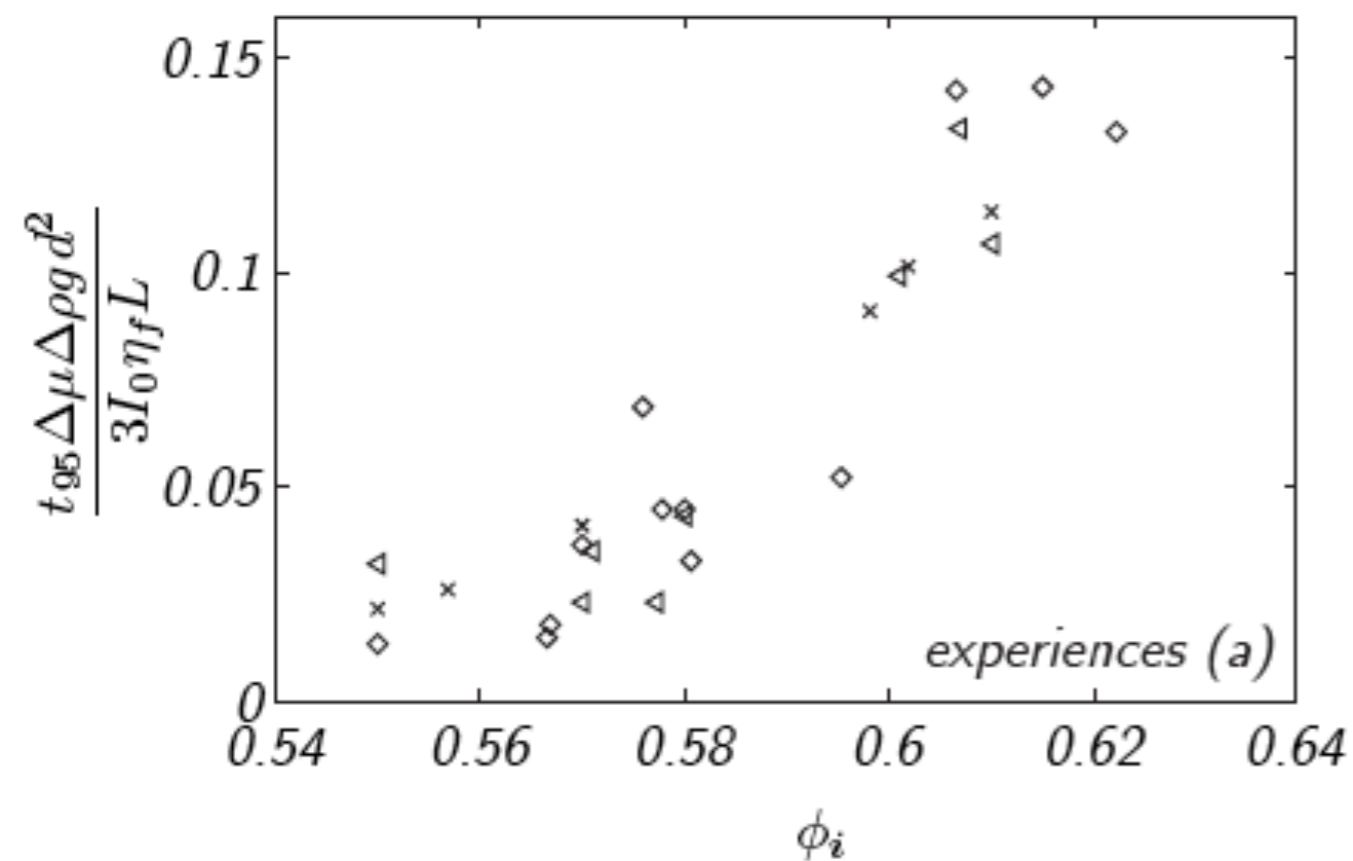
Loose initial state



# final slope of the deposit



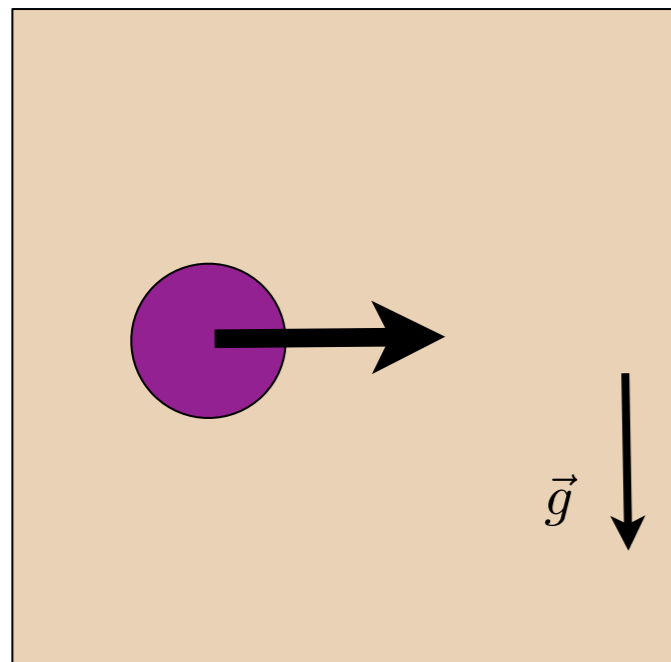
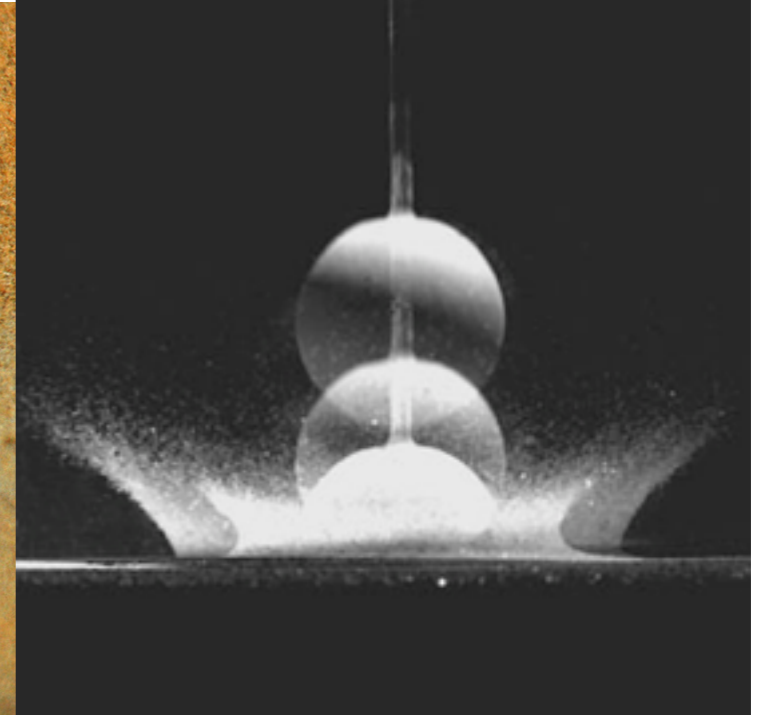
# Time of the avalanche



# Conclusions about the rheology...

- pressure imposed versus volume imposed...
- A precise measure of rheology for very dense suspensions.
- the frictional description is relevant to describe flows under gravity (avalanches, sediment transport...) ...

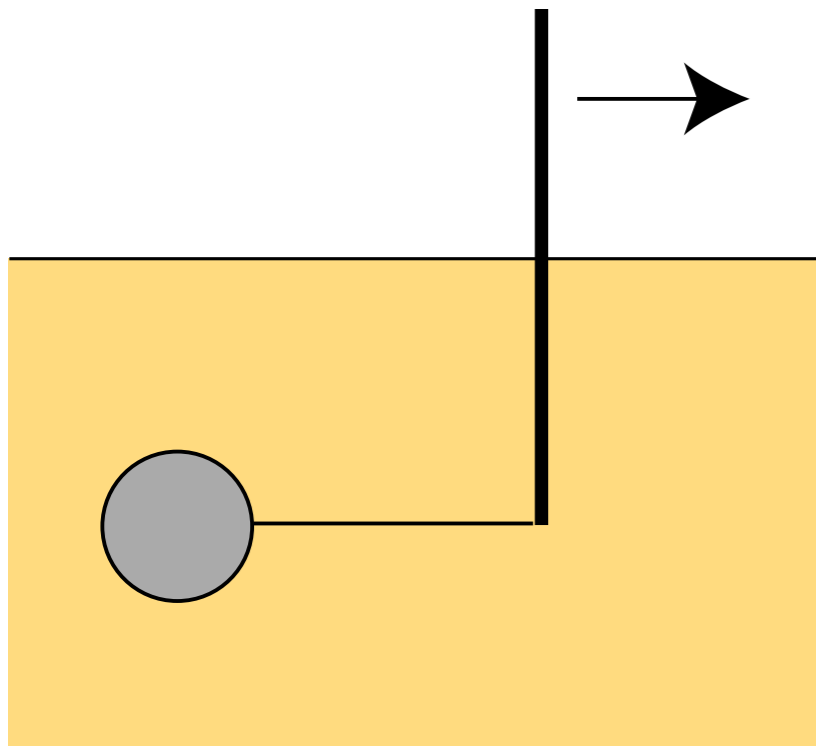
# A classical hydrodynamic problem: Forces on a moving object



drag has been studied in details,...  
what about lift ?

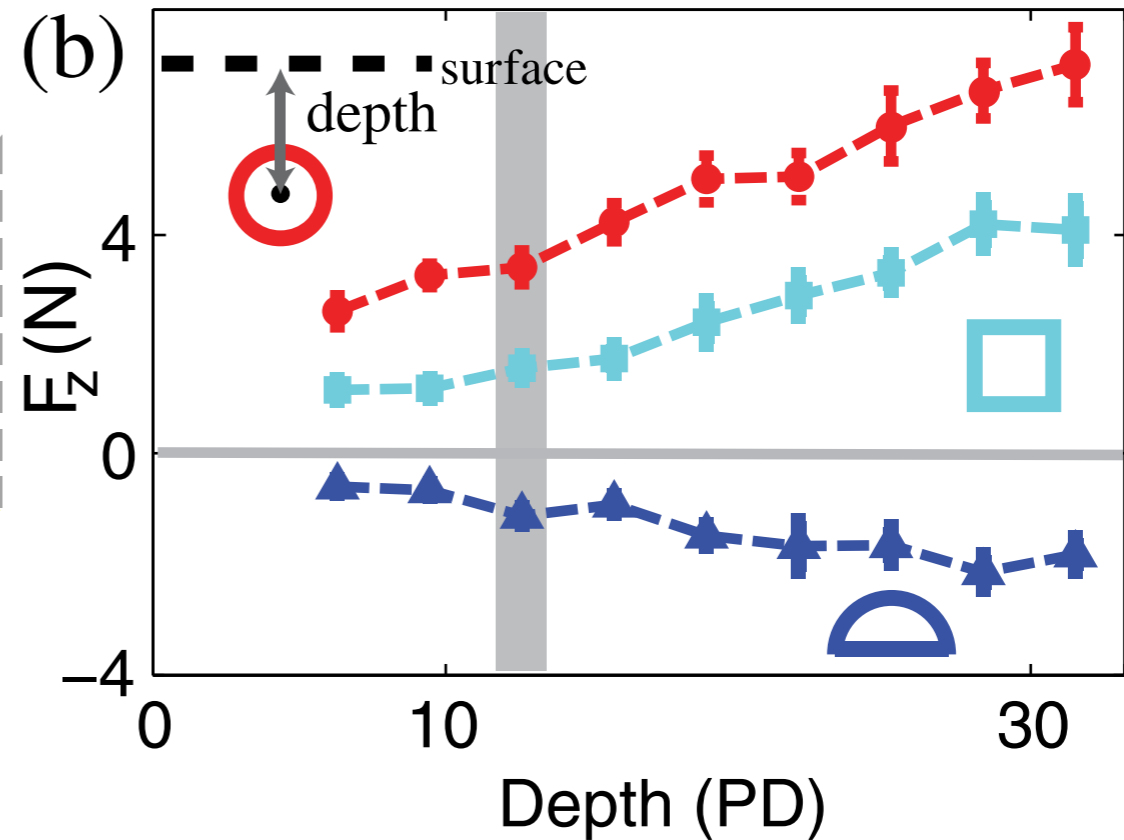
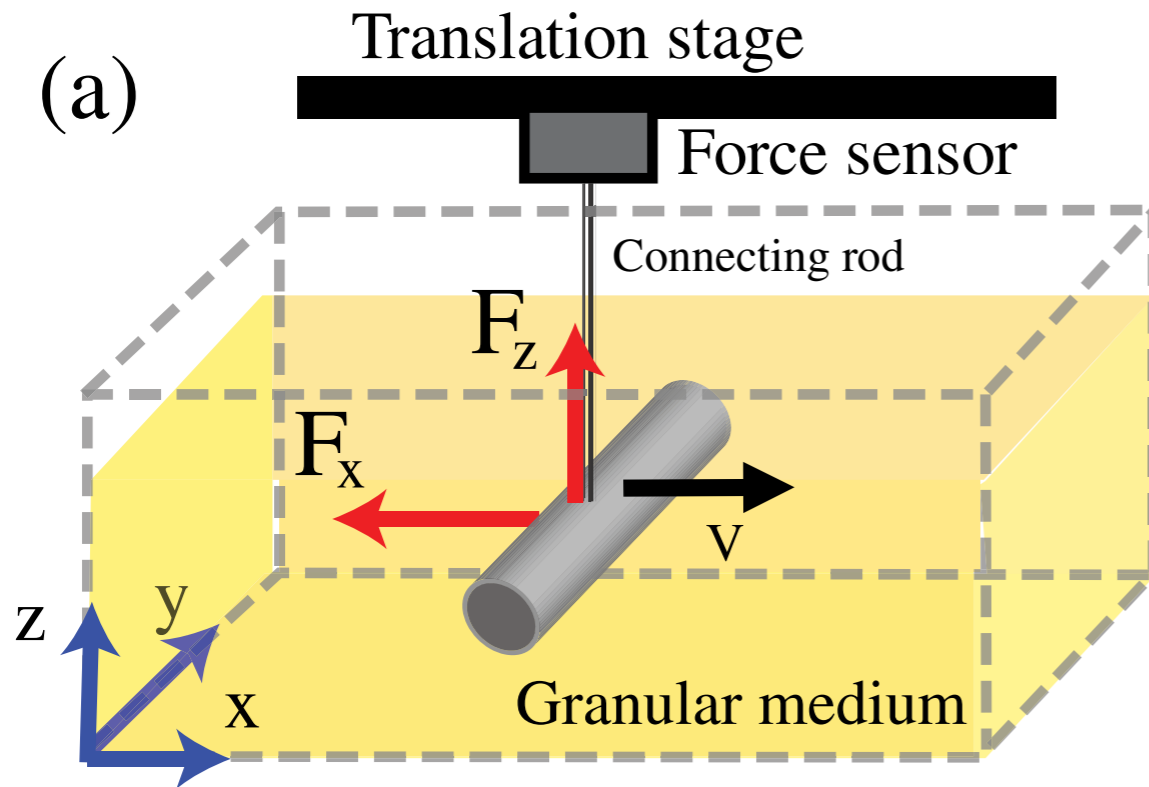


# Evidence of lift force on a symmetric object



# a recent study on lift forces...

Ding et al. Phys. Rev. Lett. 2011

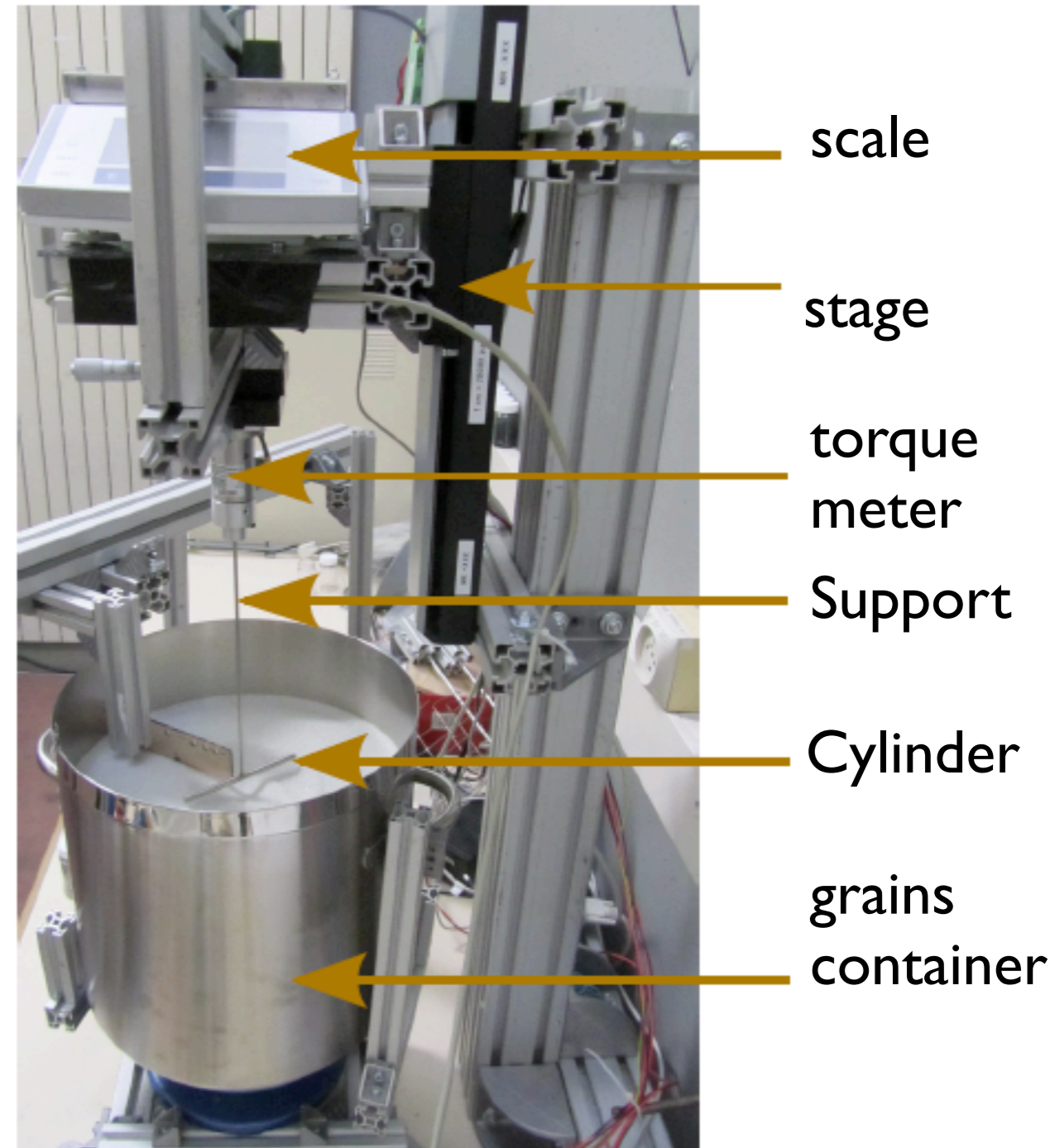
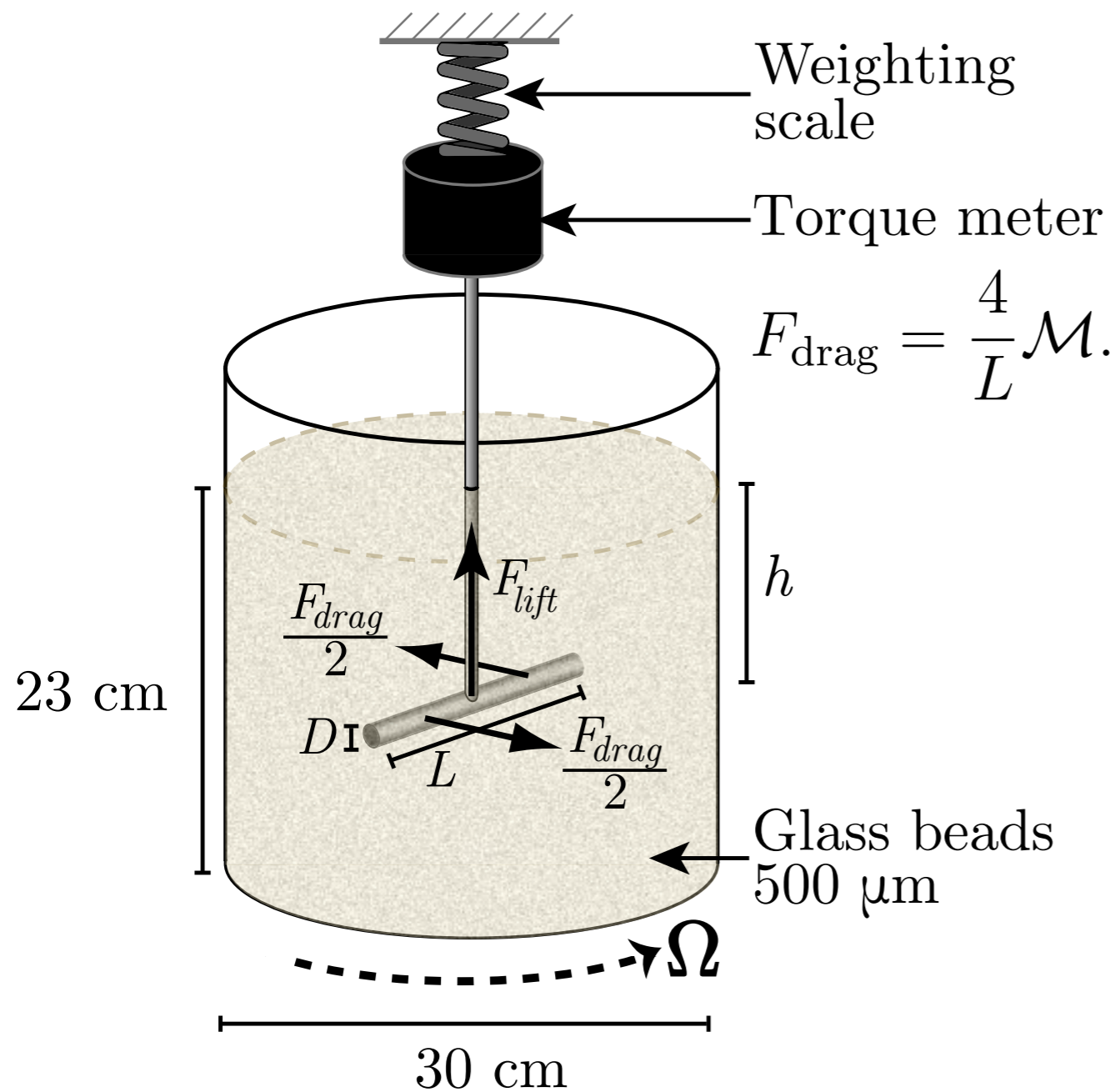


Lift on a symmetrical object which scales with depth...

physical origin, scaling ??

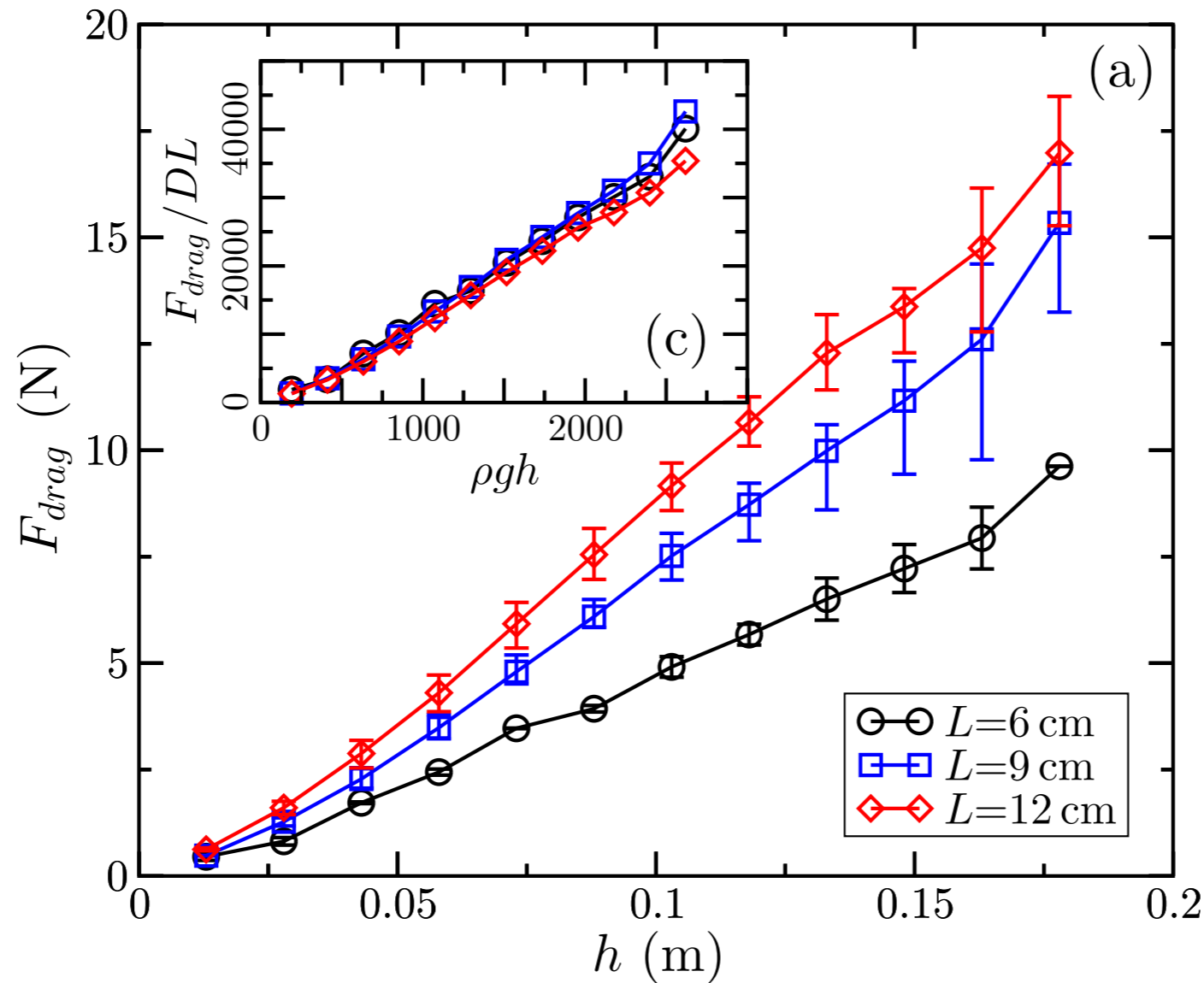
# Experimental set-up: drag and lift on a rotating cylinder

Guillard et al. PRL 2013



low velocity quasi-static regime : forces independent of velocity  
measurements first half-turn...

# Drag forces

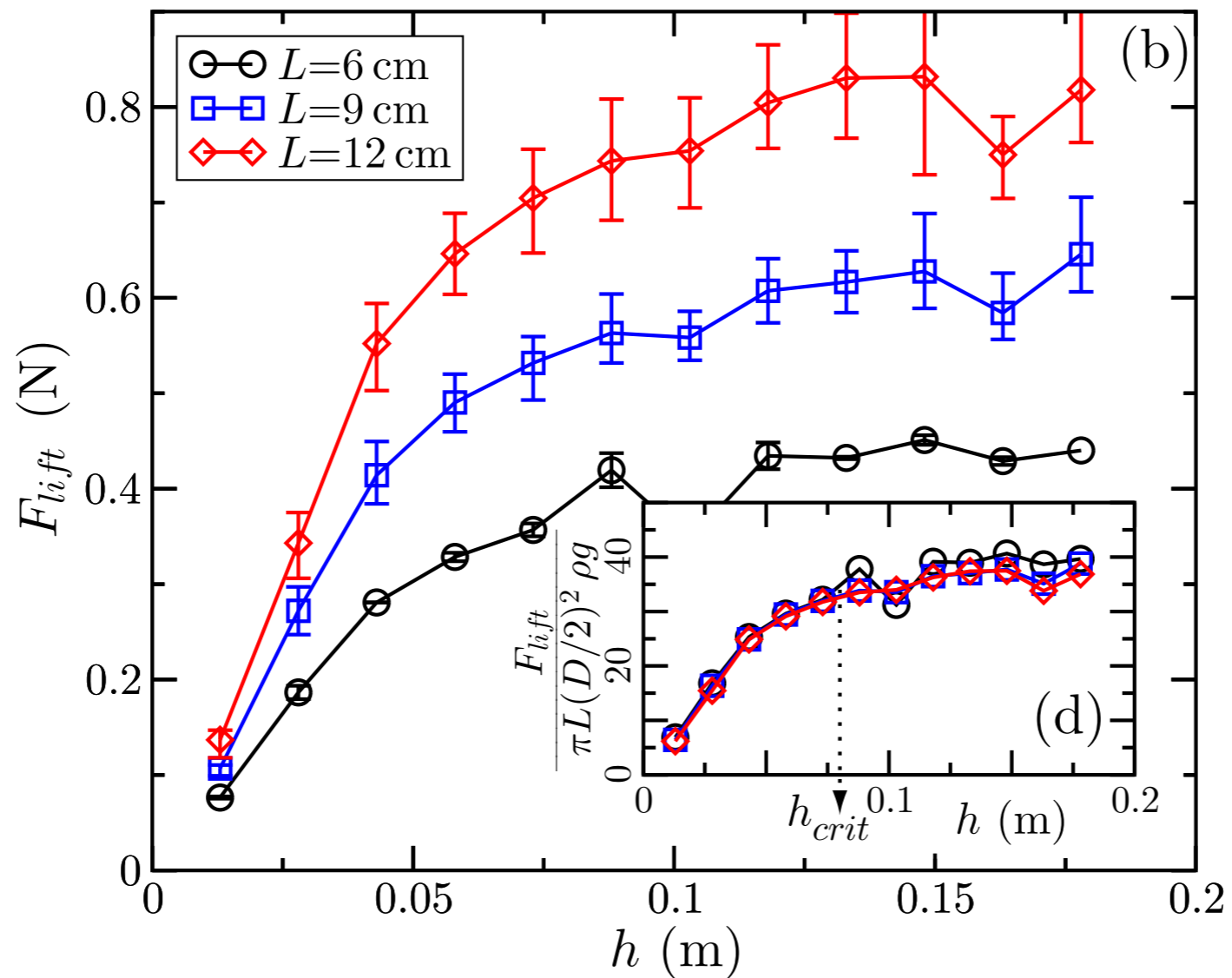


$D = 4 \text{ mm}$

Drag forces proportional to depth (pressure)

$$F_{\text{drag}} = C_d \rho gh DL \propto PS \quad C_d \simeq 13$$

# Lift forces

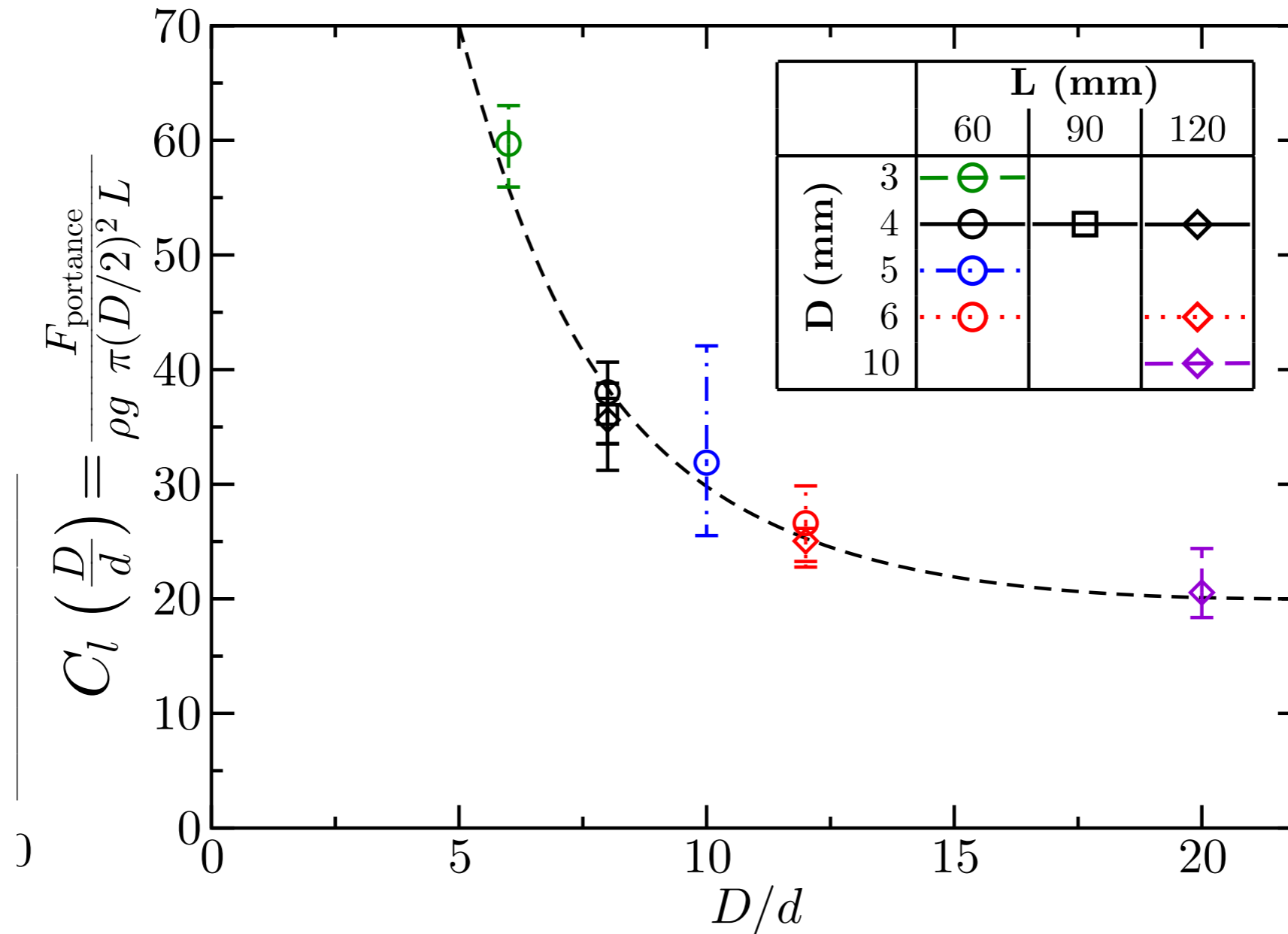


$D = 4$  mm

Saturation of the lift force with depth

$$F_{lift} \sim 20 F_{Archimede} !!!$$

# Lift force : general scaling

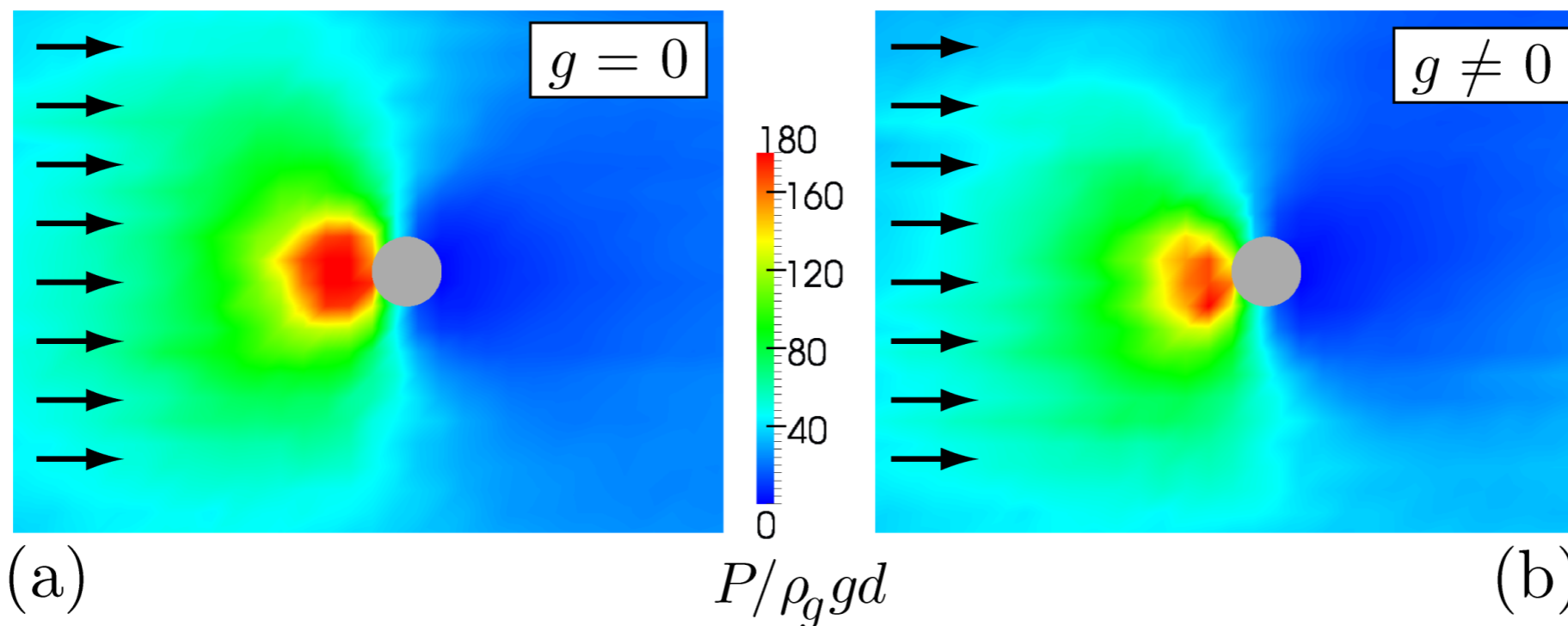
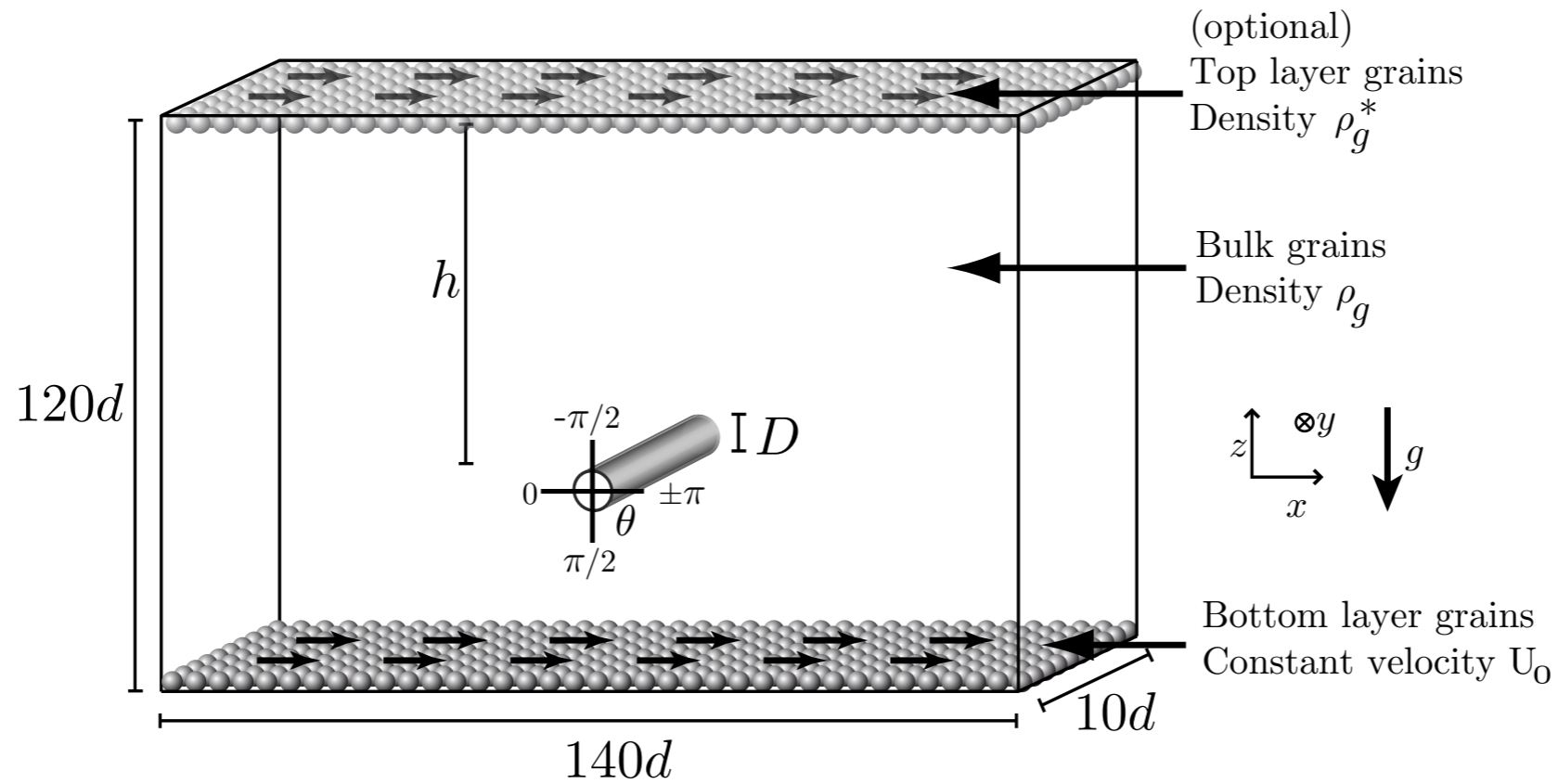


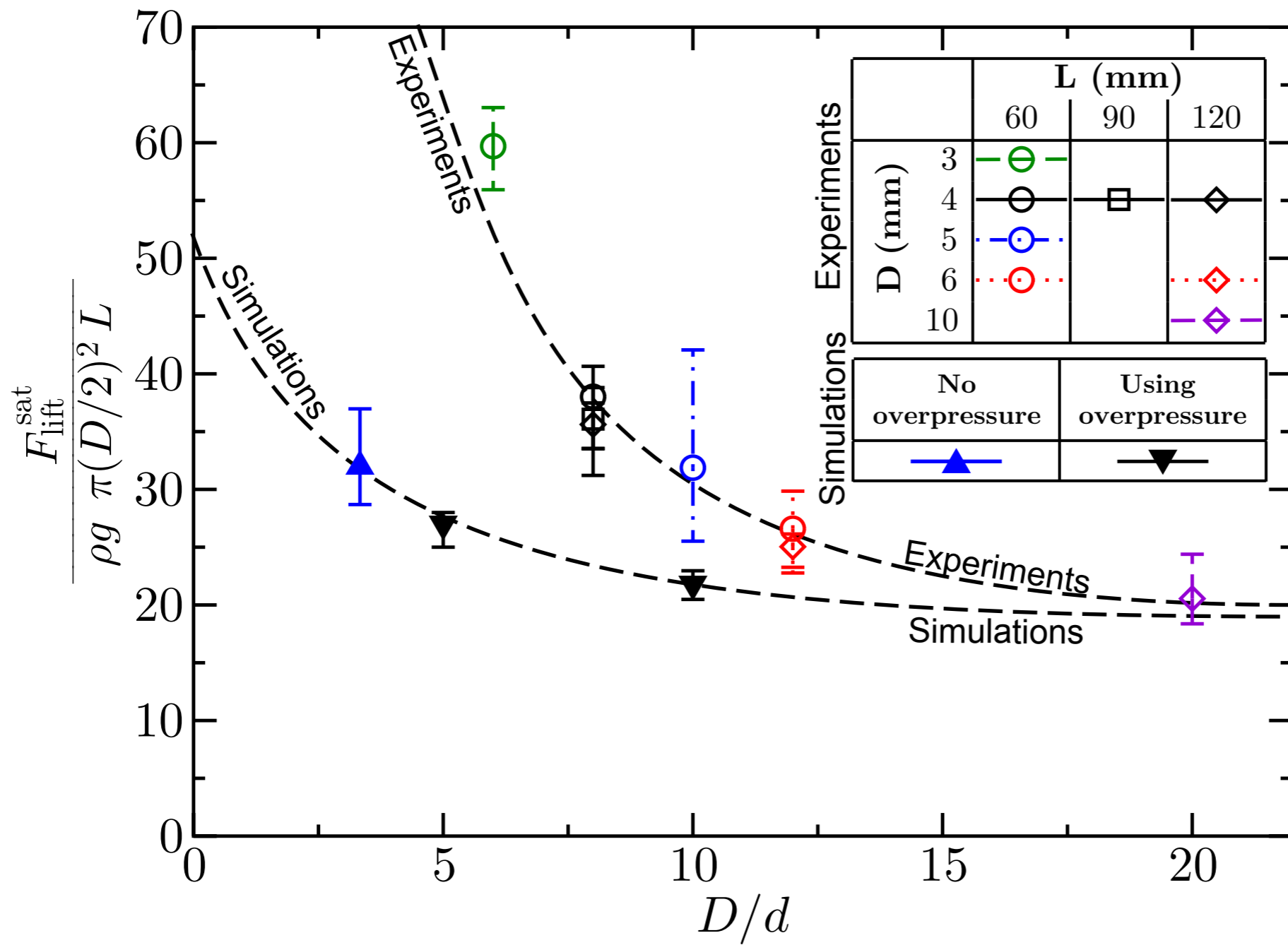
$$F_{\text{lift}} = C_l(D/d) \rho g \pi \frac{D^2}{4} L \propto \nabla P \times V$$

**Archimedes-like scaling but finite size effects**

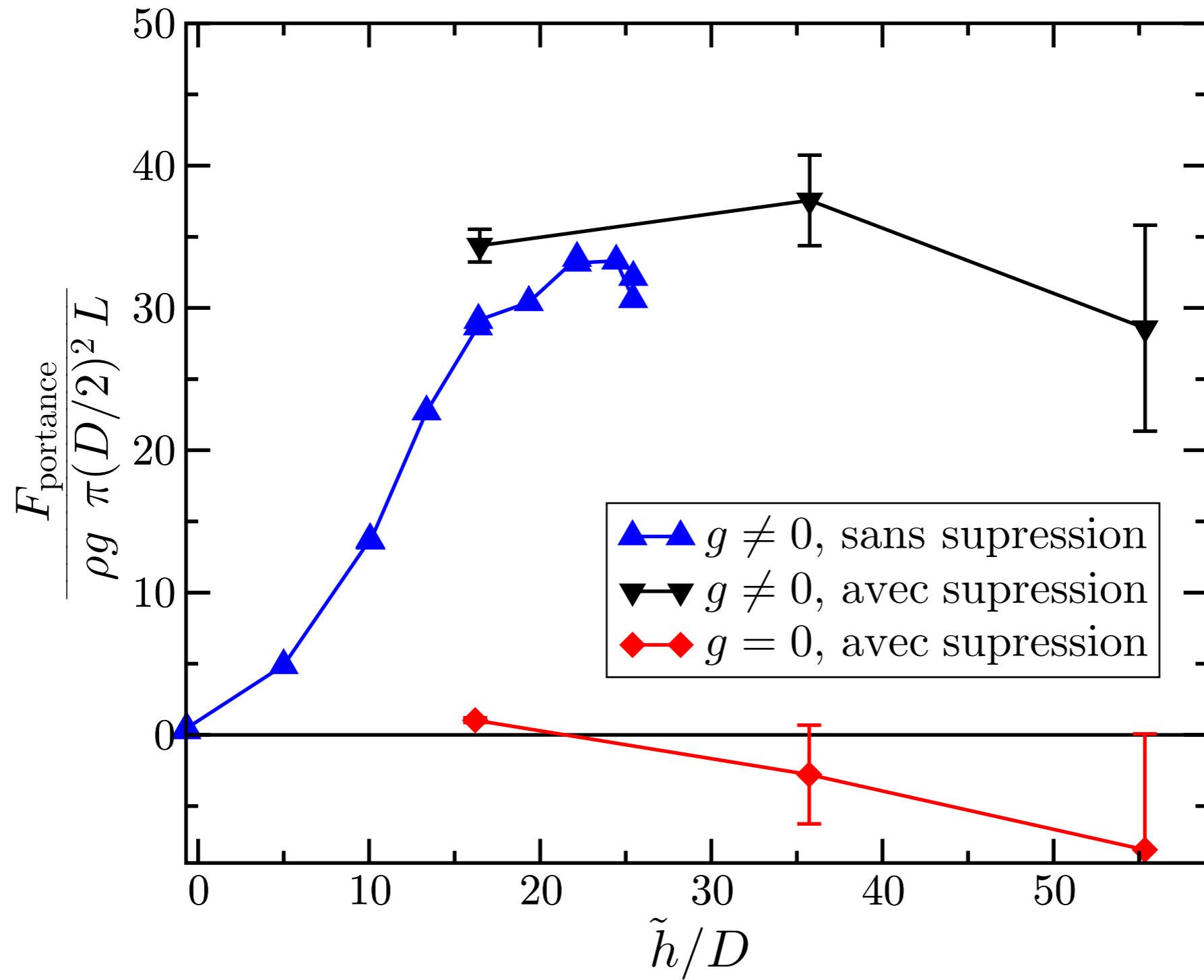
**very strong lift**  $C_l \approx 20$  when  $D/d > 15$

# Molecular dynamics simulations (LIGGGHTS)



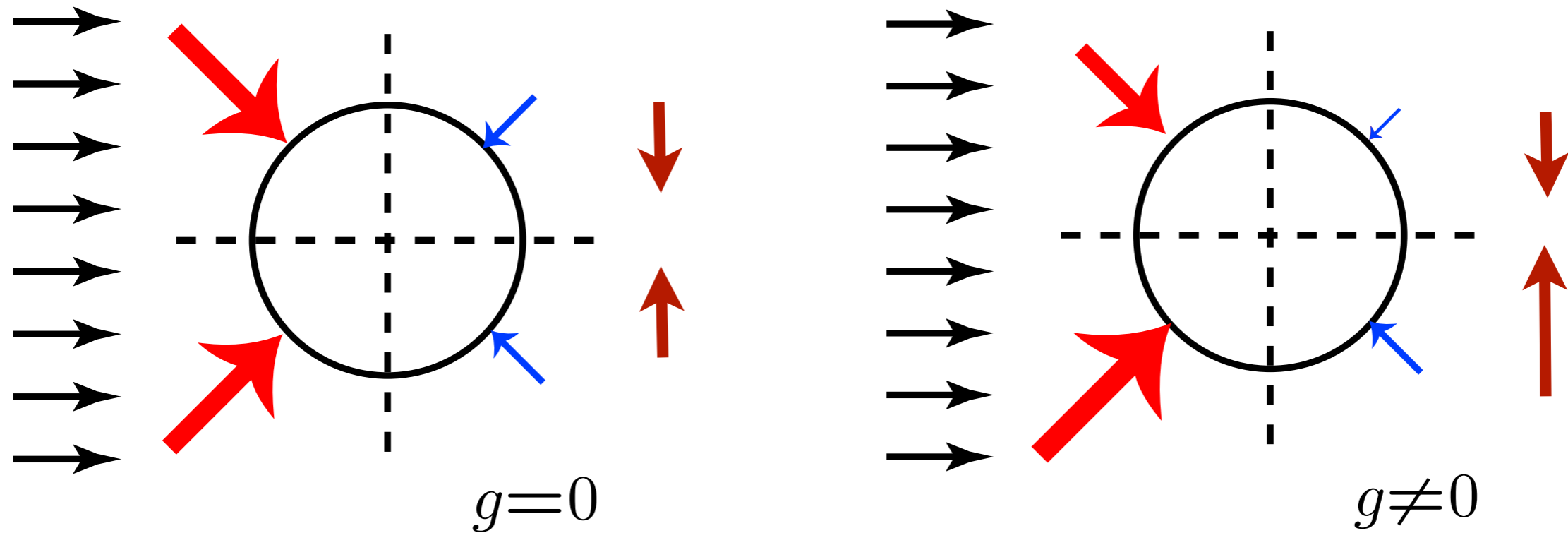






non pressure gradient => no lift

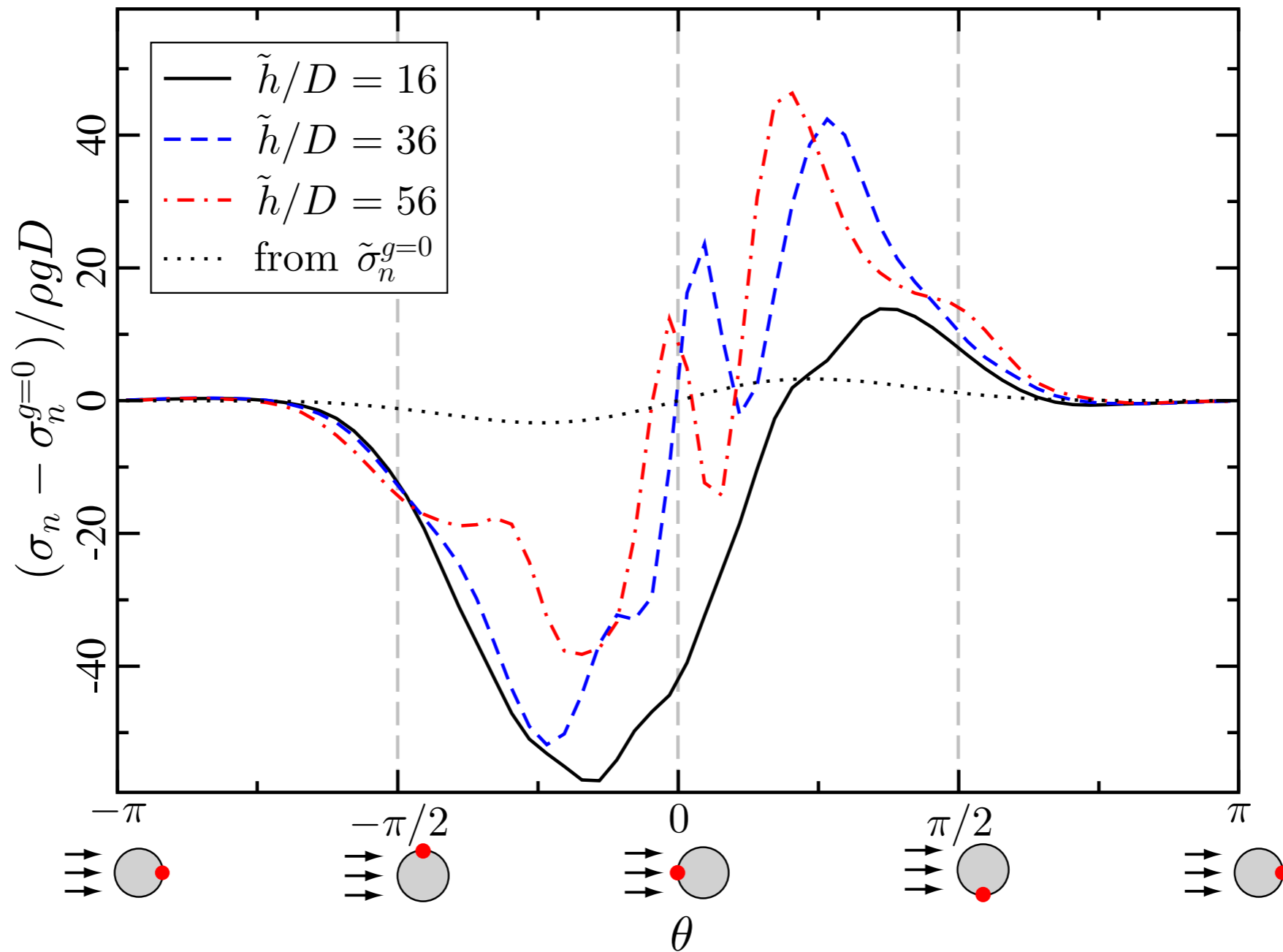
# In granular media : strong left/right asymmetry



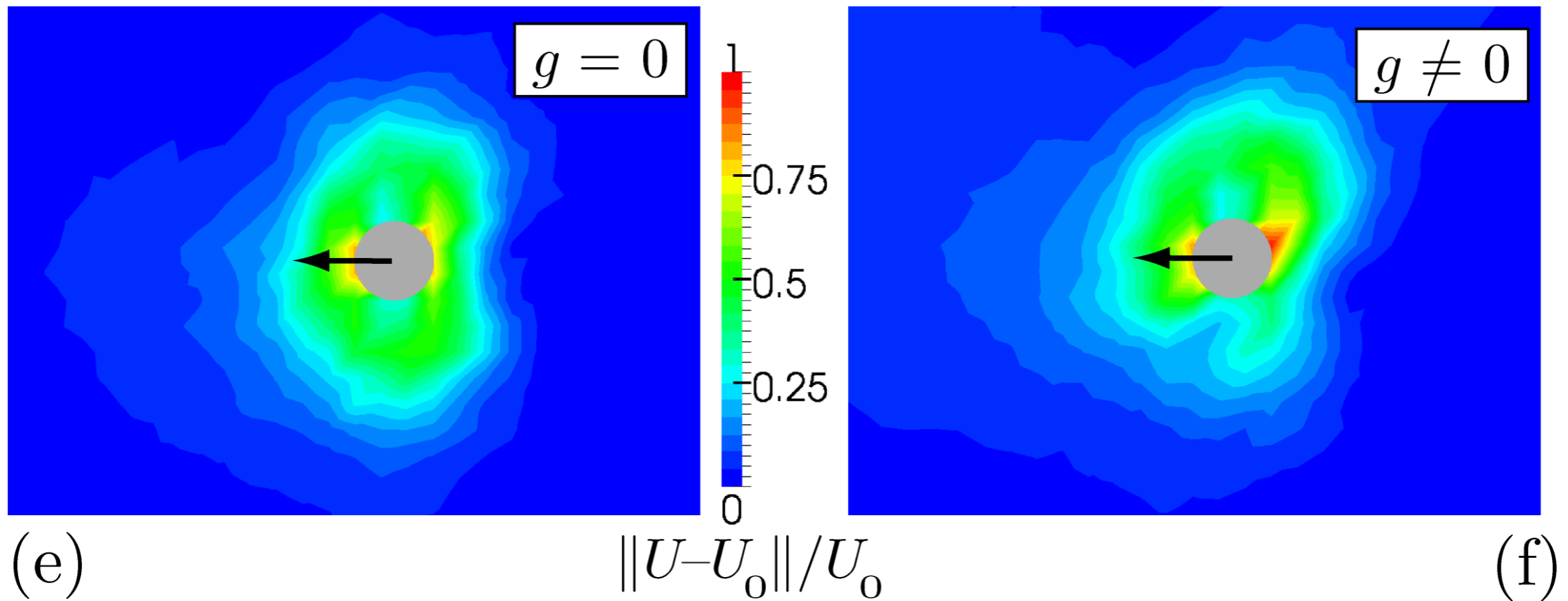
+ stresses proportional to pressure (friction)

→ in the presence of a pressure gradient : lift

$$F_{\text{lift}} = (\rho g L D^2 / 4) \int_{-\pi}^{\pi} \tilde{\sigma}_n^{g=0}(\theta) \sin^2 \theta d\theta$$



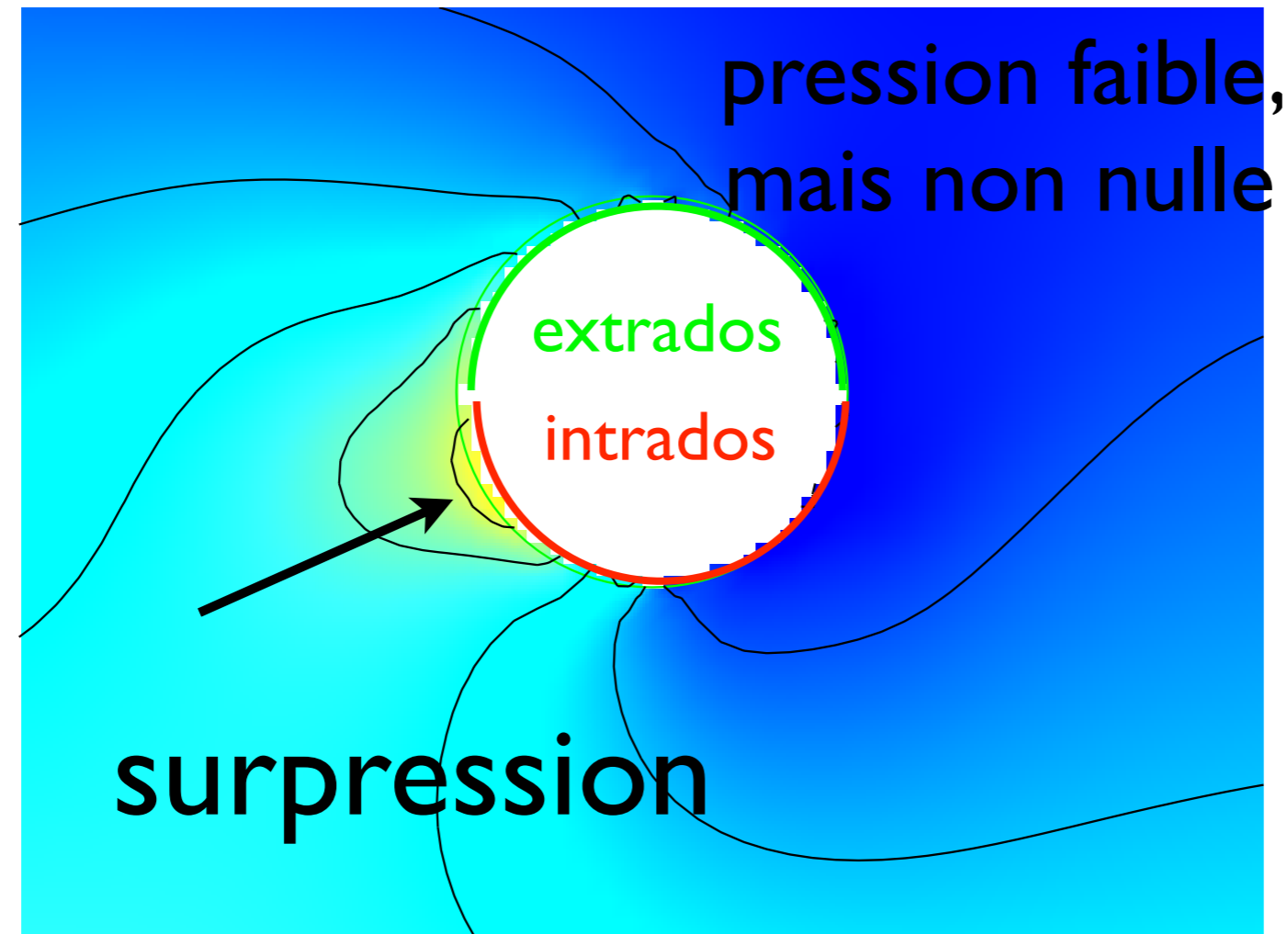
~~$$F_{\text{lift}} = (\rho g L D^2 / 4) \int_{-\pi}^{\pi} \tilde{\sigma}_n^{g=0}(\theta) \sin^2 \theta d\theta$$~~



The flow itself is strongly modified in the presence of the pressure gradient

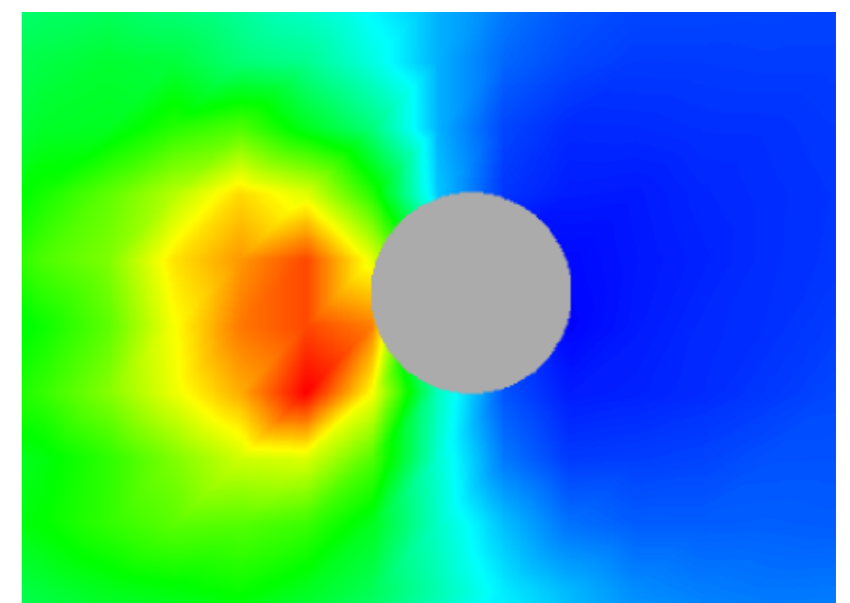
# Simulation of the visco-plastic rheology ( P. Y Lagree, M. Médale ) :

simulation



a lift exists!!

DEM simulation



## Conclusion:

- pressure imposed rheology an interesting approach
- visco-plastic frictional description captures many observed features in dense granular flows
- Hydrodynamic approach

## Questions:

- microscopic origin: role of contact? of fluctuations?
- link with quasi-static regime, and collisional regime?

thanks..

## Suspensions

François  
Boyer



## Granular media



François  
Guillard

## Granular collapse



Loic Rondon



Etienne  
Couturier

Elisabeth  
Guazzelli



Yoel Forterre



Pascale Aussillous