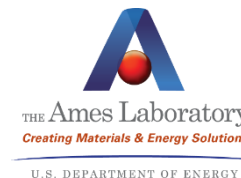


Charge Aspects of Composite Pair Superconductivity

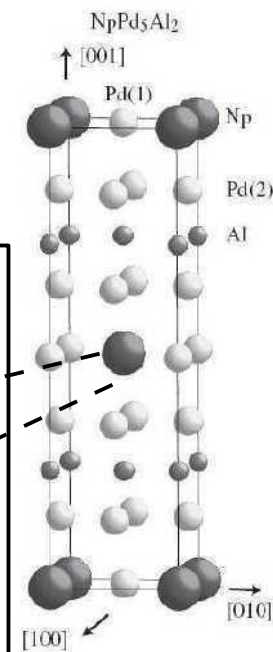
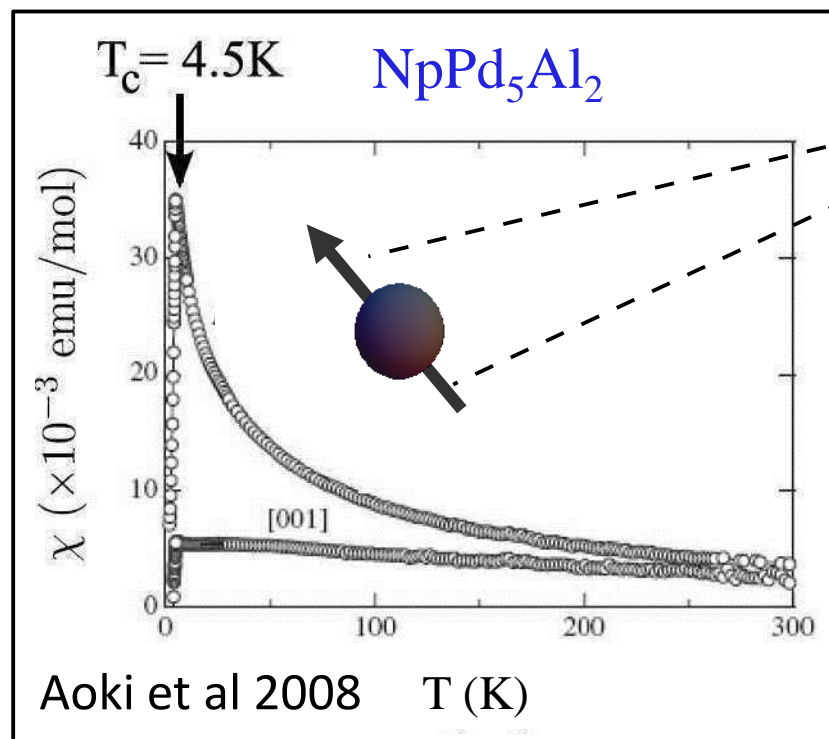
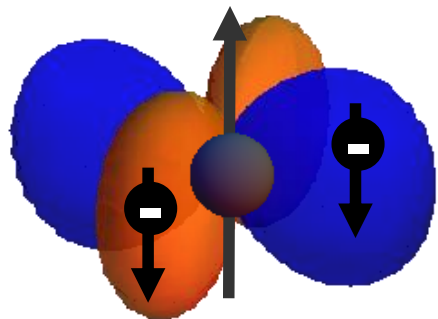
Rebecca Flint (Iowa State)



Piers Coleman (Rutgers)

Maxim Dzero (Kent State)

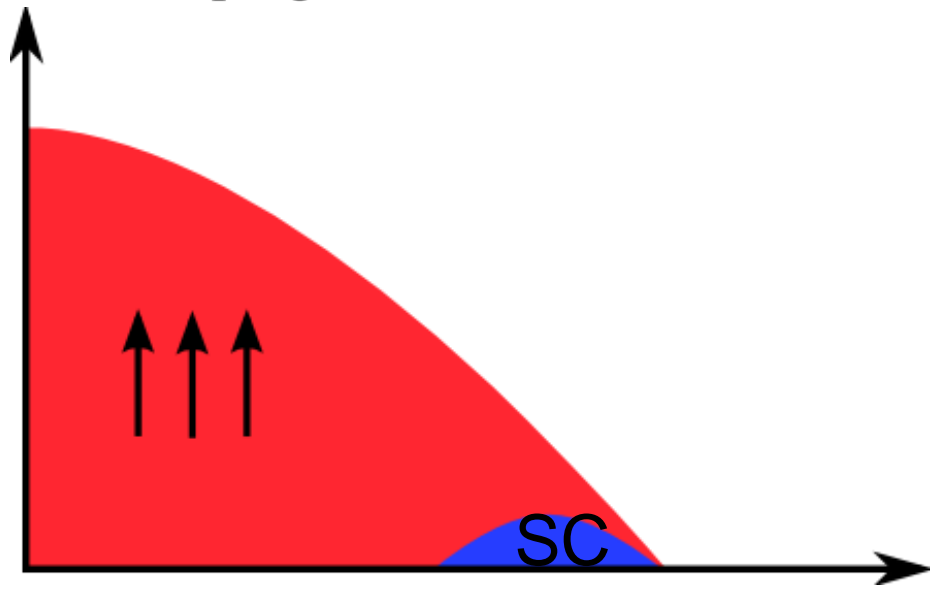
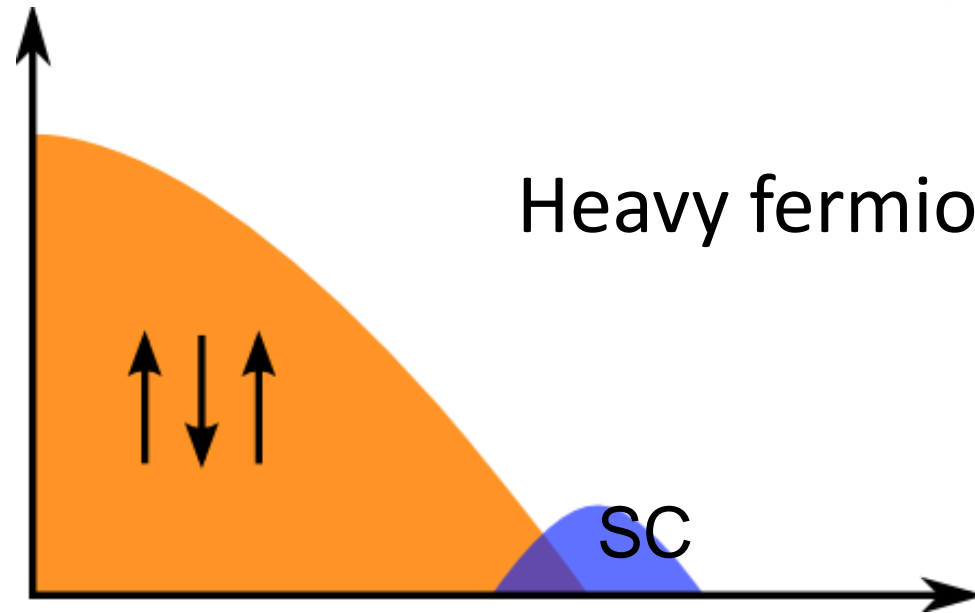
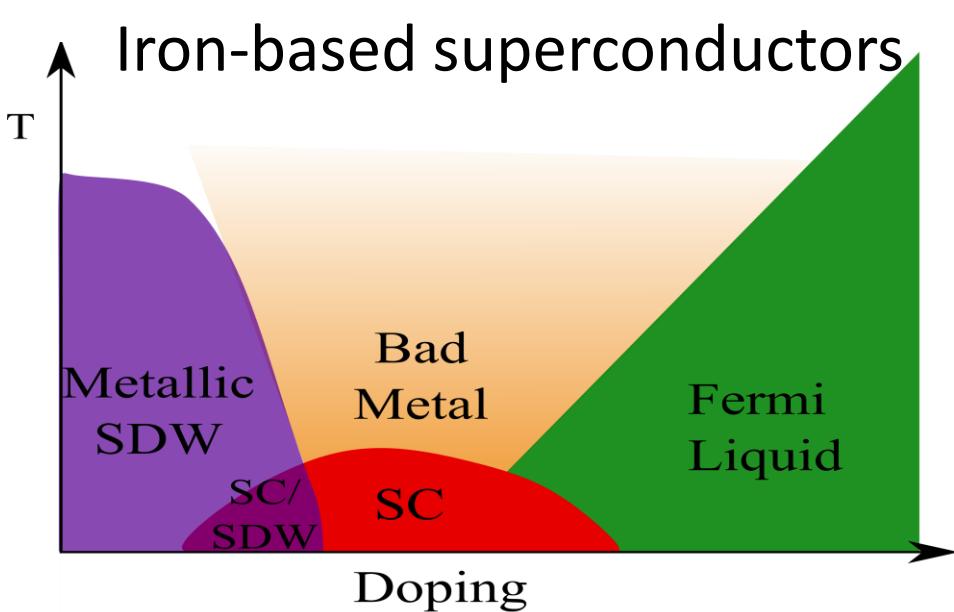
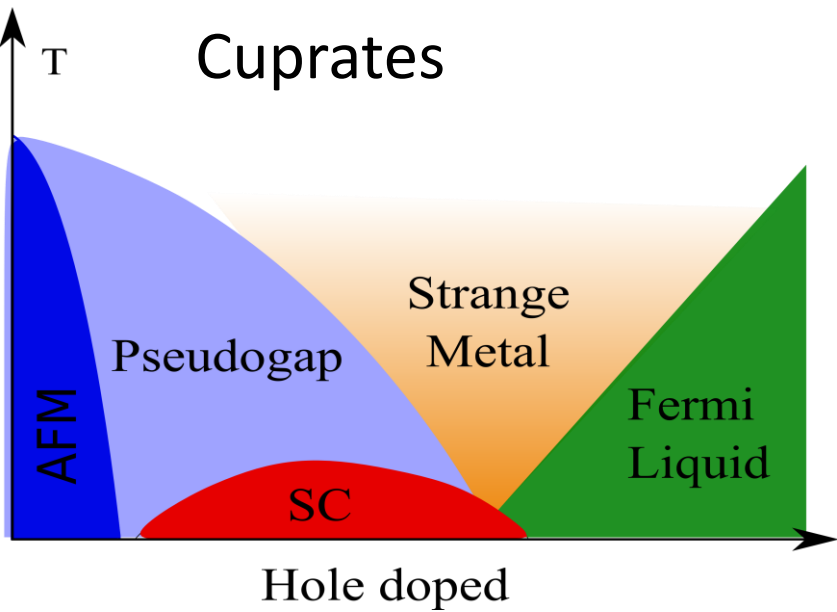
Andriy Nevidomskyy (Rice)



R. Flint, M. Dzero and P. Coleman, Nature Physics 4, 643(2008)

R. Flint, A.H. Nevidomskyy and P. Coleman, Phys. Rev. B 84, 064514 (2011)

Magnetic pairing appears ubiquitous

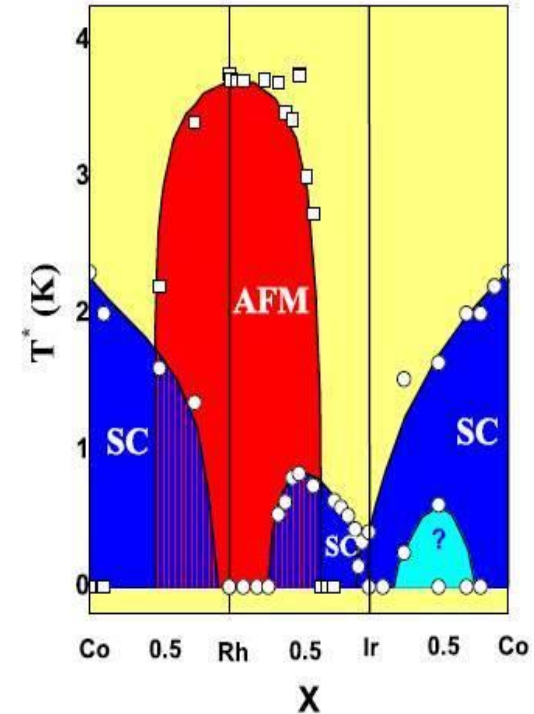


Magnetic pairing appears ubiquitous

Sarrao and Thompson, JPSJ (2007)

But...

- Two domes in CeMIn_5 ($M = \text{Co}, \text{Rh}, \text{Ir}$)
- Superconductivity without magnetism
 - NpPd_5Al_2 , PuCoGa_5
- Robust to disorder *on the f-site*
 - 5% Sn on In kills T_c , 25% La on Ce required
- Nodeless superconductivity in $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$
- Many Ce superconductors, few Yb ($T_c^{\text{Ce}} \gg T_c^{\text{Yb}}$)
 - Magnetism should not discriminate



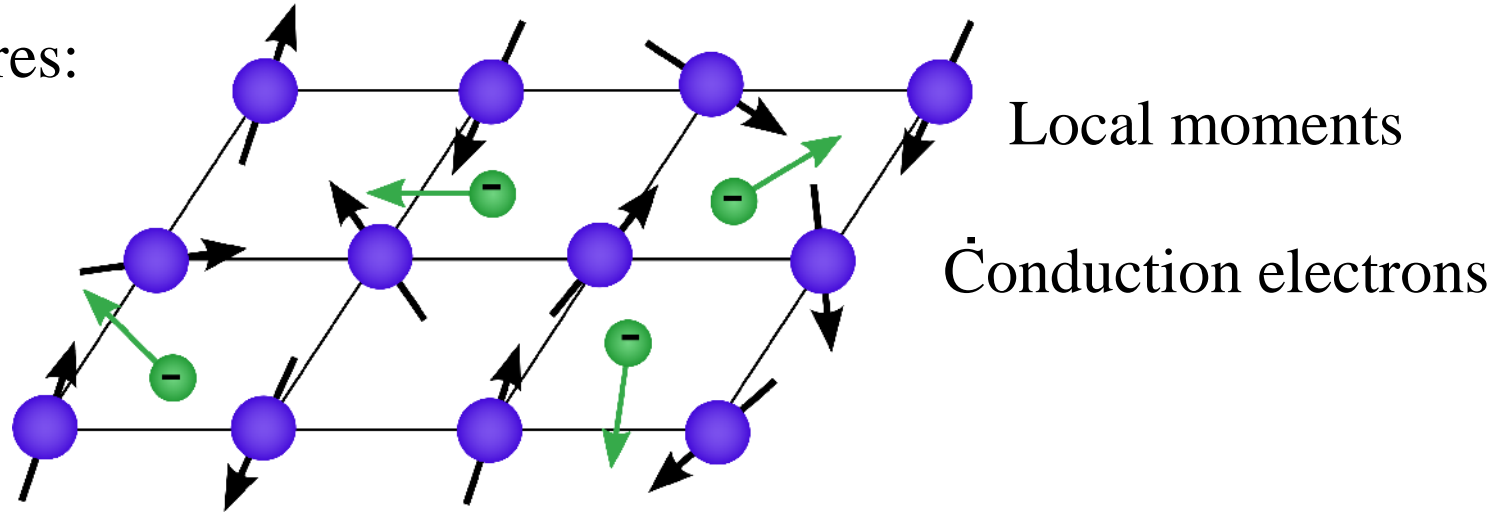
Are there other possible mechanisms?

Yes! Composite pairing

How can we tell?

“Conventional” heavy fermion superconductivity

At high temperatures:



How do we get from here to heavy Cooper pairs?

Beal-Monod, Bourbonnais and Emery (1986)

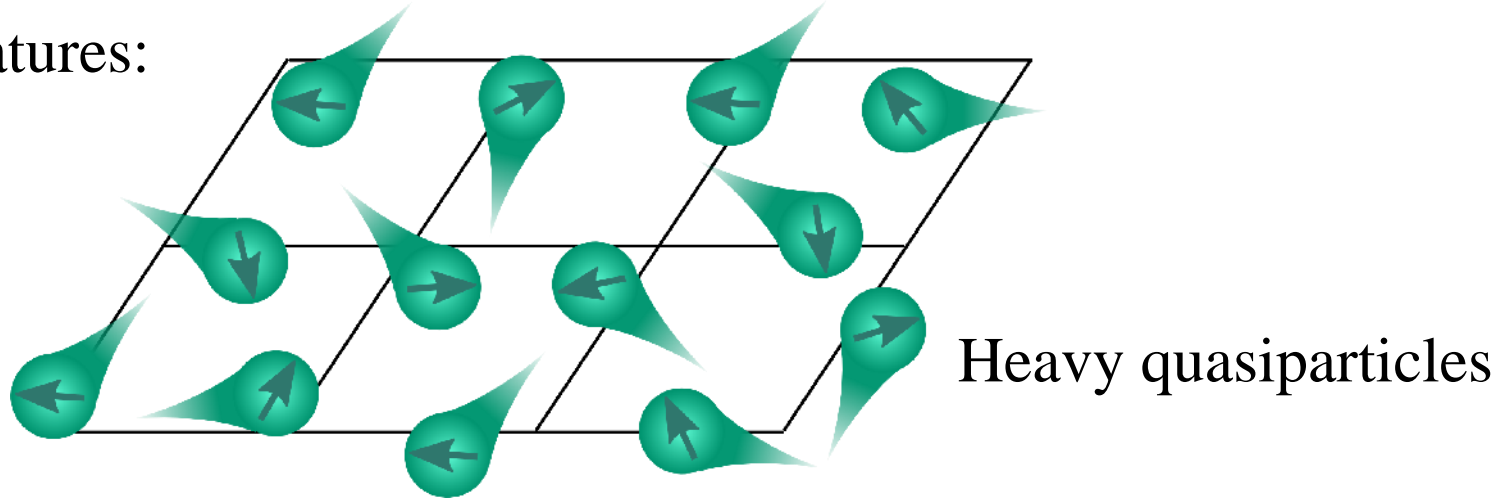
Scalapino, Loh and Hirsch (1986)

Miyake, Schmitt-Rink and Varma (1986)

“Conventional” heavy fermion superconductivity

At lower temperatures:

$$T < T^*$$



How do we get from here to heavy Cooper pairs?

1. The local moments quench [via the Kondo effect], forming heavy quasiparticles

Beal-Monod, Bourbonnais and Emery (1986)

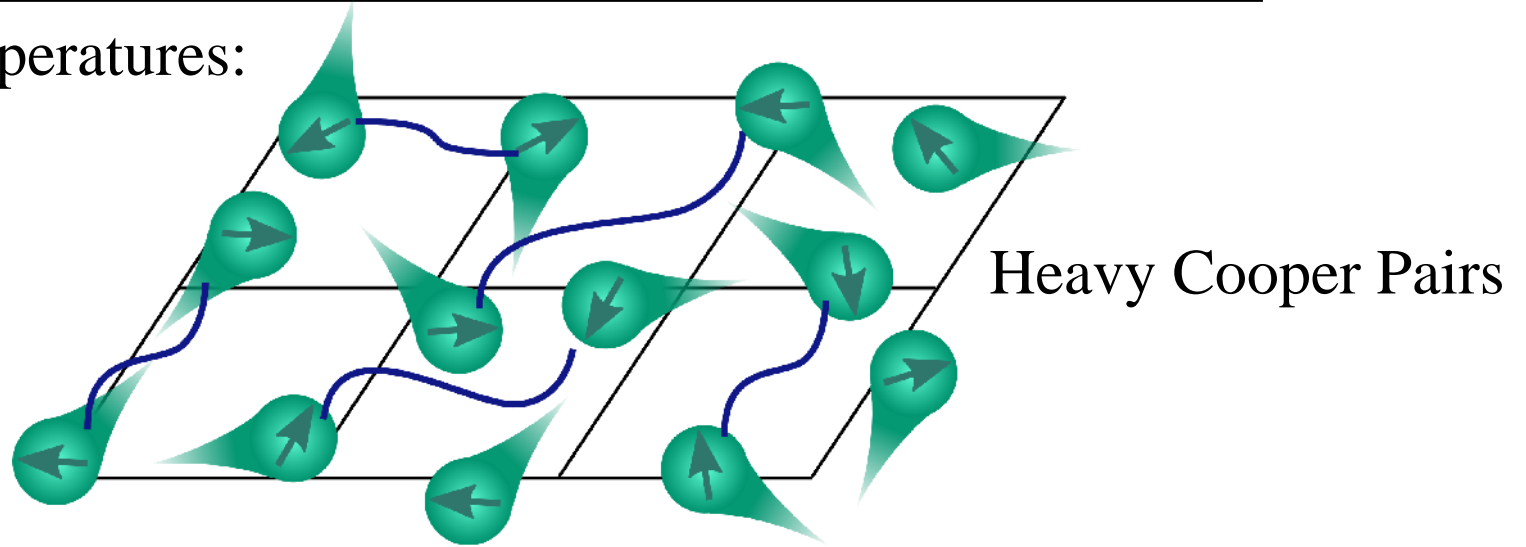
Scalapino, Loh and Hirsch (1986)

Miyake, Schmitt-Rink and Varma (1986)

“Conventional” heavy fermion superconductivity

At very low temperatures:

$$T < T_C$$



How do we get from here to heavy Cooper pairs?

1. The local moments quench [via the Kondo effect], forming heavy quasiparticles
2. The heavy quasiparticles pair [via residual spin fluctuations]

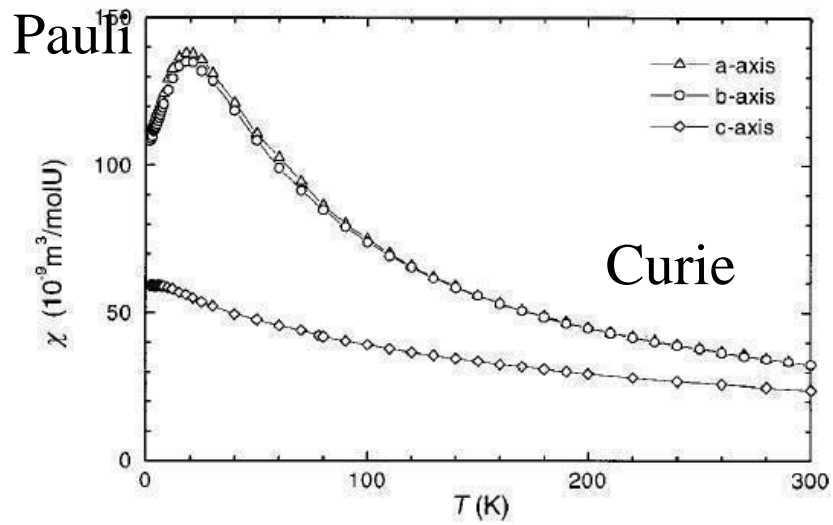
These two stages are well separated.

Beal-Monod, Bourbonnais and Emery (1986)

Scalapino, Loh and Hirsch (1986)

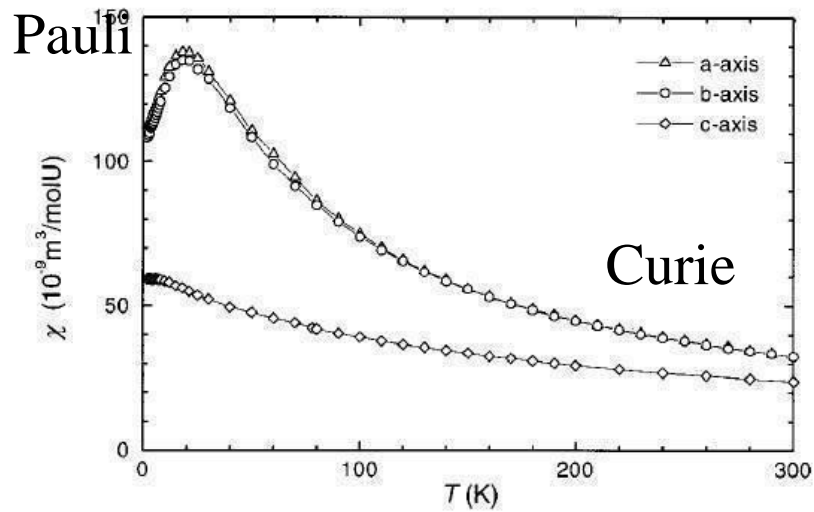
Miyake, Schmitt-Rink and Varma (1986)

UPt3



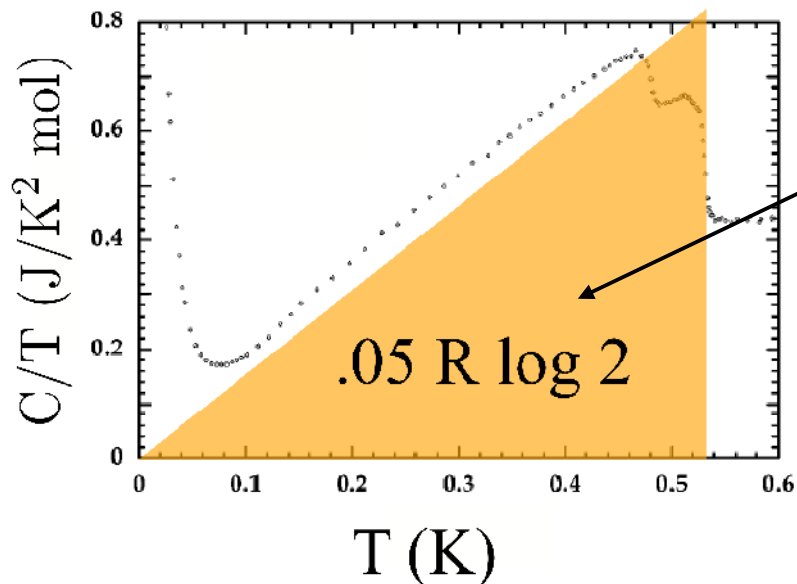
Pauli paramagnetic by 30K

UPt3



Pauli paramagnetic by 30K

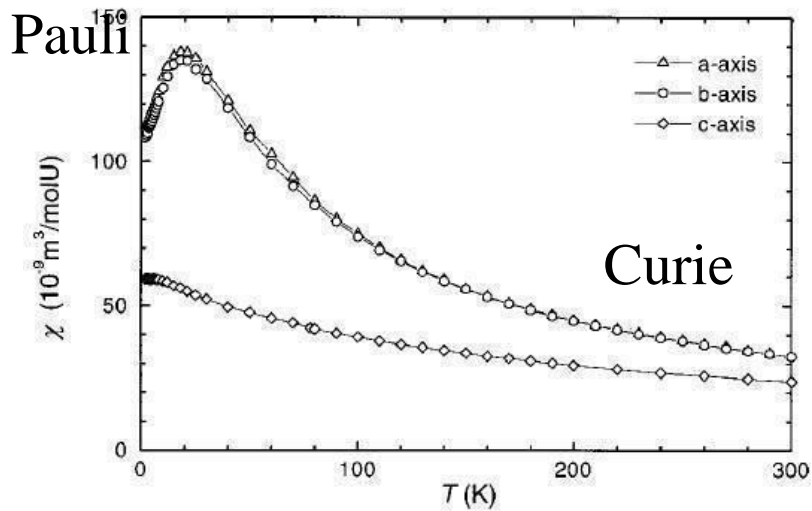
$T_c = 0.5\text{K}$



$$S = \int_0^T \frac{C(T')}{T'} dT'$$

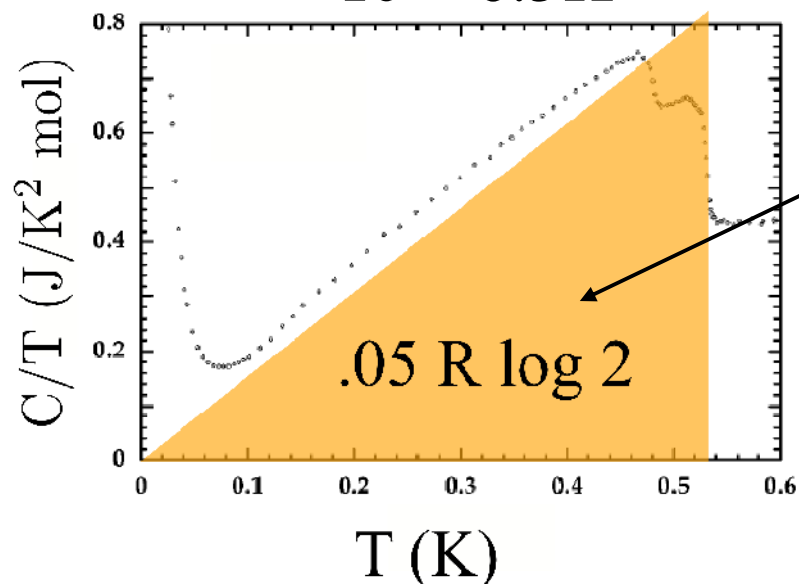
Frings *et al.* J. Magn. Magn. Mater. **31**, 240(1983)
 Brison *et al.* J. Low Temp. Phys. **95**, 145(1994)

UPt3



Pauli paramagnetic by 30K

$T_c = 0.5\text{K}$



$$S = \int_0^T \frac{C(T')}{T'} dT'$$

Small condensation entropy

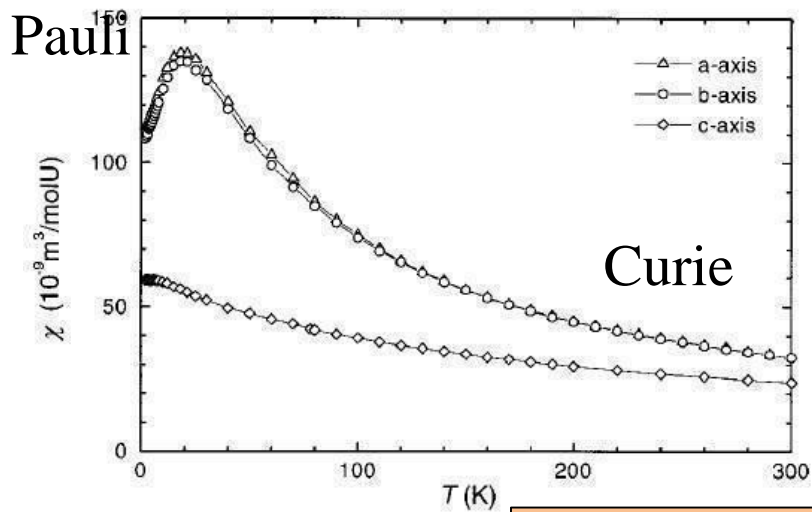
Spins are quenched

Two stages are well separated

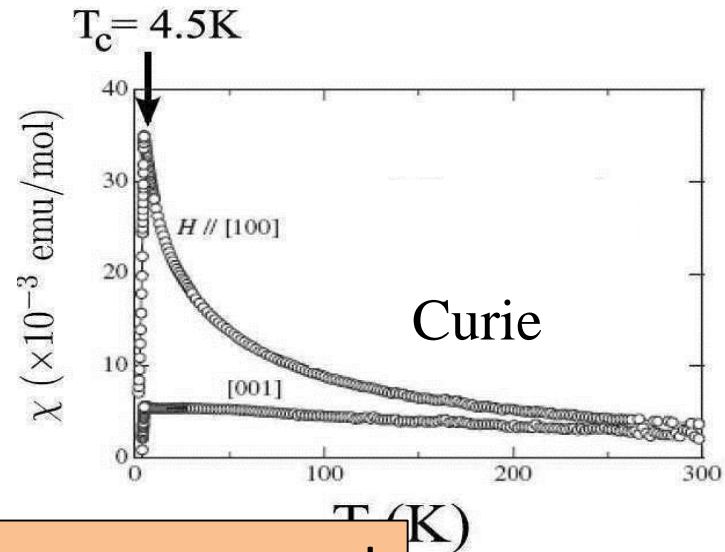
Frings *et al.* J. Magn. Magn. Mater. **31**, 240(1983)

Brison *et al.* J. Low Temp. Phys. **95**, 145(1994)

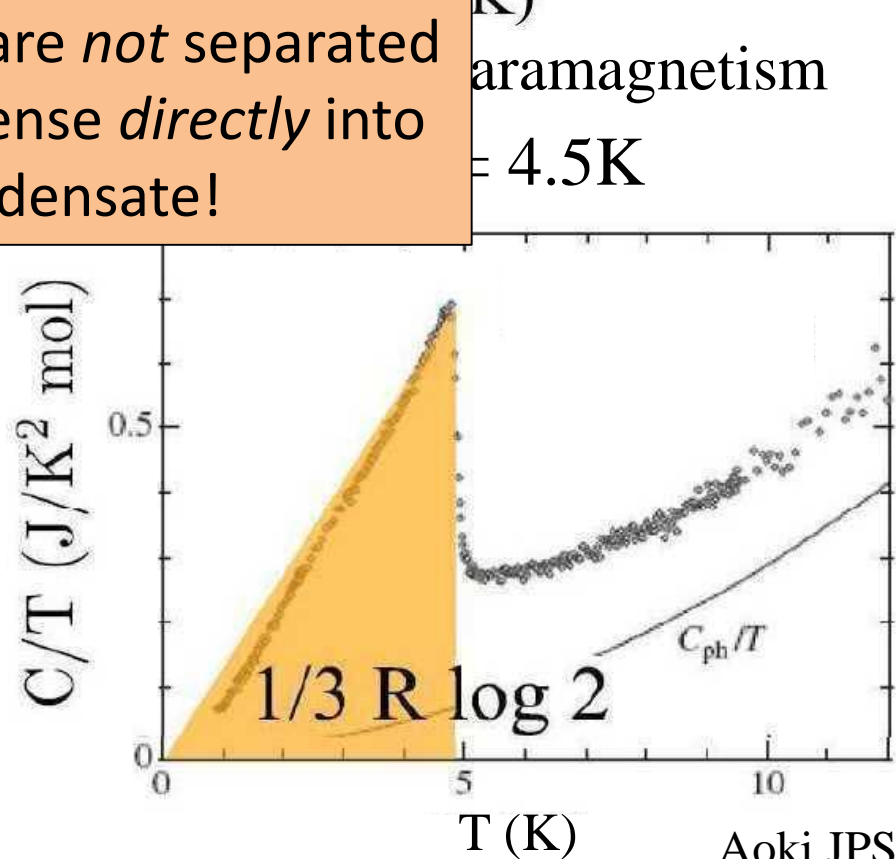
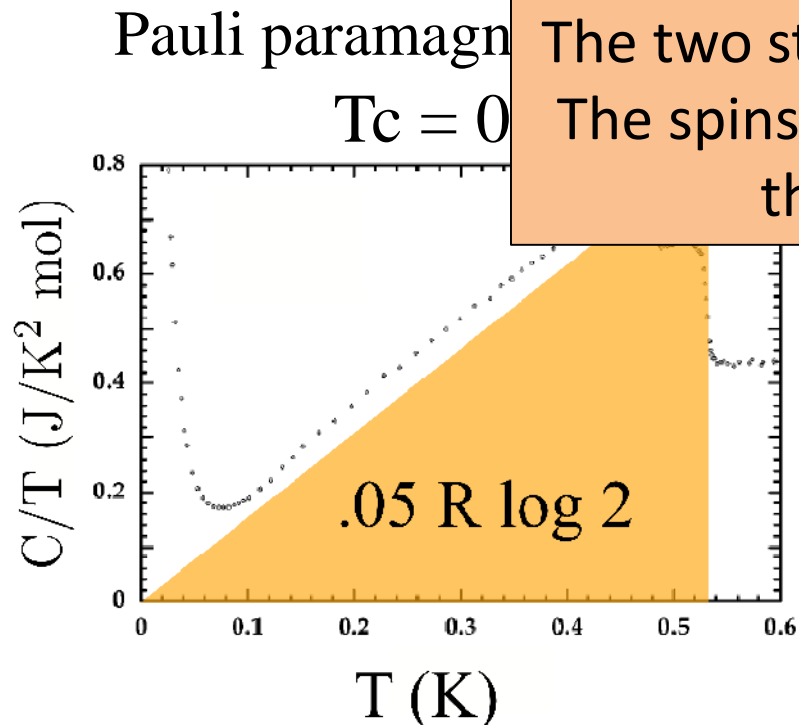
UPt₃



NpPd₅Al₂



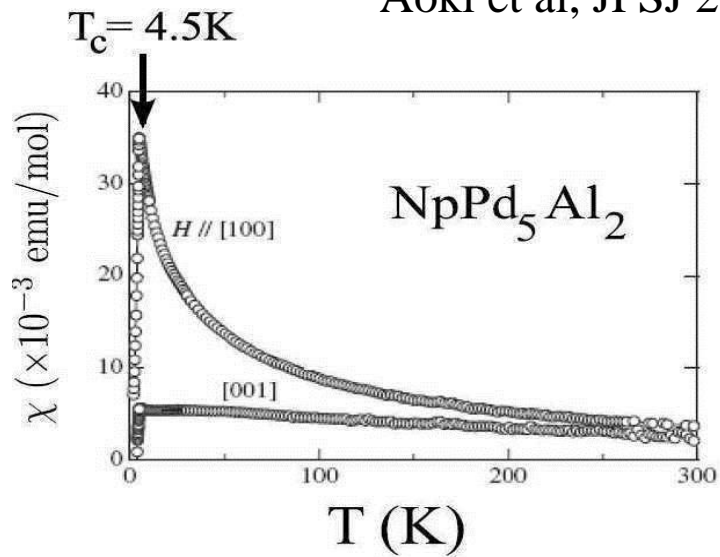
The two stages are *not* separated
 The spins condense *directly* into
 the condensate!



Frings *et al.* J. Magn. Magn. Mater. **31**, 240(1983)
 Brison *et al.* J. Low Temp. Phys. **95**, 145(1994)

Local moment superconductivity

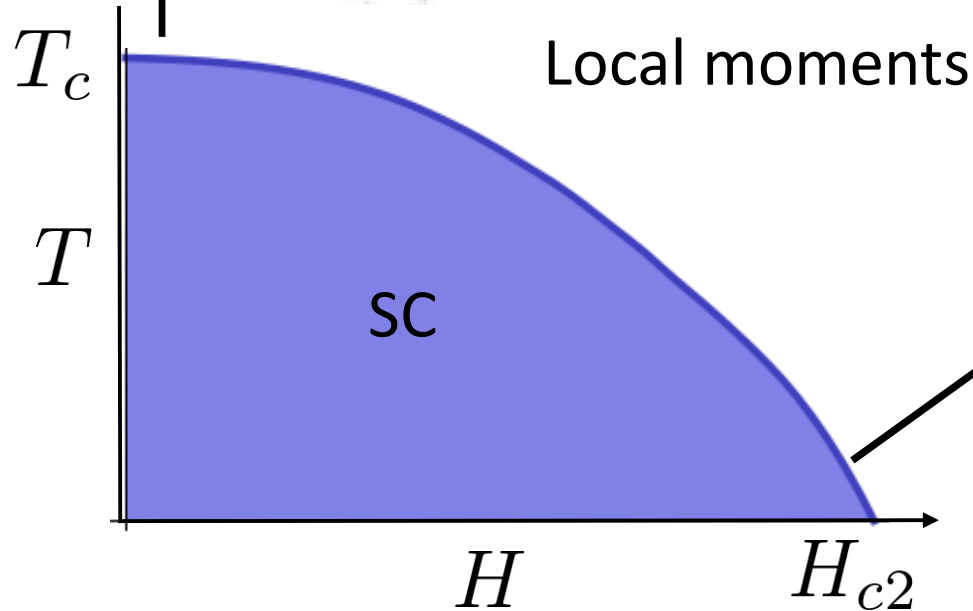
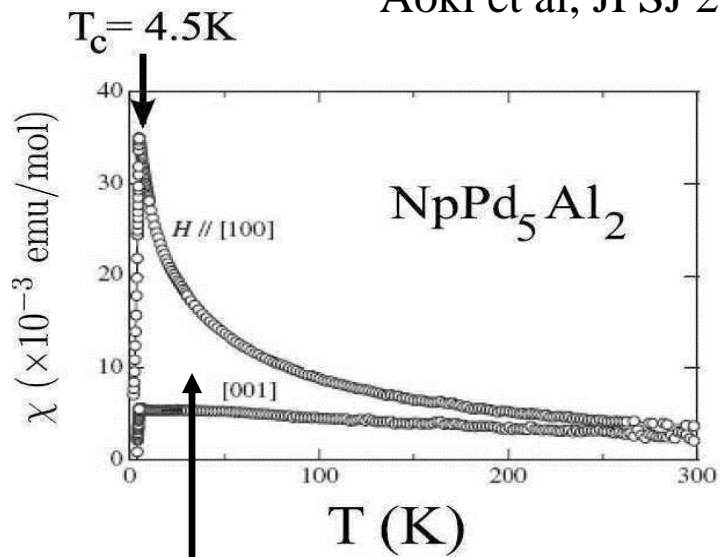
Aoki et al, JPSJ 2007



- Large condensation entropy ($\sim 1/3 R \log 2$)
 - Free local moments above T_c

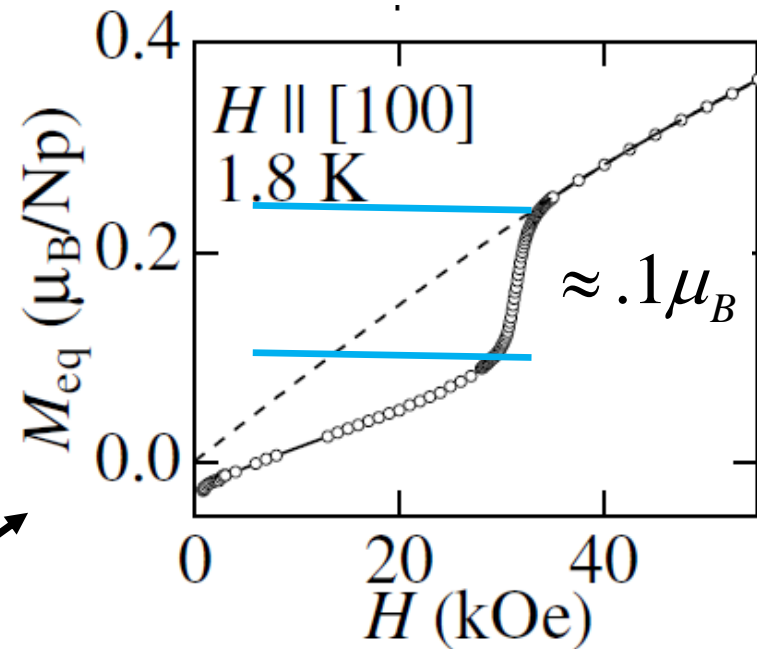
Local moment superconductivity

Aoki et al, JPSJ 2007



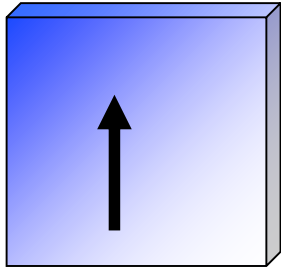
- Large condensation entropy ($\sim 1/3 R \log 2$)
 - Free local moments above T_c, H_{c2}
- Spins quench directly into the condensate!

How?



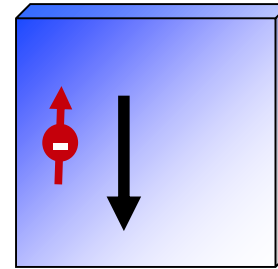
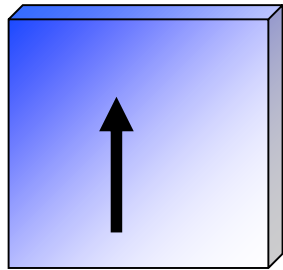
Composite Fermions (Kondo effect)

Co-tunneling



Composite Fermions (Kondo effect)

Co-tunneling



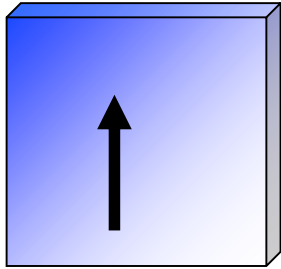
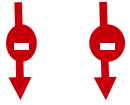
$$c_{\uparrow}^{\dagger} S_{-}$$

Composite fermion

The product of a conduction electron and a spin flip

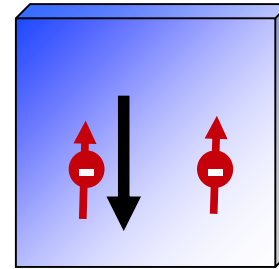
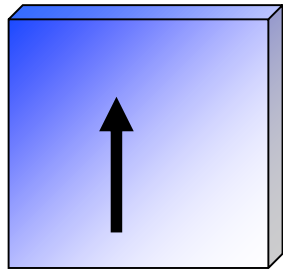
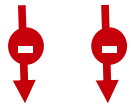
Composite Pairing

Copairing



Composite Pairing

Copairing



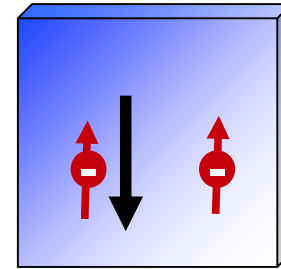
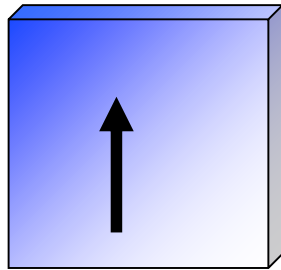
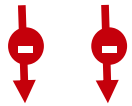
$$c_{A\uparrow}^\dagger c_{B\uparrow}^\dagger S_-$$

Composite **pair**

The product of a conduction electron **pair** and a spin flip

Composite Pairing (Two channel Kondo)

Copairing



$$c_{A\uparrow}^\dagger c_{B\uparrow}^\dagger S_-$$

Composite pair

The product of a conduction electron pair and a spin flip

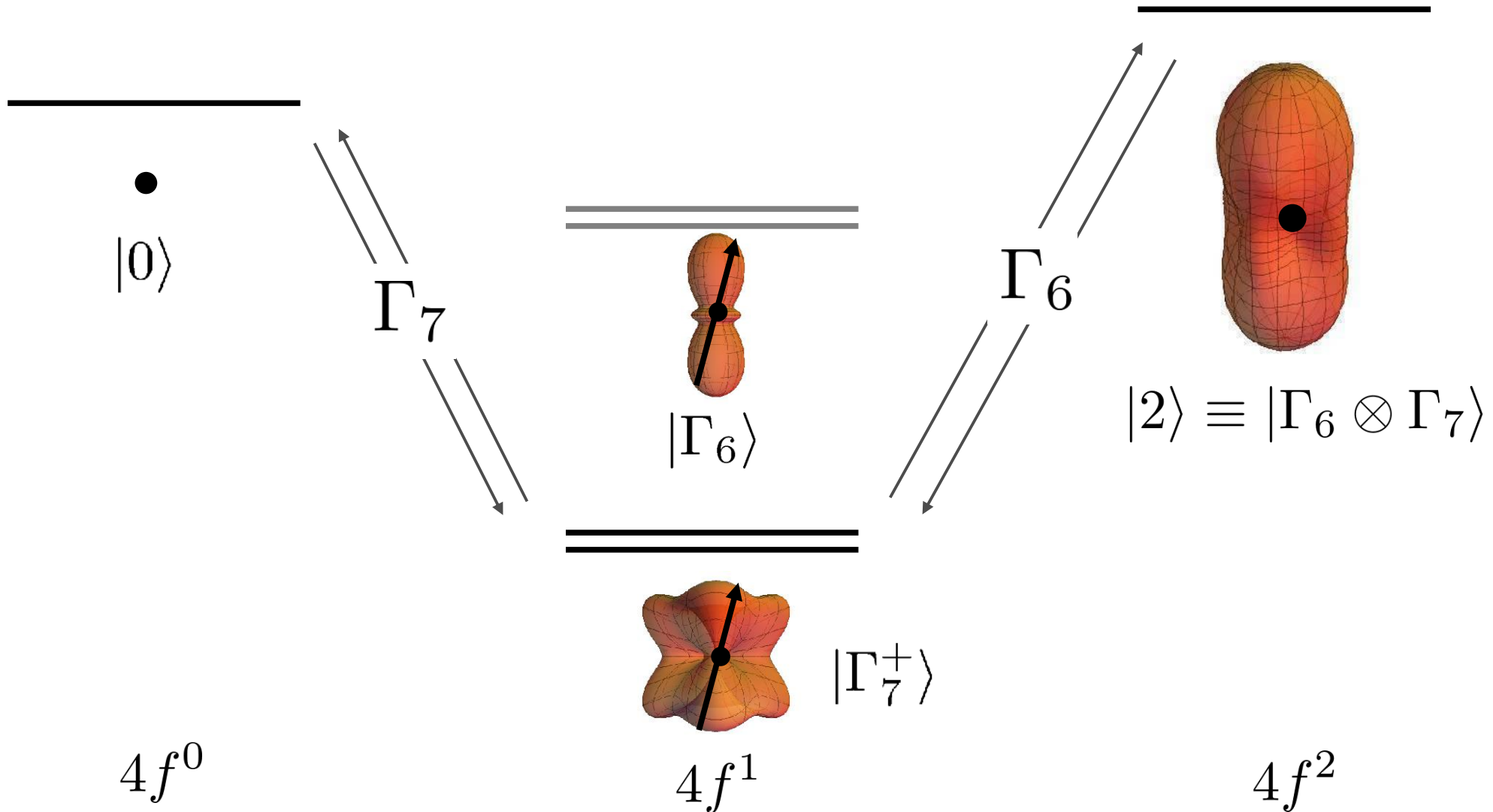
The composite pair is a **singlet**

– requiring a **triplet** pair of conduction electrons

→ Antisymmetric spatially – two electrons, two orthogonal channels

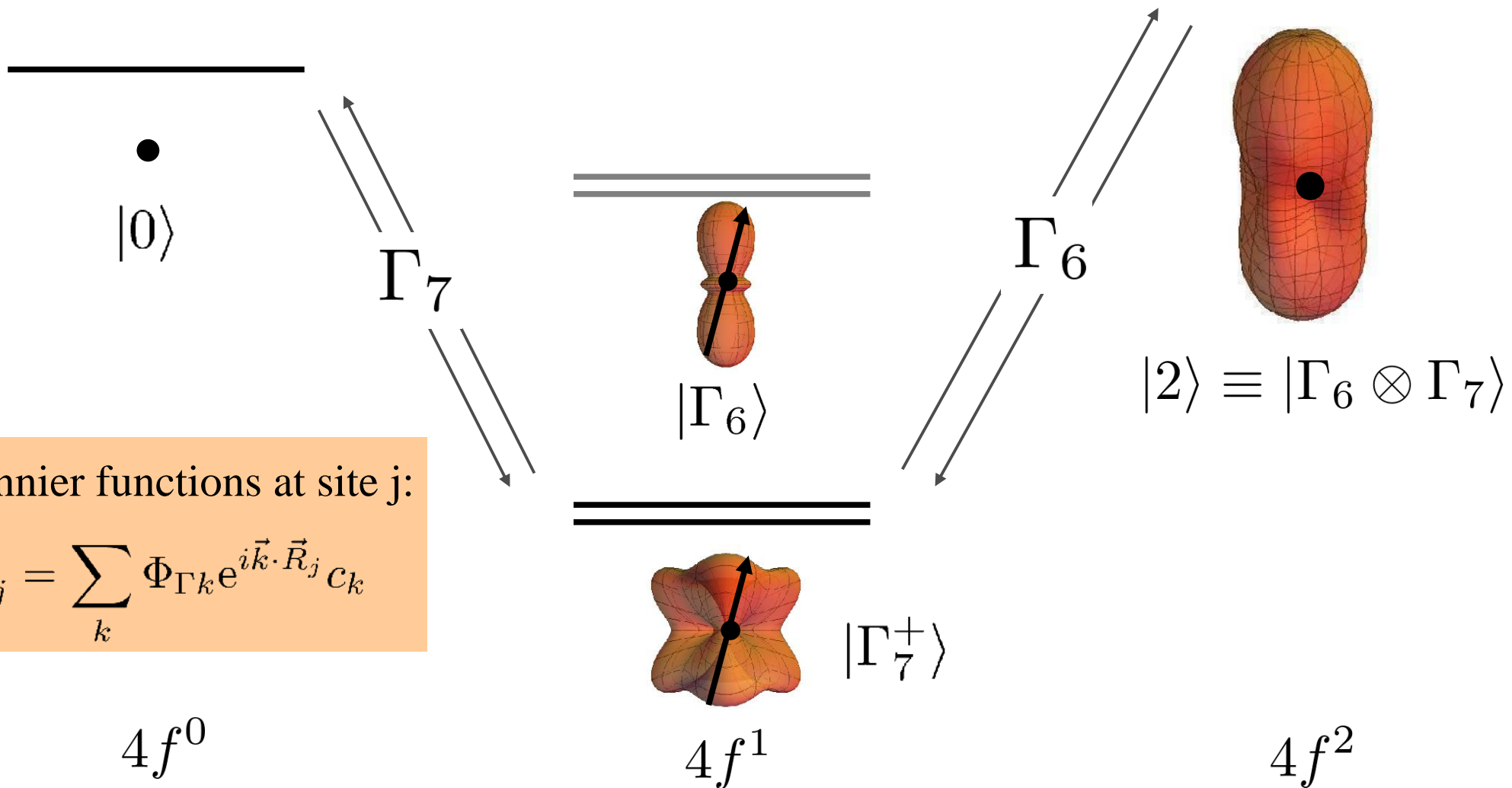
The origin of composite pairing

Requires two orthogonal Kondo channels



The origin of composite pairing

$$H = \sum_k \epsilon_k c_k^\dagger c_k + J_1 \sum_j \psi_{1j\alpha}^\dagger \vec{\sigma}_{\alpha\beta} \psi_{1k\beta} \cdot \vec{S}_j + J_2 \sum_j \psi_{2j\alpha}^\dagger \vec{\sigma}_{\alpha\beta} \psi_{2k\beta} \cdot \vec{S}_j$$



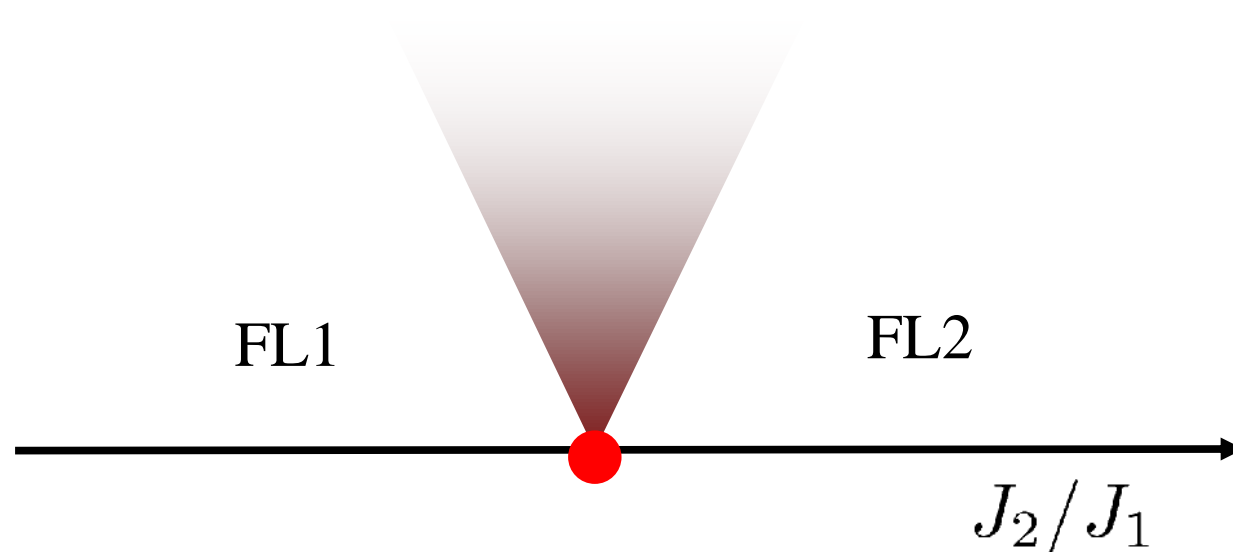
The two channel Kondo model: impurity

The impurity has a quantum critical point for $J_1 = J_2$

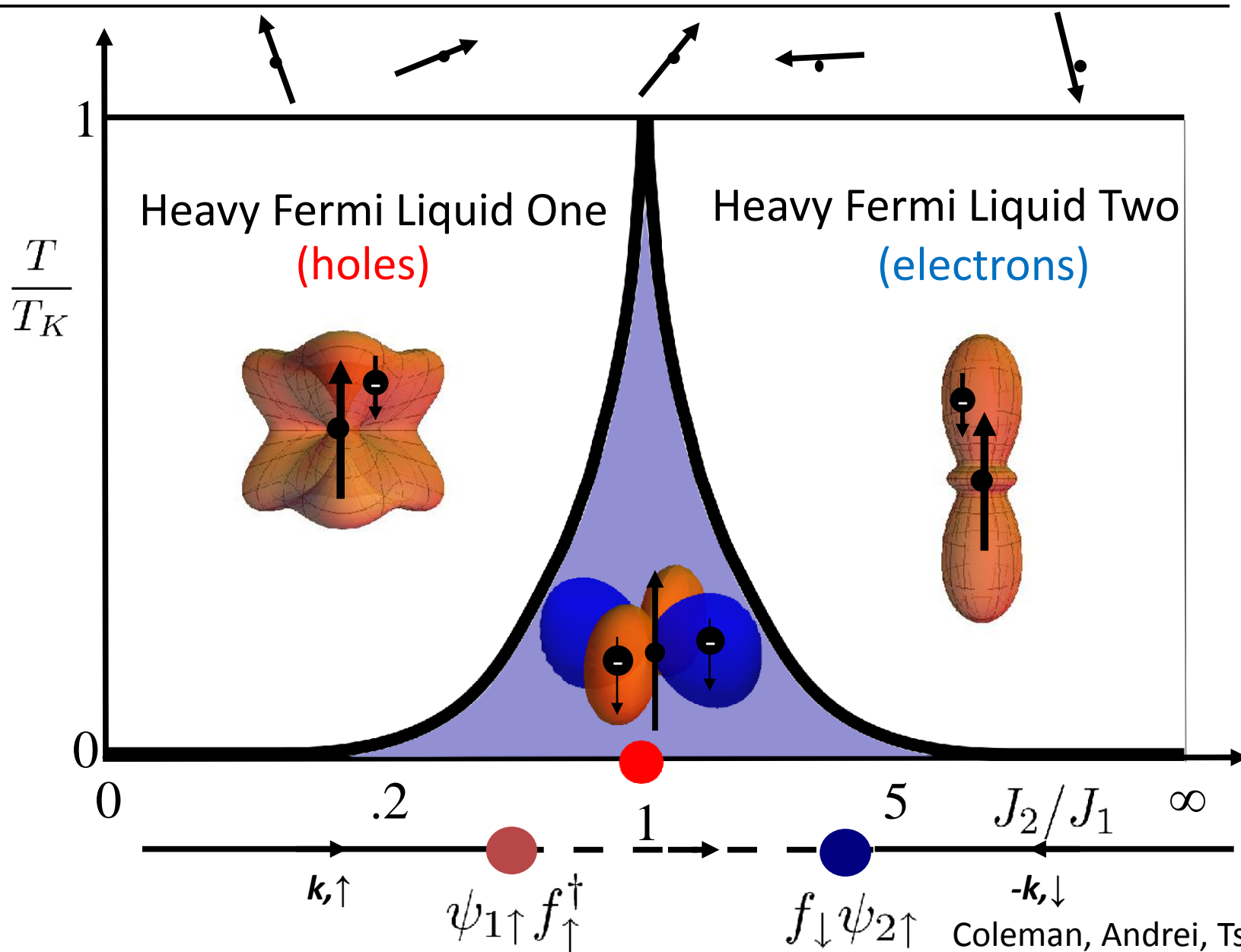
Zero point entropy $\frac{1}{2} R \log 2$

Singular composite pair fluctuations

– expect to be hidden in the lattice



The two channel Kondo model: lattice



Coleman, Andrei, Tsvetlik, Kee 1999

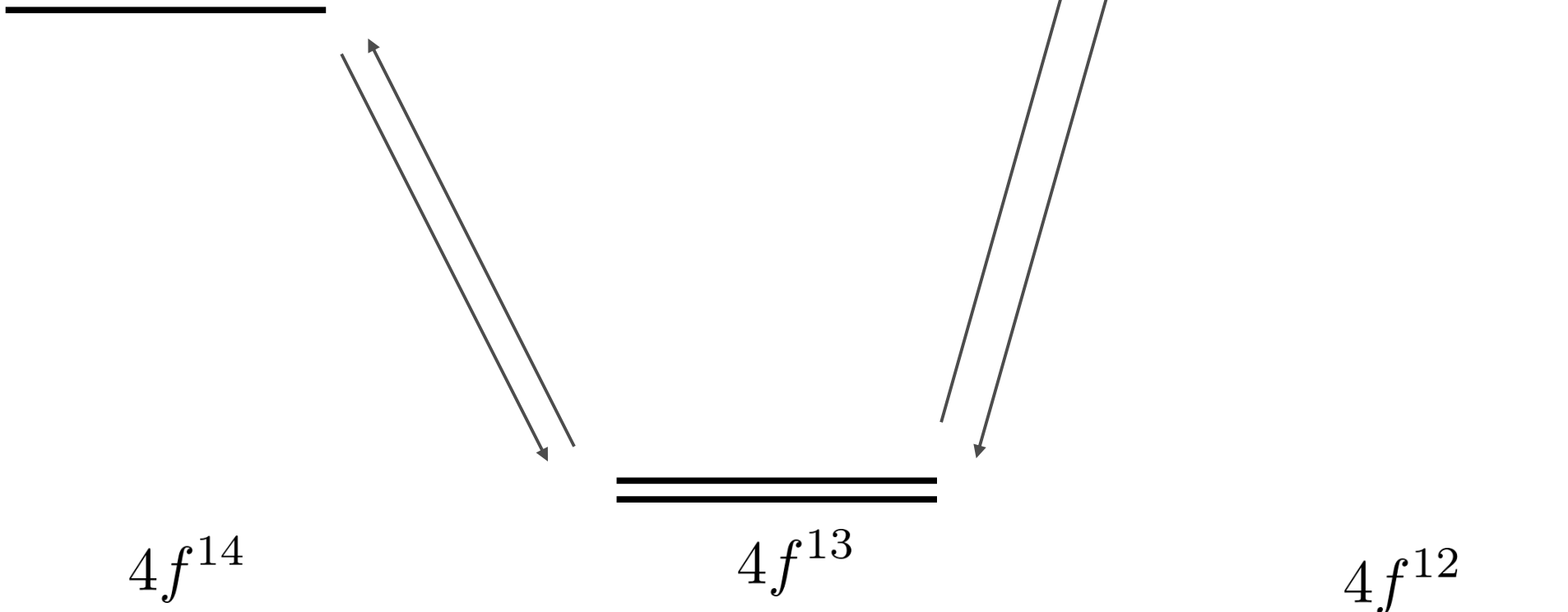
Flint, Dzero, Coleman 2008

What about Ytterbium?

Requires two orthogonal Kondo channels

$$J_1 \gg J_2$$

Composite pairing will have very low T_c



What is a composite pair?

- Total spin singlet:

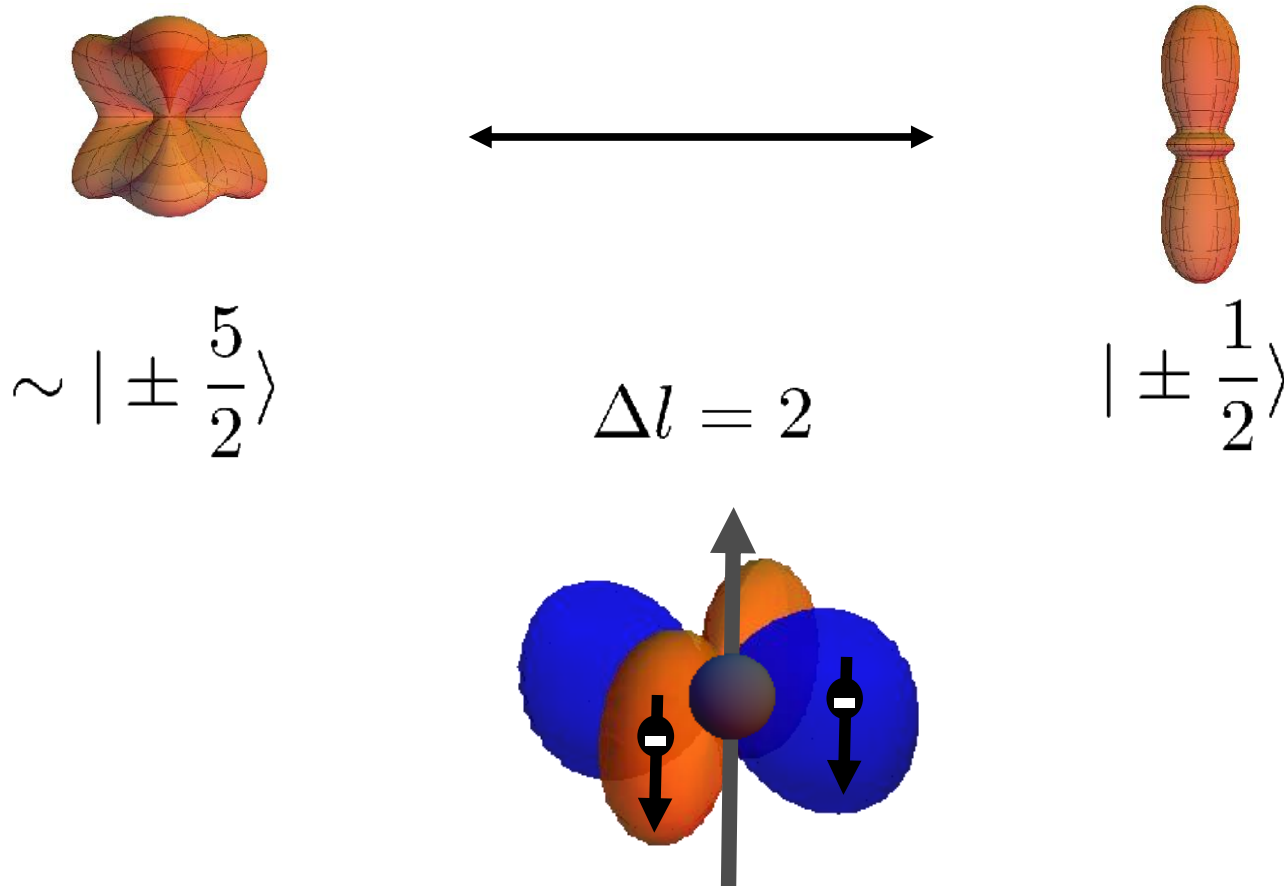
$$c_{j1\downarrow}^\dagger c_{j2\downarrow}^\dagger S_{j+}$$

What is a composite pair?

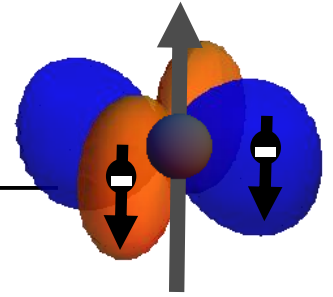
- Total spin singlet: $c_{1j}^\dagger \vec{\sigma}(i\sigma_2) c_{2j}^\dagger \cdot \vec{S}_j$

What is a composite pair?

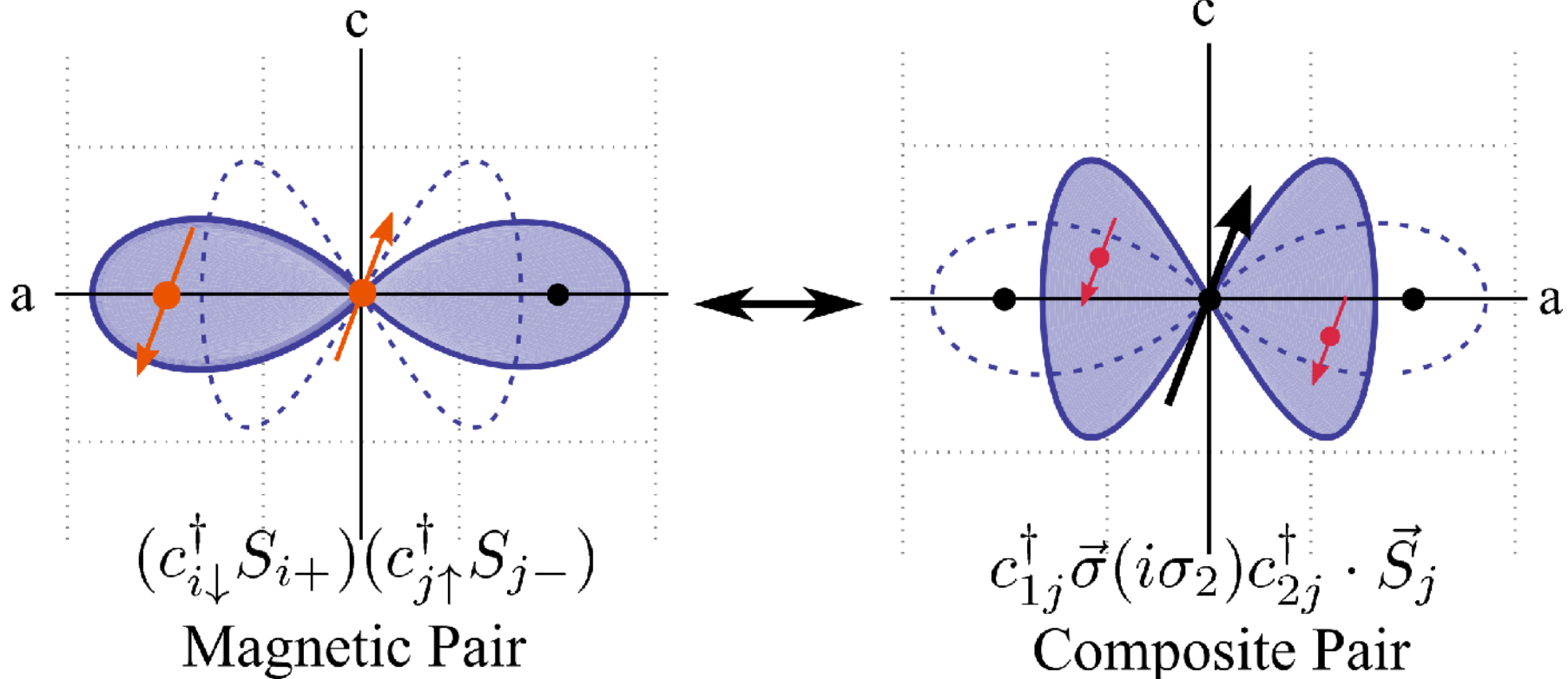
- Total spin singlet: $c_{1j}^\dagger \vec{\sigma} (i\sigma_2) c_{2j}^\dagger \cdot \vec{S}_j$
- Nodal d-wave symmetry emerges naturally from crystal fields



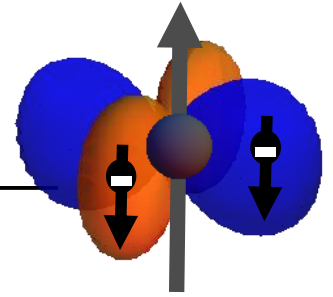
What is a composite pair?



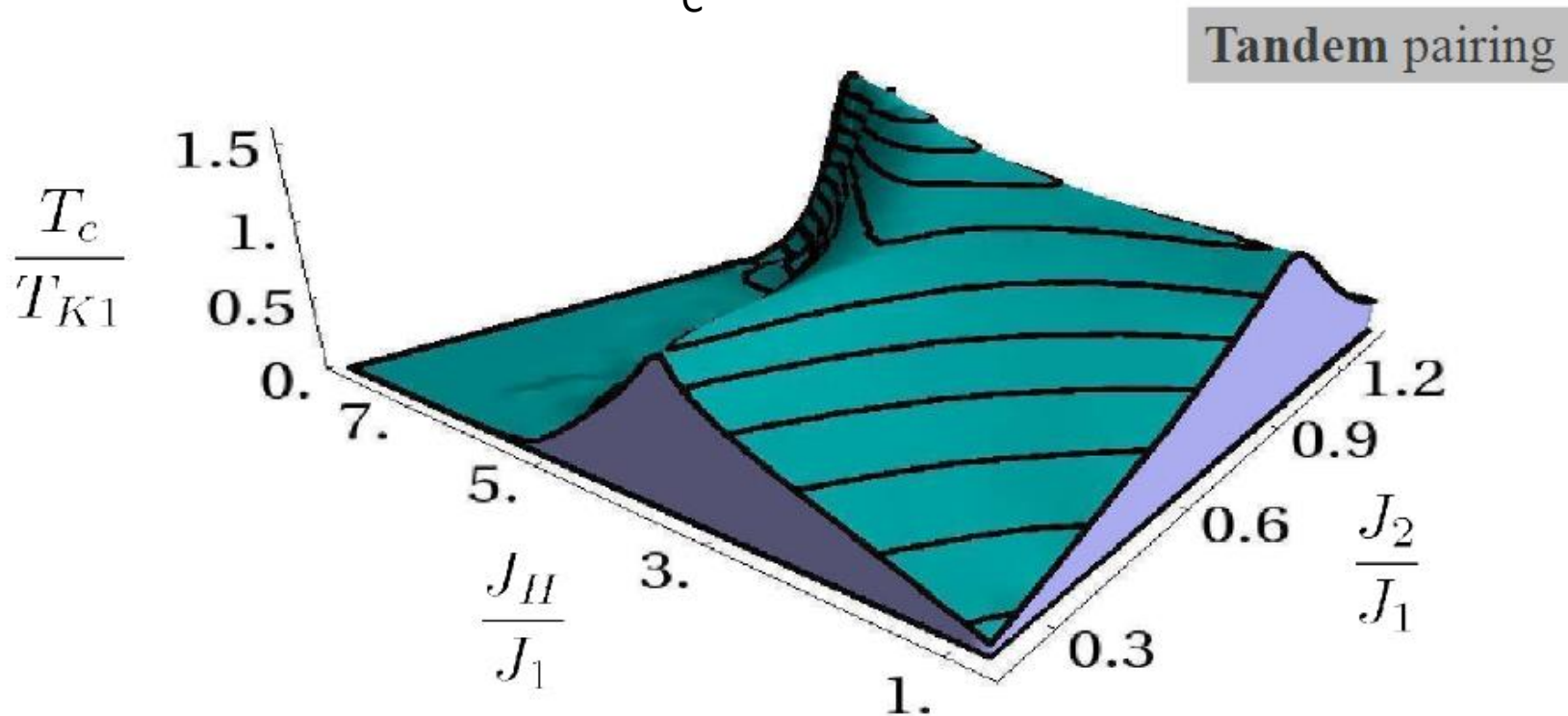
- Total spin singlet: $c_{1j}^\dagger \vec{\sigma}(i\sigma_2) c_{2j}^\dagger \cdot \vec{S}_j$
- Nodal d-wave symmetry emerges naturally from crystal fields
- These symmetries are shared with magnetic pairing, but composite pairing is *local*



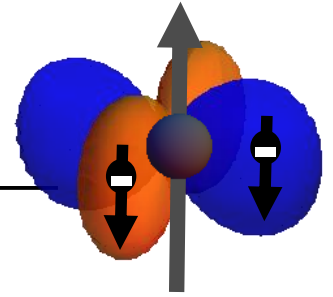
What is a composite pair?



- Total spin singlet: $c_{1j}^\dagger \vec{\sigma} (i\sigma_2) c_{2j}^\dagger \cdot \vec{S}_j$
- Nodal d-wave symmetry emerges naturally from crystal fields
- These symmetries are shared with magnetic pairing, but composite pairing is **local**
 - Can work *in tandem* to raise T_c



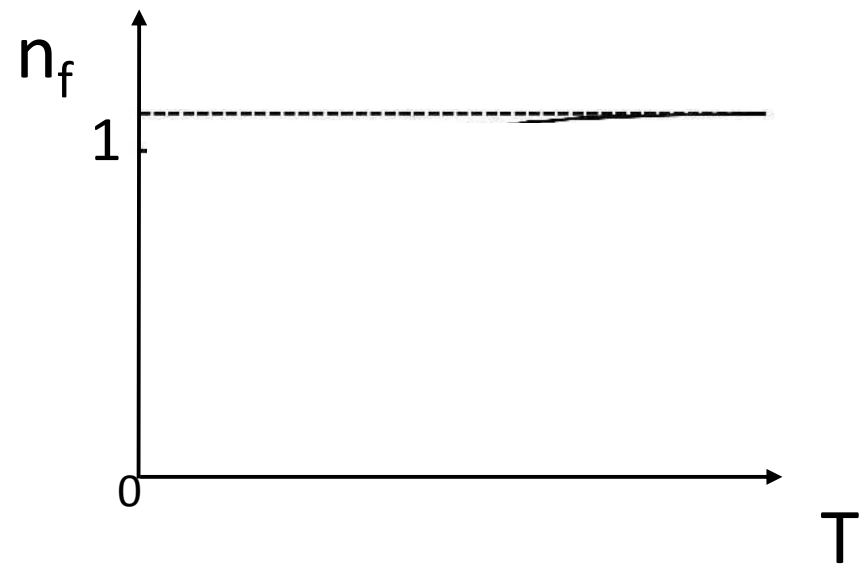
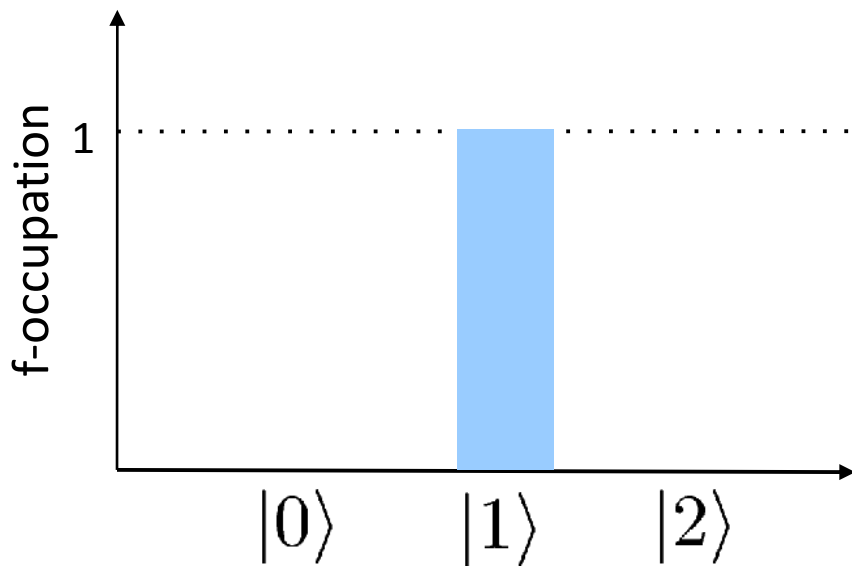
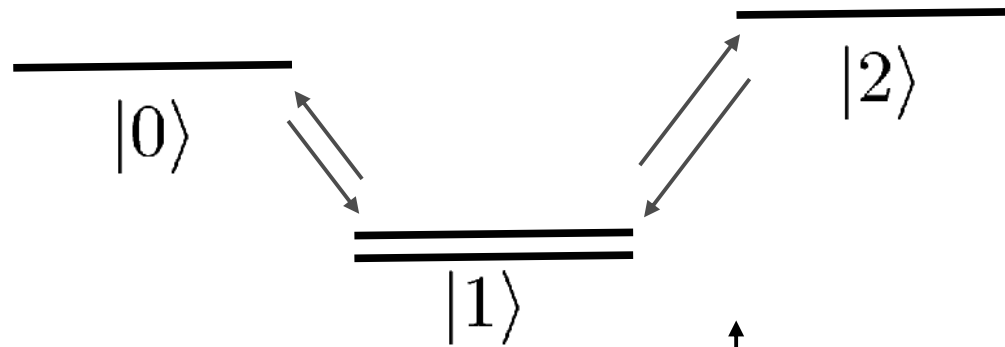
What is a composite pair?



- Total spin singlet: $c_{1j}^\dagger \vec{\sigma} (i\sigma_2) c_{2j}^\dagger \cdot \vec{S}_j$
- Nodal d-wave symmetry emerges naturally from crystal fields
- Charge aspects: Kondo effect *adds and removes* electrons

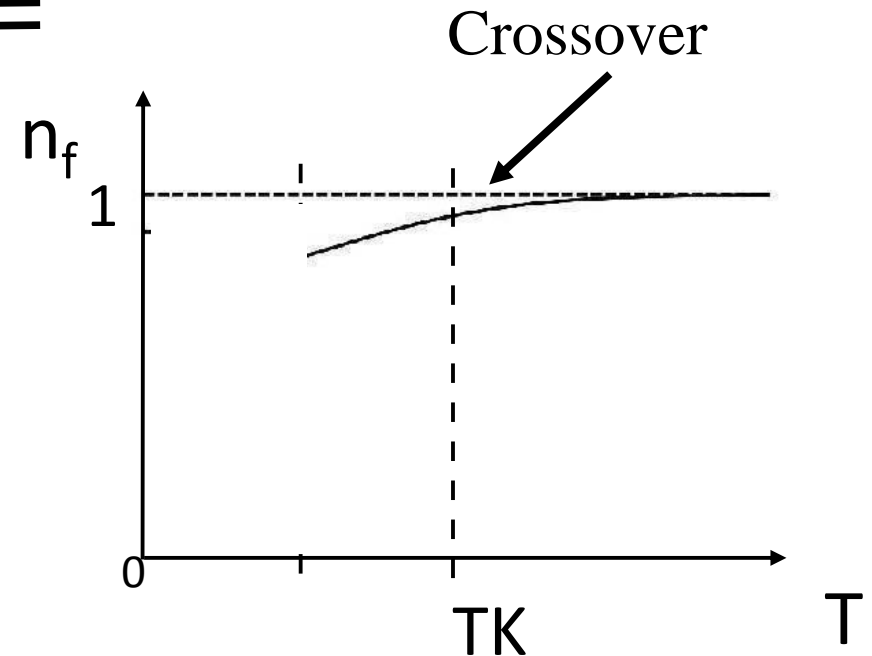
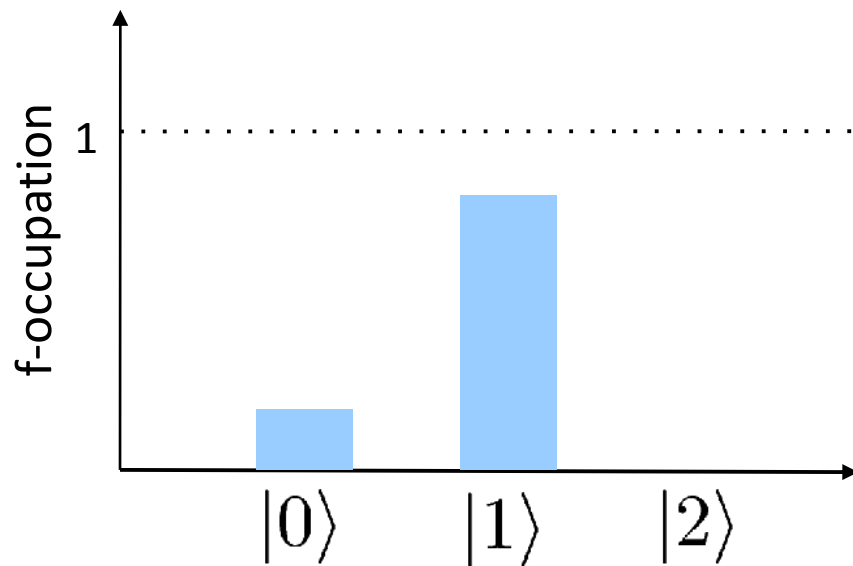
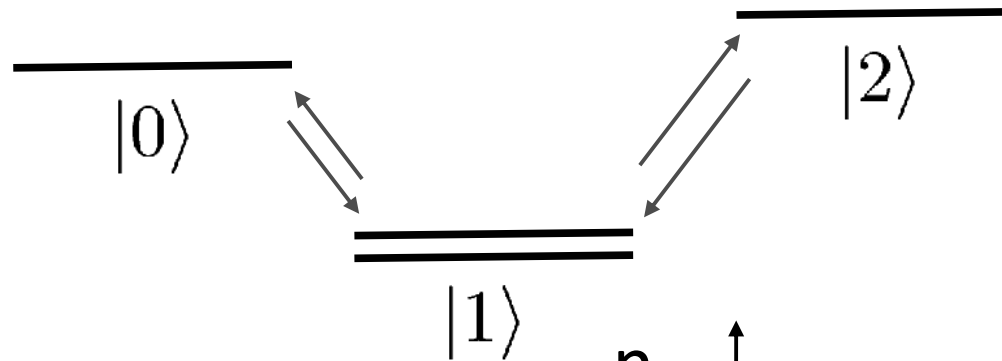
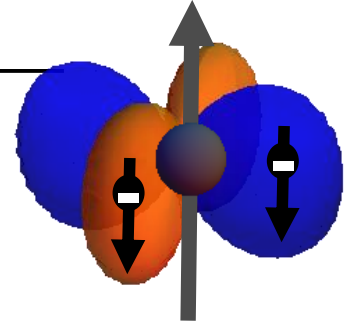
Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons



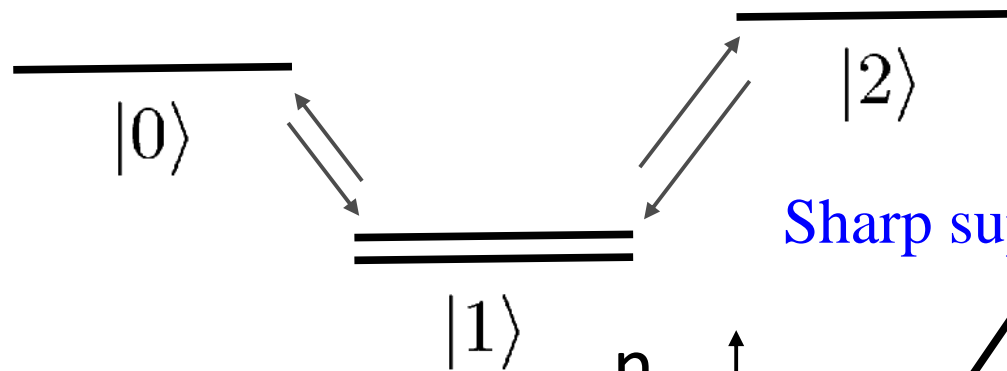
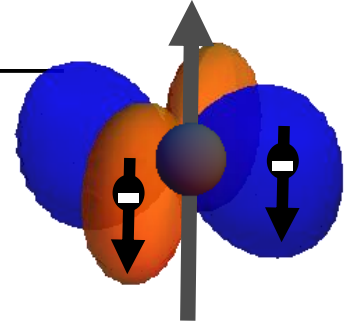
Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons

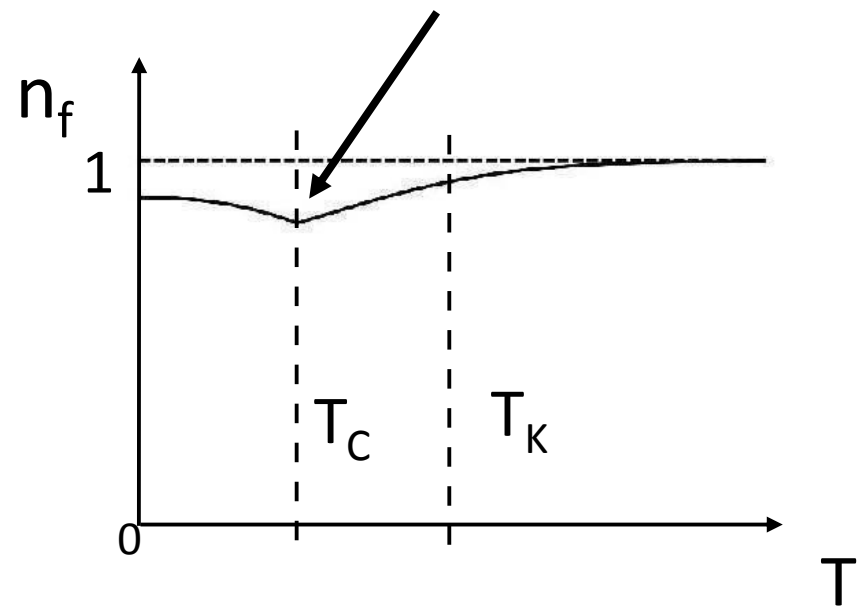
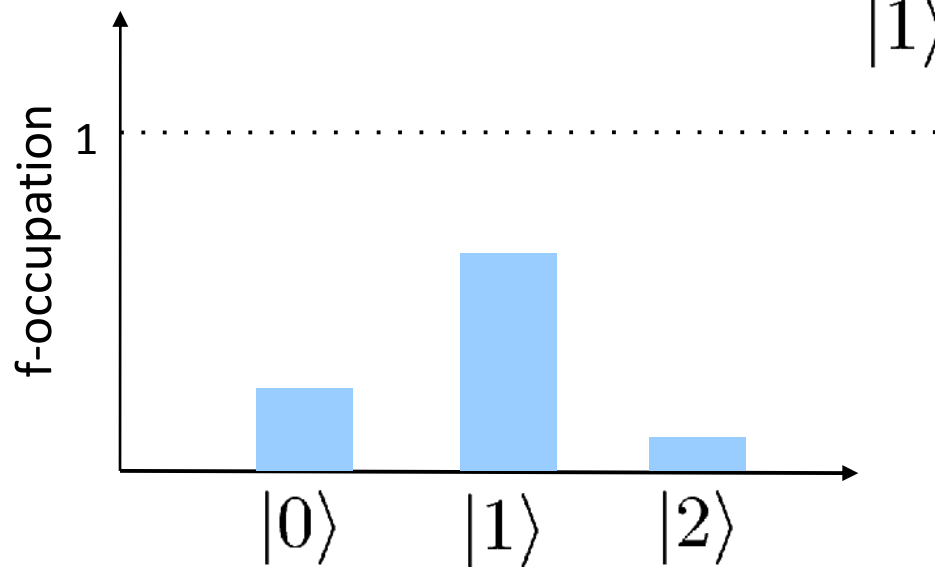


Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons



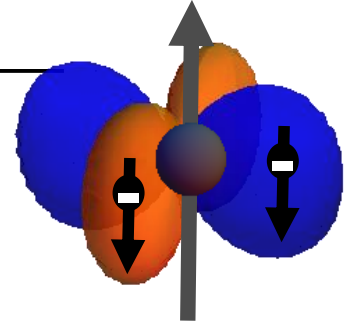
Sharp superconducting shift



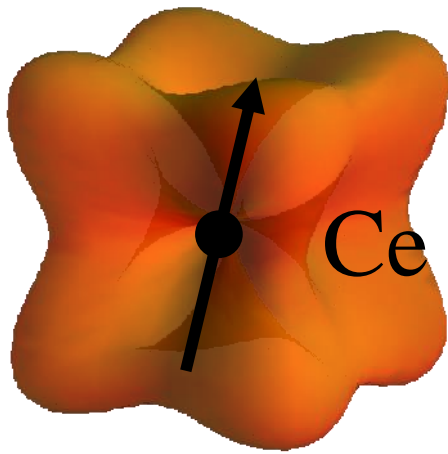
$n_f(T)$ measured by core-level x-ray spectroscopy or Mossbauer isomer shift

Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons

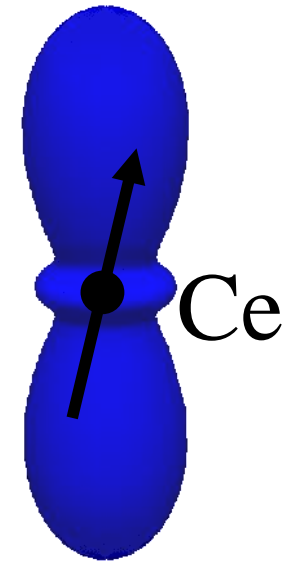


Hole Kondo effect



Γ_7^+

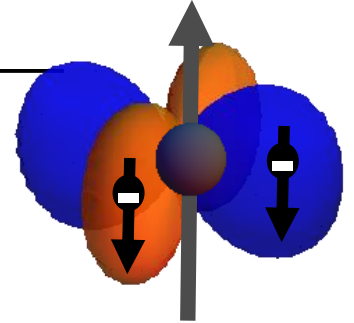
Electron Kondo effect



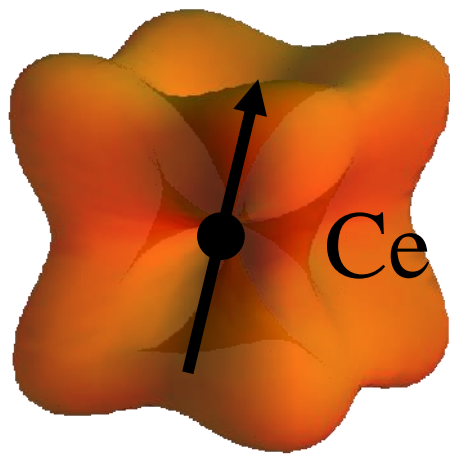
Γ_6

Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons



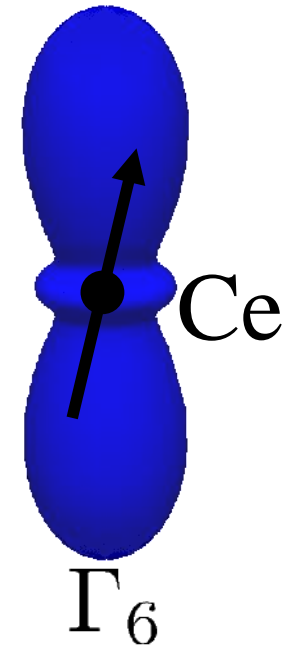
Hole Kondo effect



Transfer electrons

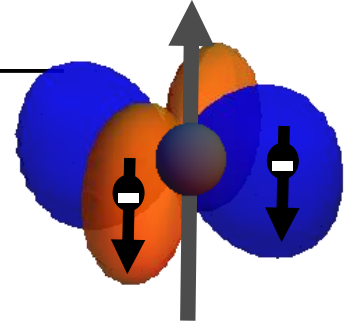


Electron Kondo effect

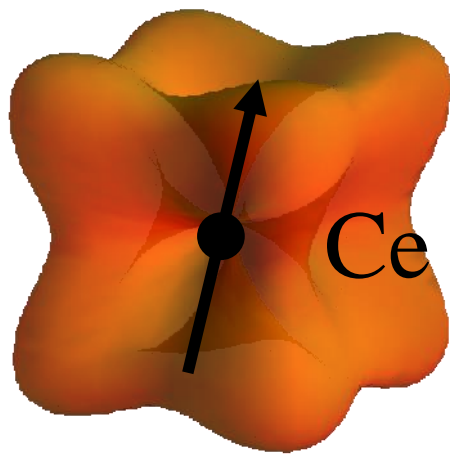


Electrostatically active composite condensate

- Charge aspects
 - Hybridization *adds and removes* f-electrons



Hole Kondo effect

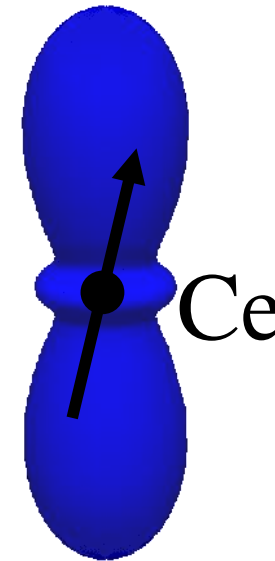


Γ_7^+

Transfer electrons



Electron Kondo effect



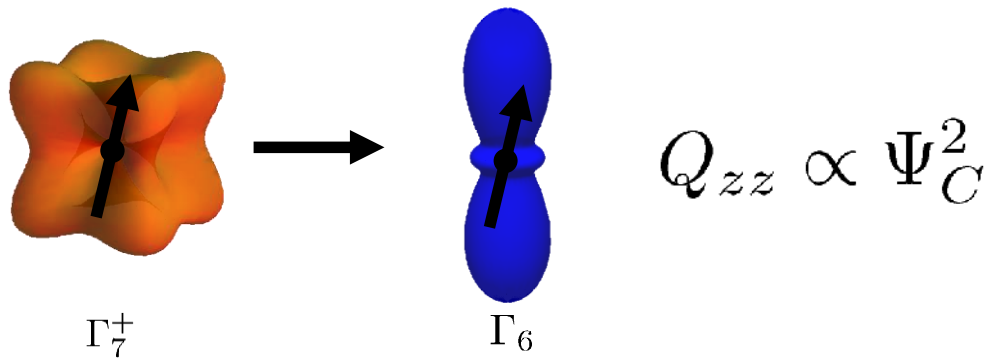
Γ_6

- Composite pair condensate acquires a quadrupole moment

$$Q_{zz} \propto \Psi_C^2$$

Electrostatically active composite condensate

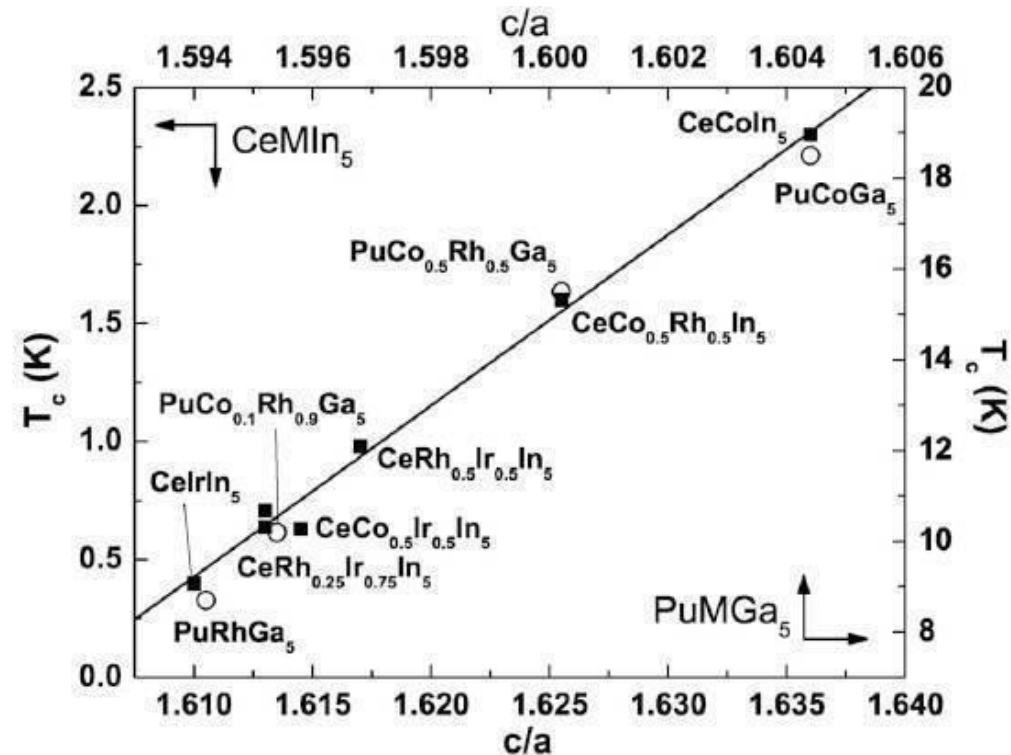
- Charge aspects
 - Composite condensate acquires a quadrupole moment



The quadrupole moment couples to tetragonal strain $\sim c/a$

$$\Delta F \propto -Q_{zz} \frac{c}{a}$$

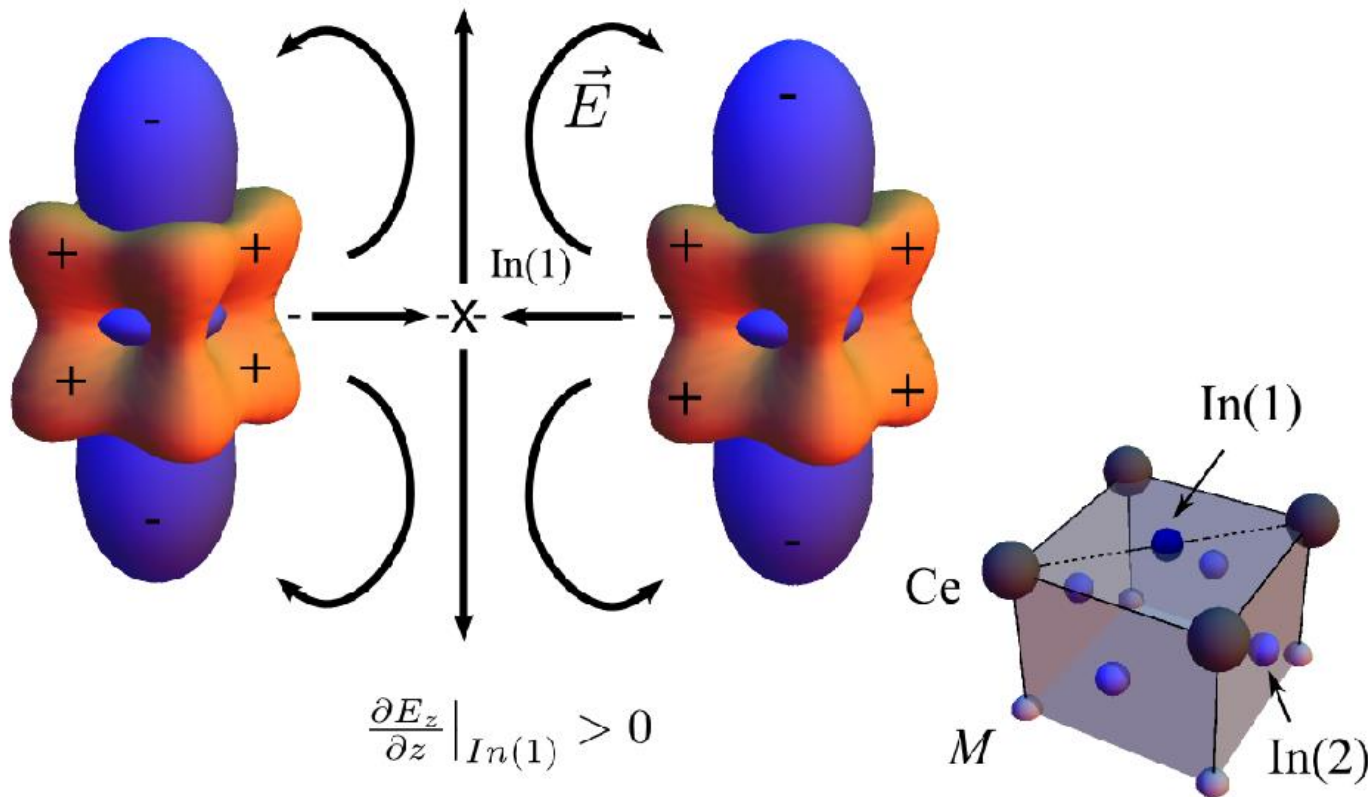
$$F = \alpha_2 \left[T - \left(T_c + \lambda \frac{c}{a} \right) \right] \Psi_C^2$$



T_c increases linearly with c/a

Electrostatically active composite condensate

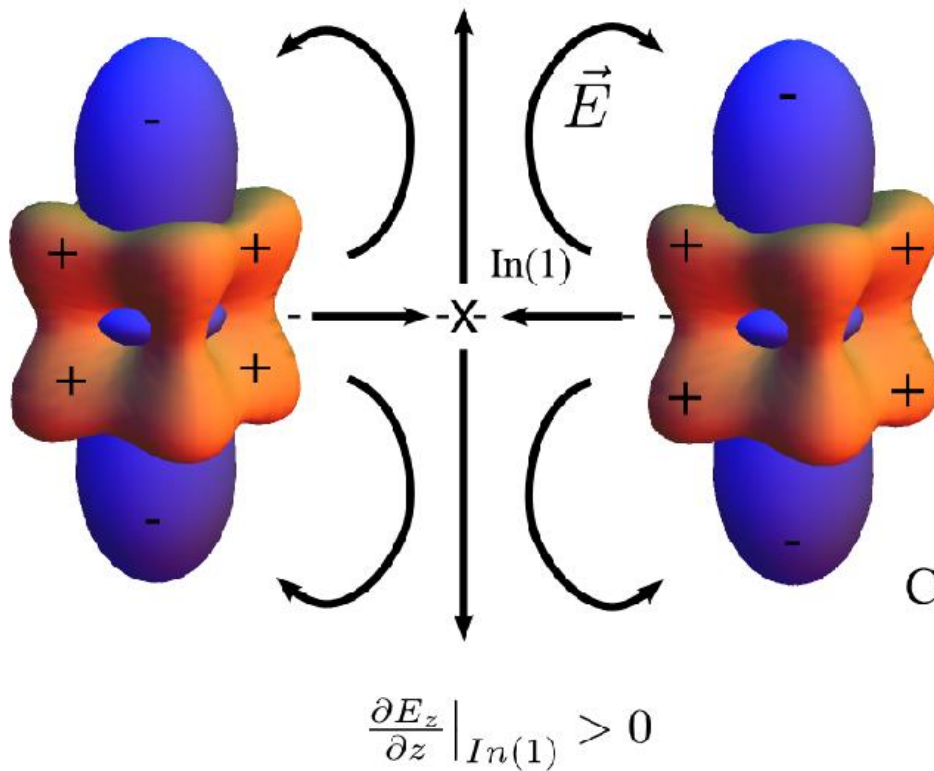
- Charge aspects
 - Composite condensate acquires a quadrupole moment



Measurable by sharp shifts in NQR or Mossbauer

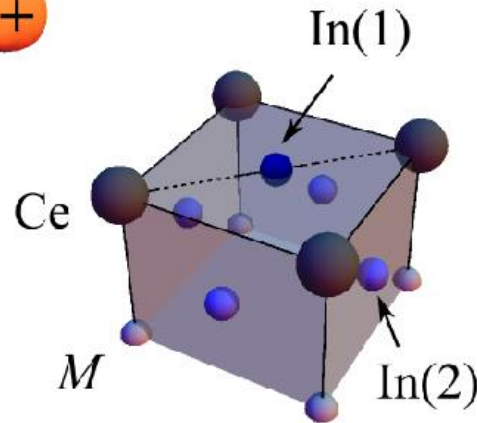
Electrostatically active composite condensate

- Charge aspects
 - Composite condensate acquires a quadrupole moment



$$\Delta\rho(\mathbf{x}) \propto V_1^2 \rho_{7+}(\mathbf{x}) - \Delta_2^2 \rho_6(\mathbf{x})$$

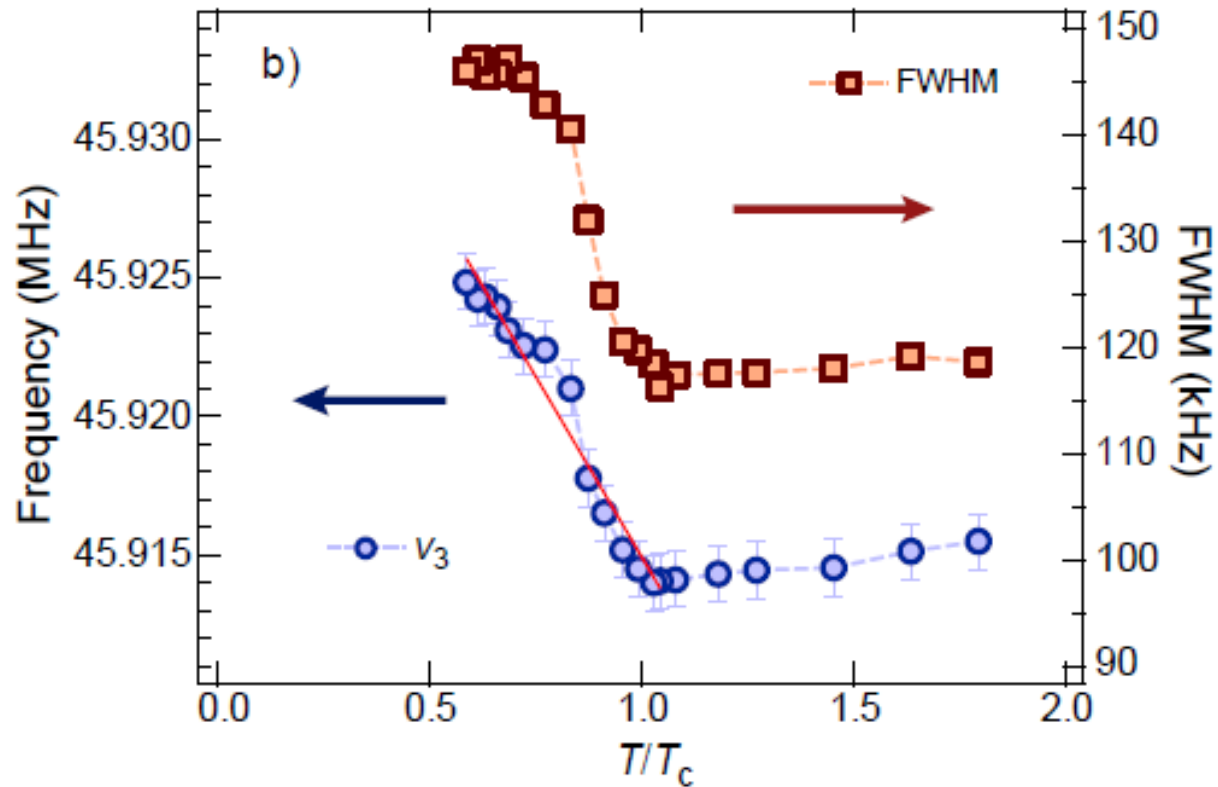
$$\Delta\nu_{NQR} \approx 5\text{kHz/K} \times (T_c - T)$$



Measurable by sharp shifts in NQR or Mossbauer

NQR shift observed in PuCoIn₅!

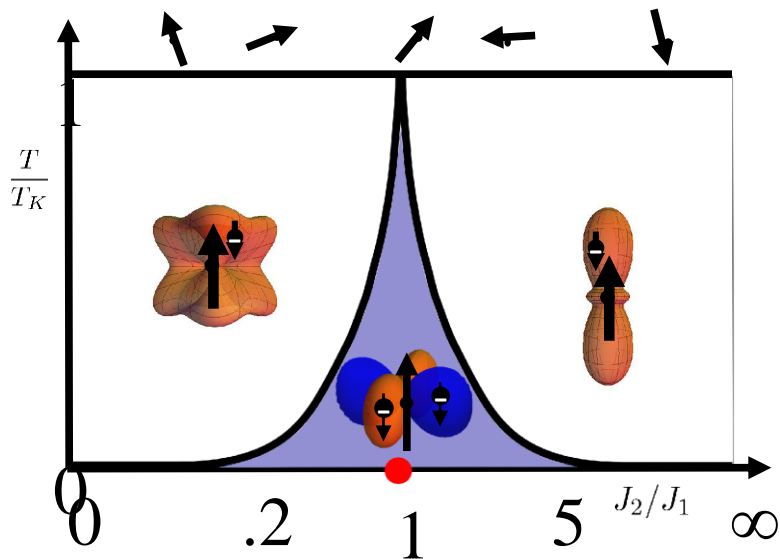
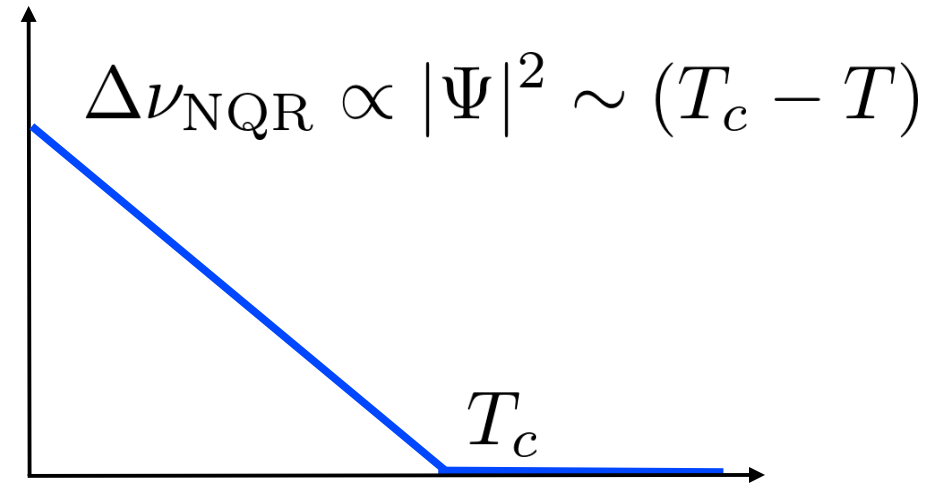
Sharp linear shift in NQR frequency seen at T_c in PuCoIn₅
Also seen in PuRhIn₅, CeCoIn₅



NQR Shift

Three possible origins of NQR shift

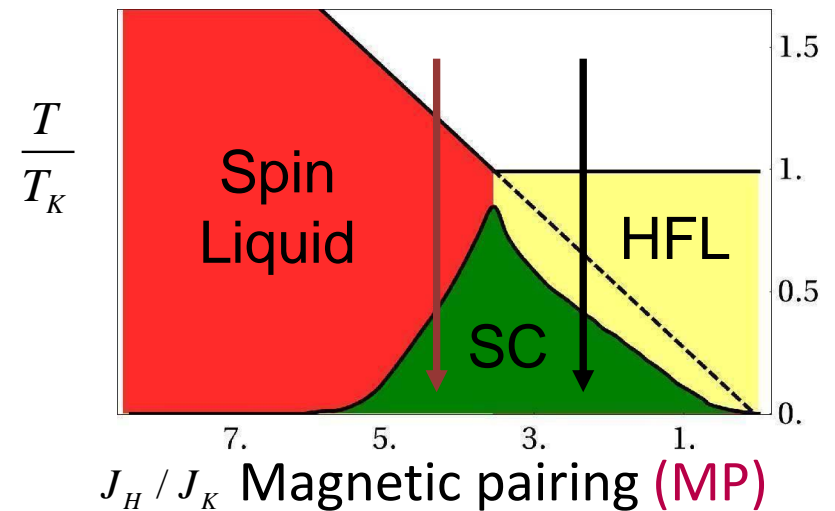
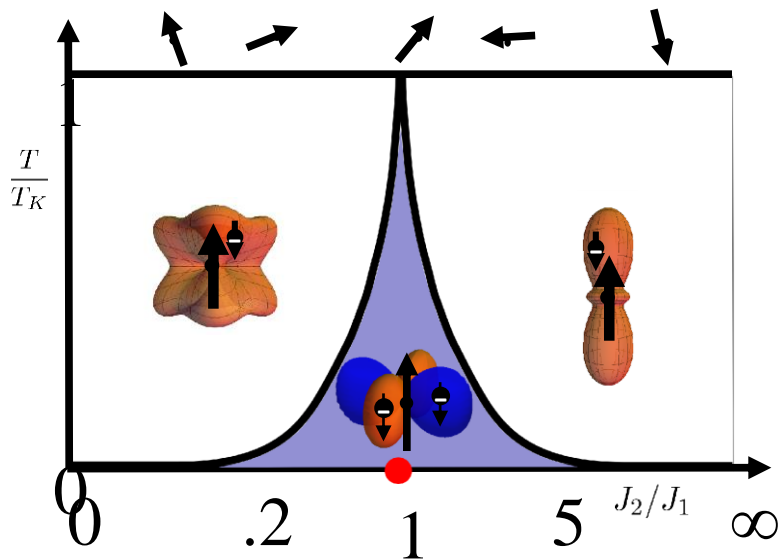
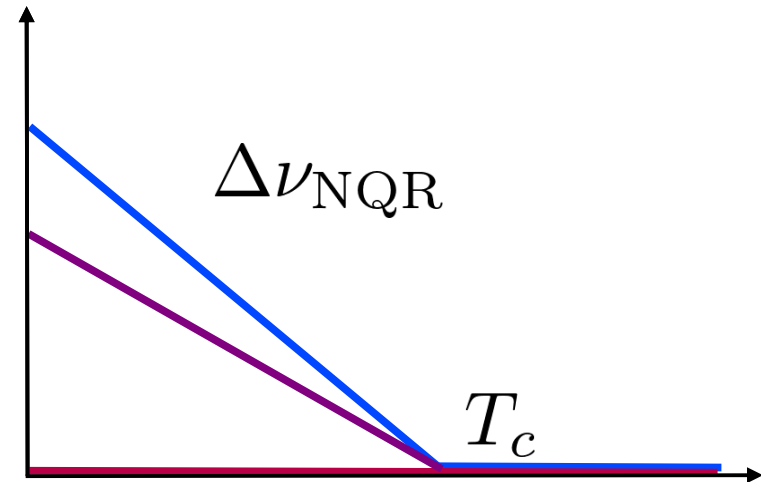
- Change in lattice parameters
- Change in valence Δn_f (CP)
- Change in on-site symmetry (CP)



NQR Shift

Three possible origins of NQR shift

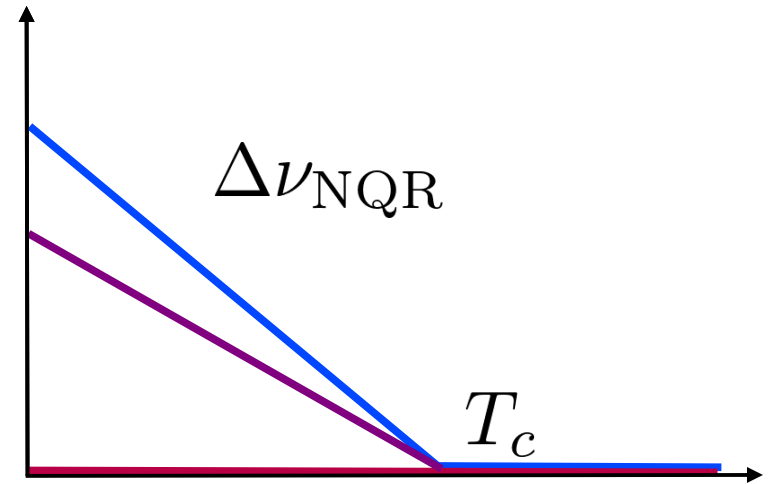
- Change in lattice parameters
- Change in valence Δn_f (CP, MP)
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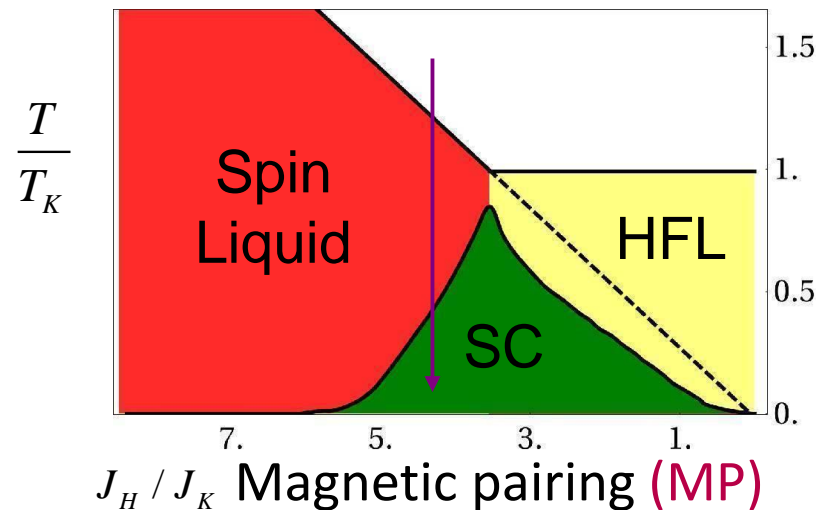
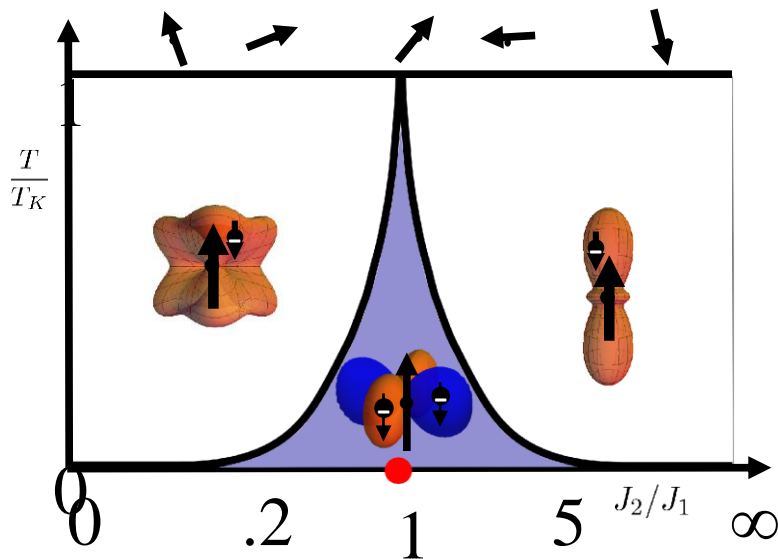
NQR Shift

Three possible origins of NQR shift

- Change in lattice parameters
- Change in valence Δn_f (CP, MP)
- Change in on-site symmetry (CP)



Magnetic pairing gives $\Delta\nu_{\text{NQR}}$ only below “spin liquid” – CeRhIn₅? **Not** CeCoIn₅
Composite pairing *always* gives $\Delta\nu_{\text{NQR}}$

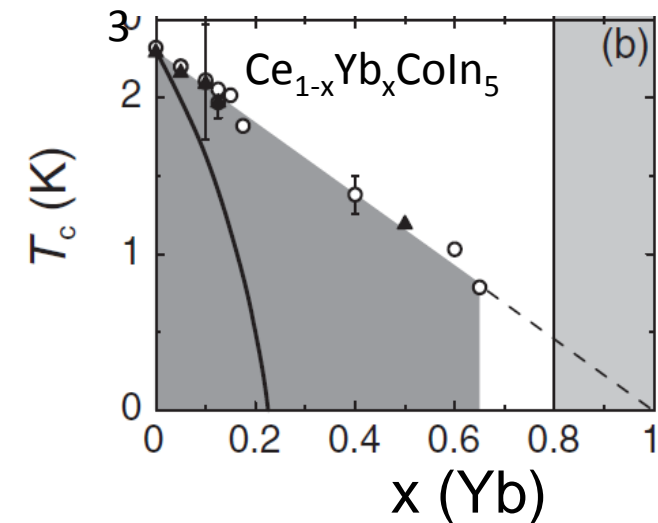
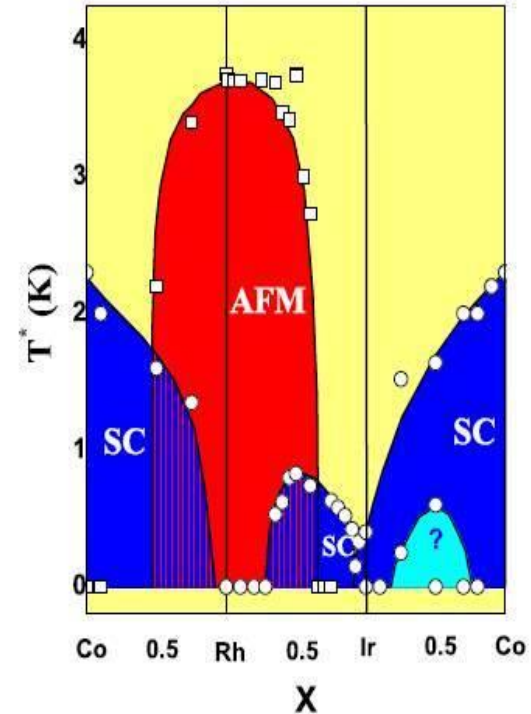


Magnetic pairing appears ubiquitous

Sarrao and Thompson, JPSJ (2007)

But...

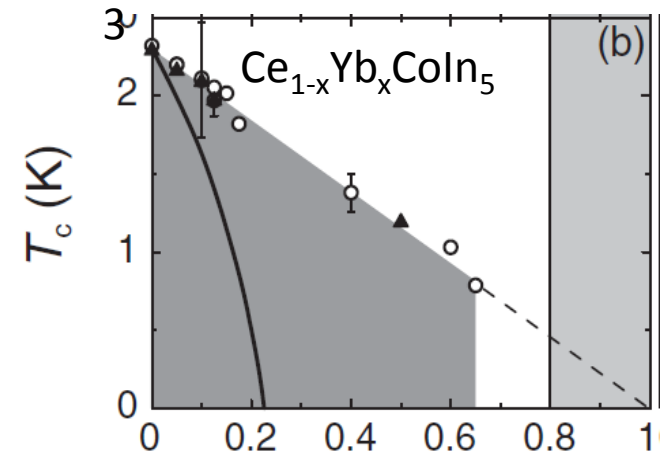
- Two domes in CeMIn_5 ($M = \text{Co}, \text{Rh}, \text{Ir}$)
 - *Tuning between composite and magnetic pairing*
- Superconductivity without magnetism
 - *Not required for composite pairing*
- Extremely robust to disorder *on the f-site*
 - *Local nature of composite pairs protects T_c*
- Nodeless superconductivity in $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$ ($x > .2$)
 - *Composite pairing doesn't require an underlying Fermi surface*



L. Shu et al PRL 2011

Superconductivity in Yb-doped CeCoIn₅

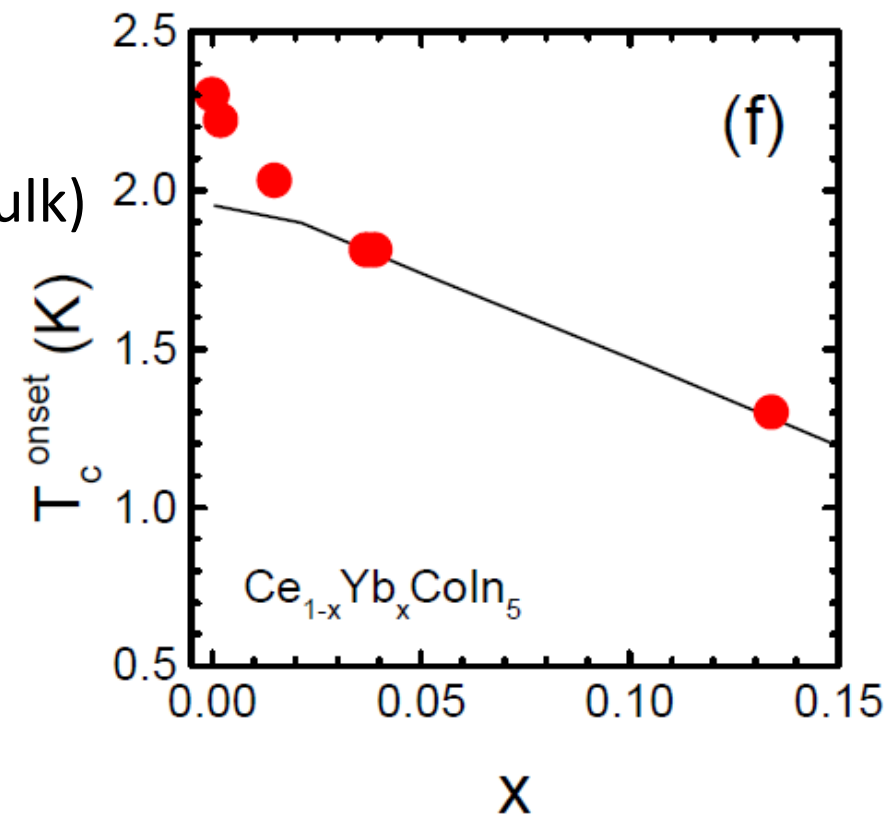
- Ce_{1-x}R_xCoIn₅ superconducts out to x ~ .25-.33
- Both T_c and T* are unexpectedly unaffected by these Kondo holes
- Yb-doped CeCoIn₅ seems to be special
 - Superconducts out to nominal x_n ~ 1 (!)



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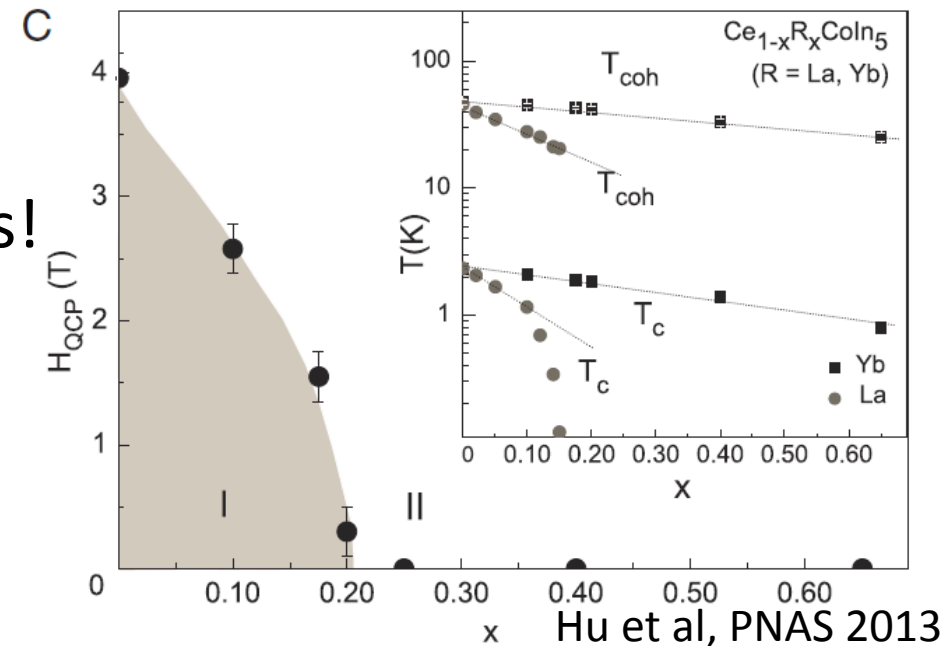
Superconductivity in Yb-doped CeCoIn₅

- Ce_{1-x}R_xCoIn₅ superconducts out to $x \sim .25-.33$
- Both T_c and T^* are unexpectedly unaffected by these Kondo holes
- Yb-doped CeCoIn₅ seems to be special
 - Superconducts out to nominal $x_n \sim 1$ (!)
 - New data shows actual $x \sim x_n/3$
(resolves discrepancy between films +bulk)



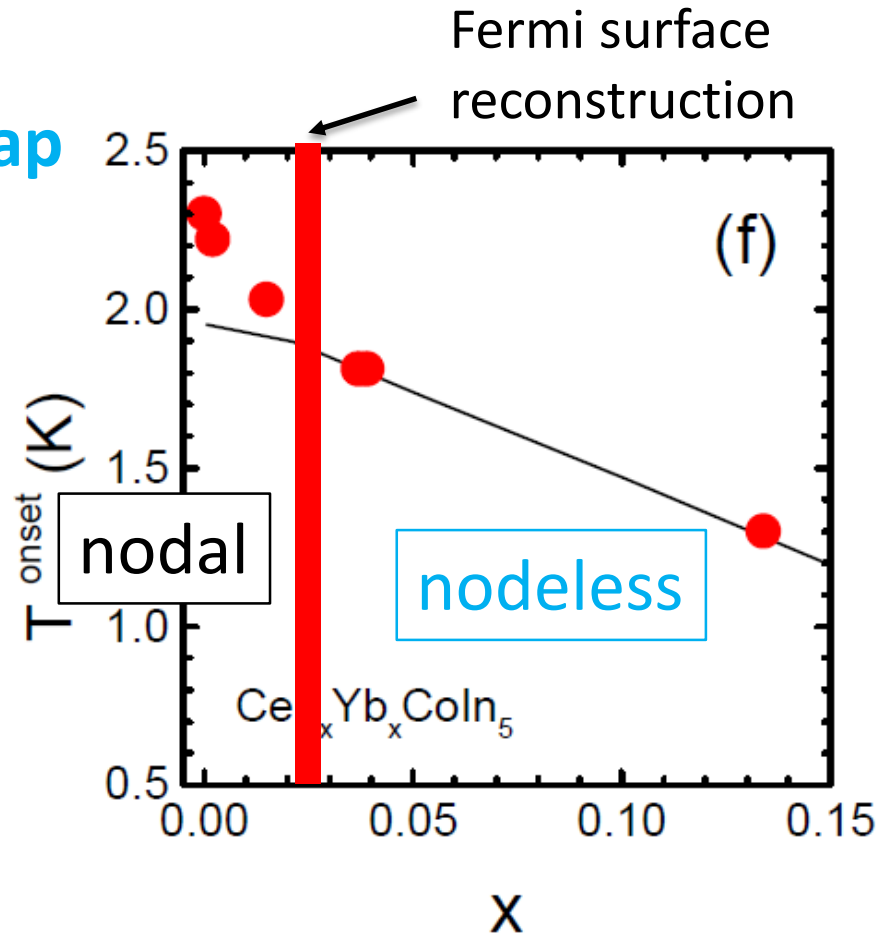
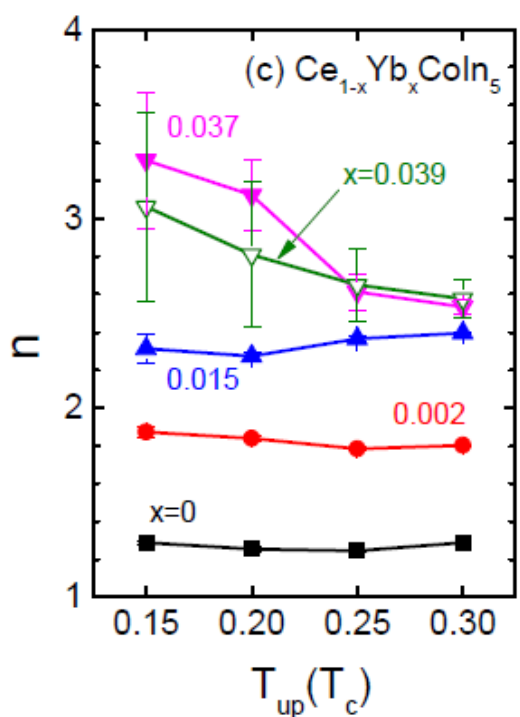
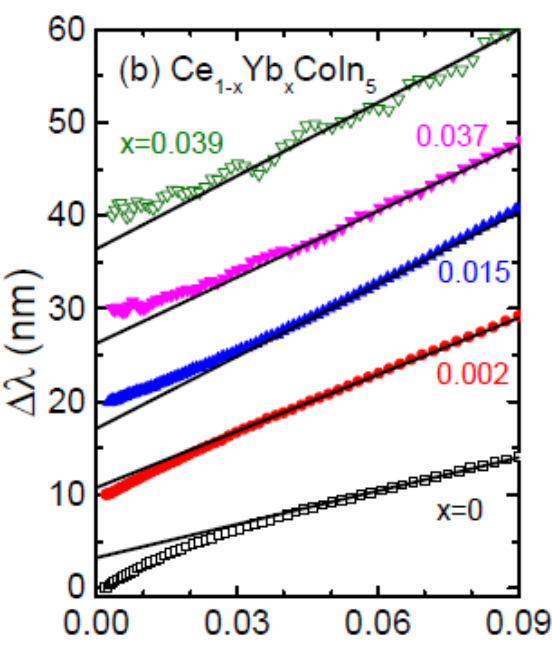
Superconductivity in Yb-doped CeCoIn₅

- Yb-doped CeCoIn₅ seems to be special
 - Yb is mixed valent ($n_f \sim 2.3$)
- Unusual quantum criticality, suppressing H_{QCP} at $x_n = .2$
- Lifshitz transition in dHvA at $x_n = .2$
 - Hole doping removes Fermi surface
- Superconductivity becomes nodeless!
 - Also at $x_n = .2$



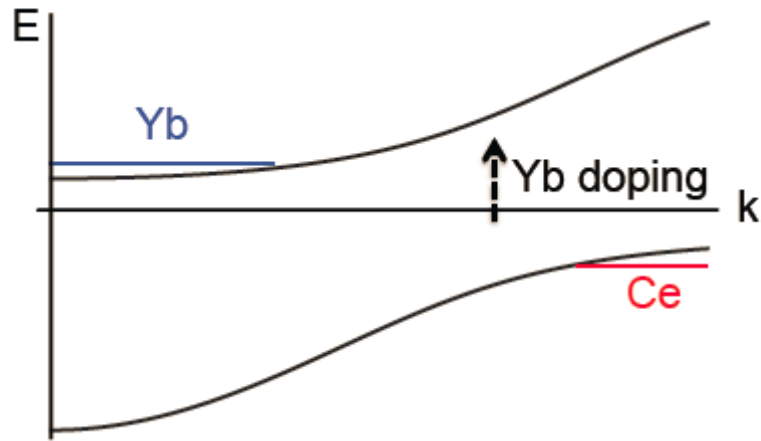
Nodeless superconductivity in $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$

- Pure CeCoIn_5 is a nodal superconductor ($d_{x^2-y^2}$)
- The penetration depth approaches $\Delta\lambda \sim T$ for very clean samples
 - $\Delta\lambda \sim T^2$ for dirtier samples (dirty d-wave)
- Beyond a critical x_c , the penetration depth goes as $\Delta\lambda \sim T^n$ $n > 3$
 - Inconsistent with nodes => **full gap**

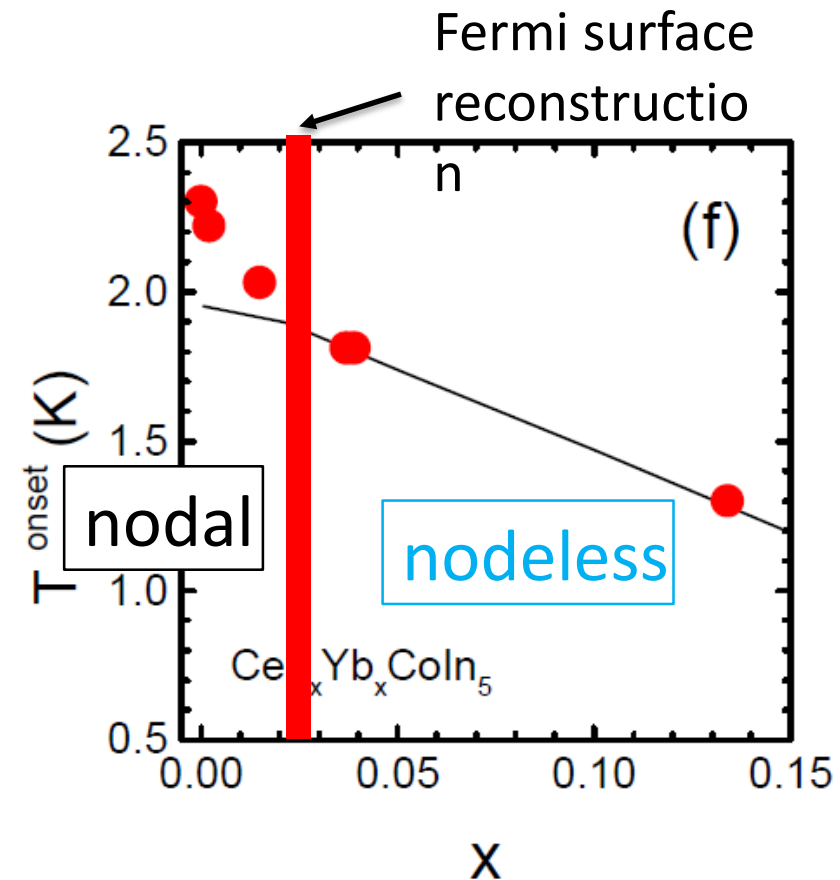


Nodeless superconductivity in $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$

- $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$ is a **fully gapped** superconductor beyond $x_n = .2$
- How can Yb doping induce nodeless superconductivity?
 - Adds holes, adjusts Fermi surface
 - Fermi sheet disappears, but T_c is smooth!

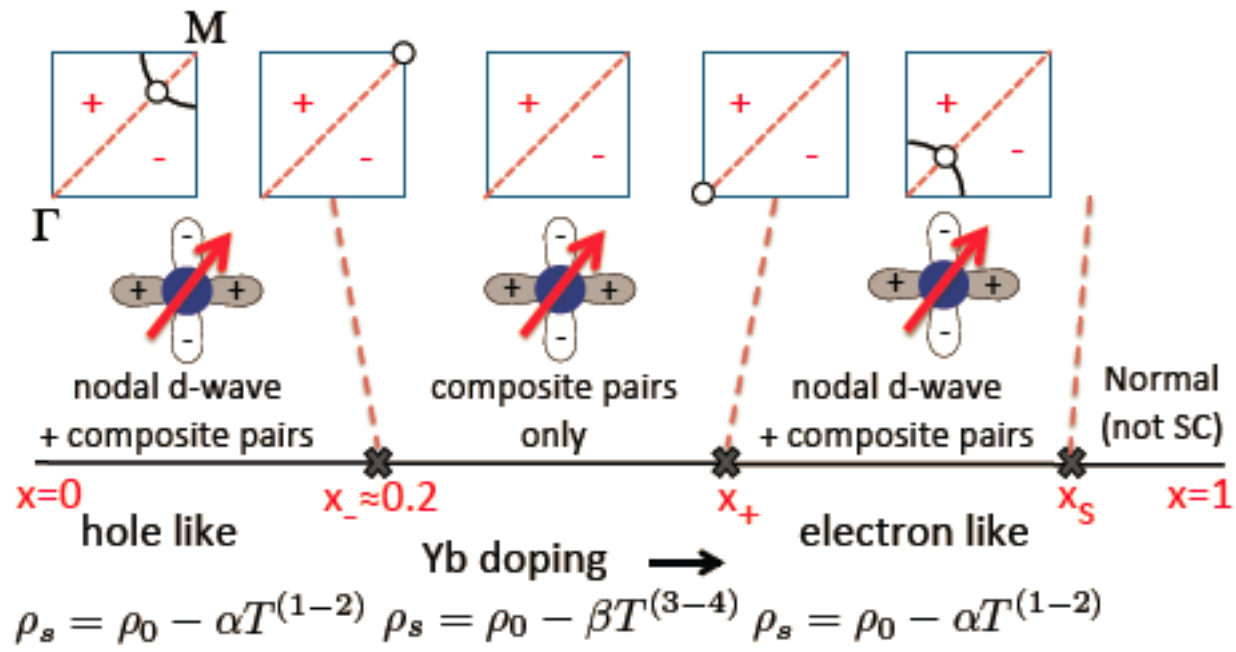


- Composite pairing does not require a FS
 - Can exist even below a Kondo insulator



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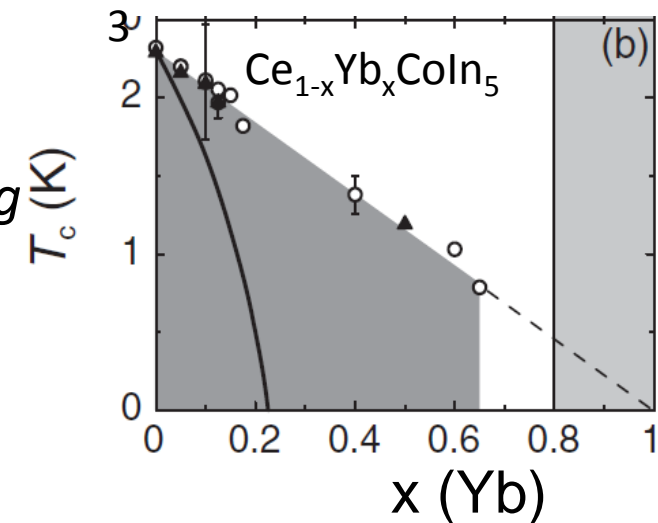
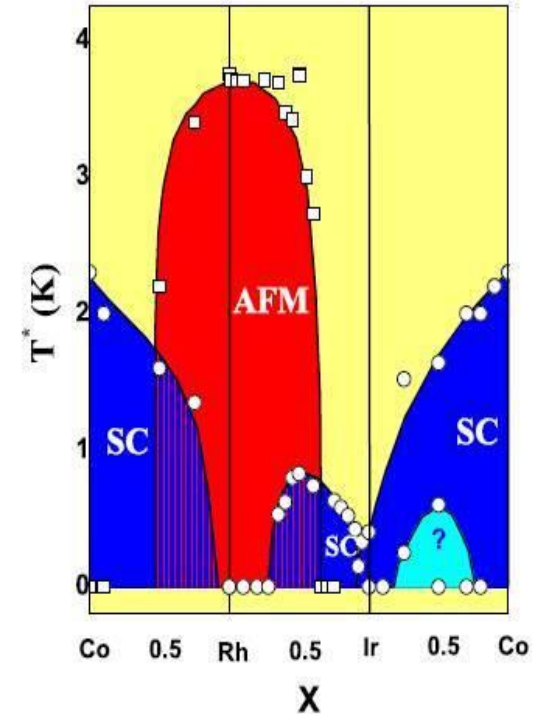
Kim et al, arXiv: 1404.3700
 Erten, Flint, Coleman, arXiv:1402.7361

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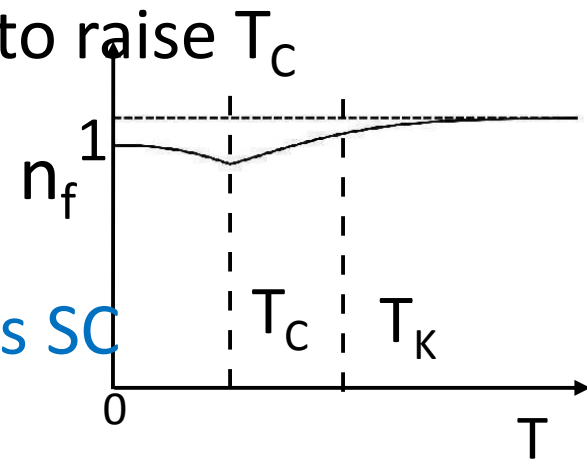
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- Nodeless superconductivity in $\text{Ce}_{1-x}\text{Yb}_x\text{CoIn}_5$ ($x > .2$)
 - *Composite pairing doesn't require an underlying Fermi surface*
- Many Ce superconductors, only one Yb
 - *Magnetism doesn't discriminate, but composite pairing is inactive in Yb compounds*



L. Shu et al PRL 2011

Conclusions

- Two electrons in orthogonal Kondo channels can screen the same local moment to form a composite pair
- Composite pairing produces d-wave, singlet pairs
 - Can work in tandem with magnetic pairing to raise T_C
- To resolve the two – look at the charge:
 - Suggestive: NQR shift
 - Suggestive: robustness to disorder, nodeless SC
 - Smoking gun: valence shift



R. Flint, M. Dzero and P. Coleman, Nature Physics 4, 643(2008)

R. Flint and P. Coleman, Phys. Rev. Lett. 105, 246404 (2010)

R. Flint, A.H. Nevidomskyy and P. Coleman, Phys. Rev. B 84, 064514 (2011)

Open questions

- How does disorder affect different pairing mechanisms?
- Quantum criticality related to two channel Kondo physics?
- Can composite pairing be extended to other materials?
(d-electron 115s?)
- Tandem pairing in other materials (eg – nematic fluctuations aiding magnetic pairing)
- Relationships to pair-density wave in Q-phase of CeCoIn_5 ?