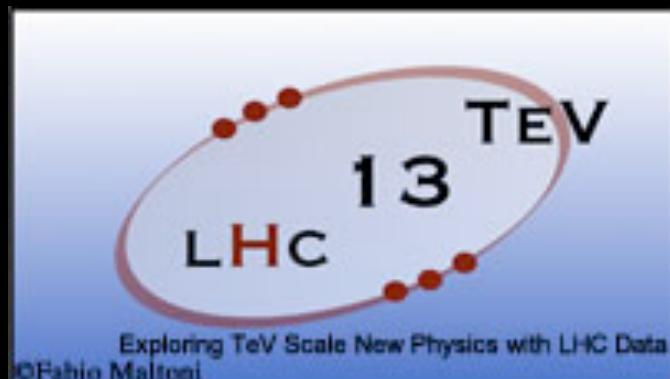


A Flavorful Top-Coloron Model

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MICHIGAN STATE UNIVERSITY



- Extended Color Dynamics
- A Top-Coloron Model
- Flavor Symmetries and Constraints
- Scalars: Same Sign Top Signature
- Flavor Independent Constraints
- Conclusions

KITP

JULY 10, 2013

EXTENDED COLOR DYNAMICS

New colored gauge bosons

Classic Axigluon: P.H. Frampton and S.L. Glashow, Phys. Lett. B 190, 157 (1987).

Topgluon: C.T. Hill, Phys. Lett. B 266, 419 (1991).

Flavor-universal Coloron: R.S. Chivukula, A.G. Cohen, & E.H. Simmons, Phys. Lett. B 380, 92 (1996).

Chiral Color with $g_L \neq g_R$: M.V. Martynov and A.D. Smirnov, Mod. Phys. Lett. A 24, 1897 (2009).

New Axigluon: P.H. Frampton, J. Shu, and K. Wang, Phys. Lett. B 683, 294 (2010).

Other color-octet states: (cf. “partial compositeness”)

KK gluon: H. Davoudiasl, J.L. Hewett, and T.G. Rizzo, Phys. Rev. D63, 075004 (2001)
B. Lillie, L. Randall, and L.-T. Wang, JHEP 0709, 074 (2007).

Techni-rho: E. Farhi and L. Susskind, Physics Reports 74, 277 (1981).

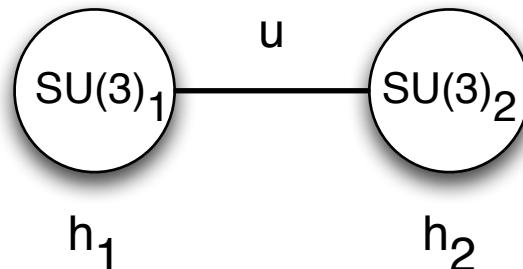
Recent catalog of colored states:

Color sextets, colored scalars, low-scale scale string resonances...

T. Han, I. Lewis, Z. Liu, JHEP 1012, 085 (2010).

GAUGE SECTOR

COLORON MODELS: GAUGE SECTOR



$SU(3)_1 \times SU(3)_2$ color sector with $M^2 = \frac{u^2}{4} \begin{pmatrix} h_1^2 & -h_1 h_2 \\ -h_1 h_2 & h_2^2 \end{pmatrix}$

unbroken subgroup: $SU(3)_{1+2} = SU(3)_{QCD}$

$$h_1 = \frac{g_s}{\cos \theta} \quad h_2 = \frac{g_s}{\sin \theta}$$

gluon state: $G_\mu^A = \cos \theta A_{1\mu}^A + \sin \theta A_{2\mu}^A$

couples to: $g_S J_G^\mu \equiv g_S (J_1^\mu + J_2^\mu)$

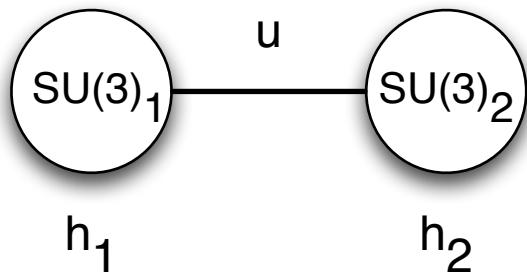
coloron state: $C_\mu^A = -\sin \theta A_{1\mu}^A + \cos \theta A_{2\mu}^A \quad M_C = \frac{u}{\sqrt{2}} \sqrt{h_1^2 + h_2^2}$

couples to: $g_S J_C^\mu \equiv g_S (-J_1^\mu \tan \theta + J_2^\mu \cot \theta)$

low-energy current-current interaction: $\mathcal{L}_{FF}^2 = -\frac{g_S^2}{2M_C^2} J_C^\mu J_C^\mu$

FERMIIONS

COLORON MODELS: QUARK CHARGES



$$g_S J_G^\mu \equiv g_S (J_1^\mu + J_2^\mu)$$

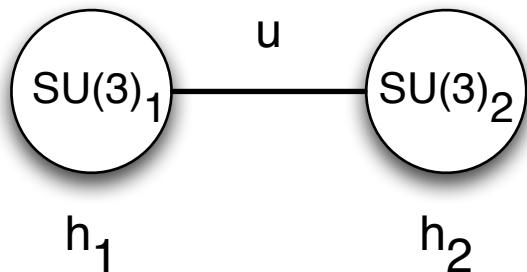
$$g_S J_C^\mu \equiv g_S (-J_1^\mu \tan \theta + J_2^\mu \cot \theta)$$

low-energy current-current interaction: $\mathcal{L}_{FF}^2 = -\frac{g_S^2}{2M_C^2} J_C^\mu J_{C\mu}$

Depending on how quarks transform under $SU(3)_1 \times SU(3)_2$ the presence of colorons may impact

- LHC **dijet** mass distribution (or angular distribution)
- kinematic distributions of **tt or bb** final states
- asymmetry in top-quark production: **A^t_{FB}**
- **FCNC** processes: $K\bar{K}, D\bar{D}, B\bar{B}$ mixing, $b \rightarrow s\gamma$
- **precision EW** observables: delta-rho, **R_b**

COLORON MODELS: QUARK CHARGES



$$g_S J_G^\mu \equiv g_S (J_1^\mu + J_2^\mu)$$

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- **precision EW observables:** $\Delta\rho, R_b$

PATTERNS OF QUARK CHARGES

SU(3) ₁	SU(3) ₂	model	pheno.
	(t,b) _L q _L t _R ,b _R q _R	coloron	dijet
q _R	(t,b) _L q _L t _R ,b _R		
t _R ,b _R	(t,b) _L q _L q _R		
q _L	(t,b) _L t _R ,b _R q _R		
q _L t _R ,b _R	(t,b) _L q _R	new axigluon	dijet, A ^t _{FB} , FCNC
q _L q _R	(t,b) _L t _R ,b _R	topgluon	dijet, tt, bb, FCNC, R _b ...
t _R ,b _R q _R	(t,b) _L q _L	classic axigluon	dijet, A ^t _{FB}
q _L t _R ,b _R q _R	(t,b) _L		

(No spectators required)

q = u,d,c,s

PATTERNS OF QUARK CHARGES

$SU(3)_1$	$SU(3)_2$	model	pheno.
	(t,b) _L q _L t _R ,b _R q _R	coloron	dijet
q _R	(t,b) _L q _L t _R ,b _R		
t _R ,b _R	(t,b) _L q _L q _R		
q _L	(t,b) _L t _R ,b _R q _R		
q _L t _R ,b _R	(t,b) _L q _R	new axigluon	dijet, A^t_{FB} , FCNC
q _L q _R	(t,b) _L t _R ,b _R	topgluon	dijet, tt, bb, FCNC, R_b ...
t _R ,b _R q _R	(t,b) _L q _L	classic axigluon	dijet, A^t_{FB}
q _L t _R ,b _R q _R	(t,b) _L		

(No spectators required)

$q = u,d,c,s$

A FLAVORFUL TOP-COLORON MODEL

R.S.C., Elizabeth Simmons, N. Vignaroli
arXiv:1302.1069

FLAVORFUL TOP-COLORON MODEL

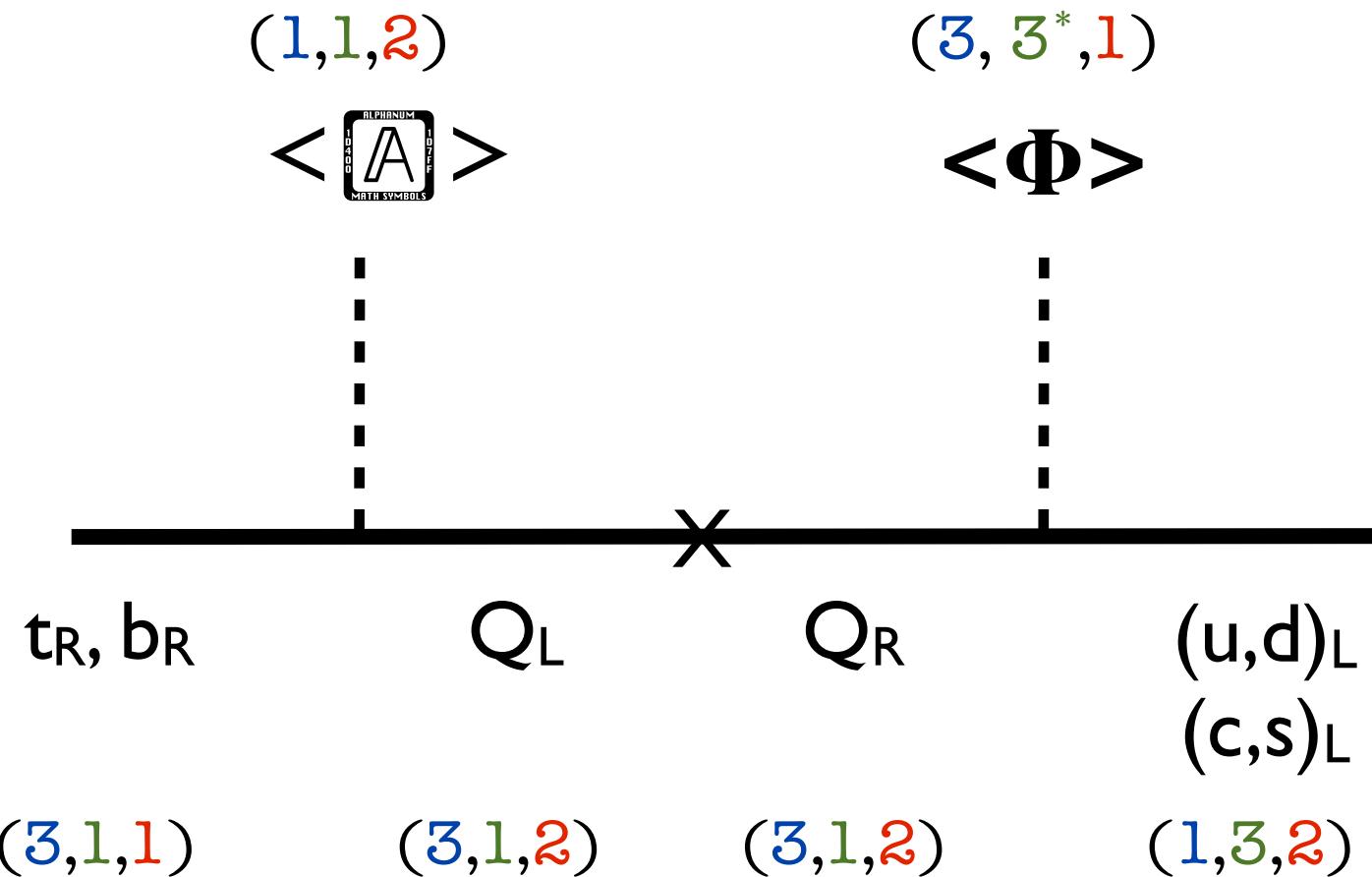
particles		SU(3) ₁	SU(3) ₂	SU(2) _W
3rd generation quarks	(t,b) _L	3	1	2
	t _R ,b _R	3	1	1
light quarks	(u,d) _L (c,s) _L	1	3	2
	u _R ,d _R c _R ,s _R	1	3	1
vector quarks	Q _L ,Q _R	3	1	2
light scalar	  	1	1	2
heavy scalar	Φ	3	3*	1

Next to minimal flavor symmetry:

$$U(2)_{\vec{\psi}_L} \times U(2)_{\vec{u}_R} \times U(2)_{\vec{d}_R} \times U(2)_{\vec{Q}_L} \times U(1)_{t_R} \times U(1)_{b_R} \times U(1)_{Q_R}$$

GENERATIONAL MIXING

$$SU(3)_1 \times SU(3)_2 \times SU(2)_W$$



Mixing to third generation occurs indirectly, through mixing with vector quarks.

GENERATIONAL MIXING

Light Generations

$$\mathcal{M}_u = M \cdot \begin{pmatrix} \Delta_u & \vec{0} & \vec{\alpha} \\ 0 & 0 & \beta_t & 0 \\ 0 & 0 & \lambda'_t & 1 \end{pmatrix}, \quad \mathcal{M}_d = M \cdot \begin{pmatrix} c\Delta_d & \vec{0} & \vec{\alpha} \\ 0 & 0 & \beta_b & 0 \\ 0 & 0 & \lambda'_b & 1 \end{pmatrix}$$

Third Generation

Vector Quarks

Weak Mixing \Rightarrow Cabibbo Matrix, \mathcal{C} , and

$$V_{ub} = \alpha_1 d = \alpha_1 \left(\frac{\lambda'_t}{\beta_t} - \frac{\lambda'_b}{\beta_b} \right) = A\lambda^3(\rho - i\eta) = 0.00131 - i0.00334$$

$$V_{cb} = \alpha_2 d = \alpha_2 \left(\frac{\lambda'_t}{\beta_t} - \frac{\lambda'_b}{\beta_b} \right) = A\lambda^2 = 0.0415 ,$$

$d = \mathcal{O}(1)$
$\alpha_1 = \mathcal{O}(\lambda^3)$
$\alpha_2 = \mathcal{O}(\lambda^2)$

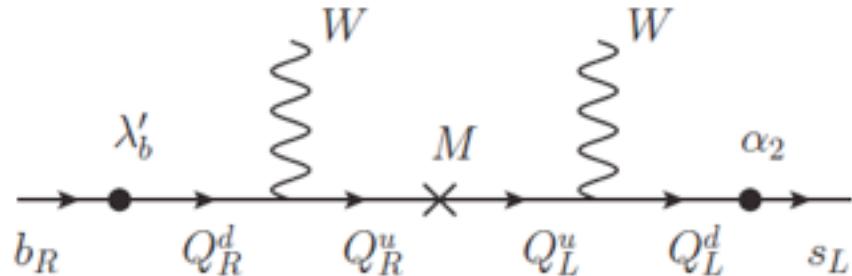
CONSTRAINTS FROM FLAVOR PHYSICS

R.S.C., Elizabeth Simmons, N. Vignaroli
arXiv:1302.1069

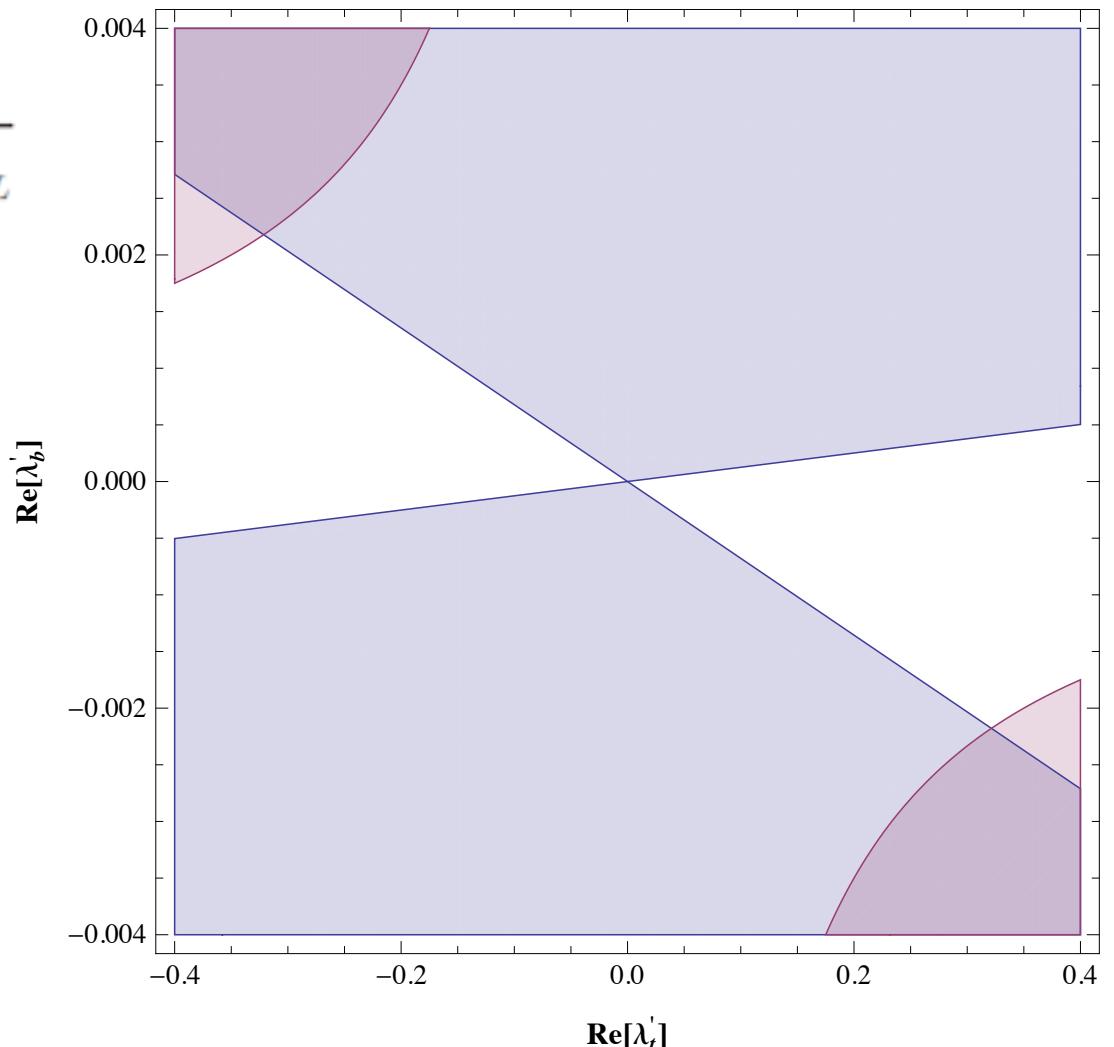
FCNC IN TOP-COLORON MODEL

- Mixing among ordinary and heavy vector quarks also leads to flavor-changing b-quark decays: $b \rightarrow s\gamma$
- Coloron exchange yields KK, DD, and BB mixing
 - ▶ quark charges under strong gauge groups are non-universal
 - ▶ the top and bottom mass eigenstate quarks are admixtures of ordinary and heavy vector gauge eigenstate quarks

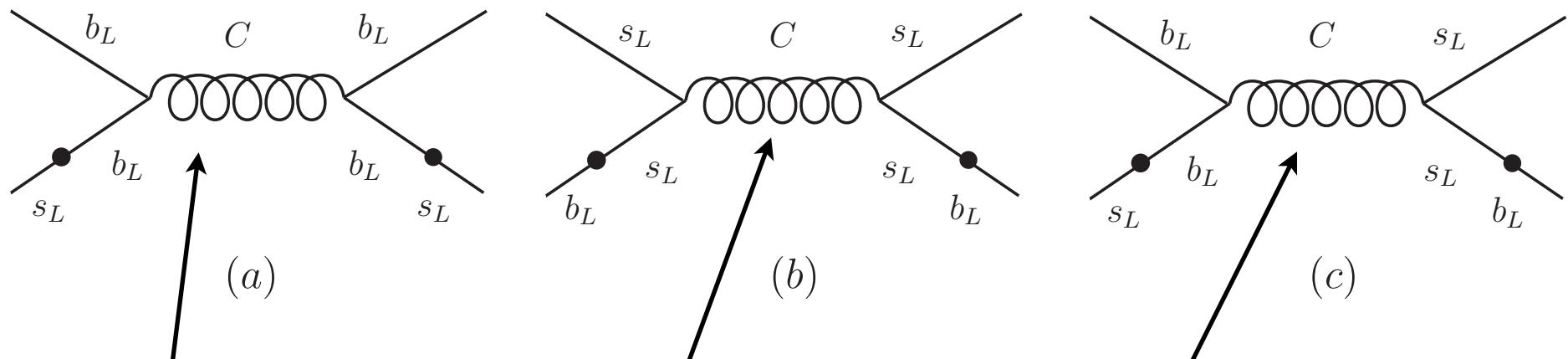
CONSTRAINTS: $B \rightarrow S\gamma$



Mixing with right-handed electroweak doublets enhances contributions to $b \rightarrow s\gamma$



CONSTRAINTS: B-BBAR MIXING



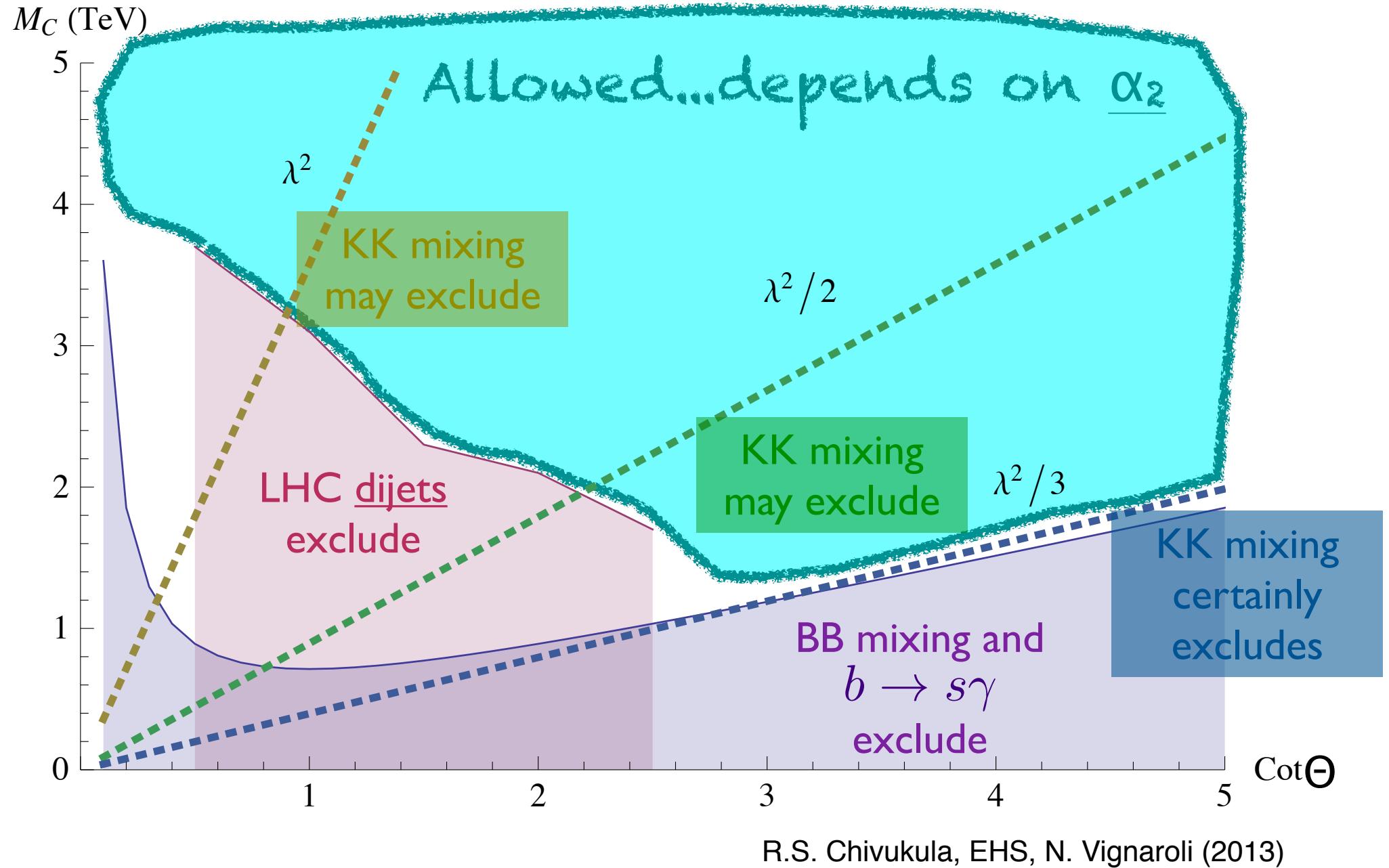
$$g_S^2 (\cot \omega J_{bL}^{a\mu})^2 (D_L)_{32}^2$$

$$g_S^2 (-\tan \omega J_{sL}^{a\mu})^2 (D_L^\dagger)_{32}^2$$

$$2g_S^2 (-\tan \omega J_{sL}^{a\mu}) (\cot \omega J_{bL}^{a\mu}) (D_L)_{32} (D_L^\dagger)_{32}$$

**Flavor-changing Effects from Coloron Exchange:
interplay between mixing and coupling strengths**

FLAVOR LIMITS ON TOP-COLORON MODEL



SCALAR BOSONS

R.S.C., Elizabeth Simmons, N. Vignaroli

arXiv:1306.2248

Bogdan Dobrescu and Yang Bai

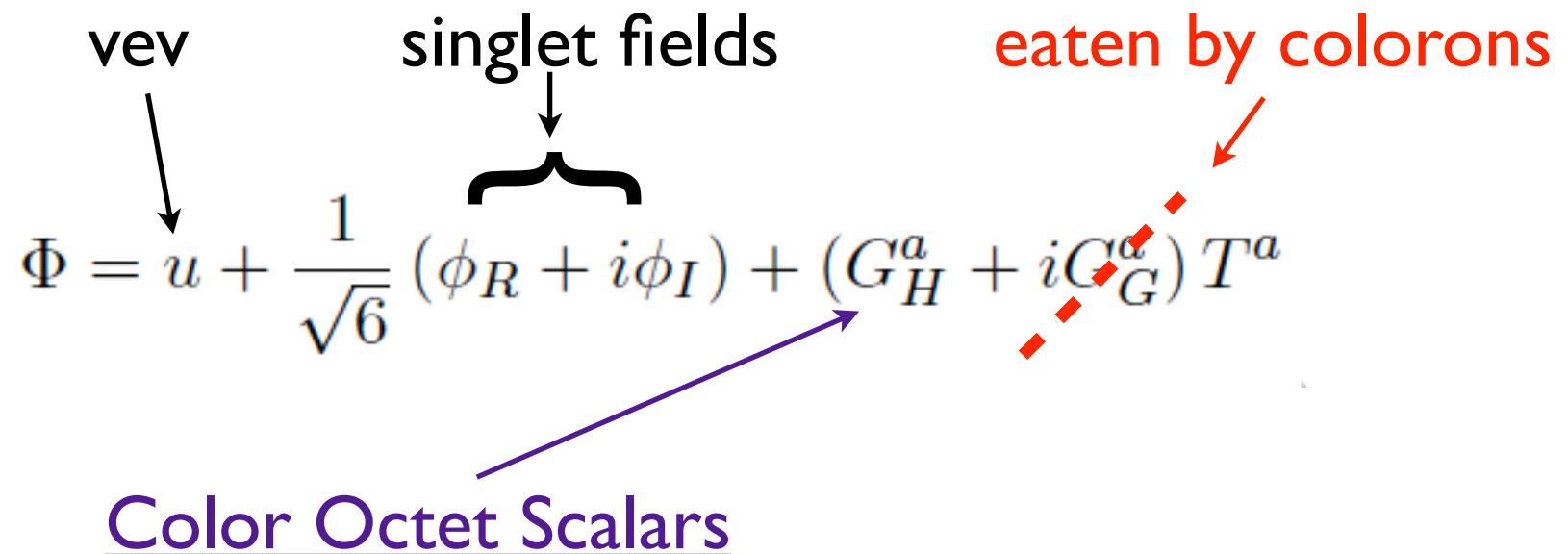
arXiv:1012.5814

COLORED SCALARS AND THEIR POTENTIAL

Most general renormalizable $(3, \bar{3})$ potential:

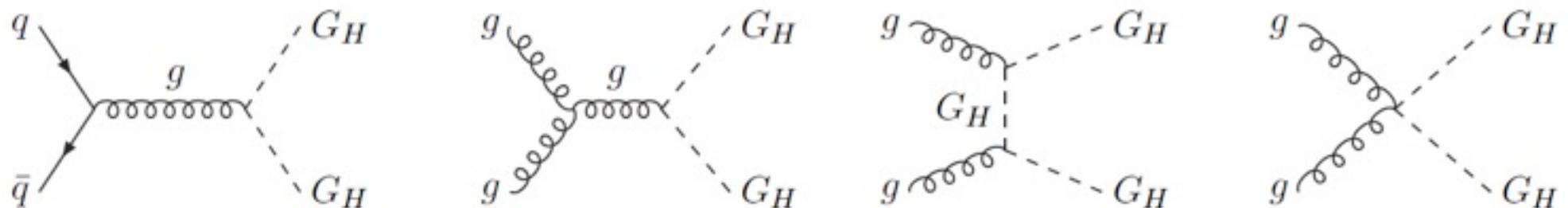
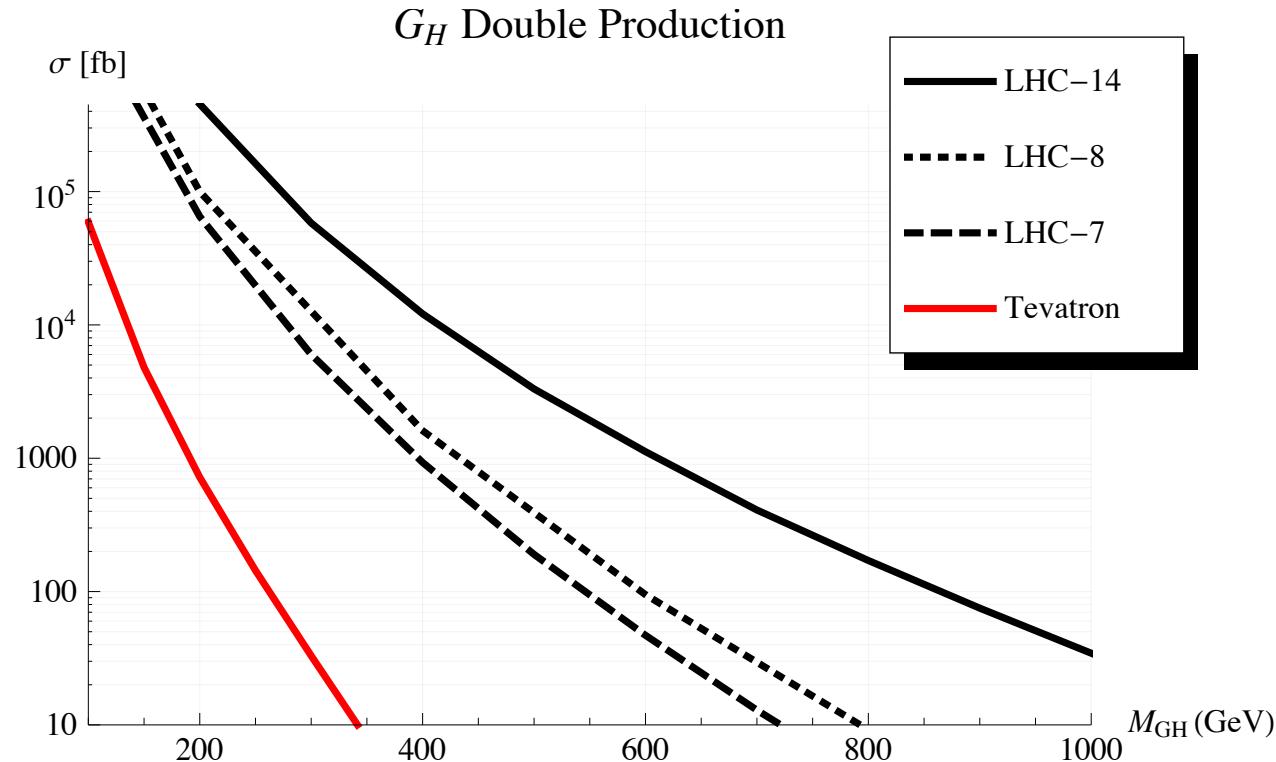
$$V(\Phi) = -m_\Phi^2 \text{Tr}(\Phi\Phi^\dagger) - \mu(\det\Phi + \text{H.c.}) + \frac{\xi}{2} [\text{Tr}(\Phi\Phi^\dagger)]^2 + \frac{k}{2} \text{Tr}(\Phi\Phi^\dagger\Phi\Phi^\dagger)$$

For an appropriate range of parameters:



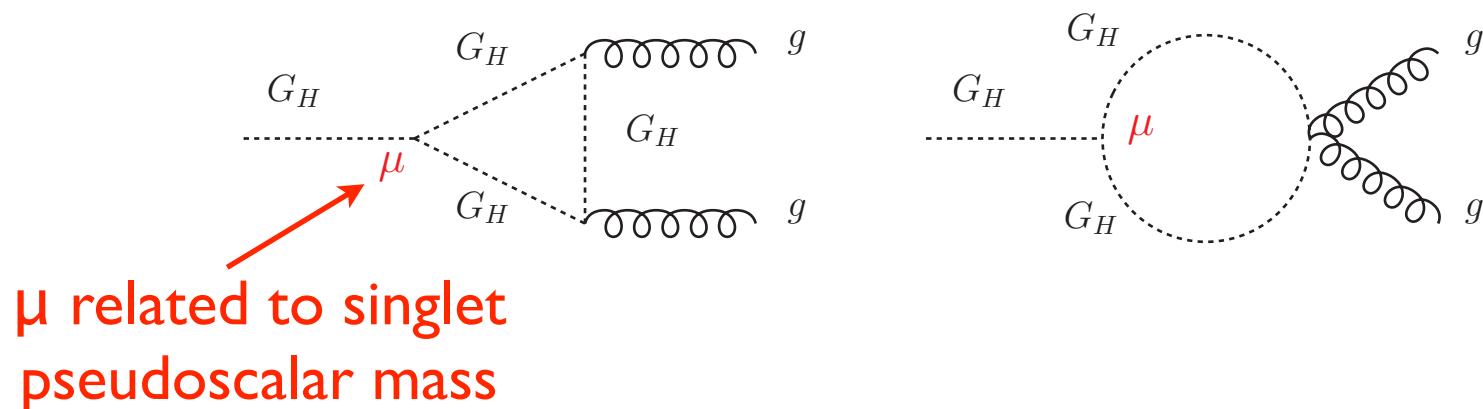
Quark couplings fixed from above!

OCTET SCALAR PRODUCTION

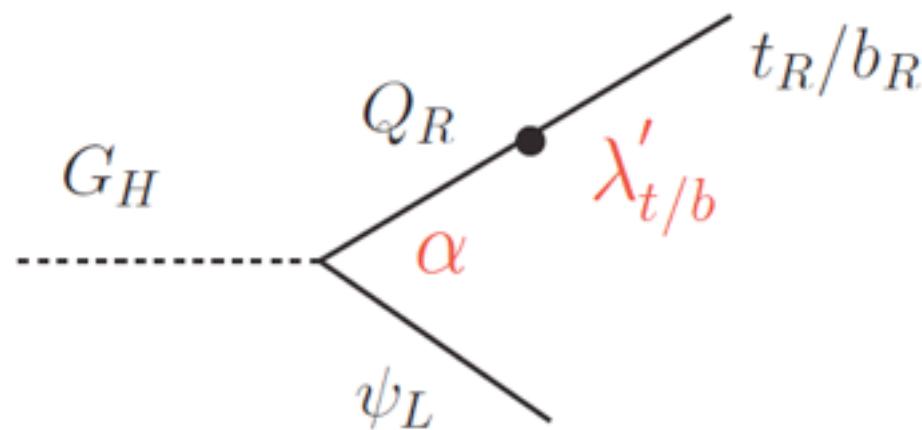


OCTET SCALAR DECAY

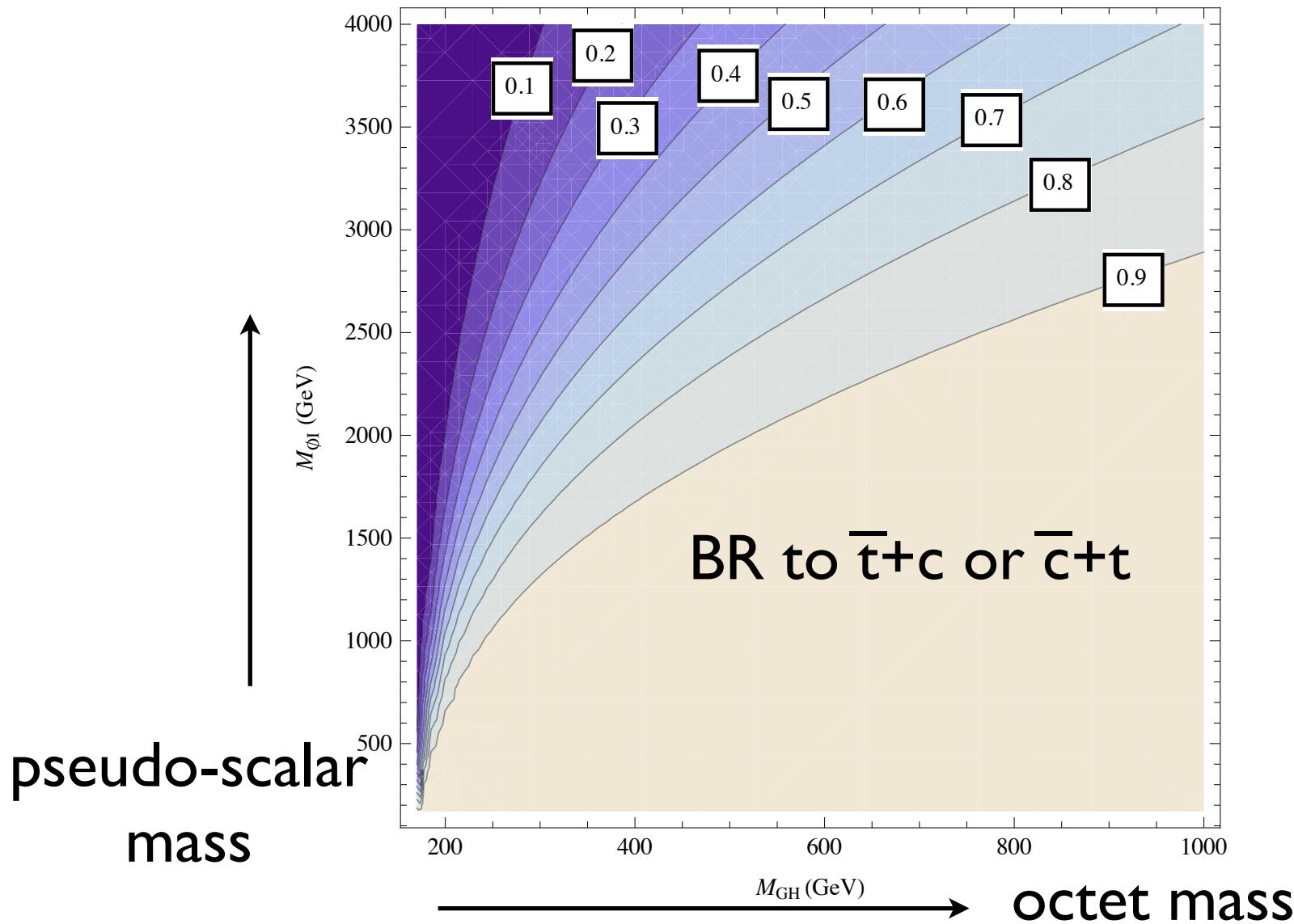
Dijets:



$\bar{c}_L t_R + \bar{t}_R c_L :$

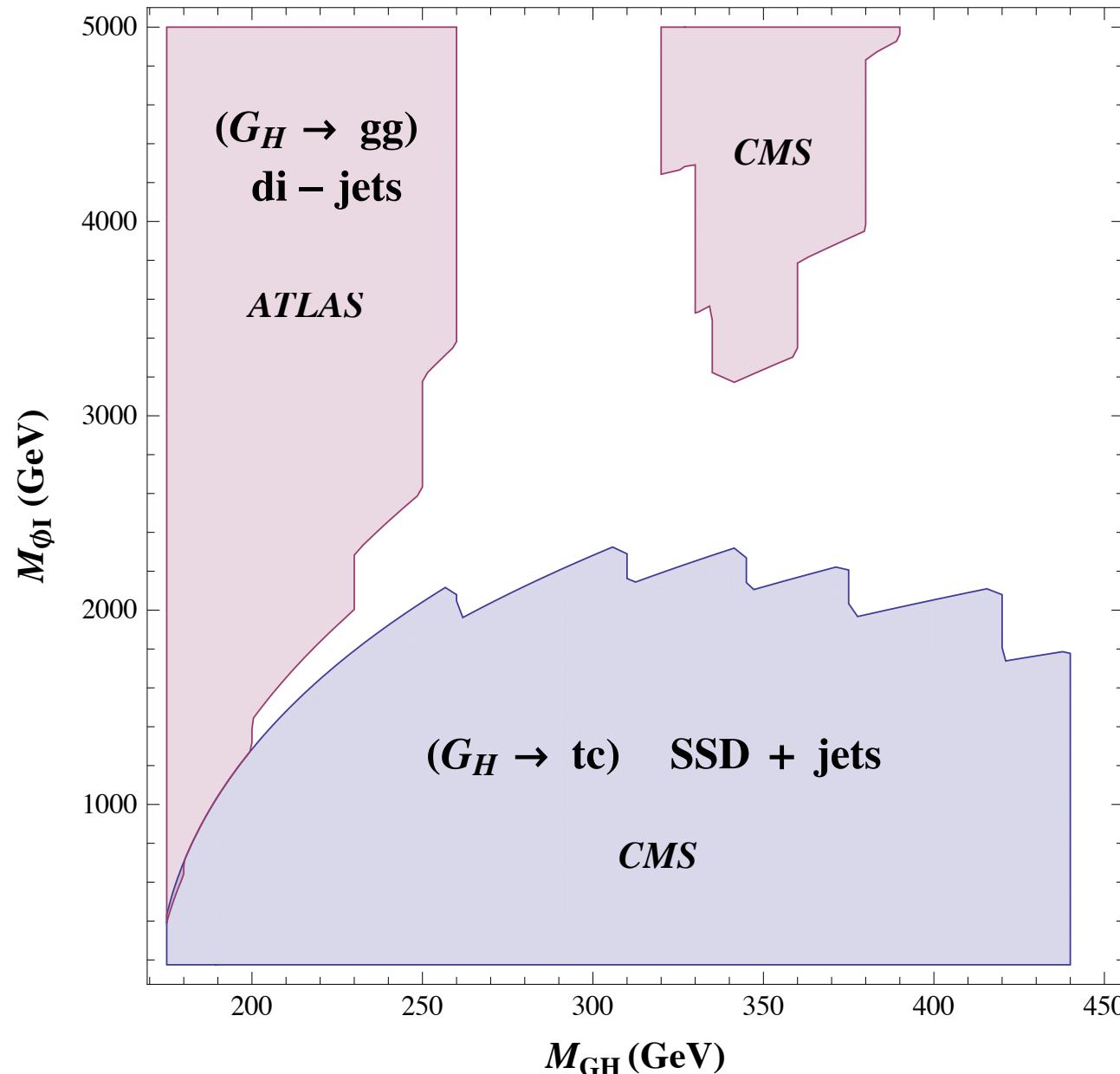


TOP + CHARM OFTEN VERY LARGE!



Octet pair production can lead to same-sign tops (dileptons)!

EXPERIMENTAL CONSTRAINTS



Singlet mass dependence from
behavior of BRs

References: CMS PAS SUS-12-029
ATLAS arXiv:1210.4826
CMS arXiv:1302.0531

FLAVOR- UNIVERSAL CONSTRAINTS ON SCALARS

R.S.C., Arsham Farzinnia, Jing Ren,
and Elizabeth Simmons
arXiv:1307.xxxx

SCALAR POTENTIAL: HIGGS AND MIXING

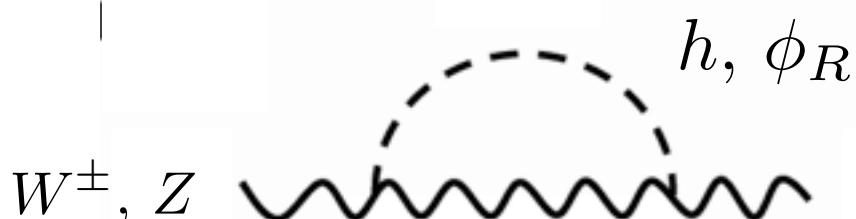
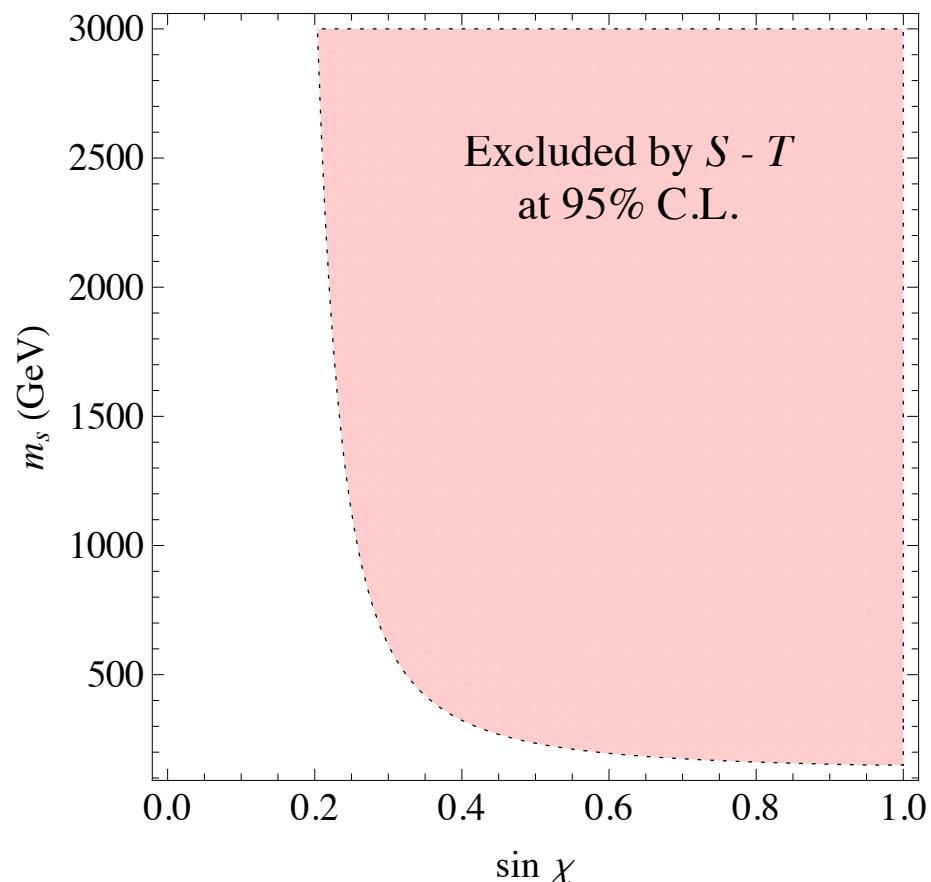
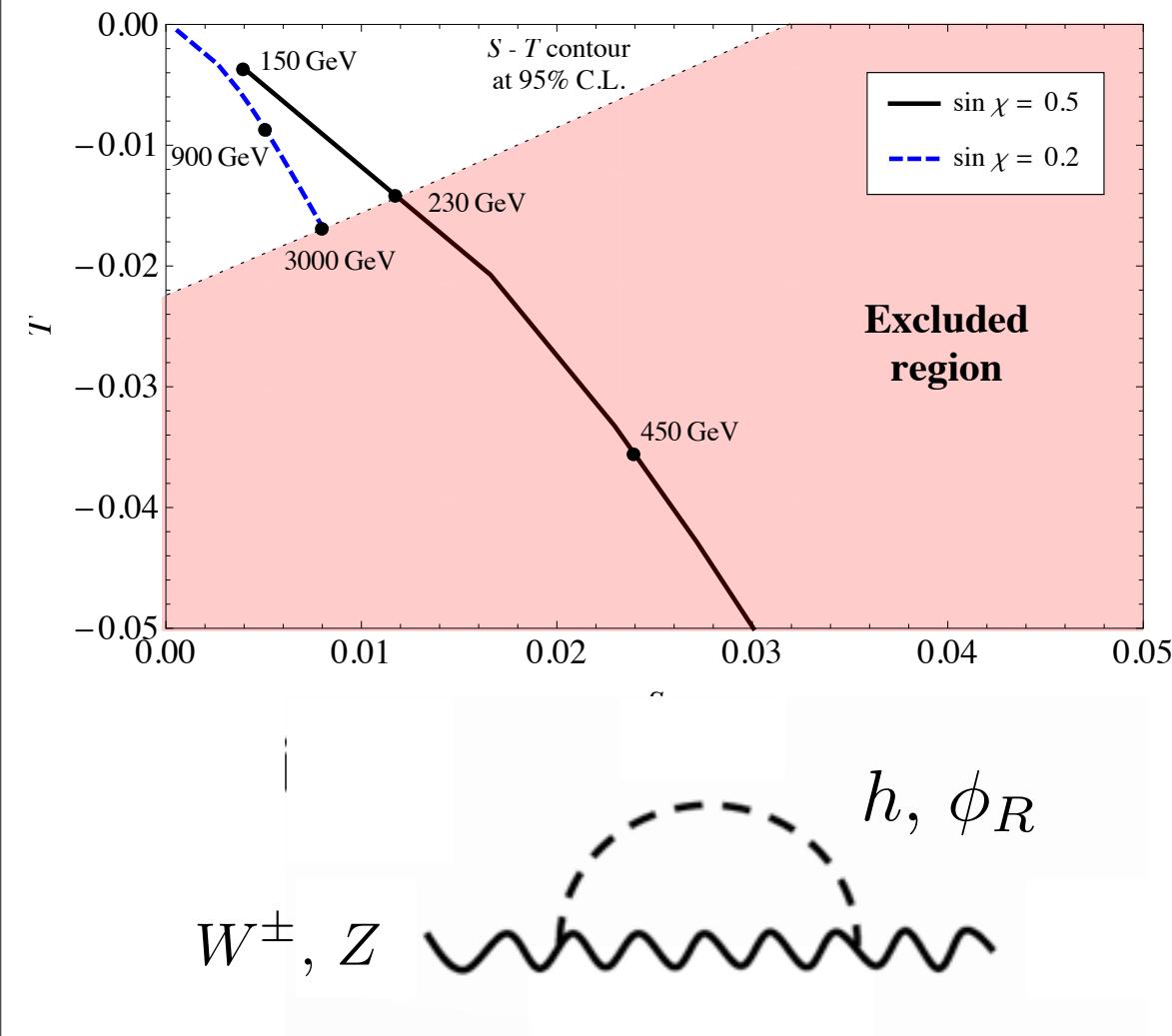
Scalar potential includes Higgs boson as well:

$$V(\phi, \Phi) \subset \frac{\lambda_h}{6} \left(\phi^\dagger \phi - \frac{v_h^2}{2} \right)^2 + \lambda_m \left(\phi^\dagger \phi - \frac{v_h^2}{2} \right) \left(\text{Tr} [\Phi^\dagger \Phi] - \frac{v_s^2}{2} \right)$$


“Higgs portal” coupling: mixing between electroweak and color sectors

$$h = \cos \chi h_0 - \sin \chi \phi_{0R}$$

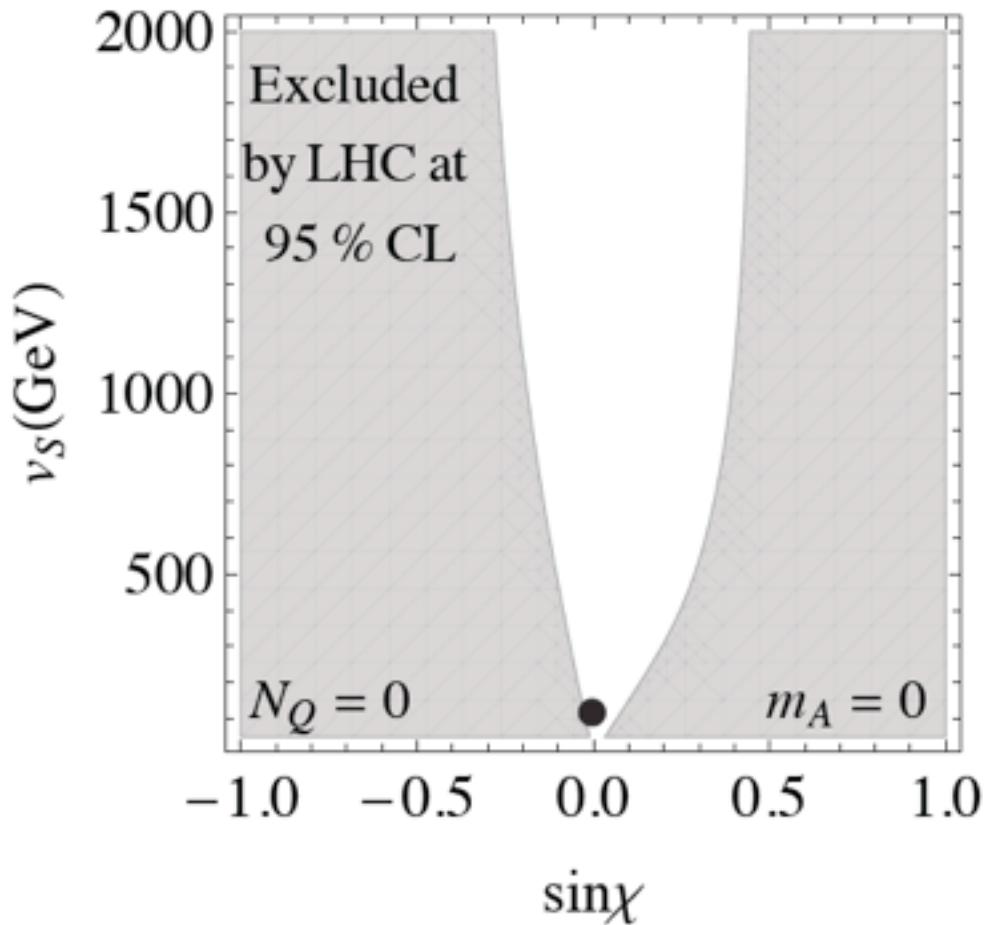
PRECISION ELECTROWEAK CONSTRAINTS



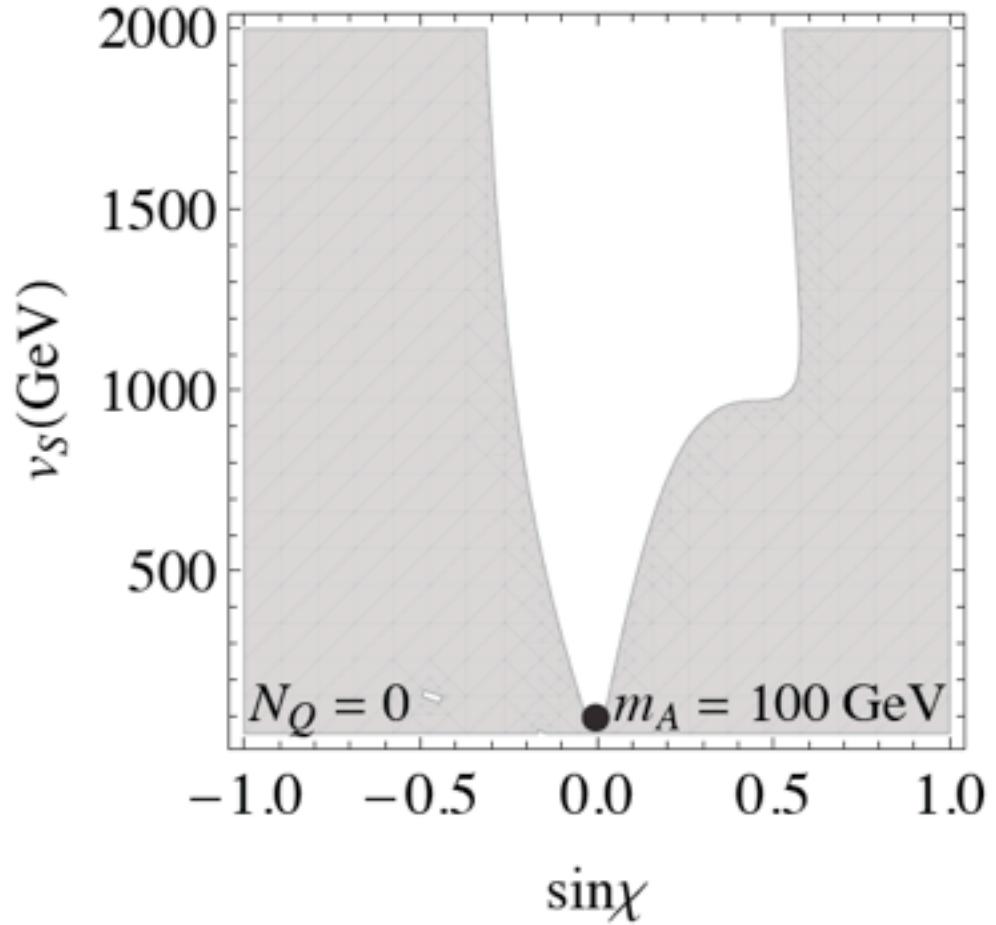
S-T contours from Gfitter, arXiv:1209.2716

CONSTRAINTS FROM HIGGS OBSERVATION

Coloron and colored scalar contributions to production...



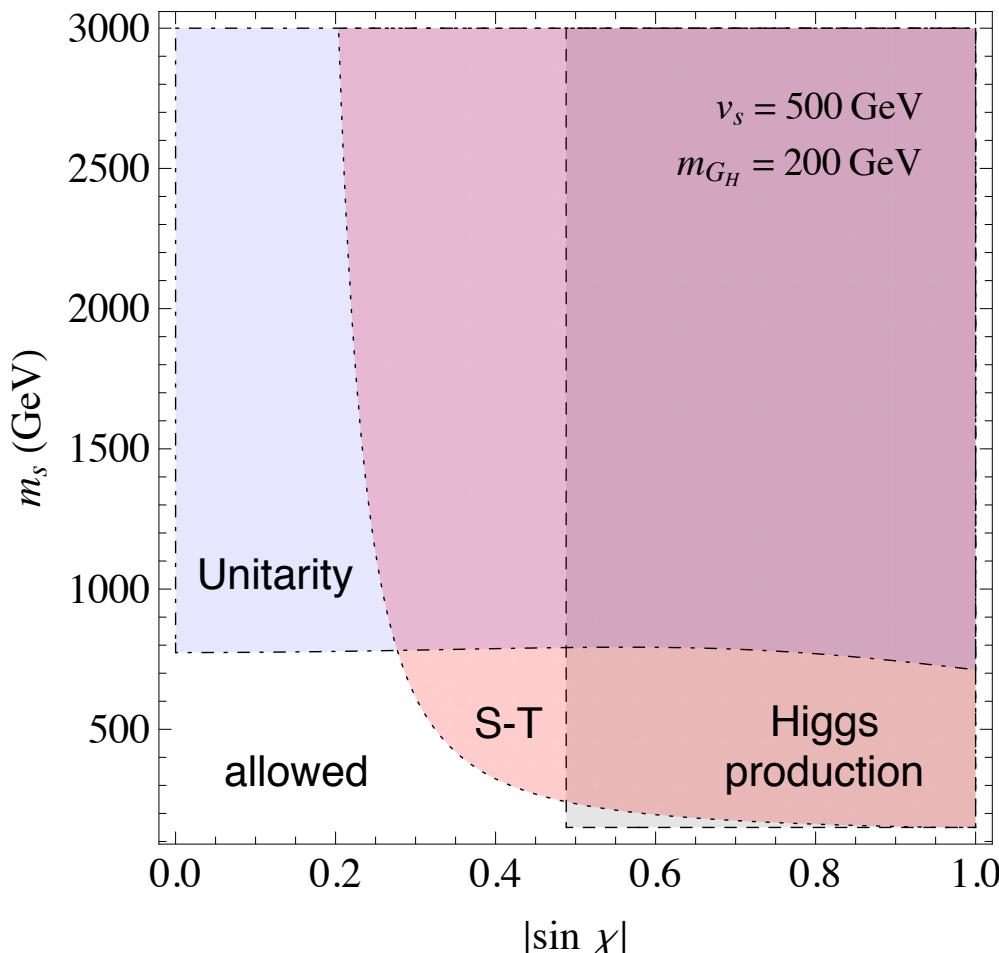
$h \rightarrow \phi_I \phi_I$ allowed



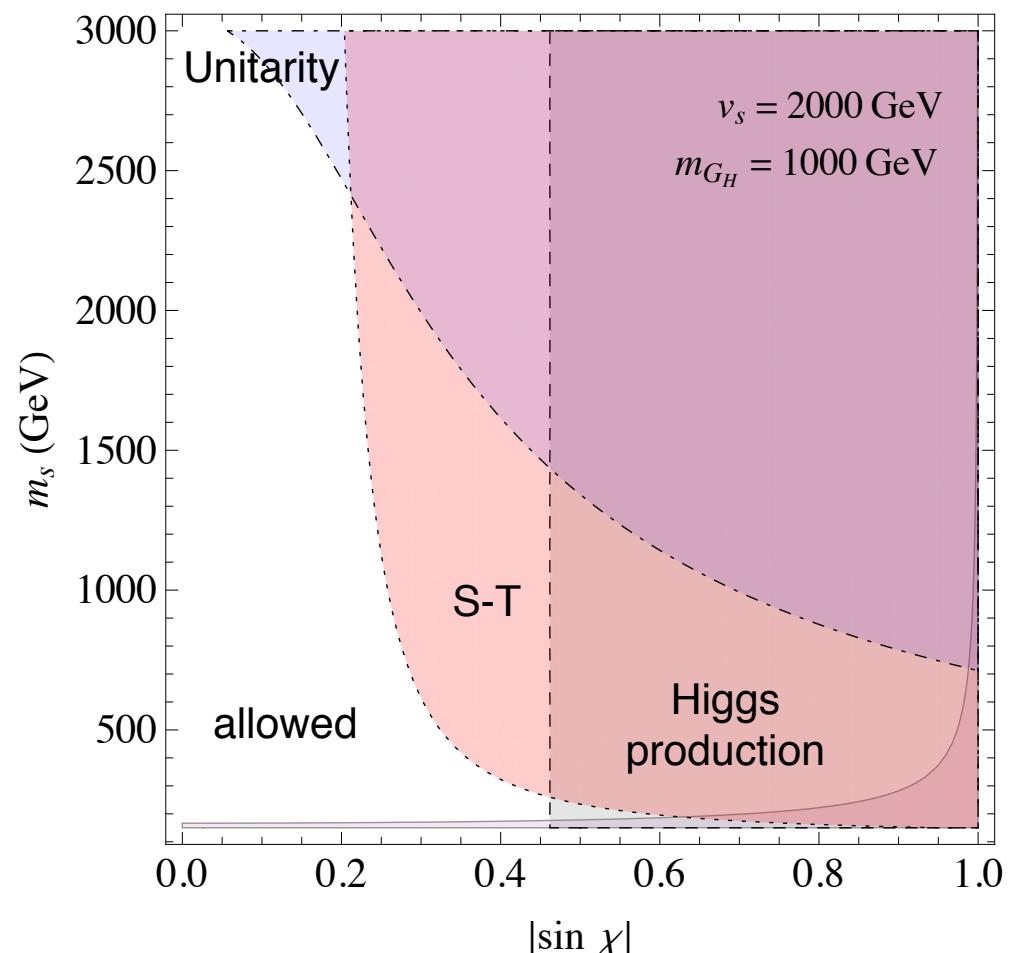
CMS-PAS-HIG-13-005
ATLAS-CONF-2013-034
Yao, Moriond EW 2013

ILLUSTRATION OF COMBINED RESULTS

$u=500 \text{ GeV}$
 $m_{GH}= 200 \text{ GeV}$



$u=2000 \text{ GeV}$
 $m_{GH}= 1000 \text{ GeV}$



Illustrates interplay of different constraints ...
and of direct and indirect bounds

CONCLUSIONS

CONCLUSIONS

Many models predict extended strong interactions

Is this extended dynamics flavor-universal or not?

- Introduced a flavorful top-coloron model
- Constraints from FCNCs favor NMFV.
- Same-sign tops, and therefore dileptons, an interesting signature for new colored scalars.

Additional effects of extended strong interactions?

- Color symmetry breaking sector can mix with EWSB
- Constraints on Higgs mixing and from observed properties of Higgs boson