# **Precision for New Physics**

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- Precise understanding of the Standard Model processes
  - $\star$  Precise predictions
  - $\star\,$  Good understanding of detector and of the modeling

\* ...

- $\rightarrow$  many talks and discussions here
- New Physics processes: Need for precise predictions?

- For finding a deviation: no?
- For interpretation, including testing and limit setting: yes?

- roughly SM like
- could be embedded in an enlarged Higgs sector (2HDM, Singlet extension,...)

### Precision for a Higgs boson

- QCD corrections (no new coloured particles):
  - $\star$  can be partly taken over from the Standard Model



easily implementable, see e.g. sHDECAY

[Costa, Mühlleitner, Sampaio, Santos, arXiv:1512.05355]

 $\star$  additional corrections for new types of particles



### Precision for a Higgs boson

Electroweak corrections: cannot be extracted from SM ones

 $\Rightarrow$  calculate:

- \* perform renormalization:
  - choice of set of independent input parameters
  - choice of renormalization conditions

help: NLOCT based on FeynArts and FeynRules

[Degrande]

[Hahn; Küblbeck, [Alloul, Christensen, Böhm, Denner] Degrande, Duhr, Fuks]

calculate diagrams

#### help: FeynArts/FormCalc, GoSam, MadGraph5\_aMC@NLO, ... [Alwall Erodoriy

[Cullon yon Dourzon Croiner

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[Hahn; Hahn,	Perez-Victoria]	Heinrich, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, Reiter, Schlenk, von Soden- Fraunhofen, Tramontano]	Frixione, Hirschi, Maltoni, Mattelaer, Shao, Stelzer, Torrielli, Zaro]

Higgs potential:

$$\begin{split} V &= m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 + m_{12}^2 (\Phi_1^{\dagger} \Phi_2 + \text{h.c.}) \\ &+ \lambda_1 (\Phi_1^{\dagger} \Phi_1)^2 + \lambda_2 (\Phi_2^{\dagger} \Phi_2)^2 + \lambda_3 (\Phi_1^{\dagger} \Phi_1) (\Phi_2^{\dagger} \Phi_2) + \lambda_4 (\Phi_1^{\dagger} \Phi_2) (\Phi_2^{\dagger} \Phi_1) \\ &+ \lambda_5 \left[ (\Phi_1^{\dagger} \Phi_2)^2 + (\Phi_2^{\dagger} \Phi_1)^2 \right] \end{split}$$

- CP conserving
- $\bullet$  invariant under  $\Phi_1 \to -\Phi_1$





Chosen set of independent input parameters:



- FeynRules for the creation of the model file
- Matrix element generation with FeynArts/FormCalc
- Implementation into PropHecy4f [Bredenstein, Denner, Dittmaier, Mück, Weber]

Azimuthal angle distribution:

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[Altenkamp, Dittmaier, HR]



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- Differences at LO (dashed) rather small
- NLO corrections sizeable
- For scenarios, see [Haber, Stål, arXiv:1507.04281]

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Azimuthal angle distribution (relative correction): [Altenkamp, Dittmaier, HR]



#### New coloured particles

QCD corrections have to be reevaluated:

• For Higgs boson production and decay within the MSSM:

a lot of work has been done

Small remark:

For Higgs boson production via VBF in the MSSM:

SUSY electroweak effects are larger than SUSY QCD effects,

[Hollik, Plehn, Rauch, HR, arXiv:0804.2676; Figy, Palmer, Weiglein, arXiv:1012.4789]

#### New coloured particles

QCD corrections for processes with new coloured particles:

Example: Squark production and decay:



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#### New coloured particles

QCD corrections for processes with new coloured particles: Example: Squark production and decay with parton shower:



[Gavin, Hangst, Krämer, Mühlleitner, Pellen, Popenda, Spira, arXiv:1305.4061]

$$m_{1/2}=550$$
 GeV,  $m_0=825$  GeV,  $A_0=0$  GeV,  ${\rm tan}(\beta)=10,$   ${\rm sgn}(\mu)=+1$ 

#### Precision input and precision observable

Example: Higgs physics:

In the Standard Model:

The Higgs boson mass is a free parameter.

In extensions of the Standard Model:

The Higgs boson mass can depend on

other parameters of the theory,

e.g. in the (Next) Minimal Supersymmetric Standard Model.

#### Precision input and precision observable

- Needed as consistent input for the calculation of cross sections and decay widths in the (N)MSSM
- The experimental measured value

ATLAS/CMS:  $m_H = 125.09 \pm 0.21$  (stat)  $\pm 0.11$  (syst) GeV

constrains the viable parameter space of the MSSM:

A precise theoretical prediction is needed to fully exploit this constraint.

#### Higgs boson masses at higher orders

Two-point-function:



Calculate the zeros of the determinant of  $\hat{\Gamma}$ 

 $\Rightarrow$  loop-corrected Higgs masses

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### Implications of a 125 GeV Higgs boson (MSSM)



- A 125  $\pm$  3 GeV mass constrains the parameter space but does not exclude the MSSM. (theory uncertainty  $\approx$  3 GeV)
- here: no known 3-loop contributions included

[Martin; Harlander, Kant, Mihaila, Steinhauser]

#### Higgs-boson self energies in the MSSM

	real parameters				complex parameters	
	DR scheme		OS/mixed schemes		OS/mixed schemes	
	$p^2 = 0$	$p^2  eq 0$	$p^{2} = 0$	$p^2 \neq 0$	$p^2 = 0$	$p^2  eq 0$
one-loop				$\checkmark$		$\checkmark$
two-loop						
$\mathcal{O}(\alpha_t \alpha_s)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$	
$\mathcal{O}(\alpha_t^2)$			$\checkmark$		$\checkmark$	
$\mathcal{O}(\alpha_b \alpha_s)$			$\checkmark$			
$\mathcal{O}(\alpha_t \alpha_b, \alpha_b^2, \alpha_\tau^2)$			$\checkmark$			
$\mathcal{O}(\alpha_{\tau}\alpha_{b})$			$\checkmark$			
full	$\checkmark$					
1st 5 rows above		$\checkmark$				
three-loop						
$O(\alpha_t \alpha_s^2)$			$\checkmark$			

Higher orders: "only" logarithmic contributions known

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### Higgs-boson self energies in the MSSM

[Chankowski, Pokorski, Rosiek, 92; hep-ph/9303309; Dabelstein hep-ph/9409375] [Frank, Hahn, Heinemeyer, Hollik, HR, Weiglein, hep-ph/0611326] [Zhang, hep-ph/9808299; Degrassi, Slavich, Zwirner, hep-ph/0105096] [Degrassi, Di Vita, Slavich, arXiv:1410.3432] [Heinemeyer, Hollik, Weiglein, hep-ph/9812472; Degrassi, Slavich, Zwirner, hep-ph/0105096] [Borowka, Hahn, Heinemever, Heinrich, Hollik, arXiv:1404.7074, arXiv:1505.03133: Degrassi, Di Vita, Slavich, arXiv:1410.3432] [Heinemeyer, Hollik, HR, Weiglein, arXiv:0705.0746] [Brignole, Degrassi, Slavich, Zwirner, hep-ph/0112177] [Hollik. Paßher. arXiv:1401.8275: arXiv:1409.1687] [Brignole, Degrassi, Slavich, Zwirner, hep-ph/0206101; Heinemeyer, Hollik, HR, Weiglein, hep-ph/0411114] [Dedes, Degrassi, Slavich, hep-ph/0305127] [Allanach, Djouadi, Kneur, Porod, Slavich, hep-ph/0406166] [Martin, hep-ph/0211366] [Martin, hep-ph/0405022] [Harlander, Kant, Mihaila, Steinhauser, arXiv:0803.0672; arXiv:1005.5709]

#### **Remark about renormalization**

In the calculation: choice of independent input parameters and of the renormalization scheme

#### Criteria:

- observables ↔ parameters

   (e.g. choice of mass parameters
   which can be related to masses)
- respect symmetry relations
- avoid unphysically large corrections

#### For illustration:

Mass of lightest CP-even NMSSM Higgs boson



#### \* Different schemes can give an error estimate.

#### Higgs boson mass for large stop masses

Prediction obtained via Feynman diagrammatic approach:

- + all log and non-log terms are taken into account at a certain order of perturbation theory
- possible appearance of large logs:

$$\Delta M_h \sim \log \frac{M_S}{m_t} \qquad M_S = \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}: \text{SUSY particle mass scale}$$
$$m_{\tilde{t}_1}, m_{\tilde{t}_2}: \text{ scalar top quark masses}$$
$$m_t: \text{ top quark mass}$$

- $\Rightarrow \bullet$  good prediction for lower SUSY mass scales
  - no reliable prediction for large SUSY mass scales

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### Higgs boson mass for large stop masses



 $\rightarrow$  Combine both approaches

#### Higgs boson mass for large stop masses



Quantum corrections can change the mixing of Higgs bosons  $\Rightarrow$  induce changes of couplings

Needed for predictions of Higgs decay widths

#### NMSSM example:

Coupling to vector bosons squared: (normalized wrt Standard Model)



#### Tools

	fix. order	RL	scheme	CPV?	RPV?	BMSSM?
SUSPECT [1]	2L	—	DR	—	—	(gen.)
SPheno [2]	2L	—	DR	$\checkmark$	$\checkmark$	gen.
SoftSUSY [3]	2L	—	DR		$\checkmark$	NMSSM
CPsuperH [4]	2L	—	MS	$\checkmark$		
FeynHiggs [5]	2L	2L	mixed	$\checkmark$		(NMSSM)
SUSYHD [6]	—	3L	DR	$\checkmark$		
H3m [7]	3L	—	$\overline{DR}$ /mixed			
NMSSMTools [8]	2L	—	DR	(√)		NMSSM
NMSSMCalc [9]	2L	—	$\overline{DR}/mixed$	$\checkmark$		NMSSM
FlexibleSUSY [10]	2L	—	DR	—		gen.

see [Draper, HR, arXiv:1601.01890]

## Tools

- [1] [Djouadi, Kneur, Moultaka, hep-ph/0211331]
- [2] [Porod, hep-ph/0301101; Porod, Staub, arXiv:1104.1573]
- [3] [Allanach, hep-ph/0104145; Allanach, Athron, Tunstall, Lewis, Voigt, Williams, arXiv:1311.7659; Allanach, Bednyakov, Ruiz de Austri, arXiv:1407.6130]
- [4] [Lee, Pilaftsis, Carena, Choi, Drees, Ellis, Wagner, hep-ph/0307377; Lee, Carena, Ellis, Pilaftsis, Wagner, arXiv:0712.2360; arXiv:1208.2212]
- [5] [Heinemeyer, Hollik, Weiglein, hep-ph/9812472; hep-ph/9812320; Degrassi, Heinemeyer, Hollik, Slavich, Weiglein, hep-ph/0212020; Frank, Hahn, Heinemeyer, hep-ph/0611326; Hollik, HR, Weiglein, Hahn, Heinemeyer, Hollik, HR, Weiglein, arXiv:1312.4937]
- [6] [Vega, Villadoro, arXiv:1504.05200]
- [7] [Kant, Harlander, Mihaila, Steinhauser, arXiv:1005.5709]
- [8] [Ellwanger, Gunion, Hugonie, hep-ph/0406215; Ellwanger, Hugonie, hep-ph/0508022]
- [9] [Baglio, Gröber, Mühlleitner, Nhung, HR, Spira, Streicher, Walz, arXiv:1312.4788]
- [10] [Athron, Park, Stöckinger, Voigt, arXiv:1406.2319]

For testing new physics:

Higher-order corrections can be important for the interpretation:

- NLO corrections for production and decay processes (QCD and electroweak)
- Higher-order corrections for masses and mixings

PhD Scholarship in Higgs Physics and Physics beyond the Standard Model

At CP3-Origins and the Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark in Odense

Application deadline: May 31, 2016

http://jobbank.sdu.dk/en/job/837488/syddansk-universitet/
phd-scholarship-in-higgs-physi

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#### **Definition of the** *Z***s**

$$Z_4 = \frac{1}{4} \sin^2(2\beta) \left[ \lambda_1 + \lambda_2 - 2(\lambda_3 + \lambda_4 + \lambda_5) \right] + \lambda_4$$
  

$$Z_5 = \frac{1}{4} \sin^2(2\beta) \left[ \lambda_1 + \lambda_2 - 2(\lambda_3 + \lambda_4 + \lambda_5) \right] + \lambda_5$$
  

$$Z_7 = -\frac{1}{2} \sin(2\beta) \left[ \lambda_1 \sin^2 \beta - \lambda_2 \cos^2 \beta + (\lambda_3 + \lambda_4 + \lambda_5) \cos(2\beta) \right]$$

$$m_t = 173.2 \text{ GeV}, \ M_{\text{SUSY}} = 1000 \text{ GeV},$$
  
 $|\mu| = 1000 \text{ GeV}, \tan \beta = 10, \ M_A = 500 \text{ GeV},$   
 $|M_2| = 200 \text{ GeV}, \ |M_3| = 1500 \text{ GeV}$ 

$$\begin{split} &\kappa = \lambda/5, \ \tan\beta = 2, \ A_{\lambda} = 500 \ \text{GeV}, \ A_{\kappa} = -10 \ \text{GeV}, \\ &v_{S} = \frac{1}{\sqrt{2}} \frac{250}{\lambda} \ \text{GeV}, \ M_{S} = 300 \ \text{GeV}, \ A_{t} = A_{b} = A_{\tau} = -1.5 M_{S}, \\ &M_{1} = M_{S}/3, \ M_{2} = 2/3 M_{S}, \ M_{3} = 2 M_{S}, \\ &m_{t} = 173.3 \ \text{GeV}, \ \mu_{\text{ren}} = M_{S} \end{split}$$

$$\begin{split} |\lambda| &= 0.65, \ |\kappa| = 0.33, \ \tan\beta = 1.94, \ M_{H^{\pm}} = 420.9 \ \text{GeV}, \\ |A_{\kappa}| &= 37.9 \ \text{GeV}, \ \text{Re}(A_{\kappa}) < 0, \ |v_{s}| = \sqrt{2} \frac{193}{|\lambda|} \ \text{GeV} \\ M_{Q_{3}} &= 557 \ \text{GeV}, \ M_{t_{R}} = 709 \ \text{GeV}, \ M_{\tilde{f} \neq \{\tilde{b}, \tilde{t}\}} = 3 \ \text{TeV}, \\ |A_{b}| &= 963 \ \text{GeV}, \ \varphi_{A_{b}} = 0, \ |A_{t}| = 1441 \ \text{GeV}, \ \varphi_{A_{t}} = \pi, \\ M_{1} &= 855 \ \text{GeV}, \ M_{2} = 269 \ \text{GeV}, \ M_{3} = 2794 \ \text{GeV}, \ m_{t} = 173.5 \ \text{GeV}, \\ \mu_{\text{ren}} &= \sqrt{M_{Q_{3}}M_{t_{R}}} \end{split}$$