



HVP contribution to the muon anomalous magnetic moment from lattice QCD

Christine Davies
University of Glasgow
HPQCD collaboration

Santa Barbara
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Work with: Bipasha Chakraborty, Gordon Donald,
Rachel Dowdall, Jonna Koponen, Peter Lepage



Using the Darwin (9600 core) Sandybridge/infiniband
cluster at Cambridge, part of STFC's DiRAC HPC
facility

Muon anomalous magnetic moment

$$\vec{\mu} = g \frac{e}{2m} \vec{S} \qquad a_{\mu} = \frac{g - 2}{2}$$

Measure using polarised muons circulating in E and B fields. At a momentum where $\beta \times E$ terms cancel, difference between spin and cyclotron frequencies:

$$\omega_a = -\frac{e}{m} a_{\mu} B$$

BNL result:

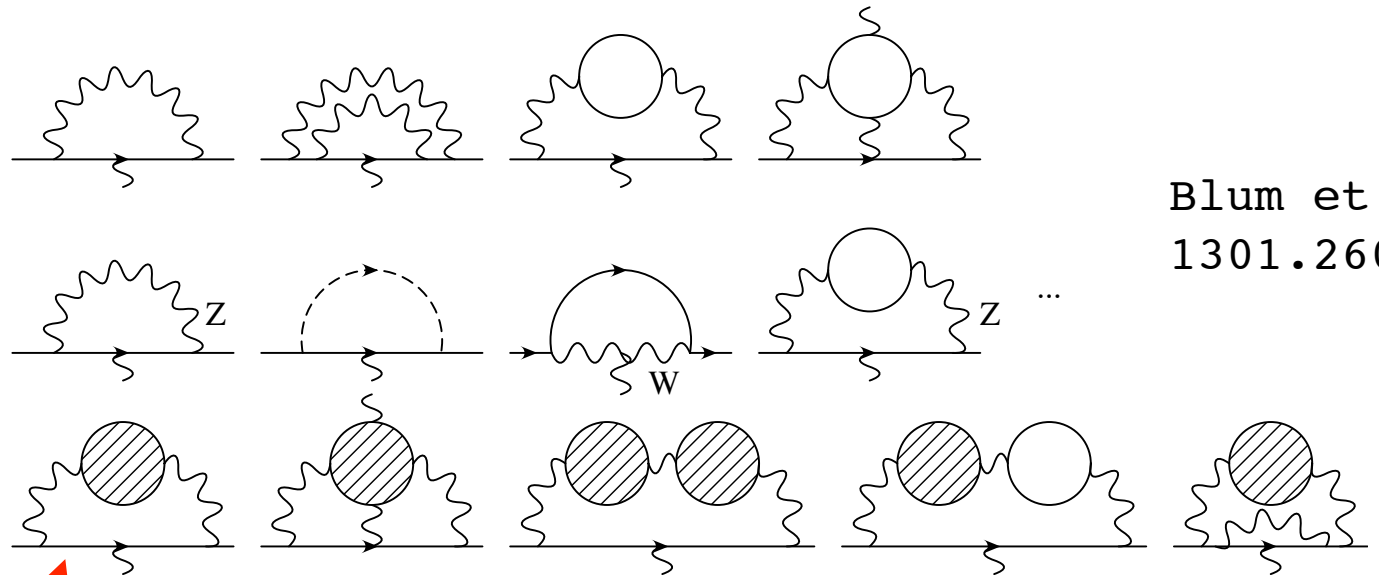
$$a_{\mu}^{expt} = 11659208.9(6.3) \times 10^{-10}$$

E989 (FNAL) will
reduce exptl uncty to
1.6, starting 2017



Standard Model theory expectations

Contributions from QED, EW and QCD interactions.
QED dominates.



Blum et al,
1301.2607

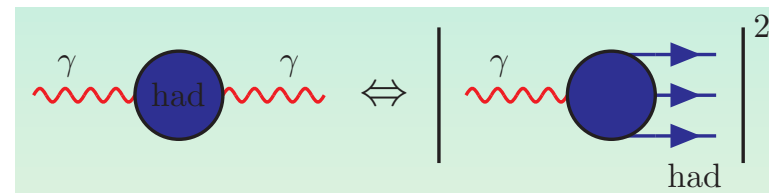
LO Hadronic vacuum polarisation (HVP)

$$\begin{aligned}
 a_{\mu}^{expt} - a_{\mu}^{QED} - a_{\mu}^{EW} &= 721.7(6.3) \times 10^{-10} \\
 &= a_{\mu}^{HVP} + a_{\mu}^{HOHVP} + a_{\mu}^{HLBL} + a_{\mu}^{new\ physics}
 \end{aligned}$$

$$a_{\mu}^{HVP, no\ new\ physics} = 721.0(6.8) \times 10^{-10}$$

Best method to date for HVP uses exptl e⁺e⁻ cross-section

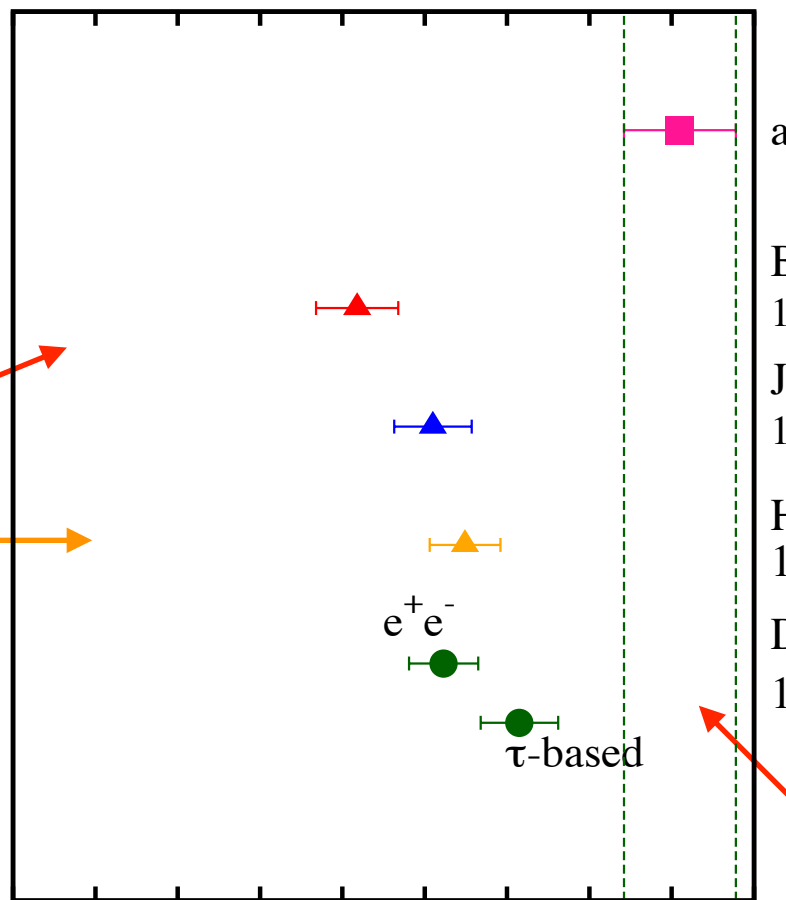
$$a_{\mu}^{HVP} = \frac{1}{4\pi^3} \int_{m_{\pi}^2}^{\infty} ds \sigma_{had}^0(s) K(s)$$



$$e^+e^- \rightarrow \gamma^* \rightarrow hadrons$$

“bare” cross-section
but inc. final-state
radiation

some “tension”
between results.
Difference is
use of BaBar radiative
return data



a_{μ}^{HVP} , no new physics

Benayoun et al
1210.7184, $e^+e^- + \tau$

Jegerlehner+Szafron
1101.2872, $e^+e^- + \tau$

Hagiwara et al
1105.3149, e^+e^- only

Davier et al
1010.4180

e^+e^-

τ -based

SM 3σ
below nnp

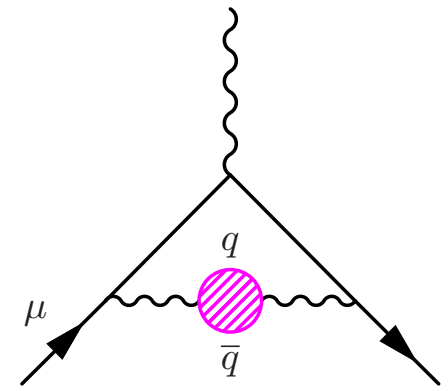
Hagiwara et al:

$$a_{\mu}^{HVP} = 694.9(4.3) \times 10^{-10}$$

Lattice calculation of HVP

Analytically continue to Euclidean q^2 .

$$a_{\mu}^{HVP,i} = \frac{\alpha}{\pi} \int_0^{\infty} dq^2 f(q^2) (4\pi\alpha e_i^2) \hat{\Pi}_i(q^2)$$



Blum, hep-lat/
0212018

connected contribution for
flavour i

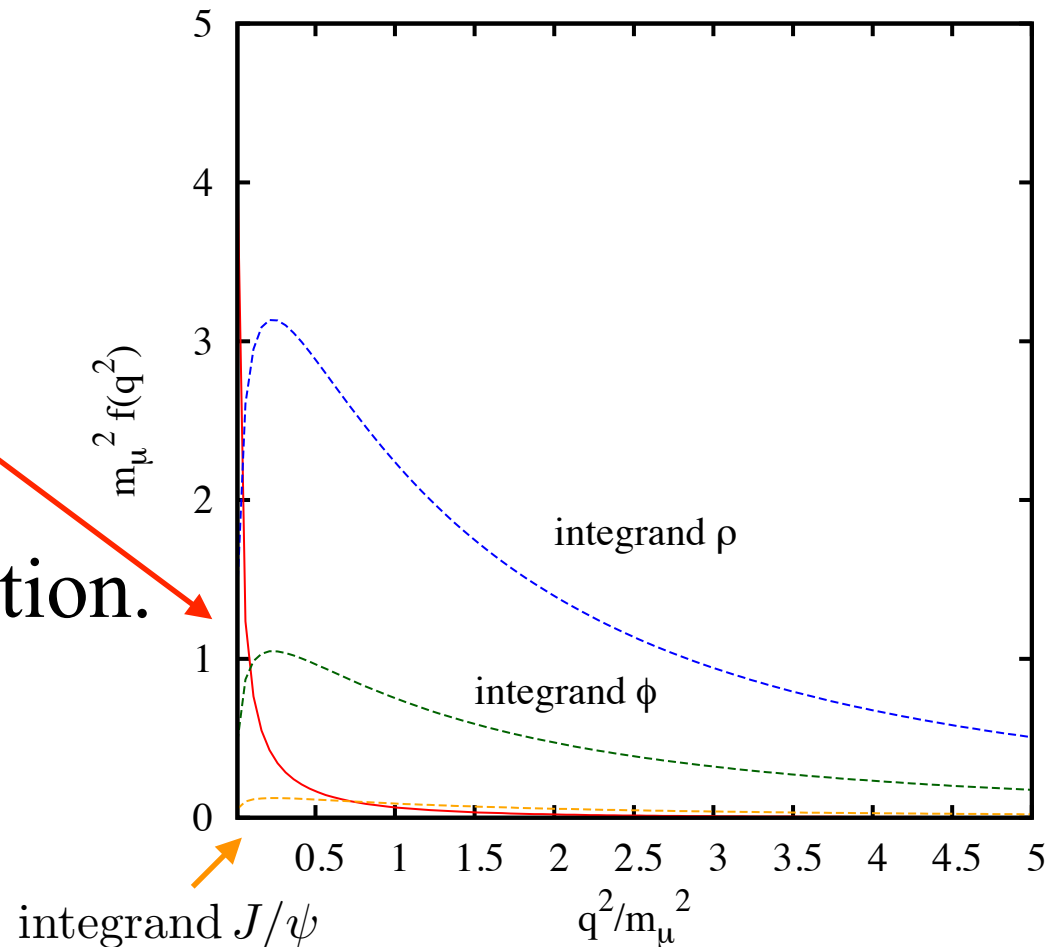
$f(q^2)$ is divergent function
with scale set by m_{μ}

$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

is vacuum polarisation function.

Test with mesons:

$$\hat{\Pi}(q^2) \propto \frac{1}{m_V^2} - \frac{1}{q^2 + m_V^2}$$



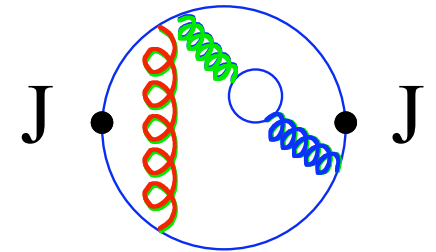
For spatial vector currents at zero spatial momentum

$$\Pi^{jj}(q^2) = q^2 \Pi(q^2) = a^4 \sum_t e^{iqt} \sum_{\vec{x}} \langle j^j(\vec{x}, t) j^j(0) \rangle$$

Time-moments of lattice current-current correlators

$$G_{2n} \equiv a^4 \sum_t \sum_{\vec{x}} t^{2n} Z_V^2 \langle j^j(\vec{x}, t) j^j(0) \rangle$$

$$= (-1)^n \left. \frac{\partial^{2n}}{\partial q^{2n}} q^2 \hat{\Pi}(q^2) \right|_{q^2=0}$$

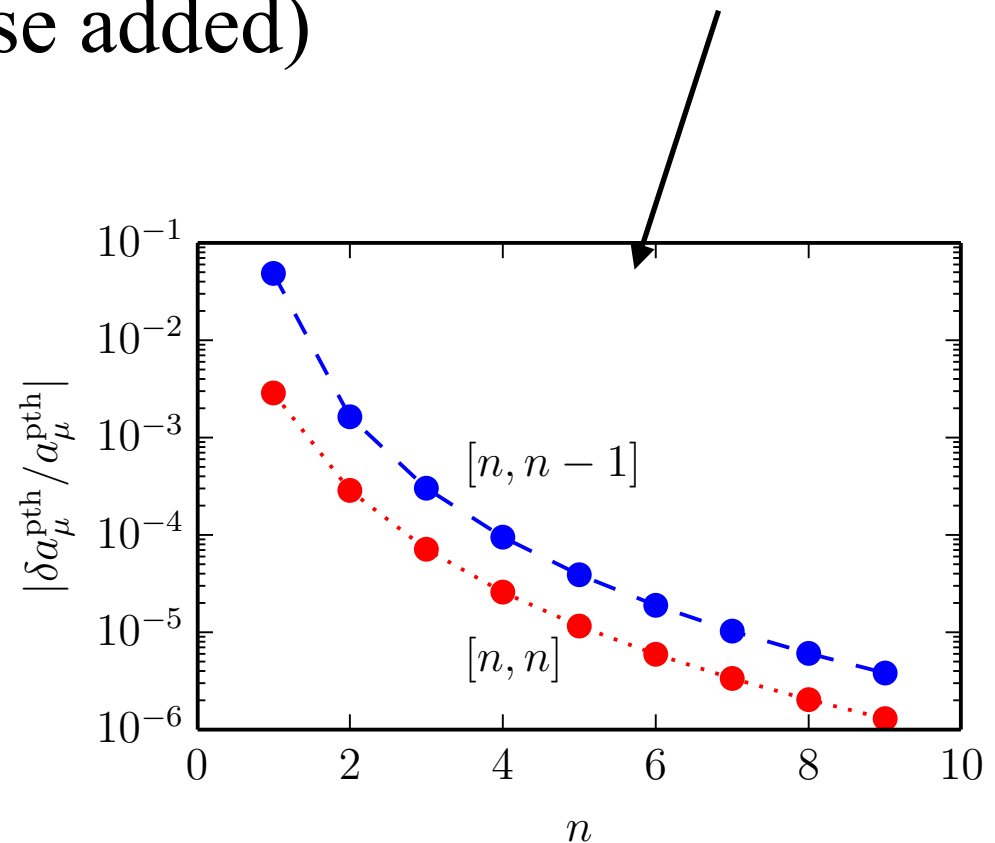
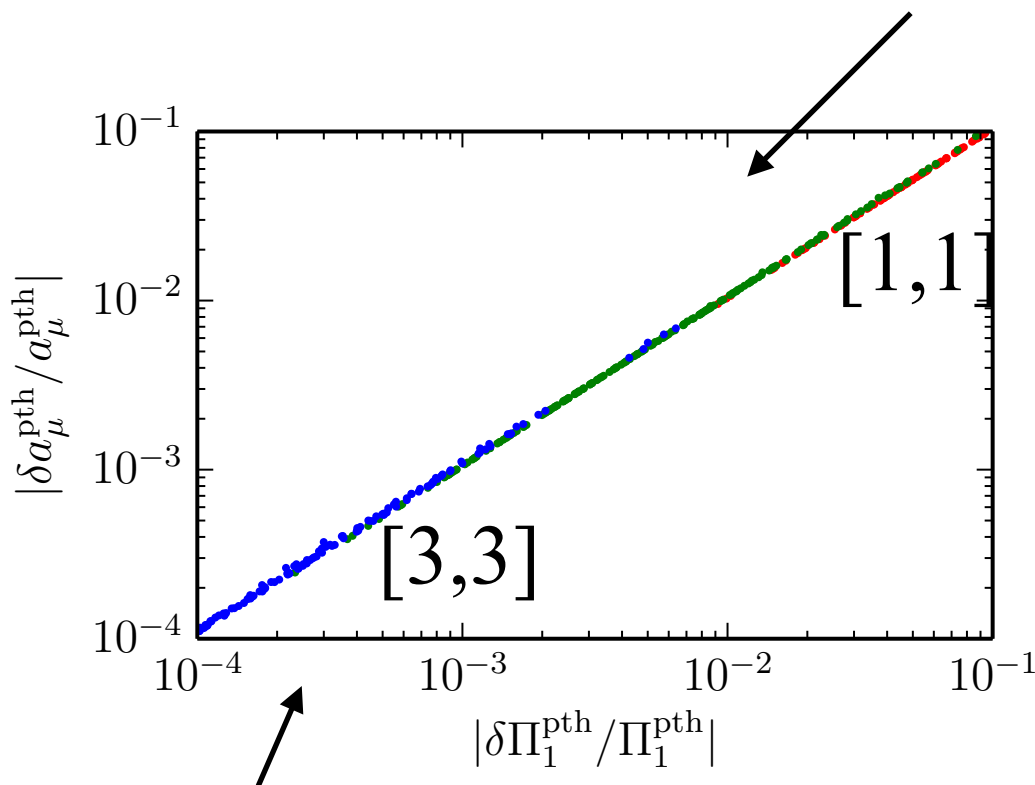


$$\hat{\Pi}(q^2) = \sum_{j=1}^{\infty} q^{2j} \Pi_j \quad \text{with} \quad \Pi_j = (-1)^{j+1} \frac{G_{2j+2}}{(2j+2)!}$$

Allows us to reconstruct $\hat{\Pi}(q^2)$ and integrate

Use Pade approximants (ratio of m/n polynomials) rather than Taylor expansion for better large q^2 behaviour.

Test Pade approximants in similar scenarios (1-loop quark vacuum polarisation, with noise added)



Improved precision allows higher order Pade - we use [2,2]

CHARM contribution

HPQCD 1004.4285,
1208.2855

Part of the set of calculations that gave

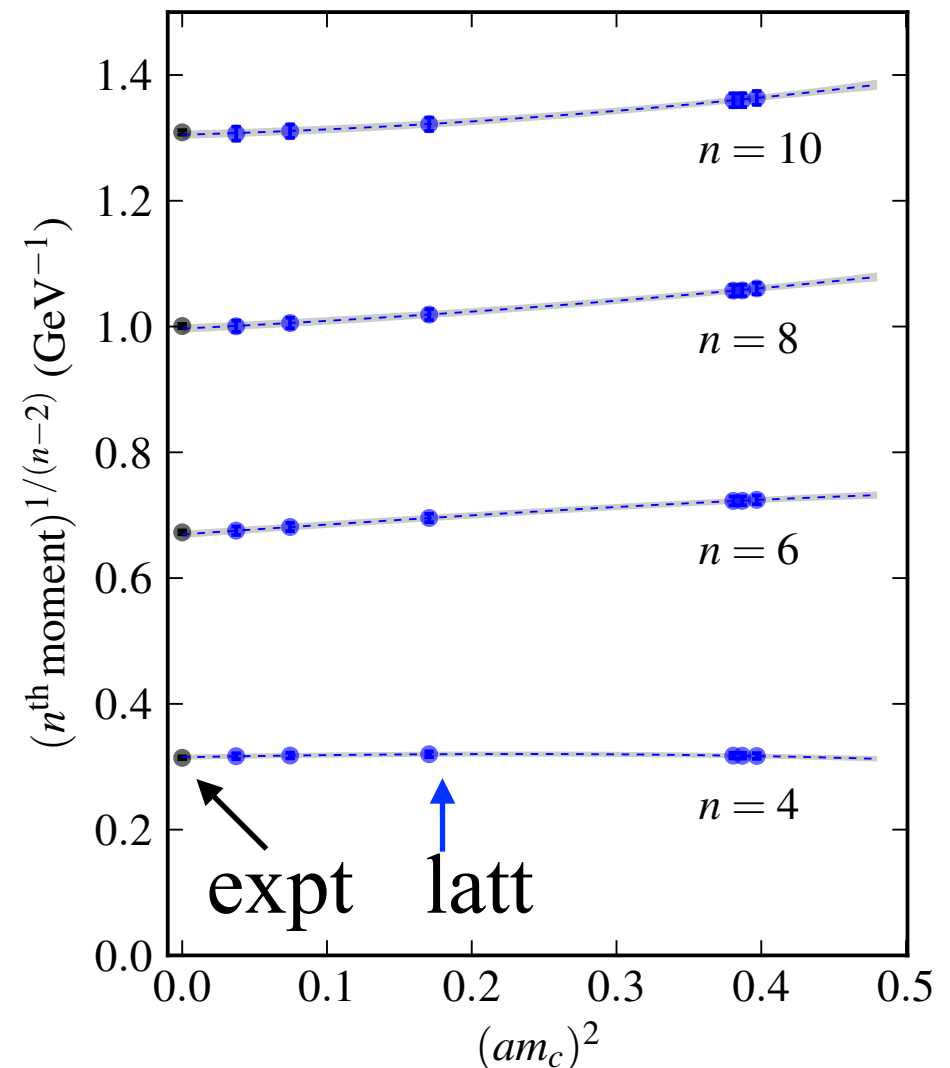
$$m_c, M(J/\psi) - M(\eta_c), \Gamma(J/\psi \rightarrow e^+e^-), \Gamma(J/\psi \rightarrow \eta_c\gamma)$$

Used HISQ valence quarks on MILC 2+1 asqtad configs. Z_V from contnm QCD pert. th.

Extrapolation to physical point allows us to compare directly to moments from e^+e^- expt. in charm region

$$a_\mu^{HVP,c} = 14.4(4) \times 10^{-10}$$

HPQCD 1403.1778



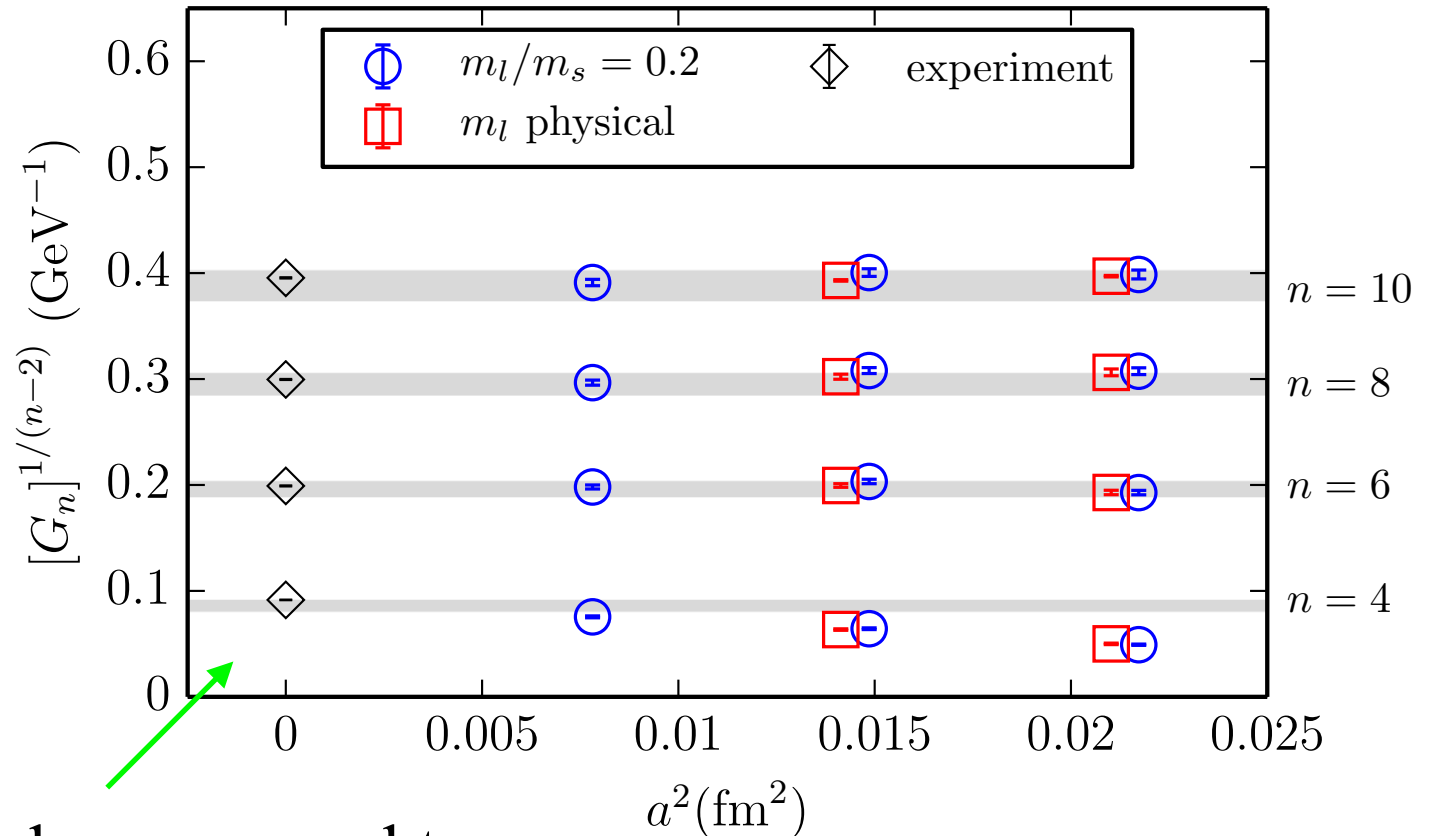
BOTTOM contribution

HPQCD 1110.6887,
1309.5797,
1408.5768

Part of the set of calculations that gave

$$m_b, M(\Upsilon) - M(\eta_b), M(\Upsilon') - M(\eta'_b), \Gamma(\Upsilon \rightarrow e^+e^-), \Gamma(\Upsilon' \rightarrow e^+e^-)$$

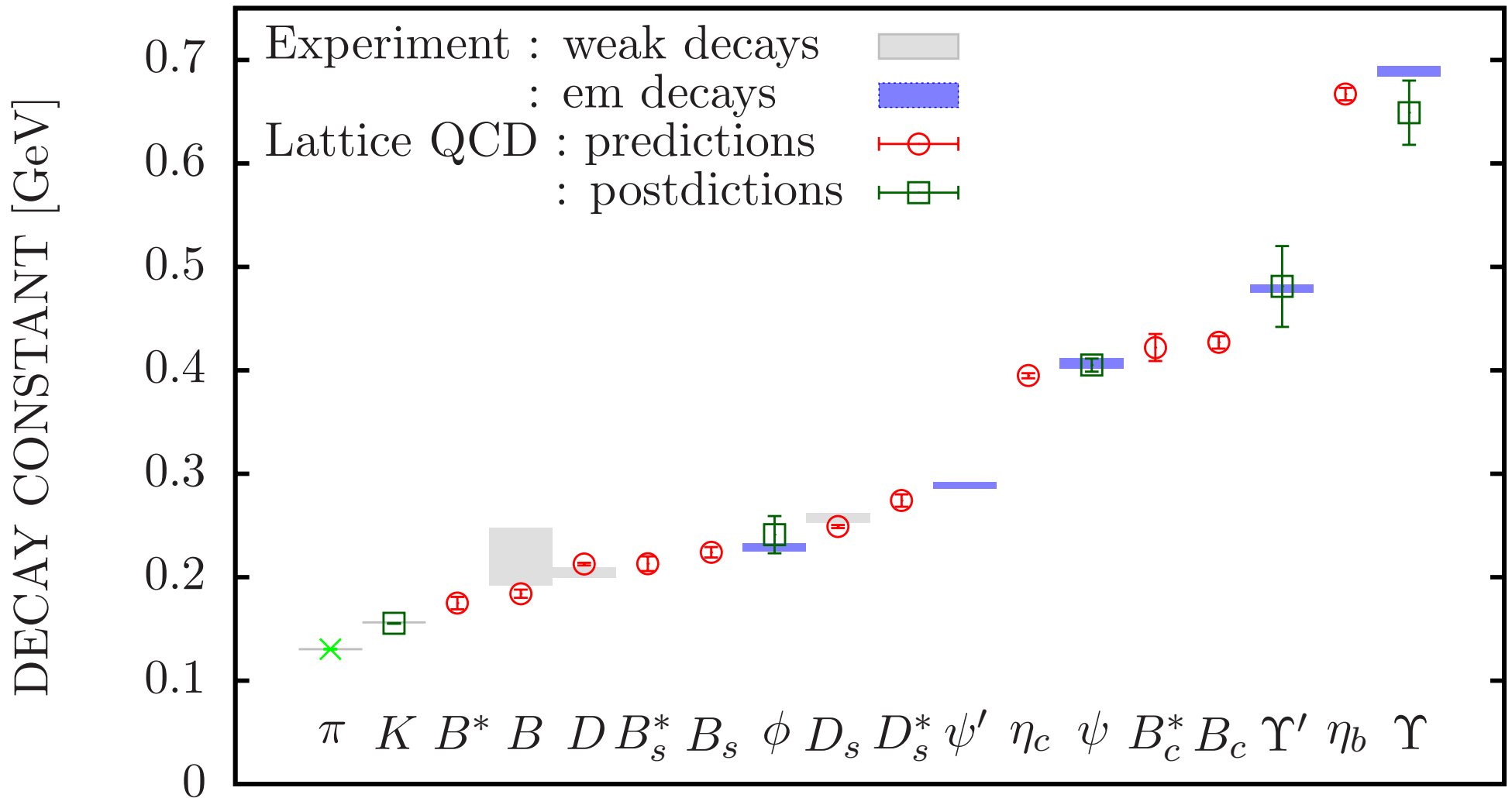
Used NRQCD
valence
quarks on
MILC 2+1+1
HISQ configs.
 Z_V from
contnm QCD
pert. th.



Again, moments can be compared to those extracted from expt.

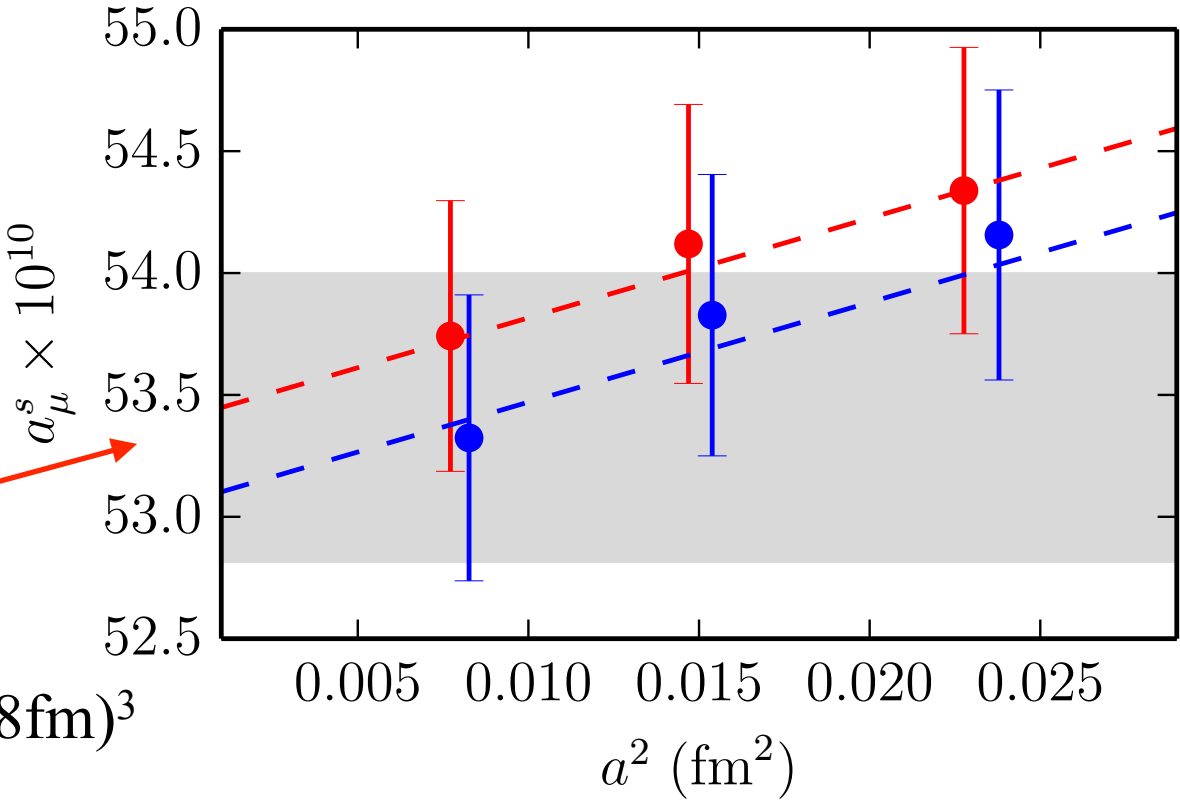
$$a_\mu^{HVP,b} = 0.27(4) \times 10^{-10}$$

Keep an eye on the ‘big’ picture whilst doing this



STRANGE contribution

HISQ valence quarks on MILC 2+1+1 HISQ configs. Local J_V - nonpert. Z_V .
 multiple a (fixed by w_0), m_l (inc. phys.), volumes. Tune s from η_s up to $(5.8\text{fm})^3$

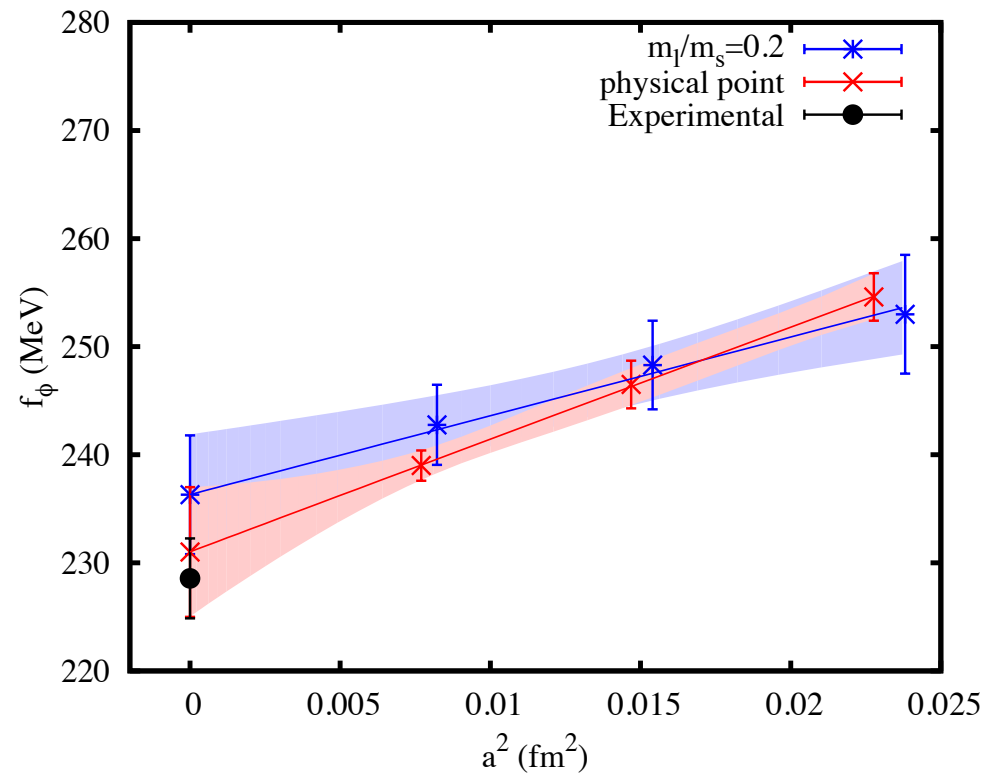
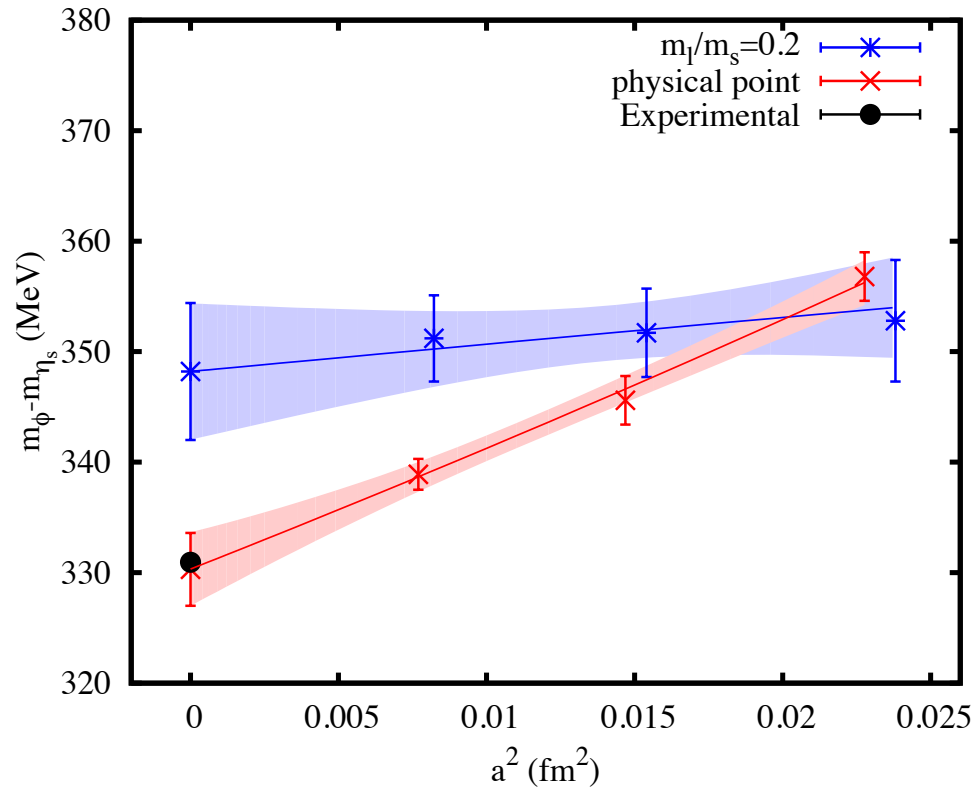


	a_μ^s
Uncertainty in lattice spacing (w_0, r_1):	1.0%
Uncertainty in Z_V :	0.4%
Monte Carlo statistics:	0.1%
$a^2 \rightarrow 0$ extrapolation:	0.1%
QED corrections:	0.1%
Quark mass tuning:	0.0%
Finite lattice volume:	< 0.1%
Padé approximants:	< 0.1%
Total:	1.1%

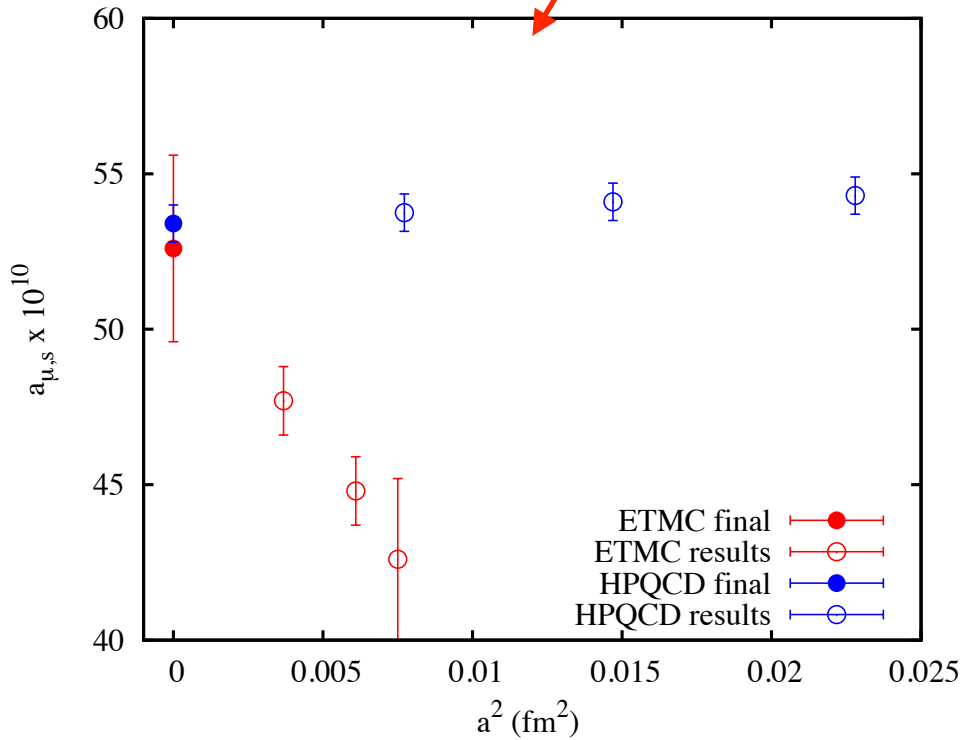
$$a_{\mu, \text{lat}}^s = a_\mu^s \times (1 + c_{a^2}(a\Lambda_{\text{QCD}}/\pi)^2 + c_{\text{sea}}\delta x_{\text{sea}} + c_{\text{val}}\delta x_{\text{val}})$$

$$a_\mu^{\text{HVP},s} = 53.41(59) \times 10^{-10}$$

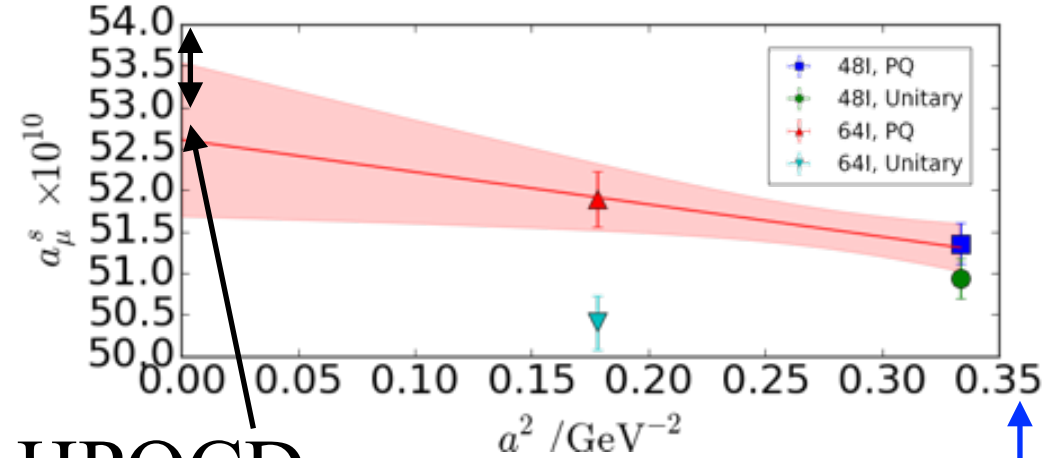
Check mass and decay constant of ϕ from these correlators against expt



Results from ETMC, RBC/UKQCD



continuum limit for strange contribution to a_μ



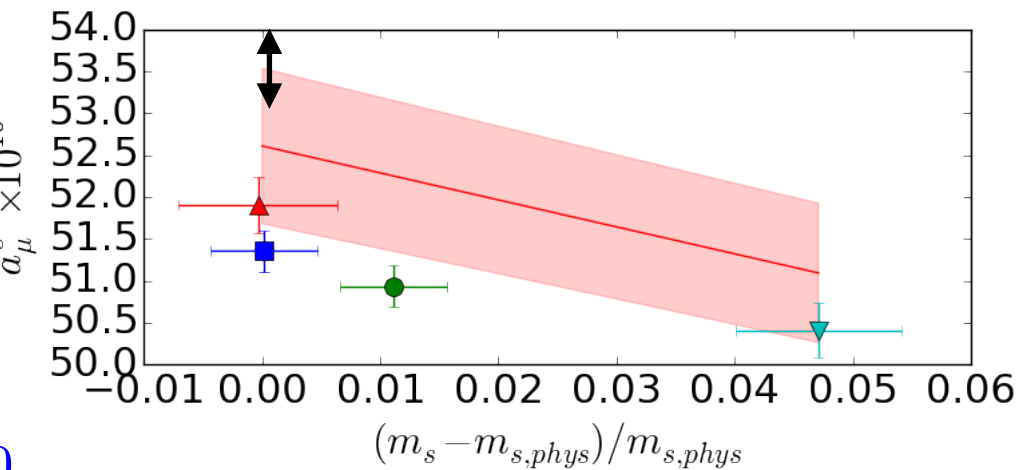
HPQCD

0.014 fm^2

$$a_\mu^{HVP,s,ETMC} = 53(3) \times 10^{-10}$$

Continuum estimate/
upper limit:

$$a_\mu^{HVP,s,cont} = 55.3(8) \times 10^{-10}$$



1403.1778

LIGHT contribution $m_u = m_d$

HISQ valence quarks on MILC 2+1+1 HISQ configs. Use Z_v from s calc.

Multiple a (use w_0), m_1 (inc. phys.), volumes (at $m_l/m_s=0.1$).

New ingredient since correlators much noisier. Use:

$$G(t) = \begin{cases} G_{\text{data}}(t) & \text{for } t \leq t^* & \leftarrow \text{from Monte Carlo} \\ G_{\text{fit}}(t) & \text{for } t > t^* & \leftarrow \text{from multi-exponential fit} \end{cases}$$

$t^* = 1.5\text{fm} = 6/m_\rho$ so 70% of result from G_{data}

- 80% of result comes from ρ meson pole, so need to understand ρ on lattice, inc. finite-volume from $\pi\pi$.
- 10% from $\pi\pi$, sensitive to finite-volume and m_π (so taste-issues for staggered quarks).

One approach is to correct Taylor coefficients

$$\hat{\Pi}_j^{latt} \rightarrow (\hat{\Pi}_j^{latt} - \hat{\Pi}_j^{latt}(\pi\pi)) \left[\frac{m_\rho^{2j+2}}{f_\rho^2} \right]_{latt} \left[\frac{f_\rho^2}{m_\rho^{2j+2}} \right]_{expt} + \hat{\Pi}_j^{cont}(\pi\pi)$$

Remove lattice $\pi\pi$
using one-loop,
staggered quark
chiral pert. theory

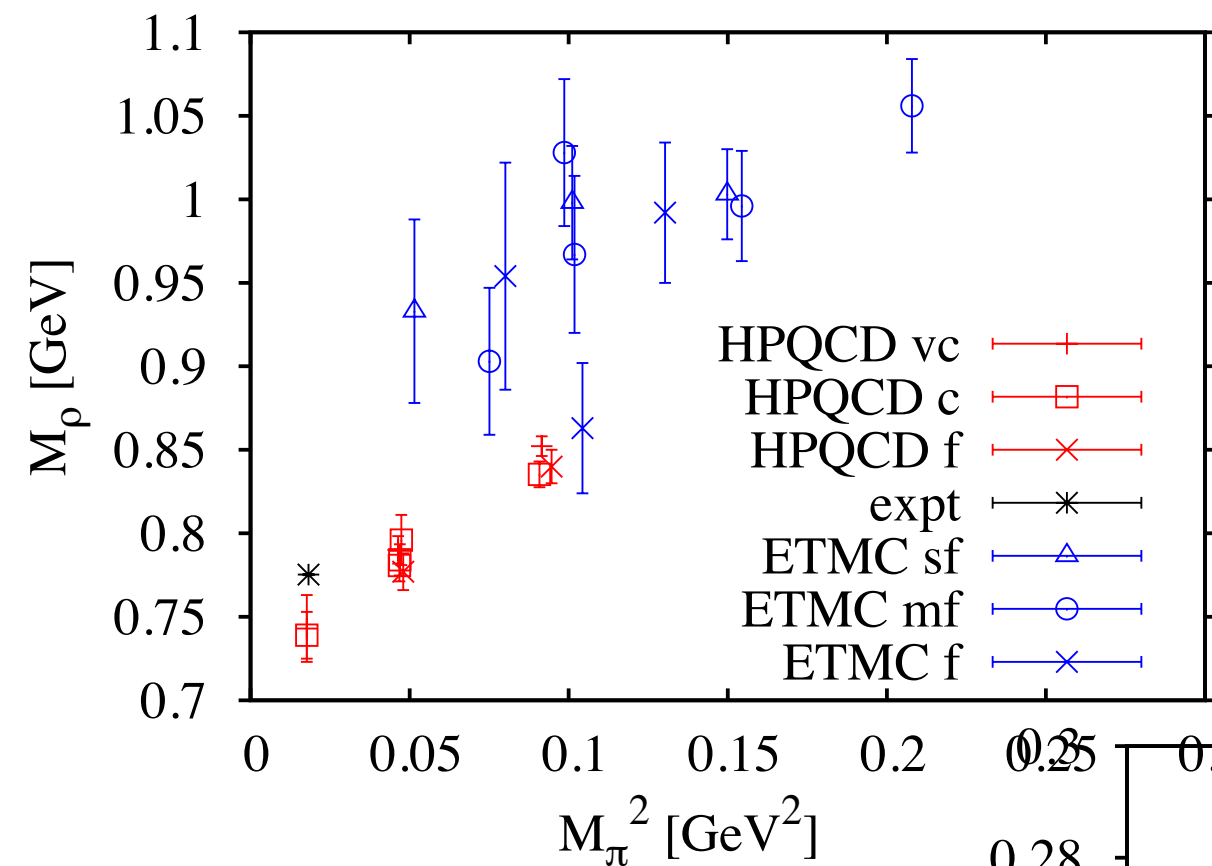
Rescale using exptl f_ρ
and m_ρ
elaborating on
ETMC : 1308.4327.
Reduces lattice
systematics, but
intro. uncties from f_ρ

Restore $\pi\pi$
from continuum
chiral pert. th.

$\pi\pi$ contribution distorted at physical point using staggered quarks on these coarse lattices. Important to inc. other masses. But note: need 7fm lattice to get this piece below 1% for contnm $\pi\pi$

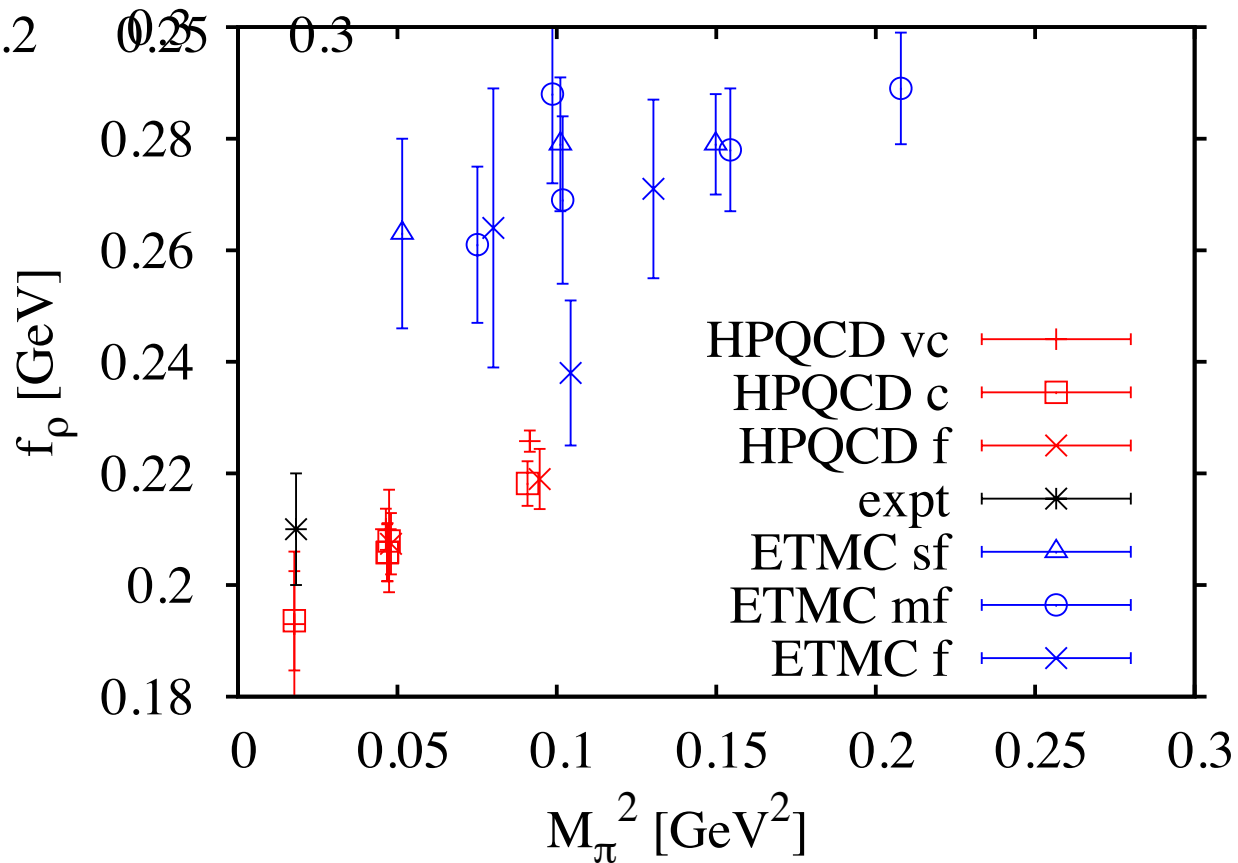
Analysis of ρ parameters

Direct comparison with ETMC possible

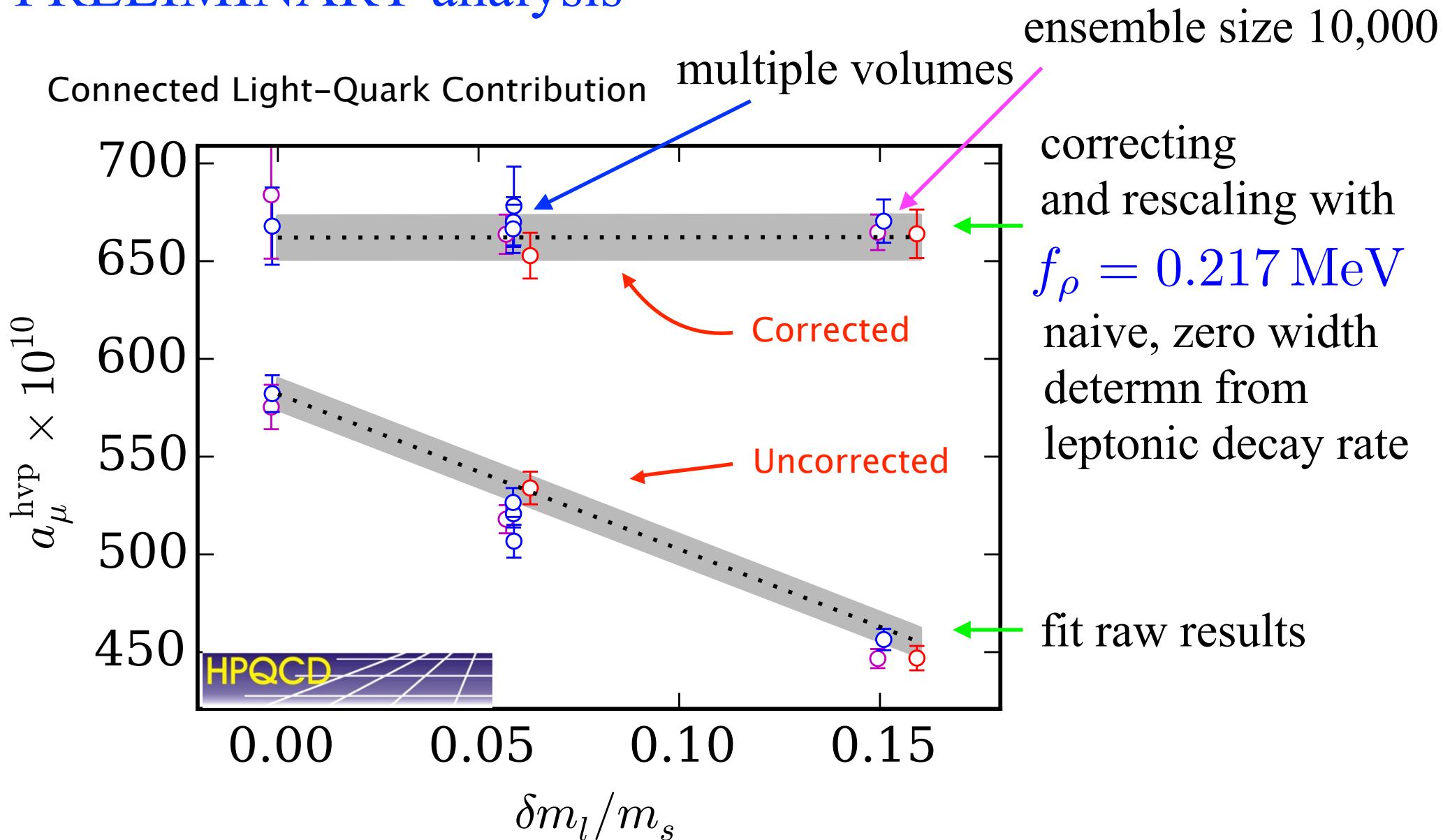


ETMC a 0.06-0.08fm
L 2.5- 2.9fm

HPQCD a 0.09-0.15fm
L 2.5-5.8fm



PRELIMINARY analysis



Improve by estimating systematics with an effective theory that includes $\rho, \gamma, \pi\pi$ e.g. Jegerlehner +Szafron, 1101.2872

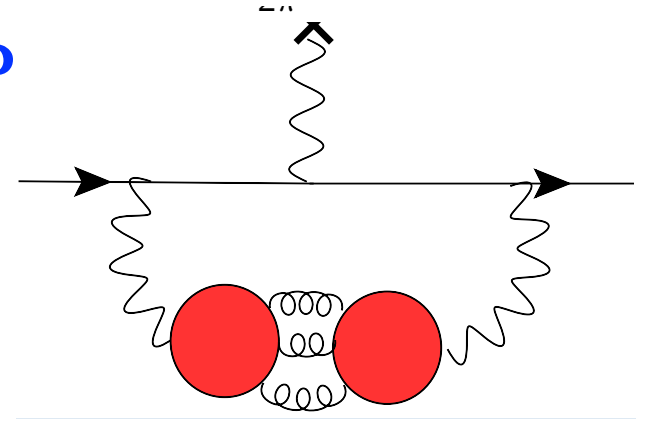
Future: improve statistics, finer lattices

DISCONNECTED contribution to HVP

Vanishes if $m_u = m_d = m_s$

since
$$\sum_i e_i = 0$$

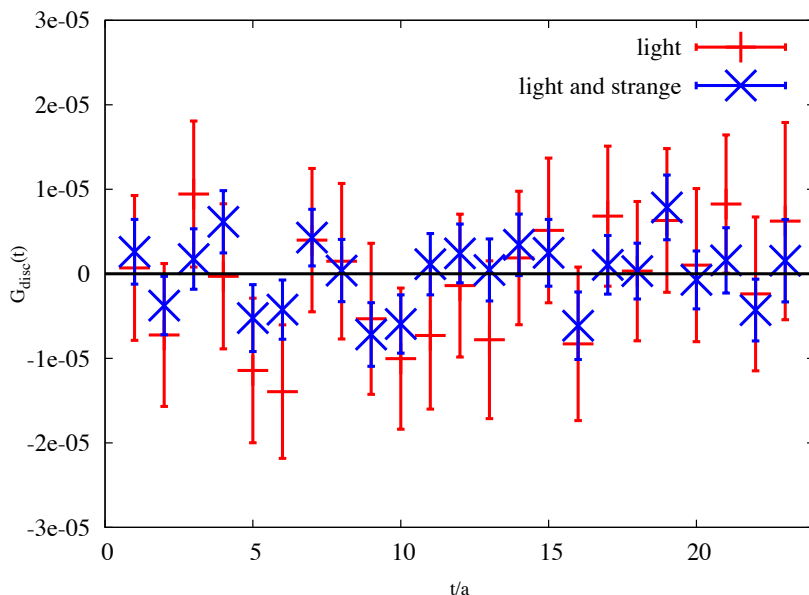
Blum, hep-lat/
0212018



For real masses, result is disconn. correlator for (1-s) current with charge 1/3 (so e^2 factor is 1/5 of connected)

Focus has been on stochastic methods. Using same source for 1 and s helps

Guelpers, Mainz, LAT14



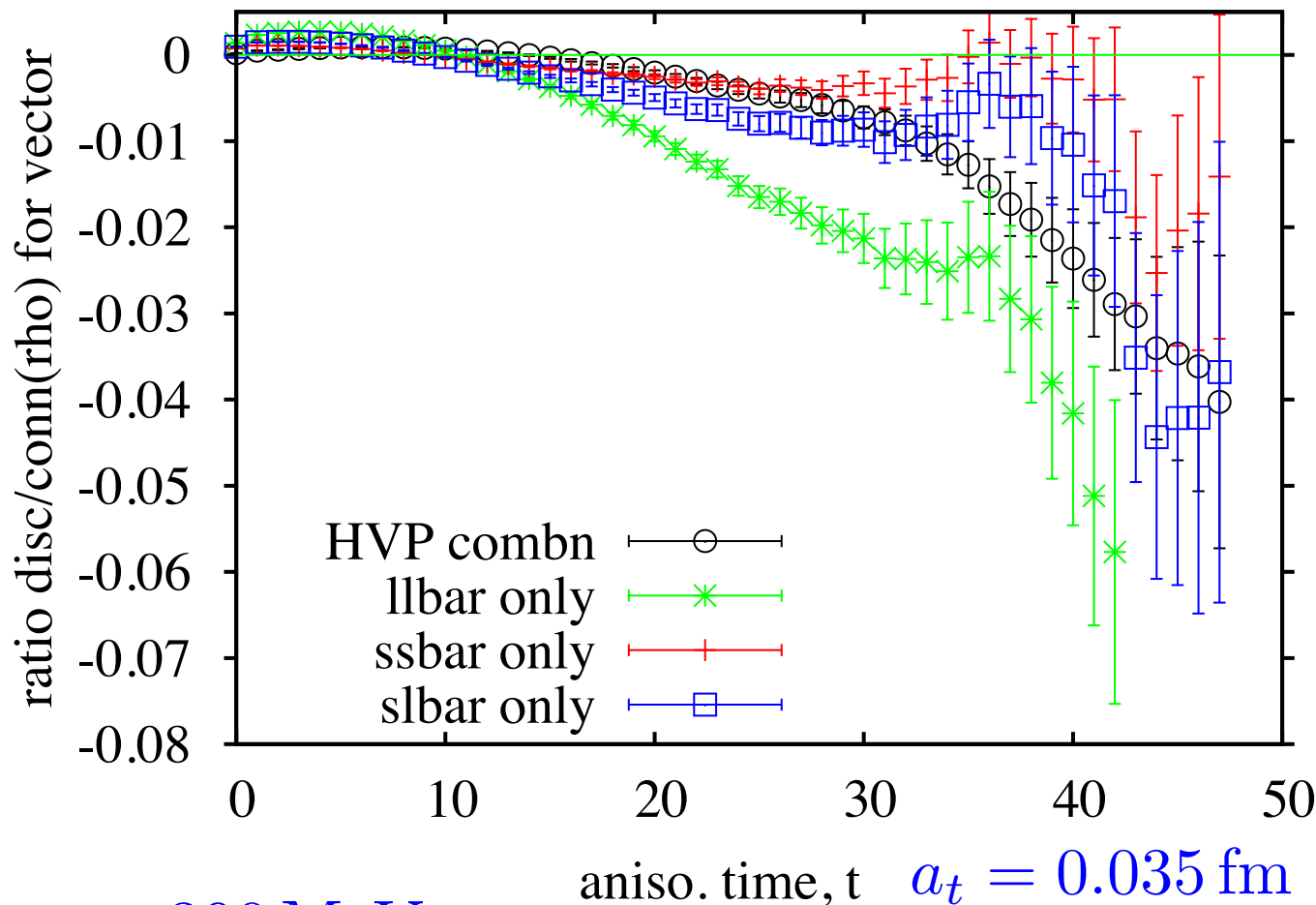
Chakraborty, HPQCD, LAT15

no signal from 50+50 sources (all-to-all) per config. One-link J

HadSpec results

e.g. Hadspec,1309.2608

Use instead many (~ 150) source vectors (eigenvectors of gauge-covariant Laplacian) for both conn. and disc. correlators to obtain good signal.



PRELIMINARY

Fitting and normalising to connected light, gives HVP disc. contribn of $\sim -1\%$

Hadspec+HPQCD, in prep.

anisotropic
clover action

$m_\pi = 390$ MeV

Simple (but conservative) argument on size of disc. pieces
 1-1 disc.pieces provide key difference between ω and ρ

$$2D_{ll} = -\frac{f_\rho^2 m_\rho}{2} e^{-m_\rho t} + \frac{f_\omega^2 m_\omega}{2} e^{-m_\omega t}$$

$$\frac{\hat{\Pi}_{j,disc}}{\hat{\Pi}_{j,conn}} = \frac{1}{2} \left[\frac{m_\rho^{2j+2} f_\omega^2}{m_\omega^{2j+2} f_\rho^2} - 1 \right]$$

We do not have accurate information on decay constants
 because of width of ρ , mixing of ω etc

Taking $f_\rho = 0.21(1)$ GeV, $f_\omega = 0.20(1)$ GeV

HVP : disc-1l/conn-1l = -1.5(1.5) %

Adding contributions to $(g-2)/2$

Largest value we can get? (since uses rescaling with largest f_ρ)

ETMC
1308.4327

HPQCD:

$\times 10^{-10}$

567(11)stat

light, connected	662(11)
strange connected	53.4(6)
charm connected	14.4(4)
bottom connected	0.27(4)
disconn. (estimate)	-25(15)
TOTAL	705(20)

preliminary

1403.1778

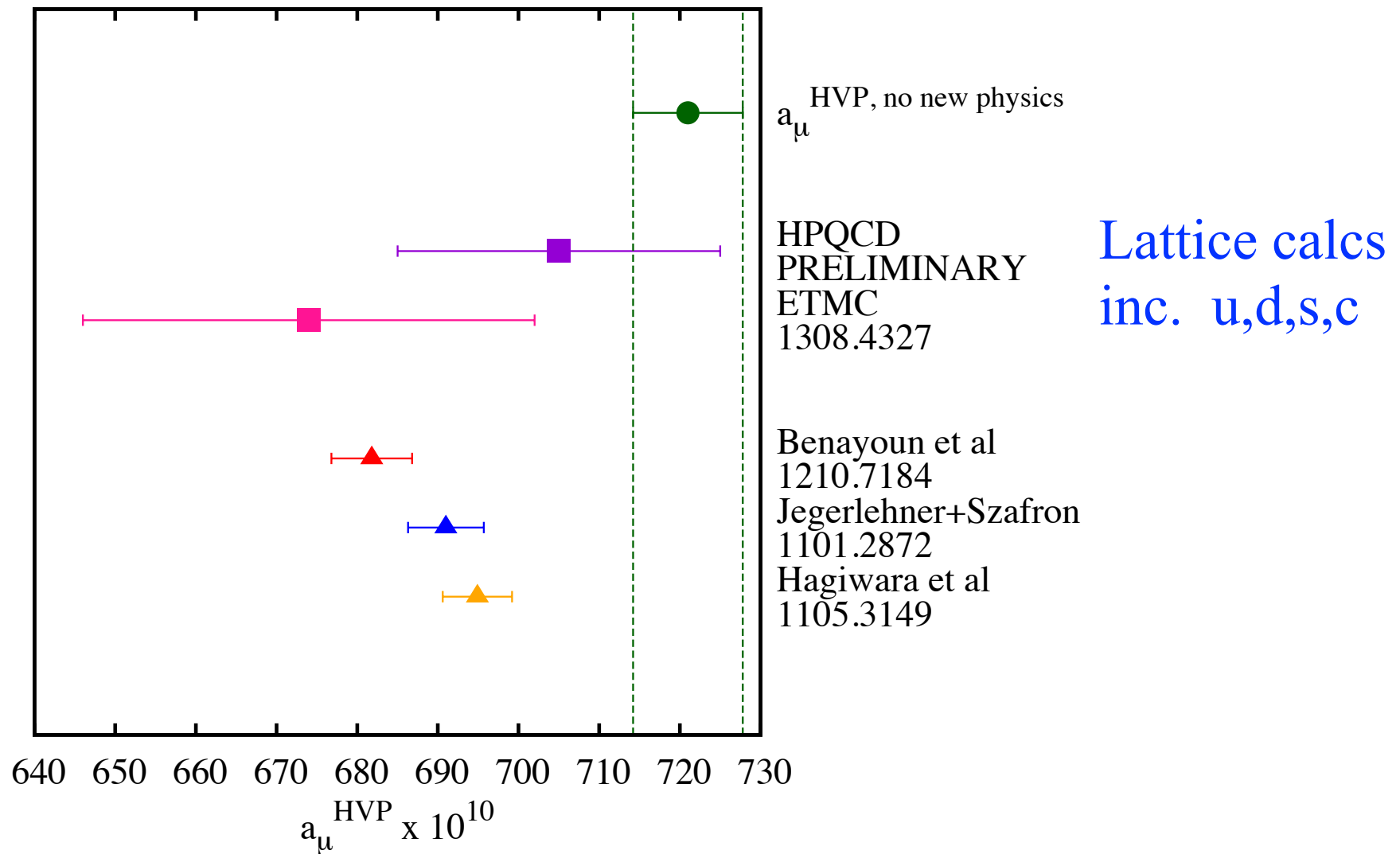
1403.1778,
1208.2855

1408.5768

add -7 from $\pi\pi$
to maximum
simple estimate

674(28)

Lattice - continuum comparison



Good agreement but lattice uncty (all from u/d) still too big.
Need to add in QED, m_u/m_d effects ($\sim 1\%$ and positive?)