

# Jamming in Dense Suspensions

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KITP Program on Granular Materials — 31 March 2005

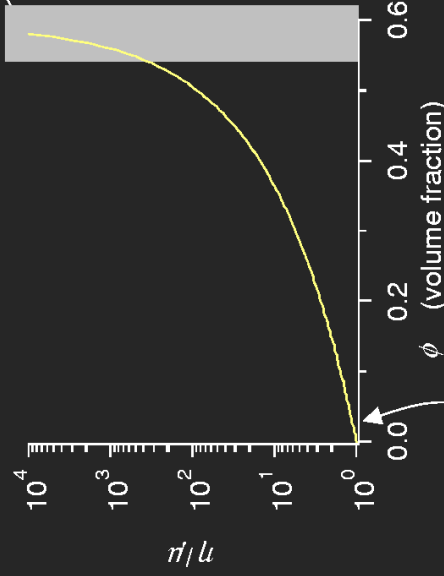
## Rheology of concentrated suspensions

- All existing data looks like this (roughly)
- Viscosity diverges at high volume fractions (<random close packing?)
- Universal behavior of viscosity?
  - Interactions — hard spheres
    - difficult to achieve
    - sensitive to short range repulsive interactions
  - Shear rate — Peclet number

$$Pe = \frac{\dot{\gamma}}{D/a^2}$$

- System eventually jams

- At what volume fraction?  $\phi_{RCP}$ ?
- Depends on shear rate (Pe)?



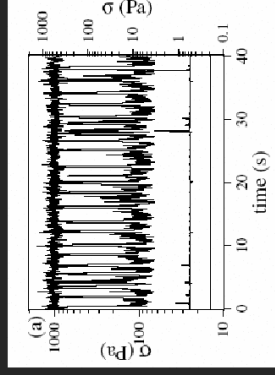
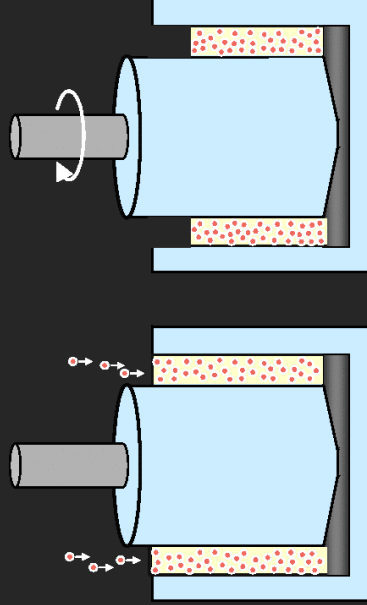
$$\eta(\phi) = \mu \left( 1 + \frac{5}{2}\phi + \dots \right)$$

(not published in 1905)

# Random loose packing

(G.Y. Onoda & E.G. Liniger)\*

RLP = loosest packing that can support an external load



Lootens, Van Damme, & Hébraud, PRL **90**, 178301 (2003).

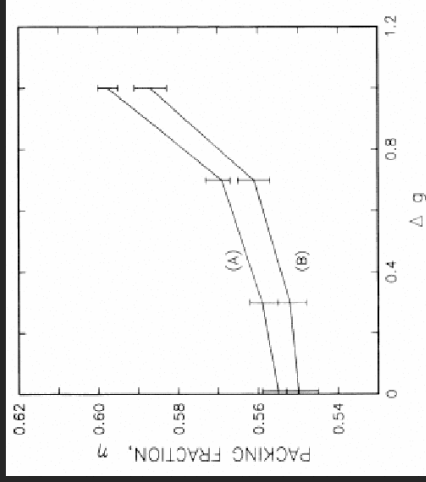
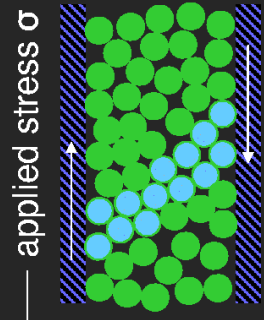
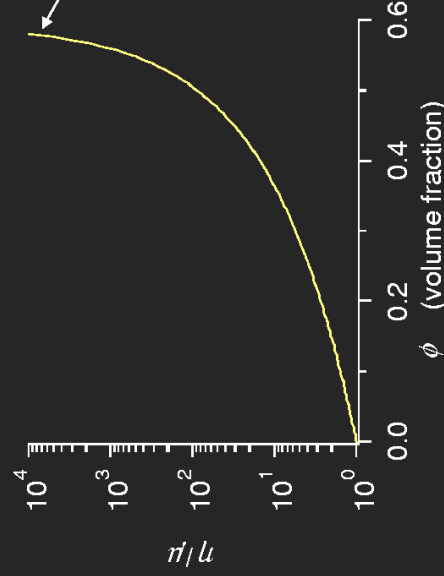


FIG. 1. Variation in packing fractions of spheres with the effective gravitational acceleration. Curve *A* is for the random loose packings obtained in the settling experiments, while curve *B* is for the expanded packings obtained in the shear-cell experiments.

\* PRL **64**, 2727 (1990) 3

# Jamming in concentrated suspensions



4

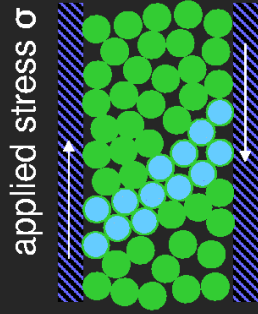
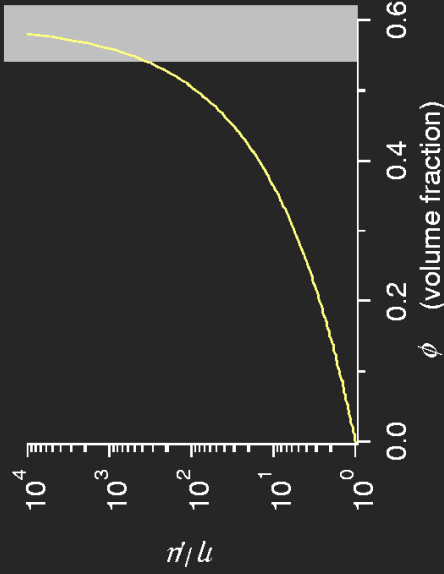
# Jamming in concentrated suspensions

- Jamming usually studied by measuring viscosity divergence
- Usually measured at constant applied shear rate  
 ⇒ Jamming is impossible under these conditions

$$\sigma = \eta \dot{\gamma}$$

- This approach *infers* jamming – does not measure it

Jamming studies should be performed at constant stress

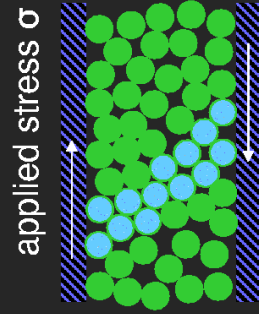
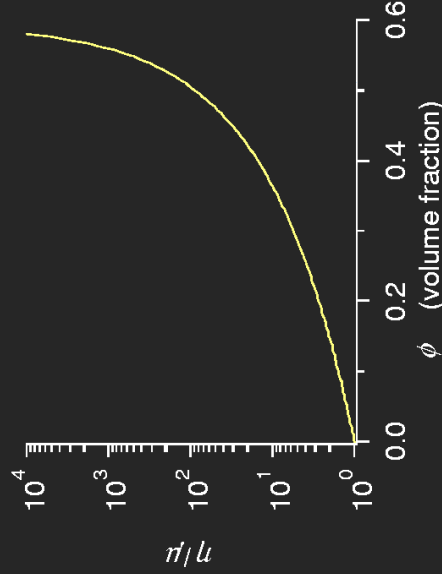


5

# Jamming in concentrated suspensions

Why study jamming in non-colloidal suspensions?

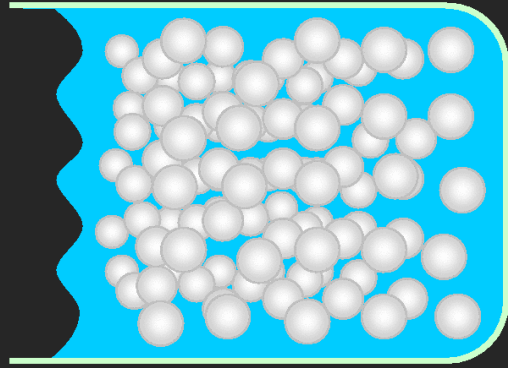
- Shear dominates diffusion:  $\dot{\gamma} \gg D/a^2$   
 diffusion frustrates jamming
- Can use *large particles* to see *inside bulk* ⇒ *microstructure*



- requires *matching refractive indices* of particles and suspending liquid
- Can minimize effects of gravity by matching density of particles and suspending liquid

6

# Our suspensions



## Spheres

- PMMA
- diameter  $\approx 210 - 250 \mu\text{m}$ 
  - *no Brownian motion* (too big)
  - *polydisperse* to avoid ordering
- concentrated  $\rightarrow 0.5 < \phi < 0.6$

## Liquid

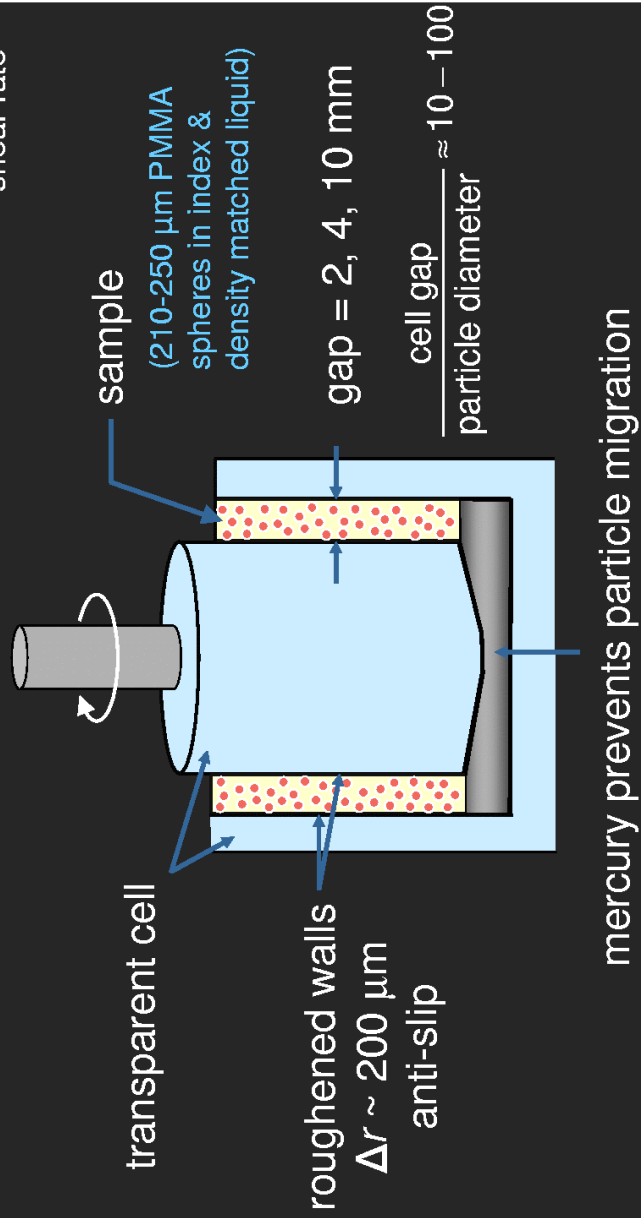
- Newtonian rheology
- refractive index matched\* to PMMA
- density matched\* to PMMA

\* Leighton formulation

# Couette cell for rheology

controlled stress  $\eta = \frac{\sigma}{\dot{\gamma}}$

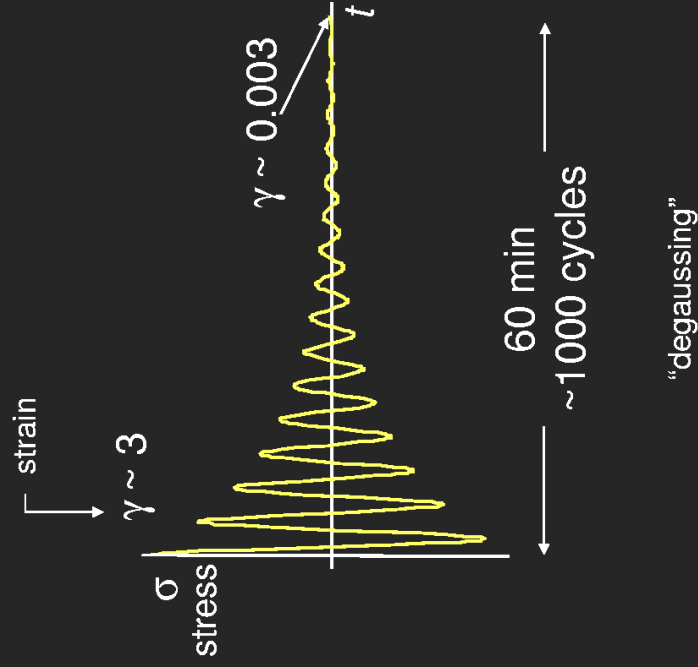
$\sigma$  ← apply stress  
 $\dot{\gamma}$  ← measure shear rate



# Randomization of Sample

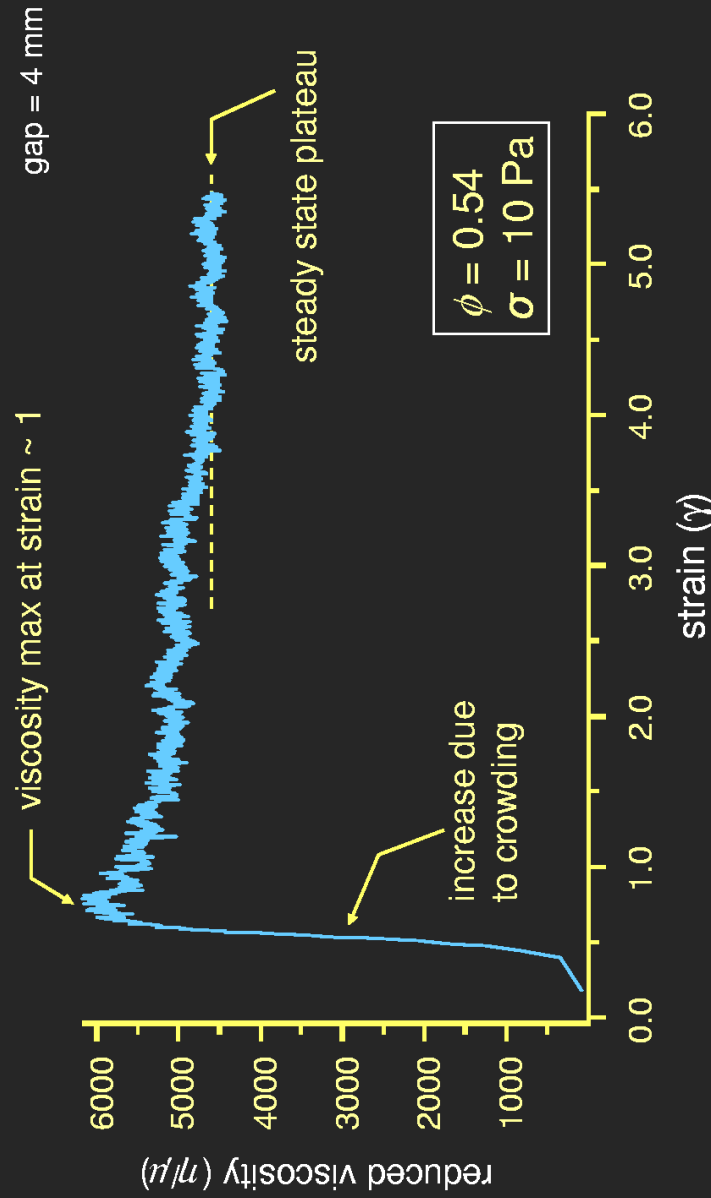
## No Brownian motion

- Initial state of system depends on sample history
- Need a protocol to produce well-defined initial conditions
- use oscillating stress (controlled shaking)
- $\Rightarrow$  loose packing (more later)



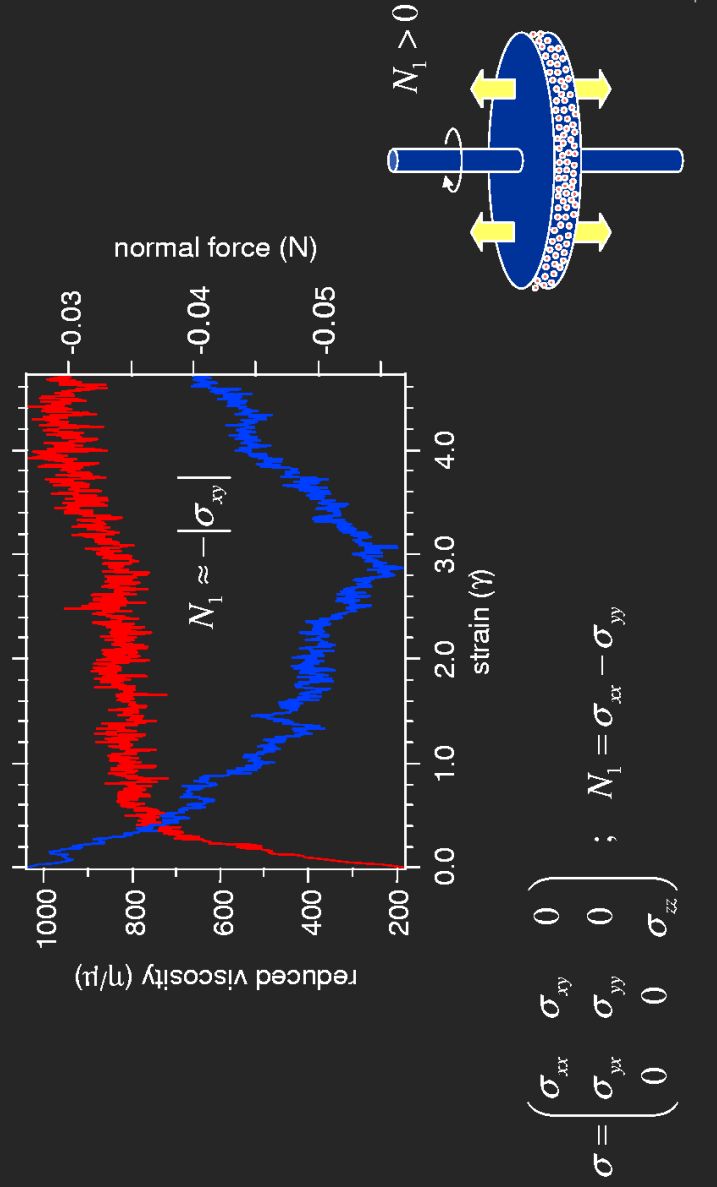
9

# Start-up flow (viscosity): no jamming



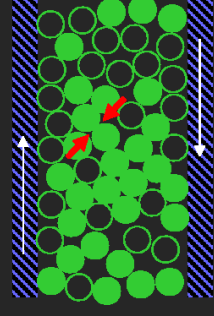
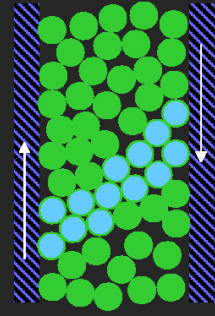
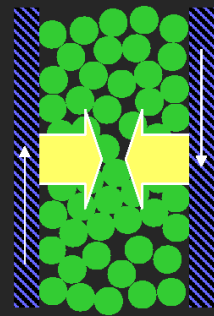
10

# First normal stress difference: no jamming



# Negative first normal stress difference

applied shear stress  $\sigma$



Negative first normal stress difference  $N_1 = \sigma_{xx} - \sigma_{yy}$  pulls plates together

Expect jamming would tend to push plates apart

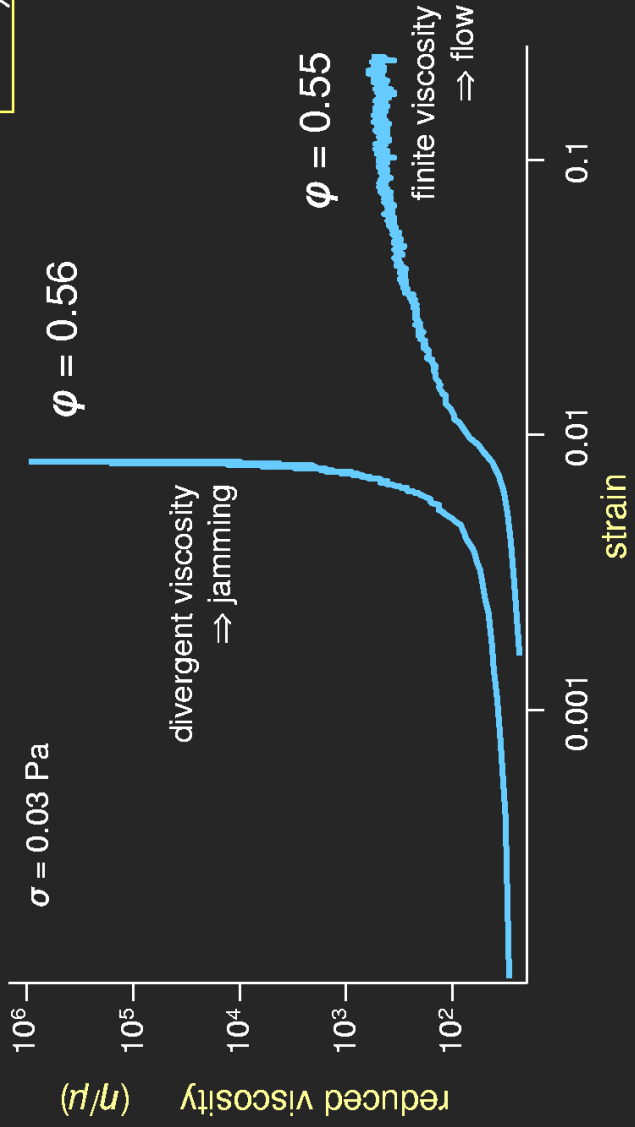
Lubrication forces can suck plates together  
(no Brownian motion)

(Laun, Brady, Morris, ...)

# Critical volume fraction for jamming

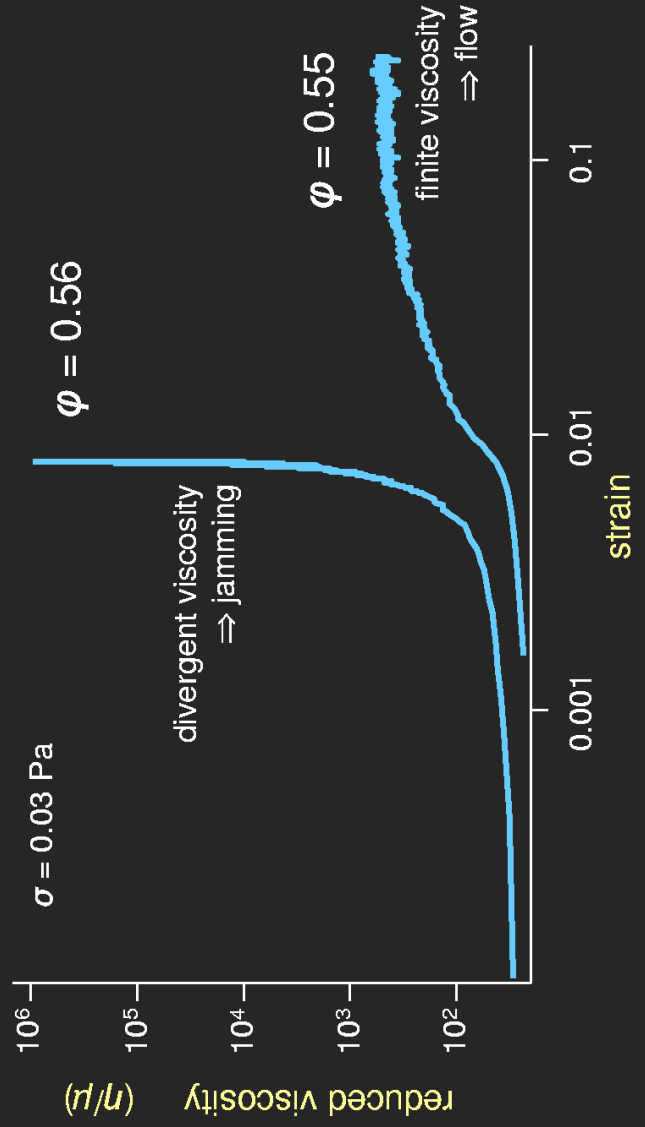
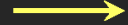
Apply constant stress  $\sigma$   
 Measure strain rate  $\dot{\gamma}$

$$\Rightarrow \eta = \frac{\sigma}{\dot{\gamma}}$$



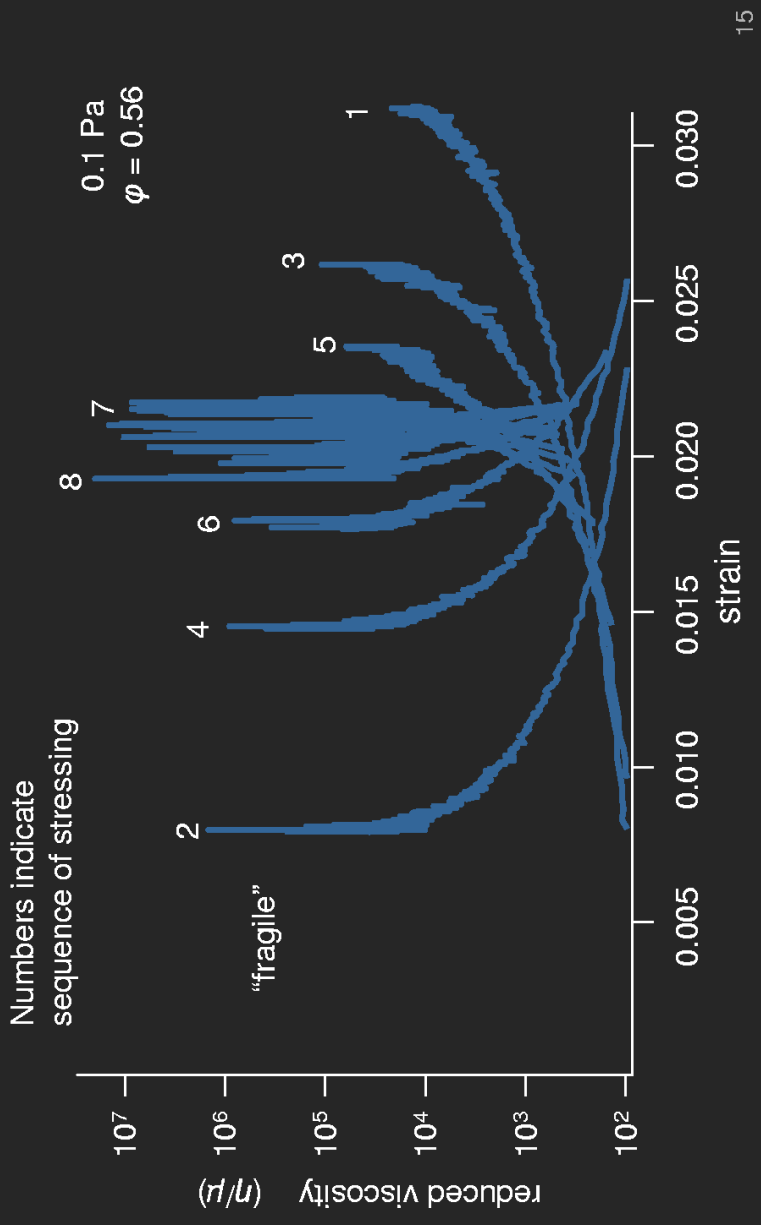
13

# Jamming case, $\phi \geq 0.56$

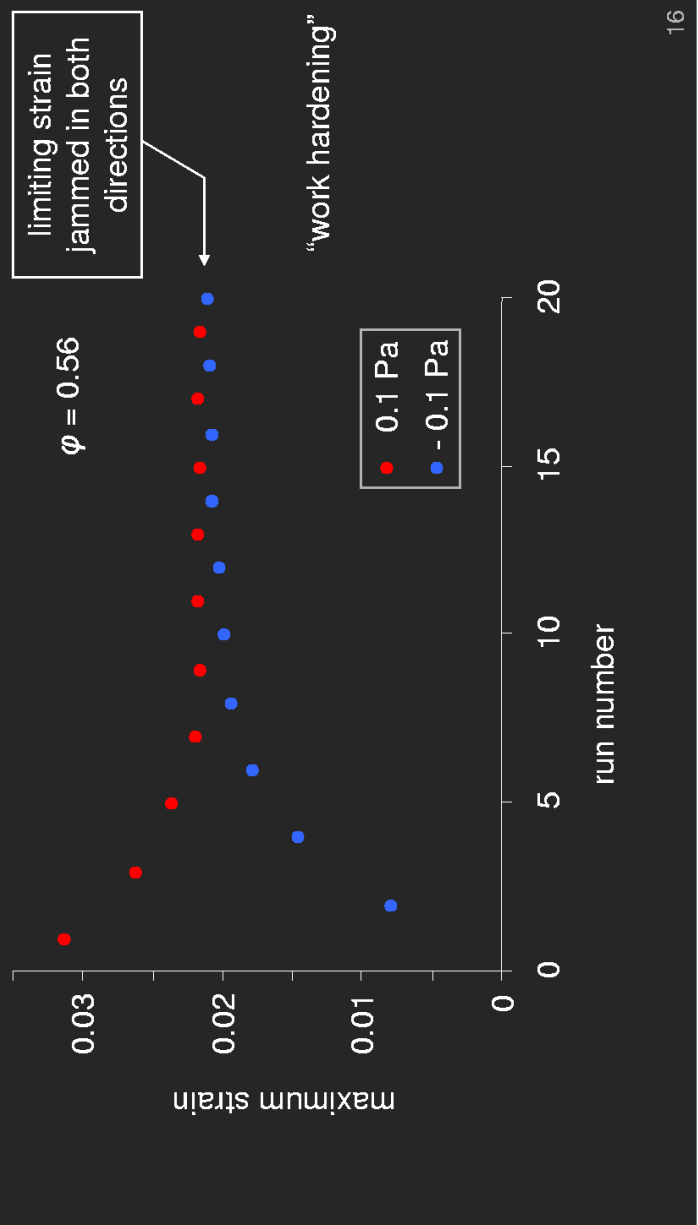


14

# Jamming reversals: low stress

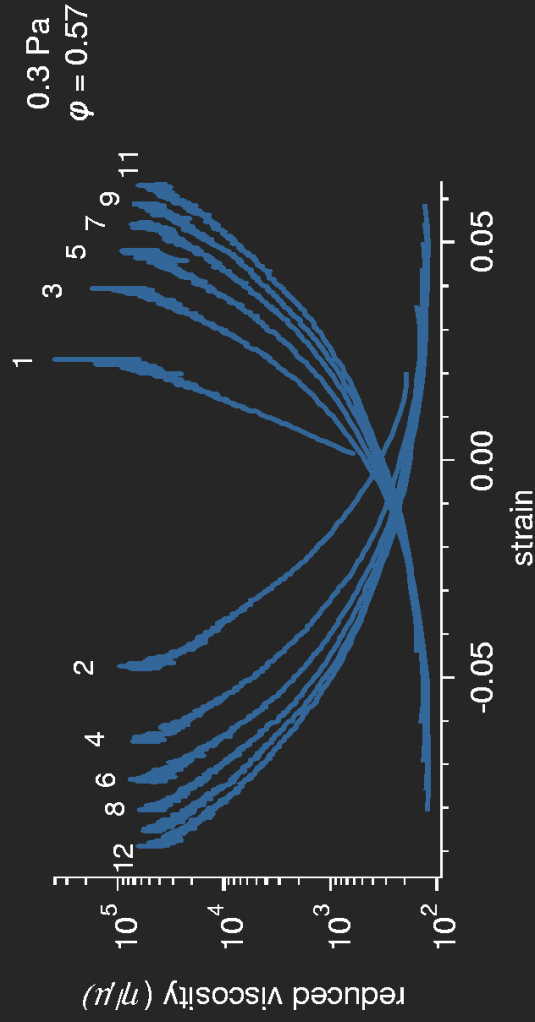


# Jamming reversals: low stress





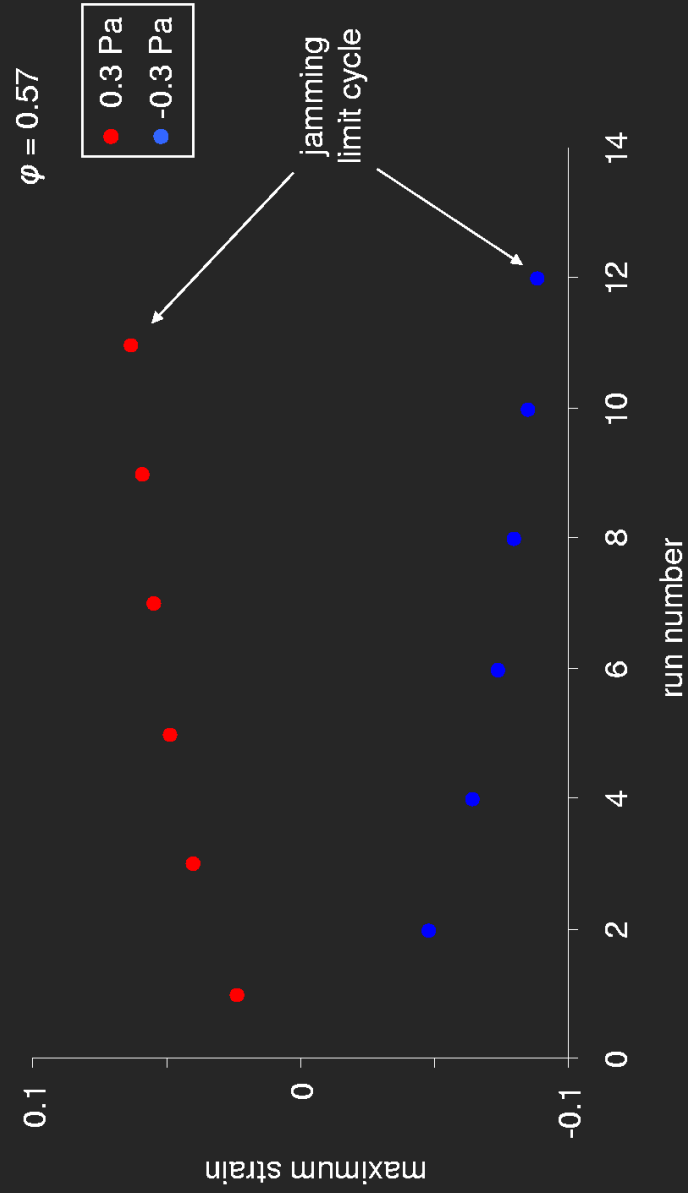
# Jamming reversals: near critical stress



Reduced viscosity rises nearly exponentially with strain  $\rightarrow \eta \sim e^{-\gamma}$

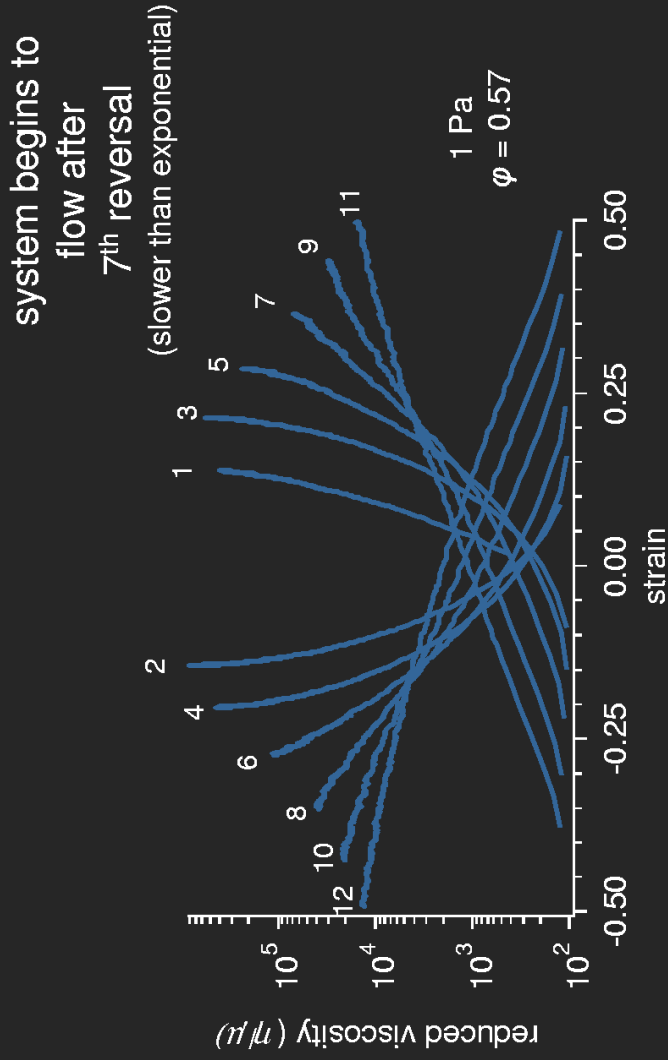
17

# Jamming reversals: $\sim$ critical stress



18

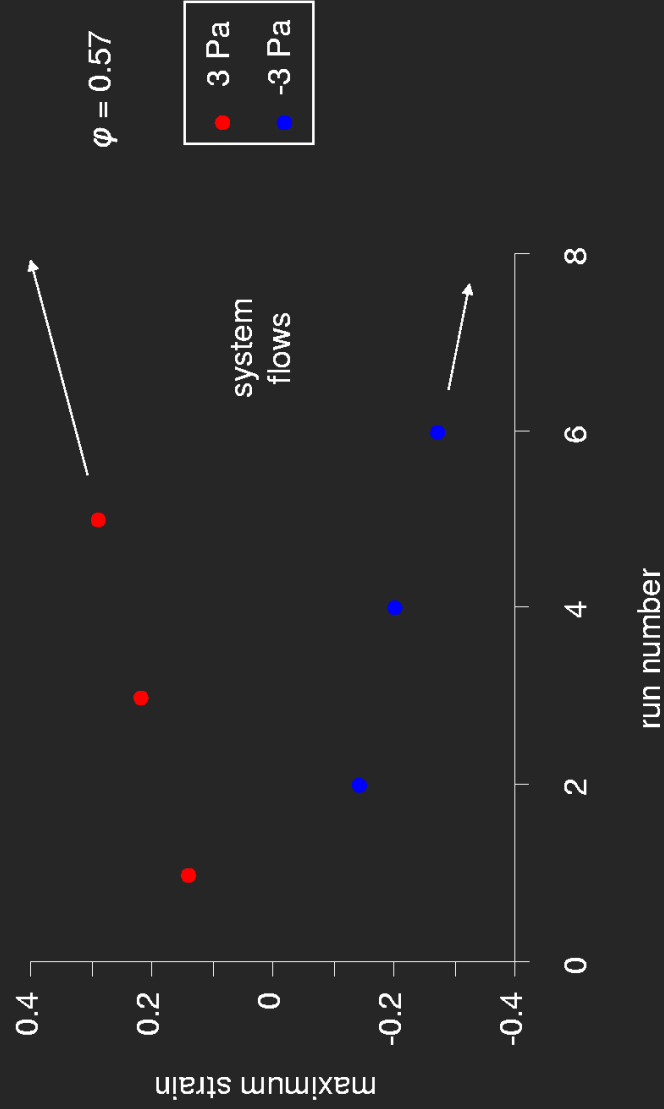
# Jamming reversals: high stress



Reduced viscosity rises slower than exponentially with strain  $\rightarrow \eta < e^{-\gamma}$

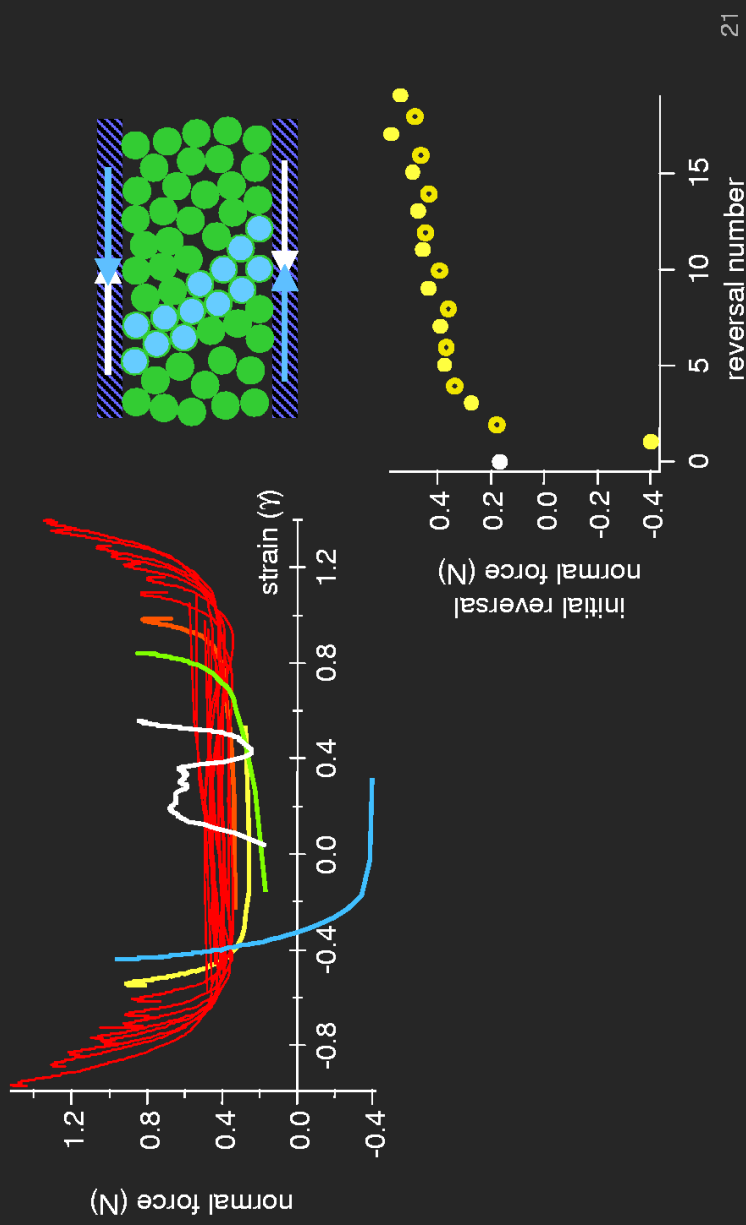
19

# Jamming reversals: high stress



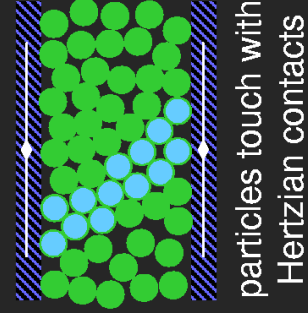
20

# Jamming normal stress



21

# Jamming both ways



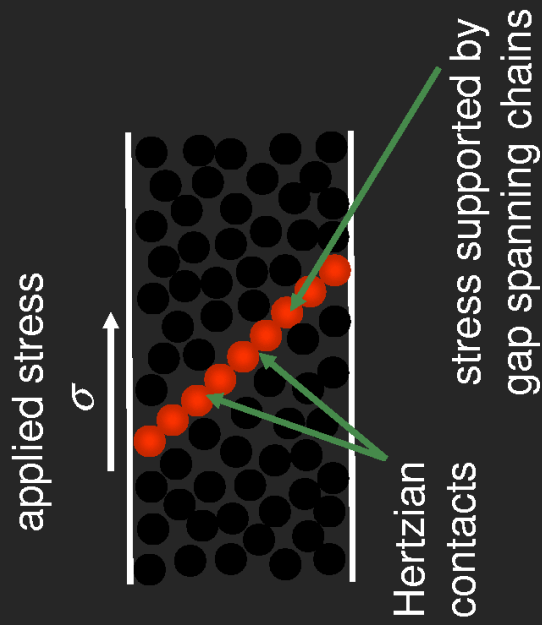
Jamming produced by flow reversal,  
not by quenching from high  $T$

*The system jams in both directions*

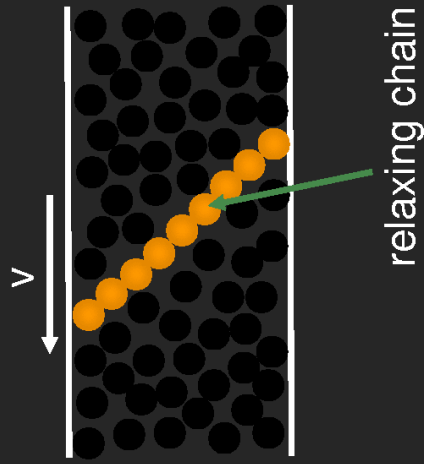
22

# Stress relaxation after jamming

stressed sample

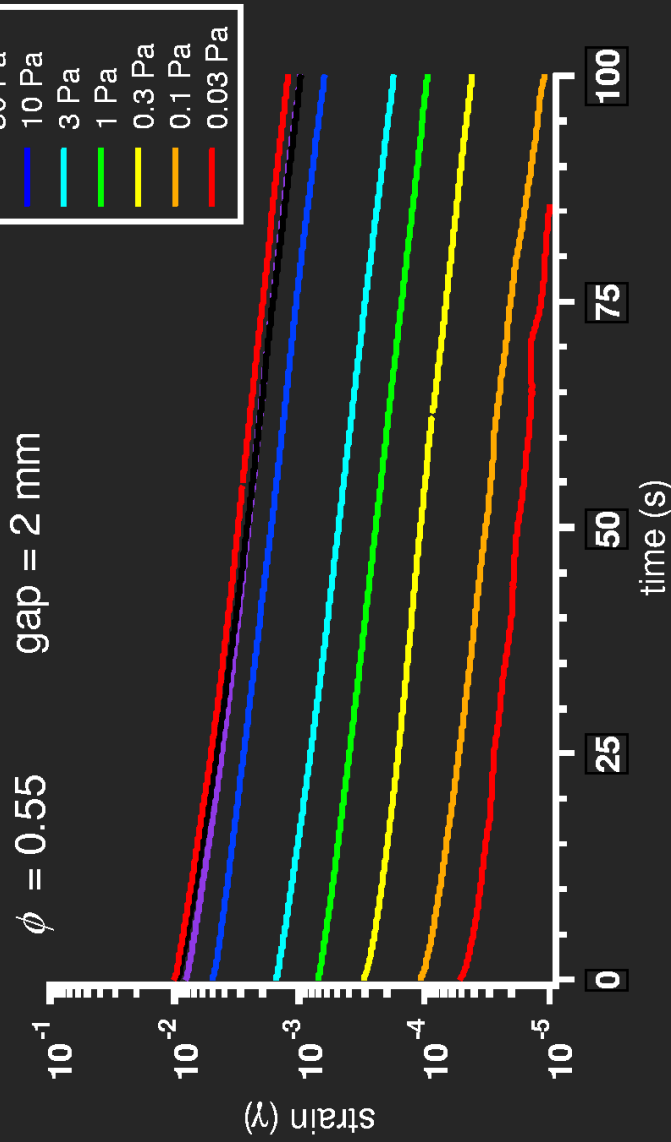


Relaxing Sample



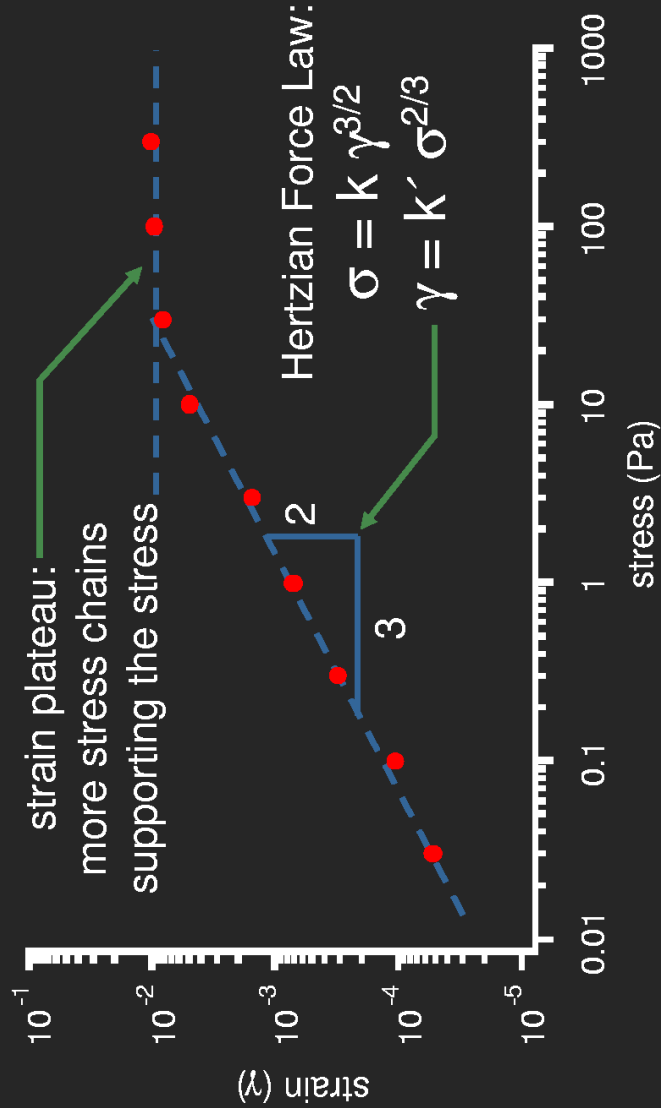
23

# Relaxation curves



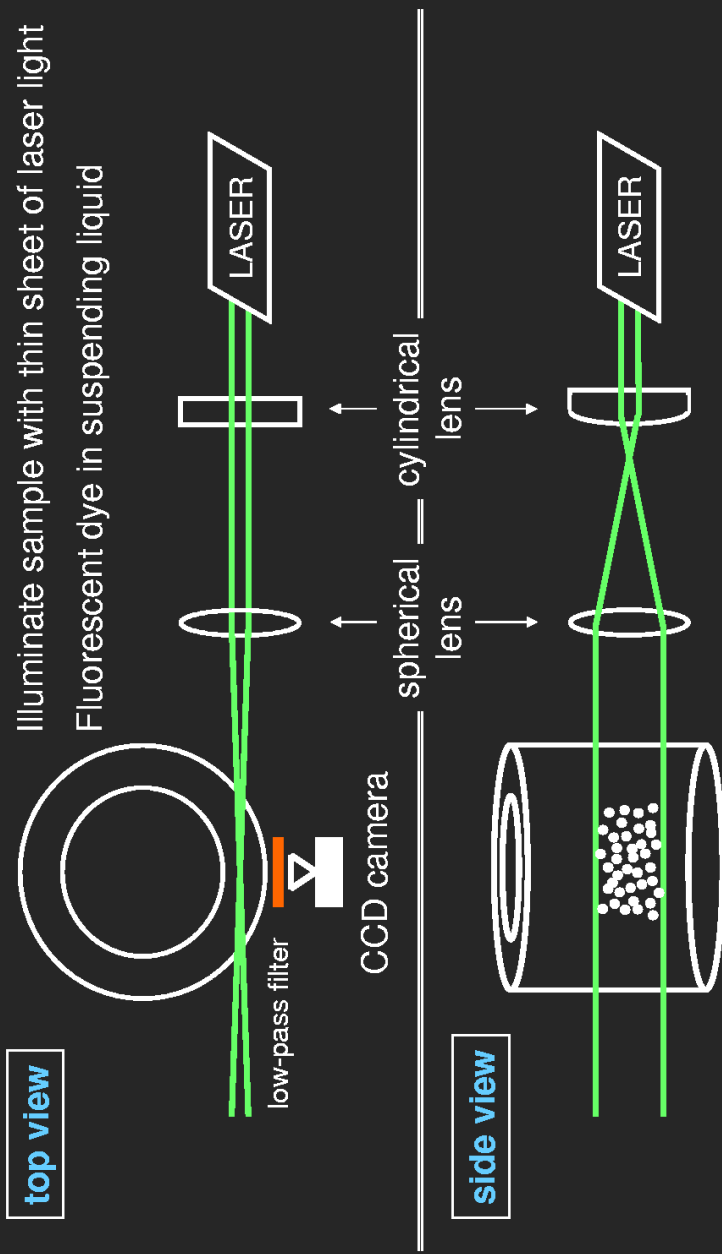
24

# Strain relaxation



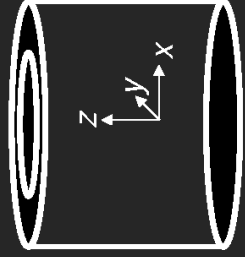
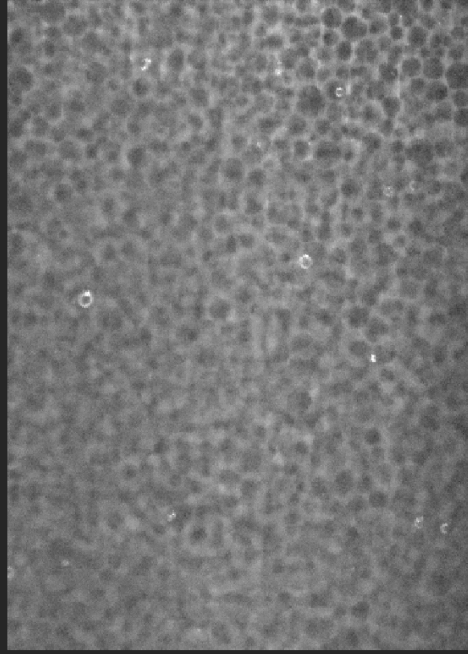
25

# Following particle trajectories



26

## 2-d image in velocity-vorticity plane

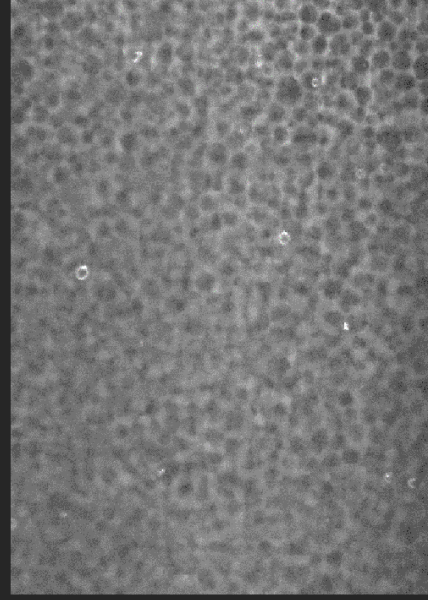
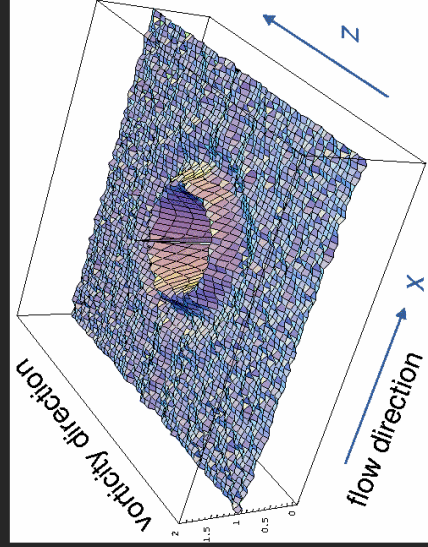


x – velocity direction  
 y – gradient direction  
 z – vorticity direction

27

## Pair correlation functions $g(r)$

jammed state - no flow

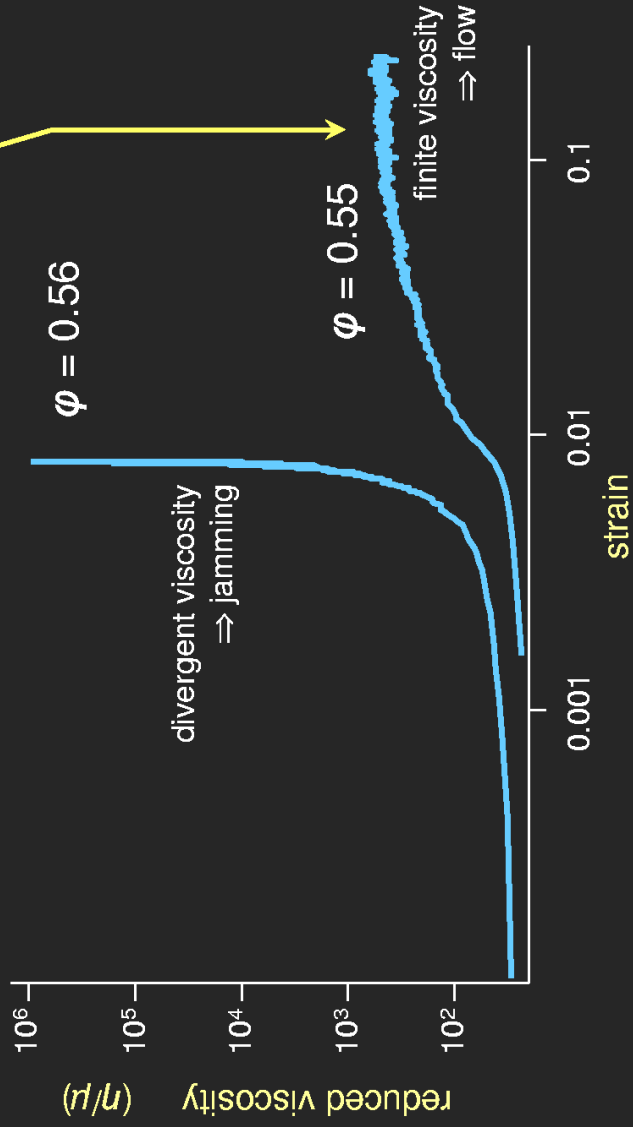


$$g(\vec{r}) = \frac{\langle n(\vec{r}_0)n(\vec{r}_0 + \vec{r}) \rangle}{\langle n(\vec{r}_0) \rangle^2}$$

random particle positions

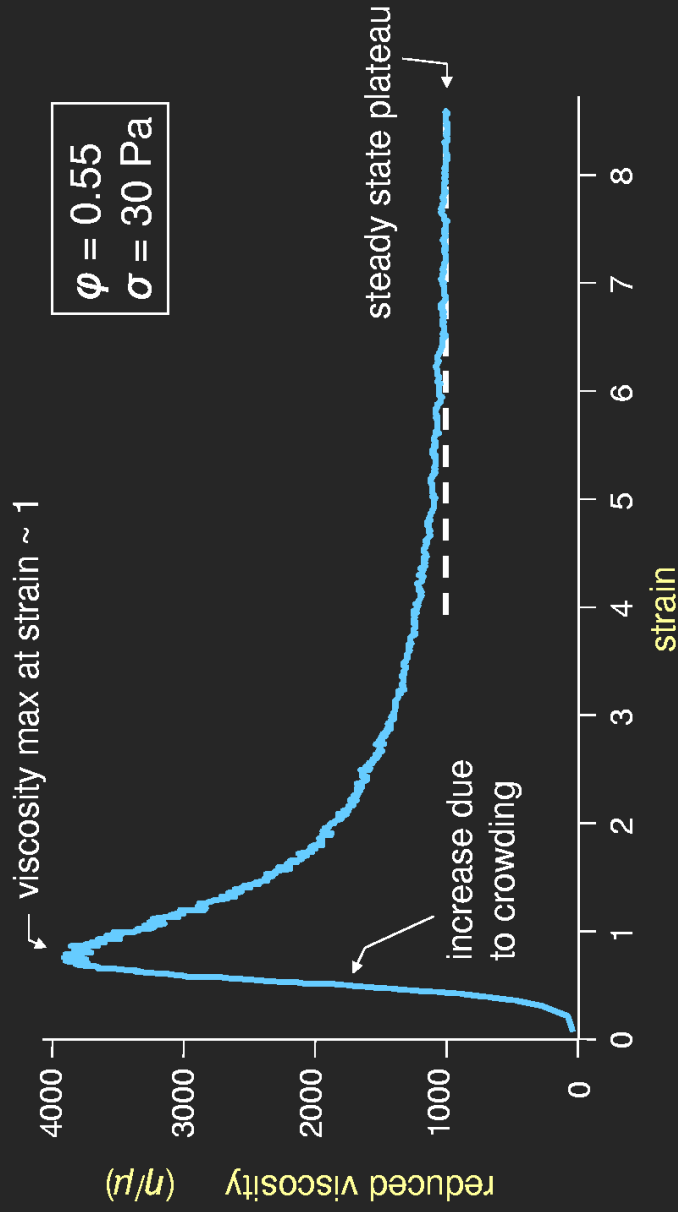
28

# Non-jamming case, $\phi \leq 0.56$



29

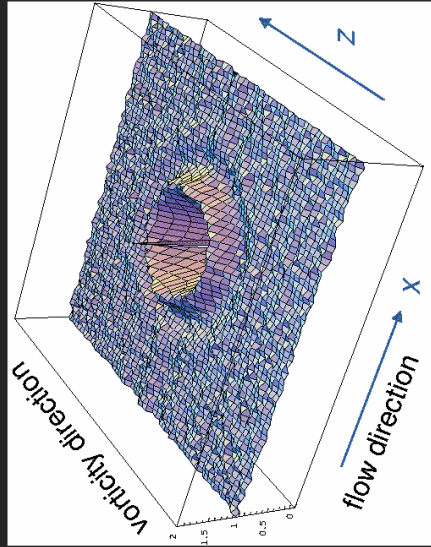
# No jamming case: transient startup response



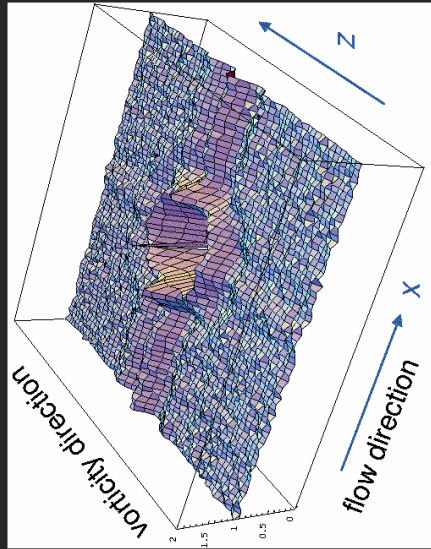
30

# Pair correlation functions $g(r)$

no flow



flow – steady state

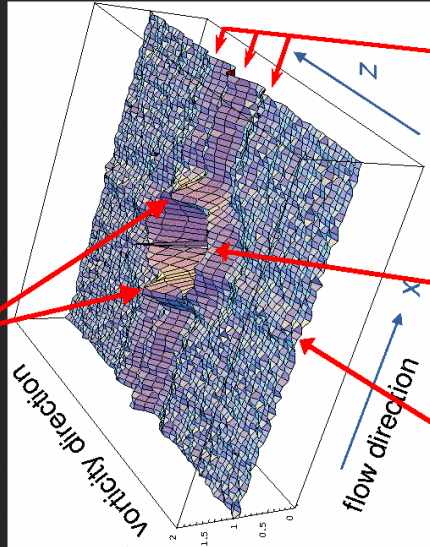


Note structure in flow direction:  
strings or layers?

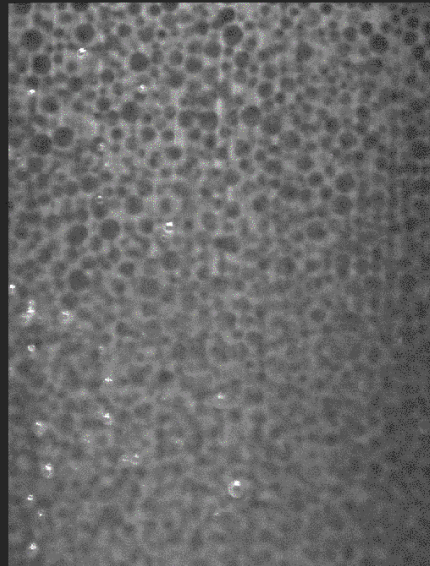
31

# Pair correlation functions $g(r)$

peaks in flow direction



flow – steady state



weak registry in vorticity plane

maxima along flow direction indicates particles line up in the flow direction

nearest neighbor ring

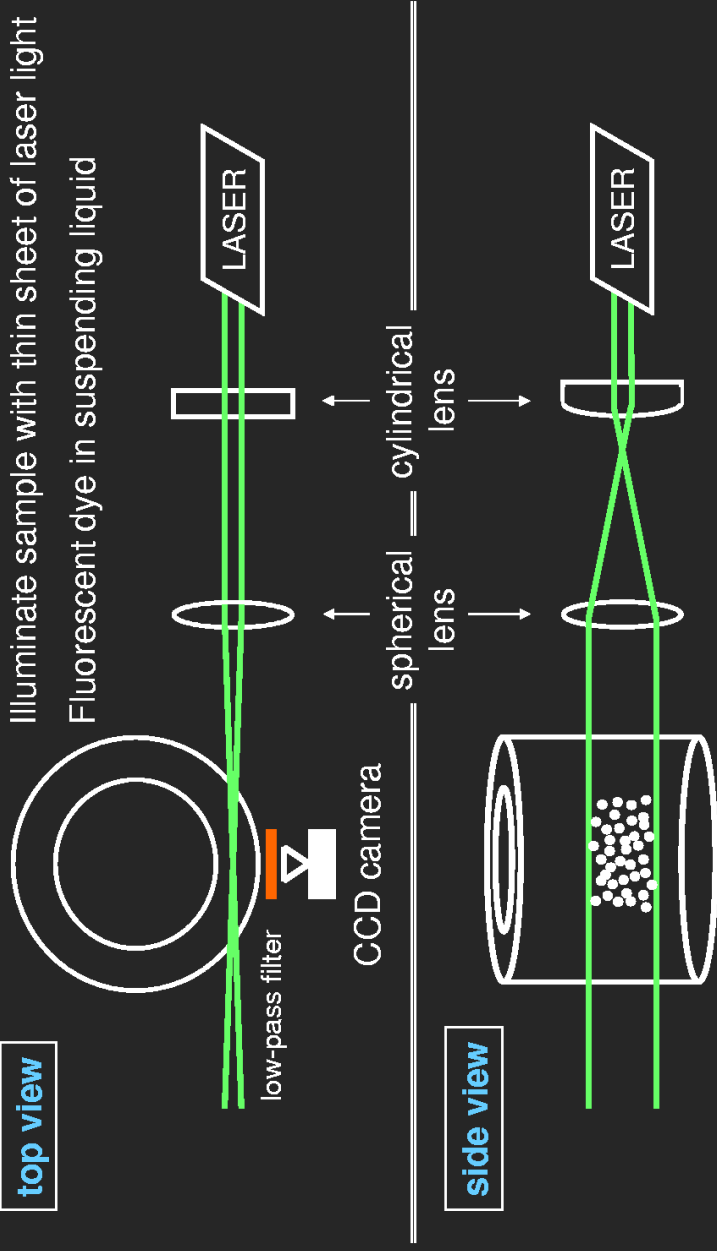
Structure in flow direction:

- strings
- layers?

32



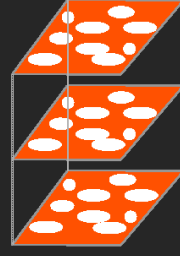
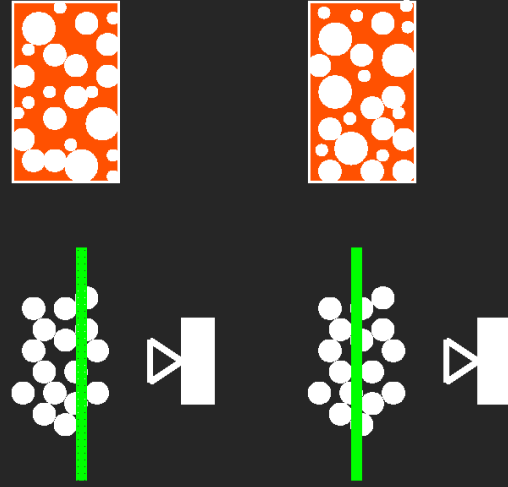
# Following particle trajectories



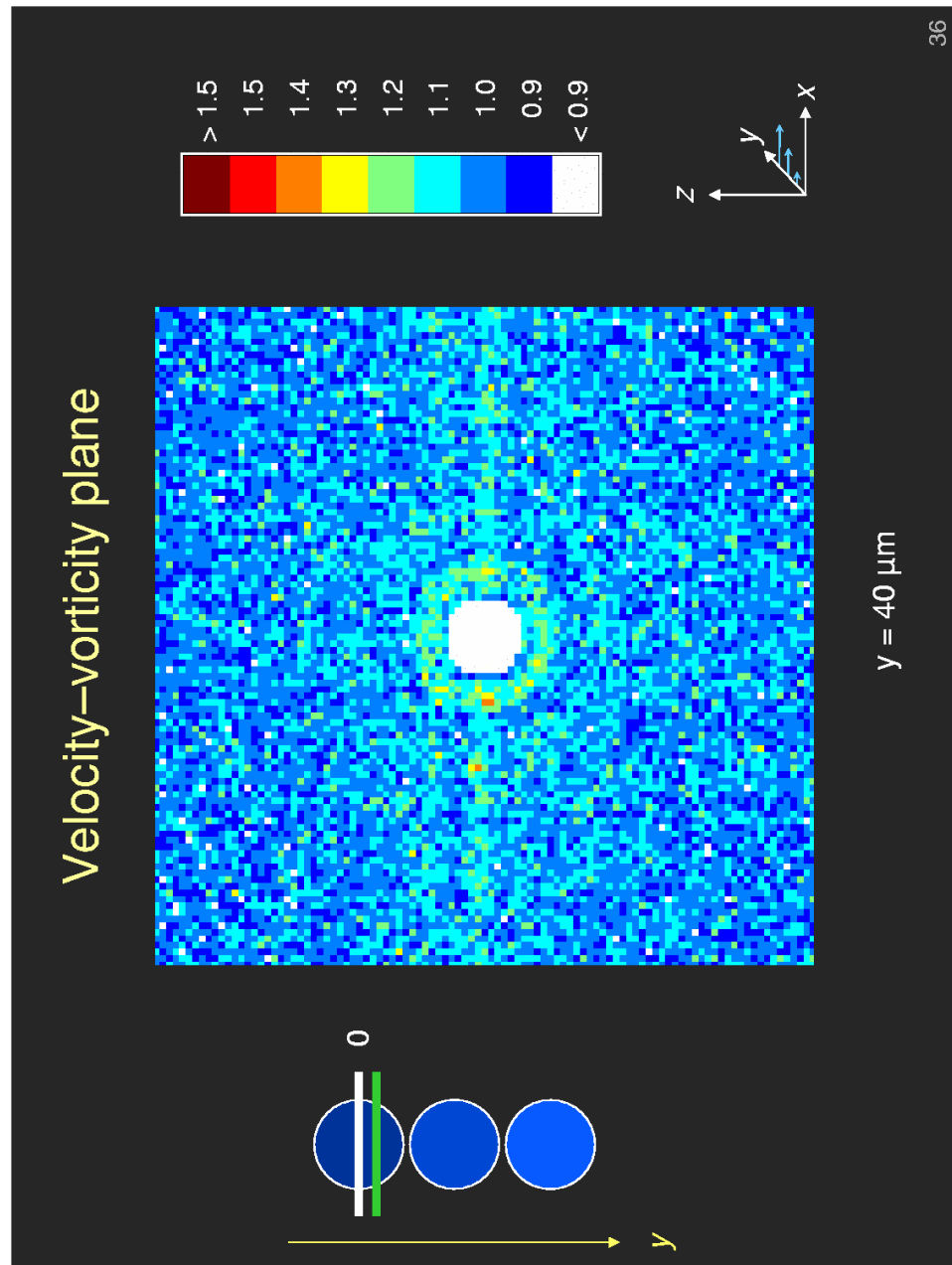
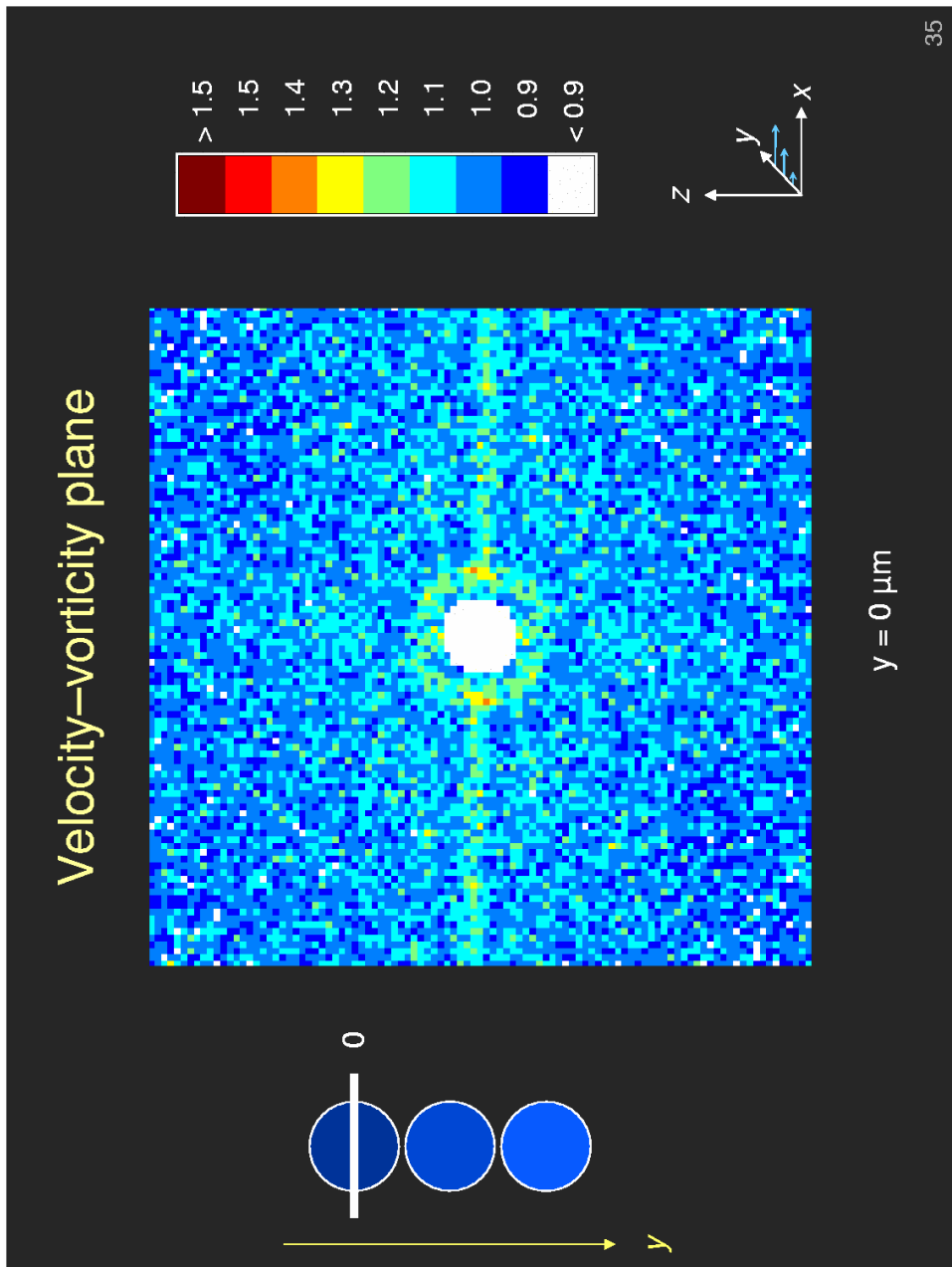
33

# 3-d: Imaging into velocity gradient direction

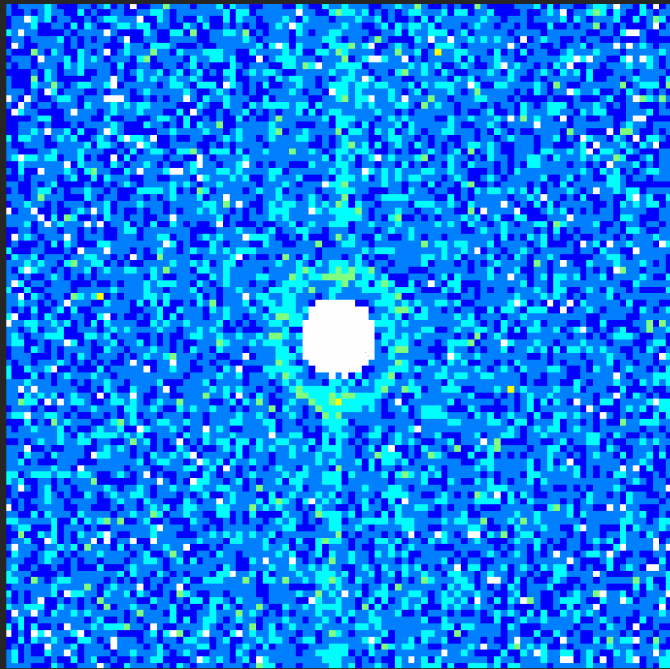
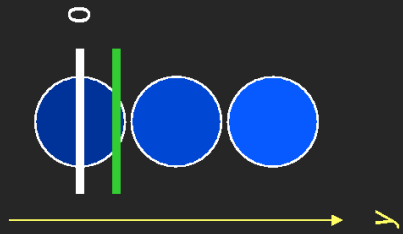
- 1) Raster the laser making a series of 2D slices
- 2) Assemble the images into a 3D vector
- 3) Find the centroid of each connected region



34



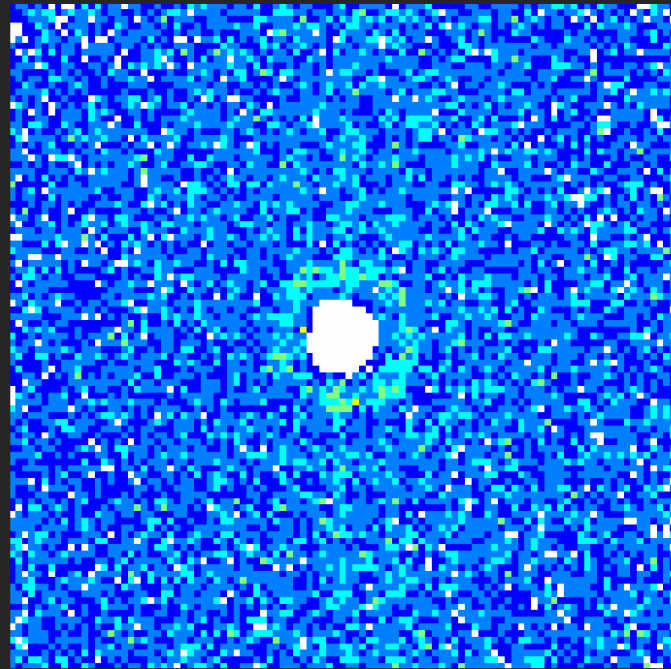
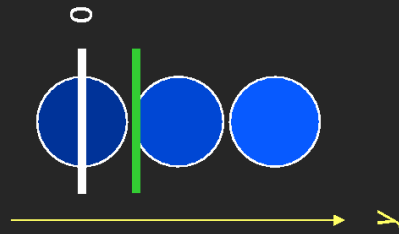
Velocity–vorticity plane



$y = 80 \mu\text{m}$

37

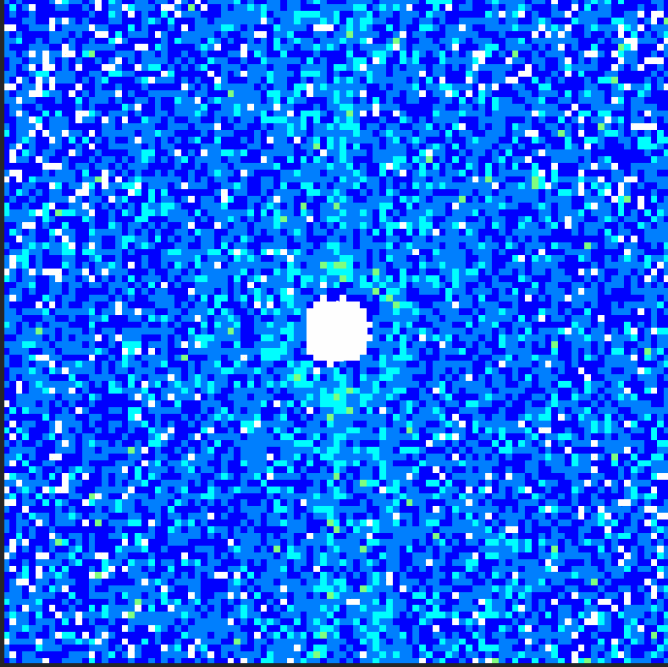
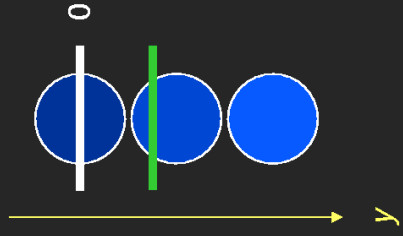
Velocity–vorticity plane



$y = 120 \mu\text{m}$

38

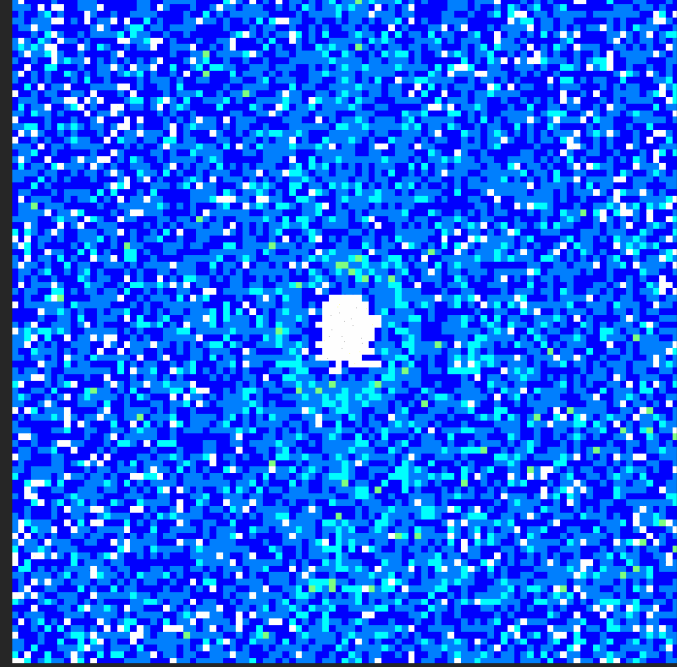
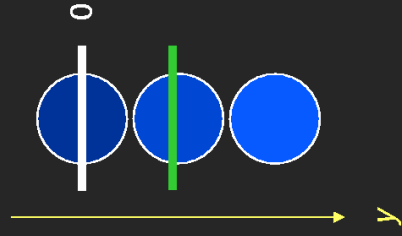
Velocity–vorticity plane



$y = 160 \mu\text{m}$

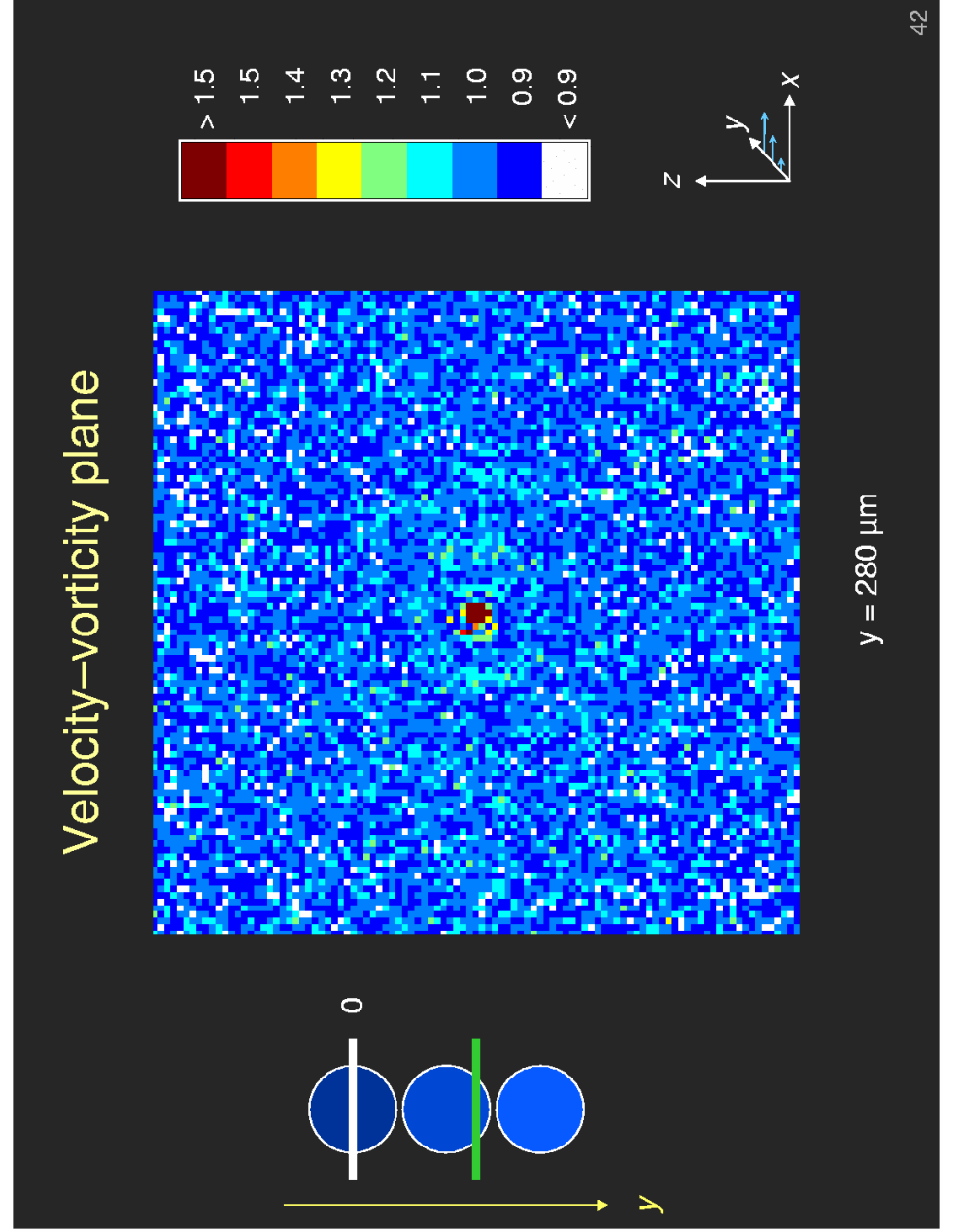
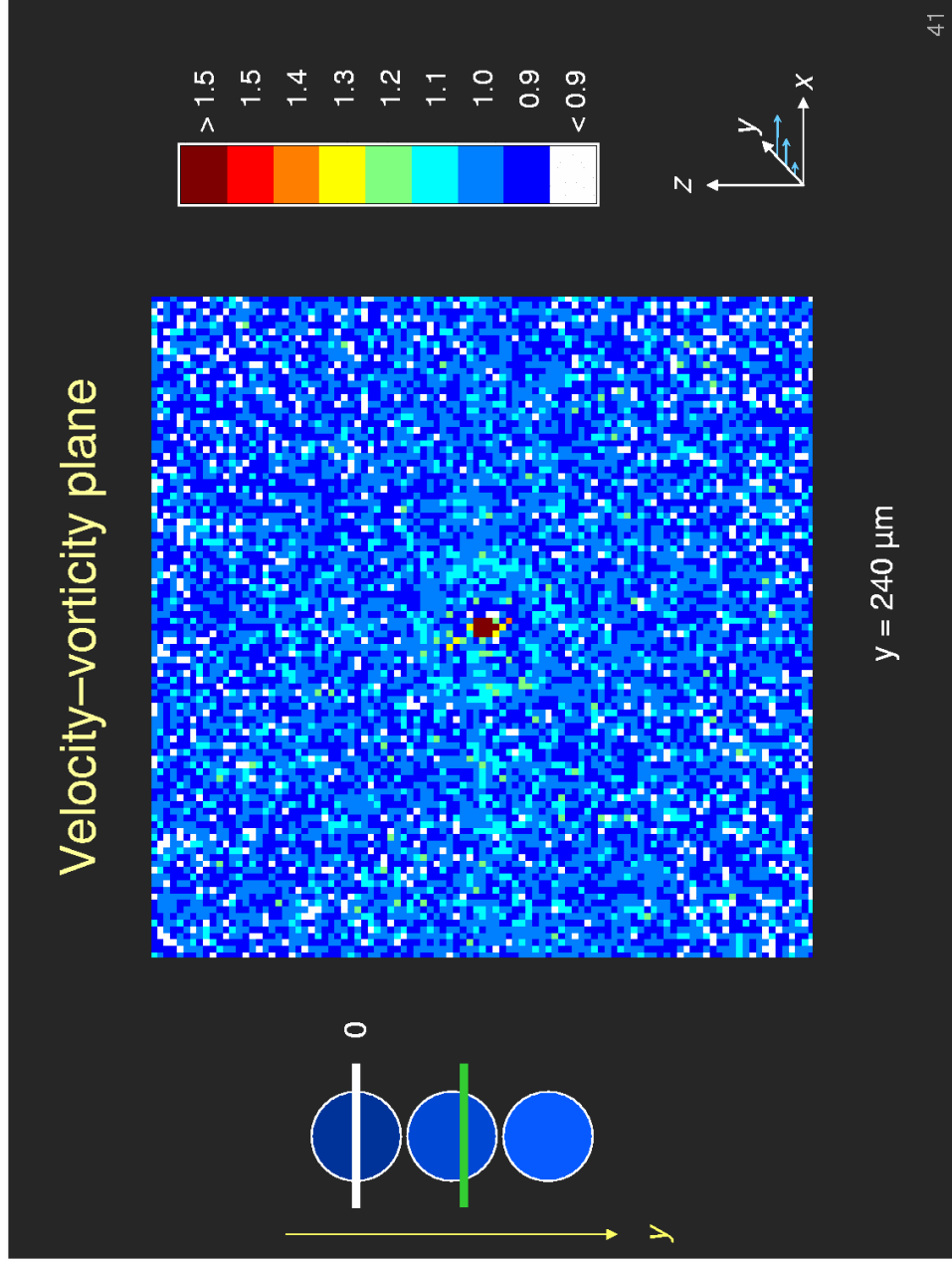
39

Velocity–vorticity plane

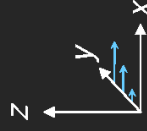
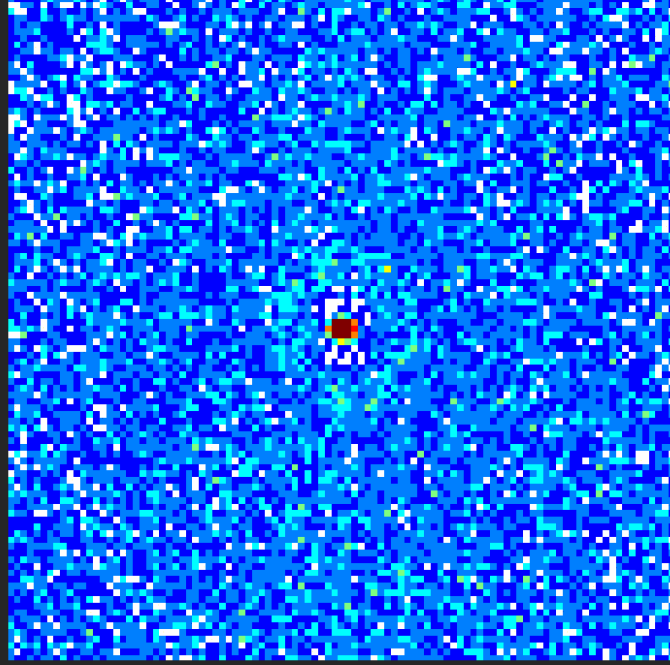
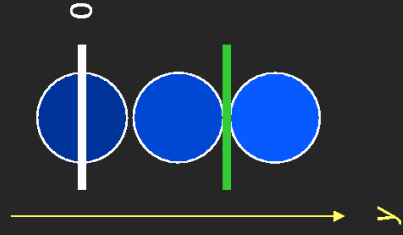


$y = 200 \mu\text{m}$

40



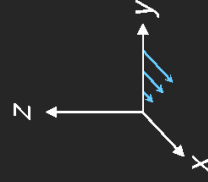
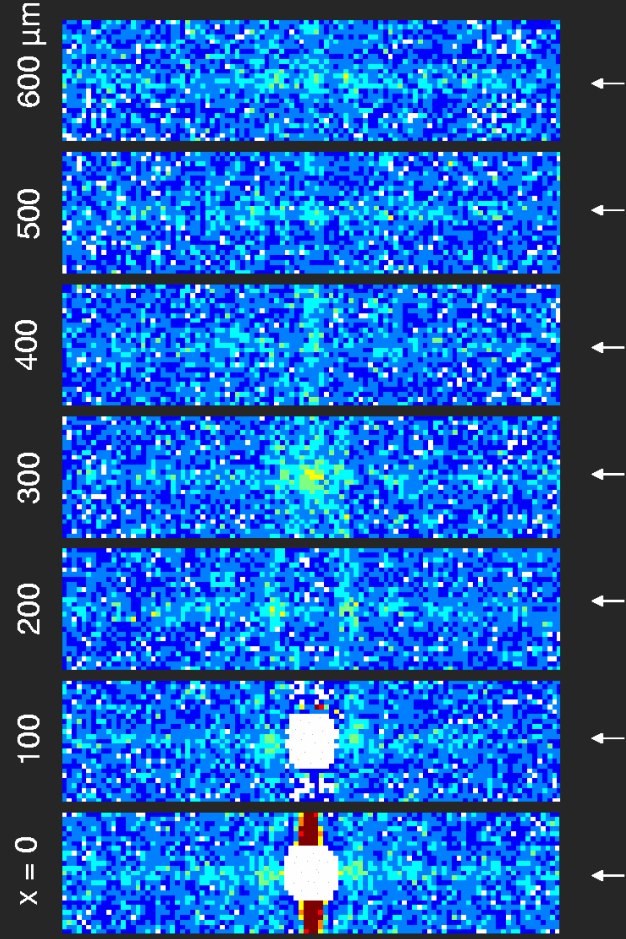
Velocity–vorticity plane



$y = 320 \mu\text{m}$

43

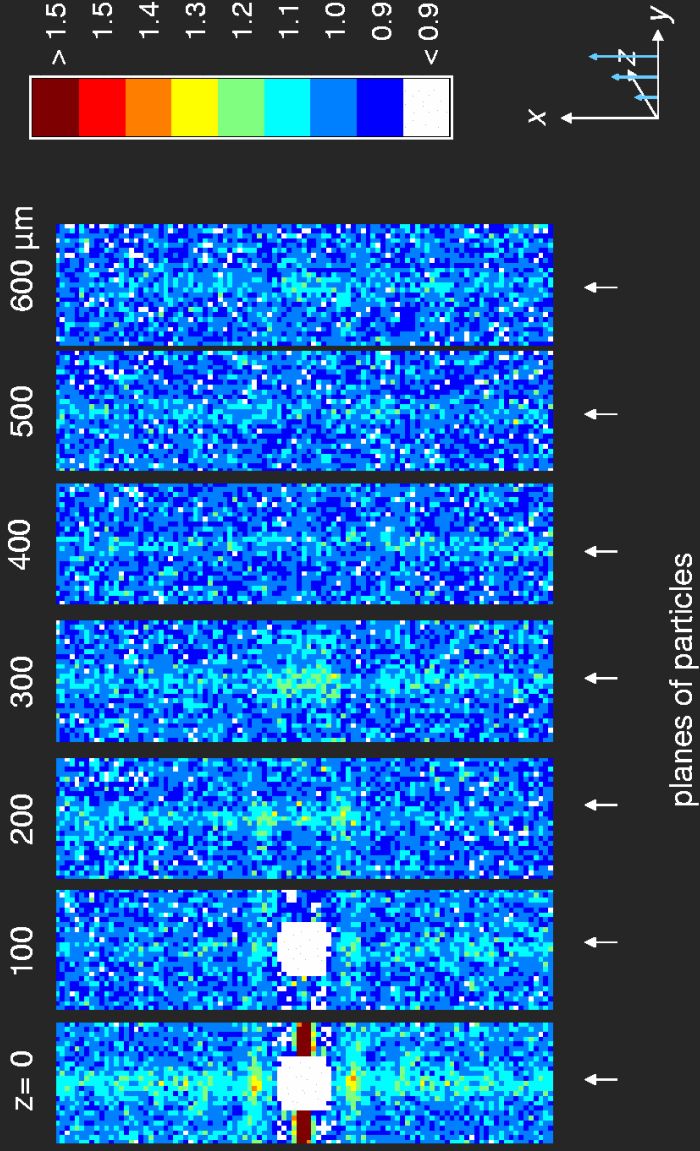
Gradient–vorticity plane



planes of particles have formed

44

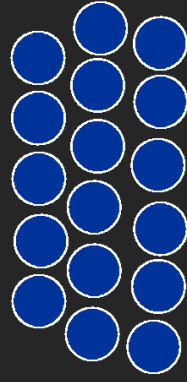
## Velocity–gradient plane



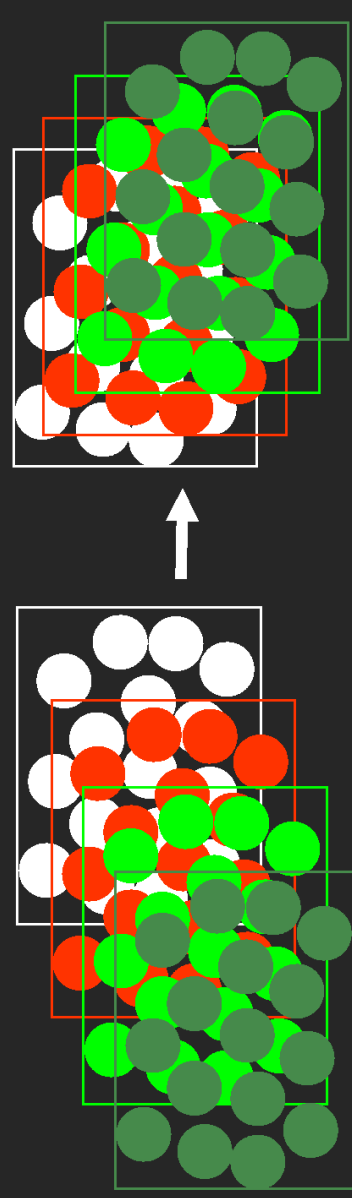
45

## Summary of flow

1) Particles tend to form strings in the flow direction

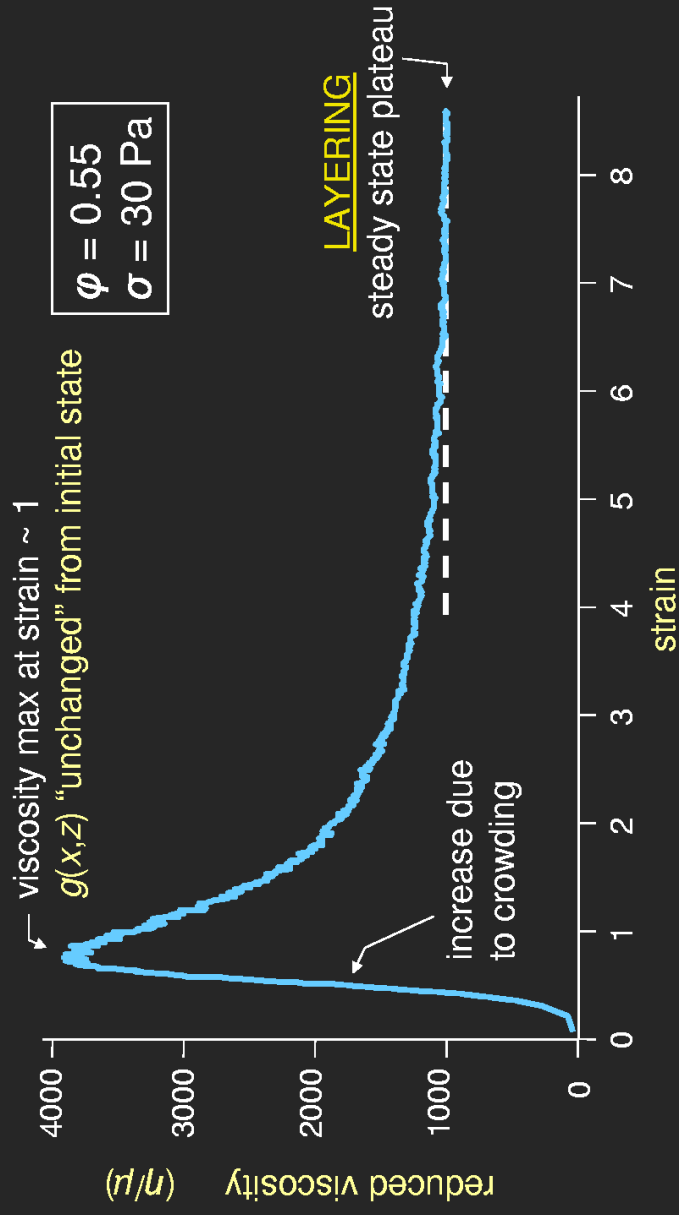


2) Planes of particles form to facilitate flow



46

## Layering $\rightarrow$ steady state viscosity



47

## Summary

- Non-colloidal suspensions jam above a (gap-dependent) volume fraction
- Suspensions jam both ways ("work hardening")
- Flowing dense suspensions organize into "strings" in the flow direction
- These "strings" are organized into sheets in the velocity-vorticity plane
- $\sigma_c$ : Jamming occurs only below a critical stress

48



## Acknowledgements

- Kurt Knipmeyer
- Keck Foundation
- NSF MRSEC Program at UCSB