

Metallic Criticality in the Charge Transport of $\text{BaFe}_2(\text{As}_{1-x}\text{P}_x)_2$

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Arkady Shekter, Ross D. McDonald, Mun K. Chan, Brad Ramshaw

NHMFL Pulsed Field Facility, Los Alamos National Laboratory

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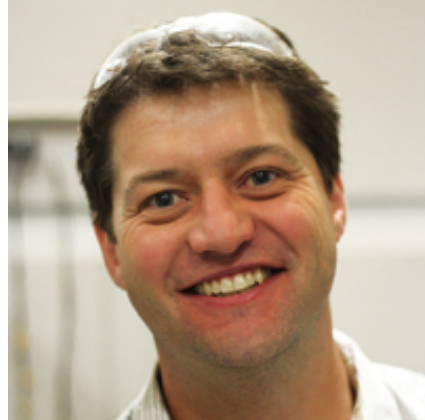
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Outline

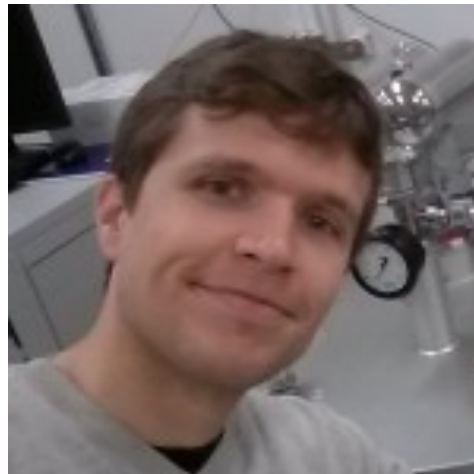
- Phenomenology of Quantum Critical metals
 - T-linear resistivity, $1/T$ Hall effect
 - Fan-like critical region
- High Field Magnetoresistance in Ba-122
 - T-linear/B-linear MR
 - Scaling between T and B, with the same exponent
- Phenomenology of R_H
 - $1/T$ enhancement at low temperatures
 - Reduction of R_H at high fields and cut-off at high dopings

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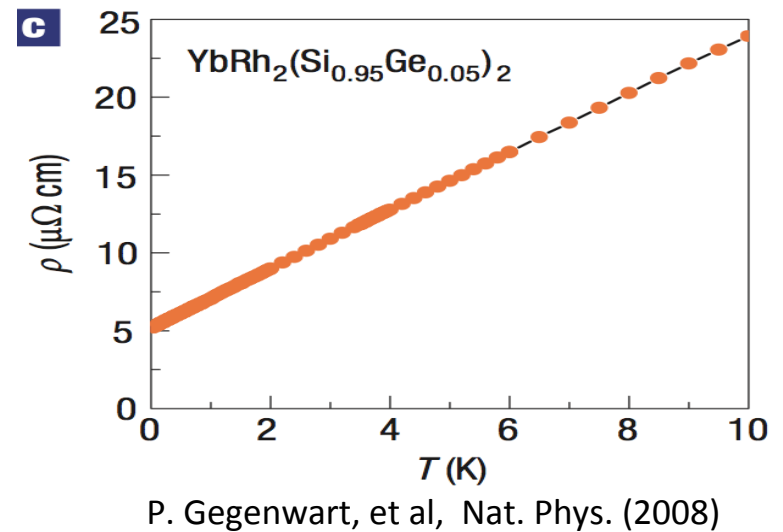
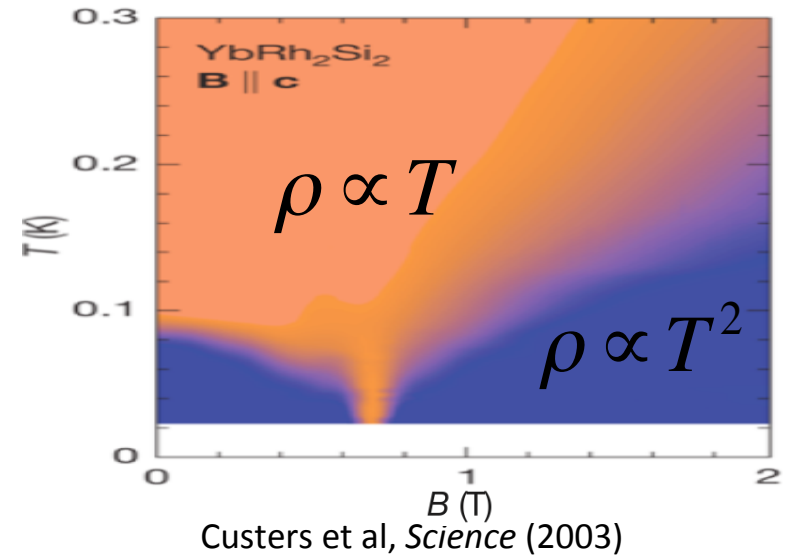
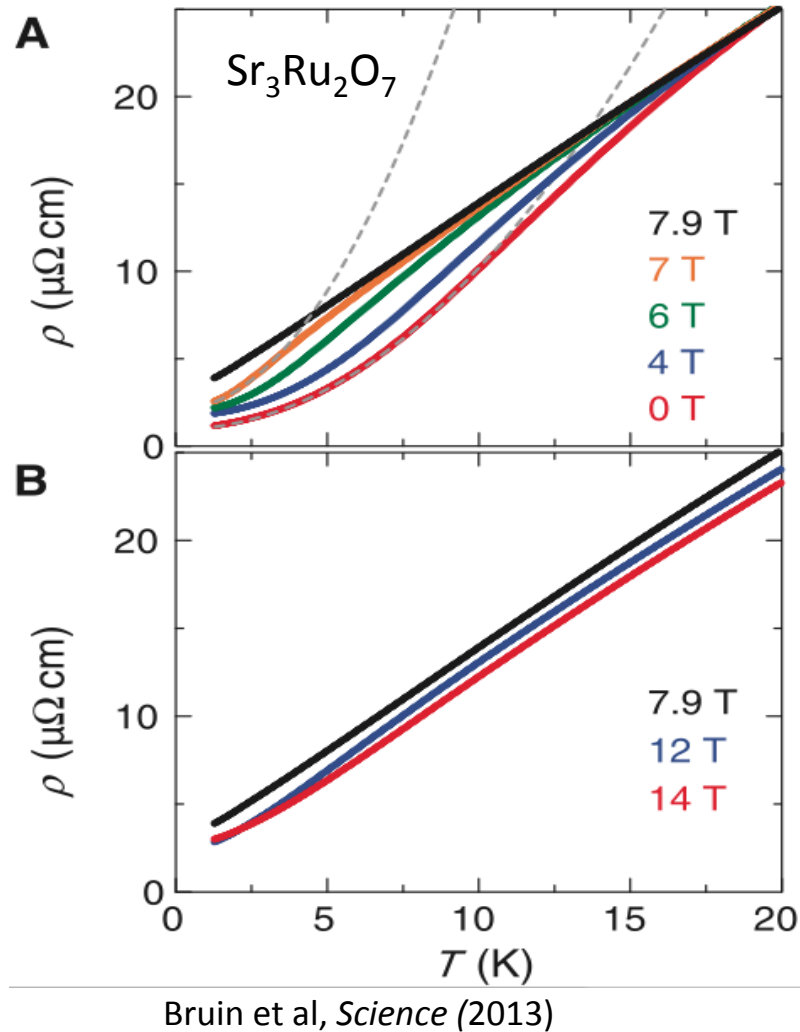
- NHMFL collaborators Arkady Shechter, Ross McDonald, Mun Chan & Brad Ramshaw



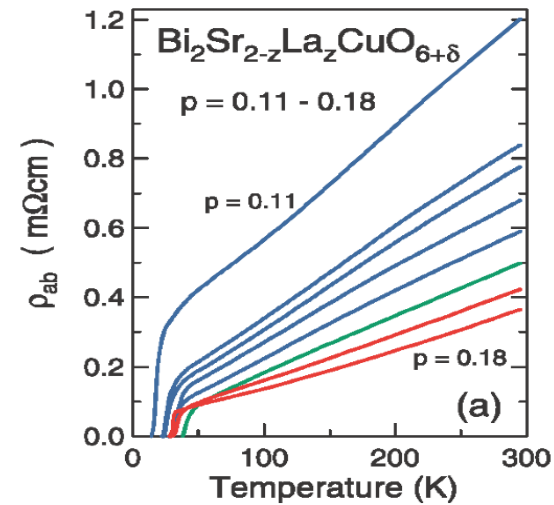
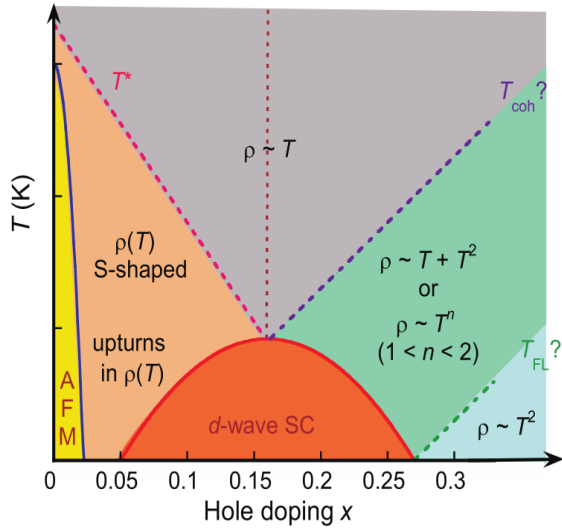
- Colleagues at Berkeley Nicholas Breznay, Toni Helm, Nikola Maksimovic & James Analytis



T-linear Resistivity in Quantum Critical Metals

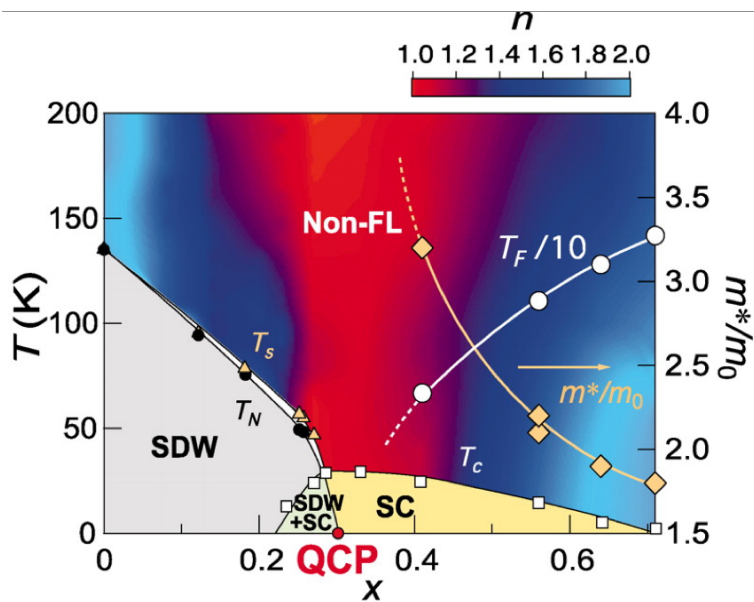


(Nonmagneto-)Transport in High- T_c s

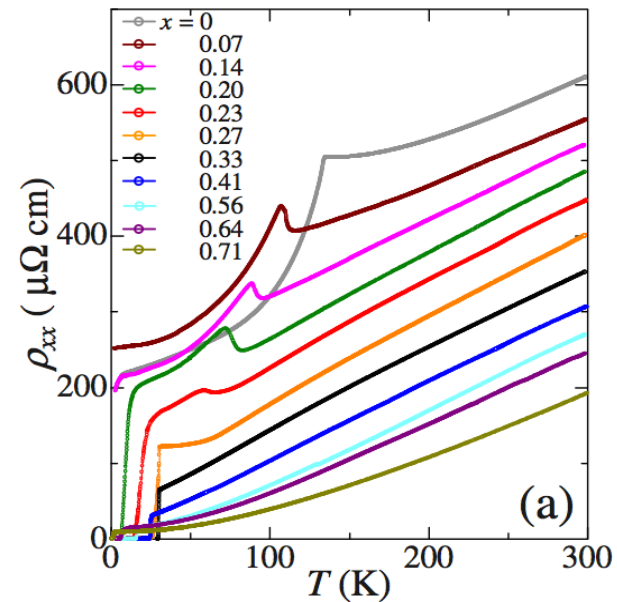


Y. Ando, PRL (2004)

N.E. Hussey, Journal of Physics: Condensed Matter (2008)

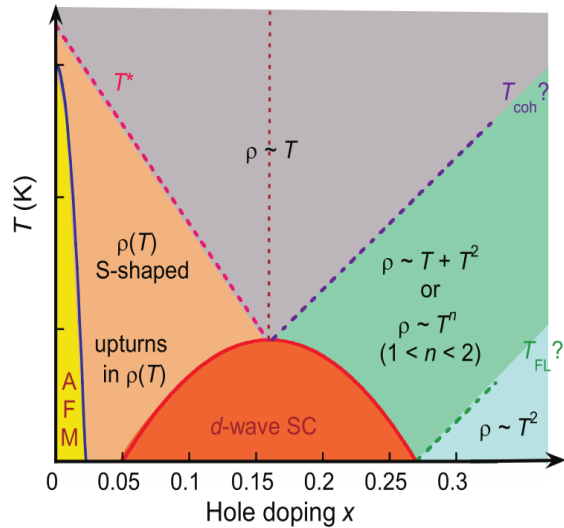


K. Hashimoto et al., Science (2012)

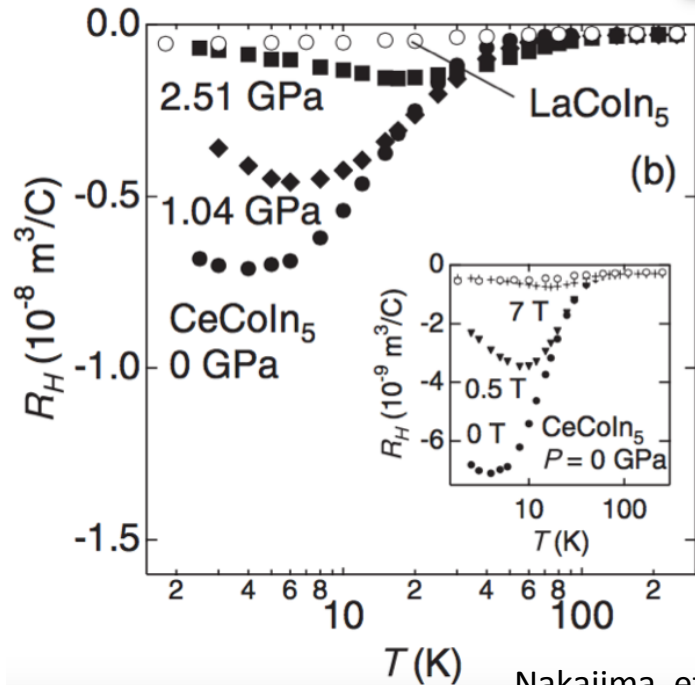


Kashahara et al., Phys. Rev. (2010)

(Nonmagneto-)Transport in High- T_c s

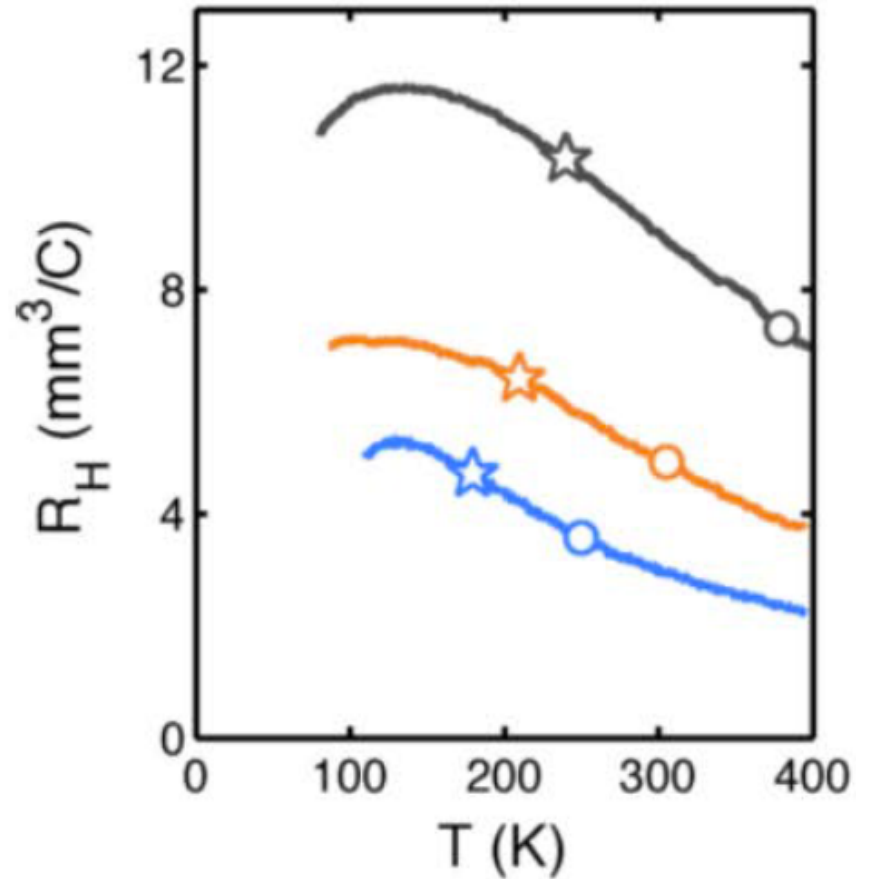


N.E. Hussey, Journal of Physics: Condensed Matter (2008)

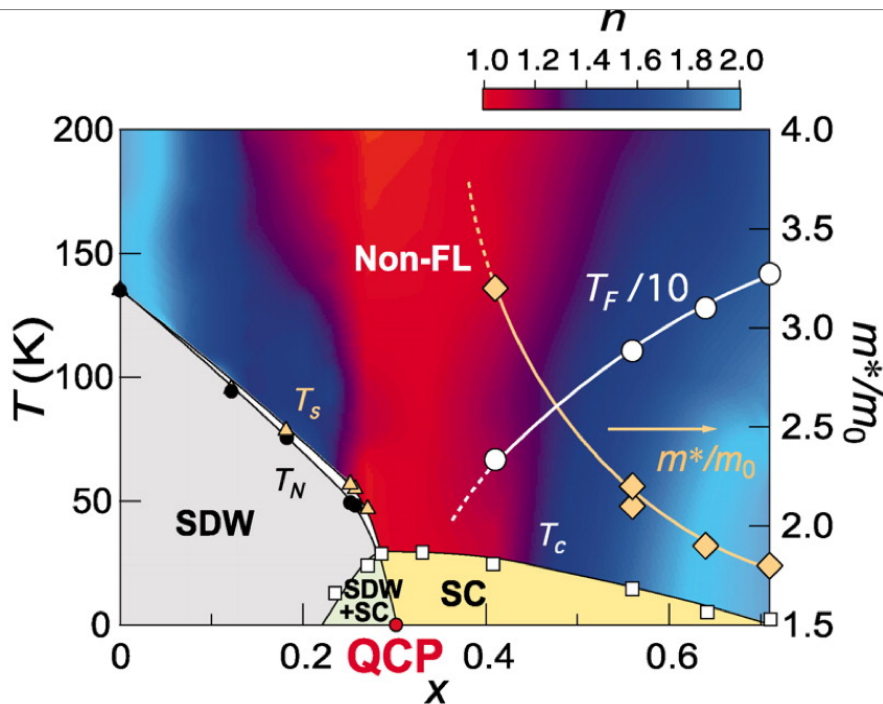


Nakajima, et al, JPSJ (2006)

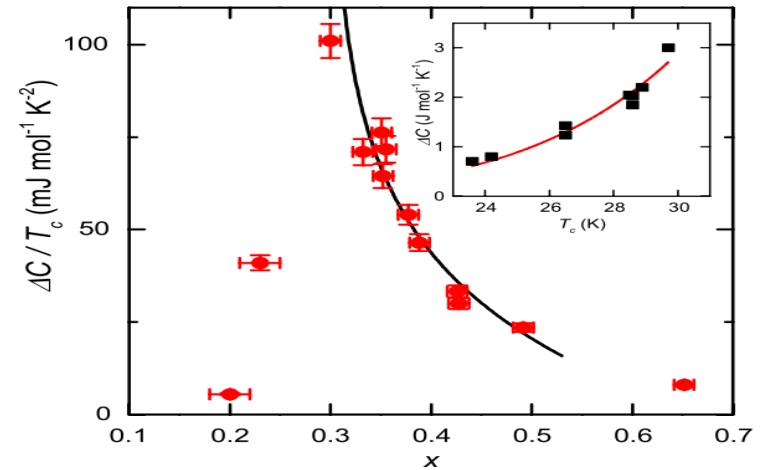
Hg-1201, from Barisic, et al, PRL (2016)



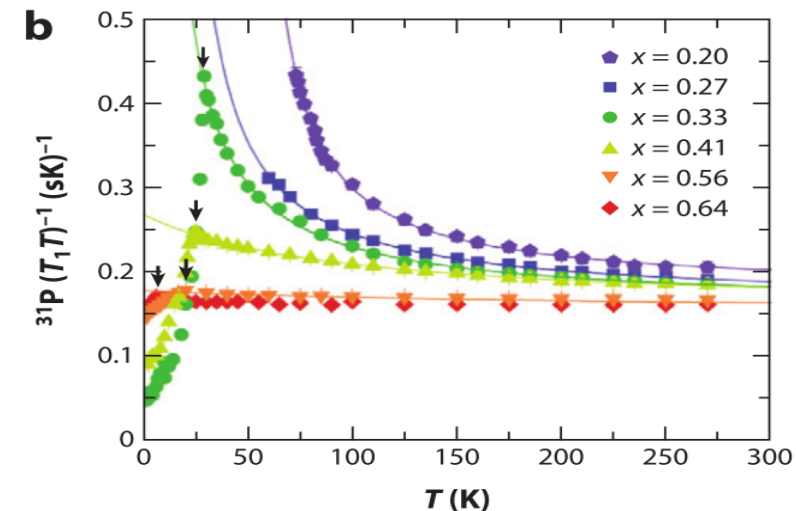
BaFe₂(As_{1-x}P_x)₂



K. Hashimoto et al., *Science* (2012)



P. Walmsley, et al, *Phys. Rev. Lett.* (2013)



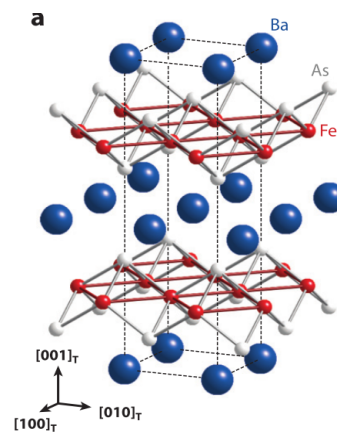
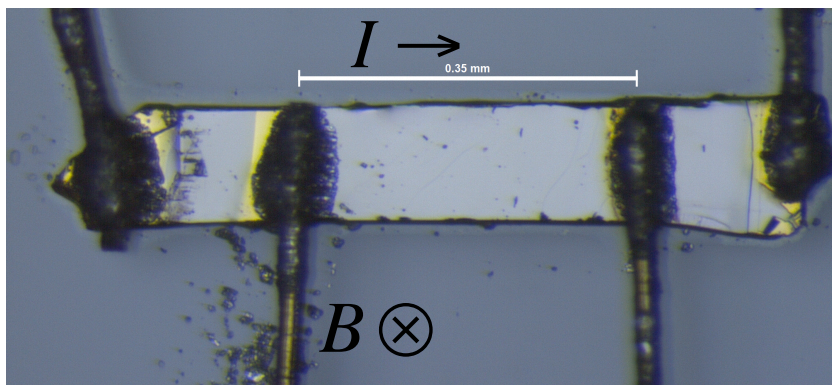
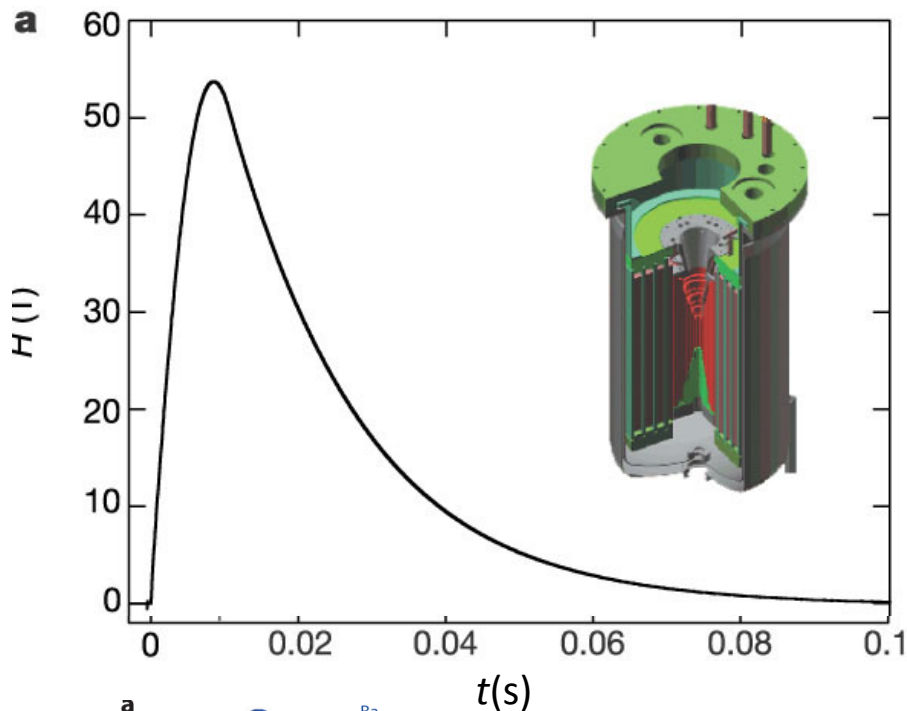
Y. Nakai, et al., *Phys. Rev. Lett.* (2010)

- Signs of strong correlations/criticality are seen in a number of probes: Quantum oscillations, heat capacity, penetration depth, NMR, charge transport.
- Normal state at optimal doping is (barely) accessible at low temperatures: $H_{c2} \sim 45$ Tesla

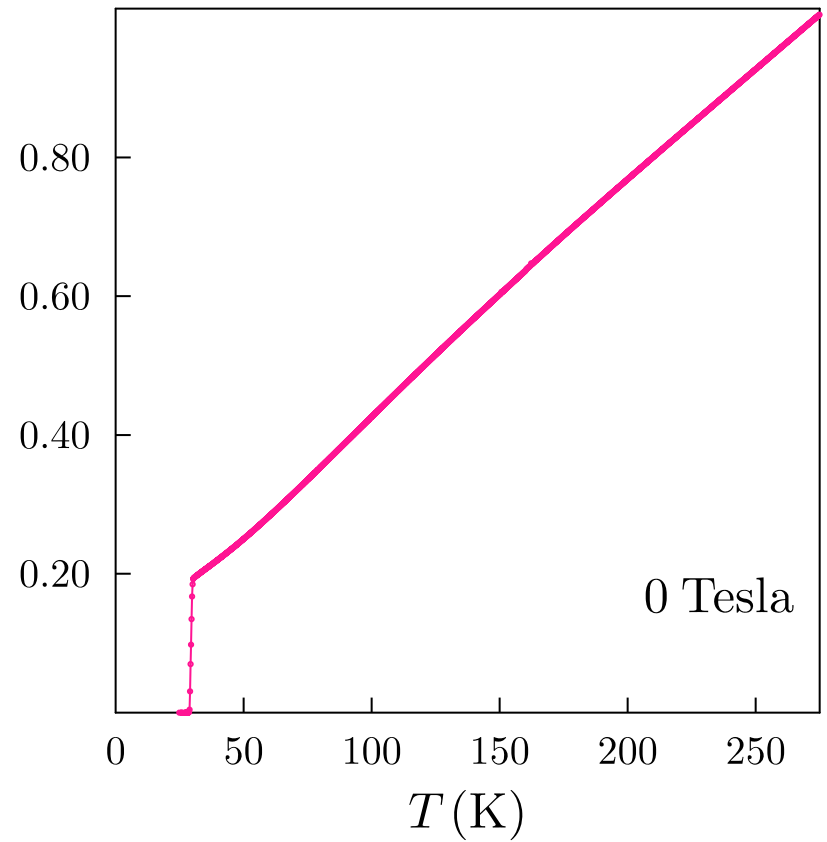
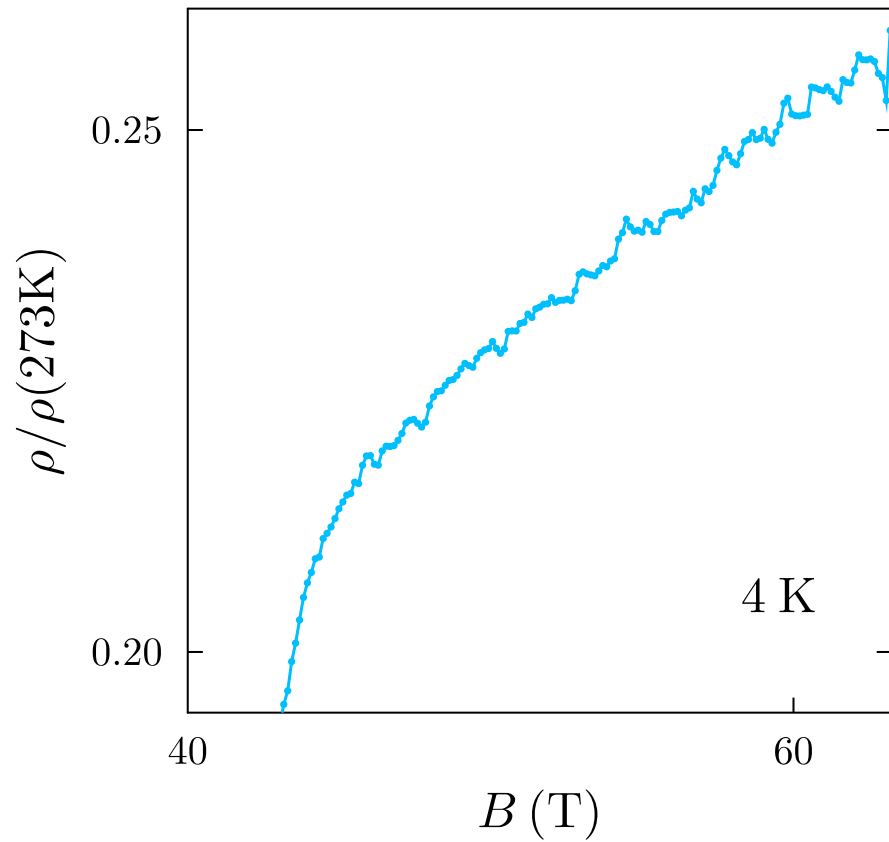
Ultra-High Magnetic fields: NHMFL Pulsed Field Facility at Los Alamos National Lab



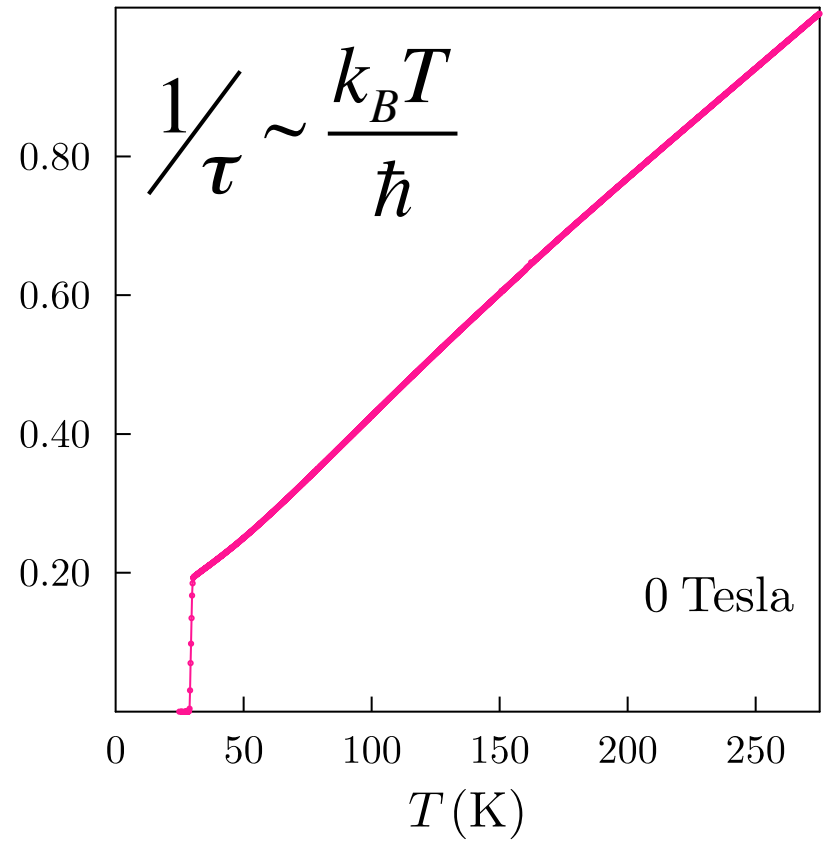
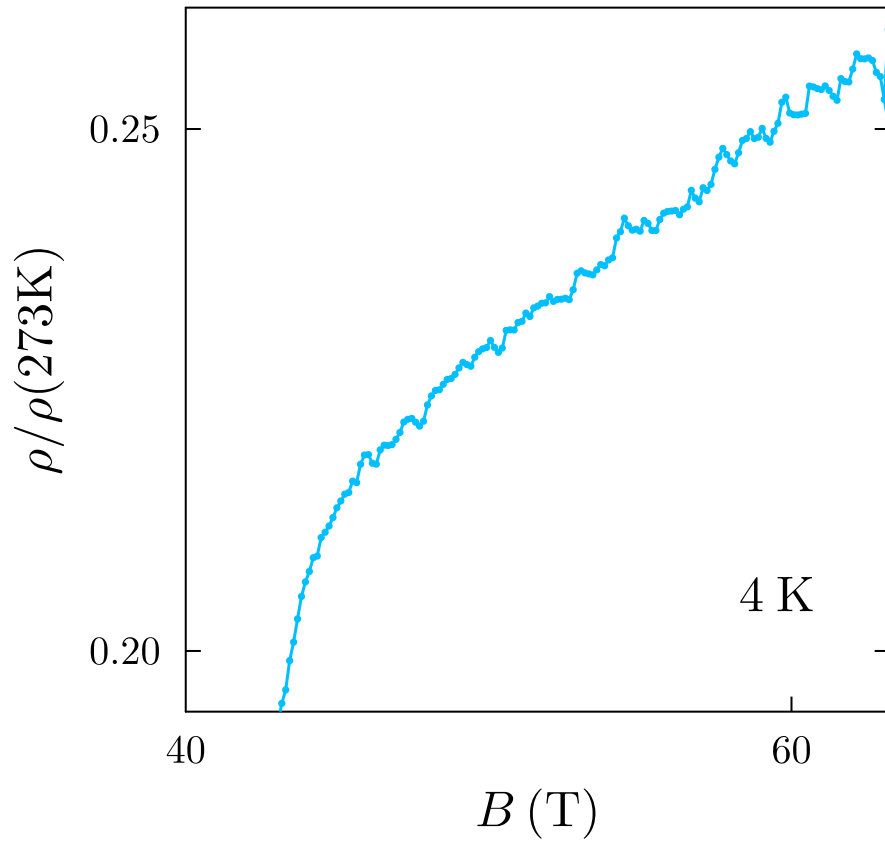
<http://www.lanl.gov>



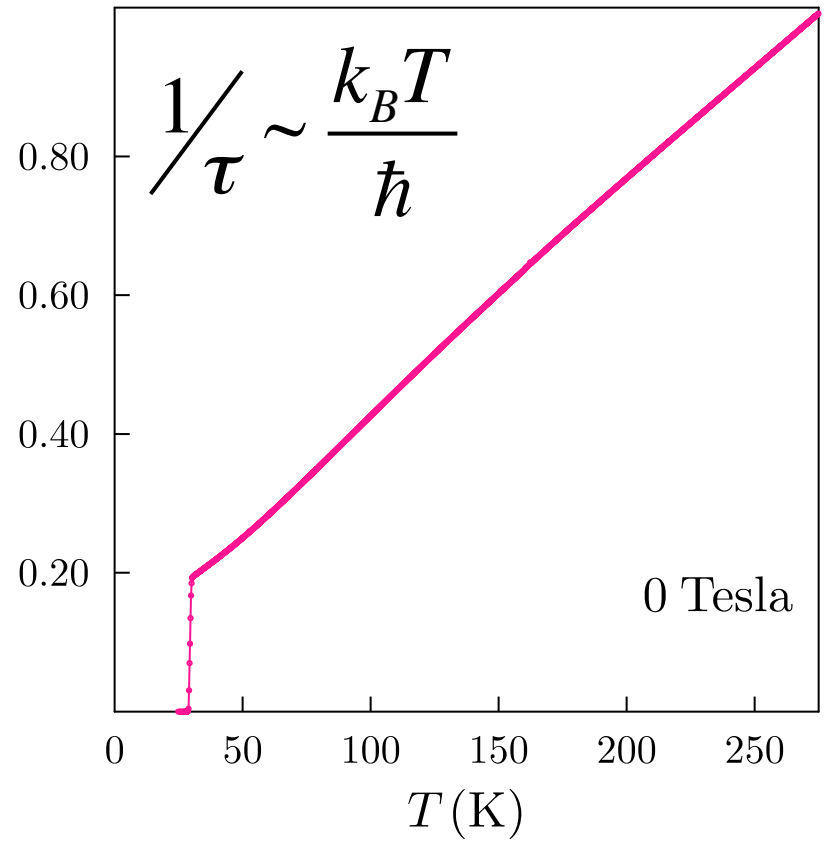
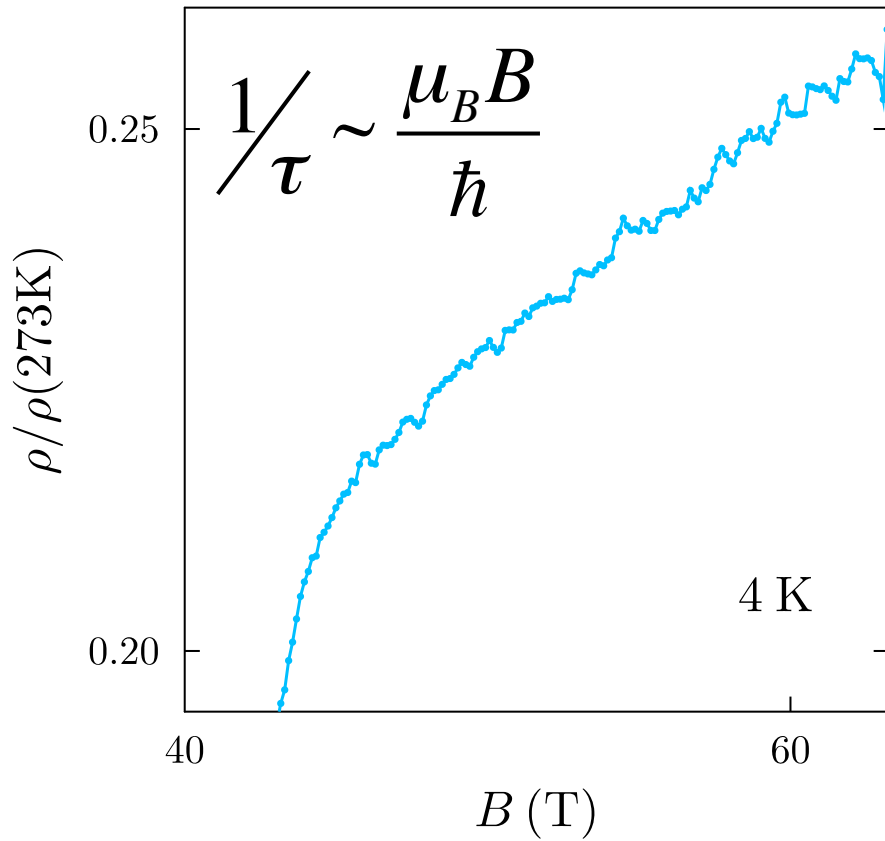
Low-T magnetotransport



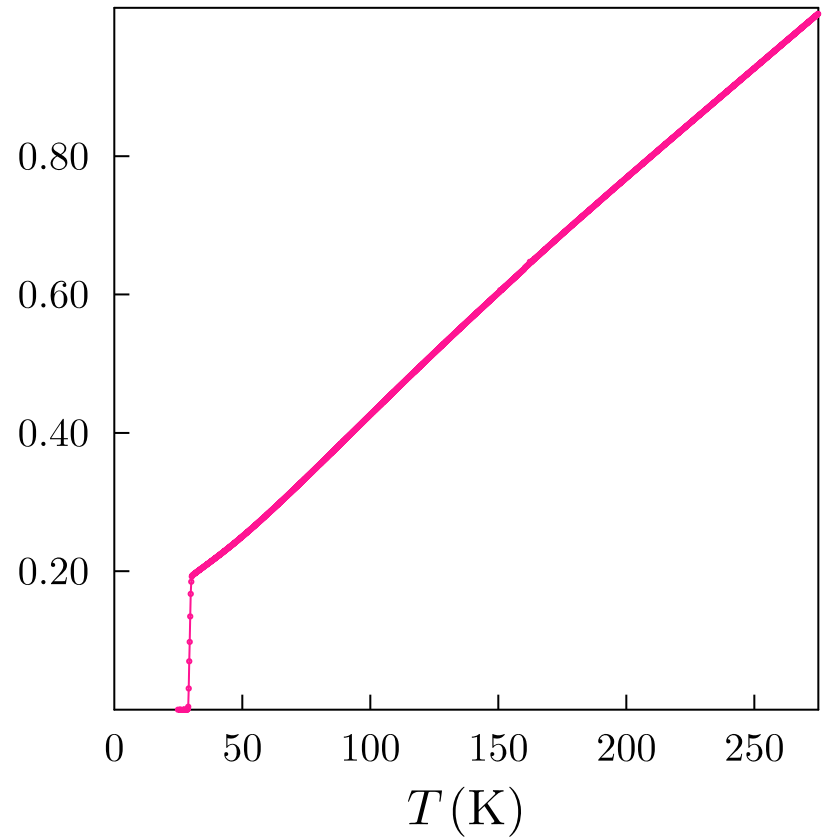
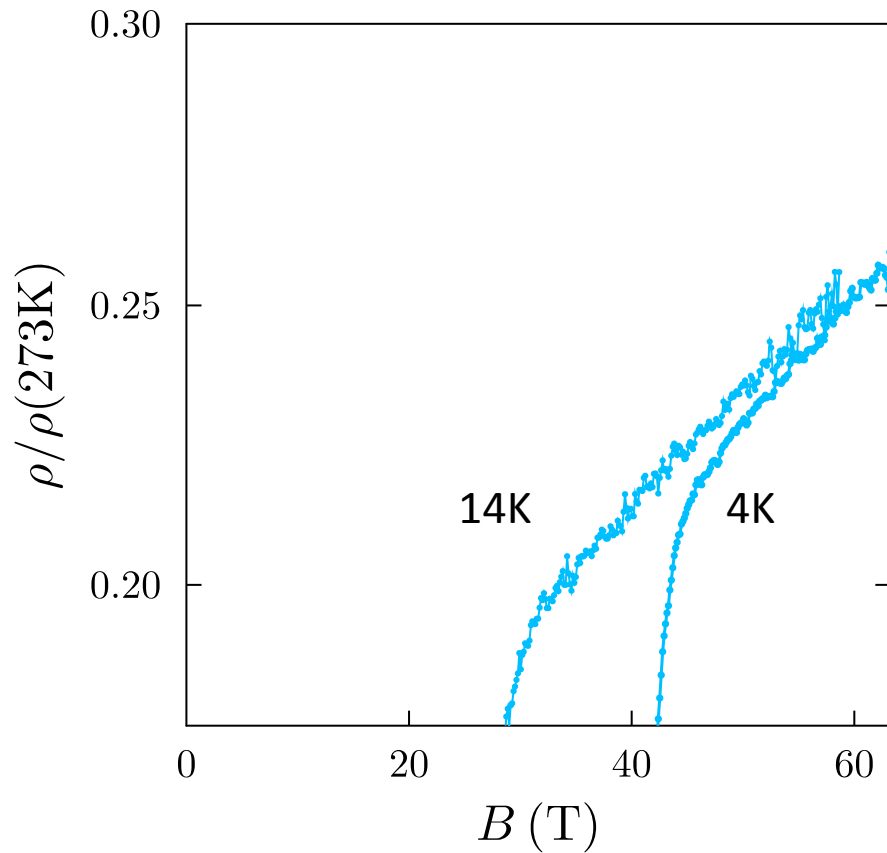
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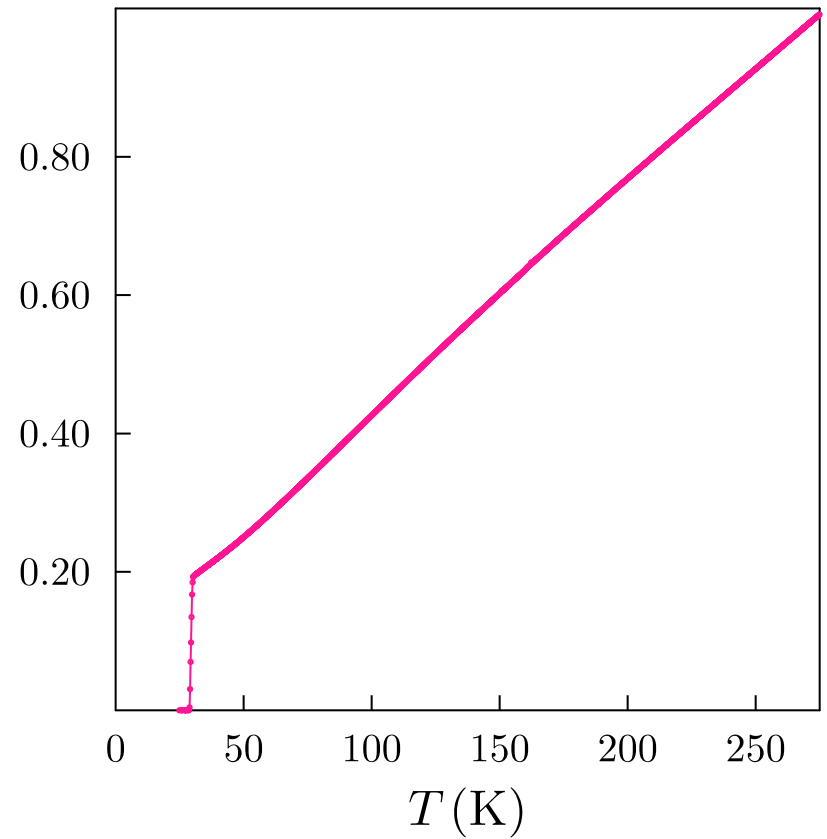
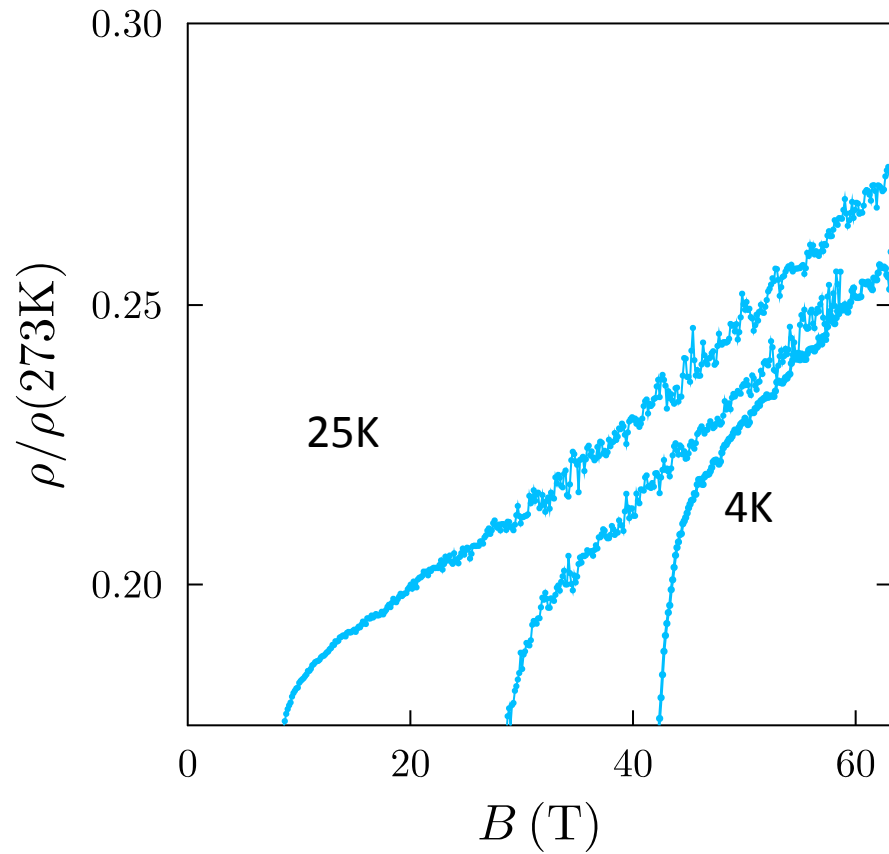
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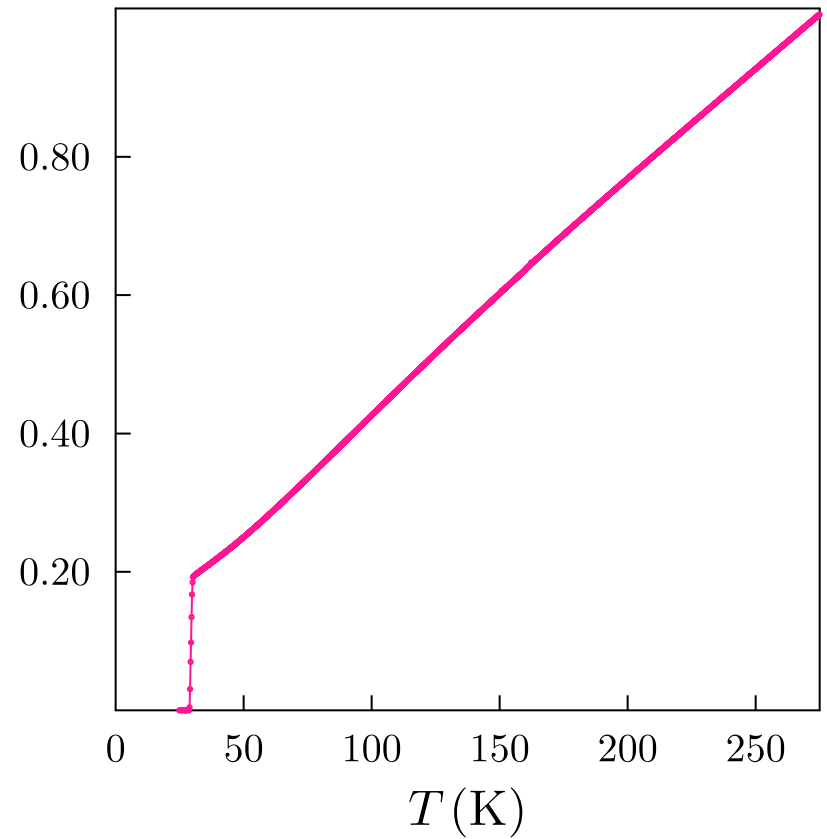
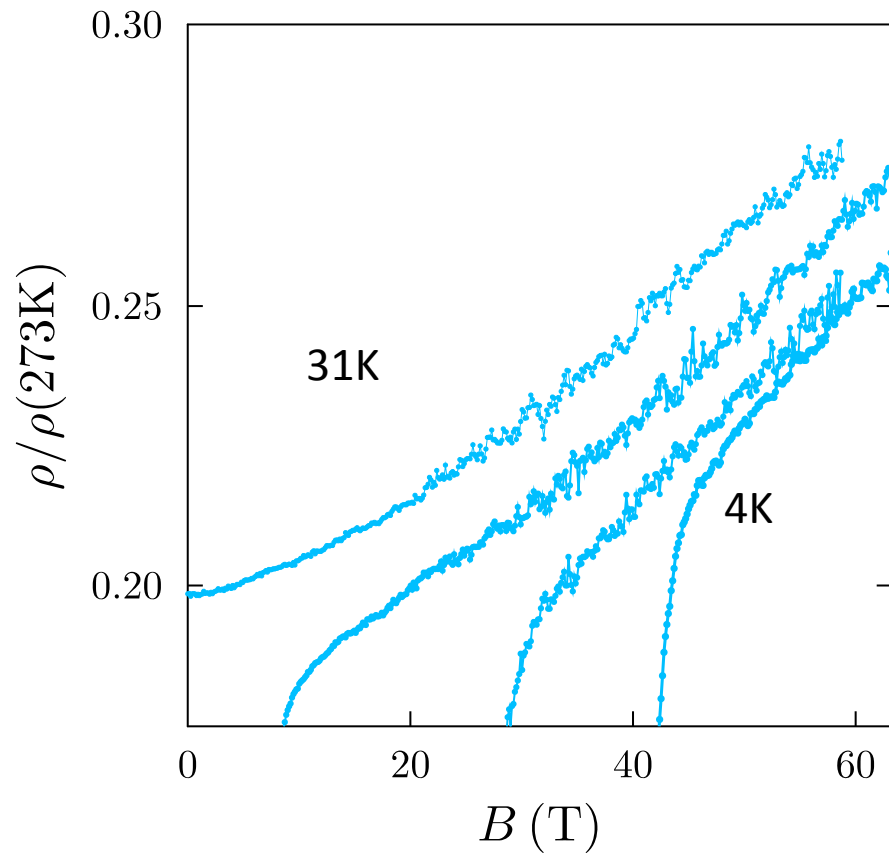
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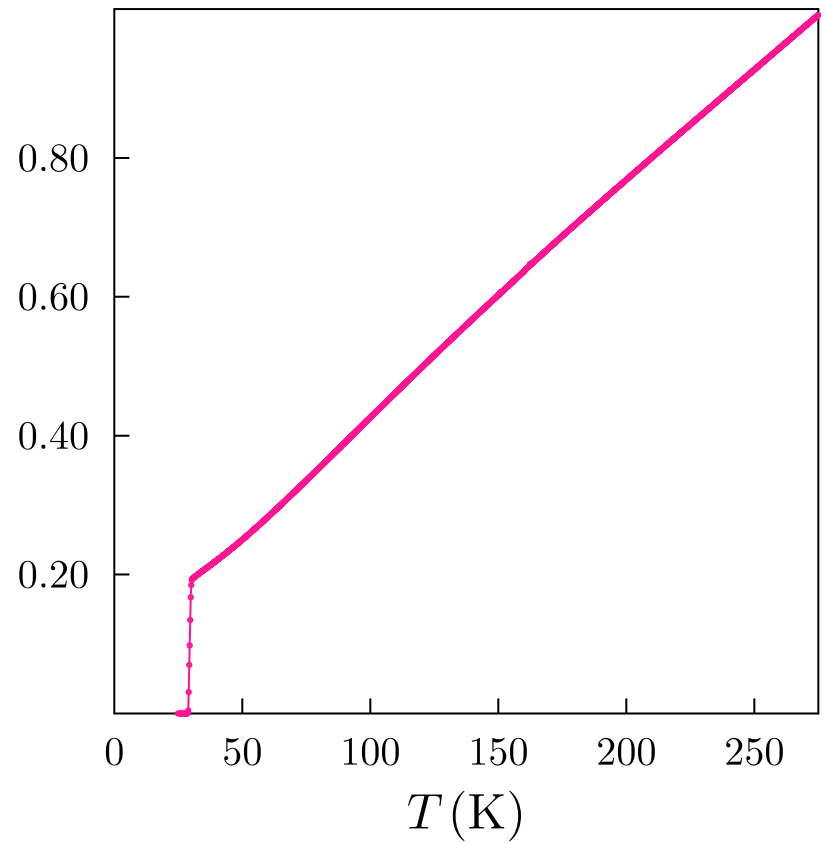
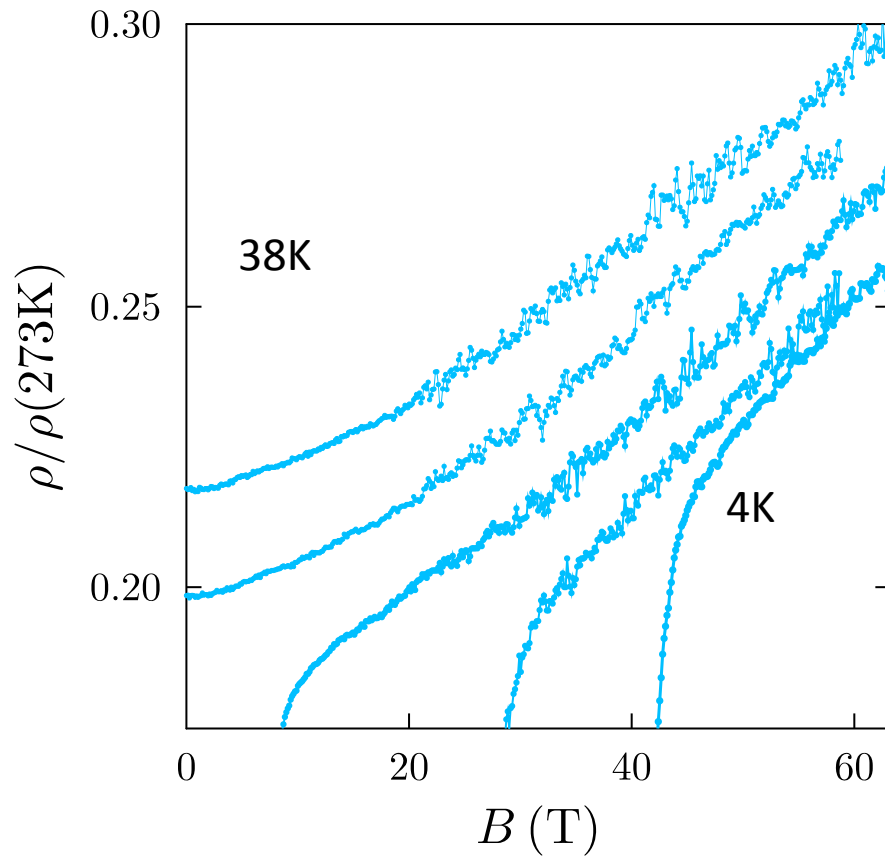
Low-T magnetotransport



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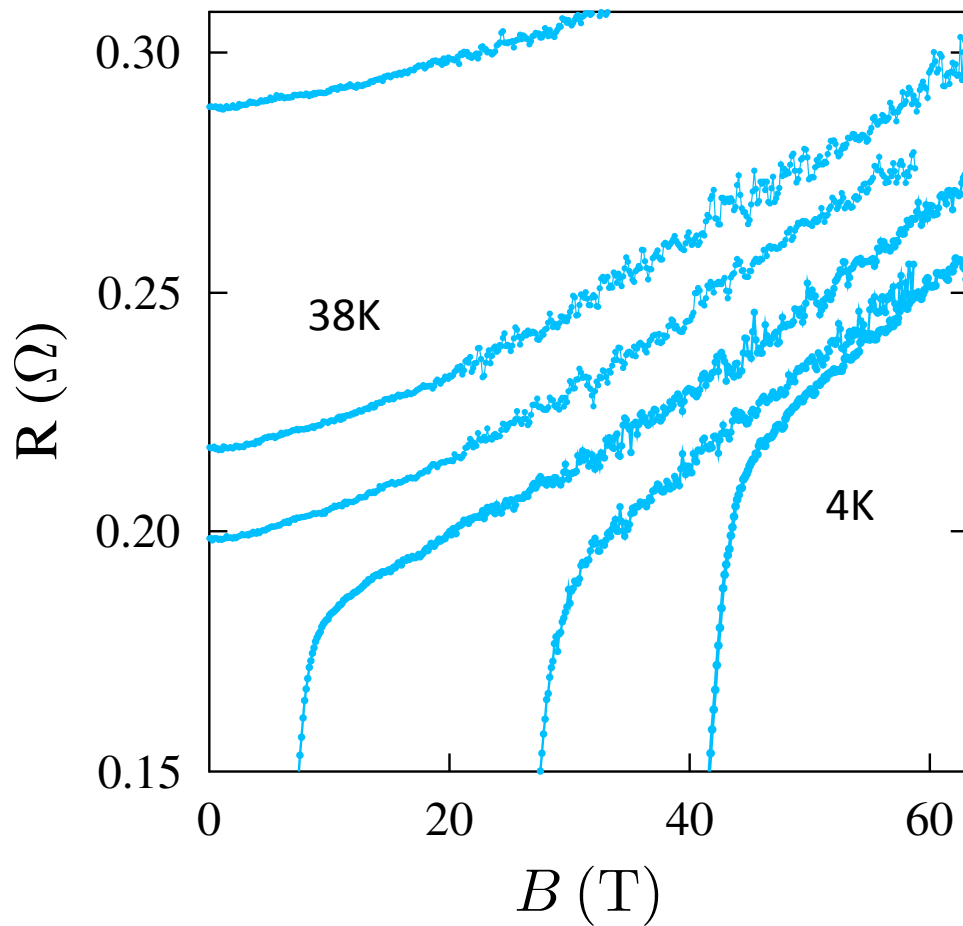


Low-T magnetotransport



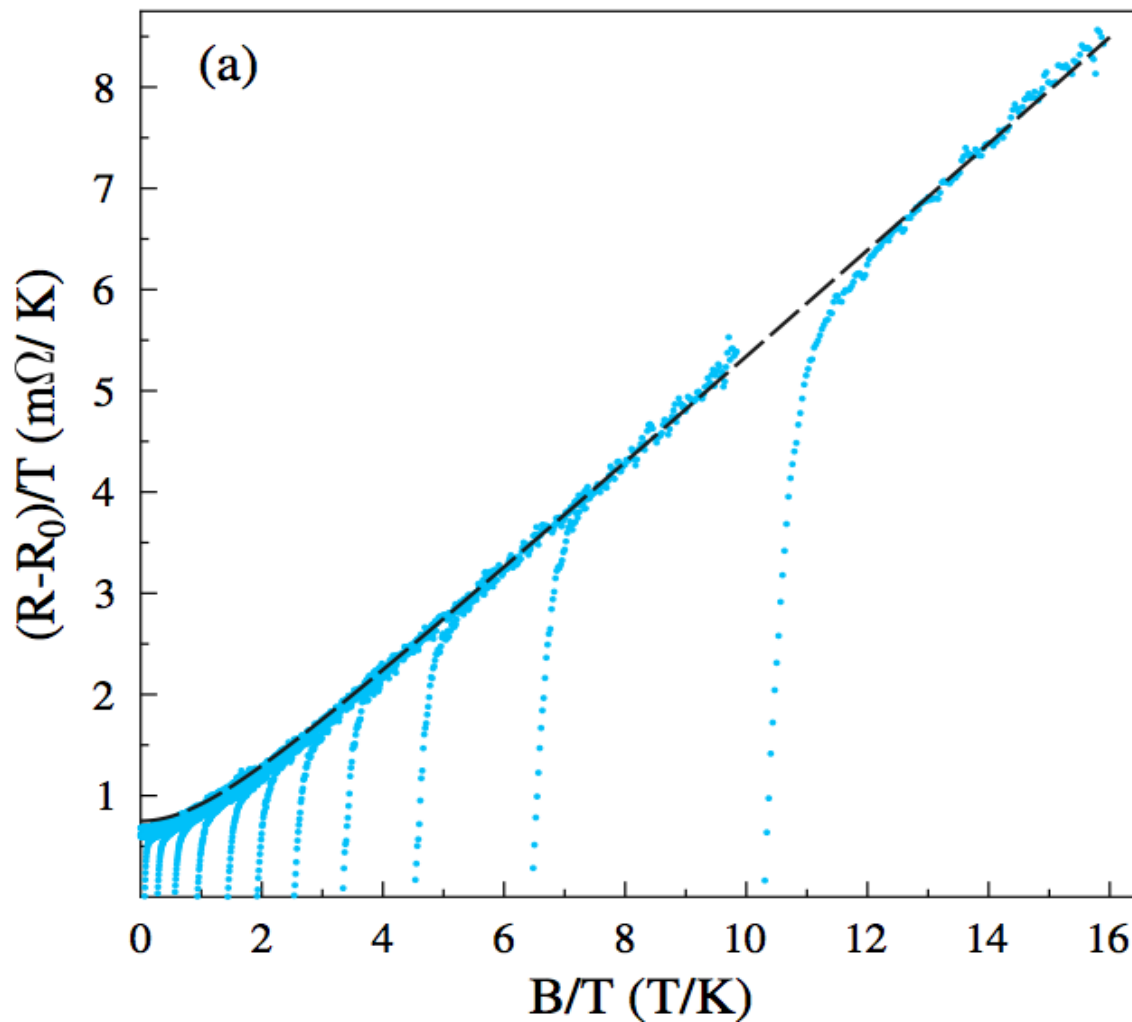
Low-T magnetotransport

$$R \propto T \times f\left(\frac{B}{T}\right)$$



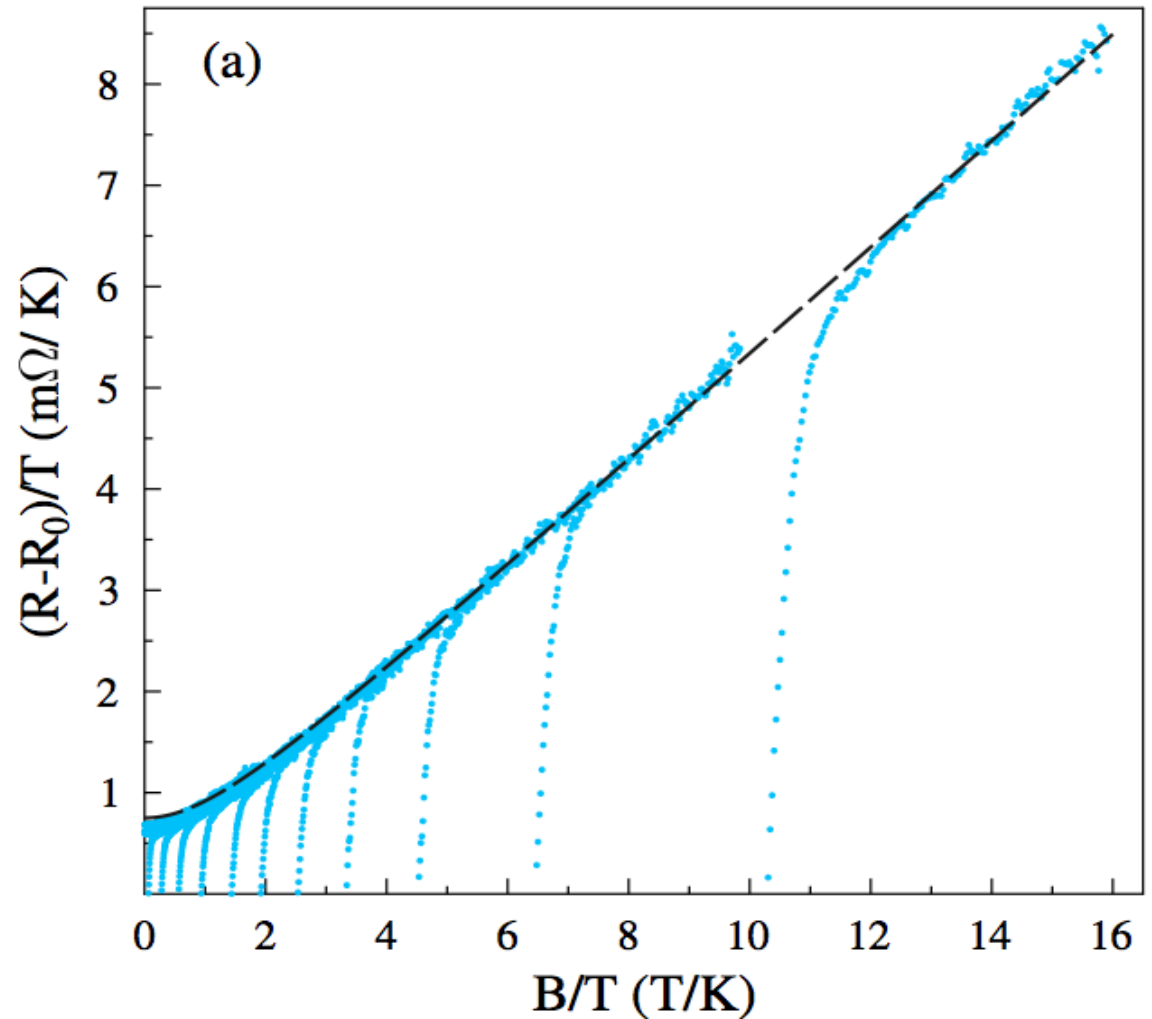
B/T Scaling in the magnetotransport

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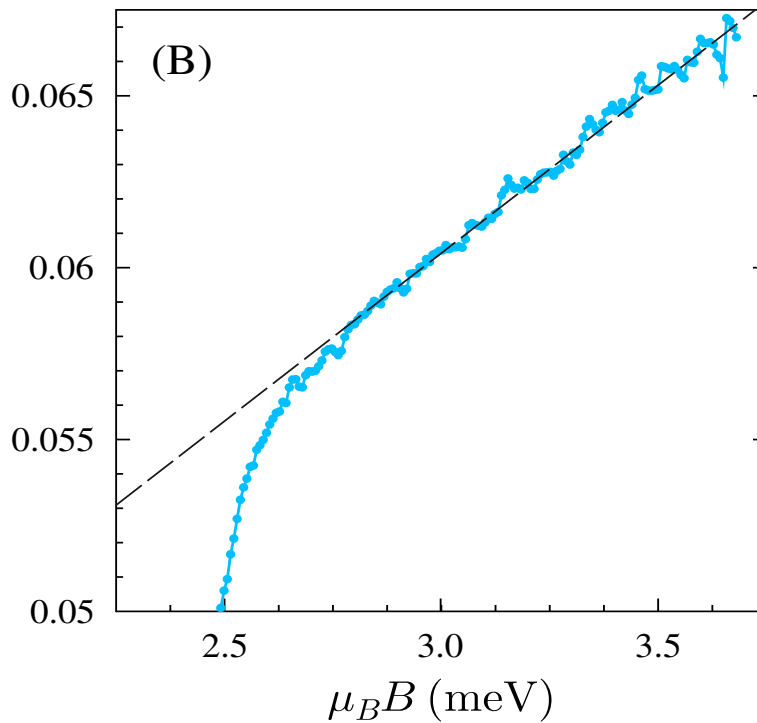
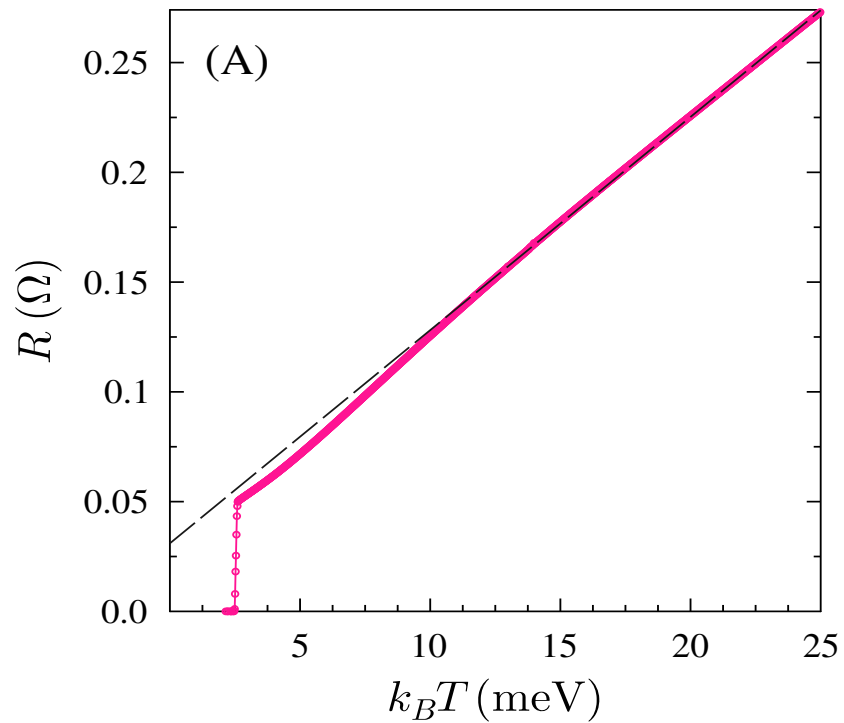
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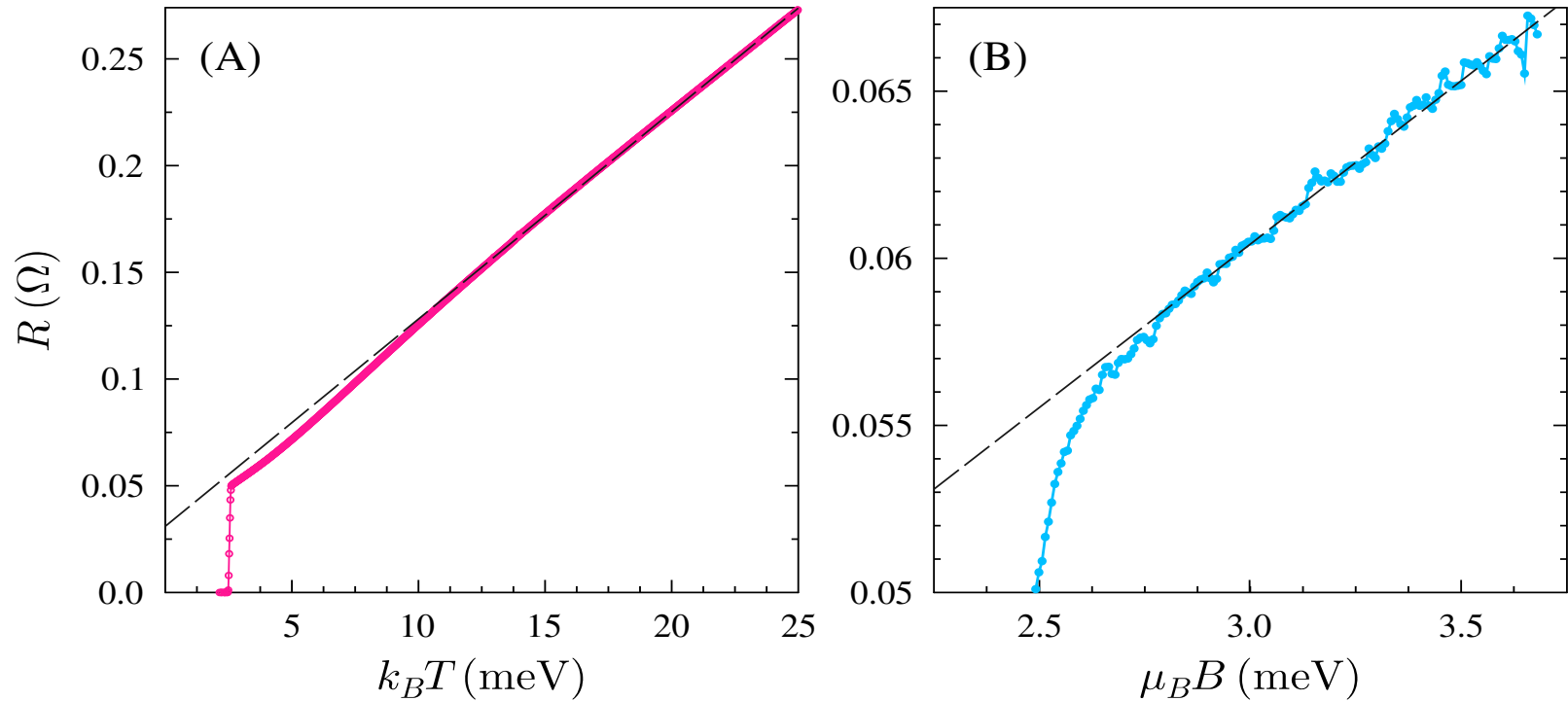


$$\rho \propto \sqrt{(\alpha k_B T)^2 + (\beta \mu_B B)^2} \equiv \Gamma$$

Intercepts and slopes



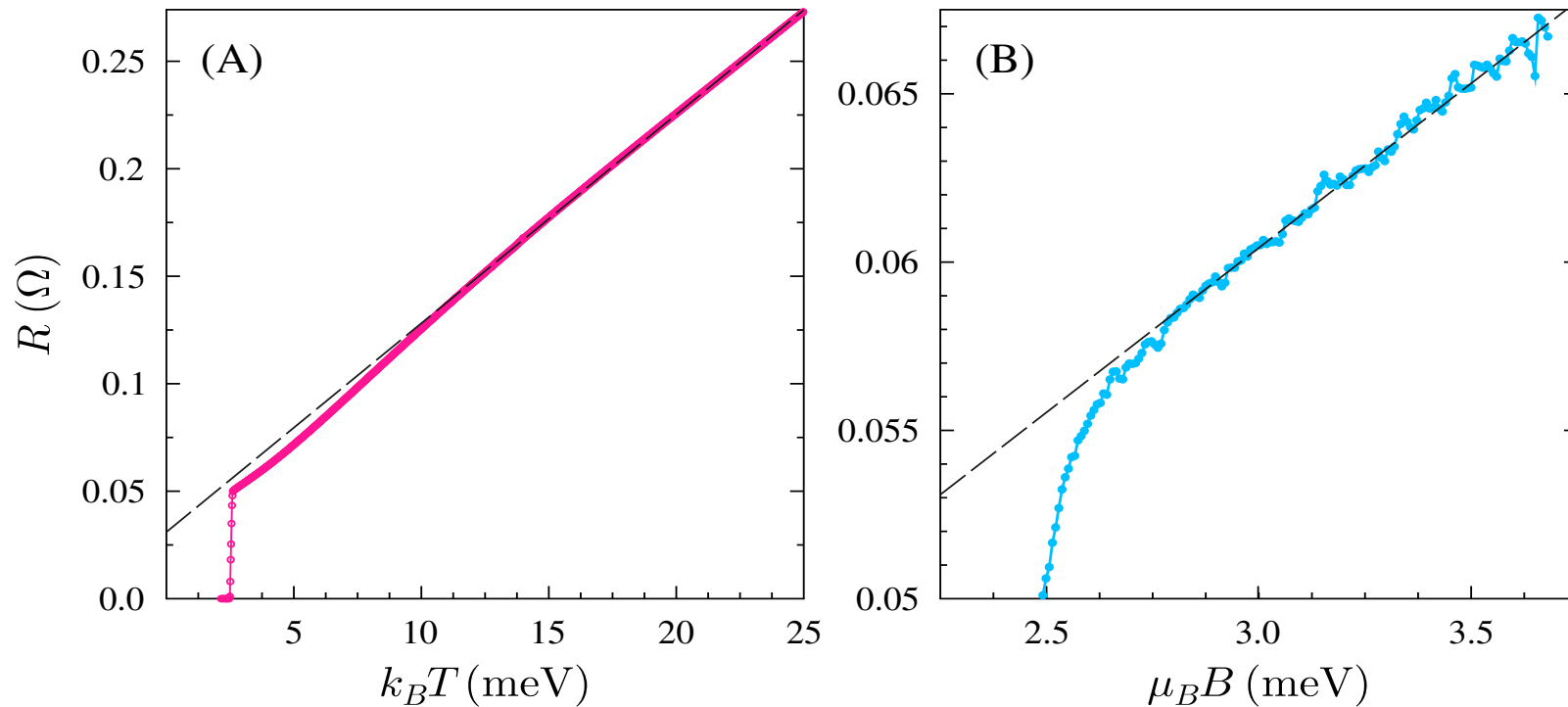
Intercepts and slopes



Within a few percent Alpha = Beta in:

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Intercepts and slopes

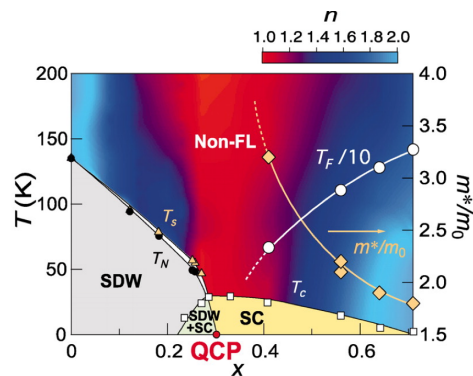
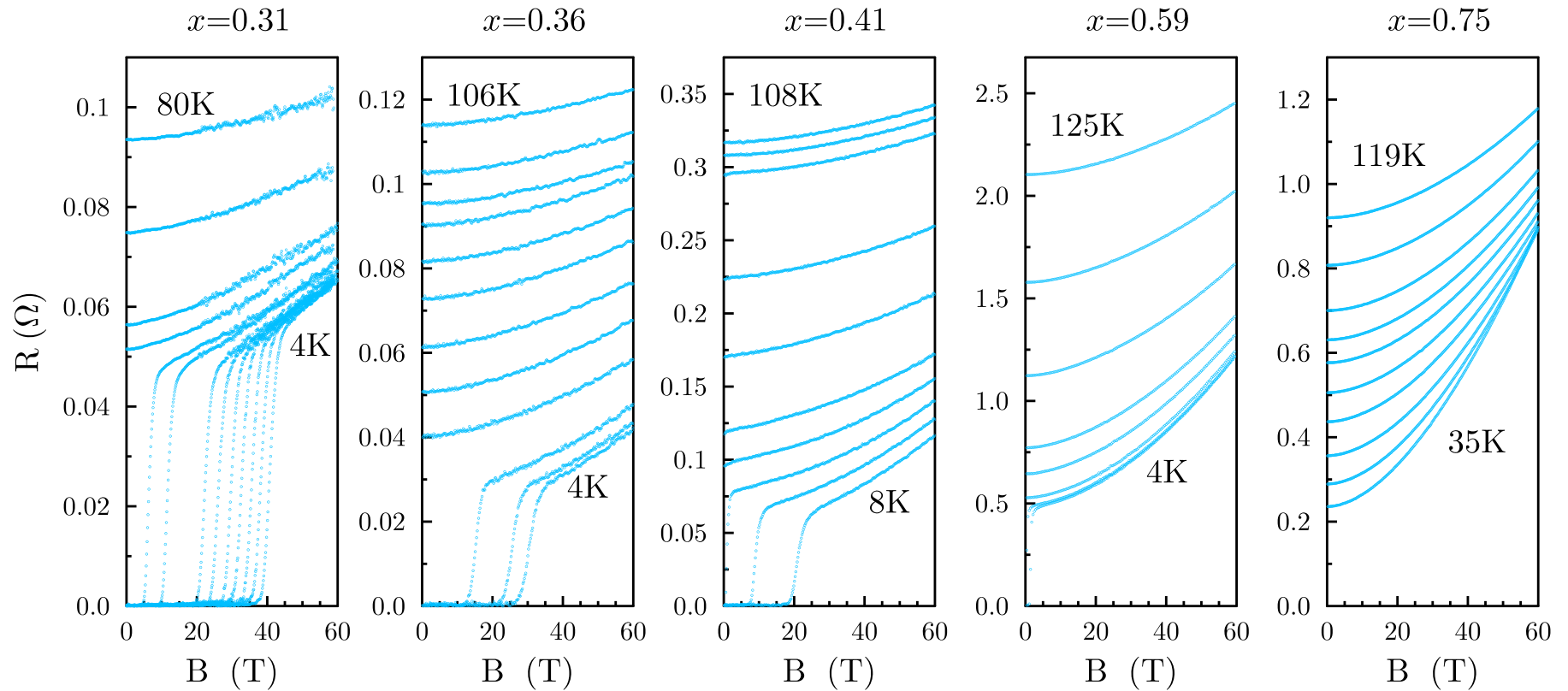


Within a few percent $\text{Alpha} = \text{Beta}$ in:

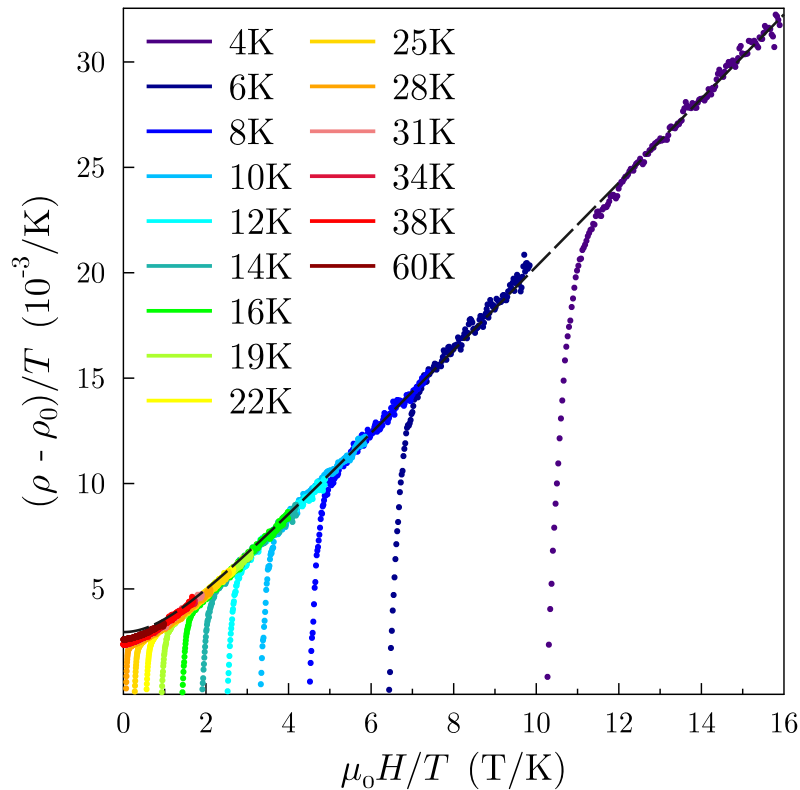
$$\rho \propto \sqrt{(\alpha k_B T)^2 + (\beta \mu_B B)^2} \equiv \Gamma$$

Intercepts at $T = 0$ or $B = 0$ are equal.

Magnetotransport across the phase diagram



Quantum Critical Magnetoresistance?



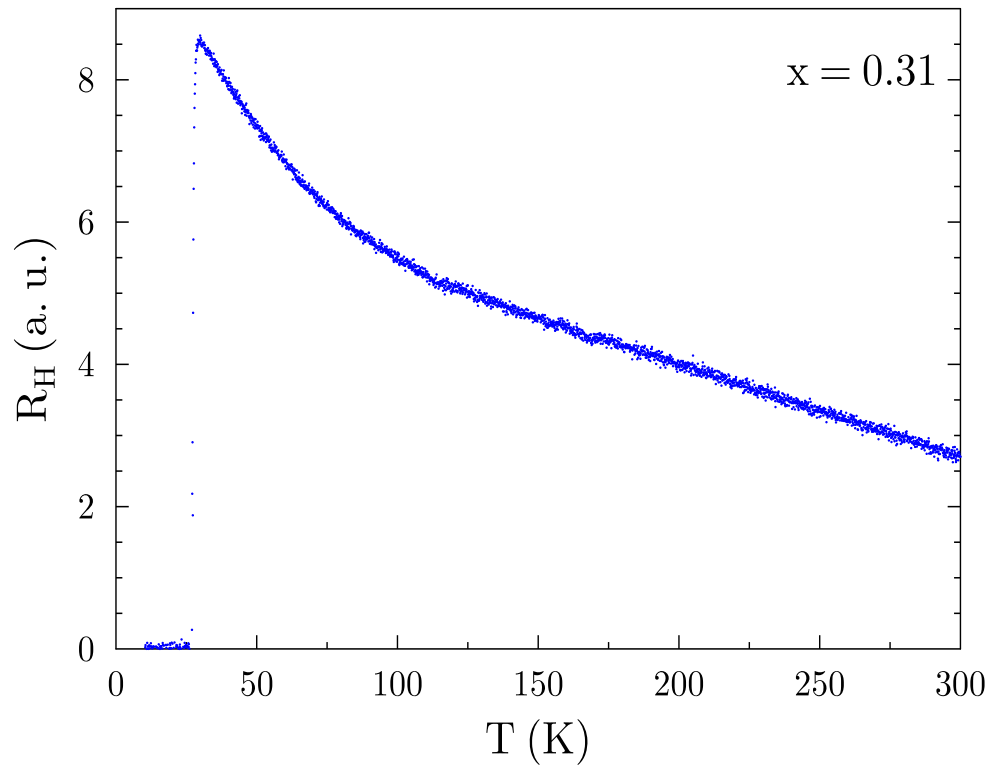
I. M. Hayes, et al, *Nature Physics* (2016)

$$\rho \propto \sqrt{(\alpha k_B T)^2 + (\beta \mu_B B)^2}$$

There is a magnetic analogue to the T-linear resistivity at optimal doping and low temperatures.

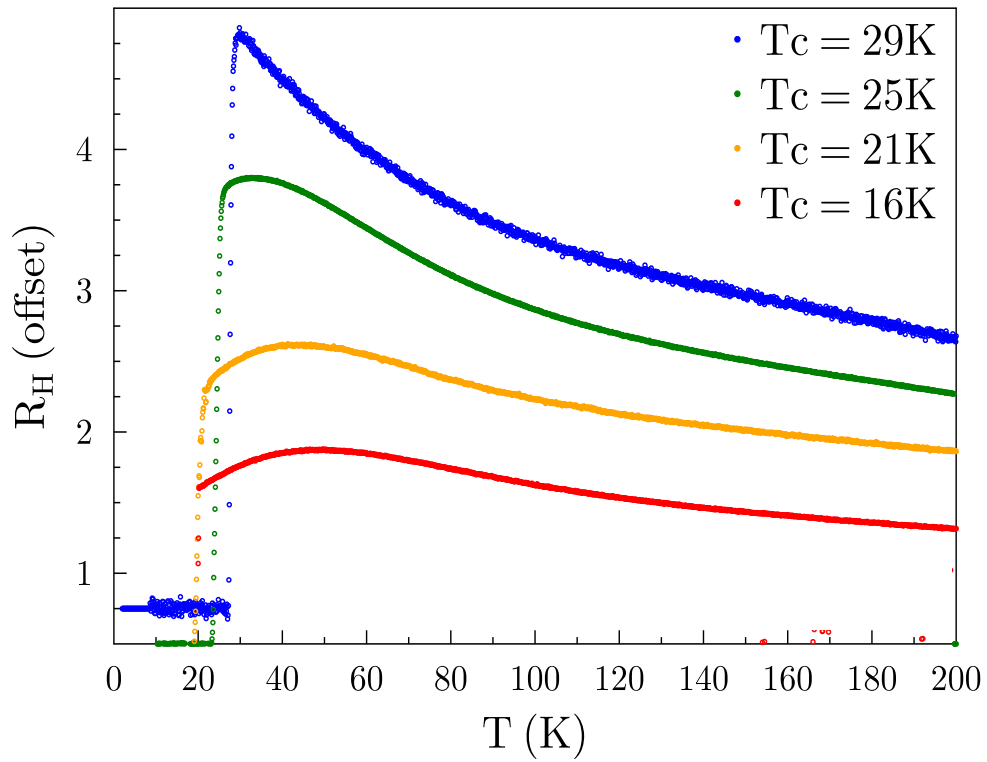
T and B have the same exponent and therefore the same scaling dimension.

Hall Effect in Ba-122



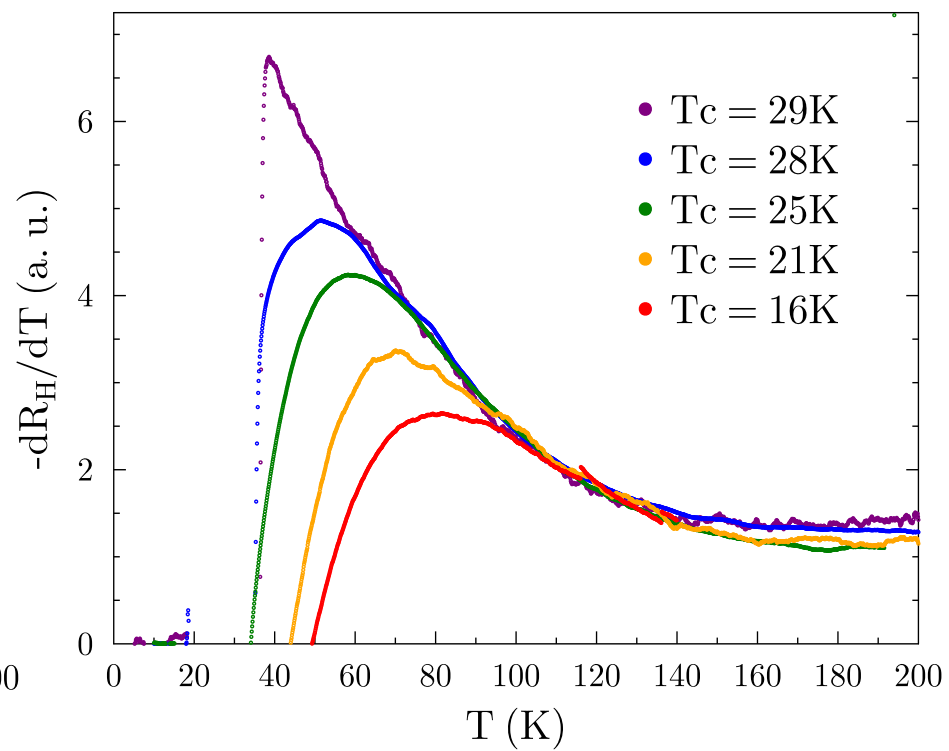
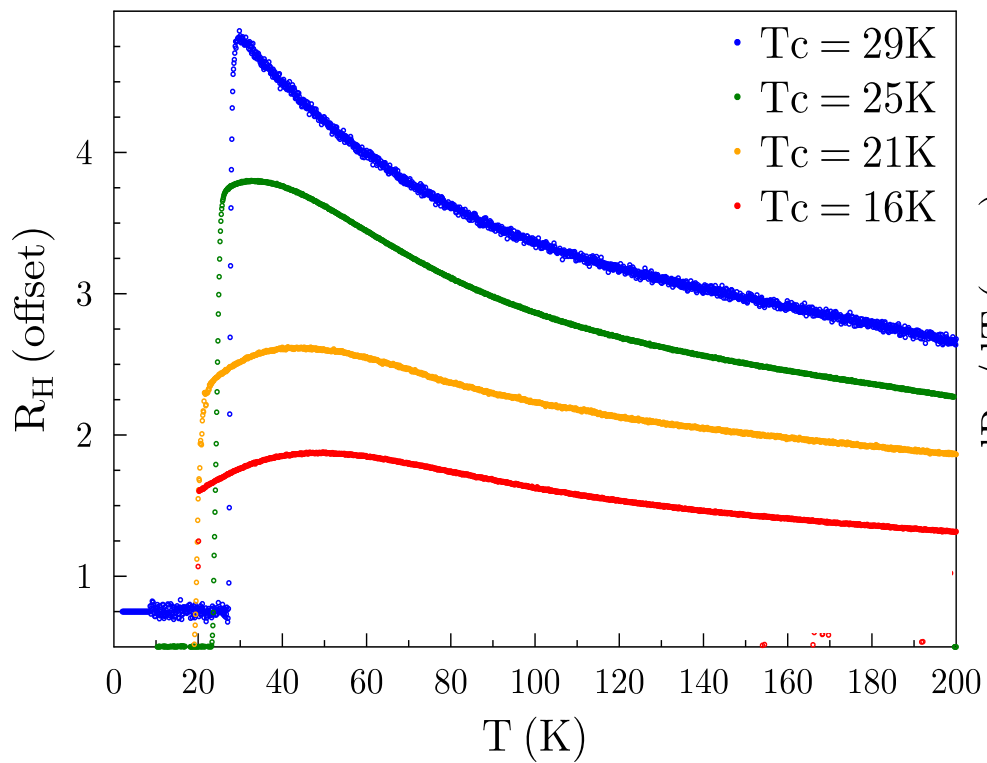
Low Field (3 Tesla)

R_H across the phase diagram

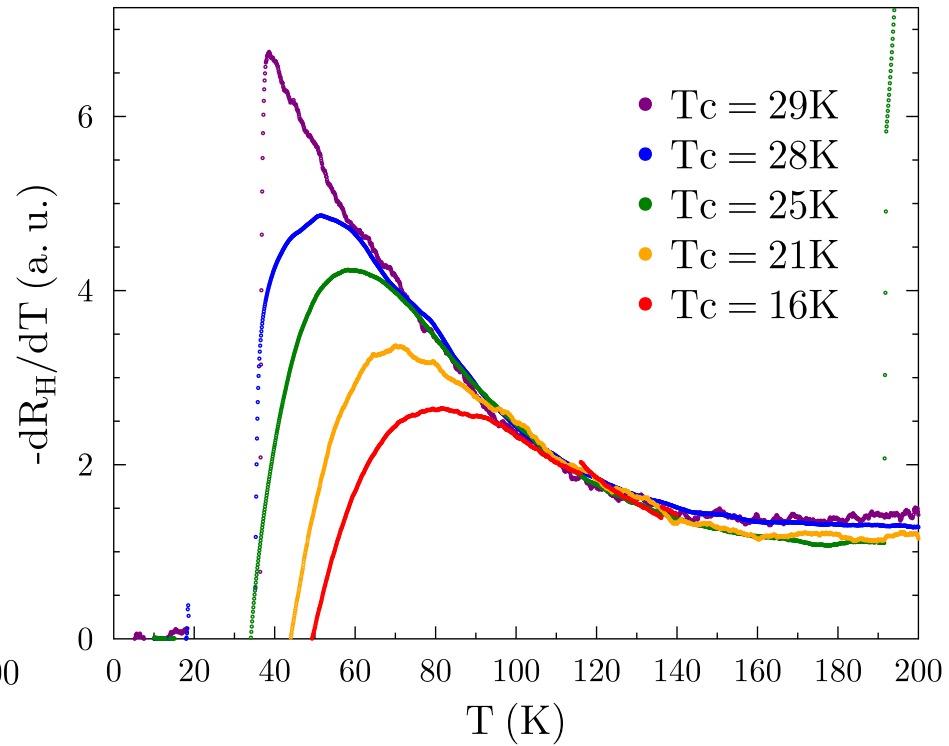
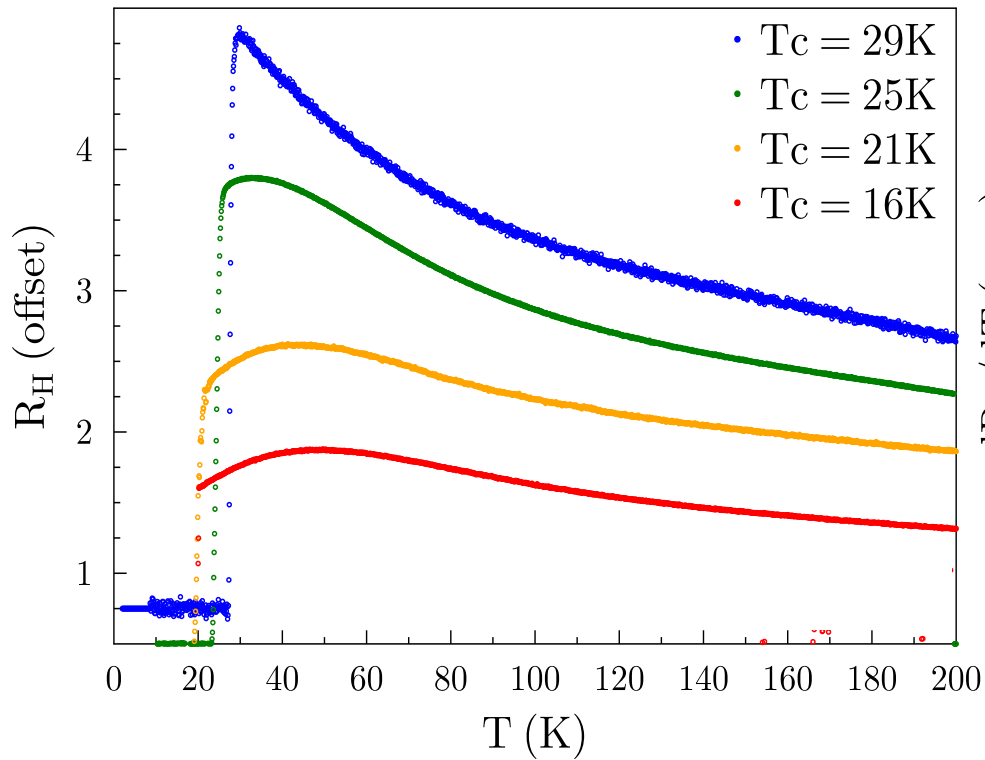


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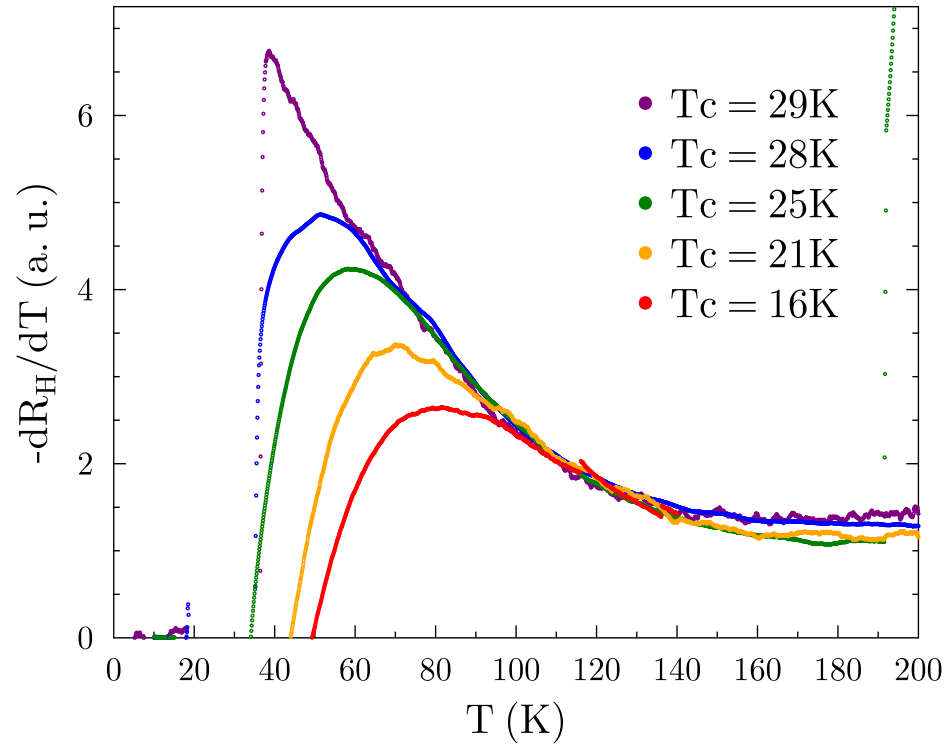
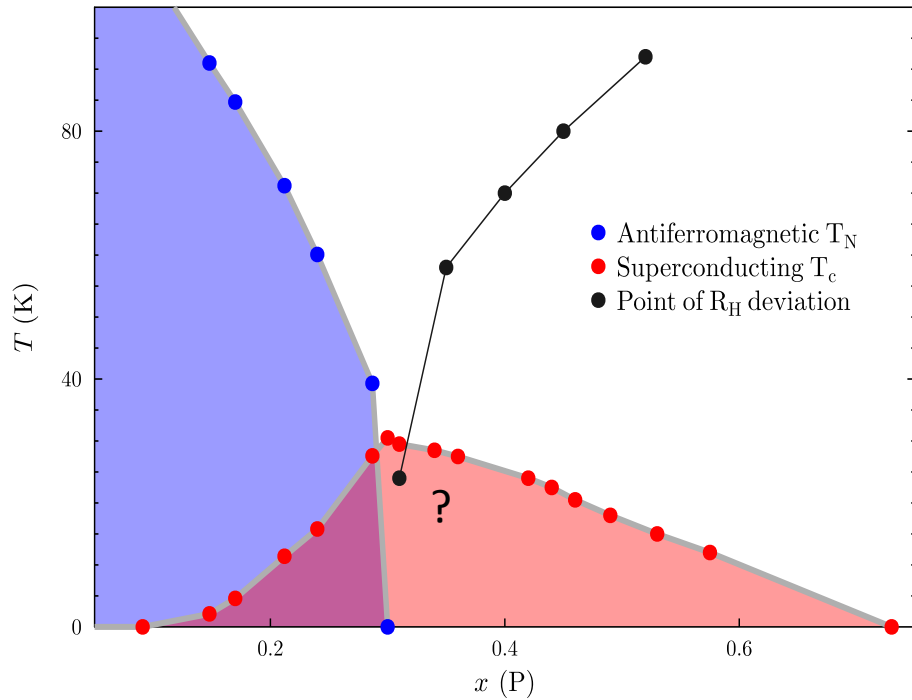


R_H across the phase diagram



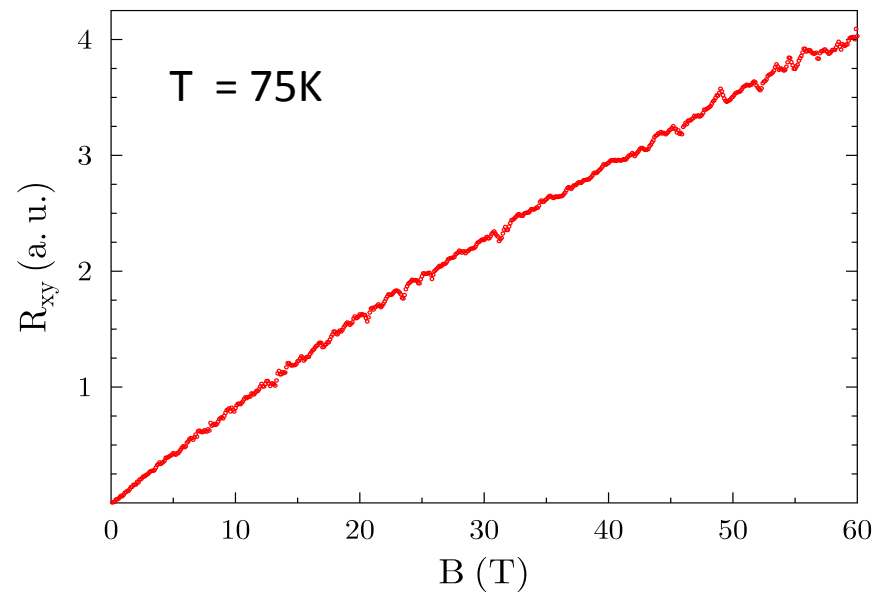
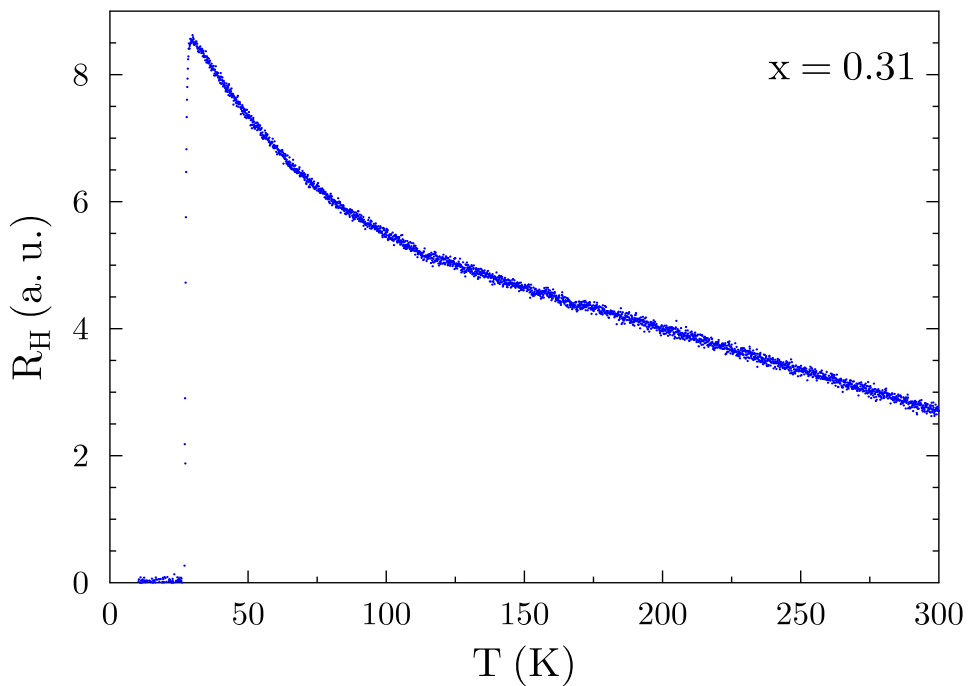
$$R_H \sim A(x)f(T)$$

R_H across the phase diagram

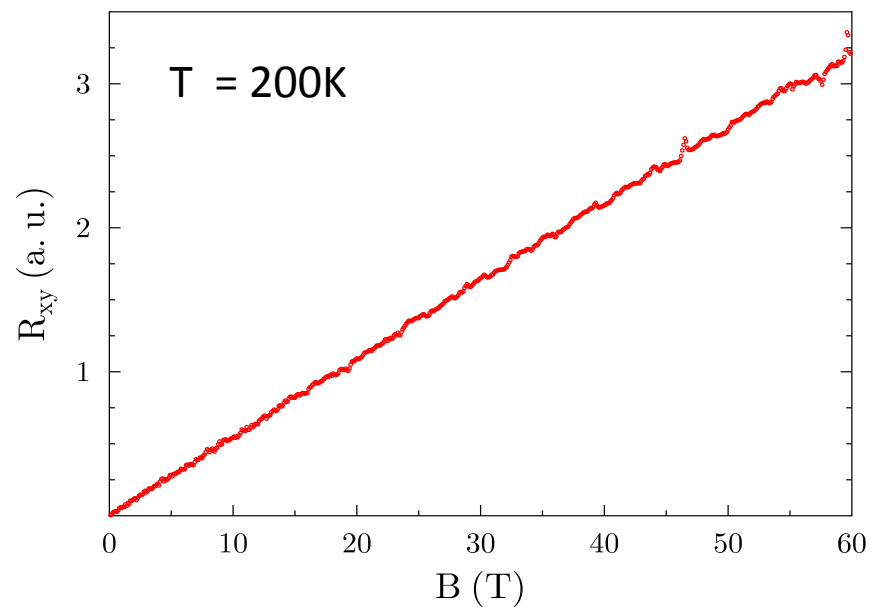
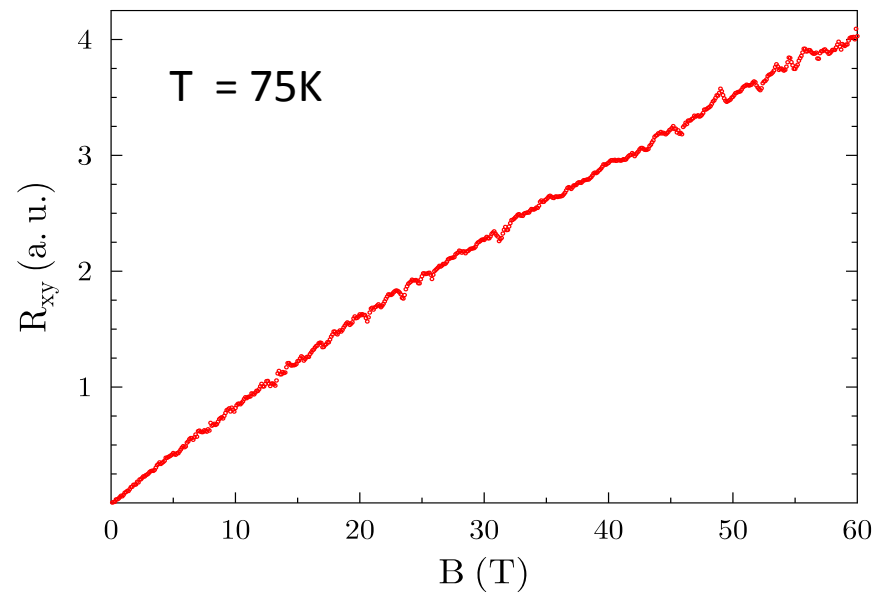
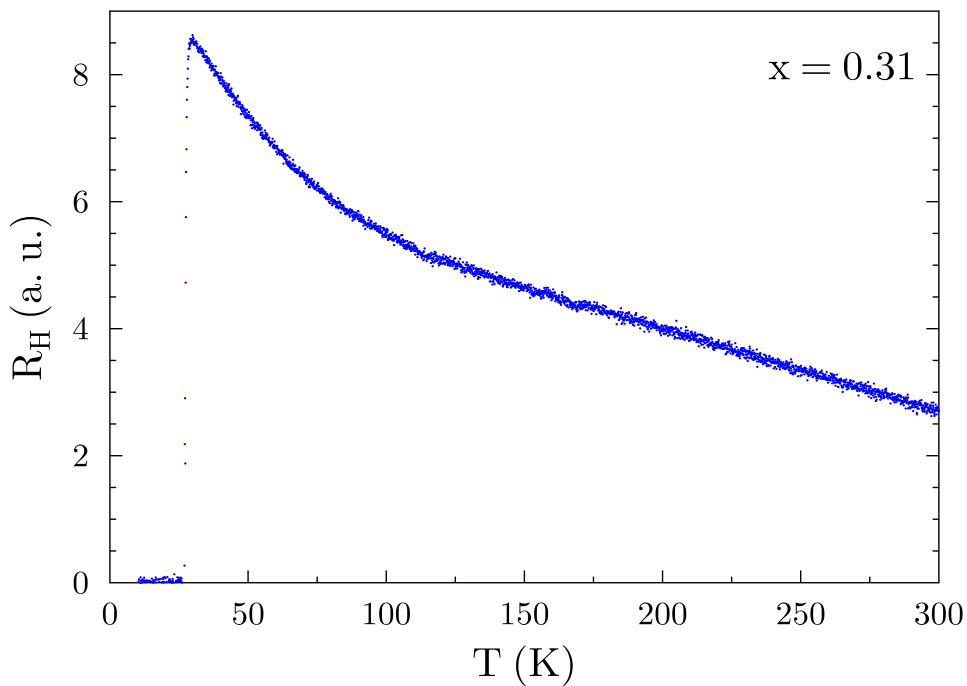


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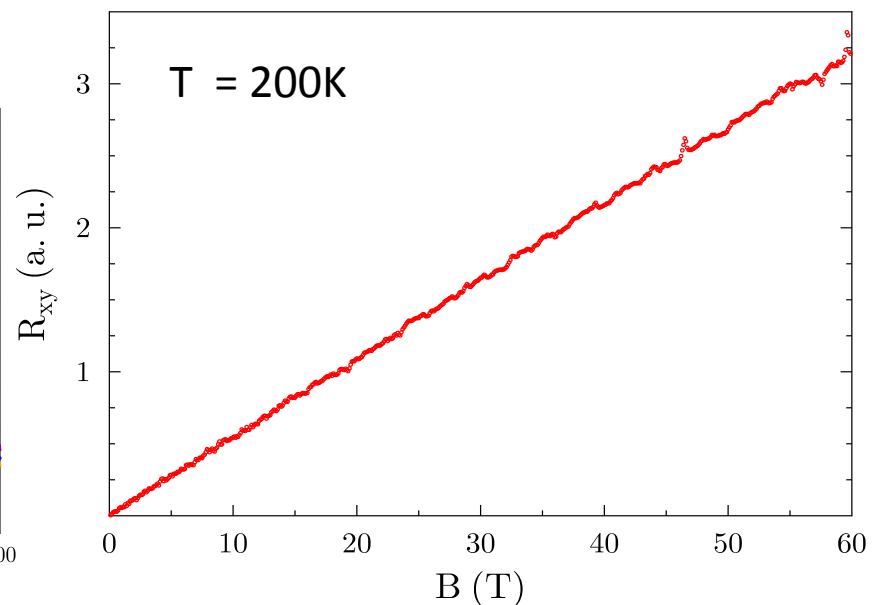
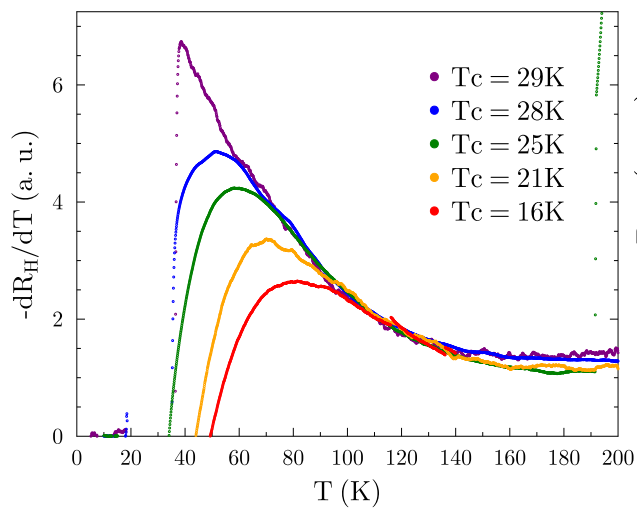
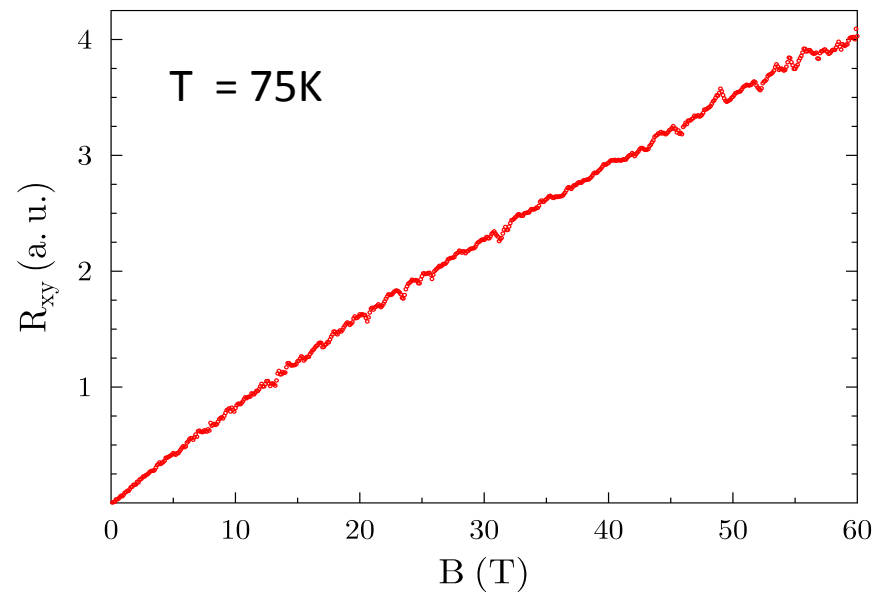
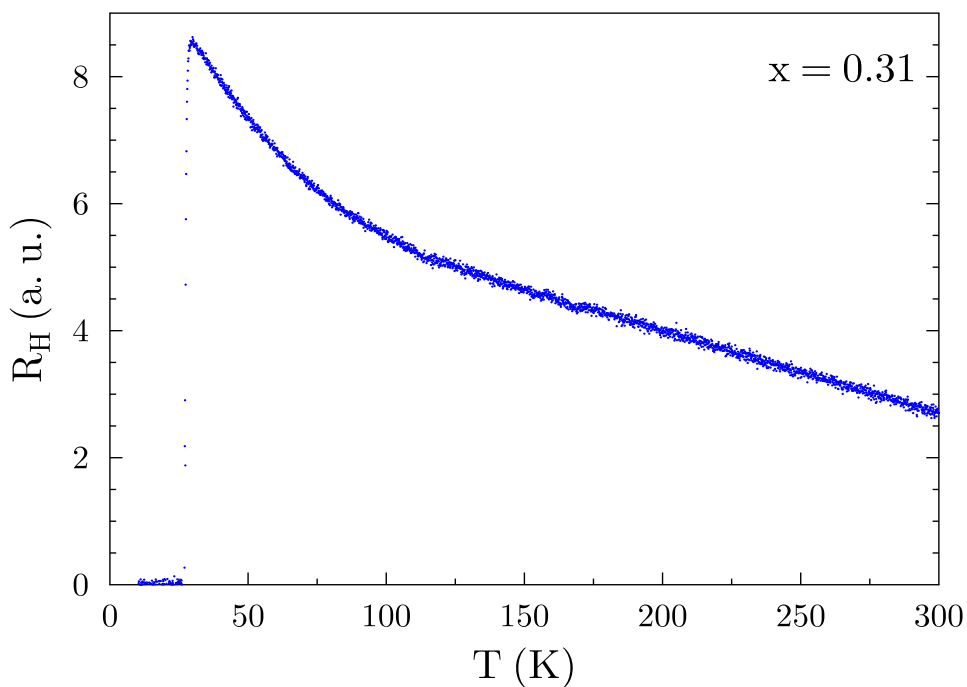
High Field Hall Effect in Ba-122



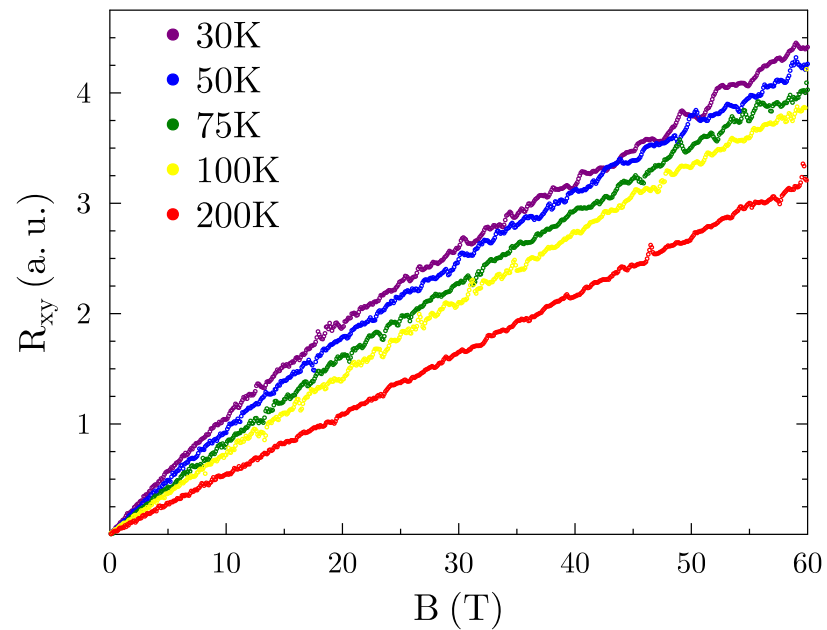
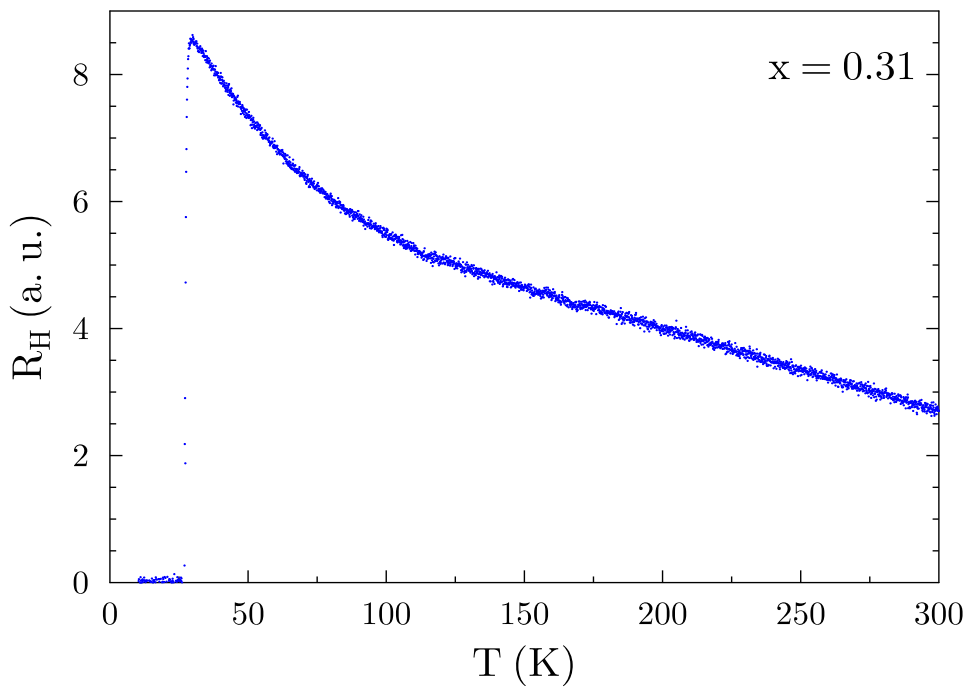
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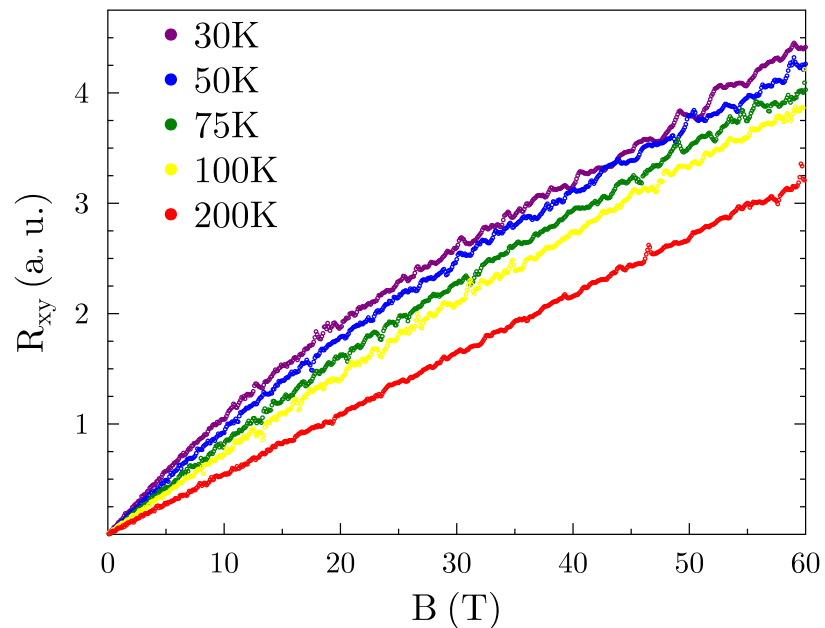
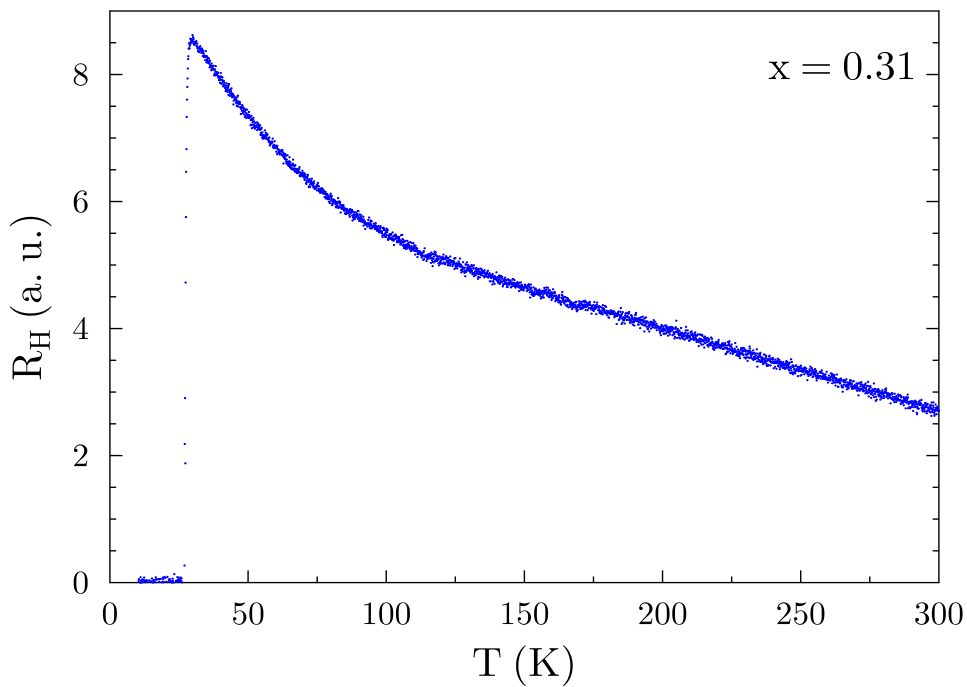
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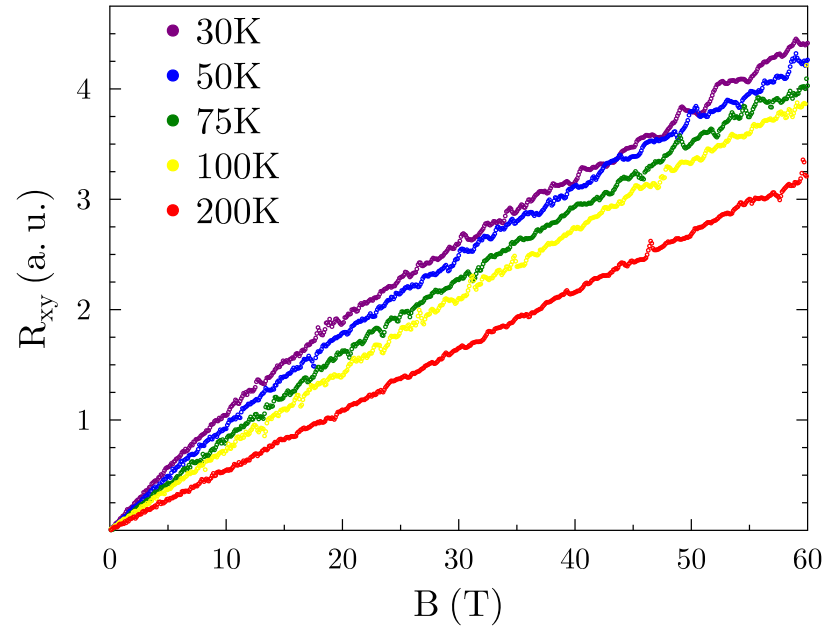
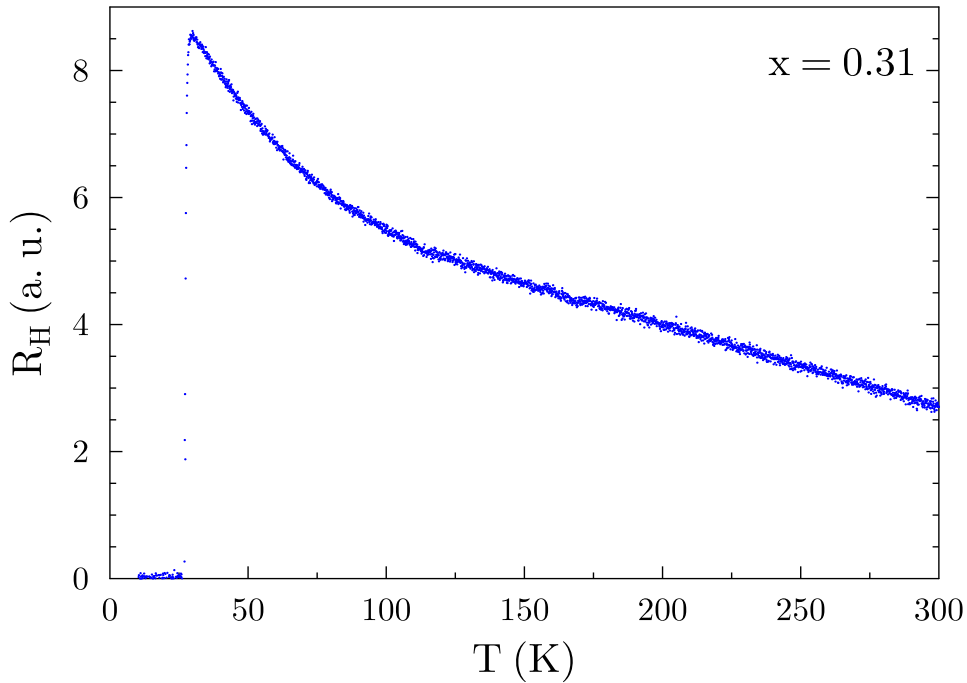


High Field Hall Effect in Ba-122



$$R_H = R_{H\text{Background}}(T) + \frac{1}{T} f(B/T)$$

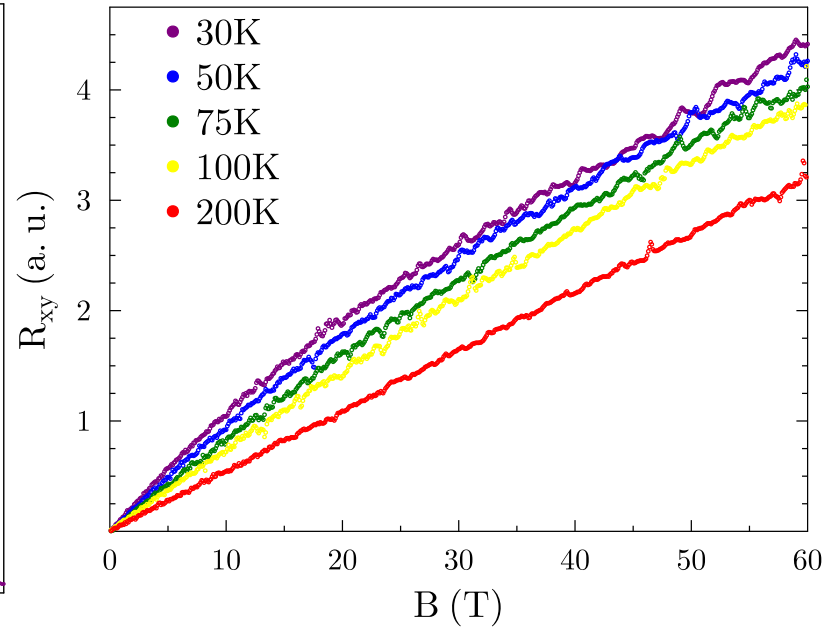
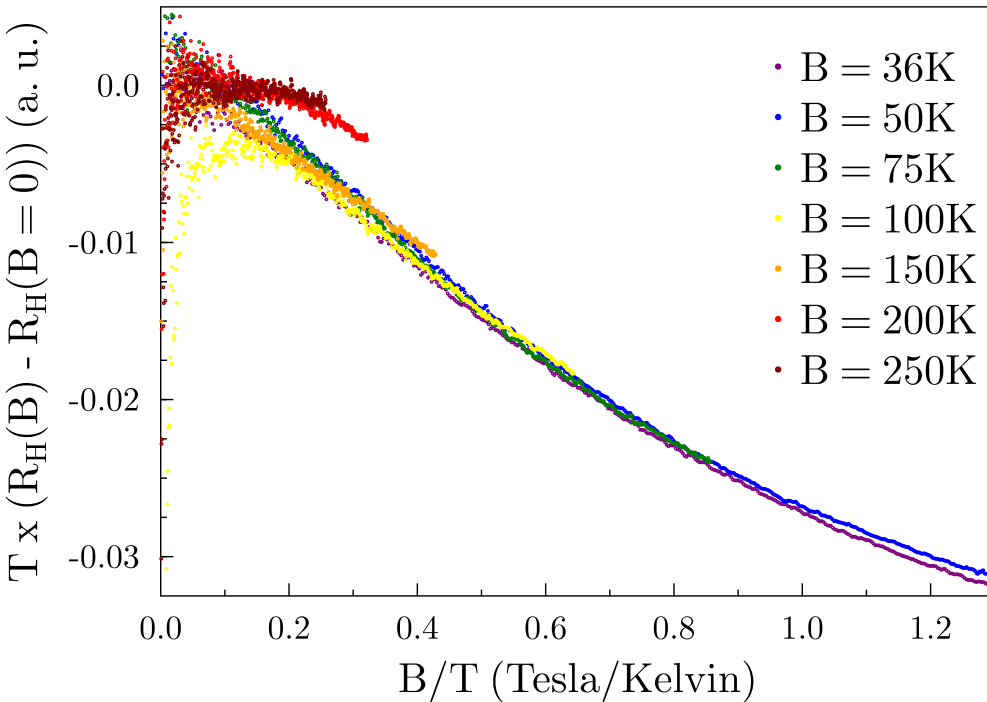
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$$R_H = R_{H\text{Background}}(T) + \frac{1}{T} f(B/T)$$

$$T(R_H - R_{H\text{Background}}(T)) = f(B/T)$$

High Field Hall Effect in Ba-122

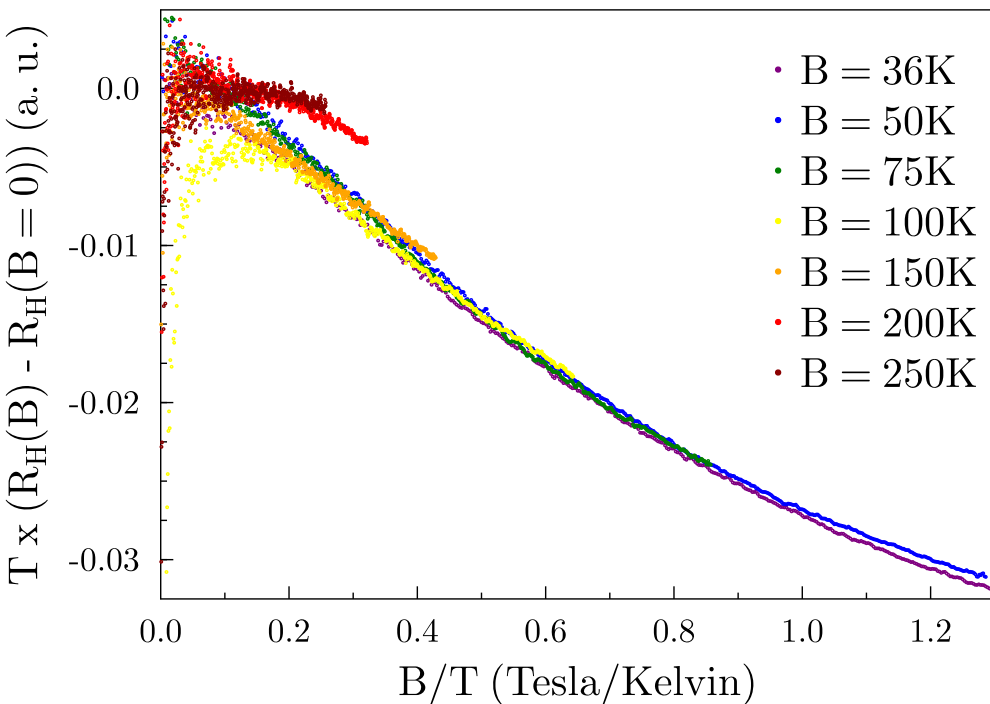


MR curves collapse for $T < 150$ K,
where the R_H upturn is pronounced.

$$R_H = R_{H\text{Background}}(T) + \frac{1}{T} f(B/T)$$

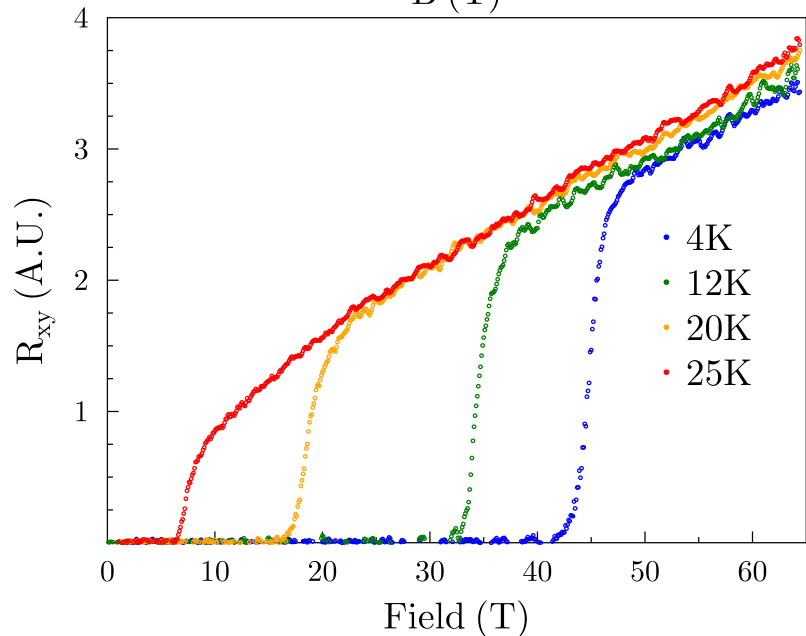
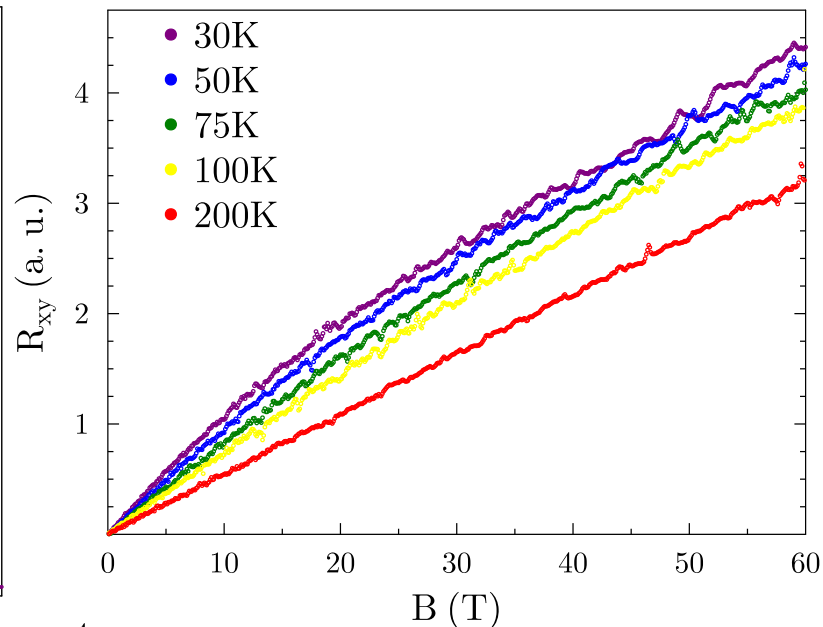
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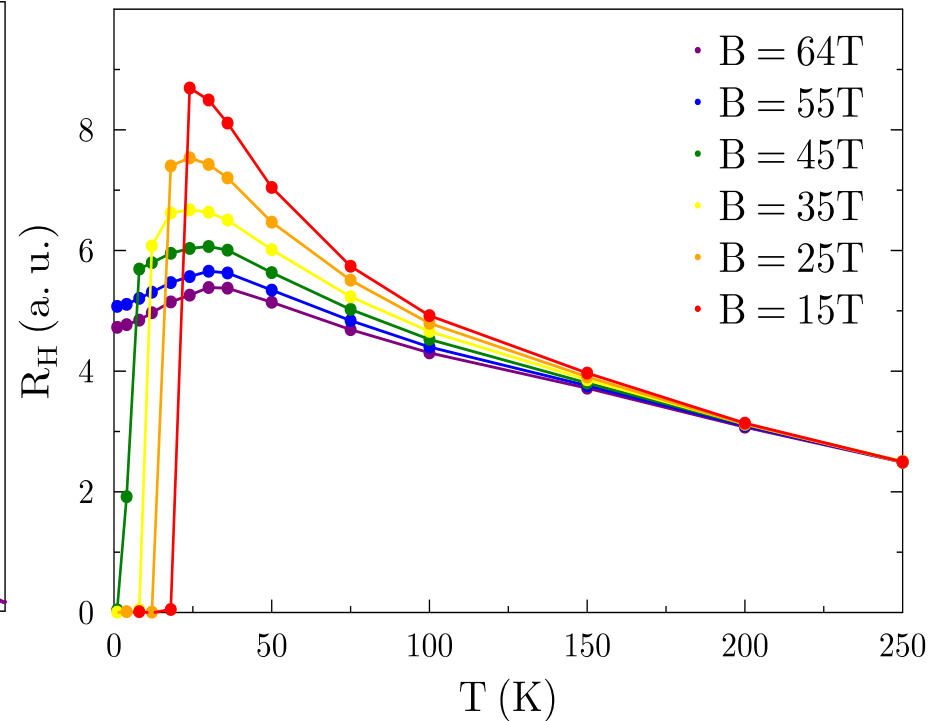
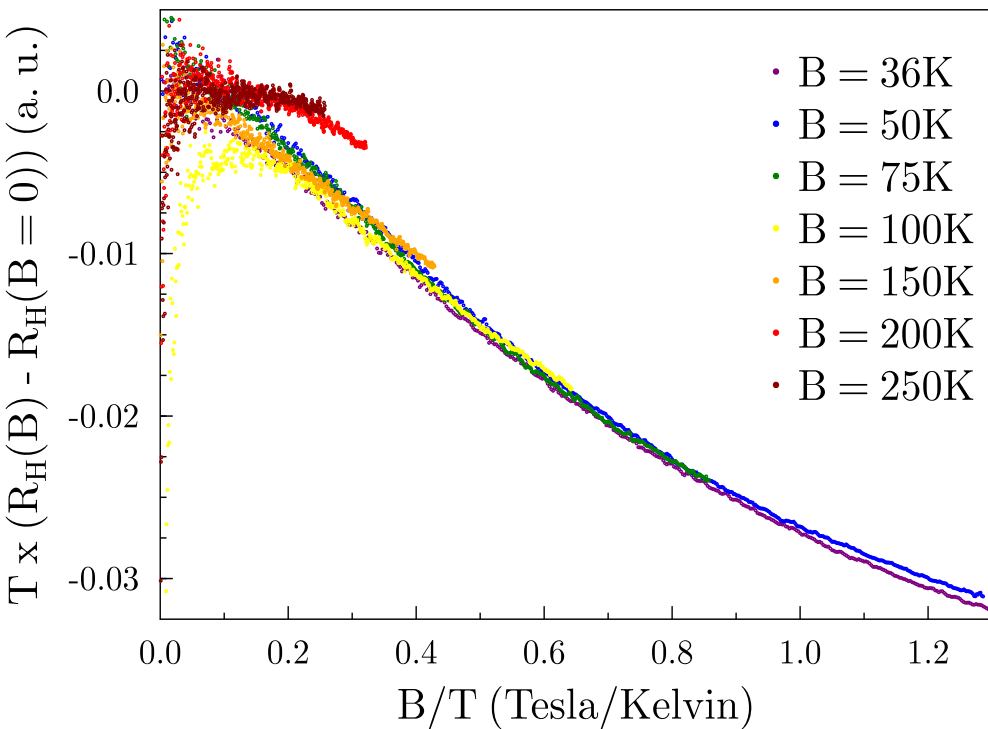


MR curves collapse for $T < 150K$, where the R_H upturn is pronounced.

Downturn in R_H below 30K means the data collapse only for $30K < T < 150K$.

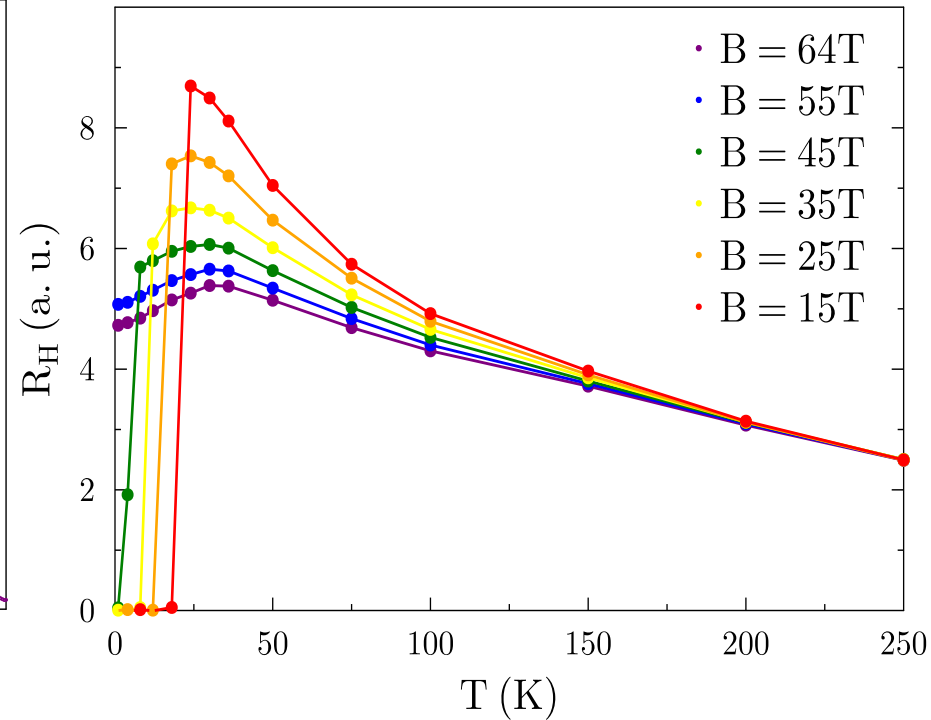
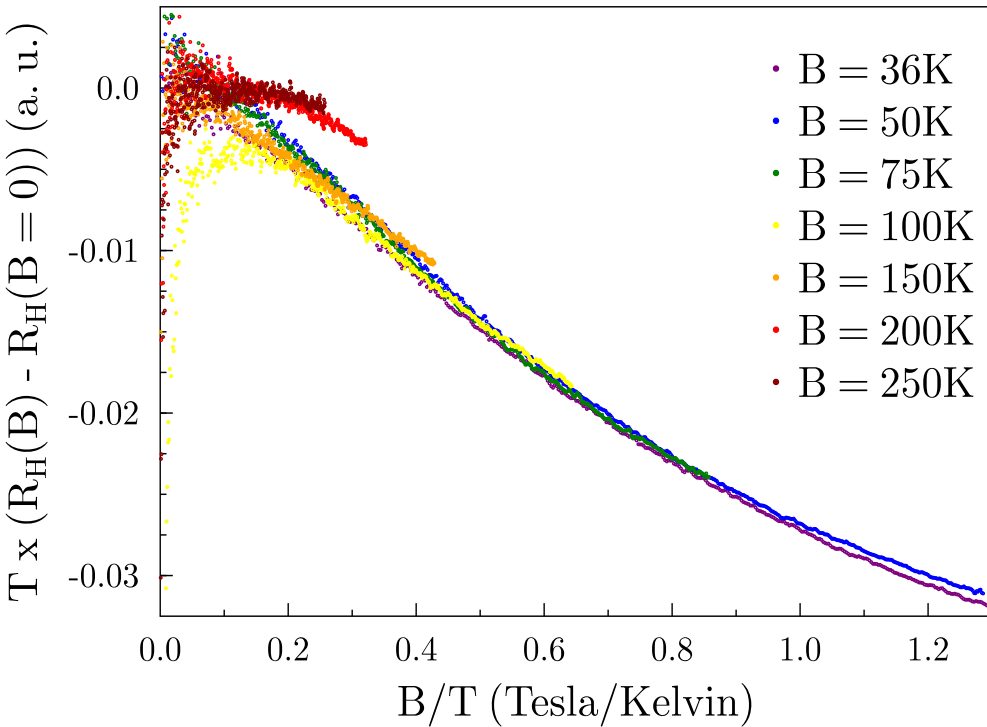


High Field Hall Effect in Ba-122



R_H versus T at high fields closely resembles R_H versus T at low fields and higher doping.

High Field Hall Effect in Ba-122



R_H versus T at high fields closely resembles R_H versus T at low fields and higher doping.

64 Tesla \sim 40K cut-off using the scale factor from the B-linear scaling

Summary and Outlook

- ρ_{xx} in Ba-122 shows scaling between B and T with both T- and B-linear resistivity, consistent with the interpretation that both tune a cut-off energy scale for charge relaxation.
- R_H shows a steep upturn at low temperatures similar to what is seen in cuprates and HF superconductors. The upturn is cut-off at successively higher temperatures as one moves away from optimal doping. R_H is reduced at high-fields as well as high temperatures
- It is an open question why field behaves as both a tuning parameter and a scaling parameter in the Hall Channel

