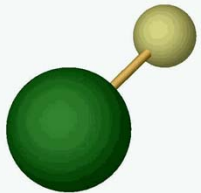


Prospects for new anapole moment measurements

- Proof of principle for the ZOMBIES nuclear anapole moment “factory”
- Near-future prospects for anapole measurements
- Questions for theorists: what measurements maximize physics impact?

DeMille



Group

Dave DeMille & Sid Cahn
Physics Department
Yale University

Past Funding
NSF

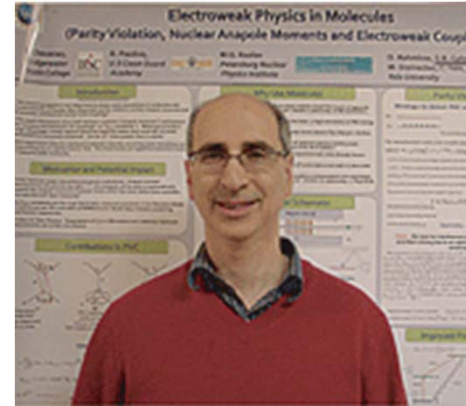


ZOMBIES @ Yale

Emine Altuntas



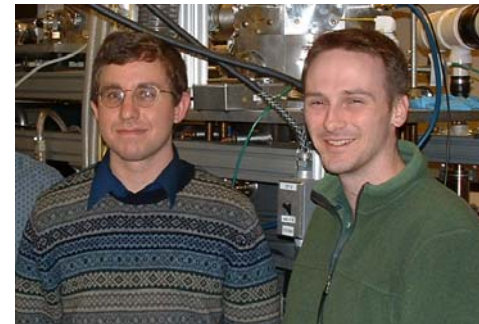
Sidney Cahn



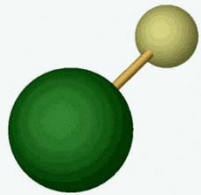
Jeffrey Ammon



David Rahmlow, Dennis Murphree



DeMille

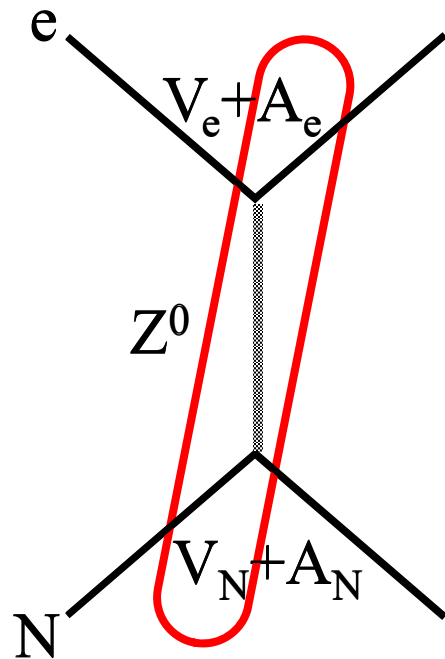


Group



Mechanisms for atomic/molecular parity violation

Axial electronic-vector nucleonic interaction



$A_e V_N$ term leads to term
in atomic Hamiltonian

$$H \propto Q_W G_F (\vec{\sigma} \cdot \vec{p}) \delta^3(\vec{r})$$

axial vector
associated with electron

short-range
Yukawa potential

Weak charge measured to 0.4%
[C. Wieman group, 1997]

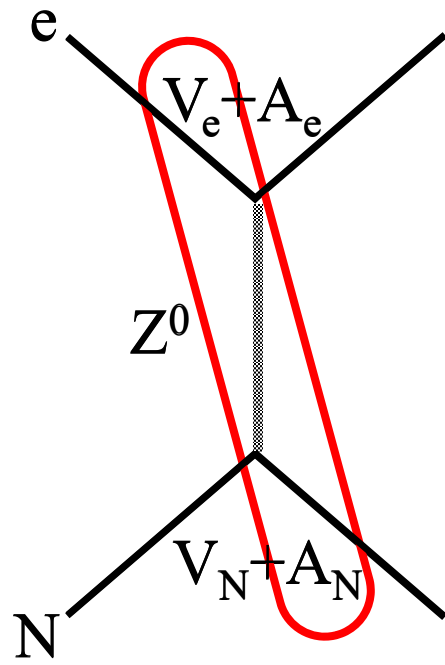
& interpreted at 0.3% level
[A. Derevianko 2010, V. Flambaum 2012]

Coherent coupling
to all nuclei =
"weak charge"
 $Q_W = -N + (1 - 4\sin^2\theta_W)Z$

\Rightarrow Running of $\sin^2\theta_W$ &
Limits on Z' bosons

Mechanisms for atomic/molecular parity violation

Vector electronic-axial nucleonic interaction



Coupling ONLY
to unpaired nucleon
coupling constant C_2

$V_e A_N$ term gives Hamiltonian:

$$H \propto C_2 G_F (\vec{\sigma} \cdot \vec{I})(\vec{\sigma} \cdot \vec{p}) \delta^3(\vec{r})$$

Nuclear spin I
= axial vector
associated with nucleon

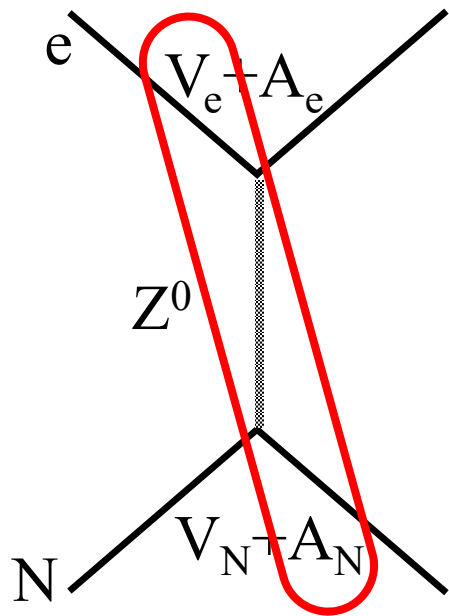
C_2 numerically small:
 $V_e/A_e = (1 - 4\sin^2\theta_W) \sim .08$

Bottom line:

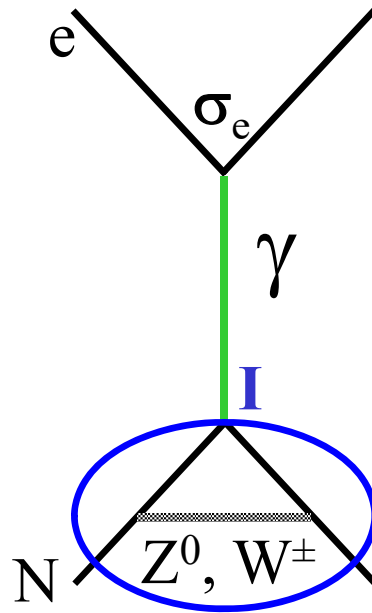
$$V_e A_N / A_e V_N \sim 10^{-3}$$

(for heavy atoms)

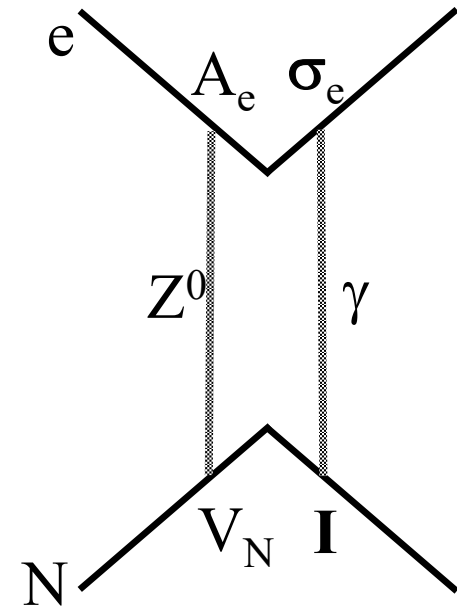
*Suppression of tree-level NSD-PV
makes radiative corrections non-negligible*



Tree-level
NSD-PV
from suppressed
 $V_e A_N$ term:
 C_2 subject to QCD
renormalization
similar to g_A



HPV interactions
inside nucleus induce
nuclear
"anapole moment":
couples to electron
magnetically

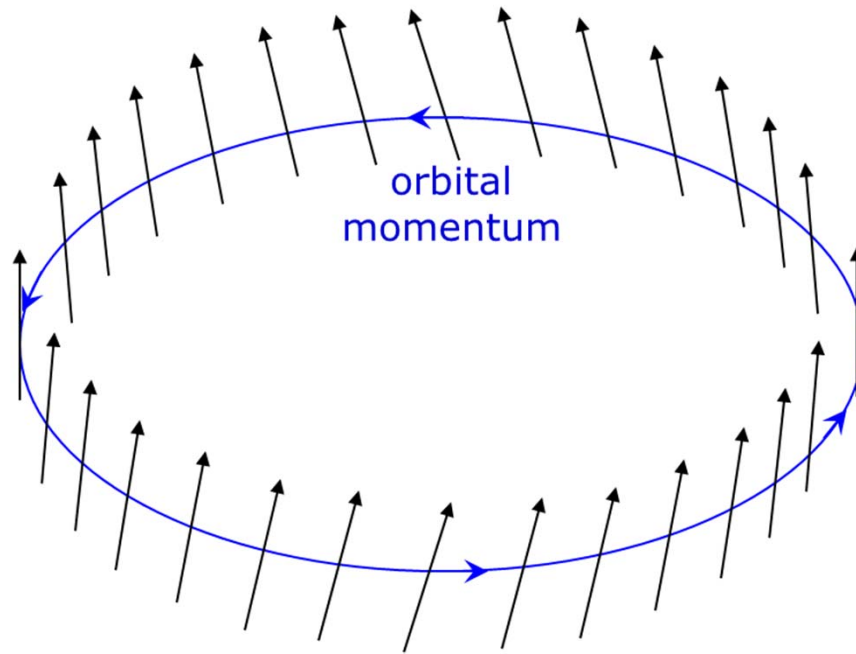


Coherent sum of
weak charge Q_W
and
electromagnetic
hyperfine interaction

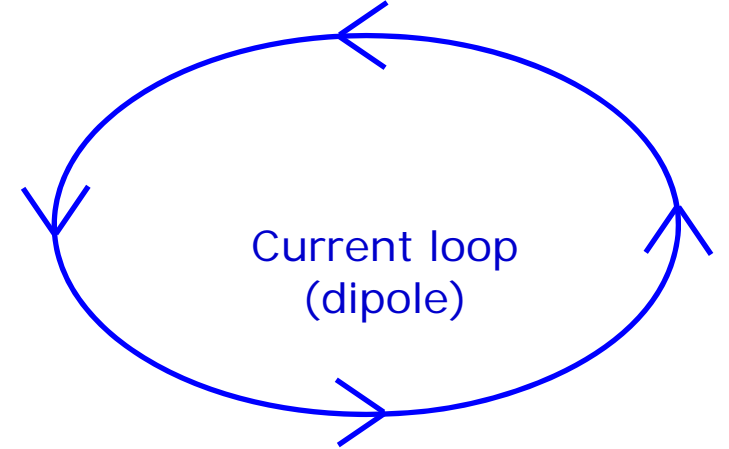
$$H_{NSD-PV} \propto (\kappa'_2 + \kappa'_a + \kappa'_Q) G_F (\vec{\sigma} \cdot \vec{I})(\vec{\sigma} \cdot \vec{p}) \delta^3(\vec{r})$$

*HPV in nucleus induces nuclear spin helix
= magnetic dipole + anapole*

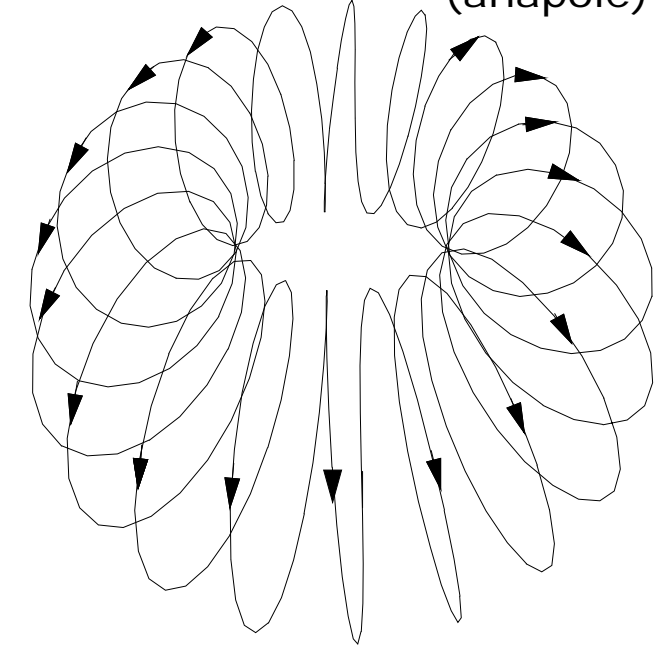
$$H_{HPV} \propto \vec{\sigma}_N \cdot \vec{p}_N \Rightarrow \begin{array}{l} \text{spin tilted} \\ \text{along momentum} \end{array}$$



=



+ Current helix (anapole)

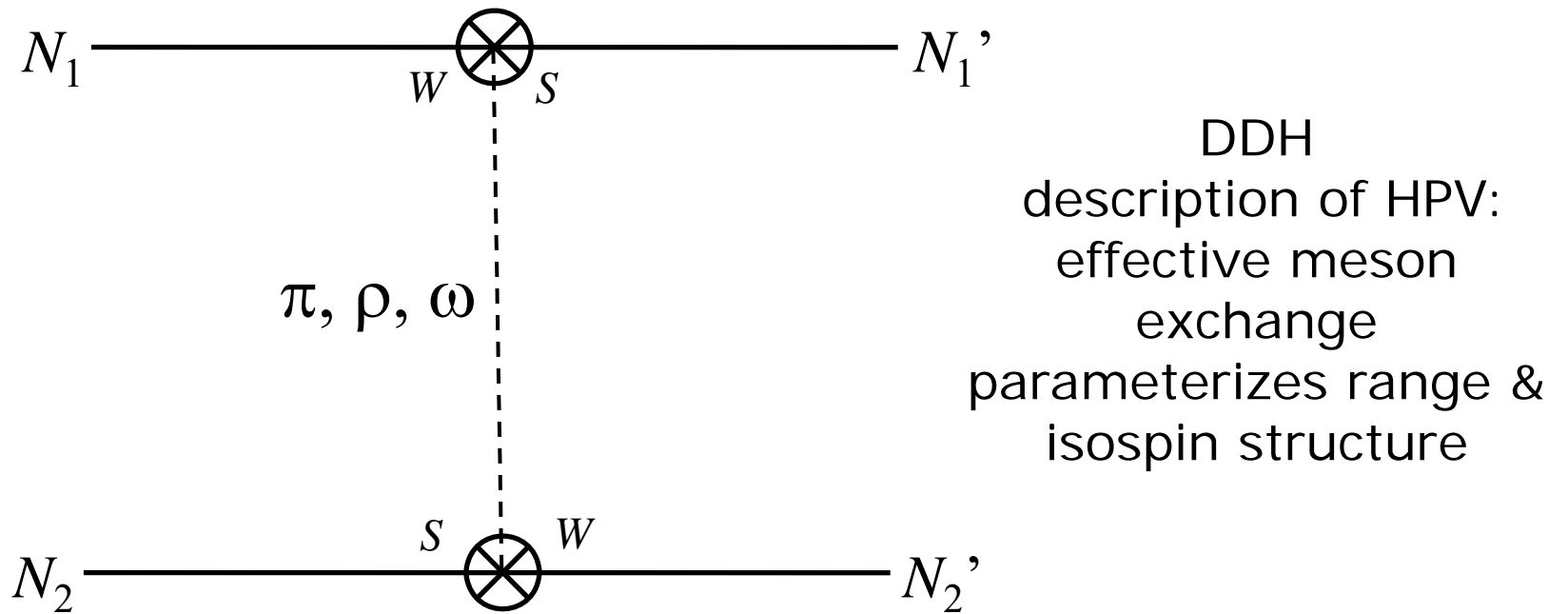


Simple model for nuclear anapole
(valence nucleon + constant-density core):

$$\vec{a} \propto g_{eff} A^{2/3} \hat{I}$$

Microscopic physics of the nuclear anapole moment

Nucleon-nucleon HPV interactions perturb nuclear structure:



Hamiltonian for unpaired nucleon interacting with paired core
gives spin-momentum correlation

$$H_{HPV} \sim G_F (\vec{\sigma}_N \cdot \vec{p}_N) \sum_i g_{\text{eff},i} F_i(\vec{r}, \vec{\tau})$$

6 terms in principle;

2 linear combinations estimated important for anapole (DDH)

Anapole moments in DDH parameterization

Nuclear anapole moments

W. C. HAXTON, C.-P. LIU, AND M. J. RAMSEY-MUSOLF

PHYSICAL REVIEW C **65** 045502

TABLE VII. PNC observables and corresponding theoretical predictions, decomposed into the designated weak-coupling combinations.

Observable	Expt. ($\times 10^7$)	$f_\pi - 0.12h_\rho^1 - 0.18h_\omega^1$	$h_\rho^0 + 0.7h_\omega^0$	h_ρ^1	h_ρ^2	h_ω^0	h_ω^1
$A_L^{pp}(13.6)$	-0.93 ± 0.21		0.043	0.043	0.017	0.009	0.039
$A_L^{pp}(45)$	-1.57 ± 0.23		0.079	0.079	0.032	0.018	0.073
$A_L^{pp}(221)$	0.84 ± 0.34		-0.030	-0.030	-0.012	0.021	
$A_L^{p\alpha}(46)$	-3.34 ± 0.93	-0.340	0.140	0.006		-0.039	-0.002
$P_\gamma(^{18}\text{F})$	1200 ± 3860	4385		34			-44
$A_\gamma(^{19}\text{F})$	-740 ± 190	-94.2	34.1	-1.1		-4.5	-0.1
$\langle A_1 \rangle / e$, Cs	800 ± 140	60.7	-15.8	3.4	0.4	1.0	6.1
$\langle A_1 \rangle / e$, Tl	370 ± 390	-18.0	3.8	-1.8	-0.3	0.1	-2.0

3 contributions to NSD-PV: scaling with Z & A

$$H_{NSD-PV} \propto (\kappa'_2 + \kappa'_a + \kappa'_Q) G_F (\vec{\sigma} \cdot \vec{I}) (\vec{\sigma} \cdot \vec{p}) \delta^3(\vec{r})$$

Overall Z^2

$\kappa'_Q \propto A^{2/3}$ small ($< \kappa'_a/4$)
& well understood
--ignore

$$\kappa'_{2P} = -\kappa'_{2N}$$

$$\cong -g_A (1-4\sin^2\theta_W)/2 \cong -.05$$

- ~independent of A
- $\mathcal{O}(20\%)$ corrections from $SU(3)_f$
- $\mathcal{O}(100\%)$ expt. uncertainty
- Quenching in larger nuclei like g_A ?

Simple shell model:
valence nucleon over closed core

$$\kappa'_a \approx \frac{9}{10} \frac{\alpha\mu}{mr_0} A^{2/3} g_{eff} \approx .05 g_{eff} \left(\frac{A}{50} \right)^{2/3}$$

$$g_{eff,p} \approx 4 - 6; \quad g_{eff,n} \approx 0.1 - 1;$$

3 contributions to NSD-PV: scaling with Z & A

$$H_{NSD-PV} \propto (\kappa_2 + \kappa_a + \kappa_Q) G_F (\vec{\sigma} \cdot \vec{I}) (\vec{\sigma} \cdot \vec{p}) \delta^3(\vec{r})$$

Overall Z^2

$$\kappa'_{2P} = -\kappa'_{2N} \approx -.05$$

$$\kappa'_a \approx .05 g_{eff} \left(\frac{A}{50} \right)^{2/3}$$

$$(g_{eff,P} \cong 4, g_{eff,N} \lesssim 1)$$

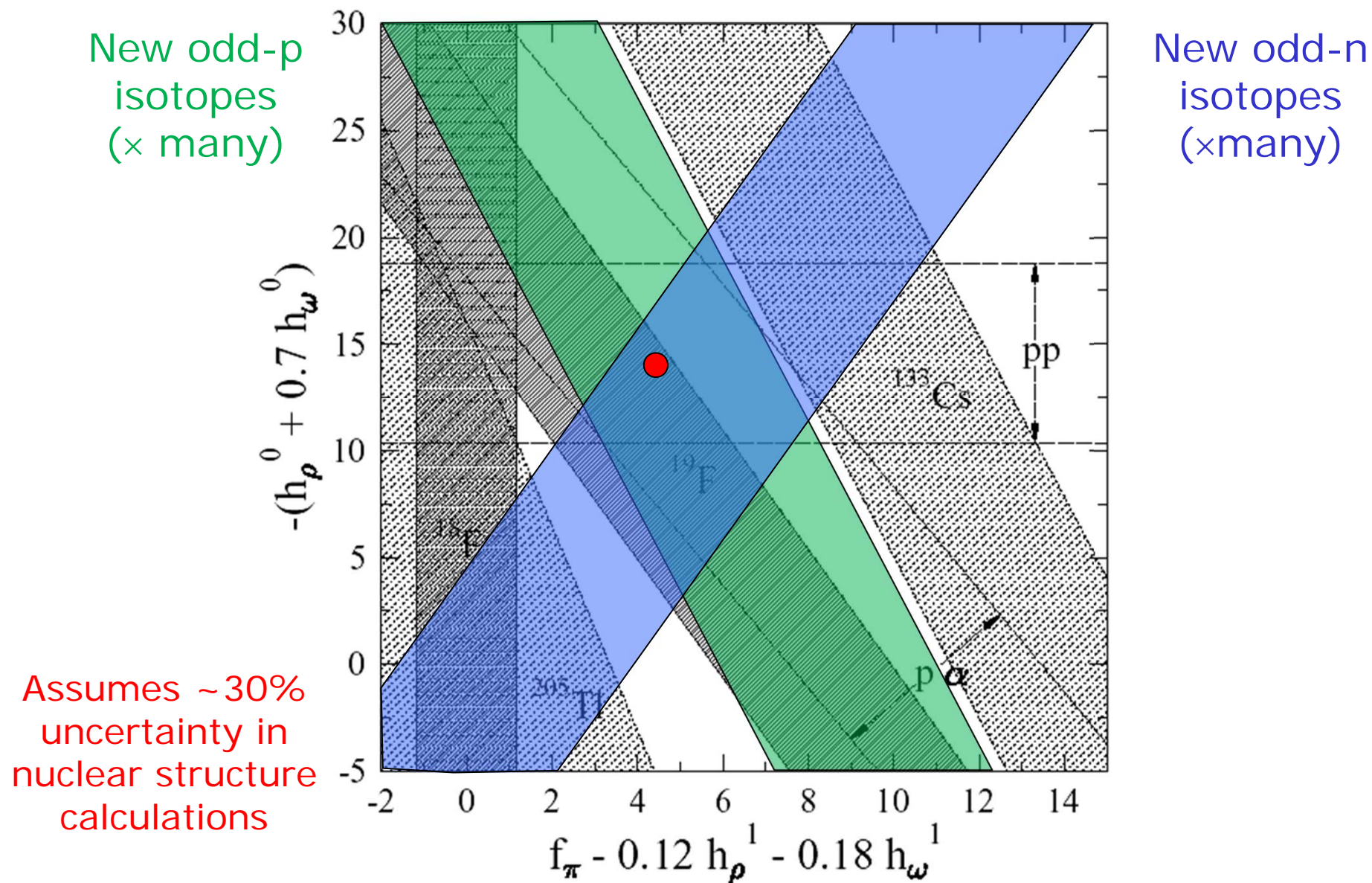
In heavy atoms, anapole term dominates: $|\kappa'_a| > |\kappa'_2|$

(Collective enhancement causes radiative correction > tree level...!)

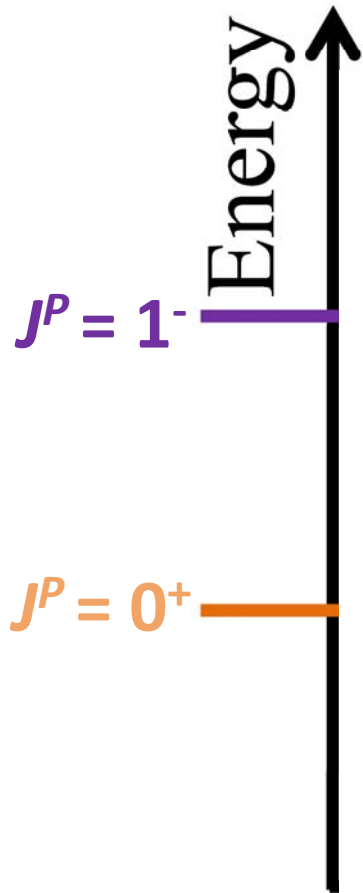
$$|\kappa_a| \approx |\kappa_2| \text{ for } A \approx 10 \text{ (odd proton)}$$

$$A \approx 100 \text{ (odd neutron)}$$

“Old style” plot of HPNC measurements including anapole measurements (past & future)



Enhanced NSD-PV mixing in simple molecules



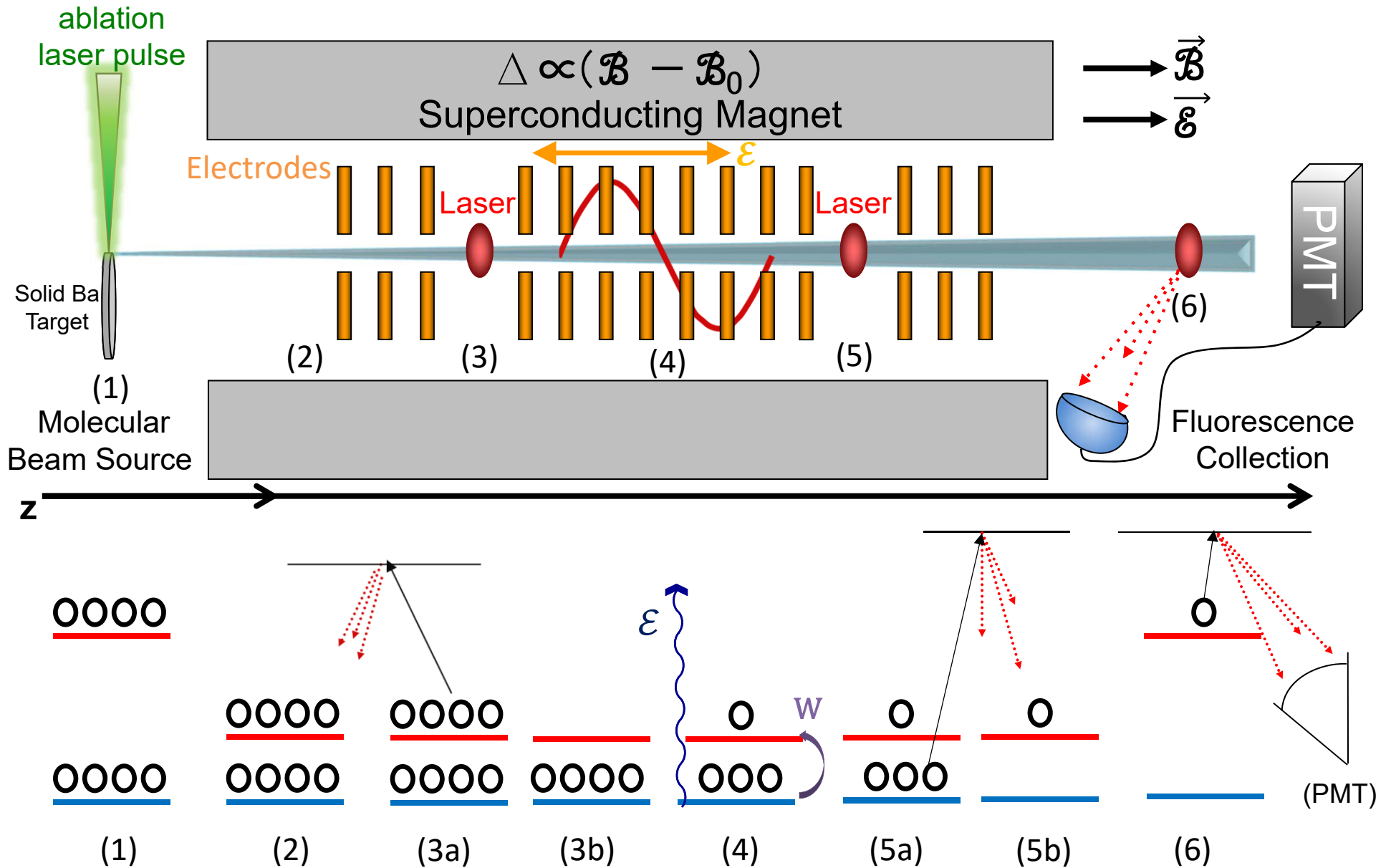
+ / - mixing

$$\eta = \frac{\langle + | H_{PV} | - \rangle}{E_+ - E_-}$$

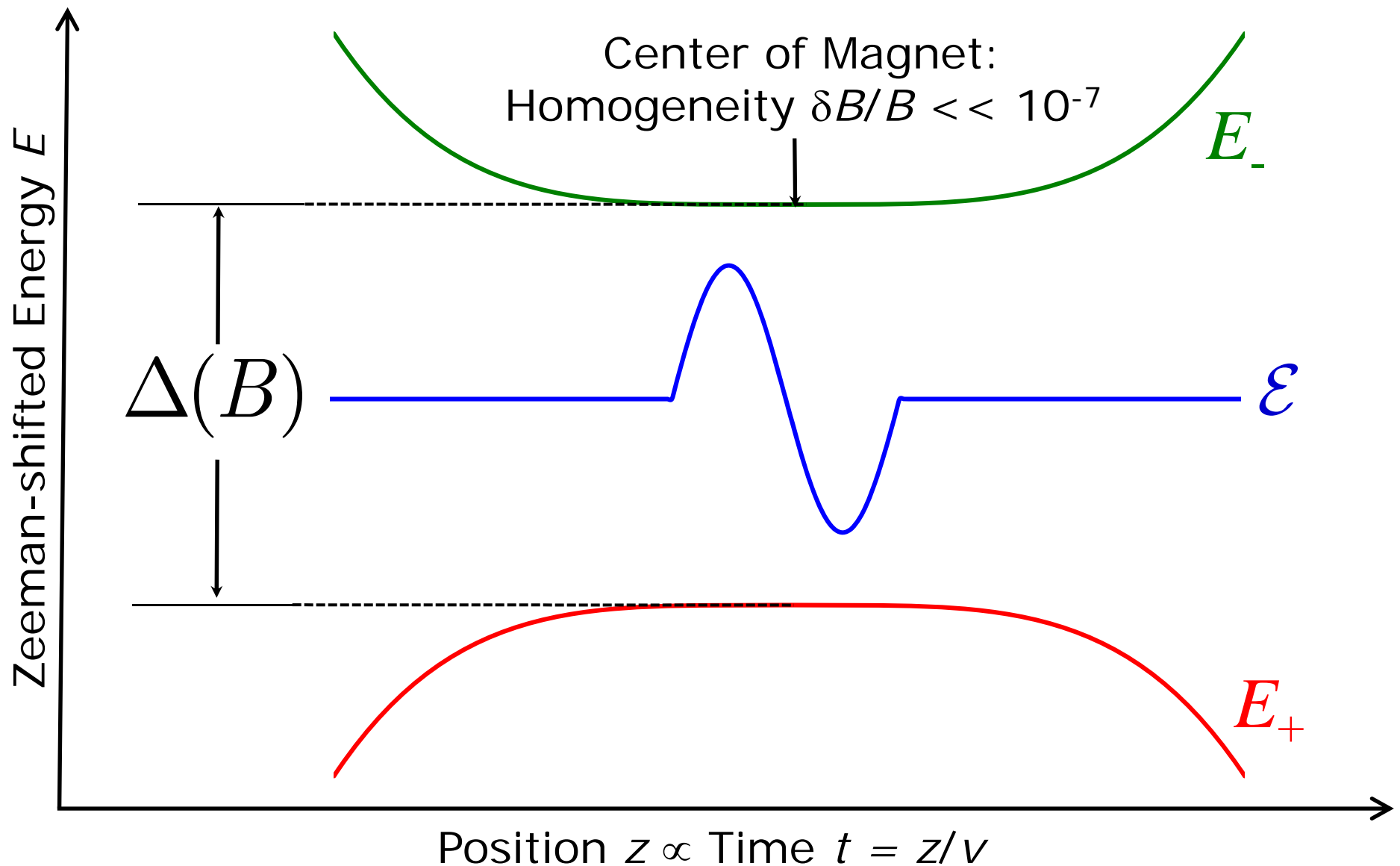
Naturally small rotational splitting ($\sim 10^{-4}$ eV vs. ~ 1 eV in atoms)
can be bridged w/Zeeman shift:

$\gtrsim 10^{11}$ enhanced PV mixing vs. classic experiments with atoms

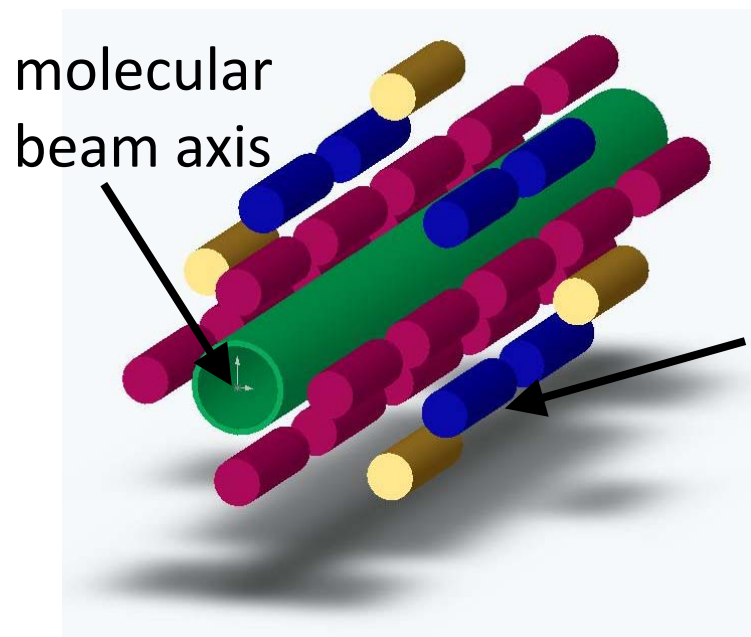
ZOMBIES experimental schematic



*Stark interference method:
apply oscillating \mathcal{E} -field to mix nearly-degenerate levels*



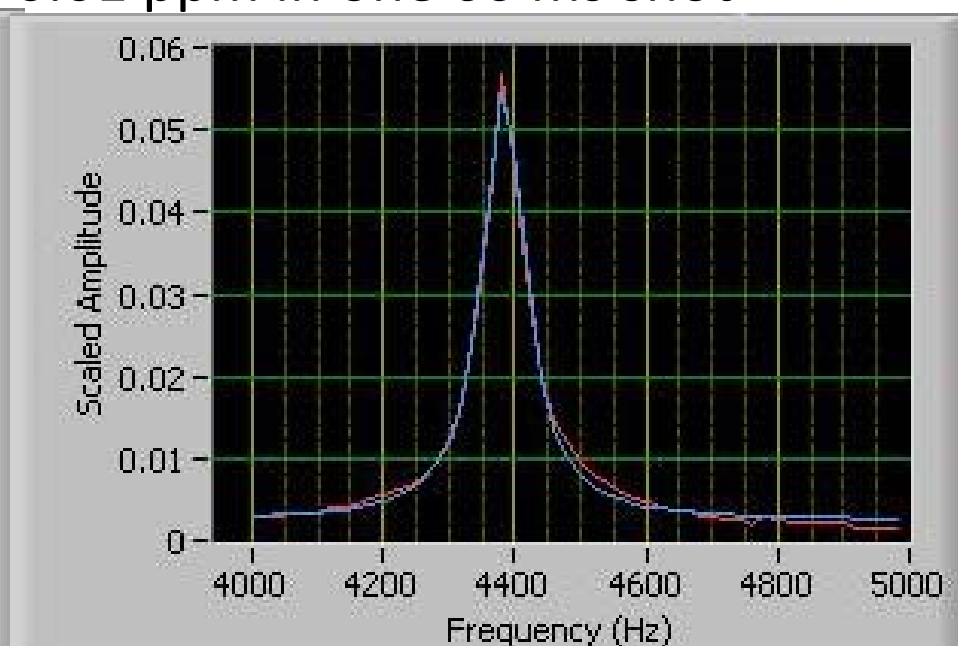
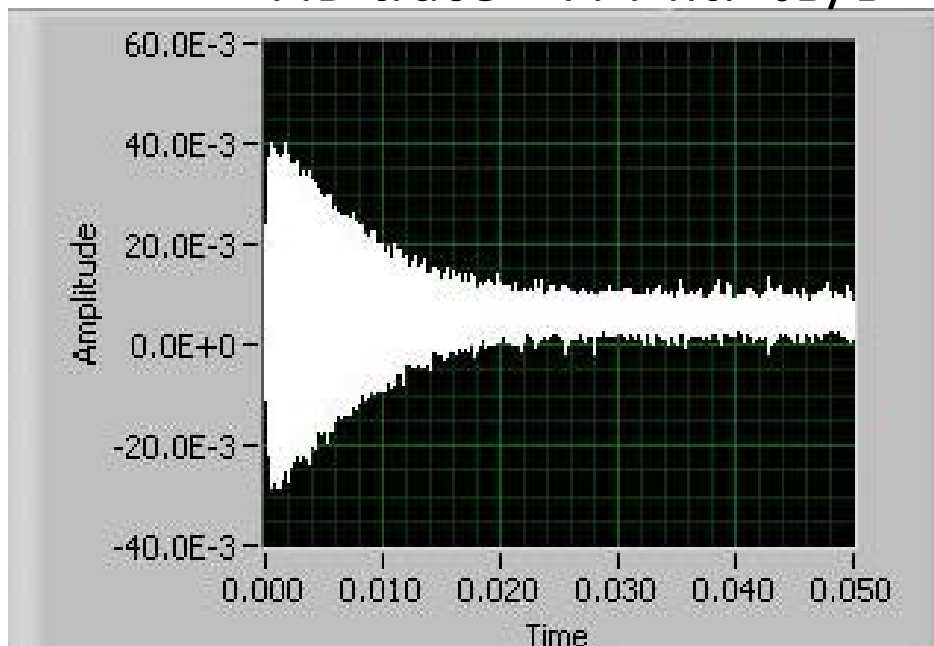
Magnetic field measurement: initial/crude



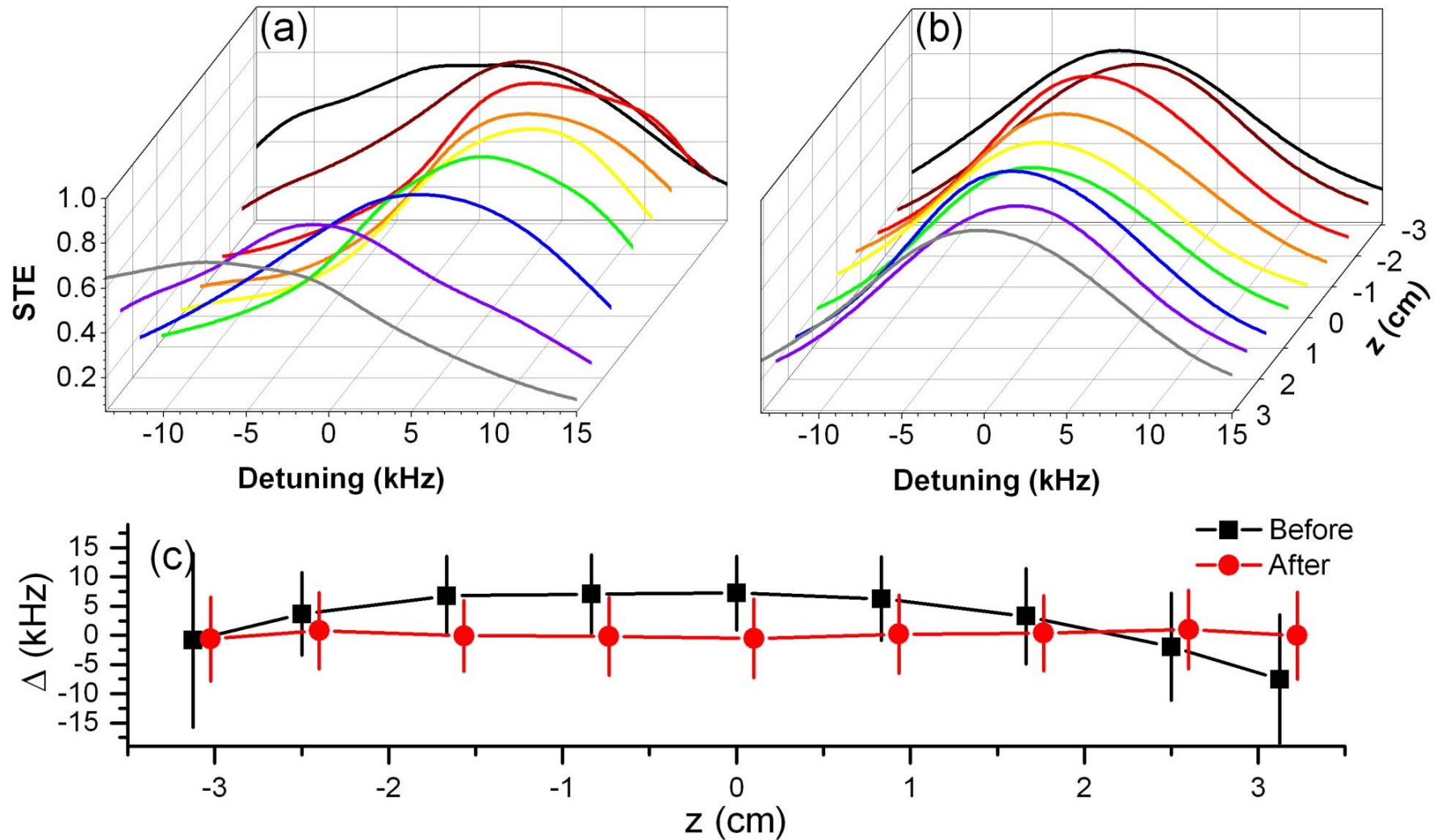
broadband probe on flex circuit



FID trace + FFT fit: $\delta B/B = 0.01$ ppm in one 60 ms shot



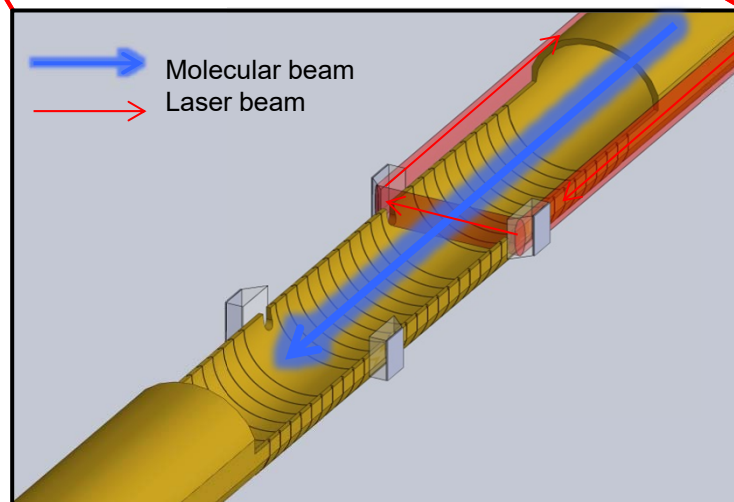
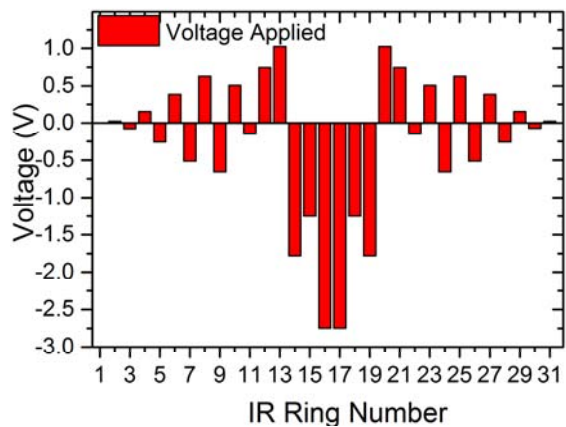
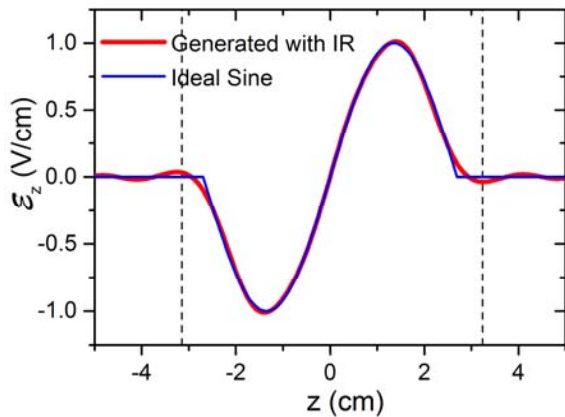
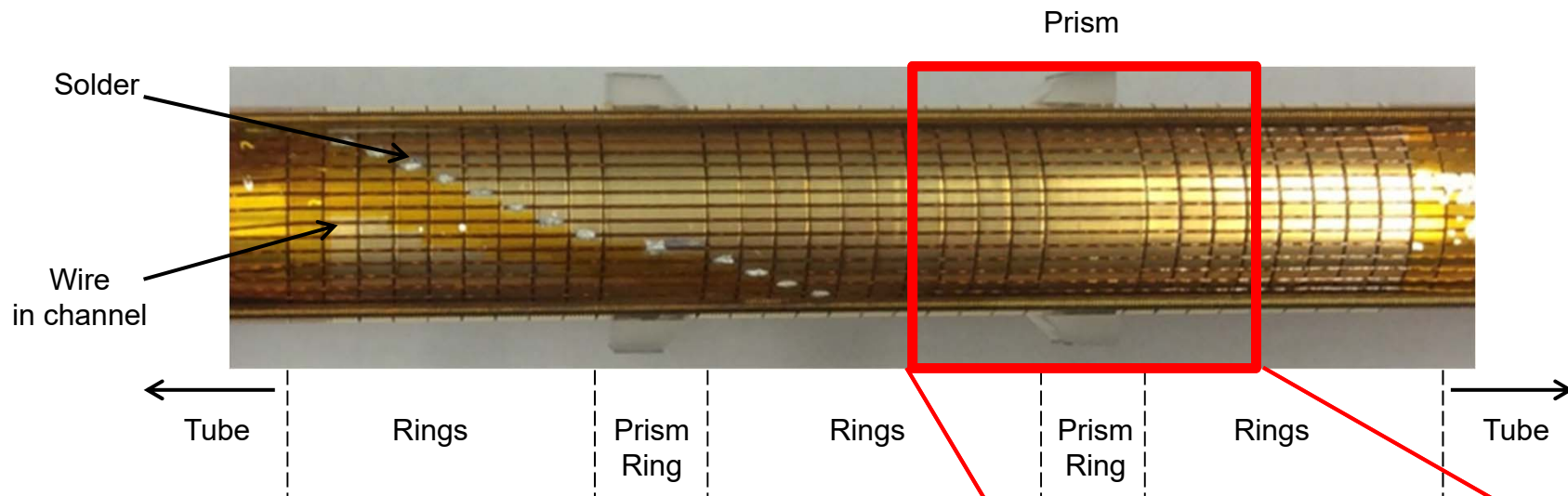
Magnetic field control: results with 52 shim coils



Using molecules for final measurement & shimming:
r.m.s. variation $\delta B/B < 20$ ppb [6 cm L. x 1 cm D. cylinder]

\mathcal{E} -field control

Ring electrodes create sine wave \mathcal{E} -field along \mathbf{z} -axis:



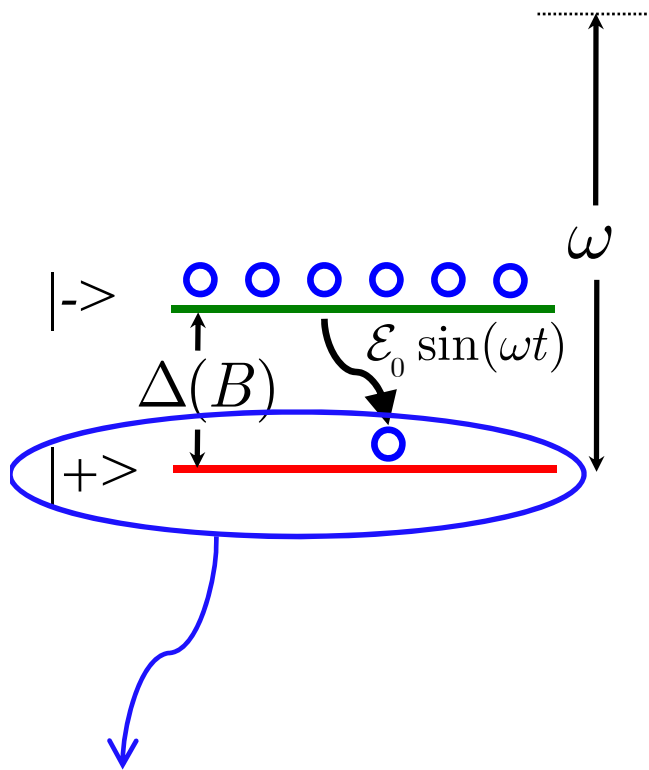
Detecting PV in near-degenerate levels: AC Stark shift

PV mixing iW encodes physics of interest

$$H = \begin{pmatrix} 0 & iW + d\mathcal{E}(t) \\ -iW + d\mathcal{E}(t) & \Delta \end{pmatrix}$$

Apply oscillating \mathcal{E} -field, 1 cycle:

$$\mathcal{E}(t) = \mathcal{E}_0 \sin(\omega t) \quad \left[\begin{array}{l} \omega \gg \Delta, d\mathcal{E}_0; \\ T = 2\pi / \omega \end{array} \right]$$



$$S = \left| \langle + | \psi(T) \rangle \right|^2 = 4 \sin^2 \left(\frac{\Delta T}{2} \right) \left[\left(\frac{d\mathcal{E}_0}{\omega} \right)^2 + 2 \frac{W}{\Delta} \frac{d\mathcal{E}_0}{\omega} \right]$$

D.D., S.B. Cahn, *et al.*
PRL **100**, 023003 (2008)

“Large”
Stark Term
Even in \mathcal{E}_0

Small
PV Term
Odd in \mathcal{E}_0

Signal, Asymmetry, Sensitivity

- - Measure signal $S(\mathcal{E}_0) \approx 4N_0 \sin^2\left(\frac{\Delta T}{2}\right) \left[\left(\frac{d\mathcal{E}_0}{\omega}\right)^2 + 2\frac{W}{\Delta} \frac{d\mathcal{E}_0}{\omega} \right]$

with opposite-sign \mathcal{E} -fields $+\mathcal{E}_0, -\mathcal{E}_0$

- - Form asymmetry to extract W in terms of known quantities :

$$\mathcal{A} = \frac{S(+\mathcal{E}_0) - S(-\mathcal{E}_0)}{S(+\mathcal{E}_0) + S(-\mathcal{E}_0)} \approx 2 \frac{W}{\Delta} \frac{\omega}{d\mathcal{E}_0}$$

Dispersion-like
function
of detuning Δ

Statistical Uncertainty

$$\delta W = \frac{1}{2\sqrt{2}} \frac{1}{\sqrt{N_0}} \frac{1}{T}$$

best sensitivity from
large interaction time T

Properties of NSD-PV asymmetry: example ^{137}BaF

Typical numbers for ^{137}BaF :

$$\Delta_0 \sim 1/T \sim 2\pi \times 1 \text{ kHz}$$

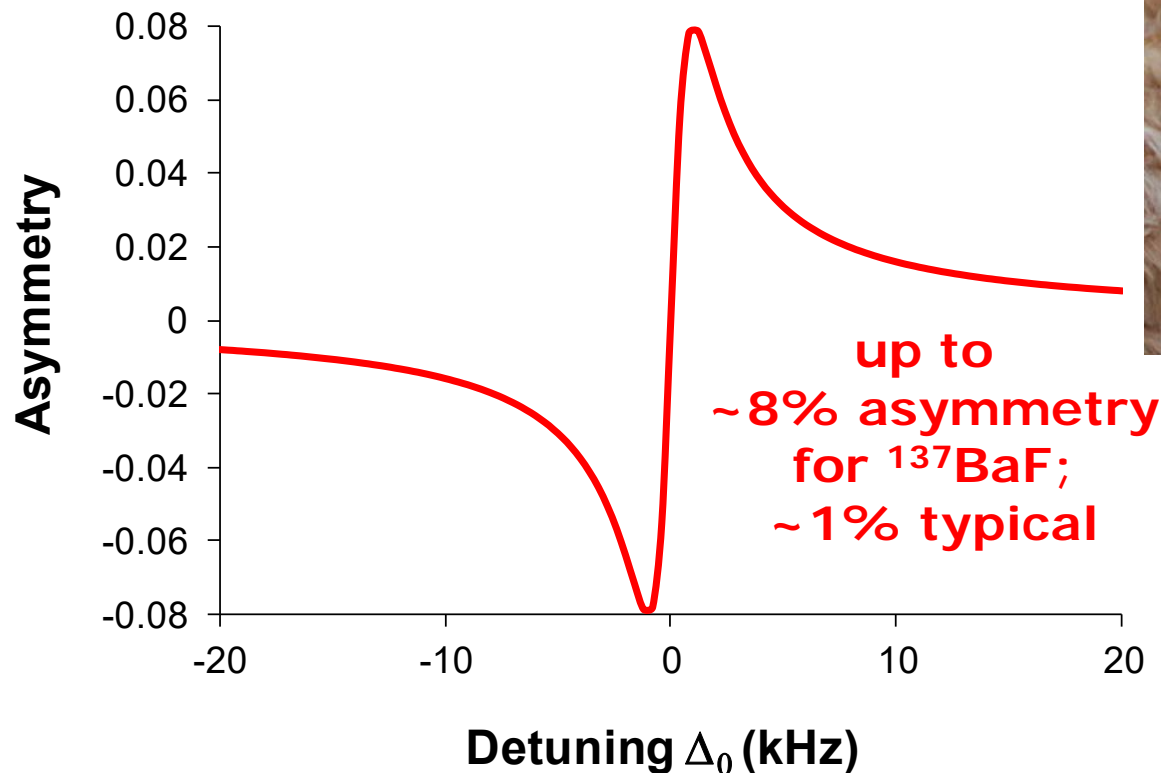
$$\omega = 2\pi \times 100 \text{ kHz}$$

$$dE_0/\omega = 0.1$$

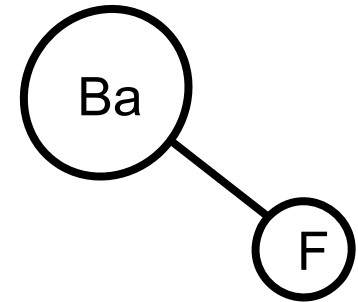
$$W = 2\pi \times 5 \text{ Hz}$$

PV Invariant

$$\left(d\vec{\mathcal{E}} / dt \right) \cdot \left(\vec{\mathcal{B}} - \vec{\mathcal{B}}_c \right)$$



NSD-PV with BaF



Initial physics goal: NSD-PV with ^{137}BaF

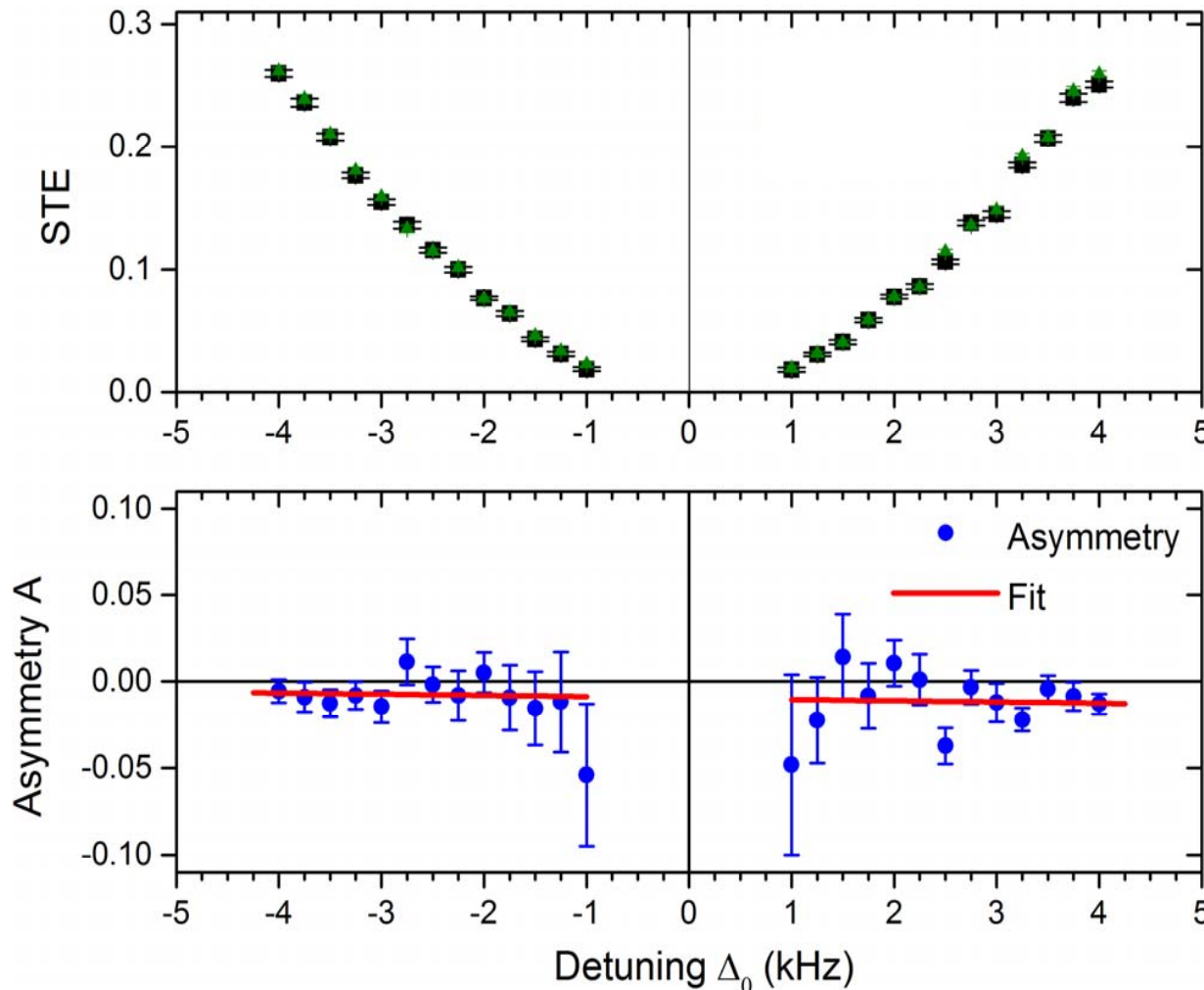
- Odd neutron (^{133}Cs had odd proton)
- Heavy \rightarrow large effect, anapole moment dominates
- Large enough natural abundance – don't need enriched source
- Required lasers = simple, cheap diodes

Proof of principle using $^{138}\text{Ba}^{19}\text{F}$: recently completed

- Larger natural abundance ($\sim 75\%$ vs $\sim 11\%$ for ^{137}Ba)
- Uses same beam source, lasers, magnet, etc. as ^{137}BaF
- $W(^{138}\text{Ba}) = 0$ Hz (no unpaired nucleons = no NSD-PV)
 $W(^{19}\text{F}) \approx 0.002$ Hz ≈ 0 (light, small electron spin density in BaF)
- **Test for systematics with known answer**

NSD-PV data with $^{138}\text{Ba}^{19}\text{F}$

- Measure, cancel, & remeasure \mathcal{B} -field gradients and non-reversing \mathcal{E} -fields to suppress possible systematics
 - Measure NSD-PV signal & asymmetry



$$S \propto \sin^2 \left(\Delta \frac{T}{2} \right) \left(\frac{d\mathcal{E}_0}{\omega} \right)^2$$

$$A = \frac{S(+\mathcal{E}_0) - S(-\mathcal{E}_0)}{S(+\mathcal{E}_0) + S(-\mathcal{E}_0)}$$

Fit to function

$$\mathcal{A}_{\text{fit}}(\Delta) = 2 \frac{W_{\text{fit}}}{\Delta} \frac{\omega}{d\mathcal{E}_0} + a_0 + a_1 \Delta$$

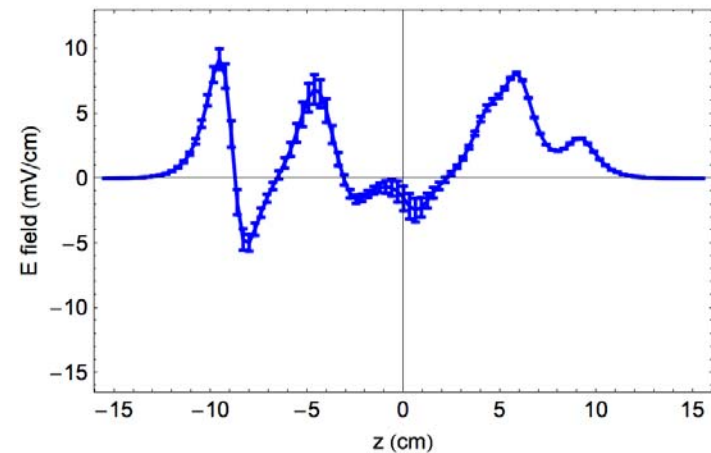
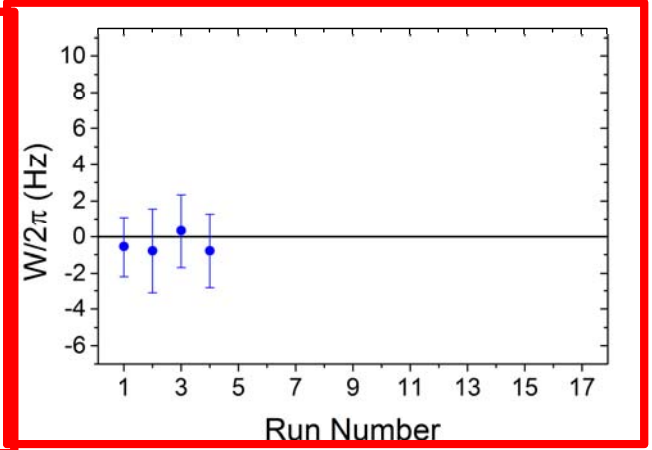
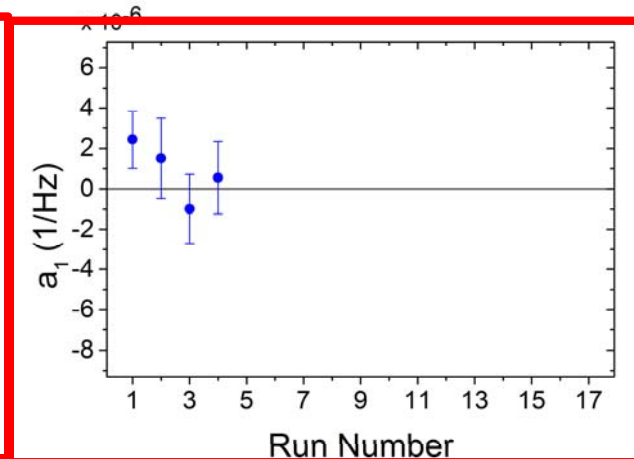
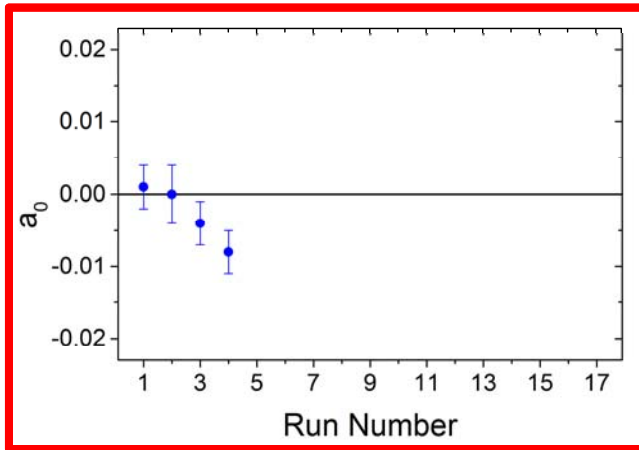
NSD-PV data with $^{138}\text{Ba}^{19}\text{F}$

Fit to $\mathcal{A}_{\text{fit}}(\Delta) = 2 \frac{W_{\text{fit}}}{\Delta} \frac{\omega}{d\mathcal{E}_0} + a_0 + a_1 \Delta$

from stray \mathcal{E} -fields alone

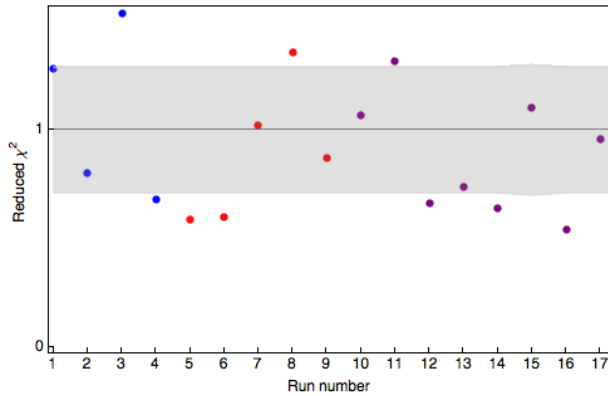
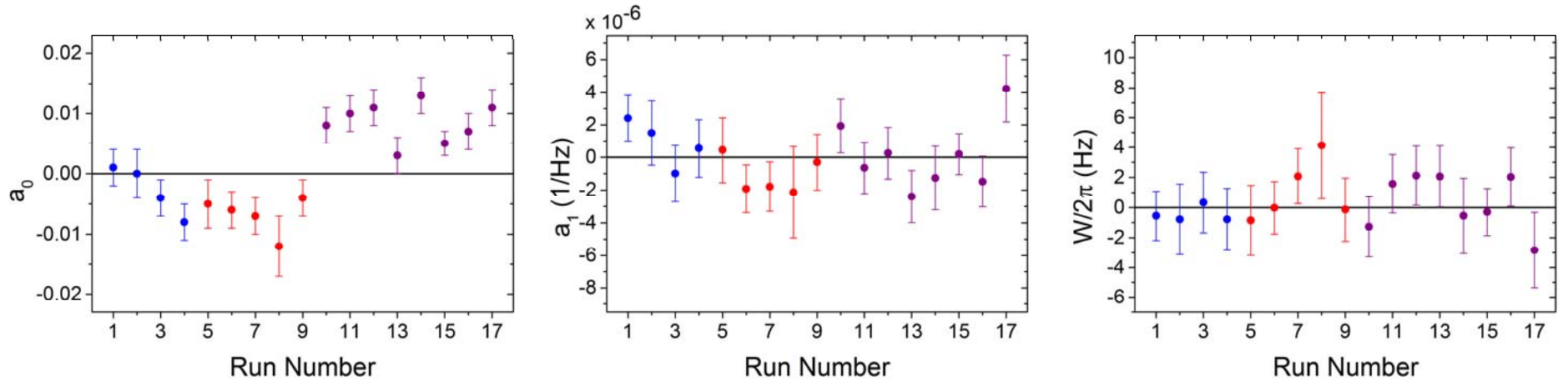
From combined
stray \mathcal{E} -fields & β -field gradients

NSD-PV



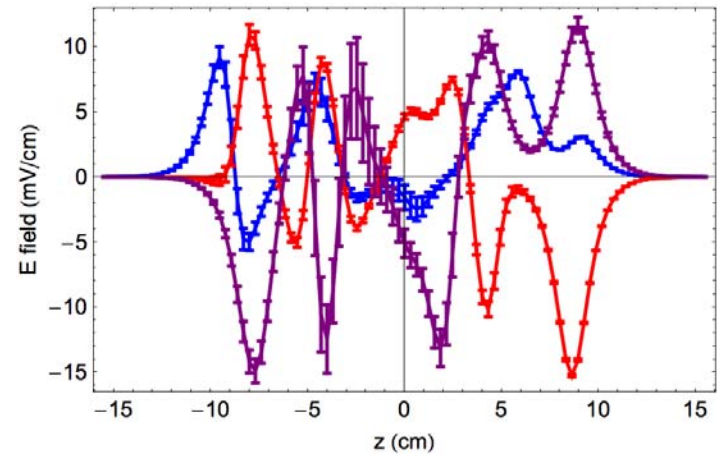
- stray \mathcal{E} -fields always below 15mV/cm

NSD-PV data with $^{138}\text{Ba}^{19}\text{F}$

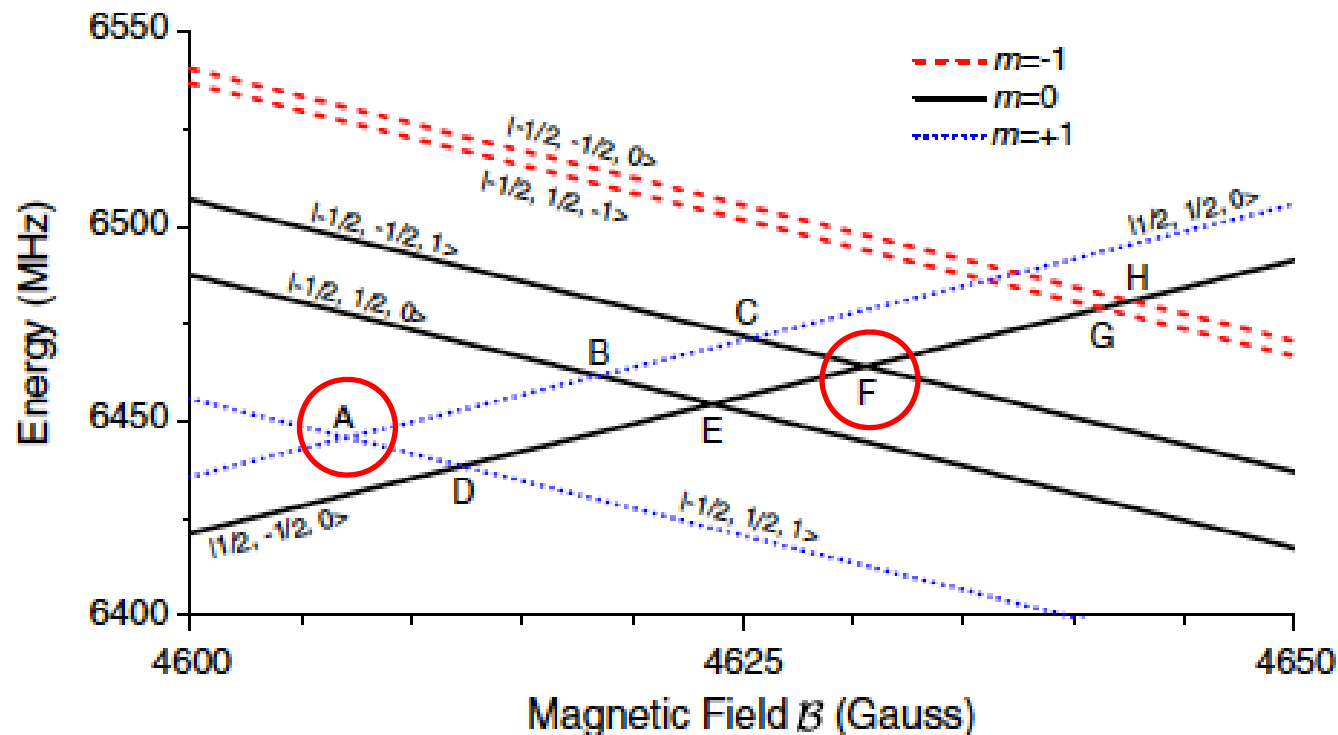


Weighted average $\rightarrow \frac{W}{2\pi} = 0.32 \pm 0.49_{\text{stat}} \text{ Hz}$
 $a_1 = \left(-1.27 \pm 4.02_{\text{stat}} \right) * 10^{-7} \text{ 1/Hz}$

- ^{138}BaF expected $W = 0$
- Measured with 3 different stray \mathcal{E} -fields (all below 15 mV/cm)
- a_1 terms consistent with zero: no systematics



Different level crossings to suppress systematics



$|m_S, m_I, m_N\rangle$

S : electron spin
 I : nuclear spin
 N : rotation
 n : molecular axis

Measured quantity, different for each crossing

$$W = W_P (\kappa'_2 + \kappa'_a) \left\langle \left(\hat{n} \times \vec{S} \right) \cdot \vec{I} \right\rangle$$

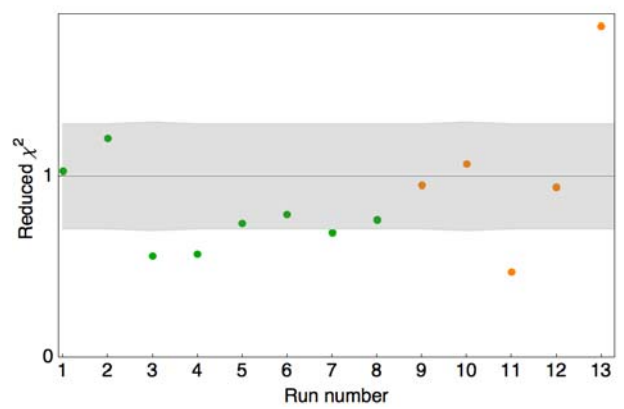
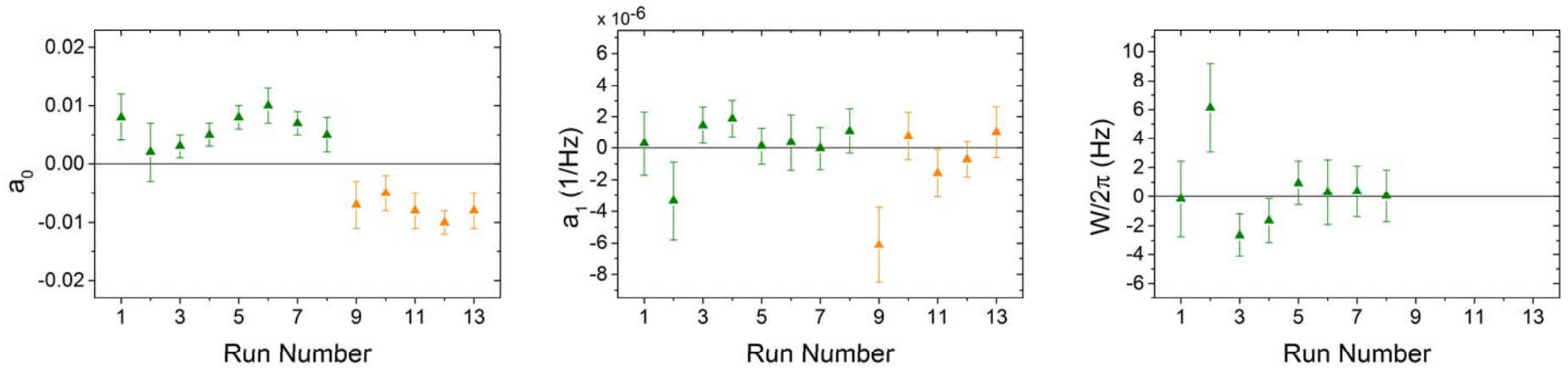
Molecular wavefunctions:
 same at all crossings,
 accurately computed

NSD-PV
 parameters:
 same at all
 crossings

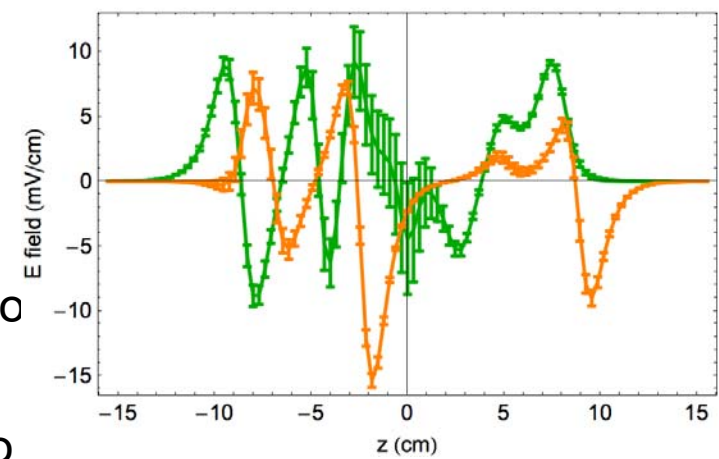
Angular factor:

- Different for each crossing (sign & magnitude)
- Analytically calculable

NSD-PV data with $^{138}\text{Ba}^{19}\text{F}$: 2nd crossing $W \rightarrow -W$



Weighted average $\rightarrow \frac{W}{2\pi} = 0.05 \pm 0.51_{\text{stat}} \text{ Hz}$
 $a_1 = \left(-1.69 \pm 3.98_{\text{stat}} \right) * 10^{-7} \text{ 1/Hz}$



- ^{138}BaF expected $W = 0$
- Measured with 2 different stray E-fields (all below 15mV/cm)
- No systematics $\rightarrow a_1$ terms consistent with zero

Systematic & total uncertainty evaluation

Strategy

- Deliberately exaggerate imperfection by known, large factor
- Measure effect on the NSD-PV matrix element W from coupling to ambient imperfections in the experiment

Parameter	Shift	Systematic δW_{sys} (Hz)	Uncertainty
Bipolar \mathcal{E}_{nr} Pulses		0.12	
Unipolar \mathcal{E}_{nr} Pulses		0.16	
\mathcal{B} -Field Inhomogeneities		0.24	
$\delta\nu_{L2}$ and \mathcal{E}_{nr} at and near Gap 22	-0.04	0.21	
Total Systematic	-0.04	0.38	

Final Error Budget with $^{138}\text{Ba}^{19}\text{F}$

Crossing	$W/(2\pi)$ (Hz)	C	d (Hz/(V/cm))	$W_{\text{mol}} = \kappa' W_P/(2\pi)$ (Hz)
A	$0.28 \pm 0.49_{\text{stat}} \pm 0.38_{\text{sys}}$	-0.41	3360	$-0.68 \pm 1.20_{\text{stat}} \pm 0.93_{\text{sys}}$
F	$0.01 \pm 0.51_{\text{stat}} \pm 0.38_{\text{sys}}$	+0.39	3530	$0.03 \pm 1.30_{\text{stat}} \pm 0.97_{\text{sys}}$
Weighted Average	-	-	-	$-0.36 \pm 0.88_{\text{stat}} \pm 0.95_{\text{sys}}$

~170 h data
~ 6×10^7 molecules total

$$W_{\text{mol}} = 2\pi \times (-0.36 \pm 1.29) \text{ Hz}$$

What does this $^{138}\text{Ba}^{19}\text{F}$ result mean?

$$W_{mol} \equiv (\kappa'_2 + \kappa'_a) W_P = 2\pi \times (-0.36 \pm 1.29) \text{ Hz}$$

\Rightarrow Limit on ^{19}F anapole + C_{2P} :

$$W_P(^{19}\text{F in BaF}) = 2\pi \times 0.05 \text{ Hz}$$

$$\kappa'(^{19}\text{F}) = -7 \pm 25 \quad \text{vs. } \kappa'(^{19}\text{F})[\text{shell model}] = 0.08$$

\Rightarrow Proves we have no unknown systematics

\Rightarrow Systematics limited by statistical power (so far)

but not really informative about HPV

....So What?

What does the $E^{138}\text{Ba}^{19}\text{F}$ result mean?

$$W_{mol} \equiv (\kappa'_2 + \kappa'_a)W_P = 2\pi \times (-0.36 \pm 1.29) \text{ Hz}$$

More useful comparison:

$$W_P \left({}^{137}\text{Ba in BaF} \right) = 2\pi \times 160 \text{ Hz}$$

Same experimental uncertainty in ${}^{137}\text{BaF}$ would mean

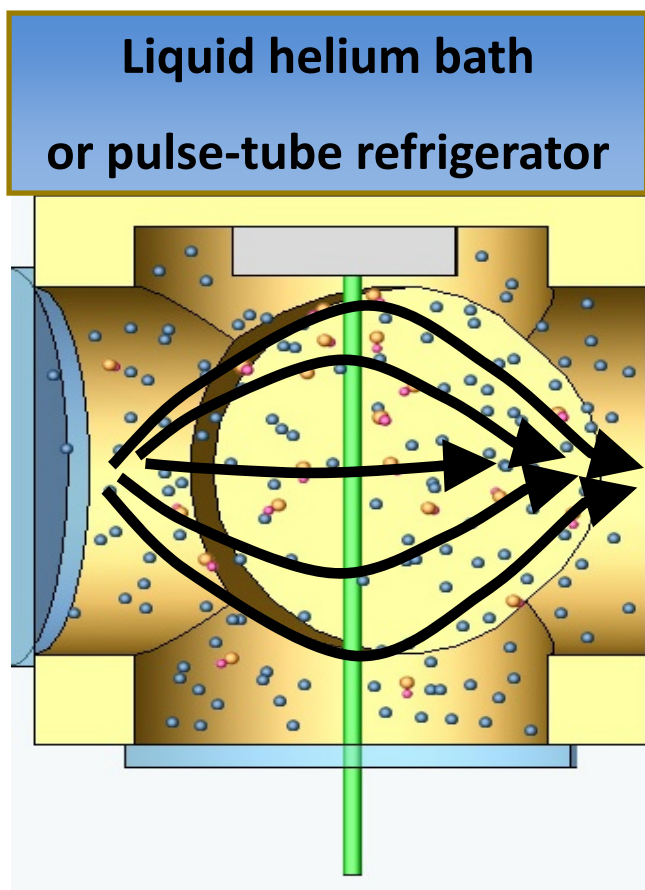
$$\underbrace{\delta\kappa' \left({}^{137}\text{Ba} \right) = 0.008 \quad \text{vs.} \quad \kappa' \left({}^{137}\text{Ba} \right) [\text{shell model}] \approx 0.07}_{\sim 10\% \text{ of predicted value}}$$

Compares favorably to JILA ${}^{133}\text{Cs}$ result: $\kappa' \left({}^{133}\text{Cs} \right) = 0.39 \pm 0.06$
C.S. Wood *et al.*, Science **275**, 1759 (1997)

- Unprecedented sensitivity to NSD-PV
- General technique enables measurements in broad range of nuclei

Cryogenic Buffer Gas-cooled Beam

[Maxwell *et al.* PRL 2005; Patterson & Doyle J Chem Phys 2007; Barry *et al.* PCCP 2011; Hutzler *et al.* PCCP 2011]



- Inject hot molecules (e.g. via laser ablation)
- Cool w/cryogenic buffer gas @ high density
- Efficient extraction to beam via “wind” in cell: $10^{-4} \rightarrow 10\%-40\%$
- “Self-collimated” by extraction dynamics
- Rotational cooling in expansion: $T \sim 4$ K
- Moderately slow: $v \sim 200$ m/s

Beam brightness $\sim 10^3 \times$ larger flux; interaction time $\sim 3 \times$ larger; enables magnetic focusing $\Rightarrow \sim 20 \times$ flux

Gain in NSD-PV statistical sensitivity: $400 \times$ (@ same total time)

Viabile nuclei for anapole/NSD-PV measurement

- 10% measurement possible with demonstrated sensitivity, $\lesssim 1$ h data
- Requires systematics ~2-10x better
- Statistics likely OK, will require systematics ~100x better

1 H																	1 H	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne	
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar	
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr	
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe	
55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn	
87 Fr	88 Ra	89 Ac	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110	111	112	114	116		118			

58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr

Other ongoing anapole-sensitive experiments

Mainz: ^{171}Yb , ^{173}Yb (similar to JILA Cs experiment)

FrPNC @ TRIUMF: $^{\text{xxx}}\text{Fr}$ (cooled & trapped Fr atoms)

Mainz: new ideas using NMR signals for light nuclei....?

ZOMBIES NSD-PV: Outlook & questions

- New lab under renovation, occupancy this summer with cryogenic beam source
- Realistic goal: ~2 years to ^{137}Ba measurement

Question for theorists: what to do next with this method?

- lightest nuclei (accessible via no-core shell model)...?
- could C_2 values be extracted reliably from light nuclei with existing HPV data & understanding?
- is consistency check among many heavier nuclei useful?
 - isotope comparisons?
 - any special cases of particular interest?
- quantitative uncertainties on shell model calculations!