

## Question for Scalar Studies

- Can Cauchy codes work using hyperboloidal slices?
- Courant condition?
- How do the boundaries work – any easier?
- Wave extraction?

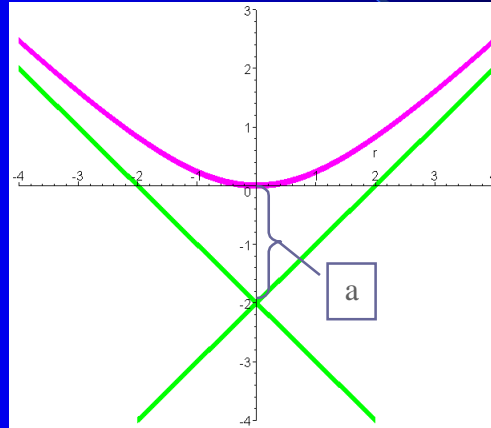
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## Familiar Future Questions

- Initial value equations for nonzero  $K$
- Slicing conditions to yield hyperboloids
- Space coordinate (shift) choices to be Courant compatible
- Boundary conditions on moving BH's
- Etc.

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## First Example: Minkowski Spacetime



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## Change the Time Coordinate

$$T = u - a + \sqrt{a^2 + R^2}$$

$a$  = 'Scri-delay' -- events at  $(R,u)=(0,0)$  influence Scri+  
no earlier than  $u=a$

$$ds^2 = -\left(1 + \frac{R^2}{a^2}\right) du^2 + \left(\frac{a dR}{\sqrt{a^2 + R^2}} - \frac{R du}{a}\right)^2 + R^2 d\Omega^2$$

Check the coordinate velocity of light:  $dR/du$  when  $ds=0$ .

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## What About the Courant Condition?

$$c_{out} = \left( \frac{R}{a} + \sqrt{1 + \frac{R^2}{a^2}} \right) \sqrt{1 + \frac{R^2}{a^2}}$$

$$c_{in} = \left( \frac{R}{a} - \sqrt{1 + \frac{R^2}{a^2}} \right) \sqrt{1 + \frac{R^2}{a^2}}$$

Trouble from  $c_{out}$  when  $r$  gets large:  $\Delta u < \frac{\Delta R}{c}$

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## Redshifted Outgoing Waves

$$\Phi(R, T) = \frac{\Lambda(T-R) - \Lambda(T+R)}{R}$$

$$T - R = u - a + \frac{a^2}{2R} - \frac{a^4}{8R^3} \dots$$

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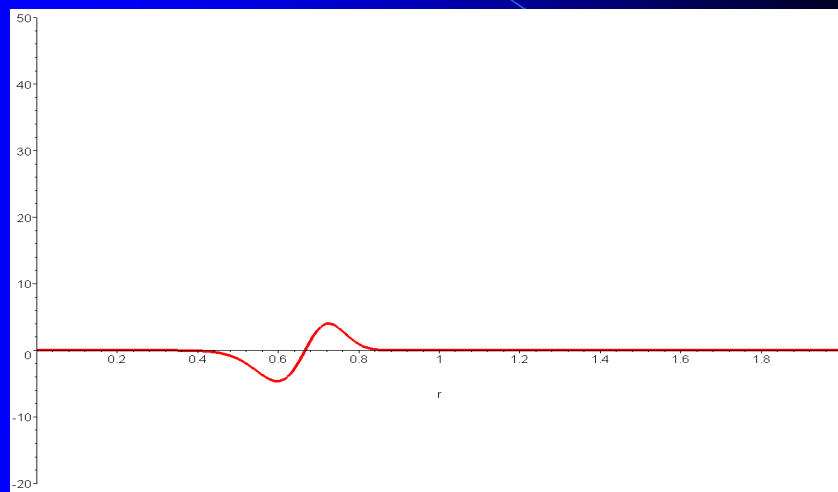
## Analytic Mesh Refinement

$$\frac{R}{a} = \frac{r}{1 - \frac{r^2}{4}}$$

Dimensionless radial coordinate compactifies the slice:  $\text{Scri}^+$  at  $r=2$ .

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## Analytic Solution



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## Working Minkowski Metric

$$ds^2 = \frac{a^2}{q^2} [-B^2 du^2 + (dr - r du)^2 + r^2 d\Omega^2]$$

$$q = 1 - \frac{r^2}{4}, \quad B = 1 + \frac{r^2}{4}$$

$$0 \leq r < 2$$

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## Working light speeds

$$c_{out} = 1 + r + \frac{r^2}{4} < 4$$

$$-1 \leq c_{in} = -1 + r - \frac{r^2}{4} < 0$$

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## Try Schwarzschild background

- Slicings from Malec and O`Murchadha's catalogue with  $K=3/a$ .
- The simplest don't cross the BH horizon (white hole  $r=2m$  instead); choose a time independent set that do.
- Still need to understand scalar wave initial conditions.

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## M-O`M Schwarzschild

$$ds^2 = -B^2 du^2 + \frac{(dR + \beta du)^2}{B^2} + R^2 d\Omega^2$$

$$B = \sqrt{1 - \frac{2m}{R} + \left(\frac{R}{a} - \frac{\kappa a^2}{R^2}\right)^2}$$

$$\beta = -B \left(\frac{R}{a} - \frac{\kappa a^2}{R^2}\right)$$

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## Light speeds

$$c_{out} = B^2 - \beta \quad , \quad c_{in} = -B^2 - \beta$$

At  $R=2m$  one finds  $B^2 = |\beta|$  so that  $c_{out} = 0$  only if  $\beta$  is positive. This requires that

$$\left(\frac{2m}{a}\right)^3 < \kappa$$

Then light going outward at  $R=2m$  can have constant  $R$ , zero speed, so BH.

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## Extrinsic Curvature (Schw)

$$K^a_b = \begin{bmatrix} -\frac{1}{a} - \frac{2a^2\kappa}{R^3} & 0 & 0 \\ 0 & -\frac{1}{a} + \frac{a^2\kappa}{R^3} & 0 \\ 0 & 0 & -\frac{1}{a} + \frac{a^2\kappa}{R^3} \end{bmatrix}$$

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### 3-curvature of slice

$$R^a_b = \begin{bmatrix} -\frac{2}{a^2} - \frac{2(m+a\kappa)}{R^3} + \frac{4a^4\kappa^2}{R^6} & 0 & 0 \\ 0 & -\frac{2}{a^2} + \frac{m+a\kappa}{R^3} + \frac{4a^4\kappa^2}{R^6} & 0 \\ 0 & 0 & -\frac{2}{a^2} + \frac{m+a\kappa}{R^3} + \frac{4a^4\kappa^2}{R^6} \end{bmatrix}$$

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### Bring Scri to r=2

$$\frac{R}{a} = \frac{r}{1 - \frac{r^2}{4}} = \frac{r}{q}$$

$$q = 1 - \frac{r^2}{4}$$

Same transformation as in flat spacetime

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## Compactified M-O`M

$$ds^2 = \frac{a^2}{q^2} [-B^2 du^2 + f^2 (dr + \beta du)^2 + r^2 d\Omega^2]$$

$$f^2 = \frac{\left(1 + \frac{r^2}{4}\right)^2}{B^2}, \quad \beta = -\left(r - \frac{\kappa q^3}{r^2}\right) \frac{1}{f}$$

$$B = \sqrt{q^2 \left(1 - \frac{2mq}{ar}\right) + \left(r - \frac{\kappa q^3}{r^2}\right)^2}$$

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## Nearly Minkowski at Scri+

Note that  $m$  and  $\kappa$  only appear in the combinations  $mq^3$  and  $\kappa q^3$ .

Since  $q \sim (2-r)$  for  $r \rightarrow 2$  (Scri+) their effects die like  $q^3 \sim \left(\frac{2a}{R}\right)^3$  toward Scri+.

In particular  $c_{\text{out}} \rightarrow 4$  and  $c_{\text{in}} \rightarrow 0$  independent of  $m$  and  $\kappa$ .

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# Wave Propagation with Hyperboloidal Slicings

