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DC, I. Sodemann & T. Senthil, arXiv:1706.00418 I. Sodemann, DC & T. Senthil, to appear

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KITP Workshop on Intertwined order and fluctuations in quantum materials

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Debanjan Chowdhury

Mixed-valence insulators with

neutral Fermi surfaces

Topological mixed-valence insulators

- Topological insulators: Spin-orbit coupled band-insulators of electrons without exotic (fractionalized) quasiparticles
- Many experimental realizations: $Bi_{1-x}Sb_x$, Bi_2Se_3 , Bi_2Te_3 ,.... Well described within band-theory.

Hasan, Kane, 2010; Qi, Zhang, 2011

Role of strong correlations? Candidate materials?



Topological mixed-valence insulators

- electrons without exotic (fractionalized) quasiparticles Topological insulators: Spin-orbit coupled band-insulators of
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Hasan, Kane, 2010; Qi, Zhang, 2011

- Role of strong correlations? Candidate materials?
- Mixed-valence systems are ideal candidates!
- SmB₆ is currently the most popular candidate correlated topological insulator

Dzero, Sun, Galitski, Coleman, 2010





Is SmB₆ a correlated topological insulator that is adiabatically related to the usual topological band insulator?

OR

Does SmB₆ host a truly exotic insulating bulk with gapless termionic excitations?





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Does SmB₆ host a truly exotic insulating bulk with gapless termionic excitations??

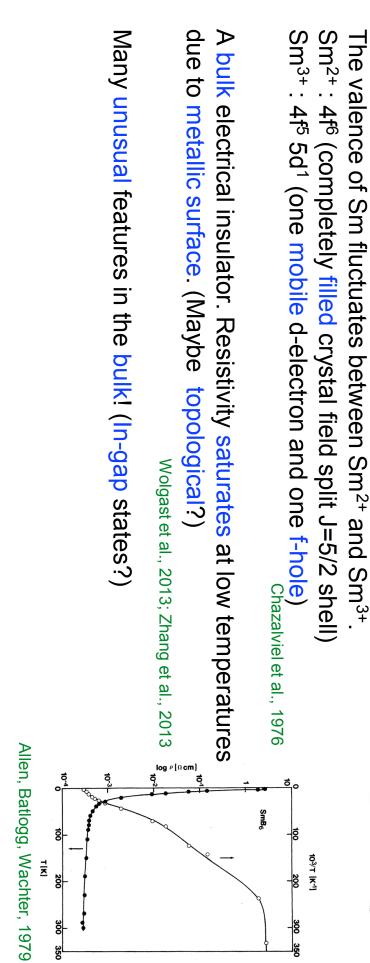
Rest of this talk: 1) Overview of the experimental puzzles

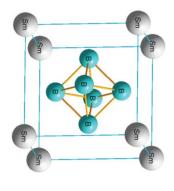
2) Theory of emergent neutral Fermi surfaces

in a three dimensional mixed-valence

insulator (in the bulk). Chowdhury, Sodemann, Senthil, 2017





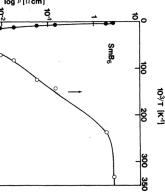


Menth, Buehler, Geballe, 1969

electrons with strong spin-orbit coupling.

Simple cubic structure. Mobile d-electrons and strongly correlated f-

SmB₆ : A classic mixed valence insulator



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SmB₆ : A classic mixed valence insulator

- Many low-temperature anomalies:
- Linear in T specific heat from bulk states.

Flachbart et al., 2006; Phelan et al., 2014; Wakeham et al., 2016

Sub-gap optical absorption.

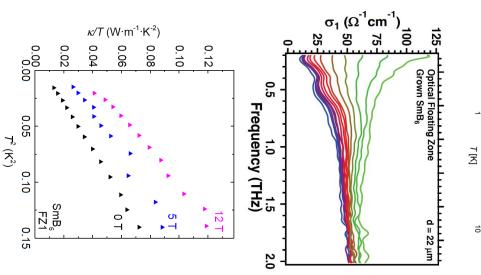
N.J. Laurita et al., 2016

NMR relaxation

V. Mitrovic et al., unpublished; T. Caldwell et al., 2007

(Field induced) thermal conductivity*

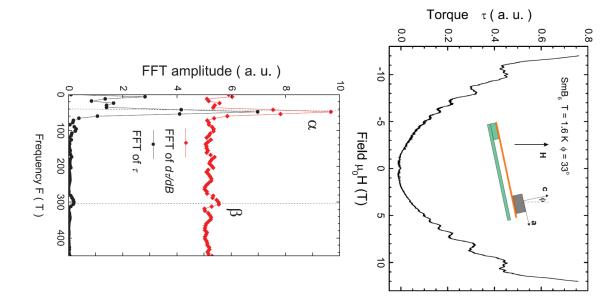
M. Hartstein, S. Sebastian et al., unpublished



* Not present universally Y. Xu et al., 2016

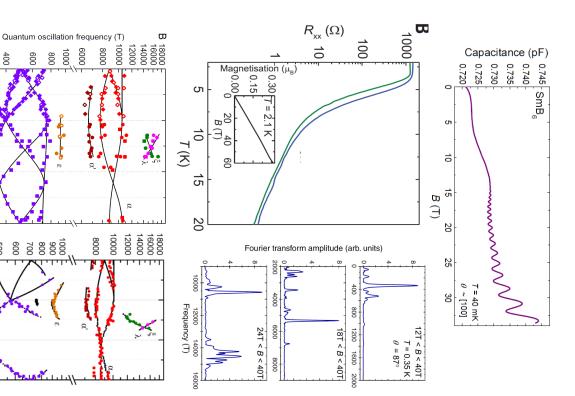
Quantum oscillations in SmB₆

- states? Oscillations from two-dimensional metallic surface
- Seen in magnetization but not in resistivity.
- Sharp (small) frequencies.



Quantum oscillations in SmB₆

- Quantum oscillations in magnetization (but not in resistivity).
- Sharp frequencies
- New high frequency oscillations (similar to observations in other metallic hexaborides).
- Angular dependence?
- from Lifshitz-Kosevich behavior Low temperature enhancement and deviation
- bulk in-gap states! Radical proposal: Oscillations arising from the





200 ∟ -45 [110]

> 0 [100]

45 [111]

90 [110]

[110]

0 45 [100] [111] θ (degrees)

90 [110] SmB₆

Exotic bulk in SmB₆?

- related theoretical considerations)? that is consistent with all of these experimental observations (and Is there a Fermi-surface of neutral Fermions in three dimensions
- field. Relevant for organics A standard example: spinon Fermi-surface coupled to U(1) gauge Motrunich, 2005; Lee, Lee 2005

Yamashita et al., 2008; Yamashita et al., 2010

- one stabilize a phase with Fermi-surface of neutral fermions in a What is the microscopic origin of the neutral fermion? How can
- mixed-valence insulator?
- Chowdhury, Sodemann, Senthil, 2017

Periodic Anderson model

Sm²⁺ : Fully filled crystal field multiplet Sm³⁺ : One mobile d-electron (charge -e) and one f-hole (charge e) The valence of Sm fluctuates between Sm²⁺ and Sm³⁺.

$$f_{\alpha} \rightarrow \varepsilon_{\alpha\beta} f_{\beta}^{\dagger} = \tilde{f}_{\alpha}$$

Simplest possible model:

$$\begin{split} H &= \sum_{k\alpha} \epsilon_d(k) d^{\dagger}_{k\alpha} d_{k\alpha} + \epsilon_f(k) \tilde{f}^{\dagger}_{k\alpha} \tilde{f}_{k\alpha} \\ &+ \sum_r (\epsilon_{\beta\gamma} V_{\alpha\beta} d^{\dagger}_{r\alpha} \tilde{f}^{\dagger}_{r\gamma} + H.c.) - U_{df} \sum_r n^{\tilde{f}}_r n^d_r \\ &+ U_{ff} \sum_r n^{\tilde{f}_r} (n^{\tilde{f}_r} - 1), \end{split}$$

Average d-electron density = f-hole density



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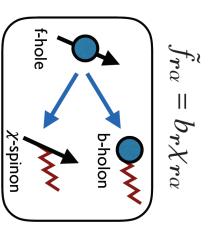
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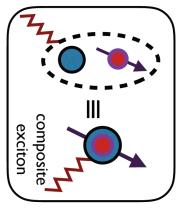


Fermionic Composite Excitons

- Spinon (χ) spin ½ Fermion that is electrically neutral Holon (b) – spinless Boson that carries physical charge eFractionalize the f-hole into:
- Spinon and holon coupled to emergent U(1) gauge-field,
- $\langle b_r^{\dagger} b_r \rangle = x$). Density of doped holes, away from 4f⁶ configuration, is xElectrical insulator when holon is gapped ($\langle b \rangle = 0$).
- charge neutral fermions carrying spin $\frac{1}{2}$. conduction electrons may lead to formation* of excitons: At large U_{df}, Coulomb attraction between holons and

$$\psi_lpha=bd_lpha$$





* 1) Similar to "molecular fermion" in Bose-Fermi mixtures (without gauge-field).
2) Fermionic excitons in multicomponent QH systems (Barkeshli et al., 2016).



Chowdhury, Sodemann, Senthil, 2017

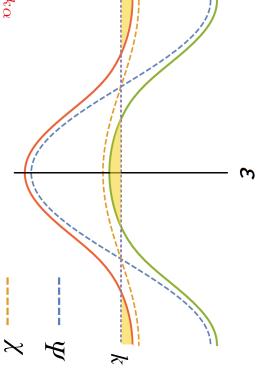
Effective model in the strong-binding limit

- electrical insulator) gauge-field. (d-electrons and holons are gapped; excitons and spinons minimally coupled to U(1) Low energy degrees of freedom : composite
- Density of excitons = Density of spinons = Density of d-electrons
- Effective Hamiltonian:

$$H = \sum_{k\alpha} \epsilon^{\psi}_{k} \psi^{\dagger}_{k\alpha} \psi_{k\alpha} + V \epsilon_{\alpha\beta} \psi_{k\alpha} \chi_{-k,\beta} + h.c + \epsilon^{\chi}_{k} \chi^{\dagger}_{k\alpha} \chi_{k\alpha}$$

- the fCE In addition, there is a strong on-site repulsion for
- "metal" of fCE A possible ground state : Compensated semi-







Katsura, Nagaosa, Lee, 2010

- Thermal conductivity: $\kappa_{xx} \sim T$
- Optical conductivity: $\sigma(\omega) \sim \omega^{\alpha}$ ($\alpha = 5/3$ or 2 depending on impurity mean-free path) Ng, Lee, 2007
- NMR relaxation rate: $\frac{1}{T_1T}$ = const.
- Specific heat :

Phenomenology of Fermionic CE liquids

 $C = \gamma T \qquad \gamma \sim \ln(1/T)$

Effects of magnetic fields?

External magnetic field induces an internal magnetic field.

Motrunich, 06

external field and sets up currents that induce an internal magnetic field. Gapped holon (charged under external and internal gauge-fields) responds to

$$\begin{array}{cccc} \mathbf{B} = \nabla \times \mathbf{A} & \mathbf{j}_b & \mathbf{B} = \nabla \times \mathbf{A} & \mathbf{b}_{\mathbf{a}} = \nabla \times \mathbf{a} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{u} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf{U} \\ \hline \mathbf{U} & \mathbf$$

- quantization Neutral fCE experience internal magnetic field and can exhibit Landau
- Energy follows from gauge-invariance

$$u(b, B) = u_0 + \frac{(b-B)^2}{2\mu_b} + \frac{b^2}{2\mu_{ce}} + \frac{B^2}{2\mu_v} + u_{osc}(b) + \dots \implies b = \alpha B, \text{ with } \alpha = \frac{\mu_{ce}}{(\mu_{ce} + \mu_b)}$$

Hard to compute α but expected to

Frequency of oscillations
$$\Delta \left(\frac{1}{B}\right) = \frac{2\pi}{S_{\perp}^{i}} \left(1 + \frac{\mu_{b}}{\mu_{ce}}\right)^{-1} = \frac{2\pi\alpha}{S_{\perp}^{i}}$$

approach 1 close to metal-insulator transition.

b B

Sodemann, Chowdhury, Senthil, to appear

Outlook

- Gapless fermionic excitations in a mixed-valence insulator?
- insulators. Low energy excitations: FERMIONIC COMPOSITE EXCITONS Theoretical description of a new phase in mixed-valence
- Immediate consequence: Weak field thermal Hall effect.

(Bound state of holon and conduction electron).

$$\kappa_{xy}^{i} = (\omega_{c,i}\tau_{i})\kappa_{xx}^{i}$$
 $\omega_{c,i} = \alpha e|B|/m_{i}$

display similar phenomenology! Other correlated mixed-valence insulators (e.g. YbB₁₂) likely

