

Elliptic - hyperbolic systems and the Einstein Equations.

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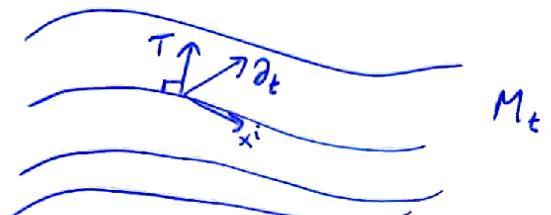
$$(\bar{M}, \bar{g}) \quad 3+1 \text{ Lorentz} \quad -+++ \\ \bar{\nabla}, \bar{R} \quad i, j, k = 1, 2, 3 \quad \alpha, \beta, \gamma = 0, 1, 2, 3, 4.$$

$t: \bar{M} \rightarrow \mathbb{R}$ time function

$$\langle \bar{\nabla}t, \bar{\nabla}t \rangle < 0$$

$\{M_t\}$ foliation by level sets of t .

$$T = \bar{N}'(\partial_t - \bar{x}).$$



$$\bar{g} = -N^2 dt^2 + g_{ij} (dx^i + \bar{x}^i dt)(dx^j + \bar{x}^j dt)$$

$$K_{ij} = -\frac{1}{2} L_T \bar{g}_{ij} \stackrel{\text{with } N \neq 0}{=} -\langle \bar{\nabla}_i T, \partial_x^j \rangle$$

Adapted frame: $(e_\alpha) = (T, e_i)$.

(1)

Backgrounds and motivation

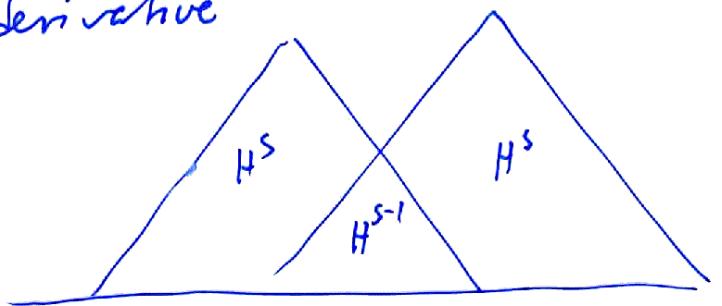
- 3+1 GR - want to understand global structure of Maximal Globally Hyperbolic Vacuum spacetimes
- Cosmic Censorship \leftrightarrow generic MGHR is C^2 -inextendible
 - Asymptotic behavior near (cosmological) singularity: generically spacelike oscillating singularity (BKL) AVTD, quiescent for models with symmetry or stiff fluid.
 - Asymptotic behavior in expanding direction
 - Global uniqueness a'la Choquet-Bruhat, Geroch. What regularity is needed?
 - Approach these problems via CMC-conjecture (spatially cpt. case).

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Results:

- 1) Elliptic-hyperbolic formulation of Einstein Eqs (CMCSH gauge)
w/p in H^s , $s > n/2 + 1$
LA, Moncrief AHP vol 4, '03
[gr-qc/0110111](#).
- 2) Future complete 3+1 vacuum spacetimes (small data global existence) [[gr-qc/0303045](#). LA, Moncrief.]
- 3) CMCSH gauge version of Einstein Eq w/p in $H^{2+\epsilon}$ in 3+1 d.
[LA, Klainerman, Rodnianski, under preparation.]

- hyperbolic gauge cond. expected to be OK locally, not globally.
For global results, need elliptic gauge like CMC.
- For global uniqueness, need elliptic gauge to gain 1 extra derivative



lose one derivative in coord. transformation
⇒ need $H^{3+\varepsilon}$ w. current results.

Long term goal: H^2 - this requires bi-linear estimates.

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(5)
Einstein Eqs in Constant Mean Curvature, Spatial Hyperbolic gauge

$\bar{M} = \mathbb{R} \times M$ 3+1 spacetime, metric $\bar{g}_{\mu\nu}$
 $\bar{R}_{\mu\nu}, \bar{\nabla}_I$, assume M compact, hyperbolic type.
Fix \bar{g} : C^∞ negative curvature background metric on M .

$$\bar{R}_{\mu\nu} = 0 \iff$$

$$(\mathcal{E}) \quad \begin{cases} \partial_t g_{ij} = -2N K_{ij} + h_{ij}^k \bar{g}_{ij} \\ \partial_t K_{ij} = -\bar{D}_i \bar{D}_j N + N(R_{ij} + 2K K_{ij} - 2K_{ik} K_{kj}) + h_{ij}^k K_{ij} \end{cases}$$

$$(\mathcal{C}) \quad \begin{cases} R - |K|^2 + (\partial_t K)^2 = 0 \\ \bar{D}_i K - \bar{D}^j K_{ij} = 0 \end{cases}$$

$$\text{Let } V^k = g^{mn} (\Gamma_{mn}^k - \hat{\Gamma}_{mn}^k) \quad (6)$$

- V^k = torsion field of $\text{Id} : (M, g) \rightarrow (M, \hat{g})$

CMSH gauge const.

$$(G) \left\{ \begin{array}{l} \text{d}K = t + t_0 \\ V^k = 0 \end{array} \right.$$

Note: 1) May consider inhomogeneous gauge const of the form

$$\left\{ \begin{array}{l} \mathcal{L}_{(t-\delta)} \text{d}K = f^0 \\ \mathcal{L}_{(t-\delta)} V^k = f^k \end{array} \right. \quad \begin{array}{l} f^\alpha \text{ given spacetime} \\ f^\alpha = (1, 0) \hookrightarrow (G) \end{array}$$

2) For AF case may use

$$\left\{ \begin{array}{l} \text{d}K = 0 \\ V^k = 0 \\ N \rightarrow 1, x \rightarrow \infty \end{array} \right.$$

$$\hat{\nabla}_m Y_n = \partial_m Y_n - \hat{\Gamma}_{mn}^k Y_k$$

$$\Delta_g h_{ij} = \frac{1}{\mu_g} \hat{\nabla}_m (g^{mn} \mu_g \hat{\nabla}_n h_{ij})$$

$$\mu_g = \sqrt{\det g}$$

$$R_{ij} = -\frac{1}{2} \Delta_g g_{ij} + S_{ij} [g, \partial g] + \delta_{ij}$$

$$\text{where } \delta_{ij} = \frac{1}{2} (\partial_i V_j + \partial_j V_i)$$

S_{ij} quadratic in ∂g .

$g_{ij} \mapsto R_{ij} - \delta_{ij}$ is quasilinear elliptic

(6b)

time differentiate (\tilde{g}):

$$(D) \quad \left\{ \begin{array}{l} -\Delta N + |K|^2 N = 1 \\ \hat{\Delta} g^k + g^{mn} \hat{R}_{men}^k \tilde{g}^l + V^l \hat{\partial}_e \tilde{g}^k = \\ = (-2NK^{mn} + 2D^m \tilde{g}^n)(\Gamma_{mn}^k - \hat{\Gamma}_{mn}^k) \\ + 2D^m N K_m^k - D^k N K_m^m \end{array} \right.$$

here \hat{R}_{men}^k is Riemann tensor of (M, \hat{g}) .Modified Evolution problem for cmcsht Einstein Eqs.

$$(E') \quad \left\{ \begin{array}{l} \partial_t g_{ij} = -2N K_{ij} + L_{\tilde{g}} g_{ij} \\ \partial_t K_{ij} = -\nabla_i \nabla_j N + N(R_{ij} + \text{tr} K K_{ij} - 2K_m K_j^m - \delta_{ij}) \\ \quad + L_{\tilde{g}} K_{ij} \end{array} \right.$$

Fact: $(E') + (D)$ is elliptic-hyperbolic:Propagation of Constraints and Gauge ⑧

$$\text{Let } \left\{ \begin{array}{l} A = \text{tr} K - t - t_0 \\ V^k = g^{mn} (\Gamma_{mn}^k - \hat{\Gamma}_{mn}^k) \\ F = R + (\text{tr} K)^2 - |K|^2 - D_i V^i \\ D_i = \nabla_i \text{tr} K - 2D^m K_{mi} \end{array} \right.$$

Let $\partial_0 = \partial_t - \tilde{g}$.

$$\left\{ \begin{array}{l} \partial_0 A \approx NF \\ \partial_0 F \approx N \Delta A \\ \partial_0 V^i \approx N D^i \\ \partial_0 D^i \approx NP_0 V^i \end{array} \right. \quad + \text{lower order.}$$

$$\text{where } P_0 V^i = \hat{\Delta} g V^i + g^{mn} \hat{R}_{men}^i V^l$$

$$\text{Let } \Sigma = \frac{1}{2} \int |A|^2 + |\Delta A|^2 + |F|^2 + \frac{1}{2} \int |V|^2 + |\hat{\Delta} V|^2 + |D|^2$$

$$\text{then } |\partial_t \Sigma| \leq C\Sigma \quad \text{so } \Sigma = 0 \quad \text{if } (E_0) = 0.$$

⑨

Thm: [LA, Moncrief, AHP 4, 2003]

The mcsH Cauchy problem $(\Sigma') + (D)$
is locally w/p in $H^{M+1+\varepsilon}$, $\varepsilon > 0$

n=3
Thm [LA, Klainerman, Rodnianski]

$$H^{2+\varepsilon}$$

locally w/p in

[under preparation]. Pf uses Strichartz
estimates:

$$\begin{aligned} \square \phi &= 0 \\ \Rightarrow \|\partial \phi\|_{L_t^p L_x^\infty} &\leq C T^\delta (\|\partial \phi\|_{L^2} + \| \partial \phi \|_{H^{5+1}}) \end{aligned}$$

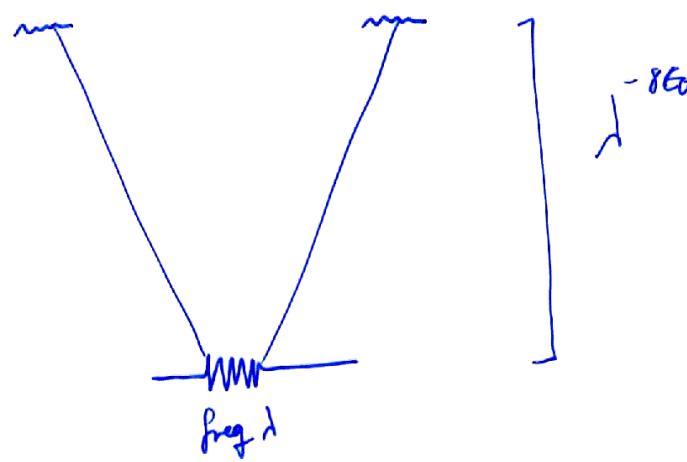
Strichartz:

$$\square g \phi = 0$$

Expansion of null cones \Rightarrow decay of waves

$\Rightarrow L_t^p L_x^\infty$ estimate for $\partial \phi$
via Littlewood-Paley.

Decay holds in rough spacetime satisfying (approximate)
Einstein Eq, after frequency localization and
rescaling.



(9a)
(9b)

Covariant wave equation for ⑩
the metric.

$$\text{Let } u_{ij} = g_{ij}, v_{ij} = -2K_{ij}$$

think of u, v as sections of bundle $Q \rightarrow M$, with covariant derivative D , defined by $D_h h_{ij} = \hat{\nabla}_h h_{ij}$.

Extend to $Q \rightarrow \bar{M}$ with cov. der

D_α , defined by

$$\begin{cases} D_0 u_{ij} = \partial_t u_{ij} \\ D_k u_{ij} = D_k u_{ij} \end{cases}$$

$$\text{let } L_h = \bar{g}^{\alpha\beta} D_\alpha D_\beta h.$$

(10a)

$$(E') \Leftrightarrow \begin{cases} \partial_t u_{ij} = N u_{ij} + \hat{\nabla}_k u_{ij} + f_{1ij} \\ \partial_t v_{ij} = N \hat{\Delta} g u_{ij} + \hat{\nabla}_k v_{ij} + f_{2ij} \end{cases}$$

$$\text{where } \begin{aligned} f_{1ij} &= u_{ej} \hat{\nabla}_i \mathcal{E}^e + u_{ie} \hat{\nabla}_j \mathcal{E}^e \\ f_{2ij} &= F_{ij} + v_{ej} \hat{\nabla}_i \mathcal{E}^e + v_{ie} \hat{\nabla}_j \mathcal{E}^e \\ F_{ij} &= F_{ij}[u, \partial u, v, \partial v, \partial^2 N] \end{aligned}$$

Then (E') can be written in the form

$$Lu = G$$

$$\begin{aligned} G = -N^{-1} f_2 - T^\alpha D_\alpha (N^{-1} f_1) + N^{-1} D^\alpha N D_\alpha u \\ + \Lambda K T^\alpha D_\alpha u \end{aligned}$$

(II)

To Do:

- AF version. \hat{g} may be chosen as Euclidean or Schwarzschild for example.
- inhomogenous gauge version (prescribed mean curvature)
- conformally split version.
- boundary value problem.

Expect: elliptic-hyperbolic formulation of E.E. \rightarrow w/p completely gauge fixed initial-boundary value problem, for ex. using Dirichlet boundary cond. for (N, Σ)