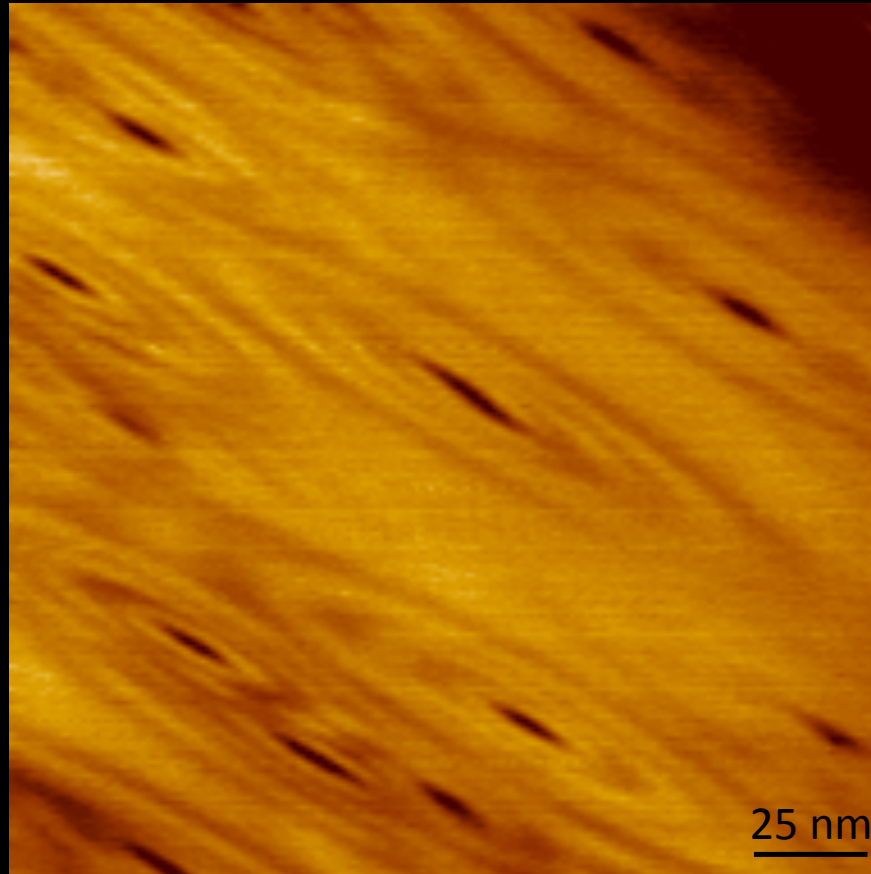


# Imaging Landau orbits to observe a nematic quantum Hall phase



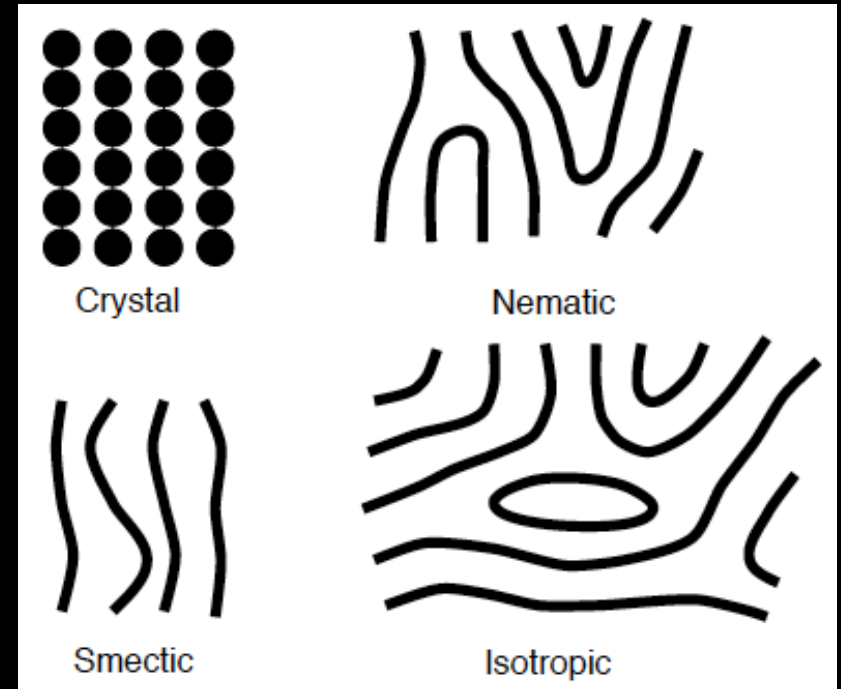
Mallika Randeria



PRINCETON  
UNIVERSITY

# Nematic electronic phases

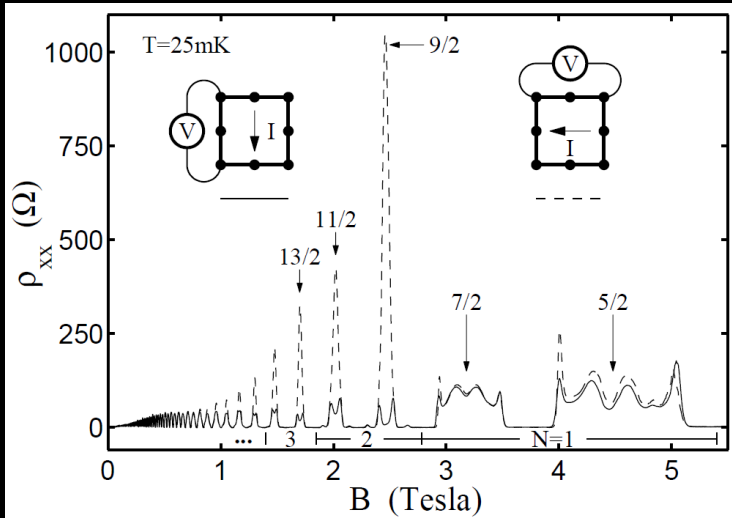
- Break the rotational symmetry of the underlying crystal lattice
- Spontaneously form, *e.g.* as a result of electron-electron interactions



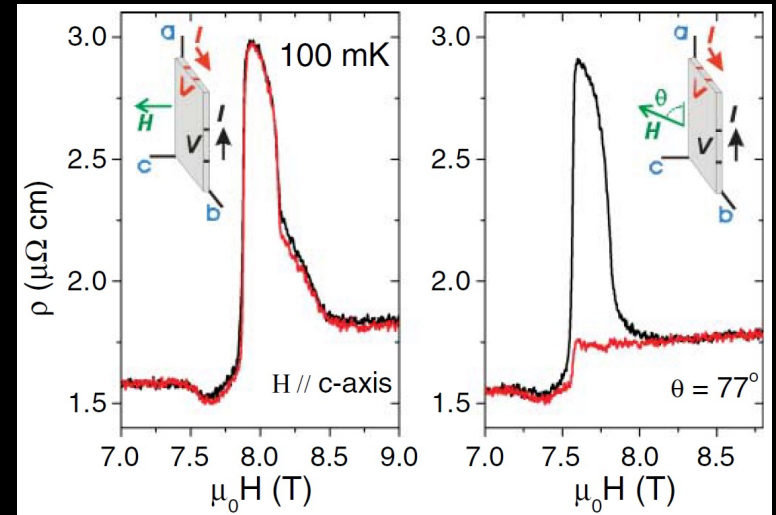
Kivelson, Fradkin & Emery, *Nature* **393**, 550-553 (1998)

# Previous evidence for nematic electronic states

2D electron systems: GaAs, AlAs, graphene

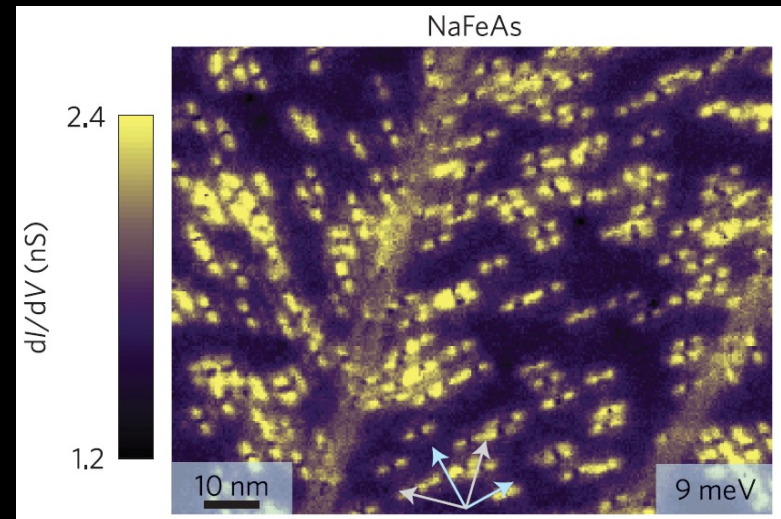


$\text{Sr}_3\text{Ru}_2\text{O}_7$



“Vestigial order” – Strain or external fields needed to align domains to enable detection

High- $T_c$ , Fe-based superconductors



Lilly, M. P. *et al.*, *Phys. Rev. Lett.* **82**, 394-397 (1999)

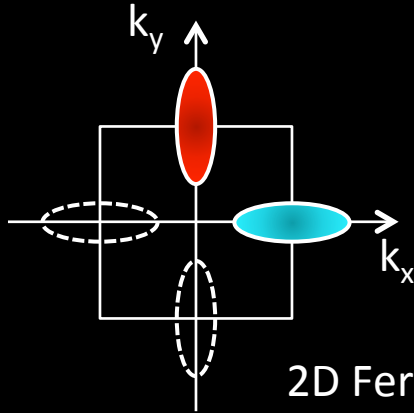
Borzi, R. A. *et al.*, *Science* **315**, 214-217 (2007)

Rosenthal, E. P. *et al.*, *Nat. Phys.* **10**, 225-232 (2014)

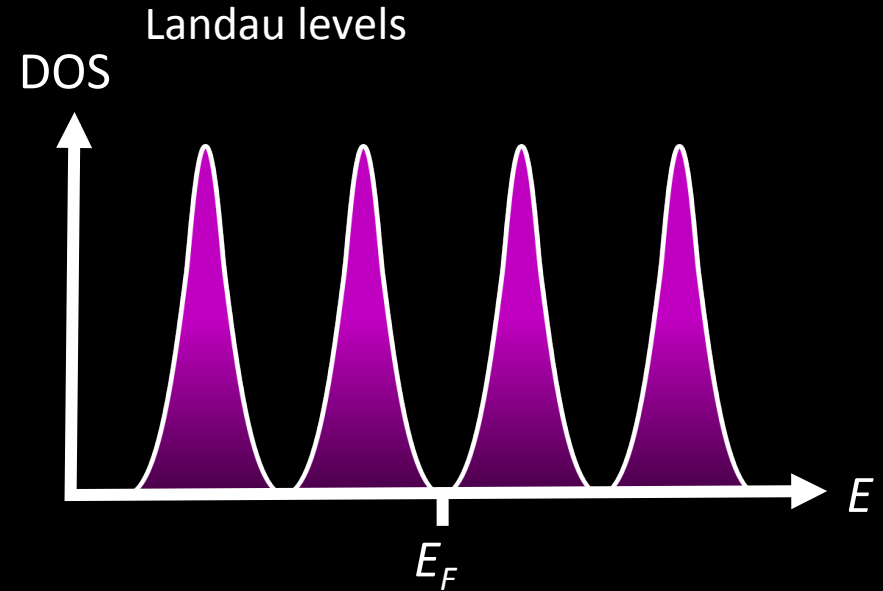
See also work from: Seamus Davis, Ian Fisher groups & many others

# A nematic quantum Hall state

Anisotropic valley  
degree of freedom



2D Fermi pockets  
(schematic for AIAs)



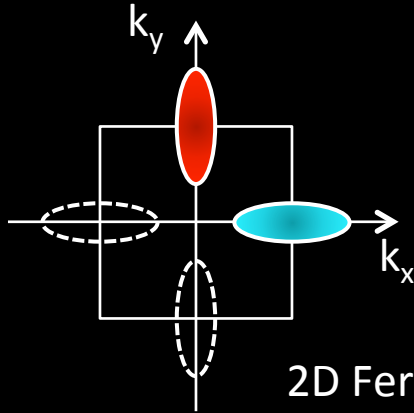
D. Abanin *et al.*, *Phys. Rev. B* **82**, 035428 (2010)

A. Kumar *et al.*, *Phys. Rev. B* **88**, 045133 (2013)

See also AIAs work from Shayegan group

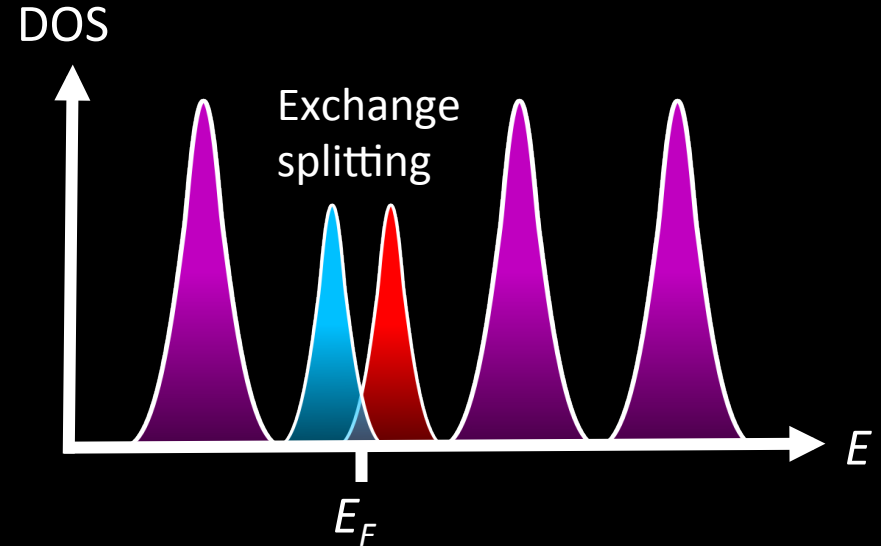
# A nematic quantum Hall state

Anisotropic valley  
degree of freedom



2D Fermi pockets  
(schematic for AIAs)

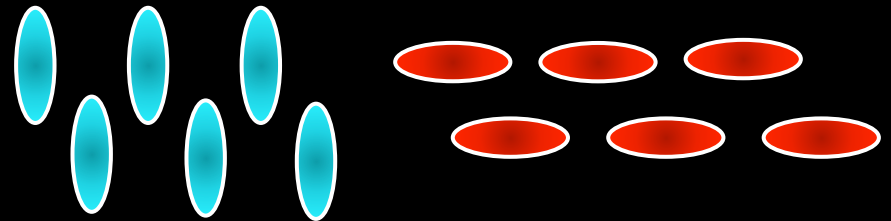
Quantum Hall ferromagnetism



Model nematic system

- Quantify role of interactions
- resulting Landau orbit anisotropy
- Locally probe individual domains

Resulting cyclotron orbit anisotropy

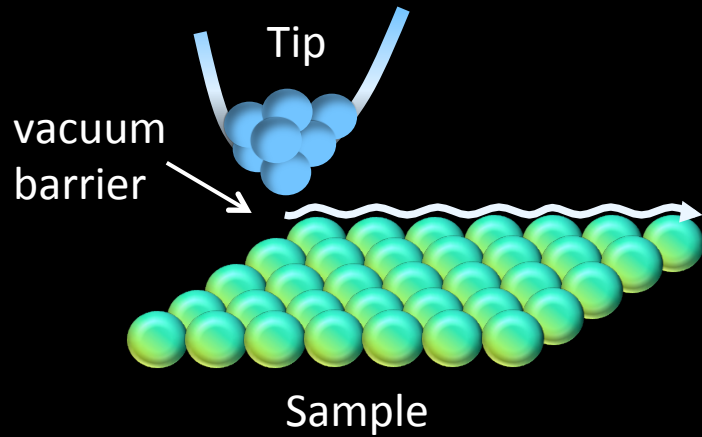


D. Abanin *et al.*, *Phys. Rev. B* **82**, 035428 (2010)

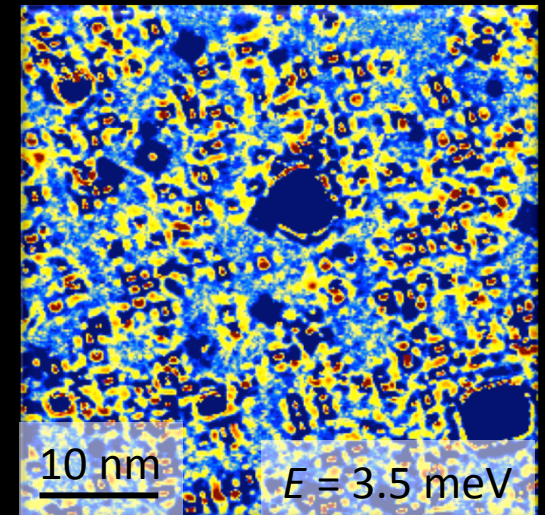
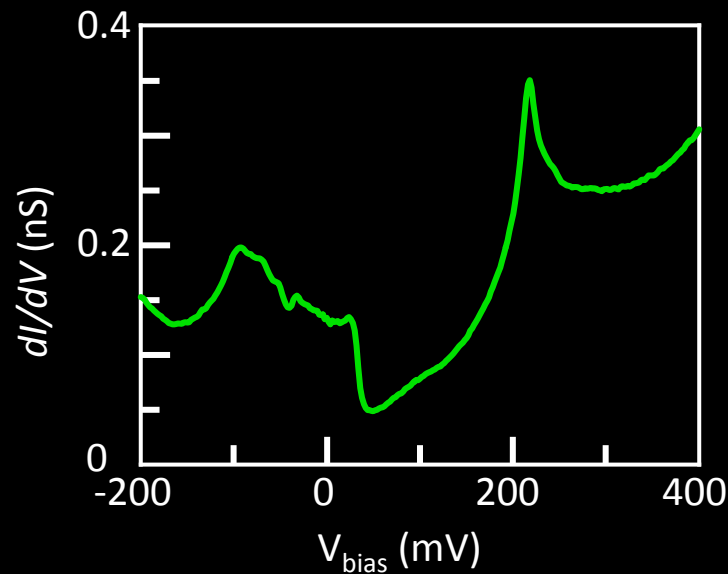
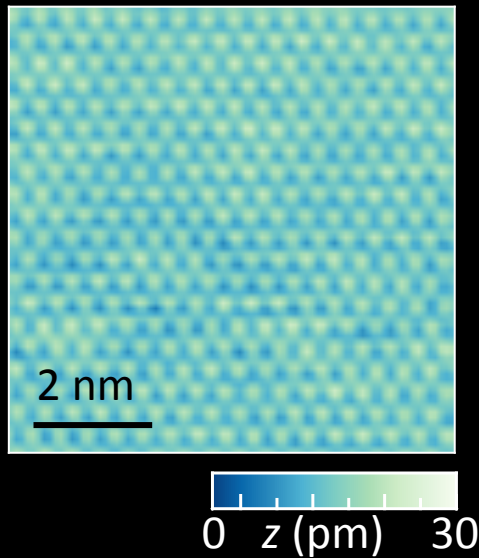
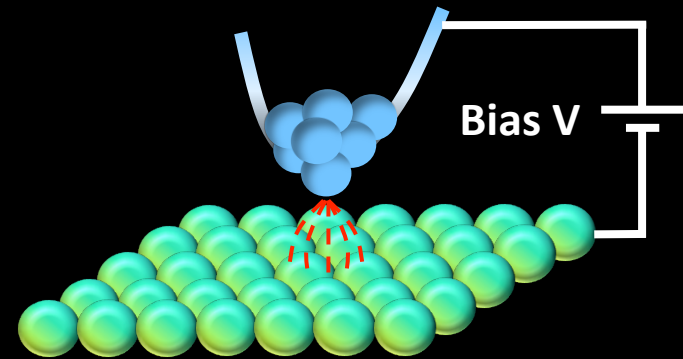
A. Kumar *et al.*, *Phys. Rev. B* **88**, 045133 (2013)

# Scanning tunneling microscope (STM)

Atomic resolution imaging



Spectroscopic mapping:  $dI/dV \propto \rho(r, V)$



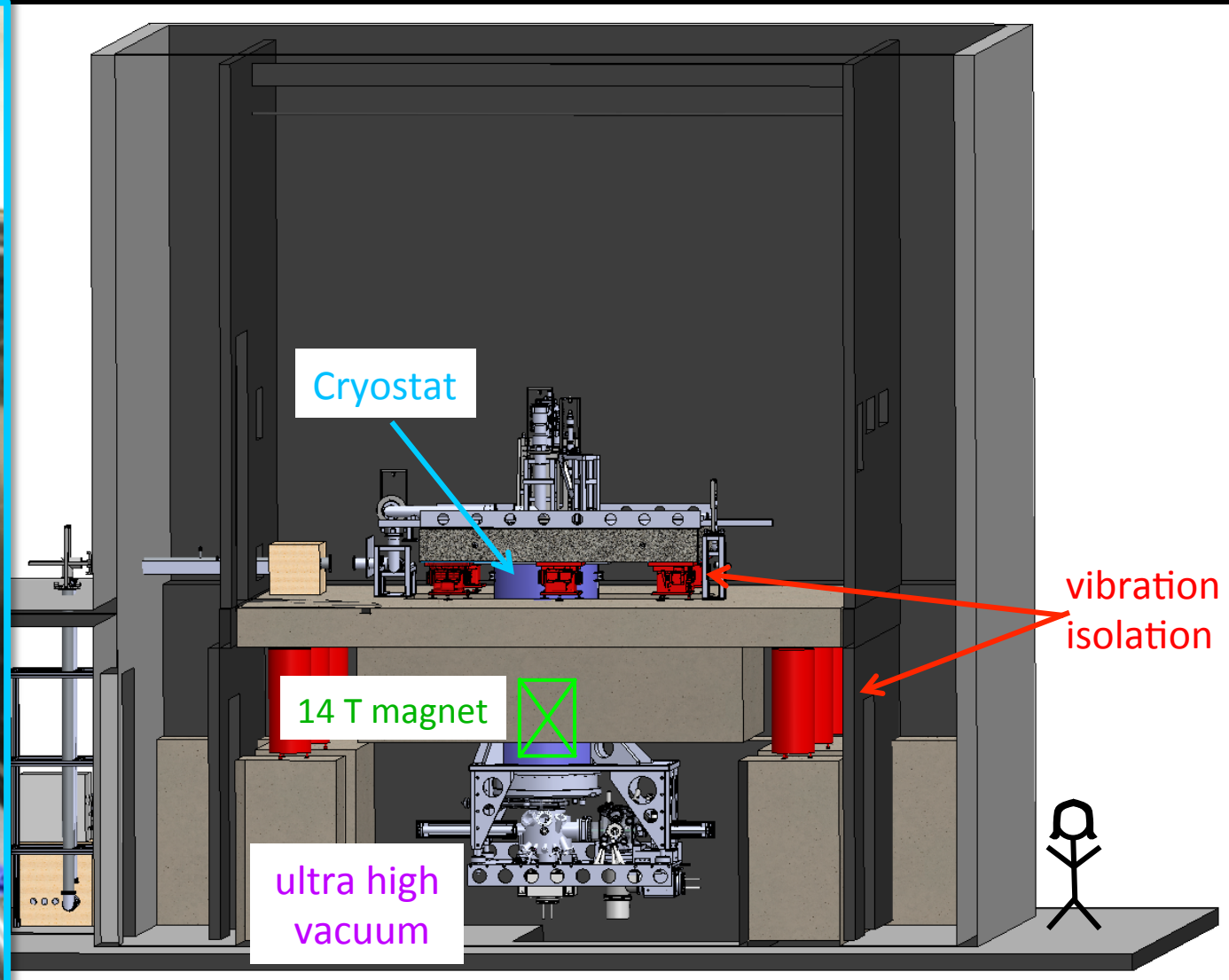
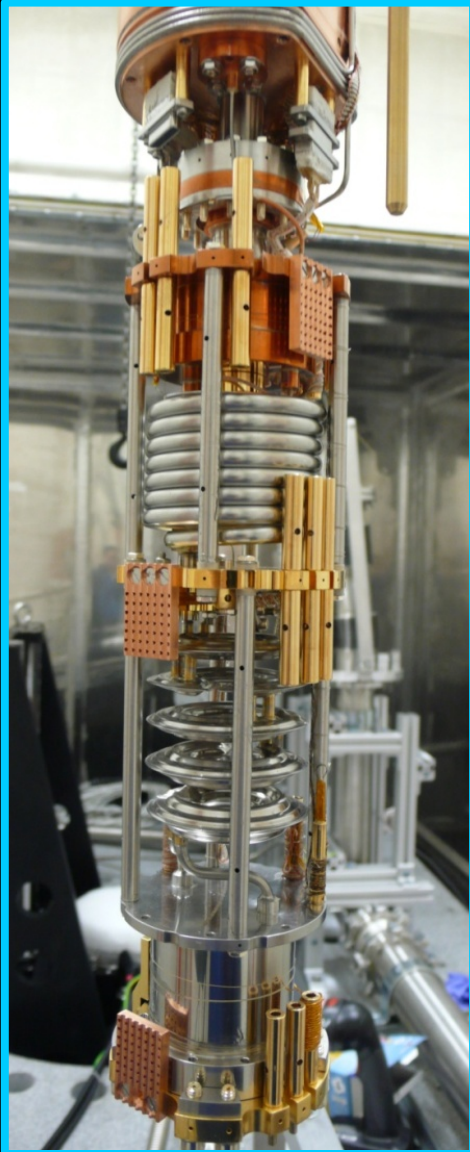
# Our STM in the lab

Ultra-low temperature high magnetic field STM

Dilution refrigerator

$$T_{\text{phonon}} = 20 \text{ mK}$$

$$T_{\text{electron}} = 250 \text{ mK}$$



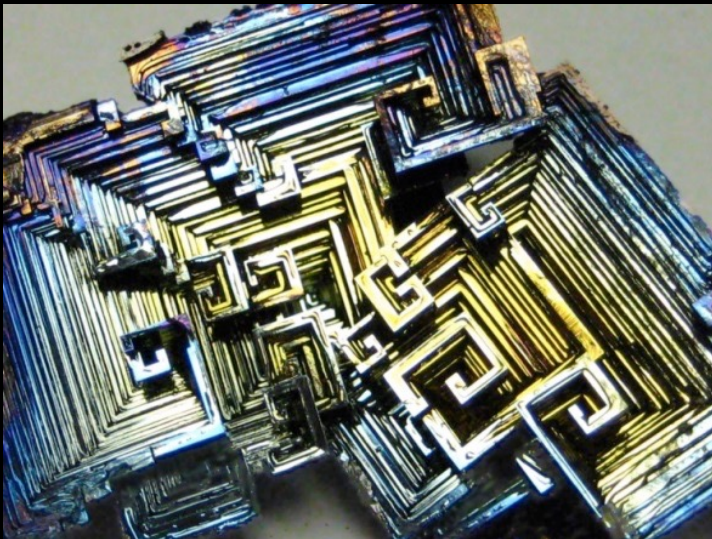
# Bismuth: the pioneering quantum material

83

**Bi**

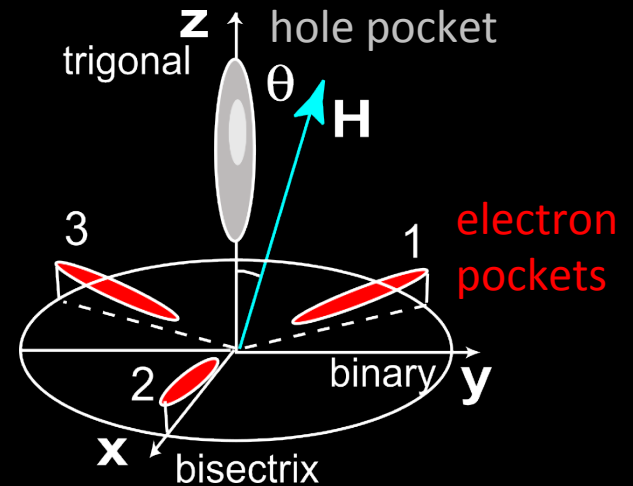
Bismuth  
208.980

- Compensated semimetal
- Extremely clean (high mobility)
- Strong spin-orbit coupling
- Multiple degenerate Fermi pockets in the bulk



Bi single crystal

Bulk Bi Fermi surface



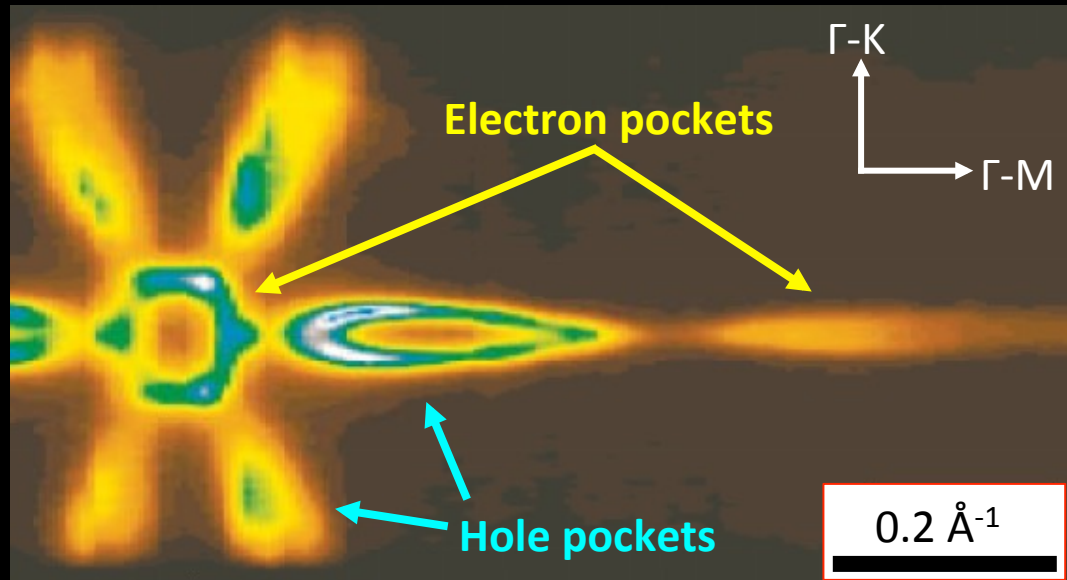
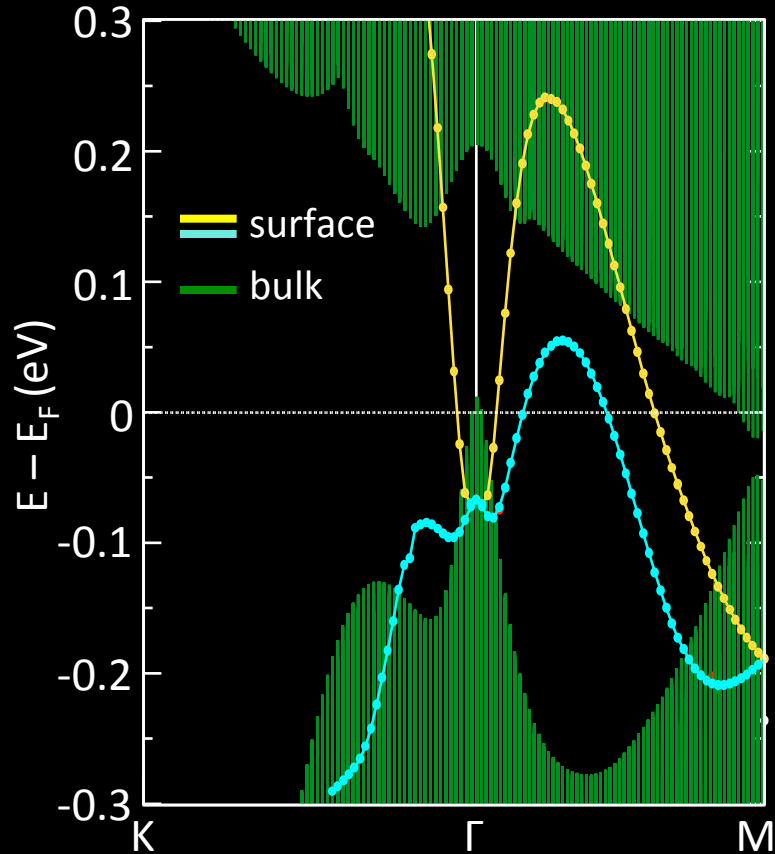
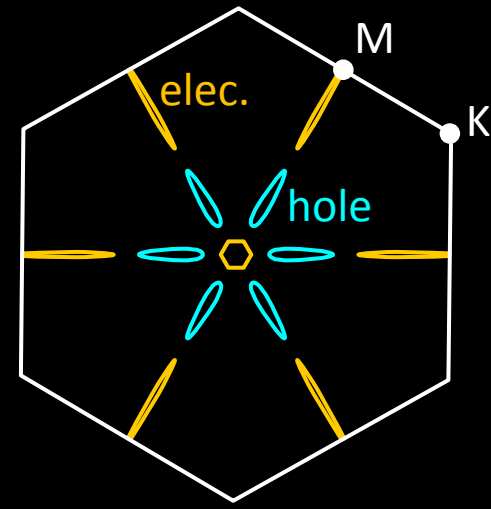
Valley polarization in bulk bismuth  
at high B, low T:

- Torque magnetometry: Li, L. *et al.*, *Science* **321**, 547-550 (2008)  
Magnetoresistance: Zhu, Z. *et al.*, *Nat. Phys.* **8**, 89-94 (2012)  
Magnetostriction: Kuchler, R. *et al.*, *Nat. Mater.* **13**, 461-465 (2014)



# Bi(111) surface states

Large spin-orbit coupling – lifts spin degree of freedom  
Multiple degenerate valleys



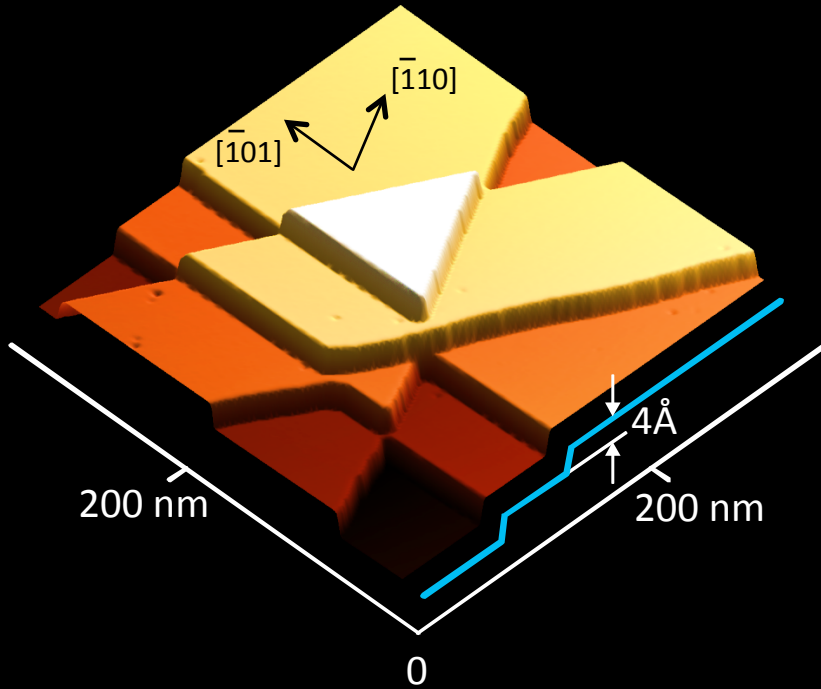
Ast, C. R. and Höchst, H., *Phys. Rev. Lett.* **87**, 177602 (2001)

Koroteev, Yu. M. *et al.*, *Phys. Rev. Lett.* **93**, 046403 (2004)

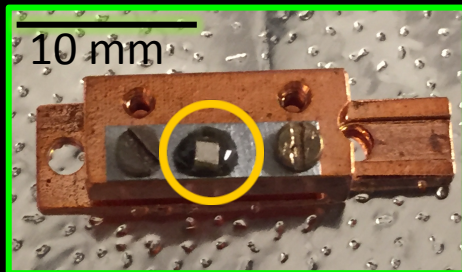
Hirahara, T. *et al.*, *Phys. Rev. Lett.* **97**, 146803 (2006)

# Bi(111) surface

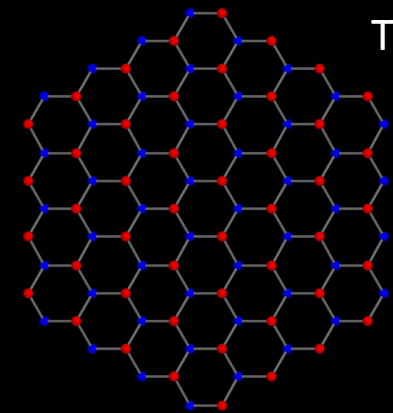
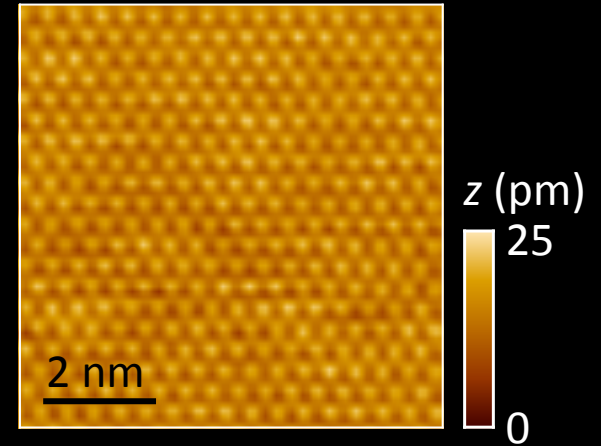
A typical cleaved surface



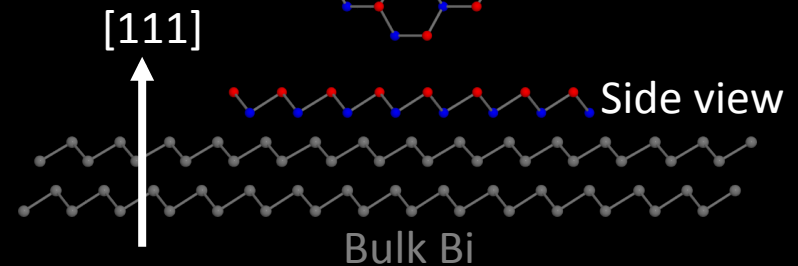
Bi crystal on sample holder



Hexagonal atomic lattice

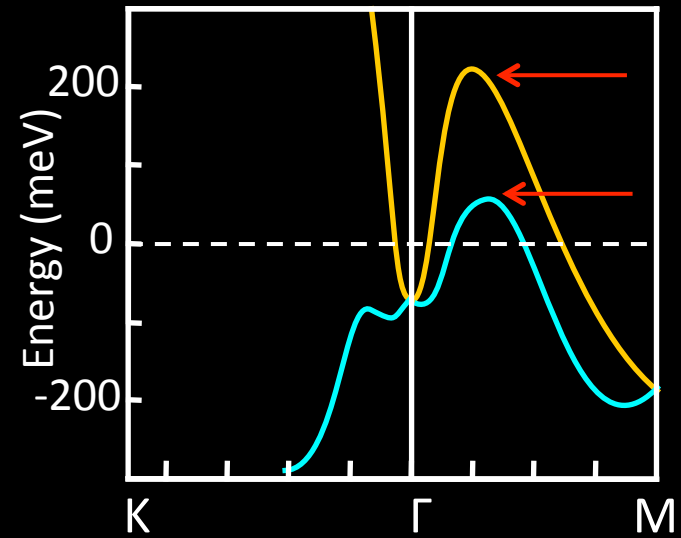
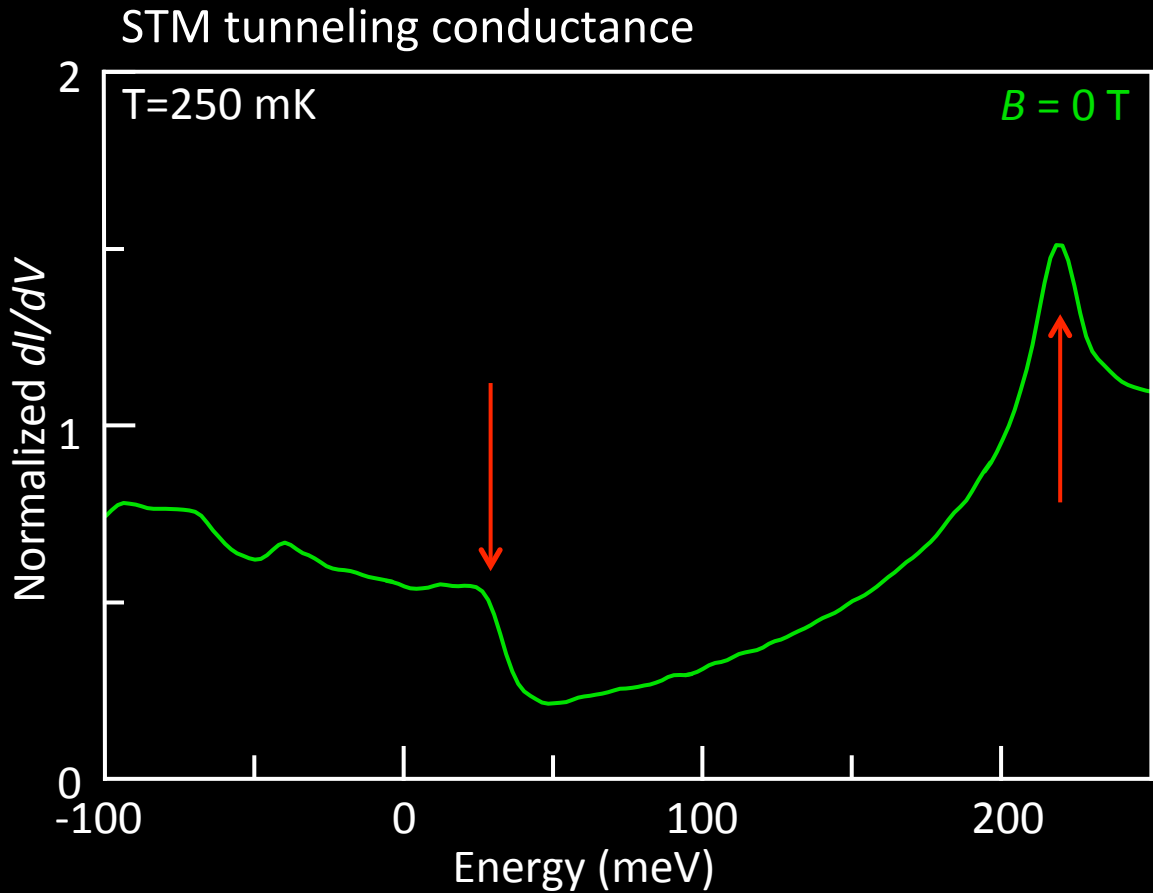


Top view

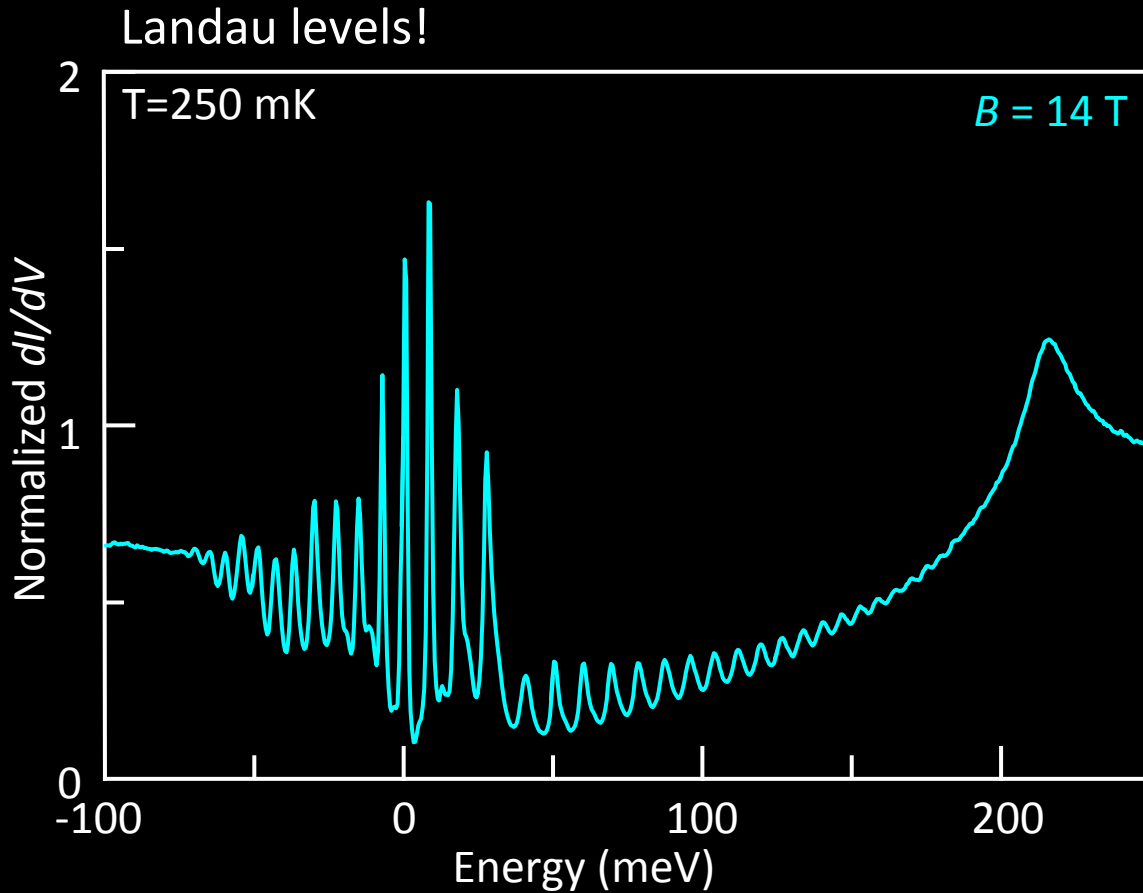


Bulk Bi

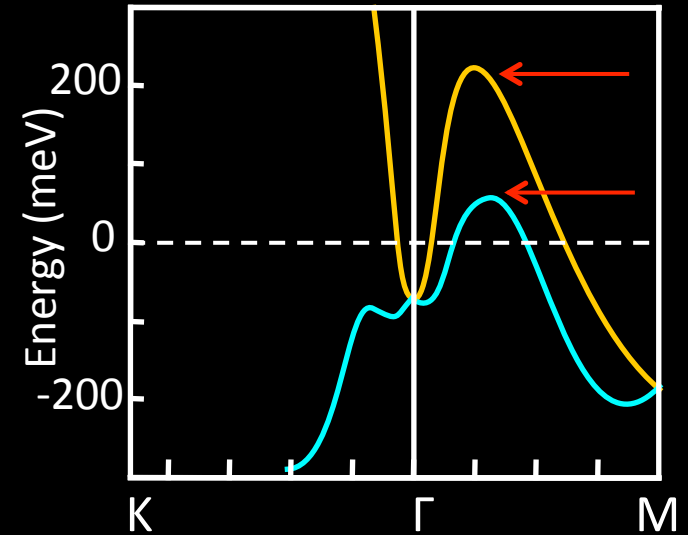
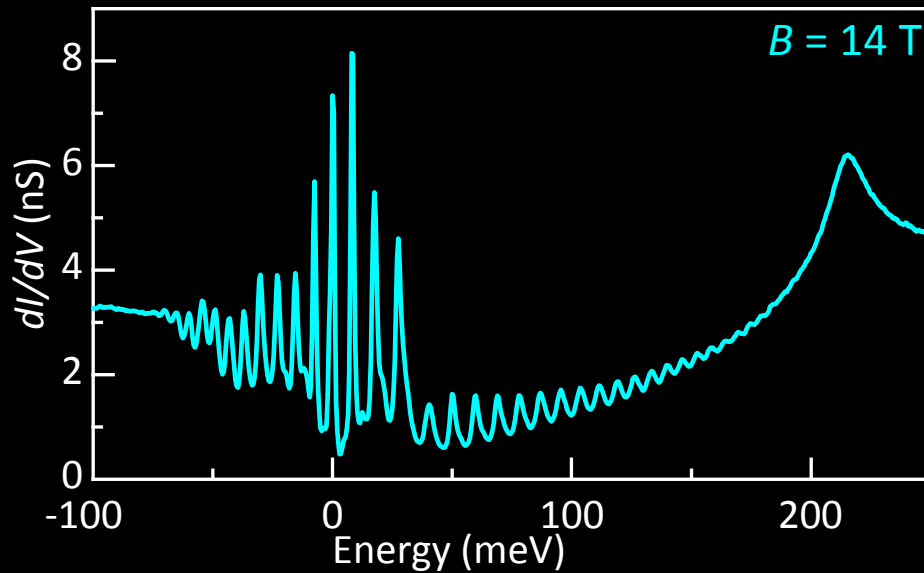
# Quantized energy levels of the bismuth surface



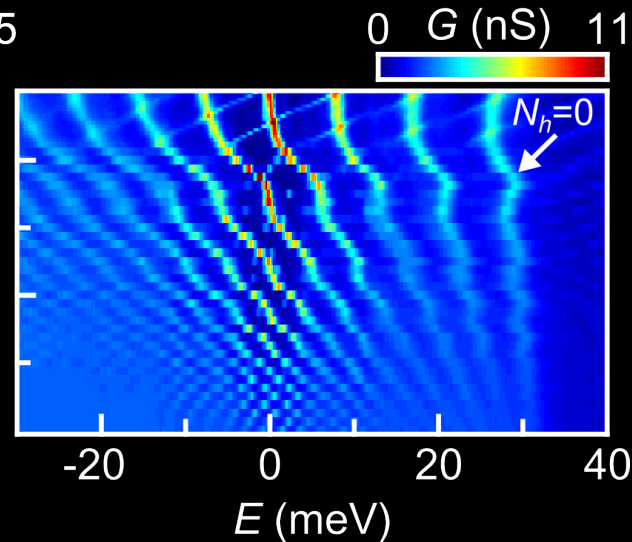
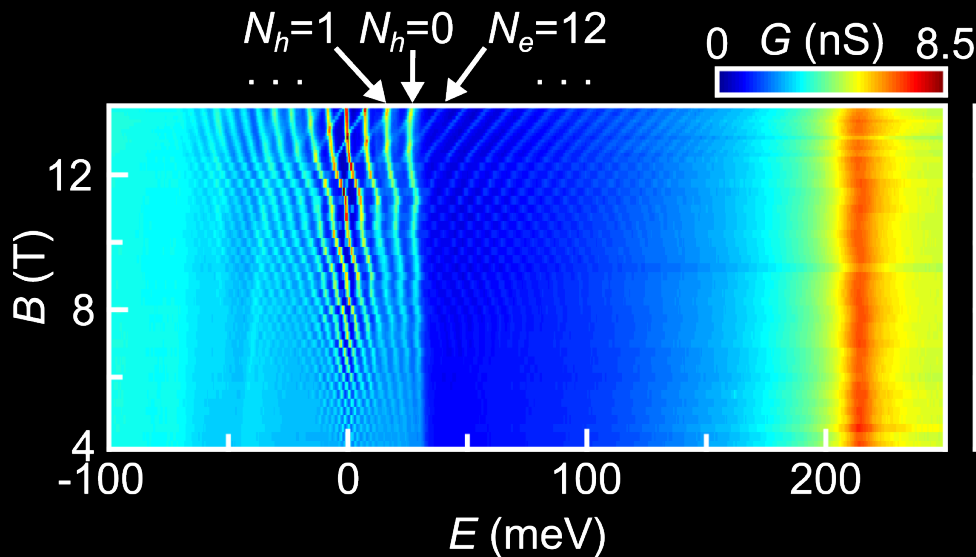
# Quantized energy levels of the bismuth surface



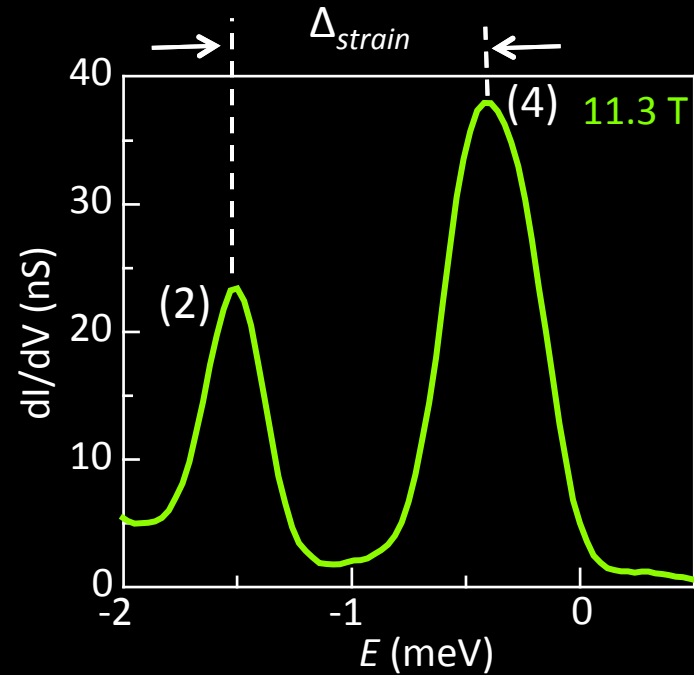
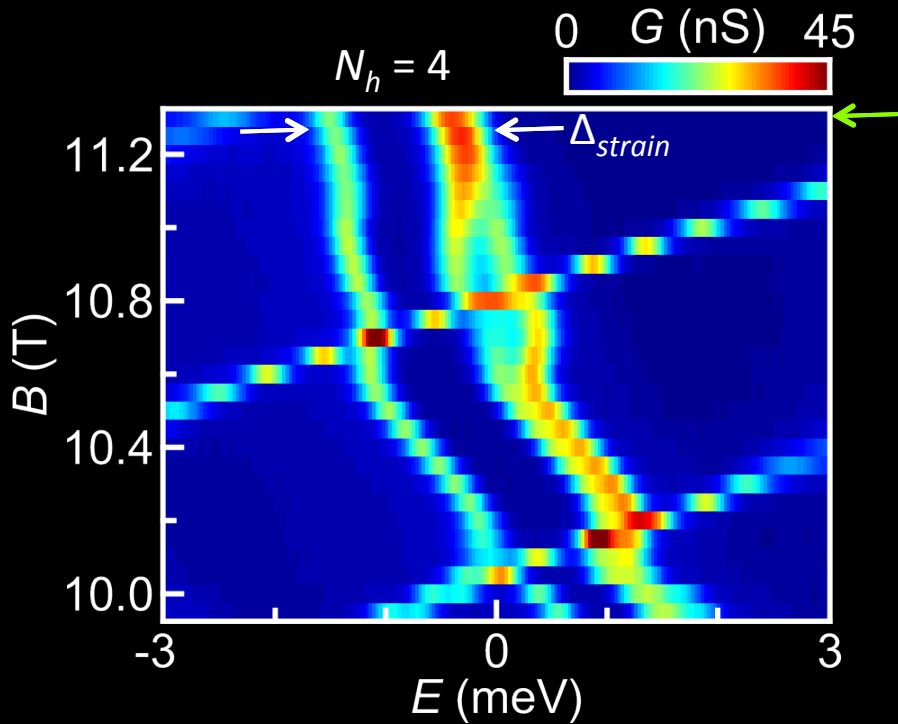
# Surface state Landau levels (LLs)



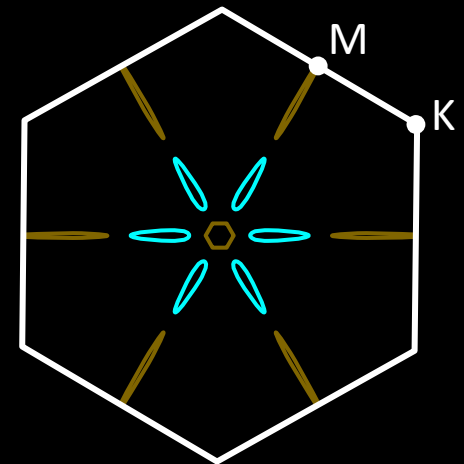
Fixed hole density:  $p \approx 7e12 \text{ cm}^{-2}$



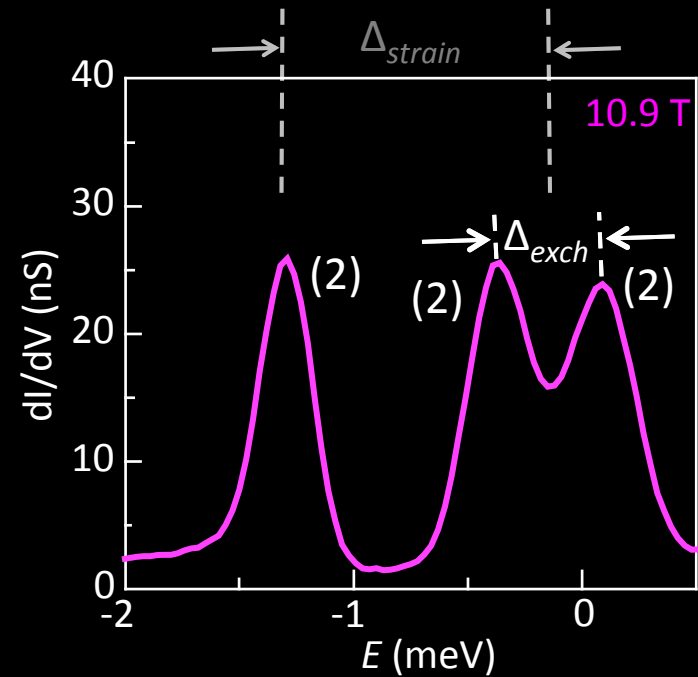
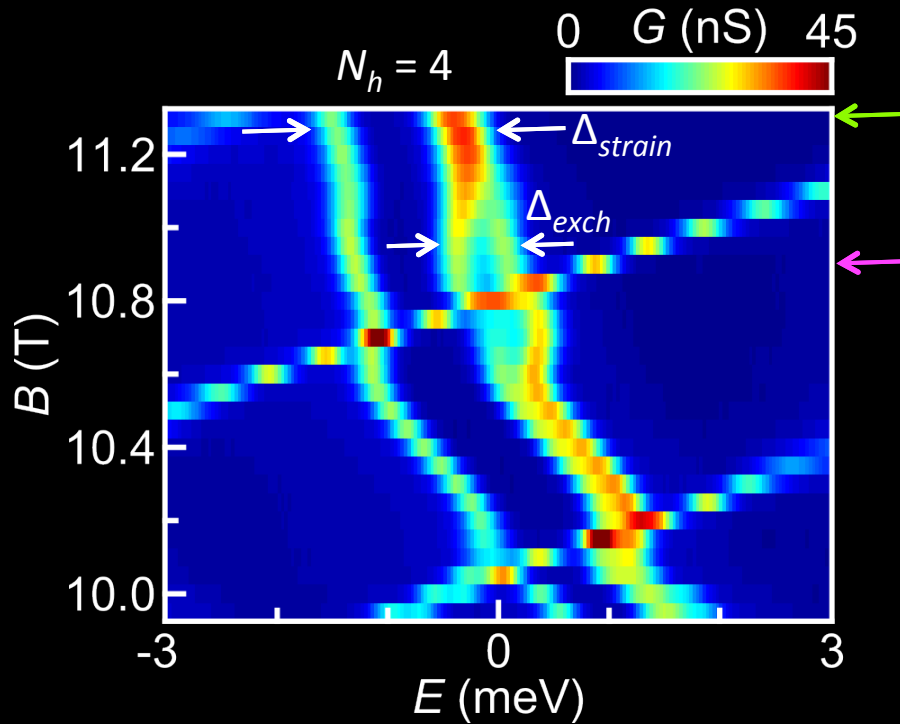
# Lifting of the six-fold valley degeneracy



**Local strain:** independent of Fermi level and orbital index, but depends on position



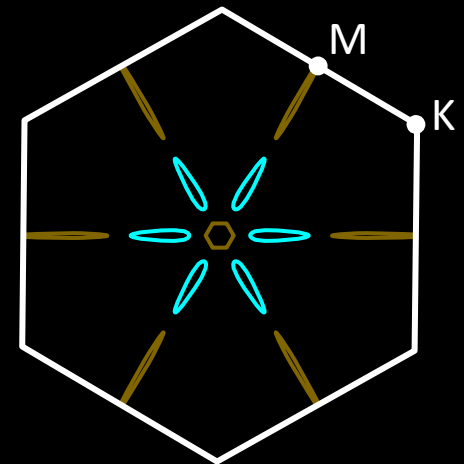
# Lifting of the six-fold valley degeneracy



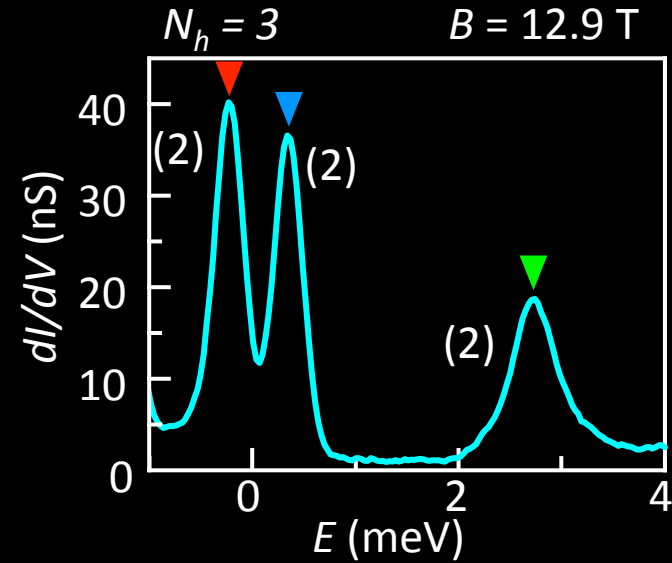
**Local strain:** independent of Fermi level and orbital index, but depends on position

**Exchange interactions:** opens gap *only* at Fermi level

Spectroscopy: quantitative probe of single-particle vs. many-body effects in the formation of these broken symmetry states

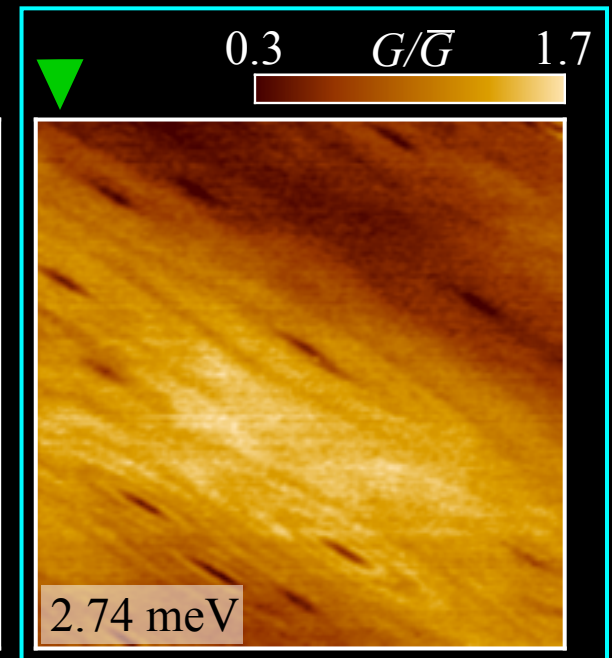
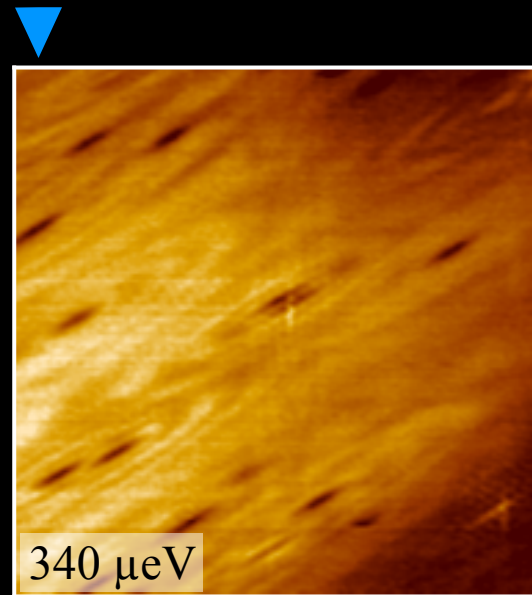
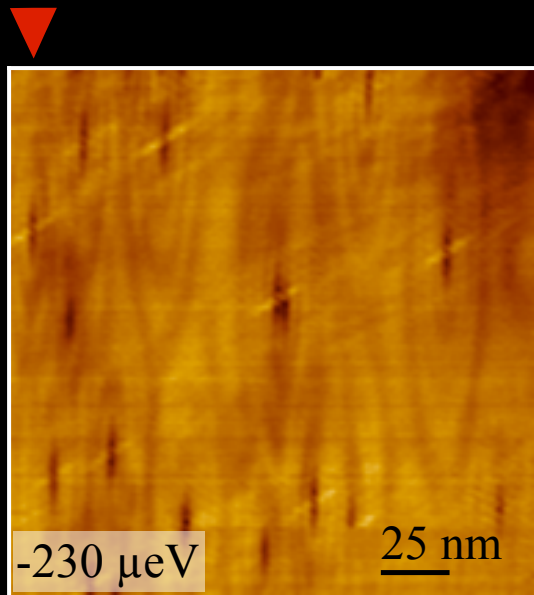
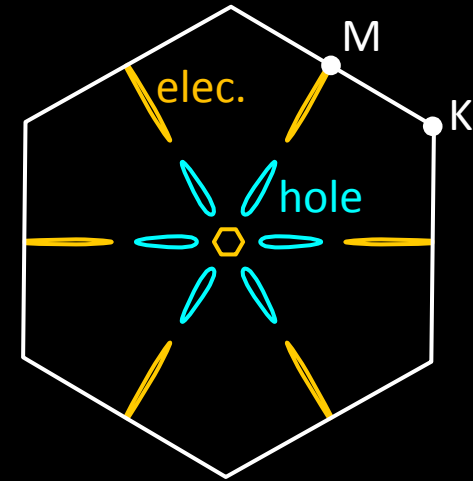


# Real-space imaging of broken symmetry states



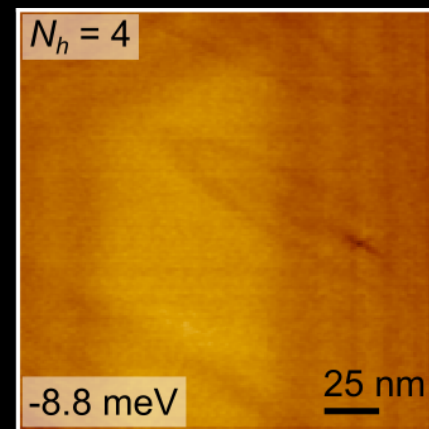
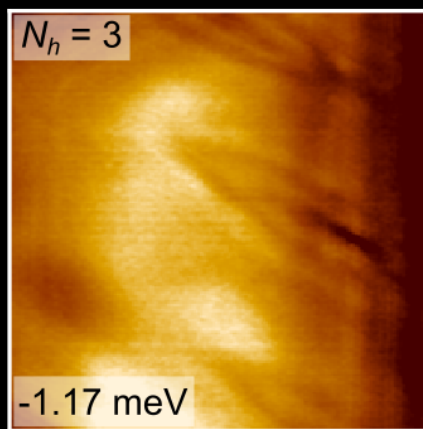
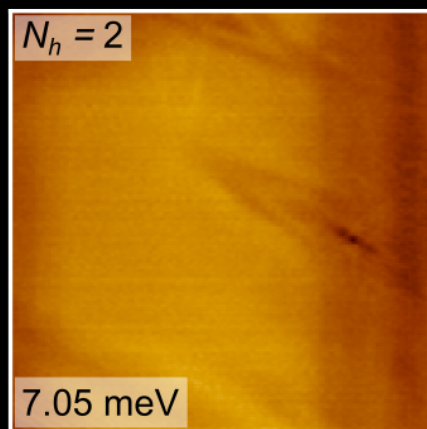
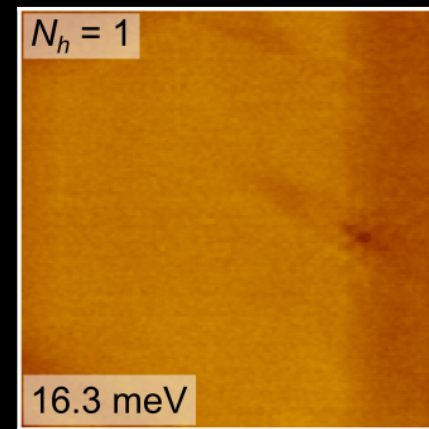
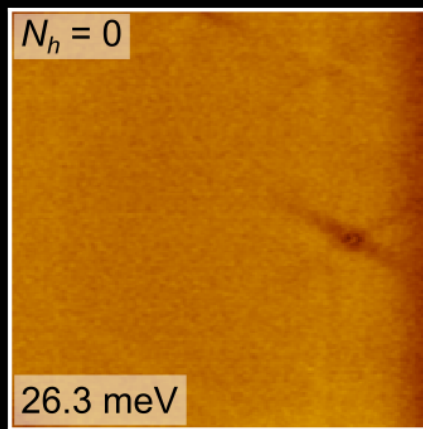
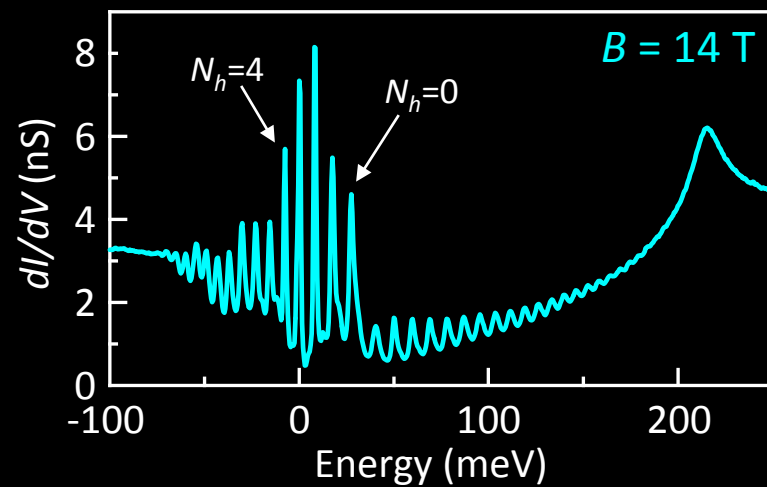
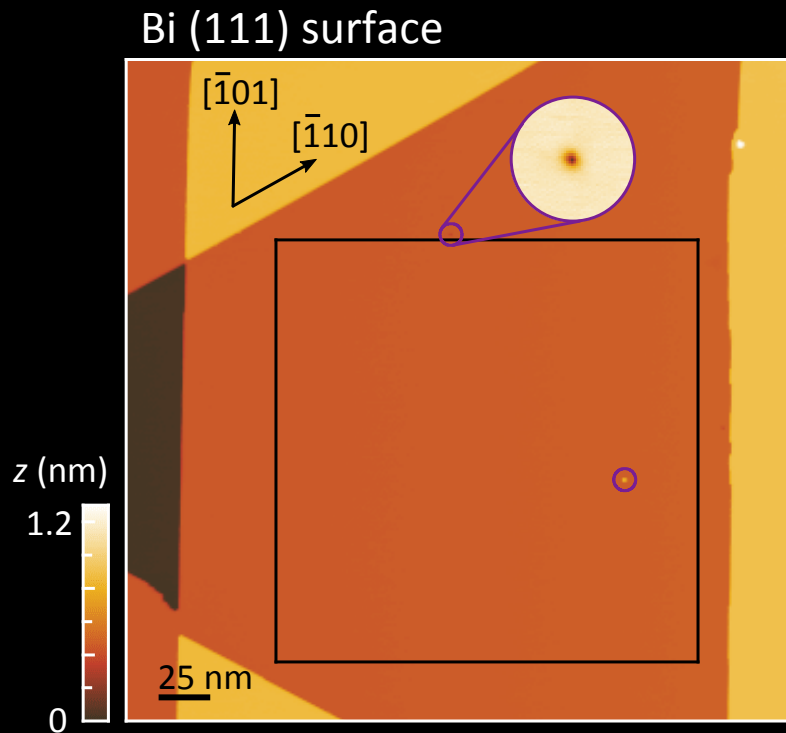
Remainder of talk:

- Imaging individual Landau level wavefunctions
- Observing nematic domains



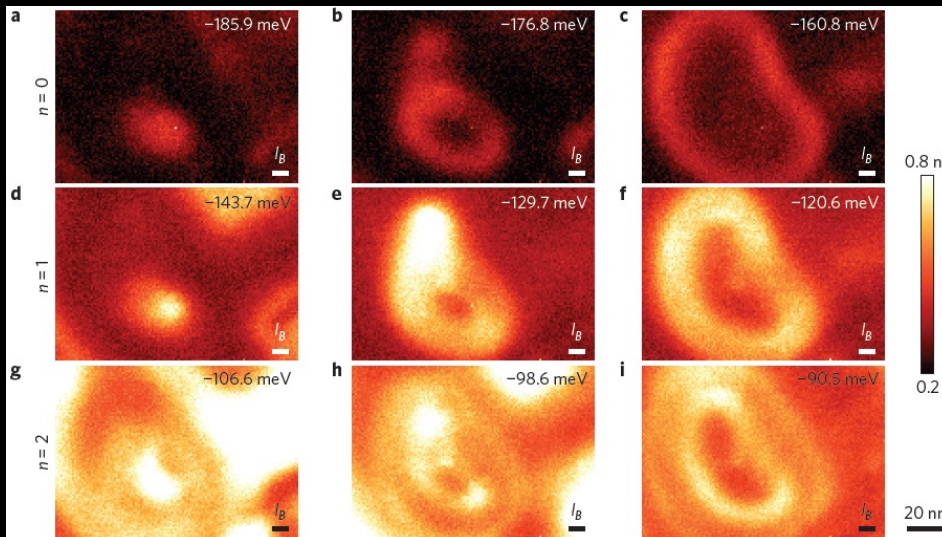
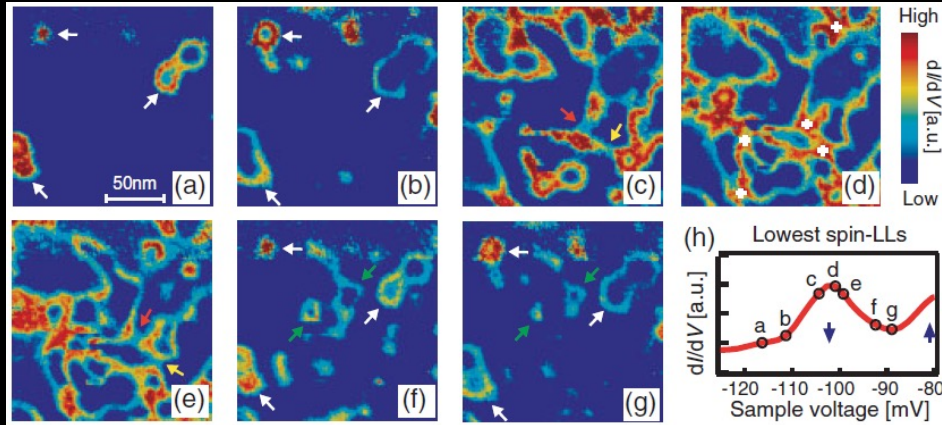


# Orbital dependence of conductance maps

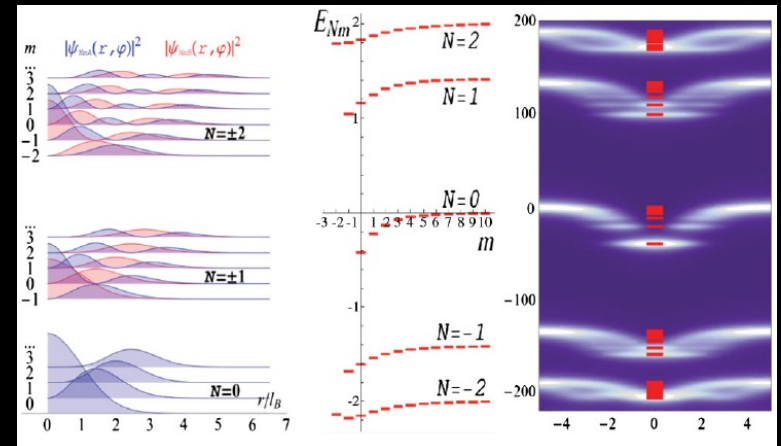
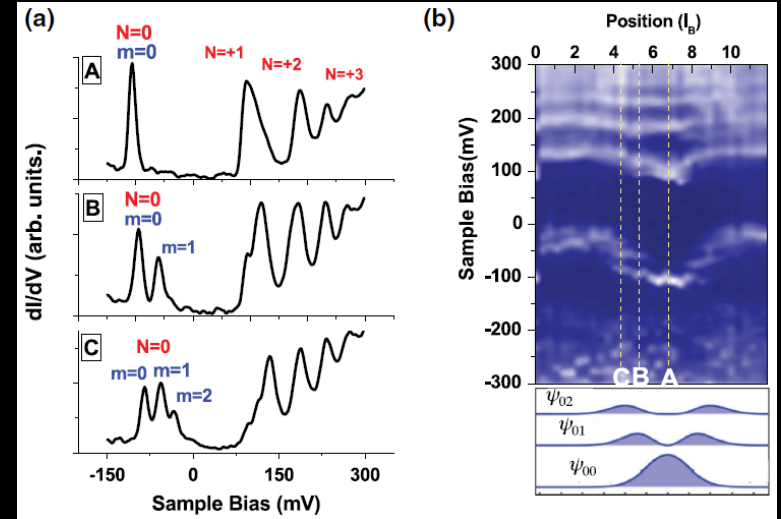


# Previous STM imaging in the quantum Hall regime

More disordered samples: drift states



Evidence of impurity-shifted Landau orbits

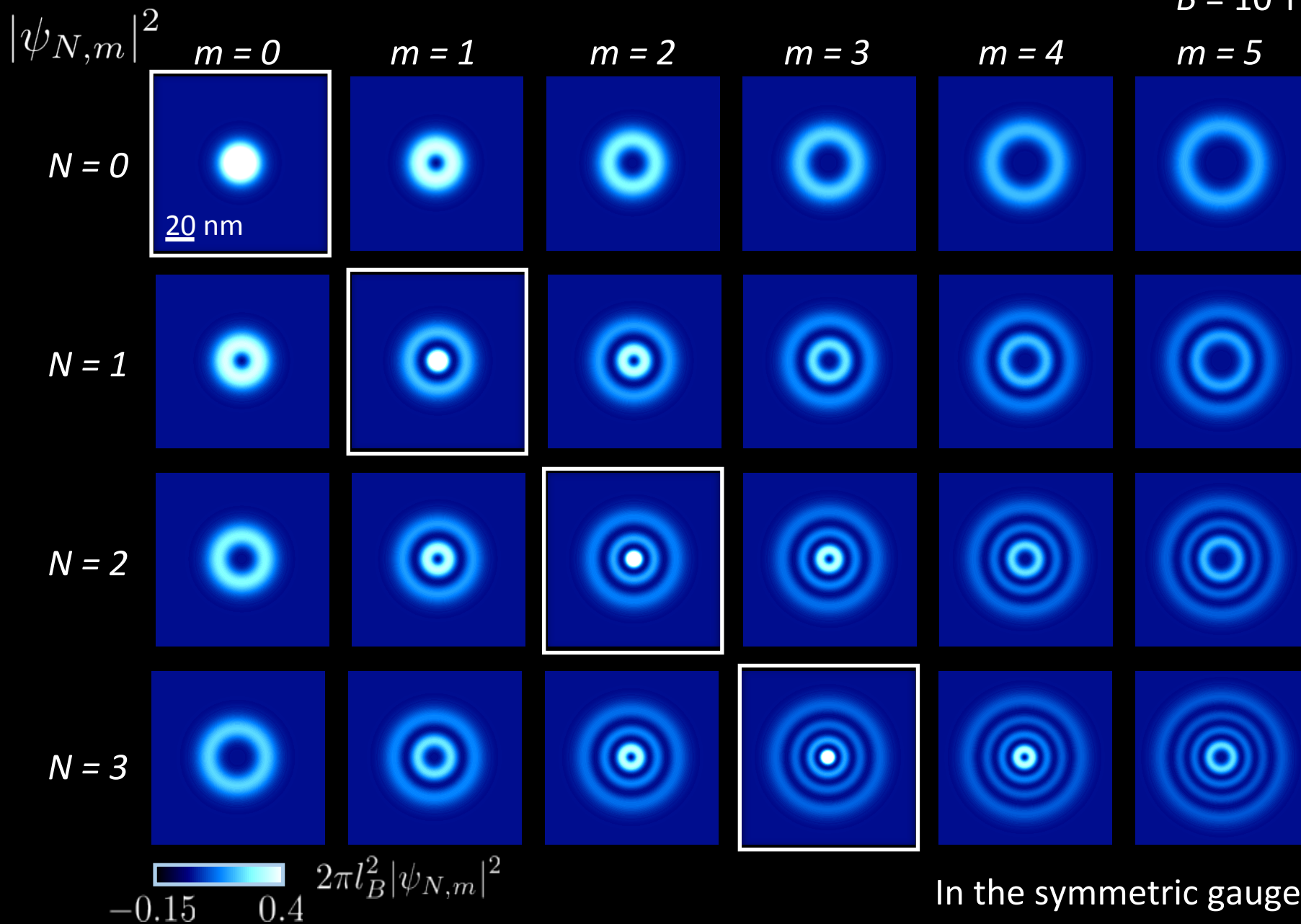


K. Hashimoto *et al.*, *Phys. Rev. Lett.* **101**, 256802 (2008)  
 Y.-S. Fu *et al.*, *Nat. Phys.* **10**, 815-819 (2014)

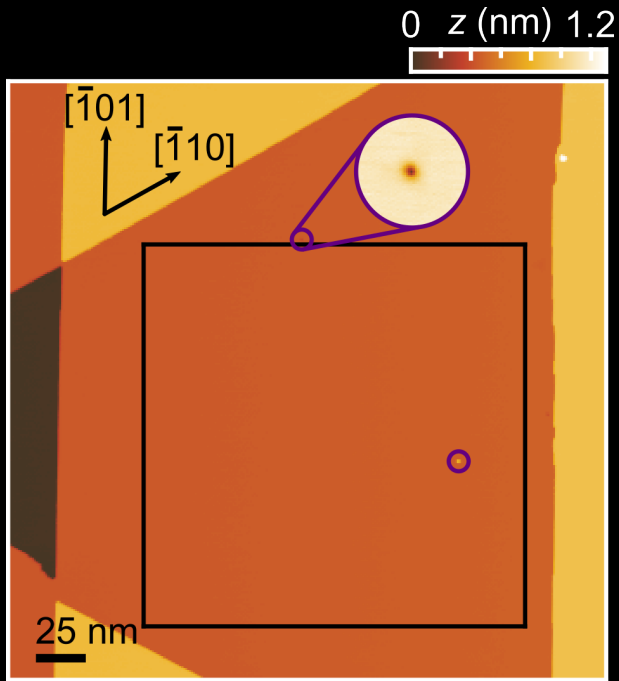
A. Luican-Mayer *et al.*, *Phys. Rev. Lett.* **112**, 036804 (2014)  
 see also: Y. Okada *et al.*, *Phys. Rev. Lett.* **109**, 166407 (2012)

# Landau Level wavefunctions

$B = 10 \text{ T}$



# Effect of a defect



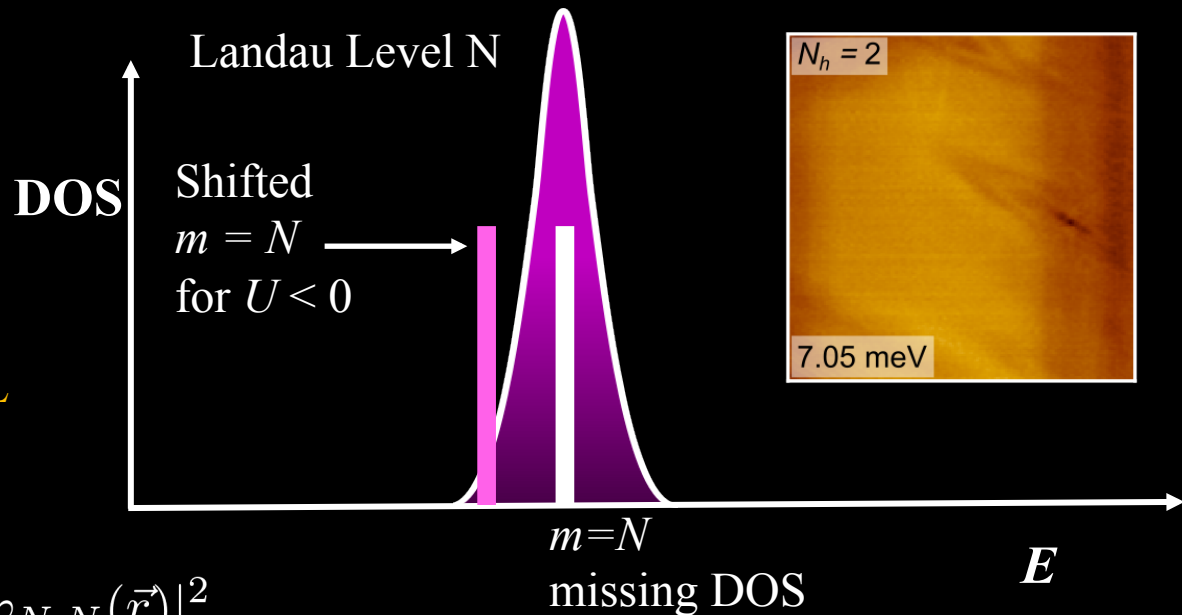
Model atomic-scale defect as a delta function potential

$$U = \alpha \delta(\vec{r})$$

Only  $m = N$  cyclotron orbit is shifted in energy

$$E_{N,N} - E_{N,m \neq N} = \alpha |\varphi_{N,N}(0)|^2 = \frac{\alpha}{2\pi l_B^2}$$

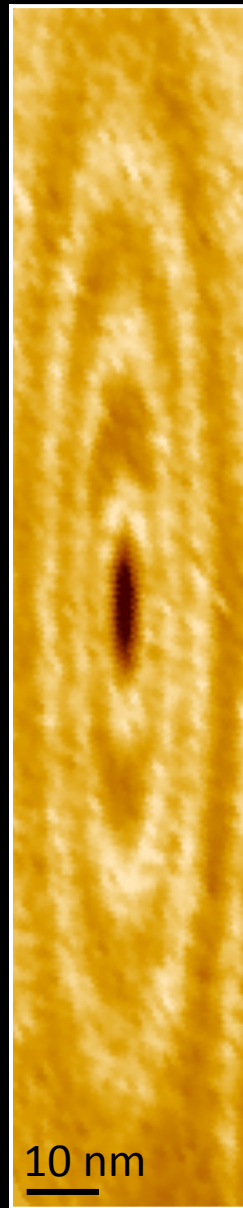
Conductance map measured at LL peak has decreased signal due to shifted cyclotron orbit.



$$\sum_{m \neq N} |\varphi_{N,m}(\vec{r})|^2 = \frac{1}{2\pi l_B^2} - |\varphi_{N,N}(\vec{r})|^2$$

# Imaging an individual Landau orbit

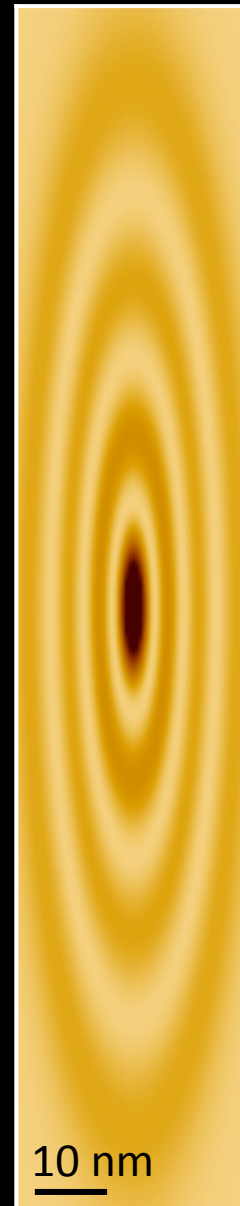
Experiment  
 $B = 12.75 \text{ T}$   
 $E = 770 \text{ } \mu\text{eV}$



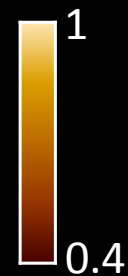
$dI/dV$  (nS)



Simulation for  
 $N=3$  Landau level  
 $\sqrt{m_x/m_y} = 5$

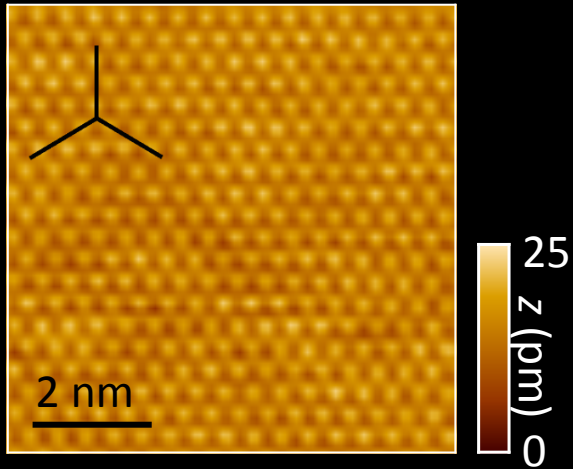


$1 - |\psi|^2$

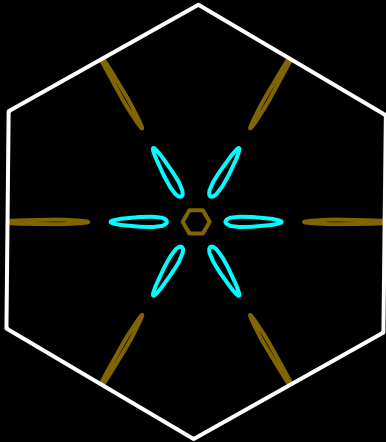


# Symmetry of the Bi(111) surface

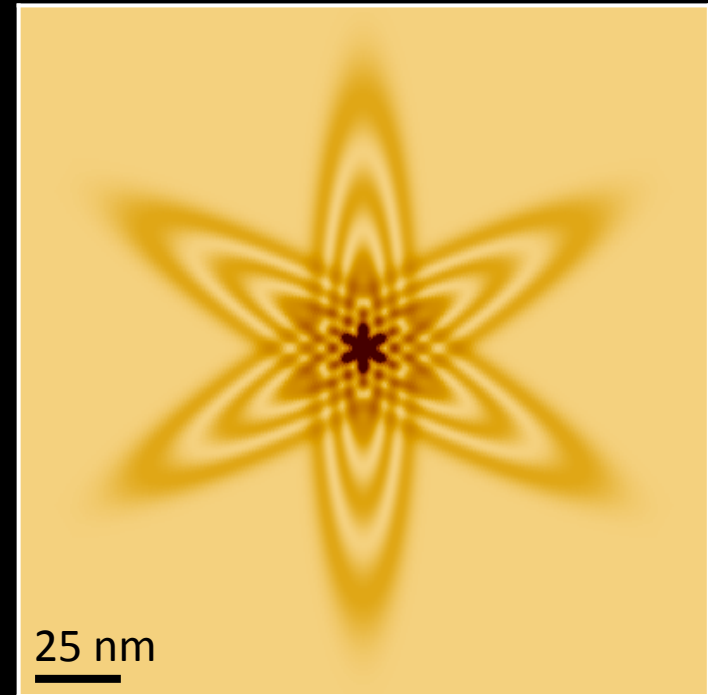
atomic lattice



Fermi surface

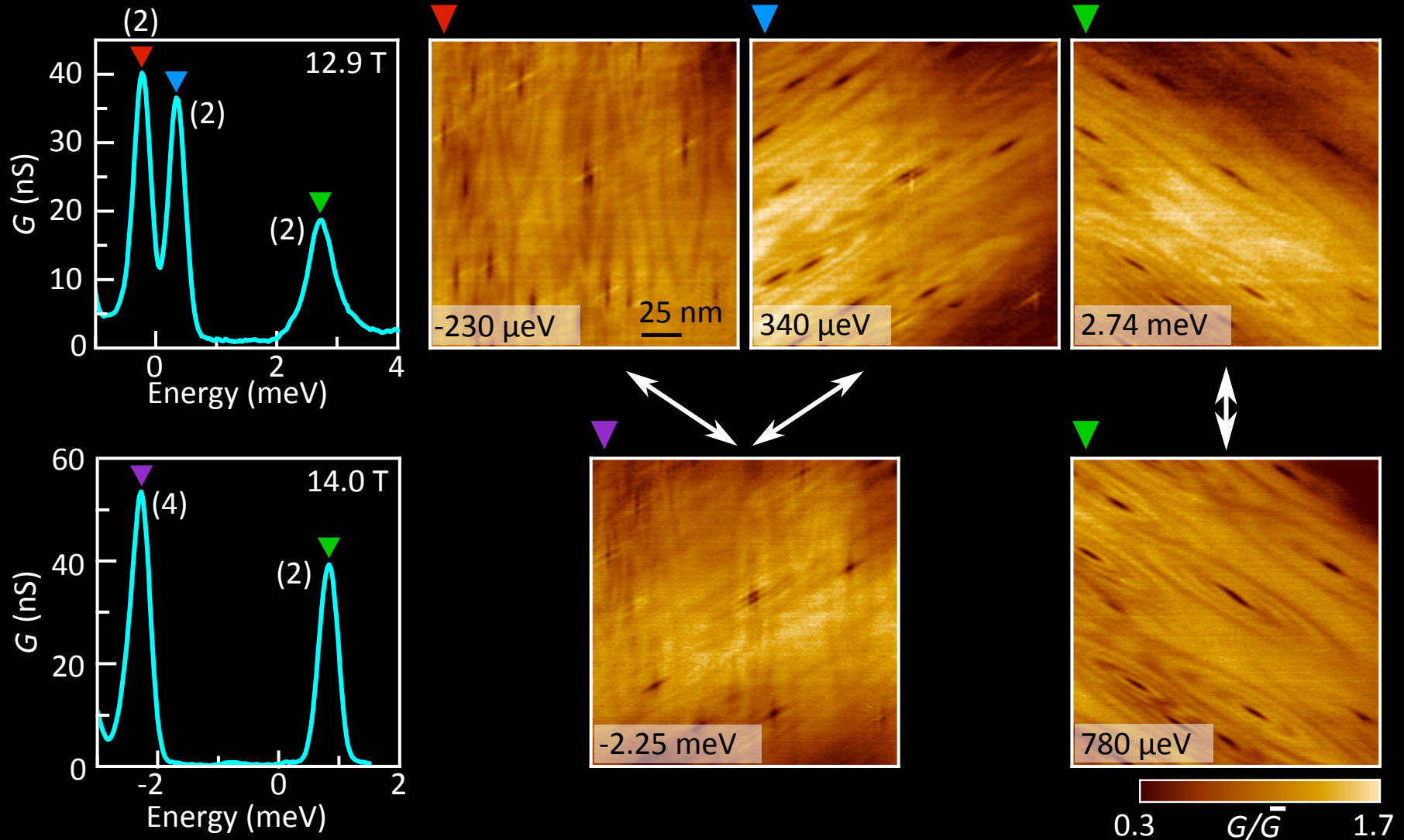


Expected electron wavefunction in the absence of symmetry breaking



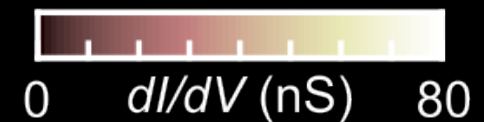
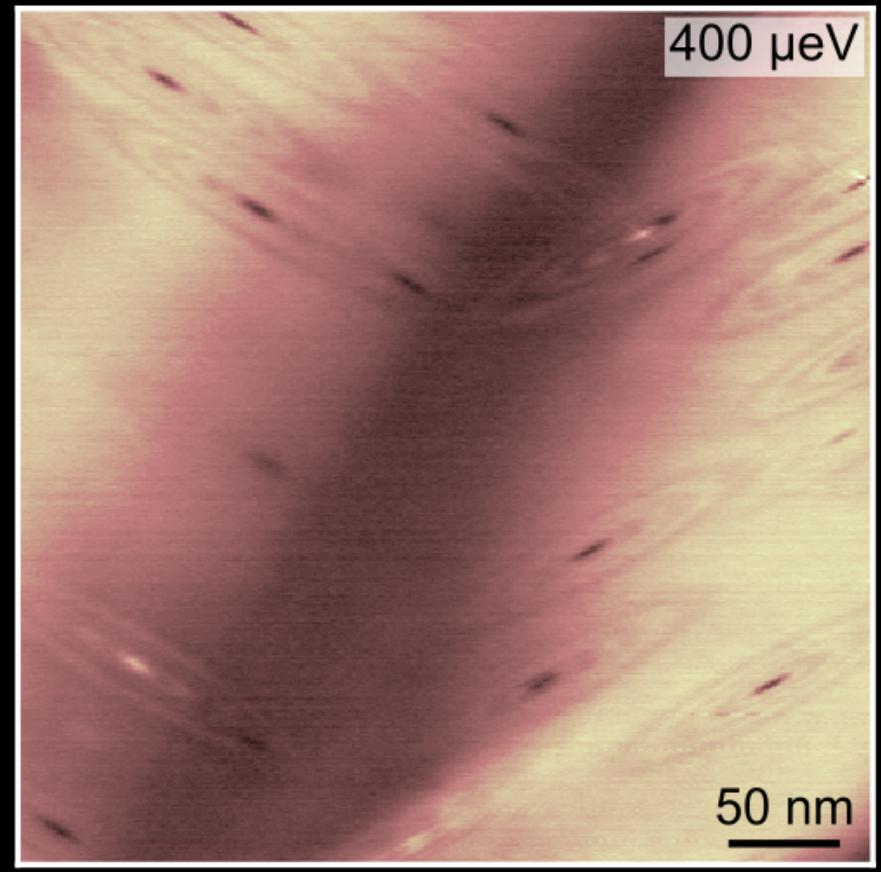
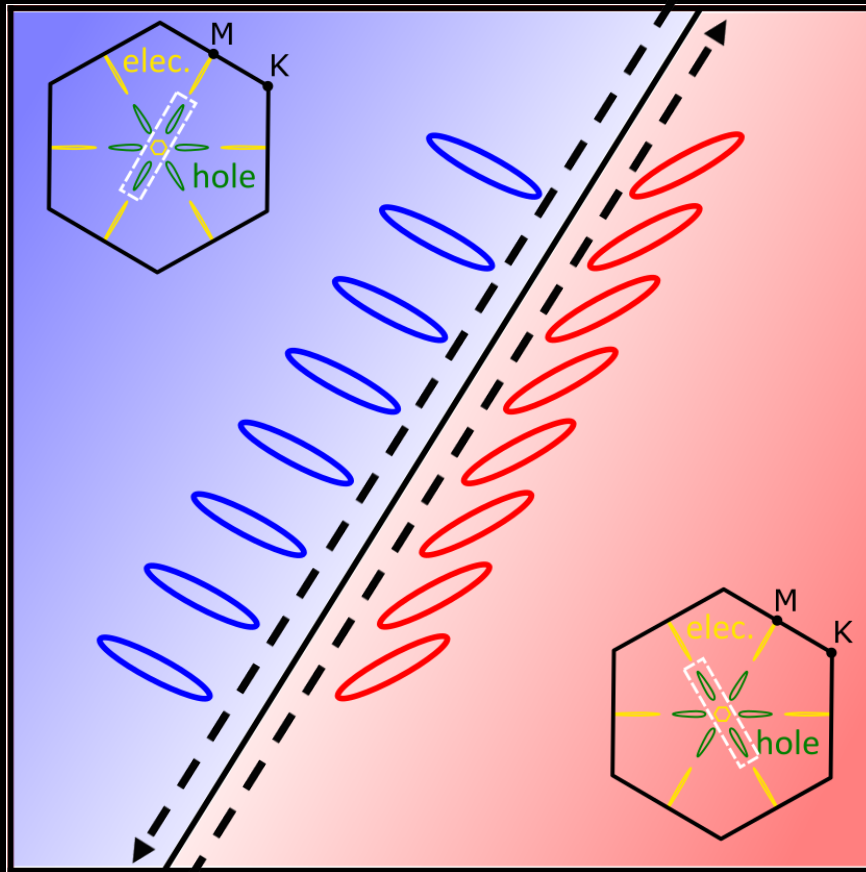
# Interaction driven nematic quantum phase

Spontaneous valley-symmetry breaking: image unidirectionality of wavefunctions  
Quantify roles of single-particle and many-body effects



# The boundary between nematic domains

- 1D counter-propagating chiral edge modes, Luttinger liquid behavior



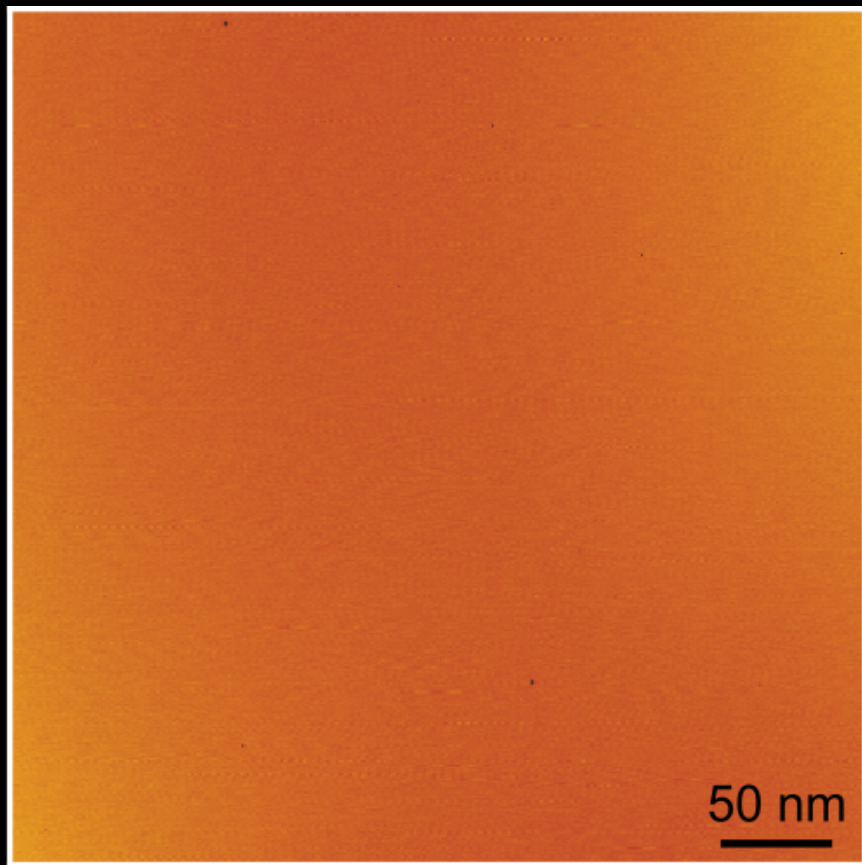
D. A. Abanin *et al.*, PRB 82, 035428 (2010)

A. Kumar *et al.*, PRB 88, 045133 (2013)

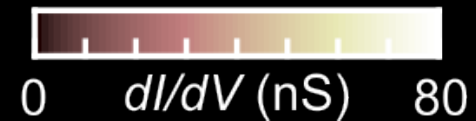
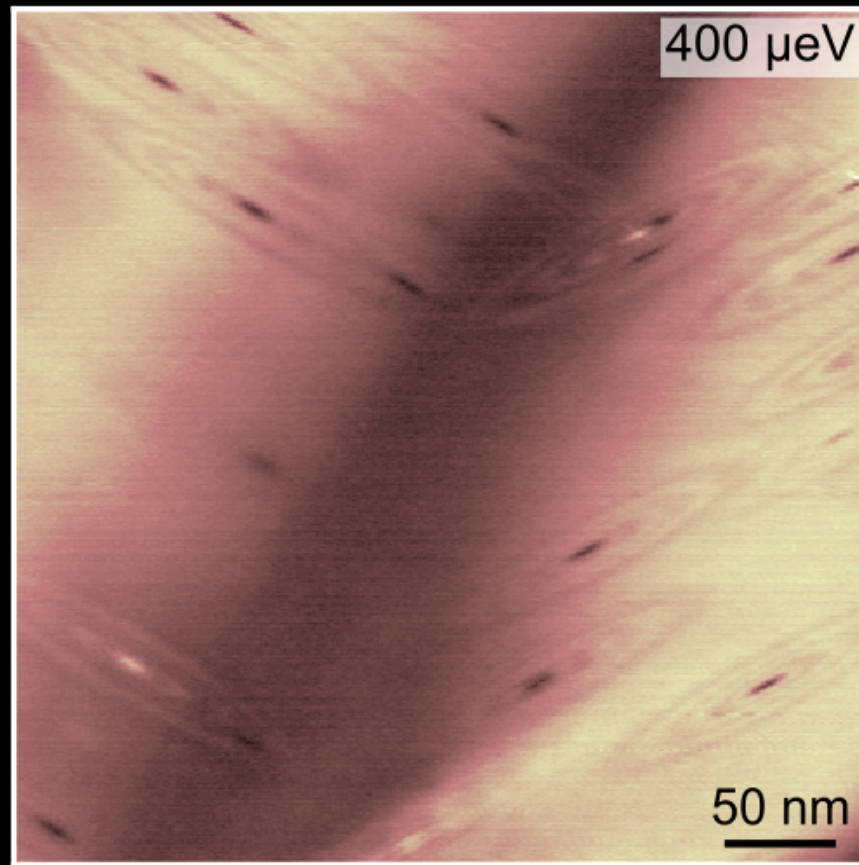


# A pristine domain wall

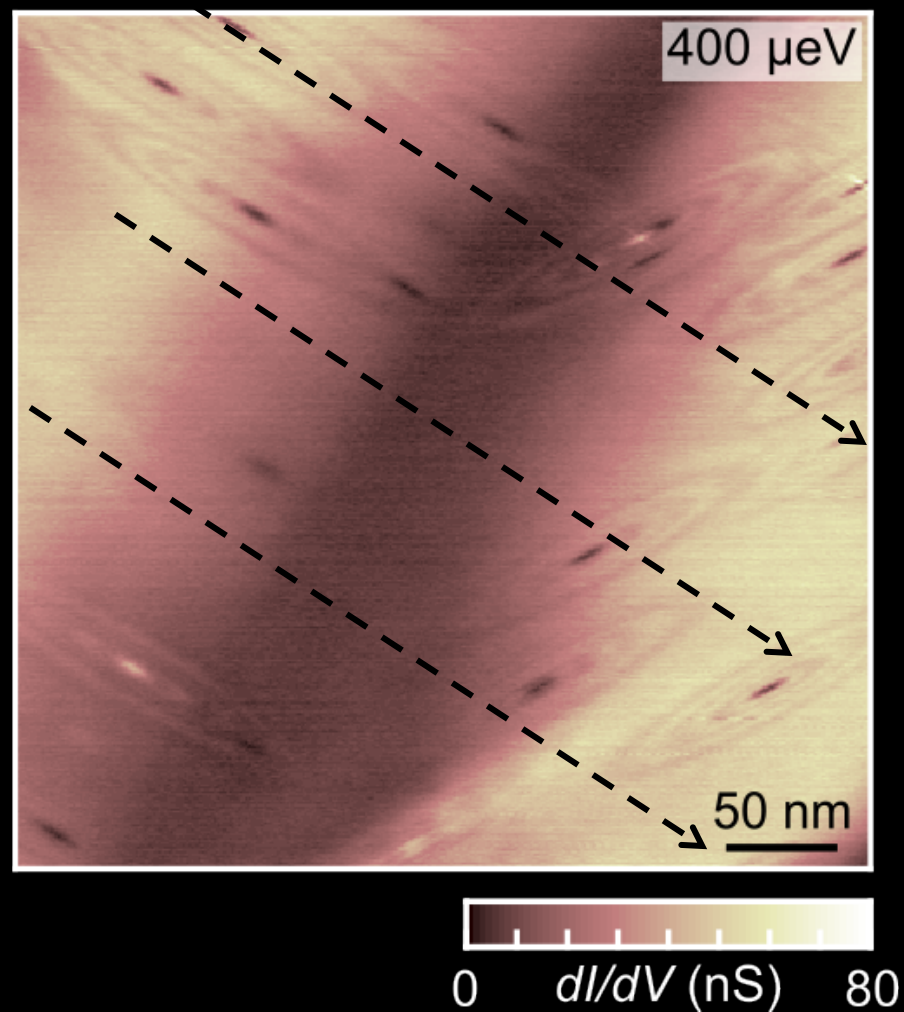
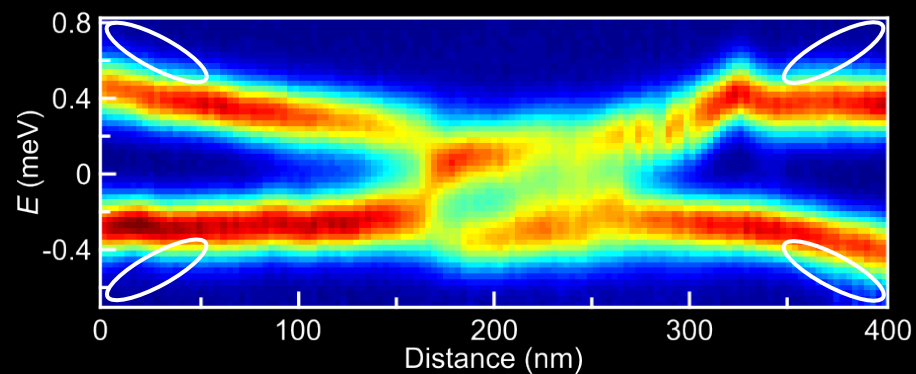
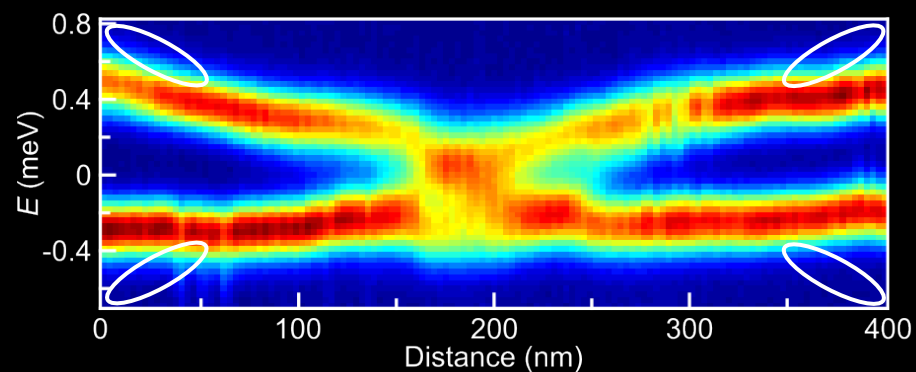
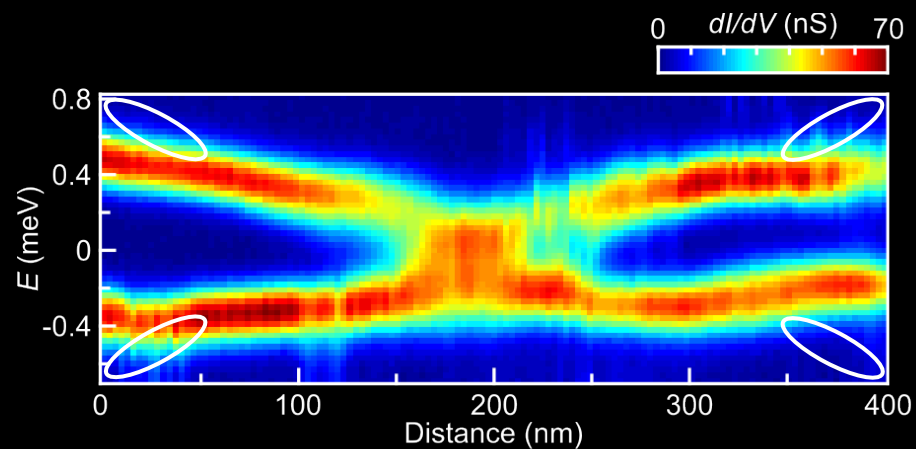
Featureless topography



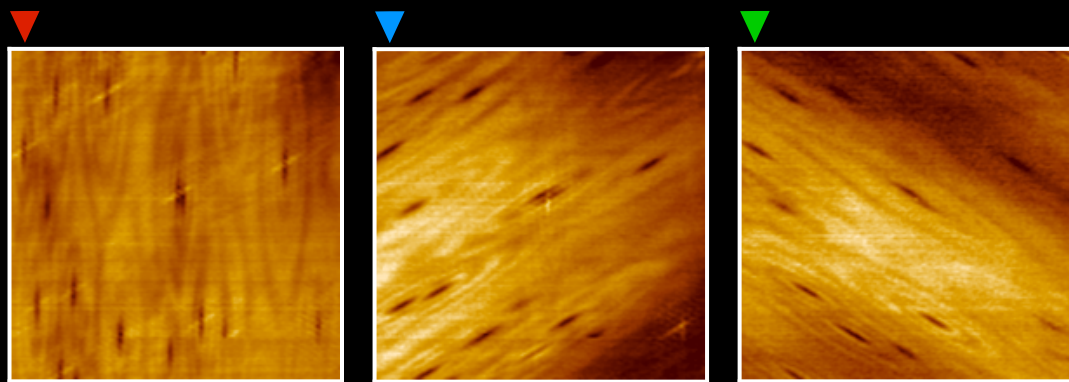
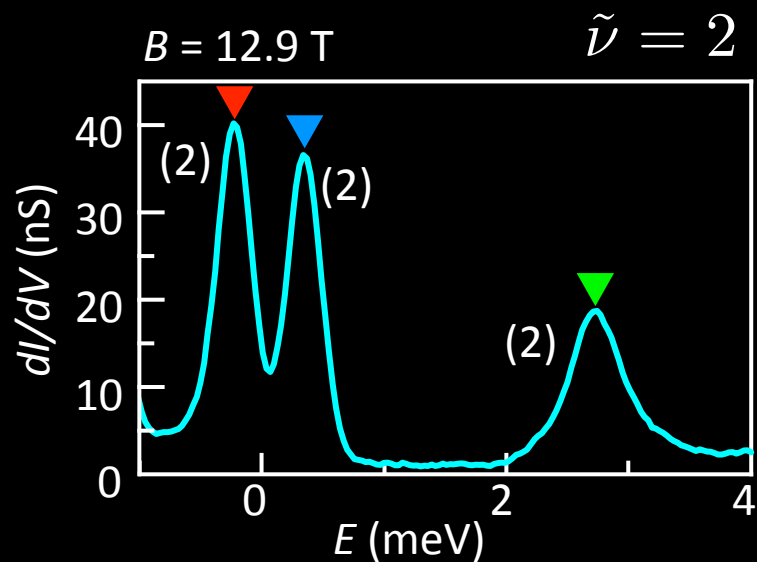
Nematic domains in conductance maps



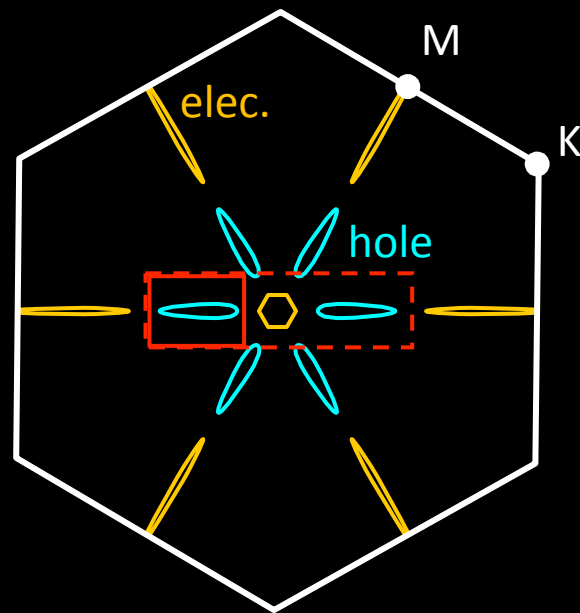
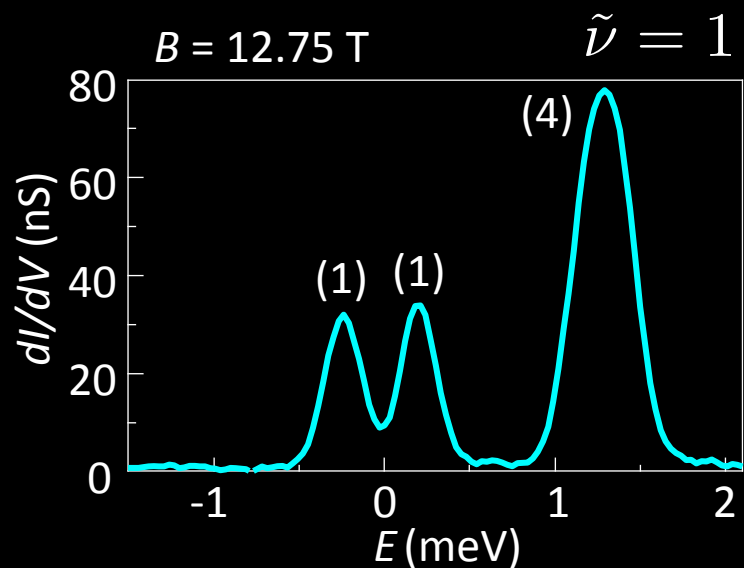
# Linecuts across nematic domain wall



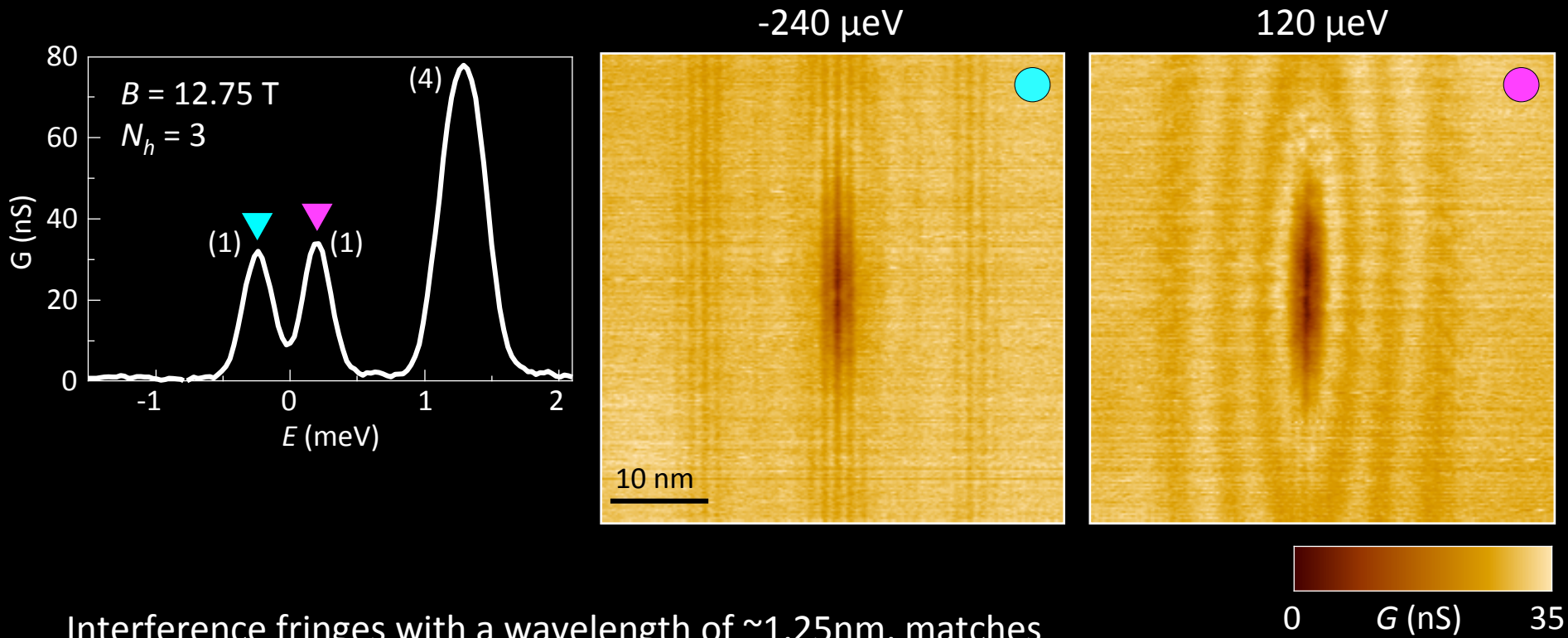
# Imaging additional broken symmetry states



Driven by electron-electron interactions

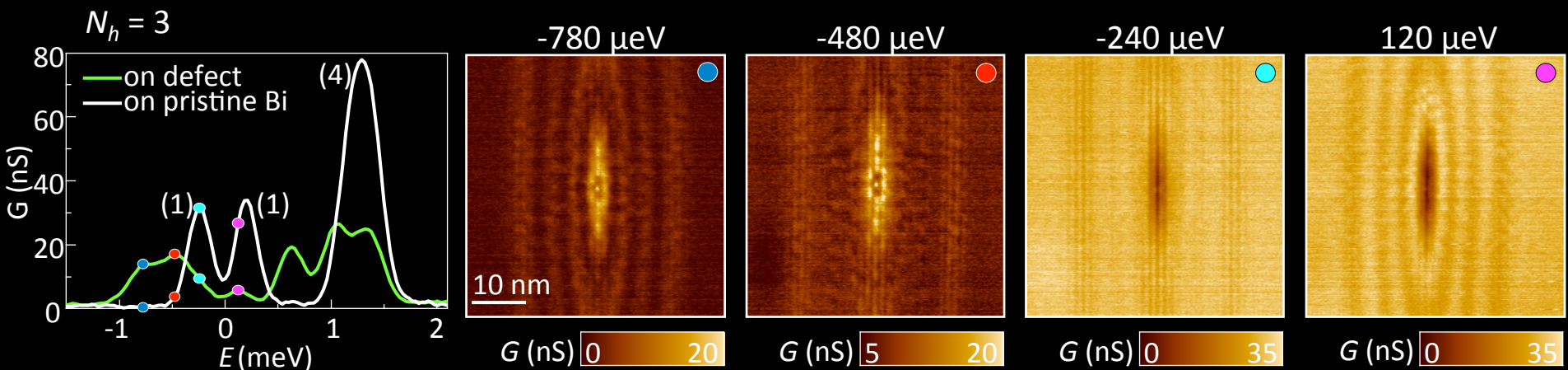


# Higher resolution wavefunction mapping

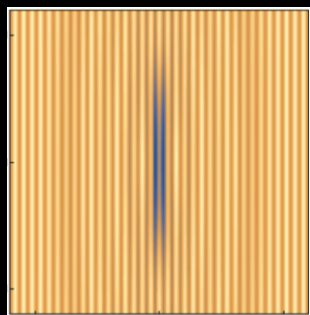


Interference fringes with a wavelength of  $\sim 1.25$  nm, matches well to corresponding valley separation in momentum space

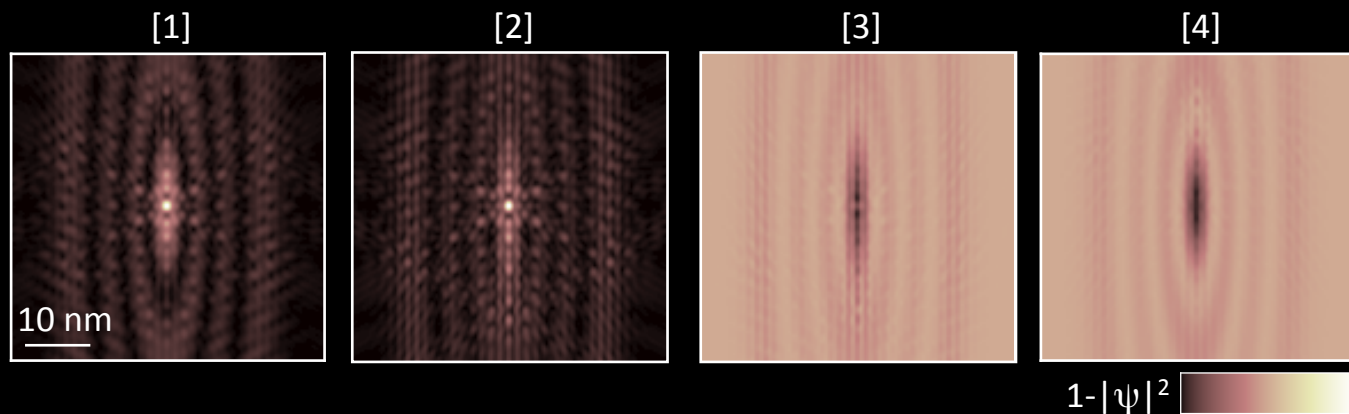
# Valley-polarized ferroelectric ground state



Theoretical simulations for a *valley-polarized* ground state, including defect induced mixing of LL wavefunctions from different valleys



Valley-coherent simulation



# Acknowledgements



## Experiment



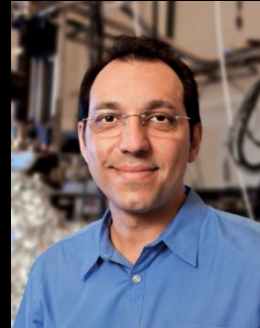
Ben  
Feldman



Andras  
Gyenis



Hao Ding



Ali  
Yazdani

## Sample growth

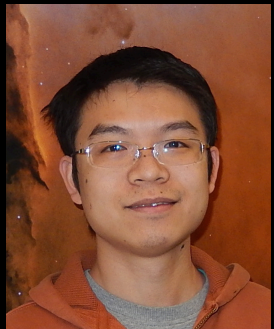


Huiwen Ji



Bob Cava

## Theoretical support



Fengcheng  
Wu



Allan  
MacDonald



Kartiek  
Agarwal



Sid  
Parameswaran



Shivaji  
Sondhi

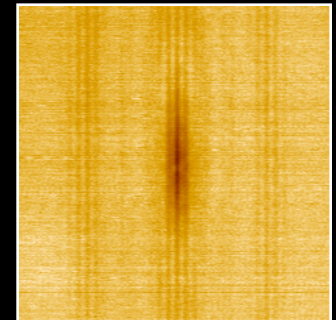
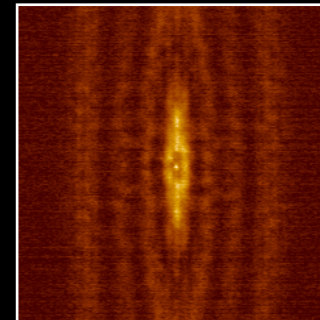
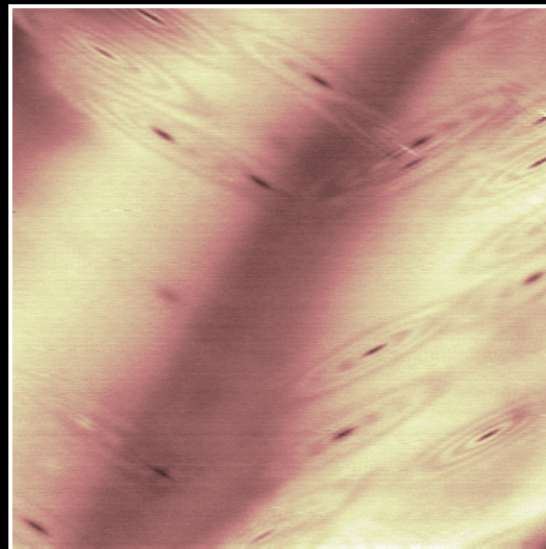
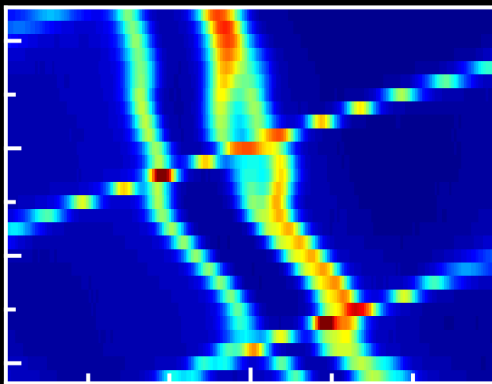
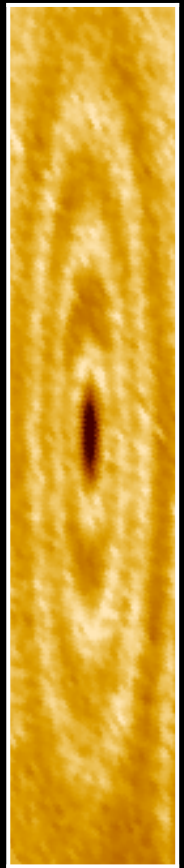
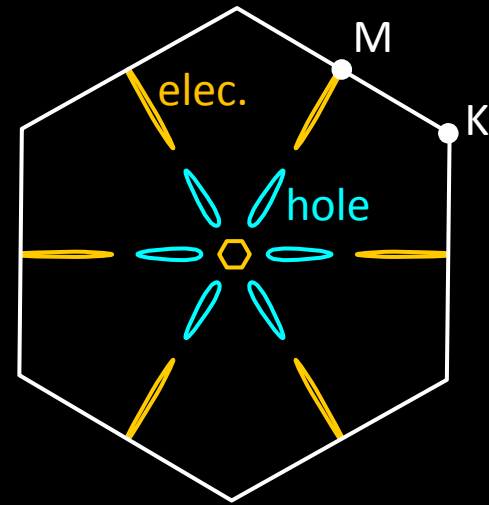
UT Austin

UC Irvine



# Summary & Outlook

- Bi(111) surface - unique platform to explore interaction driven nematic quantum Hall phase
- Real-space imaging of Landau orbits: domains of local nematic order



Thank you!

