

*Implications of charge asymmetry in the  $B_s$   
System*

July 6, 2011

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University of Washington

# OUTLINE

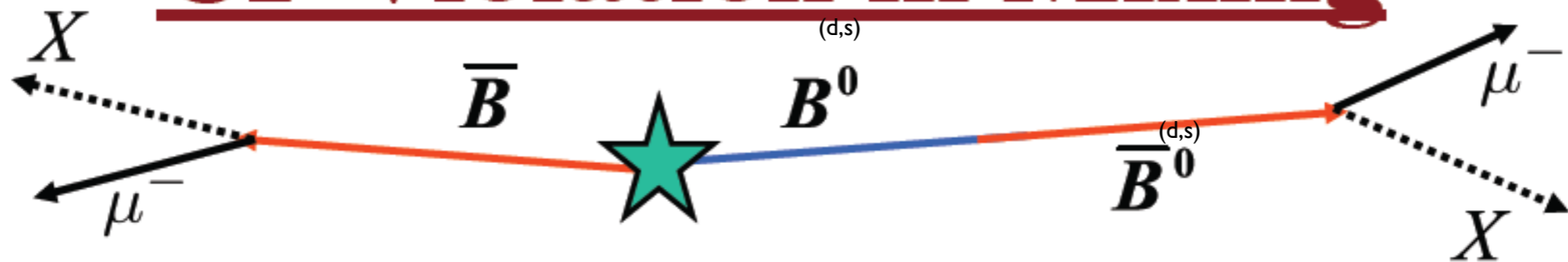
- ➔ reviewing the situation:
  - ➔ recent  $D\bar{0}$
  - ➔ upcoming LHCb
  - ➔ SM theory
- ➔ BSM theory expectations:
  - ➔ new physics in  $B_s$  mixing
  - ➔ new physics in B decay

# Experiment update

“Evidence for an anomalous like sign dimuon charge asymmetry”  
 the DØ collaboration, arXiv:1005.2757  $6.1 \text{ fb}^{-1}$

“Measurement of the anomalous like-sign dimuon charge asymmetry with  $9 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions”, arXiv:1106.6308

## CP Violation in Mixing



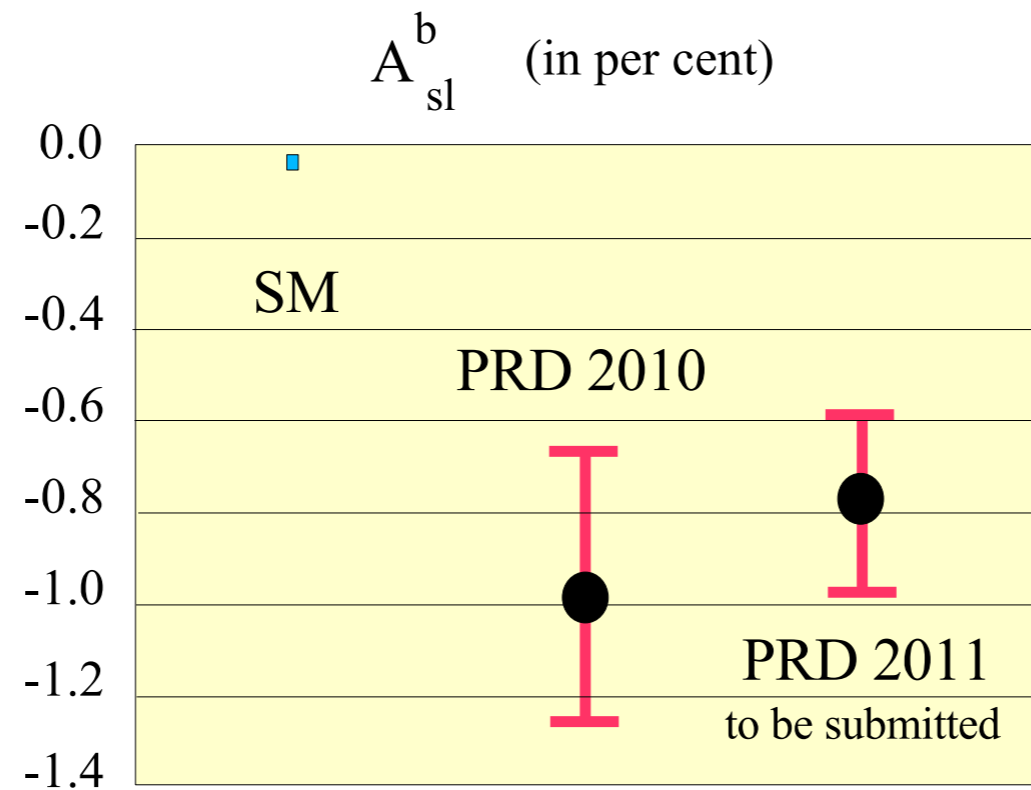
$$A_{sl}^b \equiv \frac{N_b^{++} - N_b^{--}}{N_b^{++} + N_b^{--}}$$

“wrong sign” B decay  
 from oscillation  
 gives like sign dimuon

➡ DØ:  $A_{sl}^b = (-0.00787 \pm 0.00172 \text{ (stat)} \pm 0.00093 \text{ (syst)})$

➡ differs by  $3.9\sigma$  from  $A_{sl}^b(\text{SM}) = -0.00028 \pm 0.00005$

# Comparison with last year



from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
Fermilab, 30 June 2011

Comparison of measurements of  $A_{sl}^b$ .

# Improvements (since Phys. Rev. D 82, 032001, (2010))

- To **increase the number of events**, the  $|p_z|$  cut is lowered from 6.4 GeV to 5.4 GeV.
- To **lower the  $K \rightarrow \mu$  and  $\pi \rightarrow \mu$  backgrounds**, the  $\chi^2$  of the match of track parameters obtained with the central detector and outer muon system is reduced from 40 to 12 (with 4 d.o.f.).
- The measurement of  $f_K$  is improved:  $K_S \rightarrow \pi\pi \rightarrow \mu$  (muon required for same sample composition as  $K \rightarrow \mu$ ).
- The measurement of  $R_K \equiv F_K/f_K$  is done in **two independent channels**:  $K^{*0} \rightarrow \pi^- K^+ \rightarrow \mu^+ X$  (with the **null-fit method**), and the new channel  $K_S \rightarrow \pi\pi \rightarrow \mu$ .
- The data set is increased from  $6.1 \text{ fb}^{-1}$  to  **$9.0 \text{ fb}^{-1}$** .

from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
Fermilab, 30 June 2011

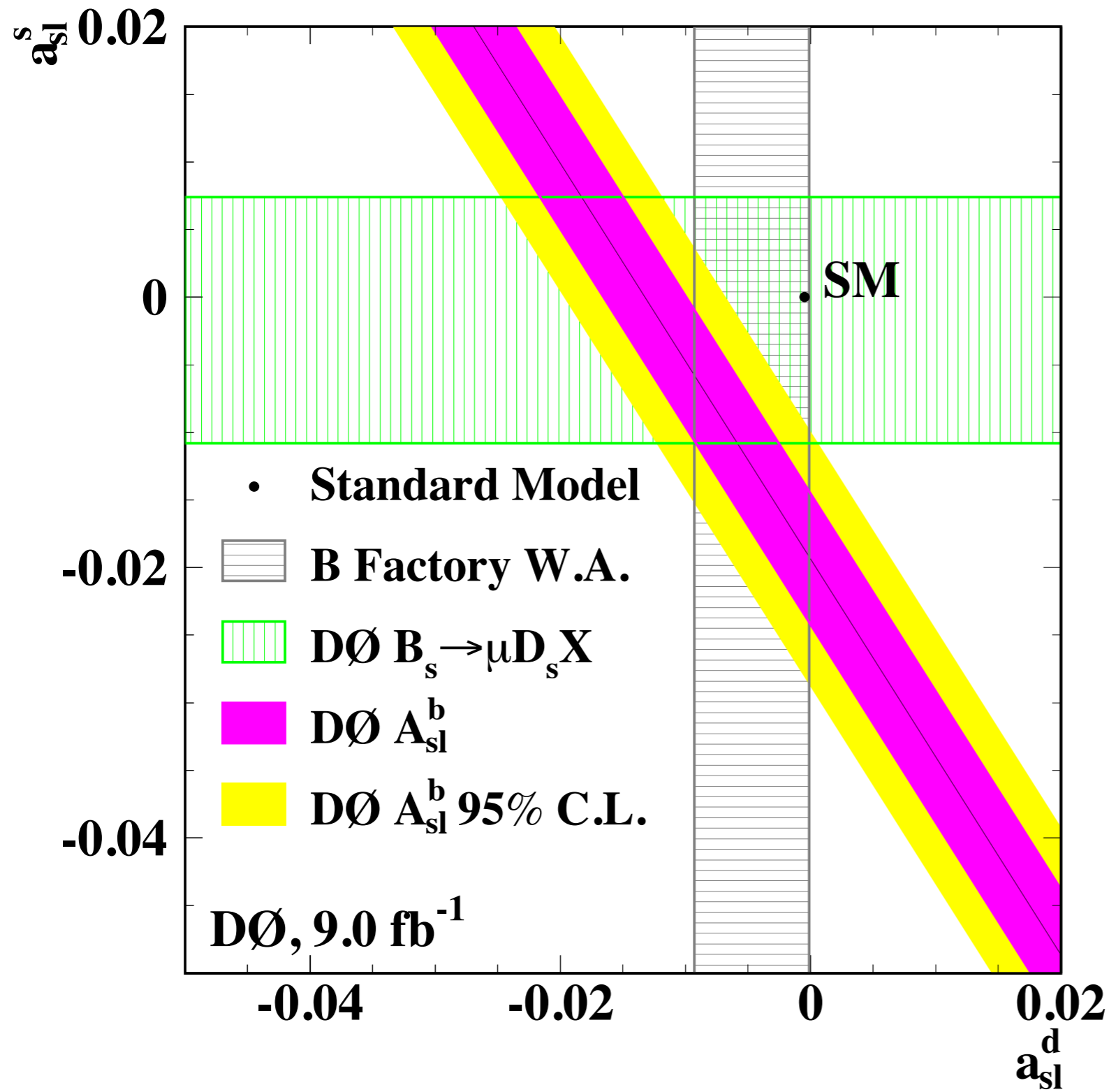
## 2. Results with $9.0 \text{ fb}^{-1}$

- From 1  $\mu$  ( $2.041 \times 10^9$  muons):  
 $A_{\text{SI}}^b = (-1.04 \pm 1.30 \text{ (stat)} \pm 2.31 \text{ (syst)}) \%$ .
- From 2  $\mu$  ( $6.019 \times 10^6$  like-sign dimuons):  
 $A_{\text{SI}}^b = (-0.808 \pm 0.202 \text{ (stat)} \pm 0.222 \text{ (syst)}) \%$ .
- $A_{\text{SI}}^b = (-0.787 \pm 0.172 \text{ (stat)} \pm 0.093 \text{ (syst)}) \%$ .  
This measurement disagrees with the prediction of the Standard Model by 3.9 standard deviations.
- The charge asymmetry of like-sign dimuon events after subtracting all background contributions from the raw charge asymmetry is:

$$\begin{aligned} A_{\text{res}} &\equiv (A - \alpha a) - (A_{\text{bkg}} - \alpha a_{\text{bkg}}) \\ &= (-0.246 \pm 0.052 \text{ (stat)} \pm 0.021 \text{ (syst)})\%. \end{aligned}$$

This quantity does not depend on the interpretation in terms of the charge asymmetry of semileptonic decays of  $B$  mesons. This measurement disagrees with the prediction of the Standard Model by 4.2 standard deviations.

from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
Fermilab, 30 June 2011



from Bruce Hoeneisen  
 representing the  
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## 4. Dependence on the impact parameter

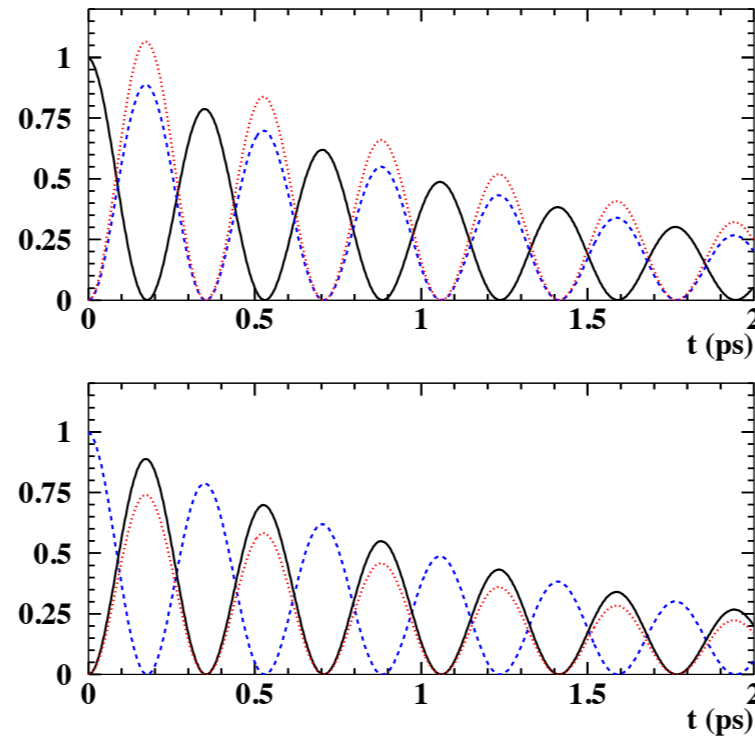
Additional measurements are made applying an impact parameter ( $IP$ ) cut on **each** muon.

$IP$  is the distance of closest approach of the muon track to the primary vertex projected onto the plane transverse to the  $p\bar{p}$  beams.

The dependence of  $A_{SI}^b = C_d a_{SI}^d + C_s a_{SI}^s$  on  $IP$  can reveal the origin of the asymmetry because  $C_d$  and  $C_s$  depend on  $IP$ .

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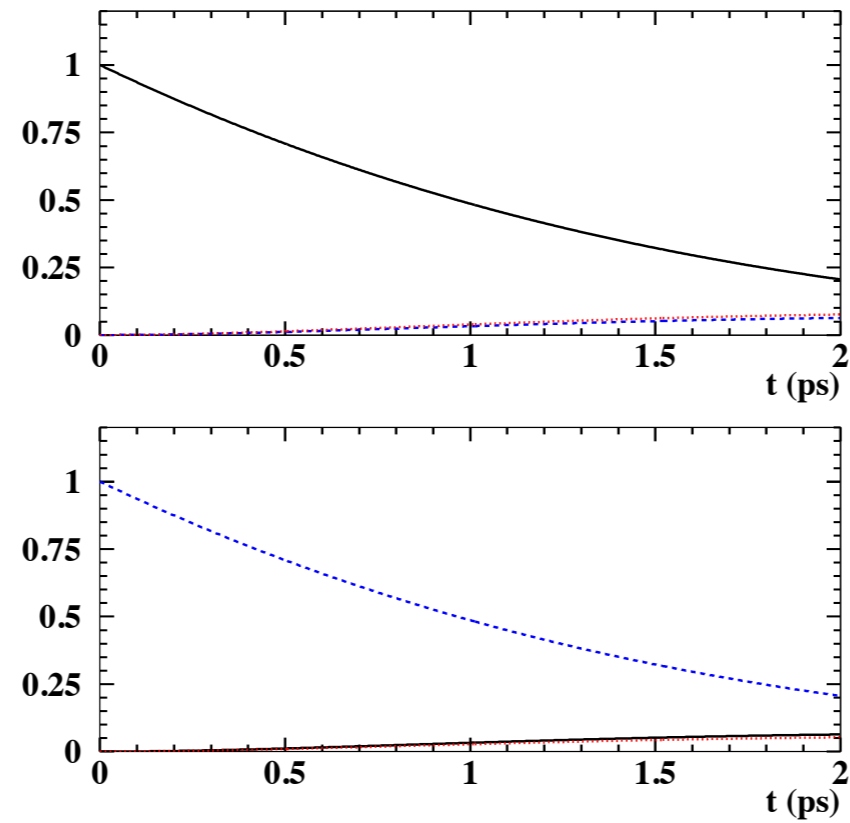
from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
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Top: Histogram of proper time of decays  $B_s^0 \rightarrow \mu^+ X$  (continuous line),  
 $B_s^0 \rightarrow \bar{B}_s^0 \rightarrow \mu^- X$  (dashed line if no CP violation, dotted red line if CP violation).  
 Bottom: The same for  $\bar{B}_s^0$  at  $t = 0$ .

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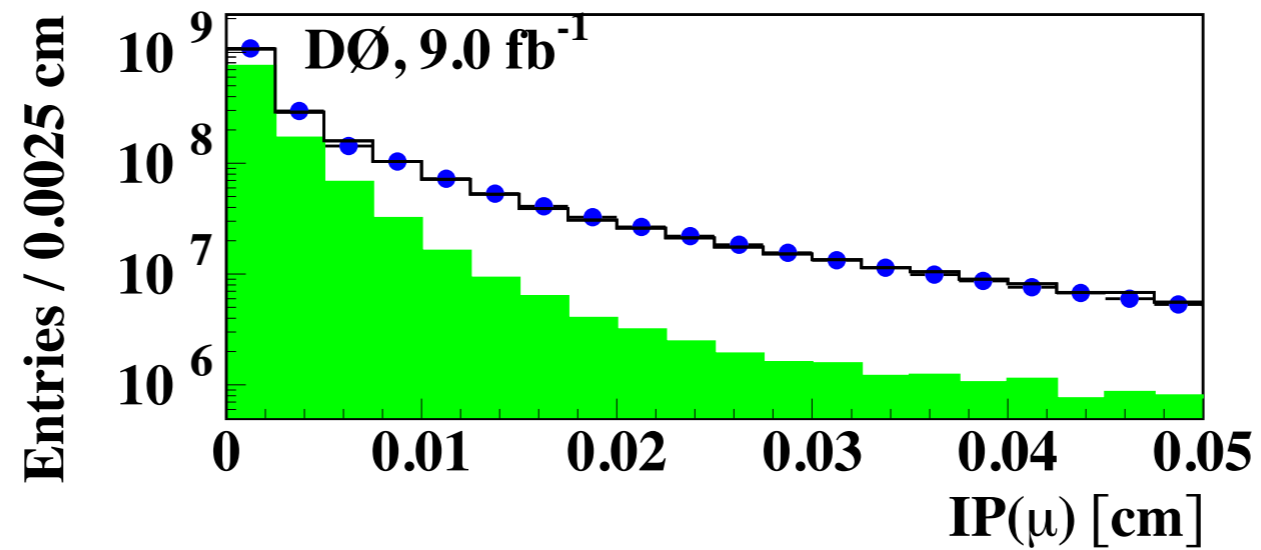
from Bruce Hoeneisen  
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Same for  $B_d^0$  (top) and  $\bar{B}_d^0$  (bottom) at  $t = 0$ . Applying an IP cut can enrich the sample in oscillating  $B_d^0$ 's (shown in red).

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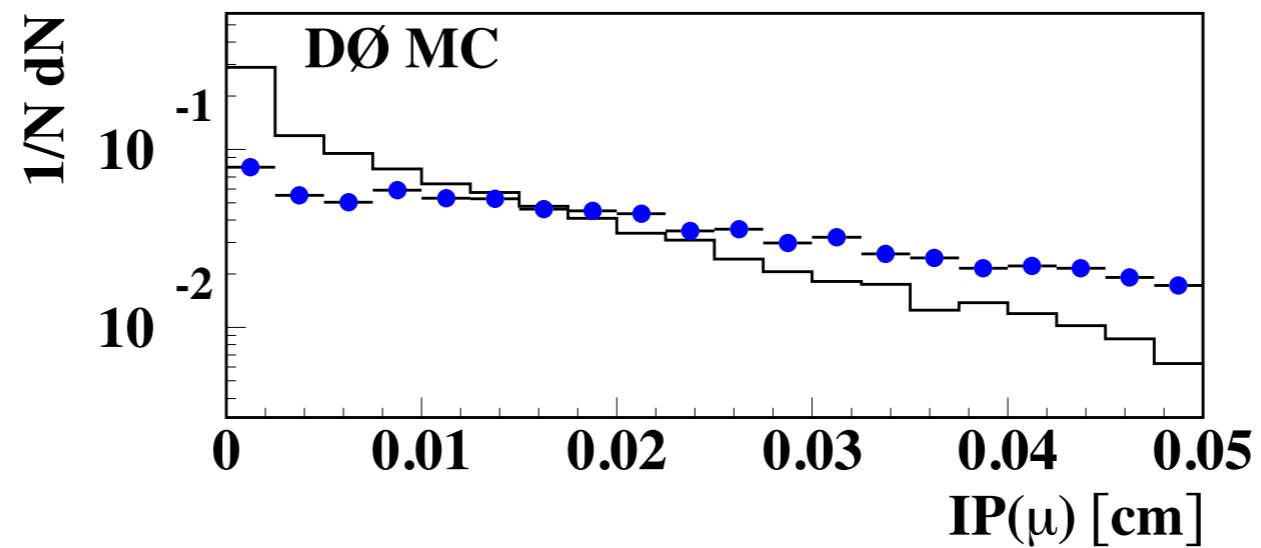
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The muon impact parameter ( $IP$ ) distribution in the inclusive muon sample (dots). The solid line represents the muon  $IP$  distribution in simulation. The shaded histogram is the contribution from  $K$ ,  $\pi$  and  $p$  background muons in simulation.

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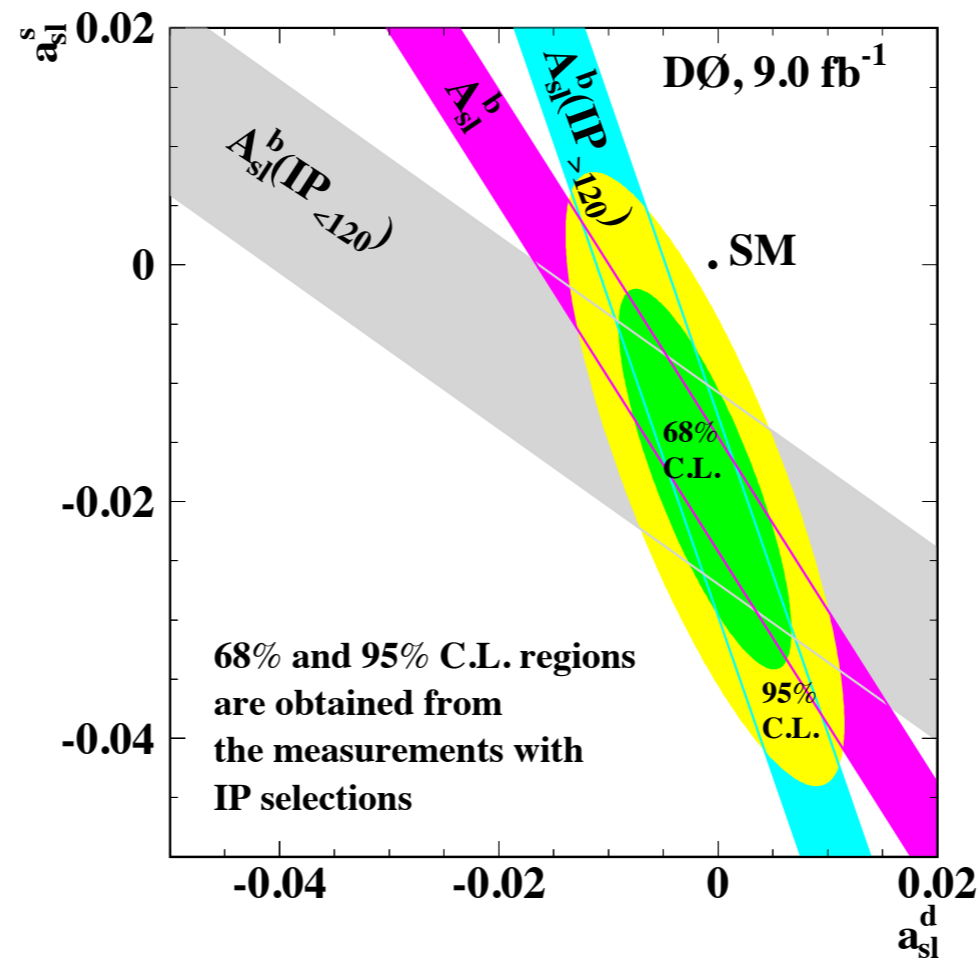
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The normalized impact parameter ( $IP$ ) distribution for muons produced in **oscillating decays** of  $B_d^0$  mesons (dots) and  $B_s^0$  mesons (solid histogram) in simulation.

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from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
Fermilab, 30 June 2011



Measurements of  $A_{sl}^b$  with  $IP > 120\mu\text{m}$  and  $IP < 120\mu\text{m}$ , and corresponding 68% and 95% confidence level regions in the  $(a_{sl}^d, a_{sl}^s)$  plane. Also shown is the measurement with no  $IP$  cut.

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from Bruce Hoeneisen  
representing the  
**DØ Collaboration**  
Fermilab, 30 June 2011

upcoming: LHCb

# LHCb and leptonic charge asymmetry

## Flavour specific asymmetry: $a_{fs}$

- ◆ D0 charge asymmetry measurement, using  $bb \rightarrow \mu\mu X$  event

$$A^b = \frac{N^{++} - N^{--}}{N^{++} + N^{--}} = (0.494)a_{fs}^s + (0.506)a_{fs}^d$$

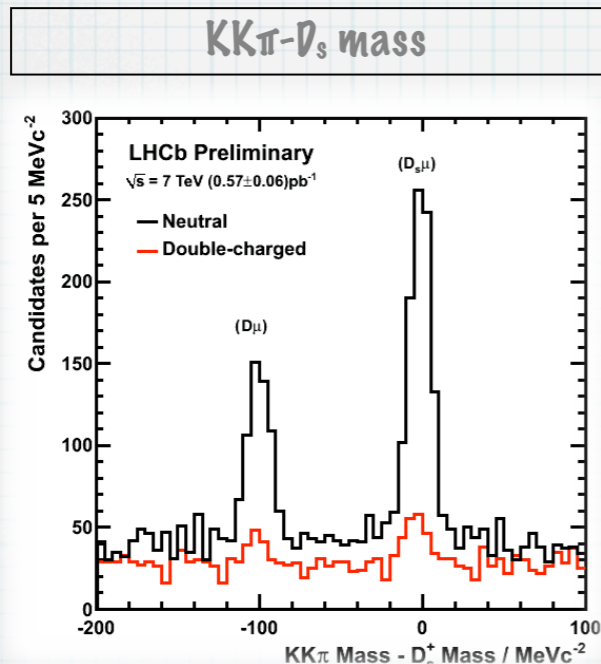
$$a_{fs}^s = \frac{\Delta\Gamma^s}{\Delta M^s} \tan \phi_s$$

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

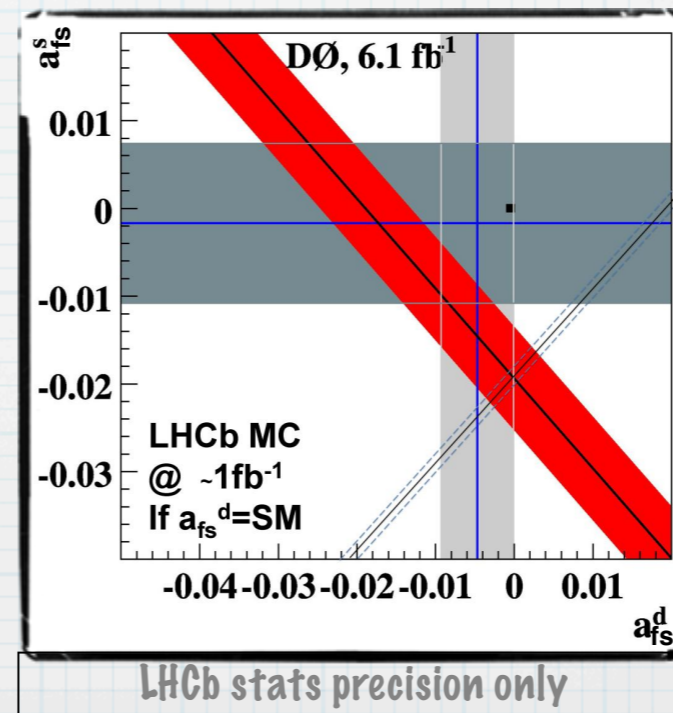
- ◆ LHCb plans to measure exclusive rates  $B_{(q)} \rightarrow D_{(q)} \mu\nu$  in pp
- ◆ Ignore time dependent part to remove production asym ( $\sim 10^{-2}$ )
- ◆ Compute the difference in the Asymmetry between  $B_s, B^0$  to remove detector asymmetries ( $\sim 10^{-2}$ )

from José Ángel  
Hernando  
Morata 2010 talk

0.57 pb<sup>-1</sup>



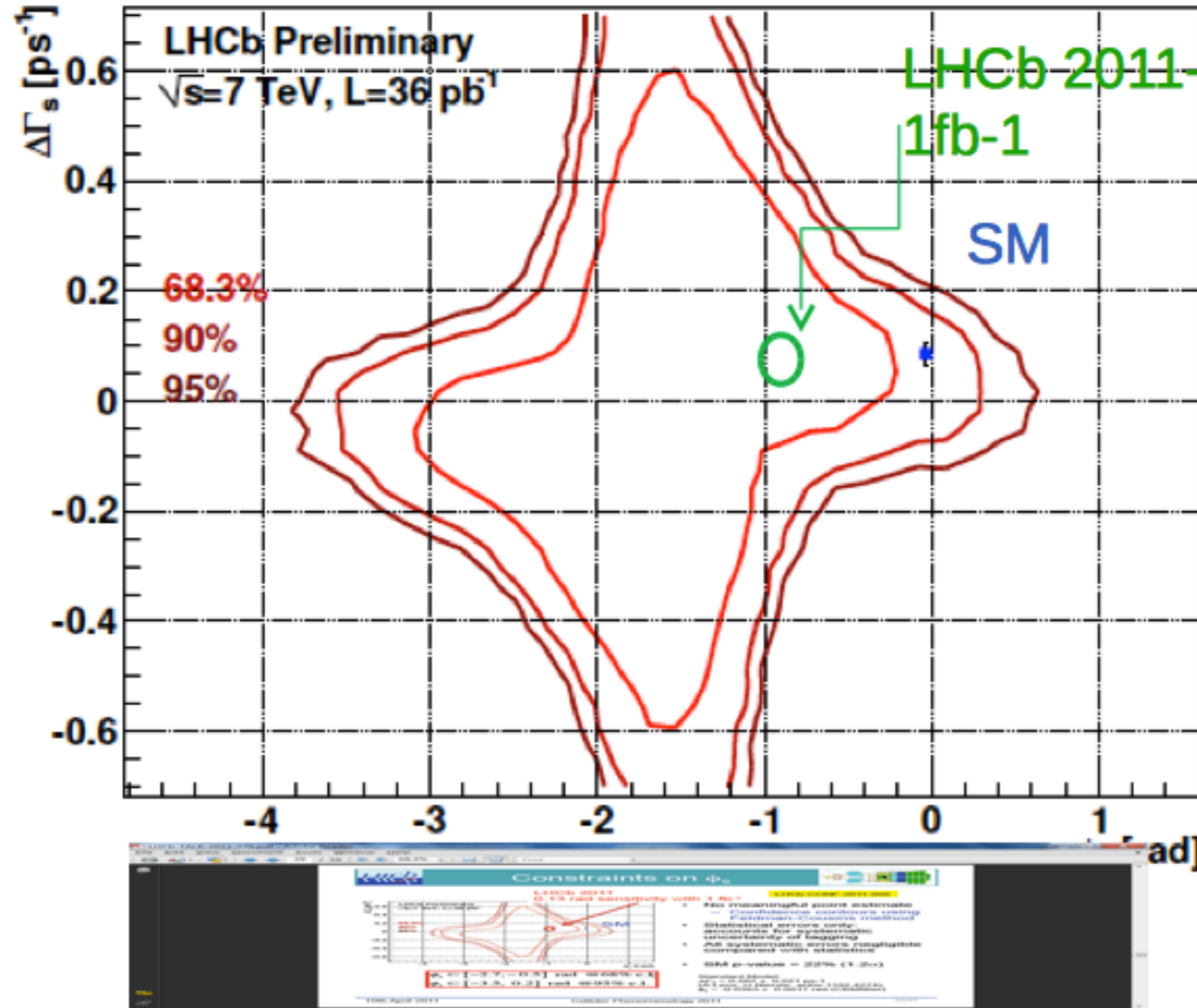
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# Flavour Tagged $\phi$ s



Flavour tagged fit to mass, time, and angular distribution

# SM Theory

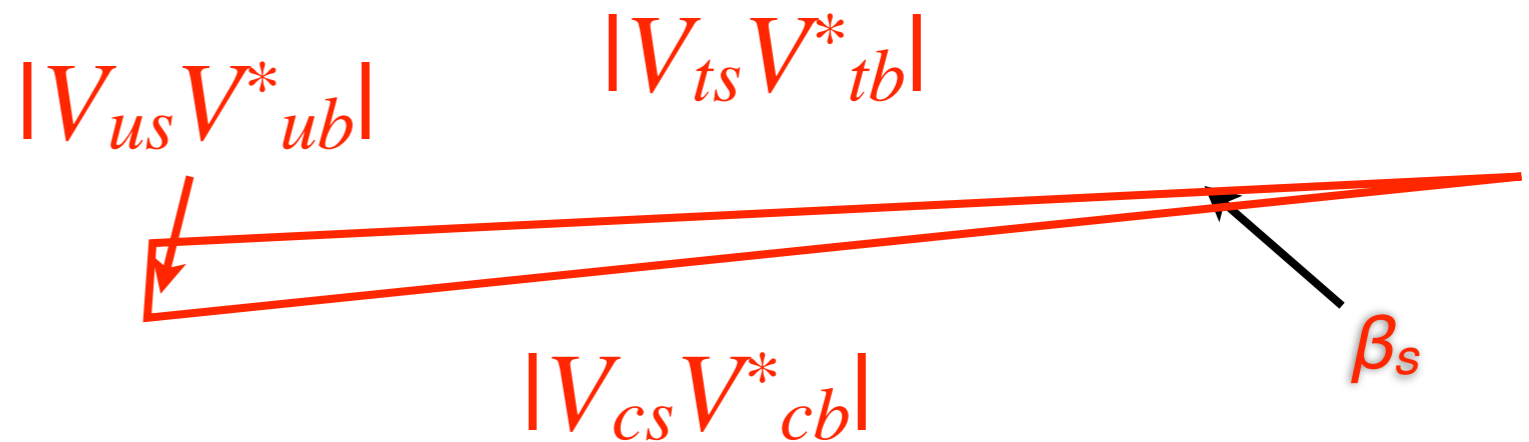
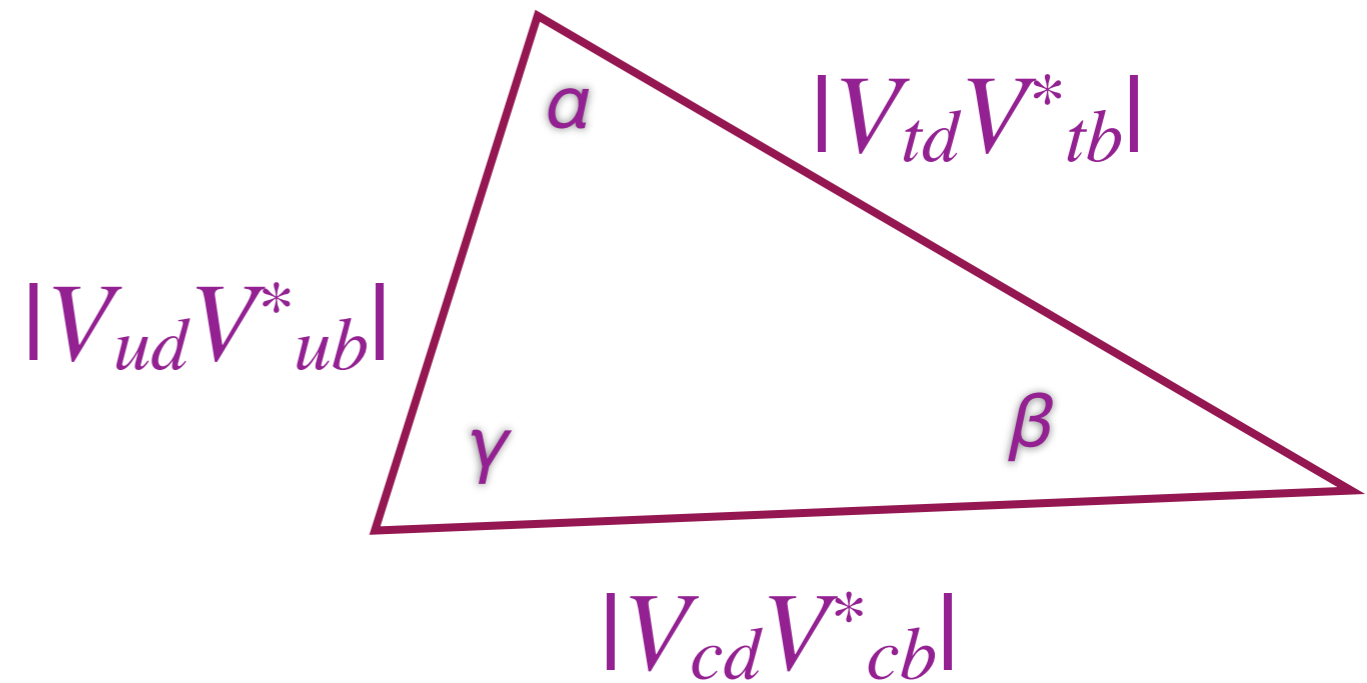
# CKM parameters

$$\beta \equiv \arg \left( -\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*} \right)$$

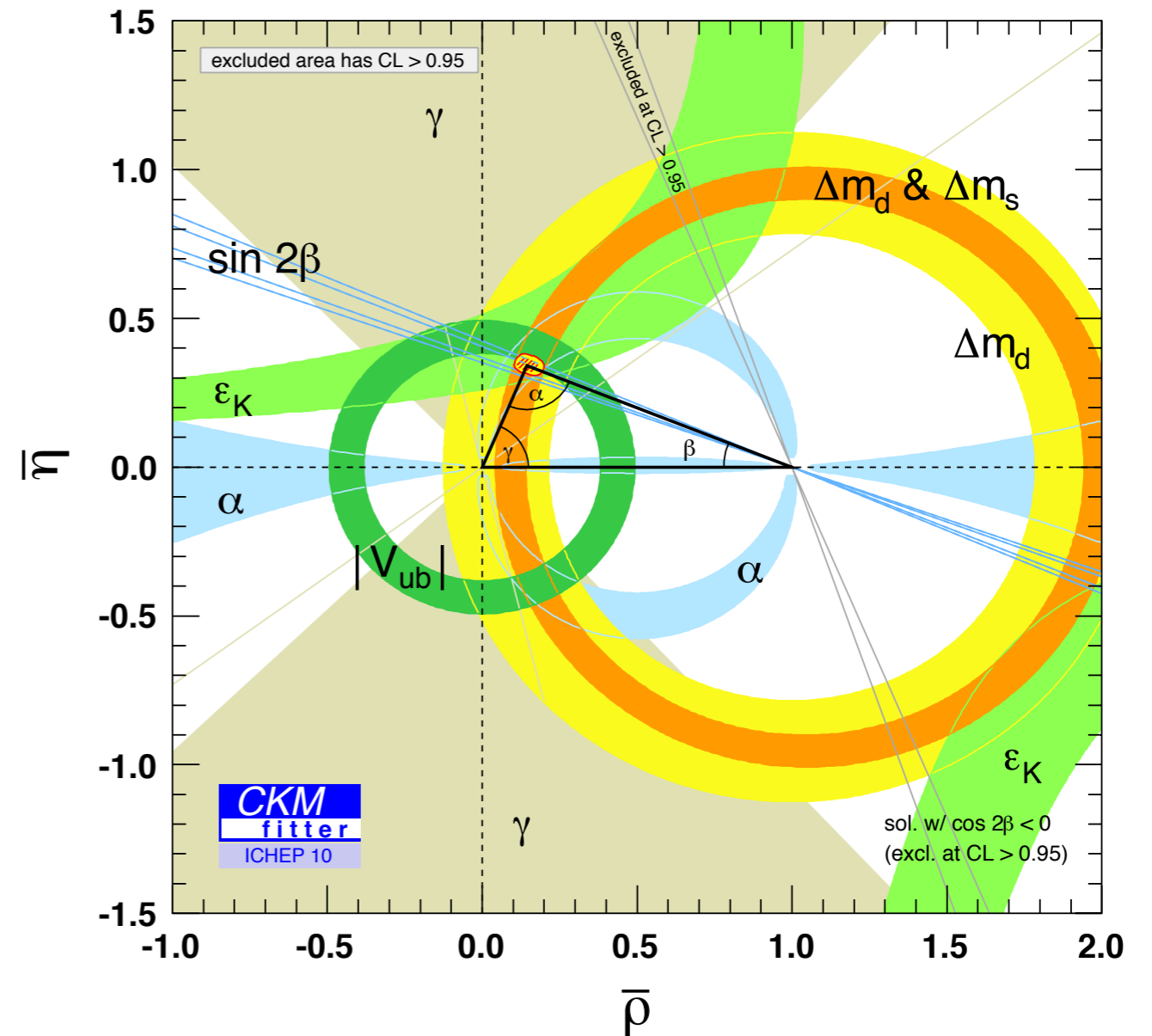
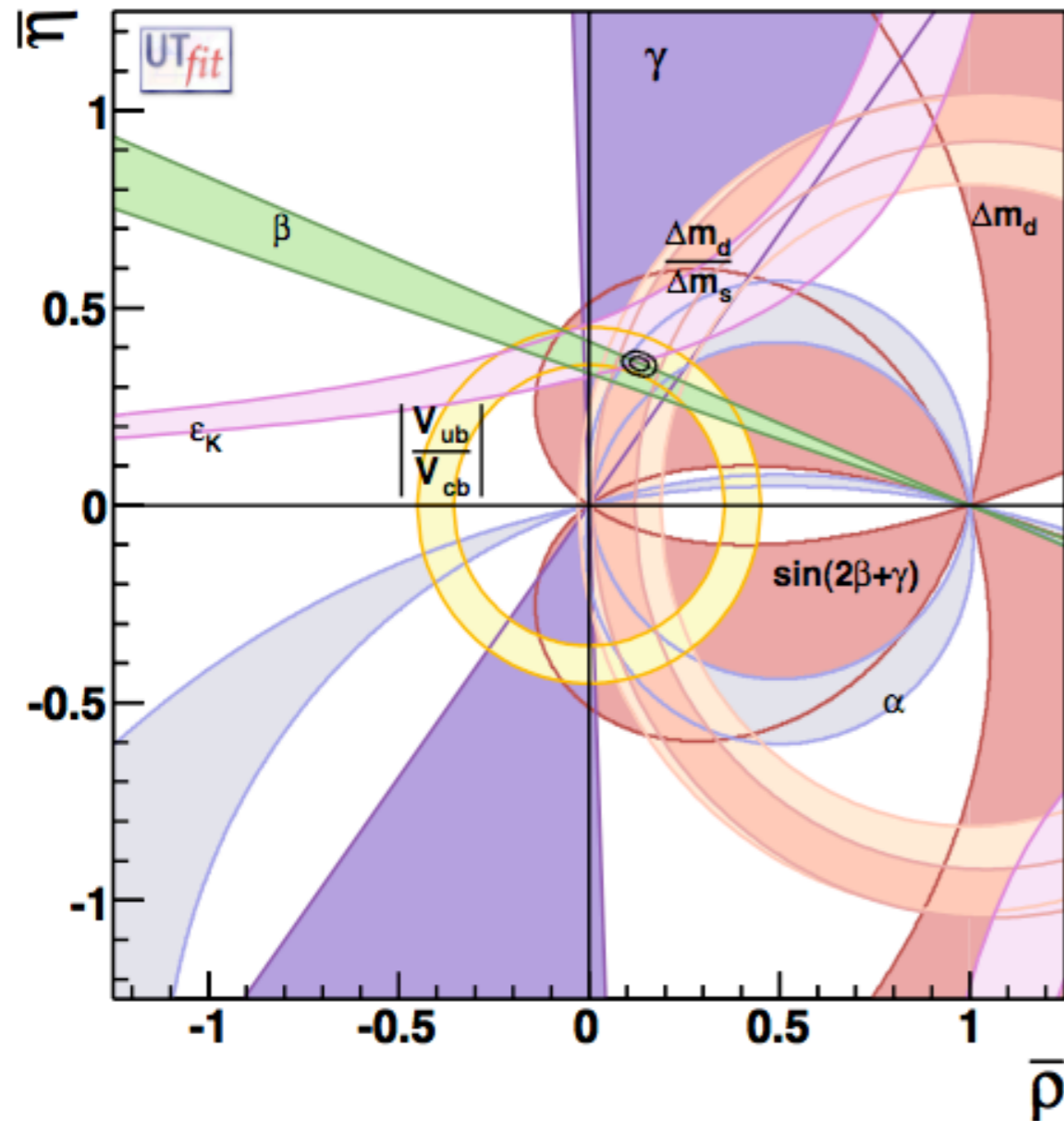
$$\alpha \equiv \arg \left( -\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*} \right)$$

$$\gamma \equiv \arg \left( -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right)$$

$$\beta_s \equiv \arg \left( -\frac{V_{cs}V_{cb}^*}{V_{ts}V_{tb}^*} \right) \approx 0$$



# The Dreaded Unitarity Triangle



Status of SM CKM parameters from CPV in  $B_d$  mixing etc. CPV in interference between mixing and decay of  $B_d$  appears to be mostly (entirely?) from SM

# Mixing Basics

$$i \frac{d}{dt} \begin{pmatrix} |B_q\rangle \\ |\bar{B}_q\rangle \end{pmatrix} = \begin{pmatrix} m^q - \frac{i\Gamma^q}{2} & m_{12}^q - \frac{i\Gamma_{12}^q}{2} \\ m_{12}^{q*} - \frac{i\Gamma_{12}^{q*}}{2} & m^q - \frac{i\Gamma^q}{2} \end{pmatrix} \begin{pmatrix} |B_q\rangle \\ |\bar{B}_q\rangle \end{pmatrix}$$

$$\phi_q \equiv \arg \left( \frac{-m_{12}^q}{\Gamma_{12}^q} \right)$$

➔ Charge asymmetry  $\Rightarrow \Gamma(B \rightarrow \bar{B}) \neq \Gamma(\bar{B} \rightarrow B) \Rightarrow \phi_q \neq 0$

$$\left| m_{12}^q - \frac{i\Gamma_{12}^q}{2} \right| \neq \left| m_{12}^{q*} - \frac{i\Gamma_{12}^{q*}}{2} \right|$$

**Note:**  $|\Gamma_{12}^q| \ll |m_{12}^q|$  in  $B_{d,s}$  systems  $\Delta m_q = 2|m_{12}^q|$

# dimuon asymmetry from $B_s$ or $B_d$ mixing?

- ➔ impact parameter analysis favors  $B_s$
- ➔  $B_d$ , FCNC in  $b \leftrightarrow d$  more constrained from B-factories
- ➔ New contribution to  $B_s$  also hinted at from  $B \rightarrow J/\psi \Phi$  time dependent CPV asymmetry
- ➔ theory can be massaged to favor sizable (relative to SM) new contribution to  $B_s$  mixing with smaller (relative to SM) contribution to  $B_d$  mixing

# New Physics vs SM backgrounds

➔ QCD uncertainty?

➔ QCD is CP symmetric (strong CPV negligible)

➔ Wolfenstein parametrization: selects basis most suitable for understanding where CPV is  $O(\lambda)$

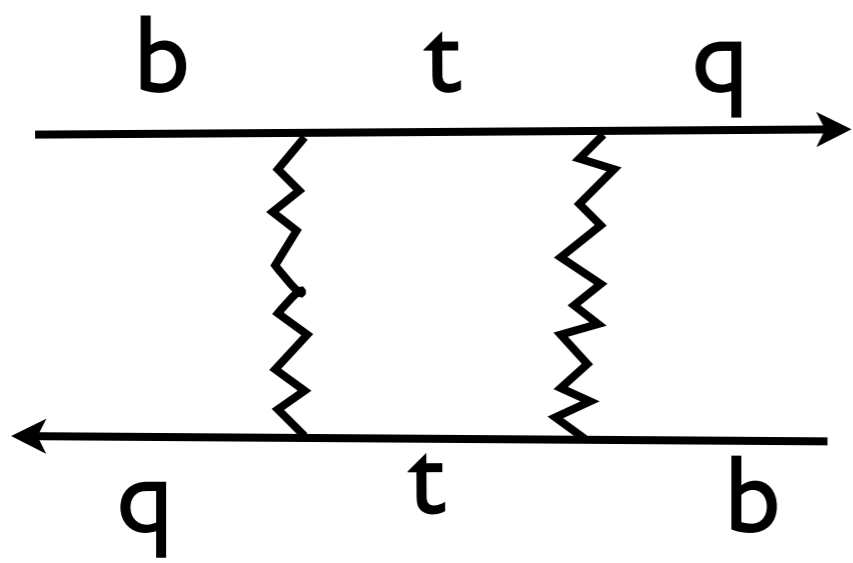
$$\begin{bmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix}.$$

➔ unsuppressed CPV only in processes dominated by  $V_{td}$  and/or  $V_{ub}$

➔ e.g.  $B_d \bar{B}_d$  mixing, not  $B_s \bar{B}_s$  mixing

# SM predicts tiny semi-leptonic asymmetry in $B_{d,s}$

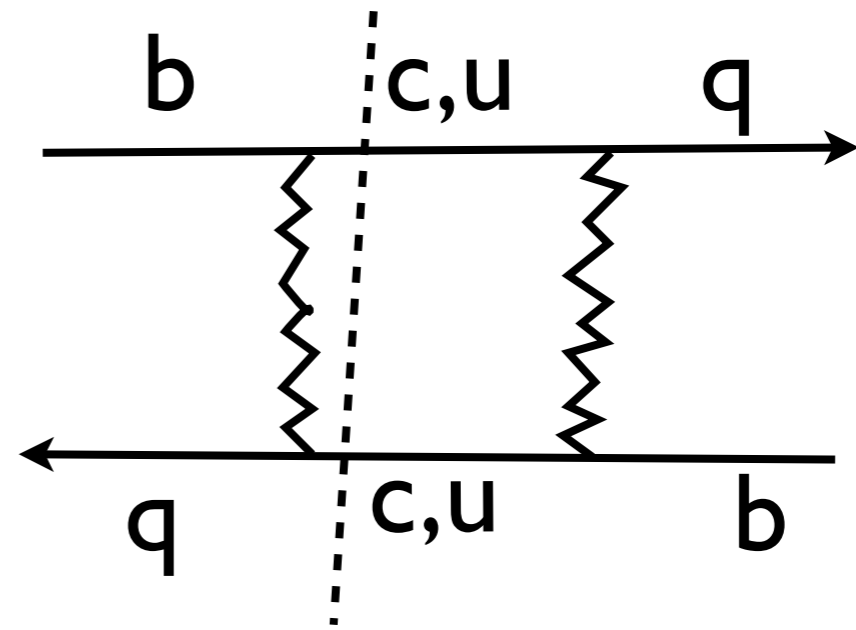
leading  $m_{12}^q \propto (V_{tb}V_{tq}^*)^2$



enhancement of mass  
mixing by heavy top

$$|\Gamma_{12}^q| \ll |m_{12}^q|$$

leading  $\Gamma_{12}^q \propto (V_{cb}V_{cq}^* + V_{ub}V_{uq}^*)^2$



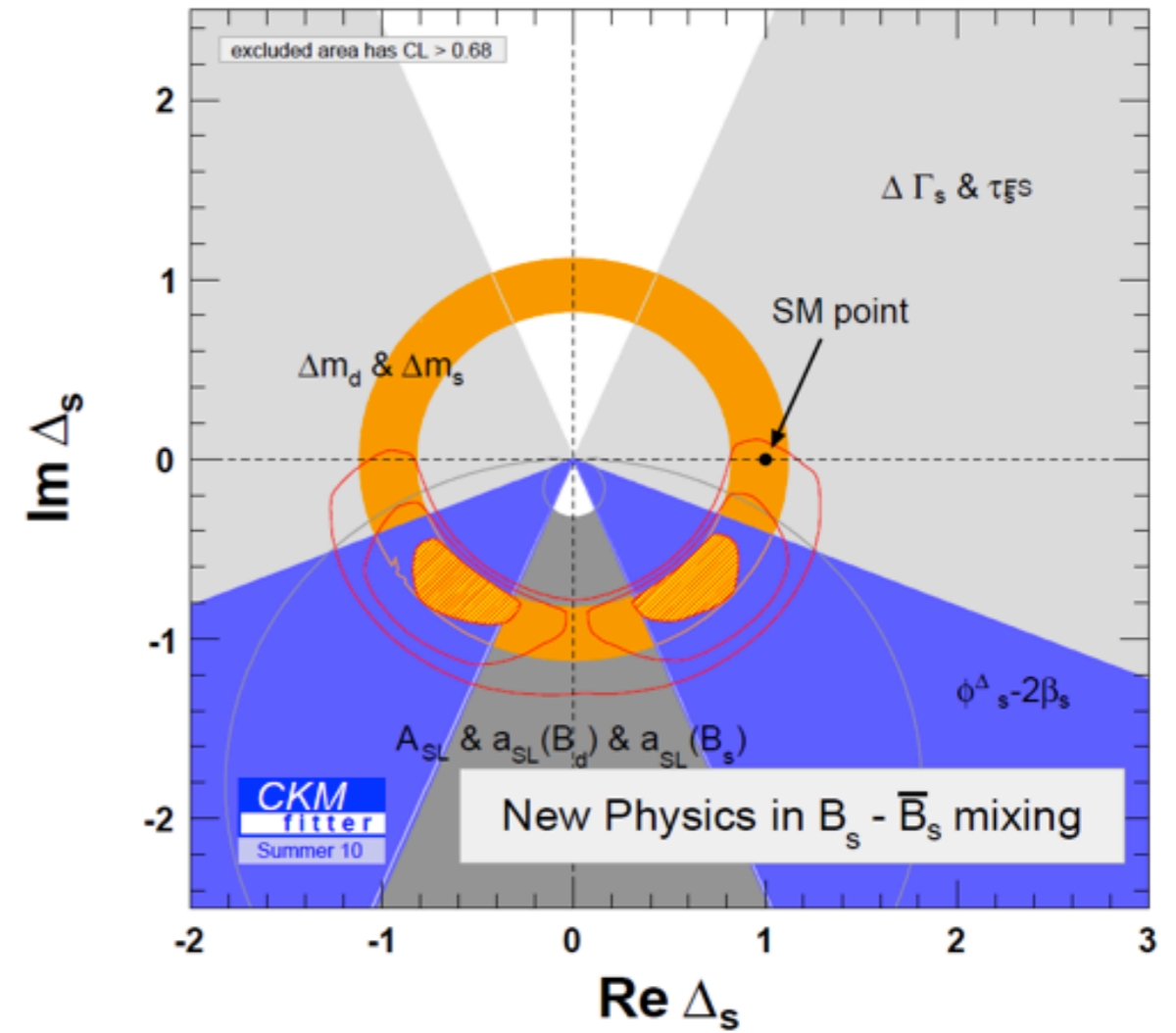
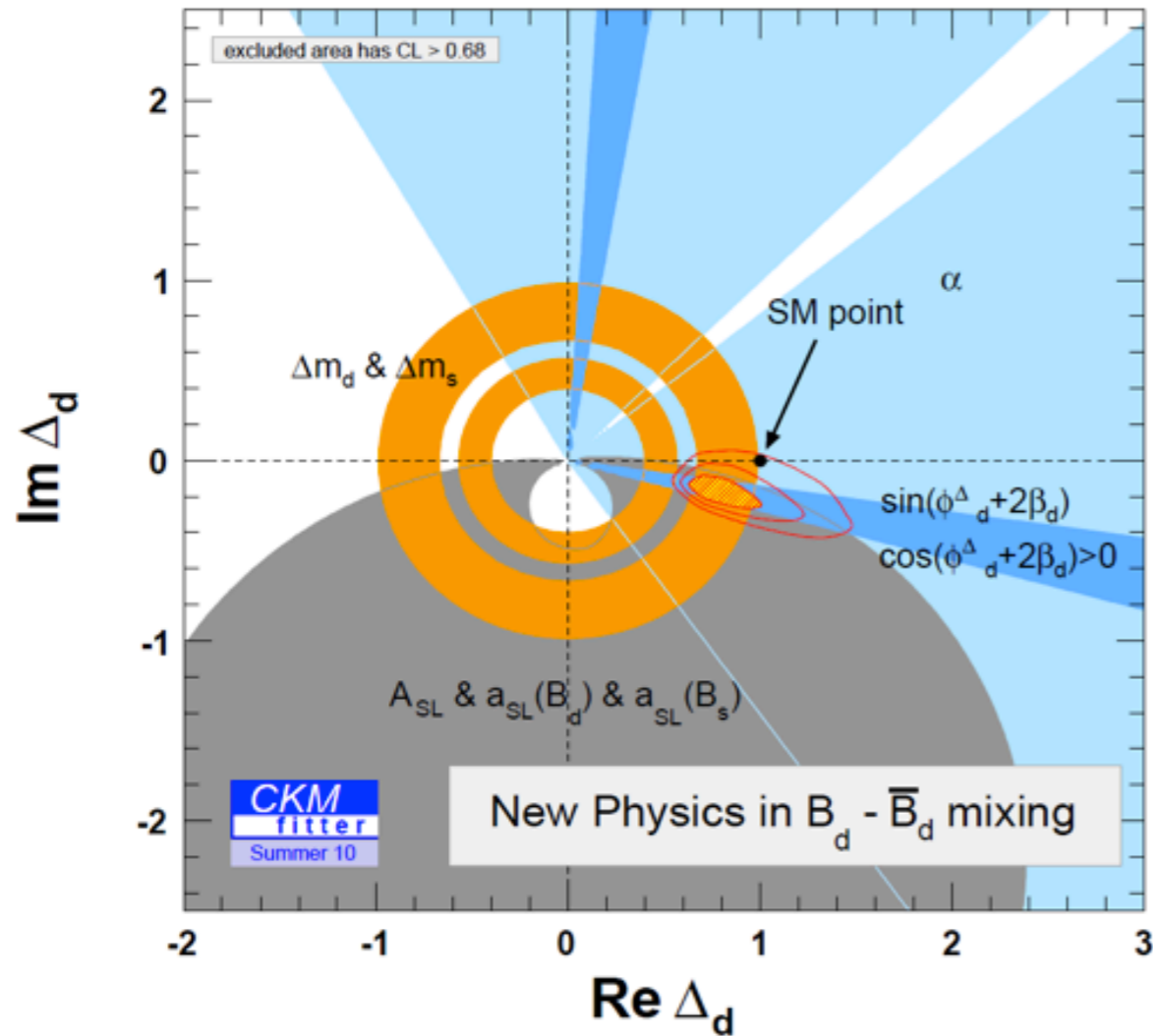
CKM unitarity +  
heavy b quark

$$\phi^q \approx 0$$



# New physics in $B_{d,s}$ mass mixing?

# Still room (indication?) for new CPV physics in mixing

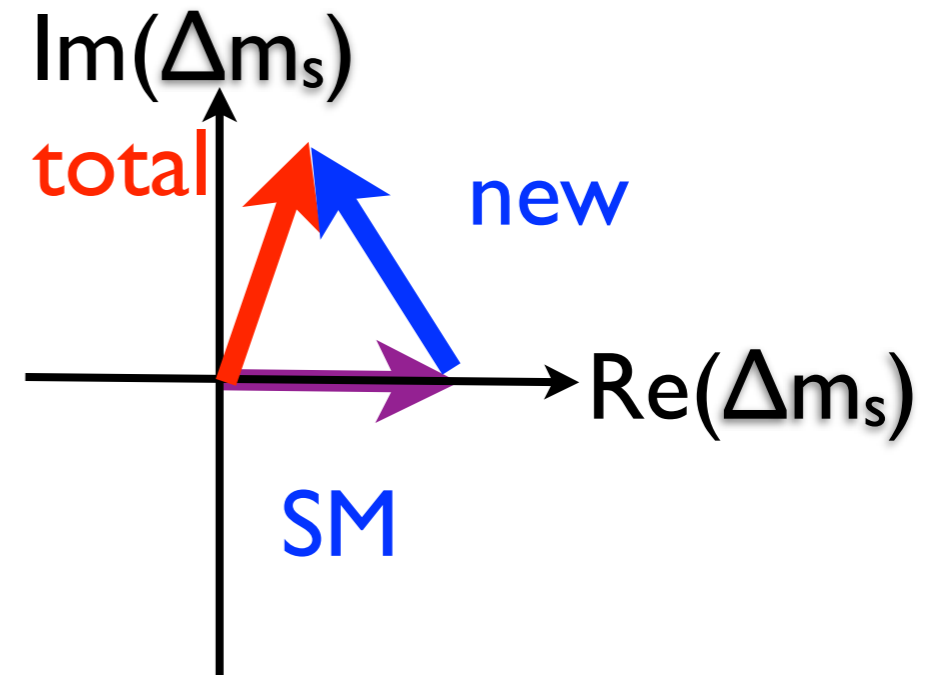


$$\Delta_q \equiv \frac{M_{12}^q}{M_{12}^{q,SM}}$$

$$\Delta_q \equiv |\Delta_q| e^{i\phi_q^\Delta}$$

## New physics in $\Delta m_{d,s}$ ?

- ➔ need a large order one new phase for dimuon asymmetry
- ➔ Don't want to change magnitude of  $|\Delta m_s|$ ,  $|\Delta m_d|$ ,  $|\Delta m_s/\Delta m_d|$  by more than 10-20% of SM
- ➔ Don't want to change phase of  $\Delta m_d$  by more than 10-20% or lose B factory CKM fit to phase
- ➔ want order one change of phase of  $\Delta m_s$  without large change of magnitude



# challenges for charge asymmetry via $\Delta m_s$

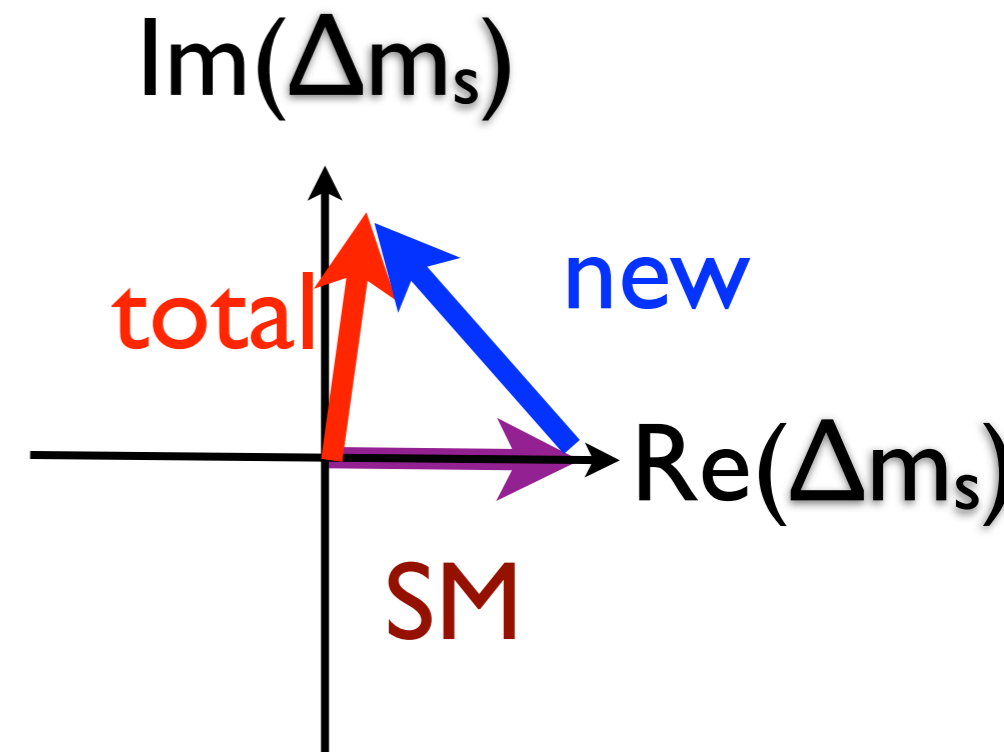
➔ Conspiracy to avoid large nonstandard magnitude

➔  $\Gamma_{12}(\text{SM})$  on small side, charge asymmetry prefers  $\sin \phi_s > 1$

➔ Width problem:  
 $\Delta\Gamma_s = \Delta\Gamma(\text{SM}) \cos \phi_s$

- $\Delta\Gamma(\text{SM})$   
 $= 0.098 \pm 0.024 \text{ ps}^{-1}$

- $\Delta\Gamma_s(\text{expt})$   
 $= 0.134 \pm 0.039 \text{ ps}^{-1}$



expt preference for  
large  $\cos \phi_s$   
and large  $\sin \phi_s$

# Standard Model Flavor Physics:

new FCNC in  $b \leftrightarrow s$  versus experiment

SM Theory

$$\Delta m_s \approx 19.6 \pm 2.2 \text{ ps}^{-1}$$

$$\text{Br}(B_s \rightarrow \mu^+ \mu^-) \approx 3.6 \times 10^{-9}$$

$$\text{Br}(B \rightarrow X_s \gamma) \approx 3.2 \pm 0.2 \times 10^{-4}$$

Current Experiment

$$\Delta m_s \approx 17.77 \pm 0.12 \text{ ps}^{-1}$$

$$\text{Br}(B_s \rightarrow \mu^+ \mu^-) < 4.3 \times 10^{-8}$$

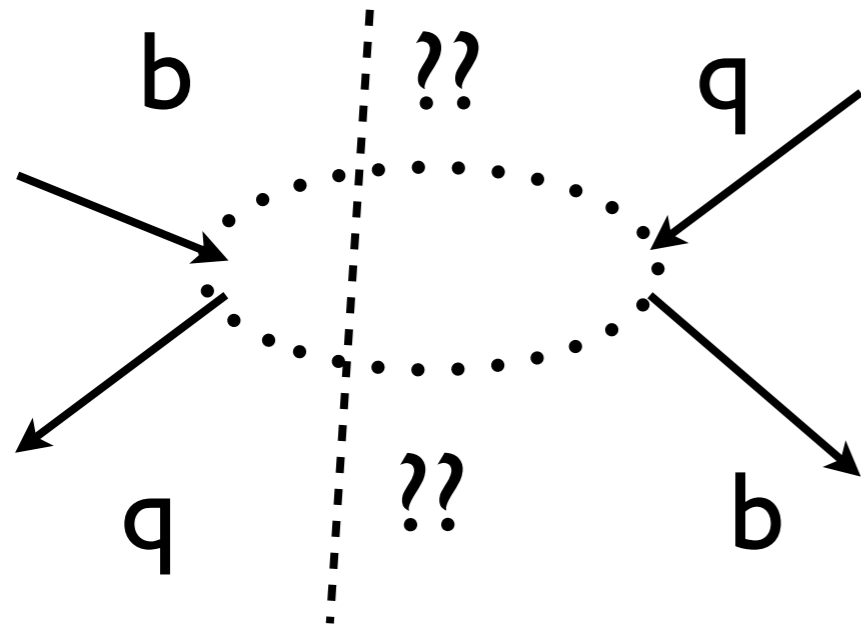
$$\text{Br}(B \rightarrow X_s \gamma) \approx 3.4 \pm 0.3 \times 10^{-4}$$

## Other than that...

- ➔ Clear sailing for model builders!
- ➔ NP at TeV scale can compete with SM loops
- ➔ similar (relative to CKM) NP contributions to  $B_d$  and  $B_s$  mixing allowed
- ➔ must violate assumption that all CPV is in Yukawas or in spurions proportional to Yukawas (minimal flavor/CPV)
- ➔ new flavor and/or CPV for third generation?

# New physics in decay?

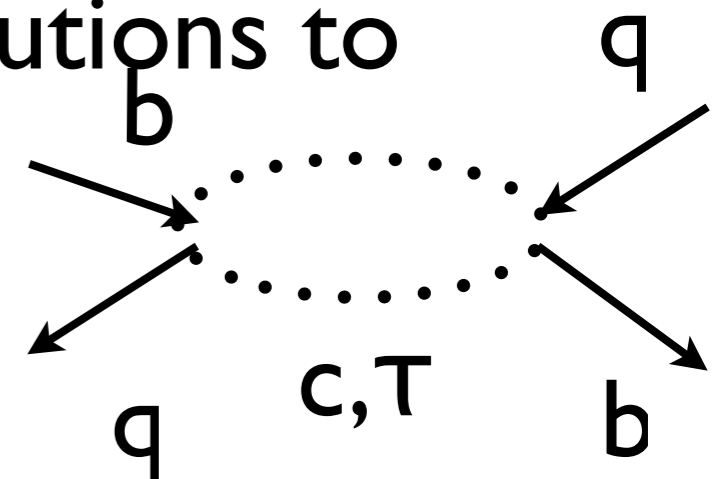
# issues with new contribution to decay



➡ affects B branching fractions, requires new physics in decay comparable to SM tree

➡ Bauer and Dunn: largish new contributions to  $B_s \rightarrow \tau^+ \tau^-, c \bar{c}$  OK

➡ potentially large contribution to  $\Delta m_s$

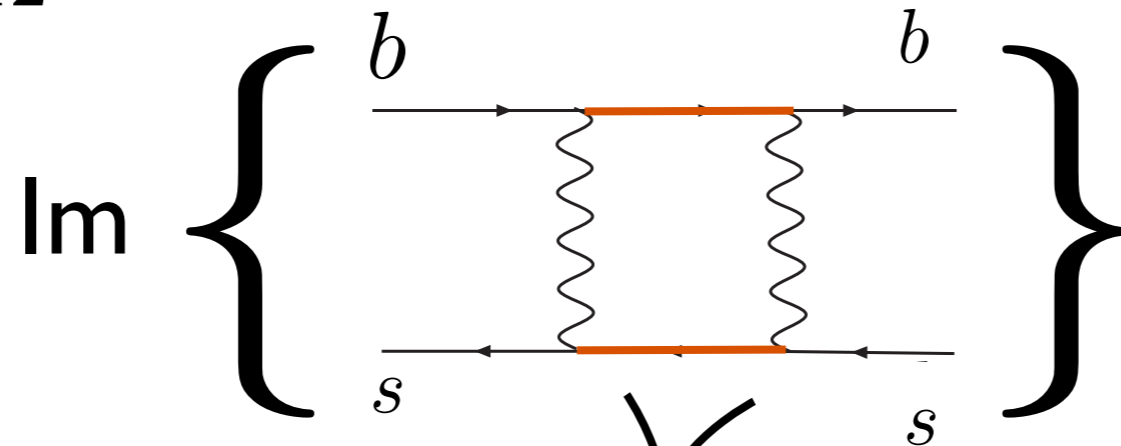




# New Physics in $a^d_{SL}, a^s_{SL}$

from Adam Martin  
Seattle, June 8 2010

more on  $\Gamma_{12}$



confusing, **but**  
rearrange slightly

so  $\Gamma_{12}^{SM} \sim \Gamma_{\bar{B} \rightarrow XX} \sim \frac{G_F^2 f_B^2 M_B^3}{16\pi}$  from tree-level calculations

Meanwhile:  $M_{12}^{SM} \sim \frac{G_F^2 M_W^2 f_B^2 M_B}{16\pi^2} \longrightarrow \frac{M_{12}^{SM}}{\Gamma_{12}^{SM}} \sim \left( \frac{M_W^2}{\pi M_B^2} \right)$

applied to new physics:

$\frac{1}{\Lambda^2} \left( b \right) \longrightarrow \frac{M_{12}^{NP}}{\Gamma_{12}^{NP}} \sim \left( \frac{\Lambda^2}{\pi M_B^2} \right)$

quickly generates large  $M_{12}$

# Our proposal for large contribution to $\Gamma_{12}^S$ : a new light pseudoscalar?

- $\zeta$  Particle with coupling  $-\frac{1}{F} \partial_\mu \zeta \bar{b} \gamma^\mu (g_V^{bs} + g_A^{bs} \gamma_5) s$
- familon?
- pseudoscalar 'Higgs' ?
- Hidden Valley?
- Mass mixing with  $B_s$   $e^{i\alpha} f^2 \zeta B_s + e^{-i\alpha} f^2 \zeta \bar{B}_s$
- Must have largish width to affect  $\Gamma$   $f = 0.0026 \times \left( \frac{F/|g_A^{bs}|}{10^6 \text{ GeV}} \right)^{-1/2} \text{ GeV}$

## $(B_s, \bar{B}_s, \zeta)$ mass matrix

$$M^2 = \begin{pmatrix} m_{B_s}^2 & \Delta m_{B_s} m_{B_s} & e^{i\alpha} f^2 \\ \Delta m_{B_s} m_{B_s} & m_{B_s}^2 & e^{-i\alpha} f^2 \\ e^{-i\alpha} f^2 & e^{i\alpha} f^2 & M_\zeta^2 \end{pmatrix}$$

- Cannot use perturbation theory to diagonalize mass due to near degeneracy of  $B_s$  system, fortunately mass matrix is simple to diagonalize exactly
- Width matrix may be diagonalized perturbatively in mass eigenstate basis
- obtain  $B_L, B_H$  eigenstates with small  $\zeta$  mixture, fit mass, width difference

# Contribution to mass, width difference

➔ Relative contribution to mass difference  
proportional to  $M_\zeta^2 - M_{B_s}^2$

➔ Order one contribution to width difference without  
order one contribution to mass difference provided  
that

$$|M_\zeta - M_{B_s}| < \frac{m_{12}^s(SM)}{\Gamma_{12}^s(SM)} \Gamma_\zeta \approx 200 \Gamma_\zeta$$

➔ either a finetuning conspiracy in the mass, or a fairly  
large  $\zeta$  width

# 2010 fit to data

	Experimental	SM prediction
$\Delta\bar{m}_s$	$(17.78 \pm 0.12) \text{ ps}^{-1}$	$(19.6 \pm 2.2) \text{ ps}^{-1}$
$\Delta\Gamma_s$	$0.134 \pm 0.031 \text{ ps}^{-1}$	$(0.098 \pm 0.024) \text{ ps}^{-1}$
$\bar{\Gamma}_s$	$0.680 \pm 0.012 \text{ ps}^{-1}$	$(0.654 \pm 0.008) \text{ ps}^{-1}$
$\tan \phi_s^{\text{sl}}$	$-1.66 \pm 0.64$	$0.0042 \pm 0.0014$
$\beta_s^{J/\psi\Phi}$	$0.21 \pm 0.12$	$0.018 \pm 0.001$

fit 5 observables to 4 variables:  $f$ ,  $\alpha$ ,  $M_\zeta$ ,  $\Gamma_\zeta$

$\chi^2(\text{SM})=14.0$ , (1.6%)  $\chi^2(\zeta \text{ best fit})=2.0$  (16%)

# Other constraints

- ➡ For light  $\zeta$ , we have 2 body decays  $b \rightarrow s$ , ruling out most of the region with  $m_\zeta < 4.8$  GeV from B width
- ➡ Other constraints depend on the decay mode of the  $\zeta$

# Allowed $\zeta$ mass for 2 widths

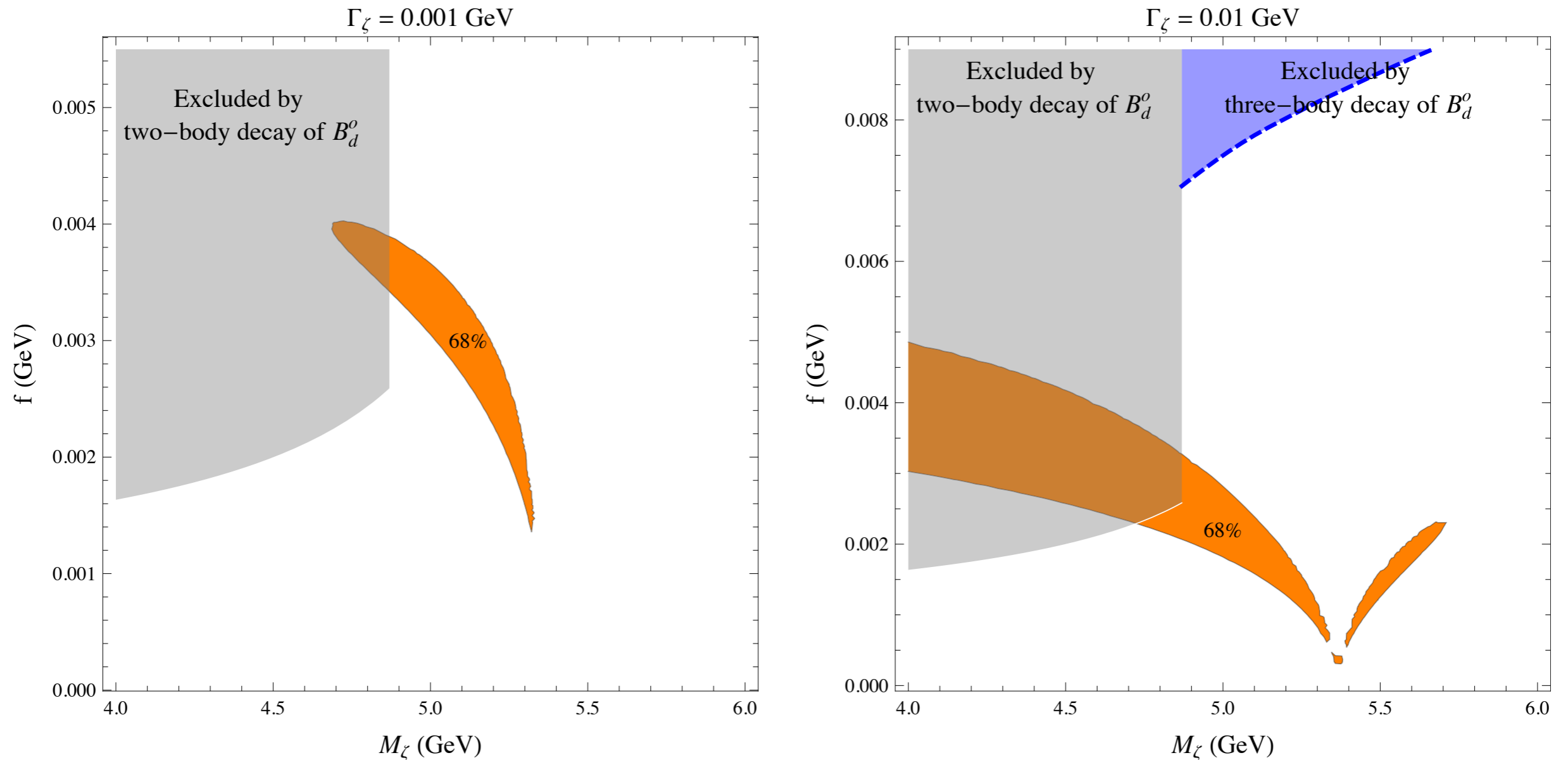


Figure 3: Left panel: the best-fit region in the  $M_\zeta$  and  $f$  space for a fixed width  $\Gamma_\zeta = 0.001$  GeV. The orange contour has 68% C.L. after minimizing  $\chi^2$  in terms of  $\alpha$ . The best fit has  $\chi^2 = 2.0$ . The gray region is ruled out by the two-body decay width of  $B_d$  when  $M_\zeta < m_{B_d} - m_K$ . Three-body decays do not rule out the best-fit region. Right panel: the same as the left panel but for  $\Gamma_\zeta = 0.01$  GeV. The best fit has  $\chi^2 = 5.4$ . The blue region is excluded<sup>39</sup> by requiring the three-body decay width to be

# Allowed Decay modes

➔  $\zeta$  can decay directly, or to other exotics which then decay back to SM particles, e.g.

	Decay Modes
Direct decay	$\tau^+ \tau^-$ , $D\bar{D}(\pi' s)$ , $D(\pi' s)X$
$\zeta \rightarrow 2a$	$2\tau^+ 2e^-$ , $2\tau^+ 2\mu^-$ , $2D^+ 2\pi^-$ , $2\pi^+ 2\pi^-$ , $2\pi^+ 2\pi^-$ , $2K^- 2\pi^+$ , $2K^+ 2K^-$
$\zeta \rightarrow a_1 + a_2$	$X + (\tau^+ e^-, \tau^+ \mu^-, D^+ \pi^-, \pi^+ \pi^-, \pi^+ \pi^-, K^- \pi^+, K^+ K^-)$

Something nonstandard accounts for  $\sim 1-3\%$  of  $B_s$  decays!!



## summary: light pseudoscalar

- ➡ data: largish new contribution to  $B_s$  decays, decay mixing without excessive new contribution to mass mixing
- ➡ weakly coupled new physics below weak scale
- ➡ economical: a new pseudoscalar
  - ➡ large width, weak coupling to SM suggests a new 'sector', large  $\Gamma(\zeta \rightarrow \text{hidden})$ , small  $\Gamma(\text{hidden} \rightarrow \text{vis})$
  - ➡ alternatively largish flavor diagonal couplings to charm or tau
    - ➡ but not 'higgs like', i.e. suppressed flavor diagonal coupling to mu ( $B_{s,d} \rightarrow \mu^+ \mu^-$  constraint), top (unitarity constraint) u
      - ➡ unless largish  $\Gamma(\zeta \rightarrow \text{other 'higgses'})$
  - ➡ also  $\Upsilon$  decay constraints on flavor diagonal b coupling

# experimental smoking guns of hidden pseudoscalar

- ➔ nonstandard contribution to  $B_s$  decays at few %
- ➔ nonstandard contribution to  $B_d, B^+$  decays from  $b \rightarrow s$  ( $\zeta$  decay products) at  $\sim 10^{-4}$
- ➔ rare  $\Upsilon$  decays (model dependent, from flavor diagonal coupling)
- ➔  $\phi_s^{sl} \neq -2\beta_s^{J/\psi\Phi}$  would indicate decay mixing, not mass mixing

# Summary

- ★  $\sim 4 \sigma$  new CPV beyond CKM indicated by  $A_{s1}^b$ 
  - ➡ theory error negligible
  - ➡ experimental systematic claimed to be small
- ★ most likely in  $B_s$  mixing and/or decay
  - ➡ new CPV in  $B_s$  mixing preferred theoretically
    - SM contribution is loop suppressed
    - many models of new TeV scale physics can do this
    - some tuning required
  - ➡ new physics in  $B_s$  decay allowed experimentally
    - hidden pseudoscalar?