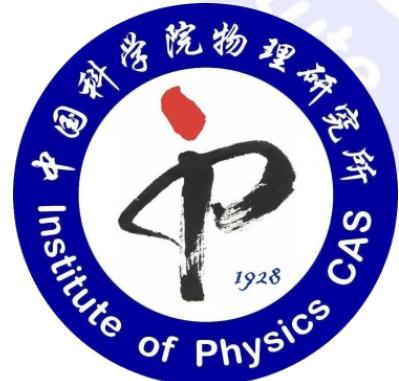


Itinerant Quantum Critical Points and Self-learning Monte Carlo Method

Zi Yang Meng

(孟子杨)

<http://ziyangmeng.iphy.ac.cn/>



Itinerant Quantum Critical Points



Xiao Yan Xu



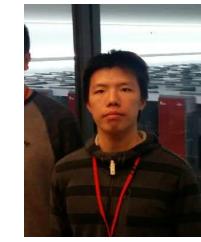
Yoni Schattner



Erez Berg



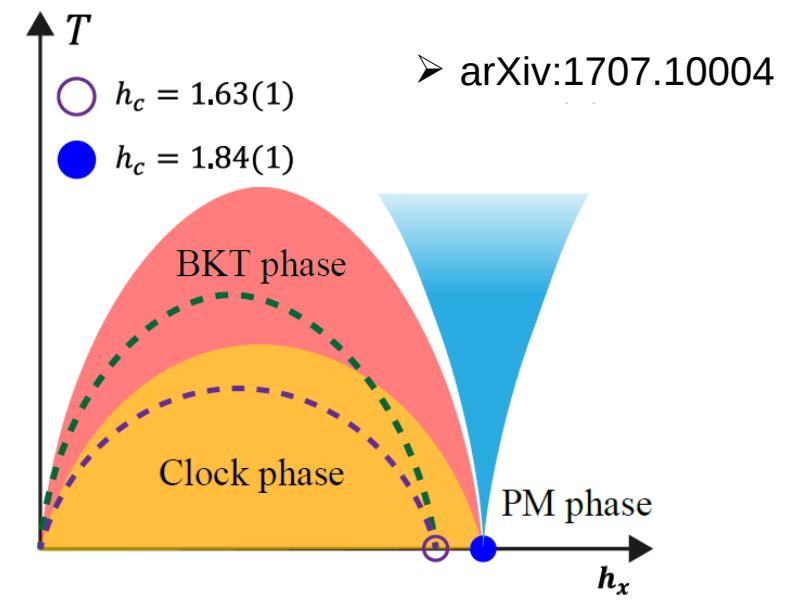
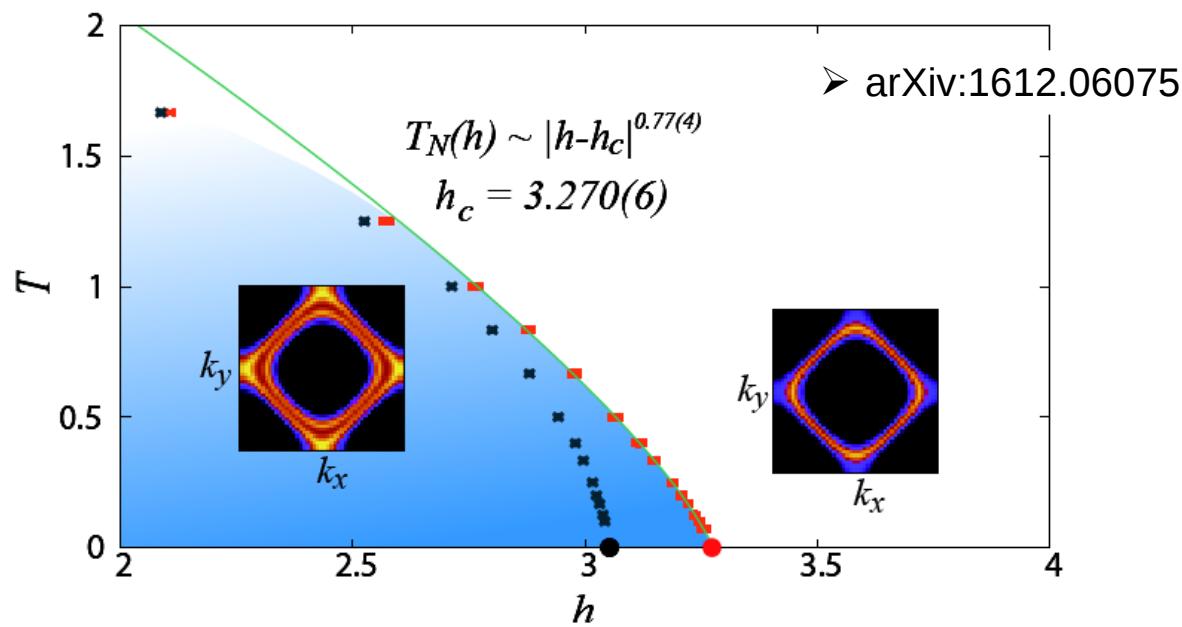
Kai Sun



Zi Hong Liu

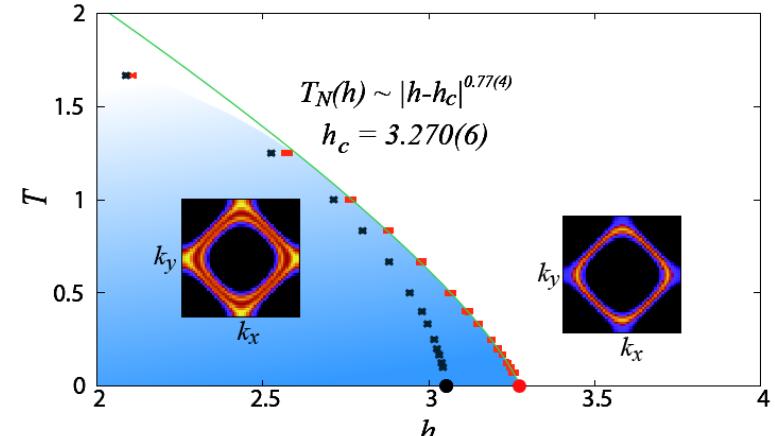


Yang Qi



Content

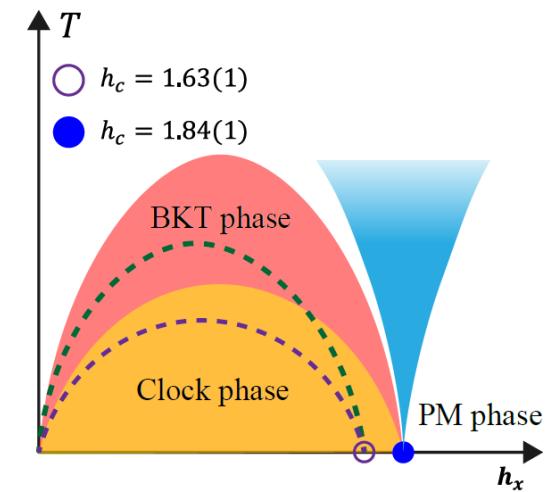
- Fermi surface coupled to critical bosonic fluctuations



➤ arXiv:1612.06075

- Results

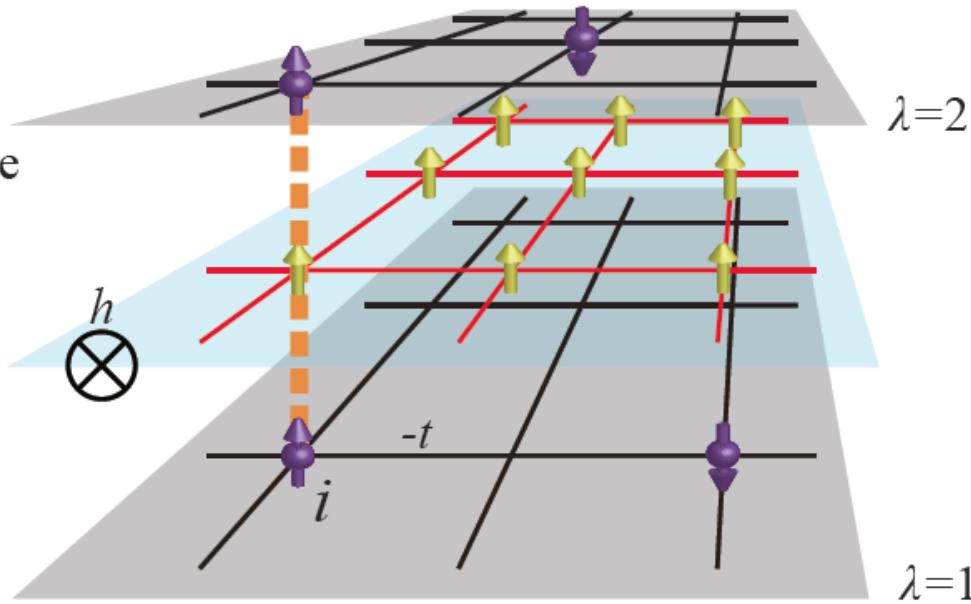
- phase diagram
- Non-Fermi liquid and itinerant FM/AFM-QCPs



➤ arXiv:1707.10004

Model

-  fermion site
-  Ising site
-  coupling
-  Ising spin
-  fermion

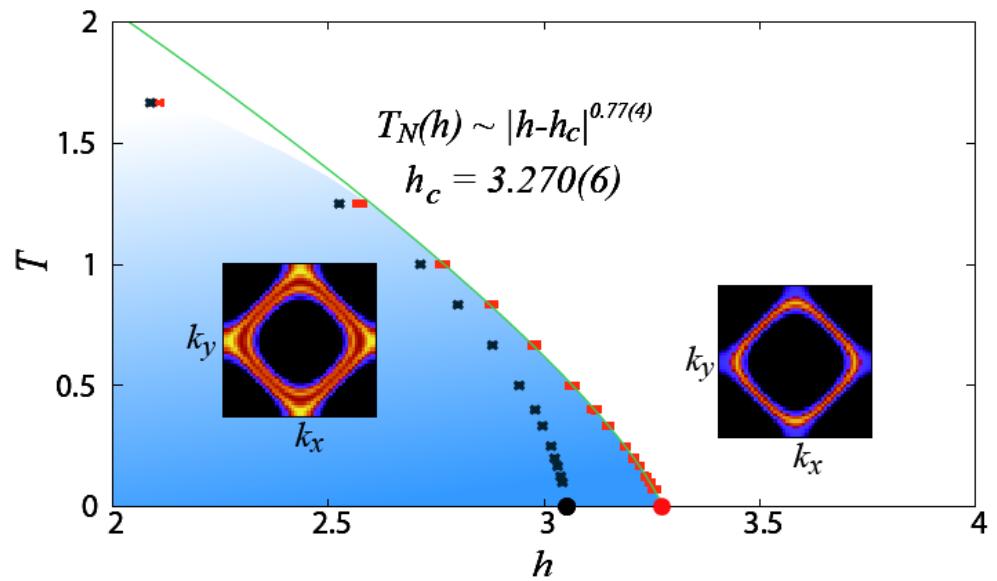


$$\hat{H} = \hat{H}_f + \hat{H}_s + \hat{H}_{sf}$$

$$\hat{H}_f = -t \sum_{\langle ij \rangle \lambda \sigma} \hat{c}_{i\lambda\sigma}^\dagger \hat{c}_{j\lambda\sigma} + h.c. - \mu \sum_{i\lambda\sigma} \hat{n}_{i\lambda\sigma}$$

$$\hat{H}_s = -J \sum_{\langle ij \rangle} \hat{s}_i^z \hat{s}_j^z - h \sum_i \hat{s}_i^x$$

$$\hat{H}_{sf} = -\xi \sum_i s_i^z (\hat{\sigma}_{i1}^z + \hat{\sigma}_{i2}^z)$$



Quantum Monte Carlo

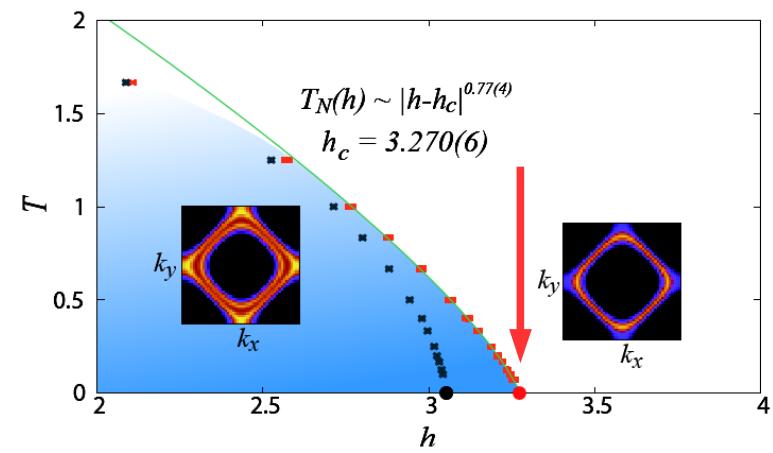
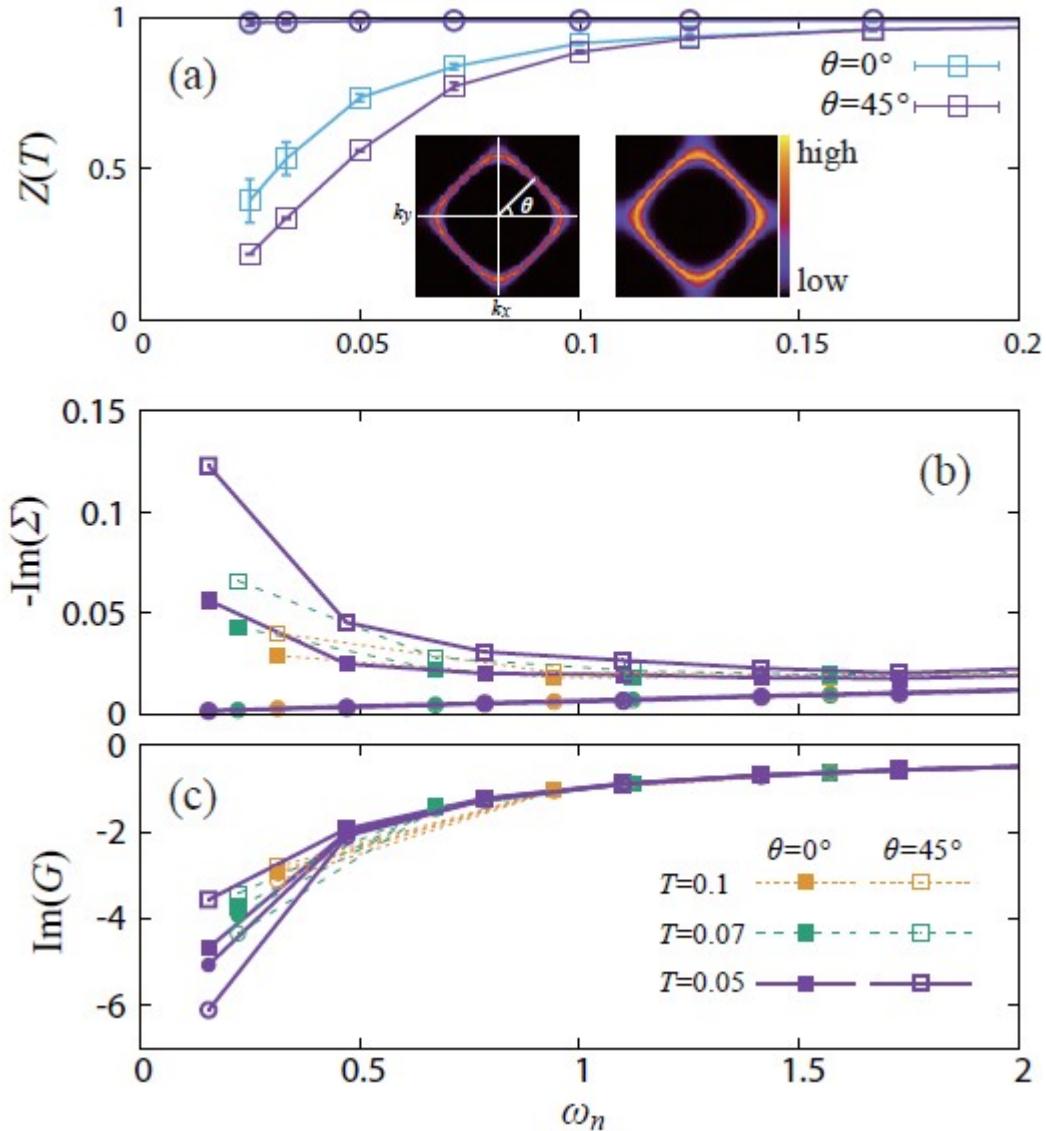
$$\begin{aligned} Z &= \mathbf{Tr} \left[e^{-\beta \hat{H}} \right] \\ &= \sum_{s_{b_1} \cdots s_{b_N} = \pm 1} \mathbf{Tr}_F \left\langle s_{b_1} \cdots s_{b_N} \left| \left(e^{-\Delta \tau \hat{H}} \right)^M \right| s_{b_1} \cdots s_{b_N} \right\rangle \\ &= \sum_{\vec{S}_1 \cdots \vec{S}_M} \mathbf{Tr}_F \langle \vec{S}_1 | e^{-\Delta \tau \hat{H}} | \vec{S}_M \rangle \langle \vec{S}_M | e^{-\Delta \tau \hat{H}} | \vec{S}_{M-1} \rangle \cdots \langle \vec{S}_2 | e^{-\Delta \tau \hat{H}} | \vec{S}_1 \rangle \\ &= \sum_{\mathcal{C}} \omega_{\mathcal{C}} \text{ with configuration } \mathcal{C} = \left\{ \vec{S}_1, \dots, \vec{S}_M \right\} \end{aligned}$$

$$\omega_{\mathcal{C}} = \left(\prod_{\tau} \prod_{\langle i,j \rangle} e^{\Delta \tau J s_{i,\tau} s_{j,\tau}} \right) \left(\prod_i \prod_{\langle \tau, \tau' \rangle} \Lambda e^{\gamma s_{\tau,i} s_{\tau',i}} \right) \left| \prod_{\sigma} \det \left(\mathbf{B}_M^{1\sigma} \cdots \mathbf{B}_1^{1\sigma} \right) \right|^2$$

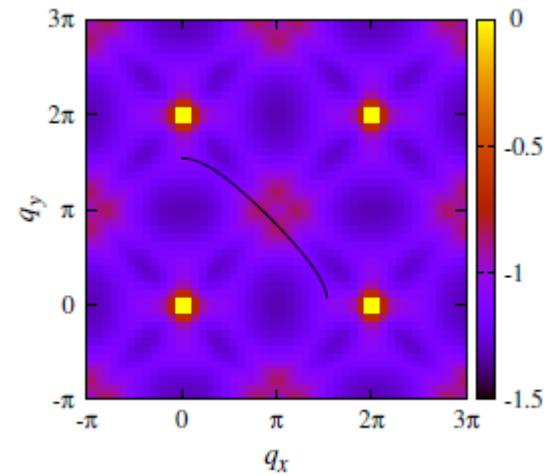
➤ Self-learning Monte Carlo Method

O(N) speedup, large lattice is possible

Non-Fermi liquid

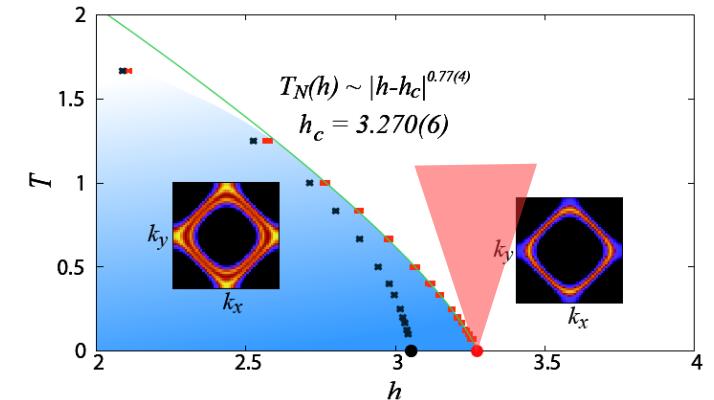


$$Z_{\mathbf{k}_F} \approx \frac{1}{1 - \frac{\text{Im}\Sigma(\mathbf{k}_F, i\omega_0)}{\omega_0}}$$



FM-QCP

$$\begin{aligned}\chi(h, T, \mathbf{q}, i\omega_n) &= \frac{1}{L^2} \sum_{ij} \int_0^\beta d\tau e^{i\omega_n \tau - i\mathbf{q} \cdot \mathbf{r}_{ij}} \langle s_i(\tau) s_j(0) \rangle \\ &= \frac{1}{c_t T^{a_t} + c_h |h - h_c|^\gamma + (c_q q^2 + c_\omega \omega^2)^{a_q/2} + \Delta(\mathbf{q}, \omega_n)}\end{aligned}$$



$$a_t = 2 \quad \gamma = 1 \quad a_q = 2 \quad \Delta(\mathbf{q}, \omega_n) = c_{HM} \frac{|\omega_n|}{\sqrt{\omega_n^2 + (v_f q)^2}}$$

$$a_q = 2 - \eta$$

$$\nu = \gamma/a_q$$

$$\chi(h=h_c, T=0, \mathbf{q}, \omega_n=0)^{-1} = c_q q^2$$

$$\chi(h=h_c, T=0, \mathbf{q}=0, \omega_n)^{-1} = c_{HM} + c_\omega \omega_n^2$$

$\eta = 0$	$\nu = 0.5$	$\gamma = 1$	$z = 3$	Hertz-Millis-Moriya
$\eta = 0.04$	$\nu = 0.63$	$\gamma = 1.24$	$z = 1$	(2+1)d Ising model
$\eta = 0.15(3)$	$\nu = 0.64(3)$	$\gamma = 1.18(4)$		Our model

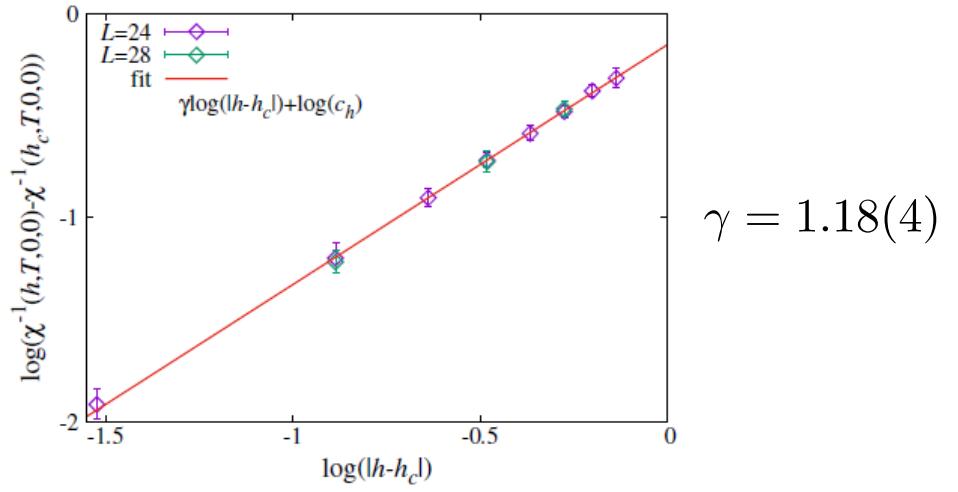
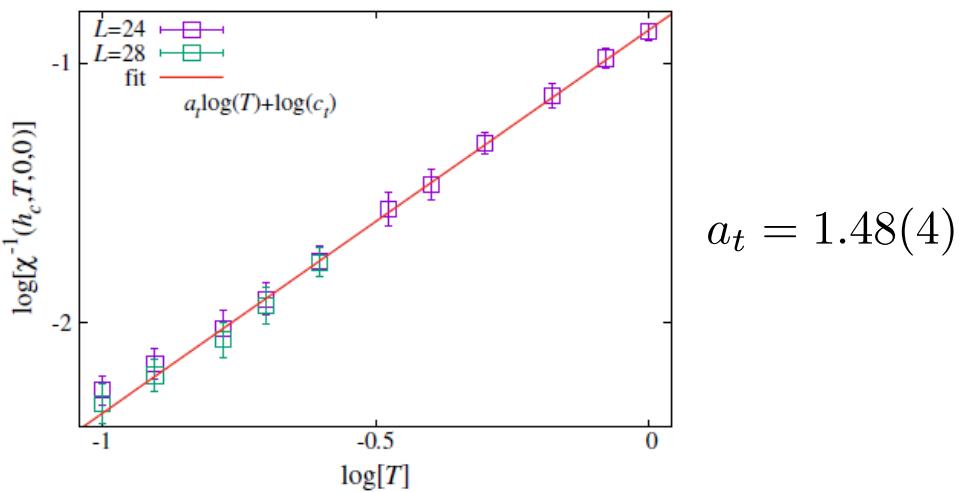
FM-QCP

$$\chi(h, T, \mathbf{q}, i\omega_n) = \frac{1}{L^2} \sum_{ij} \int_0^\beta d\tau e^{i\omega_n \tau - i\mathbf{q} \cdot \mathbf{r}_{ij}} \langle s_i(\tau) s_j(0) \rangle$$

$$= \frac{1}{c_t T^{a_t} + c_h |h - h_c|^\gamma + (c_q q^2 + c_\omega \omega^2)^{a_q/2} + \Delta(\mathbf{q}, \omega_n)}$$

$$\chi(h_c, T, \mathbf{q} = 0, \omega = 0) = \frac{1}{c_t T^{a_t}}$$

$$\chi(h, T = 0, \mathbf{q} = 0, \omega = 0) = \frac{1}{c_h |h - h_c|^\gamma}$$



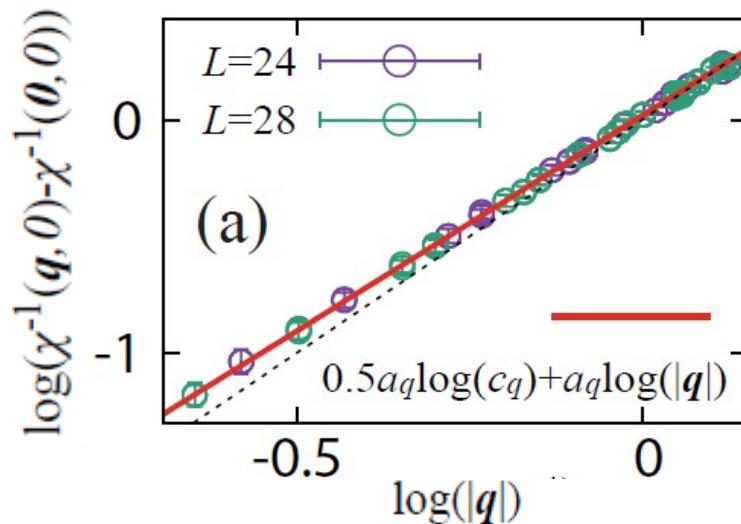
$\eta = 0$	$\nu = 0.5$	$\gamma = 1$	$z = 3$	Hertz-Millis-Moriya
$\eta = 0.04$	$\nu = 0.63$	$\gamma = 1.24$	$z = 1$	(2+1)d Ising model
$\eta = 0.15(3)$	$\nu = 0.64(3)$	$\gamma = 1.18(4)$		Our model

FM-QCP

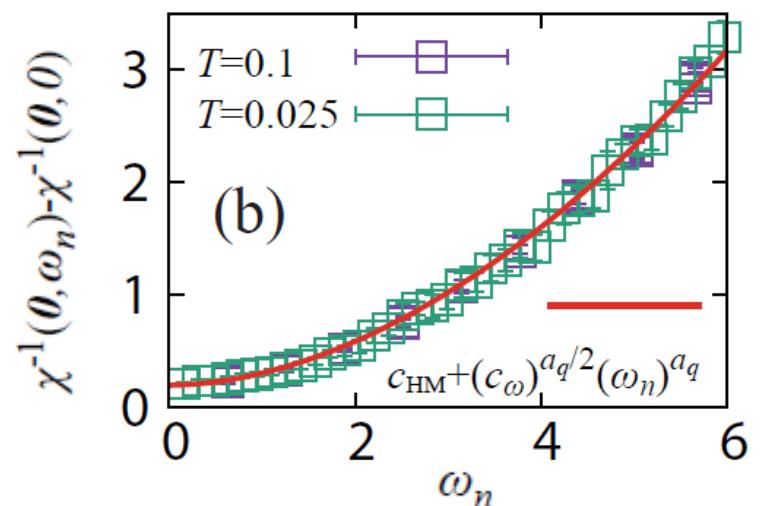
$$\begin{aligned}\chi(h, T, \mathbf{q}, i\omega_n) &= \frac{1}{L^2} \sum_{ij} \int_0^\beta d\tau e^{i\omega_n \tau - i\mathbf{q} \cdot \mathbf{r}_{ij}} \langle s_i(\tau) s_j(0) \rangle \\ &= \frac{1}{c_t T^{a_t} + c_h |h - h_c|^\gamma + (c_q q^2 + c_\omega \omega^2)^{a_q/2} + \Delta(\mathbf{q}, \omega_n)}\end{aligned}$$

$$\chi(h = h_c, T = 0, \mathbf{q}, \omega = 0) = \frac{1}{c_q^{a_q/2} |\mathbf{q}|^{a_q}}$$

$$\chi(h = h_c, T = 0, \mathbf{q} = 0, \omega) = \frac{1}{c_\omega^{a_q/2} (\omega_n)^{a_q} + c_{HM}}$$



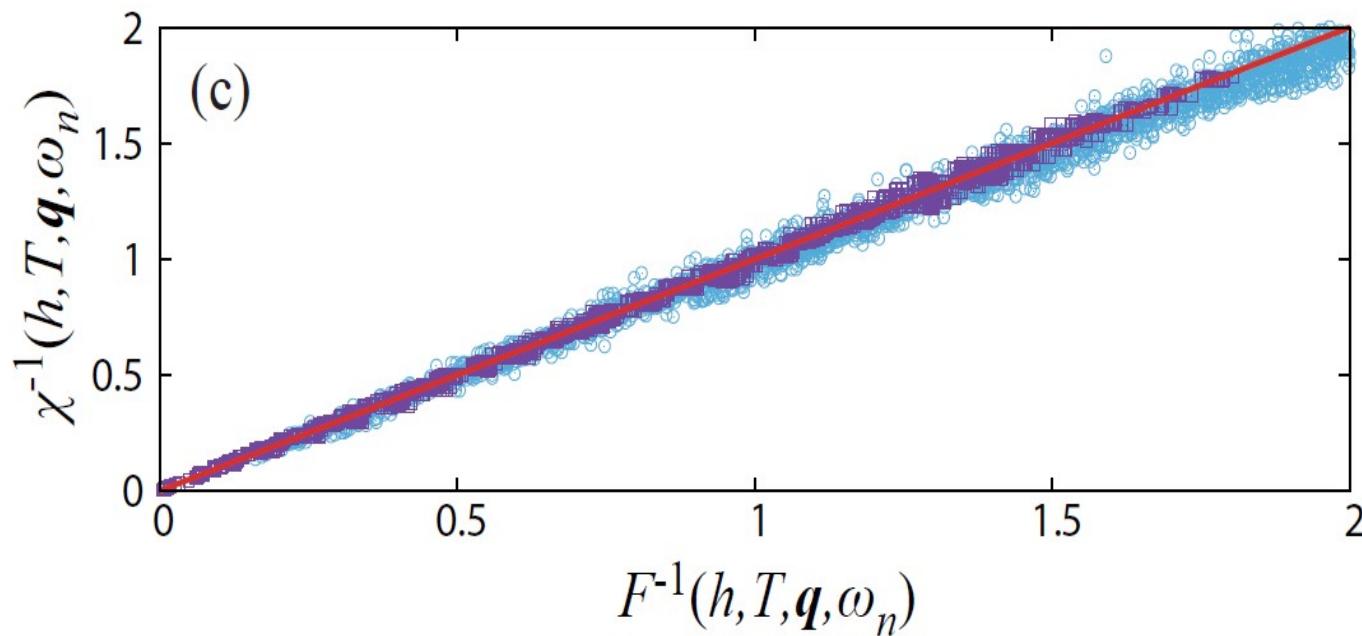
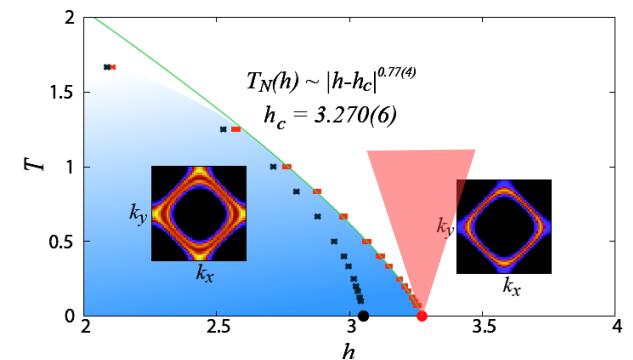
$$\begin{aligned}a_q &= 2 - \eta \\ a_q &= 1.85(3) \\ \nu &= \gamma/a_q\end{aligned}$$



$\eta = 0$	$\nu = 0.5$	$\gamma = 1$	$z = 3$	Hertz-Millis-Moriya
$\eta = 0.04$	$\nu = 0.63$	$\gamma = 1.24$	$z = 1$	(2+1)d Ising model
$\eta = 0.15(3)$	$\nu = 0.64(3)$	$\gamma = 1.18(4)$		Our model

FM-QCP

$$\begin{aligned}\chi(h, T, \mathbf{q}, i\omega_n) &= \frac{1}{L^2} \sum_{ij} \int_0^\beta d\tau e^{i\omega_n \tau - i\mathbf{q} \cdot \mathbf{r}_{ij}} \langle s_i(\tau) s_j(0) \rangle \\ &= \frac{1}{c_t T^{a_t} + c_h |h - h_c|^\gamma + (c_q q^2 + c_\omega \omega^2)^{a_q/2} + \Delta(\mathbf{q}, \omega_n)}\end{aligned}$$

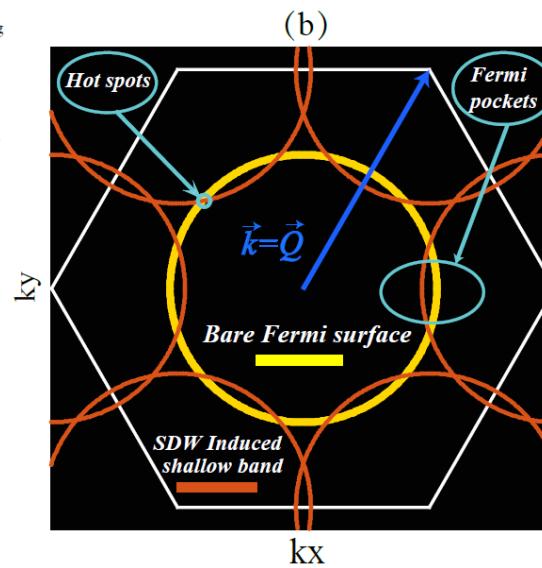
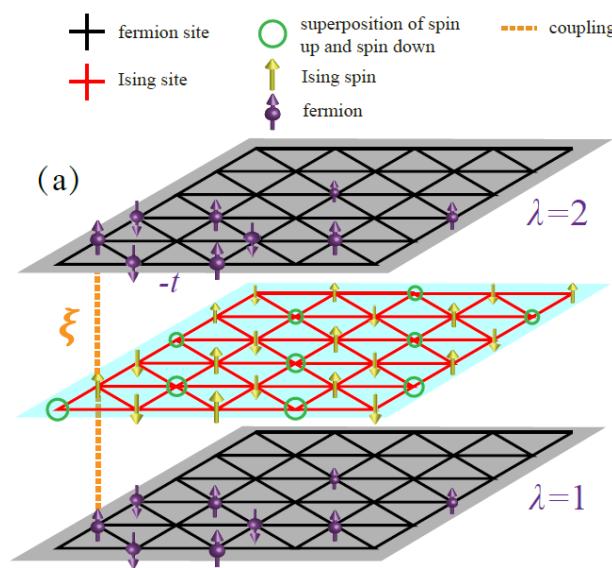


$\eta = 0$	$\nu = 0.5$	$\gamma = 1$	$z = 3$	Hertz-Millis-Moriya
$\eta = 0.04$	$\nu = 0.63$	$\gamma = 1.24$	$z = 1$	(2+1)d Ising model
$\eta = 0.15(3)$	$\nu = 0.64(3)$	$\gamma = 1.18(4)$		Our model

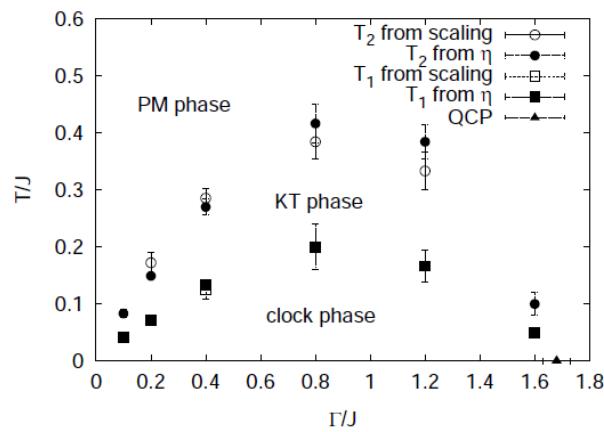
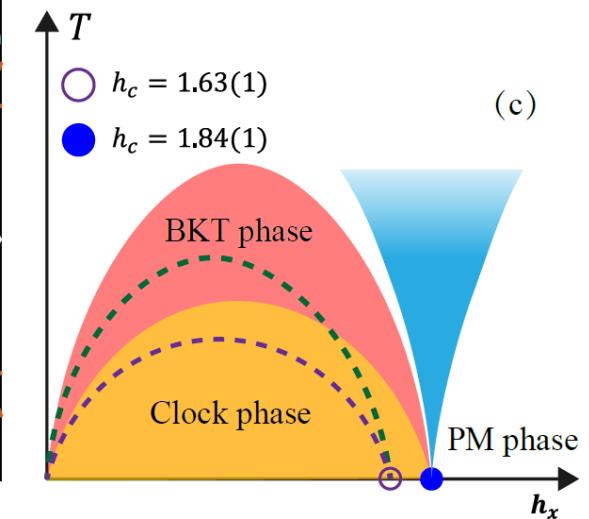
Triangle lattice



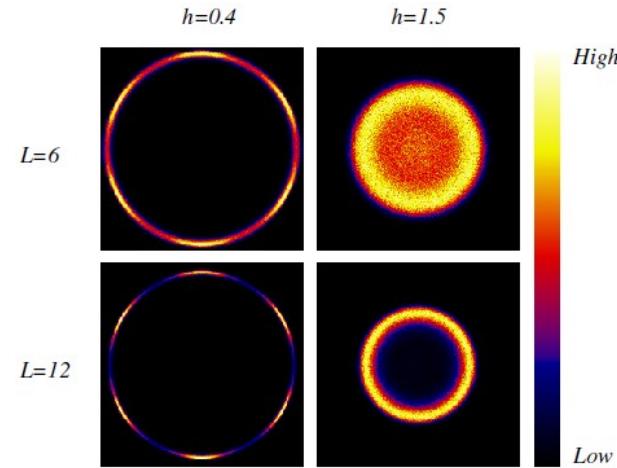
Zi Hong Liu



➤ arXiv:1707.10004



S. V. Isakov and R. Moessner, PRB 68, 104409 (2003)

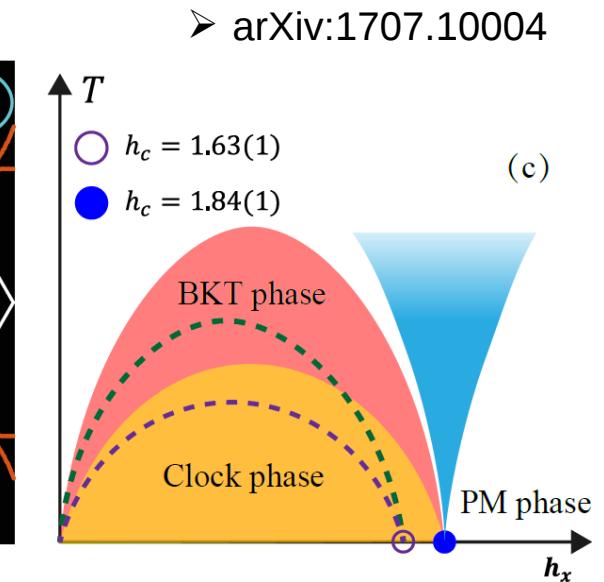
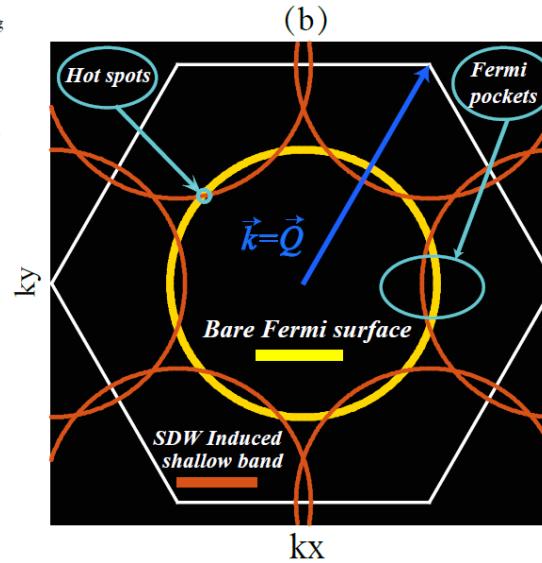
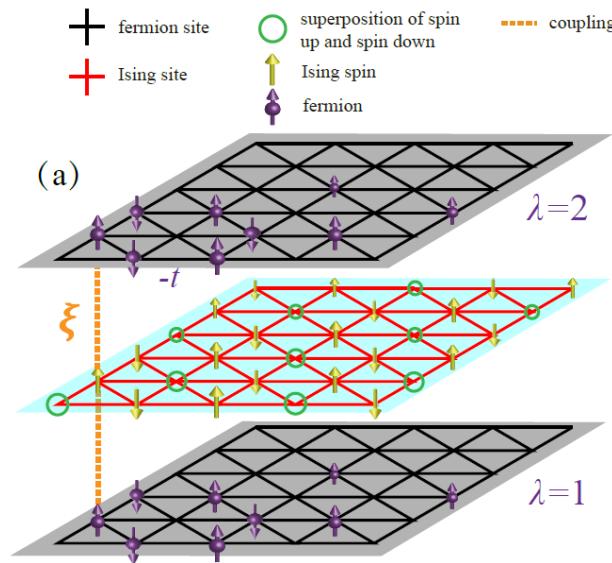


Y.-C. Wang, Y. Qi, S. Chen, Z. Y. Meng, arXiv:1602.02839

Triangle lattice

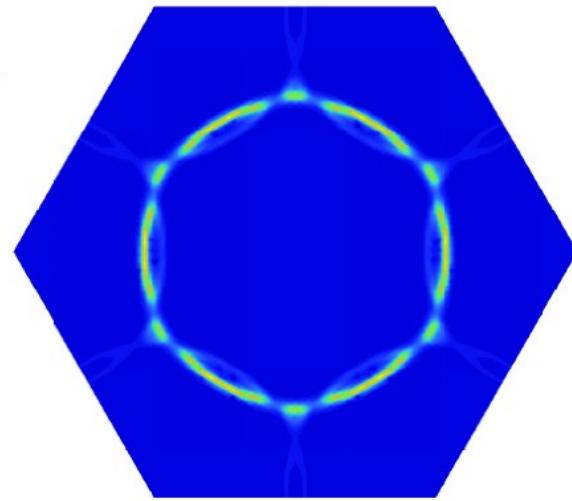


Zi Hong Liu

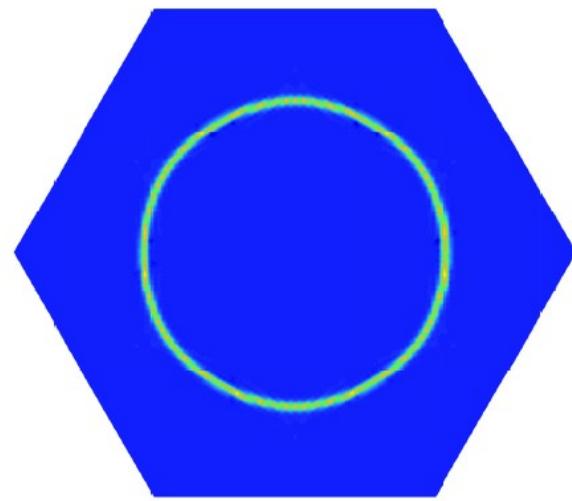


$L=30$, $\beta=30$
($30 \times 30 \times 600$)

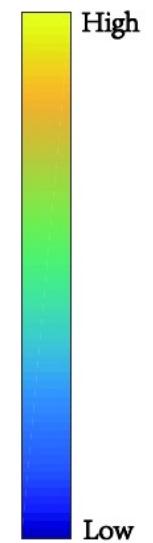
k_x
 k_y



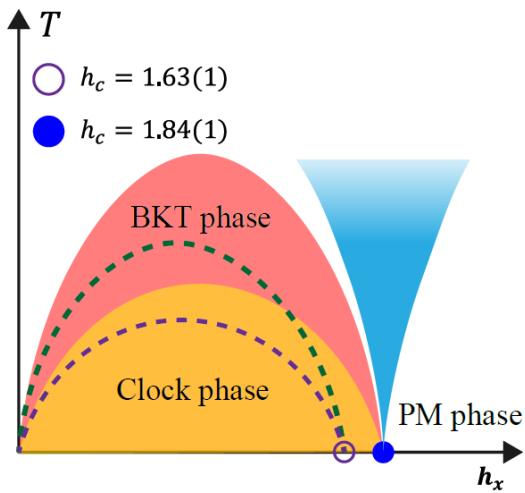
(a)



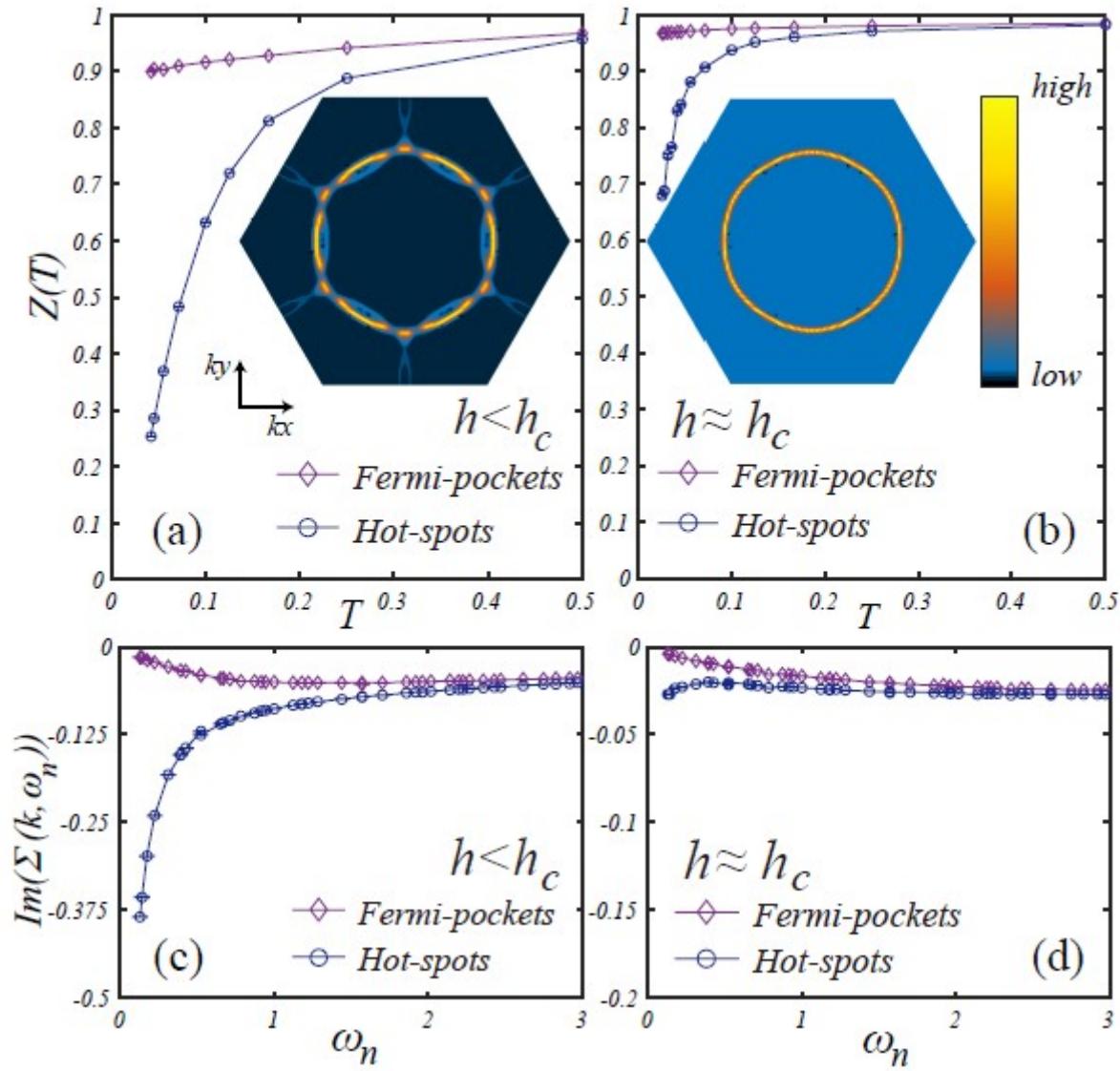
(b)



Triangle lattice

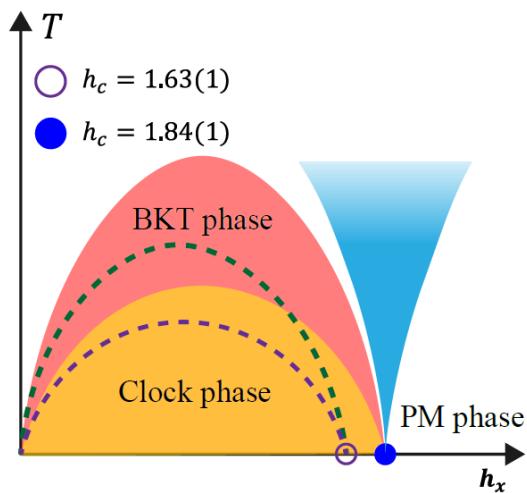


$L=30$, $\beta=30$
($30 \times 30 \times 600$)

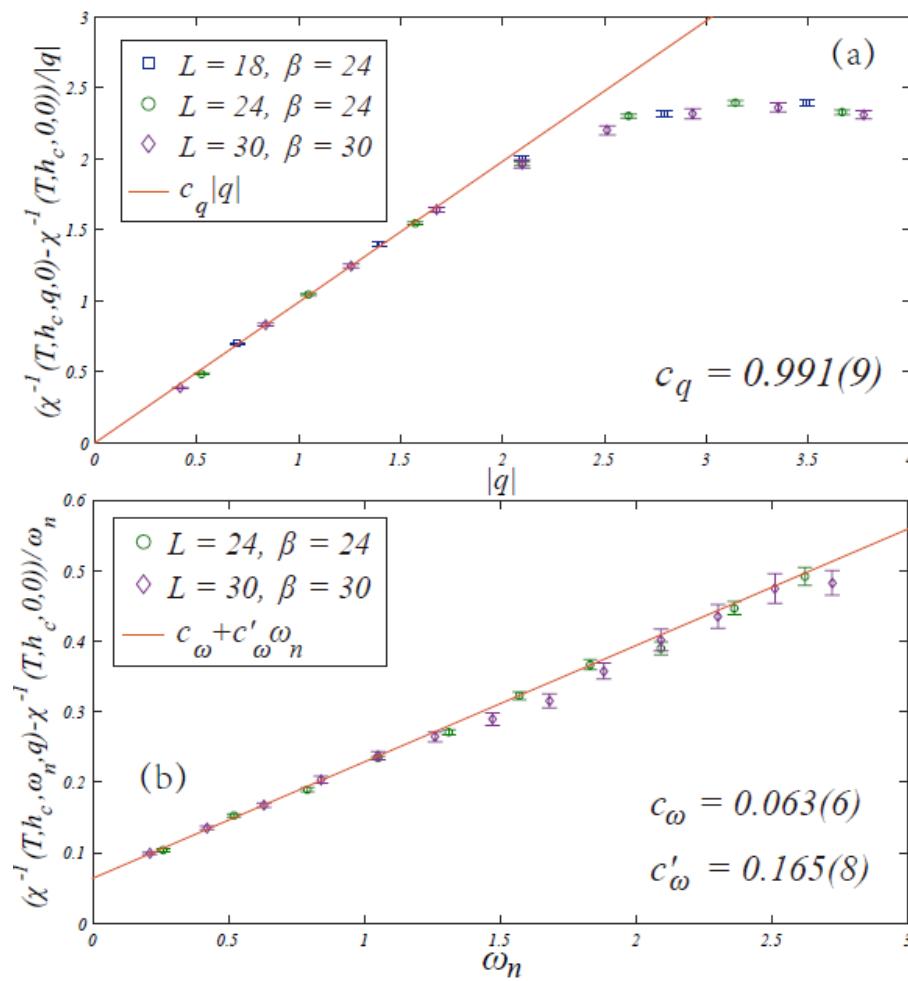


AFM-QCP

$$\chi(T, h, \mathbf{q}, \omega_n) = \frac{1}{(c_t T + c'_t T^2) + c_h |h - h_c|^\gamma + c_q |\mathbf{q}|^2 + (c_\omega \omega + c'_\omega \omega^2)}$$

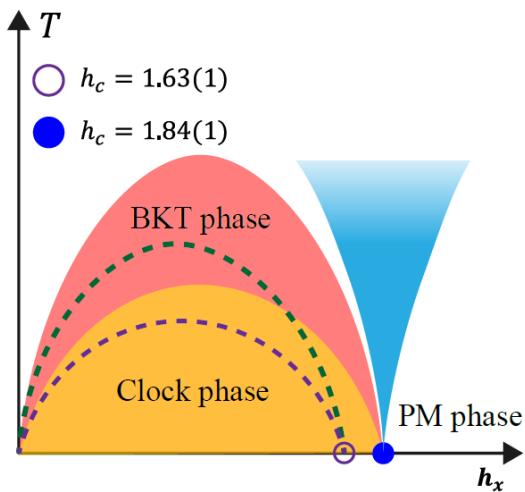


$L=30, \beta=30$
 $(30 \times 30 \times 600)$

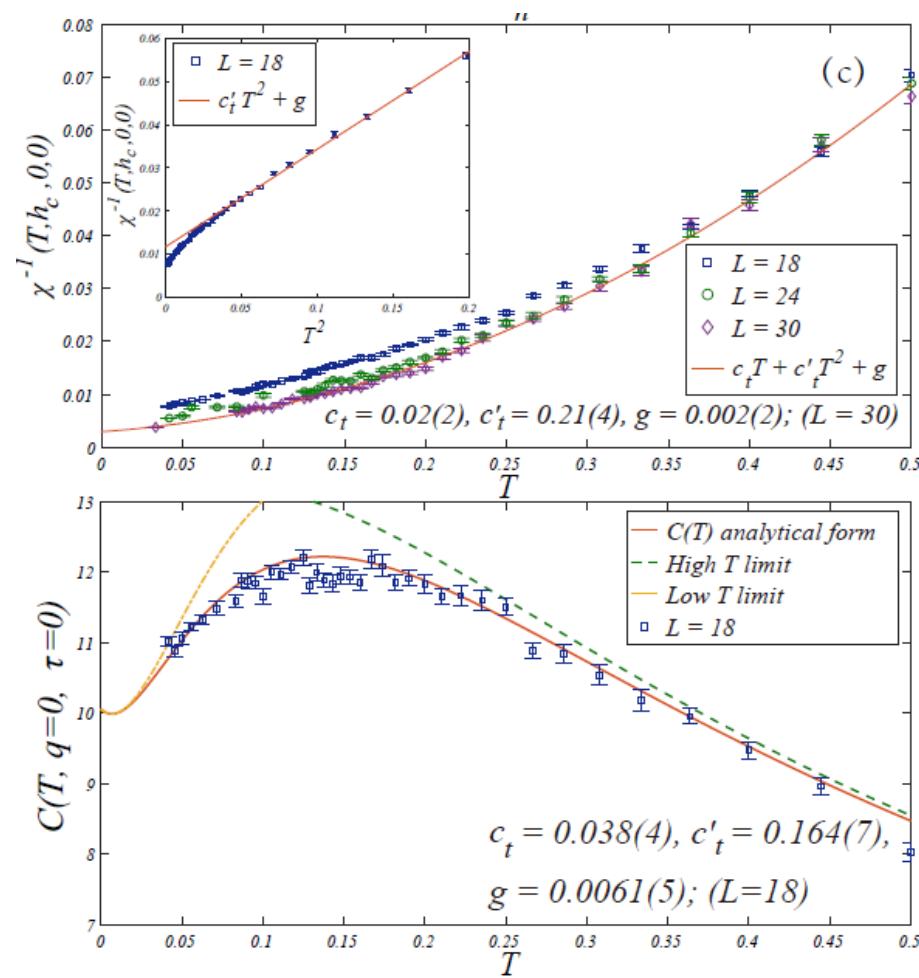


AFM-QCP

$$\chi(T, h, \mathbf{q}, \omega_n) = \frac{1}{(c_t T + c'_t T^2) + c_h |h - h_c|^\gamma + c_q |\mathbf{q}|^2 + (c_\omega \omega + c'_\omega \omega^2)}$$



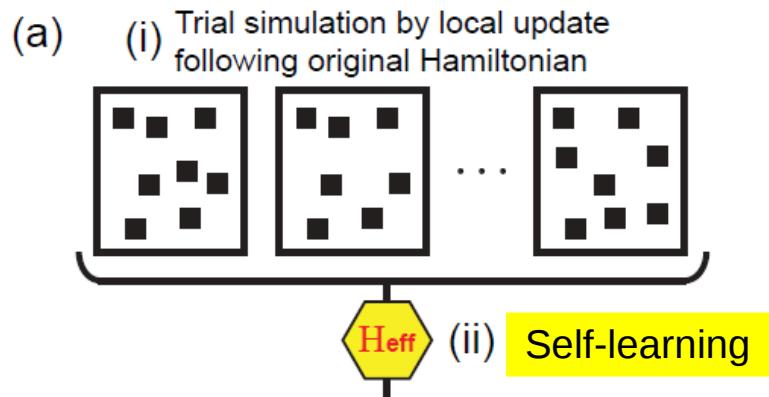
$L=30$, $\beta=30$
 $(30 \times 30 \times 600)$



A Trio of Self-learning Monte Carlo Method

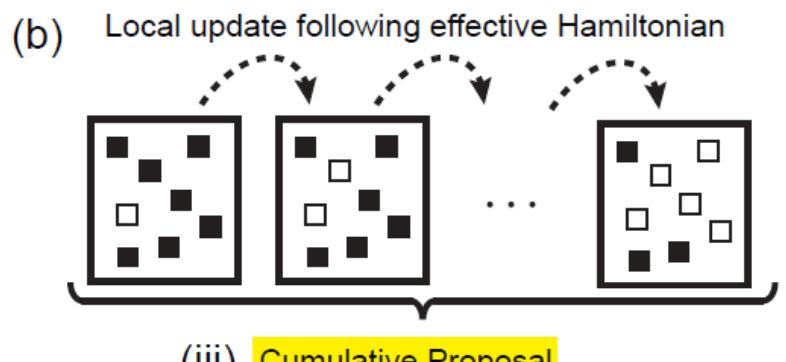
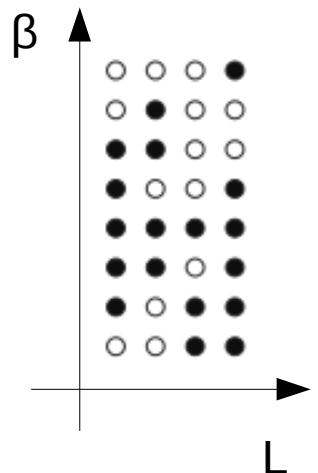
- Xiao Yan Xu, Zi Hong Liu, IOP, CAS
 - Huitao Shen, Jiuwei Liu, Liang Fu, MIT
 - Yang Qi, Fudan University
-
- Self-Learning Monte Carlo Method,
PRB 95. 041101(R) (2017)
 - Self-learning Monte Carlo method and cumulative update in fermion systems,
PRB 95, 241104(R) (2017)
 - Self-learning quantum Monte Carlo method in interacting fermion systems,
PRB 96, 041119(R) (2017)

Self-learning Monte Carlo (SLMC)



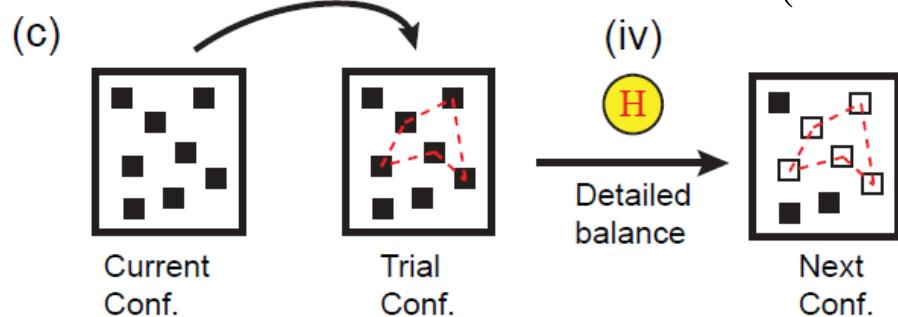
$$Z = \sum_{\{C\}} \phi(C) \det(\mathbf{1} + \mathbf{B}(\beta, 0; C))$$

$$-\beta H^{\text{eff}}[C] = \ln(\omega[C])$$

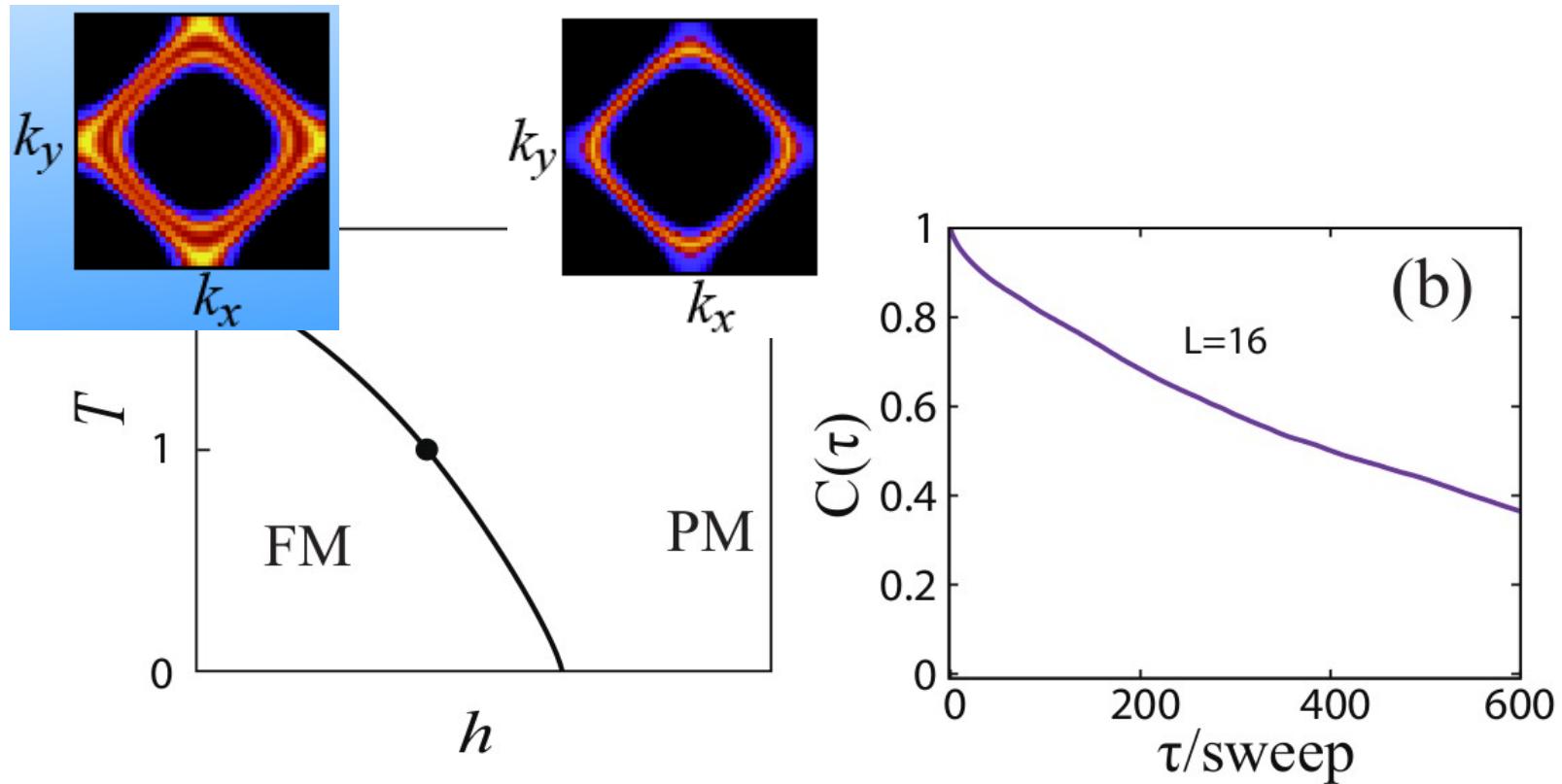


$$H^{\text{eff}} = E_0 + \sum_{(i\tau);(j,\tau')} J_{i,\tau;j\tau'} s_{i,\tau} s_{j,\tau'} + \dots$$

$$A(\mathcal{C} \rightarrow \mathcal{C}') = \min\{1, e^{-\beta((H(\mathcal{C}') - H^{\text{eff}}(\mathcal{C}')) - (H(\mathcal{C}) - H^{\text{eff}}(\mathcal{C}))}\}$$



Testing case



Complexity for getting an independent configuration: $\beta N^3 \tau_L$

Testing case

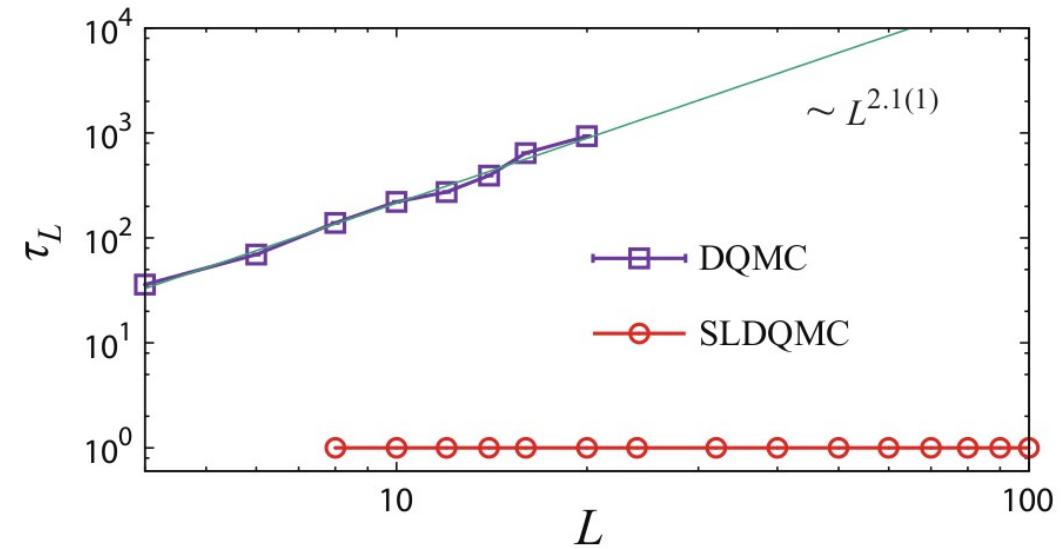
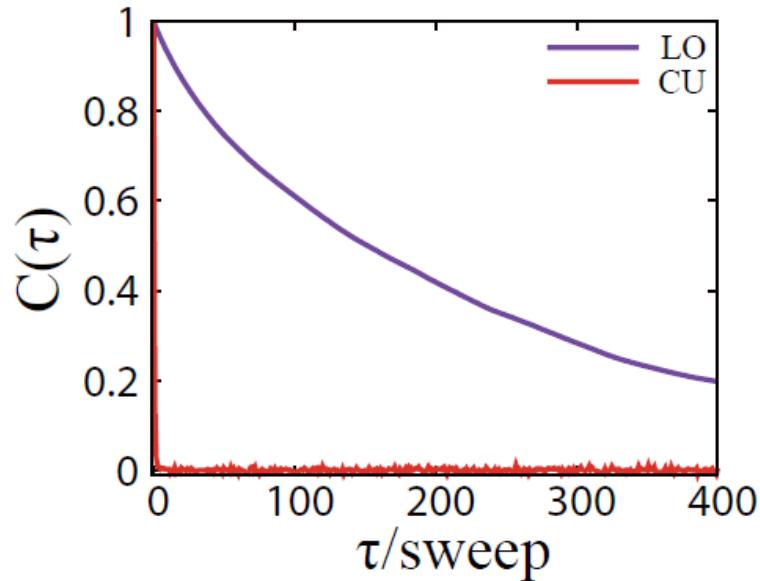
Complexity for getting an independent configuration: $\beta N^3 \tau_L$

Complexity for SLMC

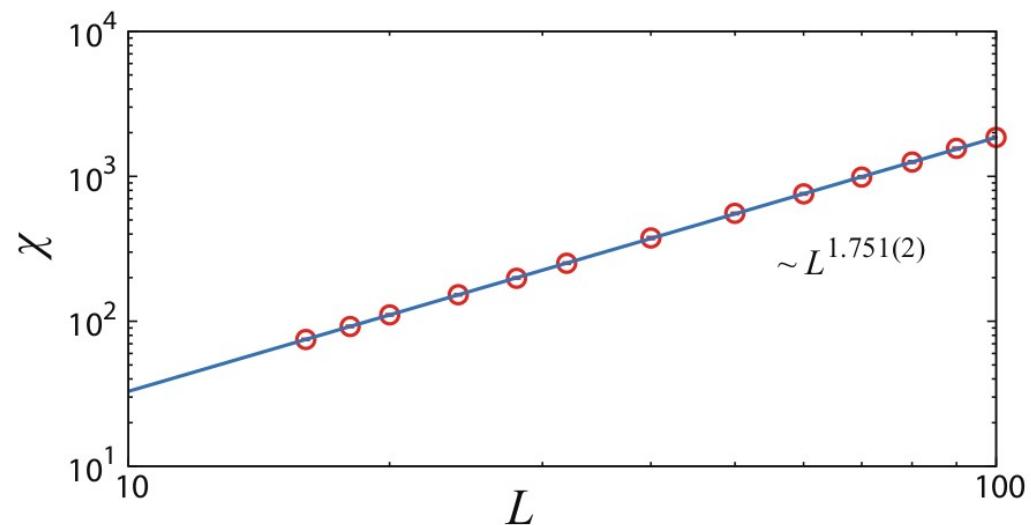
- Cumulative update: $\gamma \beta N \tau_L$
- Detail balance: N^3
$$\begin{aligned}\omega_C &= \phi(\mathcal{C}) \det (\mathbf{1} + \mathbf{B}(\beta, \tau) \mathbf{B}(\tau, 0)) \\ &= \phi(\mathcal{C}) \det (\mathbf{G}(0, 0))^{-1}\end{aligned}$$
- Sweep Green's function: βN^2

Complexity speed up $\mathcal{S} = \min\left(\frac{N^2}{\gamma}, \beta \tau_L, N \tau_L\right)$

Testing case



$$\chi(L) = \frac{1}{L^2} \sum_{ij} \int_0^\beta d\tau \langle s_{i,\tau}^z s_{j,0}^z \rangle$$



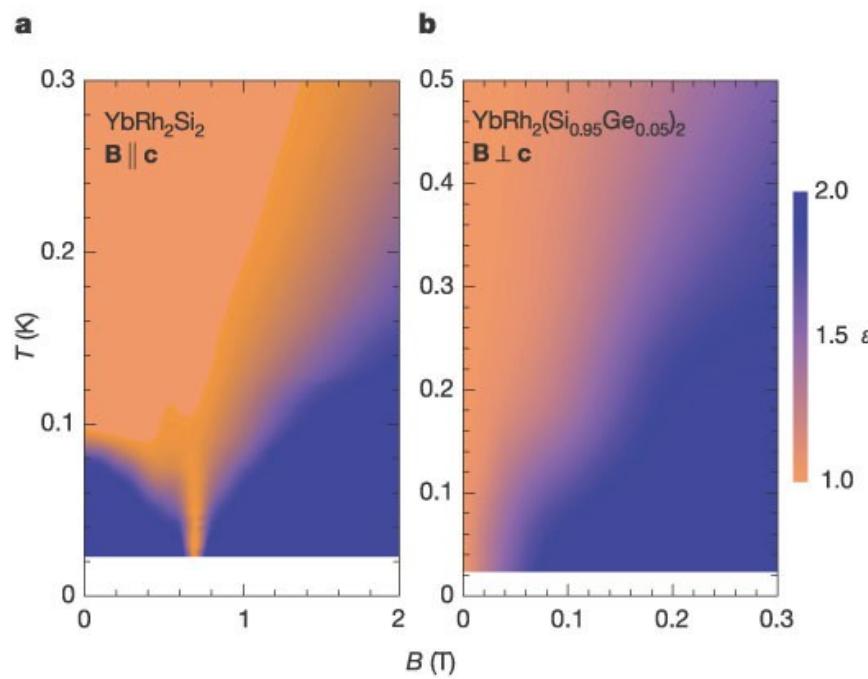
Outlook

■ Possible directions:

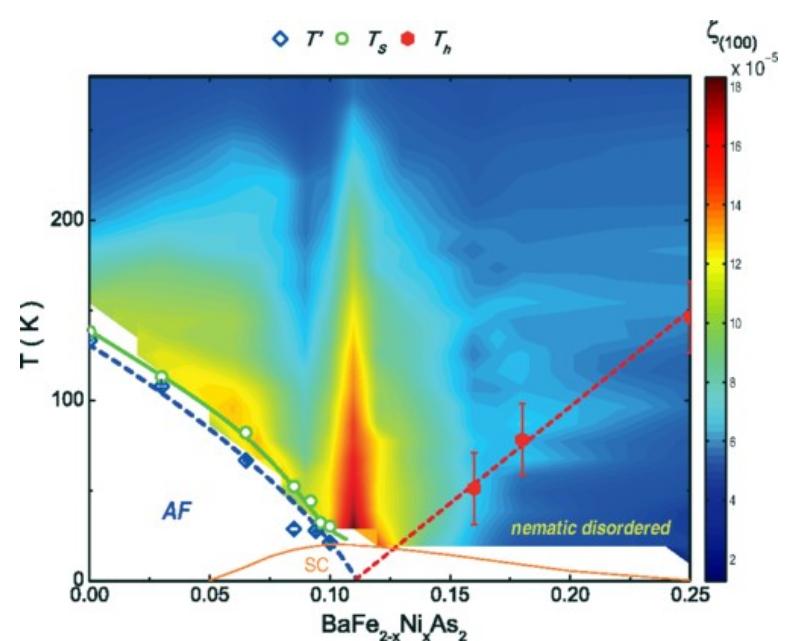
Other lattices, FM / AFM / Nematic fluctuations of itinerant electron systems

Non-Fermi liquid, fluctuation induced superconductivity

Beyond Hertz-Millis-Moriya mean-field theory, itinerant QCP



Nature 424, 524-527 (2003)



Phys. Rev. Lett. 117, 157002 (2016)