

# INTEGRABLE DEFORMATIONS AND TWISTS

BASED ON 2112.12025 AND 2207.14748

WITH SIBYLLE DRIESEN AND

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TALKING INTEGRABILITY: SPINS, FIELDS AND STRINGS  
KITP, 29/08/2022



"la Caixa" Foundation

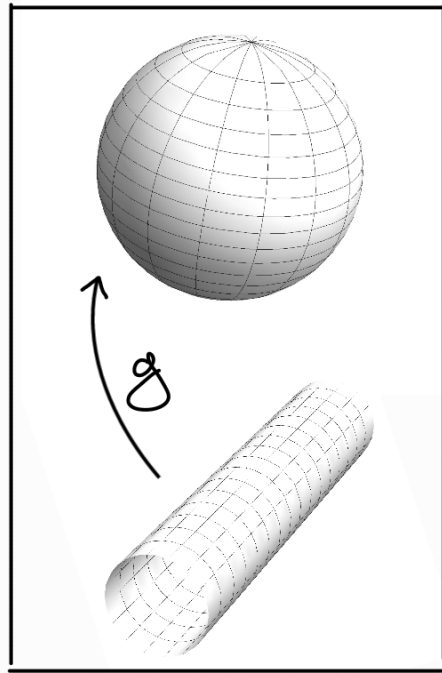
WHAT...

$\tau \in \mathbb{R}, \sigma \in [0, 2\pi], g: (\tau, \sigma) \rightarrow G$  LIE GROUP

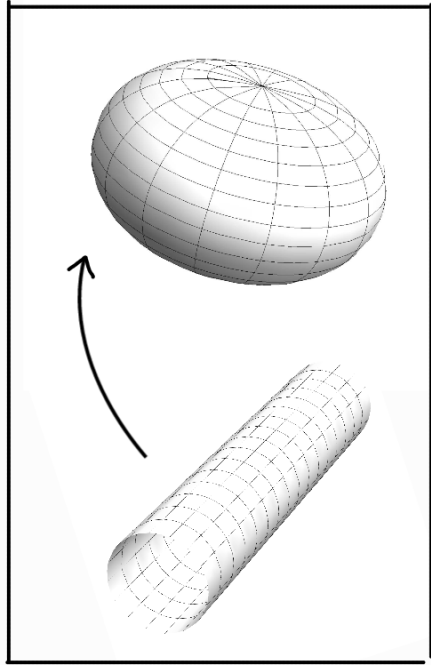
- PCM
- (SUPER) COSETS
- ...

UNDEFORMED & PERIODIC  $g(2\pi) = g(0)$

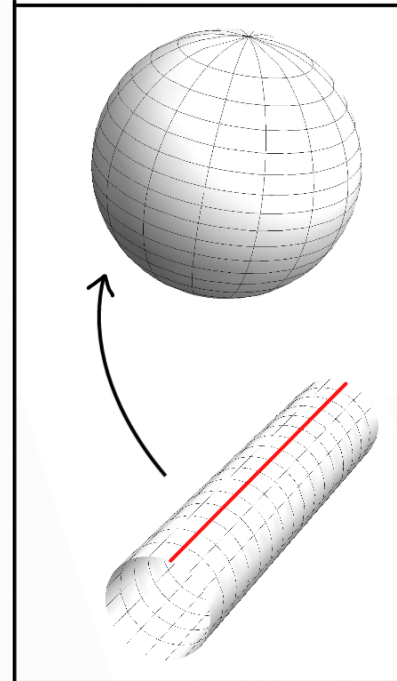
YANG-BAXTER DEF.



DEFORMED & PERIODIC



UNDEFORMED YET TWISTED



$g(2\pi) \neq g(0)$

EQUIVALENT

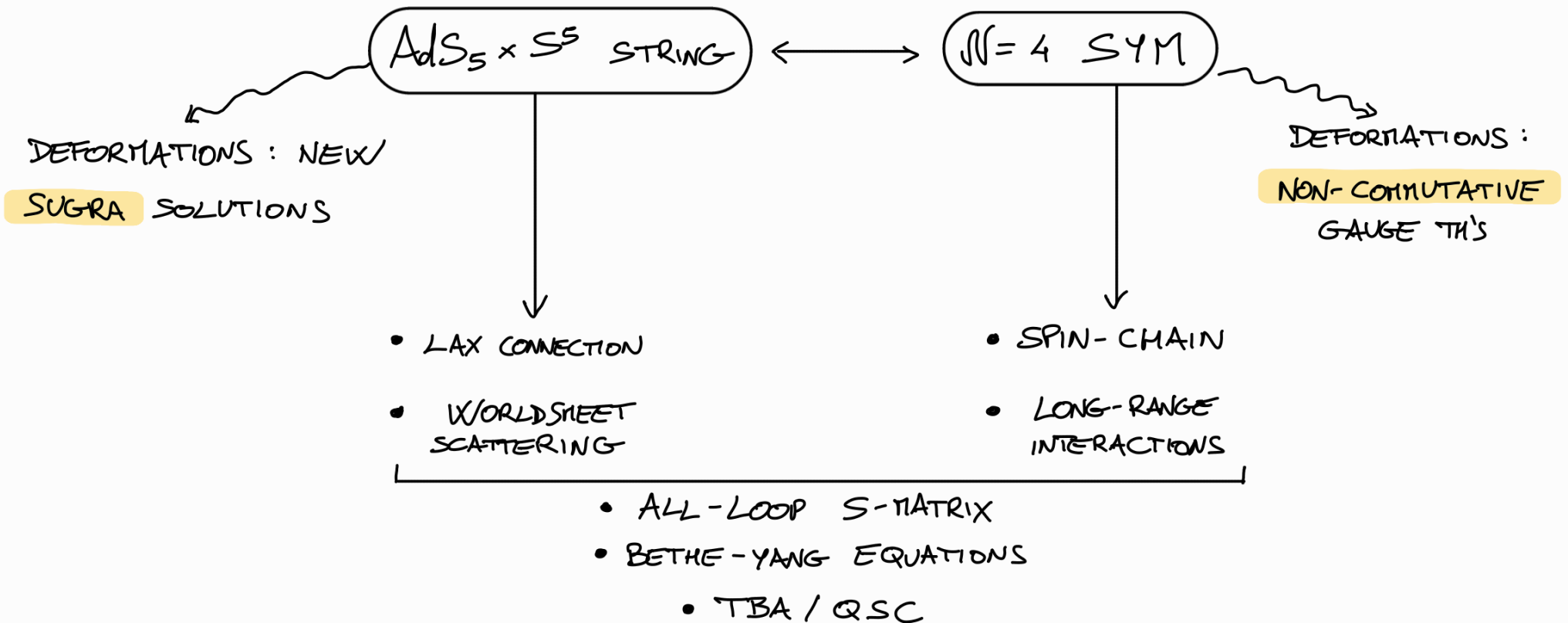
# WHY...

→ INTEGRABLE MODELS ARE PRECIOUS AND RARE

→ BREAK (SUPER) ISOMETRIES

→ UNDERSTAND ROLE OF HIDDEN SYMMETRIES

→ AdS/CFT CORRESPONDENCE

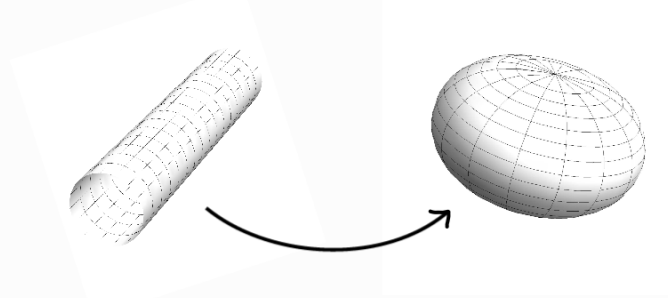


# TST DEFORMATIONS $\ni$ [LUNIN, MALDACENA '05]

SEE [FROLOV '05, FROLOV, ROIBAN, TSEYTLIN '05, ALDAY, ARUTYUNOV, FROLOV '05]

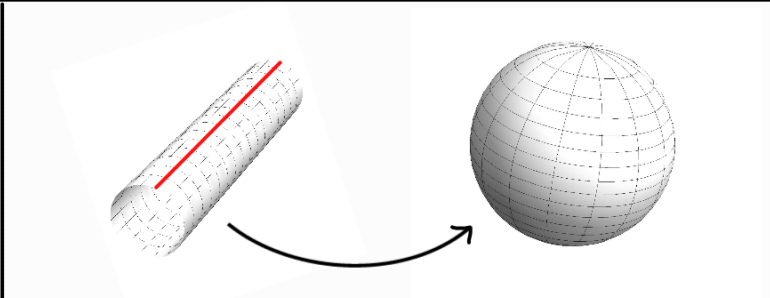
2 COMMUTING ISOMETRIES ACTING AS SHIFTS OF  $X^1, X^2$

T-DUALITY:  $X^1 \rightarrow T(X^1)$ , SHIFT:  $X^2 \rightarrow X^2 - \eta T(X^1)$ , T-DUALITY:  $T(X^1) \rightarrow X^1$



**DEFORMED**  $\sigma$ -MODEL  
WITH FIELDS  $X^1, X^2$   
**PERIODIC** BOUND. COND.  
 $X^i(2\pi) = X^i(0)$

ON-SHELL  
EQUIVALENT



**UNDEFORMED**  $\sigma$ -MODEL  
WITH FIELDS  $\tilde{X}^1, \tilde{X}^2$   
**TWISTED** BOUND. COND.  
 $\tilde{X}^i(2\pi) = \tilde{X}^i(0) + \eta \epsilon^{ij} Q_j$

$$\partial_\alpha X^i = \partial_\alpha \tilde{X}^i + \eta \epsilon^{ij} \epsilon_{\alpha\beta} J_j^\beta, \quad (\partial_\alpha J_i^\alpha = 0)$$

**YANG-BAXTER DEFORMATIONS**: **NON-LINEAR RELATIONS** FOR TWISTED B.C.

# TWISTED PCM

$g(\tau, \sigma) \in G$  LIE GROUP       $\mathfrak{g} = \text{Lie}(G)$        $\sigma^\pm = \tau \pm \sigma$

$$\boxed{S_{\text{PCM}} = - \int d^2\sigma \text{Tr}(J_+ J_-)} , \quad J = g^{-1} dg$$

EQ'S OF MOTION FROM  $\int_{\mathcal{G}} S_{\text{PCM}} \Rightarrow \partial_\alpha J^\alpha \approx 0$

VANISHING **BOUNDARY TERM**:  $\text{Tr}(g^{-1} dg J_\sigma)$  PERIODIC IN  $\sigma \in [0, 2\pi]$

$W$  CONSTANT  $\in G$   $\boxed{g(2\pi) = W g(0)}$  **LEFT TWIST** (HERE NO RIGHT)

FIELD REDEFINITION  $g' = hg$ ,  $h$  CONSTANT  $\in G$        $W' = h W h^{-1}$

LEFT TWISTS CLASSIFIED BY **ADJOINT ORBITS** OF  $G$

$G$  CONNECTED AND COMPACT  $\Rightarrow$  MAXIMAL ABELIAN SUBALGEBRA / WEYL GROUP

**YANG-BAXTER DEF'S**  $\approx W$  WRITTEN IN TERMS OF **CONSERVED CHARGES**

# YANG-BAXTER DEFORMATION (OF PCM)

[KLIMČÍK '03, DELDUC, MAGRO, VICEDO '13, KAWAGUCHI, MATSUMOTO, YOSHIDA '14, VAN TONGEREN '15...]

SEE [HOARE '21] FOR A REVIEW

$$S_{\text{YB}} = - \int d^2\sigma \text{Tr} \left( \partial_+ g g^{-1} \mathcal{O}^{-1} (\partial_- g g^{-1}) \right)$$

$$\mathcal{O}: \mathfrak{g} \rightarrow \mathfrak{g}$$

$$\mathcal{O} = \mathbb{1} - \eta R$$

$$R: \mathfrak{g} \rightarrow \mathfrak{g}, \quad t_a \in \mathfrak{g}, \quad R(t_a) = R_a^b t_b$$

$G_R$  UNBROKEN,  $G_L$  BROKEN TO  $H \subset G_L$ ,  $\text{Ad}_h^{-1} \circ R \circ \text{Ad}_h = R$

ANTI-SYMMETRY

$$\text{Tr}(x R(y)) = - \text{Tr}(R(x) y)$$

AND ...

# CLASSICAL YANG-BAXTER EQUATION (CYBE)

$$[R(x), R(y)] - R([R(x), y] + [x, R(y)]) = 0 \quad \forall x, y \in \mathfrak{g}$$

⚠ RHS = 0  $\leadsto$  "HOMOGENEOUS" YB DEF.

$$\text{Diagram 1} = \text{Diagram 2} \xrightarrow{\text{no}} [r_{12}, r_{13}] + [r_{12}, r_{23}] + [r_{13}, r_{23}] = 0$$

$$r = R^{ab} T_b \wedge T_a \in \mathfrak{g} \wedge \mathfrak{g}$$

$$\text{EX: } r = T_1 \wedge T_2$$

ABELIAN

$$[T_1, T_2] = 0 \leadsto T_S T \quad [\text{OSTEN, VAN TONGEREN '16}]$$

NON-ABELIAN

$$[T_1, T_2] = T_2 \leadsto \text{JORDANIAN DEF.}$$

## 2-COCYCLE CONDITION (2CC)

$$[R(x), R(y)] - R([R(x), y] + [x, R(y)]) = 0$$

$\mathfrak{F} = \text{Im}(R)$  IS SUBALGEBRA OF  $\mathfrak{g}$

$$\mathfrak{g} = \text{Ker}(R) \oplus \mathfrak{F}^*$$

$$R: \mathfrak{F}^* \rightarrow \mathfrak{F} \text{ SOLVES CYBE}$$

$\Downarrow$

$$\omega = R^{-1}: \mathfrak{F} \rightarrow \mathfrak{F}^* \text{ SOLVES 2CC}$$

$$\hat{\omega}: \mathfrak{F} \wedge \mathfrak{F} \rightarrow \mathbb{R}, \quad \hat{\omega}(x, y) = \text{Tr}(x \omega(y))$$

$$\hat{\omega}(x, [y, z]) + \hat{\omega}(y, [z, x]) + \hat{\omega}(z, [x, y]) = 0$$

$\omega$  IS COBOUNDARY IF  $\exists x \in \mathfrak{g}$  st.  $\omega = \text{ad}_x$  ( $\omega(y) = [x, y]$ )



# CLASSICAL INTEGRABILITY

$$[\partial_+ + \mathcal{L}_+, \partial_- + \mathcal{L}_-] = 0$$

PCM

$$\tilde{g} \in G, \quad \tilde{\mathcal{J}}_{\pm} = \tilde{g}^{-1} \partial_{\pm} \tilde{g}$$

$$\textcircled{1} \partial_+ \tilde{\mathcal{J}}_- + \partial_- \tilde{\mathcal{J}}_+ \approx 0$$

$$\textcircled{2} \partial_+ \tilde{\mathcal{J}}_- - \partial_- \tilde{\mathcal{J}}_+ + [\tilde{\mathcal{J}}_+, \tilde{\mathcal{J}}_-] = 0$$

$$\mathcal{L}_{\pm} = \frac{\tilde{\mathcal{J}}_{\pm}}{1 \mp z}$$

YB

$$g \in G, \quad A_{\pm} = \frac{1}{1 \pm \eta \text{Ad}_g^{-1} \circ R \circ \text{Ad}_g} (g^{-1} \partial_{\pm} g)$$

$$\textcircled{1} \partial_+ A_- + \partial_- A_+ \approx 0$$

$$\textcircled{2} \partial_+ A_- - \partial_- A_+ + [A_+, A_-] \approx 0$$

$$\mathcal{L}_{\pm} = \frac{A_{\pm}}{1 \mp z}$$

↑  
R<sup>T</sup> = -R  
AND  
CYBE

ON-SHELL EQUIVALENCE

$$\tilde{\mathcal{J}}_{\pm} = A_{\pm}$$

# THE TWIST

SEE [VICEDO '15] AND [VAN TONGEREN '18]

$$\text{YB} \rightarrow \left( \mathfrak{g} = \mathcal{F} \tilde{\mathfrak{g}} \right) \leftarrow \text{PCM}$$

↑  
TWIST  
FIELD

$$\tilde{\mathfrak{J}}_{\pm} = A_{\pm} \Rightarrow \partial_{\pm} \mathcal{F} = V_{\pm} \mathcal{F}, \quad V_{\pm} = \pm \eta R (\mathfrak{g} A_{\pm} \mathfrak{g}^{-1})$$

$$\mathcal{F}(\tau, \sigma) = \mathcal{P} \exp \left( \int_{(0,0)}^{(\tau, \sigma)} d\xi^{\alpha} V_{\alpha} \right) \mathcal{F}(0,0)$$

PERIODIC YB

$$\mathfrak{g}(2\pi) = \mathfrak{g}(0)$$

$$\mathbb{W} \xrightarrow{\eta \rightarrow 0} \mathbb{1}$$

TWISTED PCM

$$\tilde{\mathfrak{g}}(2\pi) = \mathcal{F}^{-1}(2\pi) \mathfrak{g}(2\pi)$$

$$= \underbrace{\mathcal{F}^{-1}(2\pi) \mathcal{F}(0)}_{\mathbb{W} \text{ TWIST}} \tilde{\mathfrak{g}}(0)$$

$\mathbb{W}$  TWIST

# A NEW TWIST

[RB, DRIESEN, MIRAMONTES '21]

$$\partial_{\pm} F F^{-1} = \pm \eta R (F \partial_{\pm} \tilde{g} \tilde{g}^{-1} F^{-1})$$

↓

$$\eta^{-1} P^T (F^{-1} \omega(\partial_{\pm} F F^{-1}) F) = P^T (\partial_{\pm} \tilde{g} \tilde{g}^{-1}) \quad (\text{s.o.v.})$$

$$F = \exp(RX) \in F$$

$$Y = \frac{1}{\eta} P^T \frac{1 - e^{-\text{ad}_R X}}{\text{ad}_R X} X$$

$$e^{-z} dz = \frac{1 - e^{-\text{ad}_z}}{\text{ad}_z} dz$$

$$\partial_{\pm} Y = \eta^{-1} P^T (F^{-1} \omega(\partial_{\pm} F F^{-1}) F)$$

$$\sim Y = \int P^T (* d\tilde{g} \tilde{g}^{-1})$$

$$\partial_{\pm} W = 0, \quad W \in F$$

$$W = \exp(-\eta RQ)$$

# JORDANIAN DEFORMATION

$$sl(2, \mathbb{R}) \subset \mathfrak{g} \quad [T_0, T_{\pm}] = \pm T_{\pm}, \quad [T_+, T_-] = 2T_0$$

$$r = T_+ \wedge T_0$$

$$X = X_0 T^0 + X_+ T^+$$

$$Y = Y_0 T^0 + Y_+ T^+$$

$$Y = \frac{1}{\eta} P^T \frac{1 - e^{-\text{ad}_R X}}{\text{ad}_R X} X$$

$$X_+ = -\log(1 - \eta Y_+)$$
$$X_0 = -\frac{Y_0}{Y_+} \log(1 - \eta Y_+)$$

$$W = \exp[\mathcal{Q}(T_0 - \eta T_+)]$$

$$\mathcal{Q} = \log \frac{1 - \eta Y_+(0)}{1 - \eta Y_+(2\pi)}$$

$$\eta = \frac{Y_0(2\pi) - Y_0(0)}{Y_+(2\pi) - Y_+(0)}$$

# MONODROMY MATRIX AND CSC

$$\Omega(z, \tau) = \mathcal{O} \exp(-\oint \mathcal{L}_\sigma), \quad \mathcal{L}_\pm = \frac{\tilde{\mathcal{J}}_\pm}{1 \mp z}$$

$$\textcircled{z \sim 0} \quad \mathcal{L}' = \tilde{g}(\mathcal{L} + d)\tilde{g}^{-1}$$

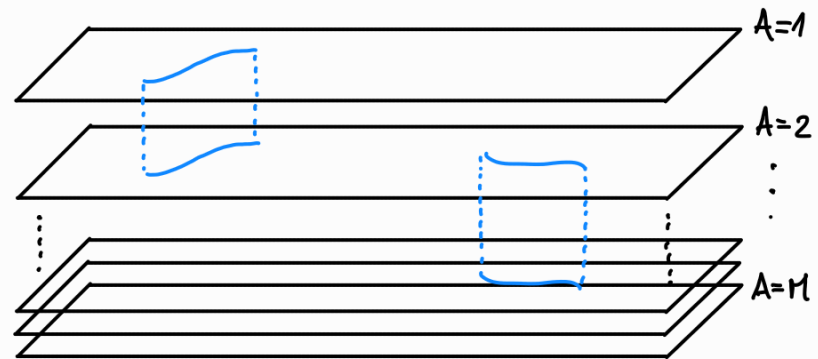
$$\Omega = \tilde{g}(2\pi)^{-1} \mathcal{O} \exp(-\oint \mathcal{L}'_0) \tilde{g}(0) = \tilde{g}(0)^{-1} W^{-1} \left( 1 - z \int_0^{2\pi} d\sigma \mathcal{L}'_0 \tilde{g} \tilde{g}^{-1} + \dots \right) \tilde{g}(0)$$

EIGENVALUES  $\lambda_A = e^{iP_A}$  "POLARIZED" BY  $W$

$$P_A \stackrel{z \rightarrow 1}{\sim} \pm \frac{q_A^+}{1 \mp z}$$

$$P(z) \stackrel{z \rightarrow 0}{\sim} P(0) + z P'(0) + \dots$$

$$P(z) \stackrel{z \rightarrow \infty}{\sim} \frac{P_{\infty}}{z} + \dots$$



$$P_A(z+i\epsilon) - P_B(z-i\epsilon) = 2\pi M_{(A,B)} \quad z \in C_{(A,B)}$$

# JORDANIAN DEFORMATION OF $AdS_5 \times S^5$

[R.B., DRIEZEN, NIETO GARCIA, WYSS '22]

$32 \rightarrow 12$  SUPERISOMETRIES

$$\text{AdS QUASIMOMENTA} \left\{ \begin{array}{l} \hat{P}_1 = -\frac{2\pi}{x}(E+S) + \mathcal{O}(1/x^2) \\ \hat{P}_2 = -\frac{i}{2}Q + \frac{2\pi}{x}S + \mathcal{O}(1/x^2) \\ \hat{P}_3 = +\frac{i}{2}Q + \frac{2\pi}{x}S + \mathcal{O}(1/x^2) \\ \hat{P}_4 = +\frac{2\pi}{x}(E-S) + \mathcal{O}(1/x^2) \end{array} \right. \quad \begin{array}{l} \hat{P}_1 \sim \frac{2\pi}{x}(+E-S_1+S_2) \\ \hat{P}_2 \sim \frac{2\pi}{x}(+E+S_1-S_2) \\ \hat{P}_3 \sim \frac{2\pi}{x}(-E-S_1-S_2) \\ \hat{P}_4 \sim \frac{2\pi}{x}(-E+S_1+S_2) \end{array}$$

GLOBAL-TIME TRASL. GENERATED BY  $P^0 - P^1 + \frac{1}{2}(K^0 + K^1) \neq \text{BIN } P^0 + K^0$

$S = \text{ANGULAR MOMENTUM IN AdS } (J_{23})$

# SEMICLASSICAL CORRECTION TO CSC

SEMICLASSICAL QUANTIZATION À LA [GROMOV AND VIEIRA '07]

$$P \rightarrow P + \delta P$$

NEW POLES FOR "MICRO CUTS"

$\delta P$  FIXED BY CONSISTENCY CONDITIONS (GLUING AT CUTS AND ASYMPTOTICS)

BMN-LIKE SOLUTION:  $E = a$ ,  $S = 0$ ,  $Q = 4\pi a \beta$

$$\hat{P}_1 = -\hat{P}_4 = \frac{2\pi a \sqrt{x^2 - \beta^2}}{1 - x^2}, \quad \hat{P}_2 = -\hat{P}_3 = \frac{2\pi a x \sqrt{1 - \beta^2 x^2}}{1 - x^2}$$

$$\delta E = \frac{a}{2} \left[ 2\beta(1-\beta) - (3-\beta^2) \log(1-\beta^2) - (1+\beta^2) \log(1+\beta^2) + (1+\beta^2) \log \frac{1-\beta}{1+\beta} \right]$$

$$\delta S = 0$$

$$\delta Q = 0$$

# SUMMARY, OUTLOOK, AND (SOME OF) WHAT I DIDN'T SAY

- EASIER TO WORK WITH TWISTED MODEL THAN DEF. ONE
- TWIST ( $\Rightarrow P$ ) HIGHLY NON-LOCAL IN TERMS OF  $g \in YB$   
↳ HIDDEN SYMMETRIES IMPORTANT
- DEFORMATION / TWIST & EXPANSION OF  $P$  DON'T COMMUTE  
 $P \sim \dots + \sqrt{Q^2 + \# / X^2 + \dots}$  NO DEF. OF BETHE EQ'S ?
- $\delta Q = 0$ , ACCIDENT OF THIS SOLUTION ?
- FACTORIZATION OF THE JORDANIAN TWIST  $W = e^{qT_+} e^{qT_0} e^{-qT_+}$
- UNIMODULAR (= SUGRA) VS NON-UNIMODULAR (= NON-SUGRA) JORD. DEF'S  
SPECTRAL EQUIVALENCE (?)



- YB  $\approx$  INTERPOLATION BETWEEN MODEL AND NON-ABELIAN T-DUALITY  
[HOARE, TSEYTLIN '16, RB WULFF '16]
  - POISSON-LIE SYMMETRIES & POISSON-LIE DUALITY TRANSFS [KLIMČÍK, SEVERA '95]
  - RELATION TO 4D CHERN-SIMONS THEORY [COSTELLO, WITTEN, YAMAZAKI '17, ...]
  - RELATION TO GENERALISED GEOMETRY AND DOUBLE FIELD THEORY [...]
- 
- DEF'S BREAKING BMN ISOMETRIES  
 $\hookrightarrow$  ALTERNATIVES TO BMN SOLUTION ?
 

}	MASSIVE SPECTRUM ?
	SYMMETRIES FOR S-MATRIX ?
	...
  - DRINFELD TWISTS ( $\Delta_F = F^{-1} \circ \Delta \circ F$ ) FOR THE QUANTUM MODELS ?  
SEE [VAN TONGEREN '15 '16]
  - NON-DIAGONALISABLE TWISTS ? [GUCA, LEVKOVICH-MASYUK, ZARETBO '17]
  - TWISTED SPIN-CHAINS ?
  - DEFORMATIONS OF  $\mathcal{N}=4$  SYM ?
  - ...

Thank  
you!