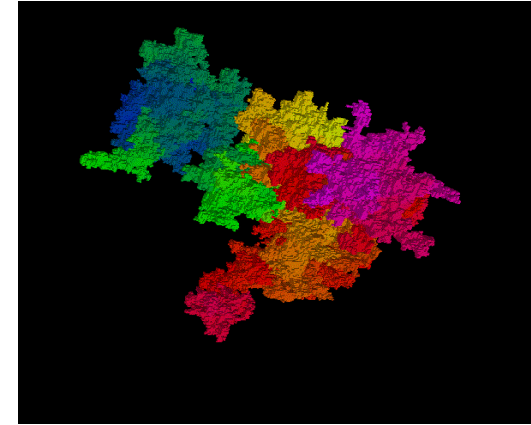


Crackling Noise, Glassiness, and Disorder Induced Critical Scaling In and Out of Equilibrium: Are they the same ???



1. Magnets:

Yang Liu, A. Mehta, J. Carpenter, S. Loverde, R. Vanderveldt, R. White, A. Travasset, *Nir Friedman*,
Braden Brinkman

Jim Sethna, M. Kuntz, O. Perkovic
(Cornell University), Gary Friedman
(Drexel.U.), Alan Middleton (Syr. U.)

Experiments:

A. Berger, O. Hellwig, Nanogune.
Gianfranco Durin (Torino, Italy).
A. Mills, Mike Weissman (UIUC),
Karen Daniels (NCSU),
Robert Behringer (Duke)

2. Plasticity, granular materials, earthquakes:

Y. Ben-Zion, J.T. Uhl,

DD Simulations: **G. Tsekenis**,
PFC: Pak Yuen Chan, Nigel Goldenfeld
superconductors:

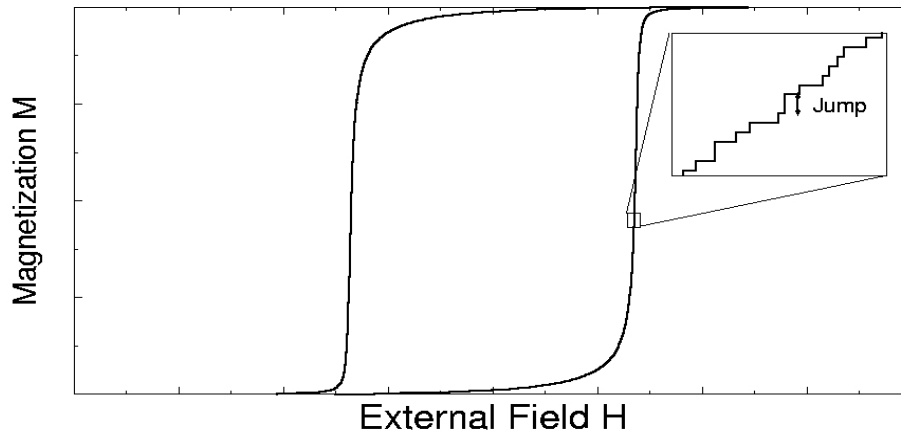
E. Carlson (Perdue), E. Fradkin (UIUC), S. Kivelson
(Stanford), D. van Harlingen, M. Weissman (UIUC)

Plastic CDW: C. Marchetti (Syr.)

Funding/Equipment:
NSF, MCC, SLOAN, UIUC, IBM, SANDIA

Magnets and Barkhausen Noise- or Martensites and acoustic emission

Hysteresis Loop Detail

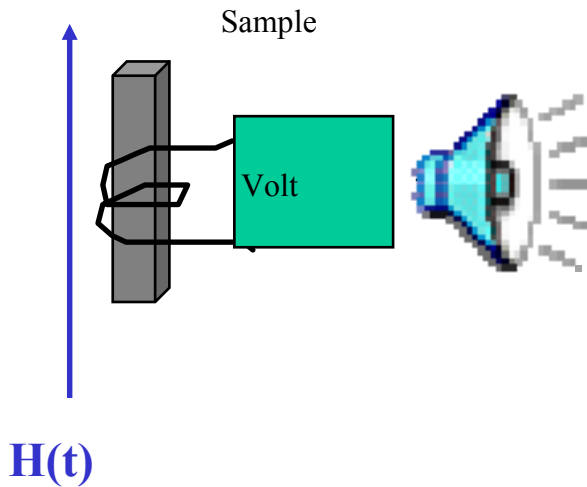
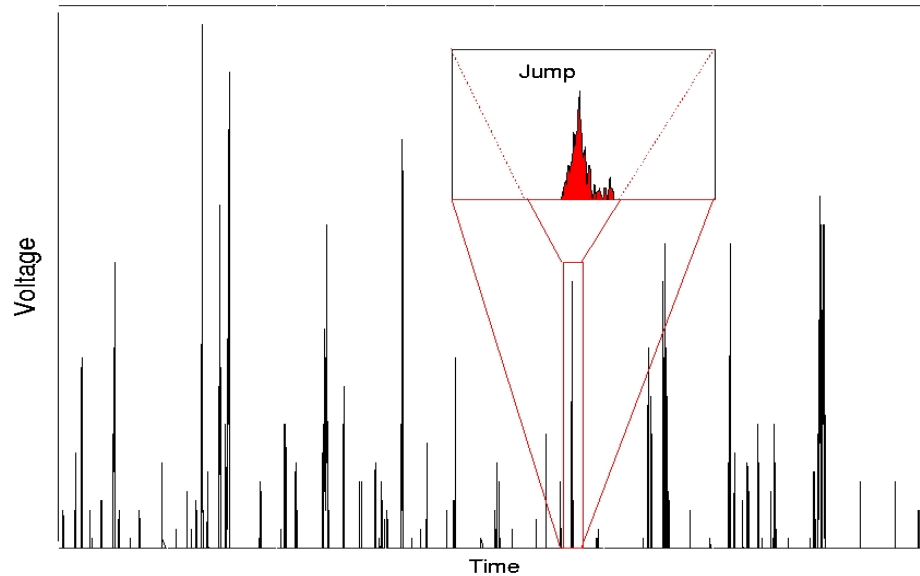


Crackling noise



Voltage vs Time

Barkhausen Jumps



Crackling Noise / Avalanches: Simple Models right!

- 1. Barkhausen Noise (magnets) (Bertotti, Durin, Weissman,...) 

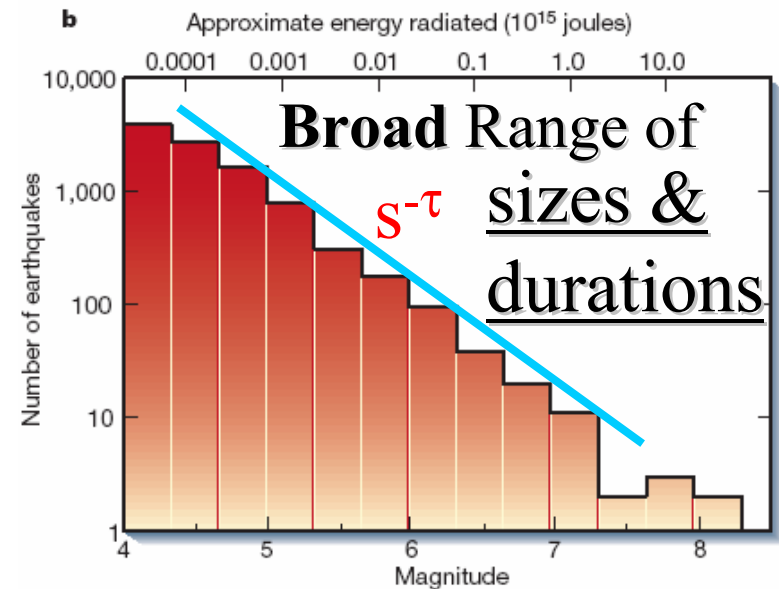
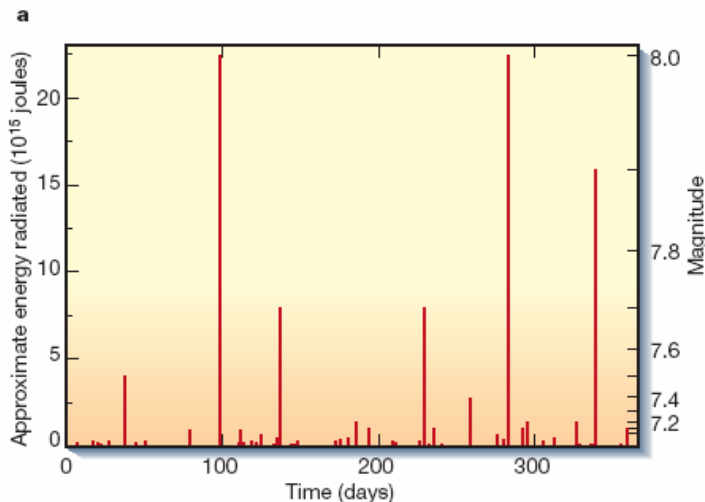
Acoustic emission (Martensites) (Ortin, Vives,...)

Liquids invading porous materials (Hallock, Lilly,..., Robbins...) 

- 2. Dislocation slips/Plasticity (Dimiduk, Weiss, Miguel, Zaiser, Zapperi,...)

- 3. Slips in Granular Materials (Behringer, Daniels,... Uhl)

- 4. Earthquakes (Ben-Zion) 



Simple (analytic) models work amazingly well!

Big open questions:

(1) Go beyond universal power laws? Universal scaling *functions* ?

Example: average temporal avalanche shape.

(2) What is the extent of the underlying universality classes ???

(3) What is the relation between crackling noise and glassiness?

Glasses: stuck in (thermal) metastable states

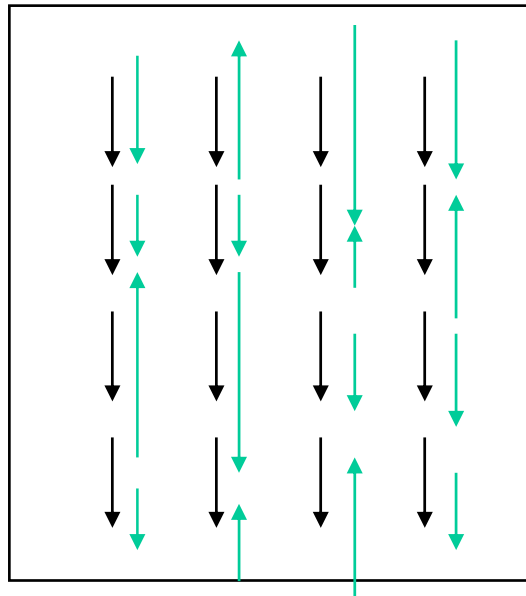
Crackling noise: (driven) transitions between metastable states

} SAME?

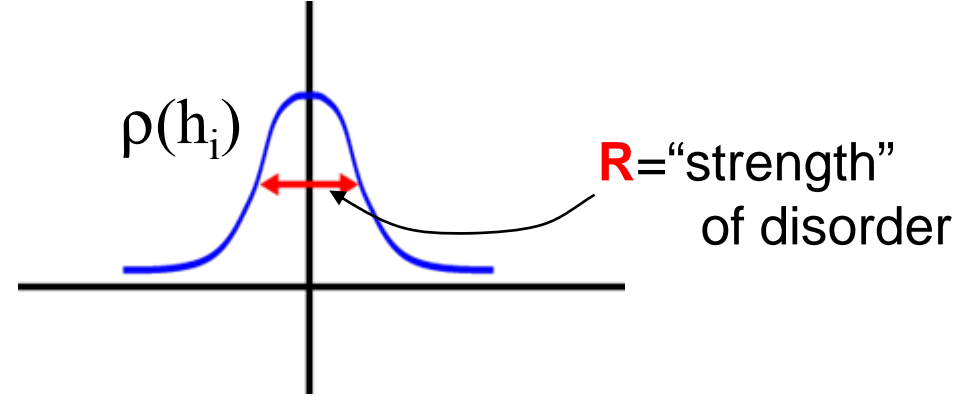
For random-field Ising model: **same scaling behavior (universality class)**

(4) Experimental tests with disorder as a tuning parameter ?!!!!

1. MAGNETS: The Zero Temperature Non-equilibrium Random Field Ising Model



Random Field Distribution

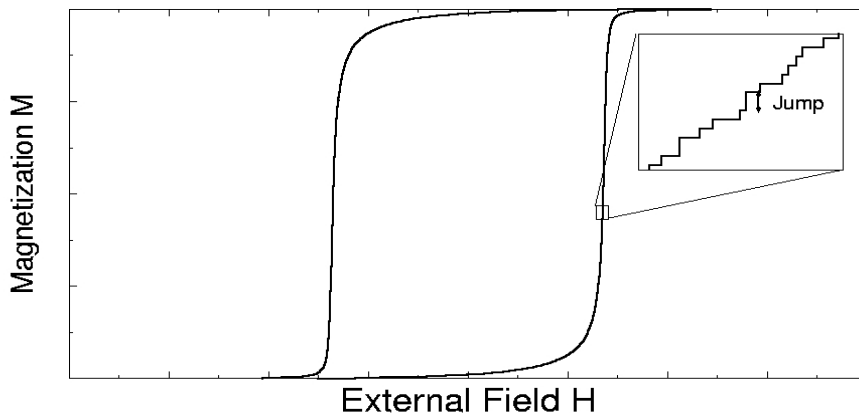


Each spin is always aligned with the direction of the

$$\text{local "force"} = H(t) + h_i + J \sum_{n.n.} S_j$$

Zero Temperature : (Equilibration time scale) \gg (Experimental time scale)

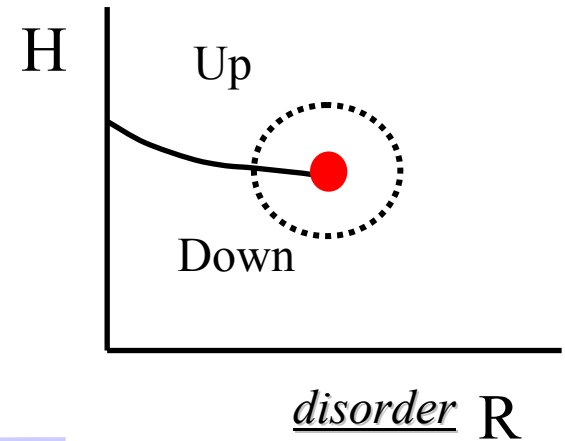
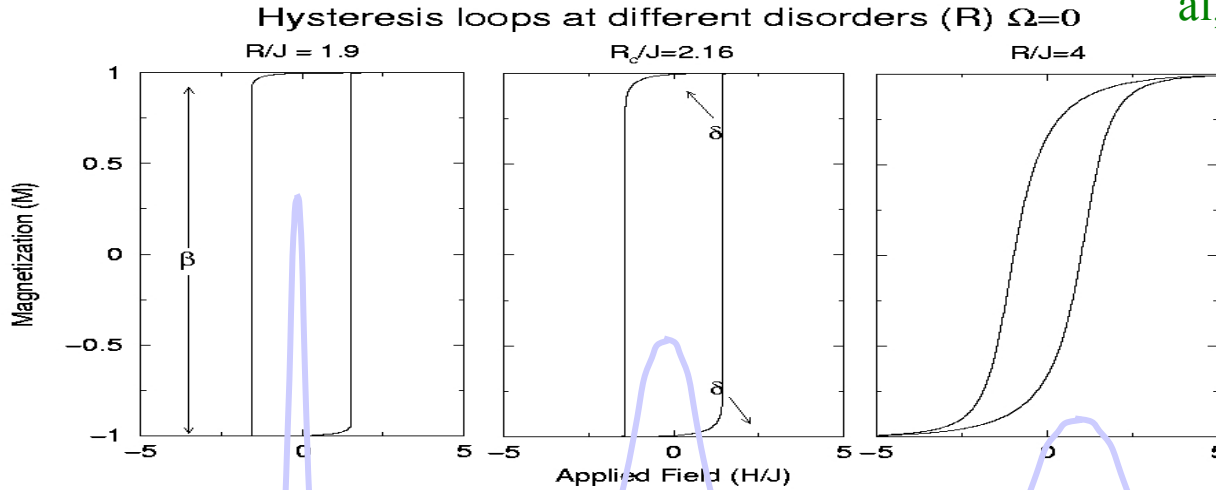
Hysteresis Loop Detail



Jumps = Avalanches =
Barkhausen noise

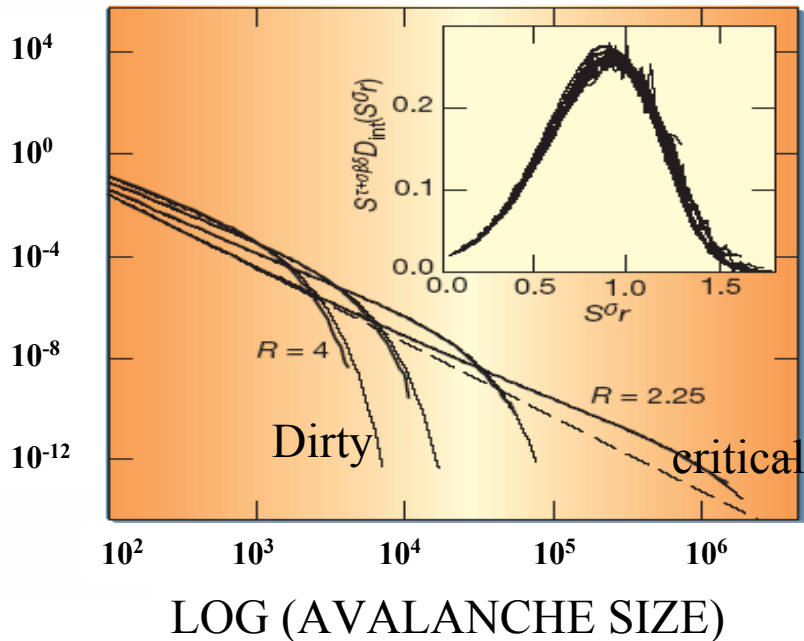
The Disorder Induced Critical Point

Seen also in experiments: Berger et al, PRL 2000, J.Appl.Phys. 2004



Clean critical Dirty

Log(#)



$$\tau + \sigma\beta\delta = 2.03 \pm 0.03, \quad 1/\sigma = 4.2 \pm 0.3, \dots$$

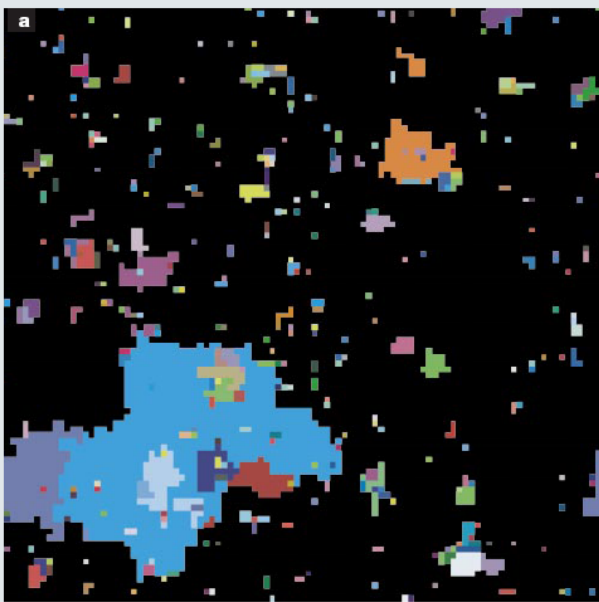
50% away from critical amount:
still find 2 decades of powerlaw
scaling!

... HUGE scaling region!!!



Self-similarity at critical disorder $R_c = 2.16J$

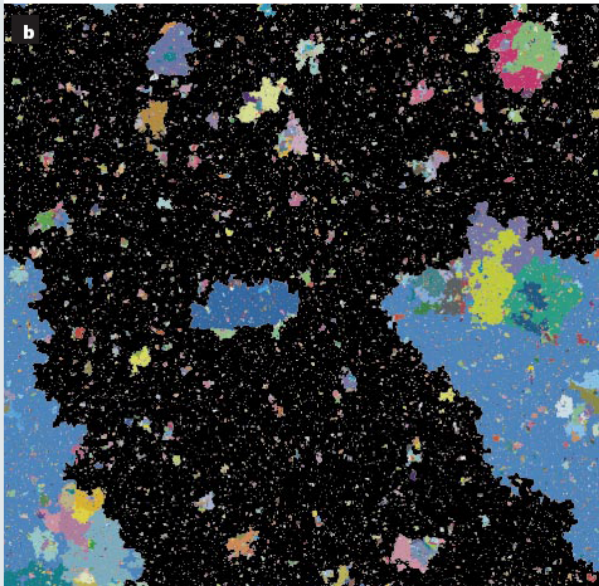
(Cross-sections of avalanches during magnetization)



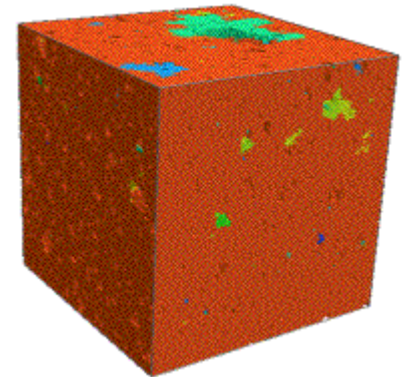
100^3

CRITICAL POINT:
system is at a fixed point
under coarse graining
transformation

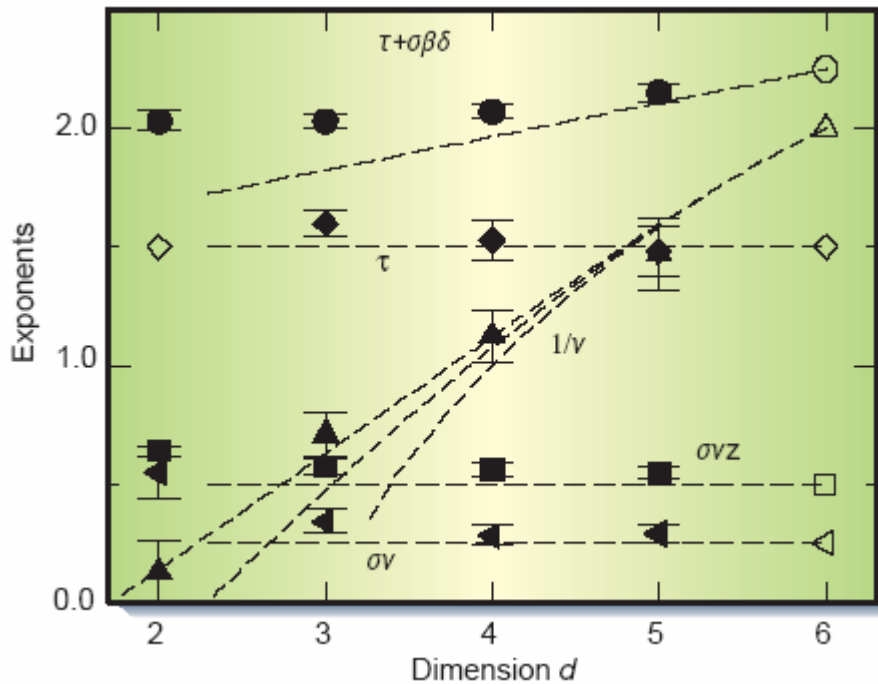
(Renormalization Group)



1000^3



RESULTS



Simulation



6- ϵ expansion

(Renormalization Group)

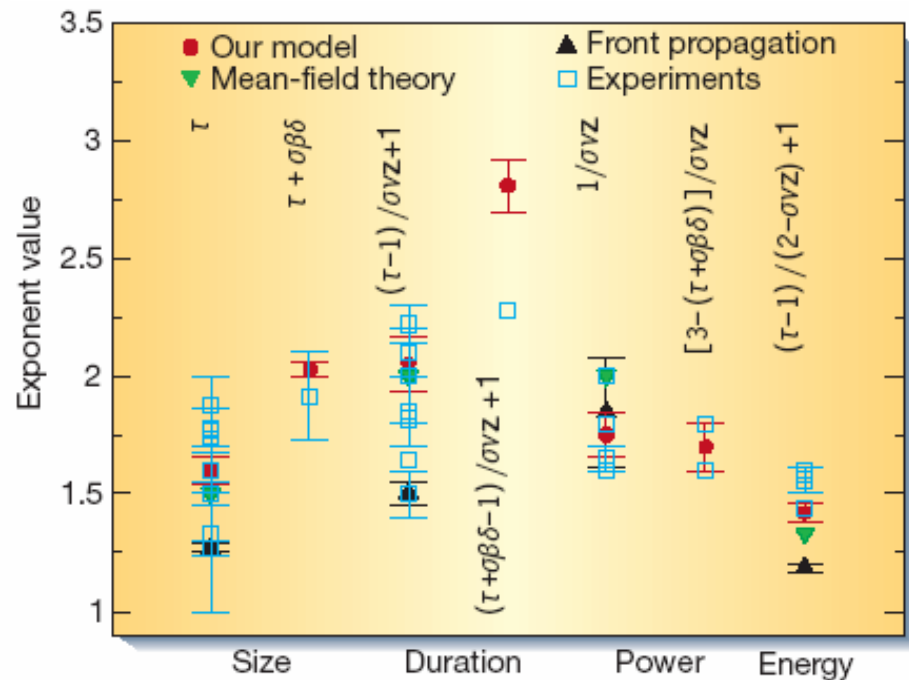
$$1/\nu = 2 - \epsilon/3 - 0.1\epsilon^2 + 0.1\epsilon^3 - 0.3\epsilon^4 + \epsilon^5 + O(\epsilon^6)$$

(PRL '93, '95, 2003 PRB '96, '99, 2002 (R))

Nature **410**, 242 (2001)

Experiments and Simulations in 3 dim. (Barkhausen Noise):

Need Noise Experiments Tuning disorder!!!



Big open questions:

(1) Go beyond universal power laws? Universal scaling *functions* ?

Example: average temporal avalanche shape.

(2) What is the extent of the underlying universality classes ???

(3) What is the relation between crackling noise and glassiness?

Glasses: stuck in (thermal) metastable states

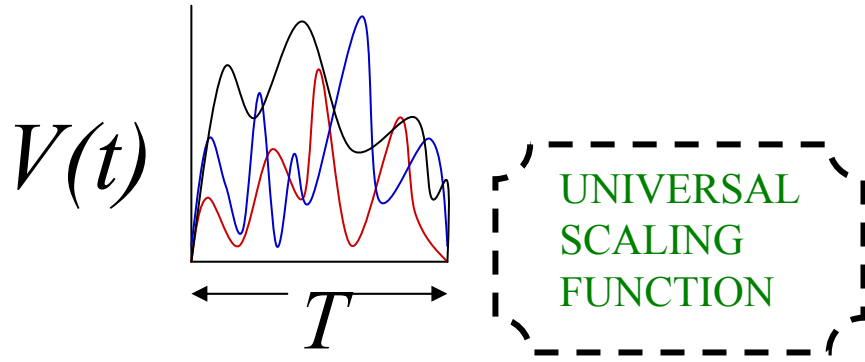
Crackling noise: (driven) transitions between metastable states

} SAME?

For random-field Ising model: same universal fixed point!

(4) Experimental tests with disorder as a tuning parameter ?!!!!

Universal Scaling Functions:



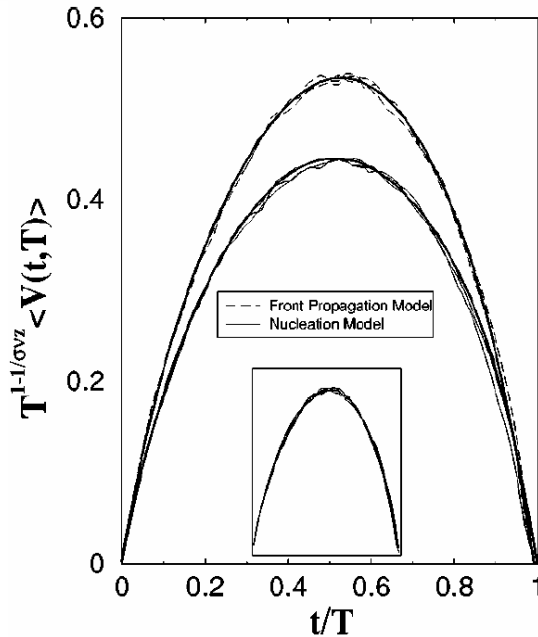
AVALANCHE SIZE

DISTRIBUTION:

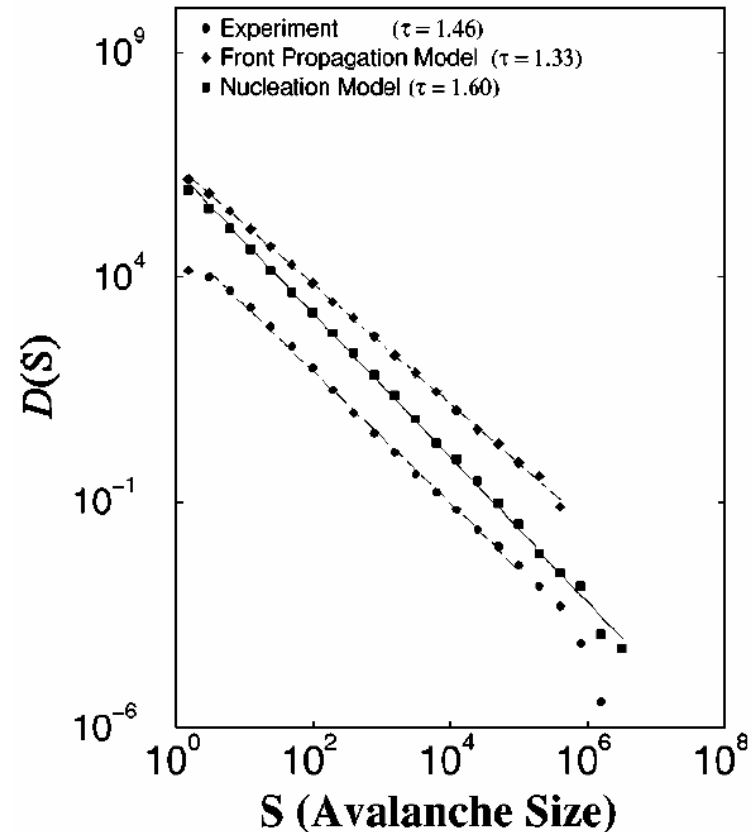
$$D(S) \sim S^{-\tau} \text{ at criticality.}$$

AVERAGE :

$$V(T, t) = T^{1/\sigma\nu z - 1} f_{shape}(t/T)$$

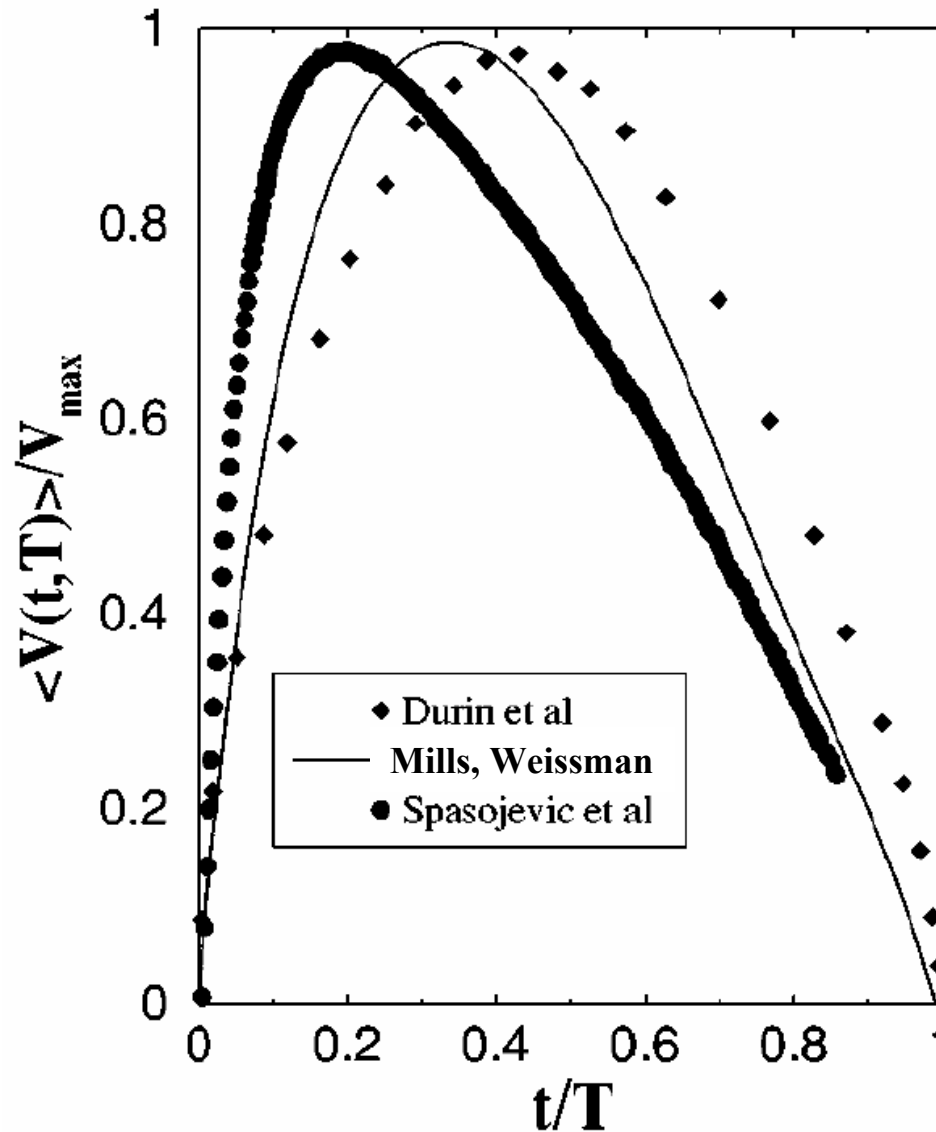


Mehta, KD,
Sethna



Kuntz
Sethna

EXPERIMENTAL SCALING FUNCTION



Asymmetric!?!

Eddy currents !

Zapperi, Durin et al.
Nature Physics.
1, 46-49 (2005);
K.D, Nature Physics
1, 13-14 (2005)

**Papanikolaou
Bohn, Sommer,
Durin, Zapperi,
Sethna (2009)**

Big open questions:

(1) Go beyond universal power laws? Universal scaling *functions* ?
Example: average temporal avalanche shape.

(2) What is the extent of the underlying universality class of this simple model ??? (I.e. class of systems with the same exponents)?

(3) What is the relation between crackling noise and glassiness?
Glasses: stuck in (thermal) metastable states

} SAME?

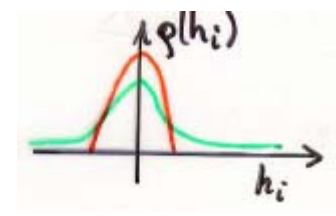
Crackling noise: (driven) transitions between metastable states

For random-field Ising model: same universal fixed point!

(4) Experimental tests with disorder as a tuning parameter ?!!!!

Huge Universality Class!!! (Details don't matter!)

$$\mathbf{H}_{system} = -J \sum_{n.n.} S_i S_j - \sum_i (H(t) + h_i) S_i$$



Replace J by random couplings

Use different distribution of h_i or replace h_i by random anisotropies

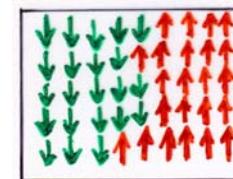
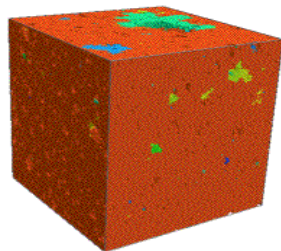
Magnets (Sethna, KD, Myers, Nature 2001), plastic charge density wave depinning (Marchetti, KD PRB 2002), maybe superconductors (Carlson, KD, Fradkin, Kivelson, 2005), same MFT: earthquakes (Fisher, KD, Ramanathan, Ben-Zion, 1997, Mehta, Ben-Zion, KD 2005), plasticity (Zaiser 2006, Miguel, Zapperi, 07, KD, Ben-Zion, Uhl 09), granular, others ?

Other Universality classes ?

$$\mathbf{H}_{system} = -J \sum_{n.n.} S_i S_j - \sum_i (H(t) + h_i) S_i + \text{long range forces}$$

2 Dynamics

With nucleation of new domains



Single domain wall.

Robbins, Ji, ...

Nattermann, Zapperi, Ciseau, Durin, Stanley, Urbach et al., Narayan, Sethna, ...

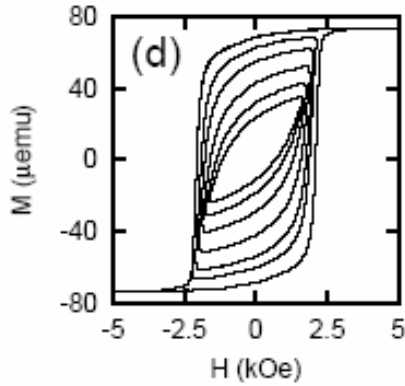
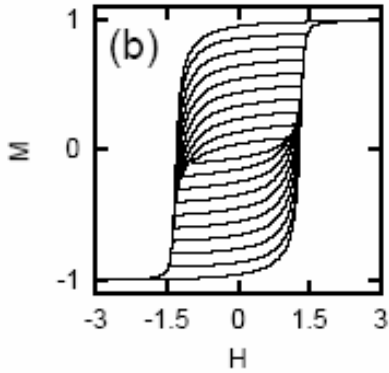
Test whether Nucleation or Domain Wall Motion?

Nucleation: avalanches (cutoff) *smaller* for inside loops!

Carpenter,
KD, 2004

Simulation

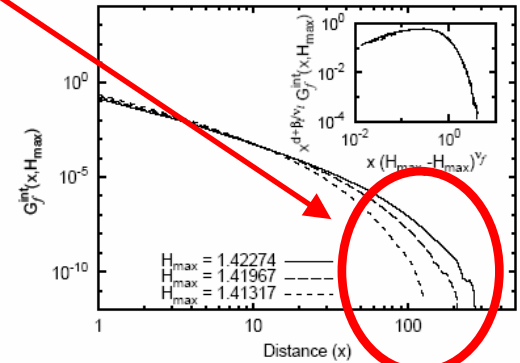
Experiment



Andreas Berger
(NanoGune, Spain)

CoPtCrB

Simulation

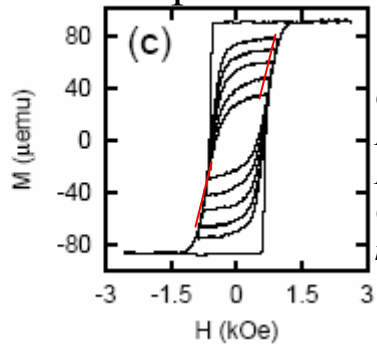
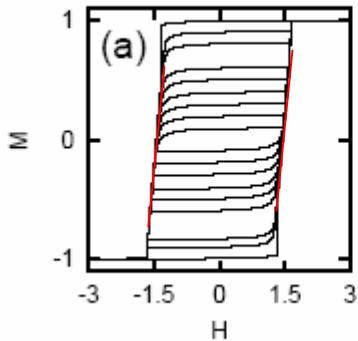


Single domain wall motion: *same* cutoff for all loops!

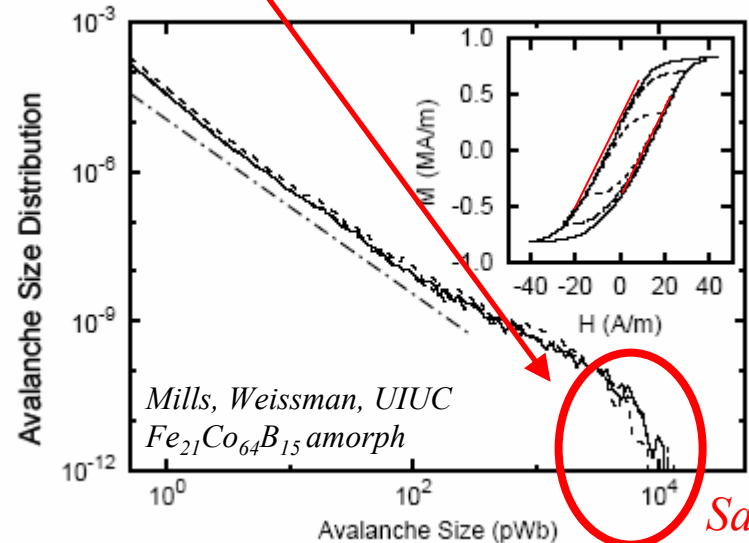
Experiment

Simulation

Experiment



Olav
Hellwig,
Hitachi
Co/Pt
multilayer



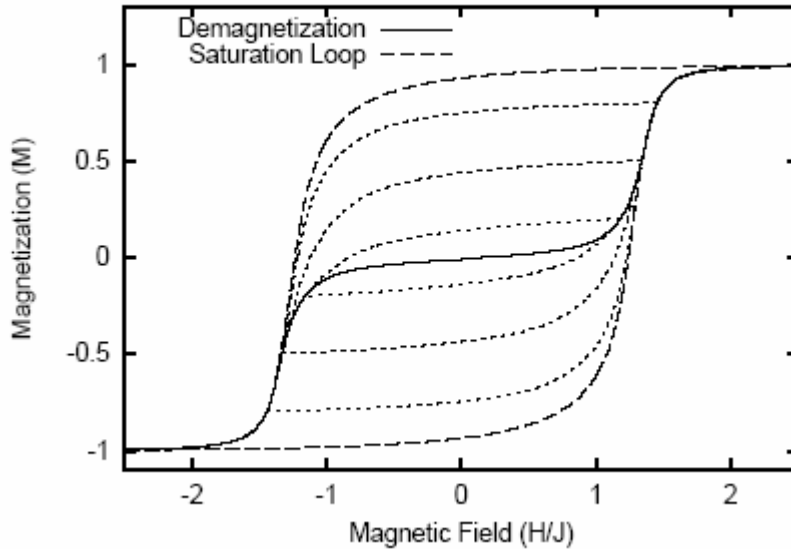
Same
Cutoff

Demagnetizing curves in the T=0 Random Field Ising Model:

also SAME scaling behavior!

John Carpenter, KD, PRB 2003 (R)

Zapperi et al., PRB 2004

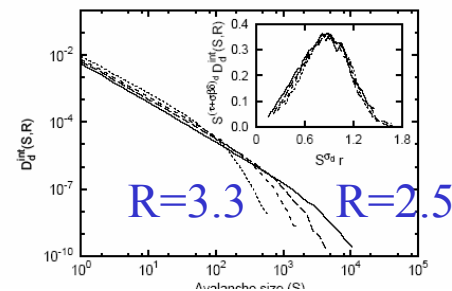


Same disorder induced critical point as saturation loop:

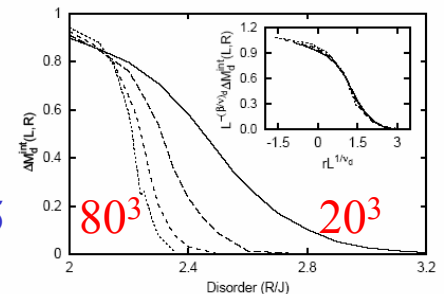
Exponent	Demagnetization	Saturation ^a
$\tau + \sigma\beta\delta$	2.10 ± 0.05	2.03 ± 0.03
$1/\sigma$	3.9 ± 0.4	4.2 ± 0.3
$d + \beta/\nu$	3.1 ± 0.2	3.07 ± 0.30
$1/\nu$	0.71 ± 0.10	0.71 ± 0.09
θ	0.01 ± 0.01	0.015 ± 0.015
β/ν	0.03 ± 0.02	0.025 ± 0.020

Need: Experiments on Barkhausen Noise (for various disorders) ???

Avalanche Size Distribution



ΔM vs R



Big open questions:

(1) Go beyond universal power laws? Universal scaling *functions* ?
Example: average temporal avalanche shape.

(2) What is the extent of the underlying universality classes ???

(3) What is the relation between crackling noise and glassiness?

Glasses: stuck in (thermal) metastable states

Crackling noise: (driven) transitions between metastable states

} SAME?

For random-field Ising model: **same universal fixed point!**

(4) Experimental tests with disorder as a tuning parameter ?!!!!

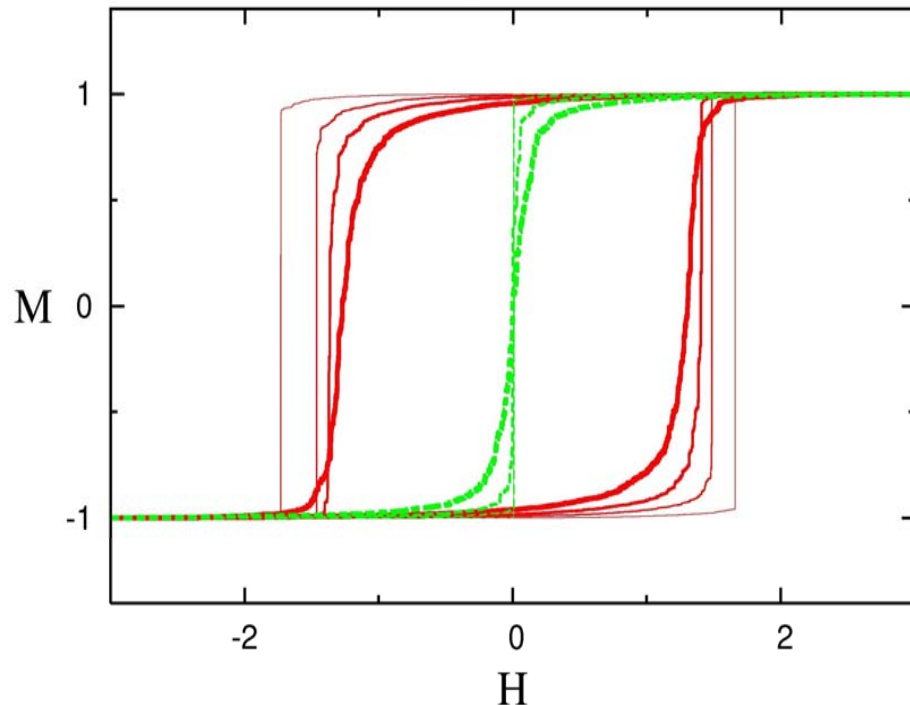
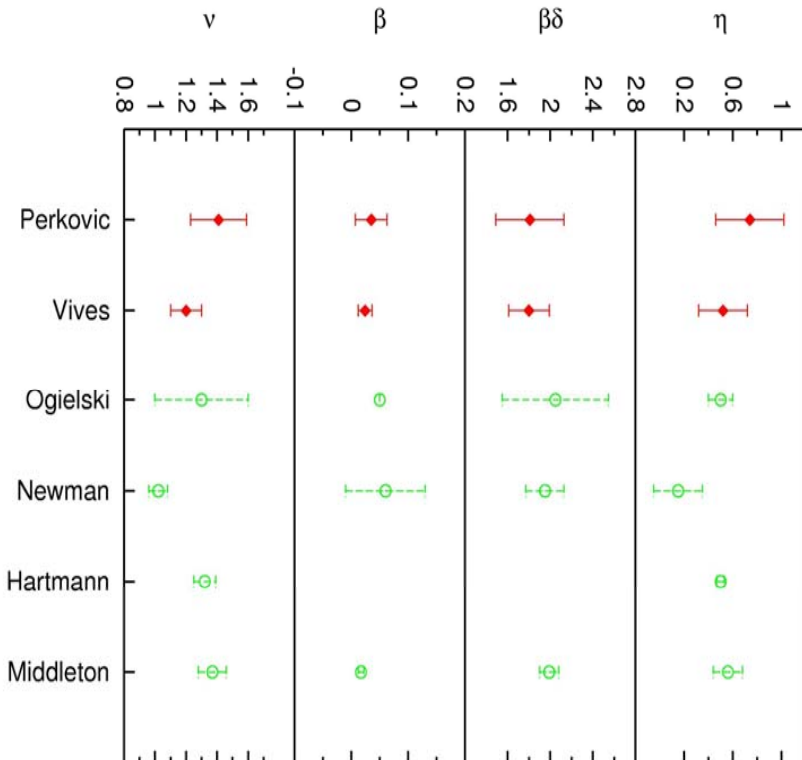
EQUILIBRIUM and NONEQUILIBRIUM RFIM:

same scaling behavior !!

(Maritan, Banavar, and Vives, Ortin, Perez Reche et al., and YANG LIU, KD):

1. SAME MEAN FIELD EXPONENTS ($\beta=1/2$, $\nu=1/2$, $\delta=3$) $\tau=3/2$ and $\sigma=1/2$ (Liu, KD)
2. ABOUT SAME SIMULATION EXPONENTS in 3D and even in 4D
3. $6-\epsilon$ expansion of noneq. mapped to all orders in ϵ to eq. expansion (KD, Sethna)

d spin can't flip back with increasing field. (Liu, KD)

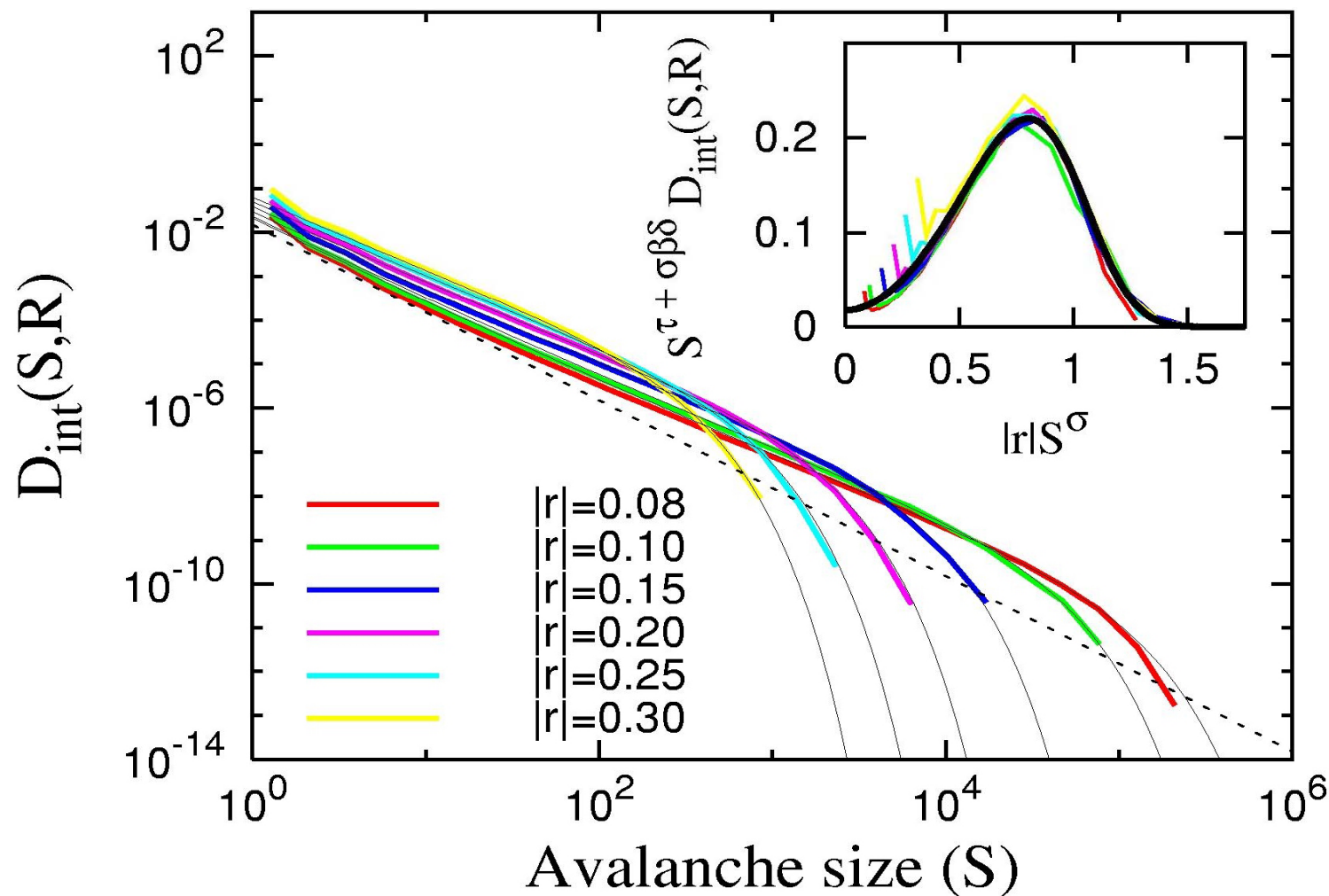


Surprising Result: same Avalanche Exponents and Scaling Function

In Equilibrium (no hysteresis) and Nonequilibrium (hysteresis):

Ground state
"Avalanches":

$$D_{\text{int}}(S, R) \sim S^{-(\tau + \sigma\beta\delta)} \bar{D}_{\pm}^{\text{int}}(S^{\sigma} |r|)$$



- Many other quantities: equilibrium and nonequilibrium

Quantities	non-equilibrium		equilibrium	
	Avalanches	Clusters	Avalanches	Clusters
d_S	2.78 ± 0.05	2.76 ± 0.04	2.77 ± 0.09	2.78 ± 0.05
d_V	2.78 ± 0.05	2.76 ± 0.04	2.77 ± 0.09	2.78 ± 0.05
d_a	2.33 ± 0.04	2.18 ± 0.04	2.16 ± 0.05	2.11 ± 0.03
A_1	0.29 ± 0.01	0.25 ± 0.01	0.30 ± 0.02	0.28 ± 0.01
A_2	0.50 ± 0.02	0.45 ± 0.02	0.50 ± 0.02	0.48 ± 0.02
Δ_D	0.16 ± 0.01	0.21 ± 0.02	0.16 ± 0.02	0.18 ± 0.02
S_D	0.06 ± 0.01	0.09 ± 0.01	0.06 ± 0.01	0.07 ± 0.01

Avalanche distributions and spatial structures (fractal dimensions and anisotropy measures), etc. all scale the same !!

=> equilibrium and non-equilibrium transitions of T=0 RFIM: same universality class... !!!???

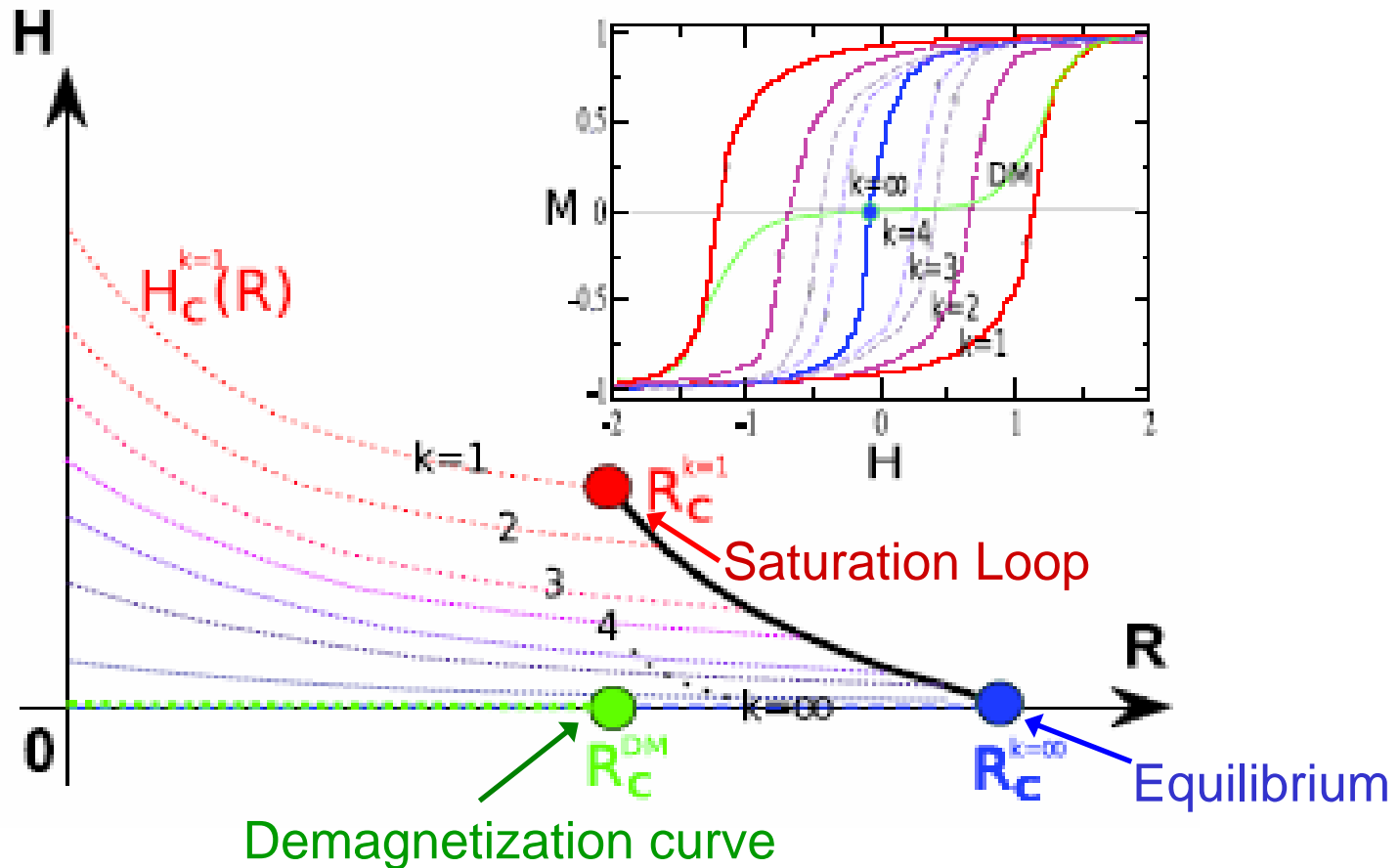


Magnets: RG and simulation, 3 equivalent critical points

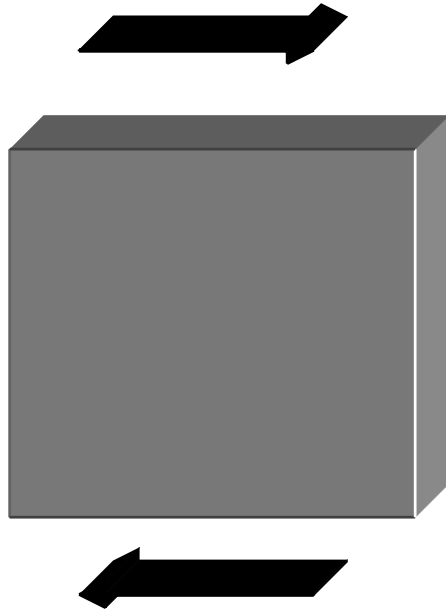
Equilibrium, Saturation Loop, Demagnetization curve

(k-spin flip dynamics may connect nonequilibrium and equilibrium critical points?)

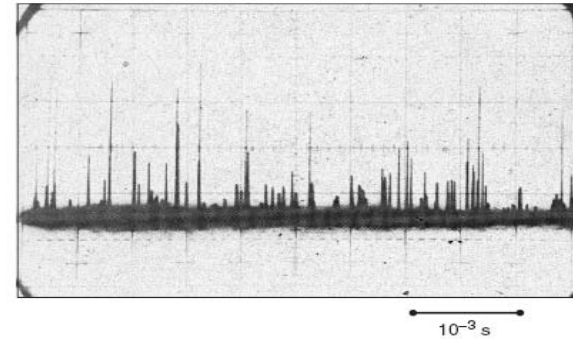
(Vives, Stein, Liu, ...)



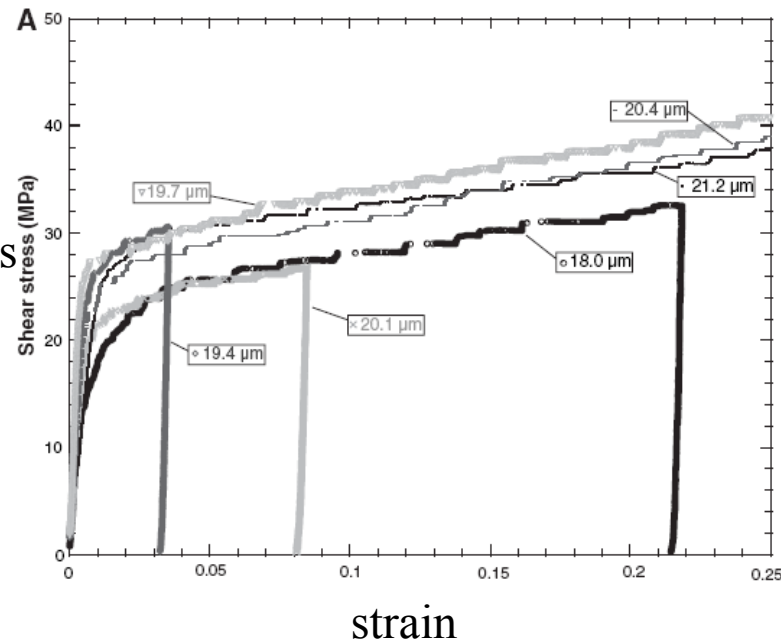
2. Slowly sheared metals: plastic response with **discrete jumps** !



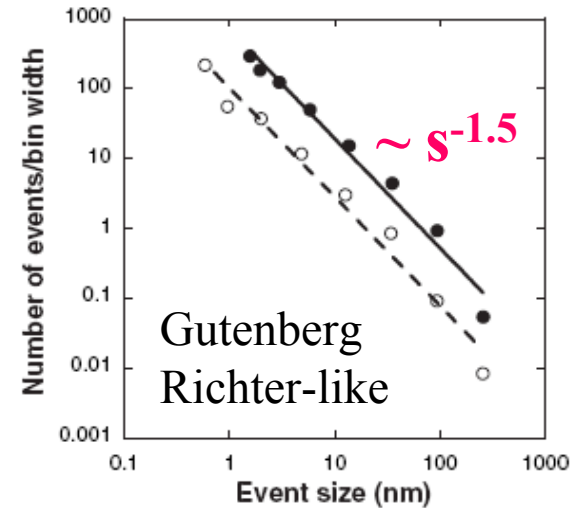
Acoustic emission



Imanaka et al.



jump size distribution



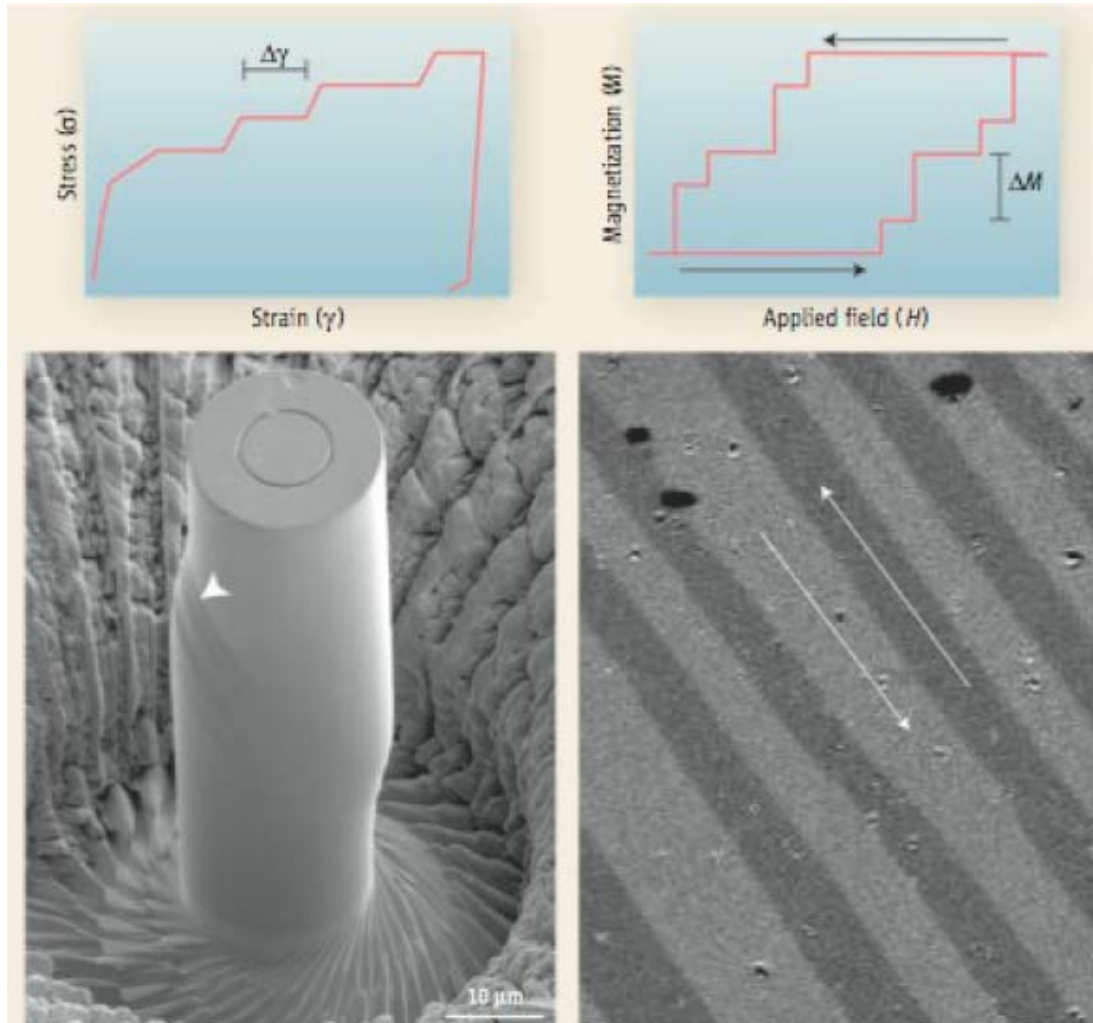
Universal!

(Dimiduk, et al
Science 2006)

Resemblance to Barkhausen Noise:

Stress-Strain

Magnetization $M(H)$



Simulations of slip avalanches in sheared crystals:

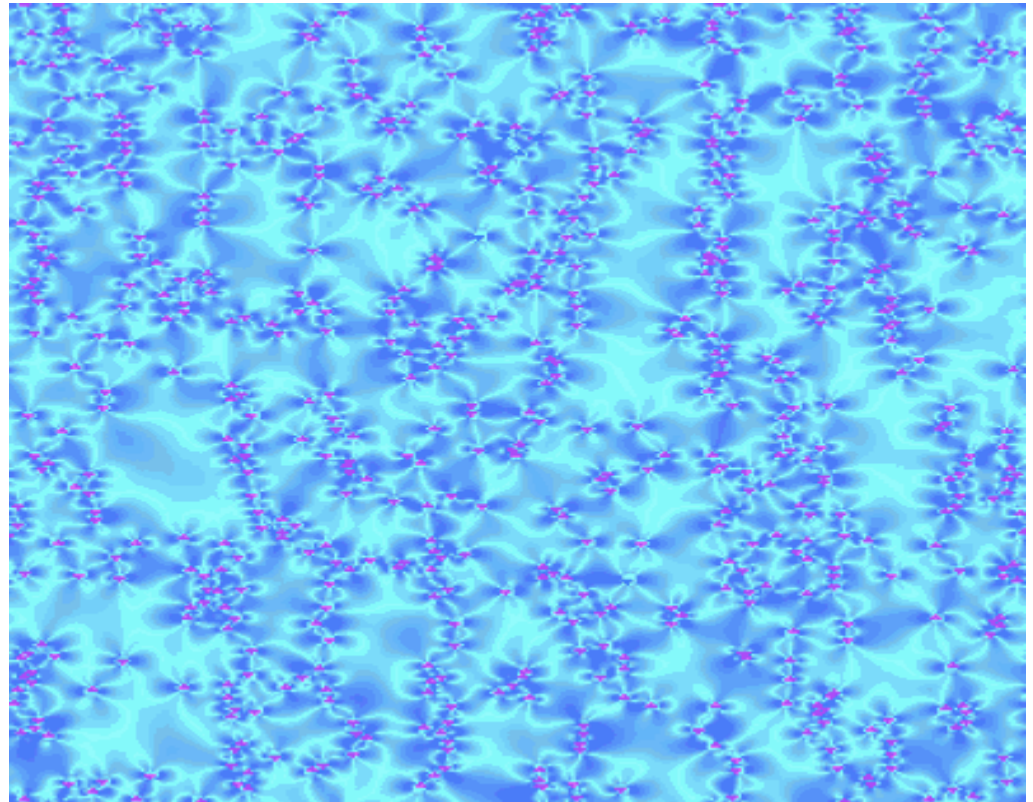
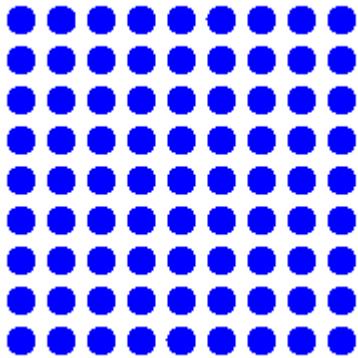
Dislocation Dynamics: Miguel et al Nature 410 (667-671) 2001, Csikor et al. Science (2007)

Georgios Tsekenis 2010;

Phase Field Crystal Models: Pak Yuen Chan, Georgios Tsekenis, J. Dantzig, KD, N. Goldenfeld (PRL 2010);

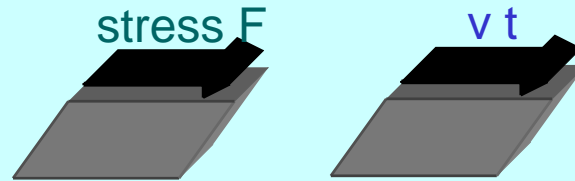
phase field models: M. Koslowski, ...

Crystal with many edge dislocations under shear (Miguel et al. Nature 2001)



Simple Analytic Model for deformation under shear:

- Simple **discrete** model (KD, Y. Ben-Zion, J.T. Uhl, PRL 2009)
- with **only 1 tuning parameter** (**weakening ε**) –
- yields ***exact analytical* scaling predictions**
- for **2** different boundary conditions:
fixed applied **stress** **OR** fixed **boundary velocity**



- Results:
- ***same* scaling exponents/functions as magnet model, (because *same* mean field theory).**
- ***same* mean field theory as for earthquakes**

stress

Dynamic equation for slip evolution in heterogeneous medium:

$$\eta \partial u(\mathbf{r}, t) / \partial t = F + \sigma_{\text{int}}(\mathbf{r}, t) - f_R[u, \mathbf{r}, \text{history}]$$

Slip velocity \sim stress + interaction + Pinning due to heterogeneities

Failure stress

Weakened failure stress

Arrest stress

ϵ

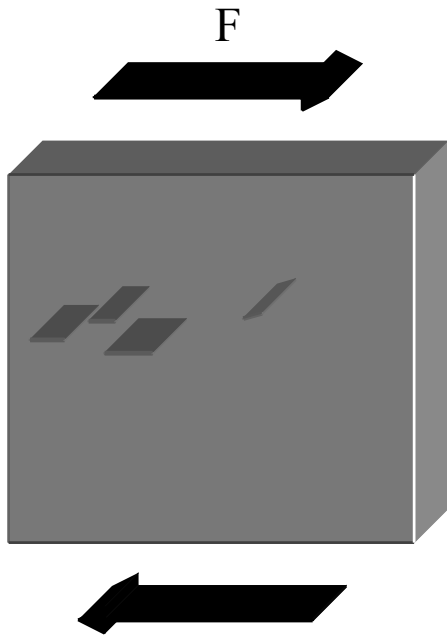
interaction:

$$\sigma_{\text{int}}(\mathbf{r}, t) = \int_{-\infty}^t dt' \int d^d r' J(\mathbf{r}-\mathbf{r}', t-t') \times [u(\mathbf{r}', t') - u(\mathbf{r}, t)]$$

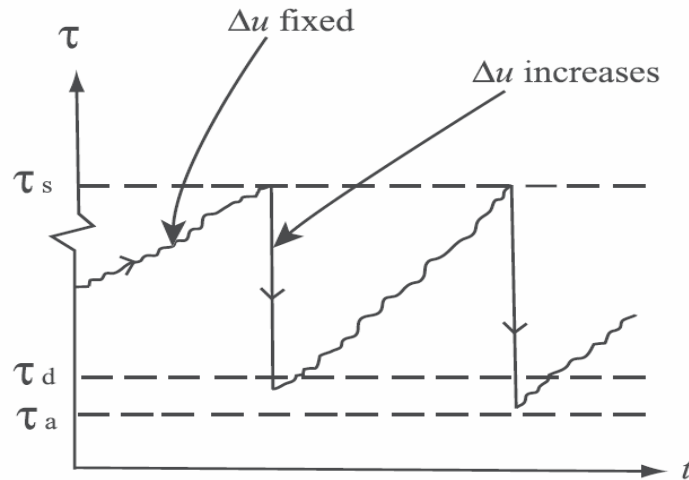
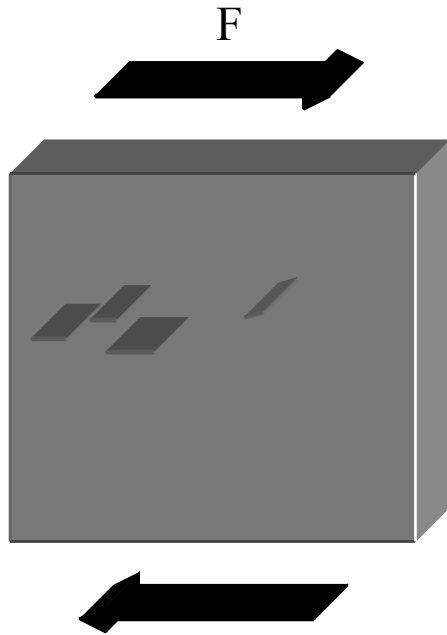
Renormalization Group:

Interaction sufficiently long range $\int dt J(\mathbf{r}, t) \sim r^{-2}$

SO THAT MEAN FIELD THEORY IS EXPECTED TO GIVE EXACT RESULTS !!!



Threshold pinning ($f_R[u, r, \text{history}]$)



$$\epsilon = (\tau_s - \tau_d) / \tau_s = \text{dynamic weakening}$$

weakening ($\epsilon > 0$)

- during failure avalanche:
- failed regions get weakened by $O(\epsilon)$
- reheal to old strength after avalanche

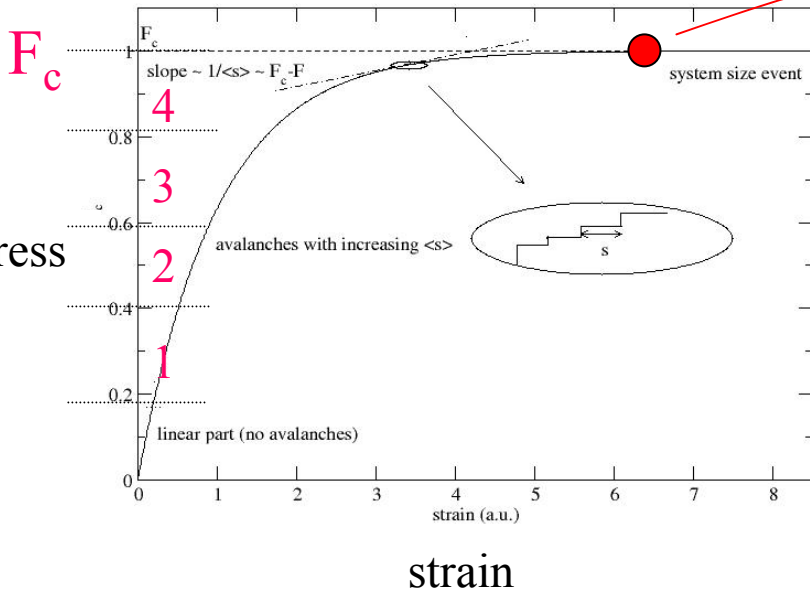
hardening ($\epsilon < 0$)

- each failure event raises failure threshold everywhere by $\sim |\epsilon/N|$.
- used to model aftershocks
- (Mehta, KD, Ben-Zion, PRE 2006)



Fixed stress boundary condition: Results for $\epsilon=0$:

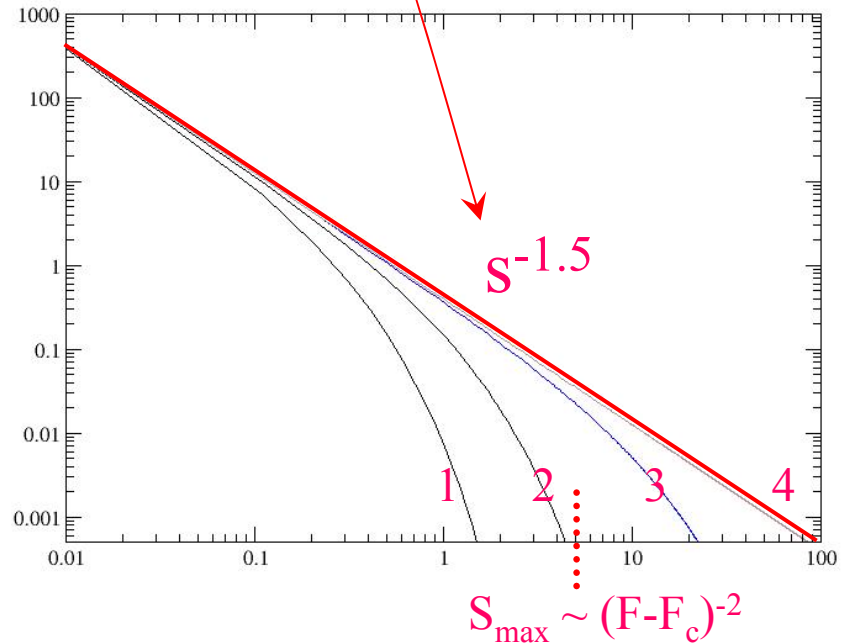
Stress strain curve (ductile)



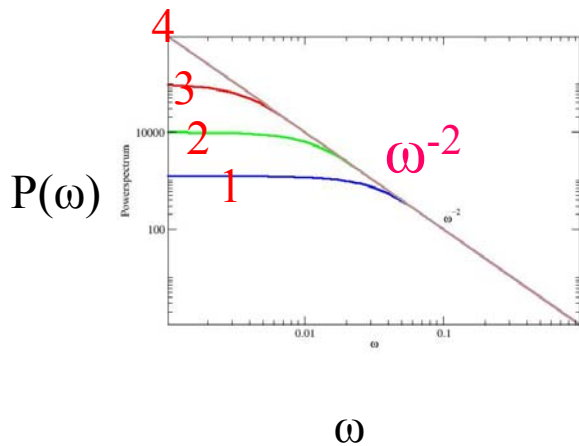
critical point

Avalanche size distribution:

$\log(D(s))$



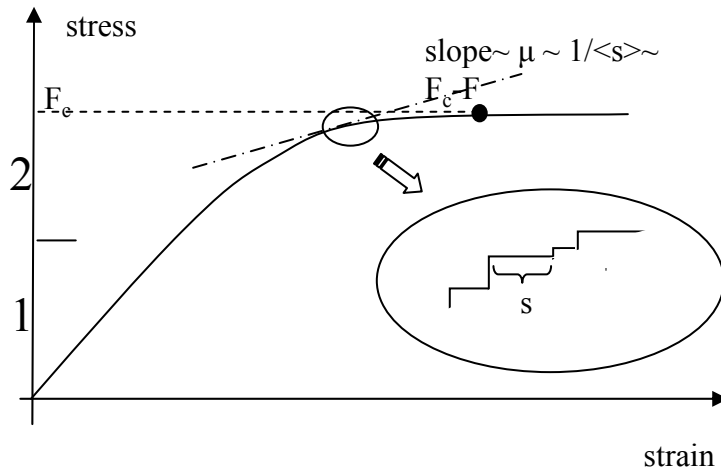
Power spectrum



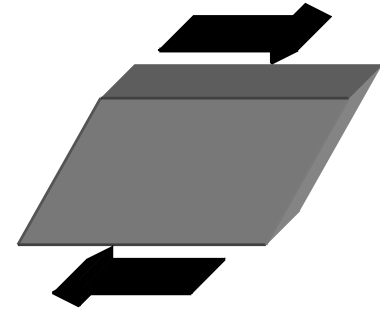
Log(size s)



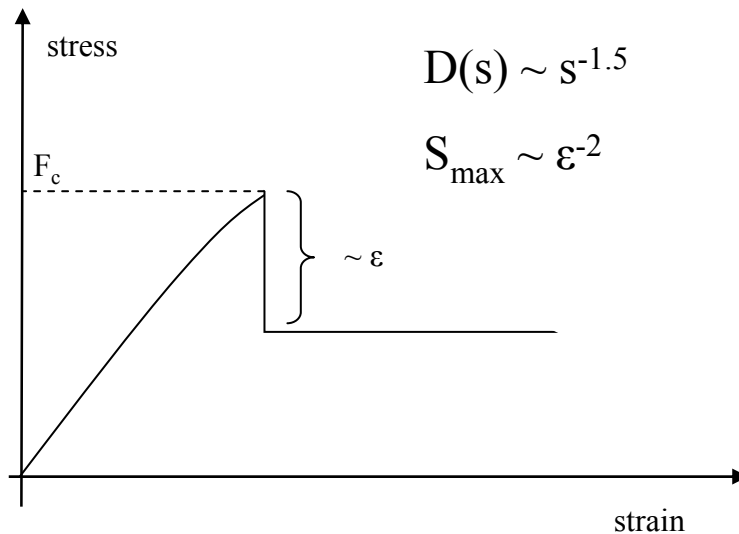
Without weakening (“ductile”): ($\epsilon=0$):



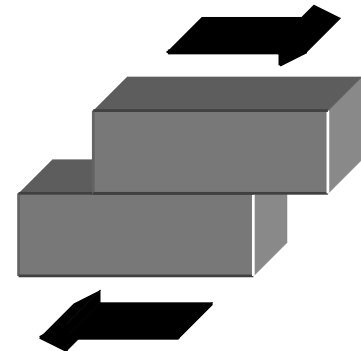
cont. depinning transition
and distributed slip



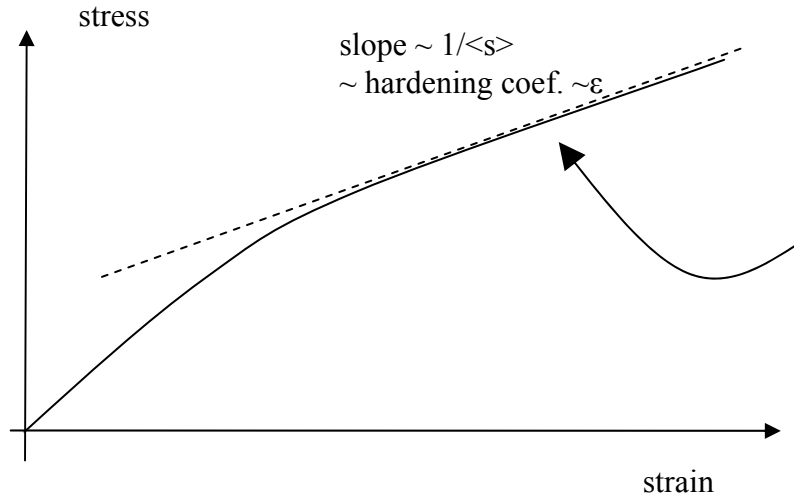
With weakening (“brittle”): ($\epsilon > 0$):



first order depinning transition
and slip localization



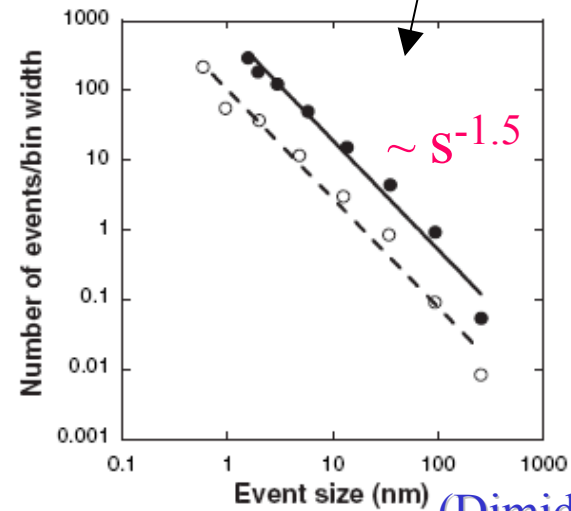
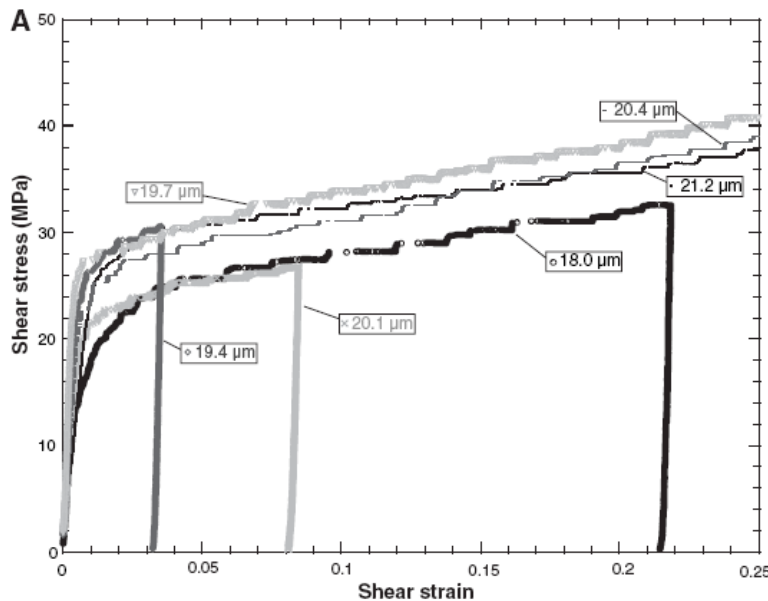
Stress Strain Curve in the presence of hardening (“ductile”): ($\epsilon < 0$)



$$\left\{ \begin{array}{l} D(s) \sim s^{-1.5} \\ P(\omega) \sim \omega^{-2} \\ D(T) \sim T^{-2} \\ \text{etc.} \end{array} \right\}$$

distributed deformation

Agrees with experiment



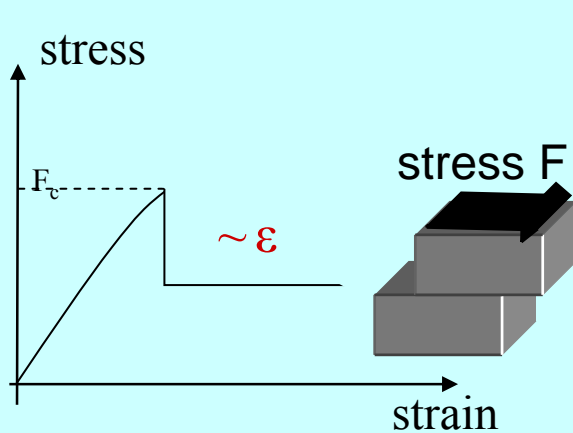
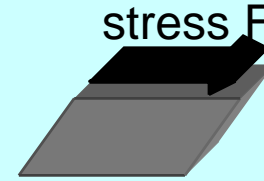
S

([Dimiduk, et al Science 2006](#))

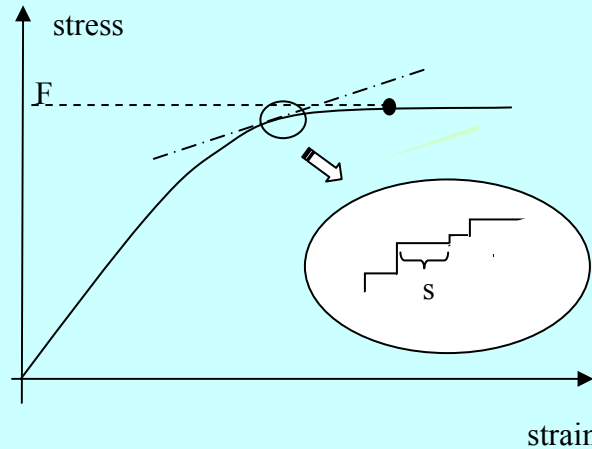
Simple Analytic Model for deformation under shear:

Main Results: (KD, Ben-Zion, Uhl, PRL 2009)

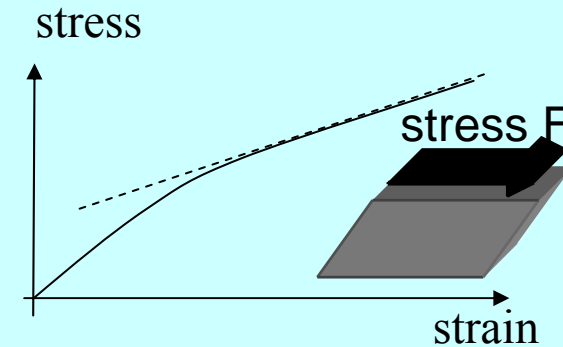
- For slowly increasing stress boundary condition:



Brittle ($\epsilon > 0$)



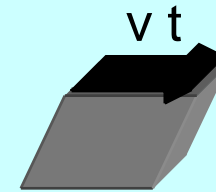
Plastic ($\epsilon = 0$)



Hardening ($\epsilon < 0$)

Avalanche-size distributions (power laws with stress dependent cutoff)

Power spectra, Scaling Functions !! Scaling behavior agrees with experiments!

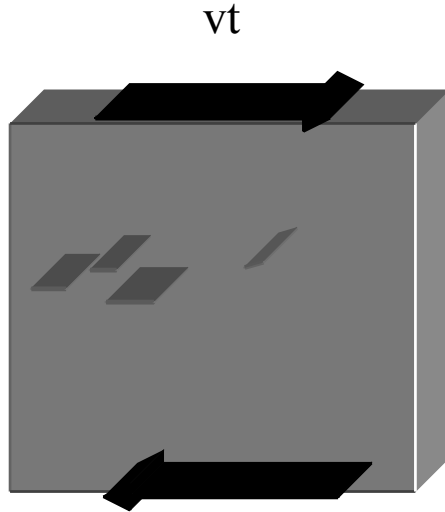


- For slow tangential velocity boundary condition:

same scaling behavior and phase diagram as Ben-Zion and Rice single earthquake fault zone model

For slowly moving boundary cond.: Same phase diagram and scaling as Ben-Zion and Rice **single** fault zone eq. model:

Ben-Zion and Rice, 1993, 1995; Ben-Zion, 1996; Fisher et al., 1997; KD et al., 1998; Mehta et al., 06; Zöller et al., 05, 07; KD and Ben-Zion, 08; Bailey and Ben-Zion, 08)



weakening ($\epsilon > 0$):

Char.Eq.Distr. + mode switching

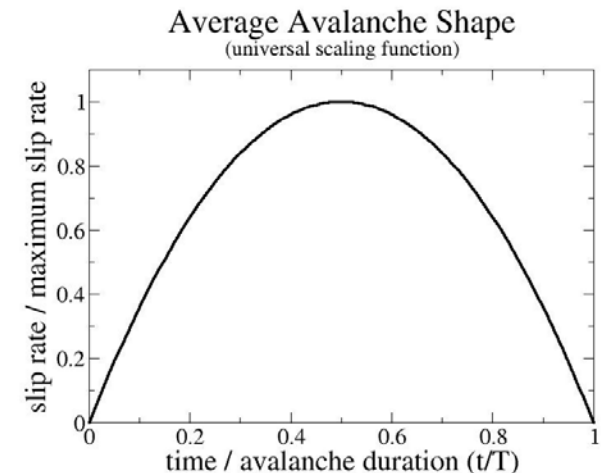
$\epsilon = 0$:

Gutenberg Richter

hardening ($\epsilon < 0$):

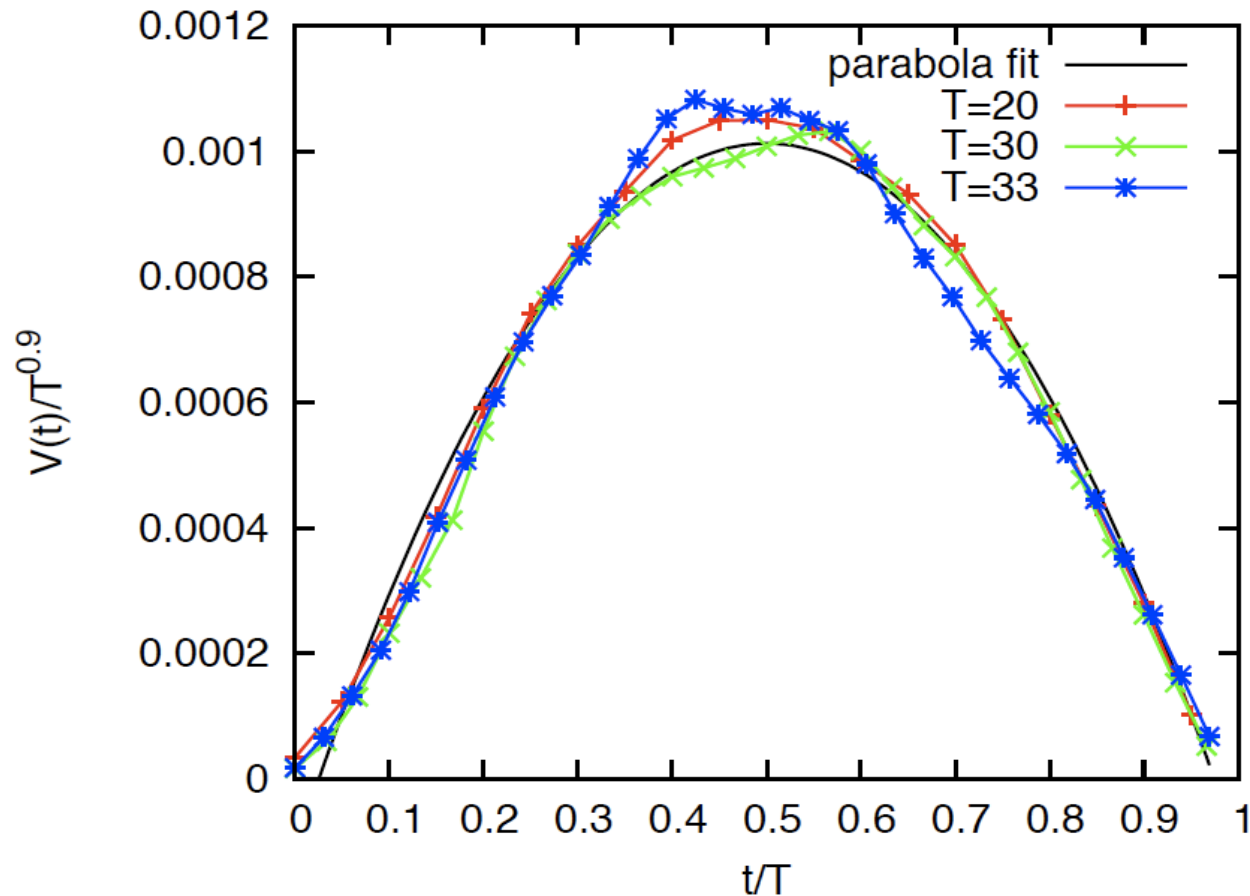
Gutenberg Richter (with Aftershocks)

There are more predictions for universal exponents, scaling functions, etc...



Georgios Tsekenis: Simulation of slip avalanche profiles

Exponents and Scaling function match Mean Field Theory Predictions!

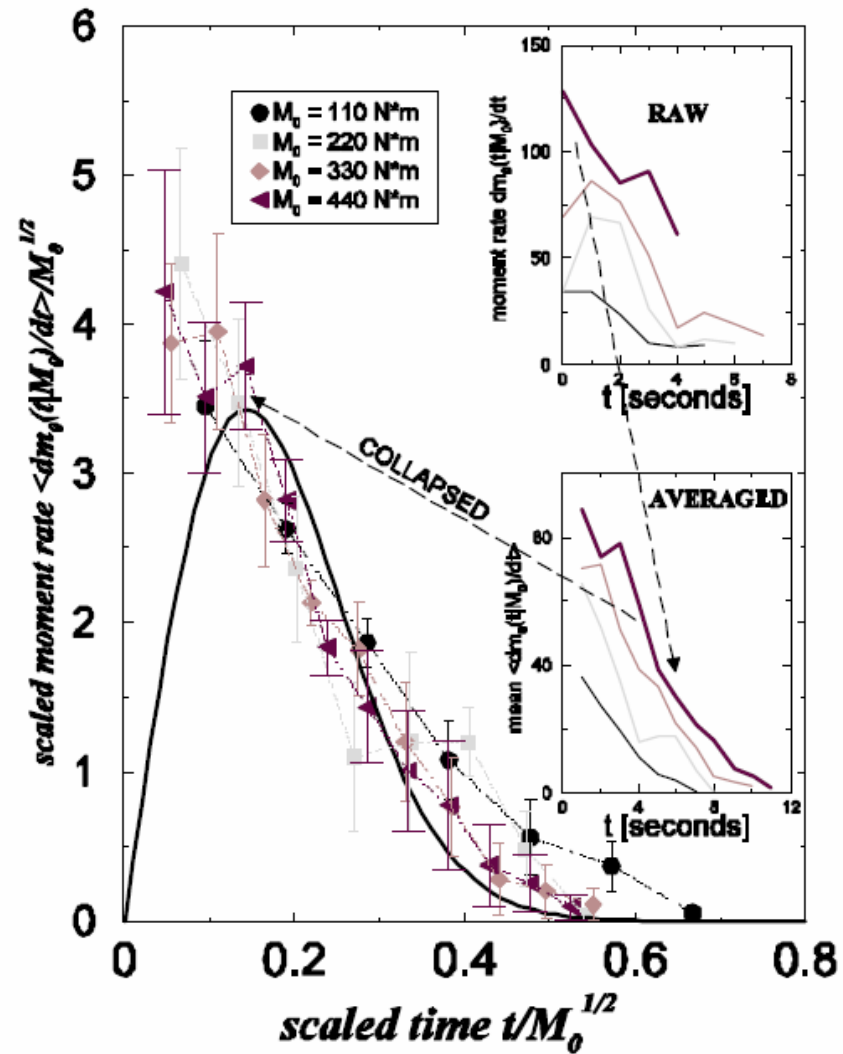
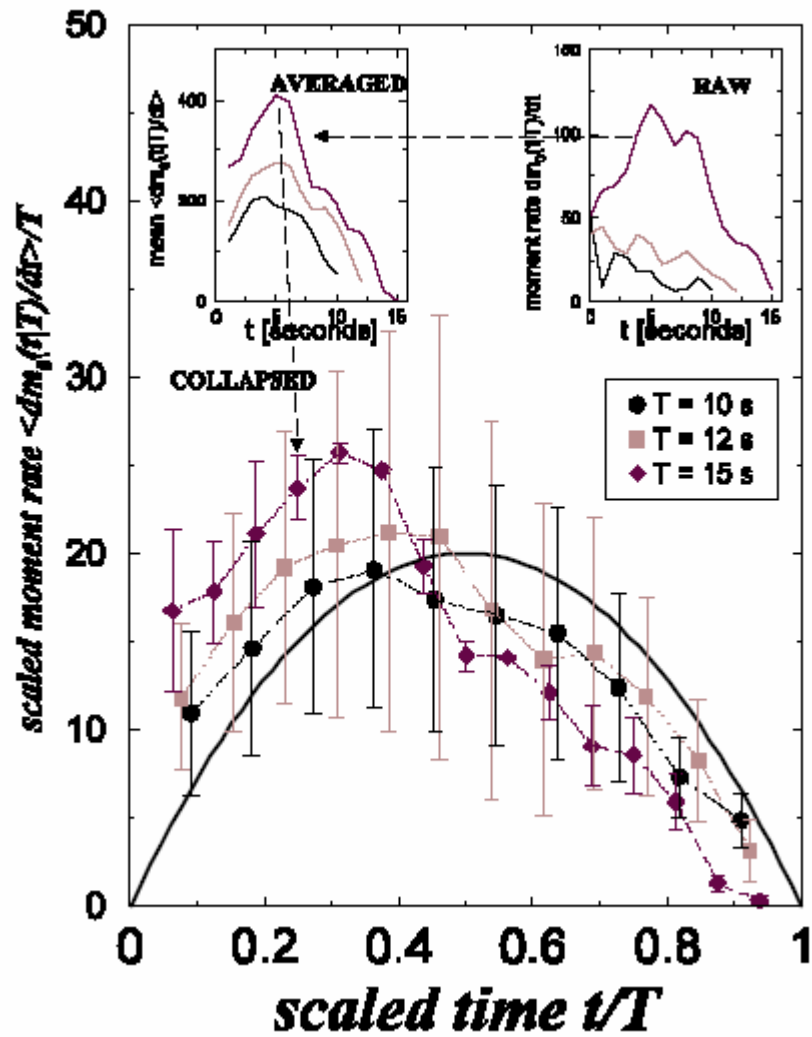


(Prediction of MFT parabola: Kuntz, Panagopoulos, Sethna)

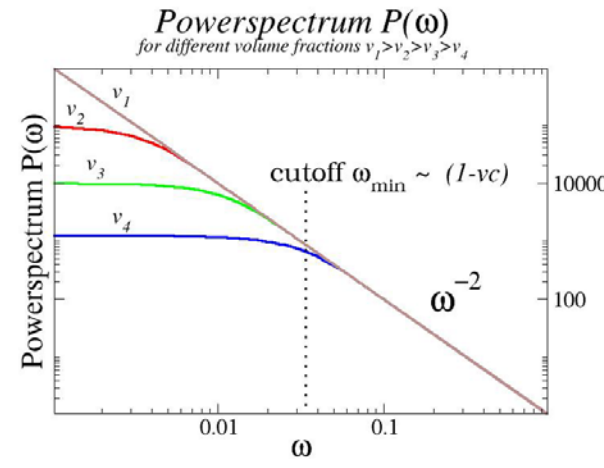
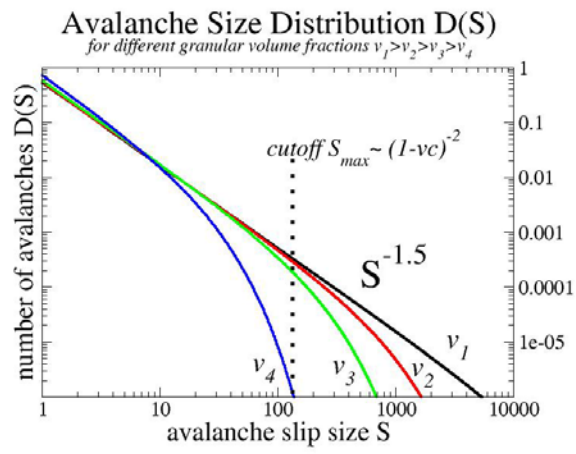
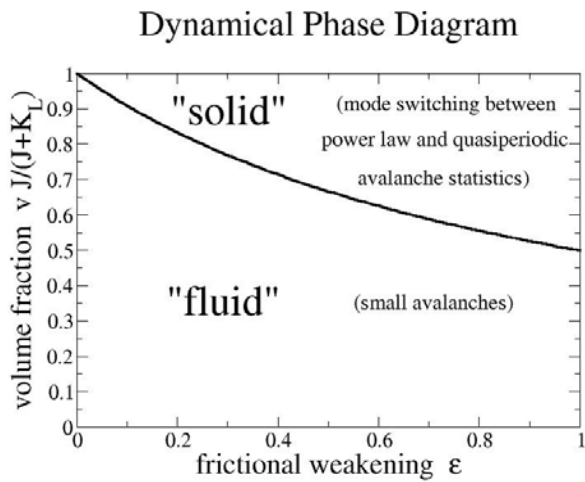
Data from Susan Bilek, see

Mehta, KD, Ben-Zion, 2005

Earthquakes: Universal Scaling Functions:

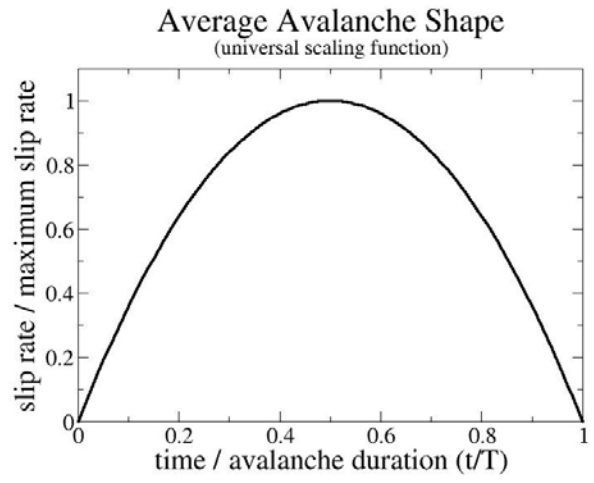


Similar SIMPLE model for avalanches in sheared granular materials:



Experiments: APS talk by Karen Daniels

Exponent or other universal quantity	Mean Field Theory	granular experiment [6,8-10,20-21]	granular simulations [2-4]
κ (size distribution)	1.5	1.5	?
$1/\rho\nu z$ (power spectrum)	2 if $v \approx 1$; 0 if $v \ll 1$	1.8-2.5, 2	2 in solid regime 0 in fluid regime
α (duration distribution) (^)	2	2 or exponential ?	?
Source time function averaged over all avalanches of same duration T .	Symmetric (parabola)	?	Symmetric: fit by sine function (?)
Quasiperiodic event statistics	Yes, if $\epsilon > 0$ and $v > v^*$	sometimes	during mode switching
Mode switching (between powerlaw and quasiperiodic)	Yes, if $\epsilon > 0$ and $v > v^*$?	Yes, in solid regime



SUMMARY

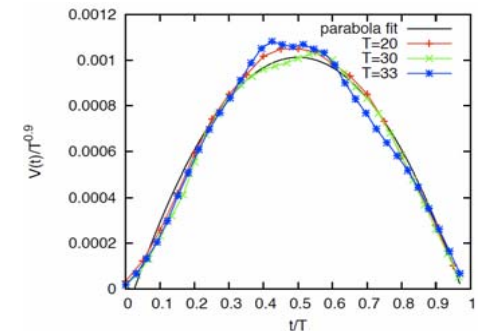
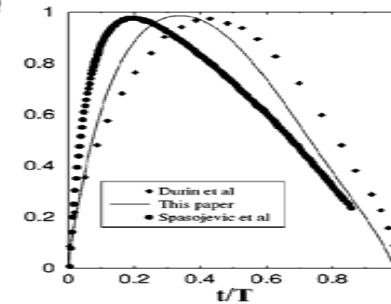
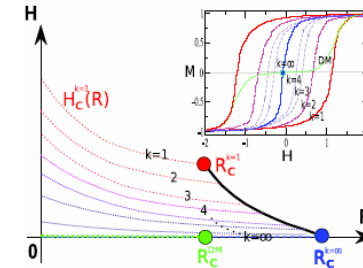
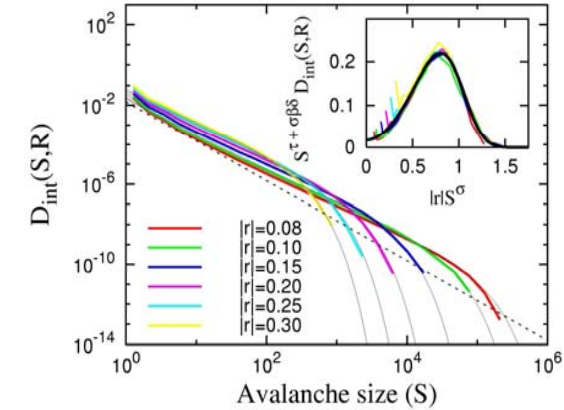
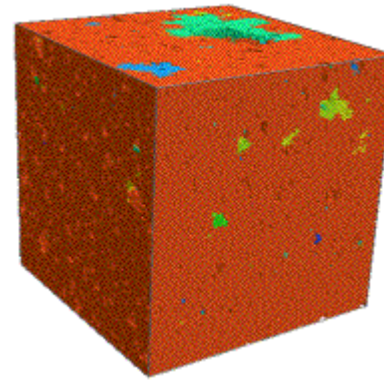
- HUGE nucleation universality class include:

- ✓ hard magnets
- ✓ equilibrium and nonequilibrium
- ✓ plastic charge density wave depinning
- ✓ high Tc superconductors(?)...

- Domain wall mean field theories exact & same scaling results for:

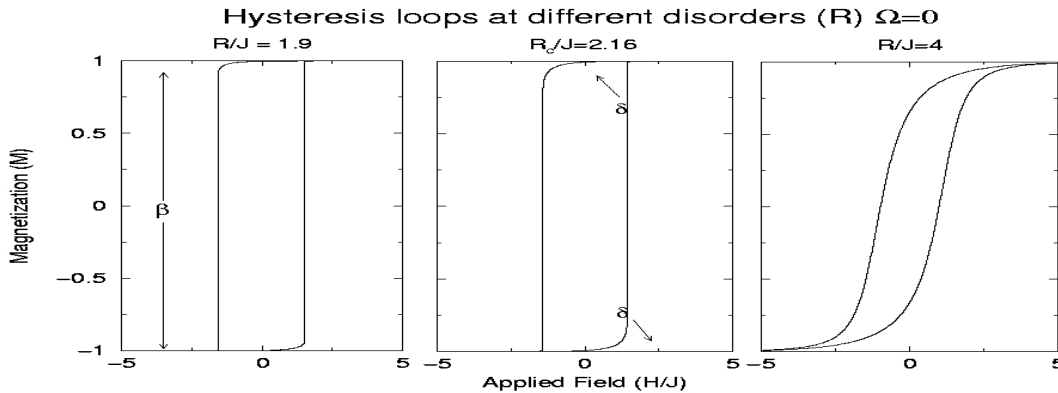
- ✓ magnetic domain wall motion
- ✓ plasticity
- ✓ granular materials
- ✓ Earthquakes

- **Experimental tests ?! Got DATA ?**

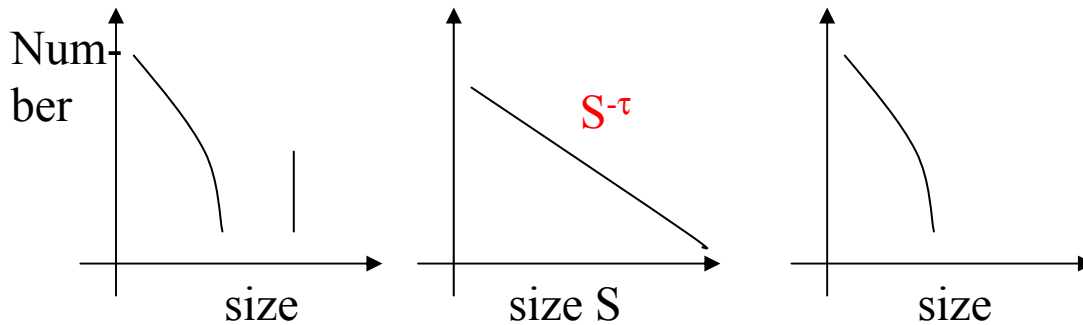


Discussion Slides:

Summary on simple magnet model & further results:



clean \rightarrow disordered



EXPERIMENTS ???!!!

Renormalization group
(huge universality
class)

equilibrium and
nonequilibrium critical
behavior the same !!

- finite sweep rate effects
(model indep. theory)

- Second Spectra (mean
field theory), universal
scaling functions, return
point memory,

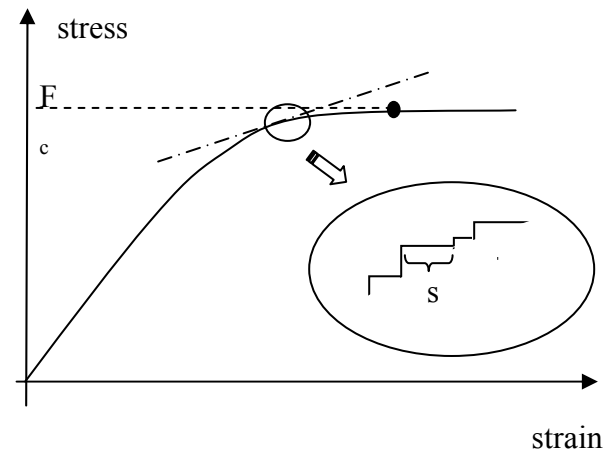
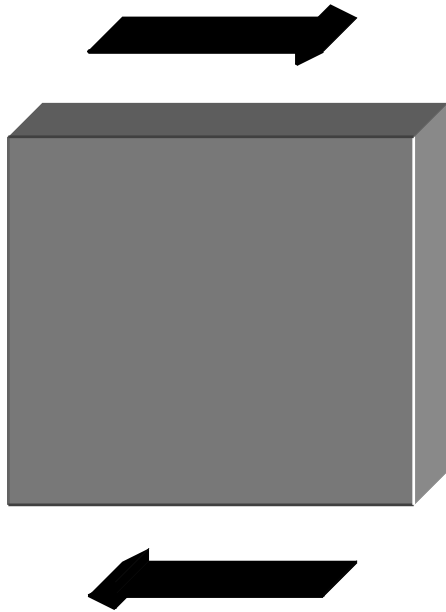
- Temperature effects...

Next:

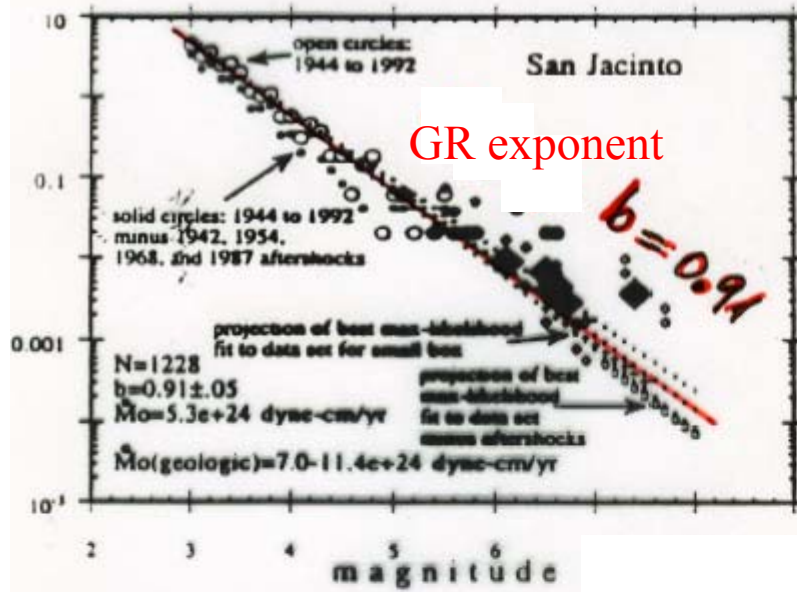
- *Plasticity,
- *Granular materials,
- *Earthquakes....

Plastic deformation and crackling noise:
Learning from Dislocations and granular materials
about Earthquakes ???

Yehuda Ben-Zion, Jonathan T. Uhl, Georgios Tsekenis, Pak Yuen Chan, Nigel Goldenfeld, Jonathan Dantzig, KD

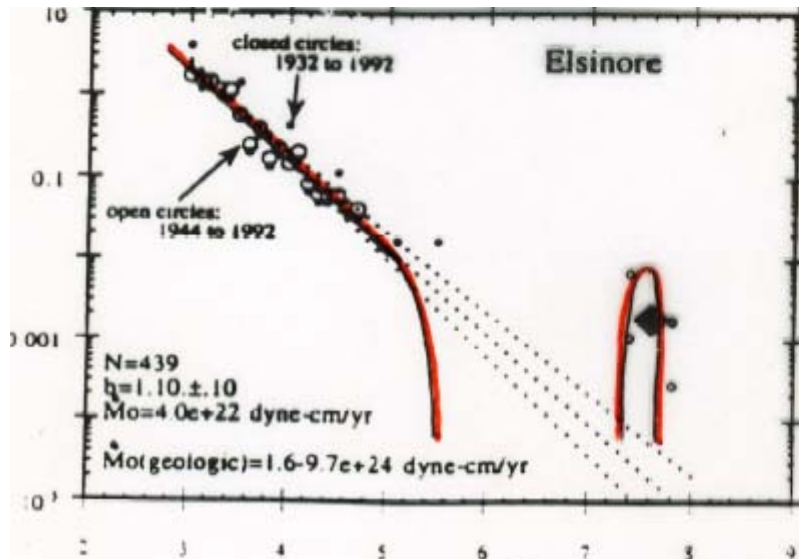


Universal Gutenberg Richter Scaling behavior



OR

Magnitude $\sim 2/3 \text{ Log}(\text{total displacement})$



Characteristic
Earthquake
distribution ?

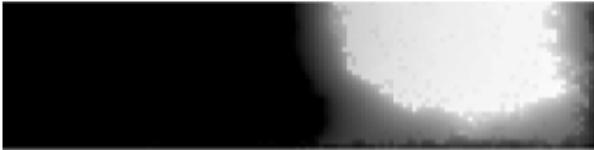
(Wesnousky '94)

Phase Diagram

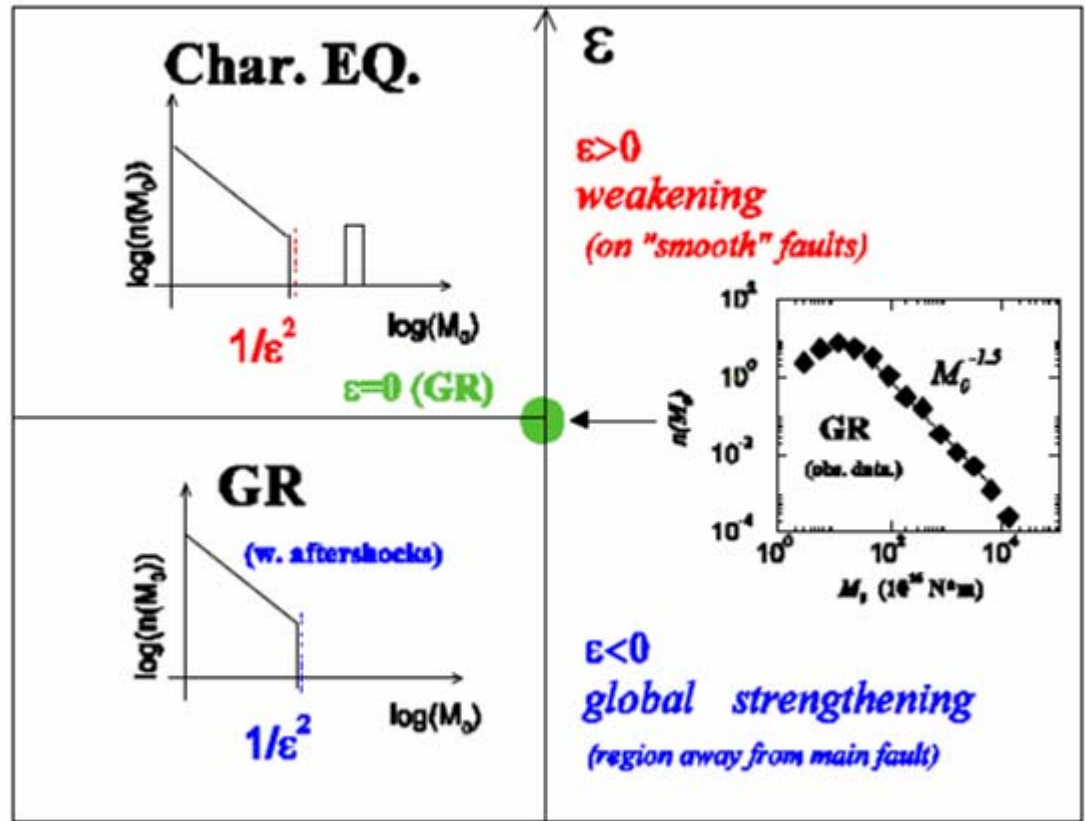
Mean field exponent: $b=0.75$

for Frequency $\sim 10^{-bM}$

cracklike large events:
moment \sim area^{3/2}

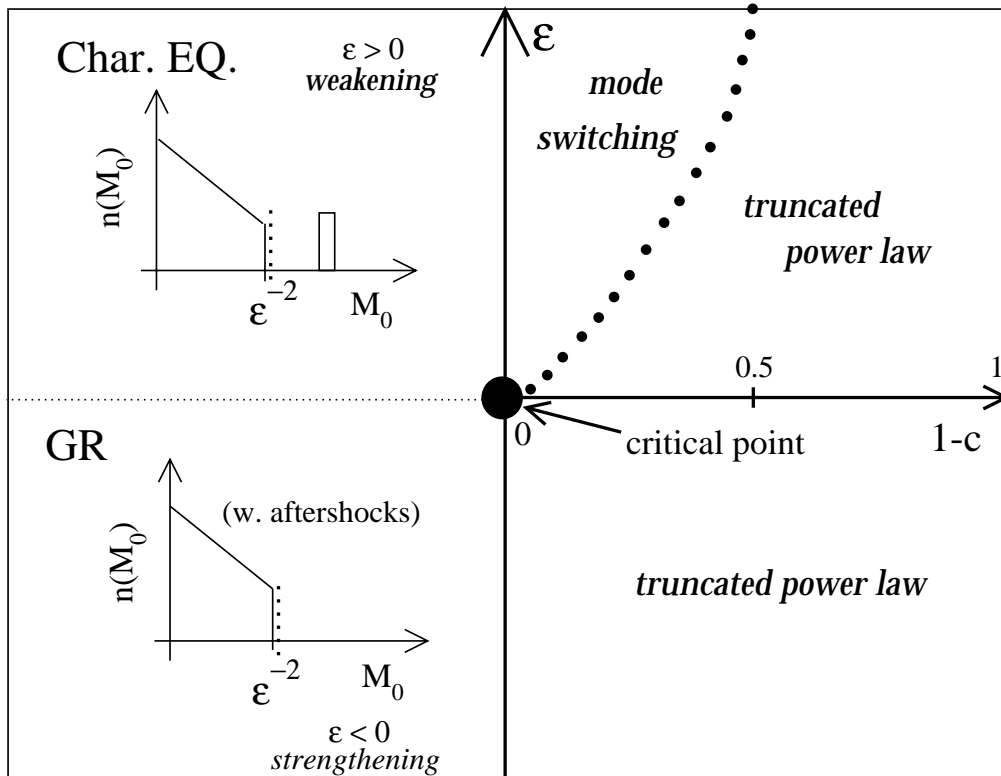
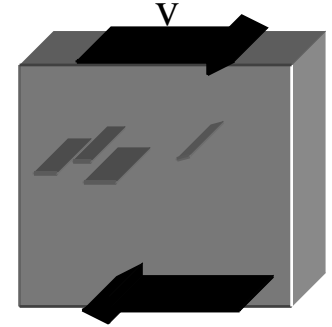


small events scale as
moment \sim area



For fixed velocity boundary conditions:

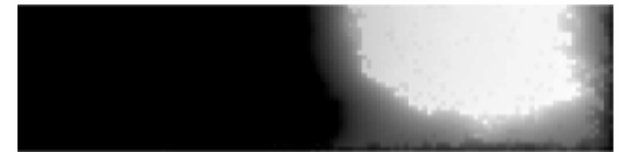
Scales like Ben-Zion and Rice earthquake model



Mean field exponent: $b=0.75$

for Frequency $\sim 10^{-bM}$

cracklike large events:
moment \sim area^{3/2}



small events scale as
moment \sim area



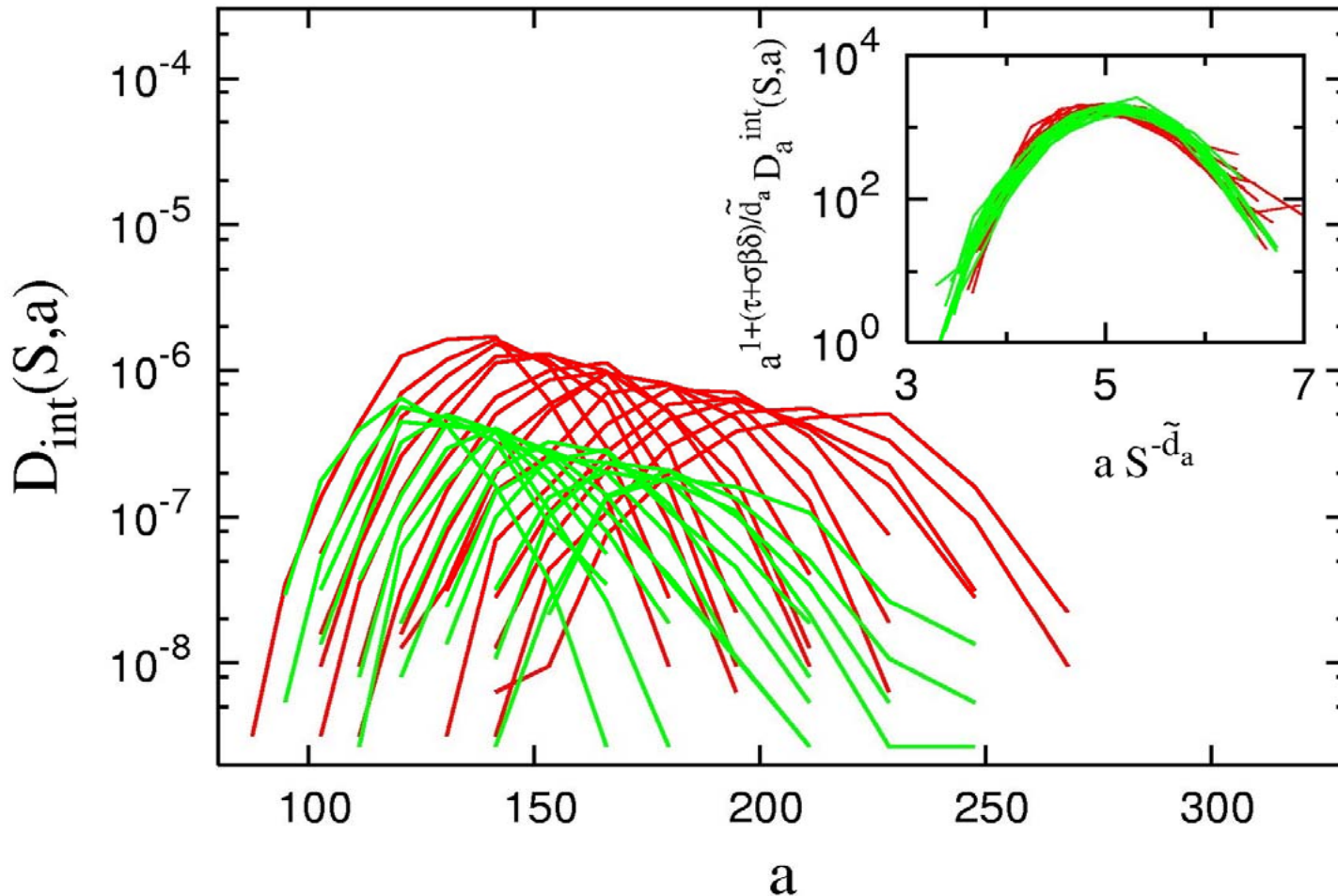
KD, Ben Zion (2008)



•Result: ~Same scaling of Avalanche Surface Area Distribution

Integrated avalanche surface distribution for avalanche size S

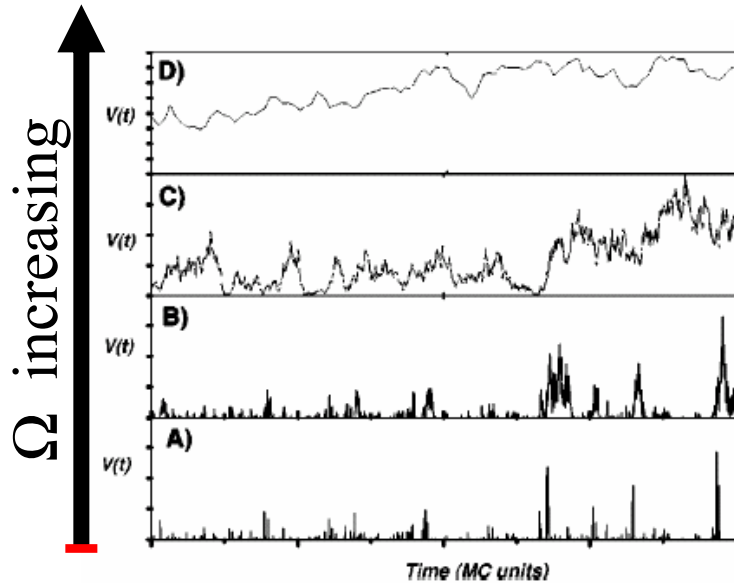
$$D_a^{(\text{int})}(S, a) \sim a^{-(\tau + \sigma\beta\delta + \tilde{d}_a)/\tilde{d}_a} \mathcal{D}_a^{(\text{int})}(a/S^{\tilde{d}_a})$$



$$\Omega \geq 0$$

Sweep Rate Regimes of Barkhausen Noise

White, KD
PRL2003



Region III (“fast”)

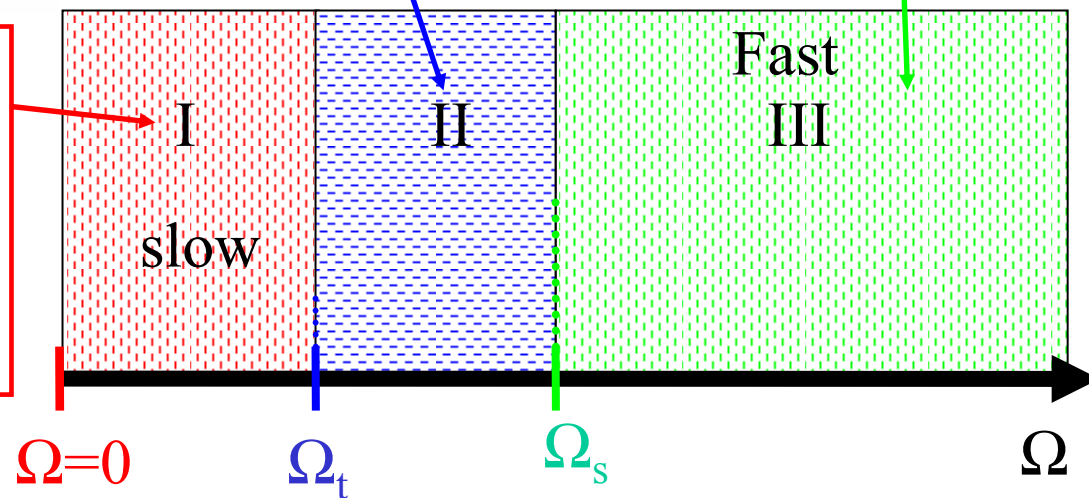
- Low Freq. PS changed due to spatial overlap

Region II (intermediate)

- Power spectra unchanged.
- Only temporal overlap

Region I (“slow”)

- power spectra unchanged
- Pulse statistics are affected in general



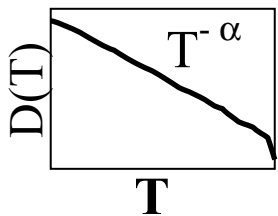
No pulse statistics
for $\Omega > \Omega_t$

How are finite sweeprate effects to be understood? (White, KD PRL 2003)

(MODEL “INDEPENDENT” THEORY):

Superposition of power law distributed avalanches

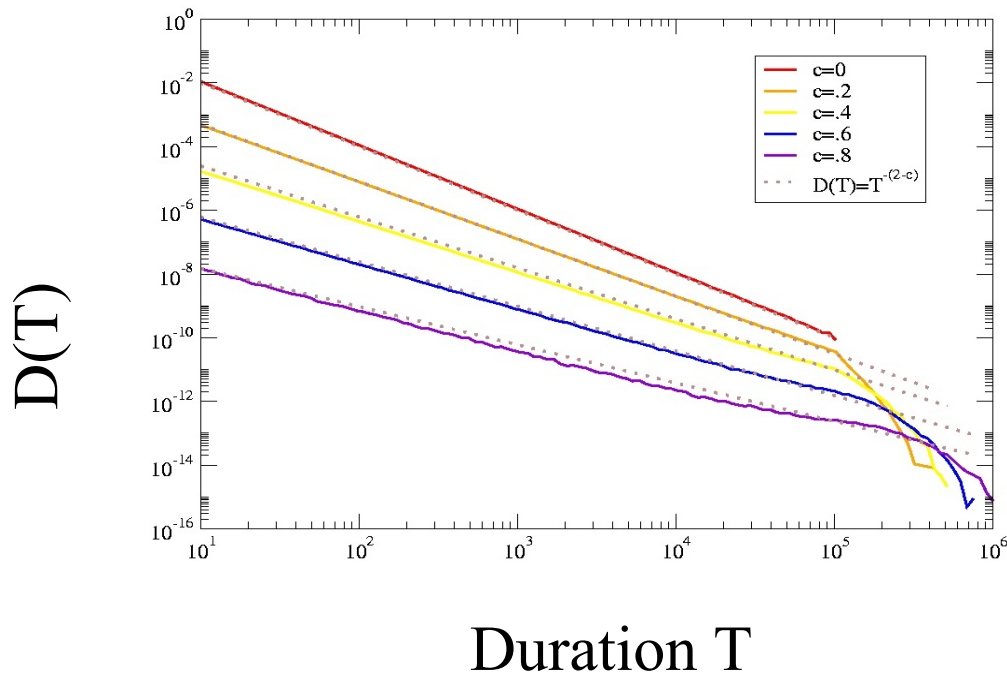
In slow regime dependence of pulse exponents on sweep rate explained by temporal overlap of pulses



$\Omega=0$	$\Omega>0$	
$\alpha_0 < 2$	α unchanged	Durin, Zapperi
$\alpha_0 = 2$	$\alpha(\Omega) = \alpha_0 - c\Omega$	ABBM, Durin
$\alpha_0 > 2$	α unchanged for large durations	???

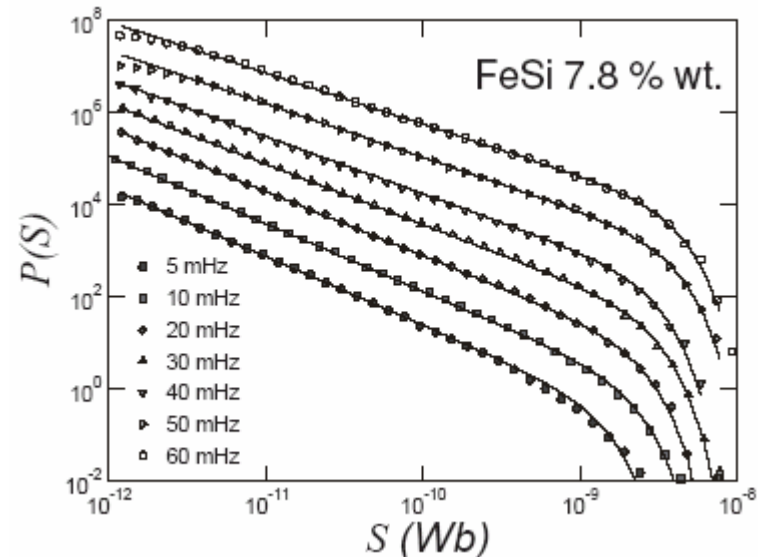
Results for $\alpha(\Omega=0) = 2$

Linear change in exponent $\alpha(\Omega) = \alpha_0 - c\Omega$



Agrees with Experiment:
Plot from Durin and Zapperi
Cond-mat/0404512v1

$$\tau(\Omega) = 1.5 - c\Omega/2$$



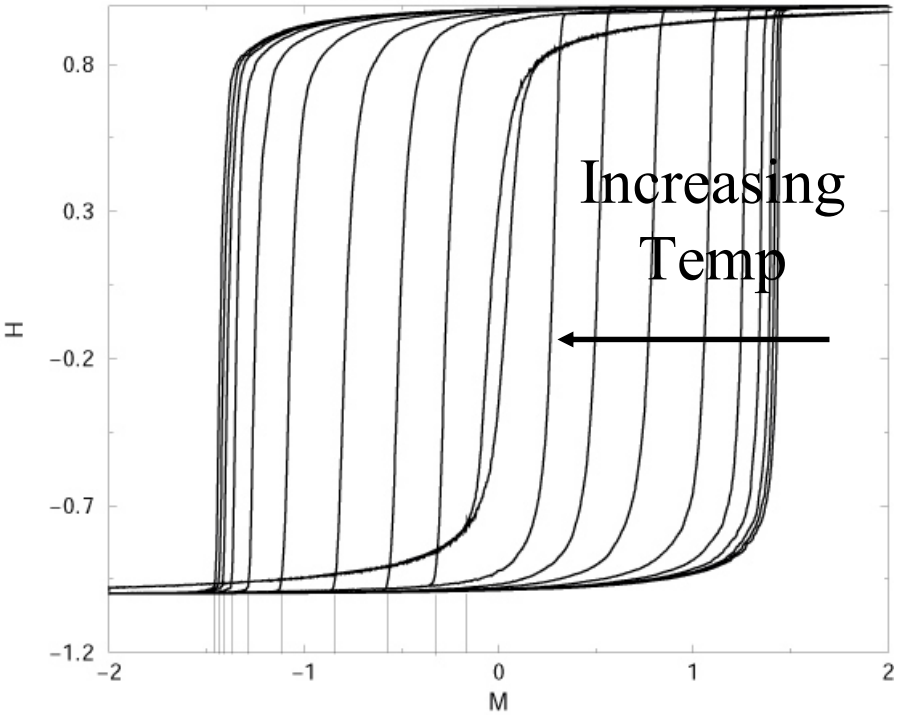
Temperature > 0

Scaling Theory and
Simulations using the random field Ising model near the
zero temperature far from eq. critical regime (in progress)

Robert White, Alex
Travasset, Yang Liu, KD,
EPL 2009,

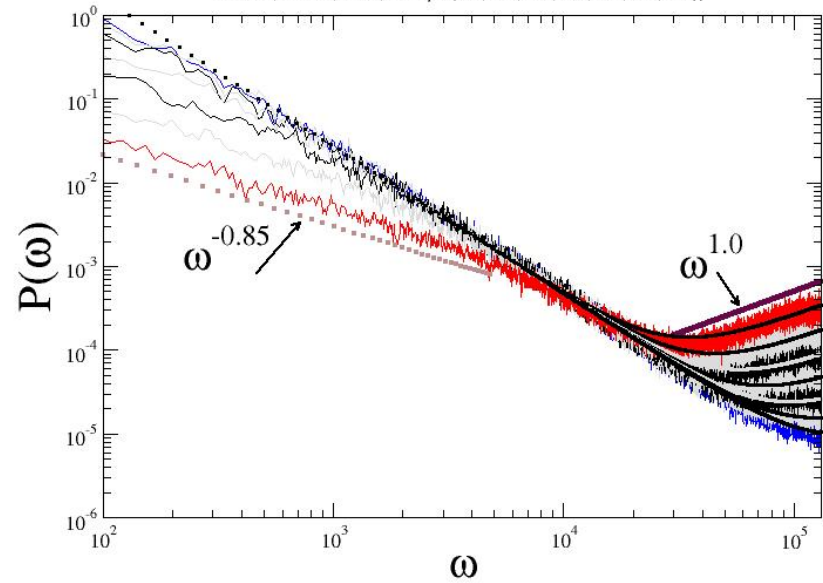
M vs H with increasing Temperature

($T=0,0.1,0.2,0.4,0.8,1.6,3.2,6.4,12.8$)



Spectra for increasing temperatures

($L=100, R=2.3, \Omega=2$, Temps (.01, .02, .04, .08, .16, .32, .64))

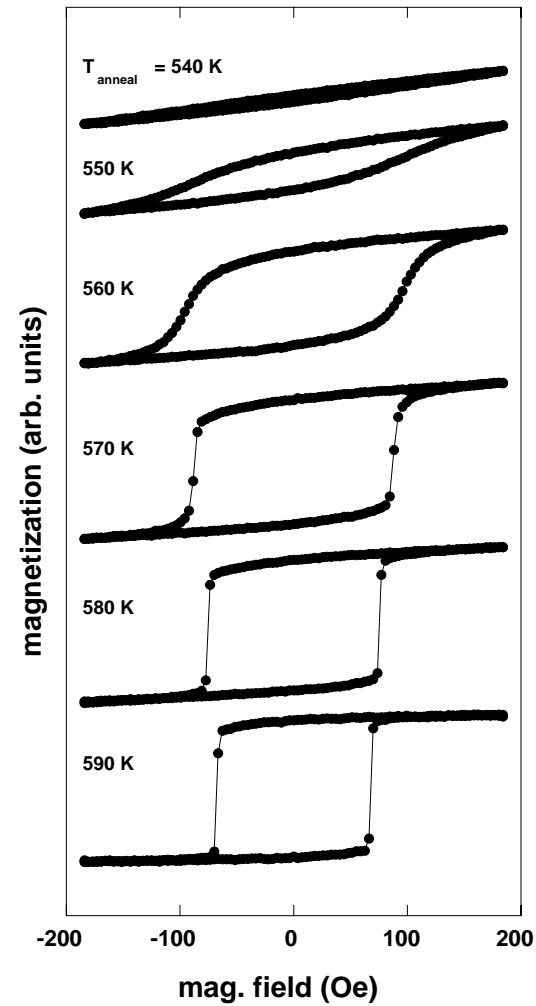
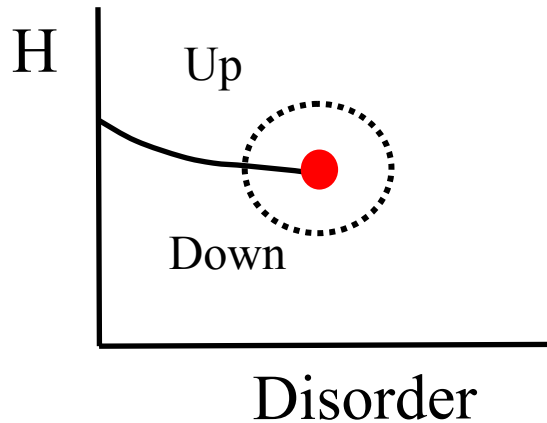


Experiments-

Phase transition?

Disorder is important!

Noise experiments?



“Dirty”



“Clean”

A. Berger et al. PRL 85 4176(2000)

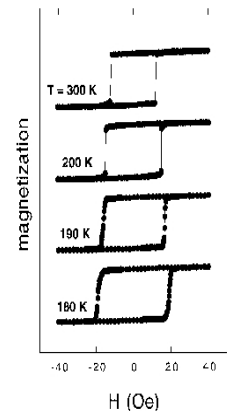
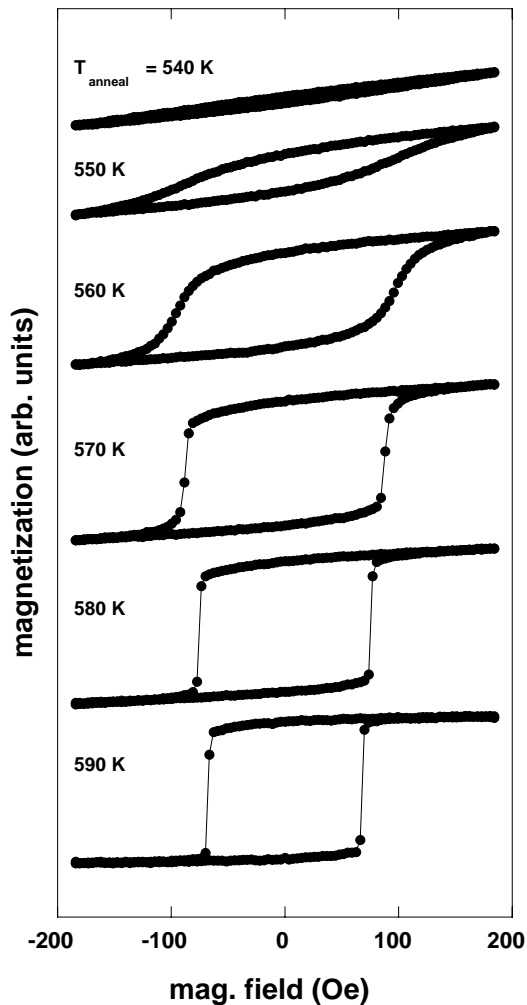


FIG. 1. $M(H)$ loops measured on a Co/CoO-bilayer structure for the temperatures indicated. The thin lines are guides to the eye.

Gd(0001)/W(110)-films

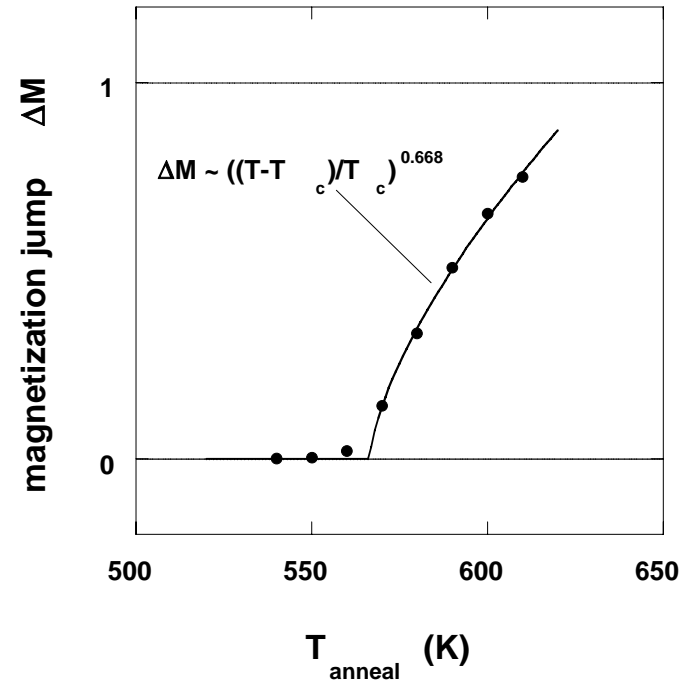
annealing study



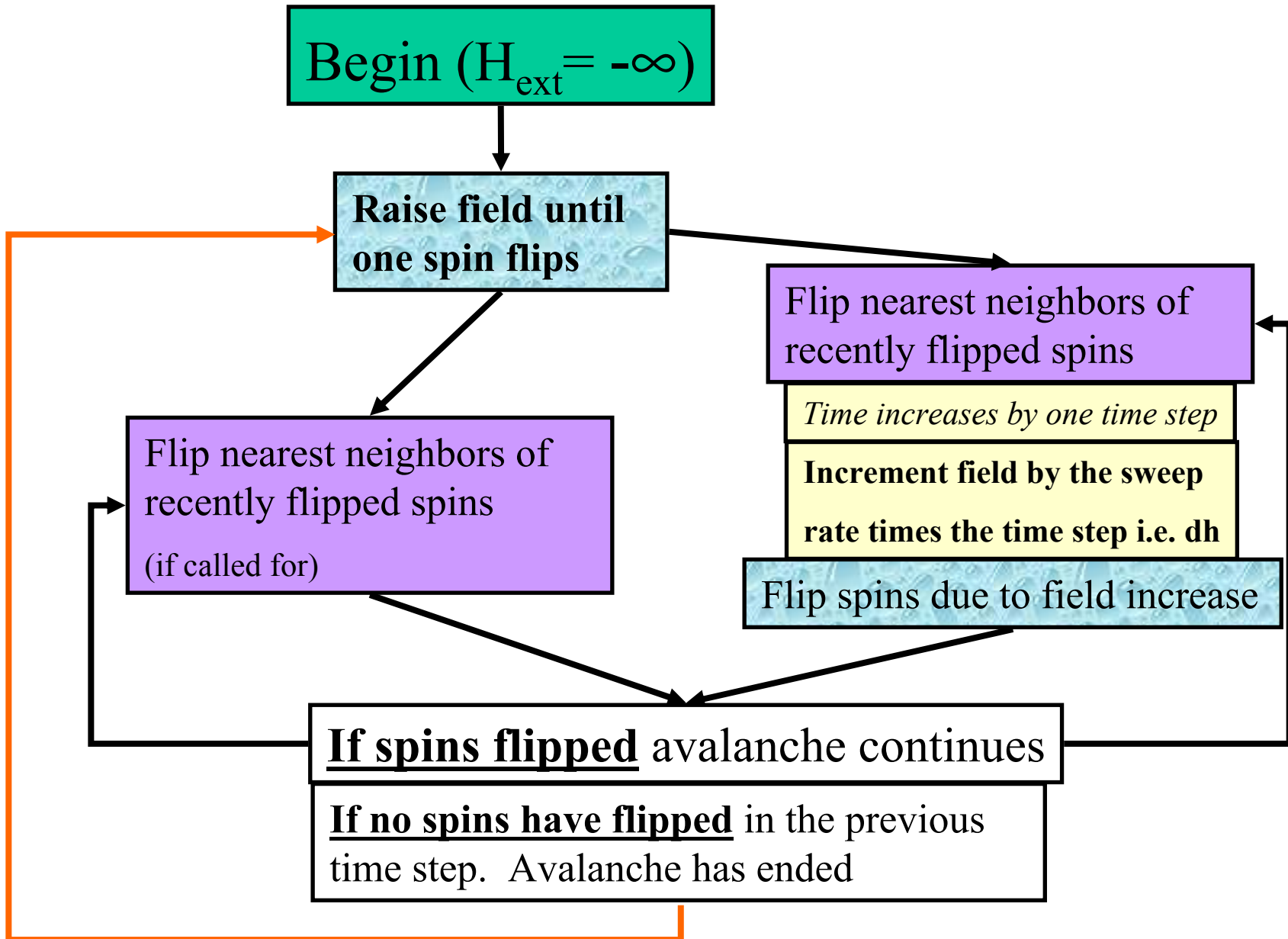
data analysis

fit of experimental data to:

$$M(H) = \Delta M \cdot \tan^{-1}\left(\frac{H - H_C}{\delta_1}\right) + A_2 \cdot \tan^{-1}\left(\frac{H - H_C}{\delta_2}\right)$$



The Algorithm- (adiabatic.... Finite sweep rate)



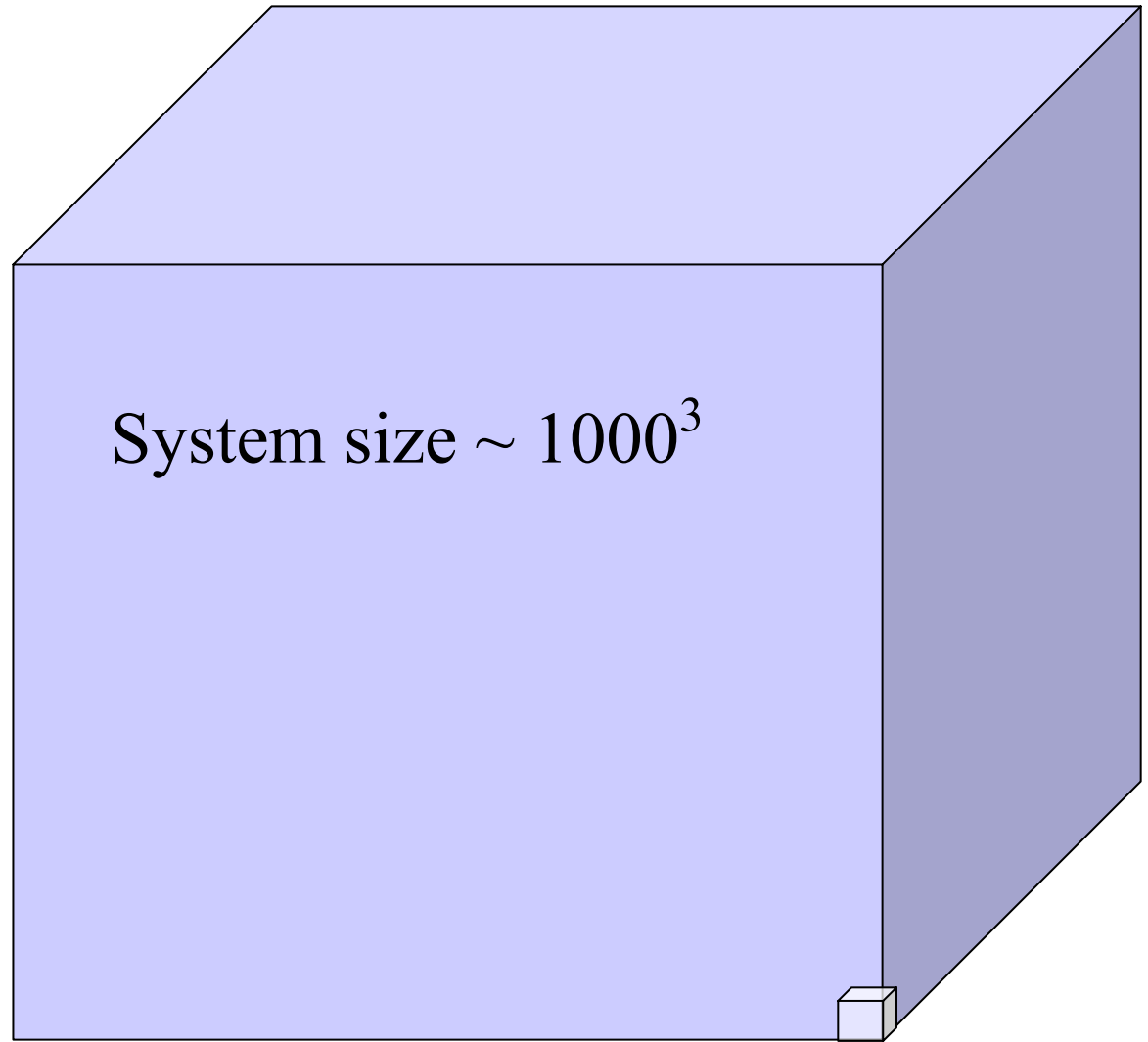
Simulation Overview-

Run Time Scaling

$O(N \log N)$

Memory Scaling

$O(N)$



Equilibrium RFIM

MC methods size $\sim 25^3$