

Power Law Correlations and Avalanche Distributions in Quasistatic Shear

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C. Maloney, Carnegie Mellon Univ.; MASM 5, 11/20/2009

How do microscopic displacements accommodate total global strain?

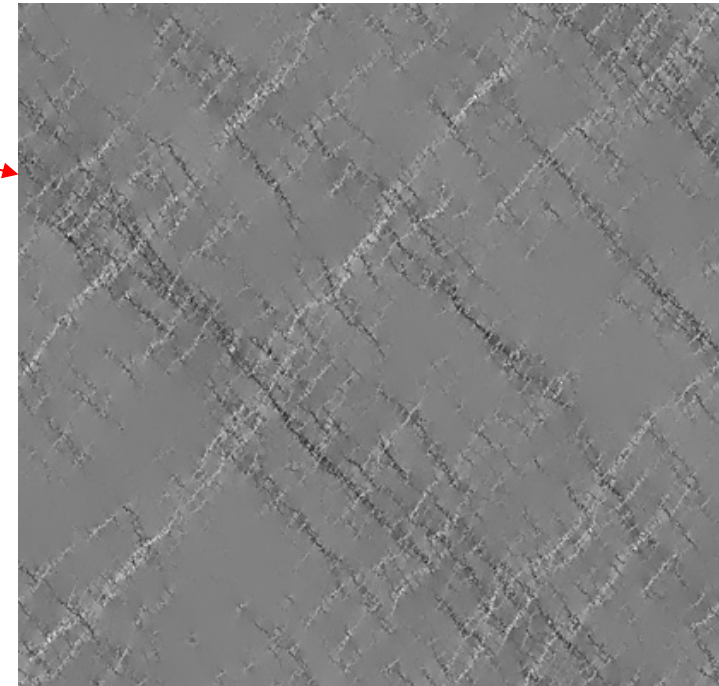
Is there a typical size of plastic event?

How are “earthquakes” distributed in energy, location and time?

How does inertia affect deformation?

Many systems critical without inertia.

Here critical with inertia.



Motivation

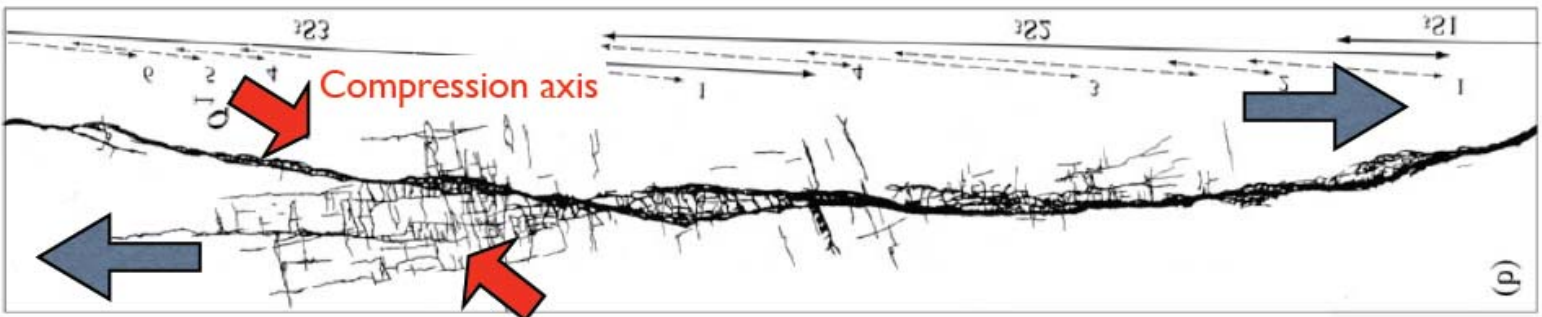
Find power law distribution of events, avalanches, earthquakes in a wide variety of systems and on wide range of scales as long as they are driven slowly

- Charge-density wave motion
- Fluid invasion of porous media, contact line motion
- Magnetic domain switching (Barkhausen noise)
- Deformation of solids
(acoustic emission, dislocation bursts, cracks, earthquakes)
- Granular media, foams,

For large objects or events temperature may be irrelevant

Dynamics are also generally underdamped at large scales

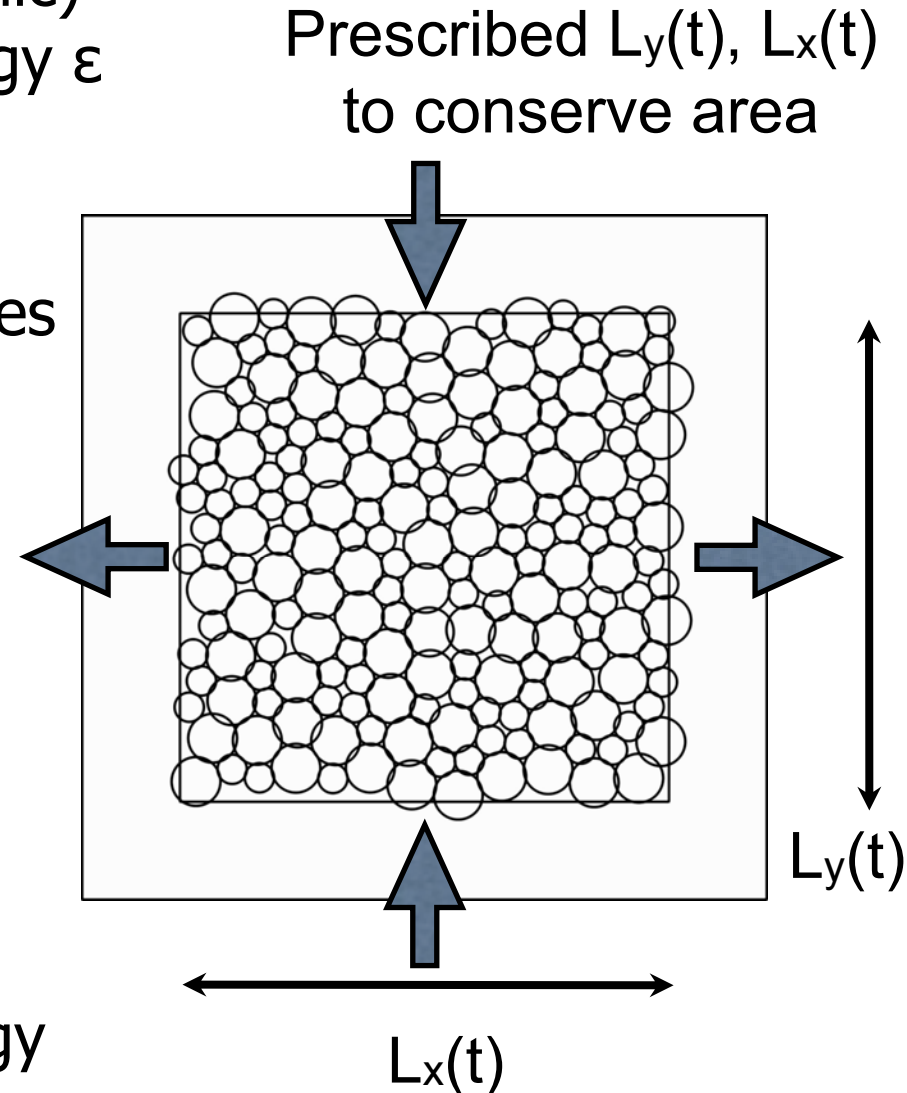
⇒ Focus on quasistatic, $T=0$, momentum conserving dynamics



Quasistatic Athermal Simulations

2D Molecular Dynamics (usually)

- Binary Lennard-Jones (or harmonic)
Mean diameter σ , binding energy ε
- Quenched at pressure $p=0$
protocol not important
- $T=0$, Underdamp relative velocities
Galilean invariant (Kelvin/DPD)
- Periodic boundaries (usually)
- Axial, fixed area strain or
simple shear, const. rate
- Quasi-static limit \Rightarrow low rate
controlled by $\Delta\gamma$ not Δt
Different than saying motion of
all atoms is always slow,
i.e. overdamp or minimize energy



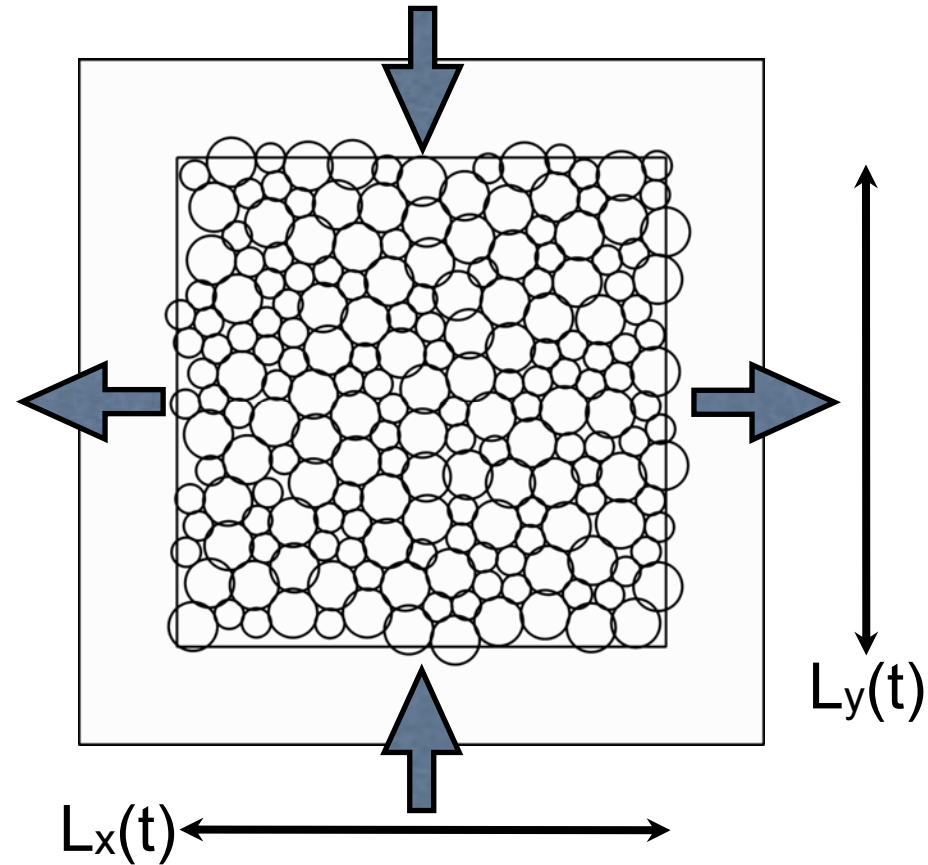
Nonaffine Particle Displacements

Non-affine displacement u =
deviation from mean motion

Integrate affine displacement
along trajectory rather than
using value for initial position.

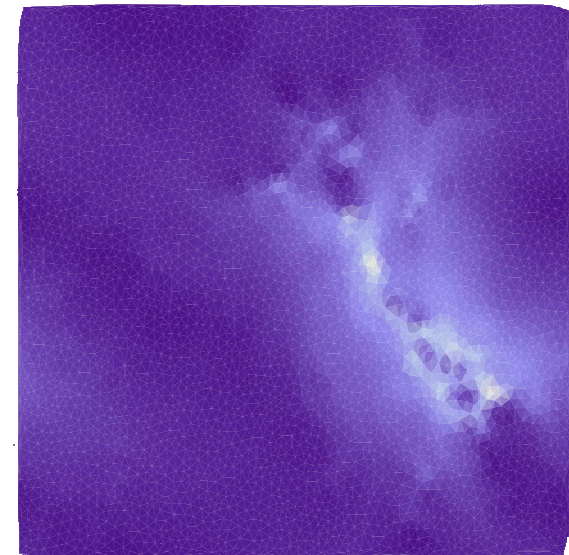
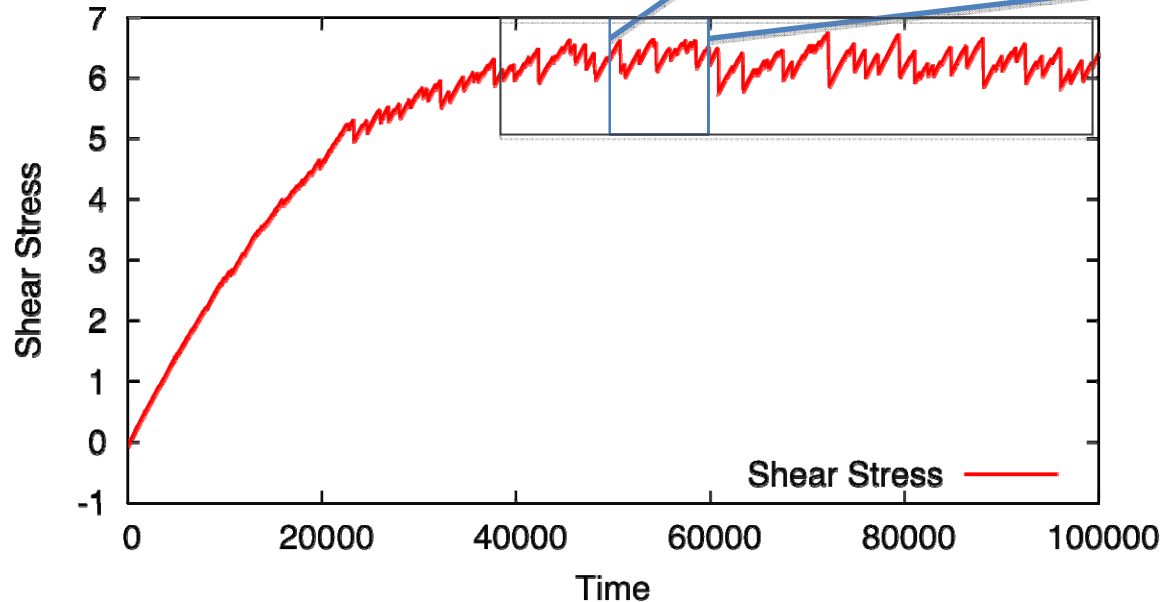
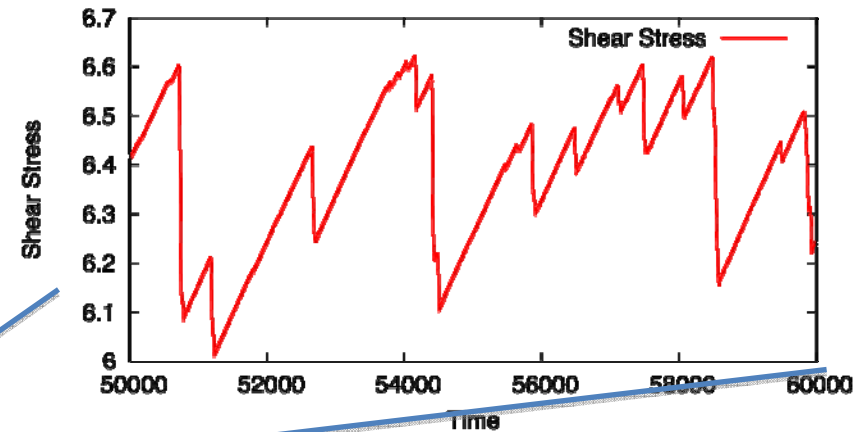
For each particle at each
time step find distance
moved relative to affine
displacement at
instantaneous position.

Eliminates terms analogous
to Taylor diffusion in simple
shear



Quasi-static, Steady-State Shear

- Shear rate low enough that independent (not causally related) events are separated in time. Typically $10^{-6} t_U^{-1}$ or less
- Number independent of rate
- Statistical properties independent of initial configuration
- Atomic velocities fast, only shear rate low



Defining Avalanches

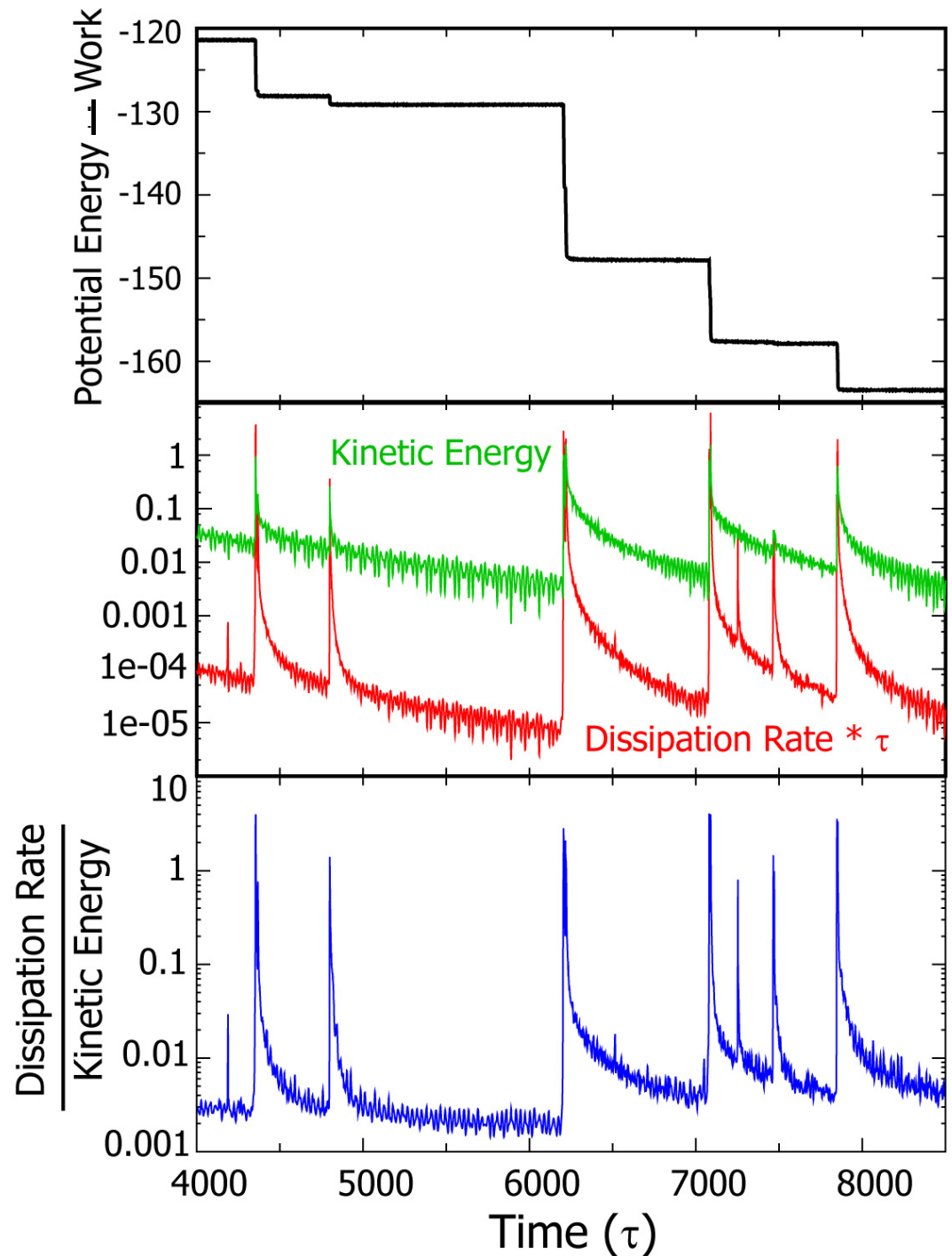
Deformation through series of avalanches

Find energy dissipated = change in potential energy - work done by system.

Drops have wide distribution, 0.03 to 20 here
Identify by sharp rises in dissipation rate.

Dissipation rate/Kinetic energy drops to constant as energy moves to longest wavelengths.

Ratio measures wavelength where have kinetic energy



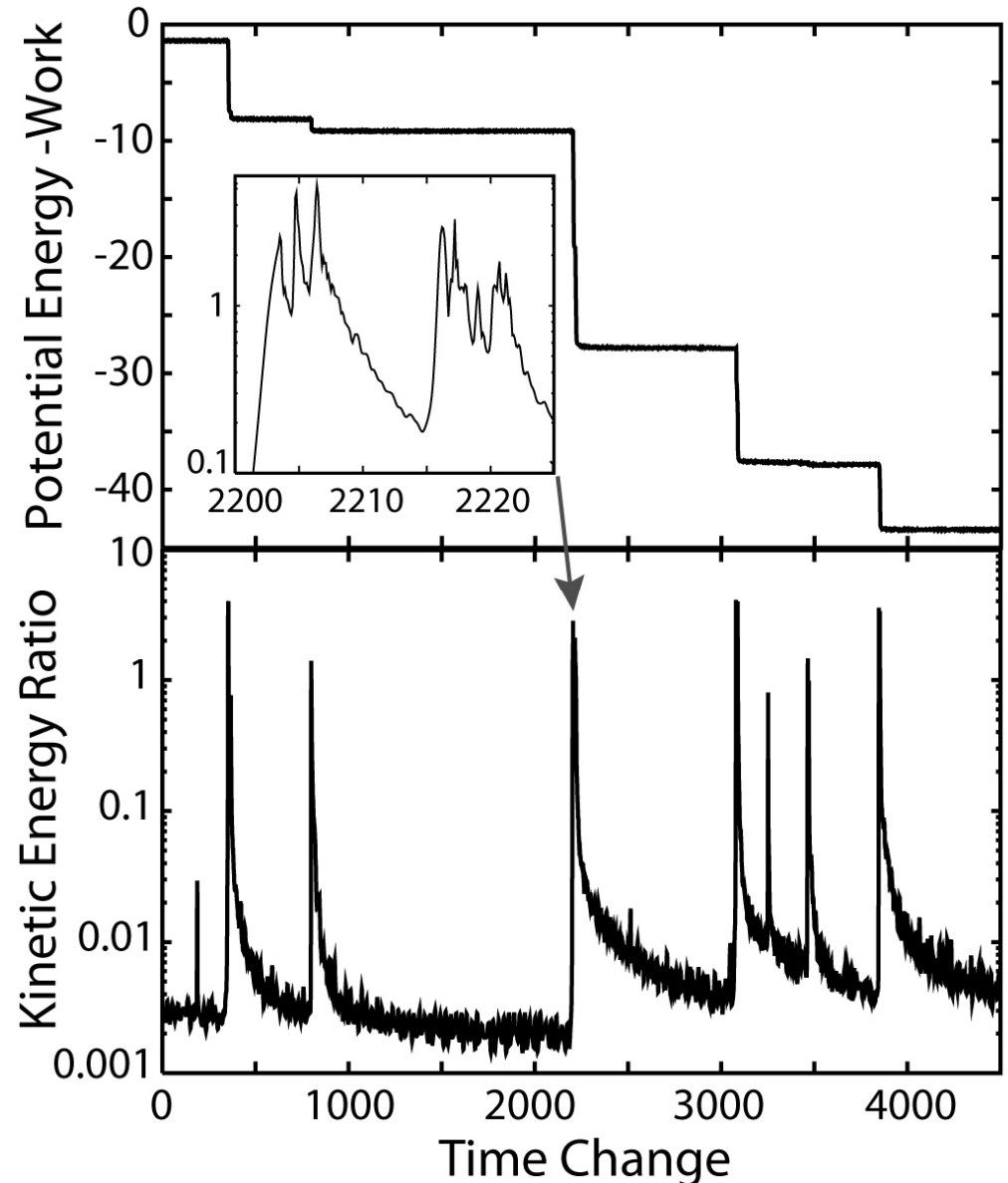
Avalanche Distribution

Amplifying scale shows earthquakes can have complex structure with many spurts of activity.

Find kinetic energy ratio gives a constant number of events per unit strain in the limit of low rates

Have also tested that stopping strain during event gives same statistics

⇒ quasistatic

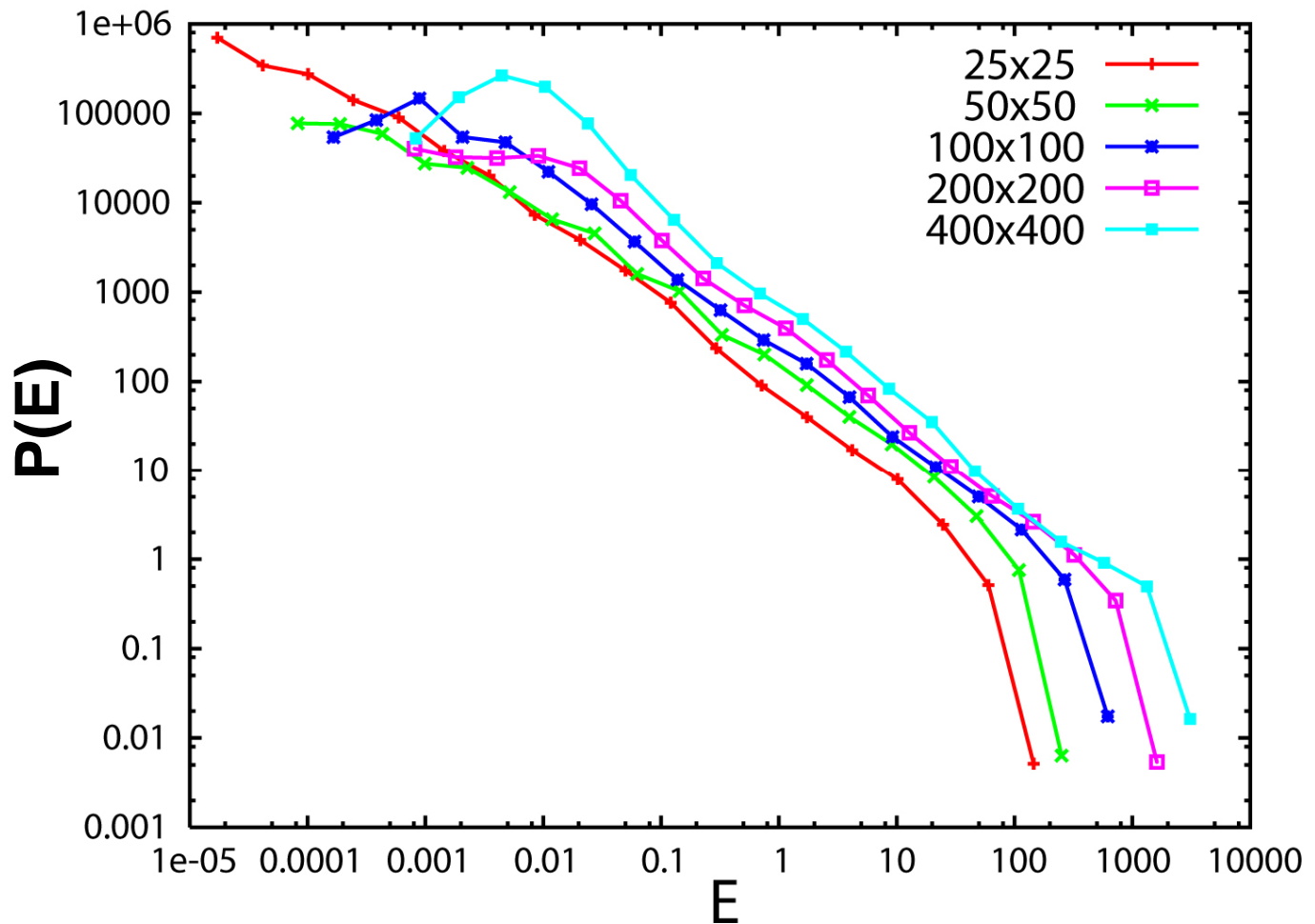


$P(E) = \text{Number of Events per Unit Strain per } E$

Reduce strain rate so quasi-static, only affects small events.

Find power law $P(E) \sim 1/E$ over 5 decades.

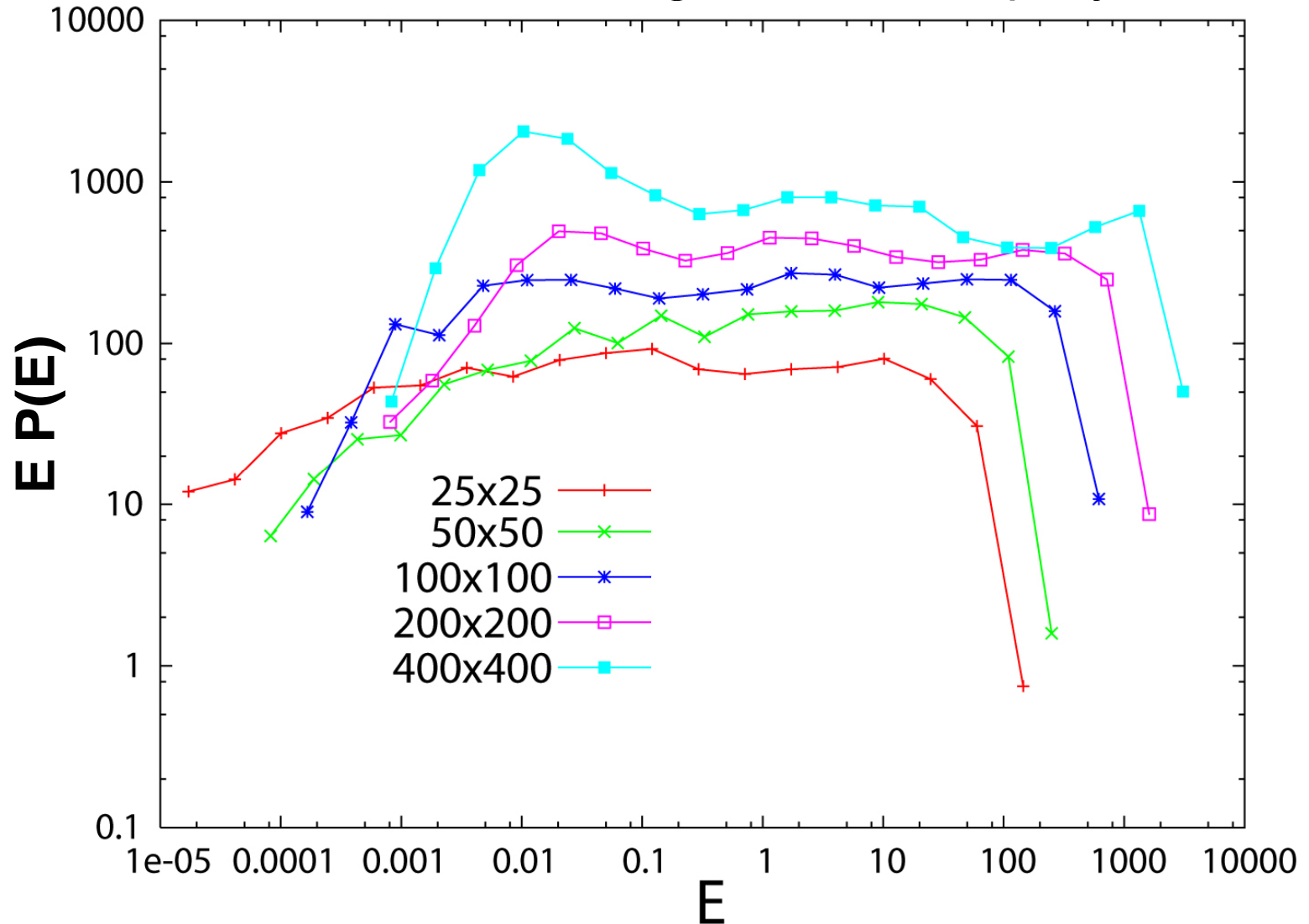
$P(E)$ and maximum E increase with system size



$E P(E) \sim \text{Constant}$

Find $P(E) \sim E^{-1}$

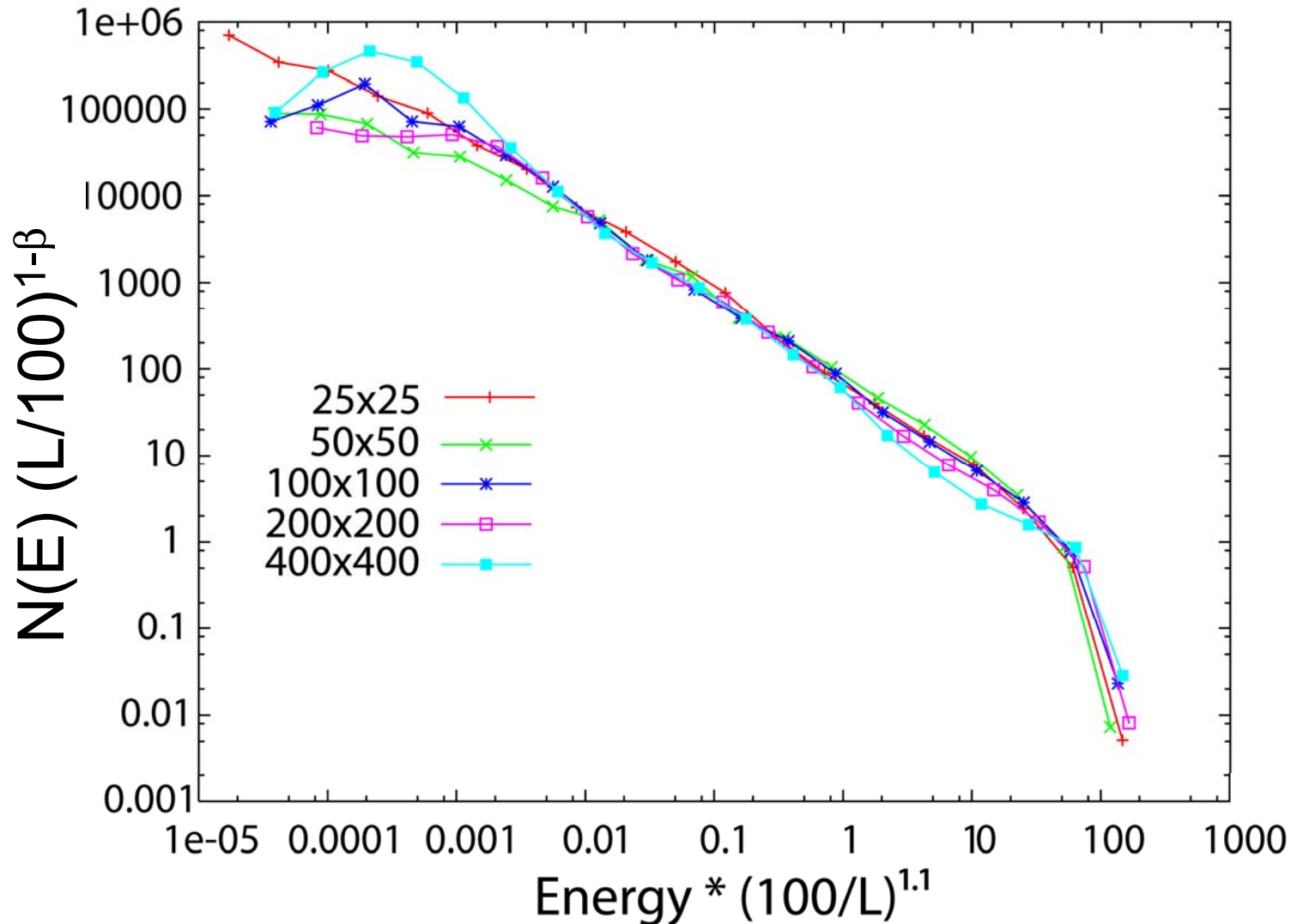
Maximum event size grows more rapidly than length L



Scaling Collapse of $N(E)$ and E

Scale E by $E_{\max} \sim L^\beta$. Find $\beta \sim 1.1$

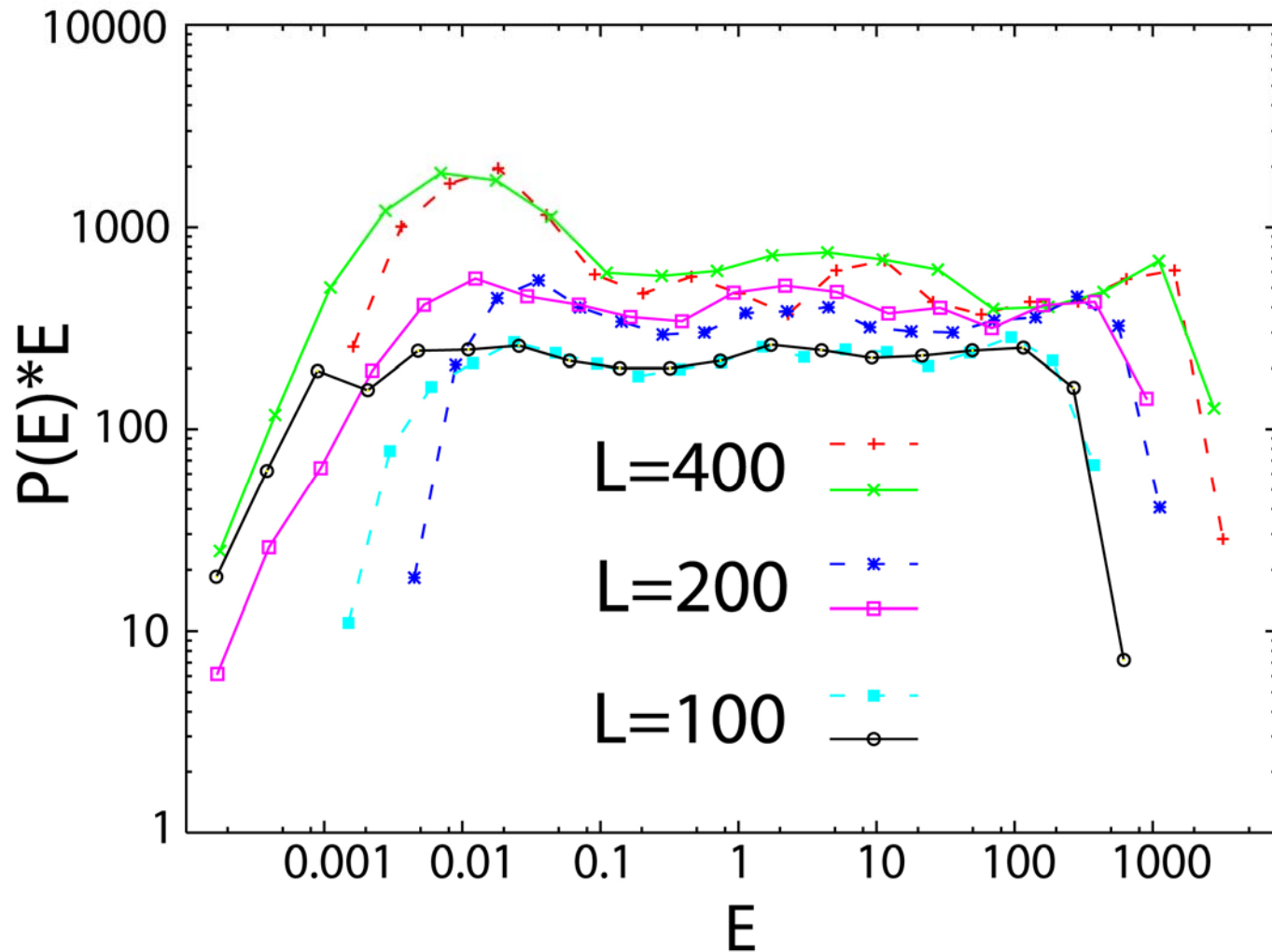
$N(E/E_{\max})$ must scale as $L^{1-\beta}$ to maintain energy balance



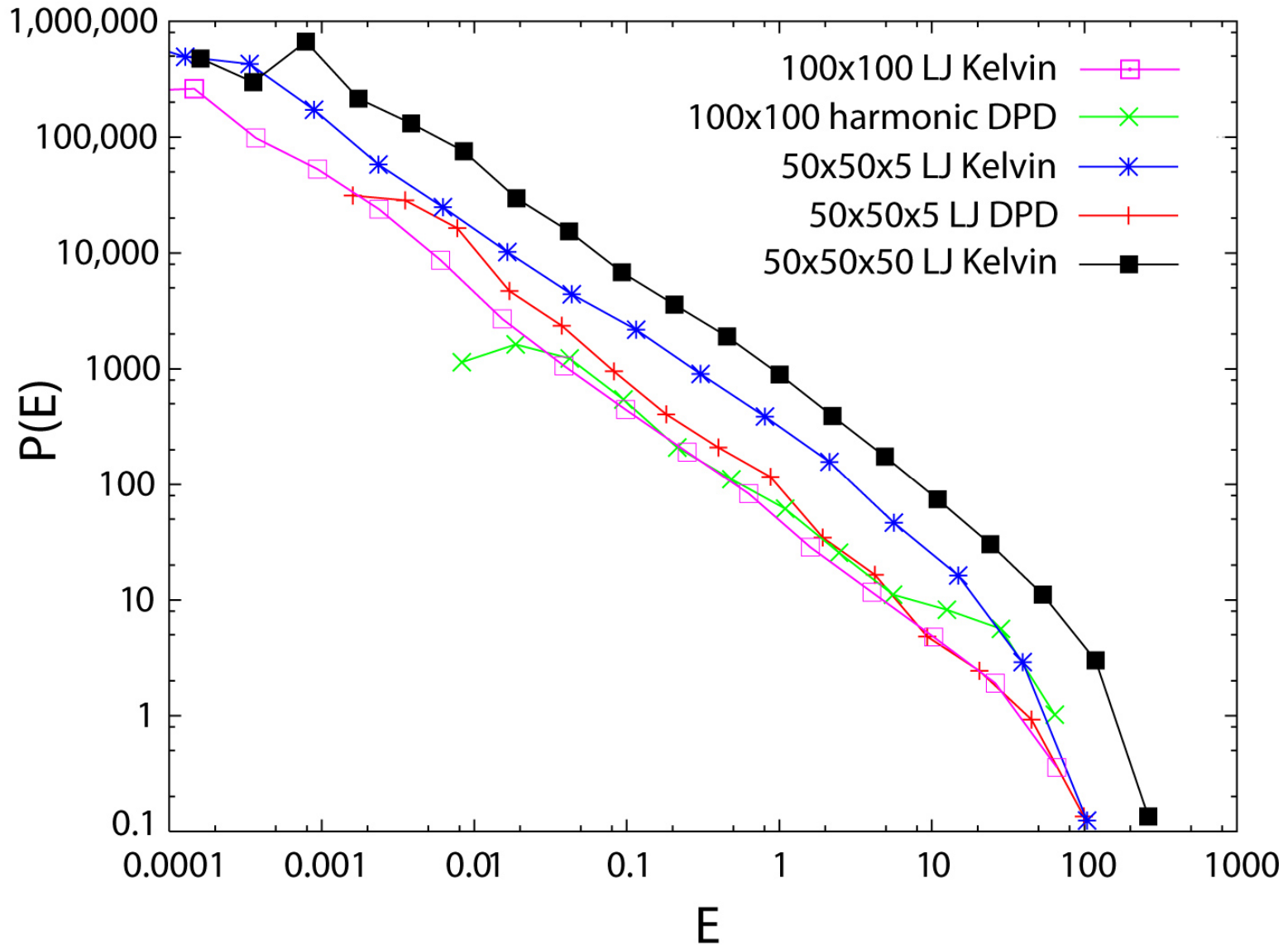
Slow Enough?

Stop straining when detect event

Same distribution (dashed curves) as continuous strain (solid)



Power Law Independent of Potential, Geometry and Thermostat

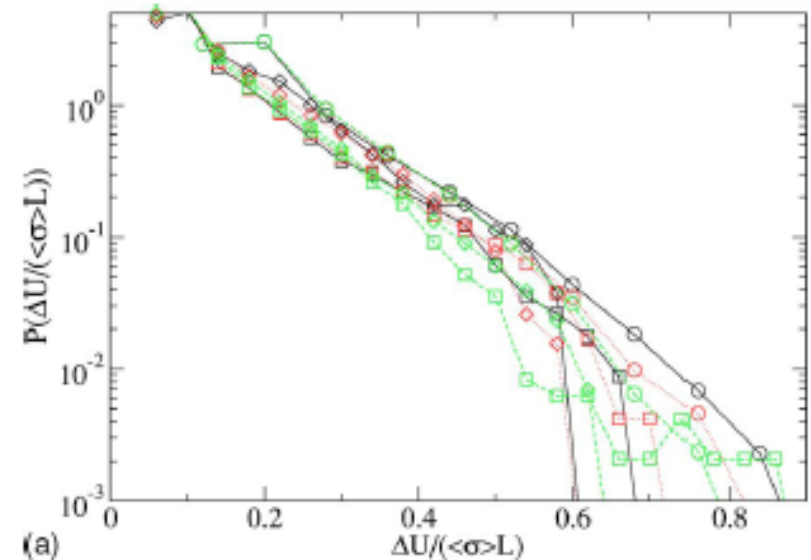
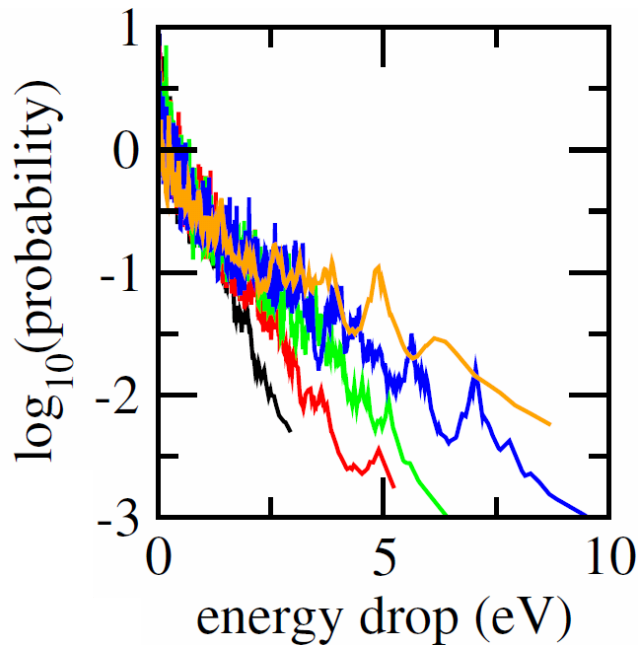


Exponential Scaling of $P(E)$ in Other Models

Same system but with energy minimization, not dynamics

Largest event $\sim 3\varepsilon$ vs. 200ε

(Maloney & Lemaître PRE 74, 016118, '06)



3D amorphous metal

$$N(E) \sim \exp[-E/L^x] \quad x=1.4$$

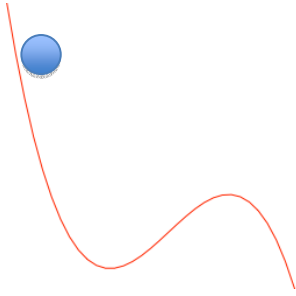
(Bailey, Schiøtz, Lemaître & Jacobsen, PRL 98, 095501, '07).

Do see power laws in overdamped systems with quenched disorder
AND

Inertial models without disorder

(Carlson, Langer, Shaw – Burrigge-Knopf)

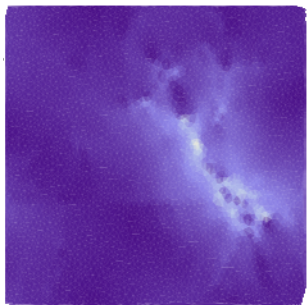
Inertia Matters



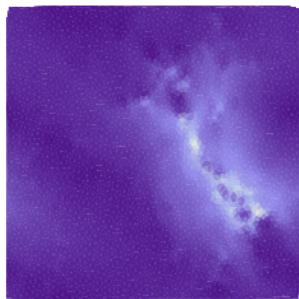
Energy minimization finds nearest minimum
Energy lower with inertia

Inertial overshoots -> weaker regions

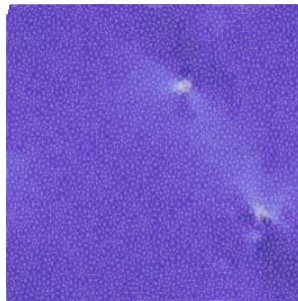
Overdamped



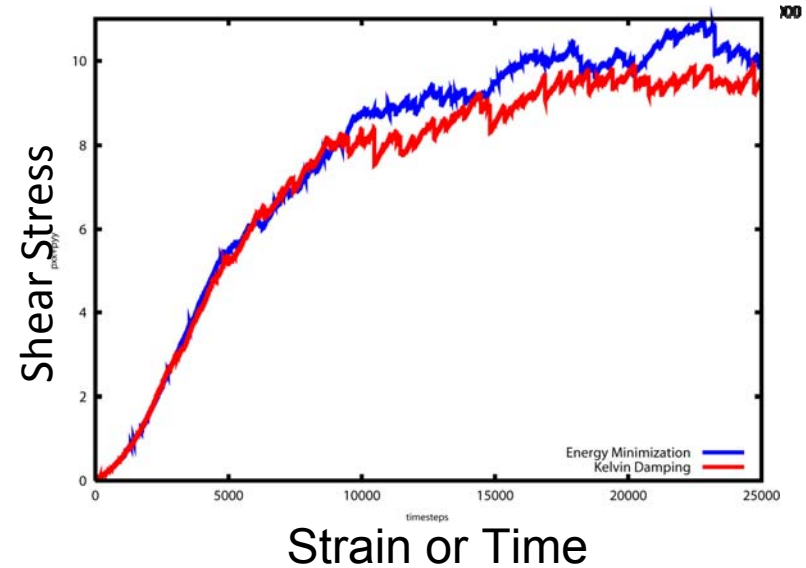
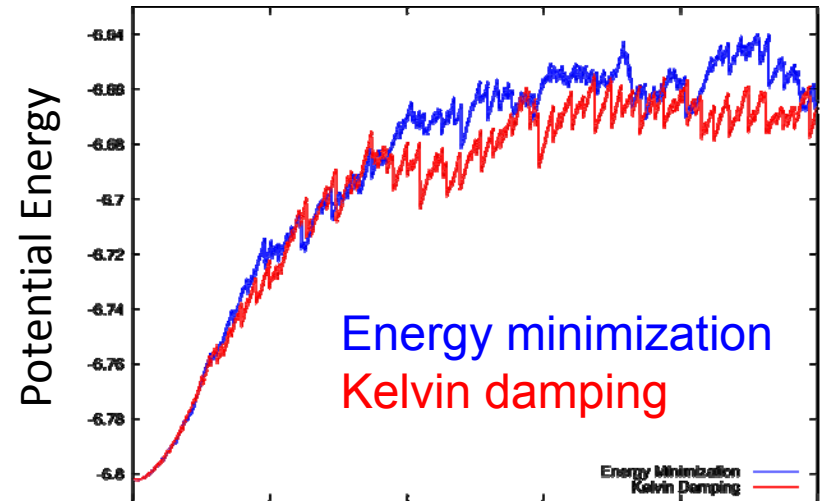
Inertial



Difference

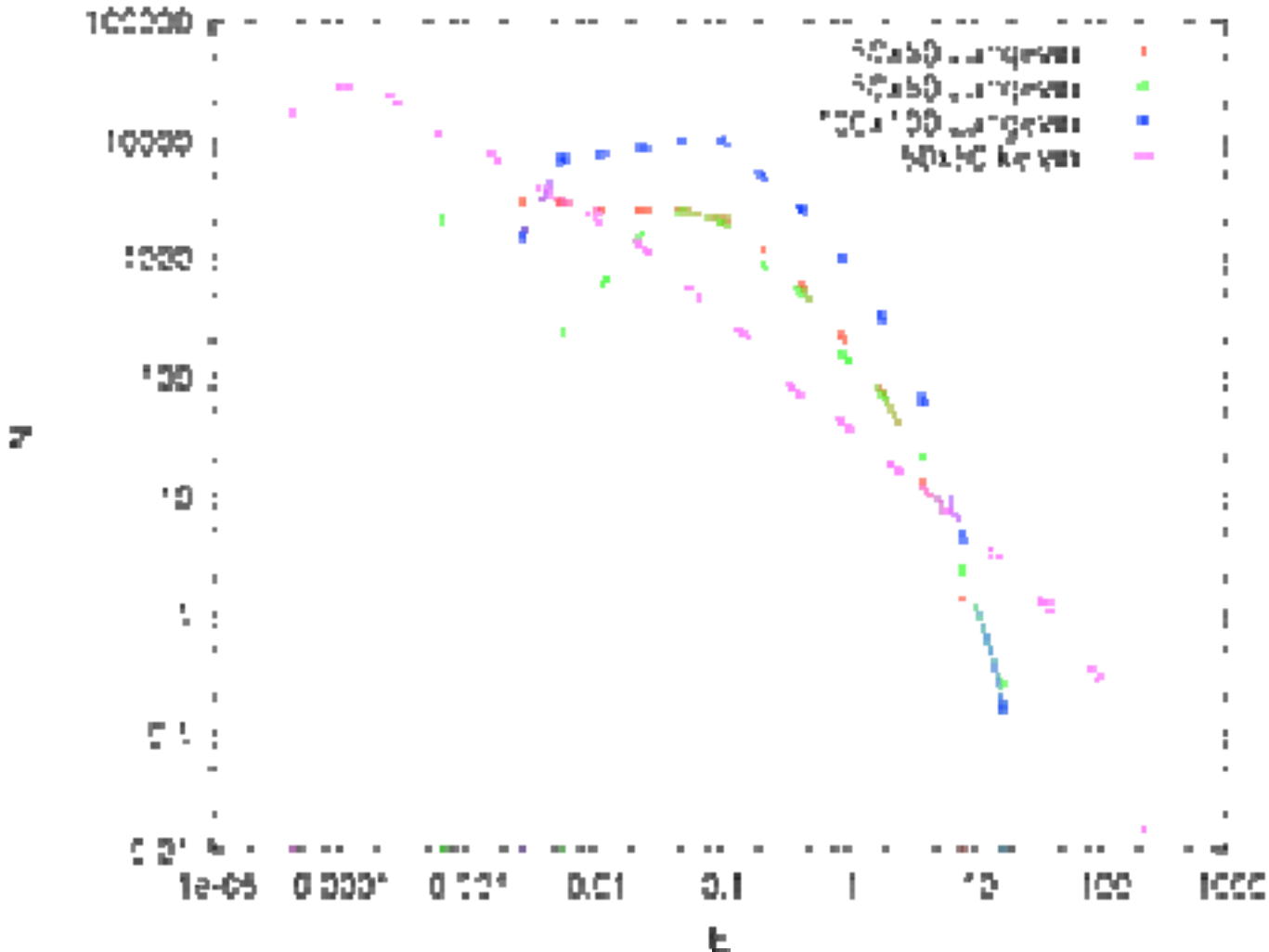


white is larger displacement of particles



Langevin Damping Changes Distribution

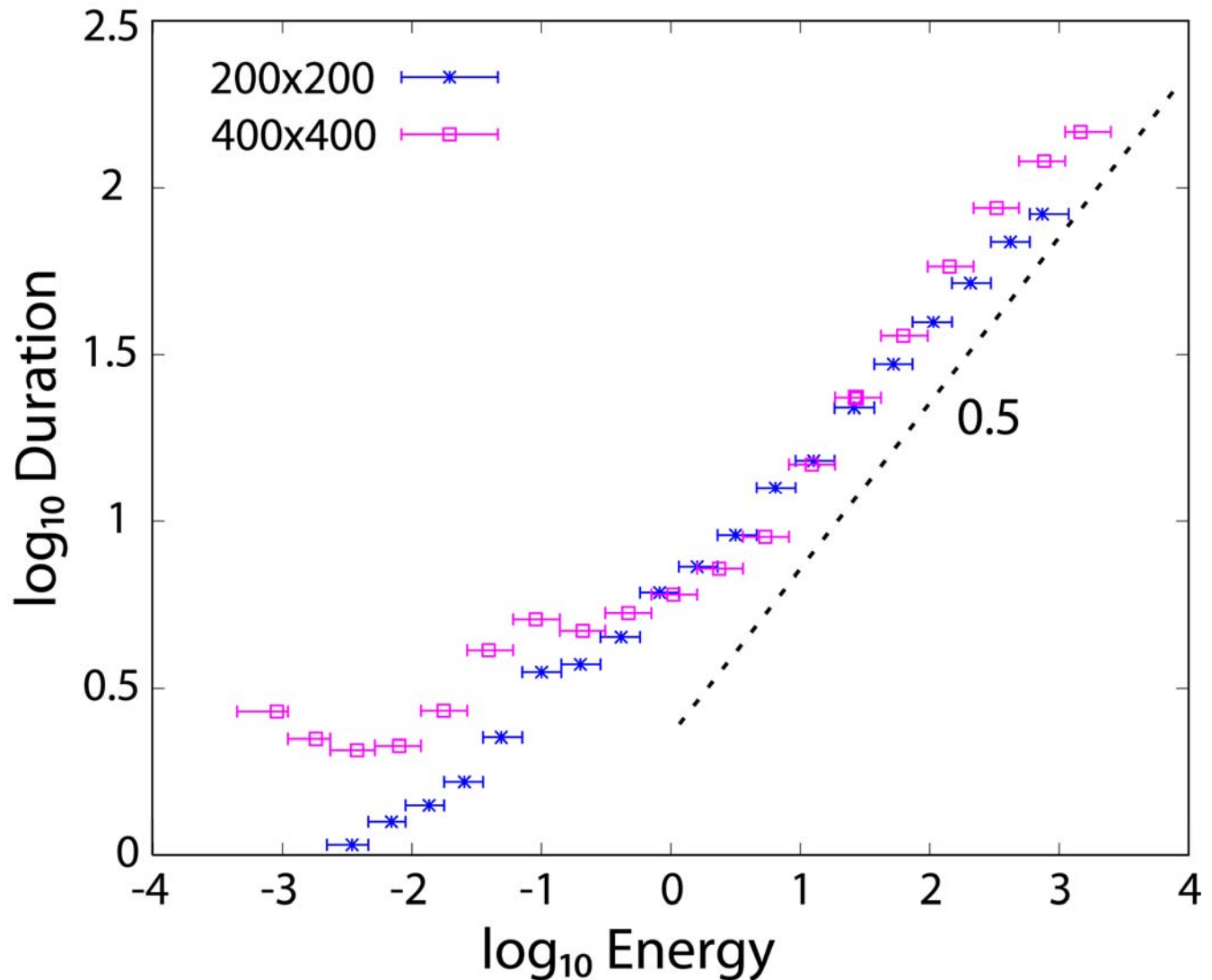
Overdamped systems \Rightarrow fewer events at small and large E
Tail at large E \sim exponential



Relation Between Duration and Energy

Duration \sim square root of energy

Consistent with $\text{area} \sim \text{energy}$ and $\text{duration} \sim \text{length} = \text{area}^{0.5}$

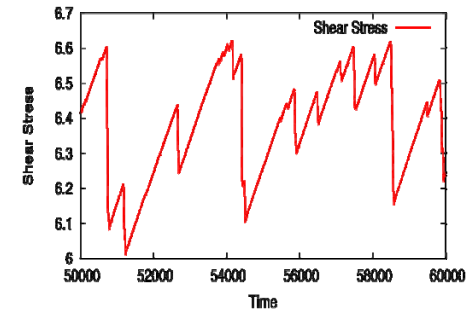
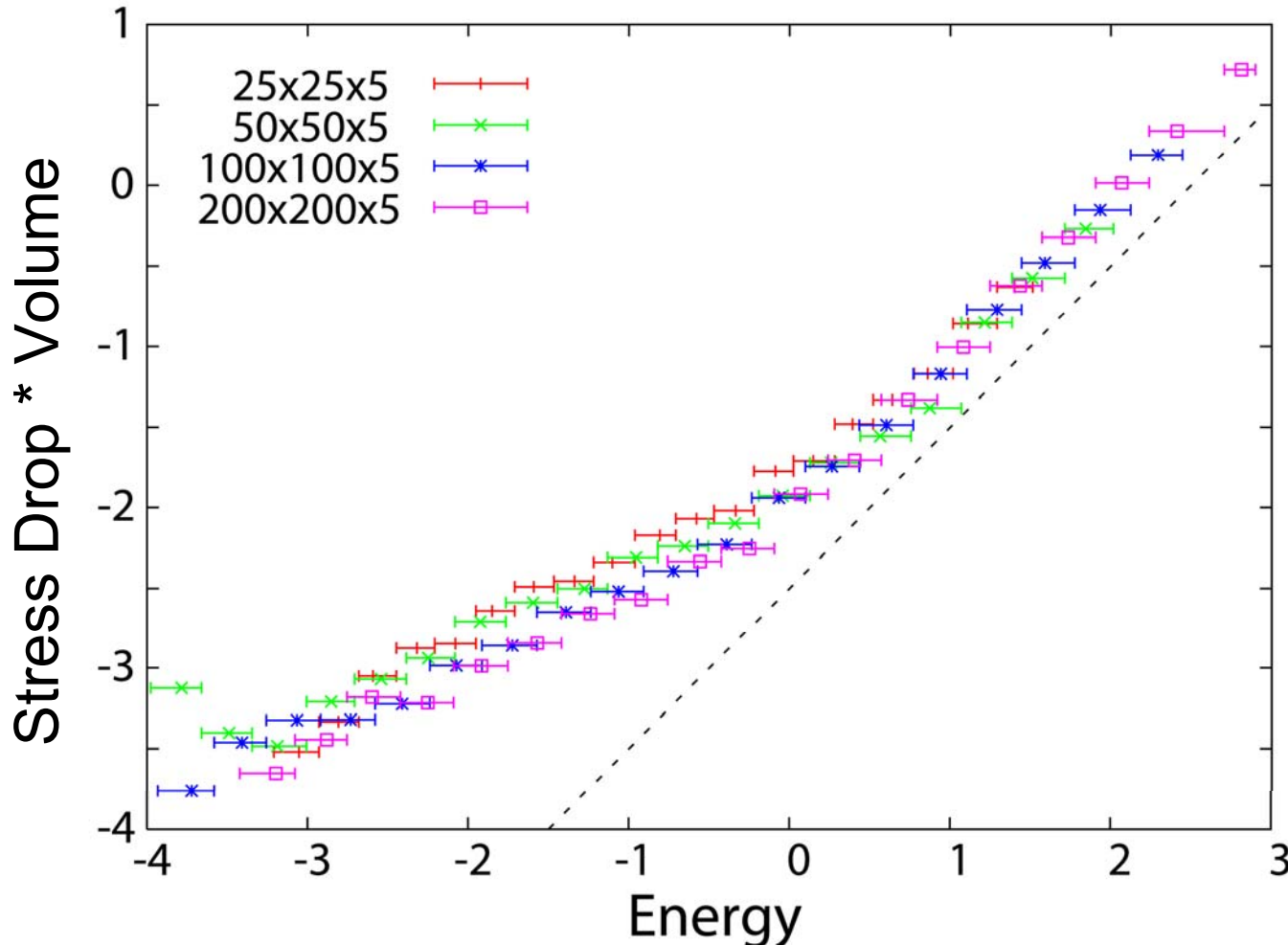


Relation Between Stress Drop and Energy

Mean stress over whole volume * volume ~ moment

Find moment ~ energy for large events (dashed line)

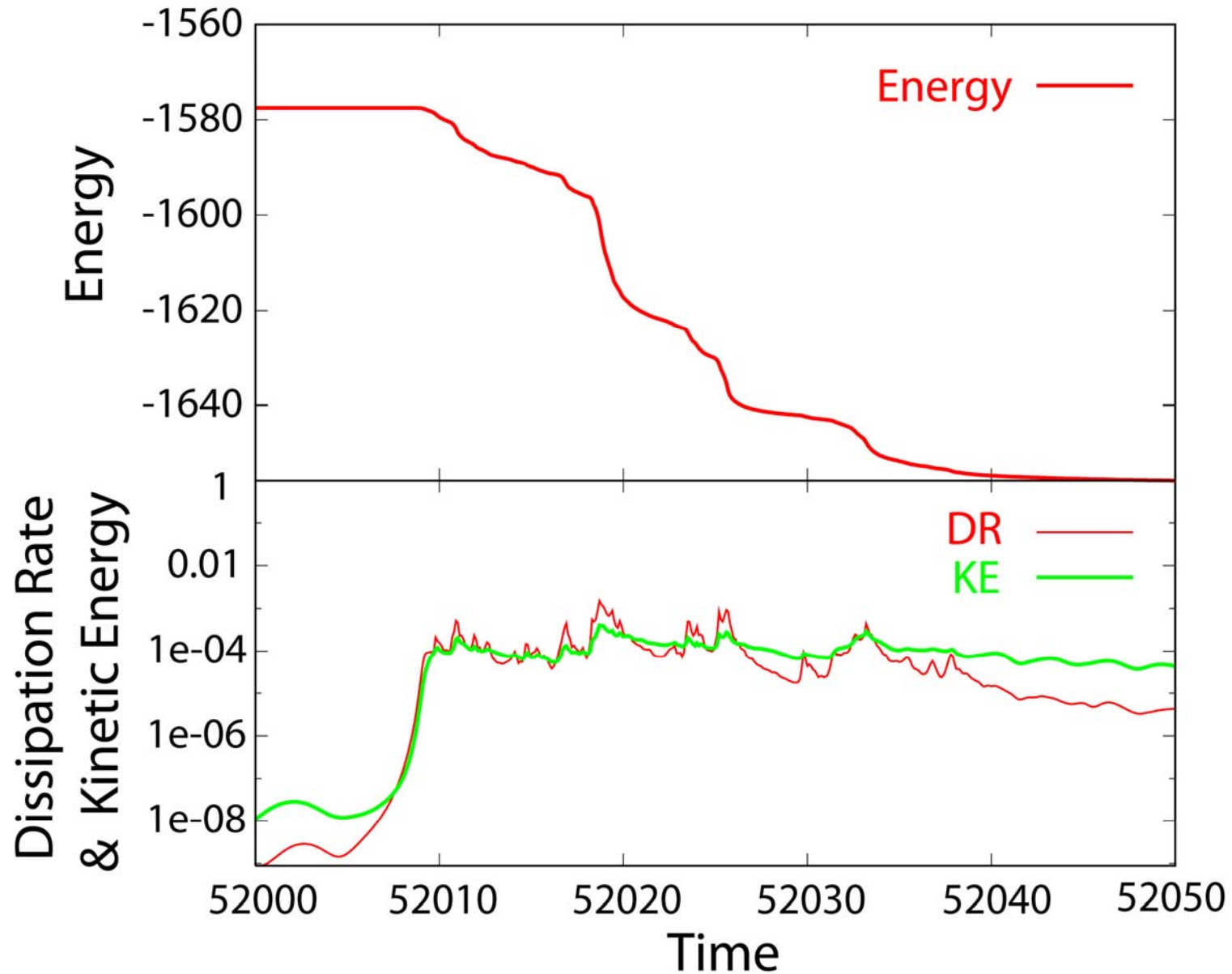
Large dispersion for small events



Different distribution if use stress minima and maxima to define event

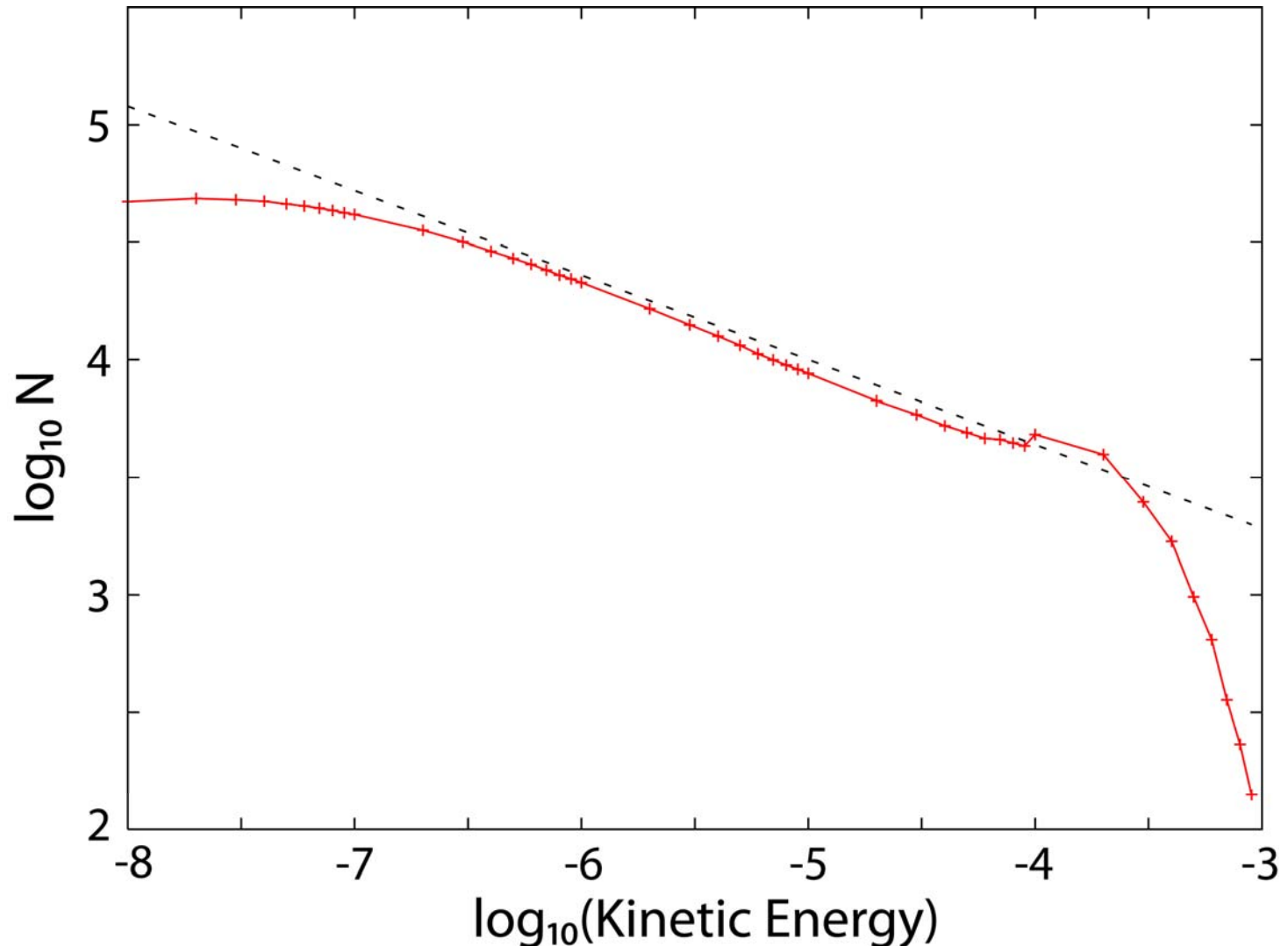
Our Quakes Include Fore and After Shocks

How Does this Affect Statistics?



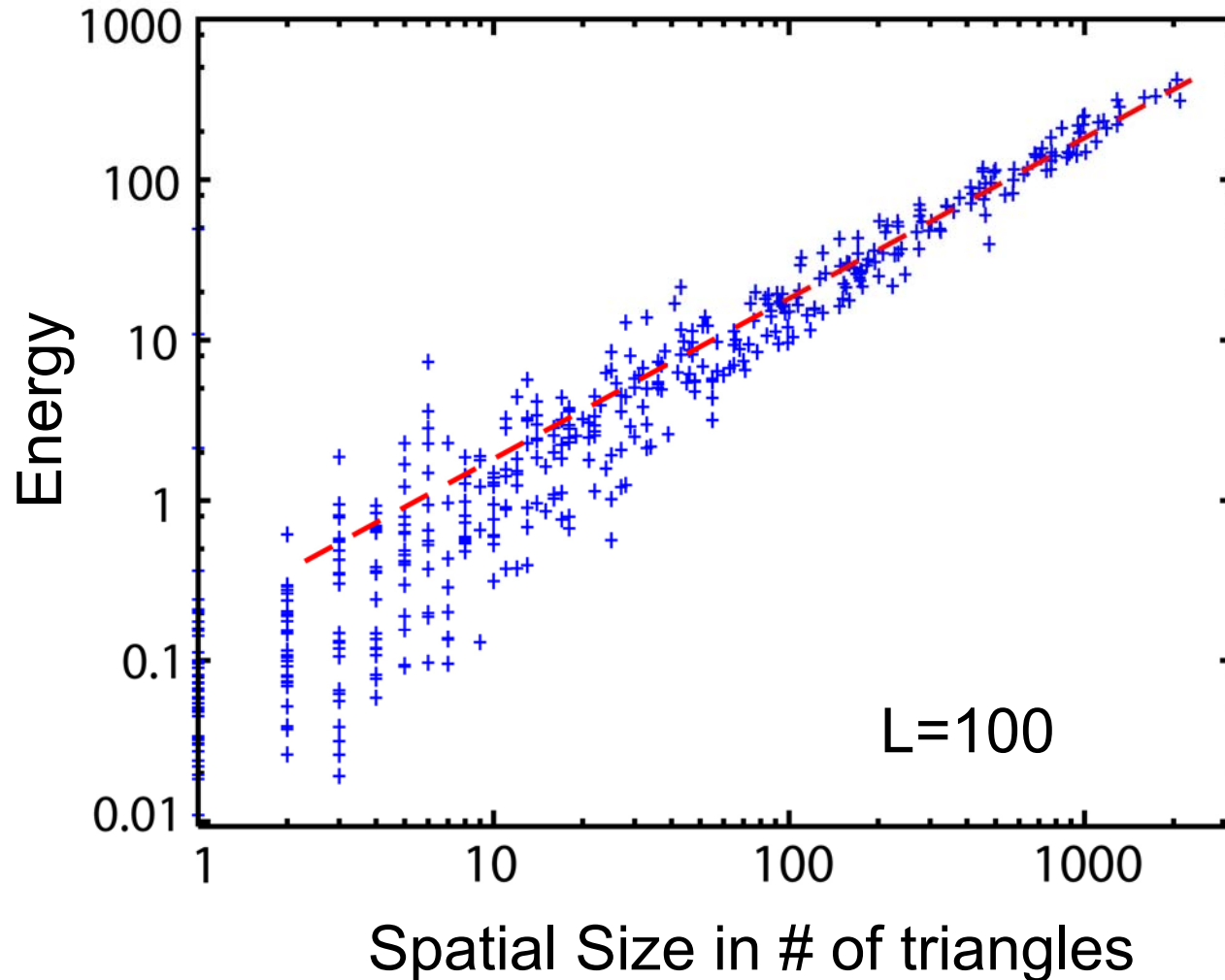
Integrated Density of Quakes Based on Kinetic Energy

Large events are broken up. Find integrated density $\sim E^{-0.4}$
Gutenberg-Richter $\sim E^{-2/3}$



Relation Between Energy and Spatial Size

Size = # triangles where curl of displacement field
has magnitude > 0.1 (limit of elastic deformations)
Find linear (dashed line) relation for large events



Defining Local Deformation

Shear related to local rotation $\omega = \nabla \times \mathbf{u}$, \mathbf{u} =displacement
 Find Delaunay triangulation for initial particle centers

For each triangle:

$$\frac{\partial u_i}{\partial x_j} = F_{ij}$$

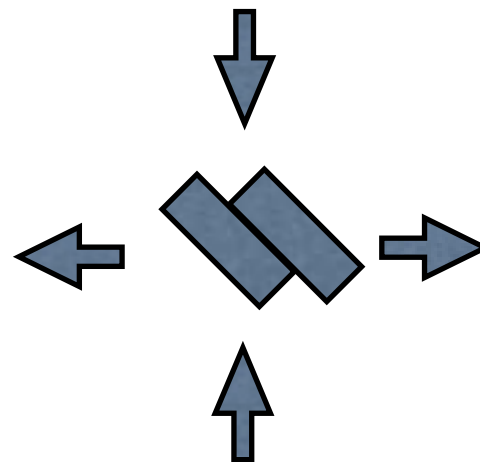
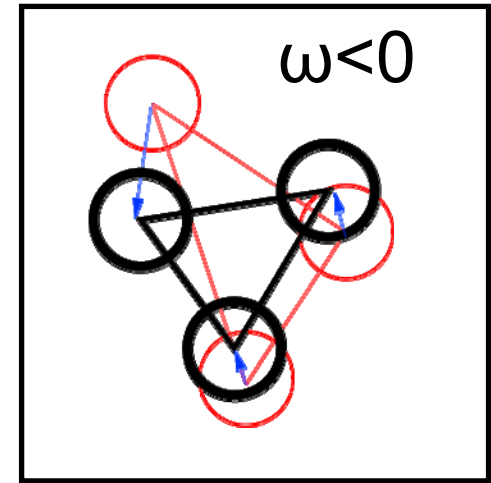
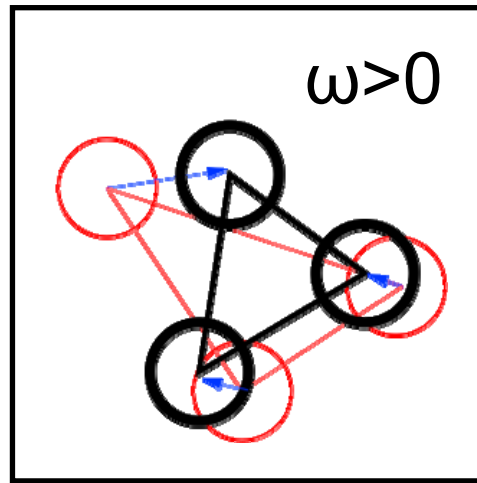
$$\epsilon_1 = \frac{F_{xx} - F_{yy}}{2}$$

$$\epsilon_2 = \frac{F_{xy} + F_{yx}}{2}$$

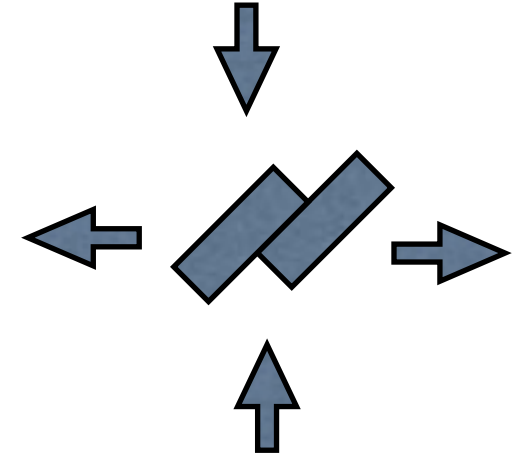
Invariants:

$$\epsilon = \sqrt{\epsilon_1^2 + \epsilon_2^2}$$

$$\omega = F_{xy} - F_{yx}$$



“Right Strain”



“Left Strain”

MR4

Mention coarse-graining property

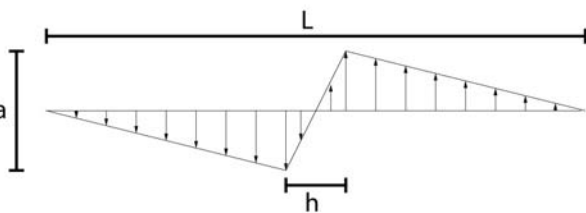
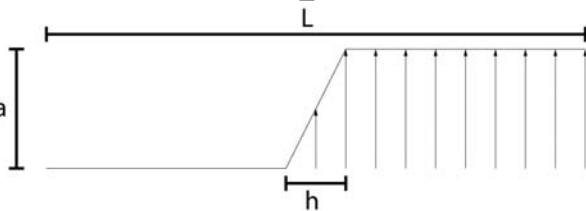
Note we 2x strain

Mark Robbins, 10/30/2008

Spatial Variation of Nonaffine Displacement u

Components of u
during 0.2% strain
 $+\sigma$ white, $-\sigma$ black

Strain localizes in
plastic bands \rightarrow step
in total displacement

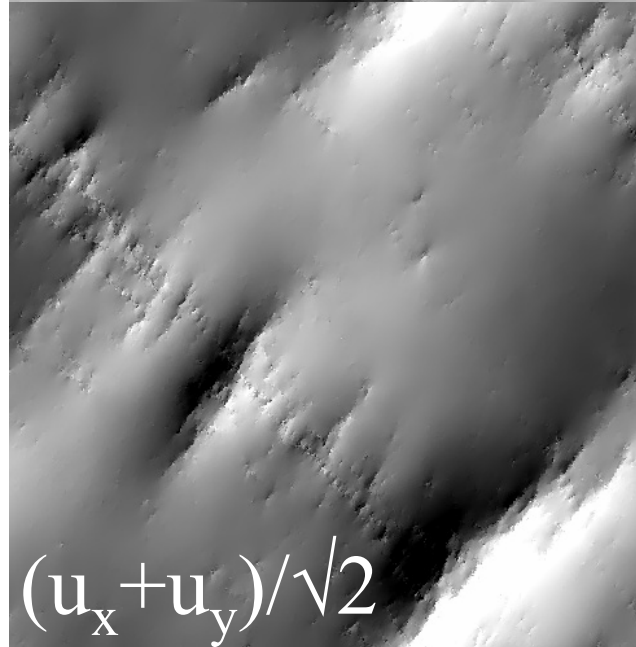
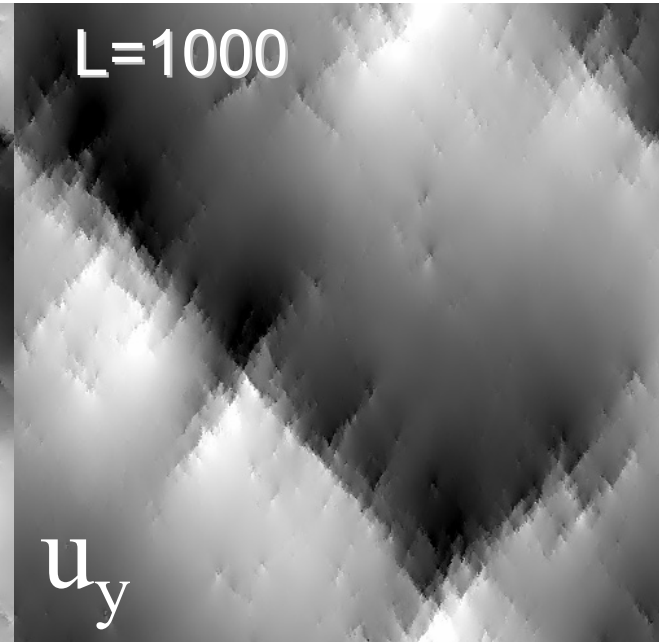
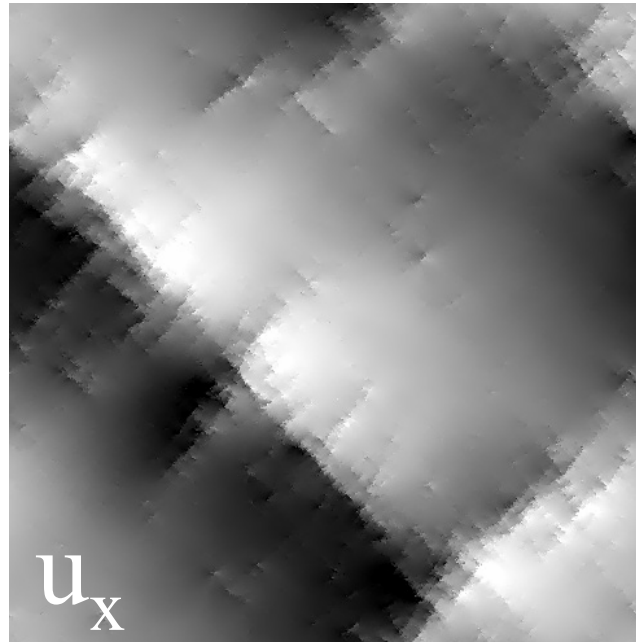


Largest projection of
 u along bands $\pm 45^\circ$

Sign \rightarrow rotation sense

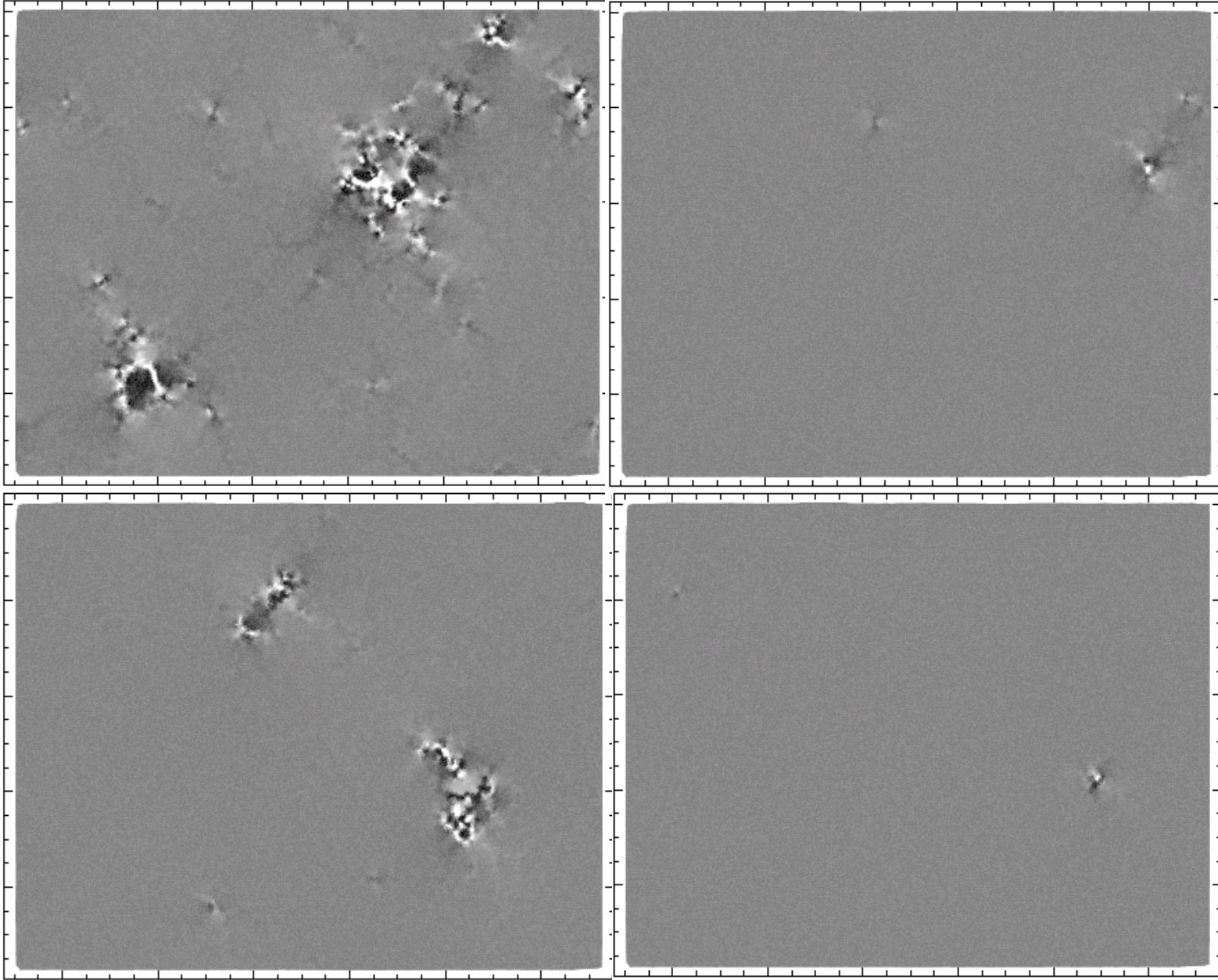
Length up to system

Typical $a \sim \sigma$



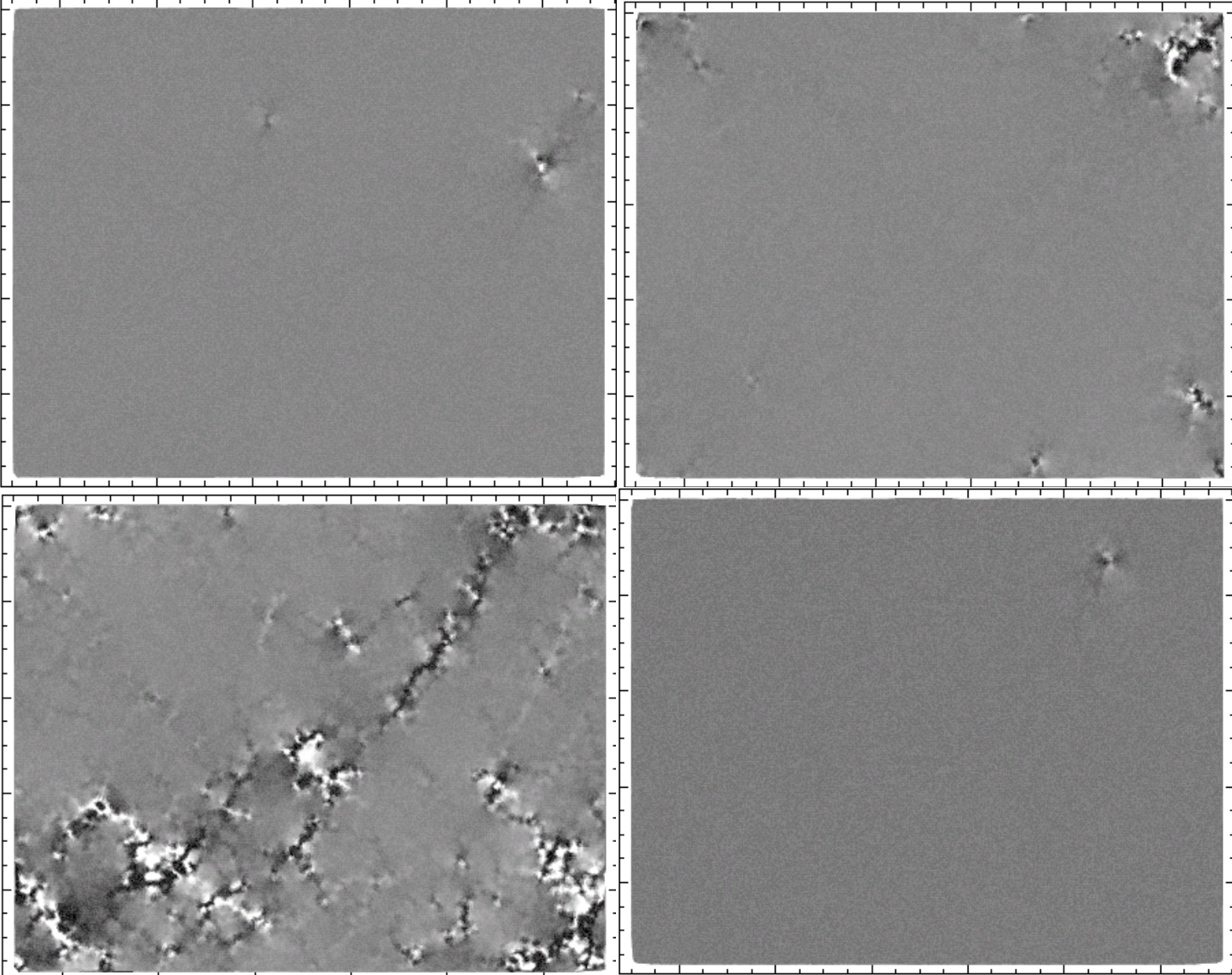
Earthquakes

Examples of events (111.065, 24.5098, $\vec{2.982}$, 1.4181)



Earthquakes

Examples of events (3.2643, 301.448, 21.3862, 0.0794)

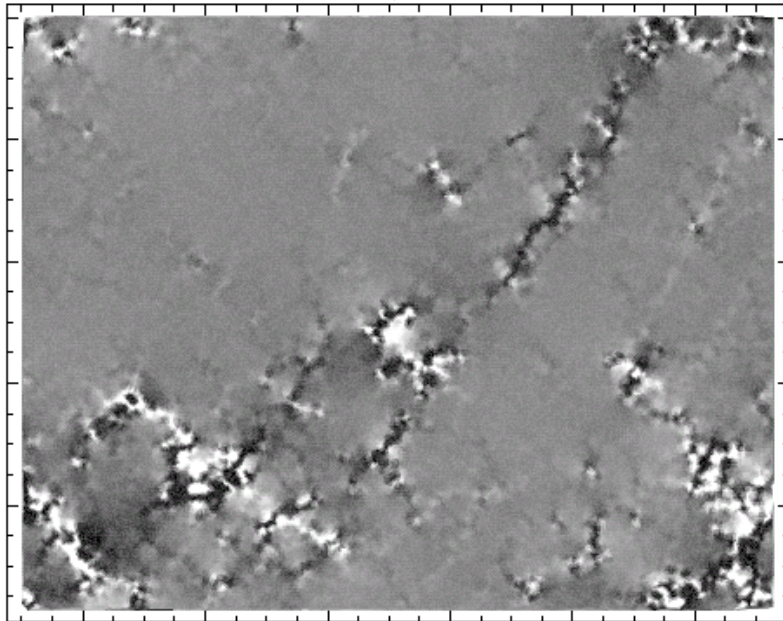
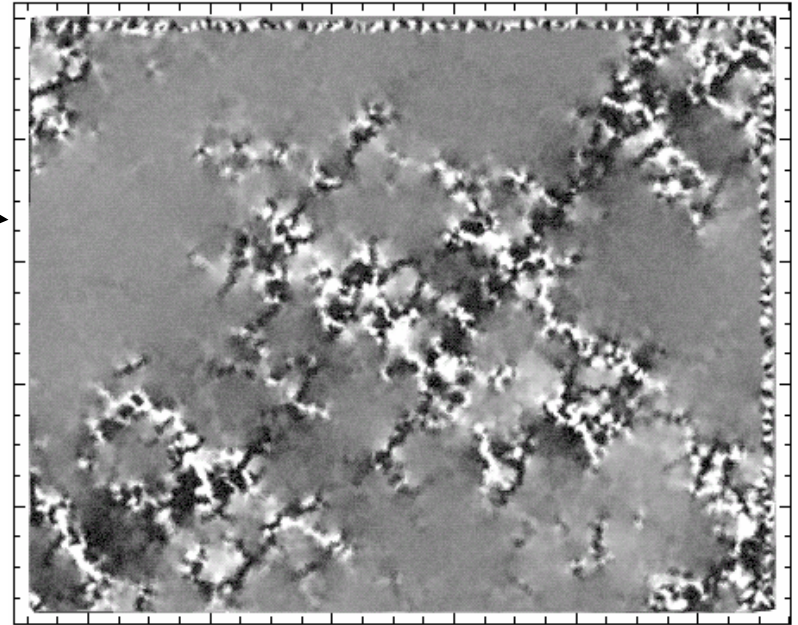


Earthquakes

Total integrated activity shows
long range correlations

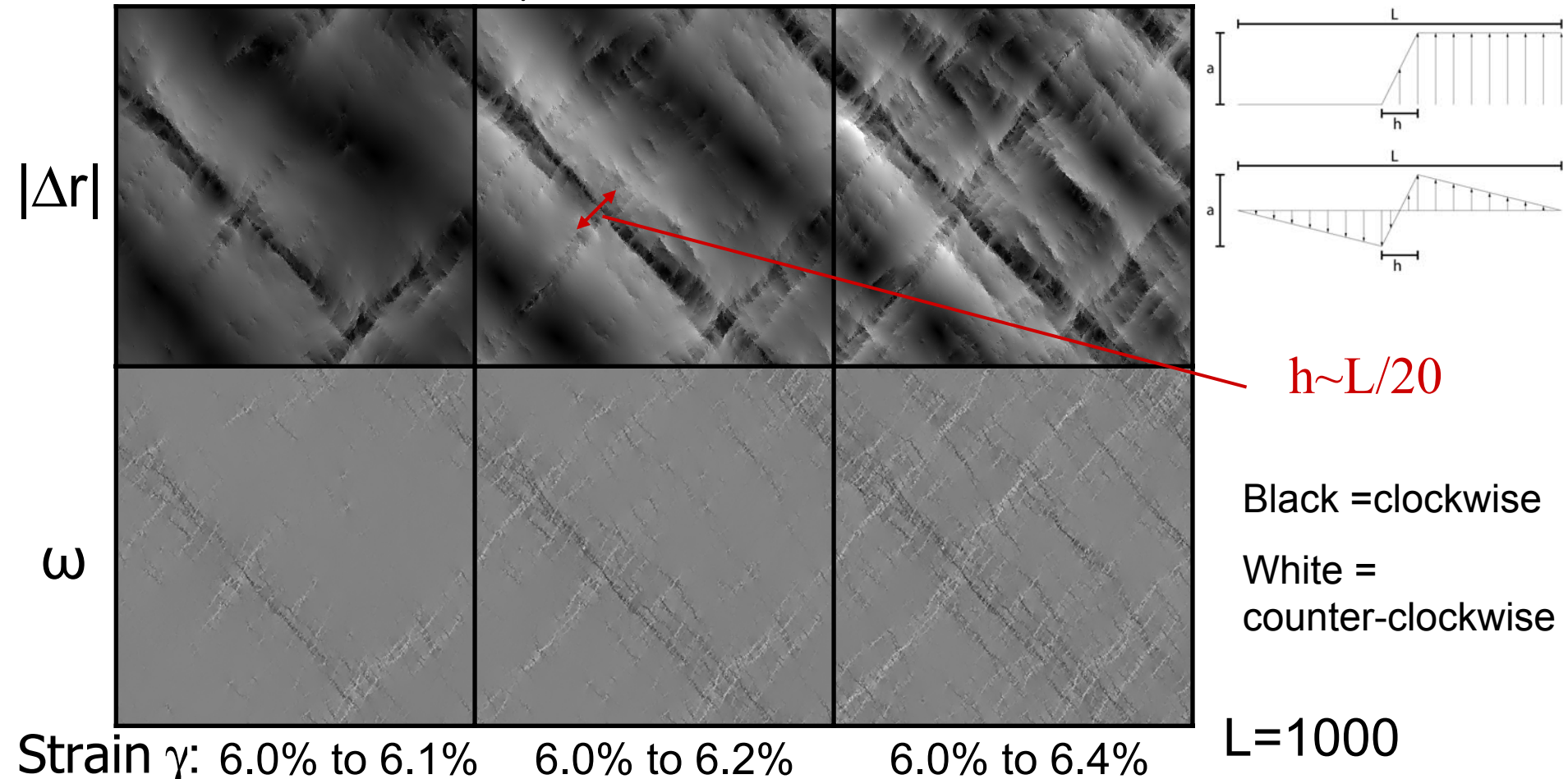
Activity over 1% strain

Single largest event



Displacement and ω in steady state

Most analysis looks at magnitude $\Delta r = |\vec{u}| \rightarrow$ smallest in plastic band
 Curl sharply localized in plastic zone, +/- regions correlate along $-/+45^\circ$
 Strain accumulates to $a \sim \sigma$ in plastic bands through many avalanches
 \rightarrow Displacement $\sim \sigma$ allows all regions to find new metastable state
 Strain over intervals $\Delta\gamma > \sim \sigma/L$ occurs in uncorrelated locations

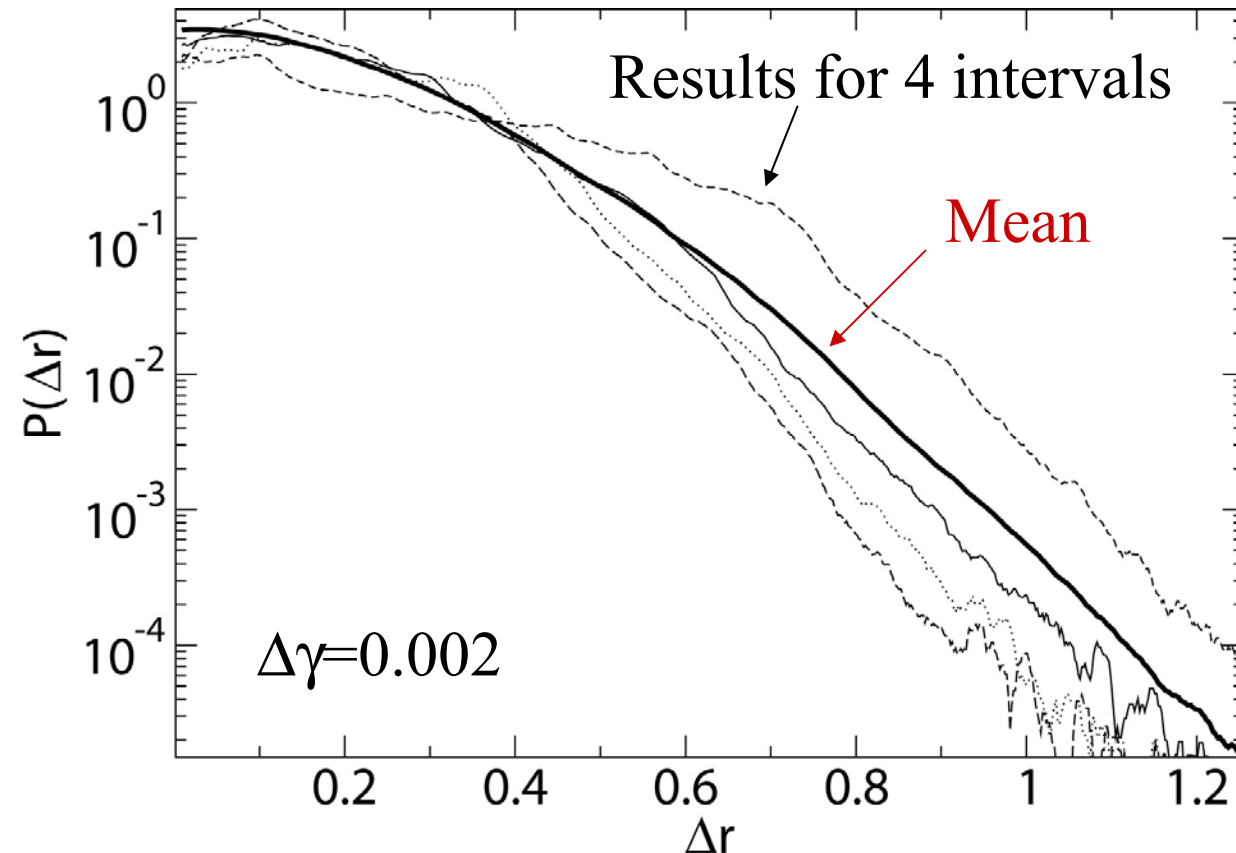
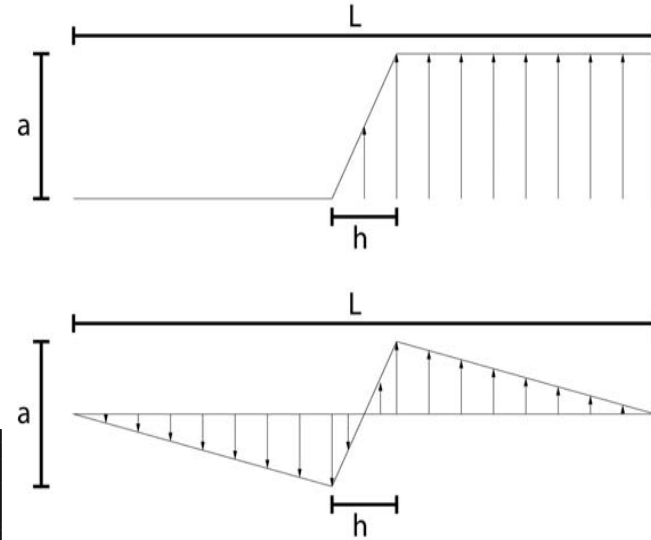


Distribution of Displacement Magnitude $\Delta r = \text{rms } u$

If displacement across band = a ,
flat distribution $P(\Delta r)$ up to $a/2$

Roughly consistent with observed P

Large fluctuations \rightarrow long-range correlations

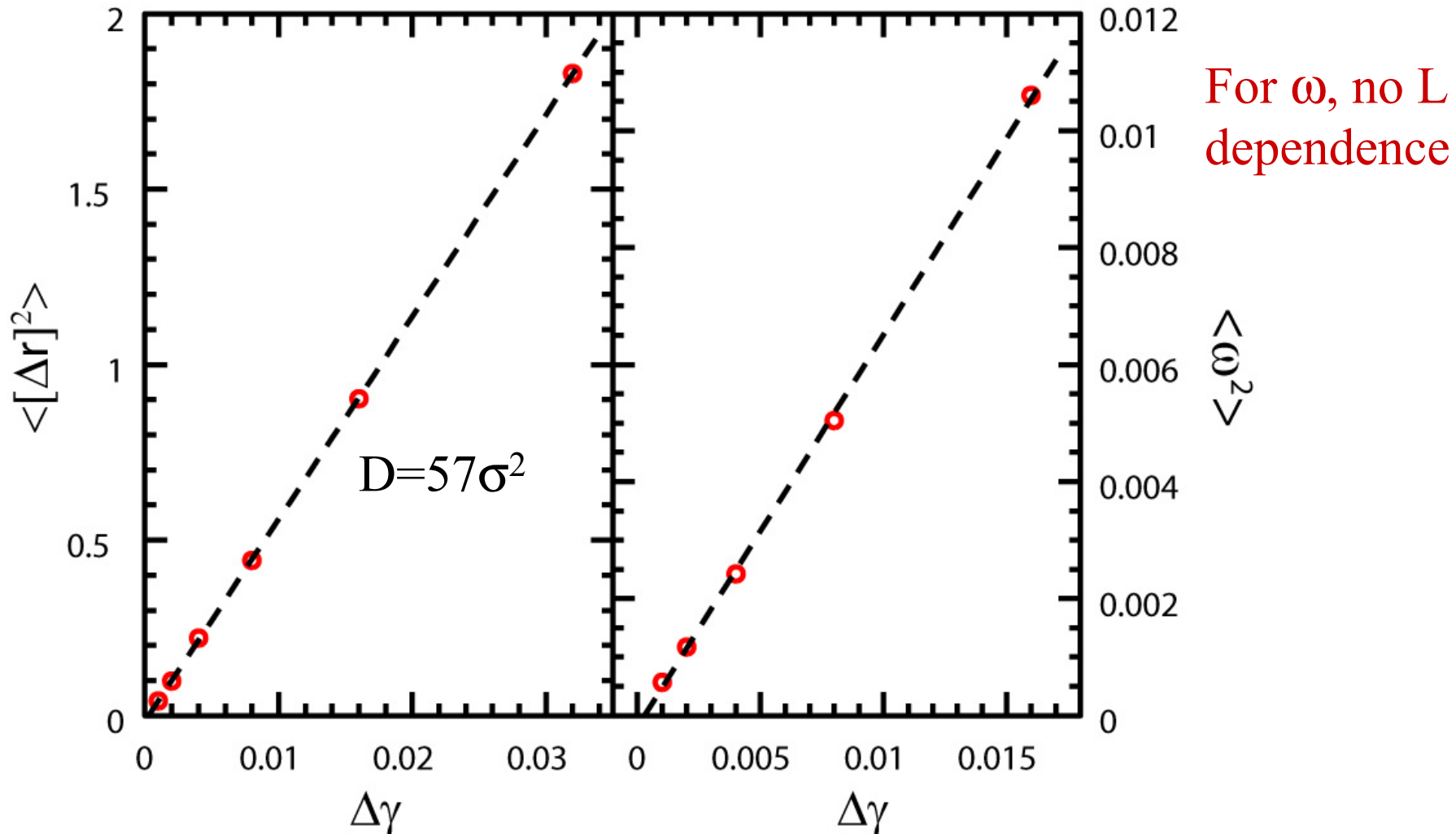


Displacements Diffusive with $D \propto L$

Plastic band formed in each strain $\sim a/L$ gives $|\Delta r| < a/2$, add incoherently

$$\Delta r^2 \sim \Delta\gamma / (a/L) \cdot a^2/12 \sim D \Delta\gamma \rightarrow D \sim L a/12$$

Consistent with observed $D = 57\sigma^2$ for $a = 0.7\sigma$ and for L down to 40



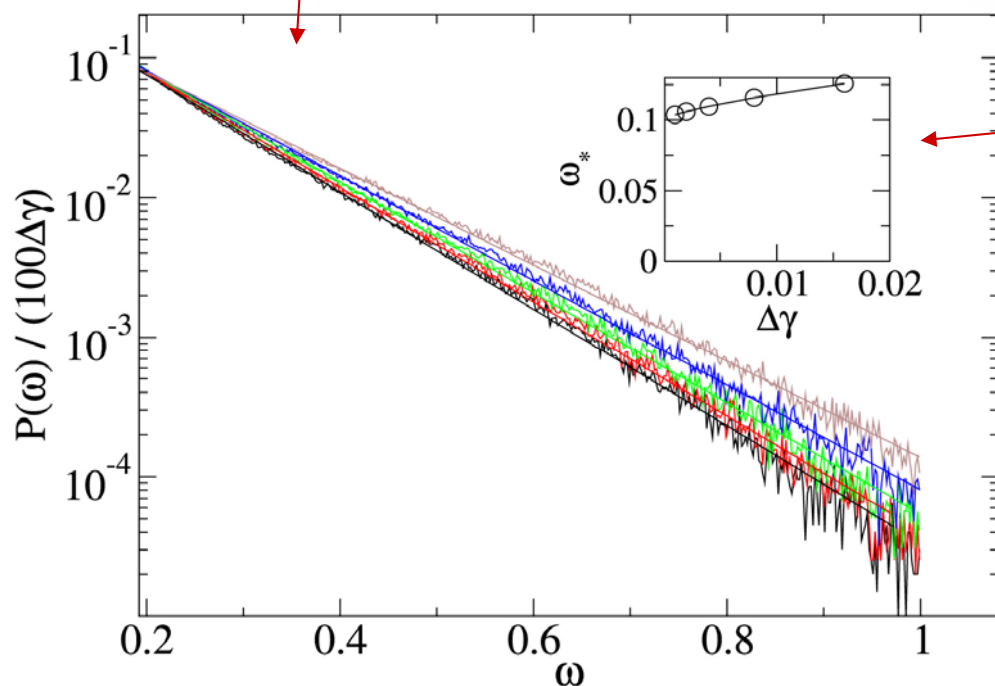
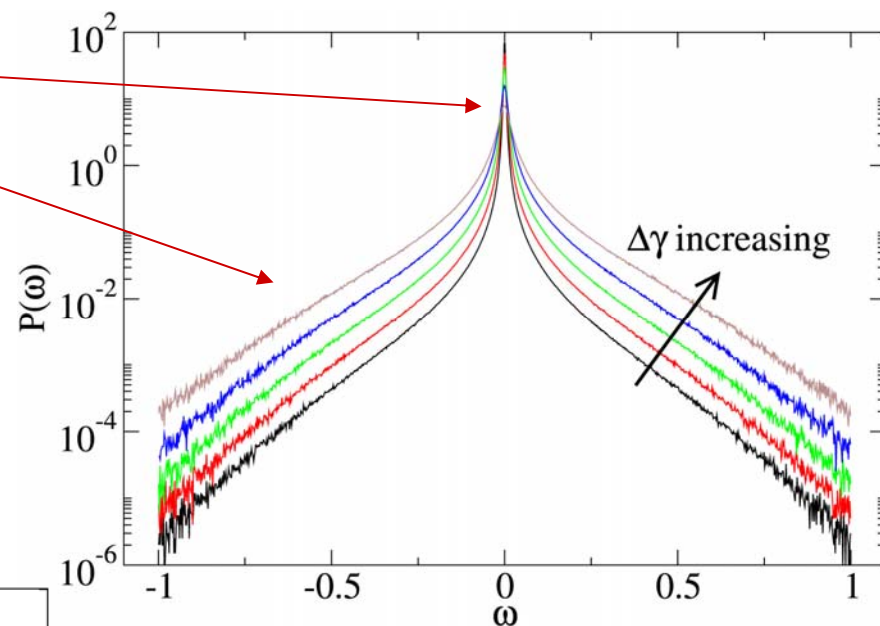
Distribution of Vorticity ω

Sharp elastic peak at small ω

Exponential tails at large ω

Weight in tails grows

\sim linearly with strain



Characteristic ω^* for decay ~ 0.1

Since $\omega \sim$ twice strain,
 $\omega^* \rightarrow 5\%$ strain
 \sim yield strain

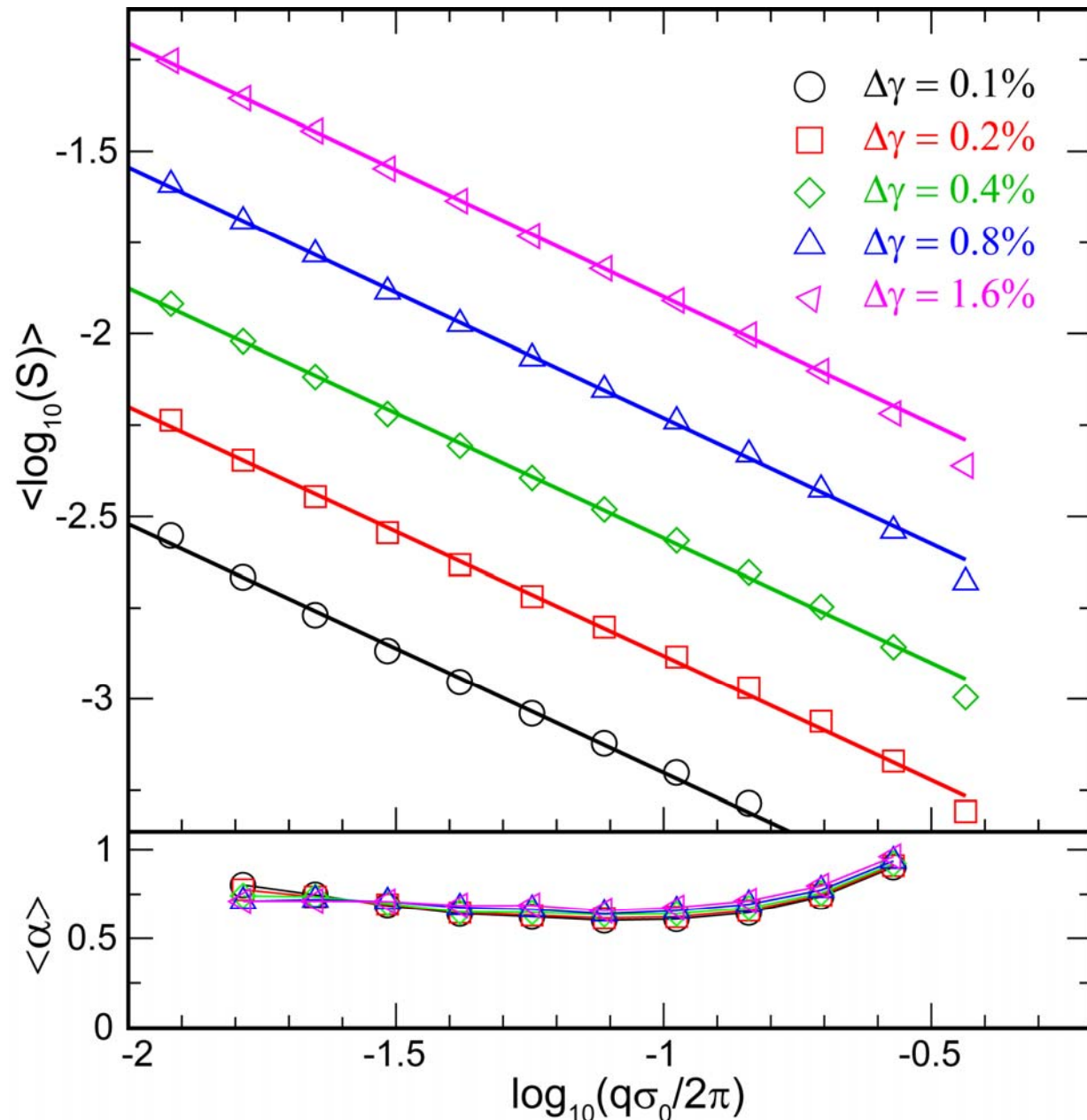
Plastic bands have ~ 10
sharper localized bands

Vorticity Correlation Function $S(\vec{q}) = \left| \int \omega(\vec{r}) \exp[i\vec{q} \cdot \vec{r}] d\vec{r} \right|^2$

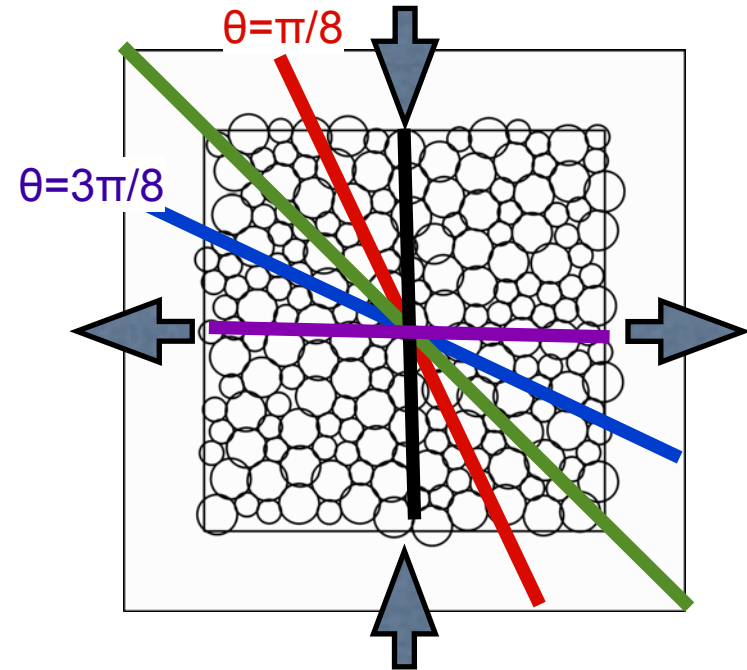
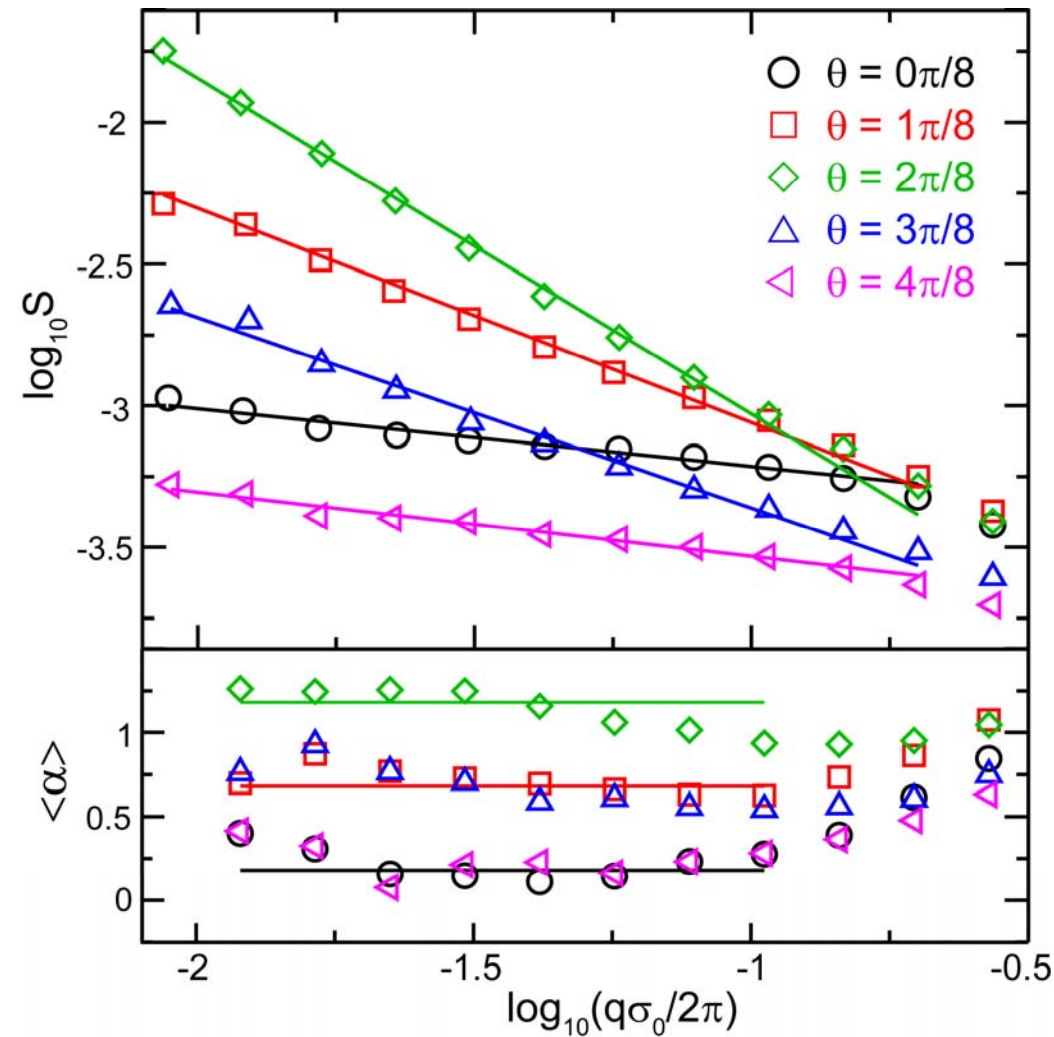
Mean of log S scales as
power of wave vector.

Prefactor linear in $\Delta\gamma$
→ Incoherent addition
of successive intervals

**BUT scaling highly
anisotropic**



Angle Dependence of Structure Factor, $\Delta\gamma=0.1\%$



$\theta = \pi/8$ and $\theta = 3\pi/8$ have same shear stress, different normal stress.
Mohr-Coulomb predicts shift to $\pi/8$

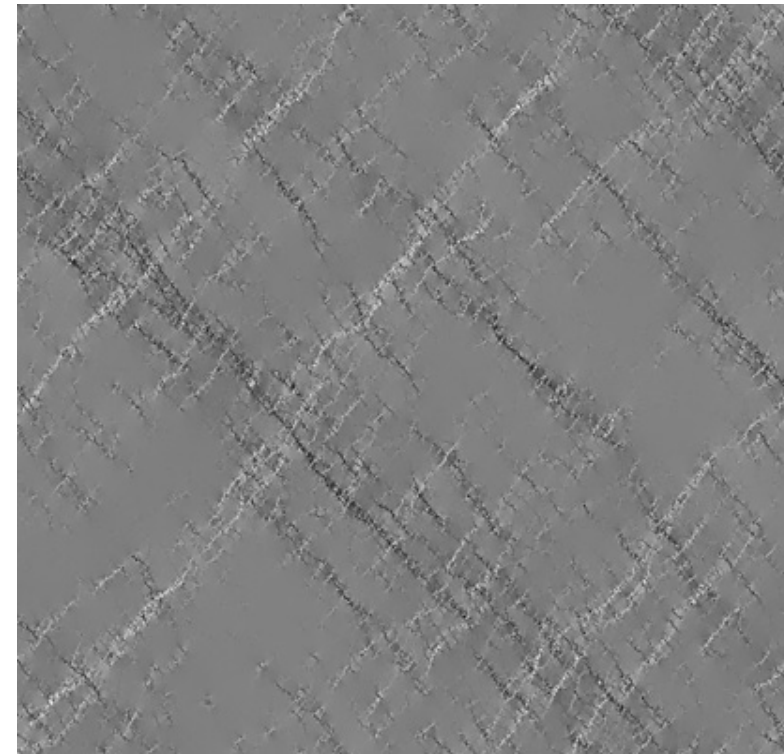
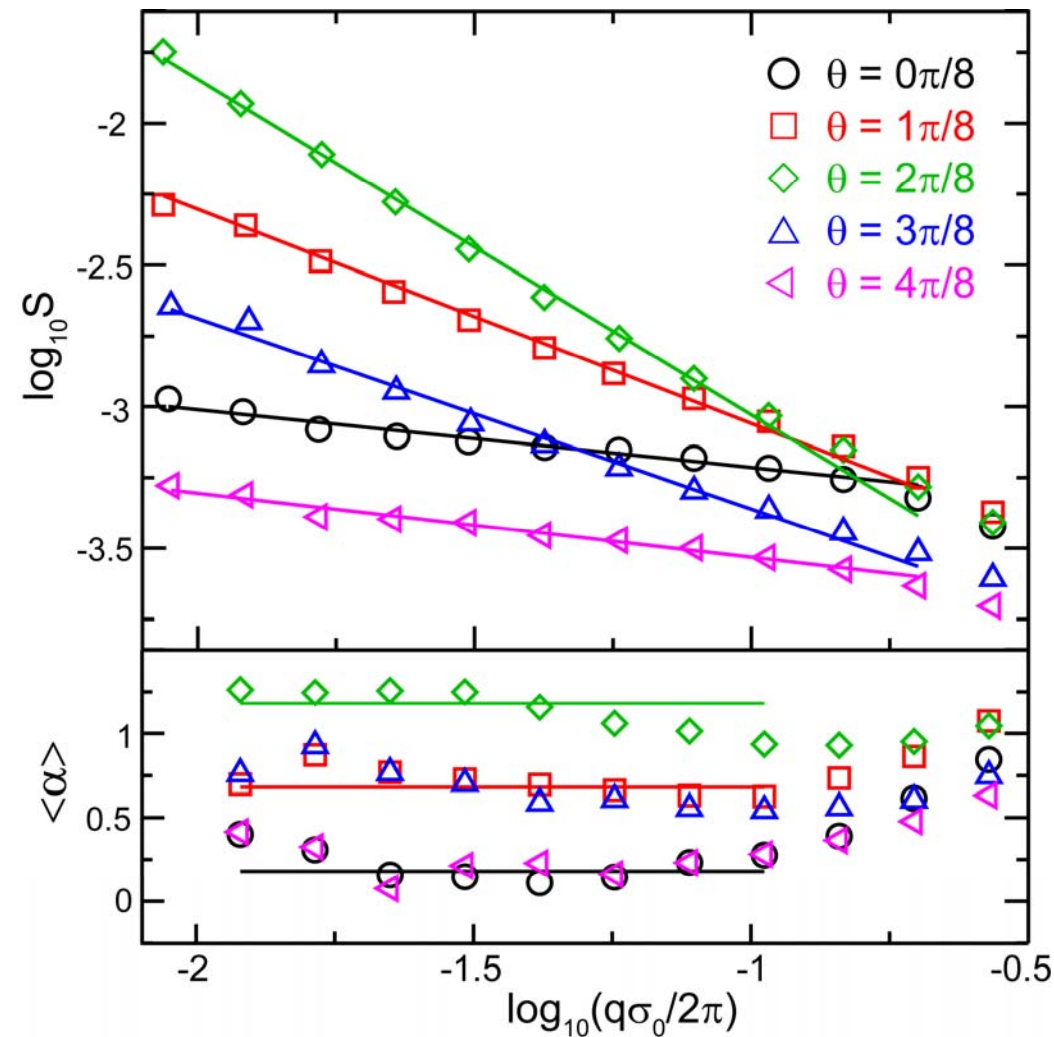
$$S(q;\theta) = A(\theta)q^{-\alpha(\theta)}$$

$$\alpha = a + b \cos(4\theta)$$

$$A = c + d \cos(2\theta)$$

A: broken shear symmetry
bigger for planes with low normal load
as predicted by Mohr-Coulomb

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Conclusions

- Earthquake probability $P(E) \sim 1/E$ over ~ 5 decades
Exponent independent of geometry, interactions
May depend on how events are broken up
- Inertia leads to qualitative changes in deformation statistics.
- Strain localizes in plastic bands that extend across system
Typical slip distance along band \sim particle diameter
Typical thickness h scales with system size $\sim L/20$
Several sharper features in h with strain $\sim 5 - 10\%$
- Non-affine displacement $\Delta r^2 \sim \Delta\gamma / (a/L) a^2/12 \sim D \Delta\gamma$, $D \sim La/12$
- Strain over short intervals has anisotropic power law correlations
 $S(|q|, \theta) \sim A(\theta) |q|^{-\alpha(\theta)}$ where $\alpha = a + b \cos(4\theta)$, $A = c + d \cos(2\theta)$
Breaking of symmetry for $A \rightarrow$ Mohr-Coulomb