

# Dynamics of 1D Quantum Magnets: High Field Fractionalization, BEC, and Transport

Wolfram Brenig

S. Grossjohann, WB, Phys. Rev. B 79, 094409 (2009)  
H. Kühne, WB, et al., Phys. Rev. B 80, 045110 (2009)

## Acknowledgements

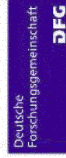
### Theory:

S. Grossjohann Braunschweig

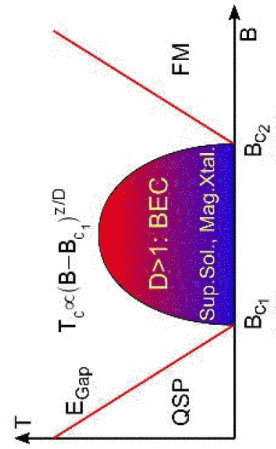
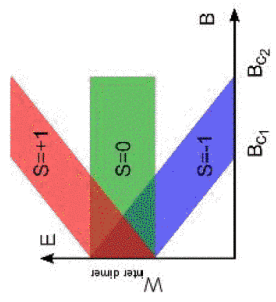
### Experiment:

H.-H. Klauss Dresden  
H. Kühne Dresden  
J.F. Litterst Braunschweig  
A.P. Reyes Tallahassee  
P.L. Kuhns "  
M.M. Turnbull Worcester MA  
C.P. Landee "

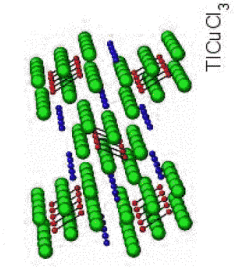
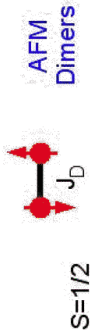
### Funding:



**Field Induced BEC in Quantum Spin Systems**



Affleck, PRB (90,91)  
Sorensen, et al., PRL (93)  
Giamarchi, et al., PRB (99)



- **D~3 S=1/2 dimer-systems** TiCuCl3  
Oosawa, et al., J. Phys., (99)  
Nikuni, et al., PRL (00)  
Ruegg, et al., Nature (03)
- **D~2 S=1/2 dimer-systems** BaCuSi2O6  
Sebastian, et al., Nature (06)

**Mapping**

• Bose liquids

$$H_{S^z=1} = \sum_{lm} (J_D - g\mu_B B) b_l^\dagger b_l + \sum_{lm} t_{lm} b_l^\dagger b_m + \sum_{lm} U_{lm}^{HC} b_l^\dagger b_m^\dagger b_m b_l$$

• Quantum magnets

$$H = \sum_{lm} J_D \tilde{S}_l \cdot \tilde{S}_m - \sum_{lm} g\mu_B B \tilde{S}_l^z + \sum_{lm, ij} \tilde{S}_l \cdot \tilde{S}_{ij}$$

bosons:  $n$

$\mu$

$U(1)$

$\langle \psi \rangle$

$\rho_S$

$\partial n / \partial \mu = \infty$

triplets:  $\langle S^z \rangle$

$B$

$O(2)_{\perp B}$

$\langle S^\pm \rangle$

$\langle S^\pm S^\mp \rangle$

$\partial \langle S^z \rangle / \partial B = \infty$

cold atoms:  $\mu$ -canonical

$U(1)$  exact

low- $T_c \sim nK$  (low dens., heavy)

quant. mag.: canonical

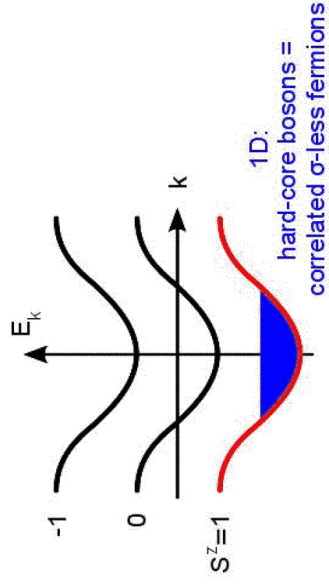
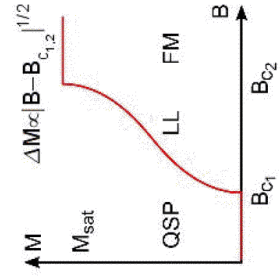
$O(2)$  ?

high- $T_c \sim 10 \dots 300K$  (high dens., light)

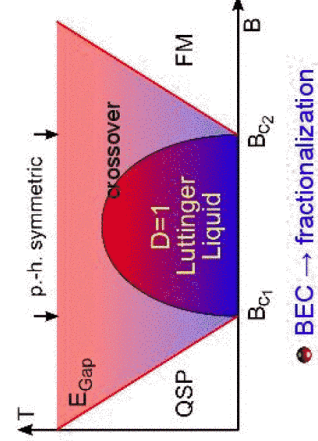
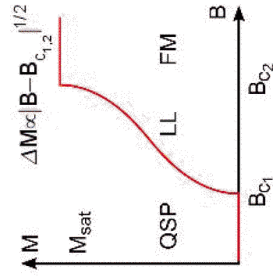
Demokritov, et al. Nature 443, (06)

Giamarchi, et al. PRB (99)  
Nikuni, et al. PRL (00)

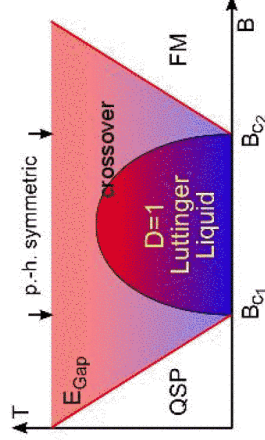
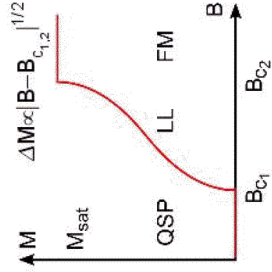
One Dimension



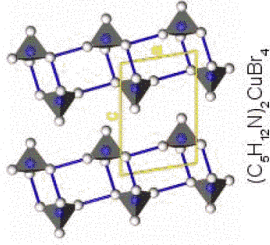
One Dimension



**One Dimension**



● BEC  $\rightarrow$  fractionalization



➤ **S=1/2 ladders**

Watson, et al., PRL (01)  
Lorenz, et al., PRL (08); Ruegg, et al. PRL (08) & (09)

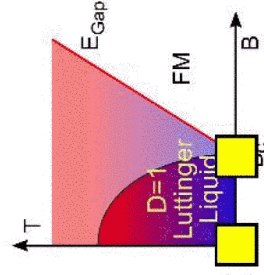
➤ **S=1/2 chains**

$Cu(C_4H_4N_2)(NO_3)_{3/2}$   
Stone, et al., PRL (03)

➤ **S=1 chains**

$Ni(C_5H_{14}N_2)_2N_3(PF_6)$   
 $NiCl_2 \cdot 4SC(NH_2)_2$   
Honda, et al., PRL (98);  
Zapf, et al., PRL (06); Zywglin, et al., PRL (07)

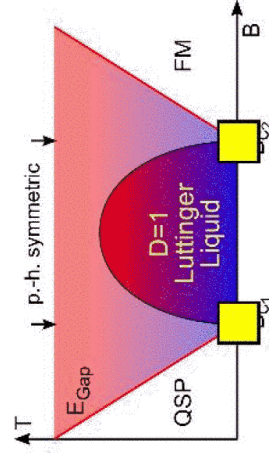
**Outline**



● Field induced critical dynamics  
of S-1/2 chain

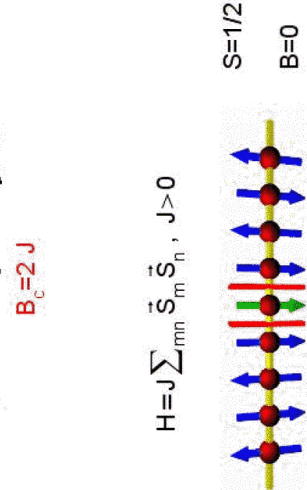
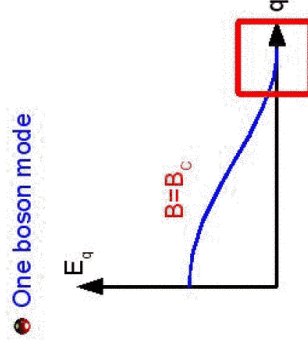
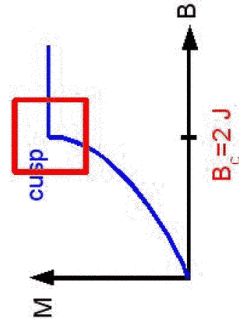
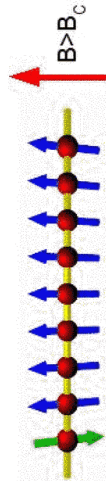
● Transport and spin diffusion in  
in Spin-1/2 chain

**Outline**

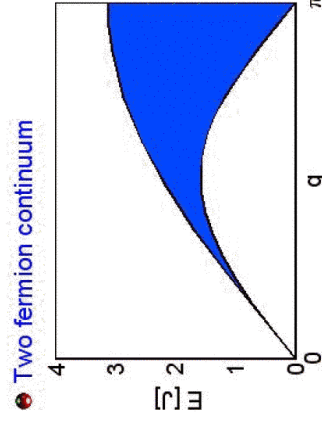


- Field induced critical dynamics of S-1/2 chain
- Transport and spin diffusion in Spin-1/2 chain
- Fractionalization in S-1 chains

**Fractionalization of Ising-Bosons**



$$H = J \sum_{mn} \vec{S}_m \cdot \vec{S}_n, \quad J > 0$$

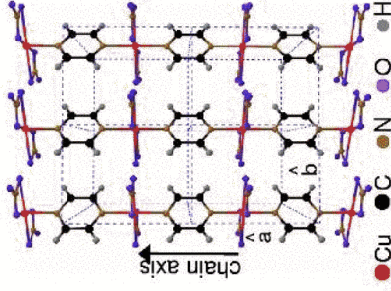


**Copper Pyrazine Dinitrate  $\text{Cu}(\text{C}_4\text{H}_4\text{N}_2)(\text{NO}_3)_2$**

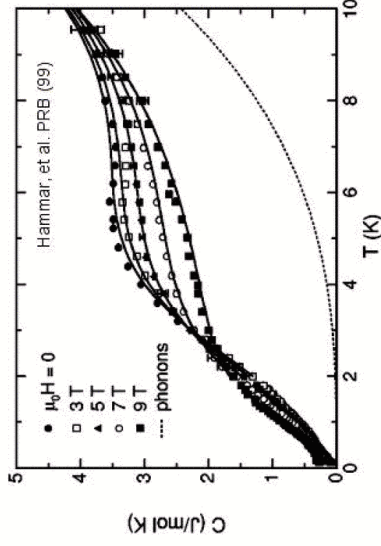
• **CuPzN: AFM  $S=1/2$  chain**

• **Thermodynamics**

Losee et al., J. Chem. Phys. (73)



$$H = \sum_{\langle ij \rangle} J_{a,ij} \mathbf{S}_i \cdot \mathbf{S}_j + J_{c,ij} \mathbf{S}_i \cdot \mathbf{S}_j$$



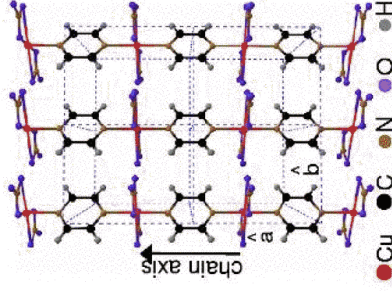
$$J/k_B \approx 10.6 \text{ K}$$

**Copper Pyrazine Dinitrate  $\text{Cu}(\text{C}_4\text{H}_4\text{N}_2)(\text{NO}_3)_2$**

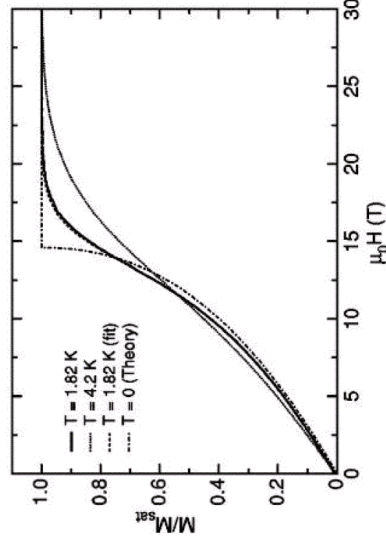
• **CuPzN: AFM  $S=1/2$  chain**

• **Thermodynamics**

Losee et al., J. Chem. Phys. (73)



$$H = \sum_{\langle ij \rangle} J_{a,ij} \mathbf{S}_i \cdot \mathbf{S}_j + J_{c,ij} \mathbf{S}_i \cdot \mathbf{S}_j$$



$$J/k_B \approx 10.6 \text{ K}$$

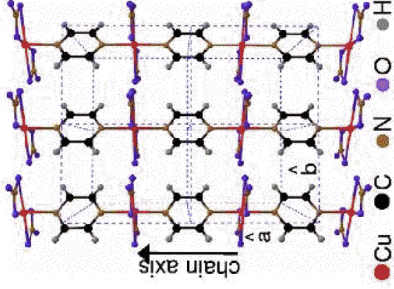
• **Saturation  $B_c = 14.5 \text{ T}$  accessible**

**Copper Pyrazine Dinitrate  $\text{Cu}(\text{C}_4\text{H}_4\text{N}_2)(\text{NO}_3)_2$**

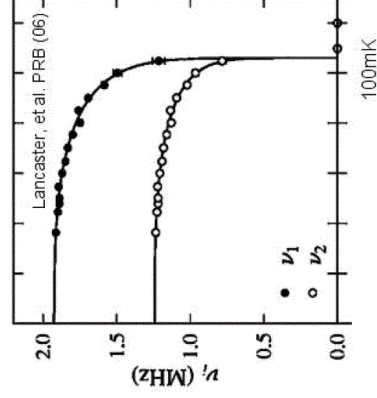
● CuPzN: AFM  $S=1/2$  chain

● Thermodynamics

Losee et al., J. Chem. Phys. (73)



$$H = \sum_{ij} J_{a,ij} \mathbf{S}_i \mathbf{S}_j + J_{c,ij} \mathbf{S}_i \mathbf{S}_j$$



● Saturation  $B_c = 14.5$  T accessible

$$J/k_B \approx 10.6 \text{ K}$$

$$|J_{\text{inter}}/J_{\text{intra}}| \sim 10^{-3} \quad T_N \approx 0.11 \text{ K}$$



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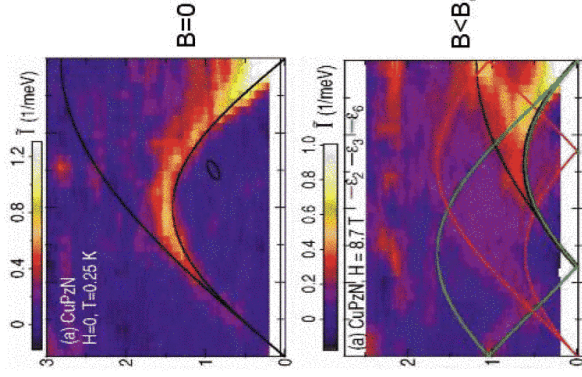


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**Copper Pyrazine Dinitrate Spin Dynamics**

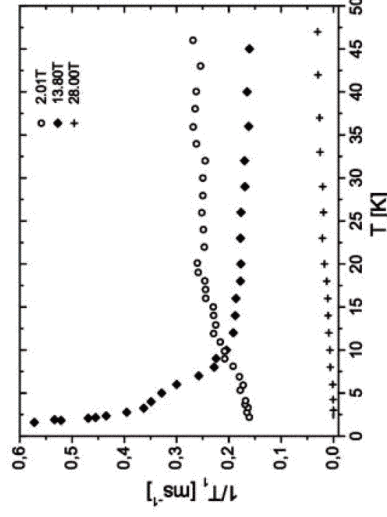
● INS

Stone, et al. PRL (03)



● NMR

H. Kühne, WB, et al., PRB 80, 045110 (2009)



fields  $0 < B < 2B_c$

temperatures  $0.25 \text{ J} < T < 5 \text{ J}$



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• Dynamical structure factor

$$S^{\alpha\beta}(\mathbf{q}, \omega) = \int_{-\infty}^{\infty} e^{i\omega t} \langle S_{\mathbf{q}}^{\alpha}(t) S_{-\mathbf{q}}^{\beta} \rangle dt$$

for spin-1/2 chain in B-field

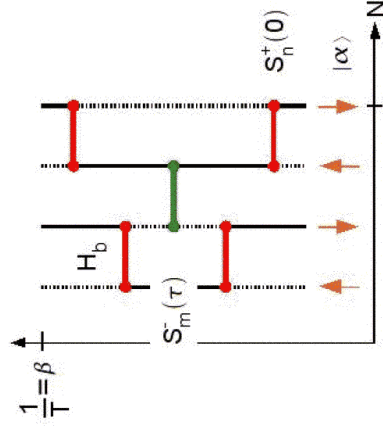
$$H = J \sum_l [S_l^z S_{l+1}^z + \Delta (S_{l+1}^x S_l^x + S_{l+1}^y S_l^y) - B S_l^z]$$

at finite T and B

- Bethe (n-spinon)  $S^{zz(x)}(\mathbf{q}, \omega)$ ,  $B \neq 0$ ,  $T = 0$
- Bosonization  $q \ll 1/a$
- ED  $N \leq 22 \dots 28$
- T(D)MRG  $S_{00}^{zz}(t)$ ,  $t < \sim 1/J$  or  $T = 0$
- QMC  $S^{zz(x)}(\mathbf{q}, \omega)$ ,  $T, B \neq 0$

• Expand partition function

$$Z = \text{Tr} \exp(-H/T) = \sum_{n,b} \text{Tr} [-H_b^n / n! T^n]$$



• Typ.:  $N \leq 10^5 \dots$ ,  $\beta \leq 1000/J$

**Quantum-Monte-Carlo**

Sandvik, PRB 69, (99)

• Spin dynamics

$$S_{mn}(\tau) = \langle S_m^-(\tau) S_n^+(0) \rangle$$

Temperature  $\exp(-H/T)$   
Time  $\exp(-i\tau H)$   
Frequency  $\omega$

$$S_{mn}(\tau) = \int d\omega K(\tau, \omega) S_{mn}(\omega)$$

Maximum Entropy Methods

• Typ.:  $N \leq 10^{2 \dots 3}$ ,  $\beta \leq 10 \dots 50/J$



**Transverse Dynamical Structure Factor**

$$\langle S_{q-q'}^x \rangle(\omega)$$

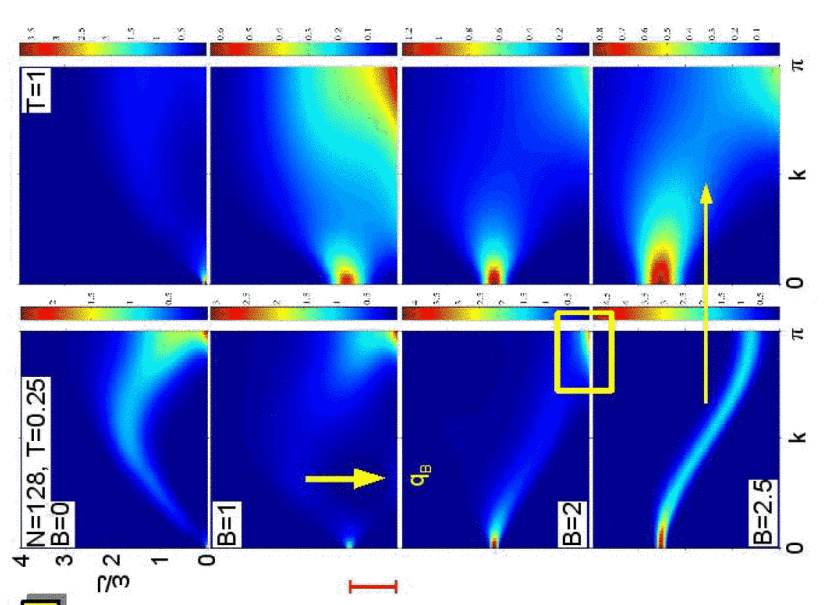
•  $B=0$ : continuum  
 $B>B_C$ : sharp mode

•  $q \sim 0, B \neq 0$ : gap

•  $B \neq 0, q=q_B$ : zone folded zero-mode

•  $B=B_C$ : critical dynamics

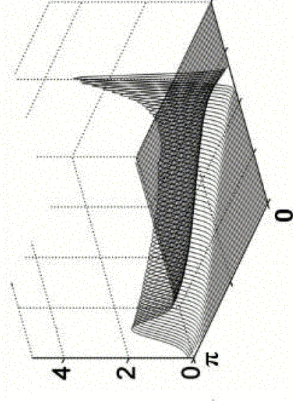
•  $T \uparrow$ : smearing



**Thermal Population of Multimagnon States**

• Momentum dependence of spectral weight

• 2-magnon contribution

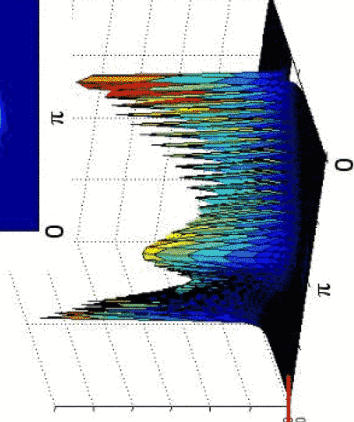
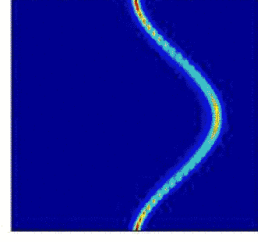


$$S^{xx}(q, \omega) = \sum_{mn} \frac{e^{-\beta E_m}}{Z} |m S_q^x n|^2 \delta(\omega - E_m + E_n)$$

$$= 1 \times n=1\text{-mag. sector}$$

$$+ e^{-\beta \Delta_B} \times n=2\text{-mag. sector}$$

$$+ e^{-2\beta \Delta_B} \times \dots$$



**Nuclear Magnetic Resonance: Critical Dynamics**

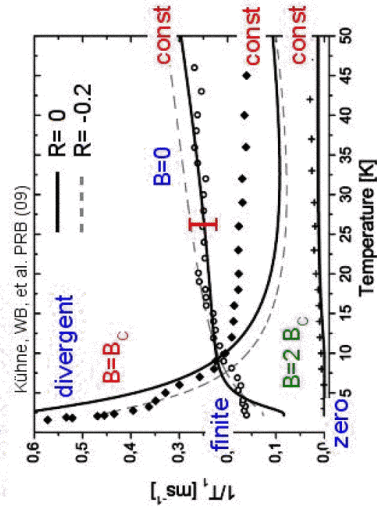
•  $T_1$ -relaxation

$$\frac{1}{T_1} = 2 \sum_{\mathbf{q}} |A(\mathbf{q})|^2 S^{\pm}(\mathbf{q}, \omega) \Big|_{\omega \rightarrow 0}$$

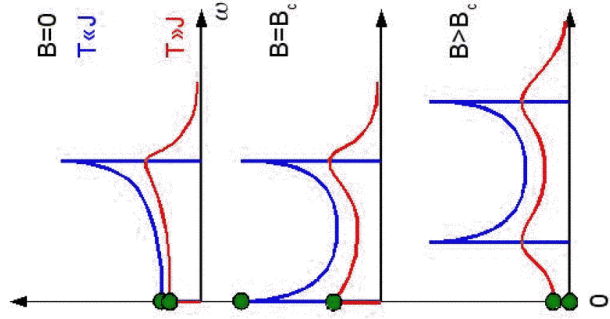
$$= \langle S_0^+(\tau) S_0^- \rangle \Big|_{\tau \rightarrow \omega, \omega \rightarrow 0}$$

~contact

• Expm. vs. QMC

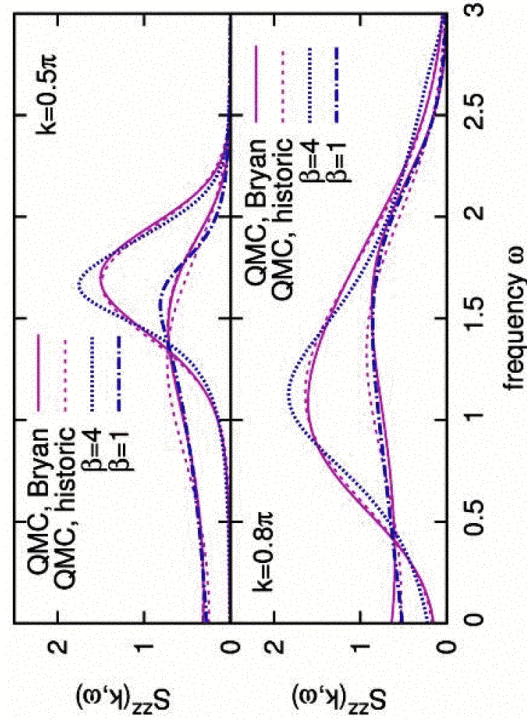


• Spin-DOS vs. B & T:



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**Comparing QMC and t-T-DMRG**



QMC

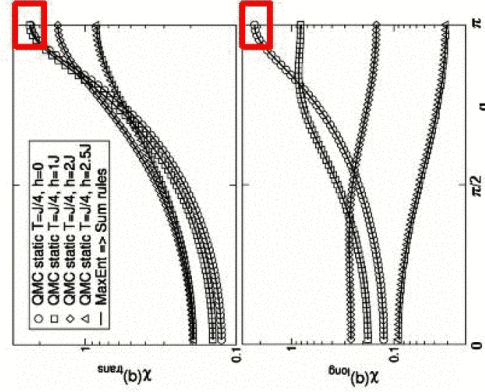
Grossjohann, Brenig  
PRB 79, 094409 (09)

t-T-DMRG

Barthel, Schollwöck, White  
PRB 79, 245101 (09)

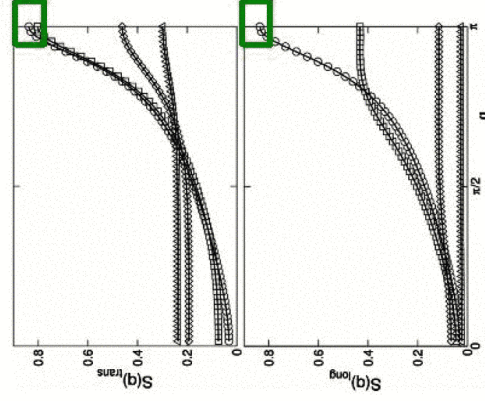
**Sum Rules**

$$\chi^{\alpha\beta}(\mathbf{q}) = \int_0^\beta d\tau \left( S_\alpha^\dagger(\tau) S_\beta^\dagger(-\mathbf{q}) \right)$$



$$S(\pi) = D_S \ln(T_S/T)^{3/2}$$

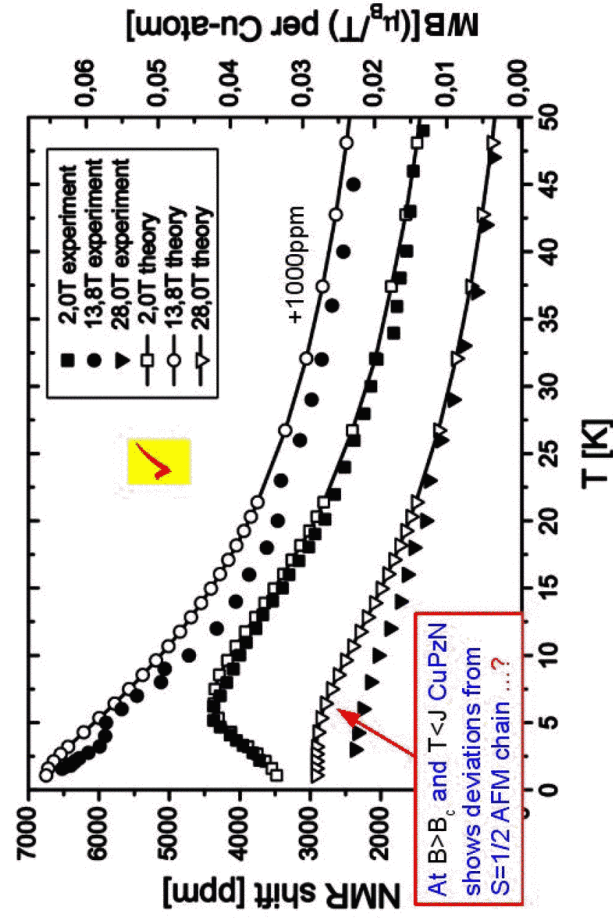
$$S^{\alpha\beta}(\mathbf{q}) = \int_{-\infty}^{+\infty} dt e^{-i\omega t} \left( S_\alpha^\dagger(t) S_\beta^\dagger(-\mathbf{q}) \right) \Big|_{t=0}$$



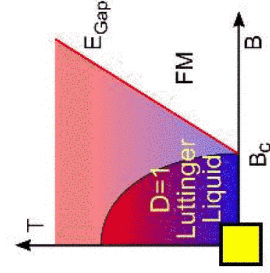
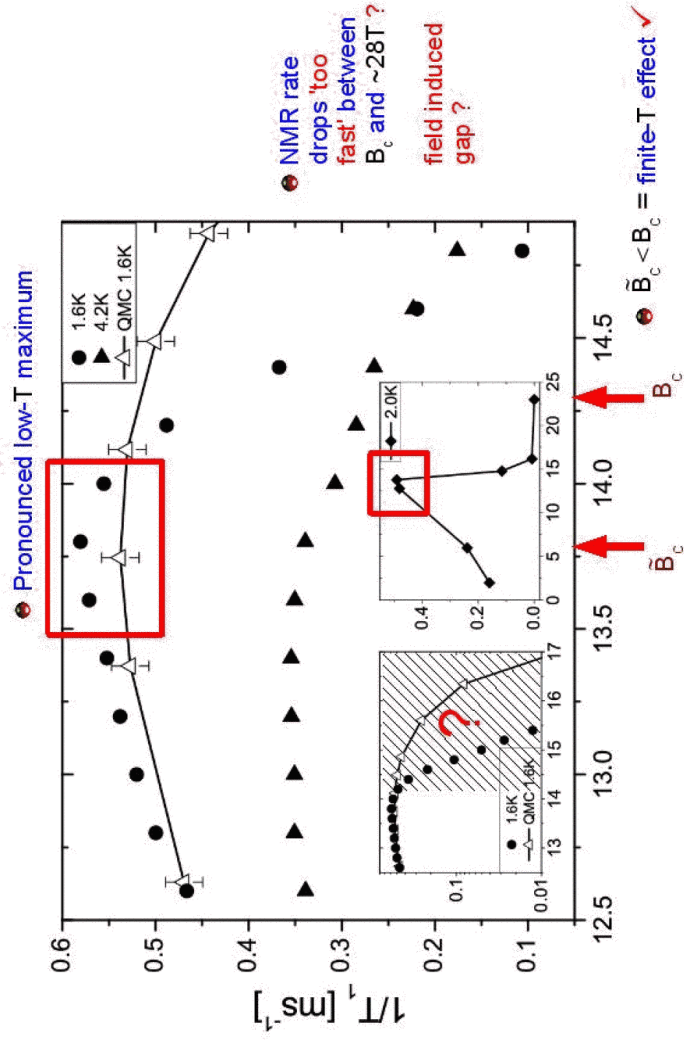
$$\chi(\pi) = (D_\chi/T) \ln(T_\chi/T)^{1/2}$$



**Knight Shift: NMR versus QMC**



Scanning the B-Field



● Field induced critical dynamics of S-1/2 chain

● Transport and spin diffusion in Spin-1/2 chain

● Fractionalization in S-1 chains

**Conductivity**

- spin current

$$\vec{J}_s(\mathbf{q}) \sim J \sum_{(m)} e^{iq} \vec{S}_m \times \vec{S}_{m+1}$$

↔ conductivity of strongly interacting spinless particles

$$H = J \sum_{\mathbf{r}} \left[ -\frac{1}{2} (c_{\mathbf{r}+1}^{\dagger} c_{\mathbf{r}} + \text{h.c.}) + \Delta (n_{\mathbf{r}} n_{\mathbf{r}+1} - n_{\mathbf{r}} + \frac{1}{4}) \right] - B \sum_{\mathbf{r}} n_{\mathbf{r}}$$

$$J_{\perp} / J$$

$$n_{\mathbf{r}} - \frac{1}{2} = S_{\mathbf{r}}^z$$

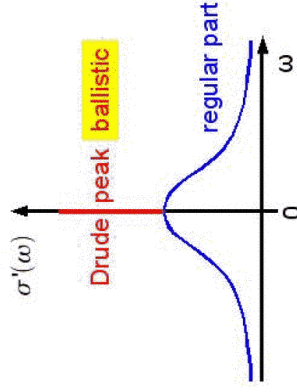
$$e^{i\pi\phi_{\mathbf{r}}} c_{\mathbf{r}}^{\dagger} = S_{\mathbf{r}}^{+}$$

Jordan-Wigner

- Optical conductivity

$$\sigma'(\omega) = D \delta(\omega) + \chi''_{\parallel}(\omega) / \omega$$

Drude regular



- regular ↔ longitudinal structure factor

$$\frac{1}{\omega} \chi''_{\parallel}(\omega) = \lim_{\mathbf{q} \rightarrow 0} \frac{\omega}{\mathbf{q}^2} \chi''_{zz}(\mathbf{q}, \omega)$$

**Transport Coefficients: Drude Weight**

- transport coefficient

Drude

$$D \sim \frac{1}{ZT} \sum_{E_m = E_n} e^{-E_m/T} |\langle m | j | n \rangle|^2$$

regular

$$\frac{\chi''_{jj}(\omega)}{\omega} \sim \frac{s(\omega)}{ZT} \sum_{E_m = E_n} |e^{-E_m/T} \times \langle m | j | n \rangle|^2 \delta(\omega - \Delta E)$$

T=0, |Δ| < 1 → D ≠ 0

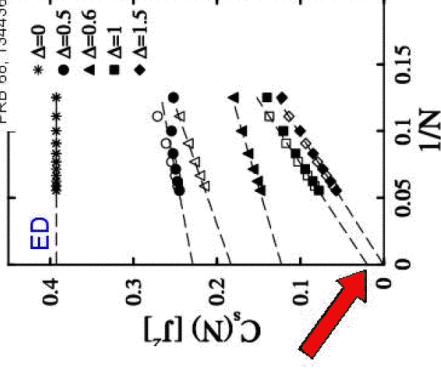
T ≠ 0, " BA: D(T) = ?

T ≠ 0, |Δ| = 1 BA: D = ?

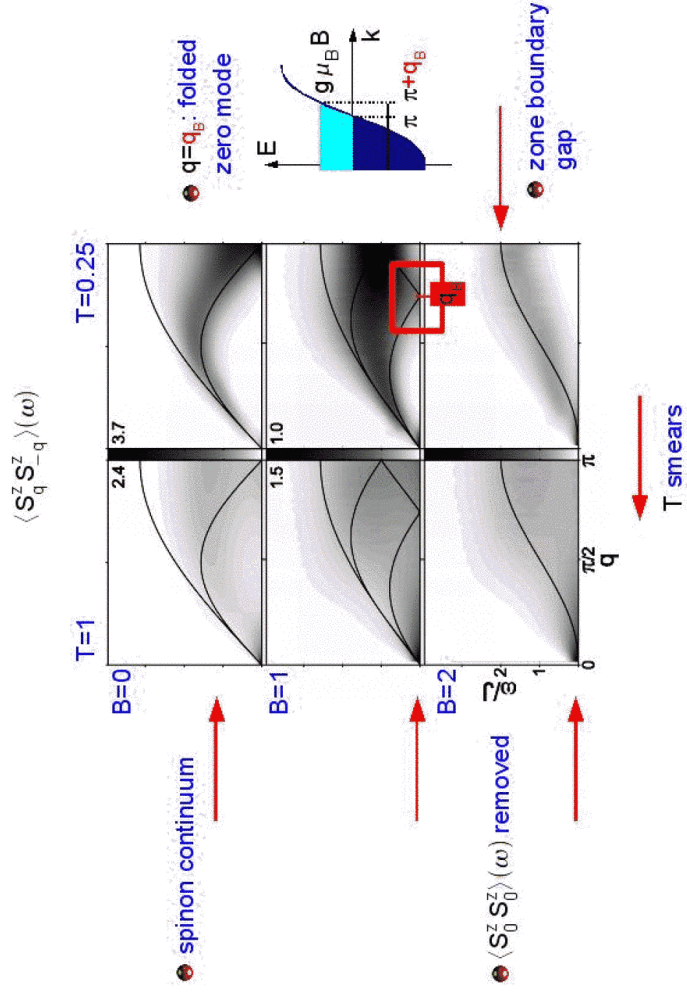
ED: D ≠ 0 Alvarez, Gros PRL (02), Heidrich-Meisner, WB; et al. PRB (03), Mukherjee, Shastya PRB (08), ...  
 tDMRG " S. Langer, et al. PRB (09)  
 MEQN " Michel, et al. PRB (08), Prosen, et al. J. StatM (09)  
 Shastya, Sutherland PRL (90)  
 Lan, Shi-Jian, Chin-PL (07); Benz, et al. EPJ (06); Zotos, PRL 82, 1764 (99)

- D: high-T residue  $\chi''_{jj} \neq 0$  !

Heidrich-Meisner, WB, et al. PRB 68, 134436 (2003)

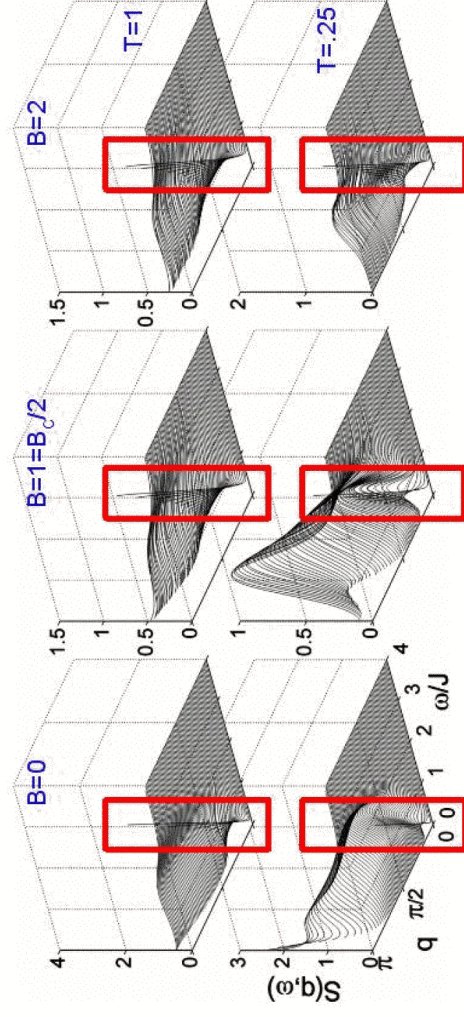


**Longitudinal Dynamical Structure Factor**



**Longitudinal Dynamical Structure Factor**

Diffusion ?



**Spin Diffusion and Susceptibility**

- Bosonization, 1<sup>st</sup> irrelevant perturbs umklapp & curvature,  $|\omega \pm vq| \ll T$

Sirker, Pereira, Affleck  
arXiv:0906.1978v1

$$\chi(\mathbf{q}, \omega) \approx \frac{Kvq^2/2\pi}{(1+b)\omega^2 - (1+c)v^2q^2 + i\pi g^2 T \omega}$$

$\uparrow$   $O(g^2)$       $\uparrow$   $g + \ln(g)/2$       $\uparrow$   $\sim \ln(J/T)$

Lukyanov, NuclePB  
522, 533 (1998)

diffusive for  $\omega \ll v^2 q^2$

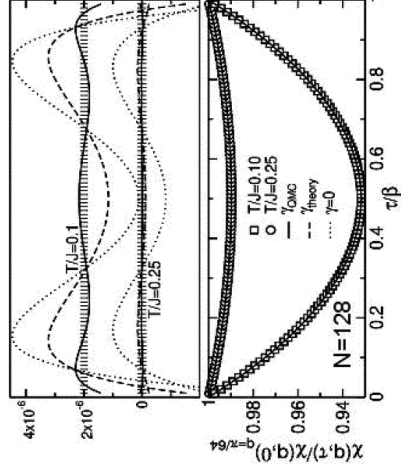
$$\Gamma = (1+c)v^2 / (\pi g^2 T)$$

$\pi g^2 T [J]$

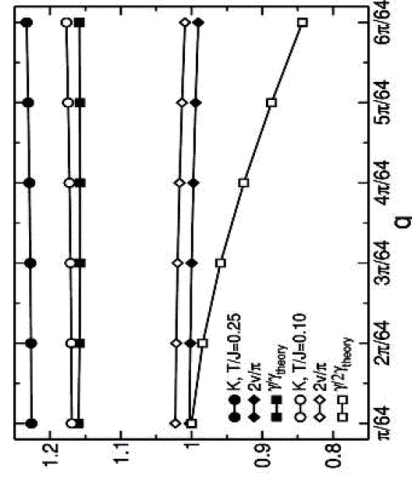
T [J]	QFT	DMRG	QMC
0.1	0.0096		0.0191
0.2	0.0297	0.0190	
0.25	0.0440		0.0507

- Compare to QMC

$$\chi(\mathbf{q}, \tau) = T \sum_n e^{i\omega_n \tau} \chi(\mathbf{q}, \omega_n)$$

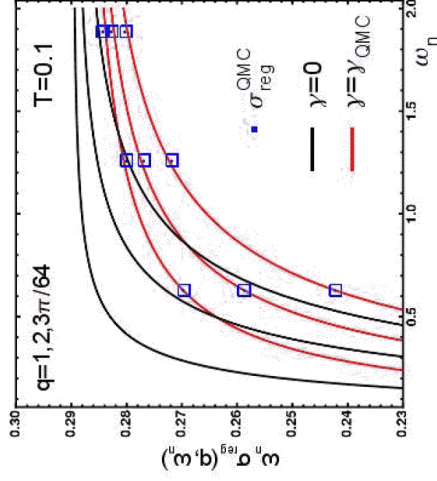


- momentum dependence



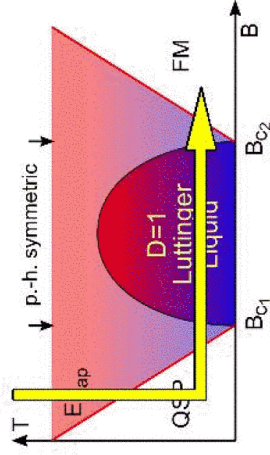
$K/v$  consistent with  $\chi(T)$  from BA ✓  
 $v \sim \pi/2$ , decreasing ✓  
 $\gamma = \Gamma(q) \rightarrow \Gamma = \Gamma(q) \leftrightarrow$  diff. to DMRG ?

- ...regular optical conductivity



compare QMC  $\rightarrow \sigma_{reg}^{QMC}(\mathbf{q}, \omega_n)$  with

$$\sigma_{reg}(\mathbf{q}, \omega_n) = \frac{\omega_n(Kv)/(2\pi)}{(1+b)\omega_n^2 + (1+c)v^2q^2 + 2\gamma|\omega_n|}$$



Field induced critical dynamics of S-1/2 chain

Transport and spin diffusion in Spin-1/2 chain

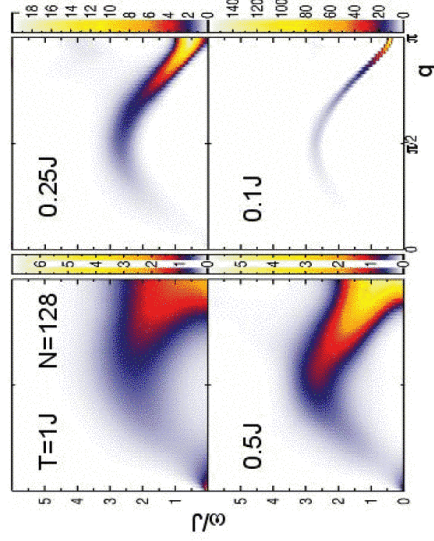
Fractionalization in S-1 chains



Cooling down at B=0



Transverse dynamic structure factor

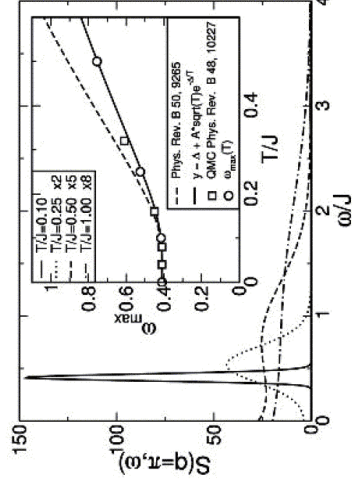


Haldane gap  $\Delta_0 \sim 0.41J$

Two-pt. continuum  $q < \pi/3$

consistent with INS Stone, et al. Nature (06)

T-dependence of  $S(q, \omega)$ -max



consistent with NLS-model

Jolicoeur, et al. PRB (94),

$$\Delta(T) \approx \Delta_0 + \sqrt{2\pi\Delta_0 T} e^{-\Delta_0/T}$$

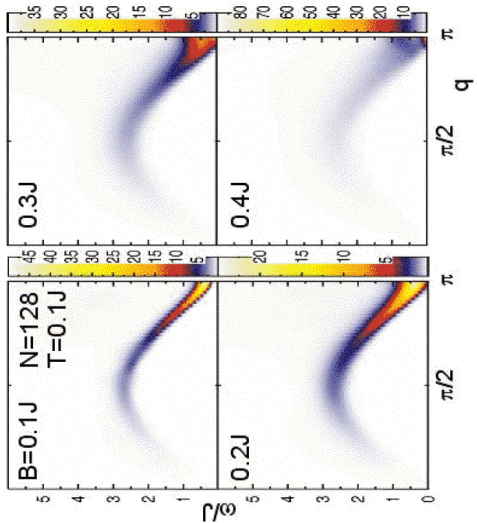
and B=0 QMC

Deiz, et al. PRB (93)

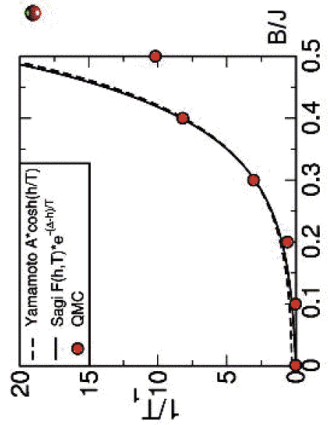
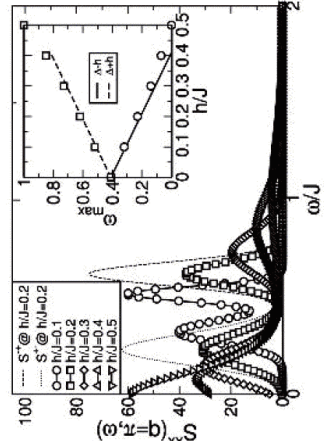


**Finite Magnetic Fields  $B < B_{c1} \approx 0.41J$**

● Transverse dynamical structure factor

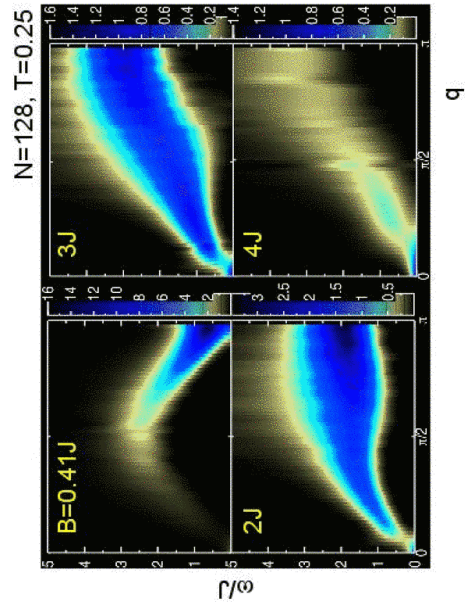


● gap closure at  $B/J = \Delta$



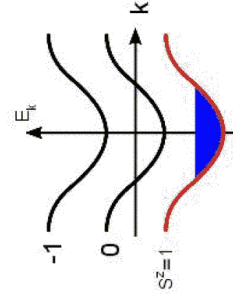
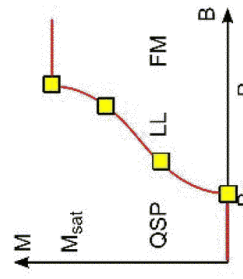
● NMR rate

● Longitudinal dynamic structure factor



- broad continua ✓
- incommensurate modes ?
- low intensity ●

**$B_{c1} \approx 0.41J < B < B_{c2} = 4J$  ... Fractionalization ... ?**



## Conclusion

- Evidence for field-induced critical dynamics of the AFHC material CuPzN due to triplet fractionalization.
- Regular part of spin conductivity is consistent with hydrodynamic limit with relaxation rate  $\ll T$ .
- First QMC results for the field induced fractionalization in the Haldane chain.