

Spin-Orbital Interplay in Double Perovskites

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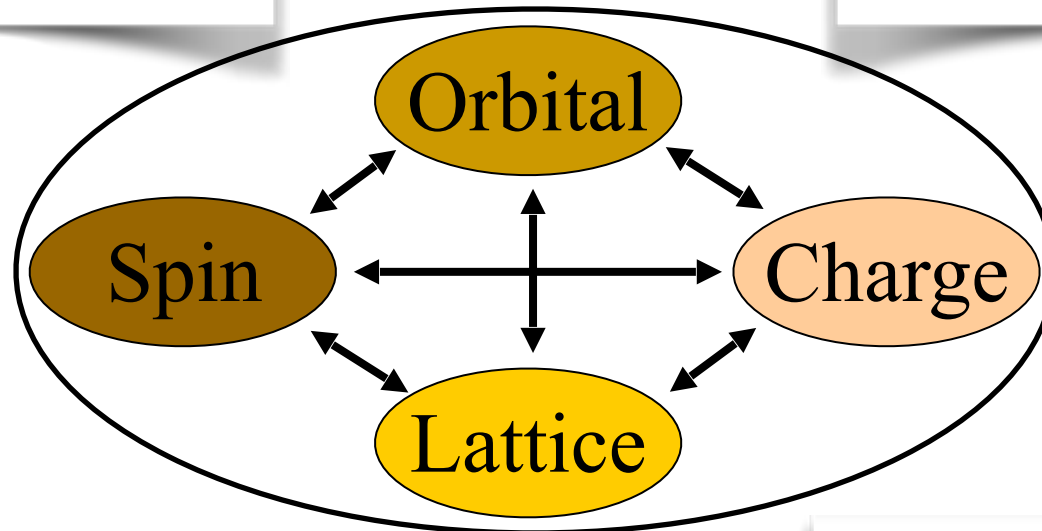
Intertwined17, KITP
Sep 26, 2017

Transition metal oxides: Plethora of Challenging Phenomena

- d- electrons are partially filled and not very extended
- Interactions dominate over kinetic energy
- Many degrees of freedom are at work

Metal to Magnetic Insulator

Colossal Magnetoresistance

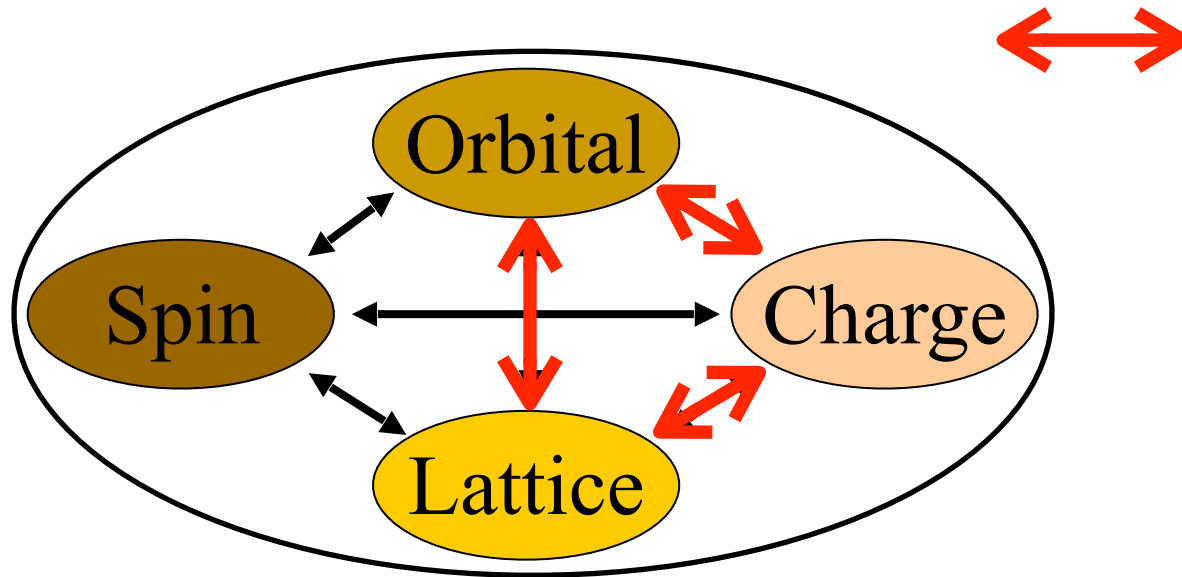


High-Tc Superconductivity

Unexpected variety of phases
and transitions between them

Transition metal oxides: Plethora of Challenging Phenomena

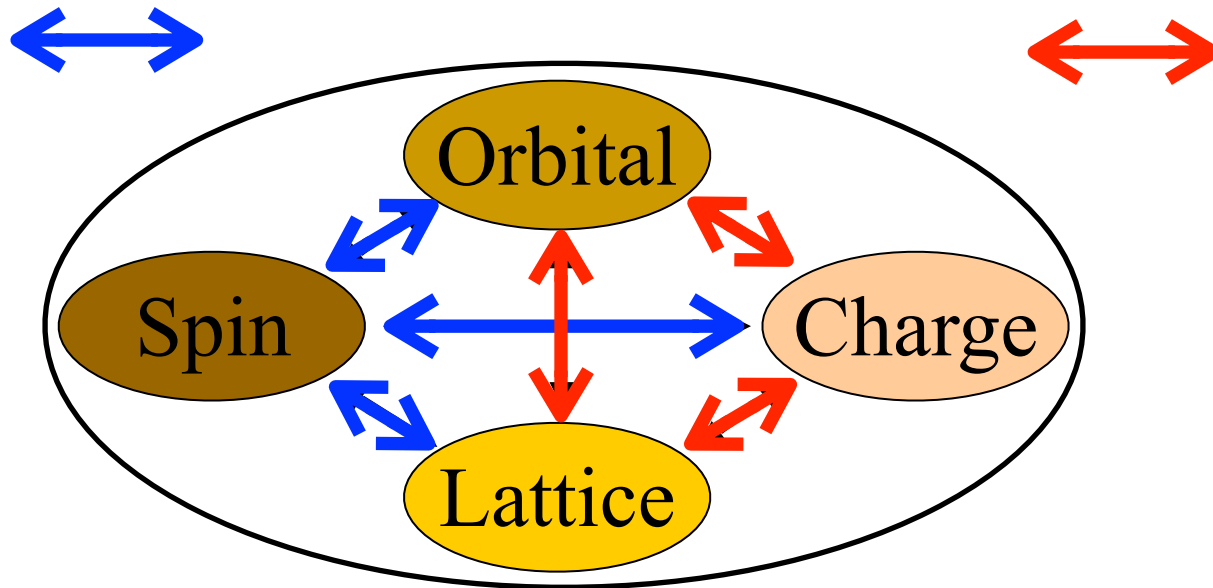
Coulomb force



Transition metal oxides: Plethora of Challenging Phenomena

Spin-orbit coupling
(relativistic)

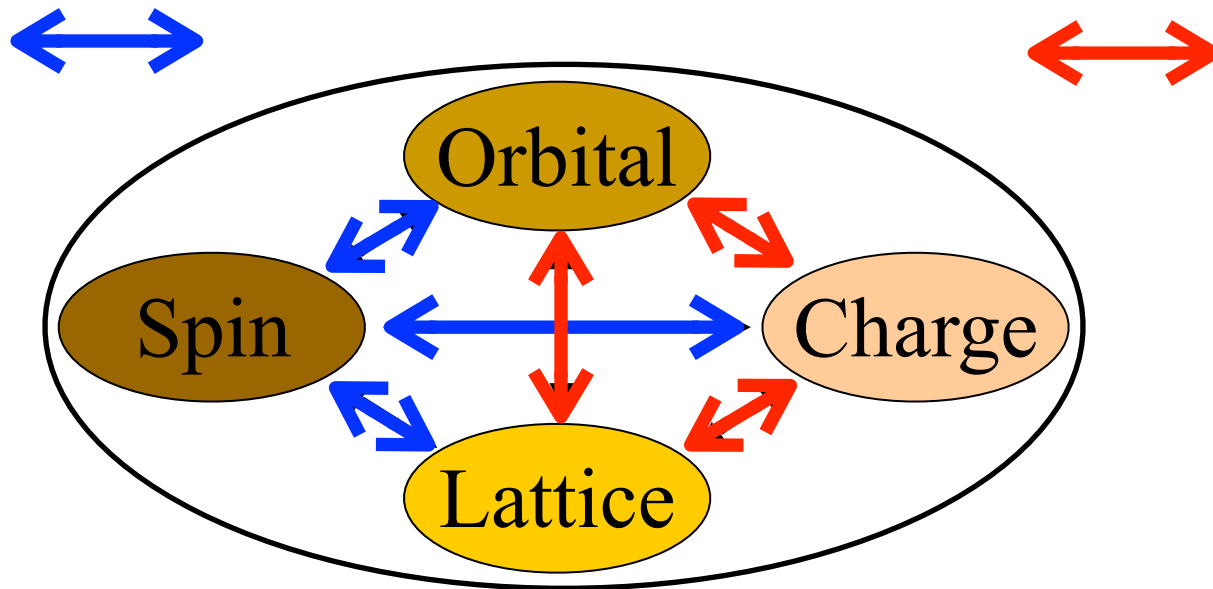
Coulomb force



Transition metal oxides: Plethora of Challenging Phenomena

Spin-orbit coupling
(relativistic)

Coulomb force



Enhanced interplay by going to heavier TM compounds:
a possible route to Novel Collective States & Phenomena

correlations

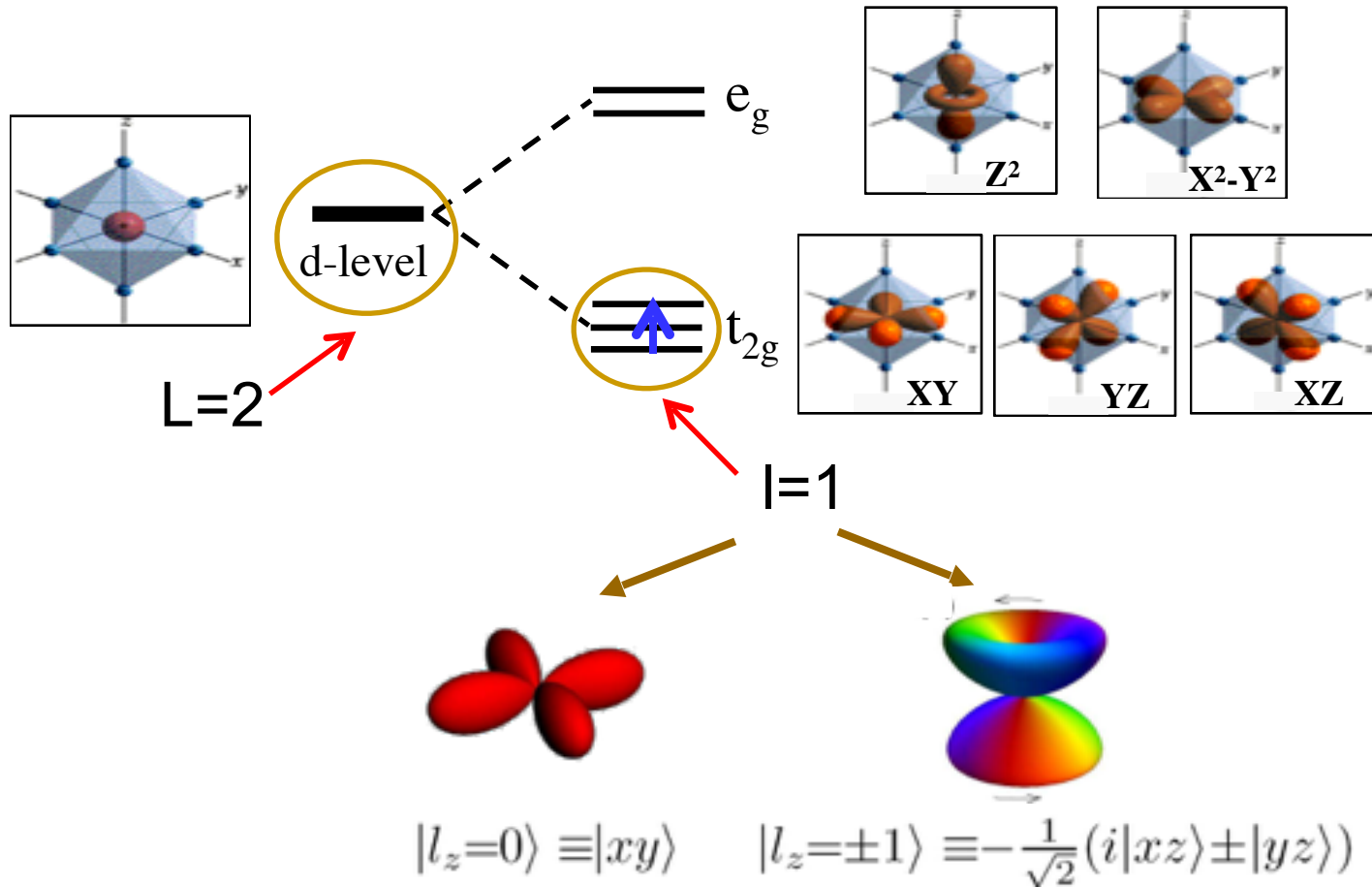


21 Sc 44.9559 Scandium	22 Ti 47.867 Titanium	23 V 50.9415 Vanadium	24 Cr 51.9961 Chromium	25 Mn 54.938 Manganese	26 Fe 55.845 Iron	27 Co 58.9332 Cobalt	28 Ni 58.6934 Nickel	29 Cu 63.546 Copper	30 Zn 65.4089 Zinc
39 Y 88.9058 Yttrium	40 Zr 91.224 Zirconium	41 Nb 92.9064 Niobium	42 Mo 85.94 Molybdenum	43 Tc 98 Technetium	44 Ru 101.07 Ruthenium	45 Rh 102.9055 Rhodium	46 Pd 106.42 Palladium	47 Ag 107.8682 Silver	48 Cd 112.411 Cadmium
71 Lu 174.967 Lutetium	72 Hf 178.49 Hafnium	73 Ta 180.9497 Tantalum	74 W 183.84 Tungsten	75 Re 186.207 Rhenium	76 Os 190.23 Osmium	77 Ir 192.217 Iridium	78 Pt 195.084 Platinum	79 Au 196.9666 Gold	80 Hg 200.59 Mercury



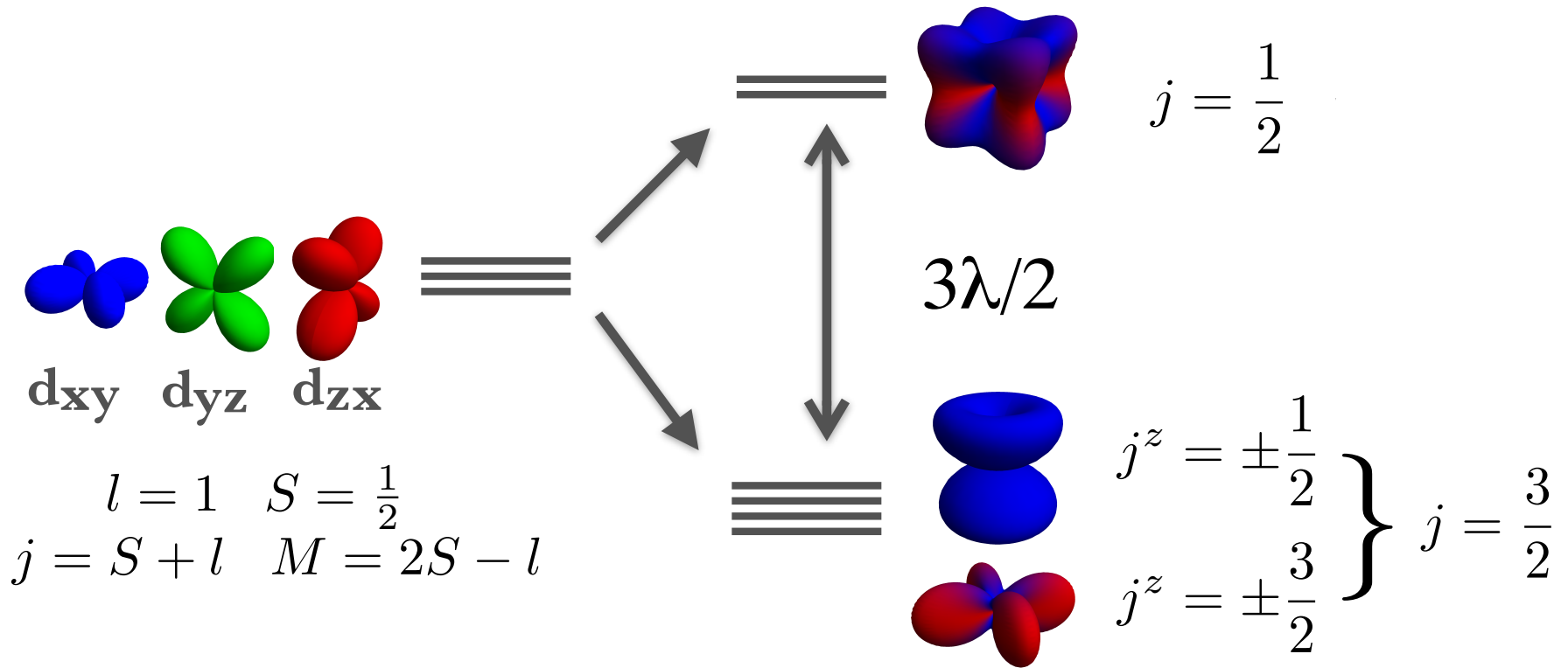
spin-orbit

Angular momentum in t_{2g} shell



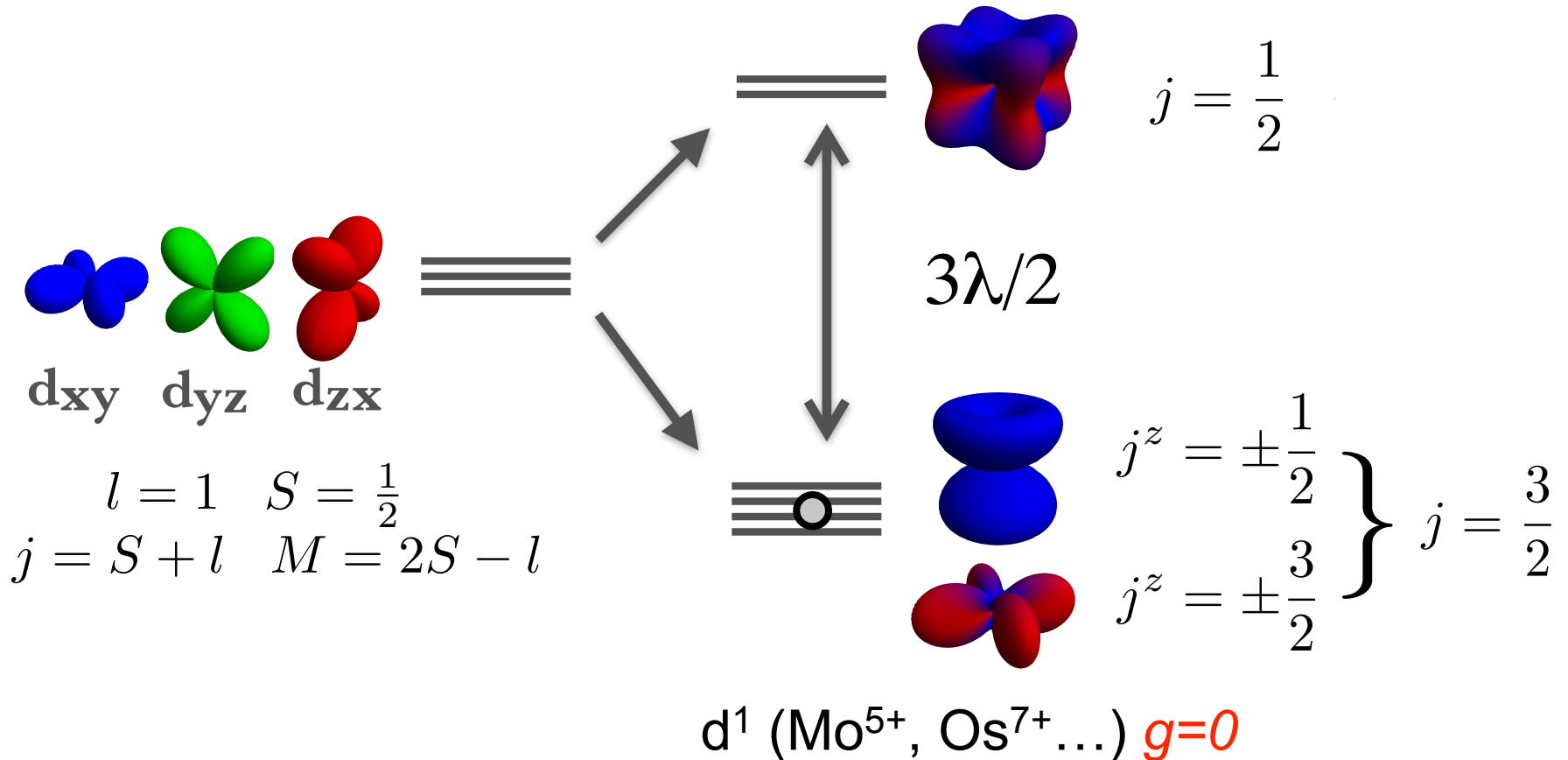
Spin-orbit coupling in t_{2g} shell

A. Abragam and B. Bleaney, "EPR of Transition Ions" (1970)



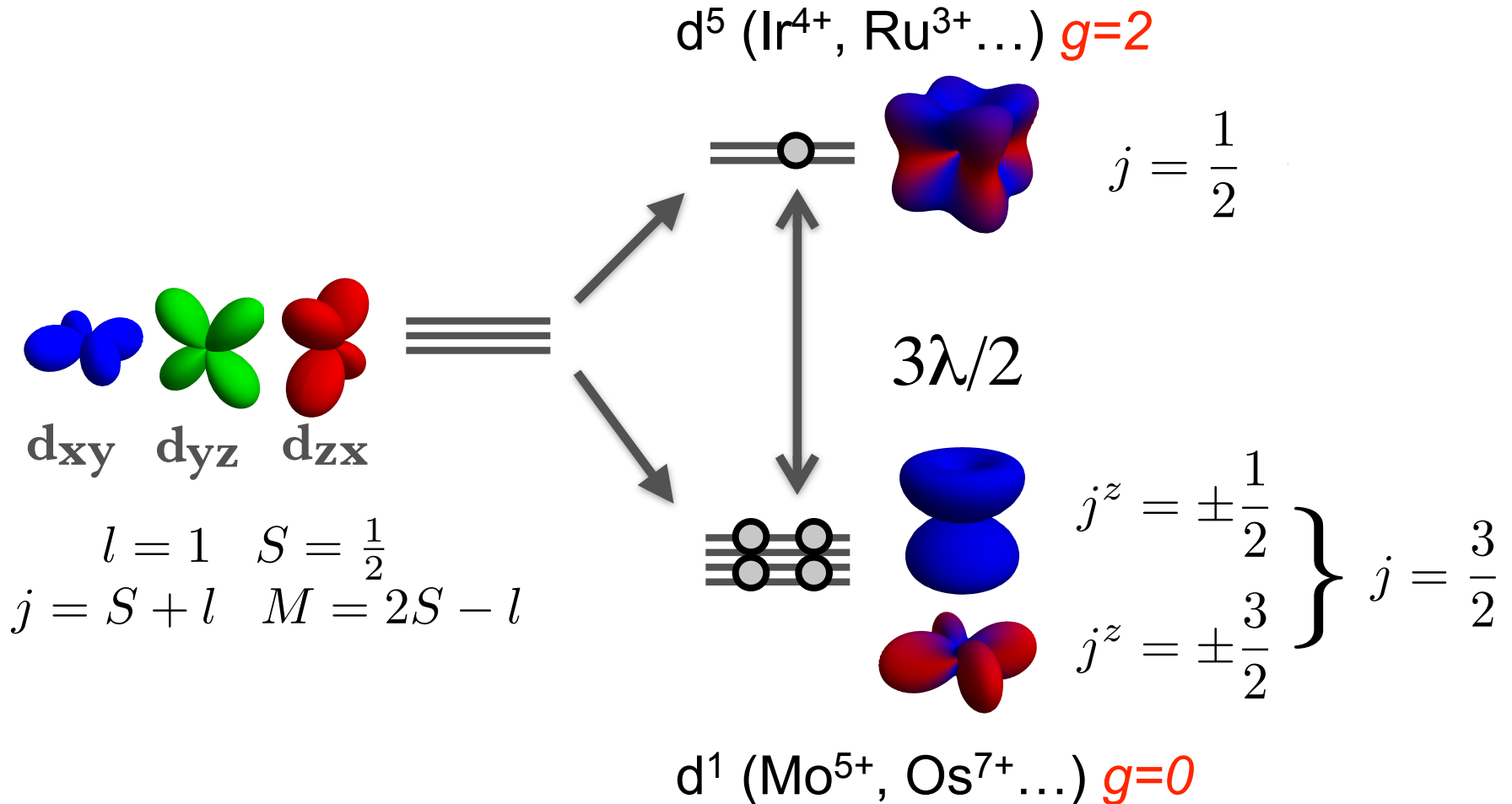
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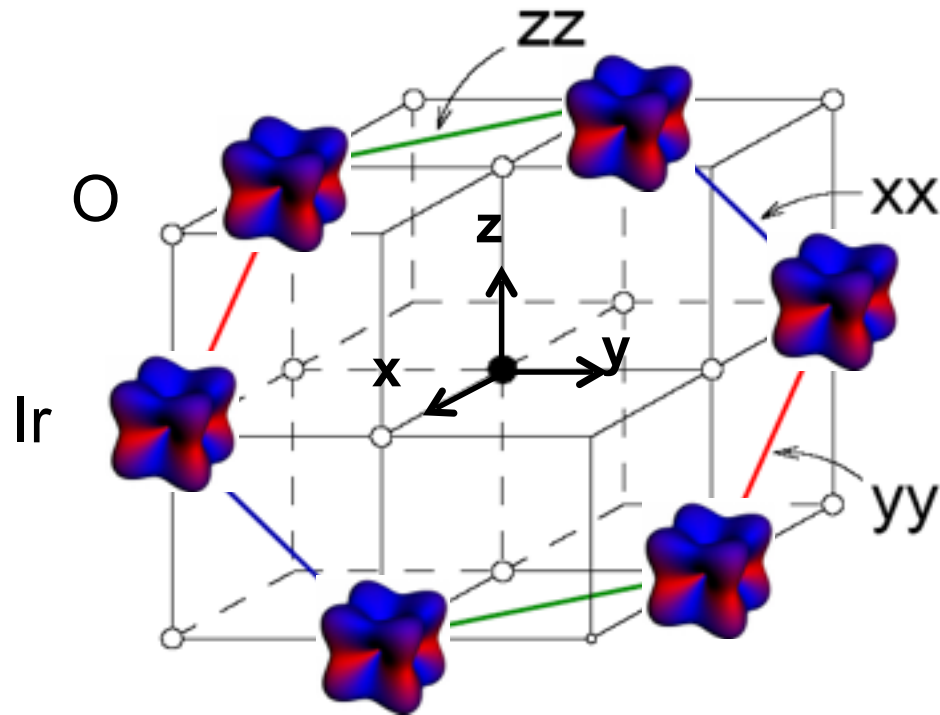
Spin-orbit coupling in t_{2g} shell

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Pseudospins on a Honeycomb layer of $A_2\text{IrO}_3$

GJ & Khaliullin, PRL (2009)



A. Kitaev, Ann. Phys'06

$$\mathcal{H}_{ij}^{(\gamma)} = JS_i^\gamma S_j^\gamma$$

Honeycomb-based Compounds do show LRO

but NON of them in Neel state

discussion on QSL

18th Aug Intertwined17@KITP

Honeycomb-based Compounds do show LRO

'Mere goodness can achieve
little against
the power of nature.'



Georg Wilhelm Friedrich HEGEL

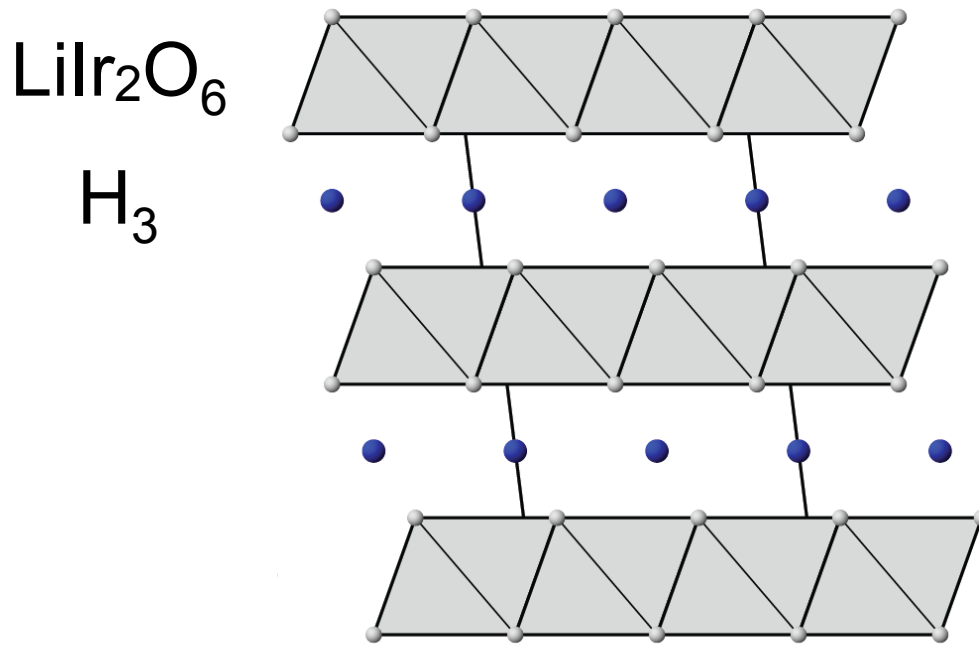
Honeycomb-based Compounds do show LRO

'Nothing great in the world
has ever been
accomplished without passion.'



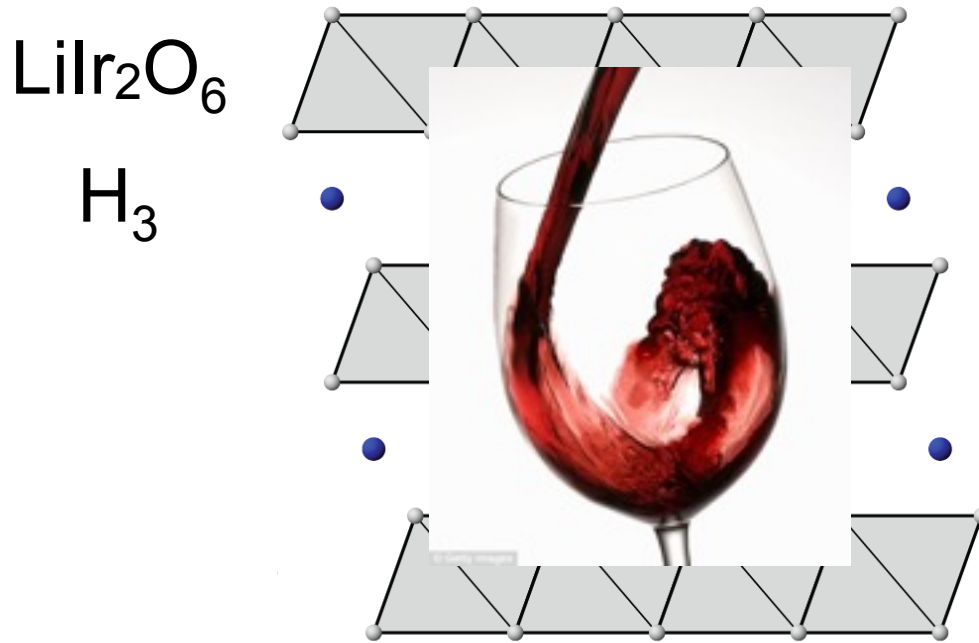
Georg Wilhelm Friedrich HEGEL

$\text{H}_3\text{LiIr}_2\text{O}_6$ (Takagi's Group)



No sign of long-range order down to 50 mK
Susceptibility, NMR, Heat capacity

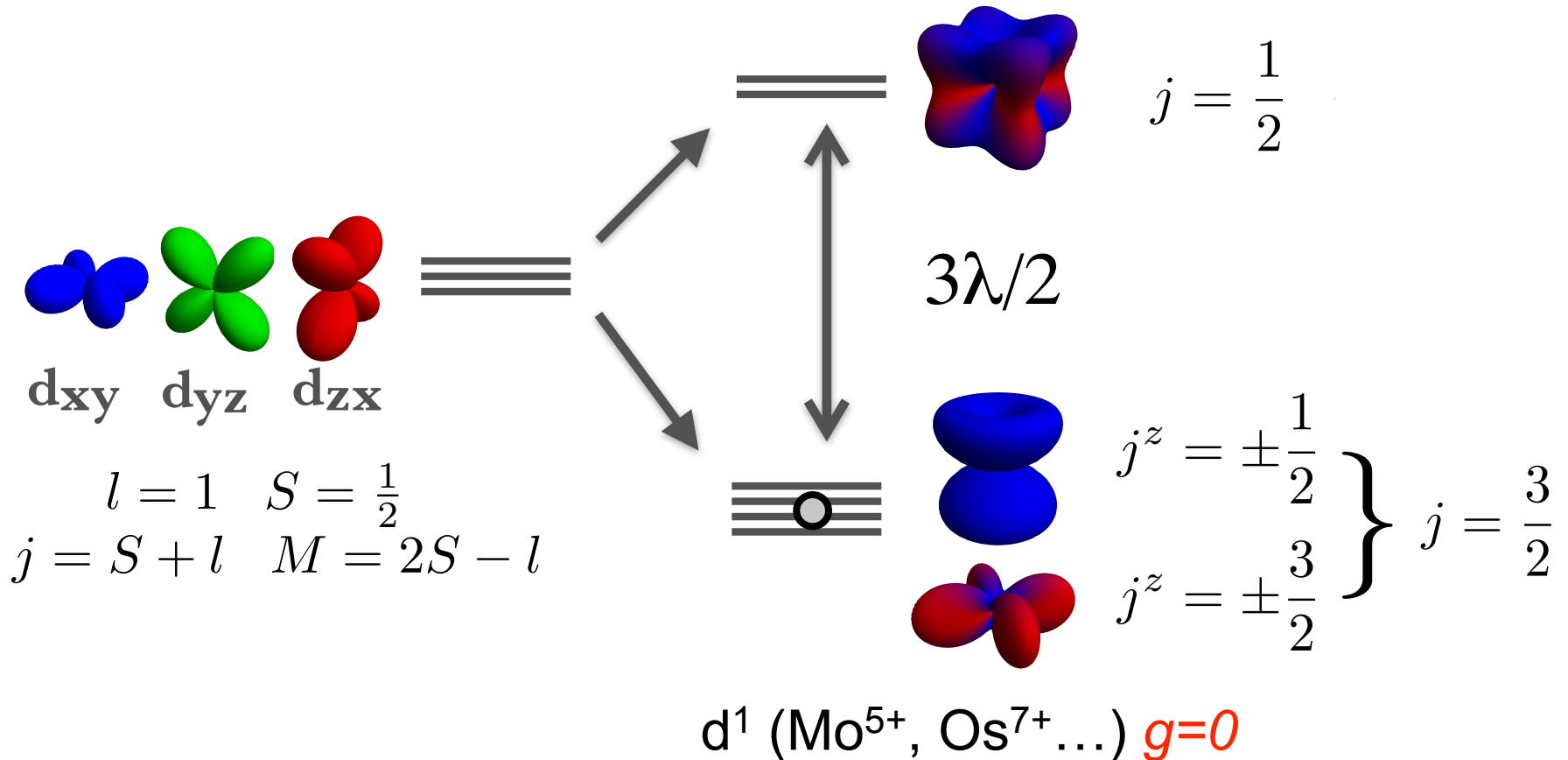
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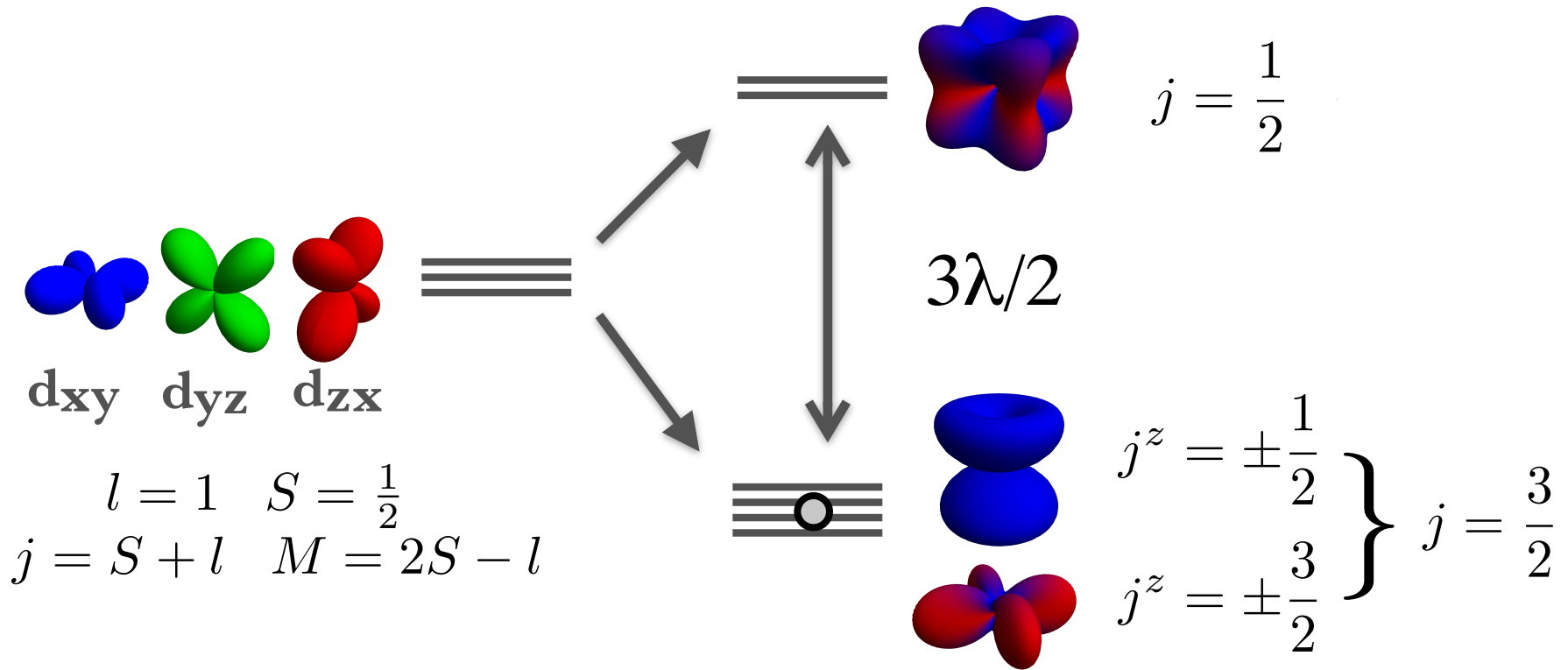
Spin-orbit coupling in t_{2g} shell

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Spin-orbit coupling in t_{2g} shell

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Magnetic field mixes $j=1/2$ with $3/2$
and induces van Vleck magnetism

Ba_2BMoO_6 ($B=Y, Lu$)
valence bond glass

↑
correlations

21 Sc 44.9559 Scandium	22 Ti 47.867 Titanium	23 V 50.9415 Vanadium	24 Cr 51.9961 Chromium	25 Mn 54.938 Manganese	26 Fe 55.845 Iron	27 Co 58.9332 Cobalt	28 Ni 58.6934 Nickel	29 Cu 63.546 Copper	30 Zn 65.4089 Zinc
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↓
spin-orbit

Ba_2BOsO_6 ($B=Na, Li$)
unusual 'spin' textures

In collaboration with



Judit Romhanyi
(MPI -> OIST)



Leon Balents
(KITP)

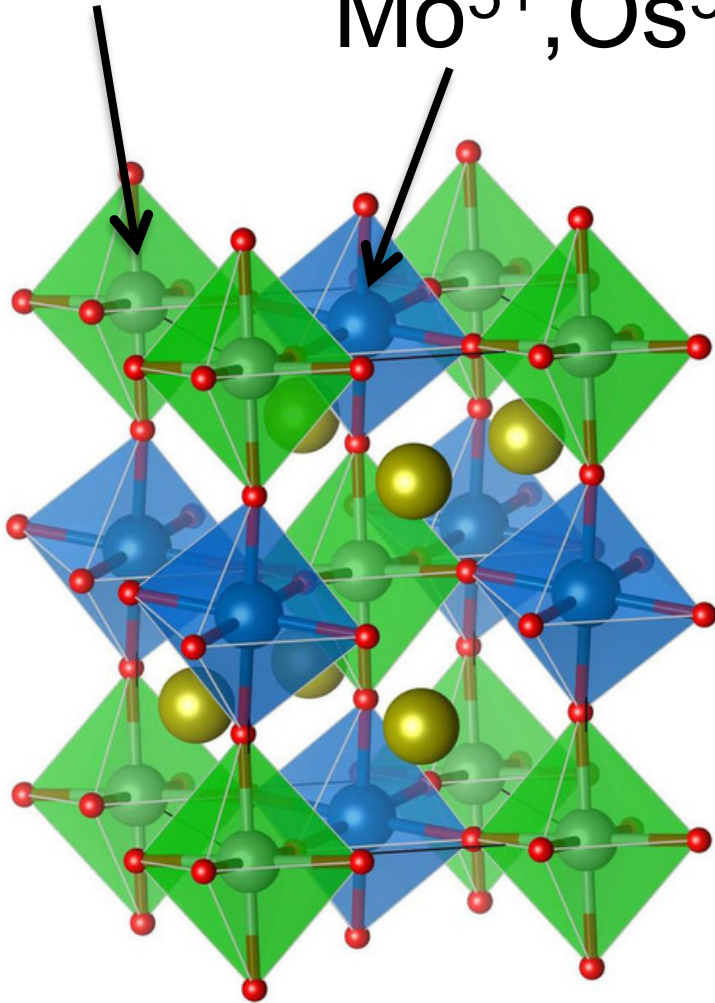
J. Romhanyi, L. Balents & GJ, [PRL \(2017\)](#)

double perovskites $A_2BB'O_6$

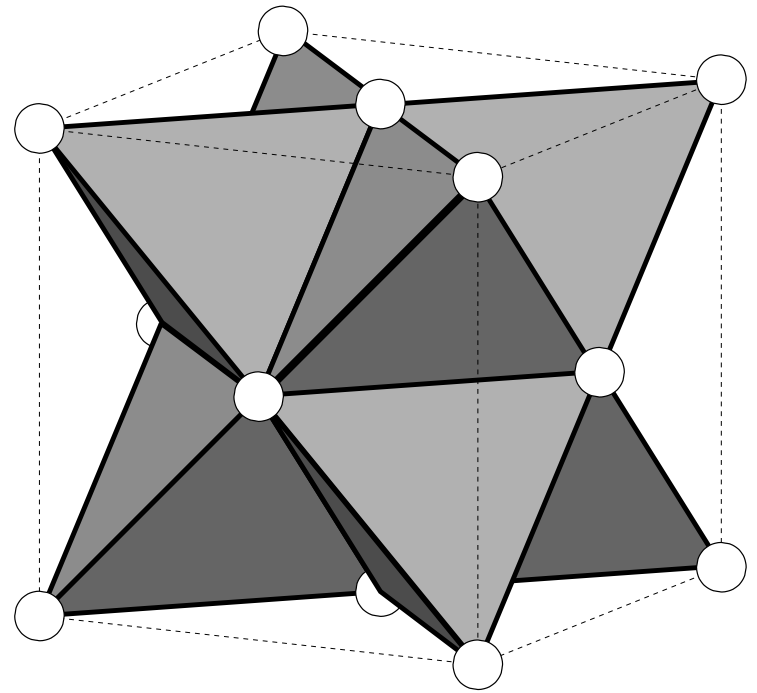
B

(nonmagnetic)

Mo^{5+}, Os^{5+}



each sublattice forms FCC





Exotic phases induced by strong spin-orbit coupling in ordered double perovskites

Gang Chen

*Physics Department, University of Colorado, Boulder, Colorado 80309, USA
and Physics Department, University of California, Santa Barbara, California 93106, USA*

Rodrigo Pereira

*Instituto de Física de São Carlos, Universidade de São Paulo, CP 369, São Carlos 13566-970, SP, Brazil
and Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California 93106, USA*

Leon Balents

Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California 93106, USA

(Received 26 September 2010; published 29 November 2010)

Compound	B' config.	Crystal structure	Θ_{CW} (K)	μ_{eff} (μ_B)	Magnetic transition	Frustration para. f	Ref
Ba_2YMoO_6	$\text{Mo}^{5+}(4d^1)$	Cubic	-91	1.34	PM down to 2 K	$f \geq 45$	11
Ba_2YMoO_6	$\text{Mo}^{5+}(4d^1)$	Cubic	-160	1.40	PM down to 2 K	$f \geq 80$	12
Ba_2YMoO_6	$\text{Mo}^{5+}(4d^1)$	Cubic	-219	1.72	PM down to 2 K	$f \geq 100$	13
$\text{La}_2\text{LiMoO}_6$	$\text{Mo}^{5+}(4d^1)$	Monoclinic	-45	1.42	Short-range AFM $T_N \sim 20$ K	$f \geq 2$	13
$\text{Sr}_2\text{MgReO}_6$	$\text{Re}^{6+}(5d^1)$	Tetragonal	-426	1.72	Spin glass, $T_G \sim 50$ K	$f \geq 8$	14
$\text{Sr}_2\text{CaReO}_6$	$\text{Re}^{6+}(5d^1)$	Monoclinic	-443	1.659	Spin glass, $T_G \sim 14$ K	$f \geq 30$	15
$\text{Ba}_2\text{CaReO}_6$	$\text{Re}^{6+}(5d^1)$	Cubic to tetragonal (at $T \sim 120$ K)	-38.8	0.744	AFM $T_N = 15.4$ K	$f \sim 2$	16
$\text{Ba}_2\text{LiOsO}_6$	$\text{Os}^{7+}(5d^1)$	Cubic	-40.48	0.733	AFM $T_N \sim 8$ K	$f \geq 5$	17
$\text{Ba}_2\text{NaOsO}_6$	$\text{Os}^{7+}(5d^1)$	Cubic	-32.45	0.677	FM $T_N \sim 8$ K	$f \geq 4$	17
$\text{Ba}_2\text{NaOsO}_6$	$\text{Os}^{7+}(5d^1)$	Cubic	~ -10	~ 0.6	FM $T_N = 6.8$ K	$f \geq 4$	18



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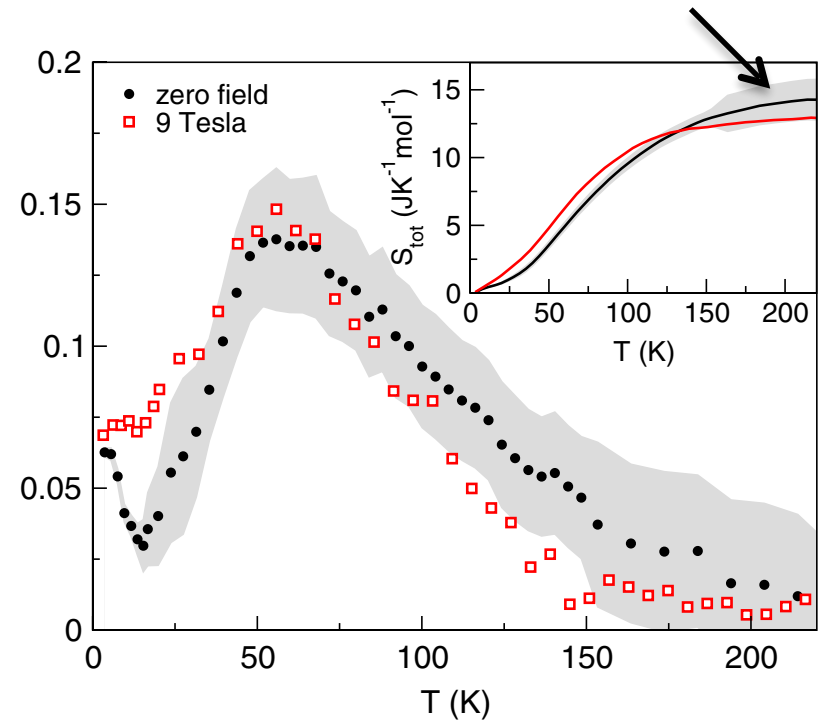
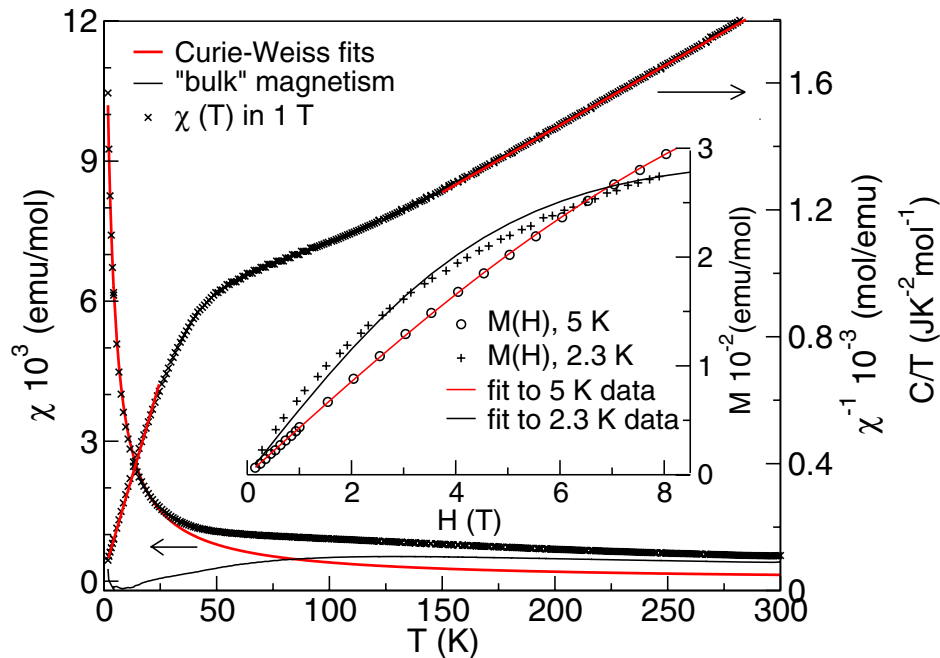
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Valence Bond Glass on an fcc Lattice in the Double Perovskite Ba_2YMoO_6 M. A. de Vries,^{1,2,*} A. C. McLaughlin,³ and J.-W. G. Bos^{4,5}

Rln4

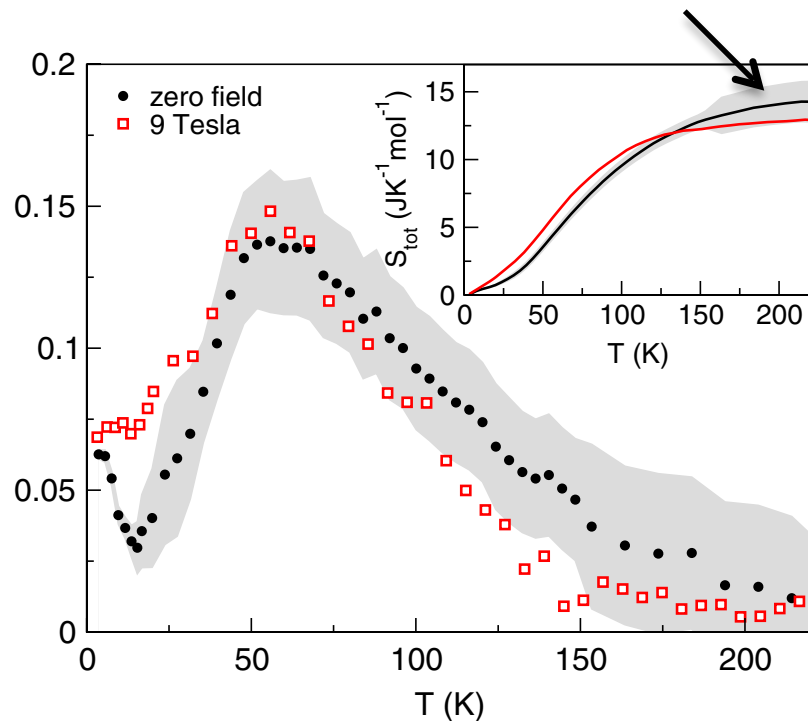
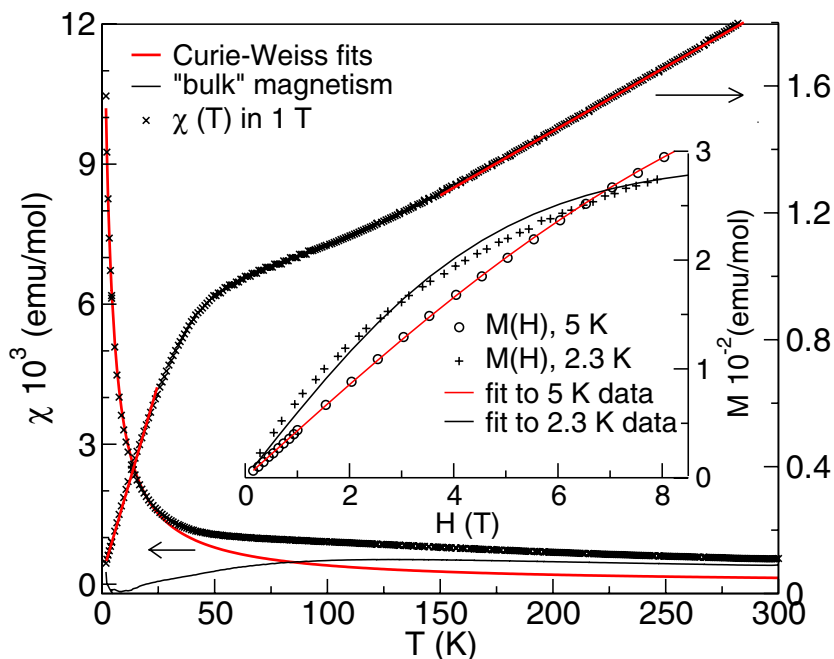


$M(H)$ at 5 (2.3)K: fit Brillouin func,
accounting for 7 (2)% of the Mo $s=1/2$.

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Rln4



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similar behaviour seen also in Ba_2LuMoO_6

**Magnetic properties of the geometrically frustrated $S = \frac{1}{2}$ antiferromagnets, $\text{La}_2\text{LiMoO}_6$ and Ba_2YMoO_6 , with the B-site ordered double perovskite structure:
Evidence for a collective spin-singlet ground state**

Tomoko Aharen,¹ John E. Greedan,^{1,2} Craig A. Bridges,¹ Adam A. Aczel,³ Jose Rodriguez,³ Greg MacDougall,³ Graeme M. Luke,^{2,3,4} Takashi Imai,^{2,3,4} Vladimir K. Michaelis,⁵ Scott Kroecker,⁵ Haidong Zhou,⁶ Chris R. Wiebe,^{6,7} and Lachlan M. D. Cranswick⁸

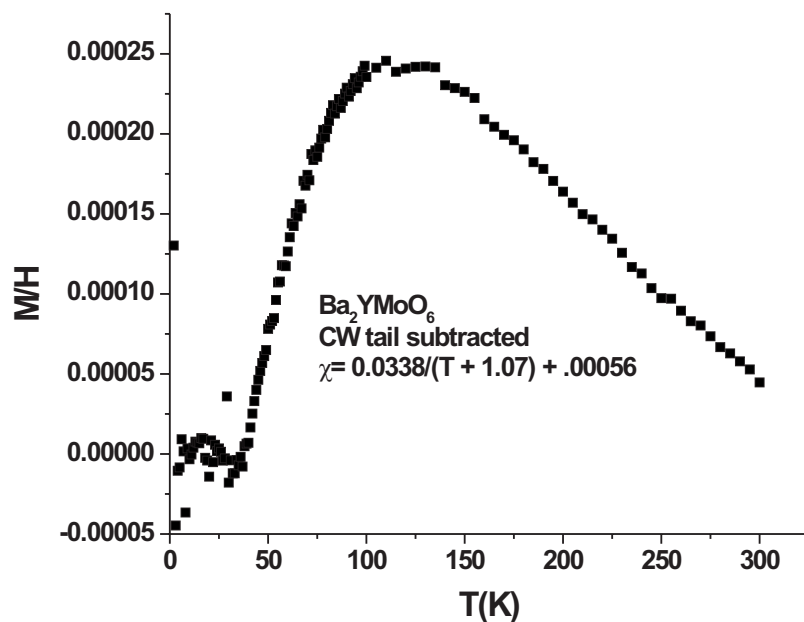
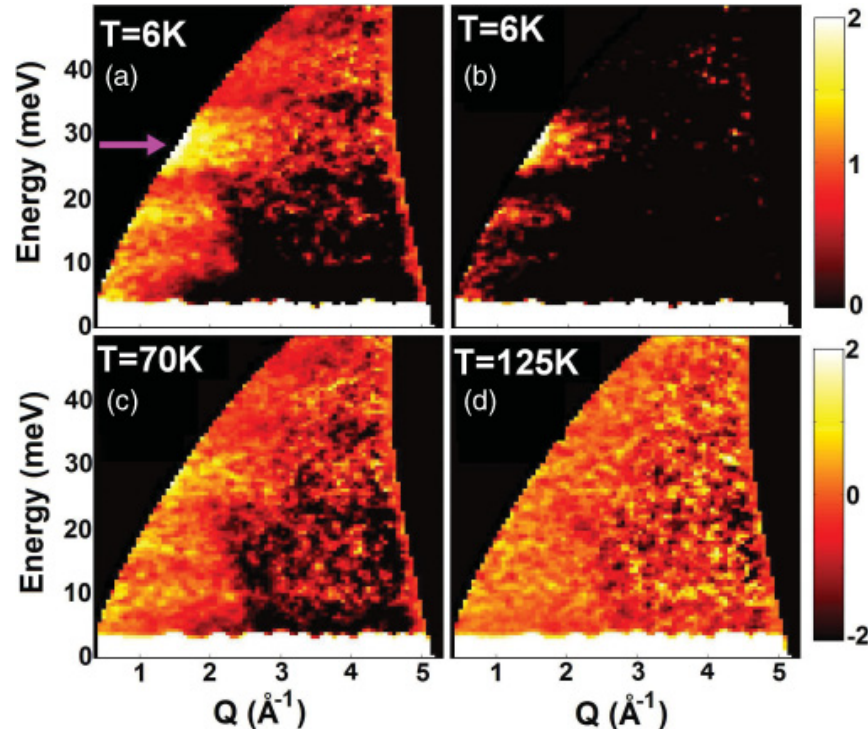


FIG. 17. The Curie-tail subtracted bulk susceptibility of Ba_2YMoO_6 .

NMR $1/T_1$ - two component response:
paramagnetic & gapped (~ 140 K)

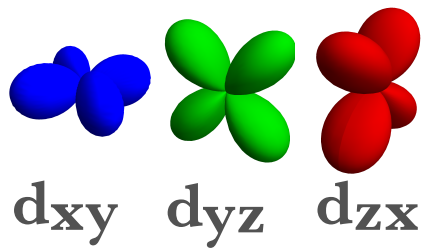
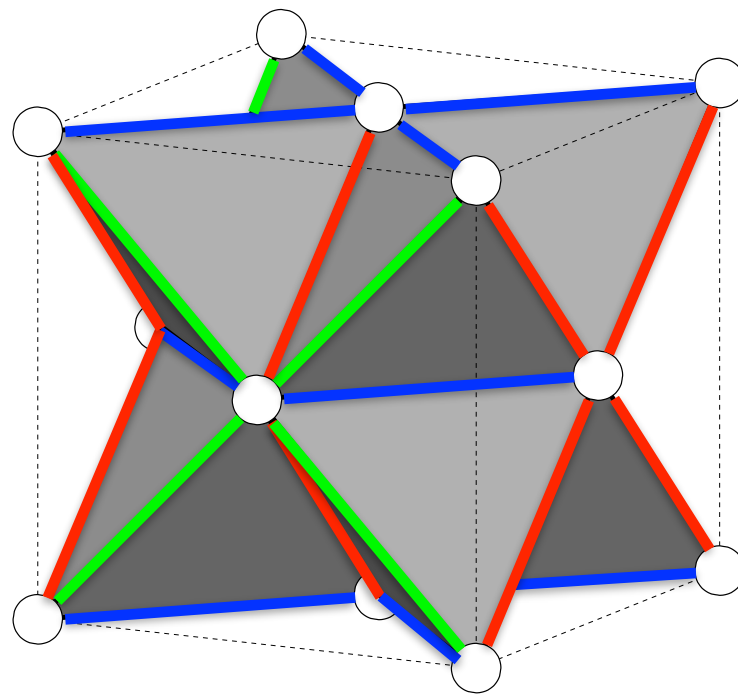
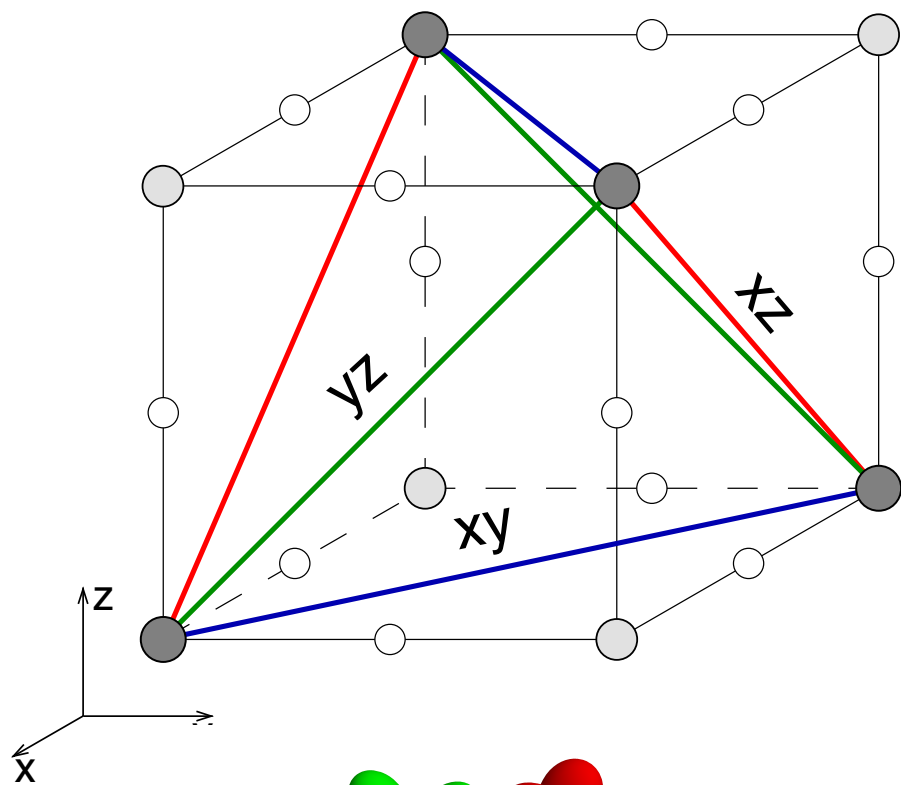
Triplet and in-gap magnetic states in the ground state of the quantum frustrated fcc antiferromagnet Ba_2YMoO_6

J. P. Carlo,^{1,2,*} J. P. Clancy,¹ T. Aharen,³ Z. Yamani,² J. P. C. Ruff,¹ J. J. Wagman,¹ G. J. Van Gastel,¹ H. M. L. Noad,¹ G. E. Granroth,⁴ J. E. Greedan,^{3,5} H. A. Dabkowska,⁵ and B. D. Gaulin^{1,5,6}



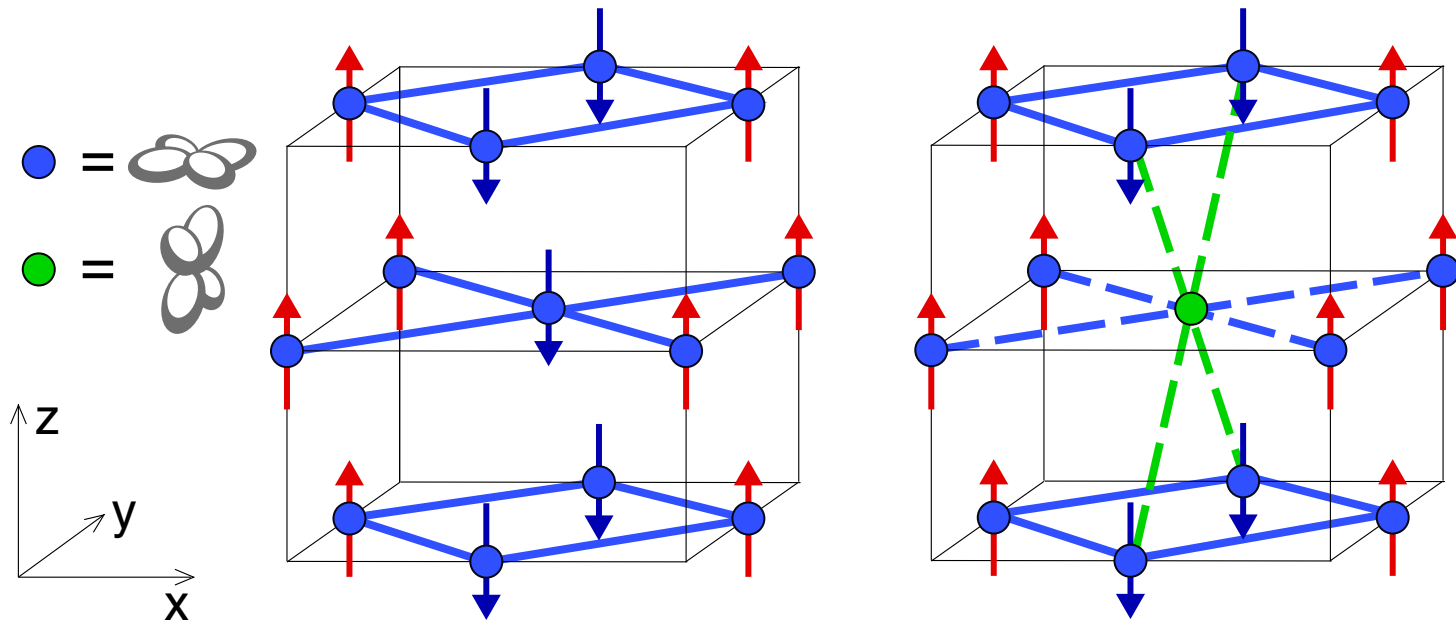
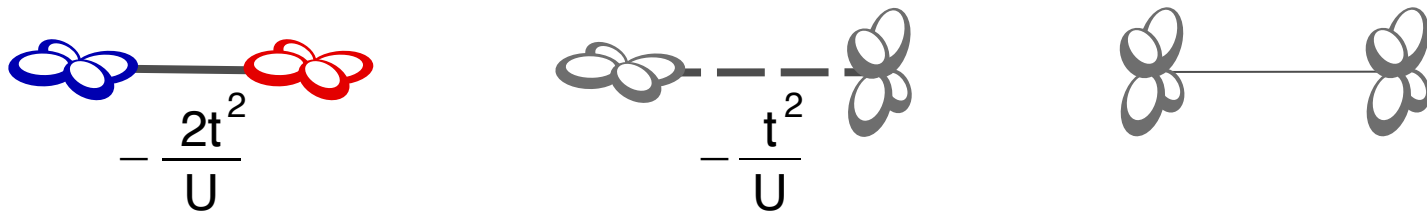
broad gapped mode ~ 28 meV (bandwidth ~ 4 meV)
 in-gap continuum 9-17 meV

Overlap of Mo - Mo t_{2g} orbitals



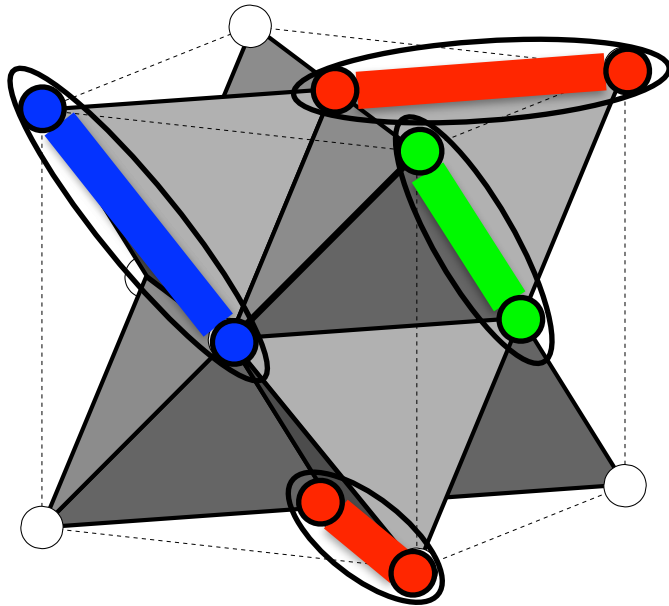
Superexchange Spin-Orbital Model

$$\mathcal{H}_{\text{ex}} = J \sum_{\langle i,j \rangle} \sum_{\alpha\beta} \left(\mathbf{S}_i \cdot \mathbf{S}_j + \frac{1}{4} \right) n_{i\alpha\beta} n_{j\alpha\beta}$$



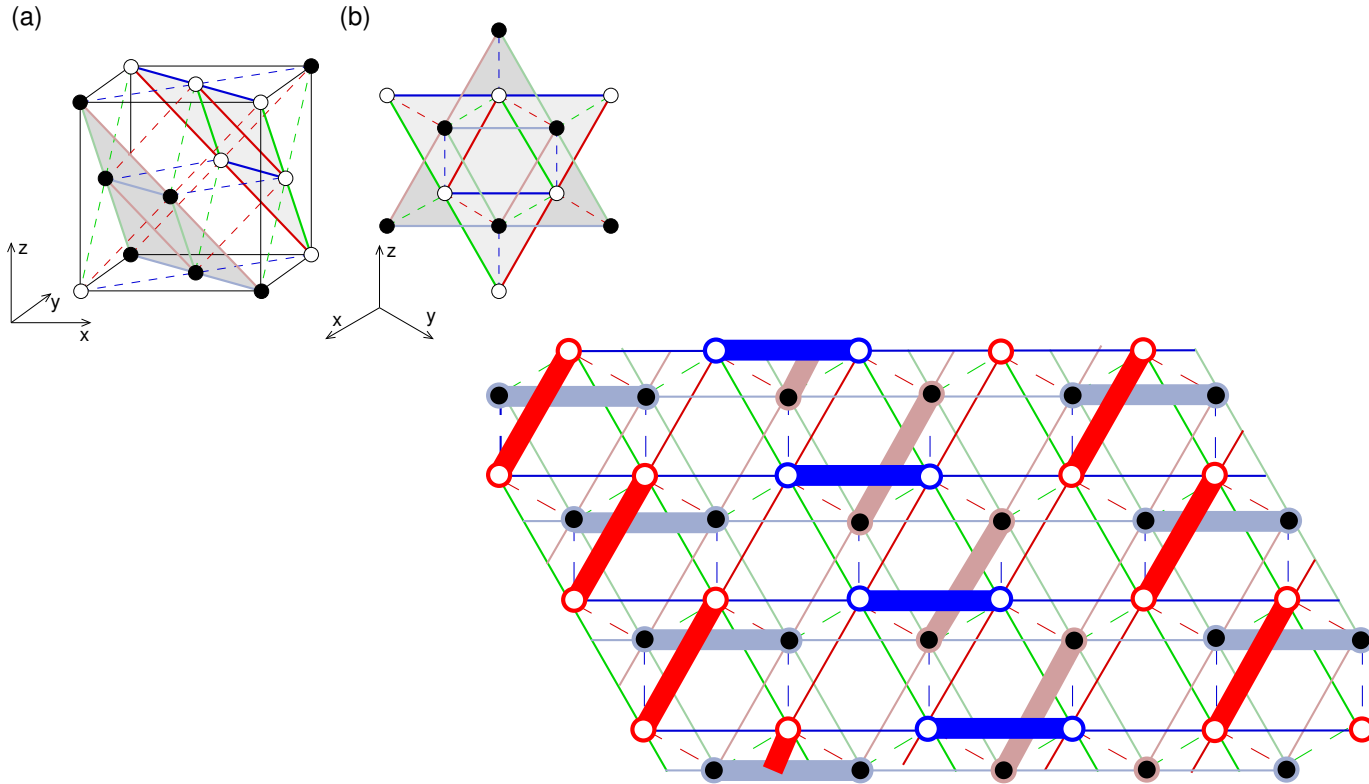
Exact ground state manifold: spanned by spin-singlet dimer coverings

$$\mathcal{H}_{\text{ex}} = J \sum_{\langle i,j \rangle} \sum_{\alpha\beta} \left(\mathbf{S}_i \cdot \mathbf{S}_j + \frac{1}{4} \right) n_{i\alpha\beta} n_{j\alpha\beta}$$

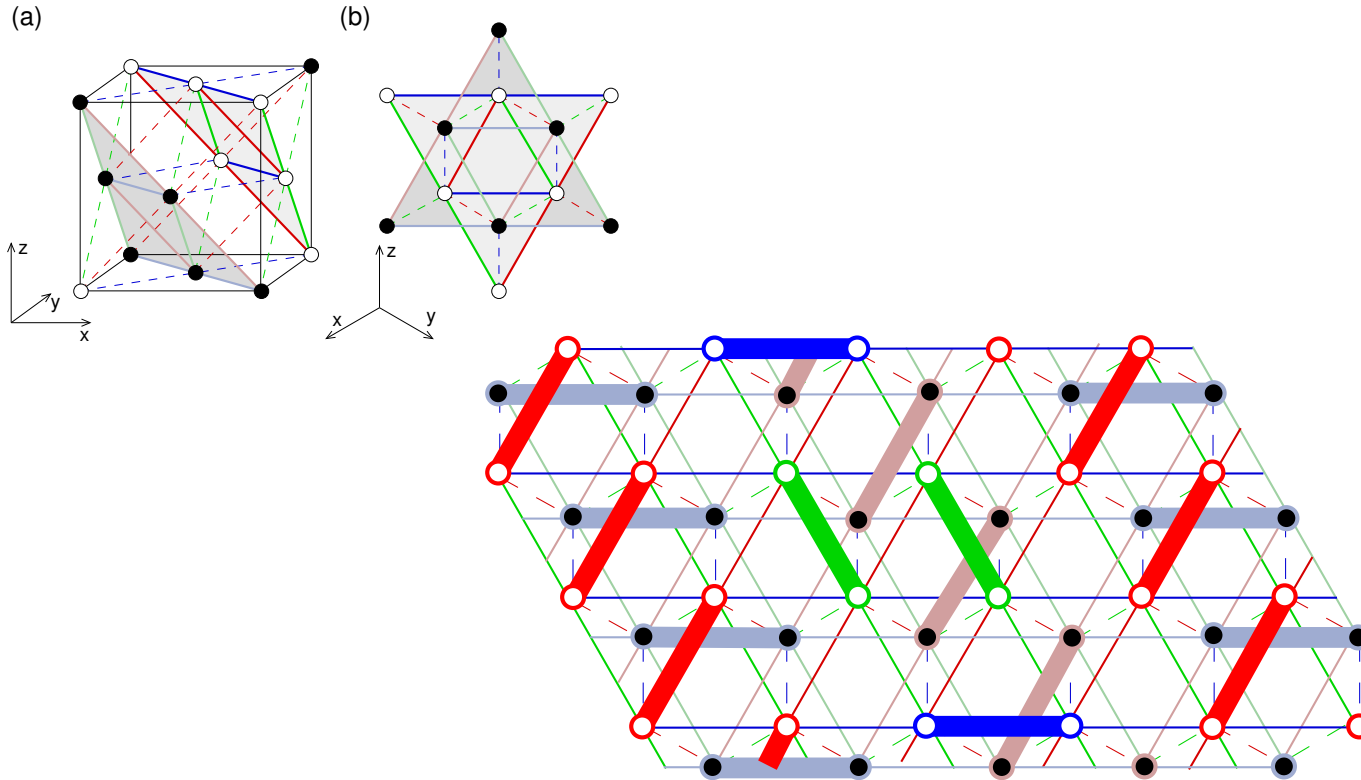


Constraint: No neighbouring dimers can lie in the same plane

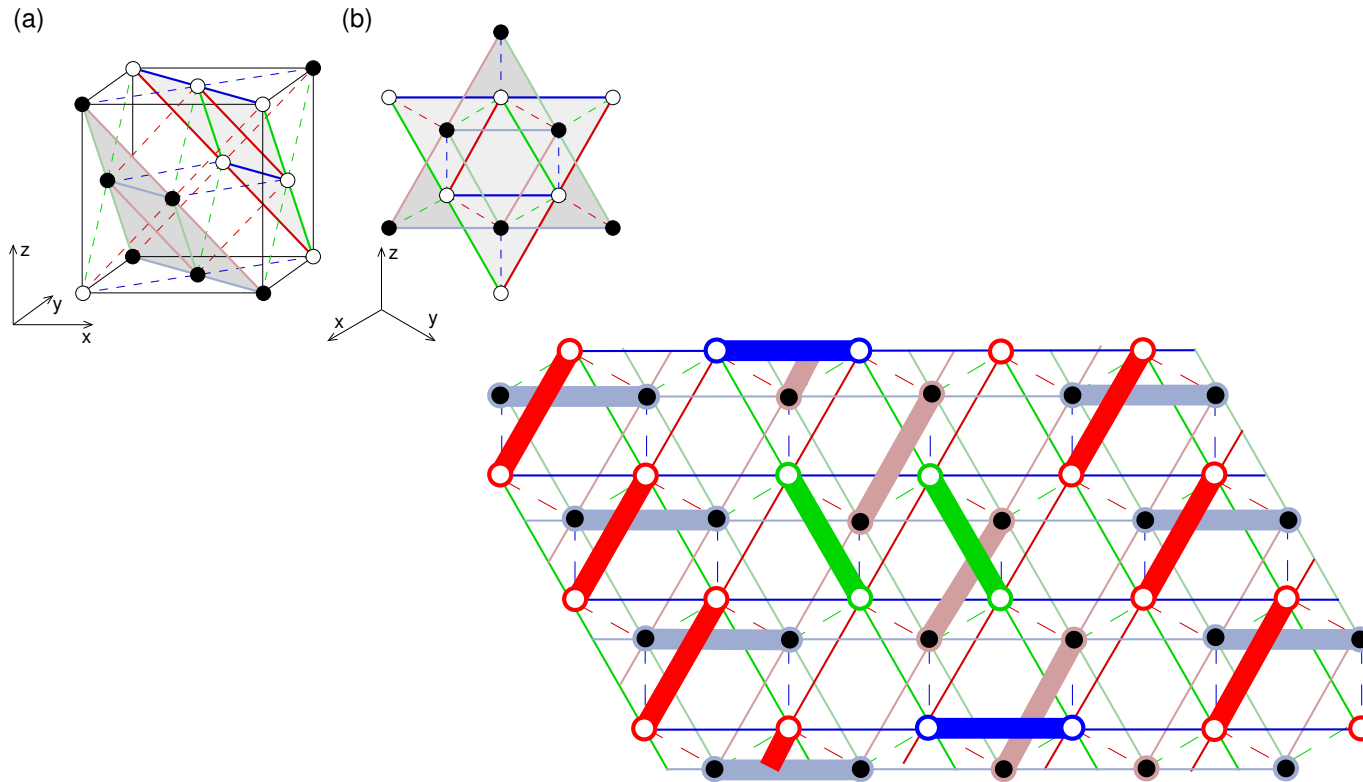
Extensive orientational degeneracy: infinitely many ways of covering



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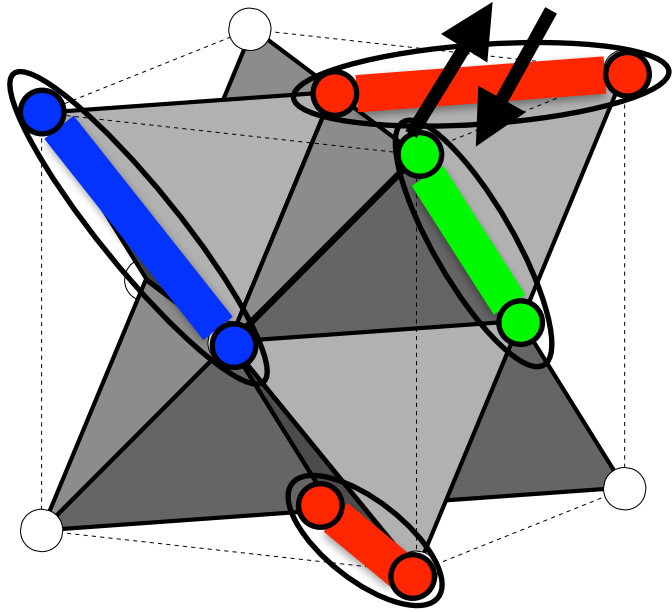


Extensive orientational degeneracy: infinitely many ways of covering

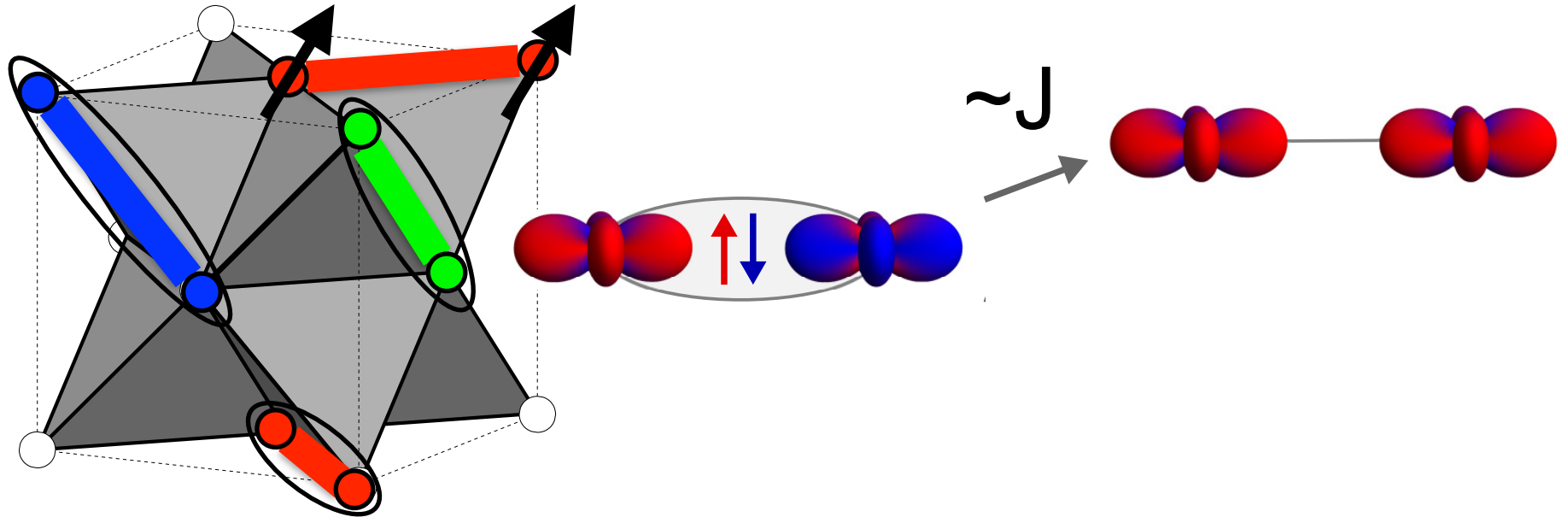


Still massively degenerate: all covering have same energy
Mis-site disorder will, most likely, choose a random pattern:
spin-orbit dimer 'glass'

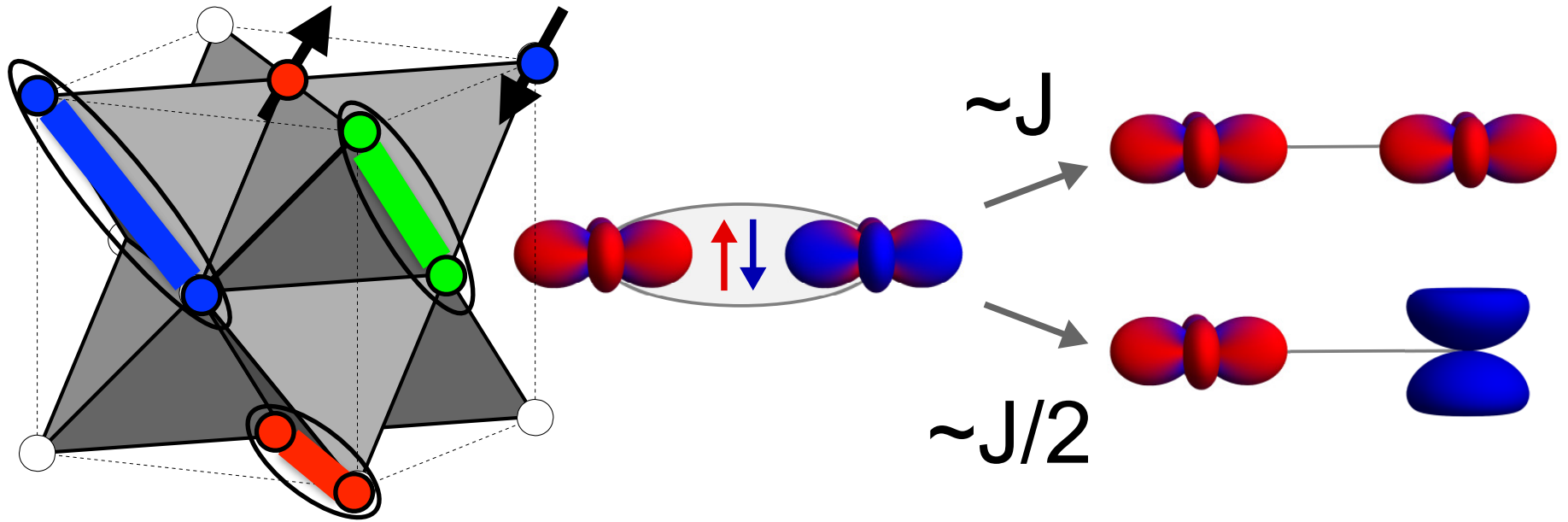
Low energy excitations & defects,



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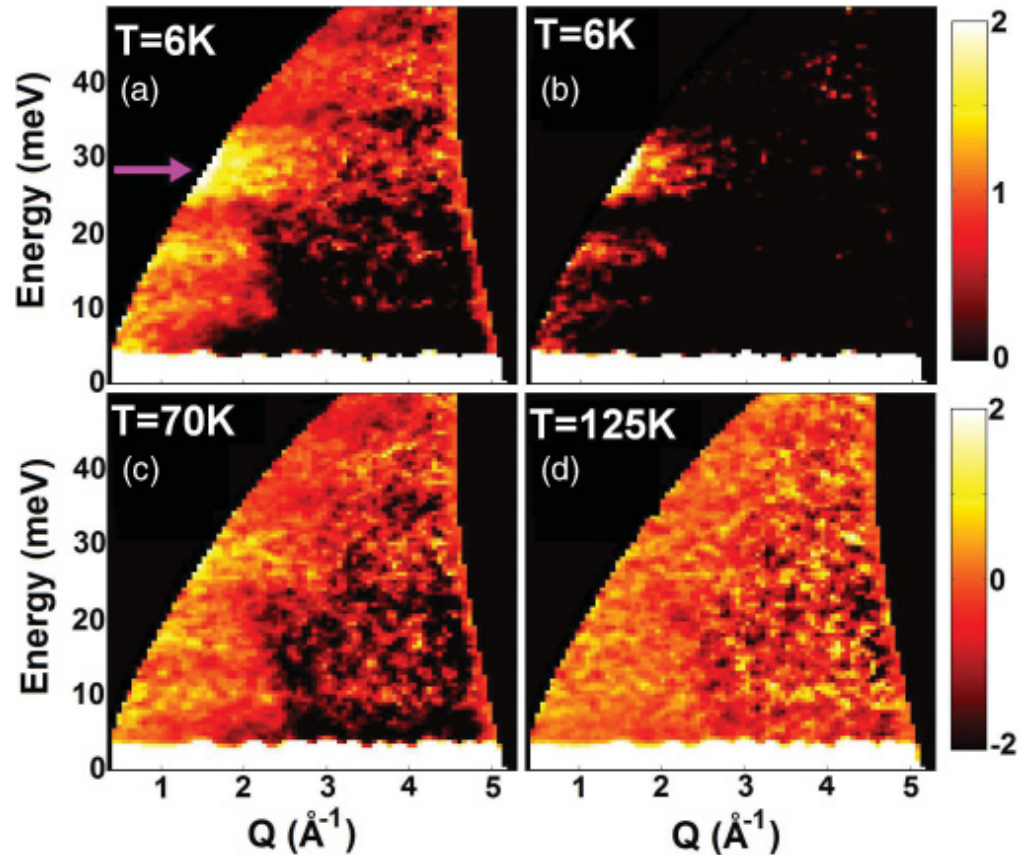
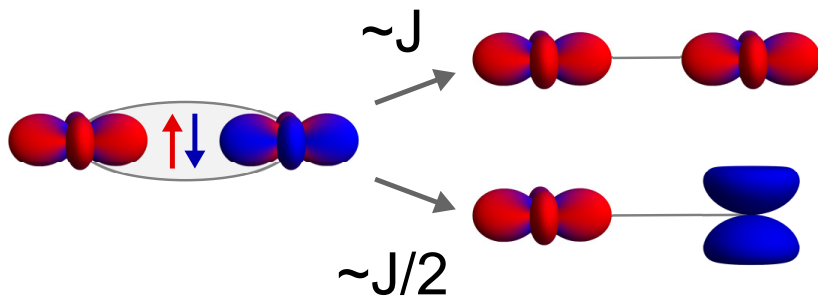


Low energy excitations & defects,

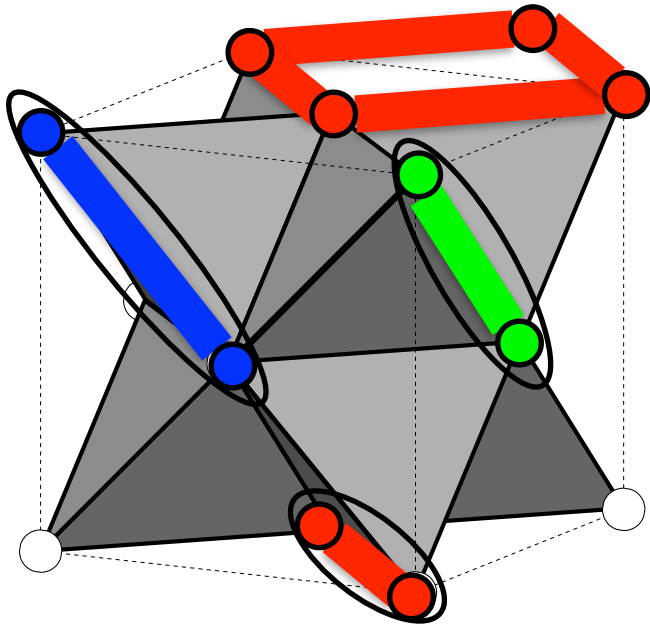


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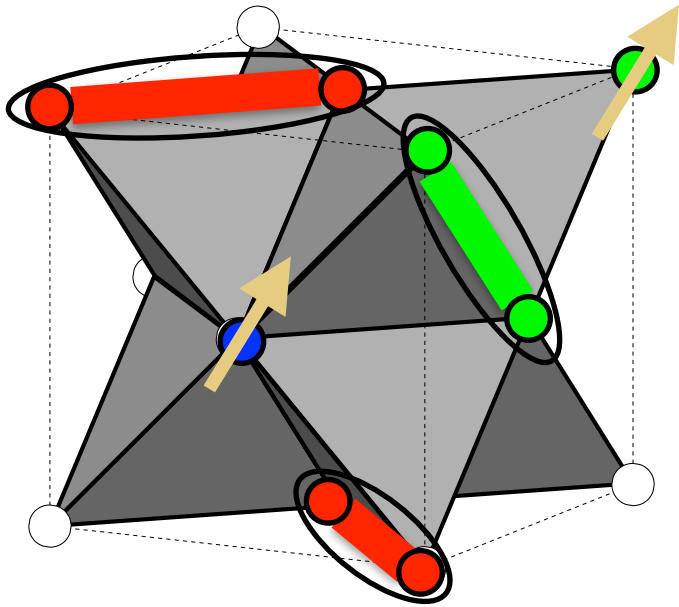
Low energy excitations & defects,



- local singlet to triplet exc: J
- local SO exc: $J/2$
- non-local SO exc: $< J/2$

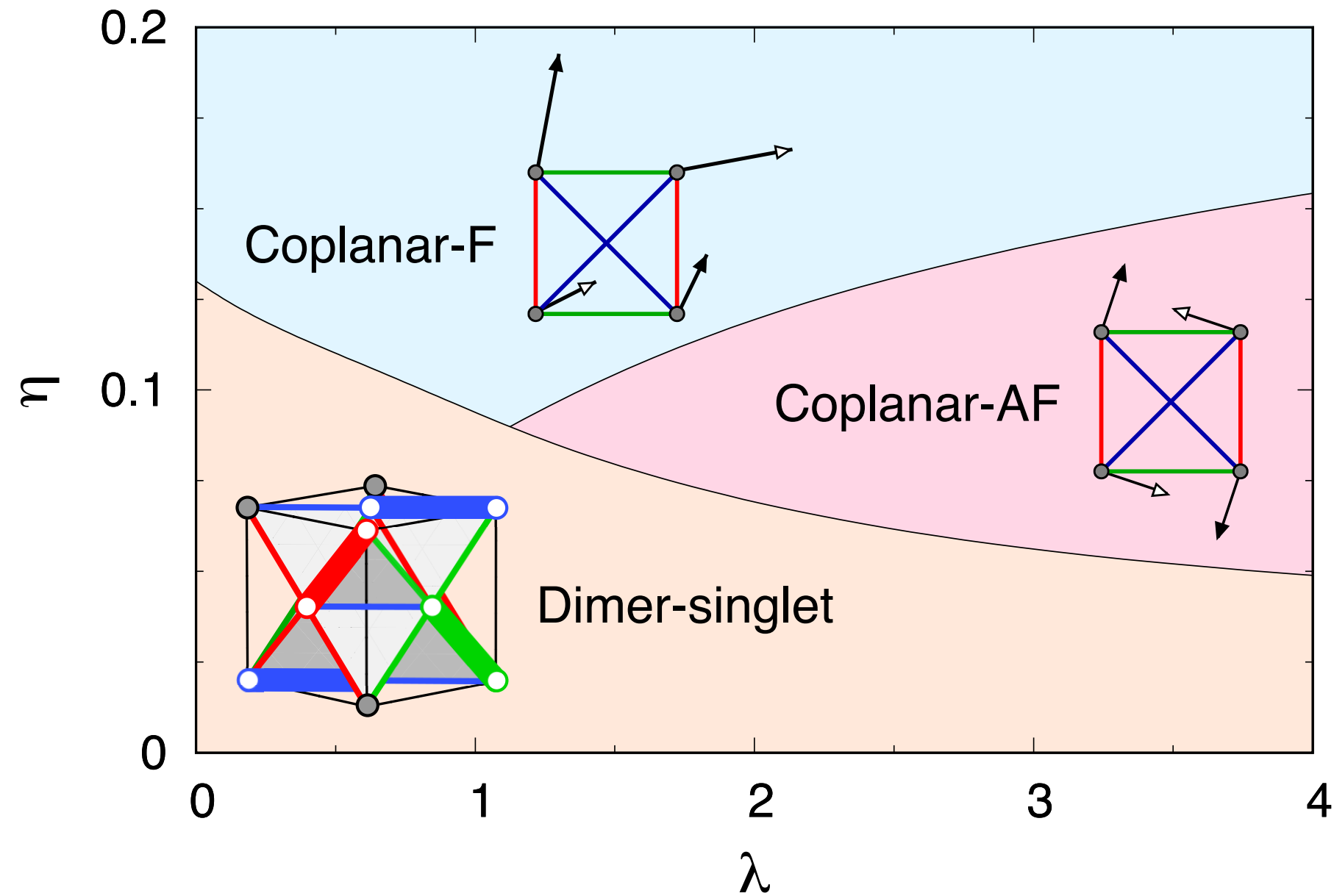
Pseudogap rather than a Hard Gap

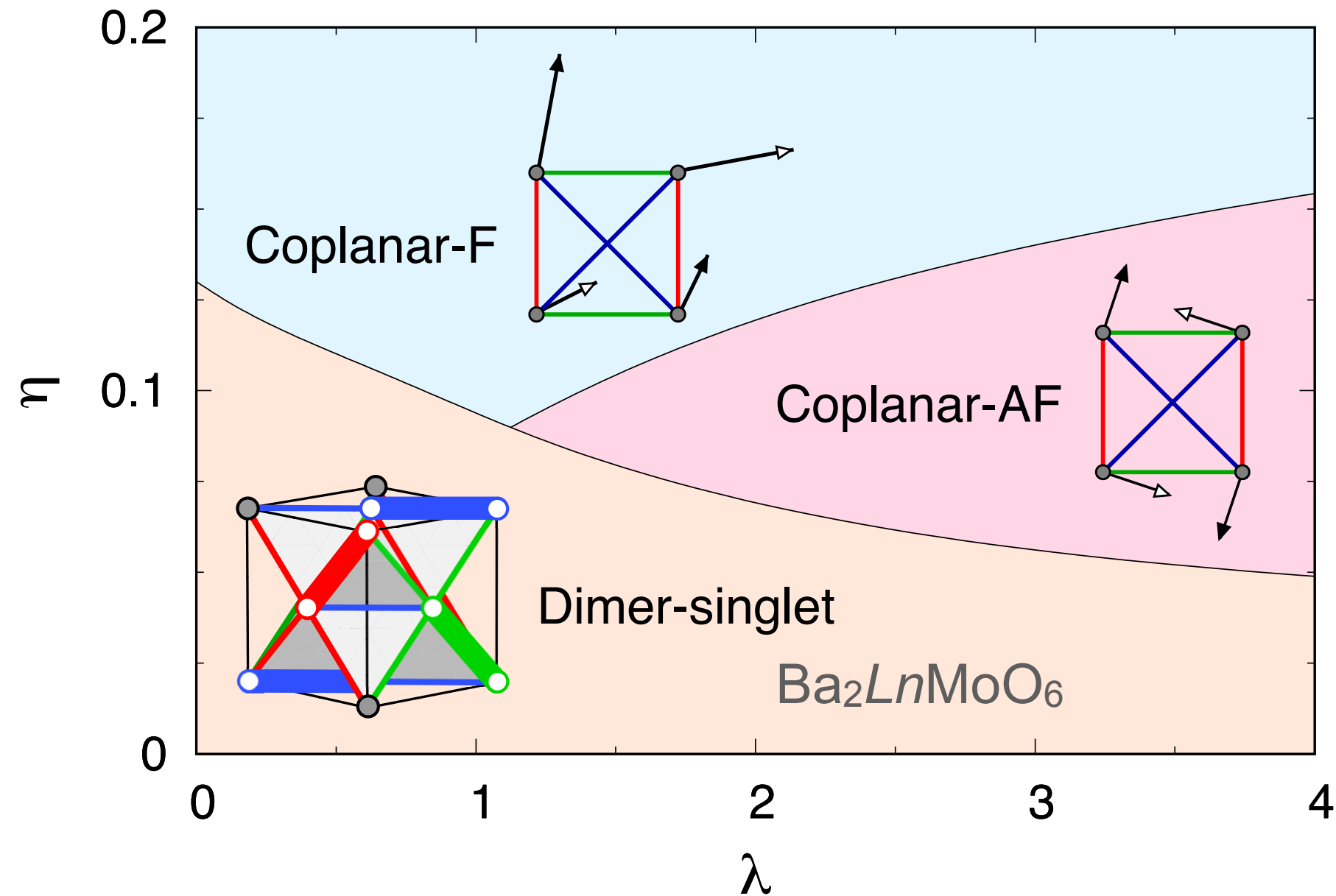
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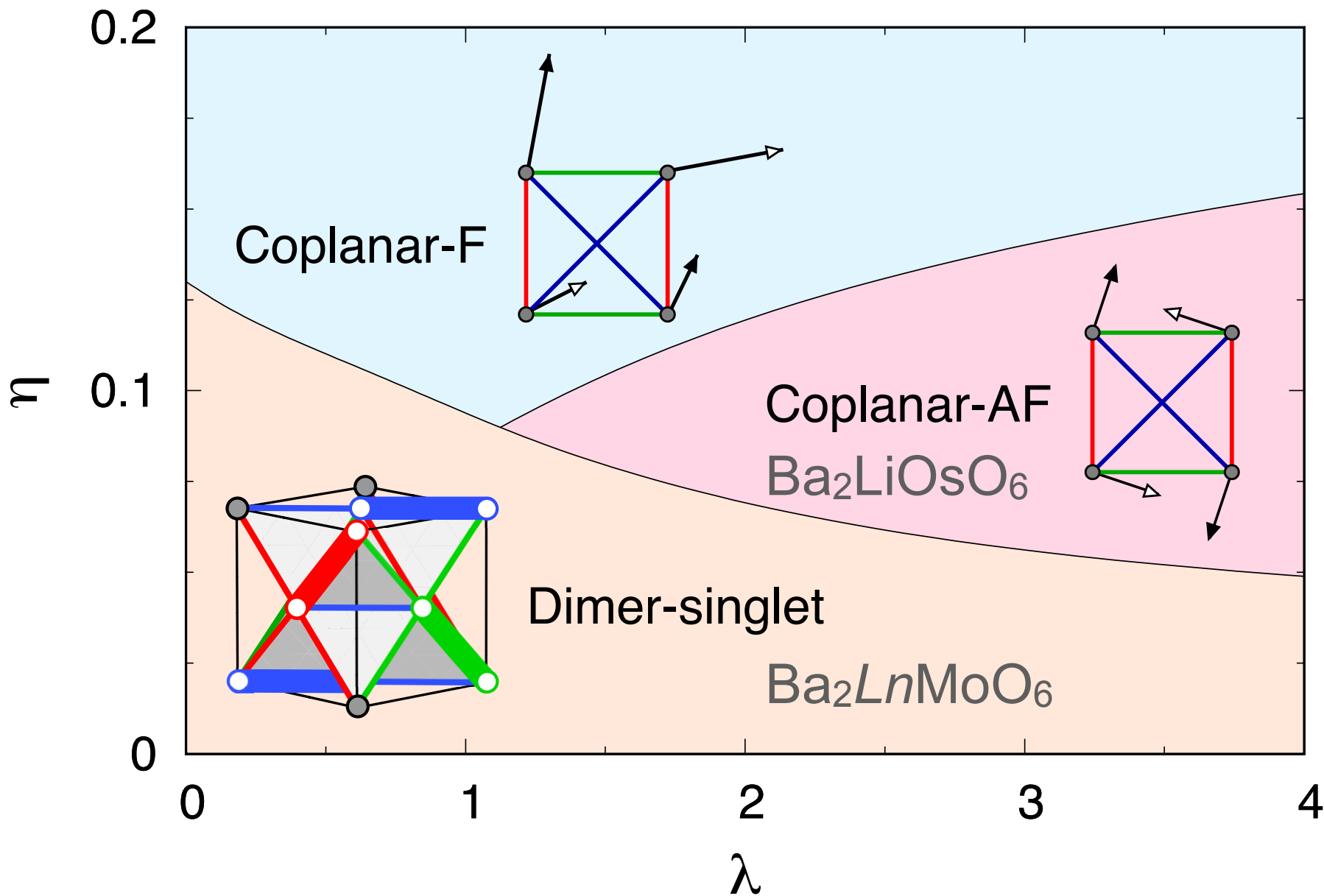


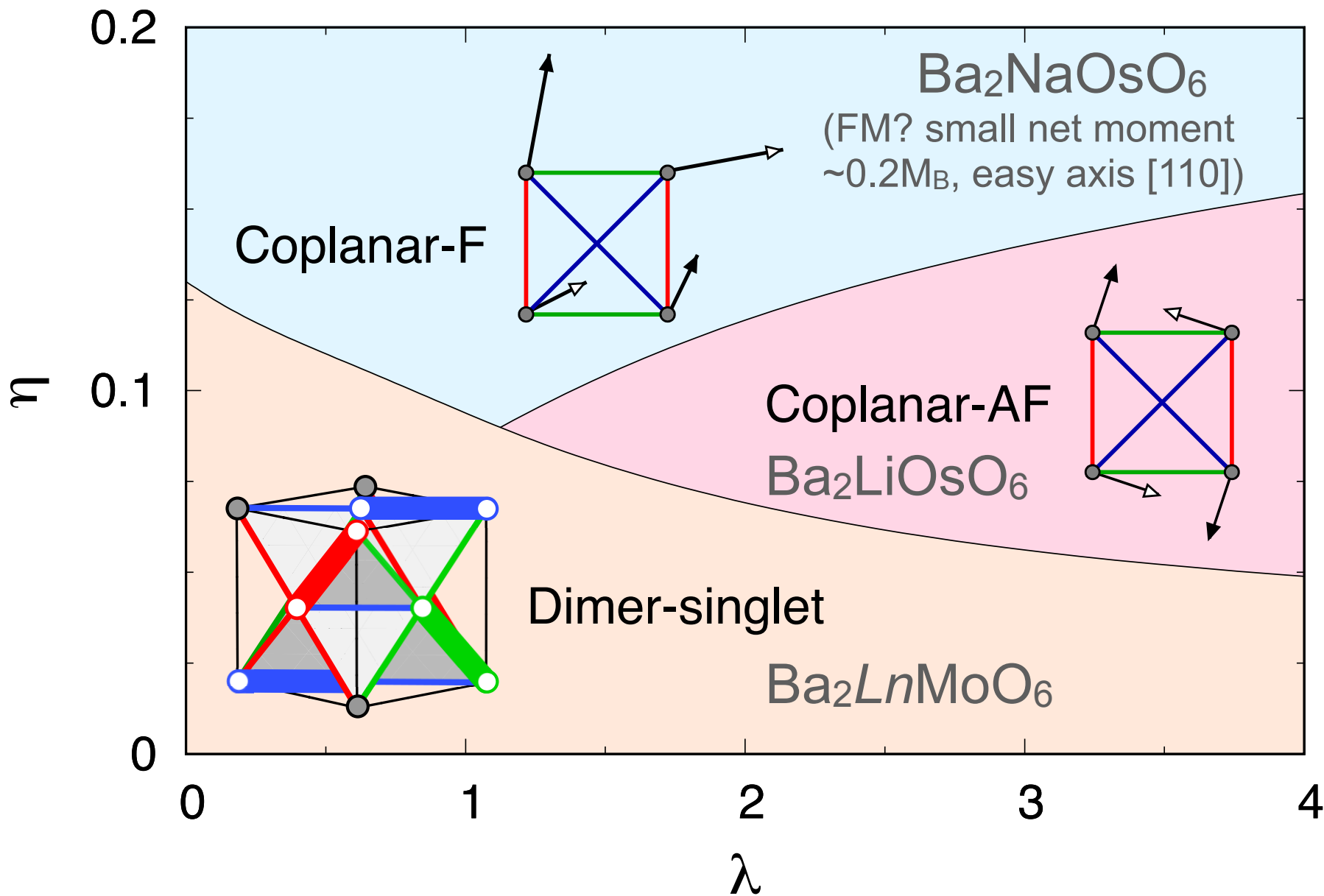
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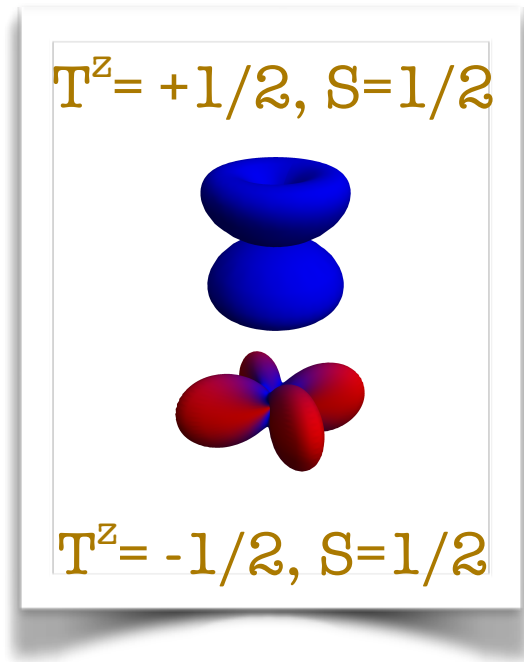






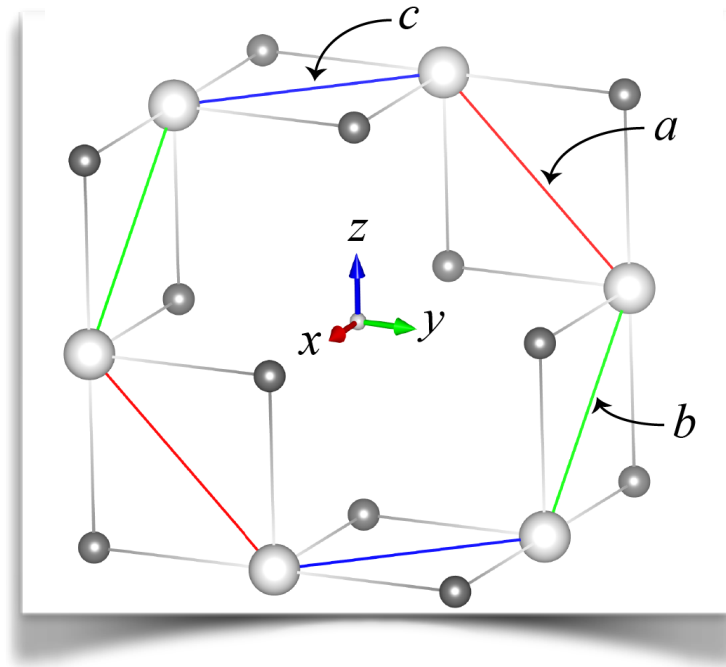
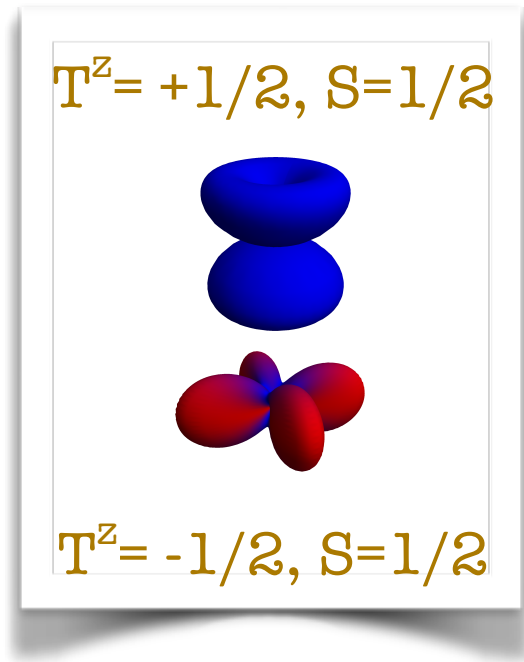
Emergent SU(4) Symmetry and More

M.G. Yamada, M. Oshikawa, & GJ, [arXiv:1709.05252](https://arxiv.org/abs/1709.05252)



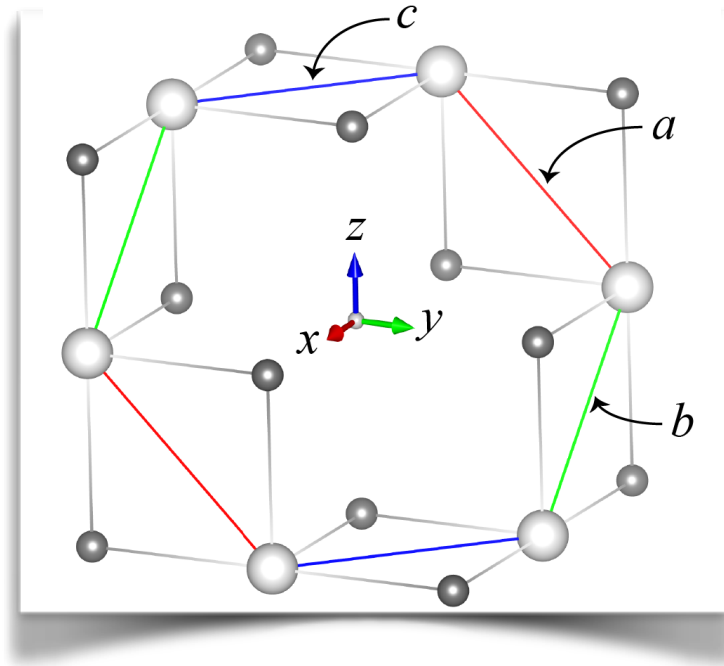
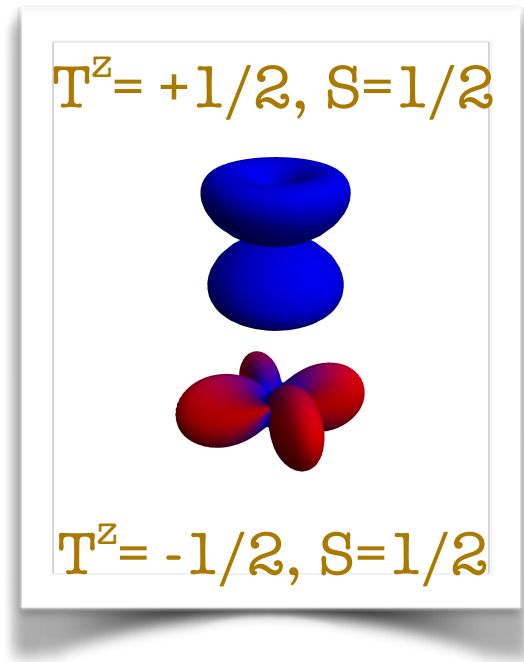
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$$H_{\text{eff}} = J \sum_{\langle ij \rangle} \left(\mathbf{S}_i \cdot \mathbf{S}_j + \frac{1}{4} \right) \left(\mathbf{T}_i \cdot \mathbf{T}_j + \frac{1}{4} \right)$$

Thank You For Listening

