Jet Substructure at the Precision Frontier

Andrew Larkoski Harvard University

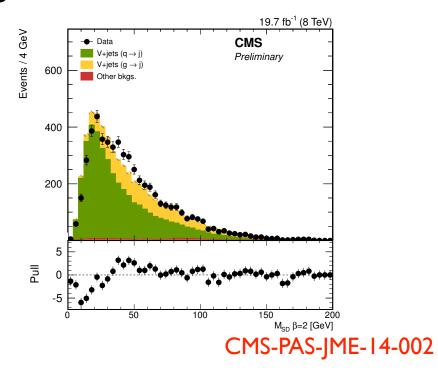
Motivation for Precision Jet Substructure

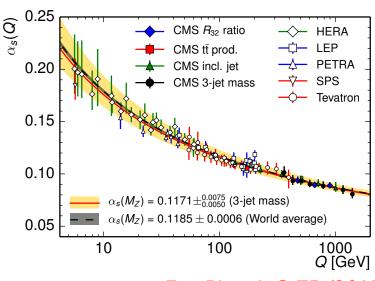
Ever increasing set of experimental measurements

Probing orthogonal regime of QCD

New α_s extractions using resummation-sensitive observables

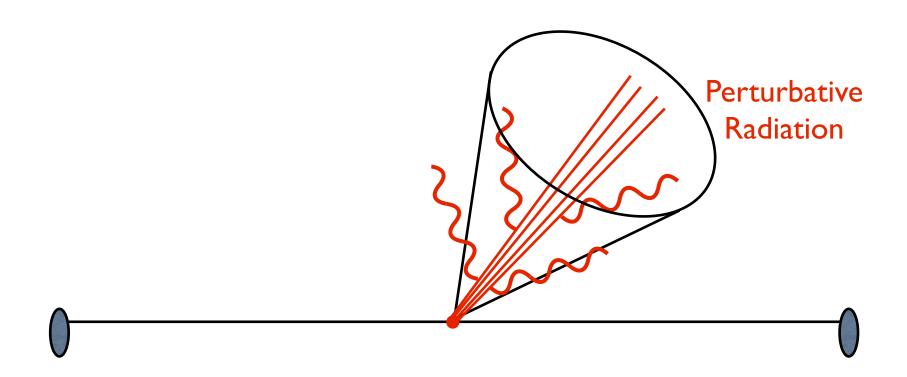
Quark and gluon jet definitions important for new physics and pdf constraints



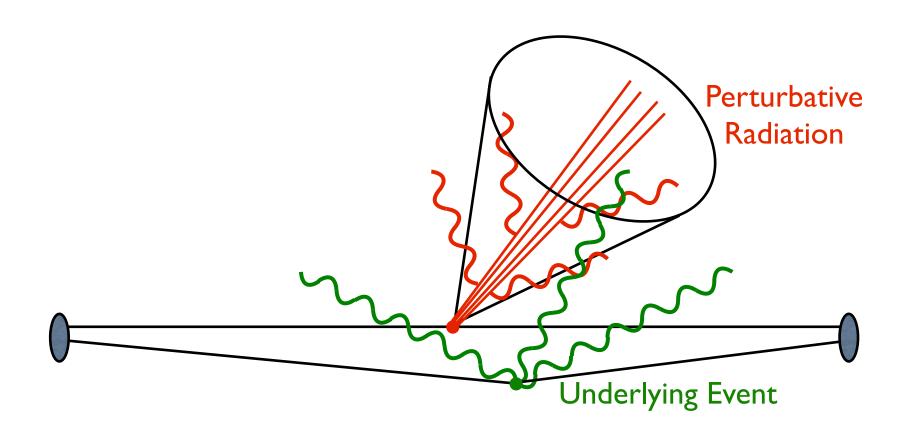


Measure m_J^2 on the jet in pp \rightarrow Z + j events

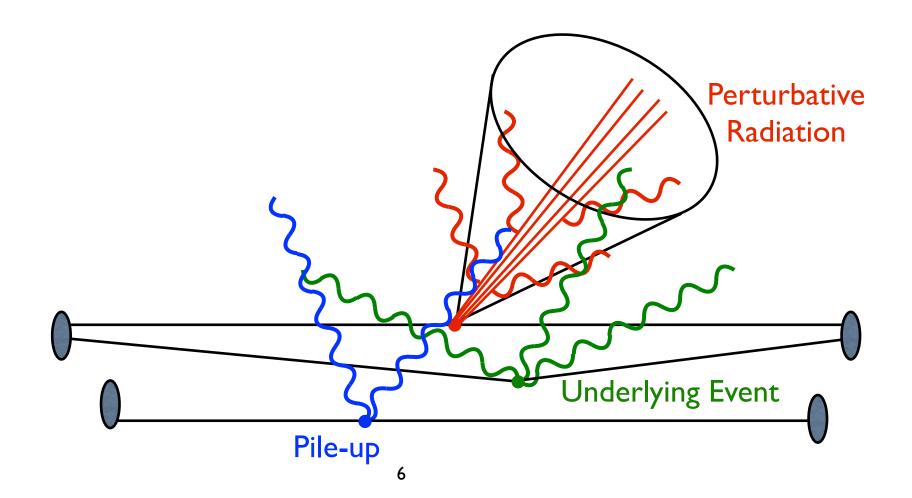
Measure m_J^2 on the jet in pp \rightarrow Z + j events



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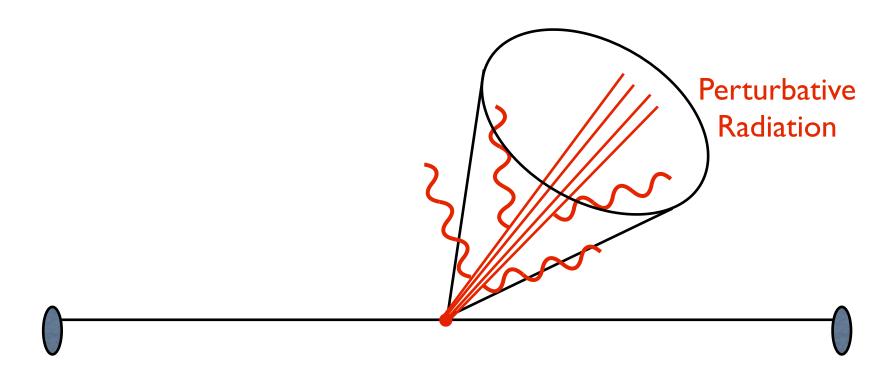


Measure m_J^2 on the jet in pp \rightarrow Z + j events

Theoretical Challenge:

Non-Global Logarithms

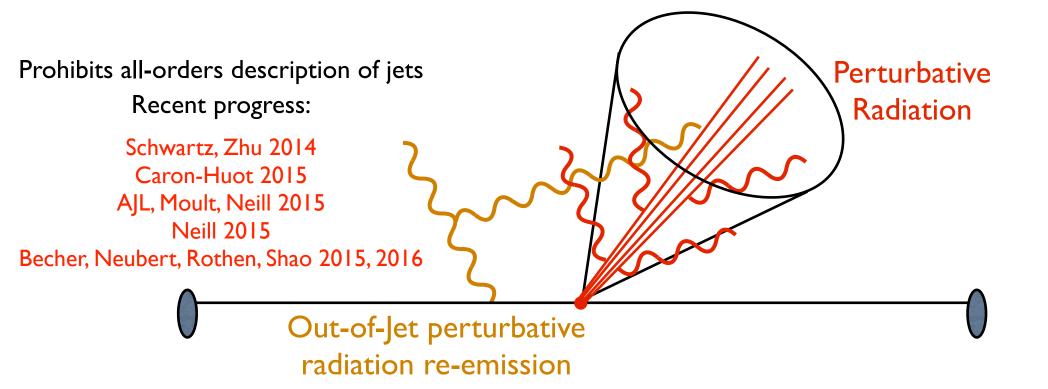
Dasgupta, Salam 2001



Measure m_J^2 on the jet in pp \rightarrow Z + j events

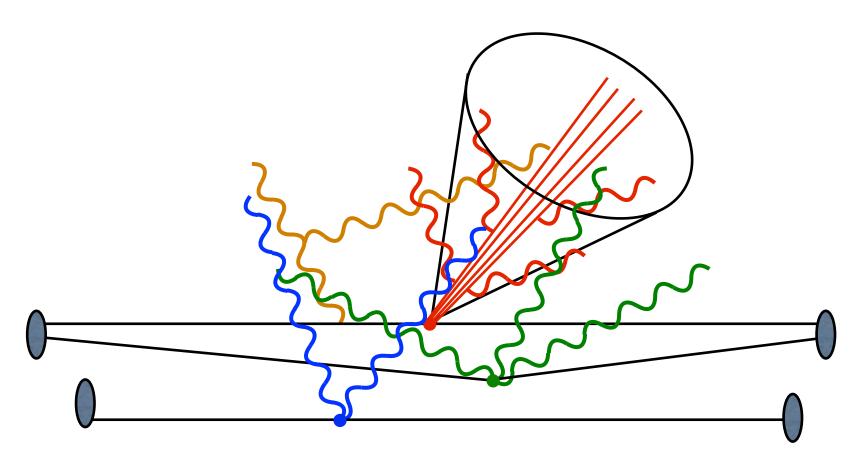
Theoretical Challenge: Non-Global Logarithms

Dasgupta, Salam 2001



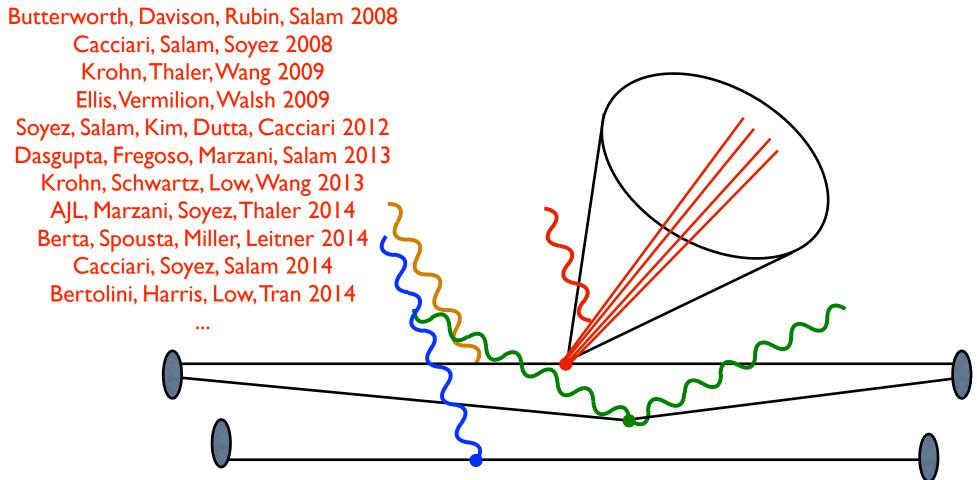
Measure m_J^2 on the jet in pp \rightarrow Z + j events

Can eliminate these problems by grooming the jet!

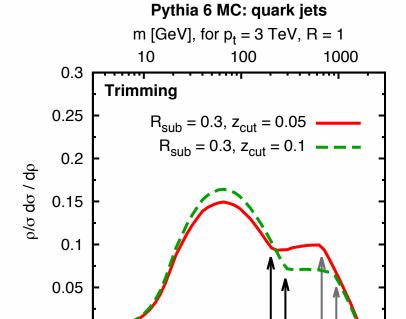


Measure m_J^2 on the jet in pp \rightarrow Z + j events

Can eliminate these problems by grooming the jet!



What has been done: NLL resummation



10⁻⁴

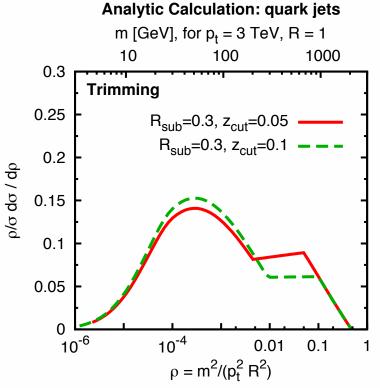
0.01

 $\rho = m^2/(p_t^2 R^2)$

0.1

0

10⁻⁶

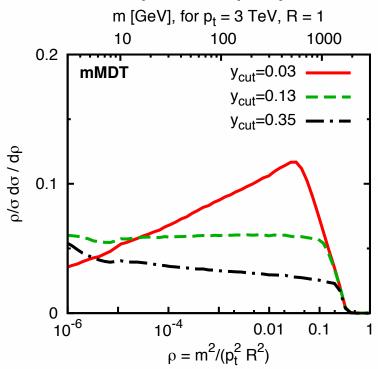


Dasgupta, Fregoso, Marzani, Salam 2013

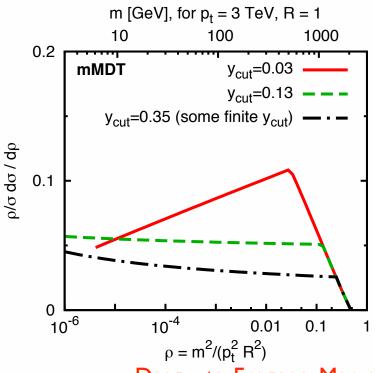
Trimming: Krohn, Thaler, Wang 2009

What has been done: NLL resummation

Pythia 6 MC: quark jets



Analytic Calculation: quark jets

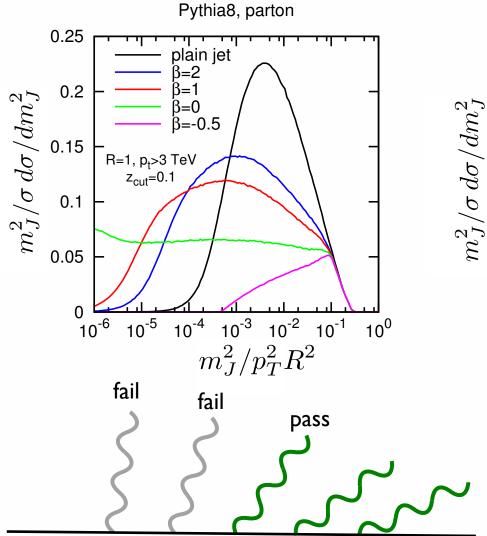


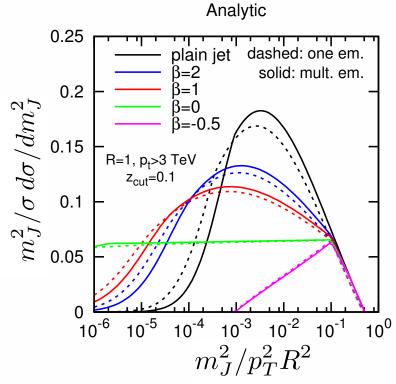
Dasgupta, Fregoso, Marzani, Salam 2013

	highest logs	transition(s)	Sudakov peak	NGLs	NP: $m^2 \lesssim$
plain mass	$\alpha_s^n L^{2n}$	_	$L \simeq 1/\sqrt{\bar{\alpha}_s}$	yes	$\mu_{\mathrm{NP}} p_t R$
trimming pruning MDT	$\alpha_s^n L^{2n}$ $\alpha_s^n L^{2n}$ $\alpha_s^n L^{2n-1}$	$z_{ m cut}, r^2 z_{ m cut} \ z_{ m cut}, z_{ m cut}^2 \ y_{ m cut}, rac{1}{4} y_{ m cut}^2, y_{ m cut}^3$	$L \simeq 1/\sqrt{\bar{\alpha}_s} - 2\ln r$ $L \simeq 2.3/\sqrt{\bar{\alpha}_s}$ —	yes yes	$\mu_{\text{NP}} p_t R_{\text{sub}}$ $\mu_{\text{NP}} p_t R$ $\mu_{\text{NP}} p_t R$
Y-pruning mMDT	$\alpha_s^n L^{2n-1}$ $\alpha_s^n L^n$	$z_{ m cut}$ $y_{ m cut}$	(Sudakov tail) —	yes no	$\mu_{\rm NP} p_t R$ $\mu_{\rm NP}^2/y_{\rm cut}$

Explicit calculations suggest better techniques!

What has been done: NLL resummation





AJL, Marzani, Soyez, Thaler 2014

Only mMDT/Soft Drop groomers eliminate NGLs!

Soft Drop:
$$\frac{\min[p_{Ti}, p_{Tj}]}{p_{Ti} + p_{Tj}} > z_{\text{cut}} \left(\frac{R_{ij}}{R}\right)^{\beta}$$

$$\beta = 0: mMDT$$

Procedure to get NNLL Resummation

Soft Drop the hardest jet in $pp \rightarrow Z + j$ events

Measure m_J^2 of the soft dropped jet:

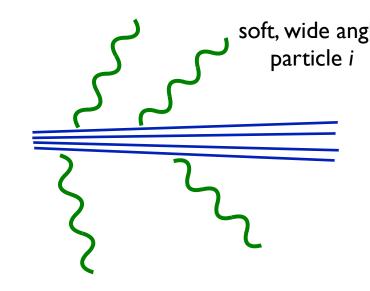
$$\frac{\min[p_{Ti}, p_{Tj}]}{p_{Ti} + p_{Tj}} > z_{\text{cut}} \left(\frac{R_{ij}}{R}\right)^{\beta}$$

$$m_J^2 \simeq \sum_{i < j \in J} p_{Ti} p_{Tj} R_{ij}^2$$

Focus on the regime where:

$$m_J^2 \ll z_{\rm cut} p_{TJ}^2 \ll p_{TJ}^2$$

All remaining particles in the jet must be collinear!



soft, wide angle particle i
 I)
$$\frac{p_{Ti}}{p_{TJ}} \sim z_{\rm cut} \longrightarrow m_J^2 \sim z_{\rm cut} p_{TJ}^2$$

2)
$$\frac{p_{Ti}}{p_{TJ}} \sim \frac{m_J^2}{p_{TJ}^2} \longrightarrow \text{groomed away}$$

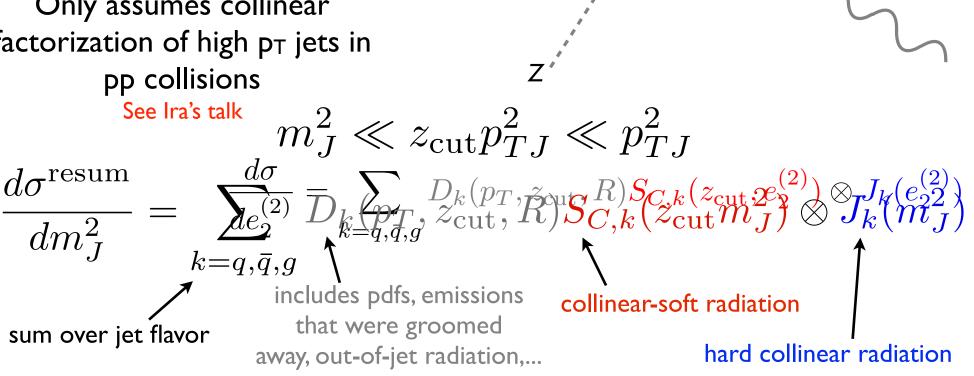
Factorization for NNLL Resummation

Effective theory for soft drop groomed jets

Frye, AJL, Schwartz, Yan 2016

Coefficient D_k can be extracted from fixed-order

Only assumes collinear factorization of high p_T jets in pp collisions



Matching NNLL to α_s^2

$$\frac{d\sigma^{\text{NNLL}+\alpha_s^2}}{dm_J^2} \equiv \frac{d\sigma^{\text{NNLL}}}{dm_J^2} + \frac{d\sigma^{\alpha_s^2}}{dm_J^2} - \frac{d\sigma^{\text{NNLL},\alpha_s^2}}{dm_J^2}$$

Use MCFM to generate relative α_s^2 cross section

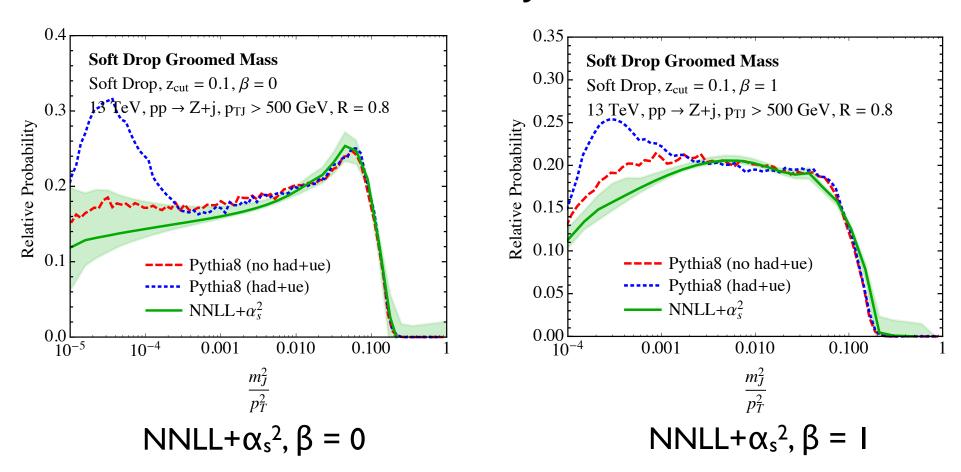
$$pp \rightarrow Z$$
 + j at NNLO with $m_J^2 > 0 = pp \rightarrow Z$ + 2j at NLO

Required extreme computing power:

To make the following plots required centuries of CPU time

The very first jet substructure calculation at high precision!

Results: NNLL+ α_s^2 Jet Substructure



Comparison with Pythia8 Monte Carlo

Hadronization and underlying event only dominate for $m_J^2/p_T^2 \lesssim 10^{-3}$

Almost three decades of perturbative control in a single jet distribution!

Summary

Precision calculations for jet substructure requires jet grooming

Only mMDT/soft drop remove contamination and eliminate NGLs

All radiation that remains in the jet is collinear

NNLL resummation of groomed jet mass is accomplished

Looking Ahead

mMDT/soft drop also makes quark/gluon jet flavor IRC safe

Improved input into pdfs?

Monte Carlo tuning to gluon jets?

Tuning Monte Carlos to precision calculations?

Precision jet substructure measurements?

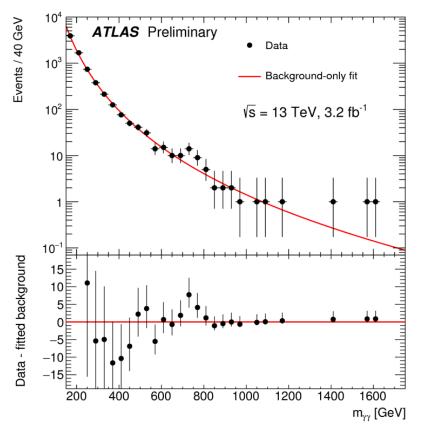
Motivation for ATLAS and CMS to make identical jet measurements? Possible for systematics at %-level?

Feedback to Fixed-Order Community

Jet substructure observables are sensitive to infrared phase space region Need very efficient methods to sample deep infrared Generic phase space reweighting à la EVENT2?

Bonus Slides

Jet Substructure



Should also see evidence in W/Z/H channels

Requires identification from decays

$$\frac{\Gamma_{Z \to l^+ l^-}}{\Gamma_Z} \simeq 10\%$$

$$\frac{\Gamma_{Z \to \mathrm{had}}}{\Gamma_Z} \simeq 70\%$$

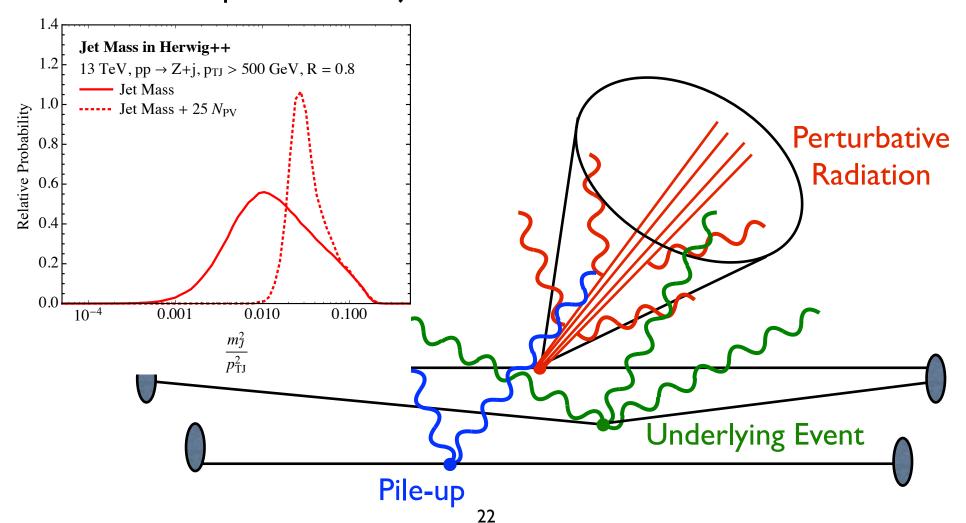
Understanding hadronic decays is essential!

ATLAS-CONF-2015-081

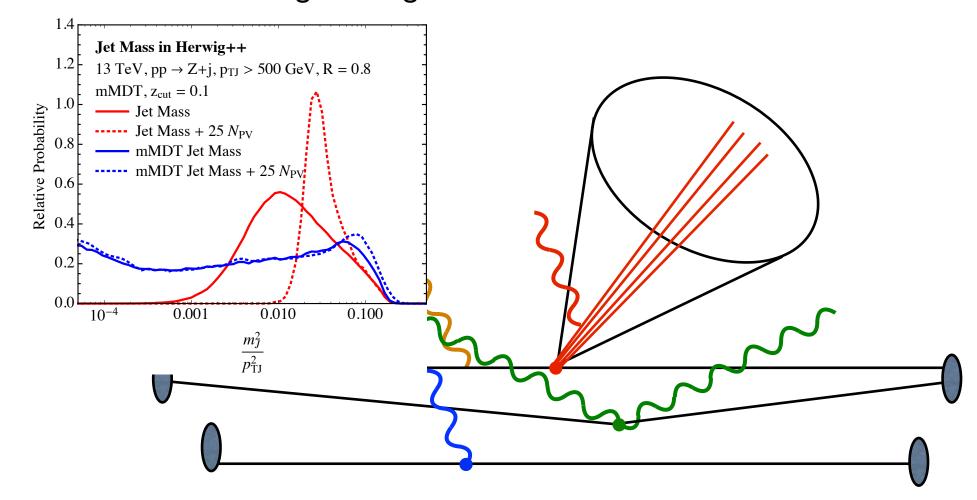
Background to boosted hadronically-decaying EW bosons are massive QCD jets!

Need precise understanding, definition, theory predictions of QCD jets at the LHC

Measure m_J^2 on the jet in pp \rightarrow Z + j events



Measure mMDT/soft drop m_J^2 on the jet in pp \rightarrow Z + j events Removes contamination radiation from the jet Eliminates non-global logarithms



Aside: Getting Collinear-Soft Function to NNLL

Factorization theorem in e⁺e⁻ collisions:

$$S_G(z_{\mathrm{cut}}Q^2) \qquad \log \frac{1}{z} \qquad \uparrow \text{ soft} \qquad \downarrow \log \frac{1}{z_{\mathrm{cut}}} \qquad \downarrow \log \frac{1}{z_{\mathrm{cut}}}$$

$$\frac{d^2\sigma}{dm_{J,L}^2 dm_{J,R}^2} = H(Q^2) S_G(z_{\text{cut}}Q^2) \left[S_C(z_{\text{cut}}m_{J,L}^2) J(m_{J,L}^2) \right] \left[S_C(z_{\text{cut}}m_{J,R}^2) J(m_{J,R}^2) \right]$$

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 $H(Q^2)$: Hard function for $e^+e^- \rightarrow qq$. Known beyond two-loops.

van Neerven 1986

Matsuura, van der Marck, van Neerven 1989

 $J(m_J^2)$: Jet function. Known at two-loops for quarks and gluons.

Bauer, Manohar 2003 Becher, Neubert 2006

 $S_G(z_{
m cut}Q^2)$: Global soft function. Related to two-loop soft function with energy veto (up to calculable clustering effects).

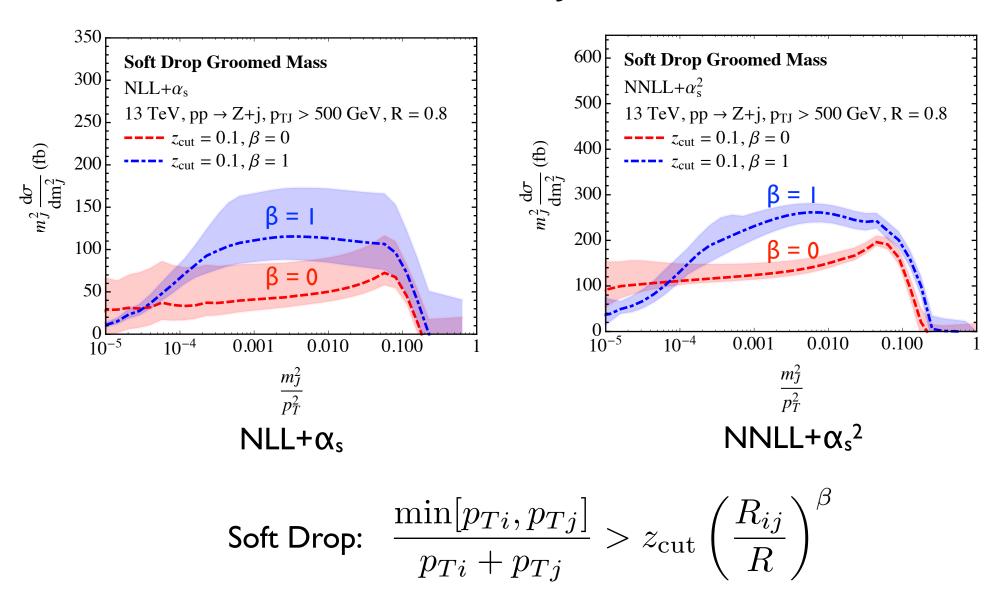
von Manteuffel, Schabinger, Zhu 2013 Chien, Hornig, Lee 2015

 $S_C(z_{
m cut}m_J^2)$: Collinear-soft function. New, no two-loop calculation exists.

Can get everything from literature and by exploiting RG invariance!

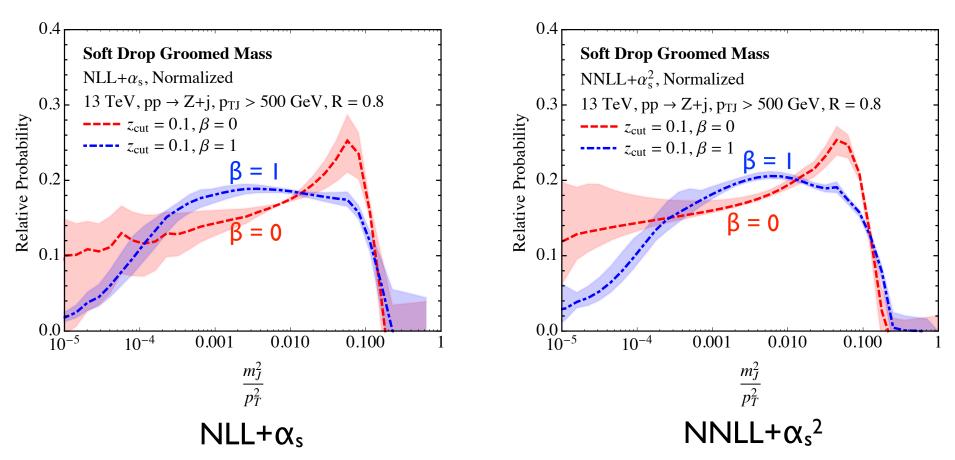
$$0 = \gamma_H + \gamma_{S_G} + 2\gamma_J + 2\gamma_{S_C}$$

Results: NNLL+ α_s^2 Jet Substructure



Significant decrease in residual scale uncertainty at NNLL+ α_s^2 !

Results: NNLL+ α_s^2 Jet Substructure



Shape of distribution only depends on collinear physics

$$\frac{d\sigma^{\text{resum}}}{dm_J^2} = \sum_{k=q,\bar{q},g} D_k(p_T, z_{\text{cut}}, R) S_{C,k}(z_{\text{cut}} m_J^2) \otimes J_k(m_J^2)$$

<10%-level residual scale uncertainties in normalized distributions!