

## Disks and Massive Planet Formation

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### Outline:

- ▶ Planet formation in the Solar System
- ▶ Observations = problems with simplest theory
- ▶ Migration: how important is disk turbulence?

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### Traditional stages of planet formation

#### 1) Dust ( $\mu\text{m}$ ) to pebbles (cm)

Small dust particles are well-coupled to the gas, settling velocity is small even *without turbulence*:

$$v_z = - \frac{\Omega^2 z a_{dust}}{C_s} \left( \frac{\rho_{dust}}{\rho_{gas}} \right) \quad t_{\text{settle}} \sim 1 \text{ Myr for } a \sim 1 \mu\text{m at } 1 \text{ AU}$$

Growth occurs via collisions + settling as size increases

Theoretically (eg Dullemond & Dominik 05) and observationally (eg Shuping et al. 03) this is not believed to be a rate limiting step

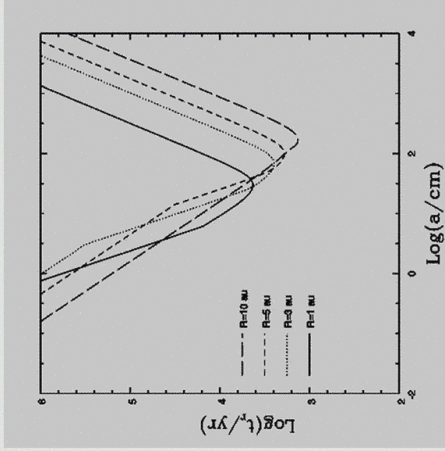
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2) Pebbles (cm) to planetesimals (>km)

Must occur rapidly - otherwise solids lost into the star due to aerodynamic drag against the gas

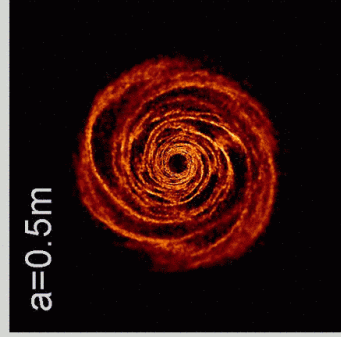
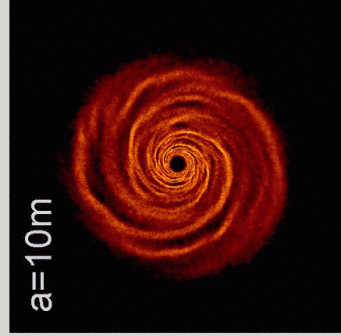
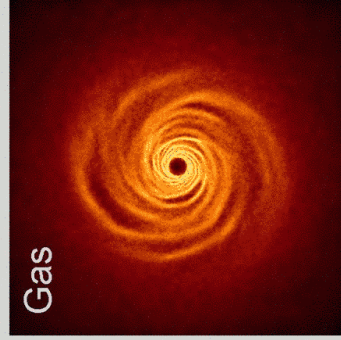
$$v_{\phi}^2 = \frac{GM_*}{r} + \frac{1}{\rho} \frac{dP}{dr}$$

'headwind' is  $O(h/r)^2$



$t_{\text{drift}} \sim 10^4$  yr or less, AND collisions of m-scale objects may not be very sticky  
 gravitational instability of a dust layer???

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$$v_{\phi}^2 = \frac{GM_*}{r} + \frac{1}{\rho} \frac{dP}{dr}$$

If the disk is very turbulent, can enhance solids at pressure maxima

e.g. self-gravity in the gas (Rice et al. 04) could locally enhance gas / dust ratio at  $a=50\text{cm}$  by  $\sim 10^2$

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### 3) Planetesimals to terrestrial planets and giant planet cores

N-body ( $N \sim 10^{11}$ ) evolution under gravity.

Growing planet accretes from an annulus of width:

$$\Delta r \approx \text{few } r_H \quad r_H = \left( \frac{M_p}{3 M_*} \right)^{1/3} a$$

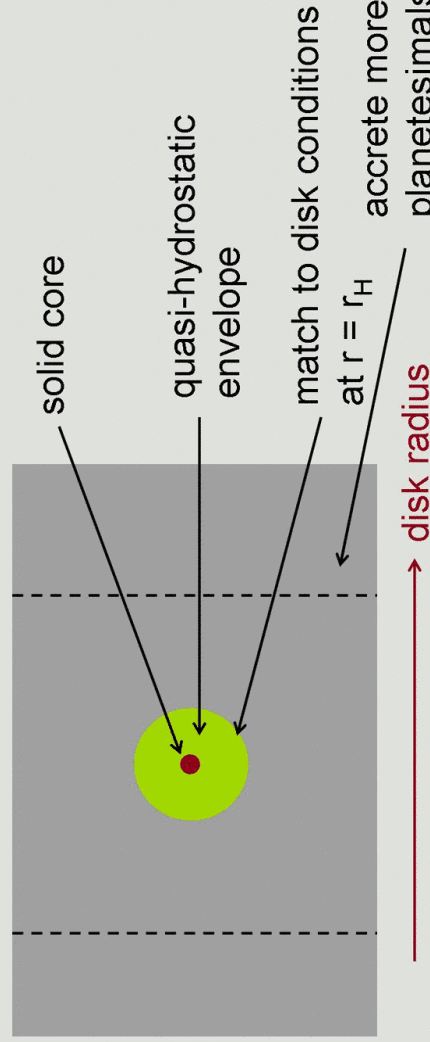
Allows rapid growth up to an isolation mass (at which point annulus is depleted of planetesimals):

$$M_{iso} \propto a^3 \sigma_p^{3/2} \quad \begin{array}{l} \text{- typically } \sim 0.1 \text{ Earth masses at } 1\text{AU,} \\ \sim \text{Earth mass at } 5\text{AU} \end{array}$$

**Growth beyond isolation is slower, but can plausibly form Earth in  $\sim 100$  Myr. Control variable for structure of the terrestrial planets is  $\sigma_p$  (eg Raymond et al. 04)**

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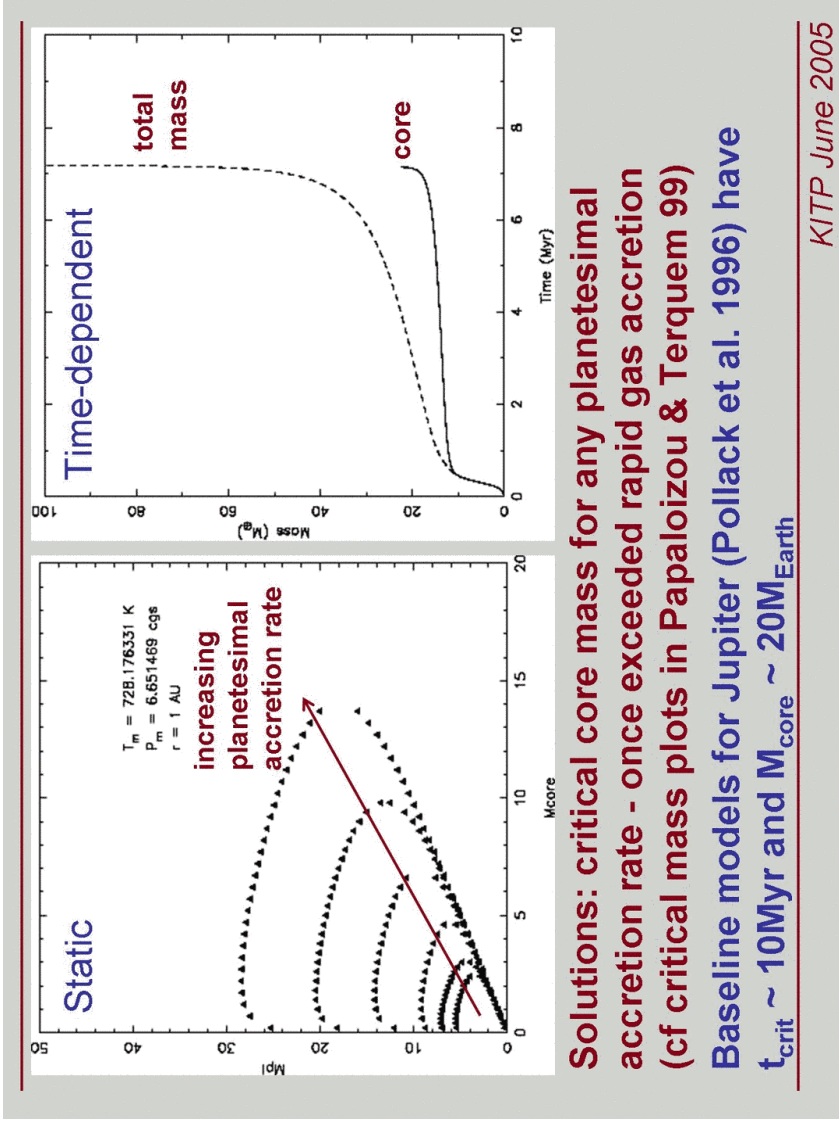
### 4) Giant planet envelopes



**Luminosity: from accretion of planetesimals (E deposited near core) + contraction of gas**

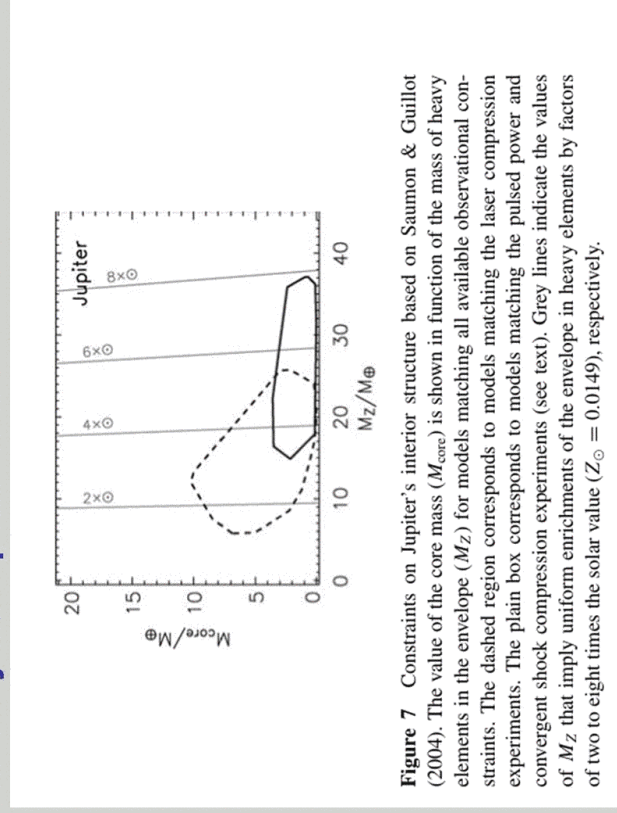
**Growth model: series of quasi-static states at fixed orbital radius. Stellar structure problem with unknown  $M$  but known  $R(M)$ .**

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## Observations

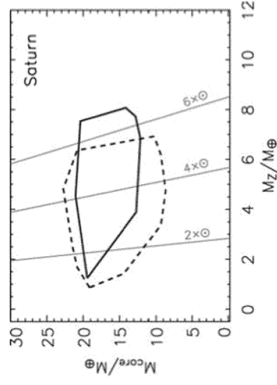
### Solar system problems:



Jupiter's core mass is thought to be less than  $10 M_{\text{Earth}}$  (review Guillot 2005)

Reducing the planetesimal accretion rate 'solves' this, but increases the formation time

**Saturn's core is OK**



**Figure 8** Same as Figure 7 in the case of Saturn. Note that smaller core masses could result either from allowing a variation of the abundance of heavy elements near the molecular/metallic transition (Guillot 1999a) or from the presence of a helium shell around the core (Fortney & Hubbard 2003).

**Formation time scale of Uranus & Neptune?**

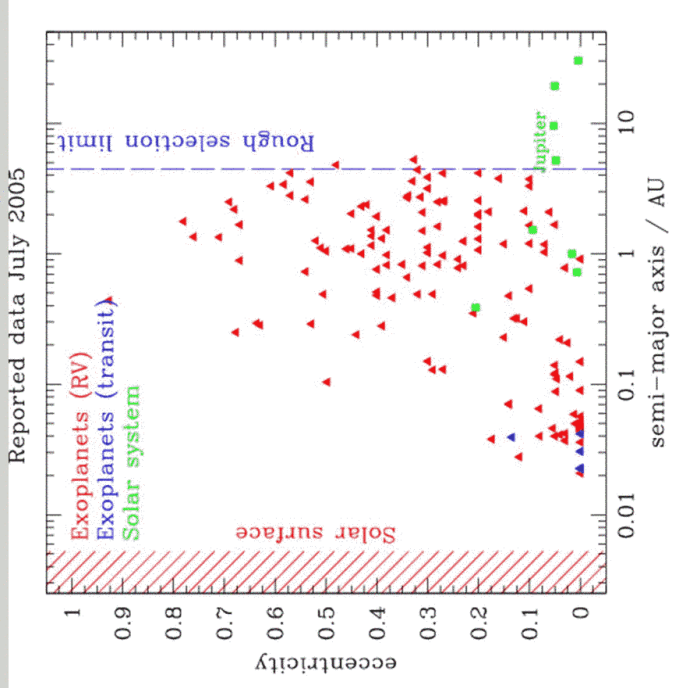
- migration (planetesimal scattering) or enhanced gravitational focusing

**Edge to the Kuiper belt at a ~ 50AU (Trujillo et al. 01)**

**Sedna: a ~ 500AU**

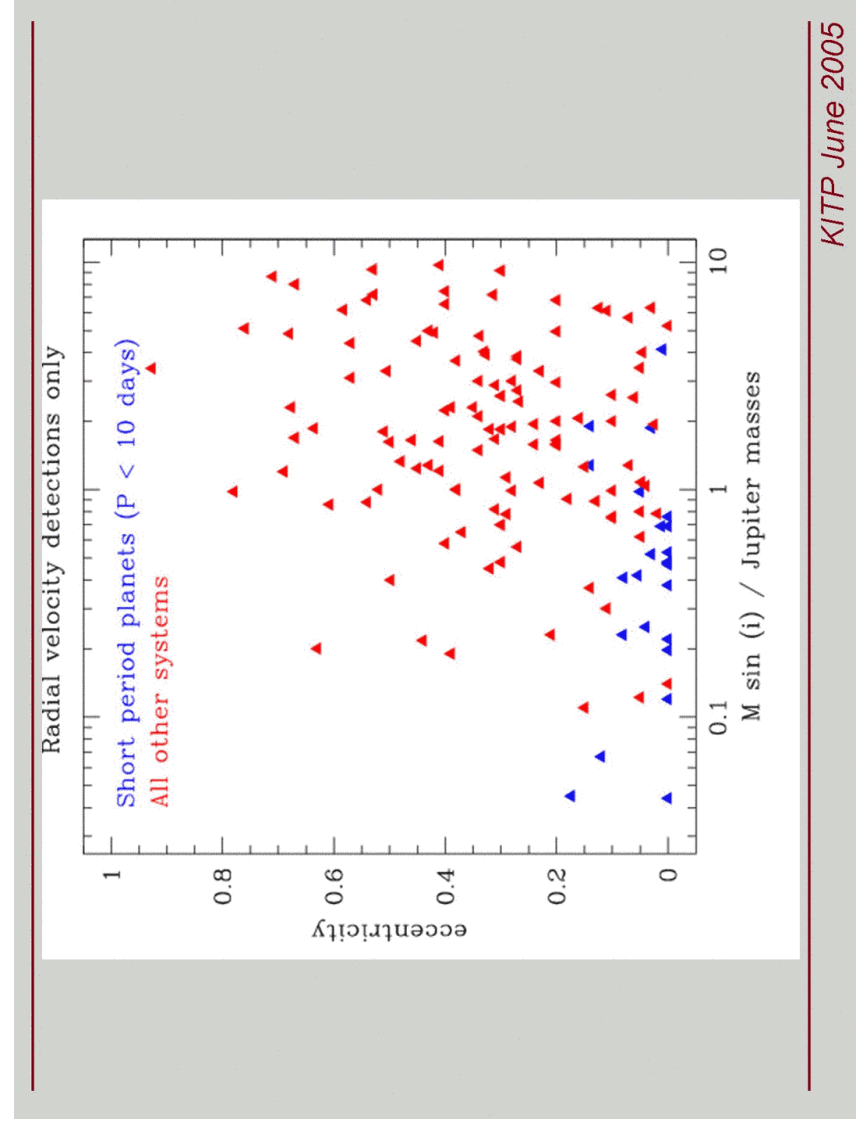
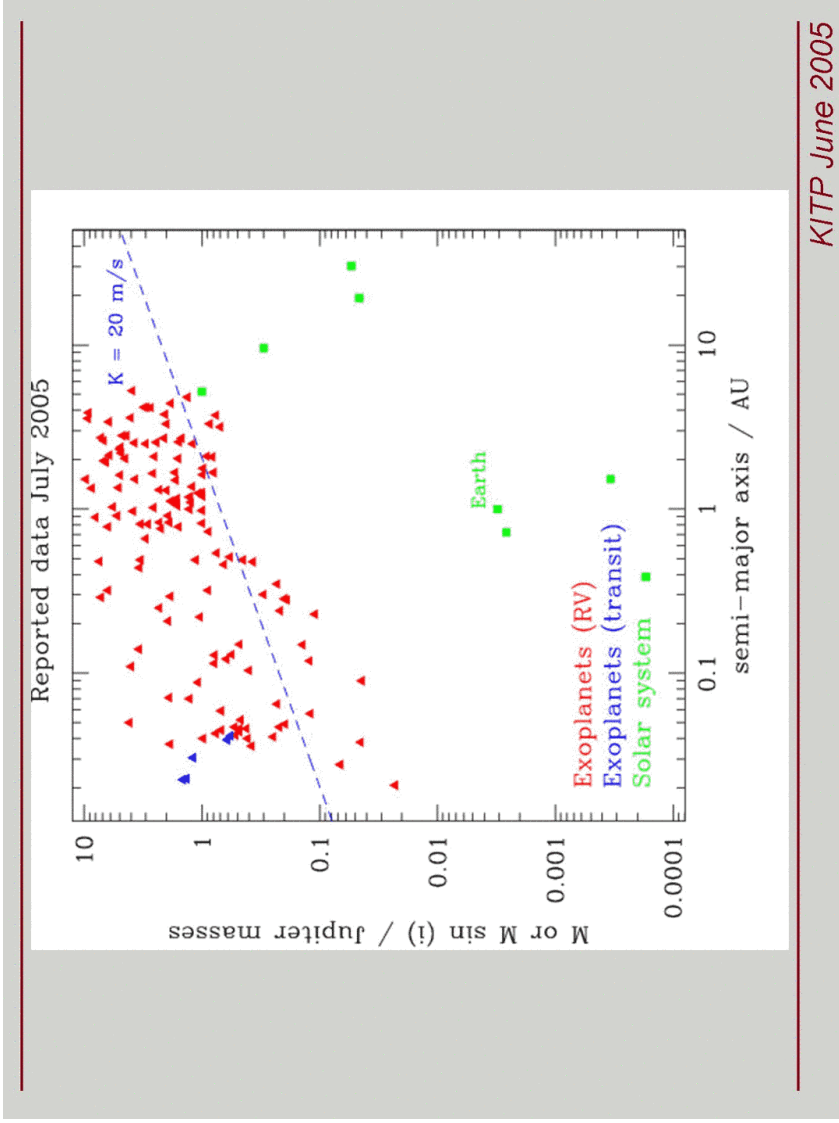
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**Extrasolar planets: know  $M_p \sin(i)$ ,  $a$ ,  $e$ ,  $[Fe/H]_{\text{host}}$  only**

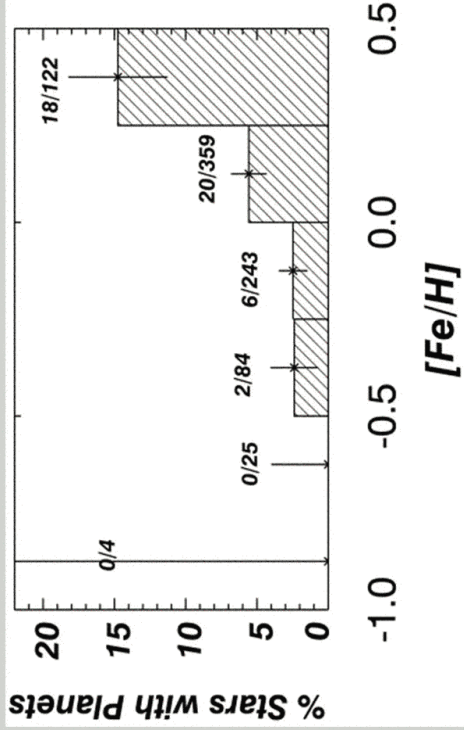


- ▶ **Massive planets in very short period orbits ( $P=1.22$  days)**
- ▶ **Larger population further out, but with ~ flat distribution of eccentricity in  $0 < e < 0.7$**
- ▶ **Raw frequency  $f \sim 5\text{-}10\%$  of ~ Solar type stars**

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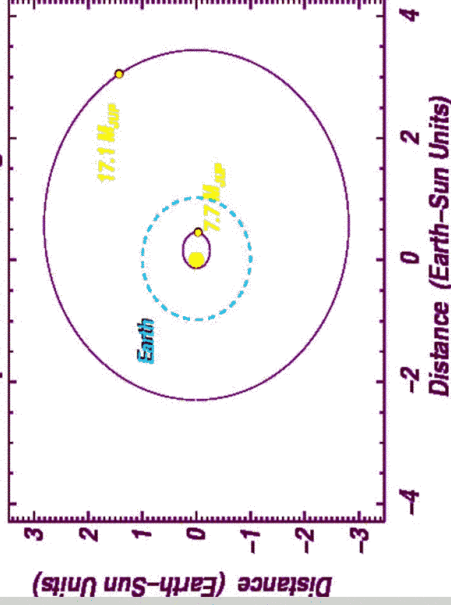
Fischer & Valenti,  
ApJ, 622, 1102  
(2005)



**BUT...  $f_{\text{planet}}([Fe/H])$  is very consistent with expectations of core accretion - most observed extrasolar planets did not form from disk instability or fragmentation**

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Two Companions Orbiting HD168443



+ odd individual systems

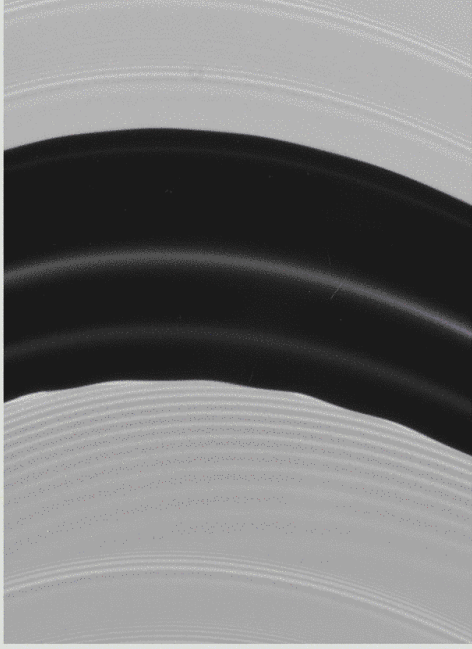
**Metallicity distribution for hosts of the most massive 'planets' is not yet well measured**

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## Migration

Gravitational interaction between planet and disk:

- ▶ Angular momentum exchange
- ▶ Planet migration
- ▶ Perturbation to  $\Sigma$



Effect on the (particle) disk can be seen in Saturn's rings

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2D, *laminar* disk simulation for illustration only!  
 Planet mass grows slowly from  $3 M_{\text{Earth}}$  to  $10 M_{\text{Jupiter}}$   
 Shown in an almost corotating frame

QuickTime™ and a YUV420 codec decompressor are needed to see this picture.

Type I: planet remains embedded in gas ( $M \sim M_{\text{Earth}}$ )

Type II: planet clears a gap ( $M \sim M_{\text{Jupiter}}$ )

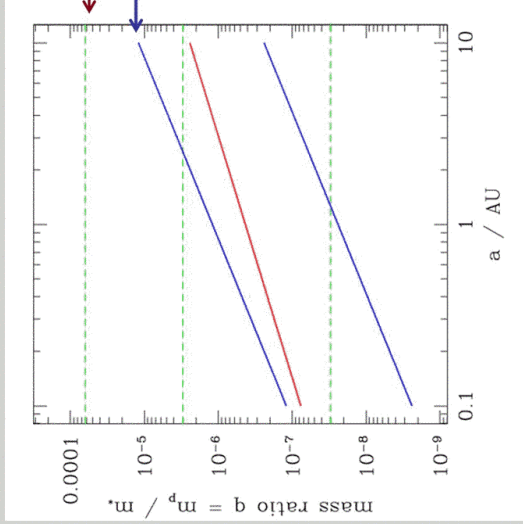
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**Type I migration and core accretion**

$$\frac{v_{\text{migrate}}}{v_{\text{Kepler}}} \approx -3q \frac{r_p^2 \Sigma}{M_*} \left( \frac{h}{r_p} \right)^{-2}$$

Type I estimate (Tanaka, Takeuchi & Ward) is about 1 cm/s per Earth mass (minimum mass Solar nebula)



critical core mass migration time scale of 1 Myr

- 1) Cores do not form in situ
- 2) Hard to reach a critical core mass before core is lost to the star

Depends on background model (Menou & Goodman 04)

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**Low mass migration in turbulent disks:**

Potential fluctuations at  $\sim 10^{-4}$  level yield a random torque on low mass planets

$$T_{\text{turbulent}} \propto M_p$$

$$T_{\text{Type I}} \propto M_p^2$$

Additive?

**Simulations:**

Laughlin, Steinacker &

Adams (2004)

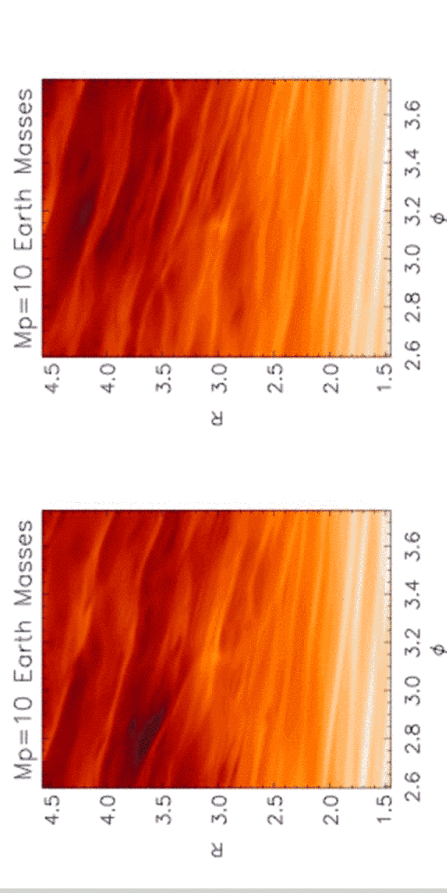
Nelson & Papaloizou

(2004)

$\Sigma$  slice of disk + boundary layer: 3D, ideal MHD, unstratified, isothermal,  $h/r \sim 0.1$

QuickTime™ and a YUV420 codec decompressor are needed to see this picture.

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*Nelson & Papaloizou (2004)*

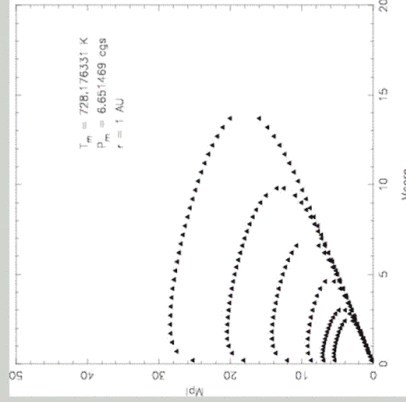
**Critical mass between stochastic / Type I migration:**

$$M_{p,crit} \sim 10M_{\oplus} \quad (\text{NP04; much higher for LSA04})$$

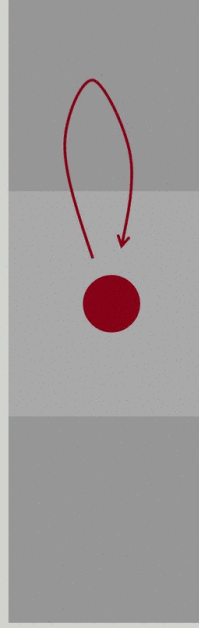
**Ideal MHD, but from simulations that realize quite inefficient transport ( $\alpha \sim 7 \times 10^{-3}$ ,  $h/r = 0.07$ )**

**Does not help survival, may alter core accretion**

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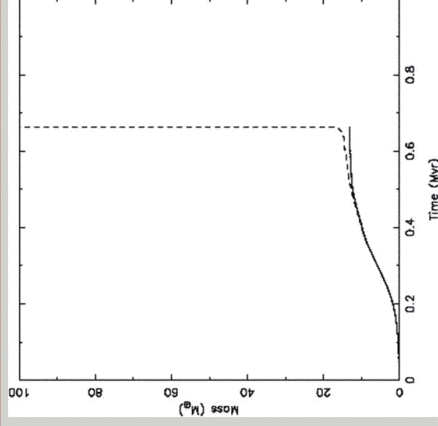
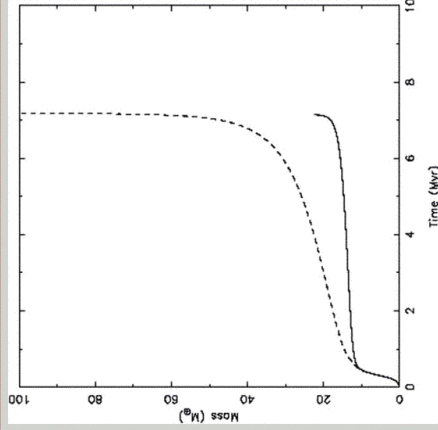
**Highly variable planetesimal accretion rate reduces the typical critical core mass**



**Grow in undepleted zones, but runaway revisiting depleted zones**

**Alternatively: multiple core formation could lead to competition for planetesimals**

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Rice & Armitage 2003

**Dominant random walk component:**

- ▶ Large reduction in time to runaway
- ▶ Can trade some of that off for smaller core mass

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**Type II migration uncertainties:**

- ▶ Disk model
- ▶ Numerical convergence at gap opening
- ▶ Turbulence

Does the interaction lead to eccentricity growth?

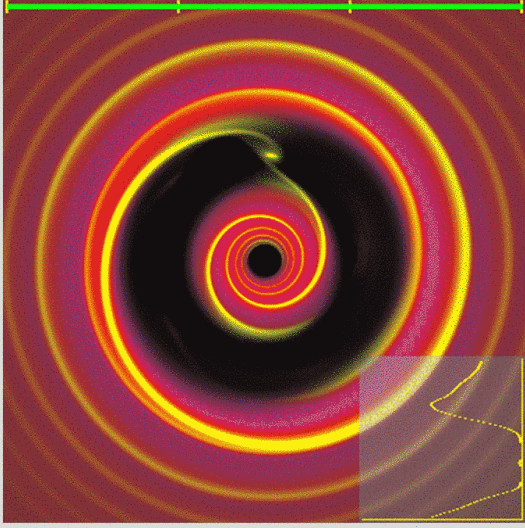
QuickTime™ and a YUV420 codec decompressor are needed to see this picture.

Winters, Balbus & Hawley '03

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**Type II migration is basis for a model for Hot Jupiters:**

- ▶ Planets form at `large' radii ( $a > 4$  AU)
- ▶ Migrate inward within gaps
- ▶ Lost to the star or stall as gas is lost



Potentially quite wasteful  
(depending on the disk  
evolution model)

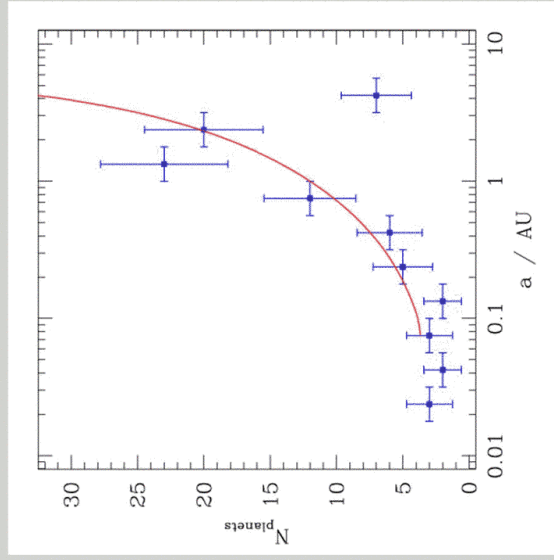
QuickTime™ and a  
GIF decompressor  
are needed to see this picture.

Armitage et al. 02, Trilling  
et al. 02

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**Predicted distribution  
of stalled planets  
from 1D disk + planet  
model**

$$\frac{\partial \Sigma}{\partial a} = \frac{1}{r} \frac{\partial}{\partial r} \left[ 3r^{1/2} \frac{\partial}{\partial r} (\nu \Sigma r^{1/2}) \right] - \frac{r^{1/2}}{\pi \sqrt{GM_*}} \frac{\partial T}{\partial r}$$



**Consistent with observed  
distribution of extrasolar  
planets (selection effects  
corrected)**

**Predicts abundance of  
planets embedded in  
gas disks (evidence for  
GM Aur, CoKu Tau/4)**

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## Summary

- ▶ Hot Jupiters / Type II migration identification is reasonably robust
- ▶ Eccentricity of many extrasolar planets remains a problem... though scattering seems to work
- ▶ Type I migration is overwhelmed by turbulent fluctuations at low masses (terrestrial planet formation?)
- ▶ If random walk migration persists to core masses, may explain Jupiter's small core

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