

G_2 Gauge Theory

Andreas Wipf

Theoretisch-Physikalisches Institut
Friedrich-Schiller-University Jena

collaboration with

A. Maas, L. von Smekal, B. Wellegehausen, C. Wozar

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- 1 Why G_2 Gauge Theory
- 2 Pure G_2 Gauge Theory
- 3 From G_2 to $SU(3)$
- 4 Spectroscopy
- 5 Finite Temperature and Density
 - Phase diagram – deconfinement transition
 - Phase diagram – baryonic matter transition



Axel Maas (Graz)



Christian Wozar (Jena)

eff. Polyakov Loop dynamics
Phys. Rev. **D80** (2009) 065028

Casimir Scaling/string breaking
Phys. Rev. **D83** (2011) 016001

phase diagram YMH
Phys. Rev. **D83** (2011) 114502

phase diagram with fermions
Phys. Rev. **D86** (2012) 111901

phase diagram, spectroscopy
Phys. Rev. **D89** (2014) 056007



Lorenz von Smekal
(Darmstadt/Gießen)



Bjoern Wellegehausen
(Jena/Gießen)

- rank $G_2 = \text{rank } SU(3) = 2$, dimension = 14
- $SU(3)$ subgroup of G_2
- $G_2/SU(3) \sim S^7 \rightarrow$ efficient parametrization
- G_2 subgroup of $SO(7)$
- smallest (simply connected) Lie group with **trivial center**
- representations are real
- fundamental representations $\{7\}$, $\{14\}$ (= adjoint)
7 **quarks** instead of 3 (cp. GUTS)
- can be **broken to $SU(3)$** with scalars in $\{7\}$

- singlet representation \rightarrow colorless states

$$\{7\} \otimes \{7\} = \{1\} \oplus \{7\} \oplus \{14\} \oplus \{27\}$$

$$\{7\} \otimes \{7\} \otimes \{7\} = \{1\} \oplus 4 \cdot \{7\} \oplus 2 \cdot \{14\} \oplus \dots$$

$$\{14\} \otimes \{14\} = \{1\} \oplus \{14\} \oplus \{27\} \oplus \dots,$$

$$\{14\} \otimes \{14\} \otimes \{14\} = \{1\} \oplus \{7\} \oplus 5 \cdot \{14\} \oplus \dots,$$

$$\{7\} \otimes \{14\} \otimes \{14\} \otimes \{14\} = \{1\} \oplus \dots$$

- branching $G_2 \rightarrow SU(3)$

$$\text{fermions:} \quad \{7\} \longrightarrow \{3\} \oplus \{\bar{3}\} \oplus \{1\},$$

$$\text{gauge bosons:} \quad \{14\} \longrightarrow \{8\} \oplus \{3\} \oplus \{\bar{3}\}.$$

- quark gluon plasma at **high** T (asymptotically free)
- confining at **low** T
- **trivial center**:
what about confinement models based on center?
- **SU(3) as subgroup**: $G_2 \xrightarrow{SSB} SU(3)$
- colorless **baryons and mesons**, different to SU(2)
- **no sign problem for any N_f** :
→ simulations at finite T and finite μ
- phase-transition at finite T and n_B
→ relevant degrees of freedom under extreme conditions?
→ proposals/speculations on **exotic phases** of cold dense matter

- 1st order **confinement/deconfinement** PT
 - Hollands, Minkowski, Pepe, Wiese 2003
 - Cossu, D'Elia, di Giacomo, Lucini, Pica 2007
- Casimir scaling
 - Greensite, Langfeld, Olejnik, Reinhardt, Tok 2007; Liptak, Olejnik 2008
- Effective Polyakov loop dynamics
 - Wellegehausen, AW, Wozar 2009
- chiral restoration in quenched theory at same T_c
 - Danzer, Gattringer, Maas 2009
- G₂ monopoles (Shnir), semiclassical dyon picture
 - (Diakonov, Petrov 2010), . . .
- string breaking, gluelump spectrum
 - Wellegehausen, AW, Wozar 2011
- SSB G₂ → SU(3)
 - Wellegehausen, AW, Wozar 2012
- topological properties, instantons
 - Maas, Olejnik, Ilgenfritz 2013
- G₂ with adjoint Weyl on $S^1 \times \mathbb{R}^3$, small L, m
 - Poppitz, Schäfer, Ünsal 2013
- equation of state, $\langle T_{\mu}^{\mu} \rangle_T \propto T^2$
 - Bruno, Caselle, Panero, Pellegrini 2015
- **ongoing**: G₂ with dynamical fermions, $\mu, T > 0$
 - Maas, Smekal, Wellegehausen, AW

approximate
order parameter:

Polyakov loop:

$$\chi_7 = \text{tr } \mathcal{P}$$

$$\chi_{14} = \text{tr } \mathcal{P}_{\text{adj}}$$

strong coupling,
mean field
→ effective models

$$S_{\text{eff}}[\mathcal{P}]$$

→ approximate
phase portrait

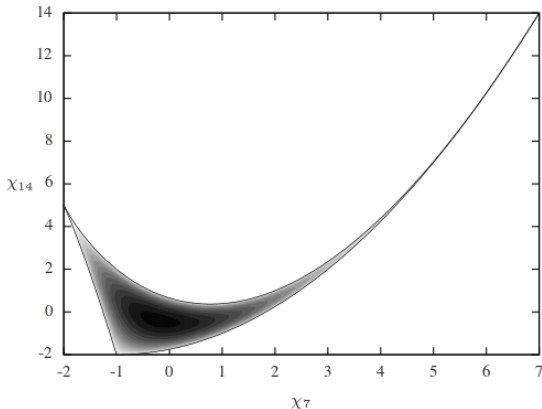
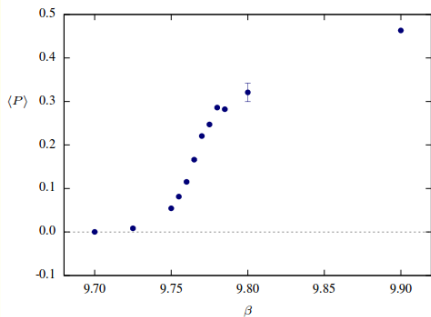
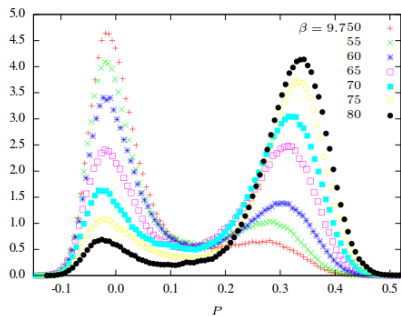


FIG. 2: Fundamental domain of G_2 . Darker regions indicate a bigger Haar measure.

- simulations: **first order PT**
- local HMC, moderate $16^3 \times 6$ lattice



rapid change with β



histogram in vicinity of $\beta_c \approx 9.765$

- static quark-antiquark potential

$$\langle P(x)P^\dagger(y) \rangle_\beta = e^{-\beta V(R)}$$

- heavy quarks in representation \mathcal{R} :

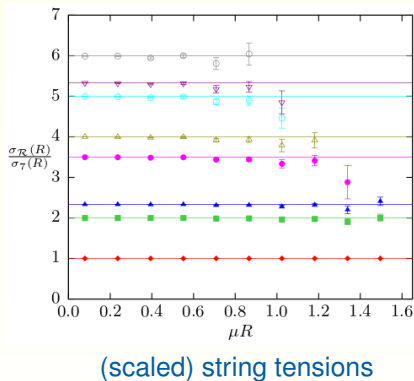
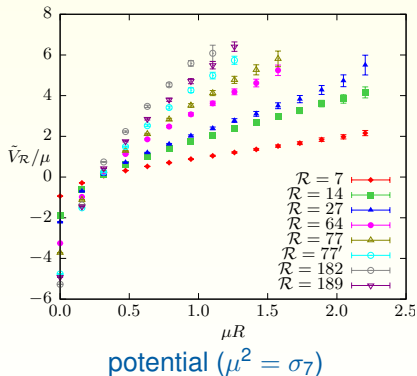
$$V_{\mathcal{R}}(R) = \gamma_{\mathcal{R}} - \frac{\alpha_{\mathcal{R}}}{R} + \sigma_{\mathcal{R}} R$$

- $e^{-\beta V(R)}$ varies over ≈ 50 orders of magnitude until string breaks
- exponential error reduction **Lüscher-Weisz multilevel algorithm**
- **Casimir scaling hypothesis** for string tensions:

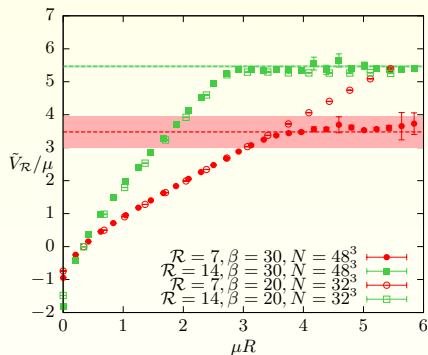
$$\frac{\sigma_{\mathcal{R}}}{C_{\mathcal{R}}} = \frac{\sigma_{\mathcal{R}'}}{C_{\mathcal{R}'}}$$

from ratios of Wilson(Polyakov) loops

- **static potential** for heavy quarks in 8 smallest G_2 representations
- multilevel, $N = 28$, $\beta = 40$
- **no sign of Casimir scaling violation** within 4% for (7,...,189) in 3d, 5% for (7,...,64) in 4d
- continuum string tension (Nair et al.), Lüscher term (CFT)



- $N = 48, \beta = 30$
- near string breaking scale:
Casimir scaling within 2.5%
- beyond: static \mathcal{R} -quarks
screened by gluons
- decay products: **gluelumps**
- cont. limit for V_7 not reached



$$\exp(-m_{\mathcal{R}} T) \propto \left\langle \bigotimes_{n=1}^{N(\mathcal{R})} F_{\mu\nu}(y)_{\mathcal{R},a} \mathcal{R}(U_{yx})_{ab} \bigotimes_{n=1}^{N(\mathcal{R})} F_{\mu\nu}(x)_{\mathcal{R},b} \right\rangle$$

Wellegehausen, AW., Wozar (2011)

string breaking in SU(2): Philipsen, Wittig; Stephenson; de Forcrand, Phillipsen; Pepe, Wiese

G_2 Higgs Model

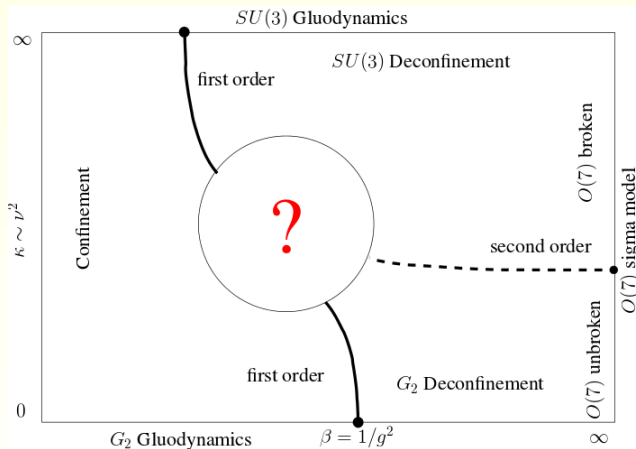
- lattice action with normalized Higgs $\varphi = (\varphi_1, \dots, \varphi_7)^T$ in $\{7\}$

$$S_{\text{YMH}}[U, \varphi] = -\beta \sum_{\square} \left(1 - \frac{1}{7} \text{tr } U_{\square} \right) - \kappa \sum \varphi_x^T U_{x,\mu} \varphi_{x+\mu}$$

- SSB-pattern for $v = \langle \varphi \rangle \neq 0$:
- $\{14\} \longrightarrow \{8\} \oplus \{3\} \oplus \{\bar{3}\}$
 - $\{8\}$: **SU(3) gluons**
 - $\{3\} + \{\bar{3}\}$: massive Vector bosons
- scalars $7 \rightarrow 1$

- $\kappa = 0$: pure G_2 gauge theory:
known 1st order deconfinement transition
- $\kappa = \infty$: 6 vector bosons decouple, pure $SU(3)$
well-known 1st order deconfinement transition
- Is there 1st order transition line connecting two gauge theories?
- observables/histograms (from up to 5×10^5 configurations)
on grid in (β, κ) -plane ($\beta = 5 \dots 10, \kappa = 0 \dots 10^4$) \implies

Phase diagram of G_2 YMH theory

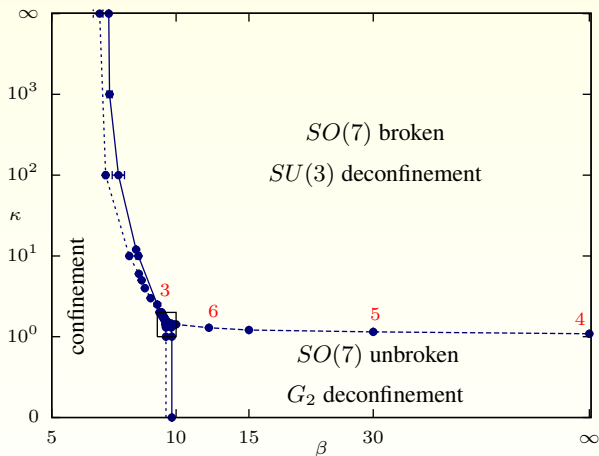


Osterwalder, Seiler; Fradkin, Shenker does not apply

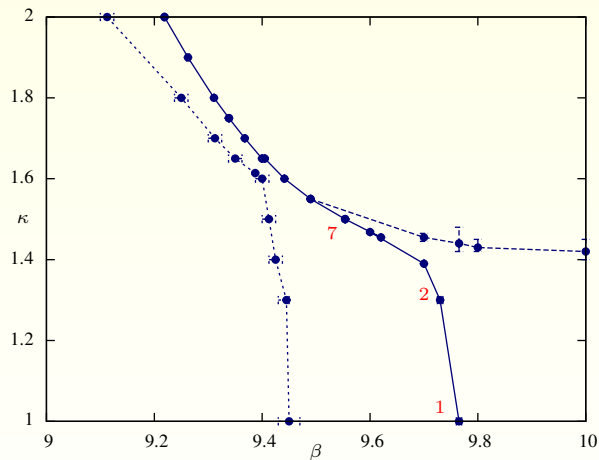
- LHMC, $N_s \leq 24$
- study of Higgs transition line (beware of bulk transition!)
- average plaquette action, Higgs action and Polyakov loop
- susceptibilities (and higher derivatives), finite size analysis
- large β : line of second order PT $O(7) \rightarrow O(6)$
- line of first order PT $G_2 \rightarrow SU(3)$ with (very) small gap
- (would be) triple point

$$\beta_{\text{crit}} = 9.55(5) \quad , \quad \kappa_{\text{crit}} = 1.50(4)$$

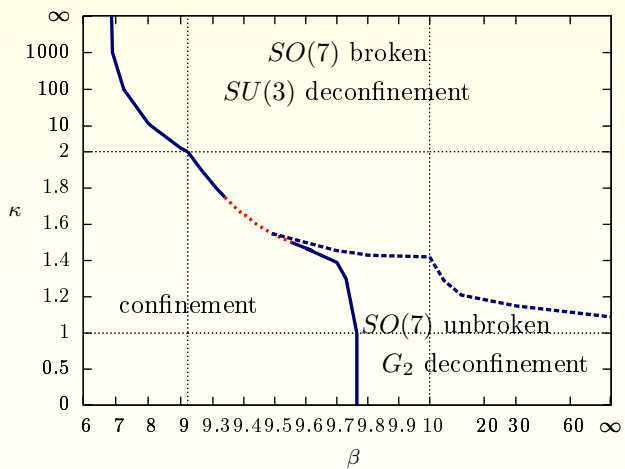
- first order (almost?) line hits second order line



$16^3 \times 6$ -lattice



$N_s \leq 24$

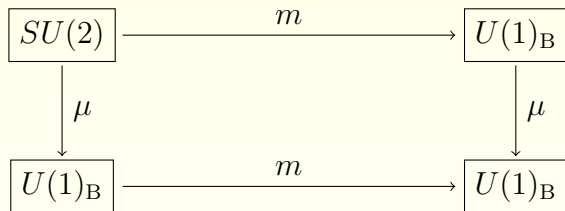


Summarizing the phase diagram of G_2 YMH-theory

Wellegehausen, Wozar, AW

- $\det D(U, \mu, m) \geq 0$
⇒ simulations at finite T and μ
- phases characterized by:
chiral symmetry, baryon-density, Polyakov loop, condensates, ...
- interpretation of phase transitions and condensates
⇒ particle spectrum (T=0)
- glueballs
- mesons
- baryons: $n_q = 1, 2, 3$

^acollaboration with Axel Maas, Lorenz von Smekal und Bjoern Wellegehausen



- enlarged chiral symmetry (Pauli-Gürsey)
- well-defined baryon number
- Goldstone bosons

$$d(0^{++}) \sim \bar{\psi}^C \gamma_5 \psi + \bar{\psi} \gamma_5 \psi^C \quad \text{and} \quad d(0^{+-}) \sim \bar{\psi}^C \gamma_5 \psi - \bar{\psi} \gamma_5 \psi^C$$

- massive state

$$f(0^{++}) \sim \bar{\psi} \psi$$

- Wilson fermions, Symanzik improved gauge action
- RHMC-algorithm
- $\beta = 0.96$, $\kappa = 0.159$ fixed (light ensemble)

Spectroscopy ensemble

16×8^3 lattice

Proton mass $m_N = 938 \text{ MeV}$

Diquark mass $m_{d(0^+)} = 247 \text{ MeV}$

Lattice spacing $a = 0.343 \text{ fm} \sim (575 \text{ MeV})^{-1}$

Finite temperature and density

Lattice size: $N_t \times 8^3$ with $N_t = 2 \dots 16$

$T = 36 \dots 287 \text{ MeV}$ (9 values)

$\mu = 0 \dots 354 \text{ MeV}$ (60 values below lattice saturation)

Name	Operator	Pos.	Spin	Colour	Flavour	J	P
η	$\bar{u}\gamma_5 u$	S	A	S	S	0	-
f	$\bar{u}u$	S	A	S	S	0	+
ω	$\bar{u}\gamma_\mu u$	S	S	S	A	1	-
h	$\bar{u}\gamma_5\gamma_\mu u$	S	S	S	A	1	+
π	$\bar{u}\gamma_5 d$	S	A	S	S	0	-
a	$\bar{u}d$	S	A	S	S	0	+
ρ	$\bar{u}\gamma_\mu d$	S	S	S	A	1	-
b	$\bar{u}\gamma_5\gamma_\mu d$	S	S	S	A	1	+

$$\{7\} \otimes \{7\} = \{1\} \oplus \dots$$

- unquenched
- partially quenched in one-flavour simulations

Name	Operator	Pos.	Spin	Col.	Flav.	J	P
Hybrid	$\epsilon_{abcdefg} u^a F_{\mu\nu}^p F_{\mu\nu}^q F_{\mu\nu}^r T_p^{bc} T_q^{de} T_r^{fg}$	S	S	A	S	1/2	\pm
$\tilde{\Delta}$	$T^{abc} (\bar{u}_a \gamma_\mu u_b) u_c$	S	S	A	S	3/2	\pm
\tilde{N}	$T^{abc} (\bar{u}_a \gamma_5 d_b) u_c$	S	A	A	A	1/2	\pm

$$\{7\} \otimes \{7\} \otimes \{7\} = \{1\} \oplus \dots$$

$$\{7\} \otimes \{14\} \otimes \{14\} \otimes \{14\} = \{1\} \oplus \dots$$

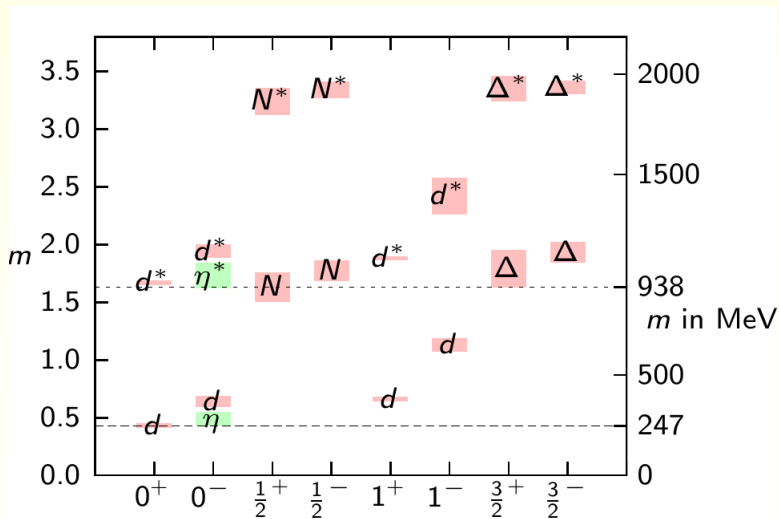
Name	Operator	Pos.	Spin	Colour	Flavour	J	P	C
$d(0^{++})$	$\bar{u}^C \gamma_5 u + \bar{u} \gamma_5 u^C$	S	A	S	S	0	+	+
$d(0^{+-})$	$\bar{u}^C \gamma_5 u - \bar{u} \gamma_5 u^C$	S	A	S	S	0	+	-
$d(0^{-+})$	$\bar{u}^C u + \bar{u} u^C$	S	A	S	S	0	-	+
$d(0^{--})$	$\bar{u}^C u - \bar{u} u^C$	S	A	S	S	0	-	-
$d(1^{++})$	$\bar{u}^C \gamma_\mu d + \bar{u} \gamma_\mu d^C$	S	S	S	A	1	+	+
$d(1^{+-})$	$\bar{u}^C \gamma_\mu d - \bar{u} \gamma_\mu d^C$	S	S	S	A	1	+	-
$d(1^{-+})$	$\bar{u}^C \gamma_5 \gamma_\mu d + \bar{u} \gamma_5 \gamma_\mu d^C$	S	S	S	A	1	-	+
$d(1^{--})$	$\bar{u}^C \gamma_5 \gamma_\mu d - \bar{u} \gamma_5 \gamma_\mu d^C$	S	S	S	A	1	-	-

$$\{7\} \otimes \{7\} = \{1\} \oplus \dots$$

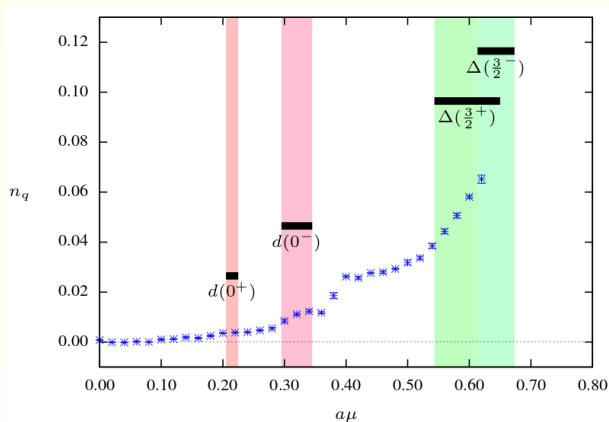
Name	Operator	Pos.	Spin	Colour	Flavour	J	P
Δ	$T^{abc}(\bar{u}_a^C \gamma_\mu u_b) u_c$	S	S	A	S	3/2	\pm
N	$T^{abc}(\bar{u}_a^C \gamma_5 d_b) u_c$	S	A	A	A	1/2	\pm

$$\{7\} \otimes \{7\} \otimes \{7\} = \{1\} \oplus \dots$$

- diquark masses are degenerate (cc)
- $m_\eta - m_{\text{diquark}} =$ disconnected contributions
- no sources for diquarks needed
- N_f complex-valued pseudo-fermions plus RHMC



Wellegehausen, Maas, Smekal, AW (2013)

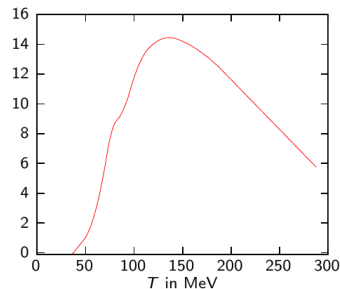
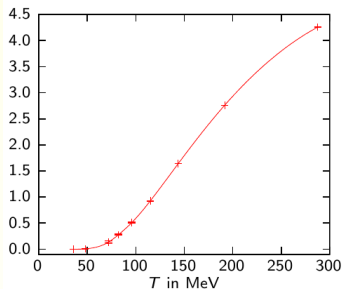


- silver blaze
- diquark BE condensation
 $m_{d(0^+)} = 247 \text{ MeV}$
- plateaus visible for larger n_q
- three transitions
- hadronic phase at $m_\Delta \approx m_N$

Wellegehausen, Maas, Smekal, AW (2013)

- low density: in accordance with **silver blaze**
- two small jumps at **diquark thresholds**
⇒ two (probably) second order PT?
- two **plateaus** after thresholds
- **Bose-condensates** of diquarks?
- one (probably) first order PT at $\approx \Delta$ threshold
- hadronic phase for higher n_q (under investigation)
- $a\mu \gtrsim 1$: lattice artifacts, e.g. saturation effects

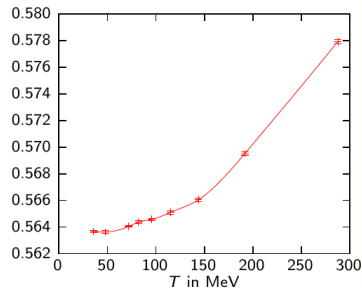
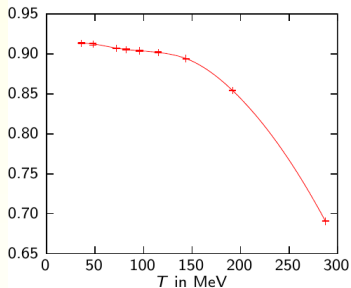
Polyakov loop $\langle P \rangle$ and susceptibility $\frac{\partial \langle P \rangle}{\partial T}$ at $\mu = 0$



Deconfinement transition at $T_c(\mu = 0) \approx 137$ MeV
(cubic splines, no thermalization at 80 MeV)

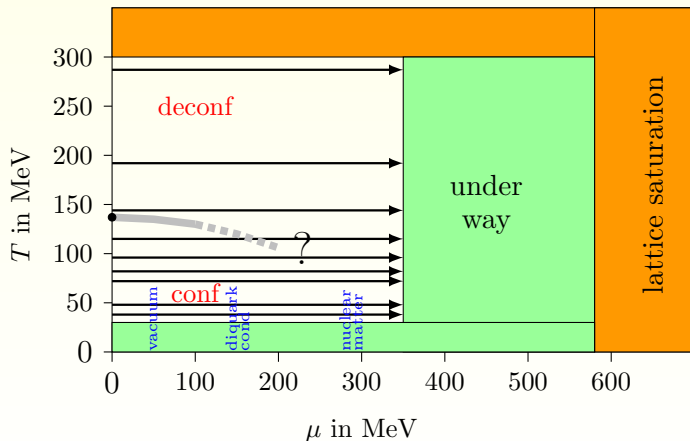
Loewe cluster, 15 000

chiral condensate $\langle \Sigma \rangle - \Sigma_{\text{SB,Latt}}$ and plaquette $\langle \text{Plaq} \rangle$ at $\mu = 0$

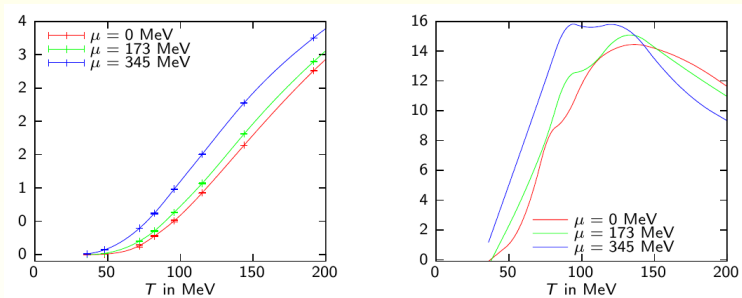


still not clear whether deconfinement
and chiral transition temperature agree

- **temperatures:** $T = 287, 192, 144, 115, 96, 82, 72, 48, 36$ MeV
- **chemical potential:** 60 μ -values between 0 and 354 MeV for each T

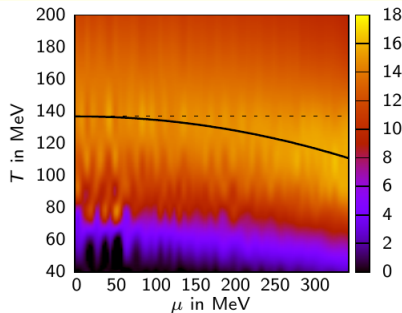
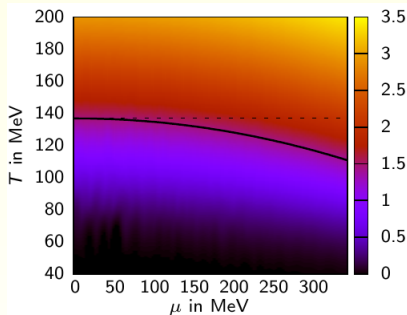


Polyakov loop $\langle P \rangle$ and susceptibility $\frac{\partial \langle P \rangle}{\partial T}$



deconfinement transition shifts to smaller T for larger μ
 peak of susceptibility increases

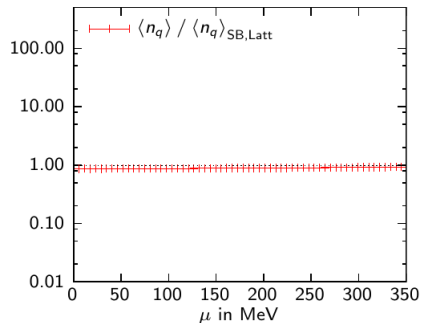
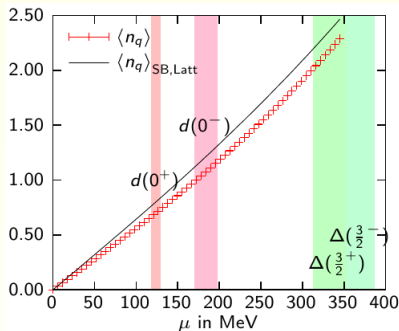
Polyakov loop $\langle P \rangle$ and susceptibility $\frac{\partial \langle P \rangle}{\partial T}$ (60μ 's)



deconfinement transition shifts to smaller T for larger μ
 peak of susceptibility increases, consistent with $T'_c(\mu)|_{\mu=0} = 0$

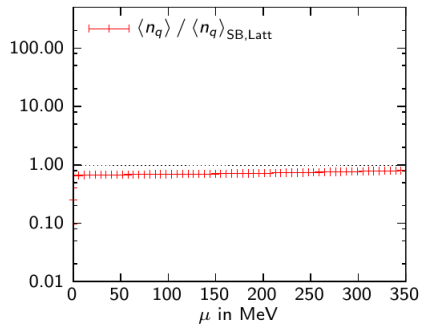
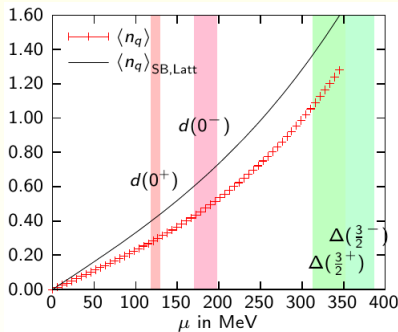
$$\frac{T_c(\mu)}{T_c(0)} \sim 1 + 0.001(8) \frac{\mu}{T_c(0)} - 0.031(6) \left(\frac{\mu}{T_c(0)} \right)^2$$

quark number density



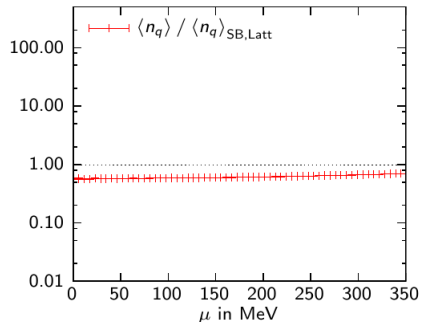
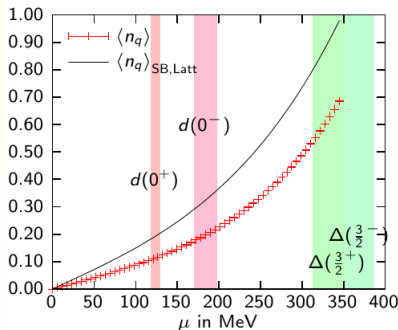
$T = 287$ MeV

quark number density



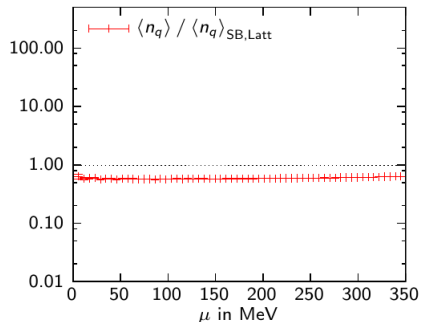
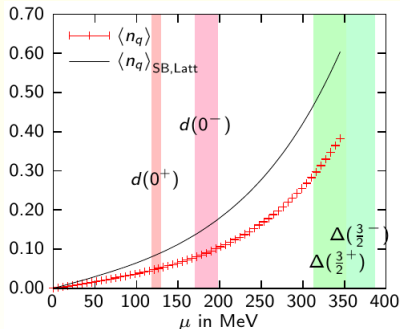
$T = 192$ MeV

quark number density



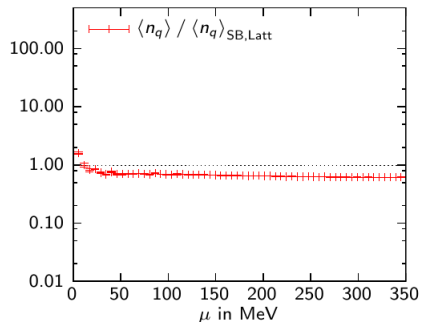
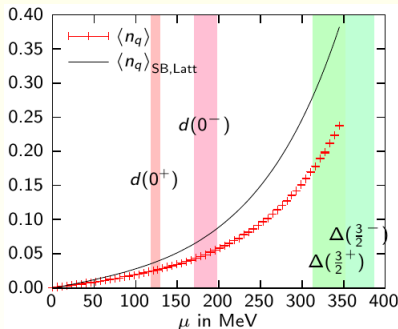
$T = 144$ MeV

quark number density



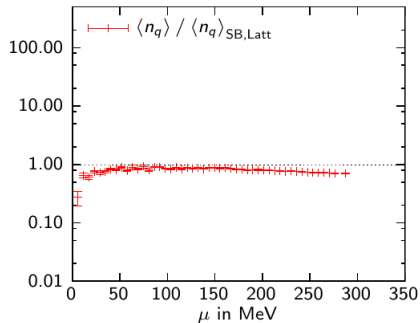
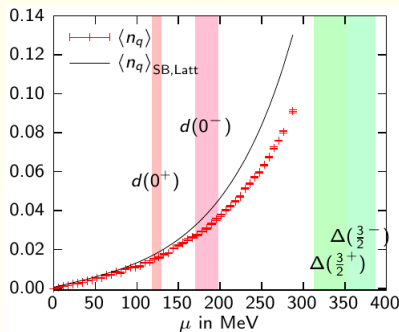
$$T = 115 \text{ MeV}, \mu_{\text{deconf}} = 317 \text{ MeV}$$

quark number density



$T = 96$ MeV, threshold effects?

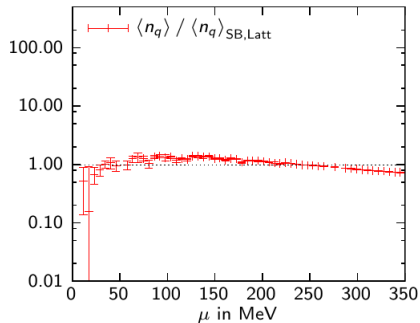
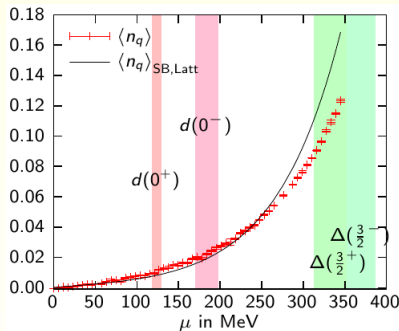
quark number density



$T = 82$ MeV

estimator problem for small μ when $\langle n_B \rangle \approx 10^{-3} - 10^{-4}$

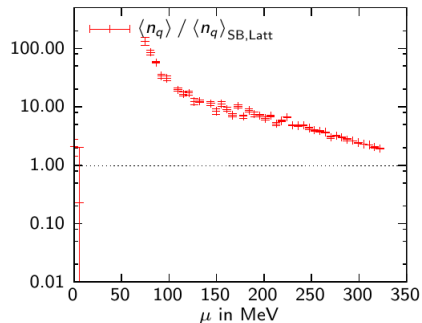
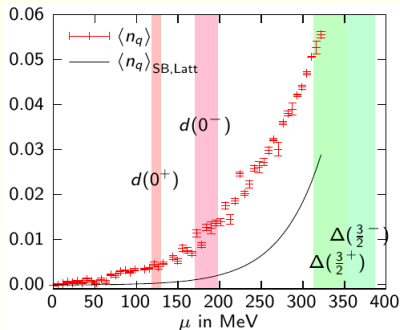
quark number density



$T = 72$ MeV

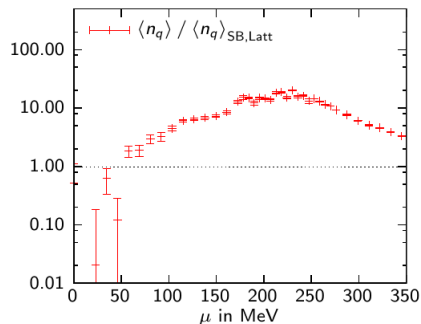
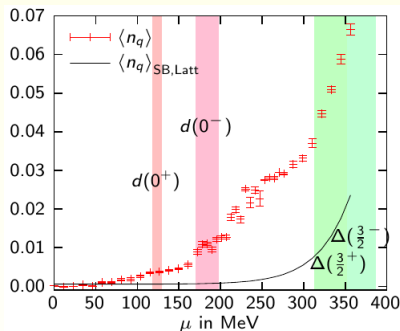
$\langle n_B \rangle$ may jump at thresholds

quark number density

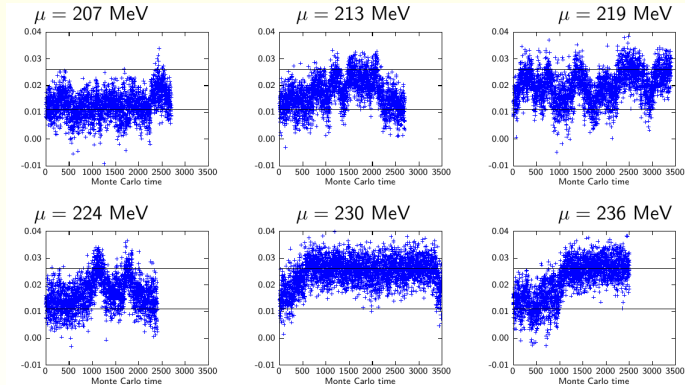


$T = 48$ MeV

quark number density

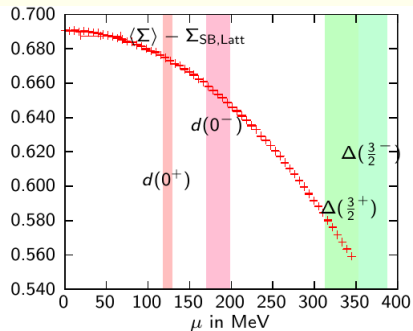
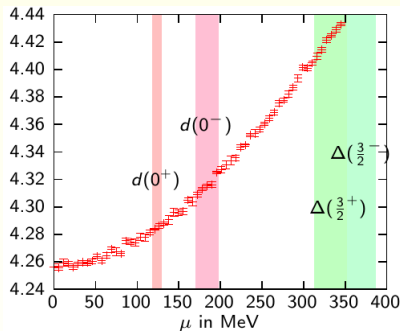


$T = 36$ MeV

quark number density at $T = 36 \text{ MeV}$ 

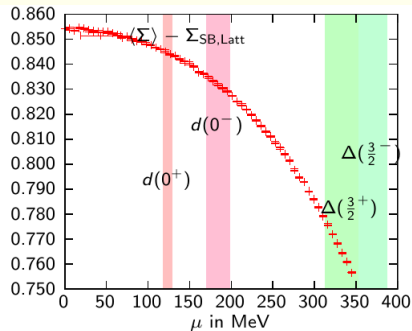
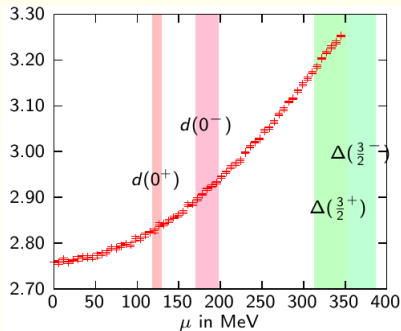
first order nuclear matter transition at $\mu_c \approx 219 \text{ MeV}$?
 10 times estimated \rightarrow wide bands
 simulation expensive near μ_c (small m_d)

Polyakov loop and chiral condensate



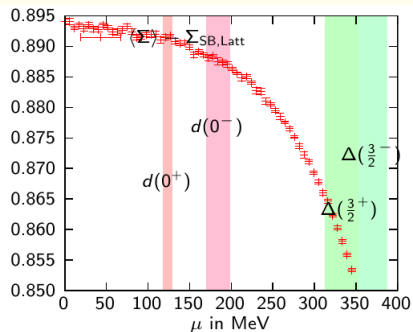
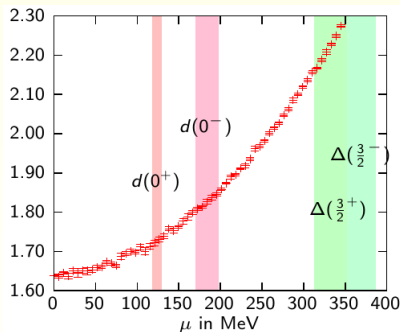
$T = 287 \text{ MeV}$

Polyakov loop and chiral condensate



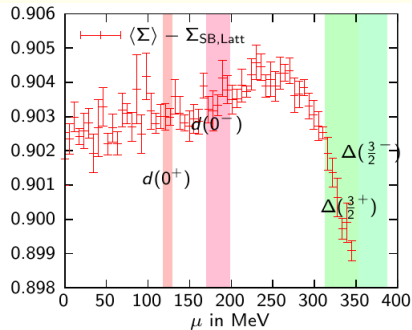
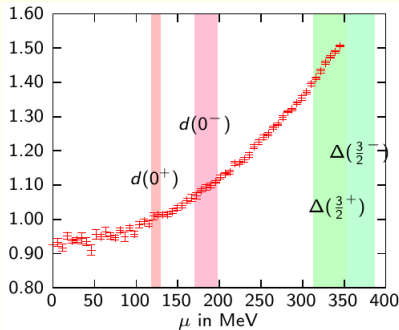
$T = 192 \text{ MeV}$

Polyakov loop and chiral condensate



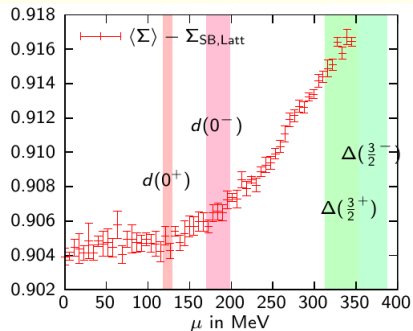
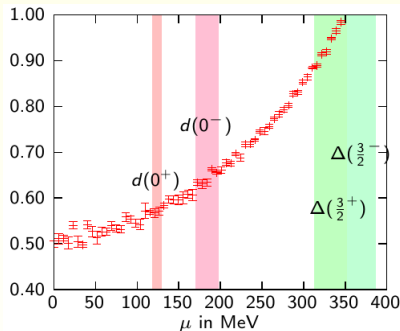
$T = 144$ MeV

Polyakov loop and chiral condensate



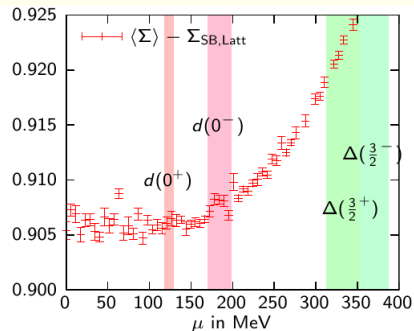
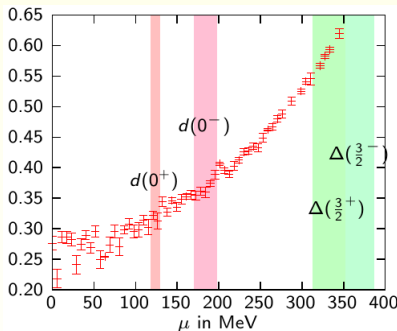
$T = 115 \text{ MeV}$, $\mu_{\text{deconf}} = 317 \text{ MeV}$
(extrapolation of χ_P)

Polyakov loop and chiral condensate



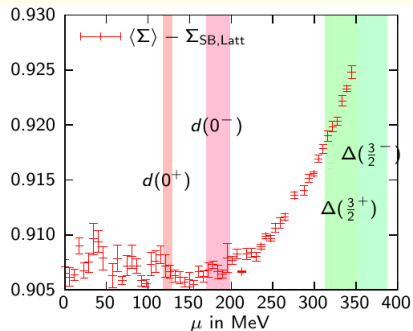
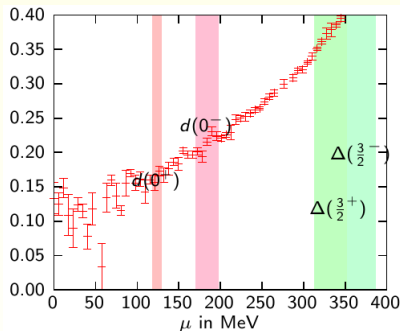
$T = 96$ MeV

Polyakov loop and chiral condensate



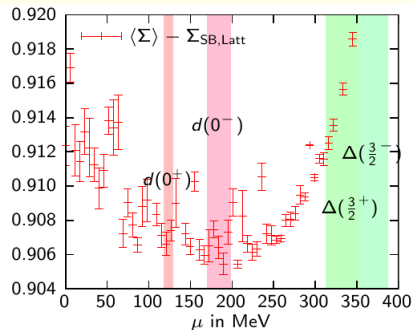
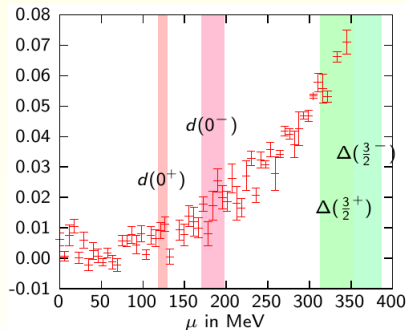
$T = 82 \text{ MeV}$

Polyakov loop and chiral condensate



$T = 72$ MeV

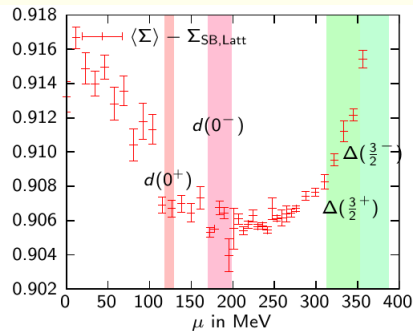
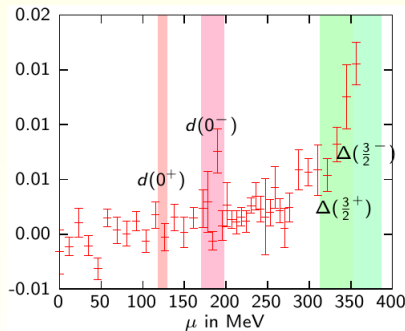
Polyakov loop and chiral condensate



$$T = 48 \text{ MeV}$$

sign of di-quark condensation
 chiral condensate \rightarrow diquark condensate

Polyakov loop and chiral condensate



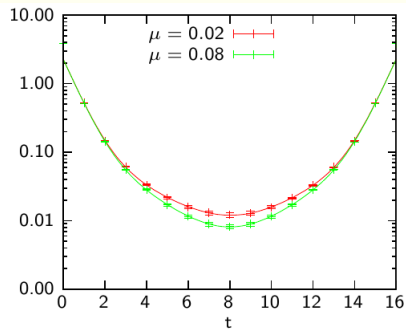
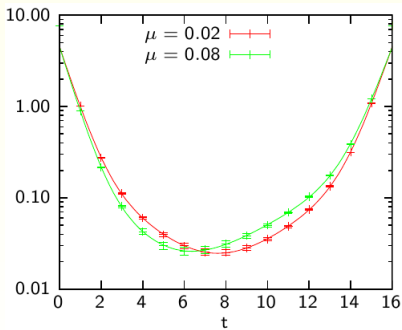
$T = 36 \text{ MeV}$

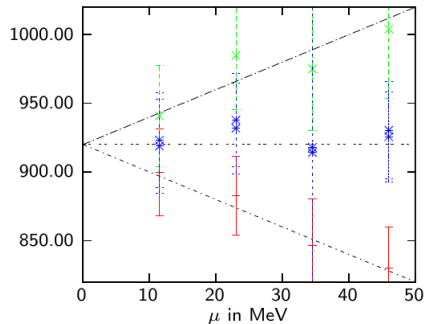
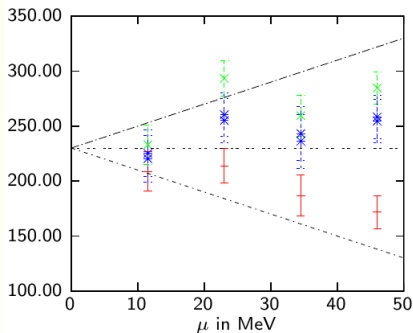
number of estimators?

Lattice correlation function for operator with quark number n_q

$$C(\mu, n_q) \sim a \exp^{-(m(\mu)+n_q\mu)t} + b \exp^{(m(\mu)-n_q\mu)t}$$

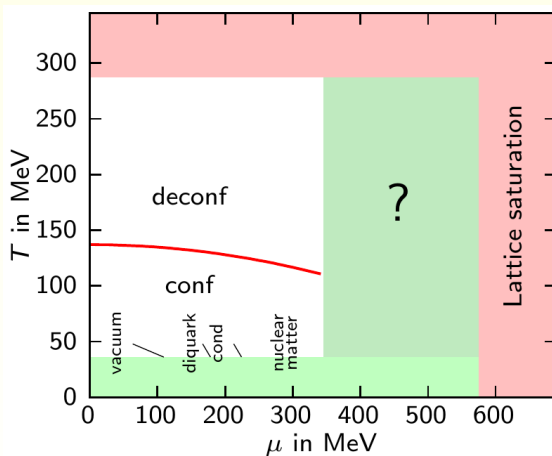
$d(0^+)$ and π correlation functions fitted with 4 exponentials for ground and excited states



ground state and first excited state masses of $d(0^+)$ and π 

below the silver blaze transition the masses do not depend on μ

G_2 -QCD phase diagram with $V = (2.7 \text{ fm})^3$ and $m_{d(0^+)} = 247 \text{ MeV}$



- finite size effects?
- evidence for first order nuclear matter transition
- chiral symmetry restoration ($\rightarrow N_f = 2$ with pions)?
- deconfinement transition at lower temperatures / critical endpoint?

- hybrid spectroscopy (difficult)
- spectroscopy at finite density (larger μ)
- diquark condensation
- $N_f = 2$ phase diagram
- equations of state

A. Maas, L. von Smekal, B. H. Wellegehausen and A. Wipf, *The phase diagram of a gauge theory with fermionic baryons*, arXiv:1203.5653 [hep-lat], 2012.

A. Maas, L. von Smekal, B. H. Wellegehausen and A. Wipf, *Hadron masses and baryonic scales in G2-QCD at finite density*, arXiv:1312.5579 [hep-lat], 2014.

- G_2 QCD is a useful laboratory
- interpolation G_2 -QCD \rightarrow QCD with Higgs-mechanism possible
- full phase diagram in principle accessible to simulations
- follow first order(?) transition line (critical end point?)
- break G_2 -QCD \rightarrow QCD with quarks via Higgs-mechanism
- deformation vs. sign problem?
- testbed for model building
- testbed for alternative approaches (expansions, renormalization group)?
- heavy ensemble vs. strong coupling / hopping parameter expansion