Fermionic atoms in optical lattices: Mott transition and metastable superconductivity



- Benedikt Binz, Rolf Helmes, David Rasch, Akos Rapp, Matthias Vojta, Achim Rosch,
- Institute for Theoretical Physics, University of Cologne
- Theo Costi

Institute for Solid State Research, Research Centre Jülich

experiments:

U. Schneider, L. Hackermüller, S. Will, Th. Best, I. Bloch

University of Mainz

- Cold atoms as quantum simulators for solid state systems: Mott transition
- New physics out of equilibrium: exotic metastable superconductivity

Quantum simulations for higher Tc?



- best: universal quantum computers
- more efficient to solve classical problems (prime number factorization, searches, ...)
- useful to simulate other quantum systems (Feynman 1981)



- simulate quantum system A (e.g. higher-Tc superconductors) by quantum system B (e.g. cold atoms)
- > useful if system B offers:
 - fewer complications (disorder, phonons, ...)
 - better control over parameters
 - other experimental probes
 - direct test of theories (!)
 - new physics (non-equilibrium)



Mott transition

Mott transition:







Sir Nevill Francis Mott 1905-1996

Hubbard model

$$H = -\mathbf{J} \sum_{\langle ij \rangle, \sigma = \uparrow \downarrow} c_{i\sigma}^{\dagger} c_{j\sigma} + \mathbf{U} \sum_{i} n_{i\uparrow} n_{i\downarrow}$$

kinetic energy () inte

interactions

1 electron per unit cell: metal for $\mathbf{U} \ll \mathbf{J}$

 $\frac{\text{Mott insulator for } \mathbf{U} \gg \mathbf{J}}{\text{Achim Rosch, Univ Cologne}}$



Mott transition

- example (V_{1-x}Cr_x)₂O₃
- but:

not described by Hubbard model orbitals, crystal fields, disorder long-range interactions



from Limelette et al. 03



trapped atoms in lattices made of light

- trap & cool atoms in lattice made of light
- optical lattice from standing waves of laser effective potential

$$egin{array}{rcl} V(r)&=&lpha(\omega)\langle E^2(r)
angle & & & & \ & & & \ & \propto & \cos^2(kx)+\cos^2(ky)+\cos^2(kz) \end{array}$$



- sufficiently high laser intensity:
 - only nearest neighbor hopping
 - only local interactions

perfect realization of Hubbard model

in external parabolic potential (Jaksch et al. 98)

- all parameters (J, U, parabolic potential) known and fully controllable !!
- bosons or fermions (or arbitrary mixtures)



Mott transition of bosons



$$H = -J \sum_{\langle ij \rangle} b_i^{\dagger} b_j + \mathbf{U} \sum_i n_i (n_i - 1) + \sum \mathbf{V_0} r_i^2 n$$

- small U: bose condensation & superfluidity
- large U: integer number of localized atom per site bosonic Mott insulator
- first realization: Greiner et al. 2002





Bloch 05

superfluid

Mott insulator

• detections: detect bose condensation!



fermionic Mott transition in optical lattices

- more fun: magnetism, superconductivity, ...
- simulate materials

- problem 1: cooling (less scattering due to Pauli principle)
- problem 2: detection

experiments:

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group of T. Esslinger (ETH)
R Jördens N. Strohmaier K. Günter H. Mori
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R. Jördens, N. Strohmaier, K. Günter, H. Moritz, T. Esslinger, Nature 08 group of I. Bloch (Mainz)

U. Schneider, L. Hackermüller, S. Will, Th. Best,

- I. Bloch, T. A. Costi, R. W. Helmes, D. Rasch,
- A. Rosch, Science 08

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theory needed !
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Theory for Mott transition?

no symmetry breaking, no obvious order parameter

method of choice: dynamical mean field theory (DMFT)

heavily used, e.g. for ab-initio description of correlated materials only approximation of **DMFT**: self-energy purely **local**

$$\Sigma_{ij}(\omega) \approx \delta_{ij} \Sigma_i(\omega)$$

naturally generalizable to inhomogeneous systems (Kotliar, Dobrosavljevic 97; Potthoff, Nolting 1999, Okamoto Millis 02, Freericks 04, Lee MacDonald 06, Snoek Hofstetter 08,....)

mapping to N single-impurity problems coupled by selfconsistency equation

solved using NRG

R. Helmes, T. Costi, A. Rosch, PRL 100, 056403 (2008)

Mott & metastable superconductivity in cold atom systems, KITP 7/09



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Mott transition of trapped atoms in optical lattices





Mott transition of trapped atoms in optical lattices



60 experimental results: Α 0.75 U. Schneider, L. Hackermül Fraction of atoms on doubly occupied sites *p*(%) 07 07 S. Will, Th. Best, I. Bloch about 10⁵ ⁴⁰K atoms 0.7 Renormalized Cloud Size $R_{
m sc}$ (d) in optical lattice initial T~0.15 E_F 0.65 0.6 0.5 1.5 Compression Et /12J \geq detection of fermionic Mott insulator 0.55 real experimental test of DMFT 0.5 without free parameters U/12J = 0> main experimental U/12J = 0.5U/12J = 1 problem: adiabaticity 0.45 U/12J = 1.50.5 1.5 2 parabolic potential · N^{2/3} Achim Rosch, Univ Cologne

Mott transition of trapped atoms in optical lattices

60



Α 0.75 action of atoms on doubly occupied sites p(%)07 07 07 **Deviations of theory** 0.7 and experiment for Size R_{sc} (d) large U? 0.65 0 م Experiment nonbud 1.5 0.5 1 6 Compression Et /12J adiabatic for large U? Renormalize 55 • U/12J = 0.5U/12J = 1 0.45 U/12J = 1.50.5 2 parabolic poten 2/3 Achim Rosch, Univ Cologne

Cold atoms as quantum simulator for Hubbard model?



- now: onset of Mott physics
- most needed: new cooling mechanism to reduce entropy by factor 4 for magnetism, factor 10 for d-wave superconductvity
- many ideas...
- likely within 2-3 years: detection of d-wave superconductivity in the 2d Hubbard model (if existent) not likely: any decent precision measurement

(Tc as function of doping, T, ...)

 already: useful as test for theories! see also: De Leo, Kollath, Georges, Ferrero, Parcollet, PRL 2008 Snoek, Titvinidze, Toke, Byczuk, Hofstetter, New J. Phys. 2008 Scarola, Pollet, Oitmaa, Troyer, PRL 2009,

more fun with cold atoms: physics out of equilibrium

simple example: metastable finite momentum superconductivity Achim Rosch, Univ Cologne

Lifetime of doubly occupied lattice sites



usually:

high energy states in interacting many body systems:

- decay rapidely due to huge phase space
- (exception: topologically stable states, solitions, domainwalls... of ordered phases)

periodic potential: formation of bands and band-gaps

regions without single-particle excitations





possiblity for (meta-) stable high energy states protected by energy conservation (not by topology)

Lifetime of doubly occupied lattice sites



Hubbard model: band-gap
$$\gg \mathbf{U} \gg J$$

$$H_h = -J \sum_{\langle ij \rangle, \sigma = \uparrow \downarrow} c^{\dagger}_{i\sigma} c_{j\sigma} + \mathbf{U} \sum_i n_{i\uparrow} n_{i\downarrow}$$

single particle states: $E_{\mathbf{k}} = -2J \sum_{i=x,y,z} \cos(k_i a)$ maximal single-particle energy: 12 **J** energy of doubly occupied site: **U**



How can doubly occupied state decay? necessary to create **many**, $\mathbf{N} \gtrsim \frac{\mathbf{U}}{12J}$, single particles excitation to absorb energy

$$t_d \gtrsim (const.)^{\mathbf{N}} \sim \frac{1}{J} \exp[const.\mathbf{U}/J]$$

exponentially large life time

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Winkler, Thalhammer, Lang, Grimm,Denschlag, Daley, Kantian, Büchler,Zoller, Nature 2006



by scattering from lattice photons

Strohmaier, Greif, Jördens, Tarruell, Moritz, Sensarma, Pekker, Altman, Demler, arXiv:0905.2963



numerics in 1d: quasistationary state within time-dependent DMRG for large U: Kollath, Läuchli, Altmann, PRL 2007

metastable superconductivity





detection of metastable superfluidity









- alternative to bose-condensation of pairs:
- phase separation (pairs clump together)
- e.g. analog situation for Bosons analyzed by
 D. Petrosyan, B. Schmidt, J. R. Anglin, M. Fleischhauer, Phys. Rev. A 76, 033606 (2007)

phase separation seems to be unavoidable for bosons

controlled calculation needed for fermionic case ! **usually** difficult, here: symmetries help

Metastable superconductivity



Simple case: only doubly occupied and empty sites:

$$H_{\text{eff}} = \frac{2J^2}{\mathbf{U}} \sum_{\langle ij \rangle} (1 - n_{i\uparrow})(1 - n_{i\downarrow})n_{j\uparrow}n_{j\downarrow} + \frac{2J^2}{\mathbf{U}} \sum_{\langle ij \rangle} c_{i\uparrow}^{\dagger} c_{i\downarrow}^{\dagger} c_{j\downarrow} c_{j\uparrow}$$

Trick:

rewrite as spin Hamiltonian: \downarrow e

↓ empty site↑ doubly occupied site

$$S_i^+ = (-1)^i c_{\uparrow i}^\dagger c_{\downarrow i}^\dagger$$

$$H_{ ext{eff}} = -rac{2J^2}{\mathbf{U}} \sum_{\langle ij
angle} \mathbf{S_i} \cdot \mathbf{S_j}$$

ferromagnetic Heisenberg model due to SU(2) symmetry in the charge sector

exact groundstates for uniform systems: magnetization in +z direction: band insulator magnetization in -z direction: no particles magnetization in x/y direction: s-wave superconductivity with momentum (π,π,π) , $(\eta$ -paring, C.N. Yang 1989) **Achim Rosch, Univ Cologne**





Is this a true superfluid with phase stiffness? **NO** (not yet)

ferromagnet: quadratic dispersion of spin-waves, $\omega \sim k^2$ superfluid: linear dispersion (necessary for superfluidity)

no superfluid ~ exact cancellation of attractions and hard-core repulsion of bosons

similarly: up to domain wall energy phase separation & superconductivity degenerate



breaking of $SU_C(2)$ symmetry of Hubbard model



Analyze leading corrections to SU_c(2) from corrections to Hubbard model

$$O(J^2/gap)$$





next nearest neighbor hopping favors phase separation but only to 2nd order $O(J^4/({\bf gap}^2 {f U}))$

assisted hopping

pair hopping

dominant: nearest neighbor interaction

$$\mathbf{U'} \sim \mathbf{U}(J/\mathbf{gap})^2 \gg J^4/(\mathbf{gap}^2\mathbf{U})$$

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breaking of $SU_C(2)$ symmetry of Hubbard model



trapping potential $lpha_t = \mathbf{V} N^{4/3} \mathbf{U} / J^2$

corrections to Hubbard model

$$\gamma = \mathbf{U'U}N^{2/3}/J^2 \gg 1$$



cloud expands exotic metastable s-wave superfluidity stabilized in strongly repulsive Hubbard model



crude estimate:

Bose condensation in trap for entropie per boson S<3.6 k_B

corresponding entropy per Fermion: S<1.8 k_B corresponding initial temperature before loading in trap: 0.22 T_F

presently reached (e.g. group of Bloch): 0.15 T_F Main question: effect of singly occupied sites!

Quantum simulations with cold atoms



- first simulations of correlated fermions in lattices: Mott transition
 - \succ next steps:

quantum magnetism (with 1/4th of entropy) d-wave superconductivity

- new phenomena out of equilibrium:
- metastable superconductivity





Helmes, Costi, Rosch, PRL **100**, 056403 (2008) Helmes, Costi, Rosch, PRL **101**, 066802 (2008) Schneider, Hackermüller, Will, Best, Bloch, Costi, Helmes, Rasch, A. Rosch, Science (2008) Rosch, Rasch, Binz, Vojta, PRL **101**, 265301 (2008)