Stacking Domain Wall Magnons in Twisted van der Waals Magnets

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PERSPECTIVE

Van der Waals heterostructures

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Stacking engineering



Cao... Jarillo-Herrero (2018)

Twistronics fundamentally is stacking engineering

Stacking DWs are ubiquitous



Ju... Wang, Nature (2015)



Edelberg...Ochoa, Pasupathy, Nat. Phys. (2020)

ID confined channels



Sunku...Basov, Science (2018)

What about 2D magnets?







Crl₃: a prototypical 2D magnet



McGuire et al, Chem. Mater. 27, 612 (2015)

Structural phase transition



Crl₃ undergoes a structural phase transition from monoclinic to rhombohedral at around 210 K

McGuire et al, Chem. Mater. 27, 612 (2015)



ween the upper and the upper domain steresis loop (about omain (orange and t is centred between seen in the resulting approximately 1-µm

an analysis of the r signal. Zero-fieldwhile cooling the perature well above upon cooling down eld ($\mu_0 H = 0.15 \text{ T}$). below $T_{\rm C}$, at which oled sweep diverge ^C for the monolayer K) for bulk samples. pportunity to inveskness. Figure 3a–c r CrI₃ samples. All stently show ferroentred at $\mu_0 H = 0 T$ nent and saturation which is an order of



Stacking and magnetic energy



Sivadas, Okamoto, Xu, Fennie, DX Nano Lett. 18, 7658 (2018)

Interlayer super-super exchange



Interlayer super-super exchange



Stacking dependent magnetism is tied to local stacking of octahedrons. Applicable to other 2D magnets as well

Sivadas, Okamoto, Xu, Fennie, DX Nano Lett. **18**, 7658 (2018); Jiang et al PRB (2019); Soriano et al, Solid State Commun. (2019); Wang et al. Nat. Comm (2018)

Experimental evidence: SHG



SHG angle dependence shows clear C₂ symmetry, consistent with monoclinic stacking (rhombohedral stacking should show C₃)

Sun, ... DX, Wu, Xu, Nature (2019)



Similar map of interlayer exchange has been found in CrCl₃



Klein et al, Nat. Phys. (2019)

What is the consequence of stacking-dependent magnetism?

Stacking dependent magnetism



the exchange interaction.

Structural domain walls



Energy of the domain wall

$$E_{\text{str}} = \int \left[\frac{V(b)}{4} + \frac{B+G}{4} (\partial_x b^x)^2 + \frac{G}{4} (\partial_x b^z)^2 \right] dx$$

B= 54.307 eV and G = 39.248 eV are bulk and shear modulus (DFT)

Structural domain walls

Minimization is performed by solving the following differential equation:

$$\frac{\partial \boldsymbol{b}}{\partial t} = -\frac{\delta E_{\rm str}}{\delta \boldsymbol{b}}$$



The RM domain wall can be described by the sine-Gordon equation

$$b^{z} = 2(b_{\text{right}}^{z} - b_{\text{left}}^{z}) \arctan[\exp(x/w)]/\pi + b_{\text{left}}^{z}$$

Variation of interlayer exchange



$$E_{\text{mag}} = \int \left[\sum_{\alpha,\beta,l} \frac{A}{2} (\partial_{\beta} m_l^{\alpha})^2 - \sum_l \frac{K}{2} (m_l^{\gamma})^2 - \sum_{\alpha} J m_1^{\alpha} m_2^{\alpha} \right] dx$$

There are two magnon branches (The Cr atoms form a honeycomb structure.). We are only looking at the low-energy long wave-length limit.

Spin configuration of RR DW



- Intralayer exchange prefers uniform spin configuration
- Easy-axis anisotropy also disfavors any spin rotation
- Inside the structural domain wall, interlayer exchange flips sign, which favors spin rotation

Magnetic structure of RR domain wall is uniform FM in the outof-plane direction

Spin dynamics of RR DW

 $\delta \boldsymbol{m}_{\pm} = \delta \boldsymbol{m}_1 \pm \delta \boldsymbol{m}_2$ $\delta \boldsymbol{m}_{\pm}^+ = \delta \boldsymbol{m}_{\pm}^x + \mathrm{i} \delta \boldsymbol{m}_{\pm}^z$

$$i\gamma^{-1}\partial_t\delta m^+_+ = [-A\nabla^2 + K]\delta m^+_+$$

$$i\gamma^{-1}\partial_t\delta m^+_- = [-A\nabla^2 + K + 2J(x)]\delta m^+_-$$



- δm₊ is blind to interlayer
 exchange coupling J(x)
- J(x) acts like a trapping potential for δm.



 δm_{-} is a trapped magnon mode (~ 0.3 K = 0.09 meV)

Universal existence of DW magnons

$$i\gamma^{-1}\partial_t\delta m^+_- = [-A\partial_x^2 + K + 2J(x) + Ak_z^2]\delta m^+_-$$

- Beyond a critical point, the magnetic ground state deviates from interlayer FM
- The magnetic energy

$$E_{\text{mag}} = \int \left[\sum_{\alpha,\beta,l} \frac{A}{2} (\partial_{\beta} m_l^{\alpha})^2 - \sum_l \frac{K}{2} (m_l^{\gamma})^2 - \sum_{\alpha} J m_1^{\alpha} m_2^{\alpha} \right] dx$$

respects continuous rotational symmetry in the out-of-plane direction

• Goldstone magnon mode exists in the DW



MM and RM domain walls





- Both domain walls support Goldstone modes.
- Numerically verified by solving the LLG equation

Magnon disperson



Twisted bilayer Crl₃



- At large twisting angle, a metastable domain (green dot) appear
- Lattice relaxation has a sudden onset as a function of twisting angle (~ I degree)

Moire Magnon network



- Moire magnon network at 0.1 degree.
- There are trapped magnon modes on the FM/FM and AFM/ AFM domain walls

Discussion



In general, honeycomb lattice allows a 2nd order DMI [Owerre, JAP 2016, Kim et al, PRL 2016, Cheng, Okamoto, DX PRL 2016], which realize a magnonic version of the Haldane model for honeycomb FM. The topological edge mode is about 15 meV



We essentially deal with ID solitons. What happens at the corners where the magnon modes merge? [Hejazi, Luo and Balents, PNAS 2020]

Summary

- All stacking domain walls of bilayer Crl₃ support ID trapped magnons. Since stacking domain walls are insensitive to external fields, these magnon modes can serve as robust ID information channels.
- Small angle twist bilayer Crl₃ hosts a ID magnon network.
- Stacking domain walls are important in understanding magnetic properties of van der Waals magnets, and offer new way of stacking engineering of magnetic dynamics.

arXiv:2007.10398