

The Effect of Small-Scale Structure on Reionization

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September 16, 2004

The Epoch of Reionization

- GP troughs detected in spectra of SDSS quasars at $z > 6 \implies$ IGM H I density high enough to suggest reionization only just ended at $z = 6$.

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The Epoch of Reionization

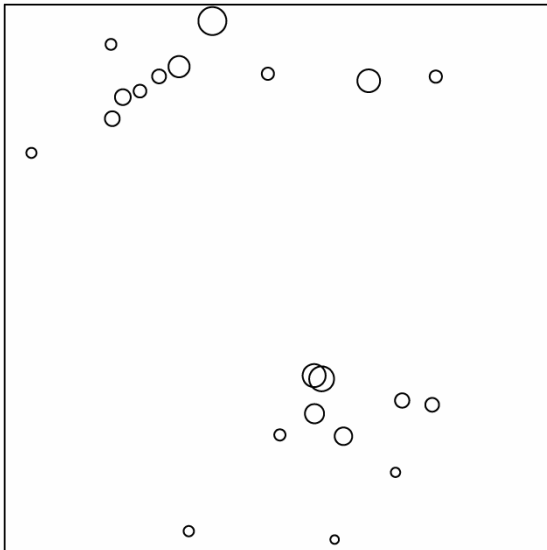
- GP troughs detected in spectra of SDSS quasars at $z > 6 \implies$ IGM H I density high enough to suggest reionization only just ended at $z = 6$.
- WMAP detection of CMB polarization fluctuations on large angular scale \implies foreground electron scattering optical depth high enough to suggest IGM mostly ionized by $z > 12$.
- Plausible explanation: reionization began by $z > 15$ but was extended in time, with final “overlap” of ionized zones at $z = 6$.

Part I: Cosmological Ionization Fronts and the Photoevaporation of Minihalos during Reionization

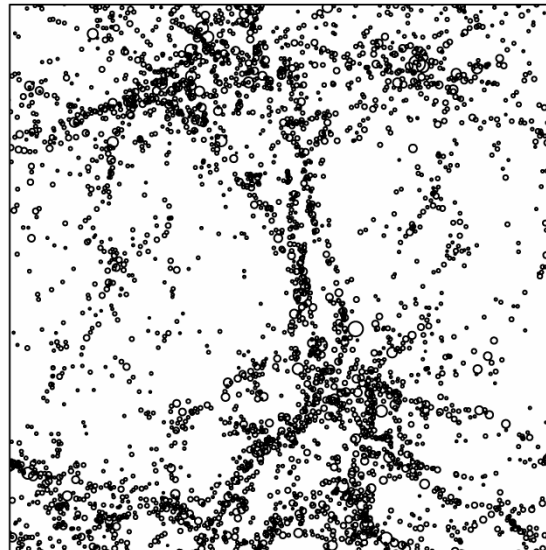
1. DWARF GALAXY MINIHALOS AT HIGH REDSHIFT

- For Λ CDM, the universe at $z > 6$ was already filled with dwarf galaxies capable of trapping a piece of the global, intergalactic I-fronts which reionized the universe and photoevaporating their gaseous baryons back into the IGM.

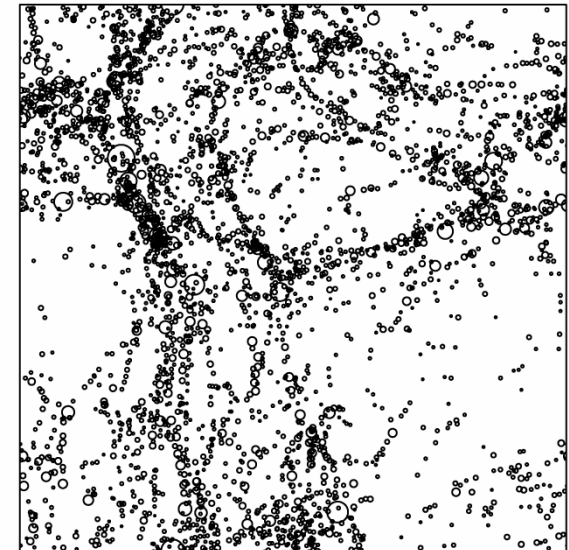
100 kpc box, $T_{\text{vir}} > 10^4 \text{K}$



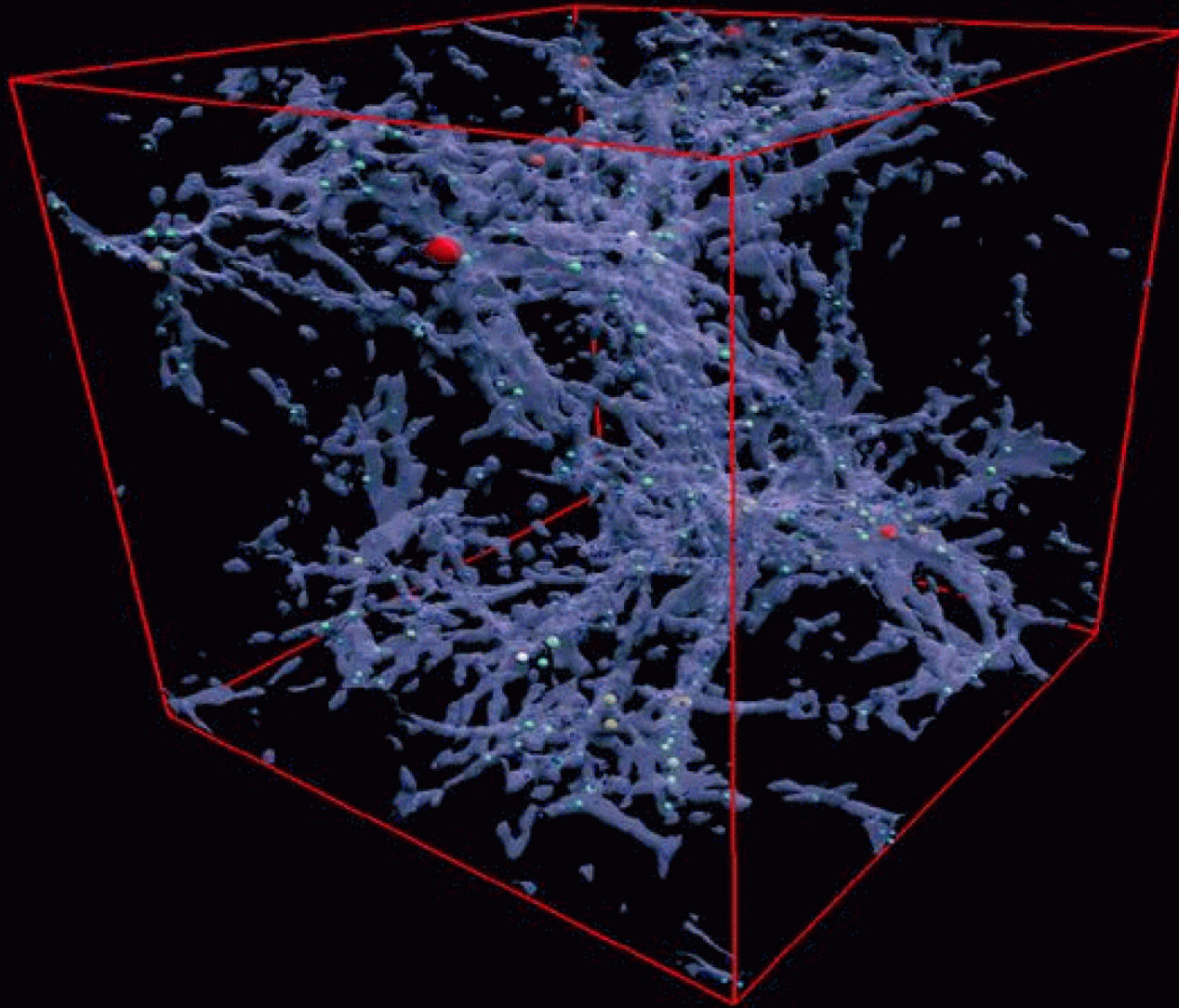
100 kpc box, $T_{\text{vir}} < 10^4 \text{K}$



50 kpc box, $T_{\text{vir}} < 10^4 \text{K}$

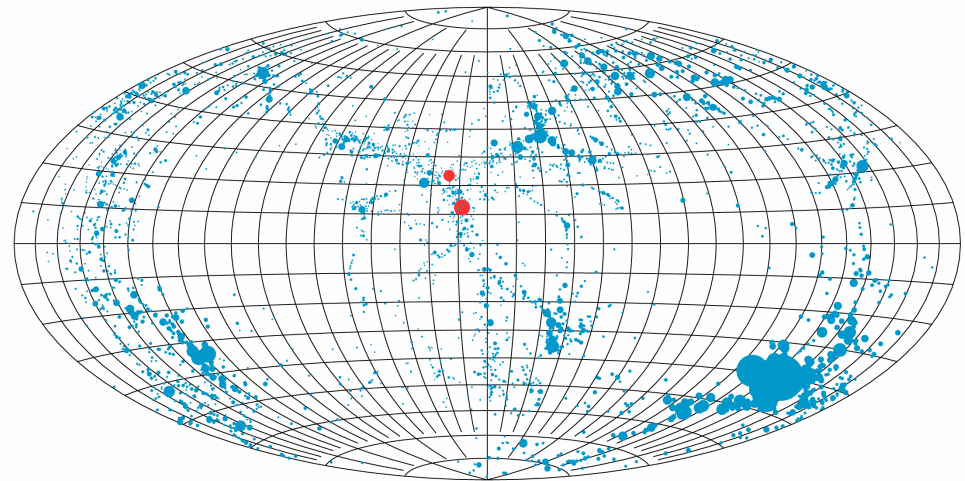


Universe at Redshift $z = 9$

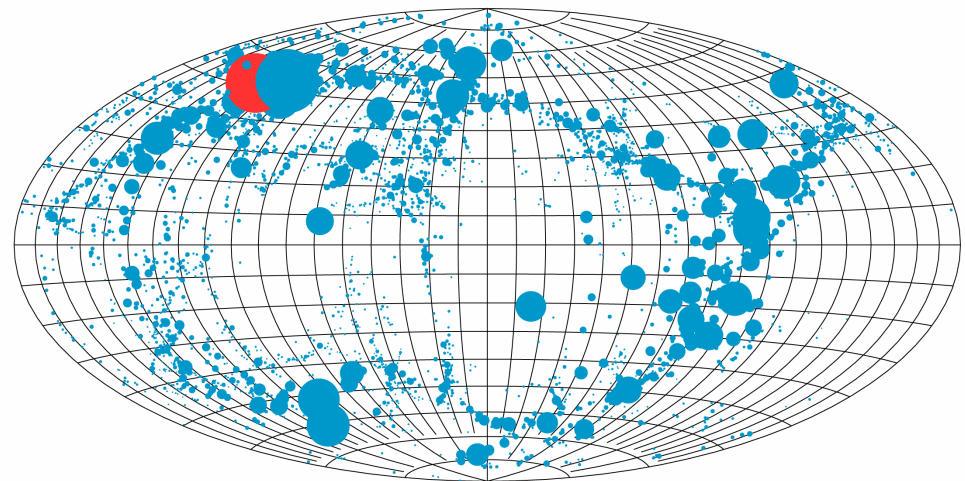


- Minihalos with $T_{\text{vir}} < 10^4$ K were common enough to cover the sky around source halos with $T_{\text{vir}} > 10^4$ K during reionization.

Λ CDM HALOS WITHIN 25 KPC AT $Z = 9$



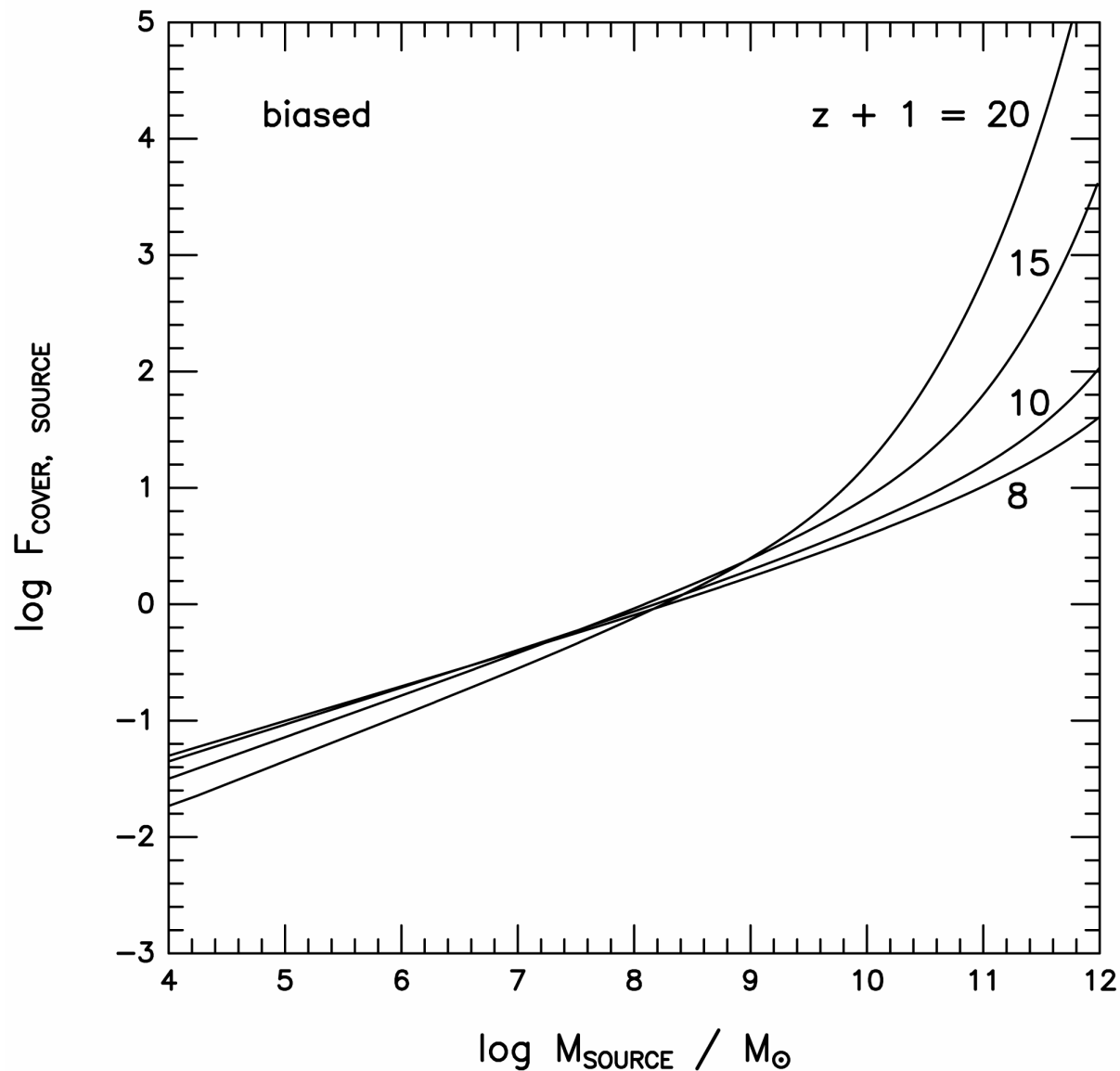
Sky as seen from a random location
Covering fraction : 10.7%



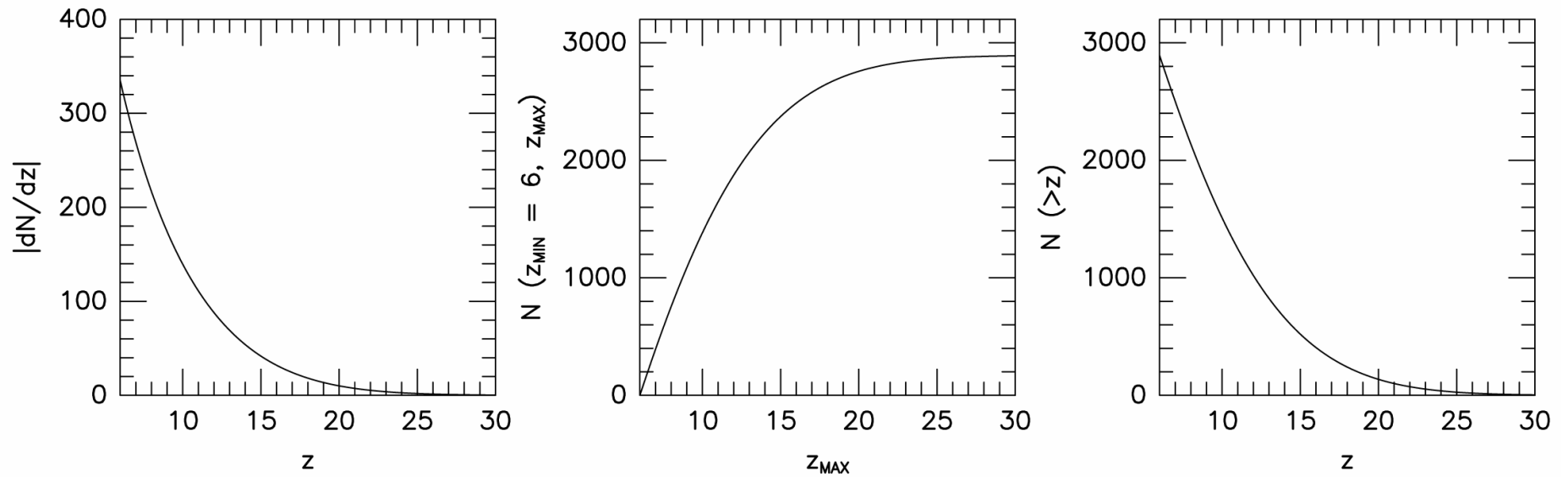
Sky as seen from a $1.1 \times 10^8 M_{\odot}$ source
Covering fraction : 23.6%

Fraction of sky covered by minihalos located within the mean volume per source halo, corrected for bias.

- If $M_{\text{source}} > 10^8 M_{\text{sun}}$, $F_{\text{cover,source}} > 1$.



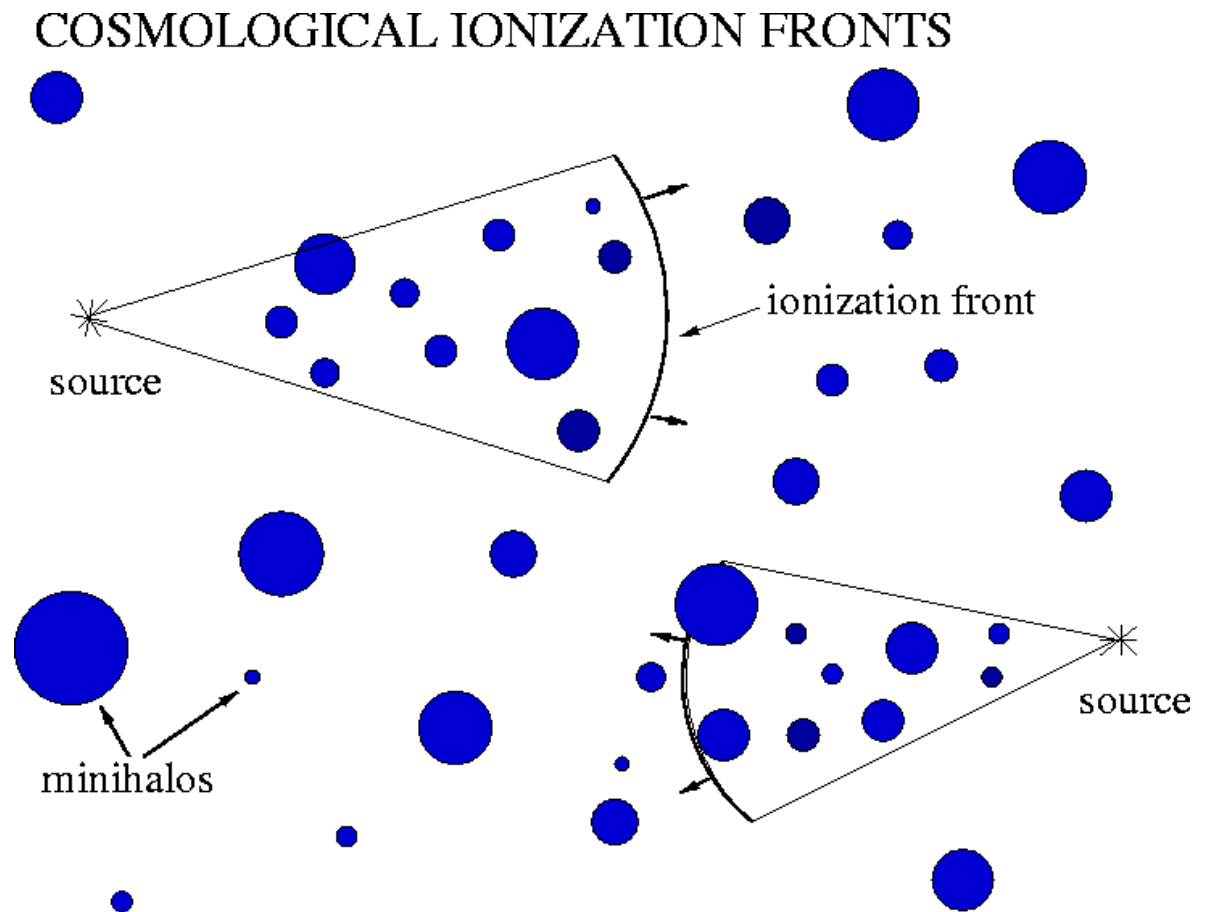
- Photons emitted by any high z source before or during reionization will typically encounter large numbers of photoevaporating minihalos at $z > 6$.



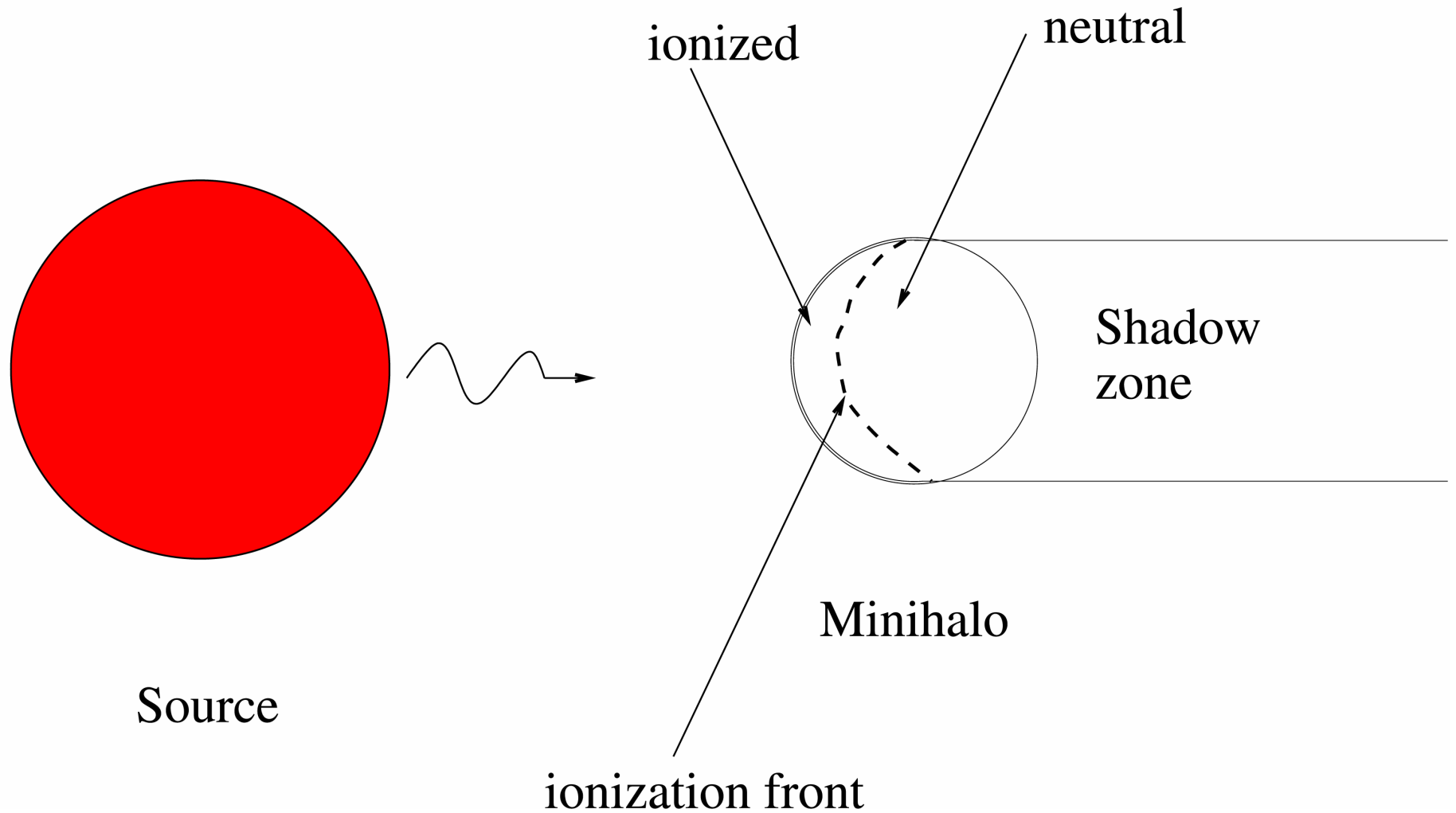
2. THE PHOTOEVAPORATION OF MINIHALOS OVERTAKEN BY COSMOLOGICAL I-FRONTS

(Shapiro, Iliev & Raga 2004)

- We have performed radiation-hydrodynamical simulations of the photoevaporation of a cosmological minihalo overrun by a weak, R-type I-front in the surrounding IGM, created by an external source of radiation.



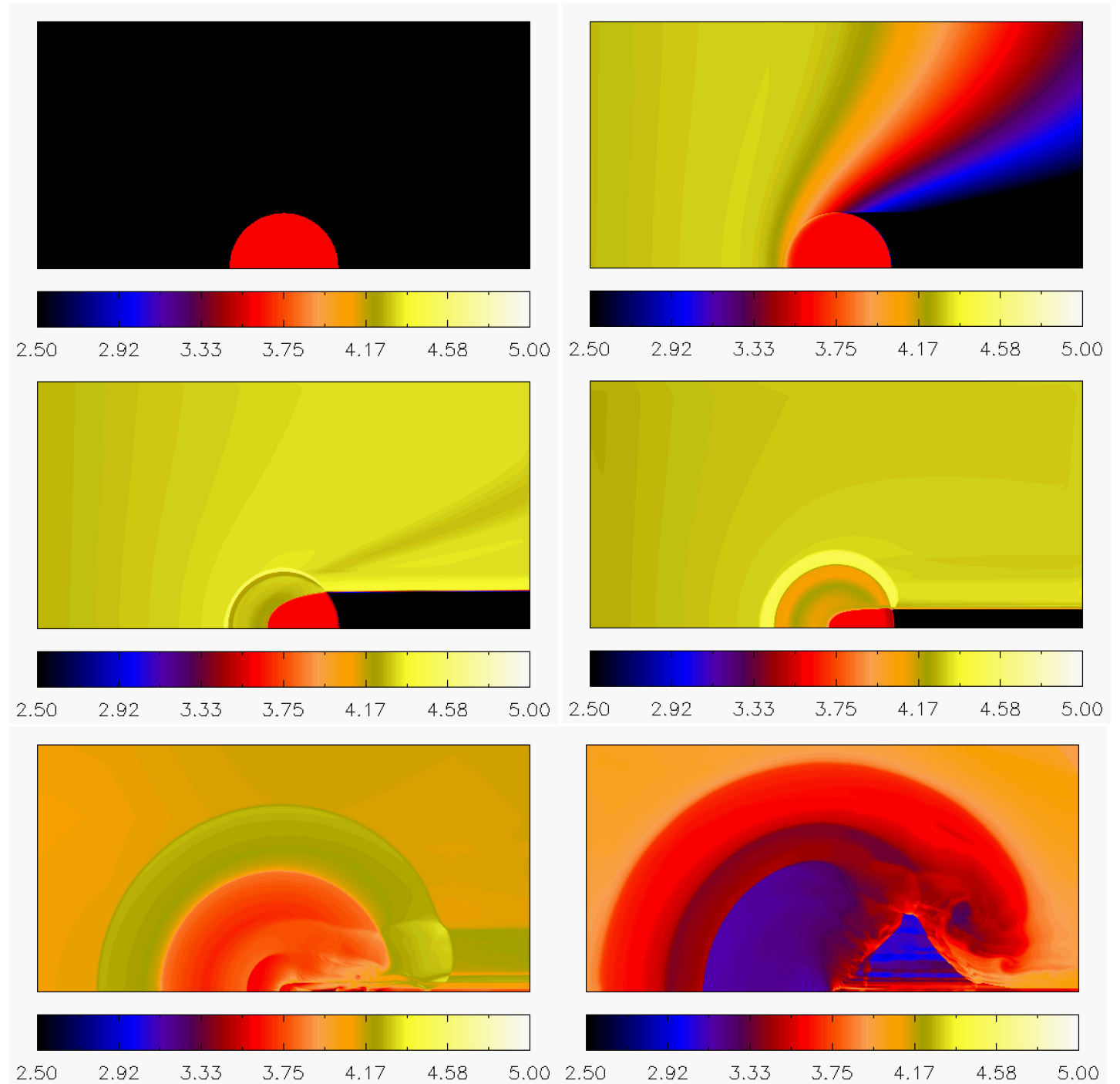
PHOTOEVAPORATION OF MINIHALO DURING COSMOLOGICAL REIONIZATION



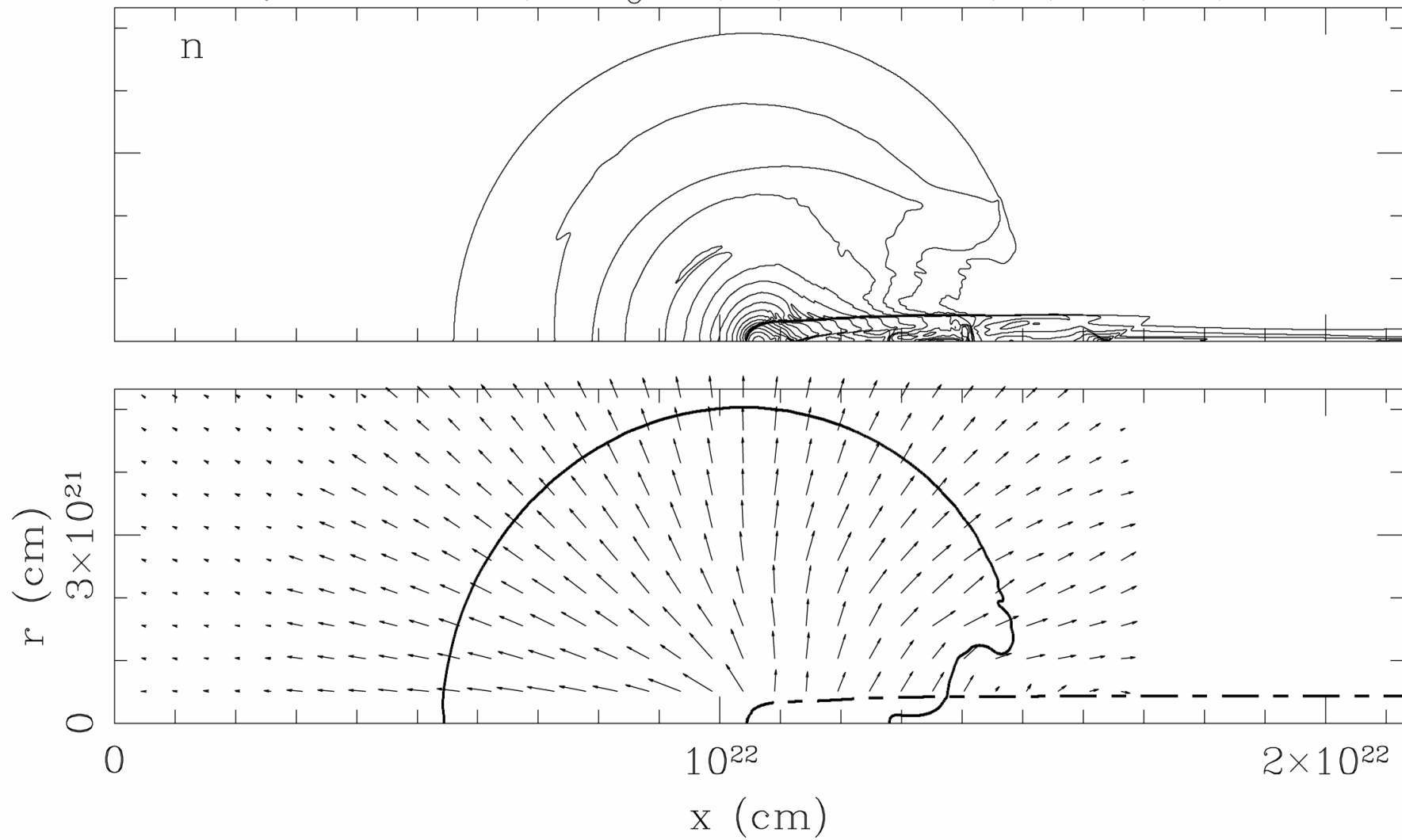
Temperature at
times $t = 0.0, 0.2,$
 $2.5, 10, 60, 150$
Myrs.

$(M_{\text{halo}}, Z_{\text{initial}}, F_0) =$
 $(10^7 M_{\text{sun}}, 9, 1).$

Pop II source.



t=60 Myr, BB 5e4, ($10^7 M_{\odot}, 9, 1$), (0.3, 0.7, 0.7), (25, 3), w/I-front



1 = IGM shock.

2 = contact discontinuity which separates shocked halo wind (between 2 and 3) from swept-up IGM (between 1 and 2).

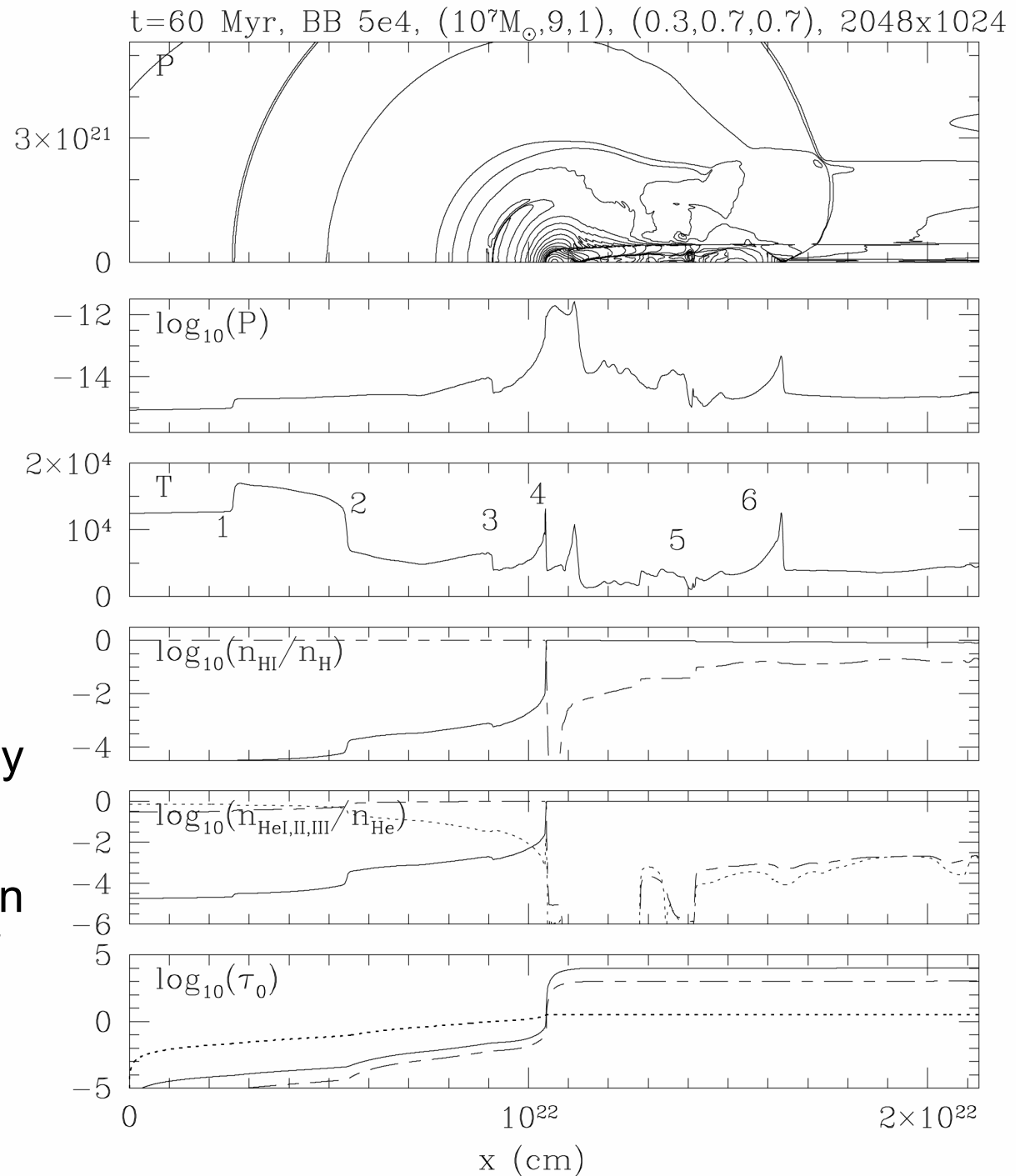
3 = wind shock.

between 3 and 4 = supersonic wind.

4 = I-front.

5 = boundary of gas initially inside minihalo at $z = 9$.

6 = shock in shadow region caused by compression of shadow gas by shock-heated gas outside shadow.



ANIMATIONS

Pop II

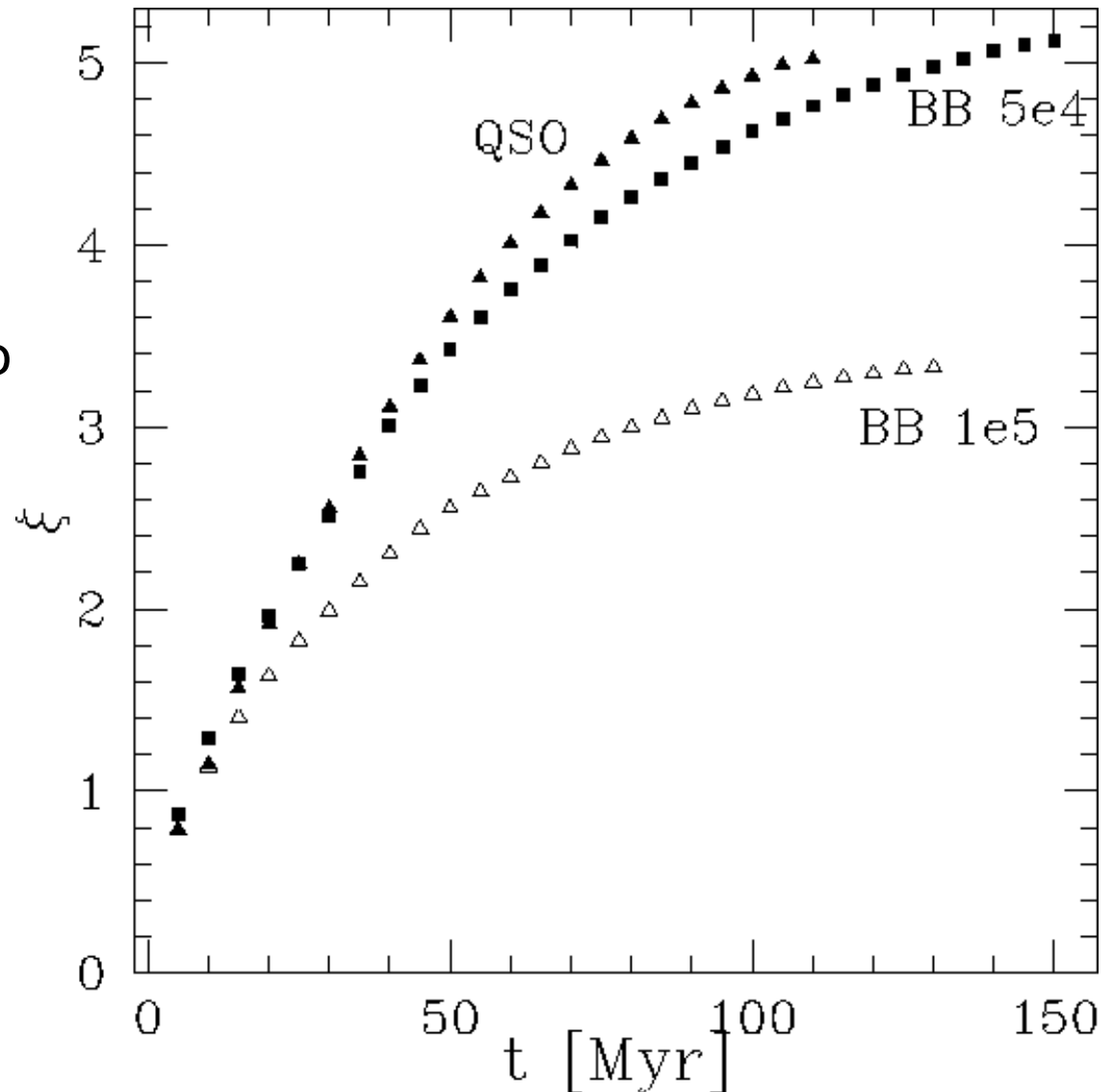
- TEMPERATURE
- DENSITY
- H I
- He II
- C IV

Pop III

- TEMPERATURE
- DENSITY
- H I
- He II
- C IV

Ionizing Photons Consumed Per Minihalo Atom

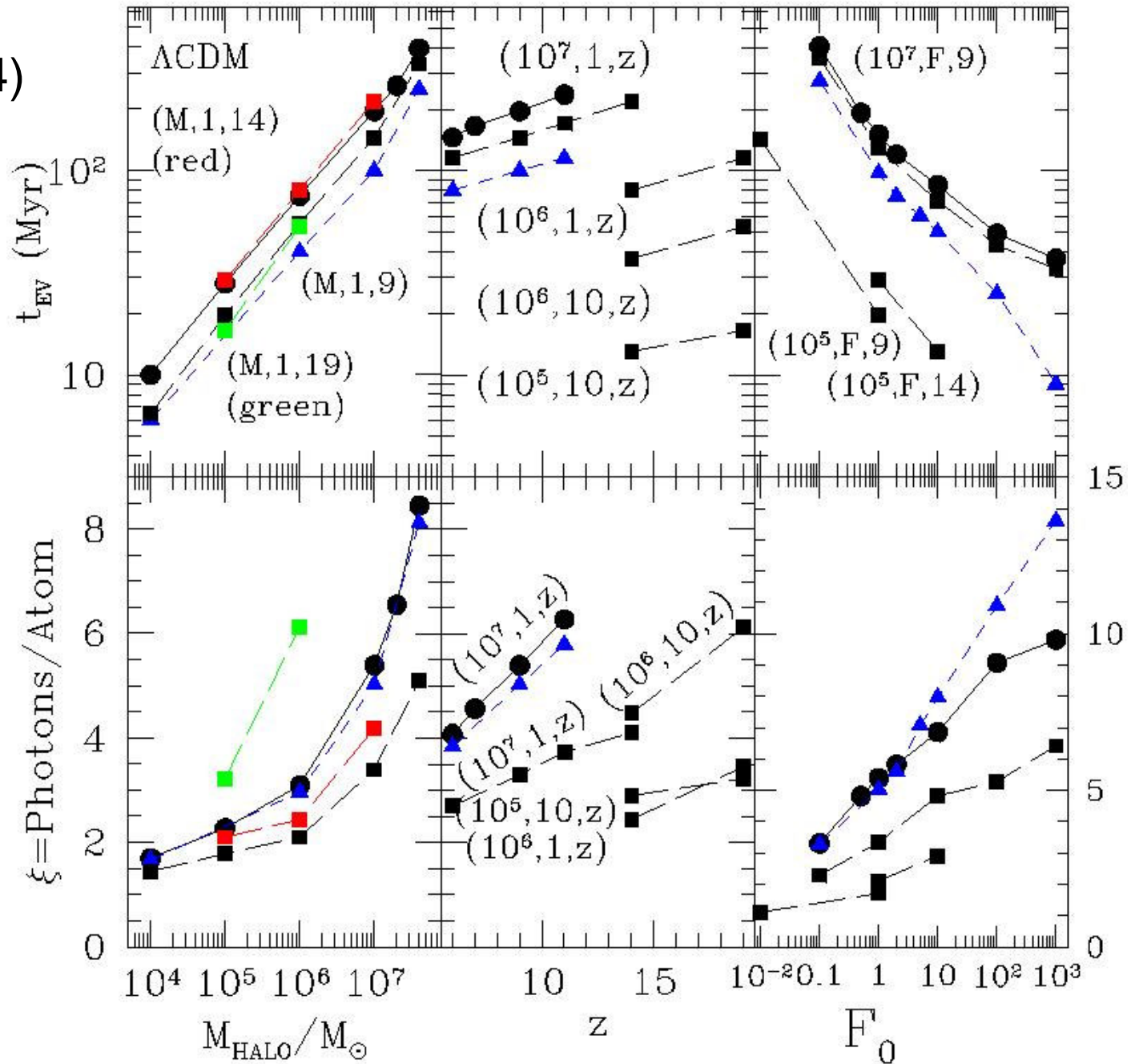
- $\xi(t)$ = # photons consumed per minihalo atom over time t .
- $(M_{\text{halo}}, z_{\text{initial}}, F_0) = (10^7 M_{\text{sun}}, 9, 1)$.



Iliev, Shapiro,
and Raga (2004)

Evaporation
times

Ionizing photon
consumption

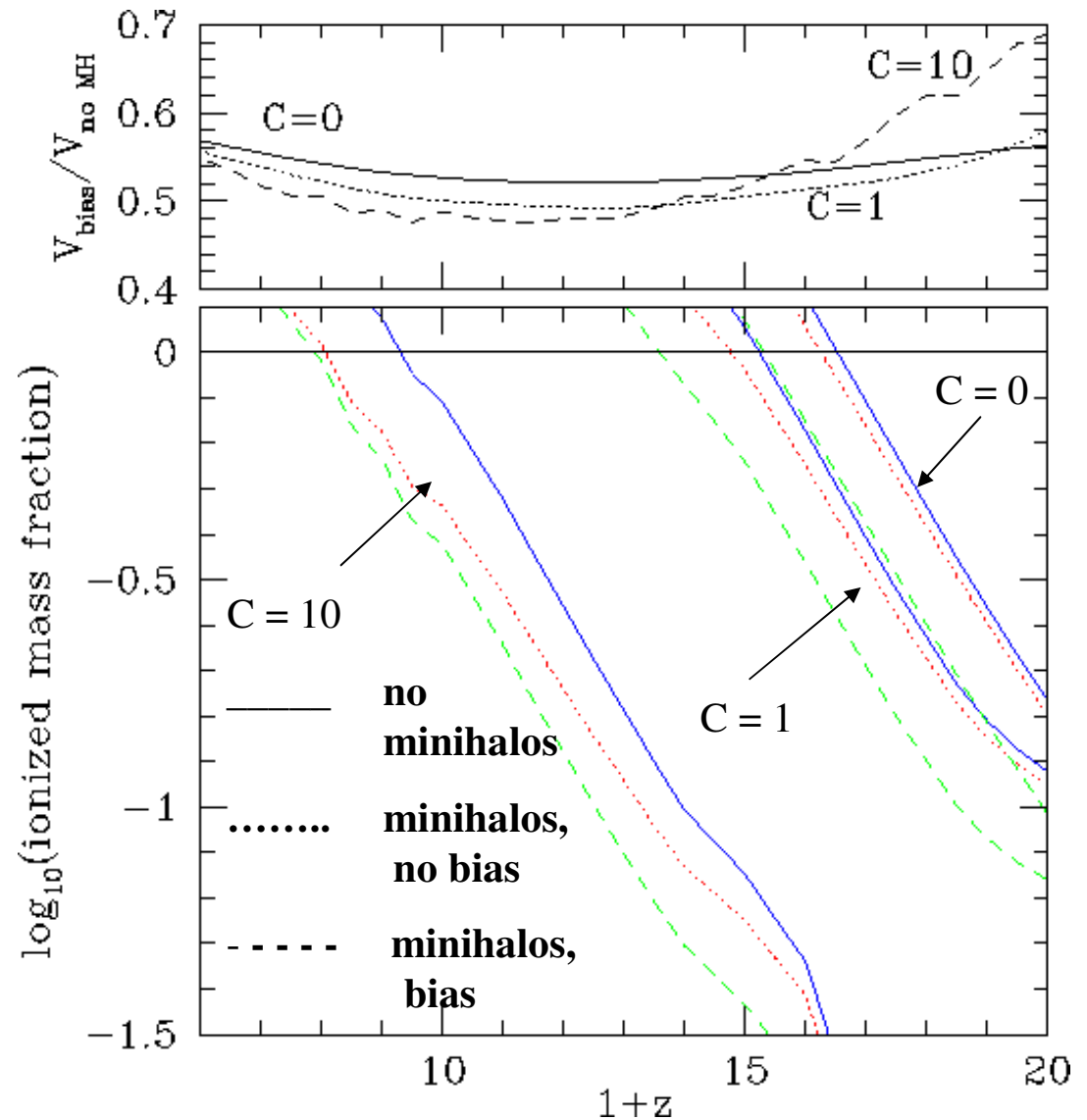


Part II:

**The Effect of Small-Scale
Structure on Reionization**

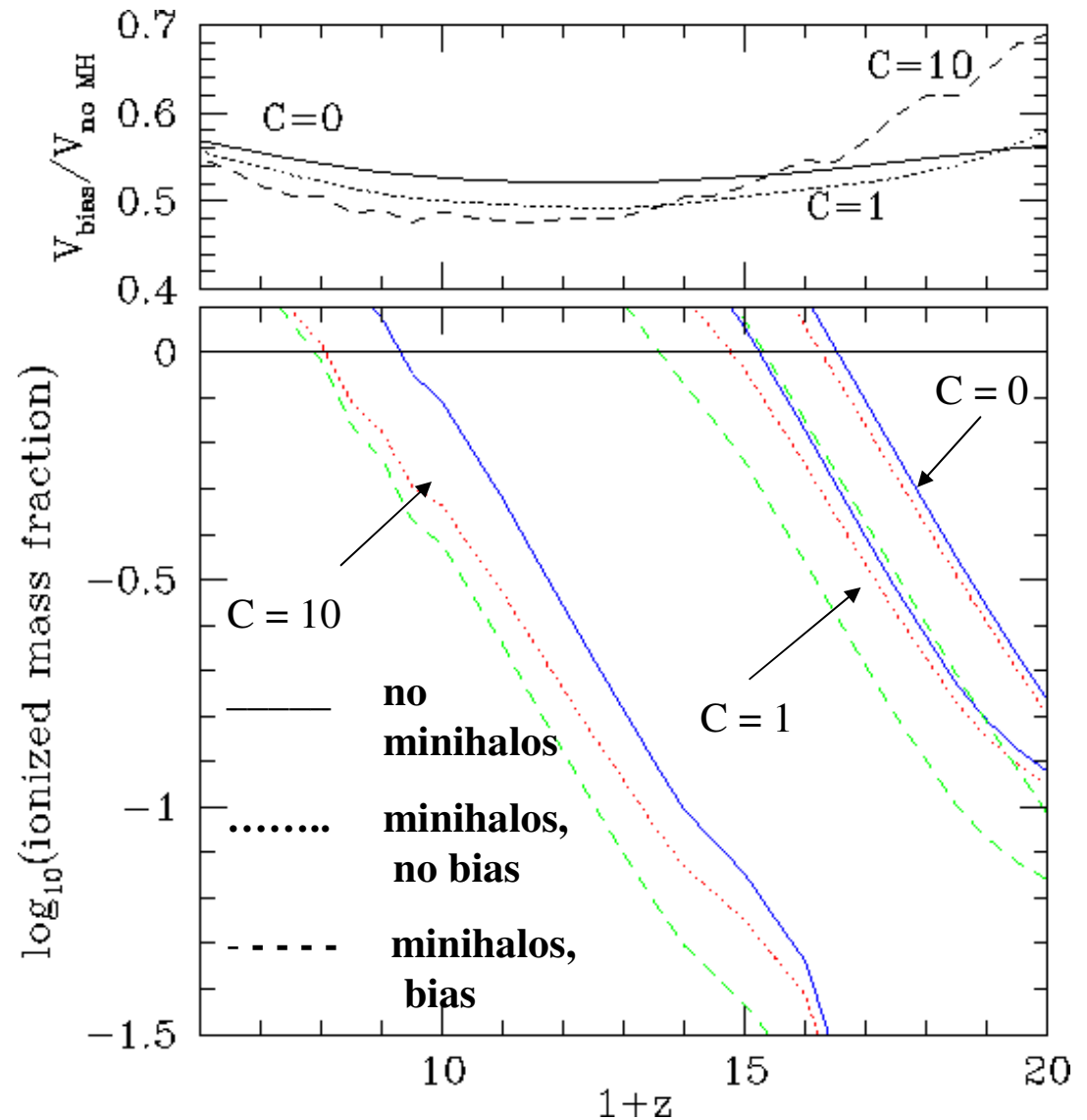
Effect of Minihalo Photon Consumption on Reionization (Iliev, Scannapieco & Shapiro, 2004)

- Let each source halo create its own expanding spherical I-front.
- I-front speed is slowed by minihalo trapping and evaporation.
- Integrate over statistical distribution of source halo masses and turn-on epochs until neighboring H II regions overlap => reionization finished.
- Minihalos increase photon consumption by factor of ~ 2 , delaying reionization.



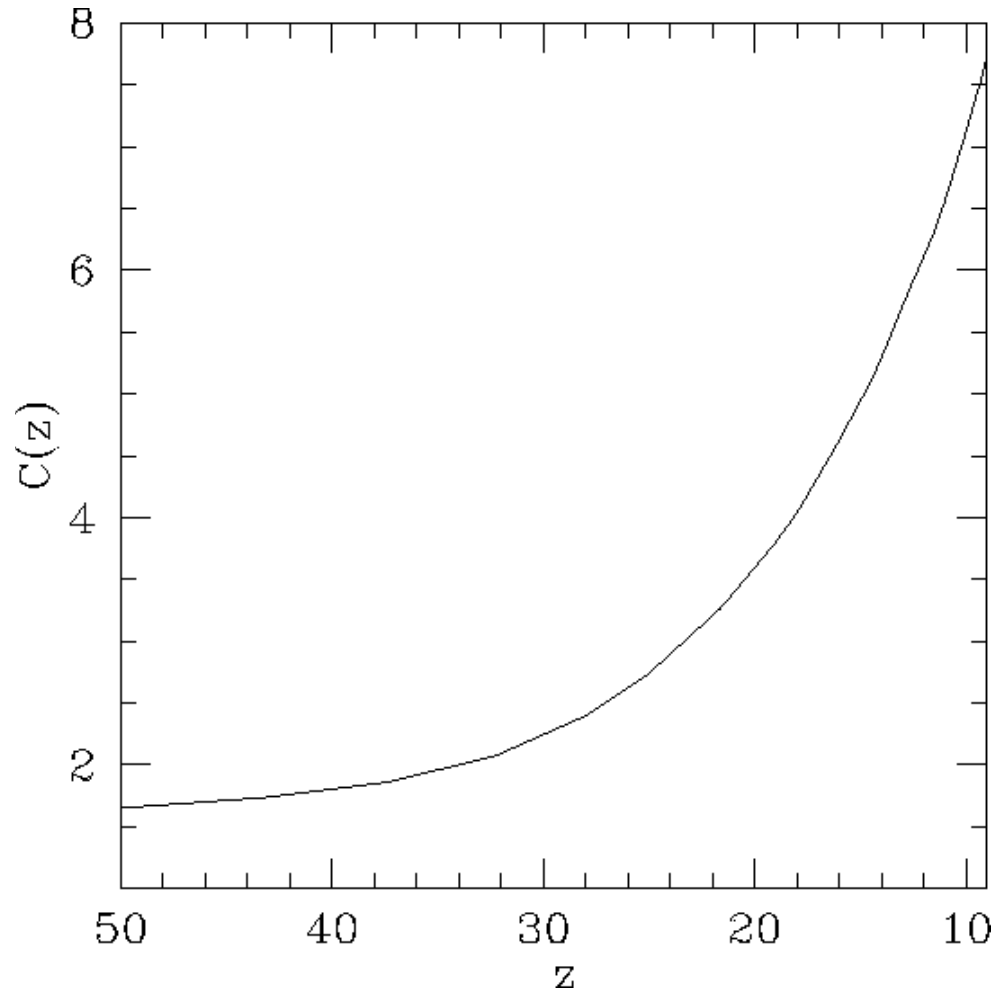
Effect of Minihalo Photon Consumption on Reionization (Iliev, Scannapieco & Shapiro, 2004)

- Delaying reionization slows the rise of electron scattering depth thru IGM.
- Enough delay to satisfy the GP constraint by ending reionization at $z \sim 6 \Rightarrow$ too little τ_{es} to explain WMAP polarization.
- To satisfy both GP and WMAP: Prolonged reionization required \Rightarrow either ionizing radiation release efficiency must drop or gas clumping outside minihalos must rise over time or both.



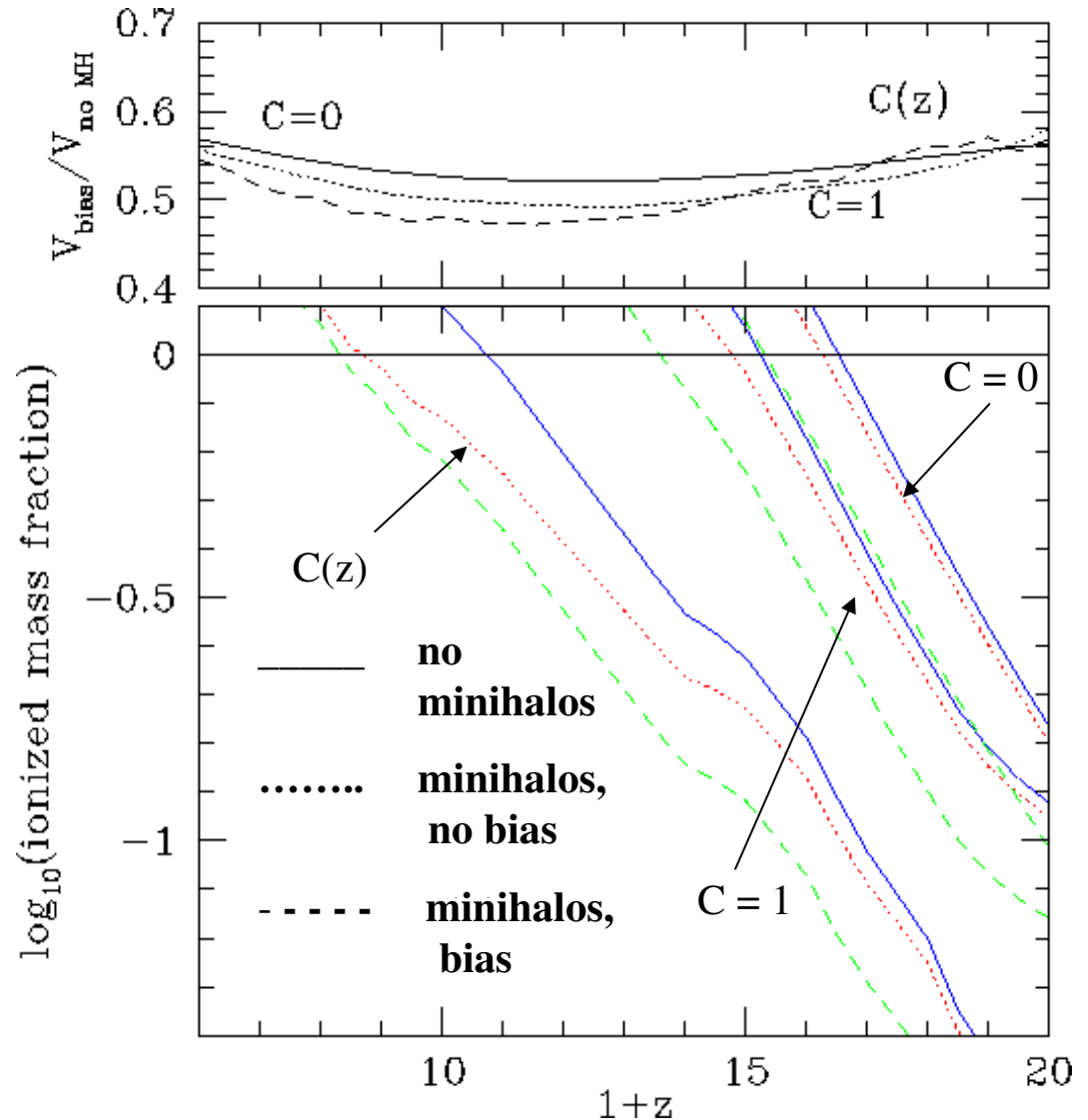
Effect of Minihalo Photon Consumption on Reionization with IGM Clumping $C(z)$

- We use N-body simulations to calculate the IGM clumping factor for gas **outside** halos.
- $C(z)$ increases with time



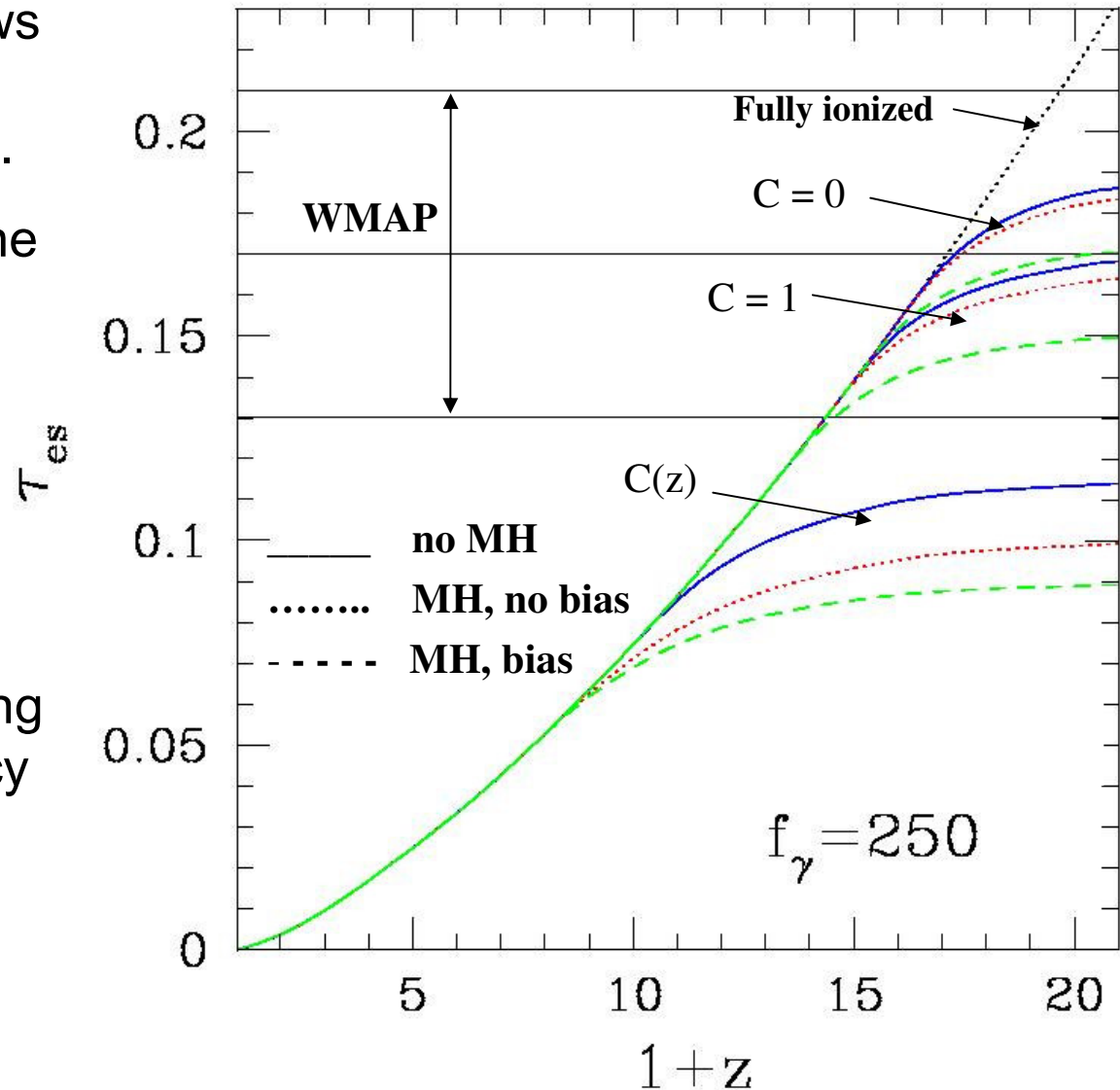
Effect of Minihalo Photon Consumption on Reionization with IGM Clumping $C(z)$

- $C(z)$ increases with time \rightarrow helps to prolong reionization epoch
- Minihalos still increase photon consumption by factor of ~ 2 , delaying reionization ($\delta z \sim 2$).
- e.g. If $z_{\text{ov}} = 9.5$ requires 14 ionizing photons/atom, then minihalos make $z_{\text{ov}} = 7.5$ with 33 photons/atom.
- Does $C(z)$ prolong reionization enough to satisfy both WMAP and GP?



Effect of Minihalos and Time-Varying Clumping of IGM on Reionization

- Delaying reionization slows the rise of electron scattering depth thru IGM.
- Enough delay to satisfy the GP constraint by ending reionization at $z \sim 6 \Rightarrow$ still too little τ_{es} to explain WMAP polarization.
- $C(z)$ helps, but to satisfy both GP and WMAP, probably also need ionizing radiation release efficiency to drop over time as well.



21-cm From Minihalos and the IGM Before Lyman- α Pumping

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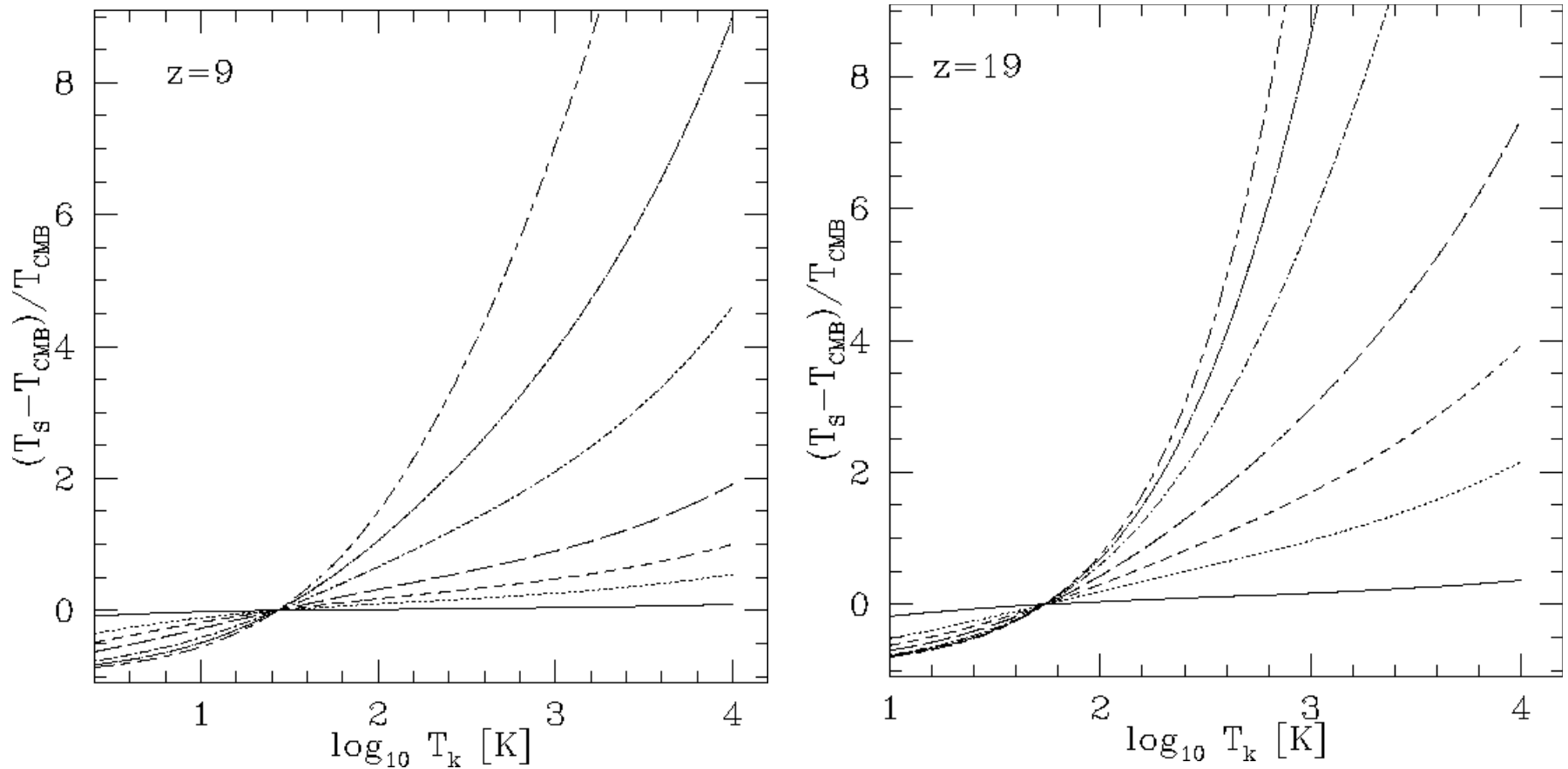
September 16, 2004

Part I:

**On the Detectability of the
Cosmic Dark Ages: 21-cm
Lines from Minihalos**

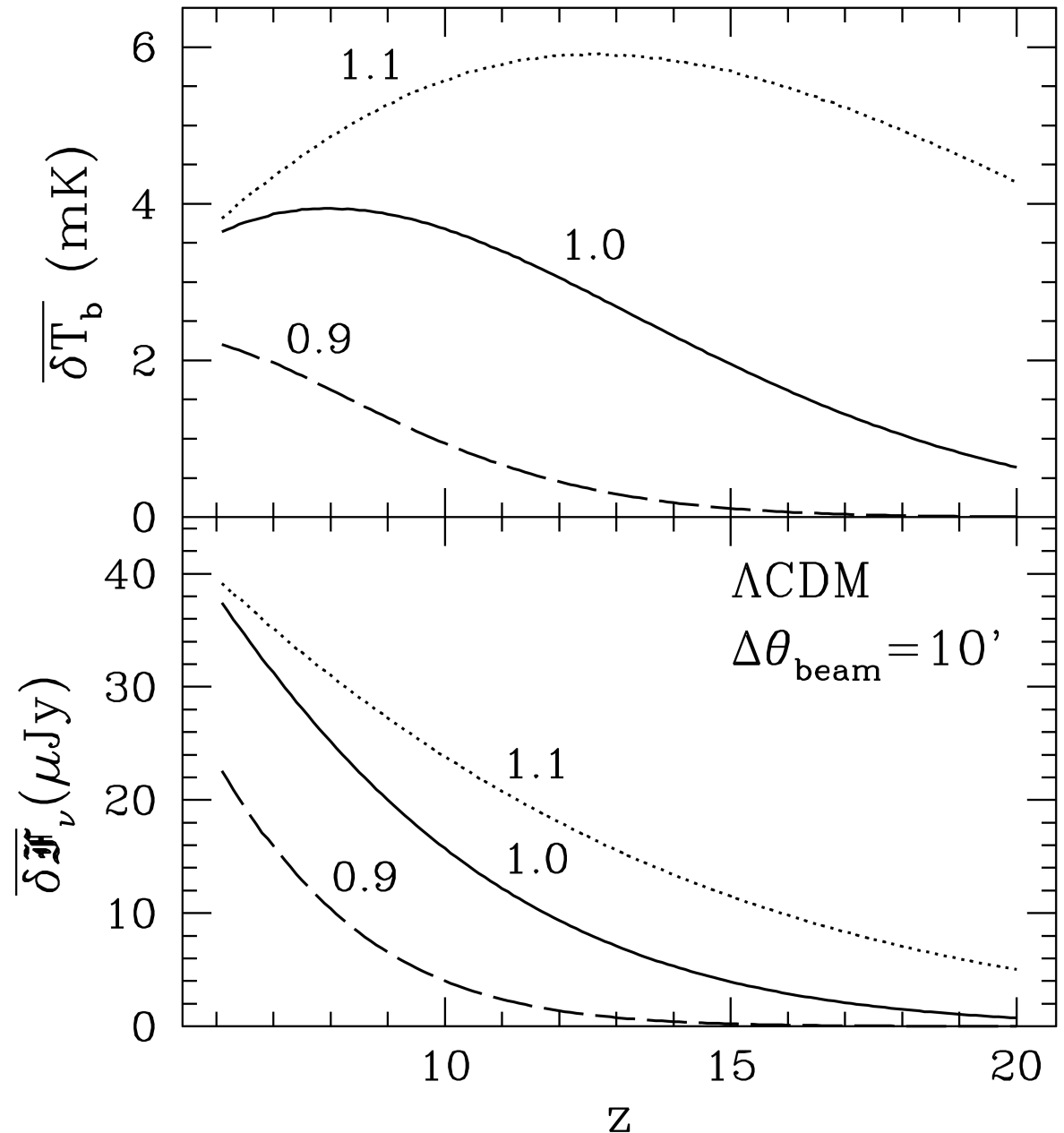
Collisional decoupling of the spin temperature T_S from T_{CMB} for different overdensities with respect to mean IGM.

bottom to top : $\delta = 0, 5, 10, 20, 50, 100$ and 200 at $z=9$ (left panel) and $z=19$ (right panel).

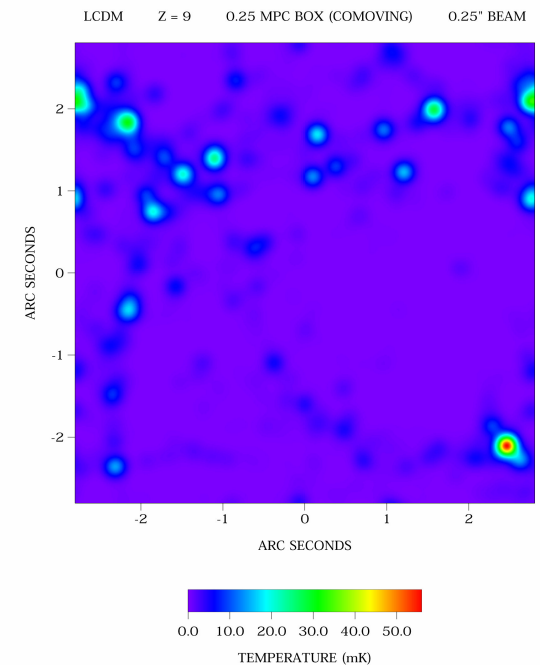
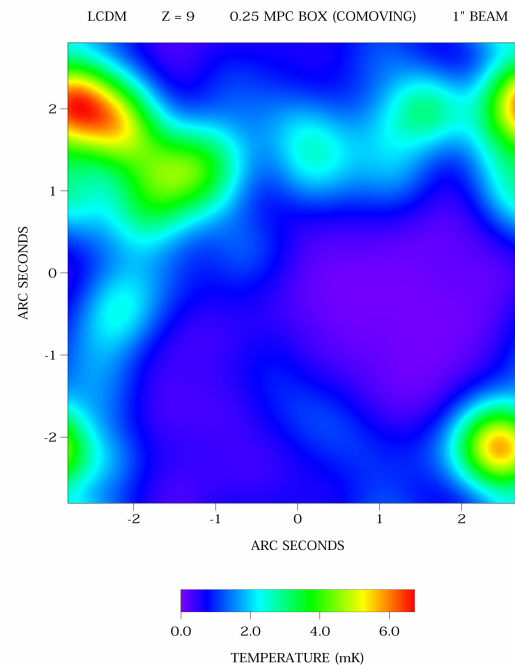
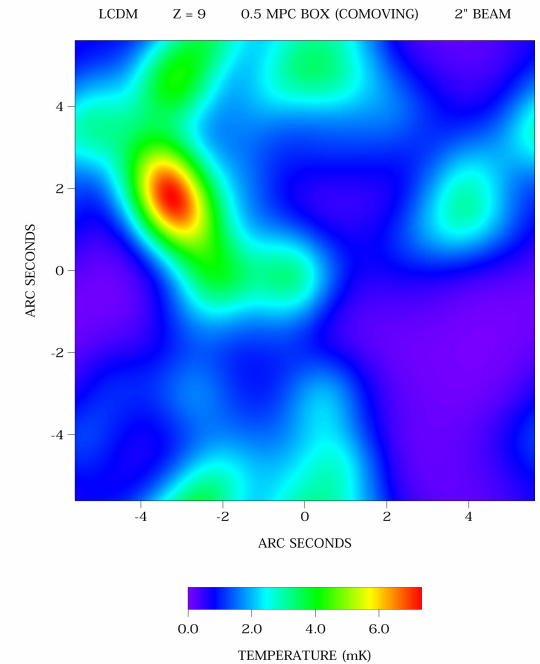
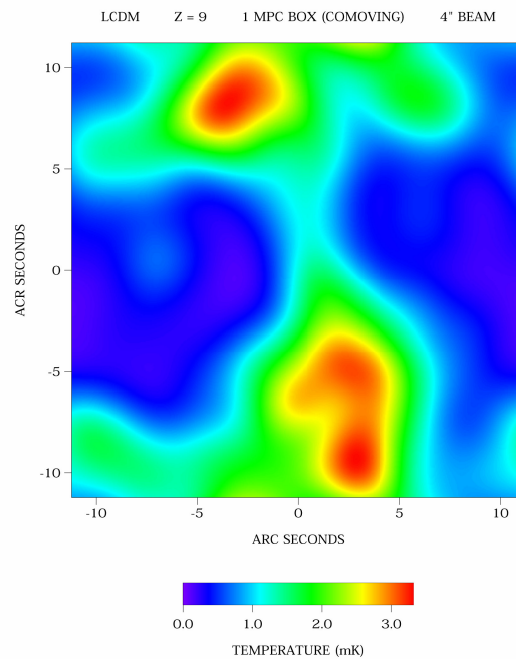


Iliev, Shapiro, Ferrara,
Martel (2002)

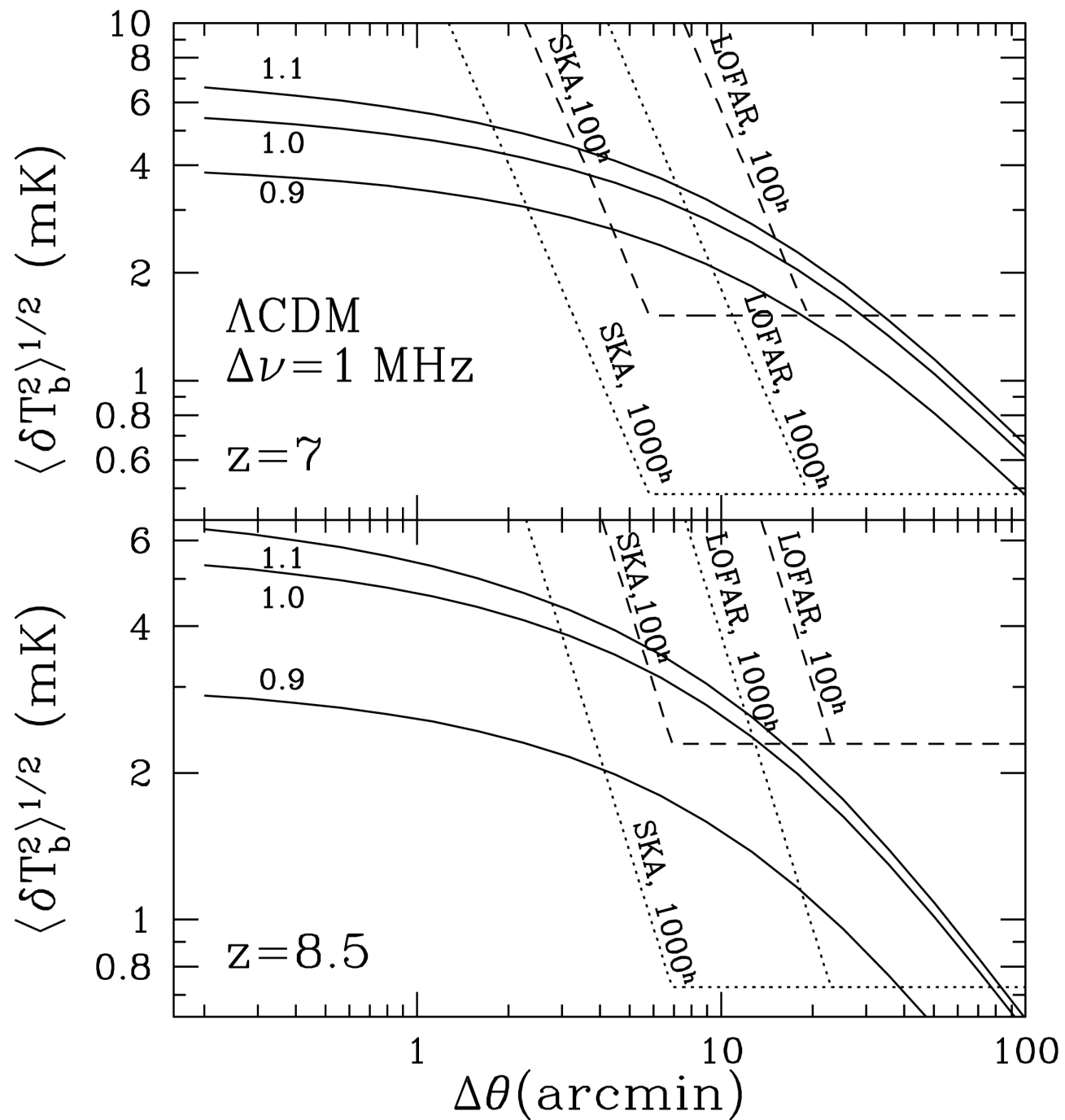
Minihalo radiation
background. Average
observed differential
antenna temperature
 $\overline{\delta T_b}$ and average
differential flux per unit
frequency $\overline{\delta \mathcal{F}_\nu}$ for beam
size of $\Delta\theta_{\text{beam}} = 10'$ at
the redshifted 21 cm
line frequency due to
minihalos vs. redshift z
for Λ CDM models with
power spectrum tilts n_p
 $= 0.9, 1.0,$ and $1.1,$ as
labeled.



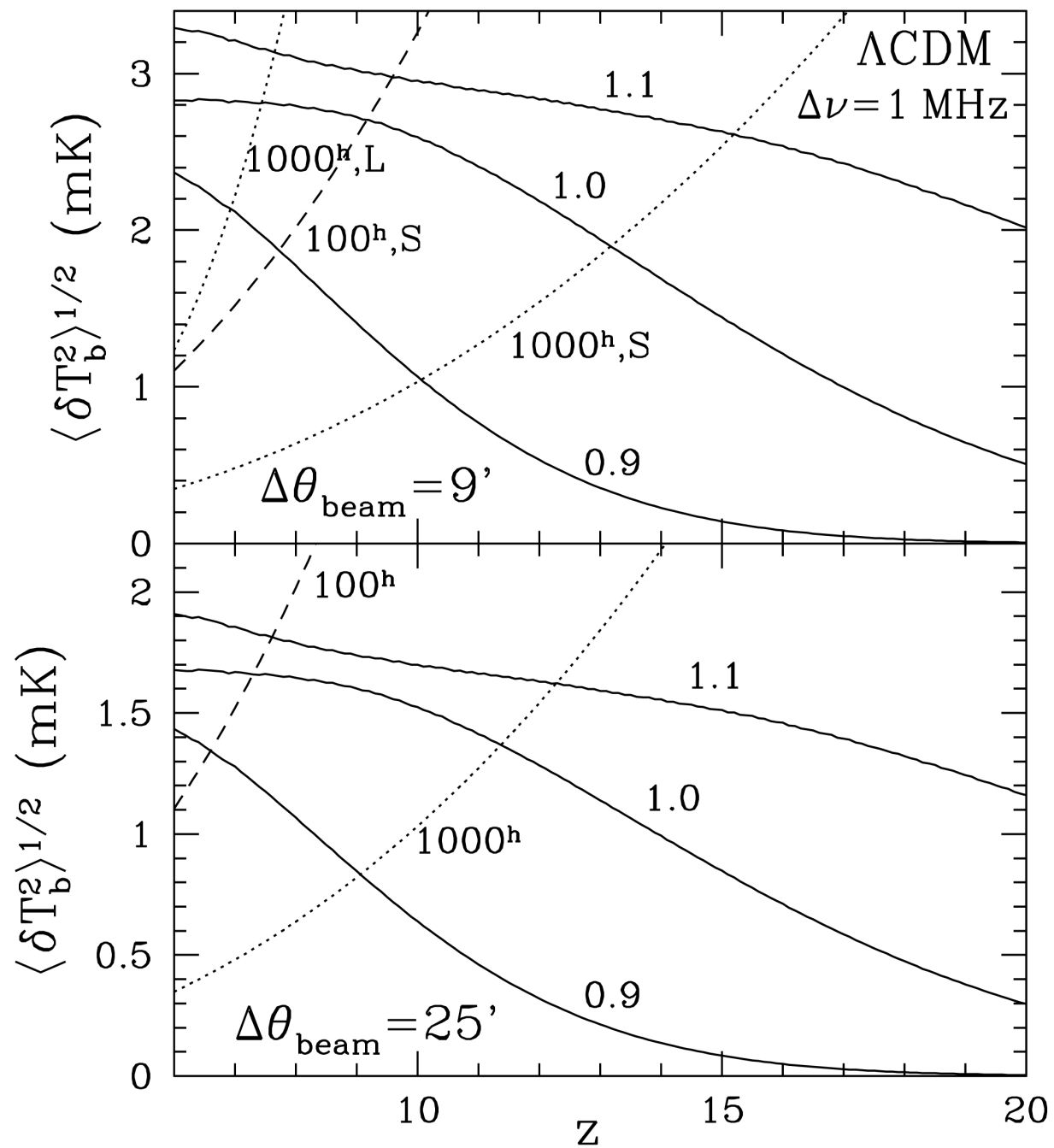
ANGULAR FLUCTUATIONS IN THE 21-CM EMISSION BACKGROUND. The clustering of minihalos as structure grows hierarchically in a CDM universe causes angular fluctuations in the 21-cm radiation background. To illustrate this, we have performed N-body simulations of halo formation in Λ CDM. Radio maps of the 21-cm emission from the simulation cubes at $z = 9$ are shown.



Predicted 3- σ
differential
antenna
temperature
fluctuations for
different beam
sizes.



Predicted 3- σ
differential antenna
temperature
fluctuations for
different redshifts.

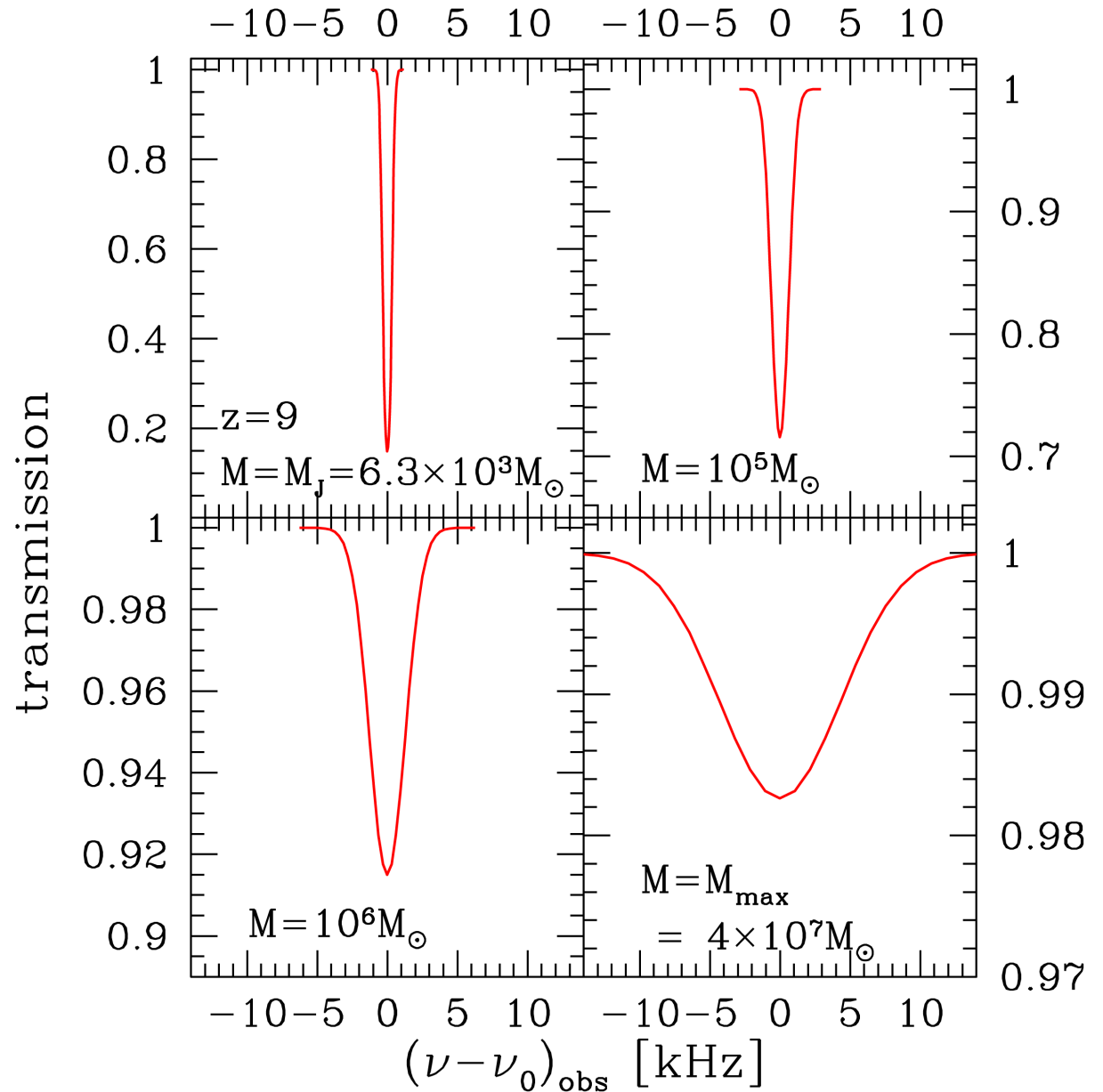


Detectability of Angular Fluctuations

- These $3\text{-}\sigma$ fluctuations should be observable with both LOFAR (300m filled aperture) and SKA (1km filled aperture) with integration times of between 100 and 1000 hours.
- For a 25' beam, for example, $3\text{-}\sigma$ fluctuations can be detected for untilted ΛCDM by both with a 100h integration for $z \sim 6 - 7.5$ and a 1000h integration for $z < \sim 11.5$, while for a 9' beam, SKA can detect them after 100h for $z < \sim 9$ and after 1000h for $z < \sim 13$.

Martel, Shapiro, Iliev,
Scannapieco, Ferrara
(2003)

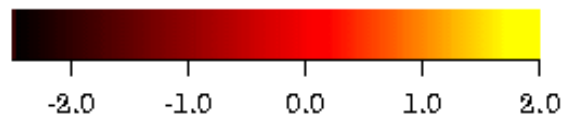
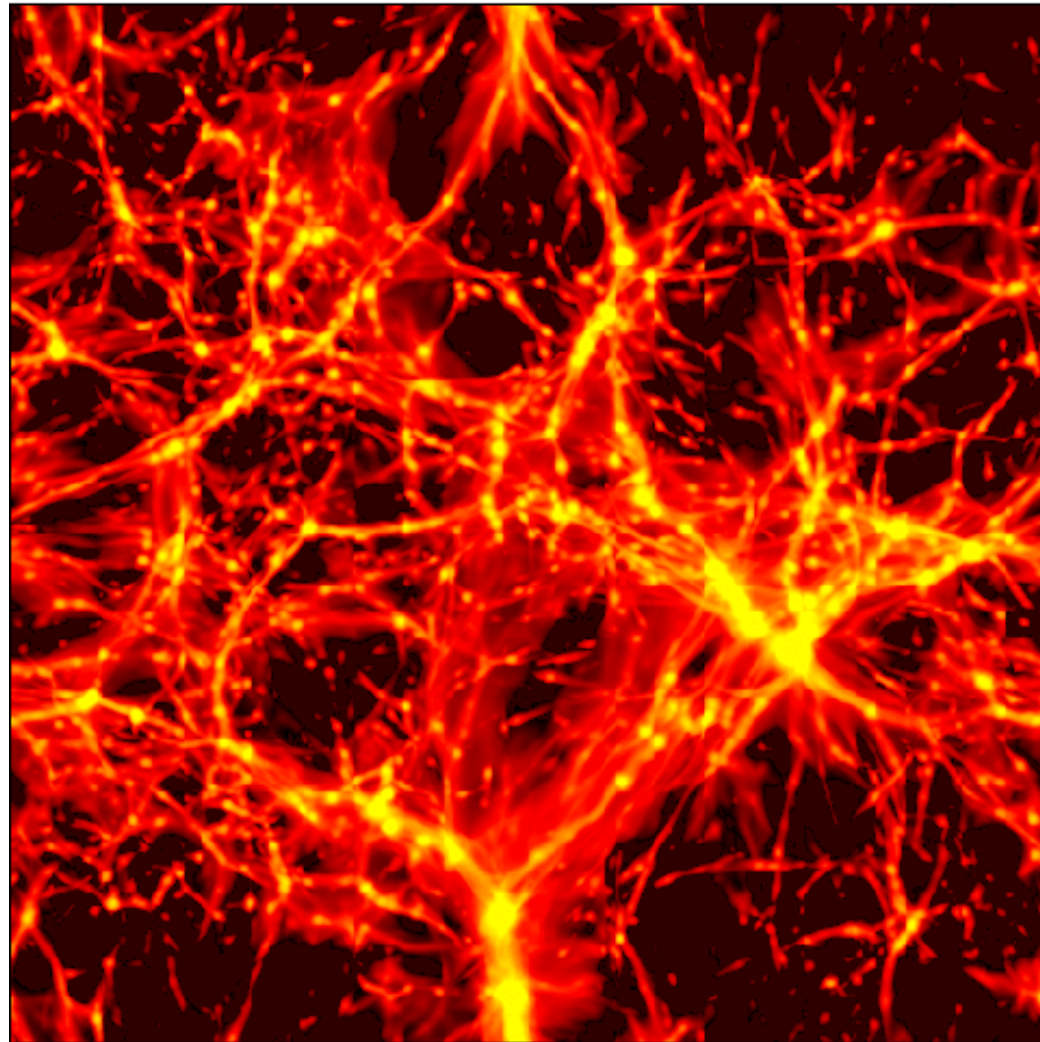
21-cm absorption
line profiles for
individual minihalos
of different mass at
 $z = 9$. Transmission
factor versus received
frequency ν_{rec} at $z = 0$,
for line of sight
through minihalo
center (i.e. zero
impact parameter).



21-cm Emission from Minihalos vs. Shock-heated Gas outside Virialized Regions (Shapiro, Ahn, Alvarez, Martel, Iliev, Ryu 2004)

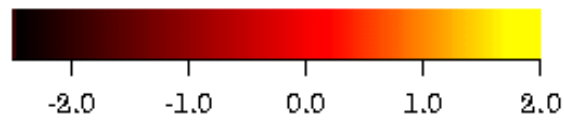
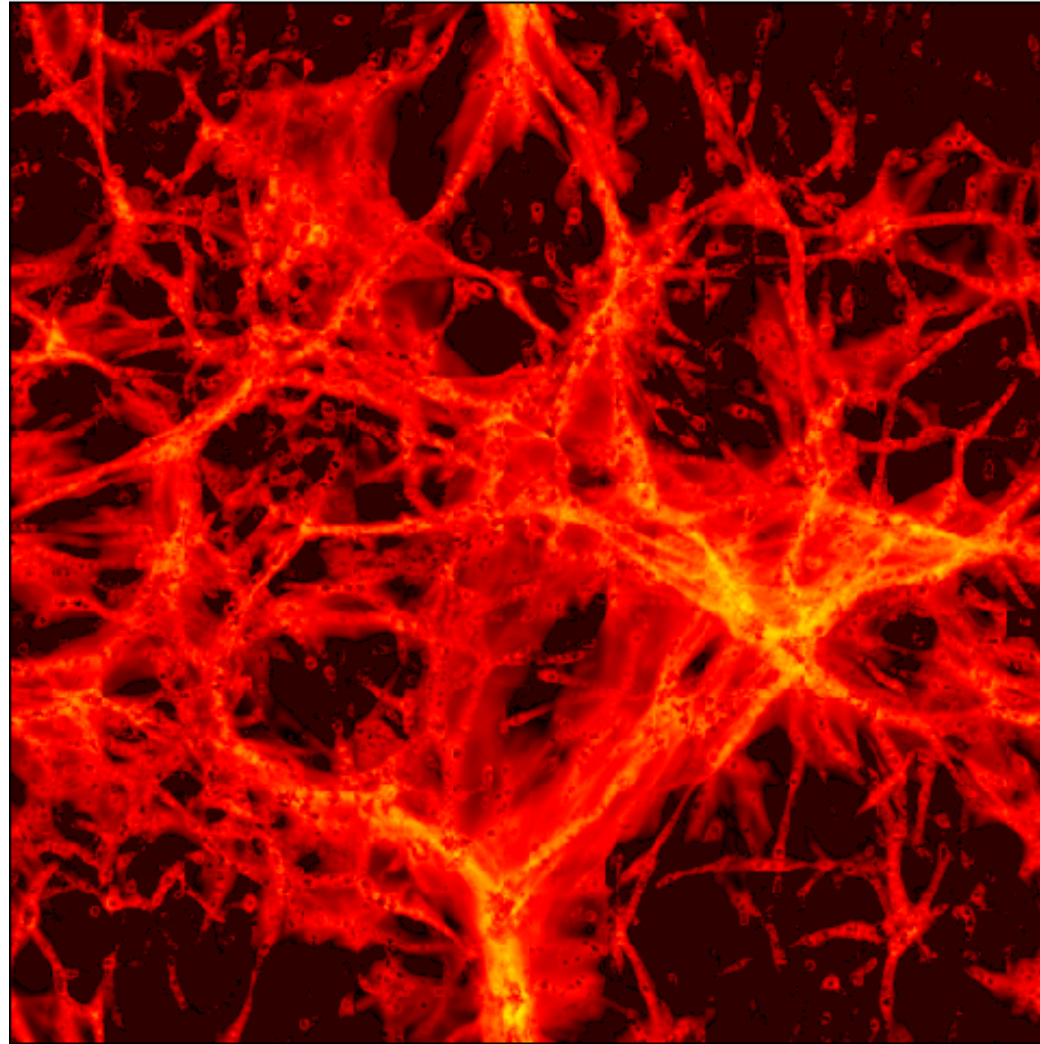
- Shock-heated IGM outside of the minihalos may also contribute to the 21-cm background, despite its less effective collisional pumping (Furlanetto and Loeb 2003).
- To quantify this, we simulated the gas and N-body dynamics of cosmic structure formation in LCDM before reionization, to compute the spin temperature at each point and map the differential brightness temperature vs. redshift.
- TVD/PM simulations of 0.7 Mpc comoving box size used 512^3 cells and 256^3 particles.

TOTAL Z=8



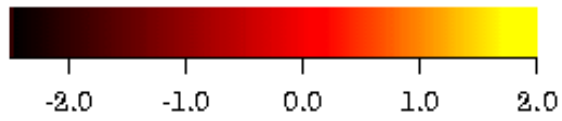
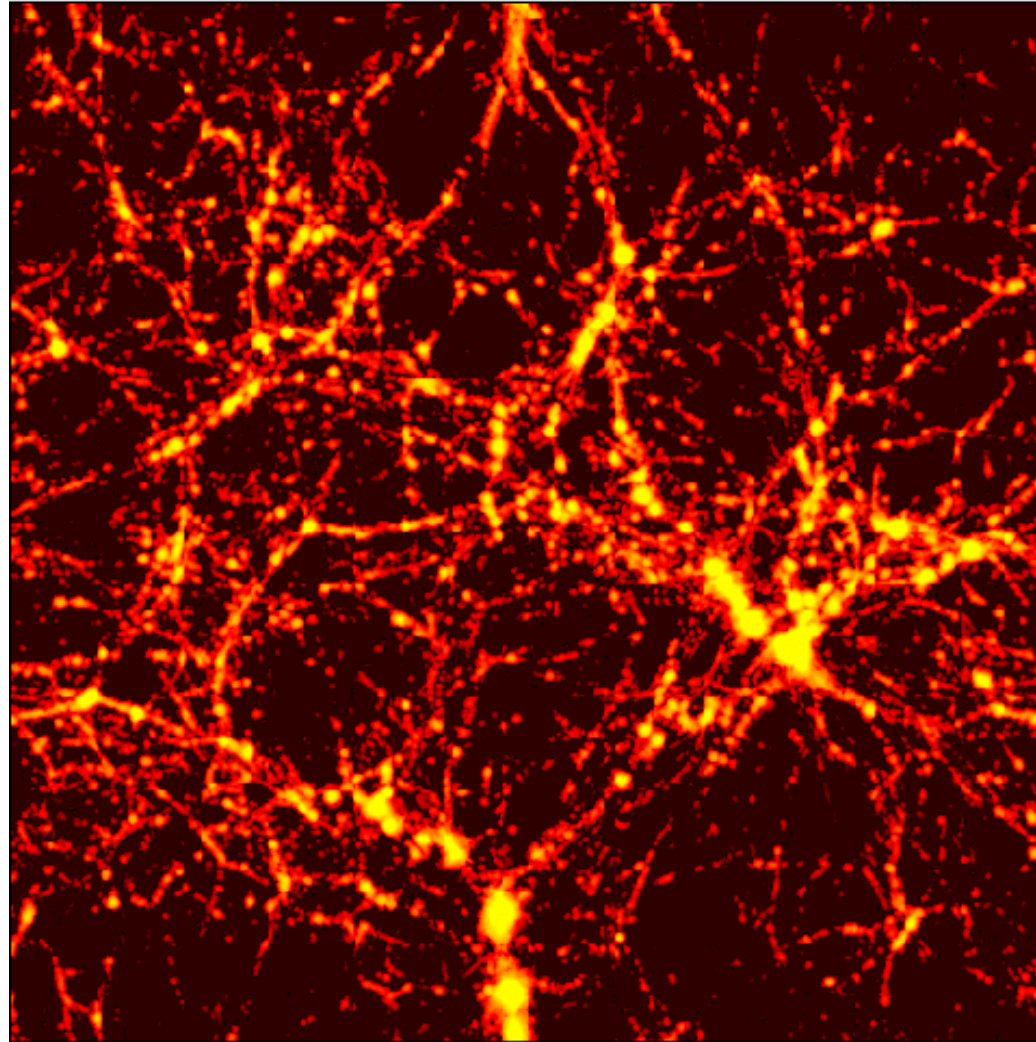
LOG 10 [DIFFERENTIAL BRIGHTNESS TEMPERATURE] (mK)

WITHOUT HALOS $Z=8$



LOG 10 [DIFFERENTIAL BRIGHTNESS TEMPERATURE] (mK)

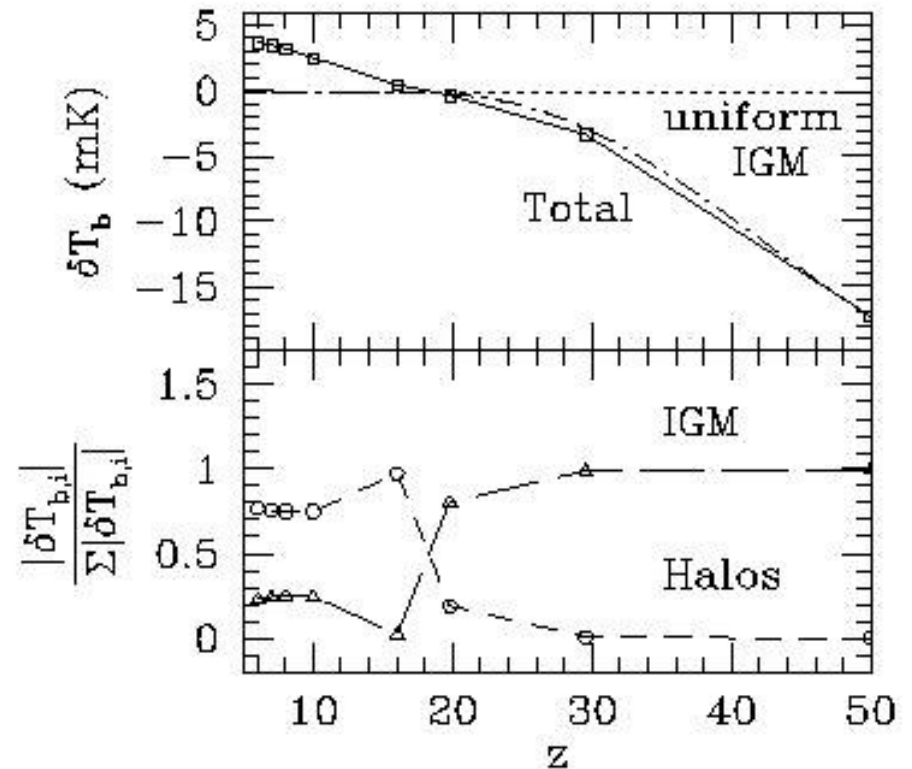
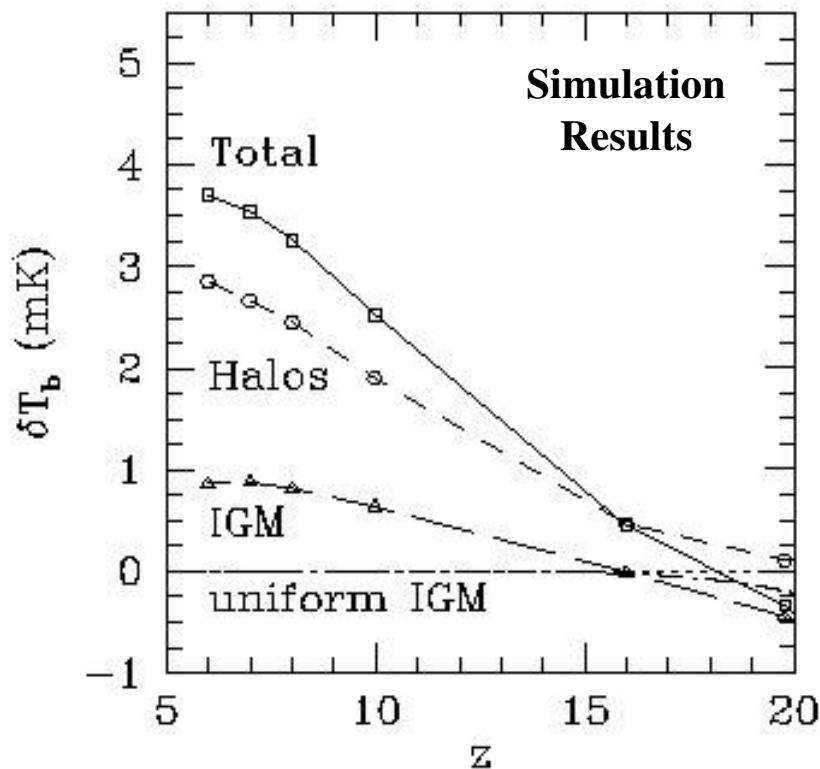
MINIHALOS ONLY Z=8



LOG 10 [DIFFERENTIAL BRIGHTNESS TEMPERATURE] (mK)

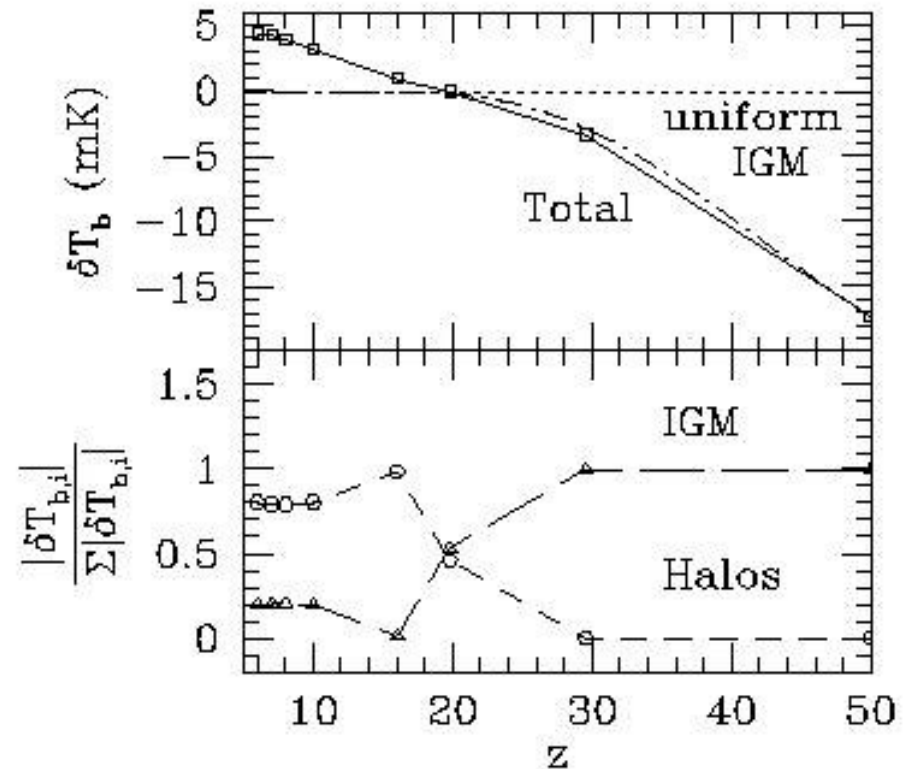
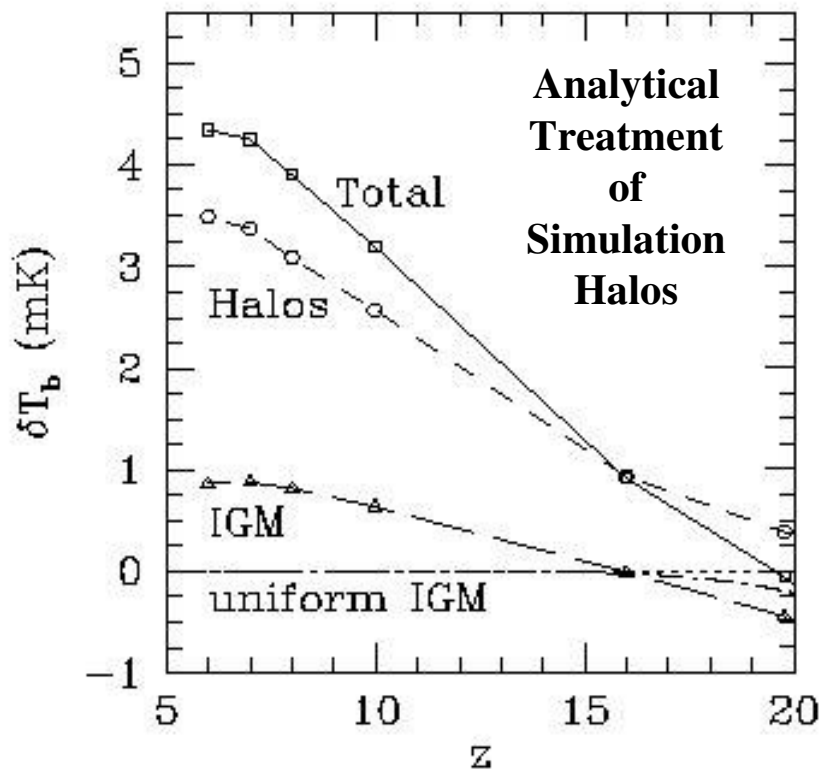
Minihalos vs. IGM

- Mean brightness temperature dominated by minihalos for $z < 20$.
- For $z > 16$, IGM contribution changes from emission to absorption.
- At $z > 20$, IGM absorption dominates over minihalo emission.
- At $z > 20$, mean brightness temperature same as uniform IGM



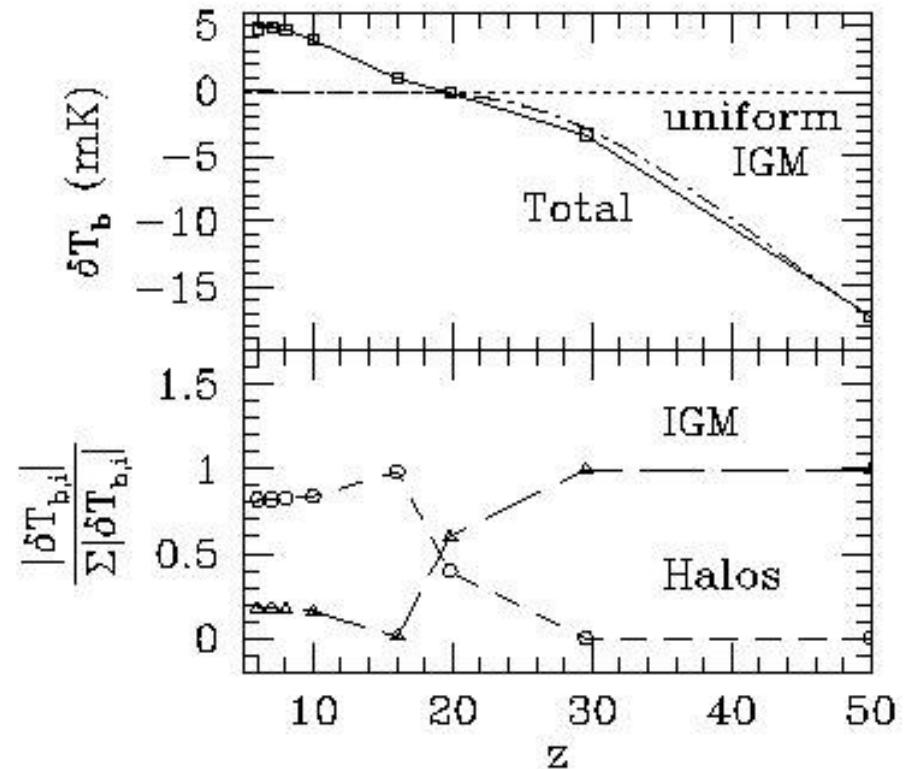
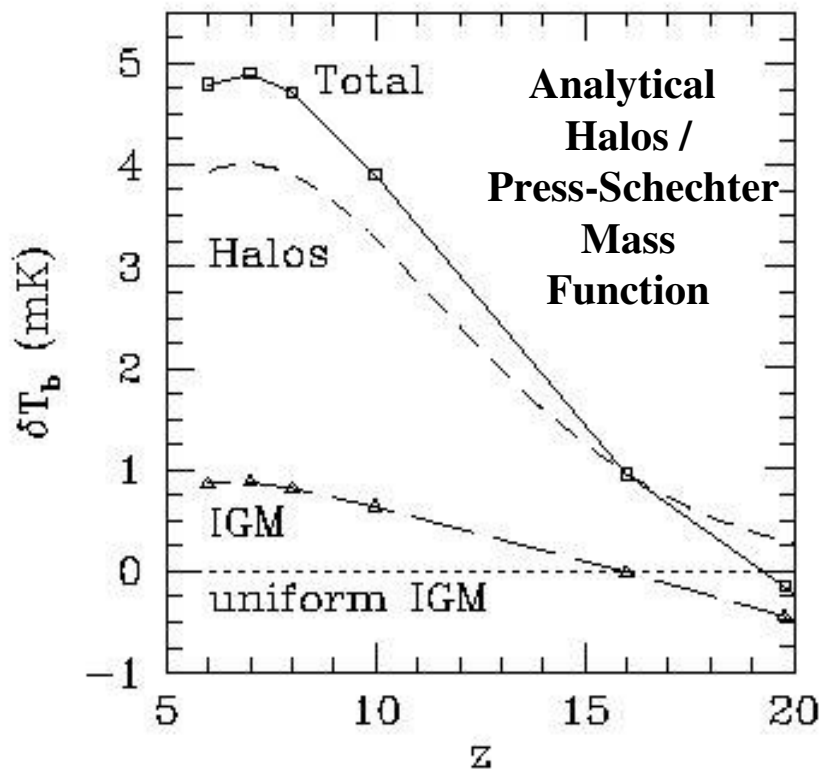
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Galaxy Formation and the Intergalactic Medium Research Group

Members:

- Paul Shapiro
- Hugo Martel
- Kyungjin Ahn
- Marcelo Alvarez

External Collaborators:

- Alejandro Raga (Mexico)
- Ilian Iliev (CITA), Andrea Ferrara (Italy), & Evan Scannapieco (KITP/UCSB)
- Hyesung Kang & Dongsu Ryu (Korea)

Computer Visualization Collaborator:

- Chandrajit Bajaj (Computer Science & ICES)