On strong hyperbolicity

Oscar Reula

reula@fis.uncor.edu

FaMAF, Córdoba, Argentina

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Collaborators:

R. GerochG. NagyO. Ortiz

OUTLINE

- Introduction, definitions, examples.
- Covariant Definitions
- Causality
- Applications
 - ADM-BSSN hyperbolicity
 - Subsidiary system hyperbolicity

First Order Constant Coefficient Systems

$$\partial_t u^{\alpha} = A^{\alpha i}{}_{\beta} \nabla_i u^{\beta} + B^{\alpha}{}_{\beta} u^{\beta}$$

Question: When is the above system well posed in the L^2 sense?

 $||u^{\alpha}(t)||_{L^{2}} \le C(t)||u^{\alpha}(0)||_{L^{2}}$

Answer: It is well posed if and only if for all co-vectors ω_i , the matrix $A^{\alpha i}{}_{\beta}\omega_i$ has only real eigen-values and a complete set of eigen-vectors.

First Order Constant Coefficient Systems II

$$\partial_t \hat{u}^\alpha = iA^{\alpha i}{}_\beta \omega_i \hat{u}^\beta$$

$$\hat{u}^{\alpha}(t) = (e^{iA^{i}\omega_{i}t})^{\alpha}{}_{\beta}\hat{u}^{\beta}(0)$$

 $||u^{\alpha}(t)||_{L^{2}} = ||\hat{u}^{\alpha}(t)||_{L^{2}} \le C(t)||\hat{u}^{\alpha}(0)||_{L^{2}} \le C(t)||u^{\alpha}(0)||_{L^{2}}$

$$C(t) = \sup_{\tau \in [0,t]} \sup_{\omega_a} ||(e^{iA^a \omega_a \tau})^{\alpha}{}_{\beta}||$$

First Order Quasi-linear Systems I

$$\partial_t u^{\alpha} = A^{\alpha i}{}_{\beta}(u, x, t) \nabla_i u^{\beta} + B^{\alpha}(u, x, t)$$

The above system is well posed (w.r.t. a Sobolev Norm) in a neighborhood of u_0^{α} if and only if for all u^{α} close enough to u_0^{α} , all co-vectors ω_i and all points, the matrix $A^{\alpha i}{}_{\beta}(u, x, t)\omega_i$ has only real eigen-values and a complete set of eigen-vectors. Plus some "technical" condition

We call such systems **strongly hyperbolic**.

First Order Quasi-linear Systems II

If a system is strongly hyperbolic then there exists a positive definite bilinear form (a metric) $H_{\alpha\beta} = H_{\alpha\beta}(u, x, t, \omega_a)$ uniformly bounded by above and away from zero in ω_a such that:

 $H_{\alpha\gamma}A^{\gamma a}{}_{\beta}\omega_a$

is also symmetric. [Kreiss Matrix Theorem] Technical condition requires H to be smooth also on ω

If there exists a $H_{\alpha\beta}$ independent of ω_a we say that the system is **symmetric** hyperbolic

If strong hyperbolicity fails it is easy to construct a sequence solutions whose initial data has norm one but whose norm at any future time tends to infi nity. Non-linear behavior can not cure this.

Examples:

Example 1: Maxwell equations:

$$W_{ij} := \partial_i A_j$$

$$\partial_t E_i = \partial^j W_{ji} - \partial^j W_{ji} - \alpha (\partial^j W_{ij} - \partial_i W_j^j)$$

$$\partial_t W_{ij} = \partial_i E_j - \frac{1}{2} \beta e_{ij} \partial^k E_k$$

This system is symmetric hyperbolic for $\alpha < 0$ and $\beta < -\frac{2}{3}$ (most general symmetrizer built out of the 3-metric). But strongly hyperbolic for all (α, β) such that $\alpha\beta > 0$.

Examples:

Example 2:

Consider the matrices,

$$\mathbf{A^{0}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & \frac{1}{2} \end{pmatrix} \quad \mathbf{A^{1}} = \begin{pmatrix} -2 & 10s_{1} & 0 \\ 0 & 1 & -2s_{2} \\ 0 & 0 & 0 \end{pmatrix} \quad \mathbf{A^{2}} = \begin{pmatrix} 0 & s_{1} & 0 \\ 0 & \frac{1}{2} & 7s_{2} \\ 0 & 0 & 1 \end{pmatrix}$$

There is no positive definite $h_{\alpha\beta}$ which would symmetrize $A^i\omega_i$ for arbitrary ω_i . Nevertheless $(A^0)^{-1}(A^1 + \lambda A^2)$ is diagonalizable.

$$(\mathbf{A^0})^{-1}(\mathbf{A^1} + \lambda \mathbf{A^2}) = \begin{pmatrix} -2 & (10 + \lambda)s_1 & 0 \\ 0 & \frac{1}{4}(2 + \lambda) & \frac{1}{2}(-2 + 7)s_2 \\ 0 & 0 & 2\lambda \end{pmatrix}$$

Covariant definitions I:

$$A^{\alpha a}{}_{\beta}(u,p)\nabla_a u^{\beta} = J_{\alpha}(u,p)$$

The sum of two symmetric matrices is symmetric

Definition: The above system is symmetric hyperbolic if there exists $h_{\alpha\beta}(u, p)$ such that:

 $h_{\alpha\beta}(u,p)A^{\beta a}{}_{\gamma}(u,p)$ is symmetric.

for some n_a , $h_{\alpha\beta}(u,p)A^{\beta a}{}_{\gamma}(u,p)n_a$ is positive definite.

One can defi ne an energy vector:

$$E^{a} := h_{\alpha\beta}(u,p) A^{\beta a}{}_{\gamma}(u,p) \delta u^{\alpha} \delta u^{\gamma}$$

 $E^a n_a \ge 0$

If n_a is as above, then $n_a + \varepsilon w_a$ is also as above afor ε is mall enough $v_{\text{Derived and only one part Objects - p.10/26}}$

Covariant definitions II:

$$A^{\alpha a}{}_{\beta}(u,p)\nabla_a u^{\beta} = J_{\alpha}(u,p)$$

The sum of two diagonalizable matrices is not necessarily diagonalizable Definition A: The above system is strongly hyperbolic if there exists n_a such that:

- $\square A^{\alpha a}{}_{\beta} n_a$ is invertible, and
- for each loop $\kappa(\lambda)_a = \lambda n_a + \omega_a \ \lambda \in [-\infty, \infty]$ where ω_a is not proportional to n_a , $dim(span\{\cup_{\lambda \in R} Kern\{A^{a\alpha}{}_{\beta}\kappa_a(\lambda)\}\}) = dim\{\text{manifold of fields}\}$

Definition B: The above system is strongly hyperbolic if there exists n_a such that for each co-vector ω_a there exists $h_{\alpha\beta}(u, p, \omega)$ satisfying:

- $= h_{\alpha\beta}(u, p, \omega) A^{\beta a}{}_{\gamma}(u, p) \omega_a \text{ is symmetric.}$
- $= h_{\alpha\beta}(u, p, \omega) A^{\beta a}{}_{\gamma}(u, p) n_a$ is symmetric and positive definite.
- If n_a is as above then $n_a + \varepsilon \omega_a$ is also as above fore similar enough compact Objects p.11/26

First Order Pseudo-Differential Systems

$$\partial_t u^{\alpha} = P^{\alpha}{}_{\beta}(u, x, t, D) u^{\beta} := \int p(u, x, t, \omega_i) \alpha_{\beta} e^{i\omega_i x^i} \hat{u}^{\beta} d\Omega$$

The above system is said to be pseudo-differential of first order if the following limit exists,

$$\lim_{\lambda \to \infty} \frac{1}{\lambda} p(u, x, t, \lambda \omega_i) \alpha_\beta := p_1(u, x, t, \omega_i) \alpha_\beta$$

If furthermore $ip_1(u, x, t, \omega_i)\alpha_\beta$ has only real eigenvalues and a complete set of eigen-vectors we say the systems is **strongly hyperbolic**.

Strongly hyperbolic pseudo-differential operators plus technical condition are well posed. [Taylor, Kreiss-Ortiz-R]

Causality

- We consider the domain of dependence of the linearized equation at a given background u_0^{α} .
- The domain of dependence of a region Σ_0 of a Cauchy surface is given by the maximal foliation of such region produced by hypersurfaces whose normal is such that:

$$E^{a}(\delta u)n_{a} \geq 0 \quad \forall \delta u \leftrightarrow \det(A^{\alpha a}{}_{\beta}n_{a}) \neq 0$$

- Surfaces with normal such that the determinant vanishes are called characteristic surfaces.
- For each co-vector k_a which is a characteristic there is a perturbation which in the high frequency limit moves along the integral lines of $V^a = \frac{\partial \det(A^c k_c)}{\partial k_a}$ at points where $\det(A^c k_c) = 0$ ($V^a k_a = 0$).
- **Question:** what happens in the case of strongly hyperbolic systems?

Holmgren's Theorem

Given an analytic coeffi cient equation system (not necessarily hyperbolic!)

$$\partial_t u^{\alpha} = A^{\alpha i}{}_{\beta}(x,t) \nabla_i u^{\beta} + B^{\alpha}{}_{\beta}(x,t) u^{\beta}$$

and assuming the solution vanishes in a hypersurface Σ_0 then the solution, if sufficiently smooth, vanishes in a whole neighborhood of it, given by the maximal foliations such that their normals do not become characteristics.

- Generalizable to the case of non-analytic coefficients for strng-hyp. systems.
 - Extend the space-time to \mathbb{R}^n or \mathbb{T}^n .
 - Approximate the system by an analytic sequence of strng-hyp. systems.
 - Use Holmgren's theorem on each one of them to conclude that the one parameter family of solutions so generated vanishes in some region Ω_n .
 - Use continuous dependence of solutions of strongly hyperbolic systems to show that the limiting solution would also vanish in a limiting set Ω .

Summary

- Strongly hyperbolic differential (and pseudo-differential) systems are well posed.
- There are global energy norms (pseudo-differential operators).
- **Ther** are covariant definitions. And open set of "space-like" hyper-surfaces.
- Strongly hyperbolic differential (and pseudo-differential) systems have fi nite propagation speeds. With domain of dependence given by their characteristic fi elds.
- **Sym**metric hyperbolic energy \leftrightarrow Summation by parts in finite differences
- **Stron**gly hyperbolic pseudo-energy \leftrightarrow Pseudo-spectral methods.

Applications:

ADM-BSSN fi rst-second order systems [Frittelli-R, Sarbach-Calabrese-Pullin-Tiglio, Kreiss-Ortiz, Nagy-Ortiz-R]

Constraint propagation. [Hyperbolicity properties of subsidiary systems of constraints.]

ADM equations I

$$G_{ab} = 0 \quad \Rightarrow \begin{cases} \mathcal{L}_{n}h_{ab} = -2k_{ab}, \\ \mathcal{L}_{n}k_{ab} = {}^{(3)}R_{ab} - 2k_{a}{}^{c}k_{bc} + k_{ab}k_{c}{}^{c} - \frac{D_{a}D_{b}N}{N}, \\ {}^{(3)}R + (k_{c}{}^{c})^{2} - k_{ab}k^{ab} = 0, \\ D^{b}k_{ba} - D_{a}k = 0, \end{cases}$$

ADM equations II

$$\begin{aligned} \mathcal{L}_{(t-\beta)}h_{ij} &= -2Nk_{ij} \\ \mathcal{L}_{(t-\beta)}k_{ij} &= \frac{N}{2}h^{kl}\left[-\partial_k\partial_l h_{ij} - \partial_i\partial_j h_{kl} + 2\partial_k\partial_{(i}h_{j)l}\right] + B_{ij} \end{aligned}$$

where

$$B_{ij} := N \left[\gamma_{ikl} \gamma_j{}^{kl} - \gamma_{ij}{}^k \gamma_{kl}{}^l - 2k_i{}^l k_{jl} + k_{ij} k_l{}^l - A_{ij} \right],$$

$$\gamma_{ij}{}^k := \frac{1}{2} h^{kl} (2\partial_{(i}h_{j)k} - \partial_k h_{ij}),$$

$$A_{ij} := a_i a_j - \gamma_{ij}{}^k a_k - 2\gamma_{ikl} \gamma_j{}^{(kl)} + \partial_i [(\partial_j N)/N],$$

$$a_i := (\partial_i N)/N.$$

ADM equations III

Hyperbolicity analysis: 1) consider only the principal part, 2) freeze coefficients, 3) substitute all derivatives by Fourier transforms $(\partial_k h_{ij} \rightarrow i\omega_k \hat{h}_{ik})$, and 4) define $\hat{\ell}_{ij} = i\omega \hat{h}_{ij}$. [Kreiss, Ortiz][Taylor]

The associated first order system is then

$$\begin{aligned} \partial_t \hat{\ell}_{ij} &\doteq i\omega \left[-2N\hat{k}_{ij} + \tilde{\omega}_k \beta^k \hat{\ell}_{ij} \right], \\ \partial_t \hat{k}_{ij} &\doteq i\omega \left[-\frac{N}{2} \left(\hat{\ell}_{ij} + \tilde{\omega}_i \tilde{\omega}_j h^{kl} \hat{\ell}_{kl} - 2\tilde{\omega}^k \tilde{\omega}_{(i} \hat{\ell}_{j)k} \right) + \tilde{\omega}_k \beta^k \hat{k}_{ij} \right] \end{aligned}$$

with $\tilde{\omega}_i = \omega_i / \omega$.

Result:

ADM equations are only weakly hyperbolic (3 eigenvectors missing).

ADM equations IV

 $N = h^b Q$ (h = determinant of h_{ij})

The associated first order system is then

$$\begin{aligned} \partial_t \hat{\ell}_{ij} &= i\omega \left[-2N\hat{k}_{ij} + \tilde{\omega}_k \beta^k \hat{\ell}_{ij} \right], \\ \partial_t \hat{k}_{ij} &= i\omega \left[-\frac{N}{2} \left(\hat{\ell}_{ij} + (1+\mathbf{b}) \tilde{\omega}_i \tilde{\omega}_j h^{kl} \hat{\ell}_{kl} - 2\tilde{\omega}^k \tilde{\omega}_{(i} \hat{\ell}_{j)k} \right) + \tilde{\omega}_k \beta^k \hat{k}_{ij} \right] \end{aligned}$$

Result:

- Modified ADM equations for b > 0 still weakly hyperbolic (2 eigenvectors missing).
- Adding Hamiltonian constraint does not change hyperbolicity, but does change characteristics.

BSSN equations I

$$f^{k} = h^{ij} \gamma_{ij}{}^{k} + dh^{kl} \gamma_{lm}{}^{m} = h^{kl} (h^{ij} \partial_{i} h_{jl} + \partial_{l} \ln h)$$

$$\begin{aligned} \mathcal{L}_{(t-\beta)}h_{ij} &= -2Nk_{ij} \\ \mathcal{L}_{(t-\beta)}k_{ij} &= \frac{N}{2}h^{kl}\left[-\partial_k\partial_l h_{ij} - b\,\partial_i\partial_j h_{kl}\right] + N\partial_{(i}f_{j)} + \mathcal{B}_{ij} \\ \mathcal{L}_{(t-\beta)}f_i &= N\left[-(2-c)D^k k_{ki} + (1-c)D_i k_k^{\ k}\right] + \mathcal{C}_i \end{aligned}$$

BSSN equations II

Hyperbolicity Analysis:

$$\begin{aligned} \partial_t \hat{\ell}_{ij} &\doteq i\omega \left[-2\alpha \hat{k}_{ij} + \tilde{\omega}_k \beta^k \hat{\ell}_{ij} \right] \\ \partial_t \hat{k}_{ij} &\doteq i\omega \left[\frac{\alpha}{2} \left(-\hat{\ell}_{ij} - b \, \tilde{\omega}_i \tilde{\omega}_j h^{kl} \hat{\ell}_{kl} + 2 \tilde{\omega}_{(i} \hat{f}_{j)} \right) + \tilde{\omega}_k \beta^k \hat{k}_{ij} \right] \\ \partial_t \hat{f}_i &\doteq i\omega \left[\alpha \left((-2+c) \hat{k}_{ik} \tilde{\omega}^k + (1-c) \tilde{\omega}_i h^{kl} \hat{k}_{kl} \right) + \tilde{\omega}_k \beta^k \hat{f}_i \right] \end{aligned}$$

Result: [Nagy-Ortiz-R]

Modified BSSN equations for b > 0 c > 0 strongly hyperbolic.

Eigenvalues: $(0, \pm 1, \pm \sqrt{b}, \pm \sqrt{c/2})$

Constraint Propagation

Evolution System:

$$\partial_t u^{\alpha} = A(u, t, x)^{\alpha a}{}_{\beta} \partial_a u^{\beta} + B(u, t, x)^{\alpha},$$

Constraints:

$$C^A = K(u, t, x)^{Aa}{}_{\beta}\partial_a u^{\beta} + L(u, t, x)^A,$$

Integrability condition (subsidiary system):

$$\partial_t C^A = S(u, t, x)^{Aa}{}_B \partial_a C^B + R(u, \partial u, t, x)^A{}_B C^B,$$

Want to study what can we say about the properties of the subsidiary system from what we know from the evolution system.

Constraint Propagation II

Problem: In general $S(u, t, x)^{Aa}{}_B$ is not unique if the constraint themselves satisfy certain identities.

For instance, if there is an $L_A(\omega)$ such that:

 $L_A(\omega)K^{An}{}_{\alpha}\omega_n=0$

we could add to $S(u, t, x)^{Aa}{}_B$

$M^{Aa}L_B$

With this addition there are easy examples where one can get any sort of badly posed systems!

Constraint Propagation III

Assume: For any ω_i , $K^{An}{}_{\alpha}\omega_n$ is surjective.

In general this is not satisfied, but in examples of interest one finds subset of constraints which do satisfy it. [Maxwell, EC].

Constraint Propagation IV

Integrability condition implies:

$$K^{A(a}{}_{\alpha}A^{|\alpha|b)}{}_{\beta} - S^{A(a}{}_{B}K^{|B|b)}{}_{\beta} = 0$$

Lemma 1: Given any fixed non-vanishing co-vector ω_a . If (σ, u^{α}) is an eigenvalue-eigenvector pair of $A^{\alpha a}{}_{\beta}\omega_a$ then $(\sigma, v^A = K^{Aa}{}_{\alpha}\omega_a u^{\alpha})$, if v^A is non-vanishing, is an eigenvalue-eigenvector pair of $S^{Aa}{}_B\omega_a$.

Constraint Propagation IV

Integrability condition implies:

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Integrability condition implies:

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- **Corollary 1:** If the evolution system is strongly hyperbolic then so is the subsidiary system. [It does not work symmetric \rightarrow symmetric].
- **Corollary 2:** The characteristics of the subsidiary system are a subset of the characteristics of the evolution system. The domain of dependence of the subsidiary system is at least as large as the domain of dependence of the evolution system.