

High-field ground state of graphene at Dirac Point

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- 1. Ground state of Dirac point in high fields**
- 2. Thermopower at Dirac point**

Zero-energy mode, $n=0$ Landau Level

For valley K

$$H = v \begin{bmatrix} 0 & \pi_x + i\pi_y \\ \pi_x - i\pi_y & 0 \end{bmatrix} = \frac{\sqrt{2}v}{\ell_B} \begin{bmatrix} 0 & a \\ a^\dagger & 0 \end{bmatrix}$$

$$\boldsymbol{\pi} = \mathbf{k} - e\mathbf{A}$$

$$a = \frac{(\pi_x + i\pi_y)\ell_B}{\sqrt{2}}$$

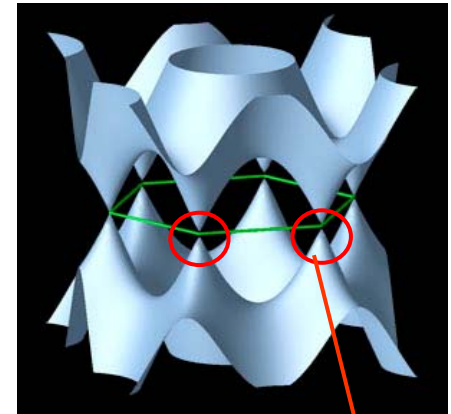
$$H^2\Psi_n = E_n^2\Psi_n$$

$$E_n = \sqrt{2n} \frac{v}{\ell_B} \quad \Psi_n = \begin{bmatrix} |n-1\rangle \\ |n\rangle \end{bmatrix}$$

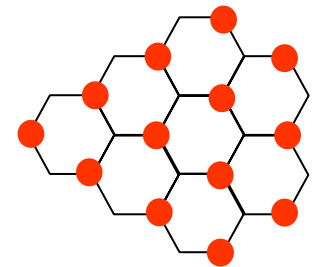
$$\Psi_0 = \begin{bmatrix} 0 \\ |0\rangle \end{bmatrix} \leftarrow$$

For $n=0$ Landau Level

- 1) K - K' valley index same as A - B sublattice index
- 2) Electrons on B sites only



2 valleys K, K'



Broken Symmetry ground states

Coulomb interaction energy in high-field state

Nomura, MacDonald, PRL 06

Goerbig, Moessner, Doucot, PRB 06

Kun Yang, Das Sarma, MacDonald, PRB 06

Ezawa, JPSJ 06

Abanin, Lee, Levitov, PRL 06, 07

Fertig, Bray, PRL 06

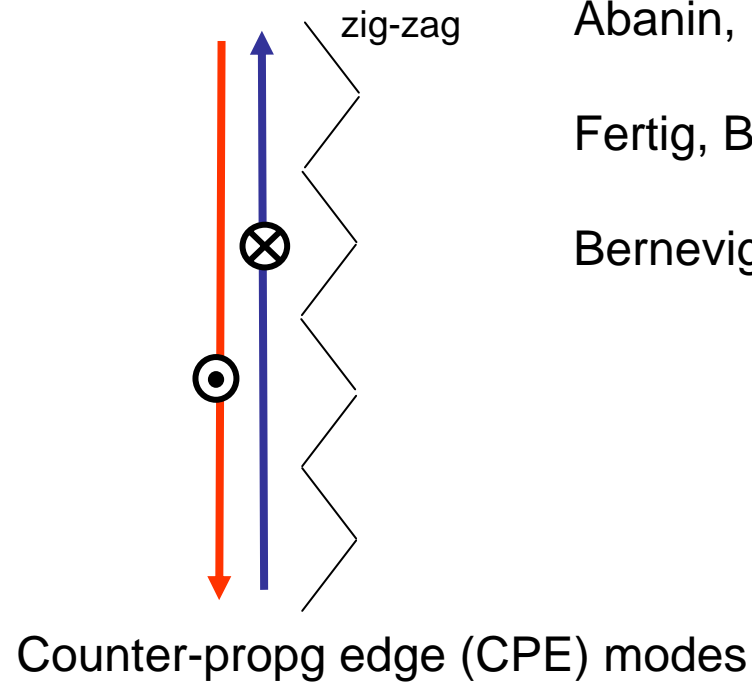
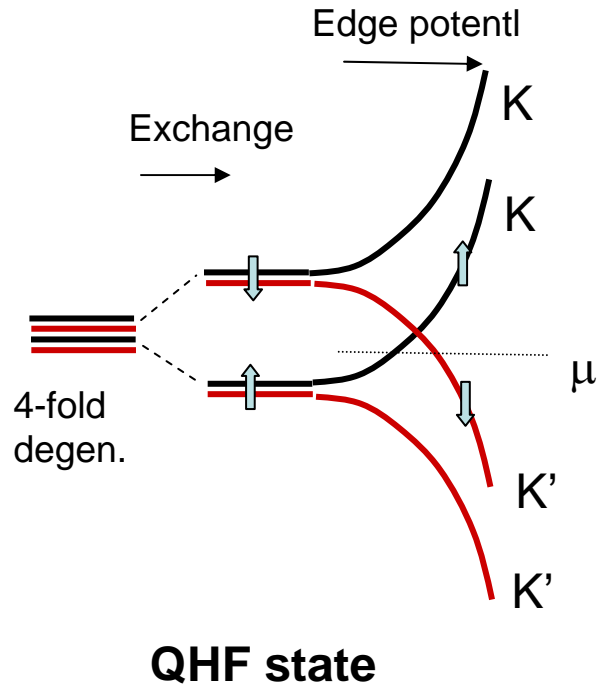
Shimshoni, Fertig, Pai, 08

Alicea, Fisher, PRB 06

Gusynin, Miransky, Sharapov, Shovkovy, PRB 06

Khveshchenko, PRL 02

Quantum Hall Ferromagnet and CPE modes



Abanin, Lee, Levitov

Fertig, Bray

Bernevig, Zhang

At Dirac point, CPE modes protected

$R_{xx} \sim 2(h/e^2)$ at intense fields (dissipative to very large H)

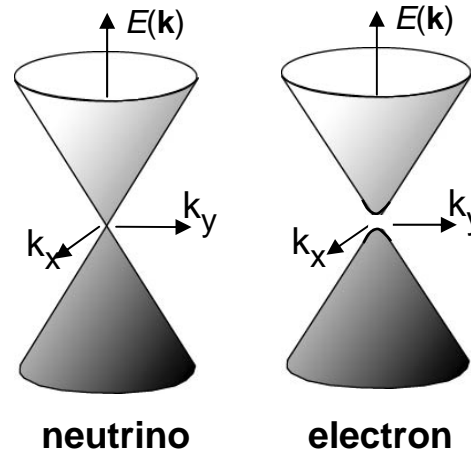
Realization of quantized spin-Hall system

Magnetic Catalysis (field induced gap)

Field H creates gap

Mass term $\Psi^+ \gamma_5 \Psi$

Chiral-symmetry breaking
In (2+1) QED₃



Appelquist et al. 1986

Miransky et al.
1995, 2003, 2005, 2008

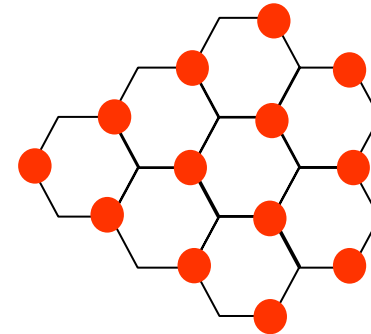
D. Khveshchenko 2002

Field-enhancement of DOS at Dirac Point
Exciton condensation of electron hole pairs

Order parameter for $H > H_c$

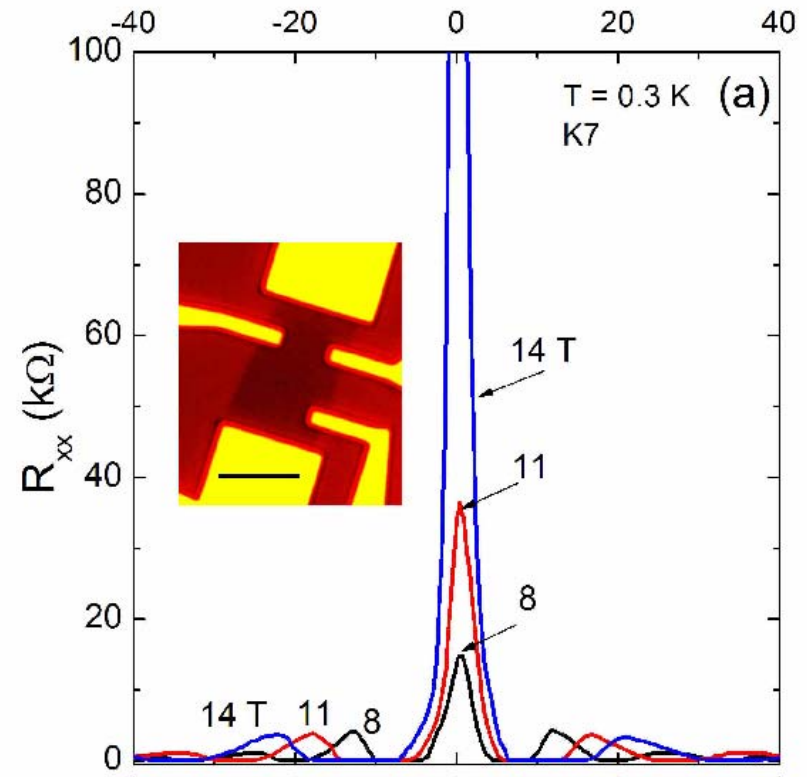
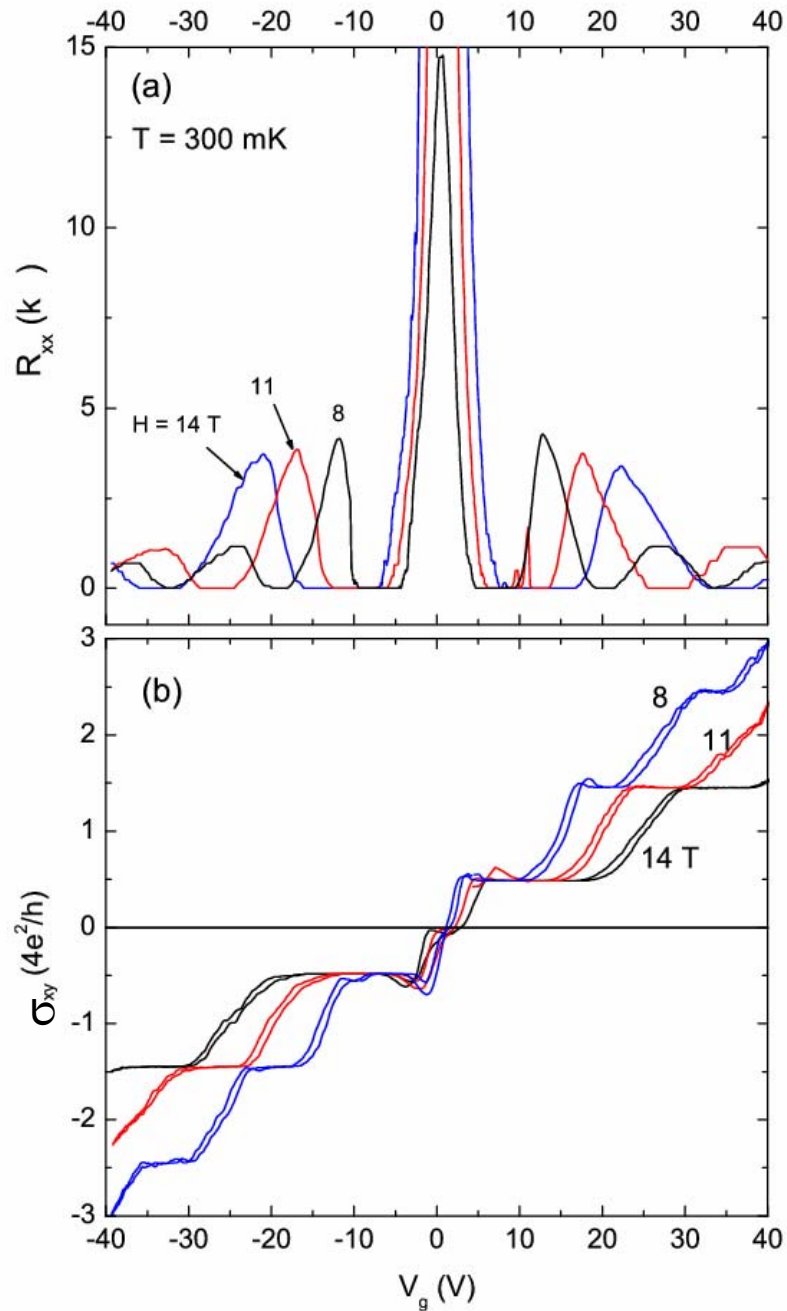
$$\langle \bar{\Psi} \Psi \rangle = \sum_{\sigma} |\psi_{A\sigma}|^2 - |\psi_{B\sigma}|^2$$

Equivalent to sublattice CDW



The resistance R_{xx} and Hall conductivity in graphene

Checkelsky, Li, Ong, PRL '08



R_{xx} diverges with H at Dirac Point

Divergent R0 and zero-Hall step

Checkelsky, Li, Ong, PRL '08

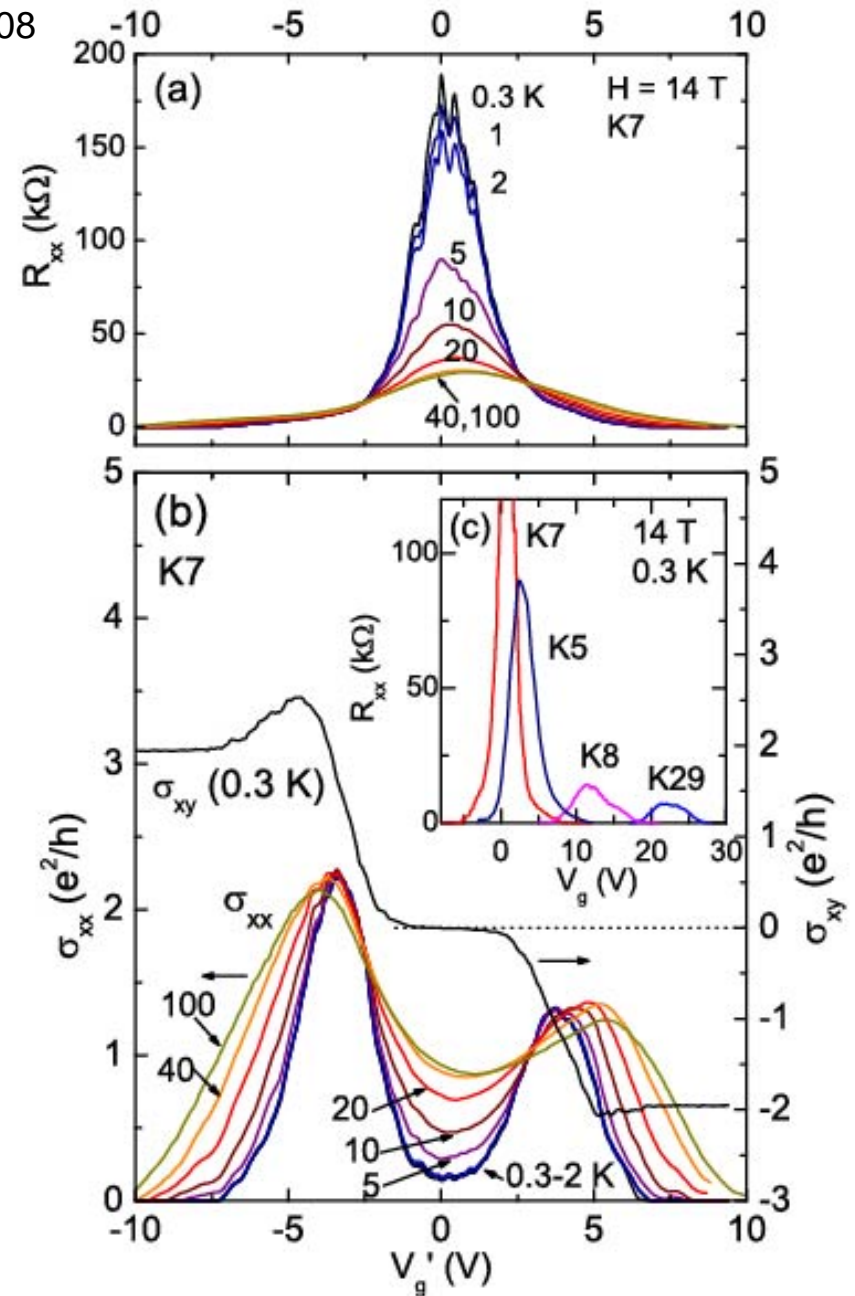
R_0 increases as T decreases

At fixed H and T , $n=0$ peak dependent on V_0 (offset voltage)

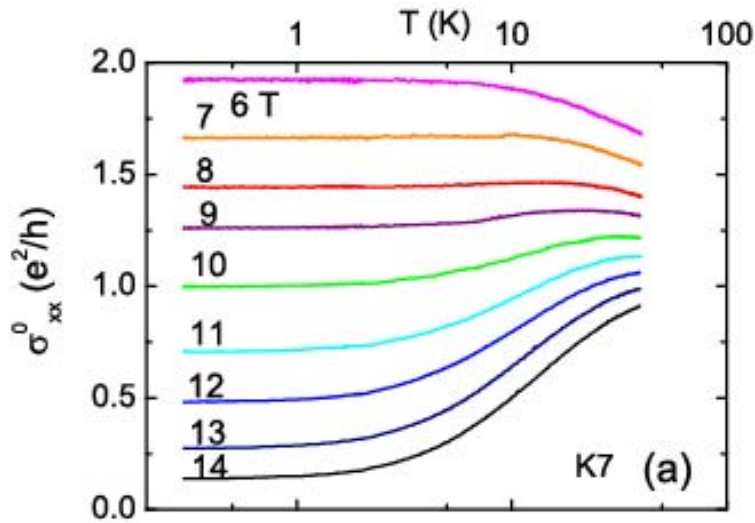
$\sigma_{xy} = 0$ at $n=0$

At 14 T, $n=0$ LL split into 2 subbands

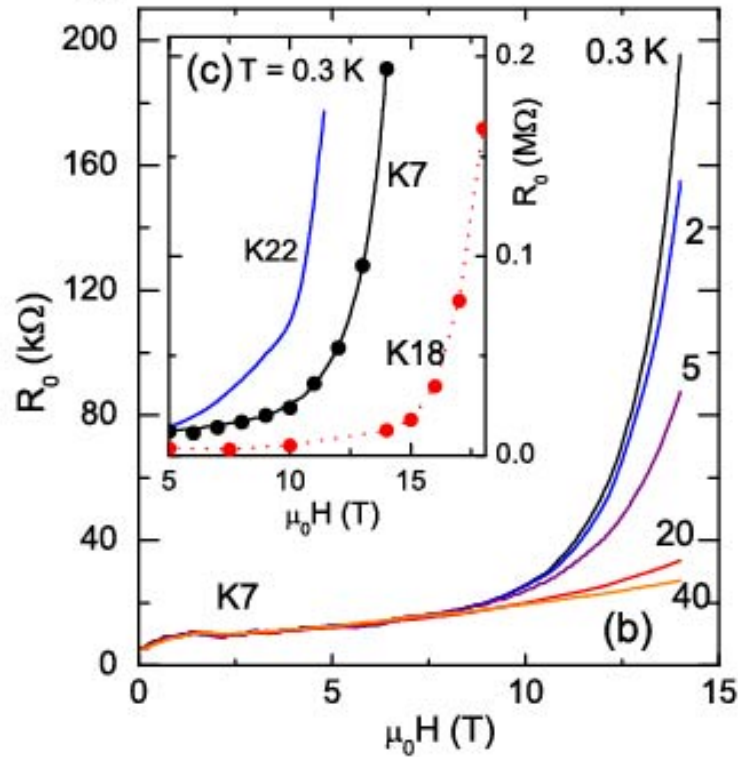
Split resolved at 100 K



R_0 saturates below 2 K

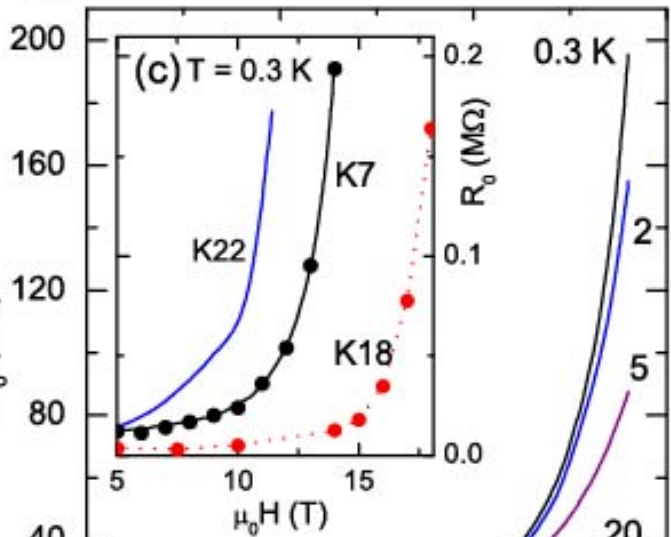


Conductance saturates
T below 2 K

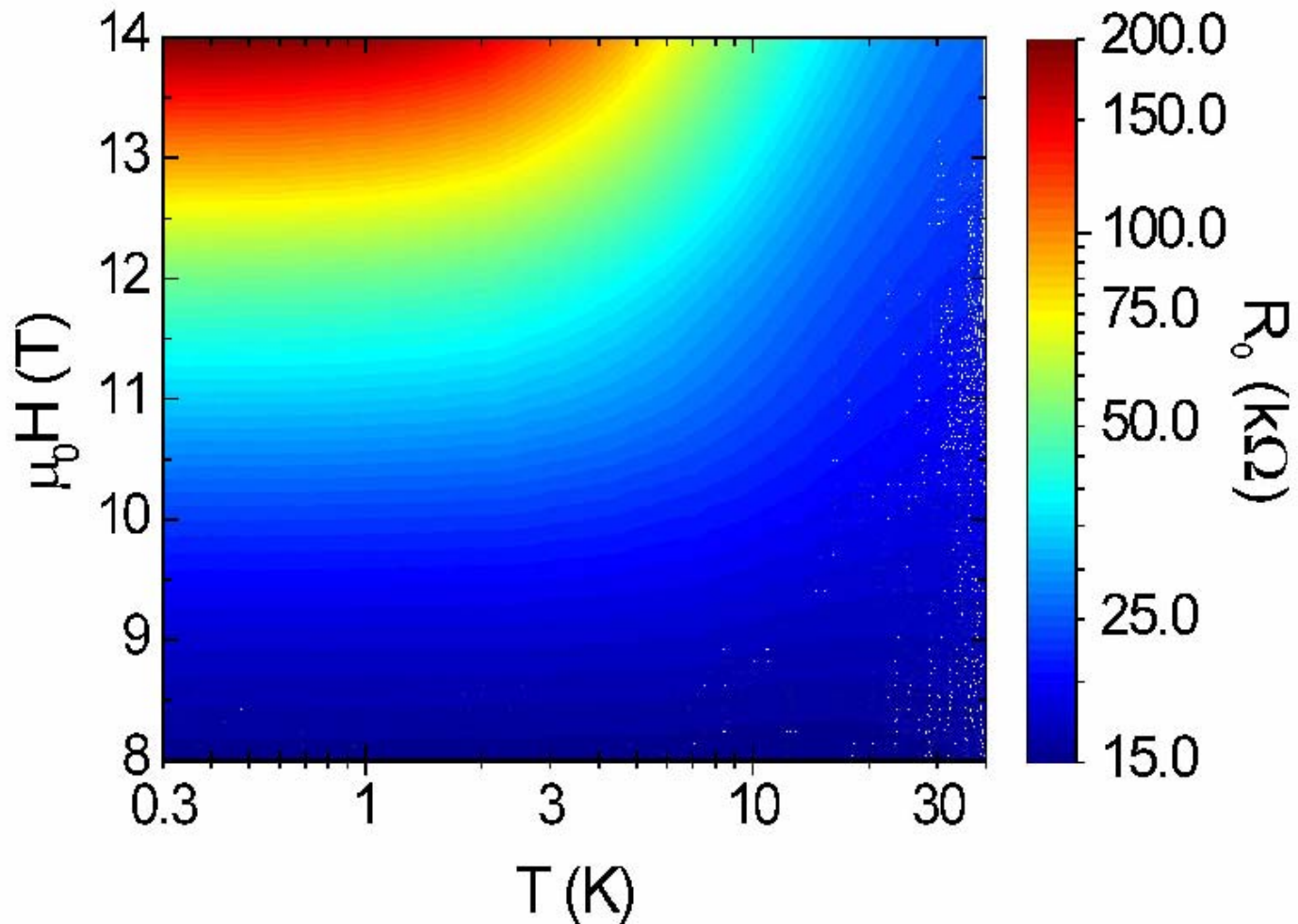


Divergence in R_0 steepest
At 0.3 K

Sample dependent

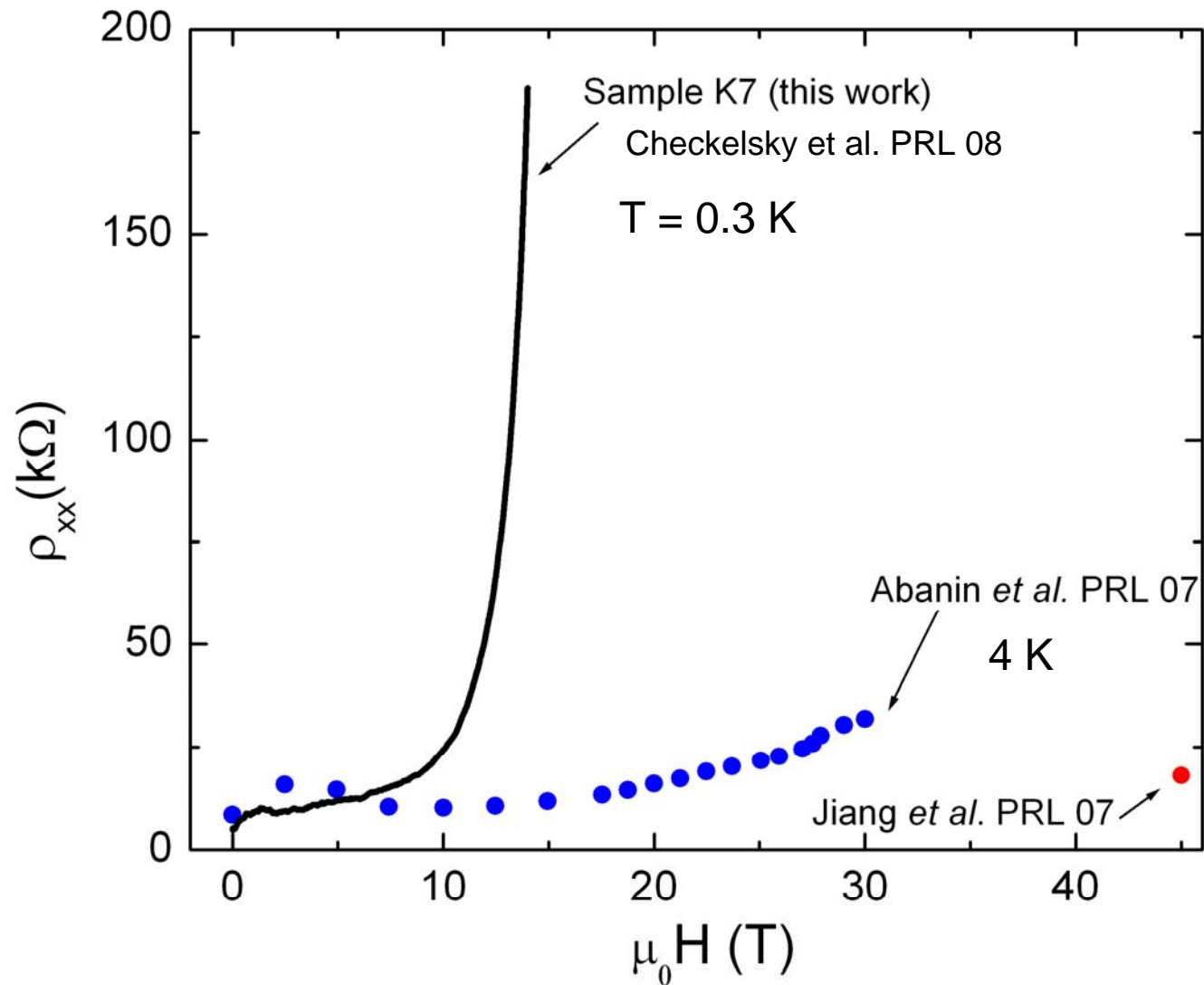


Contour map of R_0 at Dirac point in H - T plane

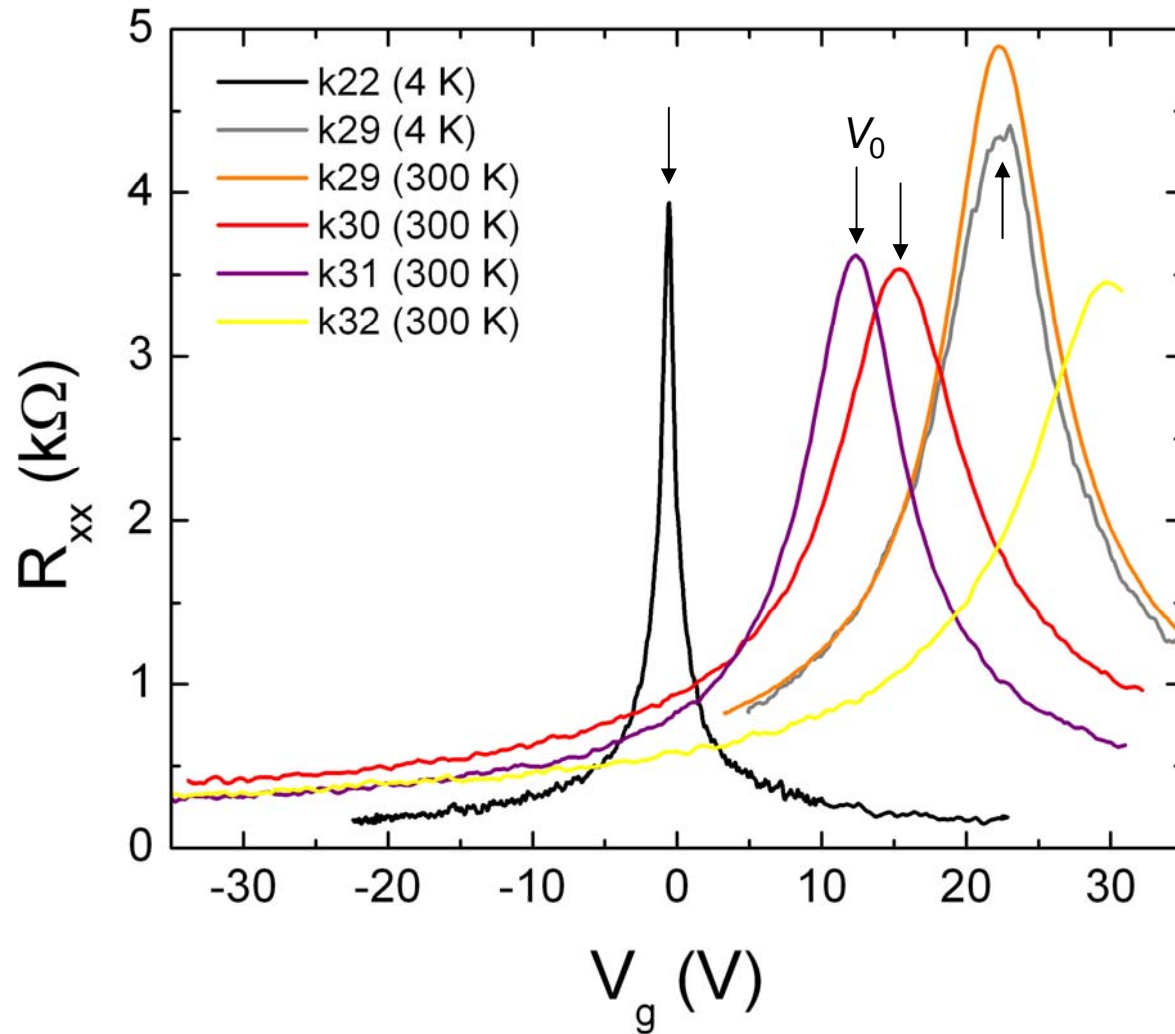


**R_0 is exponentially sensitive to H
but T -independent below 2 K**

Comparison of divergent R_0 with previous reports



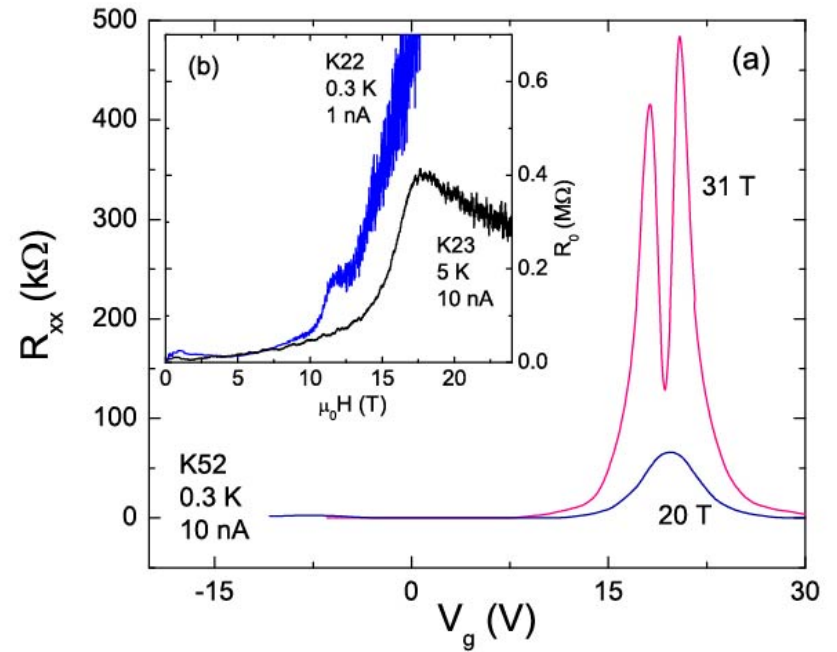
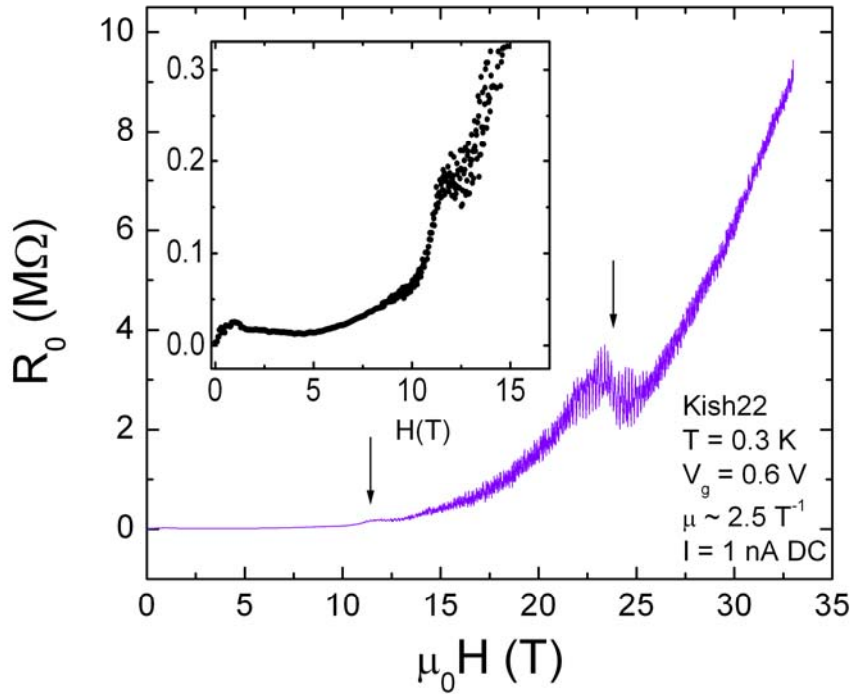
Offset Gate voltage V_0 -- an important parameter



Small offset V_0 correlates with low disorder

Spurious features from self-heating

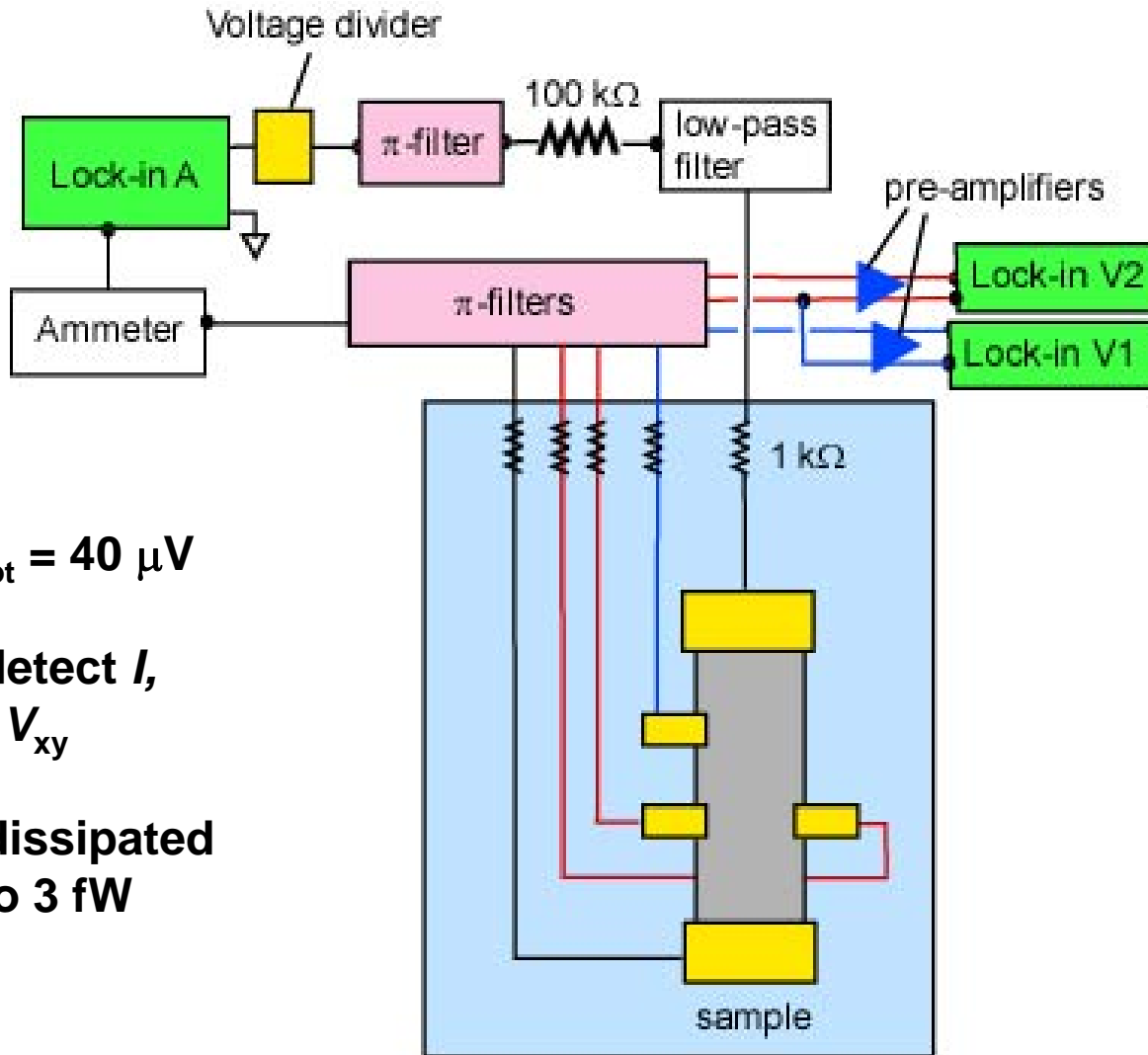
Power dissip $P > 10$ pW leads to thermal runaway at Dirac point



A major experimental obstacle

Adopt ultralow dissipation measurement circuit

Checkelsky, Li, NPO, PRB '09



Hold $V_{\text{tot}} = 40 \mu\text{V}$

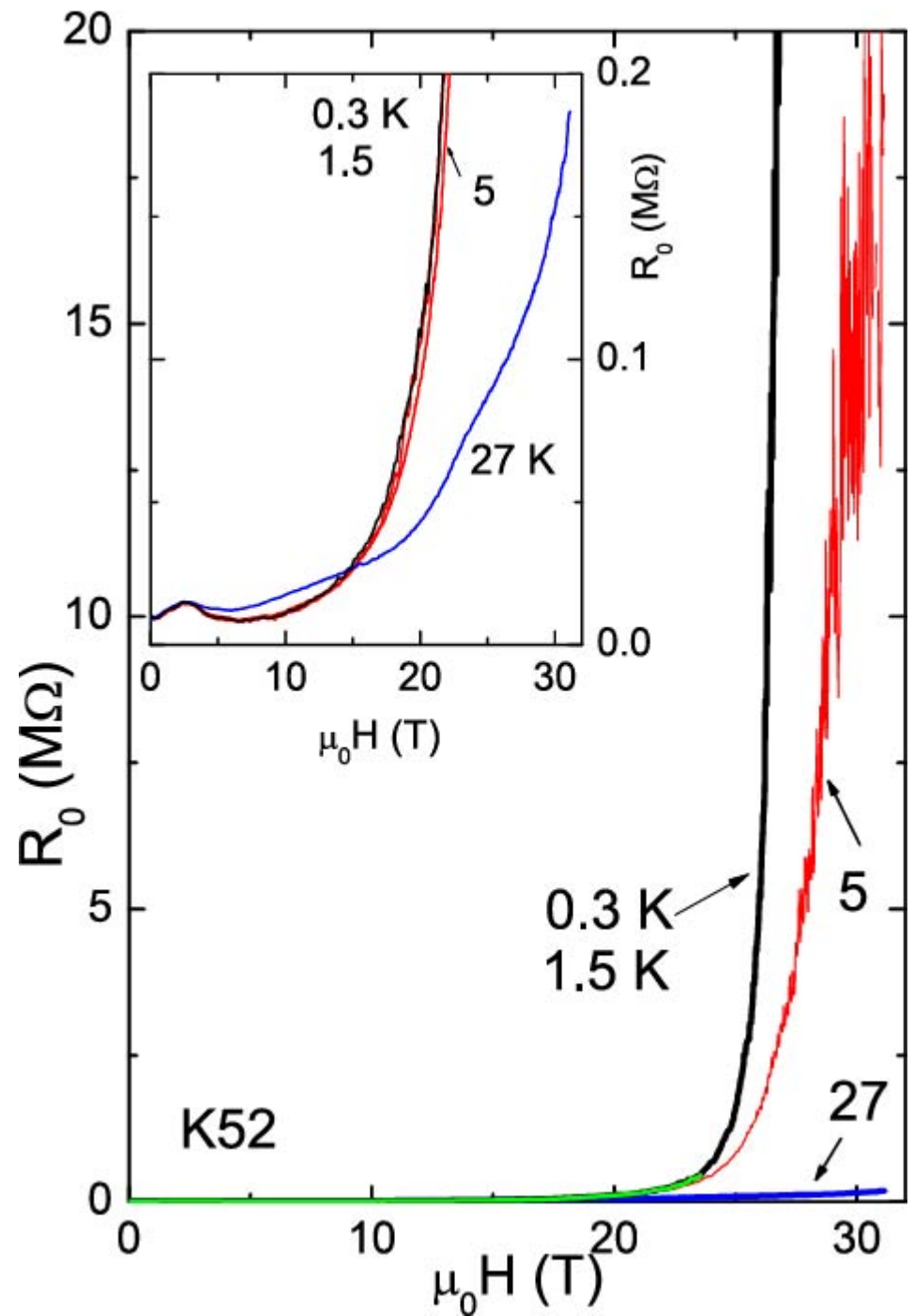
Phase detect I ,
 V_{xx} and V_{xy}

Power dissipated
10 aW to 3 fW

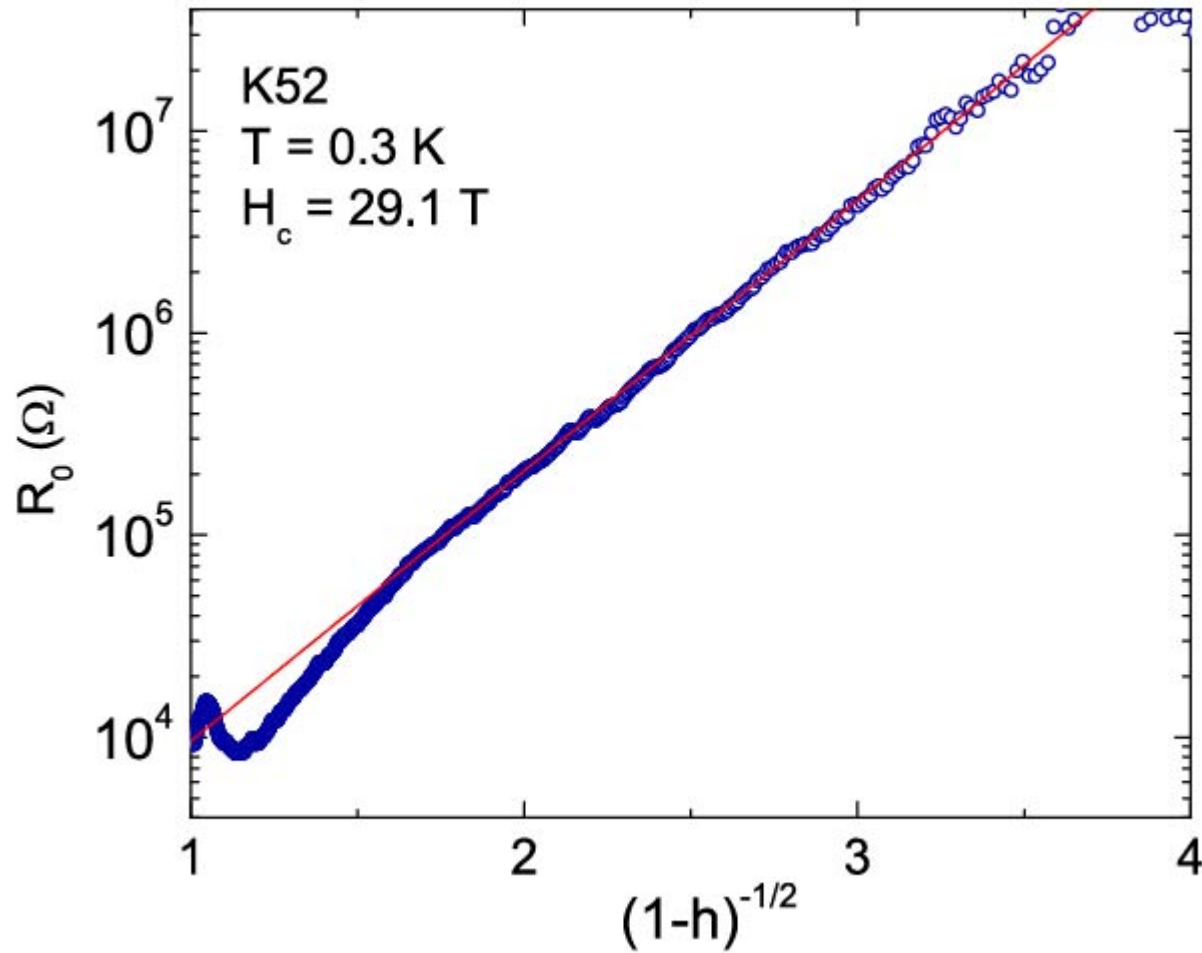
Checkelsky, Li,
NPO, PRB '09

At 0.3 K, R_0 diverges
to 40 M Ω

Divergence in R_0
seems *singular*



Fits to Kosterlitz-Thouless correlation length over 3 decades



3-decade fit to
KT correlation ξ

$$R_0 = 440 \exp \left[\frac{2b}{\sqrt{1-h}} \right]$$

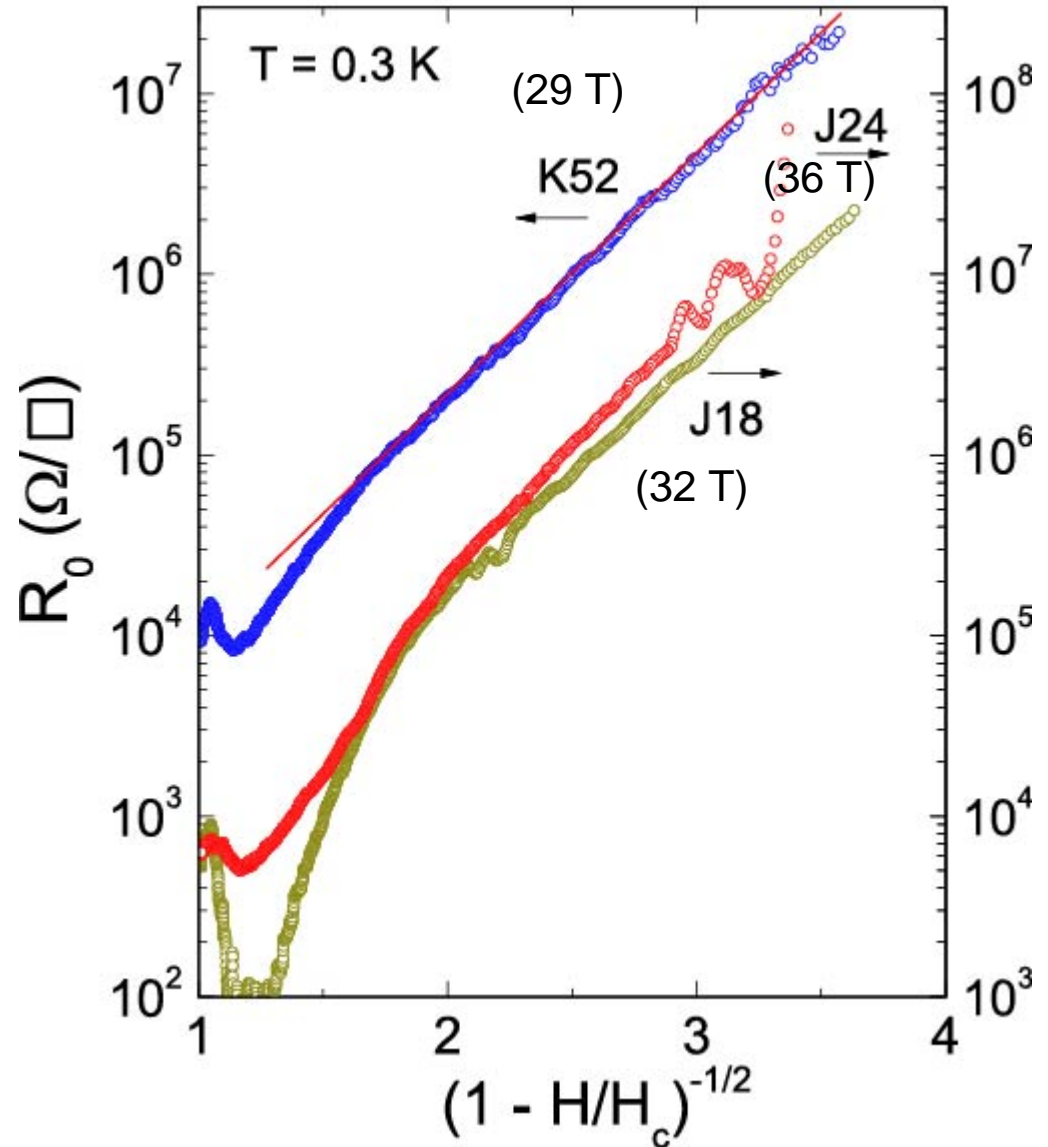
Kosterlitz Thouless fit in Samples K52, J18 and J24

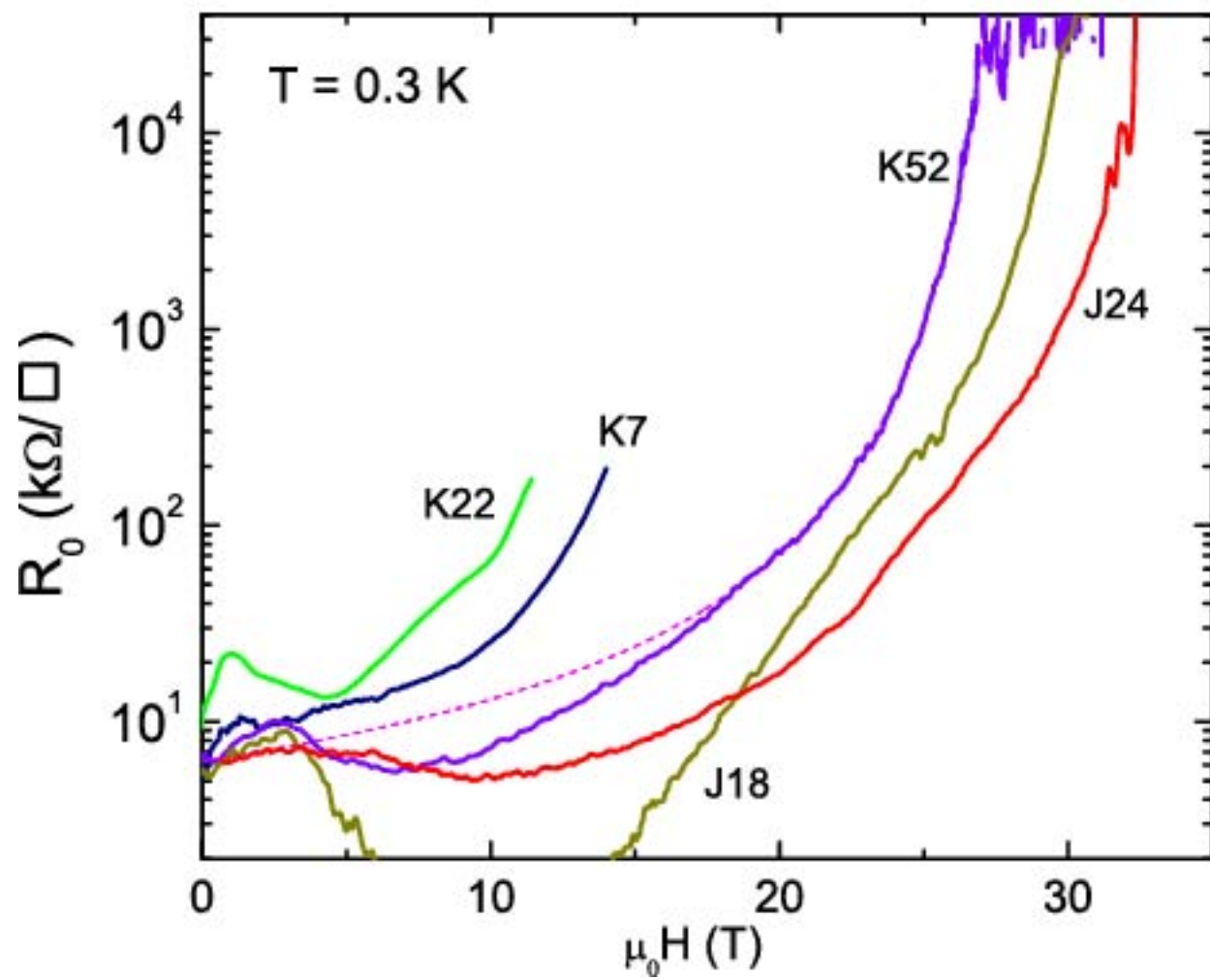
Checkelsky, Li, NPO, PRB '09

$$R_0 = A \exp \left[\frac{2b}{\sqrt{1-h}} \right]$$

Slopes b are similar

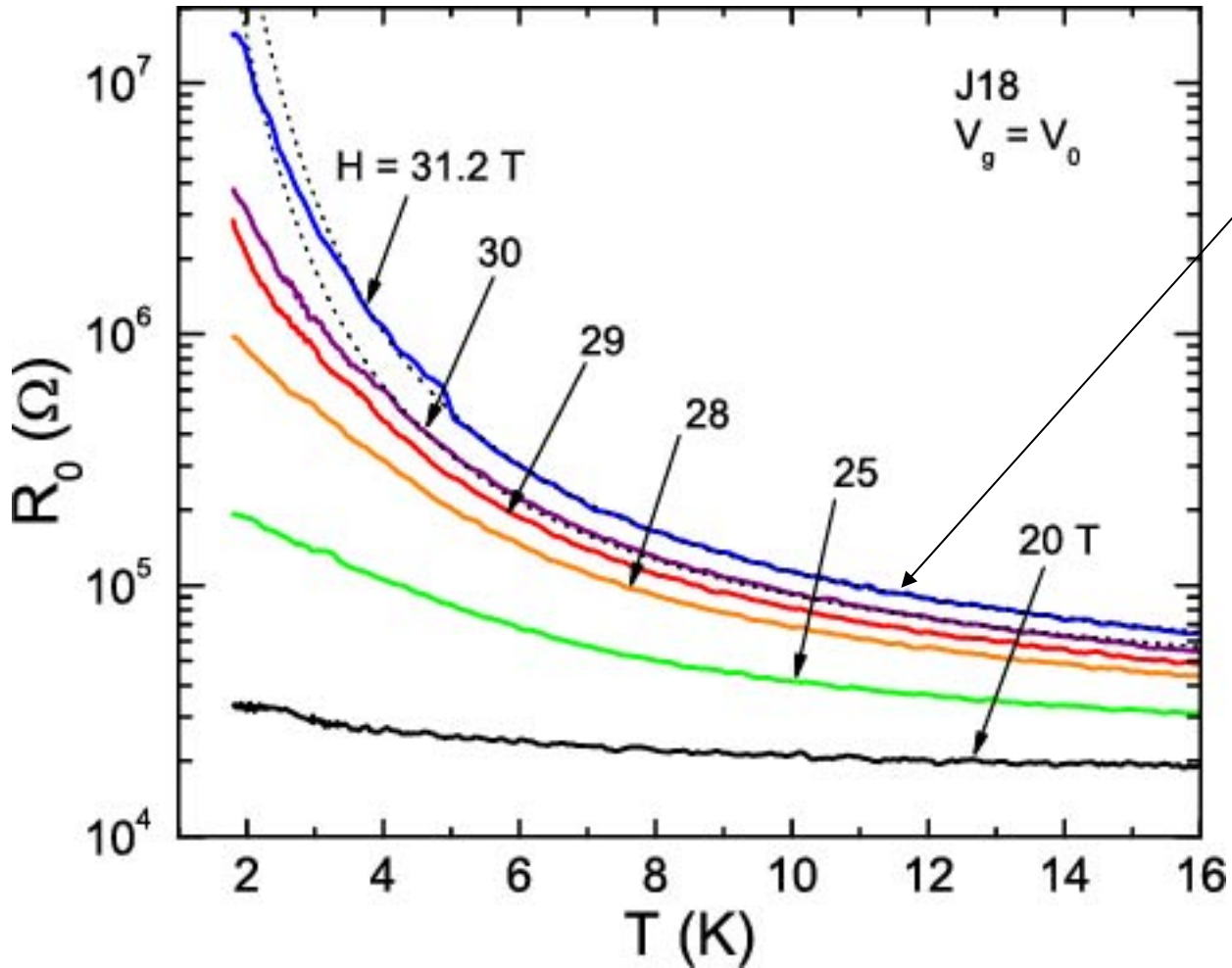
H_c sample dependt





Temperature dependences of R_0 at fields $H < H_c$

Checkelsky, Li, NPO, PRB '09

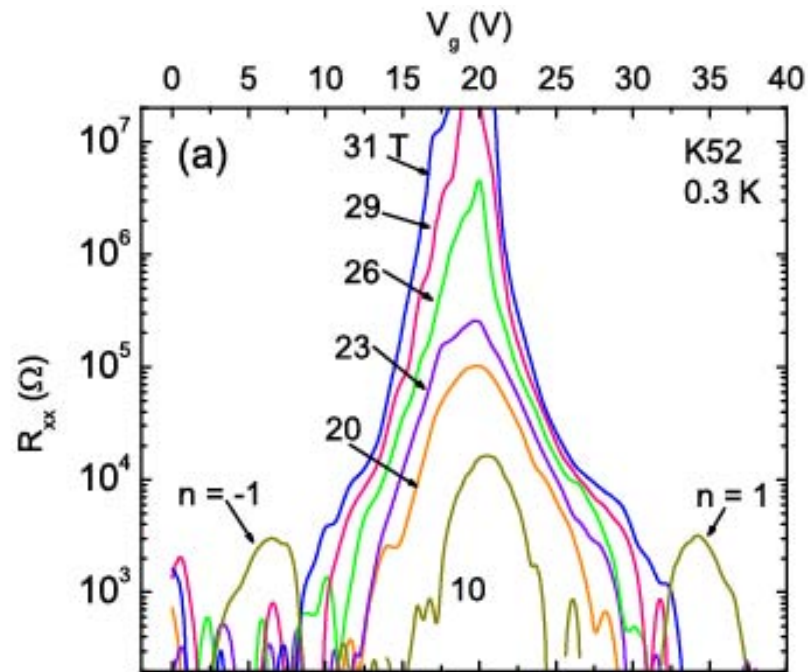


Thermally activated

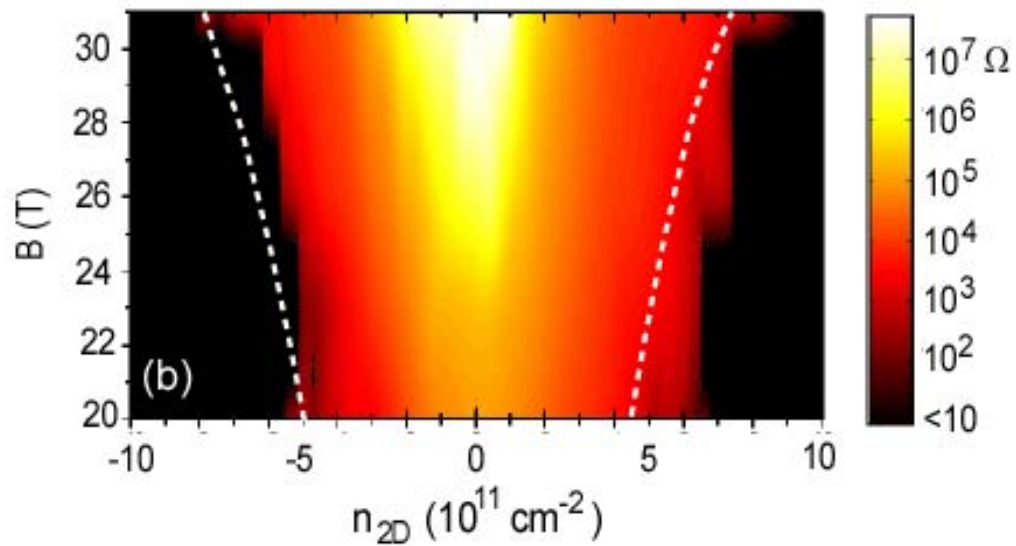
$$R_0 = A \exp(\beta\Delta)$$
$$\Delta = 14 \text{ K}$$

At "low" H , R_0
Saturates below 2 K

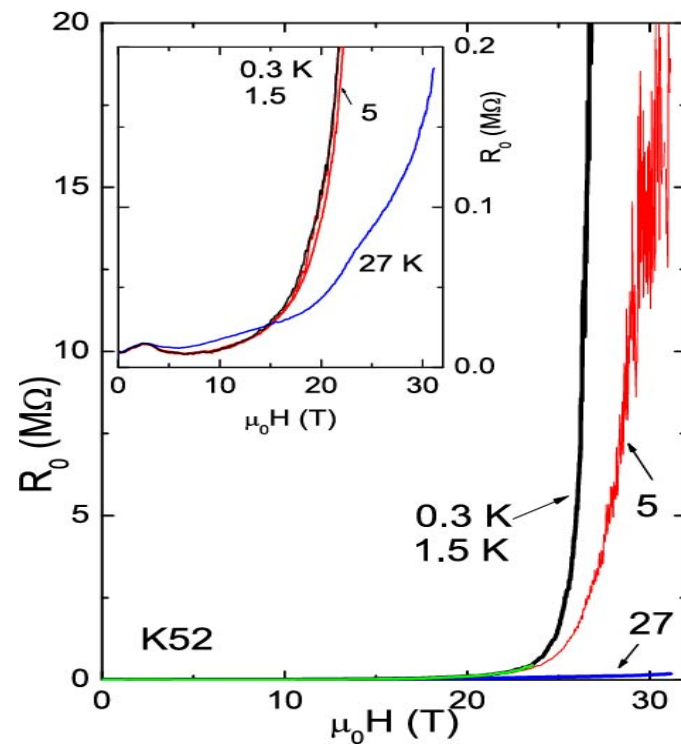
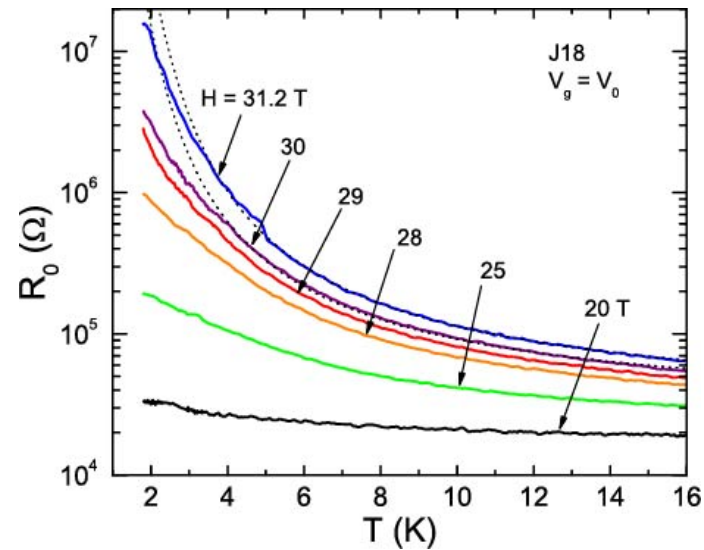
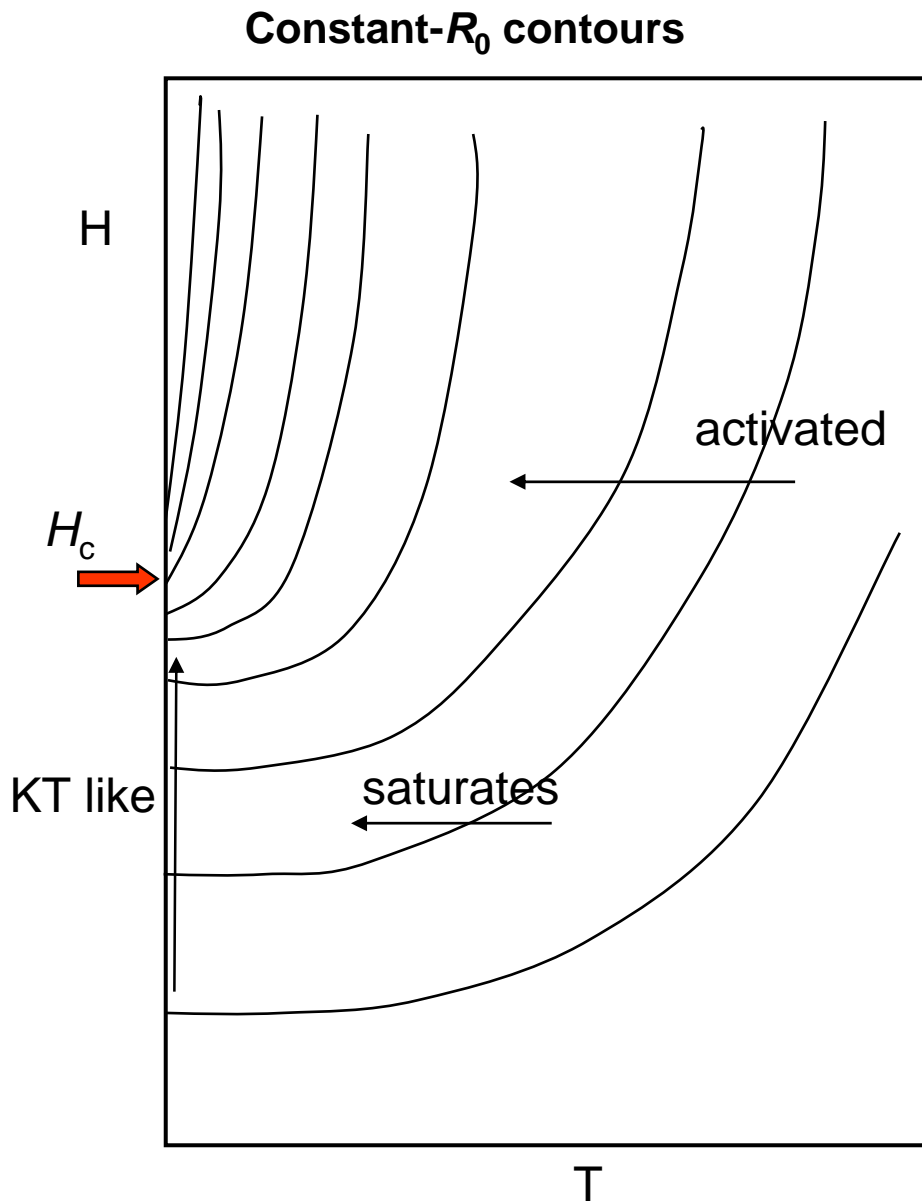
Divergence confined to $n = 0$ LL



$R_0 \sim 10^4$ larger than R_{xx} $n=1$

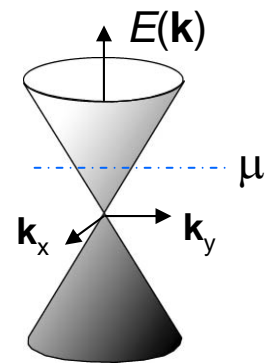
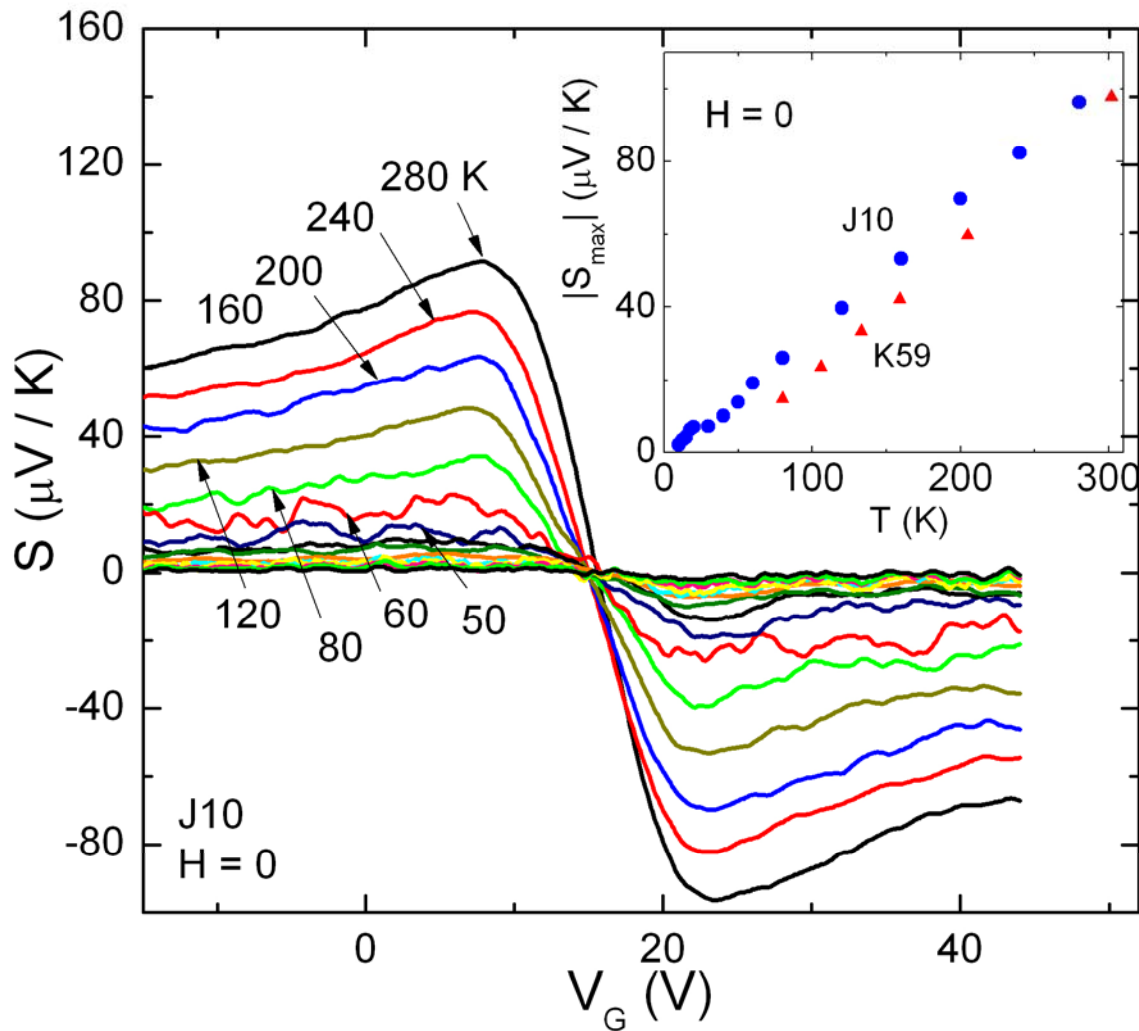


Phase diagram in T - H plane at Dirac point

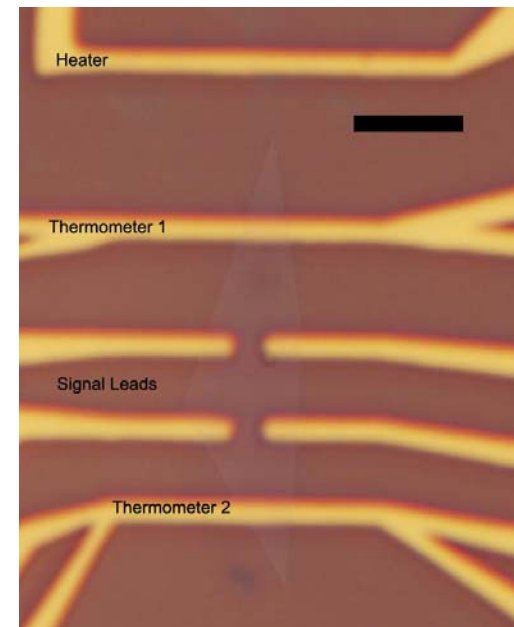


Thermopower in zero field

S is T -linear and large



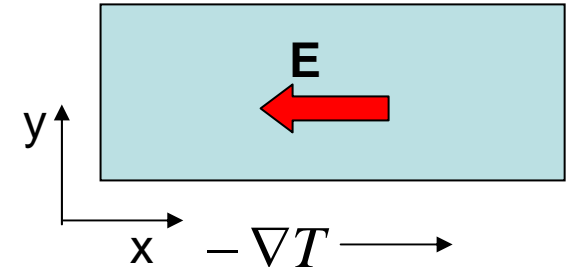
$$S = C \frac{k_B}{e} \frac{1}{\sigma} \frac{\partial \sigma}{\partial \mu}$$



Thermoelectric response functions

Current response $\mathbf{J} = \vec{\alpha} \cdot (-\vec{\nabla}T)$

Measured E field $E_i = S_{ij}(-\partial_j T)$



Total charge current is zero (open boundaries)

$$J_i = \sigma_{ij}E_j + \alpha_{ij}(-\partial_j T)$$

$$S = -S_{xx} = (\rho\alpha + \rho_{yx}\alpha_{xy}) \quad (\text{thermopower})$$

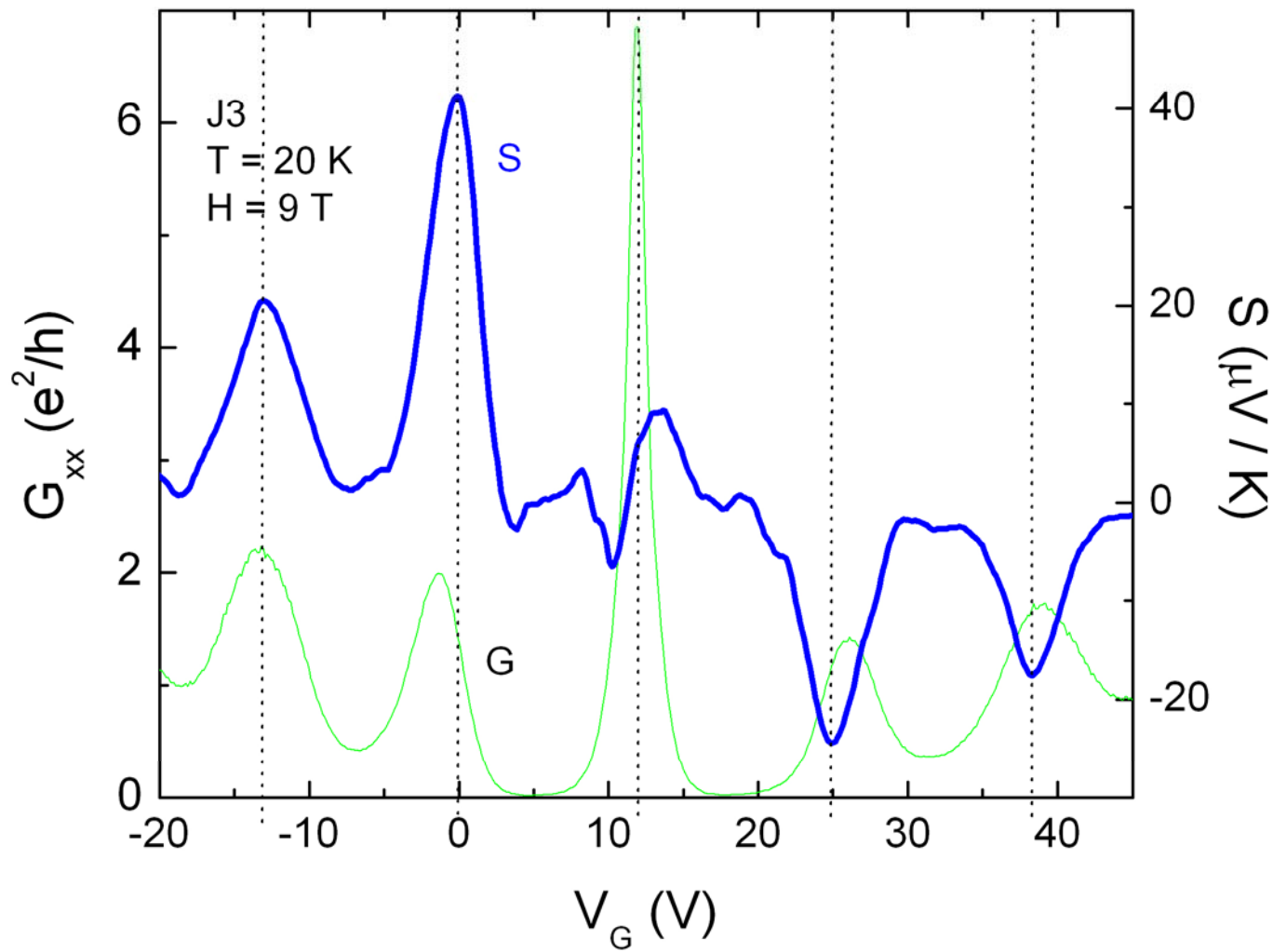
$$S_{yx} = \rho\alpha_{xy} - \alpha\rho_{yx} \quad (\text{Nernst effect})$$

Can invert to find α_{ij}

$$\alpha = -(\sigma E_x + \sigma_{xy}E_y) / |\partial T|$$

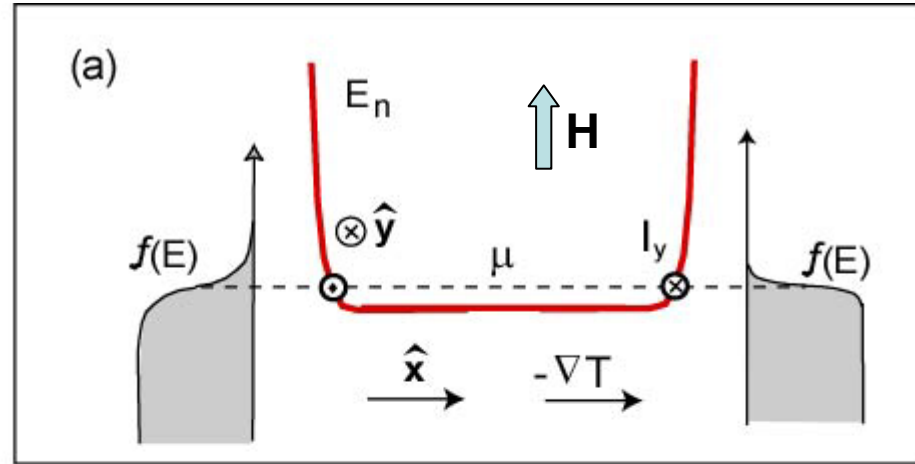
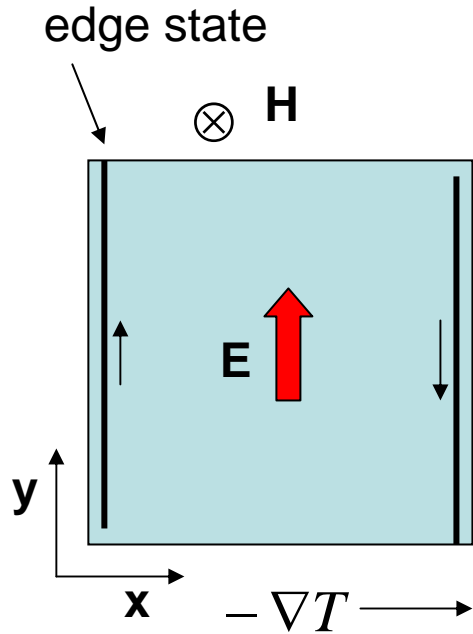
$$\alpha_{xy} = (\sigma_{yx}E_x + \sigma E_y) / |\partial T|$$

Thermopower at 9 T and 20 K



Thermopower in QHE regime ($n > 0$)

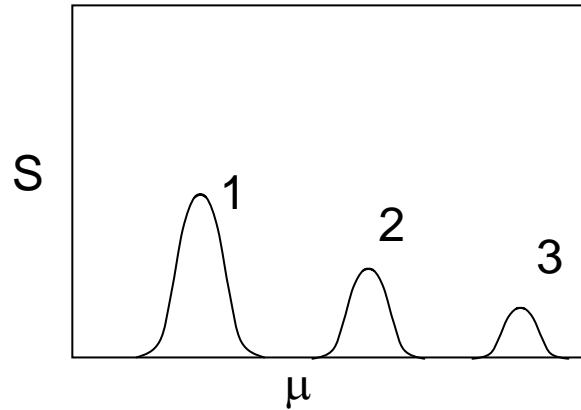
Girvin & Jonson, J Phys C 1982, PRB 1984



$$I_x = 0 \quad I_y = \frac{L}{2\pi} \sum_n \int dk \frac{e}{\hbar L} \hbar v_k f(\varepsilon_{nk})$$

$$\alpha_{xy} = \frac{e}{hT} \sum_n \int_{E_n}^{\infty} d\varepsilon (\varepsilon - \mu) \left(-\frac{\partial f}{\partial \varepsilon} \right)$$

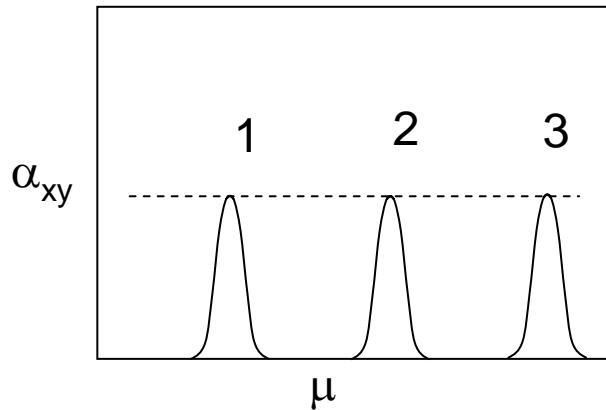
Quantization of thermoelectric current



1) Current response to ΔT is off-diagonal

$$S = \rho_{yx} \alpha_{xy} \quad S_{\max}(n) = \frac{k_B \ln 2}{e(n + 1/2)}$$

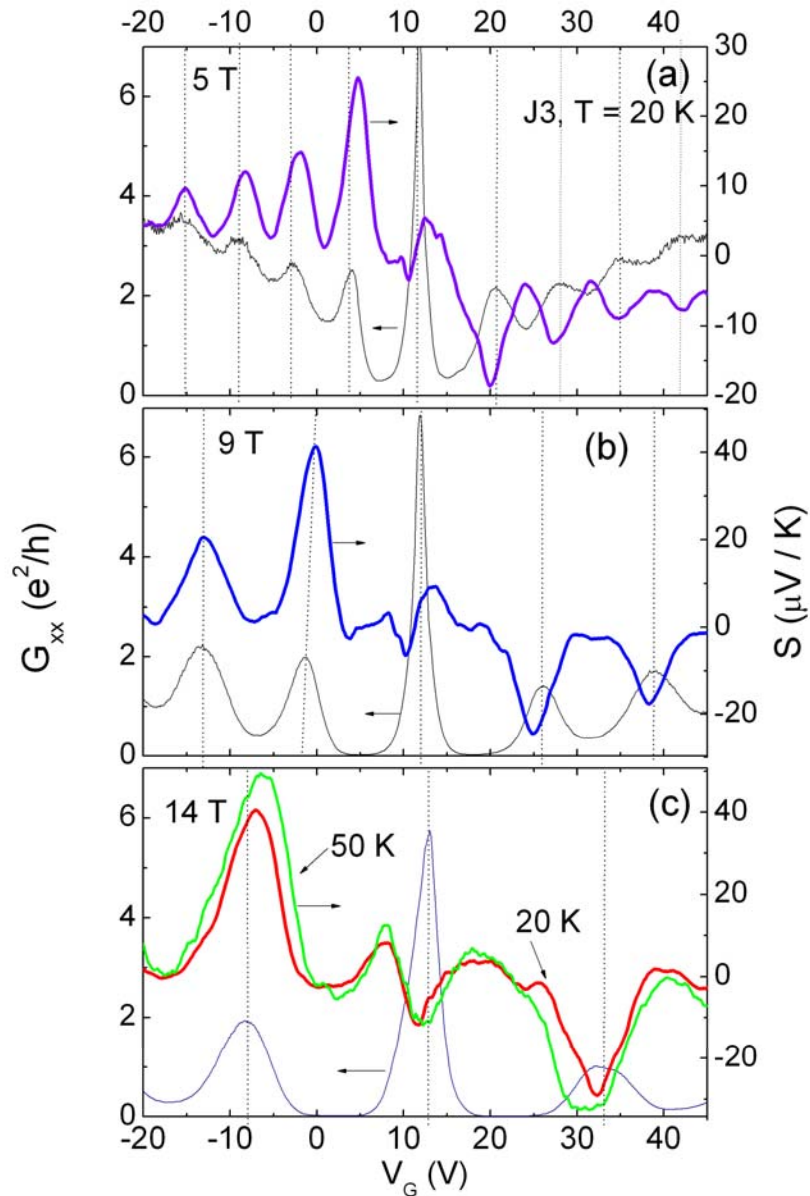
\uparrow
 Dirac



α_{xy} is quantized indpt of n

$$\alpha_{xy} = \frac{k_B e}{h} \ln 2 = 93 \text{ nA/K}$$

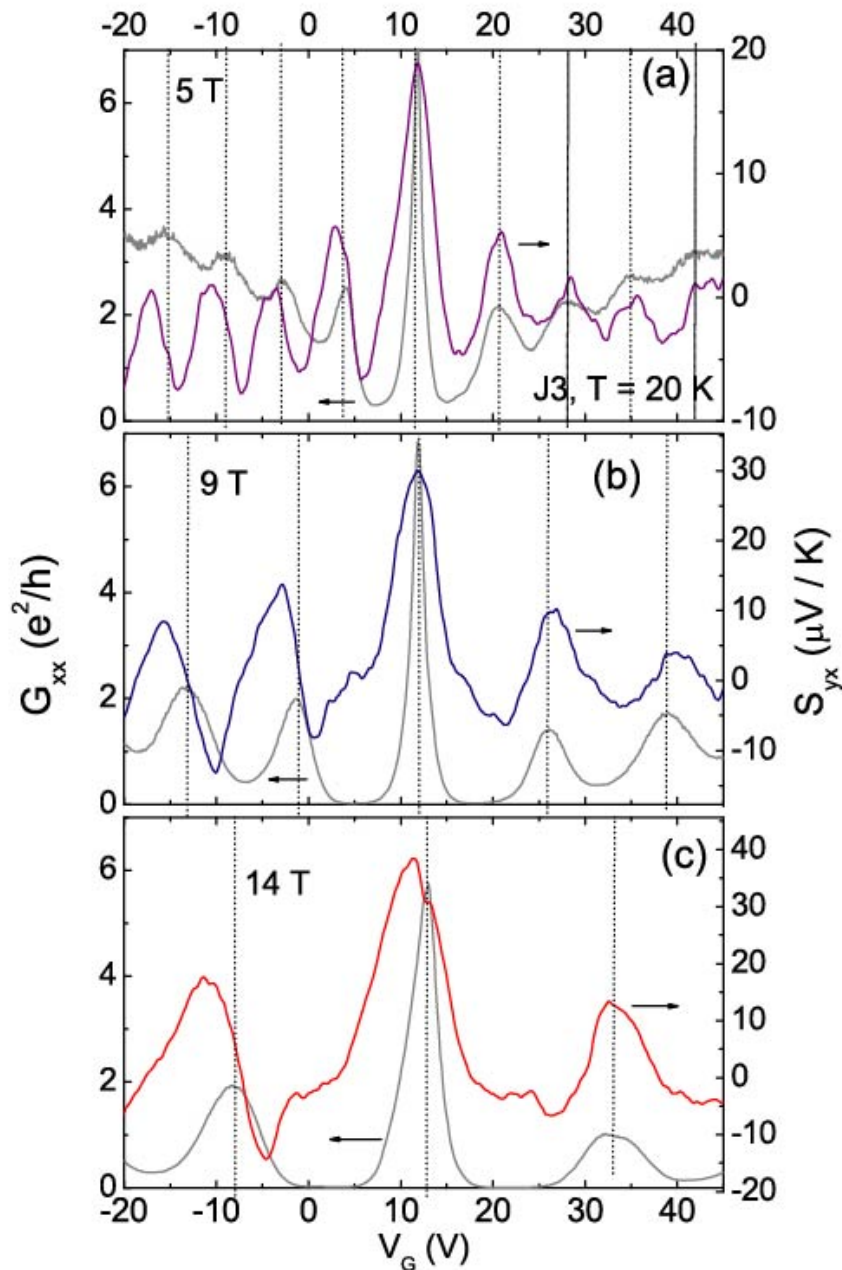
Thermopower at 5, 9 and 14 T



Features of thermopower

1. Charge antisymmetric
2. Peak at $n=1$ nearly indpt of H
3. Very weak T dependence below 50 K

Nernst Signal (off-diagonal S_{yx})



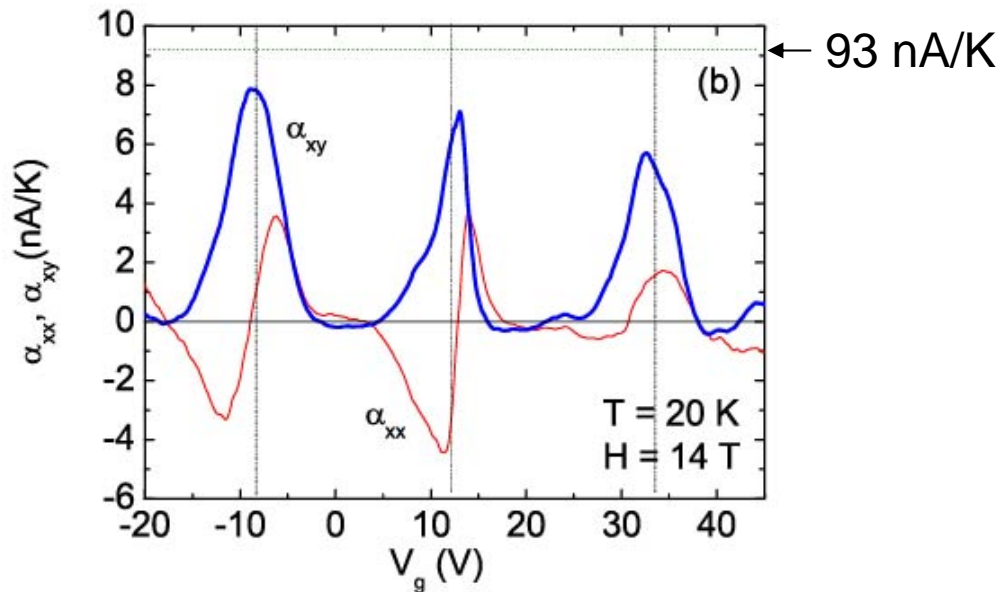
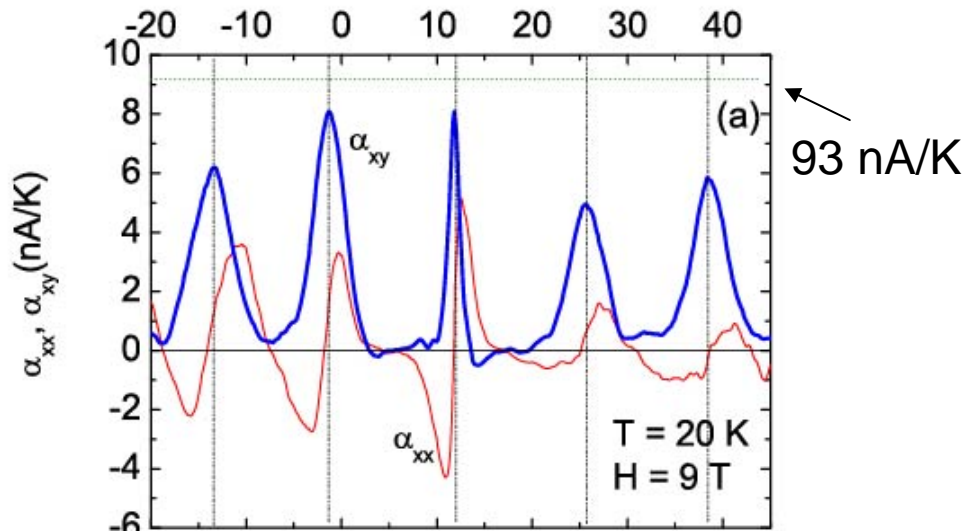
Nernst signal in graphene

- Charge symmetric
 - Nernst signal is anomalously large at Dirac Point
3. Peaks at $n=0$, but dispersive for $n=1, 2, 3, \dots$
 4. At $n=0$, sign is **positive** (similar to vortex-Nernst signal in superconductors)

The thermoelectric tensor elements α_{xx} and α_{xy}

$$\alpha = -(\sigma E_x + \sigma_{xy} E_y) / |\partial T|$$

$$\alpha_{xy} = (\sigma_{yx} E_x + \sigma E_y) / |\partial T|$$



- 1) Peaks well defined
- 2) α_{xx} is antisym. In e
- 3) α_{xy} is symmetric

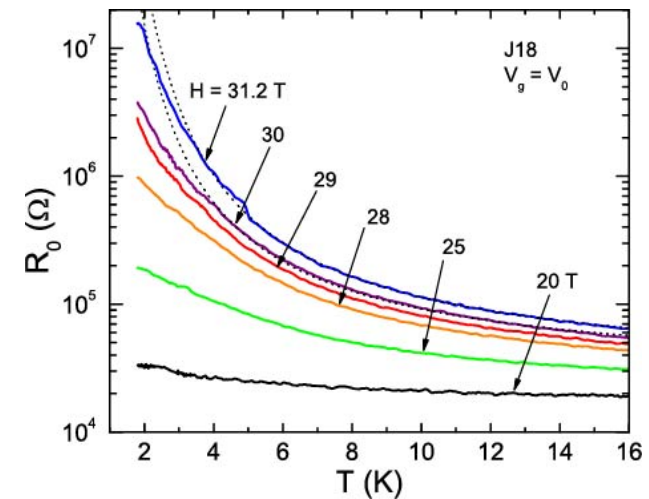
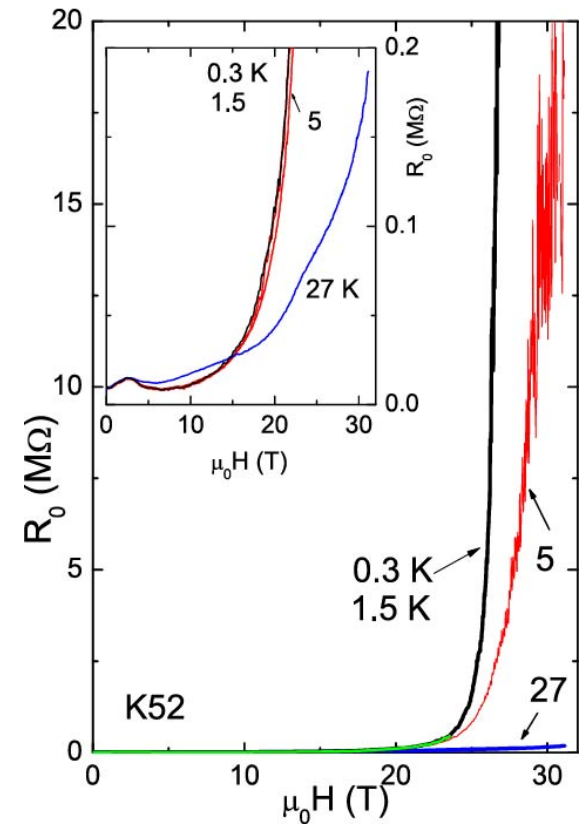
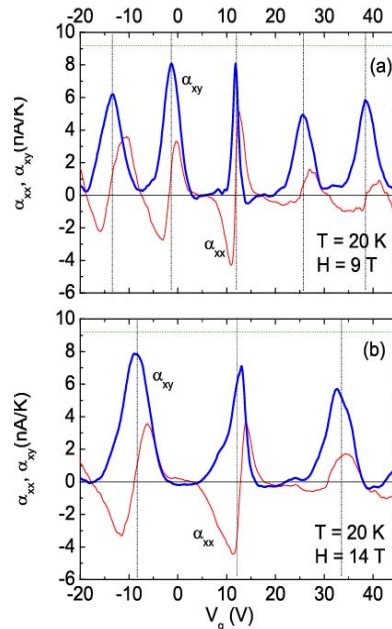
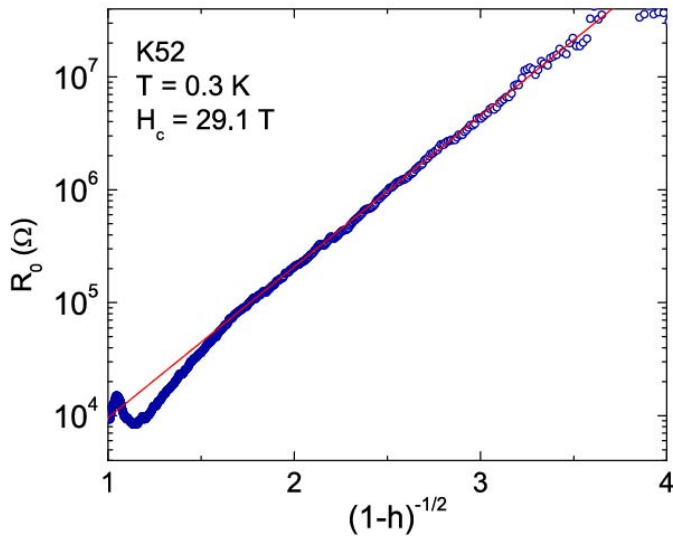
4) A surprise:

α_{xy} peak very narrow at $n=0$

Within 20% of quantized value
(need to estimate gradient
more accurately)

In graphene

1. Field induces transition to insulating state at H_c that correlates with V_0 .
2. For $H > H_c$, R_0 is thermally activated ($\Delta \sim 14$ K), $H < H_c$, R_0 saturates below 2 K.
3. $H < H_c$, divergence fits KT form $\exp[2b/(h-1)^{1/2}]$.
4. Large Nernst signal at Dirac point.
5. α_{xy} consistent with quantization ~ 90 nA/K indpt of T , H and n .



END