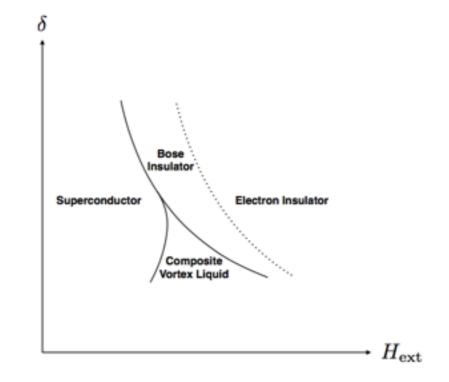
Self-duality and metallic behavior near 2d superconductor-insulator transitions

Srinivas Raghu (Stanford)





Acknowledgments



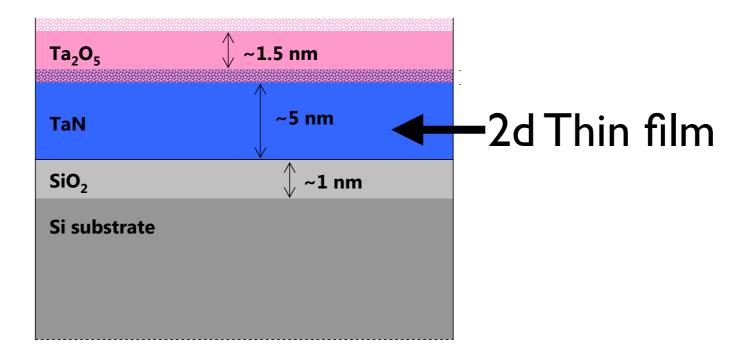
Michael Mulligan (UC Riverside)

(Collaborators on related issues):

Jing-Yuan Chen, Alfred Cheung, Hart Goldman, Pallab Goswami, Prashant Kumar, Jun Ho Son, Gonzalo Torroba, Max Zimet

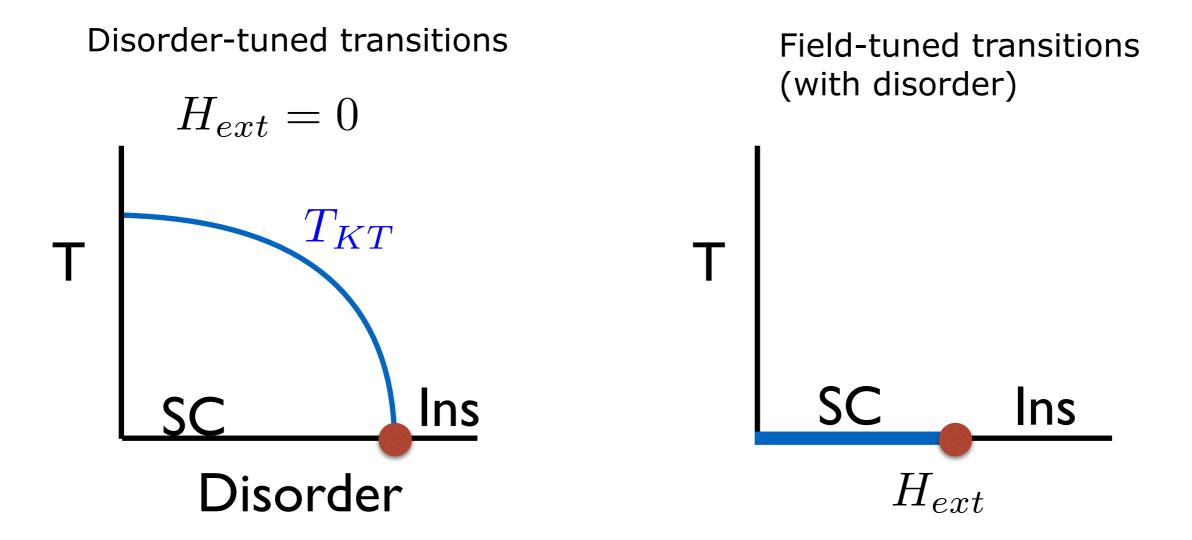
Basic setting: 2d disordered, superconducting thin films

N. Breznay thesis (Stanford)



Destroy superconductivity by tuning disorder or perpendicular magnetic field

Magnetic field vs. disorder tuning in 2d

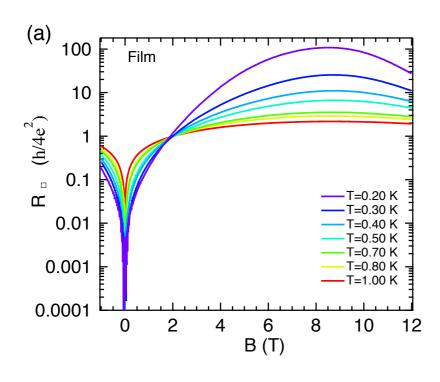


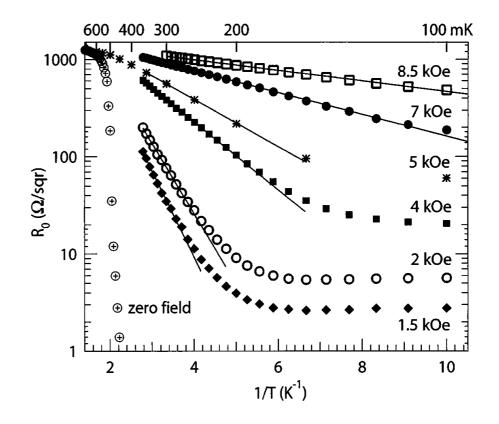
This talk: field-tuned transitions

Motivation: 2 sets of experimental observations:

I) Self-duality near the SIT

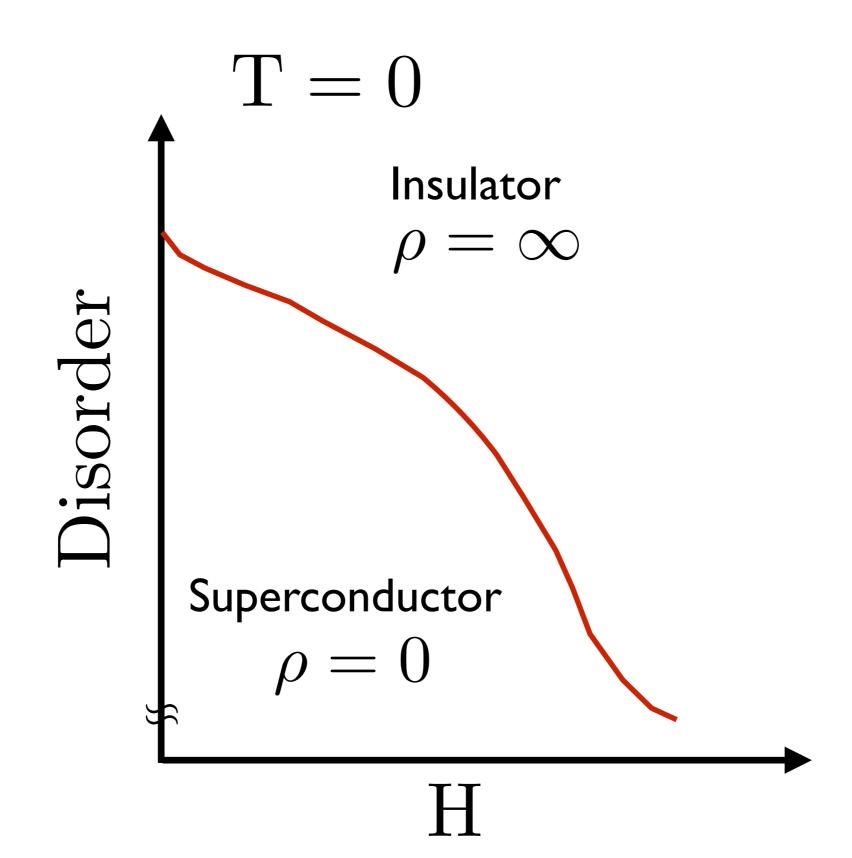
2) metallic phases near the SIT





Are these phenomena related?

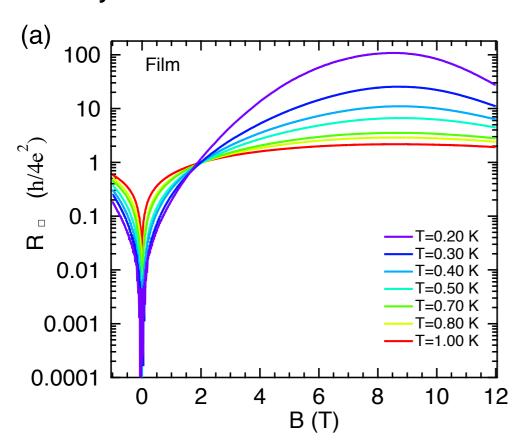
In 2d, we expect a T=0 phase diagram like this:



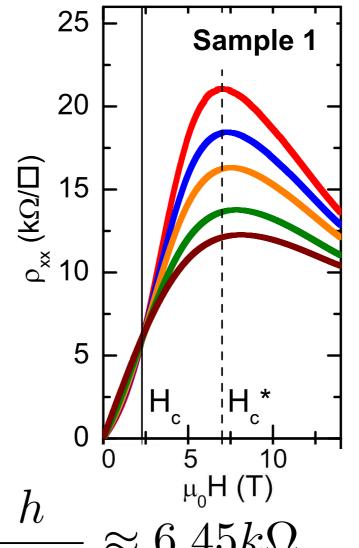
Self-duality: experiment

Paalanen, Hebard, Ruel, PRL (1992)

J. Seidemann et al, arXiv1609.07105

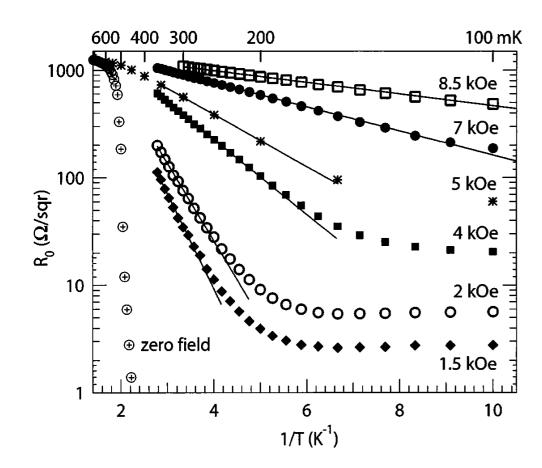


Breznay et al., PNAS (2016)

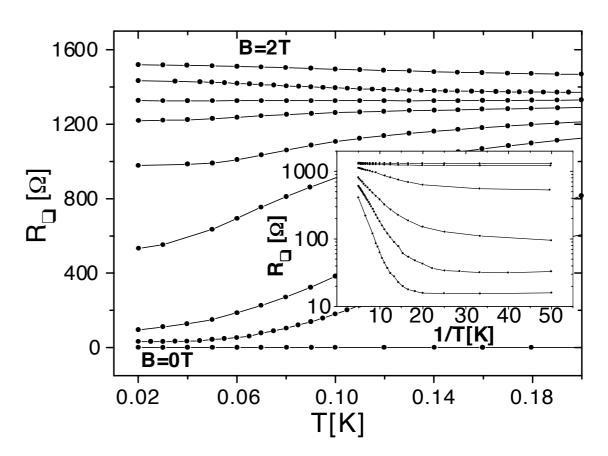


Critical resistance: $R_c = R_Q = \frac{n}{(2e)^2} \approx 6.45 k\Omega$

Metallic behavior: experiment

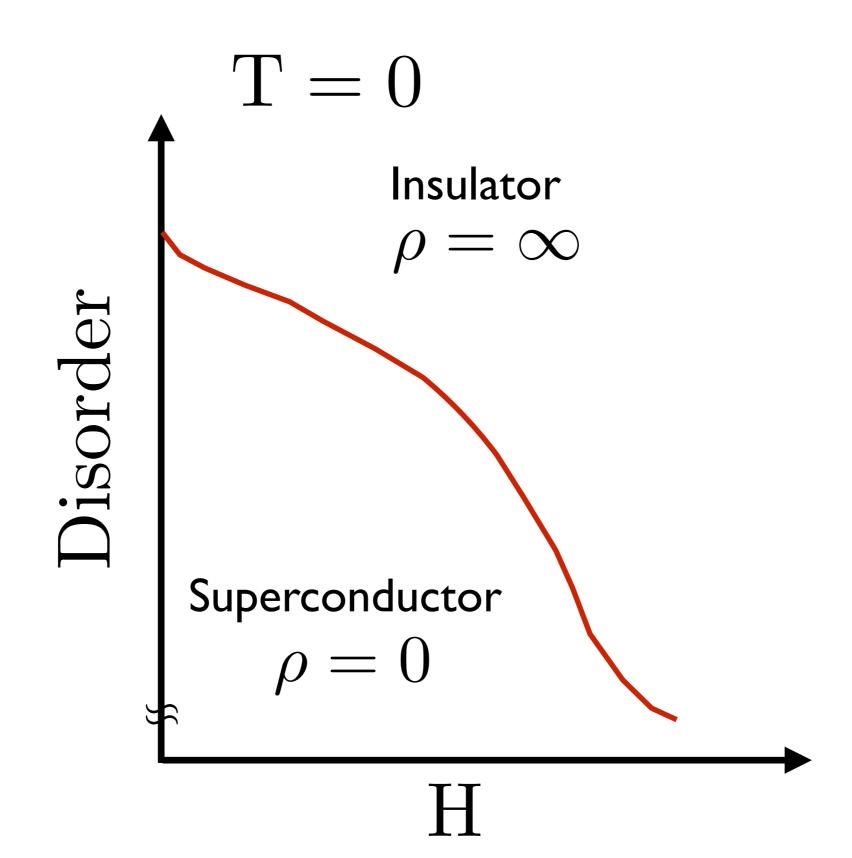


D. Ephron et al, PRL 76, 1529 (1996)

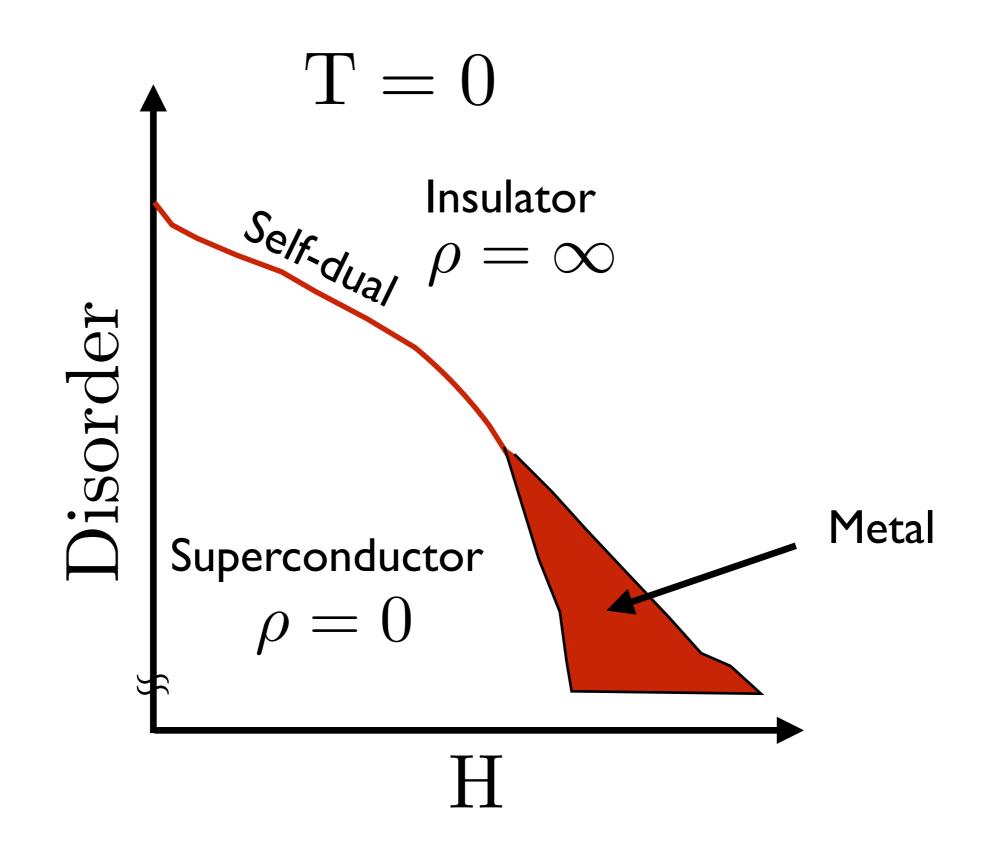


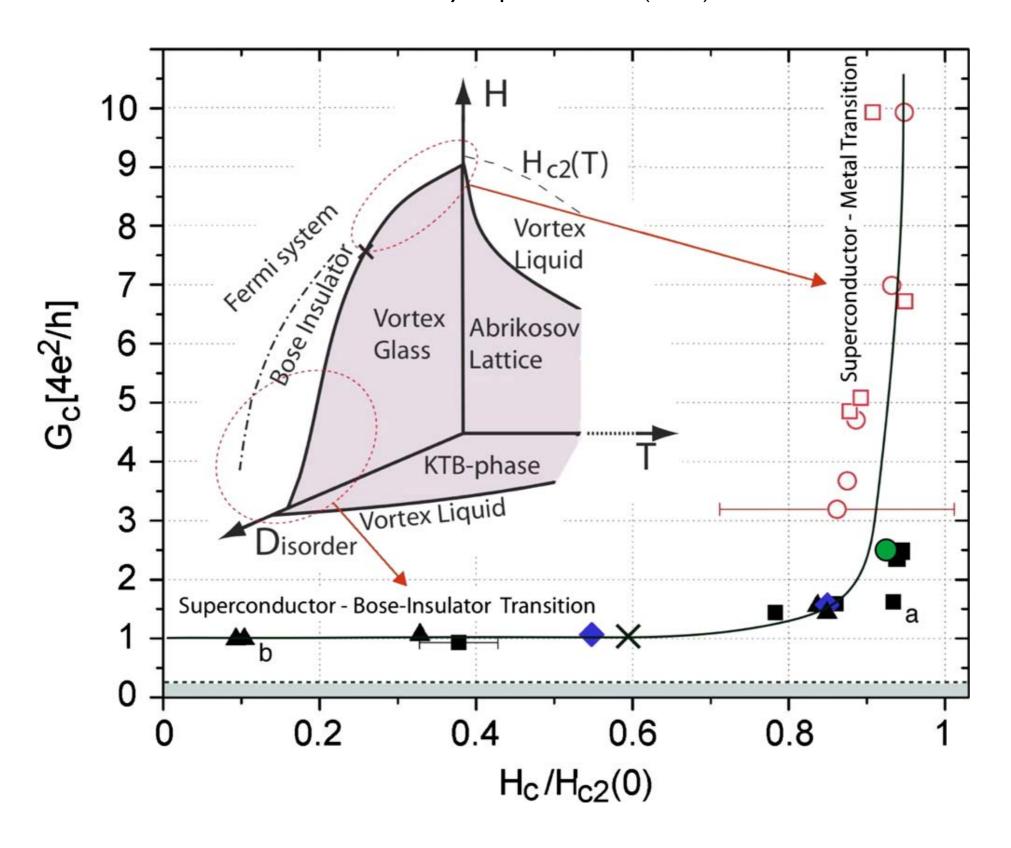
N. Mason et al, PRL 82, 5341 (1998)

In 2d, we expect a T=0 phase diagram like this:



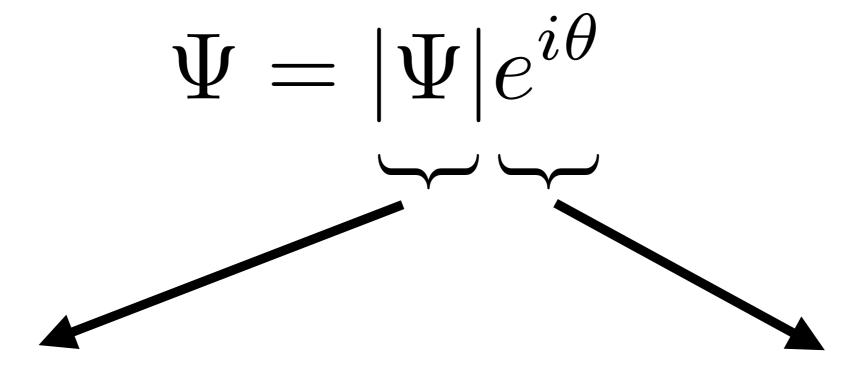
Instead, experiments suggest a phase diagram like this:





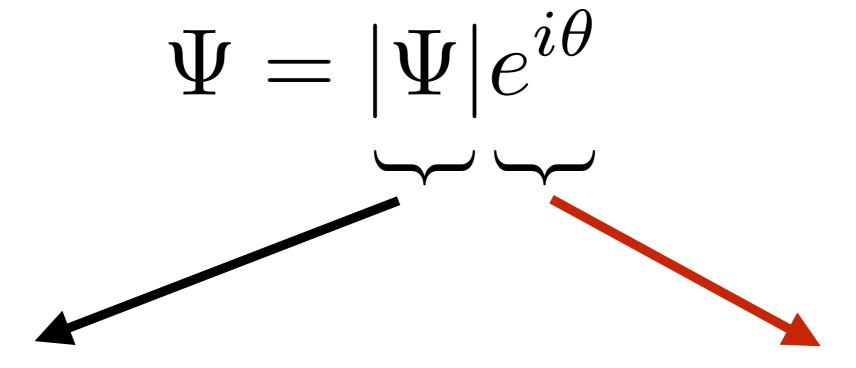
Basic ideas

Superconducting order parameter:



Maekawa, Fukuyama (1981) Finkelstein (1987) Belitz (1989) MPA Fisher (1989, 1990,...) Wen, Zee (1990)

Superconducting order parameter:

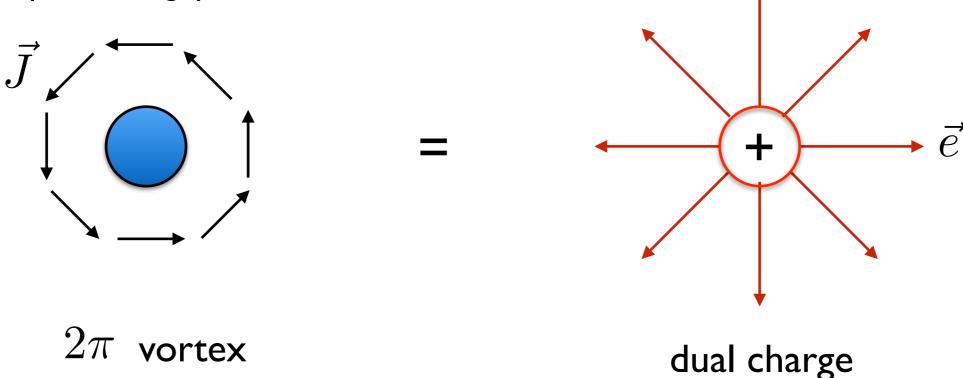


Maekawa, Fukuyama (1981) Finkelstein (1987) Belitz (1989) MPA Fisher (1989, 1990,...) Wen, Zee (1990)

Basic framework: particle-vortex duality + dirt

Particle-vortex duality

Peskin, Halperin, Dasgupta, Fisher, Lee



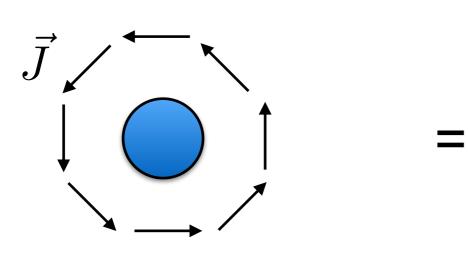
$$J_{\mu} = \frac{1}{2\pi} \epsilon_{\mu\nu\lambda} \partial_{\nu} a_{\lambda} \qquad \qquad j_{\mu} = -\frac{1}{2\pi} \epsilon_{\mu\nu\lambda} \partial_{\nu} A_{\lambda}$$

 A_{μ} : External gauge field (electromagnetism)

 a_{μ} : Emergent gauge field (2d photon)

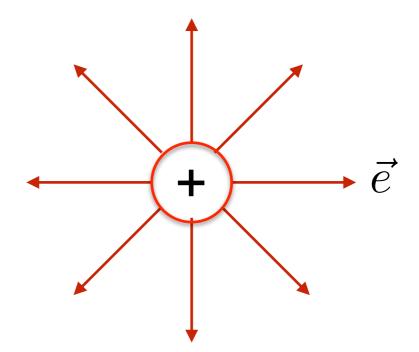
Particle-vortex duality

Peskin, Halperin, Dasgupta, Fisher, Lee



 2π vortex

$$J_{\mu} = \frac{1}{2\pi} \epsilon_{\mu\nu\lambda} \partial_{\nu} a_{\lambda}$$
$$J_{0} = -\frac{1}{2\pi} \nabla \times \vec{a} = b/2\pi$$
$$\vec{J} = \frac{1}{2\pi} \hat{z} \times \vec{e}$$



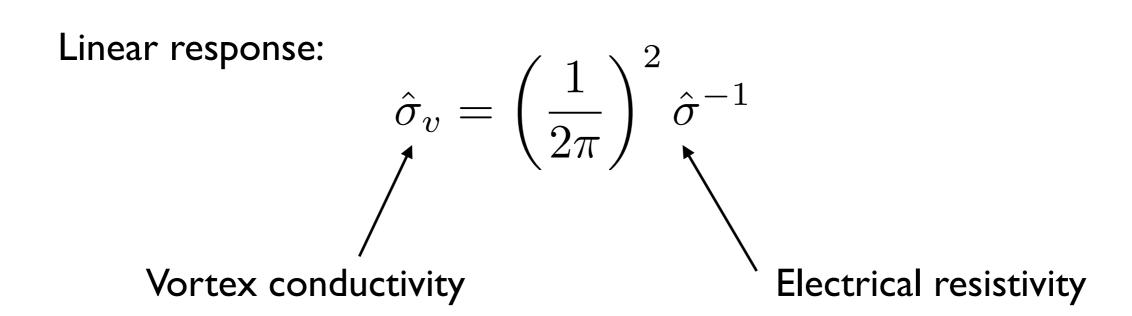
dual charge

$$j_{\mu} = -\frac{1}{2\pi} \epsilon_{\mu\nu\lambda} \partial_{\nu} A_{\lambda}$$

$$j_{0} = -\frac{1}{2\pi} \nabla \times \vec{A} = -B/2\pi$$

$$\vec{j} = -\frac{1}{2\pi} \hat{z} \times \vec{E}$$

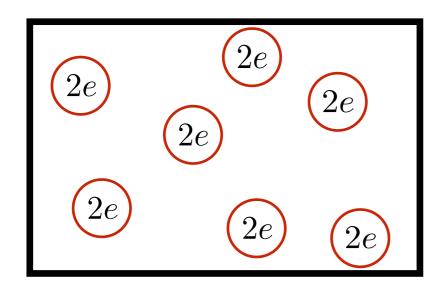
Particle-vortex duality



Phases

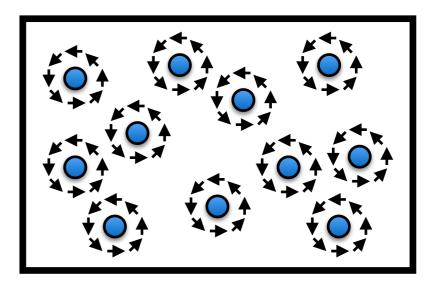
- 1) Superconductor: Cooper pair condensate, gapped vortices
- 2) Insulator: vortex condensate (Higgs), gapped Cooper pairs

Superconductor $(H \ll H_c)$



Condensate of mobile Cooper pairs

Insulator $(H >> H_c)$

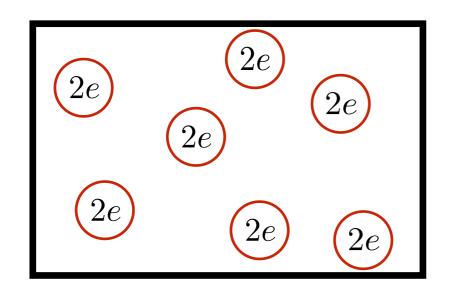


Condensate of mobile vortices

What about the critical point (SIT)?

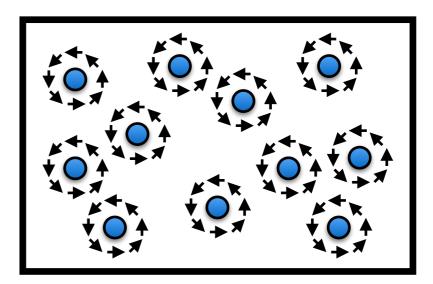
Disorder: most of the Cooper pairs are localized (frozen).

Superconductor ($H << H_c$)



Condensate of mobile Cooper pairs

Insulator $(H >> H_c)$



Condensate of mobile vortices

What about the critical point (SIT)?

Self-duality at the SIT

Duality:

$$\hat{\sigma}_v = \left(\frac{1}{2\pi}\right)^2 \hat{\sigma}^{-1} \qquad [\hbar = c = 2e = 1]$$

Self-duality:

$$\hat{\sigma}_v = \hat{\sigma}^T$$

Consequence:

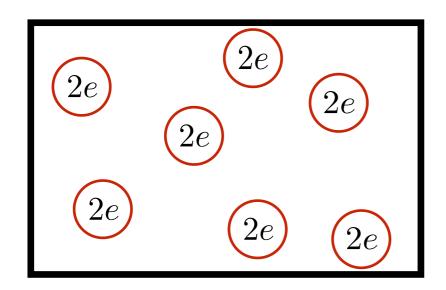
$$R_{xx}^2 + R_{xy}^2 = R_Q^2$$
 $R_Q = \frac{h}{(2e)^2} \approx 6.45k\Omega$

$$R_Q = \frac{h}{(2e)^2} \approx 6.45k\Omega$$

Requirement for self-duality: Cooper pairs and vortices both have the same response Lagrangian

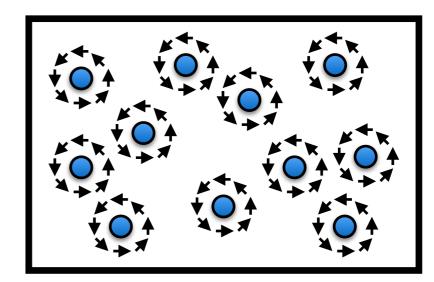
Disorder: most of the Cooper pairs are localized (frozen).

Superconductor ($H \le H_c$)



Condensate of mobile Cooper pairs

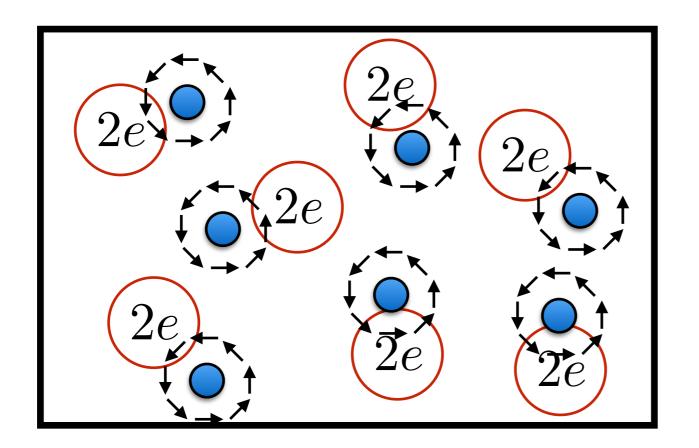
Insulator $(H >> H_c)$



Condensate of mobile vortices

Our proposal for the SIT with self-duality

Self-dual SIT
$$(H = H_c)$$

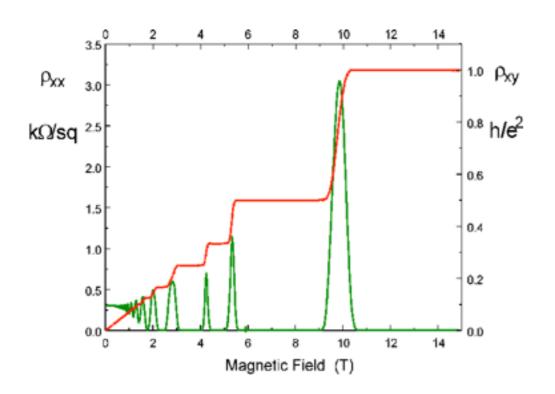


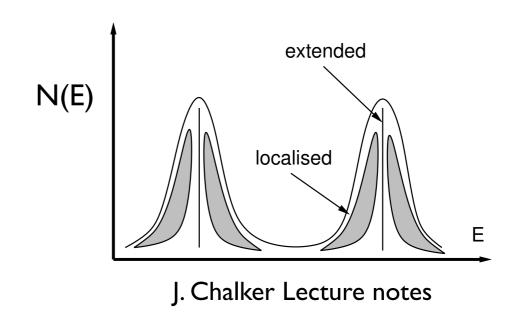
$$2e\sqrt{2}$$
 = A fermion

(composite vortex)

Finite resistance at the SIT due to emergent fermions

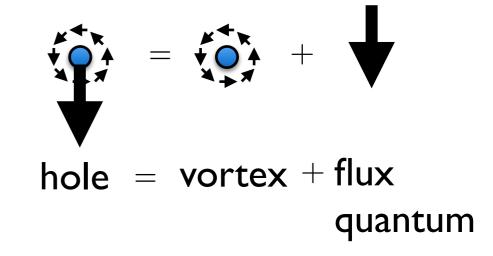
Analogous story: integer quantum Hall-insulator transition (QHIT)





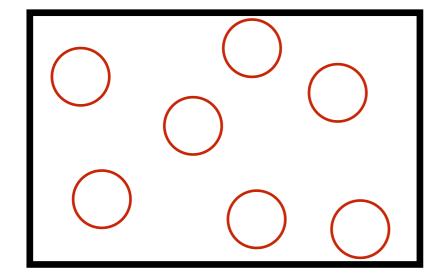
Bosonic description:

electron = boson + flux quantum



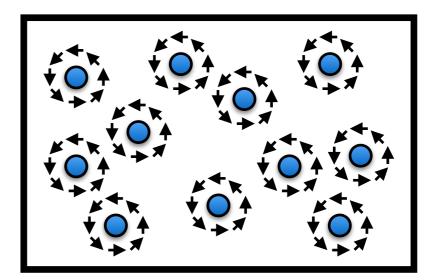
S.-C. Zhang, T. Hansson, S. Kivelson

Integer QH $(\nu=1)$



Condensate of mobile Composite bosons

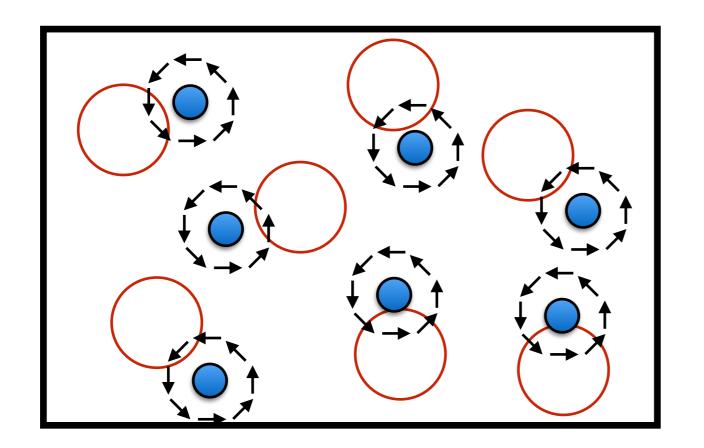
Insulator $(\nu = 0)$



Condensate of mobile vortices of comp. bosons

Integer quantum Hall - Insulator transition (QHIT)

QHIT with particle-hole symmetry



Finite resistance at the QHIT due to emergent fermions

Dictionary

QHIT

Integer QH
$$(\nu=1)$$

Insulator
$$(\nu = 0)$$

Particle-hole symmetric QHIT

$$(\nu = 1/2)$$

SIT

Superconductor

Insulator

Self-dual SIT

In both cases the critical point involves composite fermions

Slightly more quantitative discussion

Cooper pairs:

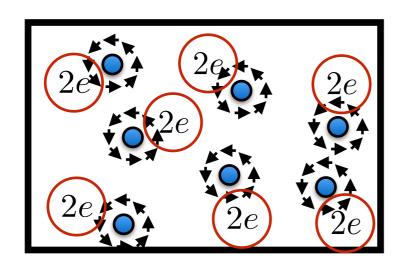
$$L_A = |D_A \Phi|^2 - V(|\Phi|) + \text{dirt}$$

 Φ : Cooper pair

 ϕ : Vortex

Vortices:

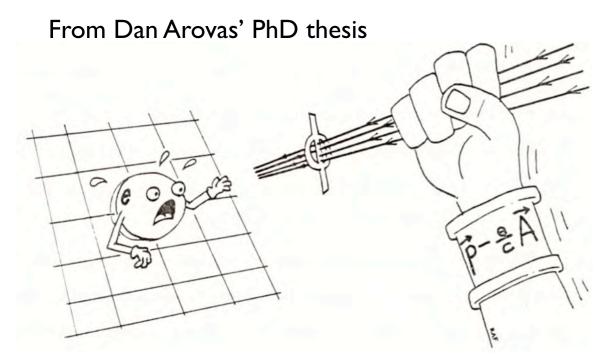
SIT: vortices at unit filling ($\nu = 1$).

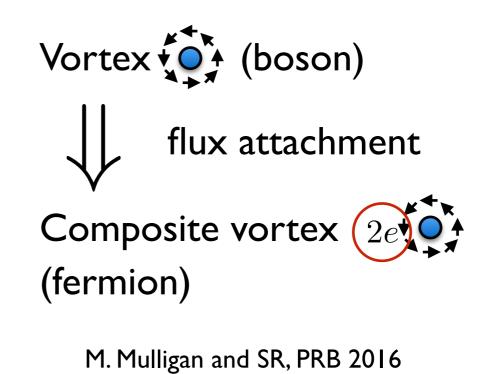


Composite vortex metal

$$2e^{-\frac{1}{2}}$$
 = A fermion

This can be implemented with flux attachment





2 choices for theoretical description of composite vortex metal:

The SIT then maps onto the problem of a half-filled Landau level

- I) Halperin-Lee-Read (HLR) composite Fermi liquid
 M. Mulligan and SR, PRB 2016
- 2) D. Son: Dirac composite Fermi liquid
 M. Mulligan, PRB 2017

QHIT: HLR theory seems to lack particle-hole symmetry

SIT: HLR theory for composite vortices seems to lack self-duality

QHIT: Son theory manifestly preserves particle-hole symmetry

SIT: Son theory for composite vortices is manifestly selfdual

M. Mulligan, PRB 2017

"Modern" discussion of self-dual theory

M. Mulligan, PRB 2017

Cooper pairs:

$$L_A = |D_A \Phi|^2 - V(|\Phi|) + \text{dirt}$$

Vortices:
$$L_B = |D_a \phi|^2 - \tilde{V}(|\phi|) + \frac{1}{2\pi} a dA + \text{dirt}'$$

composite vortices:

$$\downarrow$$

$$L_C = \bar{\chi} \not \!\!\!D_c \chi - \frac{1}{8\pi} c dc - \frac{1}{2\pi} c dA - \frac{1}{4\pi} A dA + \text{dirt}'$$

2nd-3rd Lines: make use of a conjectured (and recently wellestablished) boson-fermion duality

Boson-fermion "duality" (t < 2015)

```
Polyakov (1988),
Zhang, Hansson, Kivelson (1989),
Chen, Fisher, Wu (1993),
Fradkin, Kivelson (1996),
Barkeshli, Mcgreevy (2014)
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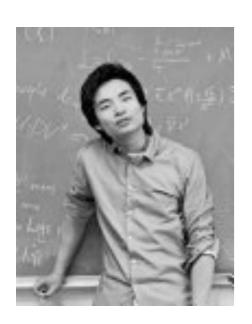
Boson-fermion duality (t > 2015)

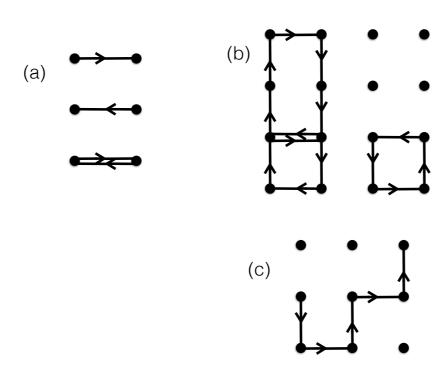
```
Karch, Tong (2016),
Seiberg, Senthil, Wang, Witten (2016),
Kachru, Mulligan, Torroba, Wang (2016),
Metlitski, Vishwanath, Xu (2016),
Mross, Alicea, Motrunich (2017),
Chen, Son, Wang, SR (2017),
```

arXiv:1705.05841

Exact Boson-Fermion Duality on a 3D Euclidean Lattice

Jing-Yuan Chen¹, Jun Ho Son¹, Chao Wang¹ and S. Raghu^{1,2}





Another dictionary

Vortices:

$$L_B = |D_a \phi|^2 - \tilde{V}(|\phi|) + \frac{1}{2\pi} a dA + \text{dirt}'$$

composite vortices:

$$\dot{L}_C = \bar{\chi} \not \!\!\!D_c \chi - \frac{1}{8\pi} c dc - \frac{1}{2\pi} c dA - \frac{1}{4\pi} A dA + \text{dirt}'$$

Vortices:

composite vortices:

Superconductor: insulator insulator

Insulator: condensate Integer QH

Self-dual SIT: ??? composite vortex metal

Self-duality = time-reversal for composite vortices

$$L_C = \bar{\chi} \not \!\! D_c \chi - \frac{1}{8\pi} c dc - \frac{1}{2\pi} c dA - \frac{1}{4\pi} A dA + \text{dirt}'$$

ct equation of motion:

$$\langle \chi^{\dagger} \chi \rangle = \frac{1}{4\pi} \langle \nabla \times \vec{c} \rangle + \frac{B}{2\pi}$$

Self-duality when

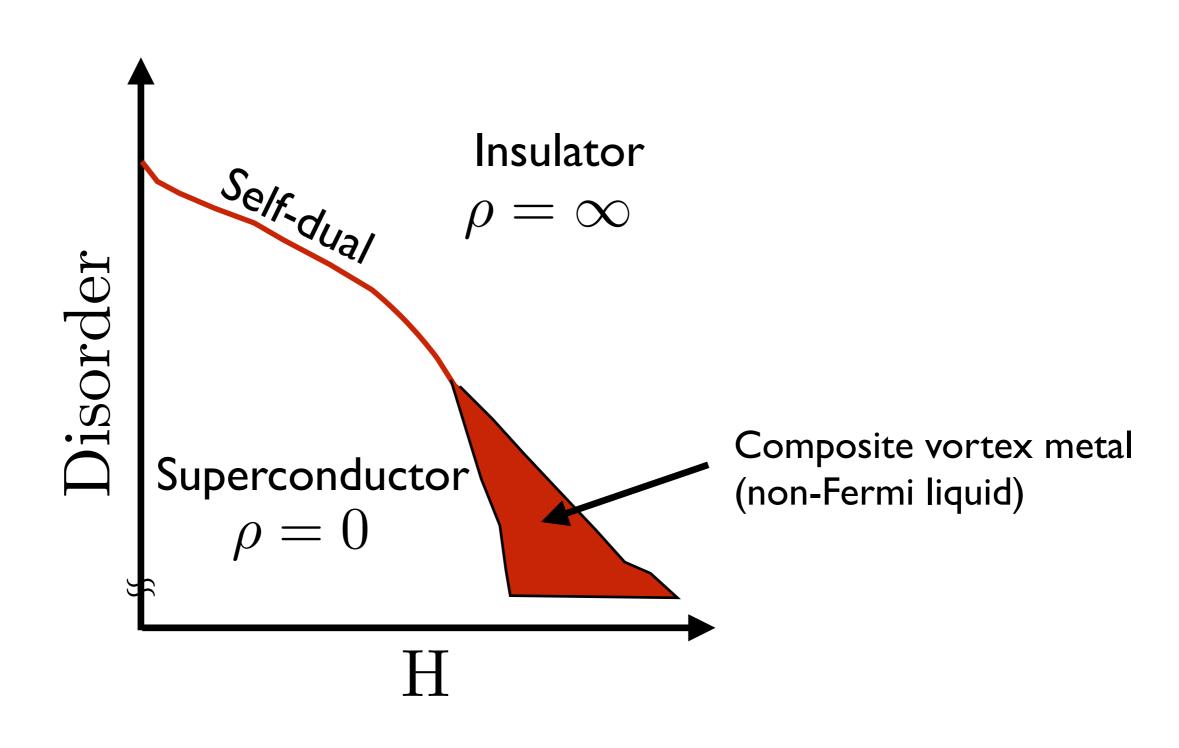
$$\langle \nabla \times \vec{c} \rangle = 0$$

Note:
$$\langle J_0 \rangle = -\frac{1}{2\pi} \langle \nabla \times \vec{c} \rangle - \frac{B}{2\pi}$$
 Mismatch between number of Cooper pairs and vortices = flux for composite vortices

Self-duality = time-reversal for composite vortices

$$\rho_{xx}^{2} + \rho_{xy}^{2} = (2\pi)^{2} \frac{(\sigma_{xx}^{\chi})^{2} + (\frac{1}{2} + \sigma_{xy}^{\chi})^{2}}{(\sigma_{xx}^{\chi})^{2} + (\frac{1}{2} - \sigma_{xy}^{\chi})^{2}}$$

Guess: composite vortex metal broadens into a phase



A prediction based on composite vortices

Borrow an old QHE trick:

Composite vortices "see" a beff << B

Huge magnetic length >> mean-free path
Therefore, very large cyclotron radii in the metal.

cleaner systems: quantum oscillations depend on beff, not B.

What we don't (yet) understand

- 1) Metallic phases have zero cyclotron resonance.
- 2) Metallic phases have zero Hall effect (previous talk).
- 3) Neighboring insulating phases: Hall insulator?
- 4) Analogous phenomena in quantum Hall context?

The existence of 2d metallic phases and neighboring transitions: challenging problem of effective field theory

Important consistency checks:

- 1) Stability of such metals to disorder?
- 2) Can they exhibit T=0 SC-metal transitions?

The existence of 2d metallic phases and neighboring transitions: challenging problem of effective field theory.

Important consistency checks:

1) Stability of such metals to disorder?

Yes in certain cases

P. Goswami, H. Goldman, SR, PRB (2017)

A. Thomson, S. Sachdev, PRB (2017)

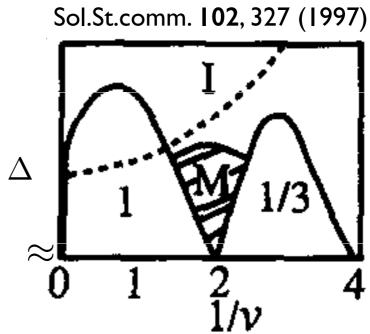
(see posters)

Perturbative

Sol.St.c.

Non-Perturbative

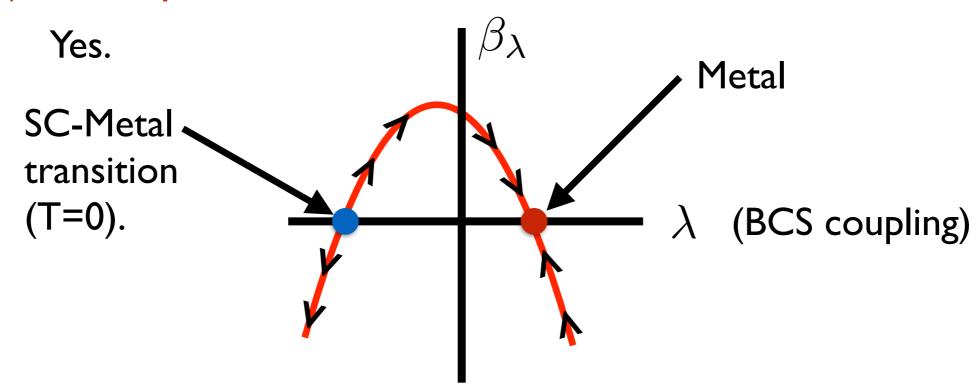
M. Mulligan, SR, G. Torroba, M. Zimet (unpublished)



The existence of 2d metallic phases and neighboring transitions: fundamental conceptual problem.

Important consistency checks (effective field theory):

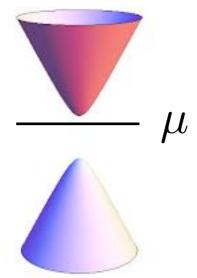
- 1) Stability of such metals to disorder?
- 2) Can they exhibit T=0 SC-metal transitions?



BCS beta function for Fermi surface + U(I) gauge field

Summary

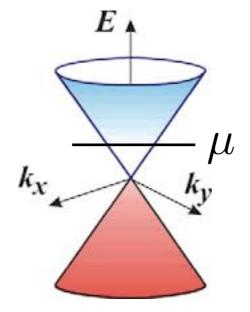
$$m, \nabla \times c > 0$$



superconductor

IQHE of fermions

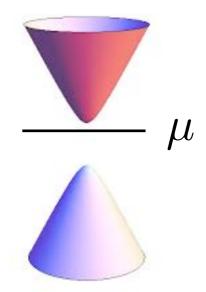
$$m, \nabla \times c = 0$$



SIT/metallic phases

Composite Fermi liquid (NFL)

$$m, \nabla \times c < 0$$



Insulator

Fermi insulator