

# Higgs phases and zero-energy states in graphene

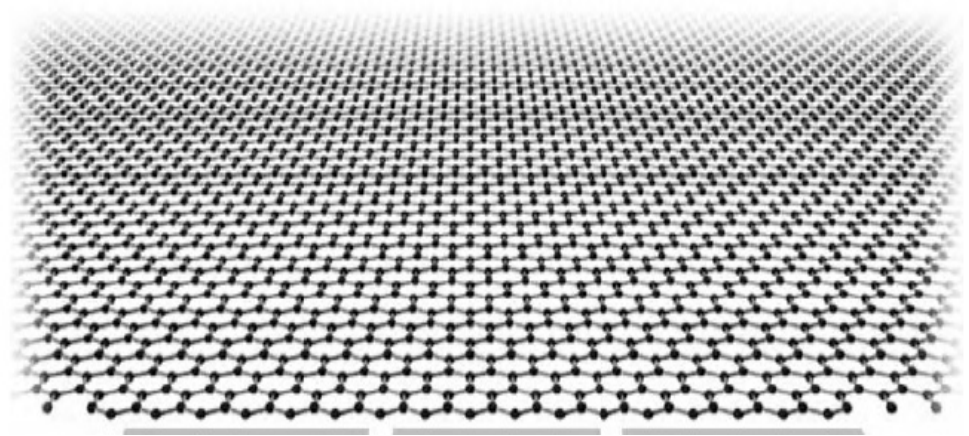
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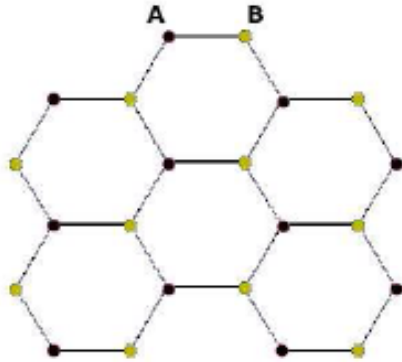
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## Graphene: 2D carbon (1s<sup>2</sup>, 2s<sup>2</sup>, 2p<sup>2</sup>)



Two triangular sublattices: A and B; one electron per site (half filling)

Tight-binding model (  $t = 2.5$  eV ):

$$H_0 = -t \sum_{\vec{A}, i, \sigma = \pm 1} u_{\sigma}^{\dagger}(\vec{A}) v_{\sigma}(\vec{A} + \vec{b}_i) + H.c.$$

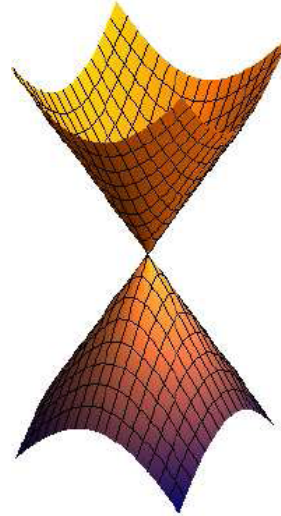
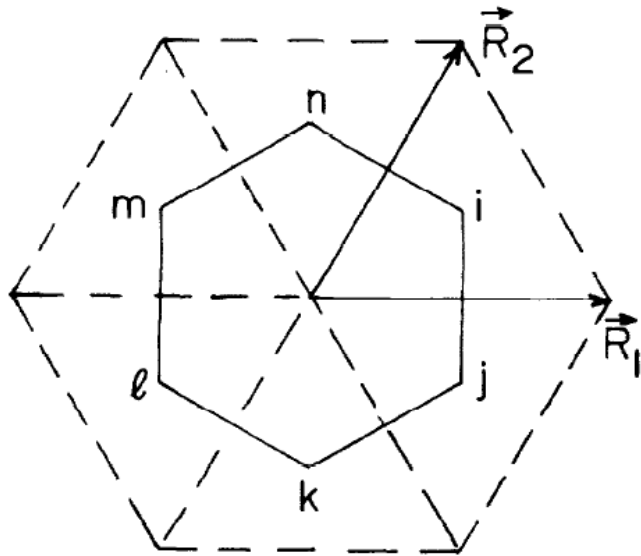
(Wallace, PR , 1947)

$$E(\vec{k}) = \pm t \left| \sum_i \exp[\vec{k} \cdot \vec{b}_i] \right|$$

The sum is complex  $\Rightarrow$  two equations for two variables for zero energy

$\Rightarrow$  **Dirac points** (no Fermi surface)

**Brillouin zone:**



Two inequivalent (Dirac) points at :

**+K** and **-K**

**Dirac fermion:**

$$\Psi_{\sigma}^{\dagger}(\vec{x}, \tau) = T \sum_{\omega_n} \int^{\Lambda} \frac{d\vec{q}}{(2\pi a)^2} e^{i\omega_n \tau + i\vec{q} \cdot \vec{x}} (u_{\sigma}^{\dagger}(\vec{K} + \vec{q}, \omega_n), v_{\sigma}^{\dagger}(\vec{K} + \vec{q}, \omega_n), u_{\sigma}^{\dagger}(-\vec{K} + \vec{q}, \omega_n), v_{\sigma}^{\dagger}(-\vec{K} + \vec{q}, \omega_n))$$

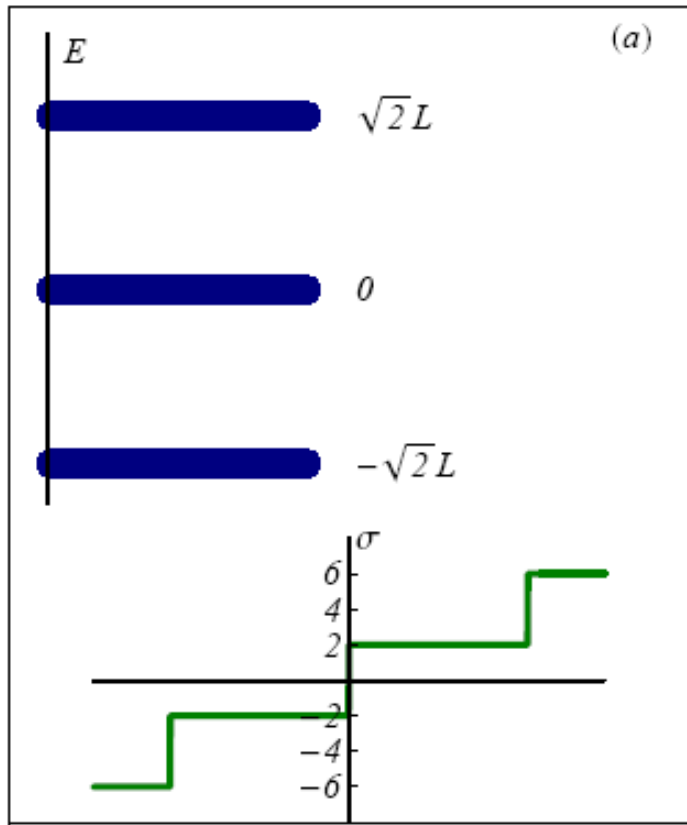
“Low - energy” Hamiltonian:  $H_0 = i\gamma_0 \gamma_i (-i\partial_i - A_i) \quad i=1,2$

$$\{\gamma_{\mu}, \gamma_{\nu}\} = 2\delta_{\mu\nu} \quad \nu, \mu = 0, 1, 2$$

( $v = c/300 = 1$ , in our units)

## Experiment: how do we detect Dirac fermions?

**Quantum Hall effect** (for example):  $L = \sqrt{\hbar v_F^2 |eB|/c}$



Landau levels: each is

$2$  (spin)  $\times$   $2$  (Dirac)  $\times$   $eB$  (Area)/ $hc$

degenerate  $\Rightarrow$  quantization in steps of

**four!**

(Gusynin and Sharapov, PRL, 2005)

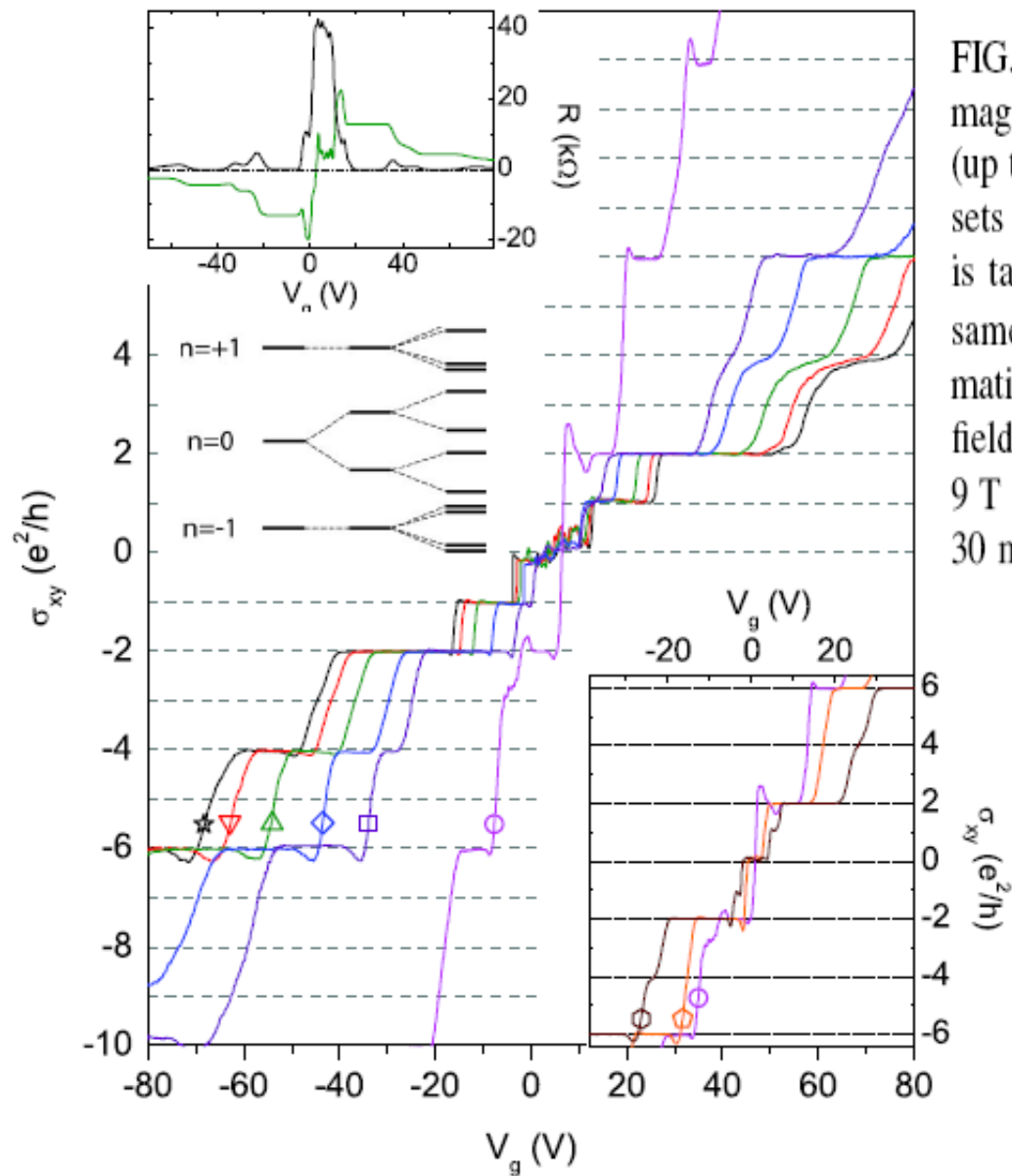


FIG. 2 (color online).  $\sigma_{xy}$ , as a function of  $V_g$  at different magnetic fields: 9 T (circle), 25 T (square), 30 T (diamond), 37 T (up triangle), 42 T (down triangle), and 45 T (star). All the data sets are taken at  $T = 1.4$  K, except for the  $B = 9$  T curve, which is taken at  $T = 30$  mK. Left upper inset:  $R_{xx}$  and  $R_{xy}$  for the same device measured at  $B = 25$  T. Left lower inset: a schematic drawing of the LLs in low (left) and high (right) magnetic field. Right inset: detailed  $\sigma_{xy}$  data near the Dirac point for  $B = 9$  T (circle), 11.5 T (pentagon), and 17.5 T (hexagon) at  $T = 30$  mK.

(Y. Zhang, PRL, 2006)

# Symmetries: exact and emergent

## 1) Lorentz

(microscopically, only  $Z_2$  (A  $\leftrightarrow$  B)  $\times$   $Z_2$  (K  $\leftrightarrow$  -K) = **D2**, dihedral group)

2) **Chiral**:  $SU(2) = \{\gamma_{35}, \gamma_3, \gamma_5\}$ ,  $\gamma_{35} = i\gamma_3\gamma_5$

Generators commute with the Dirac Hamiltonian (in 2D). Only two are **emergent!**

3) **Time-reversal (exact)**:  $I_t = U_t K$ ,  $U_t = i\gamma_1\gamma_5$

( + K  $\leftrightarrow$  -K and complex conjugation )

(IH, Juricic, Roy, PRB, 2009)

4) **Particle-hole (supersymmetry)**: anticommute with Dirac Hamiltonian

$$\vec{M} = (\gamma_0, i\gamma_0\gamma_3, i\gamma_0\gamma_5) , \quad \tilde{M} = i\gamma_1\gamma_2$$

$$\Rightarrow \psi_{-E} = M \psi_E$$

and so map zero-energy states, when they exist, into each other!

The zero-energy subspace is invariant under both symmetry (commuting) and supersymmetry (anticommuting) operators.

## “Masses” = supersymmetries

- 1) **Broken chiral symmetry, preserved time reversal**

$$H_0 \rightarrow H_0 + \vec{m} \cdot \vec{M}$$

- 2) **Broken time reversal symmetry, preserved chiral**

$$H_0 \rightarrow H_0 + \tilde{m} \tilde{M}$$

In either case the spectrum becomes **gapped**:

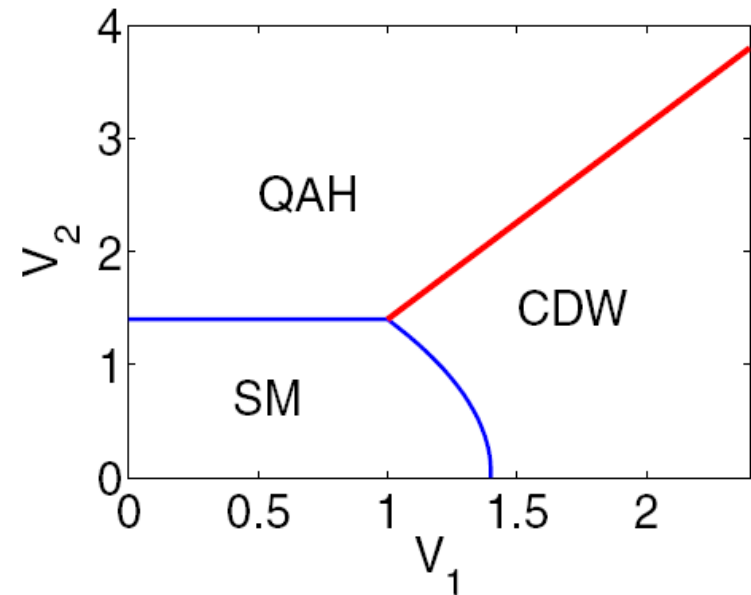
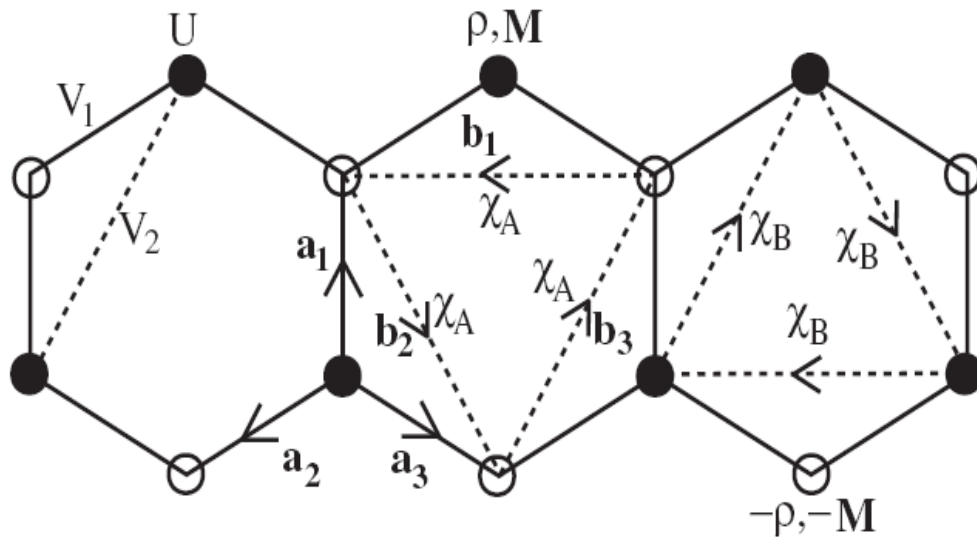
$$\varepsilon_{\pm}(\mathbf{p}) = \pm \sqrt{|\mathbf{p}|^2 + |\Delta_0|^2}, \quad \Delta_0 = m, \tilde{m}$$



## On lattice?

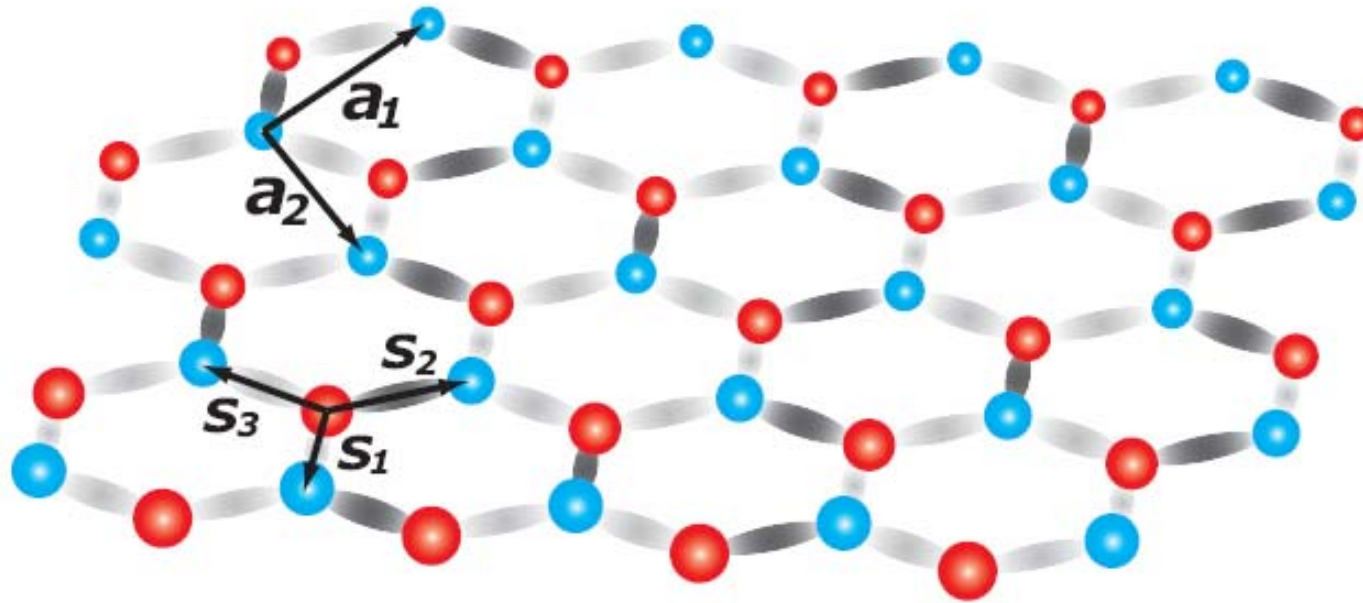
1)  $m \gamma_0$  **staggered density, or Neel (with spin); preserves translations (Semenoff, PRL, 1984)**

2)  $\tilde{m} \tilde{M}$  **circulating currents (Haldane, PRL, 1988)**



( Raghunathan et al, PRL, 2008, generic phase diagram IH, PRL, 2006 )

3)  $m_1 i\gamma_0\gamma_3 + m_2 i\gamma_0\gamma_5$  **Kekule** hopping pattern

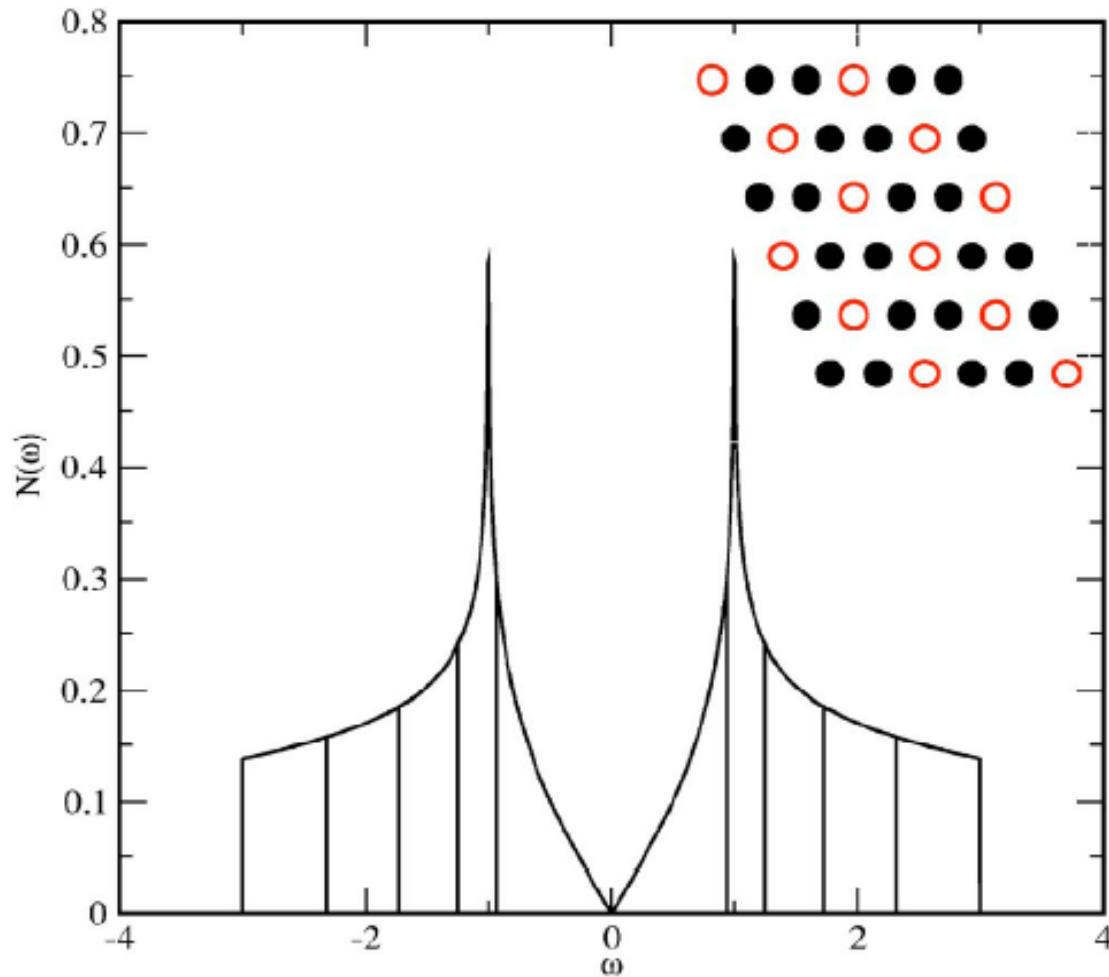


( Hou, Chamon, Mudry, PRL, 2007)

# “Catalysis” of order: magnetic and otherwise

Strong interactions are needed for the gap because there are very few states near the Fermi level:

$$U_c/t \approx 5.5$$



Density of states is linear near Dirac point:

$$N(\omega) \propto \omega^{(2-z)/z}$$

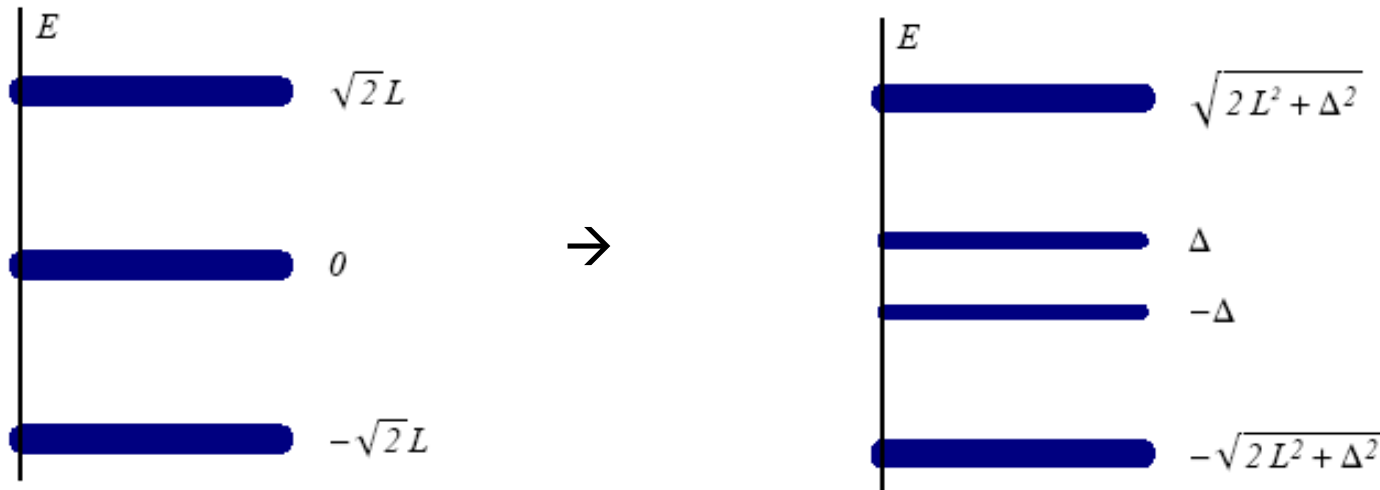
and  $z = 1$

In the magnetic field: Landau quantization

=>

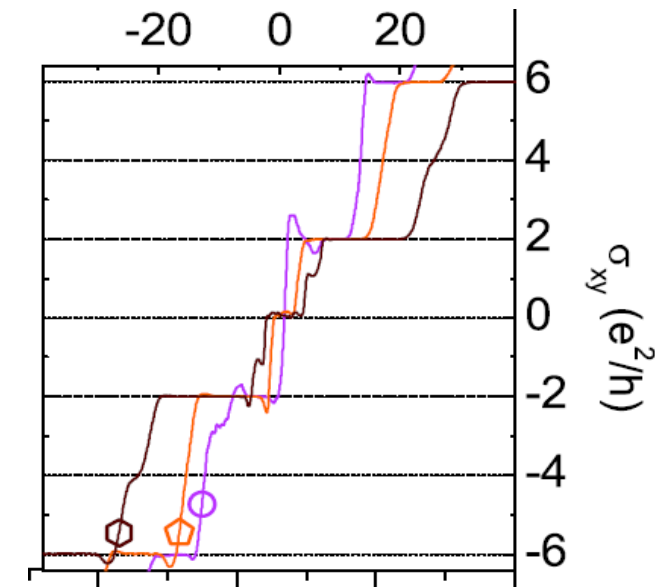
DOS infinite

=> **critical interaction infinitesimal**  
(Gusynin, Miransky, Shovkovy, PRL, 1994)



Zero-energy level is split, others only shifted  
=> quantum Hall effect at new filling factor at  
**zero!**

IH, PRB 2006, PRB 2007, IH and Roy, PRB 2008



## But which order is catalyzed?

The Hamiltonian is :  $H[A^0] = i\gamma_0\gamma_i(p_i - A_i^0)$

so that,  $H^2[A^0] = (p_i - A_i^0)^2 + \tilde{M}\epsilon_{ij}\partial_i A_j^0$

where,  $\tilde{M} = i\gamma_1\gamma_2$  is TRS breaking mass.

so all zero energy states have the same eigenvalue (+1 or -1)  $\tilde{M}$   
of  
for a uniform magnetic field.

For a non-uniform magnetic field:

$$H[A^0, 0] = e^{-\chi(\vec{x})\tilde{M}} H[0, 0] e^{-\chi(\vec{x})\tilde{M}},$$

where  $A_i^0 = \epsilon_{ij} \partial_j \chi$ . (In Coulomb gauge,  $\partial_i A_i^0 = 0$ .)

Zero-energy states with and without magnetic field are simply related:

$$\Phi_{0,n}[A^0](\vec{x}) \propto e^{\chi(\vec{x})\tilde{M}} \Phi_{0,n}[0](\vec{x})$$

since at large distance, for a localized flux  $F$

$$\chi(\vec{x}) = F \ln|\vec{x}|$$

normalizability requires them to be **-1** eigenstates of  $\tilde{M}$  !!

Since, however,

$$\text{Tr}_0 \vec{M} = 0$$

half of zero-energy states have +1, and half -1 eigenvalue of  $\vec{M}$

For any anticommuting traceless operator, such as  $\vec{M}$

$$\langle \vec{M} \rangle = \frac{1}{2} \left[ \sum_{n, \text{occup}} - \sum_{n, \text{empty}} \right] \Phi_{0,n}^\dagger(\vec{x}) \vec{M} \Phi_{0,n}(\vec{x}),$$

(IH, PRL, 2007)

So:

**TRS broken explicitly  $\Rightarrow$  CS broken spontaneously**

## Is the opposite also true?

Consider the **non-Abelian** potential:  $H = i\gamma_0\gamma_i(p_i - A_i)$

$$A_i = A_i^3\gamma_3 + A_i^5\gamma_5 + A_i^{35}\gamma_{35}$$

which manifestly **breaks CS**, but preserves TRS. Since,

$$\gamma_{35} = \sigma_z \otimes I_2$$

$A_i^{35}$  for example, represents a **variation in the position of the Dirac point**, induced by height variations or strain.

$\tilde{M}$  still **anticommutes** with  $H$  but also with  $I_t$  so  $Tr_0\tilde{M} = 0$

**CS broken explicitly  $\Rightarrow$  TRS broken spontaneously**

(IH, PRB, 2008)



# Pseudo-magnetic catalysis:

Flux of non-abelian pseudo-magnetic field



Subspace of zero-energy states (Atiyah-Singer, Aharonov-Casher)



Equally split by TRS breaking mass



**With next-nearest neighbor repulsion, TRS spontaneously broken**

In a non-uniform pseudo-magnetic field (**bulge**):

$$\langle \Psi^\dagger(\vec{x}) \tilde{M} \Psi(\vec{x}) \rangle \propto B^{35}(r)$$

**local TRS breaking!**

## In sum:

- 1) Dirac Hamiltonian in 2D has plenty of (emergent) symmetry
- 2) Chiral symmetry + Time reversal  
=> variety of Mott (“Higgs”) insulators
- 3) Interactions need to be strong for Higgs, **but**
- 4) Ubiquitous zero-energy states => catalyze Higgs  
(QHE at filling factors zero and one in uniform magnetic field)
- 5) **Non-abelian flux catalyzes time-reversal symmetry breaking!** (IH, PHYSICAL REVIEW B 78, 205433 (2008))