

Asymptotic simplicity:  
completing the picture

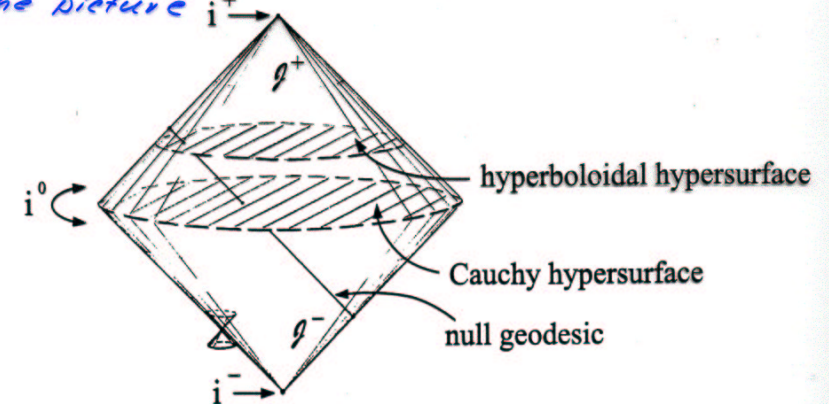
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(For refs cf. H.F. gr-qc/0304003)

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- Completing, not complete!
- the picture



- "...non-trivial existence questionable..." ?
- sharp analysis provides:
  - help with interpretation, concepts, ....
  - insight into equations, gauge cond's, ....
  - support for numerical analysis
- Rchruściel, E. Delay CQG 19(2002) L71 :  
 $\exists$  non-triv. as. simple solutions to  $R_{\mu\nu} = 0$
- 'How many' exist?  $\exists$  characterisation in terms of Cauchy data?

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- Need complete information on an arbitrary hohd of space-like infly  $i^0$ !

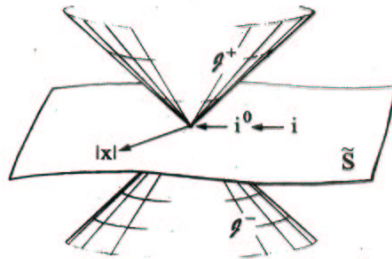
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- Also for generalizations!

- Conformal data singular at space-like infity  $i$
- remove singularity (time symmetric case) by

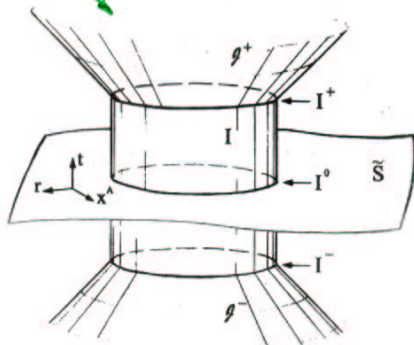
BLOW-UPS  $i \rightarrow I^0 \simeq S^2, \quad i^0 \rightarrow I \simeq ]-1, 1[ \times S^2$

Seen from inside:



not imposed!

consequence of conf. geometry and field equ's



- in suitable gauge of  $\Theta, x^a, e_a, f$ : data smooth on  $S = \tilde{S} \cup I^0$
- equ's symmetric hyperbolic near  $\tilde{S}$

- Unexpected. Specific for Einstein equ's. Most concise realization of the problem

- Reduced equations for  $u = (v, \varphi)$  where  $v = (e^{\hat{a}k}, \hat{r}^{\hat{j}k}, \hat{R}_{\hat{j}k})$ ,  $\varphi = \Theta^{-2} C^{\hat{j}k}$ :

$$\begin{cases} \partial_t v = F(v, \varphi) \\ A^{\hat{a}} \partial_{\hat{a}} \varphi + A^{\hat{a}} \partial_{\hat{a}} v + A^{\hat{a}} \partial_{\hat{a}} v = C \varphi \end{cases}$$

- sol's exist, are smooth up to  $I$  near  $\tilde{S}$
- cylinder  $I$  very special:  $A^{\hat{a}}|_I = 0$
- $\rightarrow$  Inner equ's on  $I$  for  $u^{\hat{a}} \equiv \partial_{\hat{a}} u|_I$
- $\rightarrow$  calc. of  $u^{\hat{a}}$  reduced to expressions

$$y(t) = X(t) \tilde{x}^i(t) y_0 + X(t) \int_0^t \tilde{x}^i(t') b(t') dt'$$

$$y \in \mathcal{O}^2, \quad b = b[u^0, \dots, u^{p-1}], \quad X \in M_{2 \times 2} \text{ 'known'}$$

- 'critical sets'  $I^{\pm}$  very special:  $\det(A^{\hat{a}})|_{I^{\pm}} = 0$

$\rightarrow$  break-down of hyperbolicity at  $I^{\pm}$

in general:  $u^{\hat{a}} \sim (1-t)^k \log^i(1-t)$  as  $t \rightarrow 1$

( $u^{\hat{a}}$  has polyhomogeneous expansion at  $I^{\pm}$ )

There are 2 sources of log-singularities:

- Some  $X(t)$  behave as  $(1-t)^k \log(1-t)$  as  $t \rightarrow 1$   
in that case  $b \equiv 0$ , regularity depends on  $\varphi_0$
- $\rightarrow$  no log-terms occur from this if the data  $(\bar{S}, h)$  satisfy the 'regularity condition'

(\*)  $\mathcal{R}(D_{\alpha_1} \dots D_{\alpha_p} \bar{g}_{ab})(i) = 0$ ,  $0 \leq p \leq p_*$  with  $p_* = \infty$

(non-trivial after linearization)

- in the other cases  $X(t) = Y(t)S(t)$  smooth with  $\det(Y(t)) \neq 0$ ,  $|t| \leq 1$ ,  $S(t) = \begin{pmatrix} (1-t)^{p-2} & 0 \\ 0 & (1-t)^{p-2} \end{pmatrix}$   $p \geq 2$

the regularity depends on  $b = b[u^0, \dots, u^{p-1}]$

the integrals can, in principle, be calculated recursively

(these log-terms disappear under linearization)

Implications for null infinity

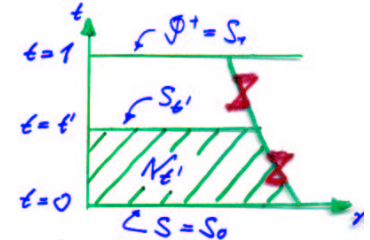
- Linearize (with  $\mathcal{F} = \{t = \pm 1, r > 0\}$ ) at Minkowski space
- Bianchi equ's then imply

$$\begin{aligned} (1+t) \partial_t \varphi_0 - r \partial_r \varphi_0 + X_+ \varphi_1 &= -2\varphi_0 \\ 2 \partial_t \varphi_1 + X_+ \varphi_2 + X_- \varphi_0 &= -2\varphi_1 \\ 2 \partial_t \varphi_2 + X_+ \varphi_3 + X_- \varphi_1 &= 0 \\ 2 \partial_t \varphi_3 + X_+ \varphi_4 + X_- \varphi_2 &= 2\varphi_3 \\ (1-t) \partial_t \varphi_4 + r \partial_r \varphi &+ X_- \varphi_3 = 2\varphi_4 \end{aligned}$$

- Standard energy estimates

$$\|\varphi\|_{L^2(S_{t'})} \leq \frac{C}{(1-t')^{3/2}} \|\varphi\|_{L^2(S_0)}$$

useless as  $t' \rightarrow 1$ .



- Useful energy estimates

$$\|\partial_r^p \varphi\|_{H_+^m(N_{t'})} \leq C \|\partial_r^p \varphi\|_{H_+^m(S_0)}, \quad p \geq m+2, \quad 0 \leq t' < 1, \quad C = \text{const.}$$

$\rightarrow \partial_r^p \varphi \in C^j(N_{t'})$ ,  $p \geq j+6$ . Two integrations give:

$$\varphi = \sum_{p=0}^{p-1} \frac{1}{p!} \varphi^{(p)} r^p + \hat{\varphi}, \quad p \geq j+6 \quad \text{with } \hat{\varphi} \in C^j(N_{t'}).$$

Solutions  $C^\infty$  if lin. reg. cond. hold with  $p_* = \infty$  otherwise  $\exists$  log-terms on  $\mathcal{F}$  (expected but not obvious!)

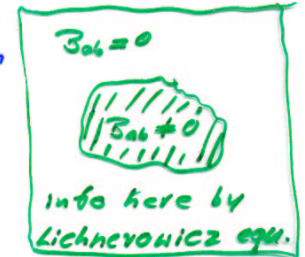
- The log-terms are determined uniquely by the data  $(\tilde{S}, \tilde{h})$ , they are not generated by the setting
- the nature of the degeneracy at  $I^\pm$  pinpoints the origin of the non-smoothness on  $\mathcal{I}$
- the transition: linear  $\rightarrow$  non-linear will not improve the situation, we can expect log-terms in general
- main tasks

i) analyse the behaviour of  $u^p$  on  $\bar{I} = I \cup I^- \cup I^+$   
derive nec. & suff. regularity conditions

ii) analyse the consequences of the behaviour of the  $u^p$  near  $I^\pm$  for the smoothness of  $\mathcal{I}^\pm$   
(generalize the linear estimates...)

- The reg. cond. (\*) is nec. + suff. for  $p = 0, 1, 2, 3$ .

- J. Valiente Kroon gr-qc/0211024 uses algebraic computer program studies data which are conformally flat near  $i$ :



$\Rightarrow u^4$  analytic on  $\bar{I}$

$\Rightarrow u^5 = u_*^5 + c^5 (1-t)^{k_5} \log(1-t)$  as  $t \rightarrow 1$   
 $u_*^5$  analytic on  $\bar{I}$ ,  $c^5 \sim NP$  const.

require  $c^5 = 0$

$\Rightarrow u^6 = u_*^6 + c^6 (1-t)^{k_6} \log(1-t)$  as  $t \rightarrow 1$

require  $c^6 = 0$  & axial symmetry

$\Rightarrow u^7 = u_*^7 + c^7 (1-t)^{k_7} \log(1-t)$  as  $t \rightarrow 1$

etc. etc.

- global conditions? things have changed!

- sequence of conditions suggests:  
regularity at all orders  $\Rightarrow$  asympt. Schwarzschild  
whence Schwarzschild near  $i$  (only static sol. in class!)

- $m=0$ : the reg. cond. (\*) are nec. + suff. for the smoothness of  $\mathcal{D}$
- static data with  $m \neq 0$  satisfy the reg cond. (\*) at all orders.
- conformal flatness,  $m \neq 0$ , and regularity seem to imply asymptotic staticity.
- $m \neq 0$ : reg. cond. (\*) is not sufficient for the smoothness of  $\mathcal{D}$

- what is the role of the mass here?
- does asymptotic staticity play a particular role here?
- is the setting with the cylinder at space-like infinity smooth at  $I^\pm$  in the static case?

static sol's :  $\tilde{g} = v^2 dt^2 + \tilde{h}_{ab} dx^a dx^b$ ,  $v(x^c)$ ,  $\tilde{h}_{ab}(x^c)$

static equ's :  $\mathcal{R}_{ab}[\tilde{h}] = \frac{1}{v} \tilde{\nabla}_a \tilde{\nabla}_b v$ ,  $\Delta_{\tilde{h}} v = 0$

For static solutions with  $m \neq 0$  everything is real analytic (in a suitable sense) in a mhd of  $\bar{I} = I \cup I^- \cup I^+$ .

Requires static field equ's in detail!

with:  $v_0 = \partial_t$ ,  $v_a = r \partial_r$ ,  $v_B = \partial_{x^B}$   
 $\alpha^0 = dt$ ,  $\alpha^1 = \frac{1}{r} dr$ ,  $\alpha^B = dx^B$  we have

$$g = g_{ik} \alpha^i \alpha^k, \quad \nabla_{v_i} v_k = \gamma_i^j \alpha_j v_k \quad \text{with}$$

- $g_{ik}$  analytic & Lorentzian near  $\bar{I}$
- $\gamma_i^j$  analytic near  $\bar{I}$
- all tensor fields derived from  $g$  are analytic in the frame above
- $u^p$ ,  $p=0, \dots$ , are analytic on  $\bar{I}$
- for data which are asympt. static up to order  $k$  the  $u^p$  are analytic on  $\bar{I}$  for  $p \leq p_*(k)$
- (cf. R. Chrusciel, E. Delay gr-qc/0301073 for such data)

$(m=0)$ -reg. cond. →  $(m \neq 0)$ -reg. cond. → asympt. static

↑  
where precisely?