Photon Inhibited and Enabled Topological Transport

Tami Pereg-Barnea McGill University

August 14, 2015 - KITP arXiv:1505.05578 arXiv:1505.05584



צדיקים מלאכתם נעשית בידי אחרים

• Saints, their work is done by others



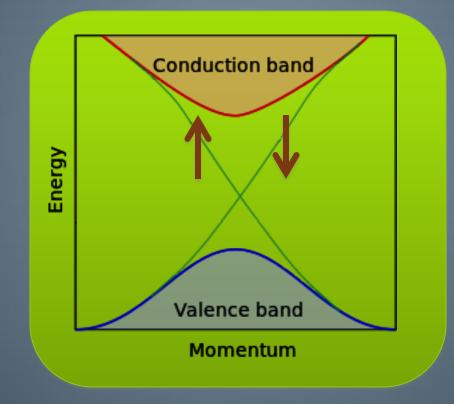
Aaron Farrell

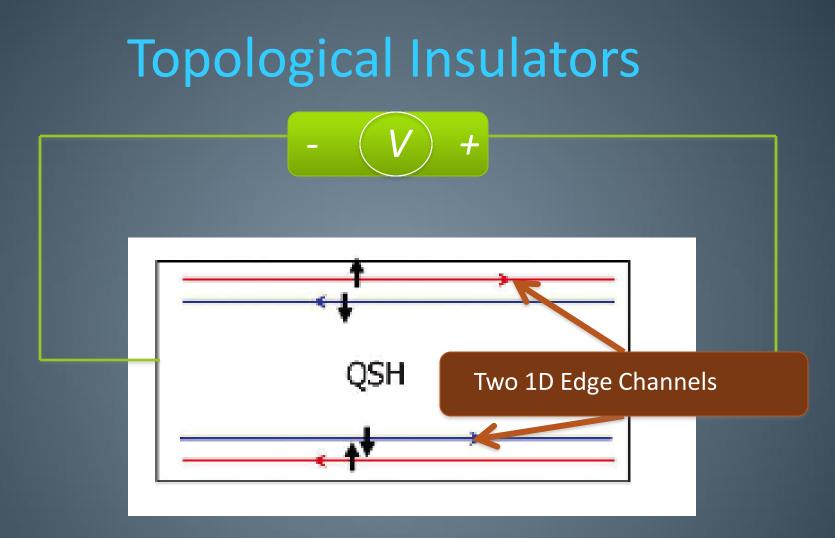
• Acknowledgment: Aash Clerk and Jean-Rene Soquet

Outline

- Floquet topological insulators
- Nonequilibrium transport results
- Photon inhibition of edge-states
- Photon enabling of edge states
- Summary and outlook

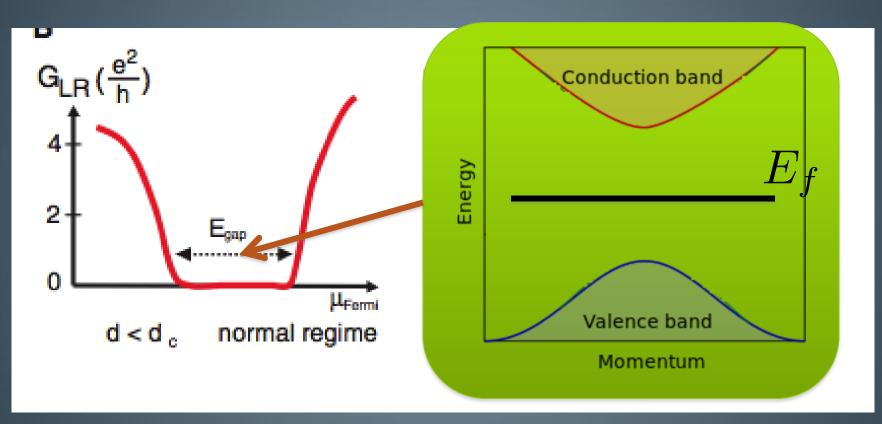
Topological Insulators Quantum Spin Hall





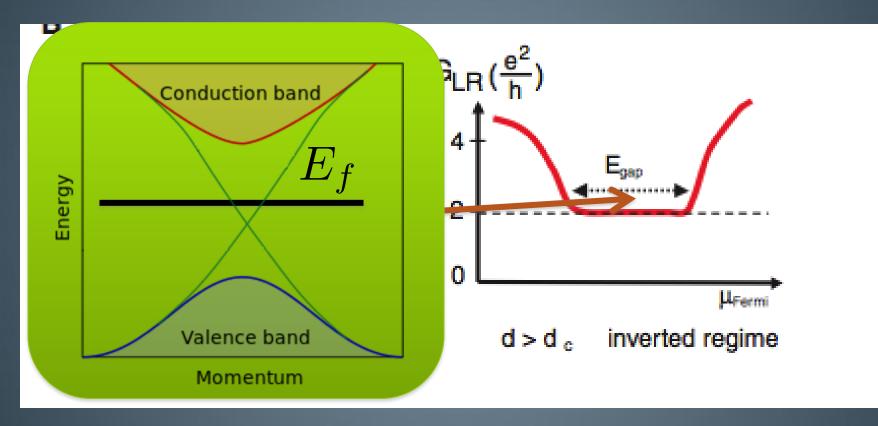
Qi and Zhang, RMP, **83** (2011)

Probes of Topological Insulators



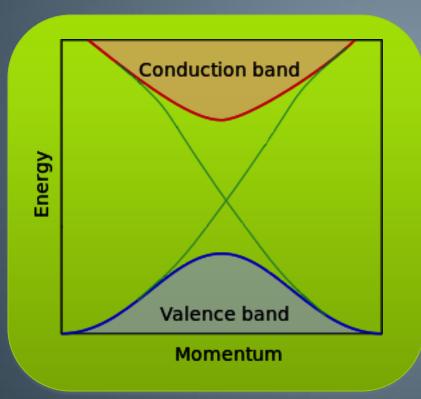
Bernevig et al, Science, 314 (2006)

Probes of Topological Insulators



Bernevig et al, Science, **314** (2006)

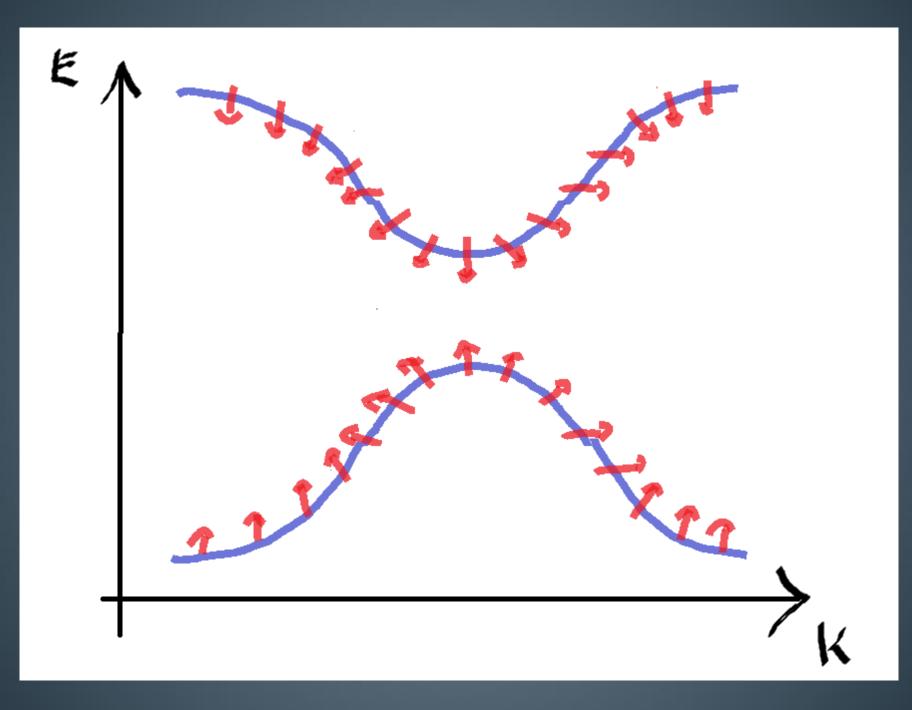
Floquet Topological Insulators





$V(t) = V(t+\tau)$

Lindner, Refael and Galitsky, Nature Physics **7** (2011)



 $i\hbar\partial_t |\psi(t)\rangle = H(t) |\psi(t)\rangle$



Time Periodic Quantum Mechanics $i\hbar\partial_t |\psi(t)\rangle = H(t)|\psi(t)\rangle$ $|\psi(t)\rangle = e^{-i\epsilon t/\hbar}|\phi\rangle$

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Additional Time Dependence

 $i\hbar\partial_t |\psi(t)\rangle = H(t)|\psi(t)\rangle$ $|\psi(t)\rangle = e^{-i\epsilon t/\hbar} |\phi(t)\rangle$ $|\phi(t+T)\rangle = |\phi(t)\rangle$

Time Periodic Quantum Mechanics $|i\hbar\partial_t|\psi(t)\rangle = H(t)|\psi(t)\rangle$ $|\psi(t)\rangle = e^{-i\epsilon t/\hbar} |\phi(t)\rangle$ $|\phi(t+T)\rangle = |\phi(t)\rangle$ Bloch: $\psi(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}}u_{\mathbf{k}}(\mathbf{r})$ $u_{\mathbf{k}}(\mathbf{r} + \mathbf{R}) = u_{\mathbf{k}}(\mathbf{r})$

Time Periodic Quantum Mechanics $i\hbar\partial_t |\psi(t)\rangle = H(t)|\psi(t)\rangle$ $|\psi(t)\rangle = e^{-i\epsilon t/\hbar}|\phi(t)\rangle$ $|\phi(t+T)\rangle = |\phi(t)\rangle$

 $(H(t) - i\hbar\partial_t)|\phi(t)\rangle = \epsilon|\phi(t)\rangle$

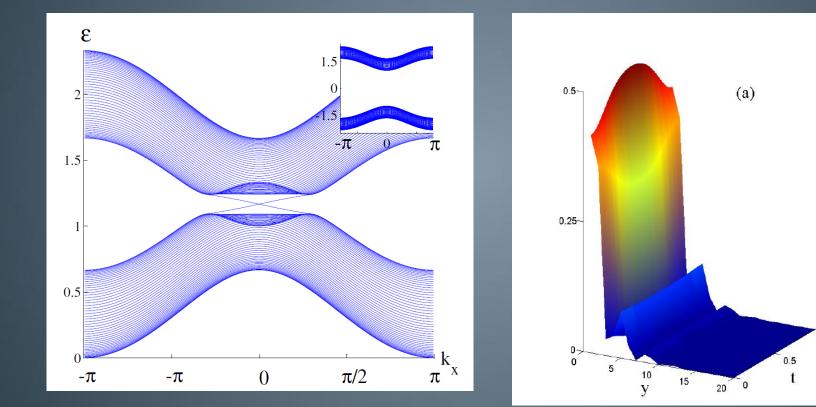
Time Periodic Quantum Mechanics $i\hbar\partial_t |\psi(t)\rangle = H(t)|\psi(t)\rangle$ $|\psi(t)\rangle = e^{-i\epsilon t/\hbar} |\phi(t)\rangle$ $|\phi(t+T)\rangle = |\phi(t)\rangle$ $\overline{(H(t) - i\hbar\partial_t)}\phi(t) = \epsilon |\phi(t)\rangle$

Quasi-energy

Floquet Topological Insulators

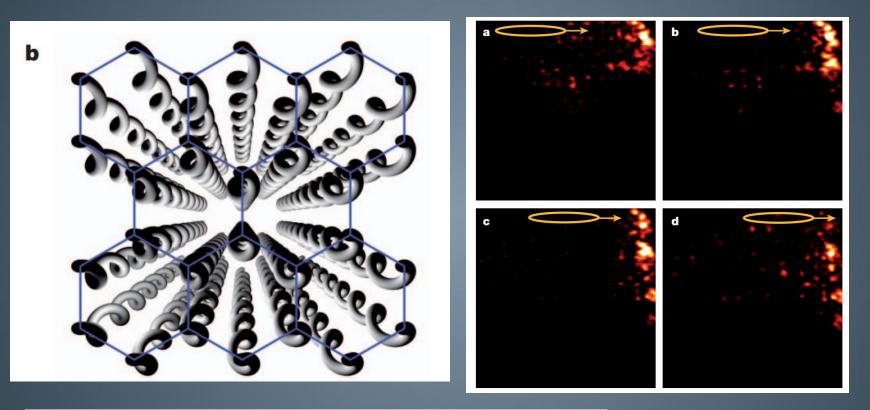
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Lindner *et al*, Nature Physics **7** (2011)

Photonic Floquet TI



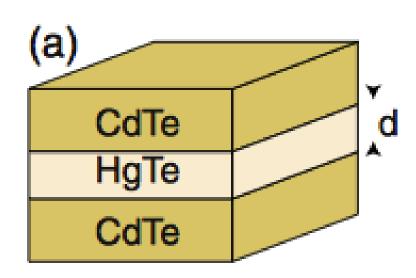
$$i\partial_z\psi(x,y,z) = -\frac{1}{2k_0}\nabla^2\psi(x,y,z) - \frac{k_0\Delta n(x,y,z)}{n_0}\psi(x,y,z)$$

Rechtsman et al., Nature 496, 196 (2013)

Goal

- Study transport properties of topological edge states in the presence of time periodic potential
- Study both forms of edge states: naturally occurring and driven.

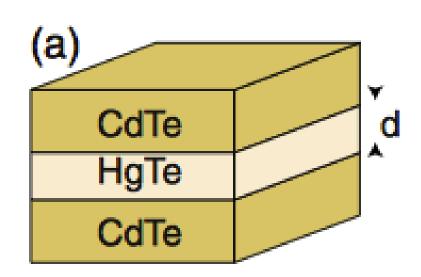
 For concreteness we study quantum well heterostructures



 $egin{pmatrix} H(\mathbf{k}) & 0 \ 0 & H^*(-\mathbf{k}) \ \end{pmatrix}$ $\tilde{H}(\mathbf{k})$

Hasan and Kane, RMP 82 (2010)

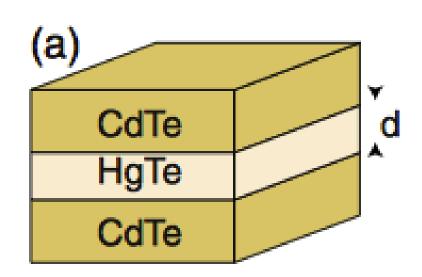
 For concreteness we study quantum well heterostructures



 $\tilde{H}(\mathbf{k}) = \begin{pmatrix} H(\mathbf{k}) & 0\\ 0 & H^*(-\mathbf{k}) \end{pmatrix}$ $H(\mathbf{k}) = \epsilon(\mathbf{k})I + \vec{d}_{\mathbf{k}} \cdot \vec{\sigma}$

Hasan and Kane, RMP 82 (2010)

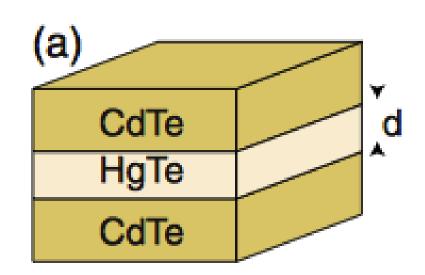
 For concreteness we study quantum well heterostructures



Hasan and Kane, RMP 82 (2010)

$$\tilde{H}(\mathbf{k}) = \begin{pmatrix} H(\mathbf{k}) & 0\\ 0 & H^*(-\mathbf{k}) \end{pmatrix}$$
$$H(\mathbf{k}) = \epsilon(\mathbf{k})I + \vec{d}_{\mathbf{k}} \cdot \vec{\sigma}$$
$$V(t) = 2V_{ext}\sigma_z \cos \Omega t$$

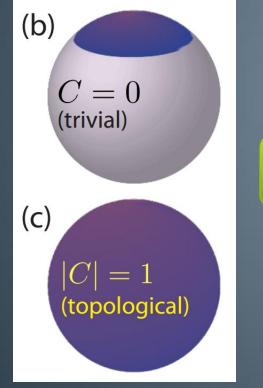
 For concreteness we study quantum well heterostructures



Hasan and Kane, RMP 82 (2010)

$$\tilde{H}(\mathbf{k}) = \begin{pmatrix} H(\mathbf{k}) & 0\\ 0 & H^*(-\mathbf{k}) \end{pmatrix}$$
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$$H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$$

 For concreteness we study quantum well heterostructures

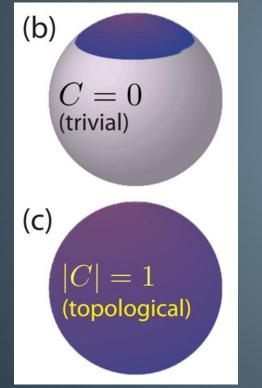


$$H(\mathbf{k}) = \epsilon(\mathbf{k})I + \vec{d}_{\mathbf{k}} \cdot \vec{\sigma}$$

 $\vec{d}_{\mathbf{k}} = (A\sin k_x, A\sin k_y, M + 2B(\cos k_x + \cos k_y - 2))$

Alicea, Rep. Prog. Phys. 75 (2012) 076501

 For concreteness we study quantum well heterostructures



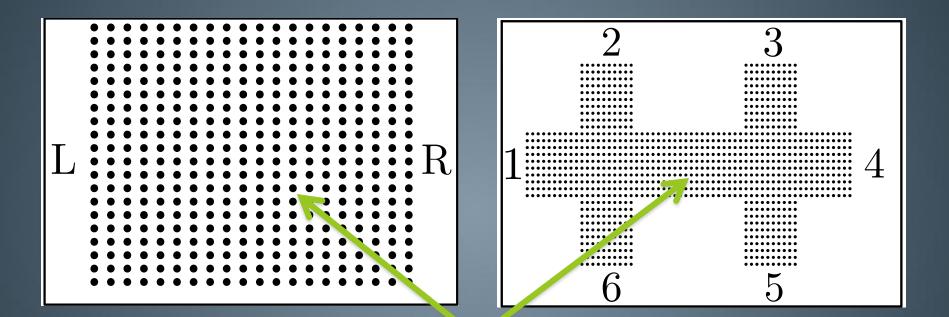
$$H(\mathbf{k}) = \epsilon(\mathbf{k})I + d\mathbf{k} \cdot \vec{\sigma}$$

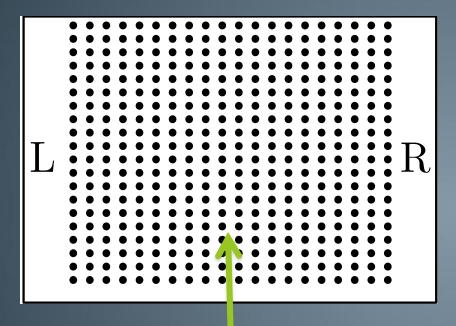
 $\vec{d}_{\mathbf{k}} = (A\sin k_x, A\sin k_y, M + 2B(\cos k_x + \cos k_y - 2))$

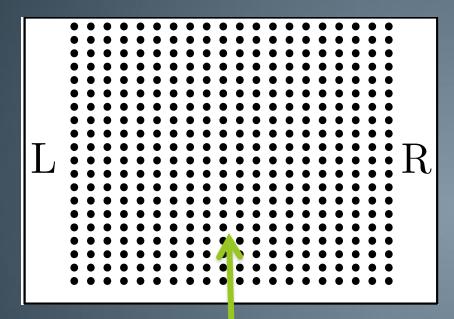
sign(*MB*)>0 \rightarrow Topological sign(*MB*)<0 \rightarrow Trivial

Alicea, Rep. Prog. Phys. 75 (2012) 076501

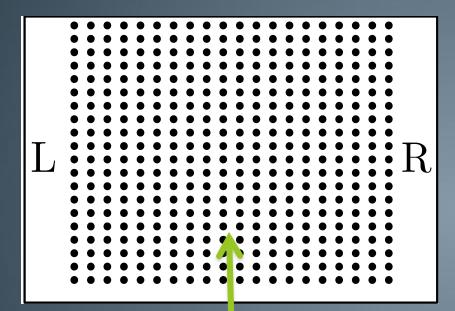
- Focus on topological parameter regime to gain insight.
- Look for signature transport values and robustness





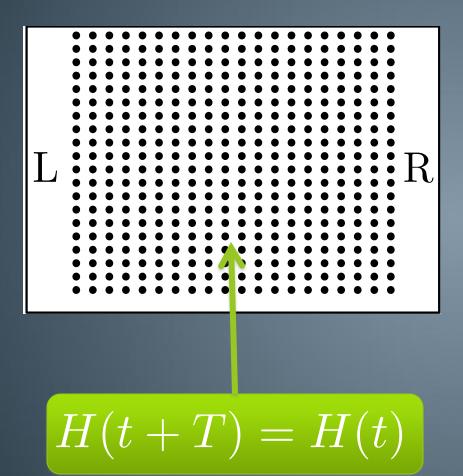


Lead $\alpha \to \Sigma_{\alpha}$



Lead $\alpha \to \Sigma_{\alpha}$

 $H(t) \to H(t) + \Sigma$



Lead $\alpha \to \Sigma_{\alpha}$

 $\overline{H(t)} \to H(t) + \Sigma$

 $\Sigma = \sum \Sigma_{\alpha} = \frac{i}{2}\Gamma$

Landauer Formalism

$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$

Landauer Formalism

$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$

$$G(\epsilon) = \int dt' e^{i\epsilon(t-t')} G(t,t')$$

Landauer Formalism

$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$

 $G(\epsilon) = \int dt' e^{i\epsilon(t-t')} G(t,t')$

G(t, t') = G(t - t')

$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$

$$G^{(n)}(\epsilon) = \frac{1}{T} \int_0^T dt e^{-in\Omega t} \int dt' e^{i\epsilon(t-t')} G(t,t')$$

Time Periodic Landauer Formalism

$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$

$$G^{(n)}(\epsilon) = \frac{1}{T} \int_0^T dt e^{-in\Omega t} \int dt' e^{i\epsilon(t-t')} G(t,t')$$

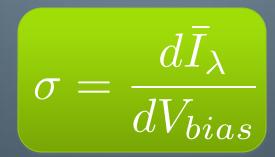
Green's function of $H(t) + \Sigma$

Time Periodic Landauer Formalism

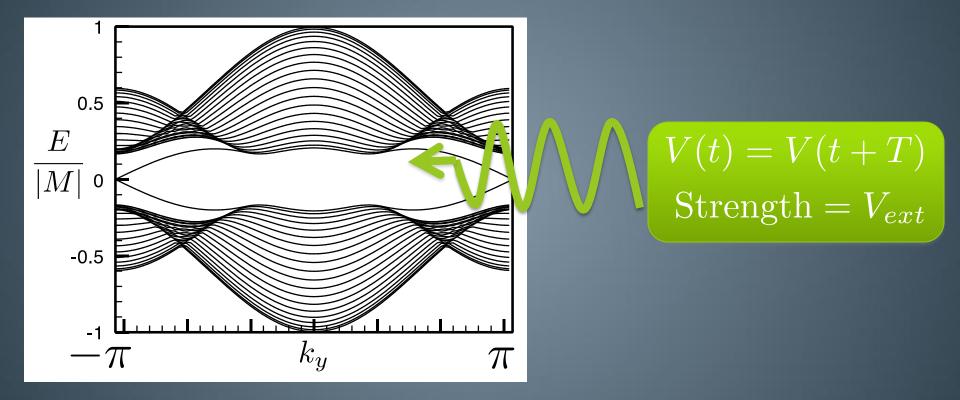
$$\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(T_{\lambda,\lambda'}(\epsilon) f_{\lambda'}(\epsilon) - T_{\lambda',\lambda}(\epsilon) f_{\lambda}(\epsilon) \right)$$
$$G^{(n)}(\epsilon) = \frac{1}{T} \int_{0}^{T} dt e^{-in\Omega t} \int dt' e^{i\epsilon(t-t')} G(t,t')$$
$$T_{\lambda,\lambda'}(\epsilon) = \sum_{n} \operatorname{Tr} \left[\Gamma_{\lambda} G^{(n)}(\epsilon) \Gamma_{\lambda'}(G^{(n)}(\epsilon))^{\dagger} \right]$$

Time Periodic Landauer Formalism

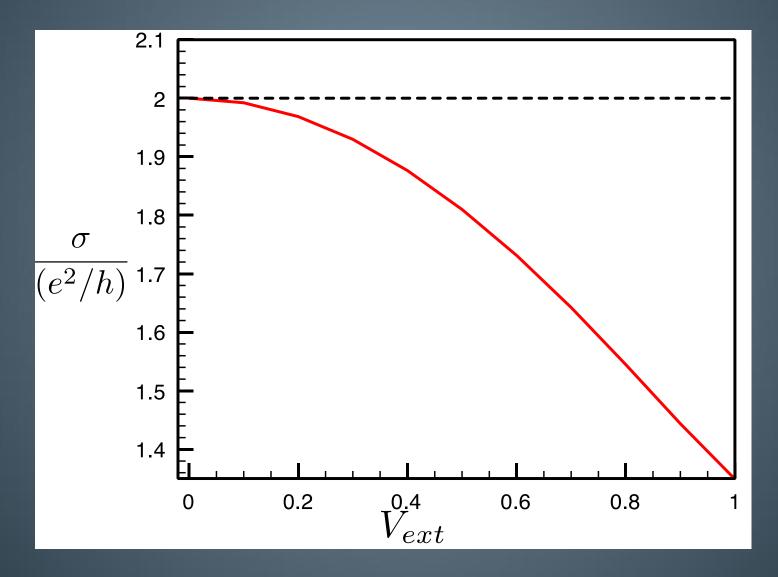
 $\bar{I}_{\lambda} = \frac{e}{h} \sum_{\lambda'} \int d\epsilon \left(\overline{T_{\lambda,\lambda'}(\epsilon)} f_{\lambda'}(\epsilon) - \overline{T_{\lambda',\lambda}(\epsilon)} f_{\lambda}(\epsilon) \right)$ $G^{(n)}(\epsilon) = \frac{1}{T} \int_0^T dt e^{-in\Omega t} \int dt' e^{i\epsilon(t-t')} G(t,t')$ $T_{\lambda,\lambda'}(\epsilon) = \sum_{n} \operatorname{Tr} \left[\Gamma_{\lambda} G^{(n)}(\epsilon) \Gamma_{\lambda'} (G^{(n)}(\epsilon))^{\dagger} \right]$



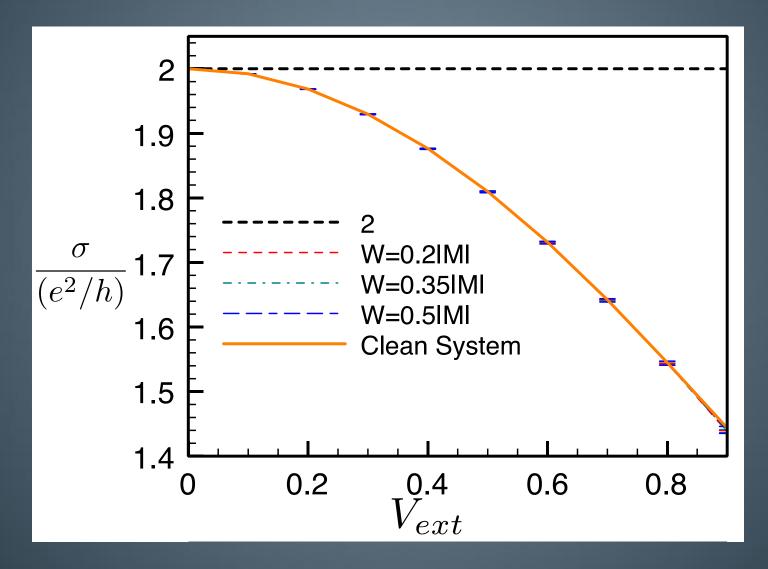
Numerical Calculation



Numerical Calculation



Insensitivity to Disorder



Periodic Topological Insulator

- No quantized conductivity
- Insensitive to disorder and material parameters
- Still "topologically robust"
- Where did the lost conductivity go?

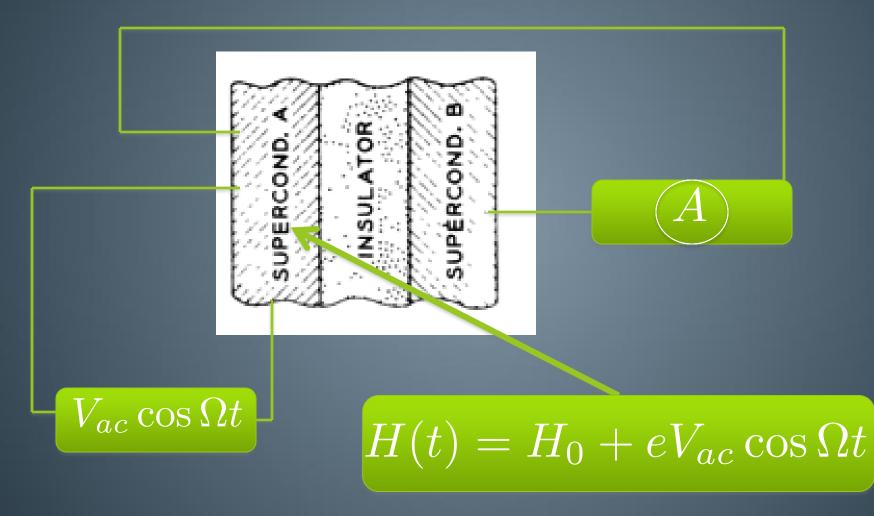
PHYSICAL REVIEW

VOLUME 129, NUMBER 2

15 JANUARY 1963

Multiphoton Process Observed in the Interaction of Microwave Fields with the Tunneling between Superconductor Films

P. K. TIEN AND J. P. GORDON Bell Telephone Laboratories, Murray Hill, New Jersey (Received 28 August 1962)

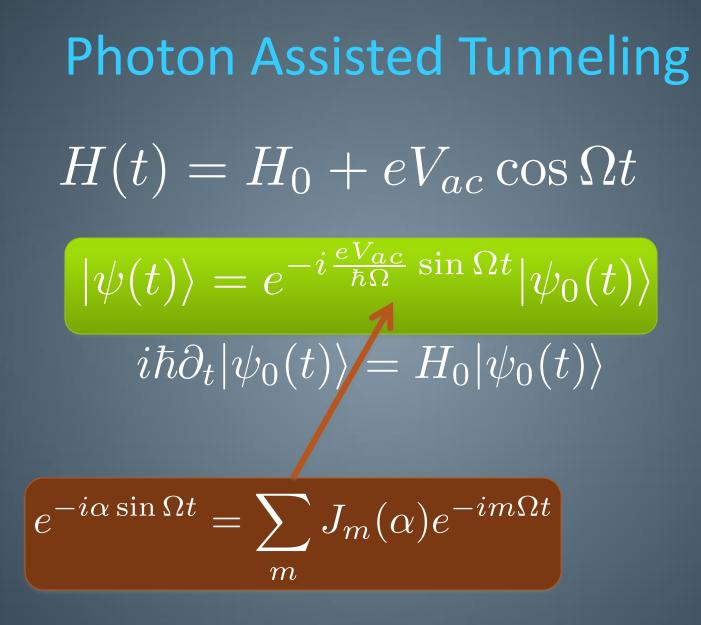


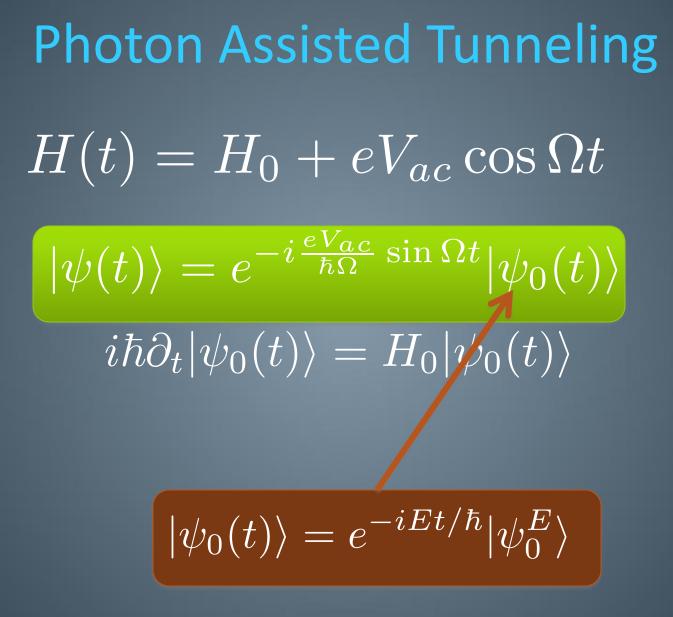
 $H(t) = H_0 + eV_{ac}\cos\Omega t$

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$$|\psi(t)\rangle = e^{-i\frac{eV_{ac}}{\hbar\Omega}\sin\Omega t}|\psi_0(t)\rangle$$

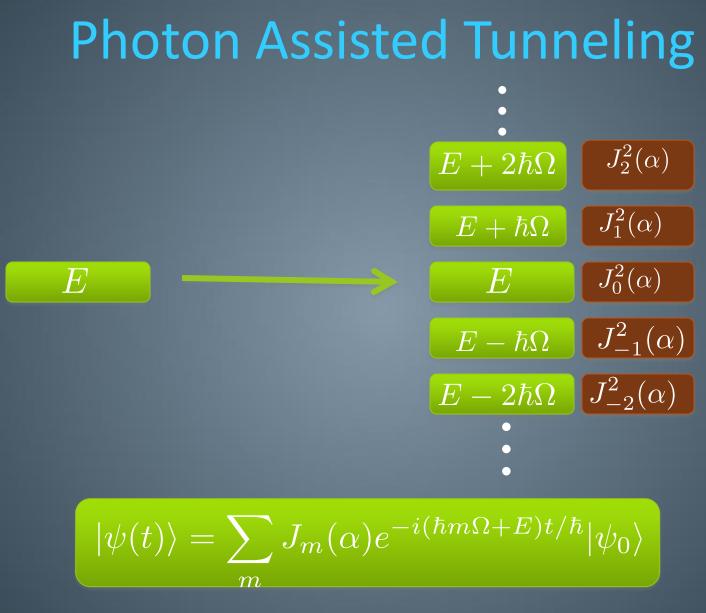
Photon Assisted Tunneling $H(t) = H_0 + eV_{ac}\cos\Omega t$ $|\psi(t)\rangle = e^{-i\frac{eV_{ac}}{\hbar\Omega}\sin\Omega t} |\psi_0(t)\rangle$ $i\hbar\partial_t |\psi_0(t)\rangle = H_0 |\psi_0(t)\rangle$

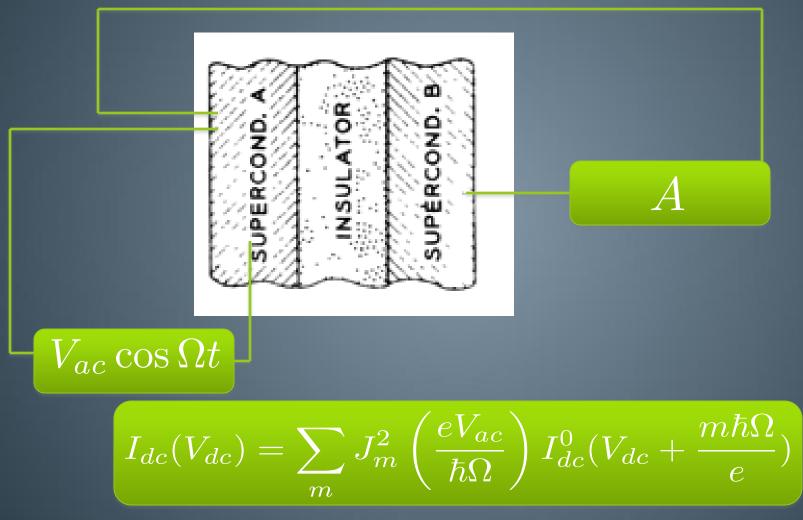




Photon Assisted Tunneling $H(t) = H_0 + eV_{ac} \cos \Omega t$ $|\psi(t)\rangle = e^{-i\frac{eV_{ac}}{\hbar\Omega}\sin\Omega t}|\psi_0(t)\rangle$ $i\hbar\partial_t|\psi_0(t)\rangle = H_0|\psi_0(t)\rangle$

$$|\psi(t)\rangle = \sum_{m} J_m(\alpha) e^{-i(\hbar m \Omega + E)t/\hbar} |\psi_0\rangle$$





$$H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$$

$$H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$$
$$H(\mathbf{k}) = \epsilon(\mathbf{k})I + \vec{d}_{\mathbf{k}} \cdot \vec{\sigma} \quad V(t) = 2V_{ext}\sigma_z \cos \Omega t$$

$$H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$$

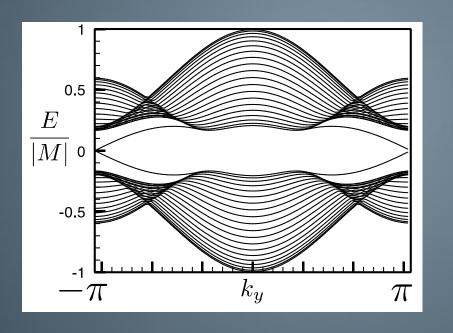
$$|\psi(t)\rangle = e^{-i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}|\tilde{\psi}(t)\rangle$$

• How does this help us....?

$$H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$$

$$|\psi(t)\rangle = e^{-i\frac{\pi}{\hbar\Omega}\sin\Omega t}|\psi(t)\rangle$$

 $\tilde{H}_{\mathbf{k}}(t) = e^{i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}H(\mathbf{k})e^{-i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}$





 $H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$ $|\psi(t)\rangle = e^{-i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}|\tilde{\psi}(t)\rangle$ $\tilde{H}_{\mathbf{k}}(t) = e^{i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}H(\mathbf{k})e^{-i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}$

 $H_{\mathbf{k}}(t) = H(\mathbf{k}) + V(t)$ $|\psi(t)\rangle = e^{-irac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}|\tilde{\psi}(t)\rangle$ $\tilde{H}_{\mathbf{k}}(t) = e^{i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}\sin\Omega t}H(\mathbf{k})e^{-i\frac{\mathbf{V}\cdot\vec{\sigma}}{\hbar\Omega}}\sin\Omega t$ $\tilde{H}_{\mathbf{k}}(t) \to H_{\mathbf{k}}^{eff} = \frac{1}{T} \int_{0}^{T} dt \tilde{H}_{\mathbf{k}}(t)$

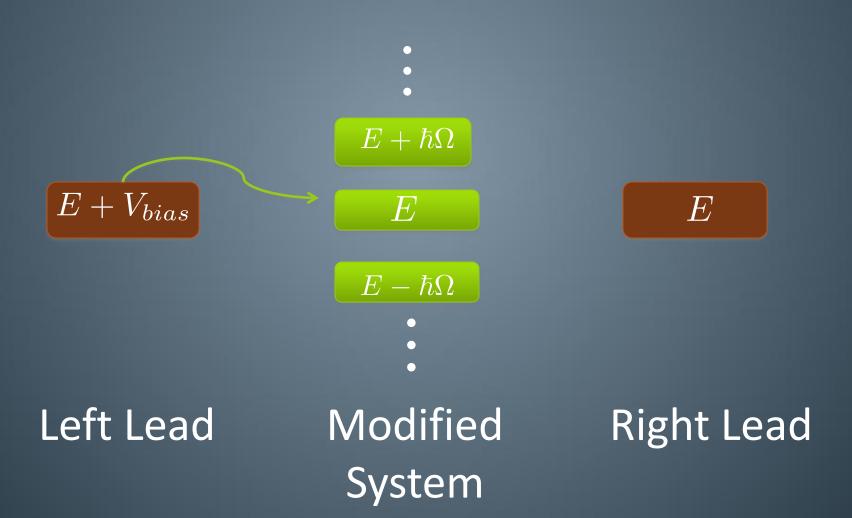
- For off-resonant light the underlying static Hamiltonian is modified
- Modified Hamiltonian is split into sidebands



Left Lead

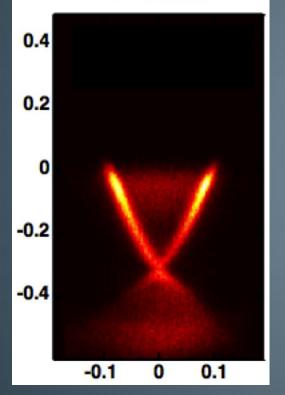
Modified System

Right Lead



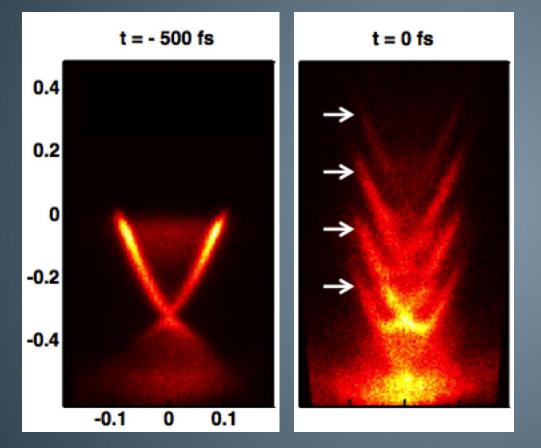
Side bands in time resolved ARPES

t = - 500 fs

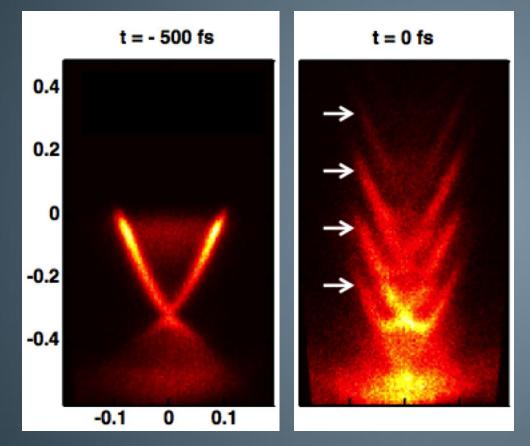


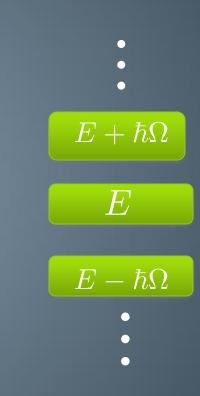
Y.H. Wang et al, Science, 342 (2013)

Side bands in time resolved ARPES



Y.H. Wang et al, Science, 342 (2013)

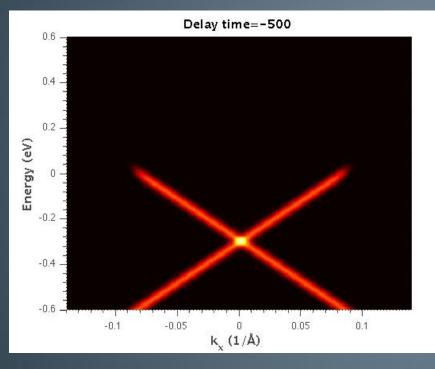


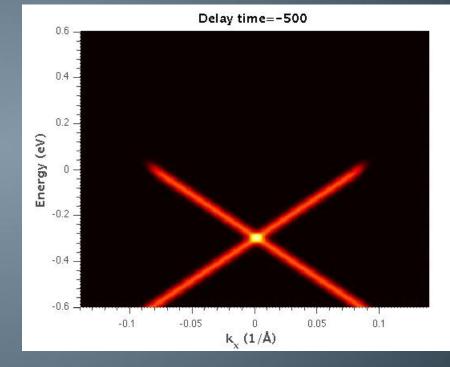


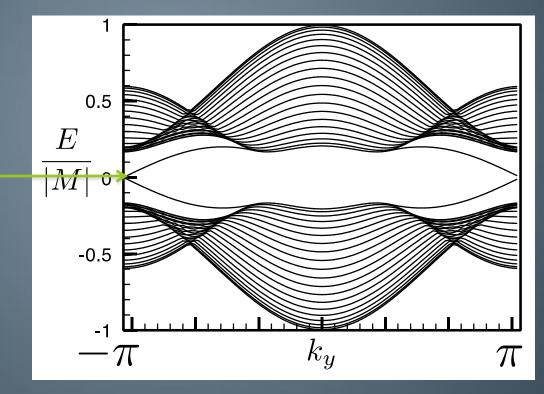
Y.H. Wang et al, Science, 342 (2013)

Aaron Farrell and TPB Work in Progress

Time Resolved ARPES simulation

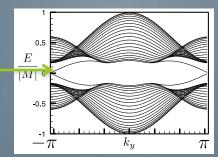






Bands of $H_{\mathbf{k}}^{eff}$

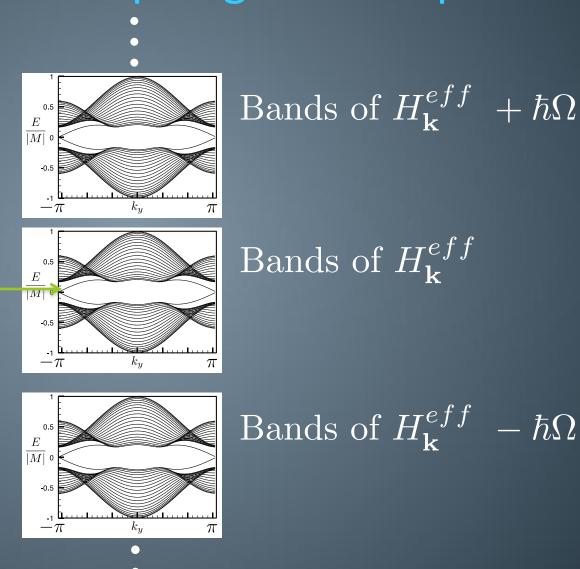
e



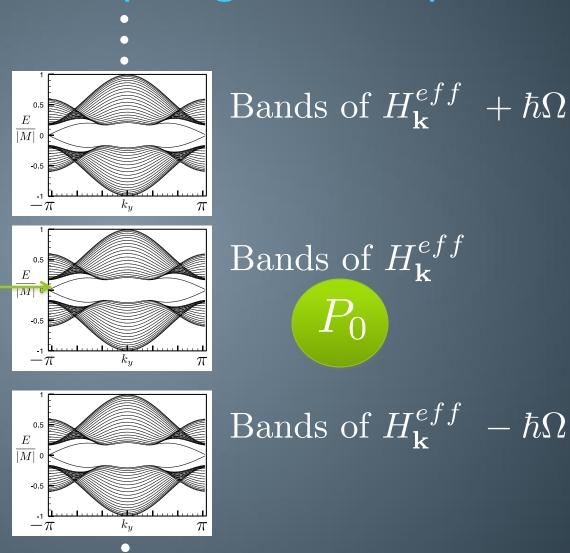
e

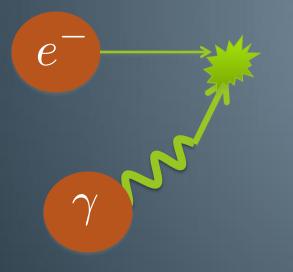


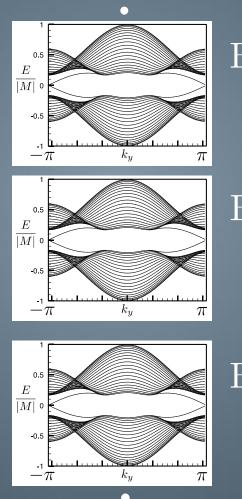
 e^{-}



 e^{-}



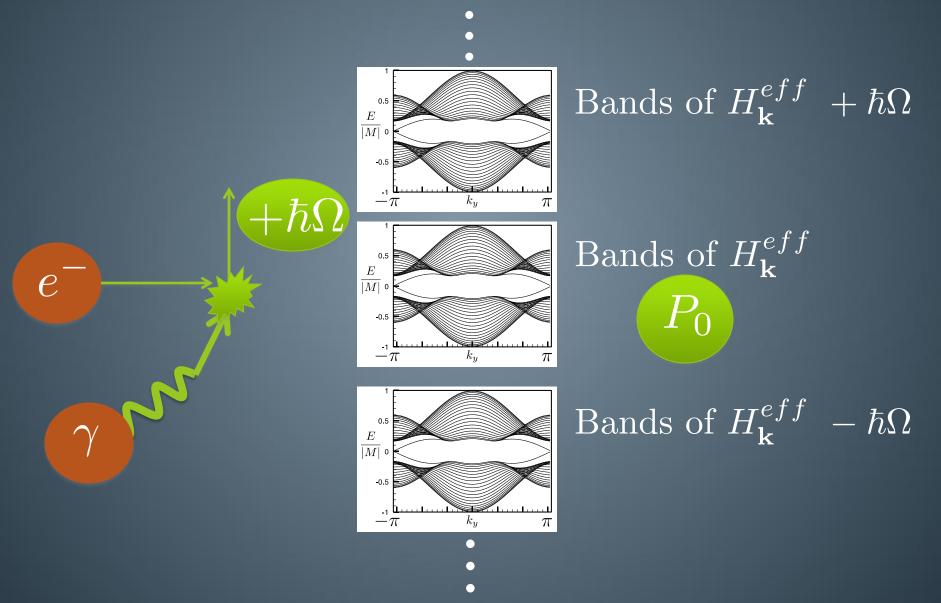




Bands of
$$H_{\mathbf{k}}^{eff} + \hbar \Omega$$



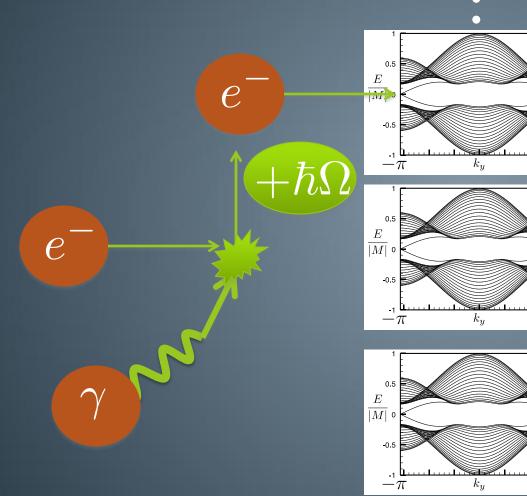
Bands of
$$H_{\mathbf{k}}^{eff} - \hbar \Omega$$



 π

 π

 π

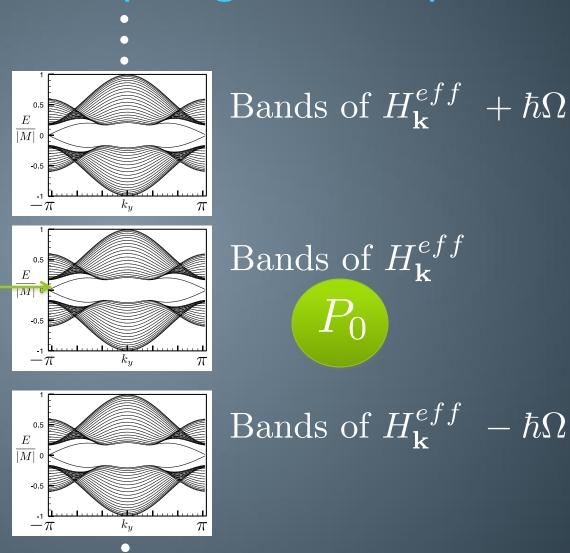


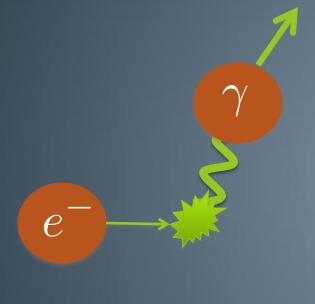
Bands of
$$H^{eff}_{\mathbf{k}} + \hbar \Omega$$

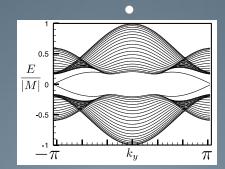


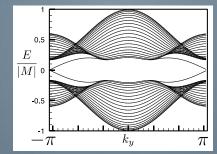
Bands of
$$H_{\mathbf{k}}^{eff} - \hbar \Omega$$

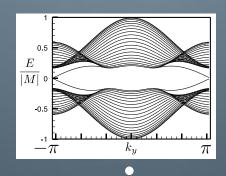
 e^{-}







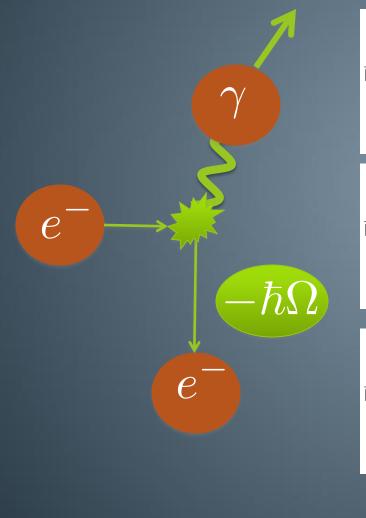


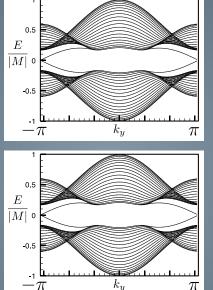


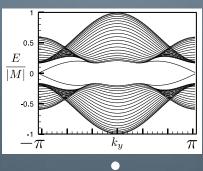
Bands of
$$H^{eff}_{\mathbf{k}} + \hbar\Omega$$



Bands of
$$H_{\mathbf{k}}^{eff} - \hbar \Omega$$







Bands of
$$H^{eff}_{\mathbf{k}} + \hbar \Omega$$



Bands of
$$H_{\mathbf{k}}^{eff} - \hbar \Omega$$

 k_u

 k_u

 π

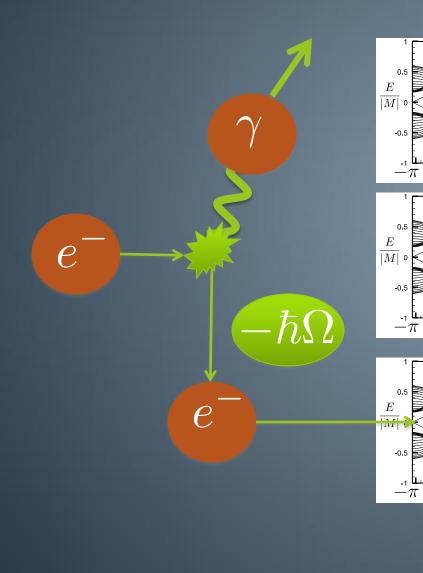
 π

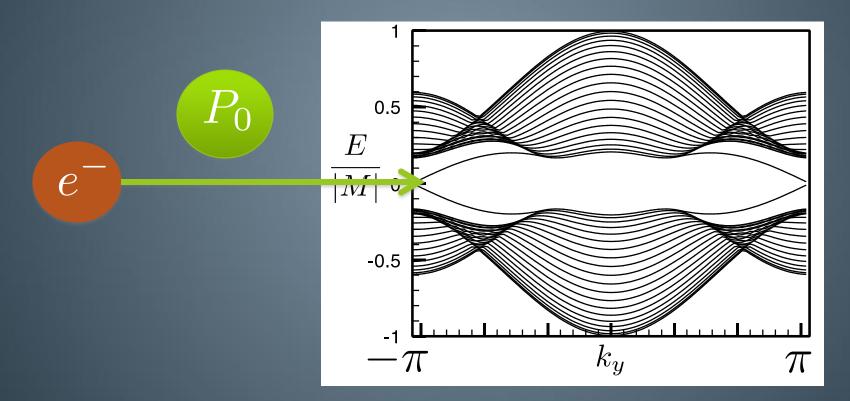
Bands of $H^{eff}_{\mathbf{k}} + \hbar \Omega$

Bands of $H^{eff}_{\mathbf{k}} - \hbar \Omega$

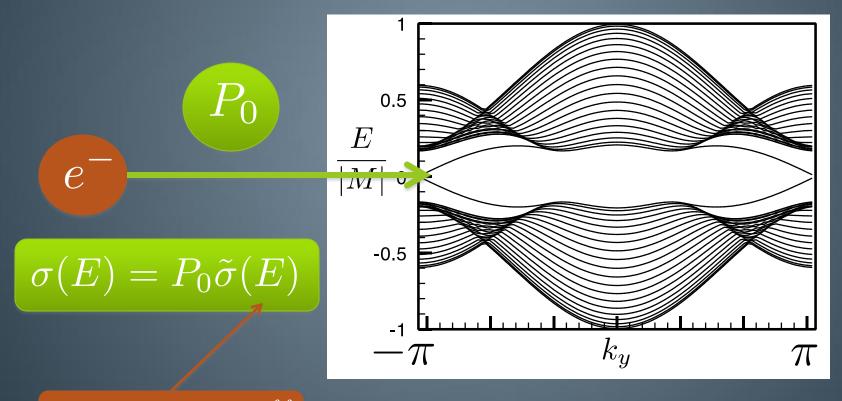
Bands of $H^{eff}_{\mathbf{k}}$

 P_0



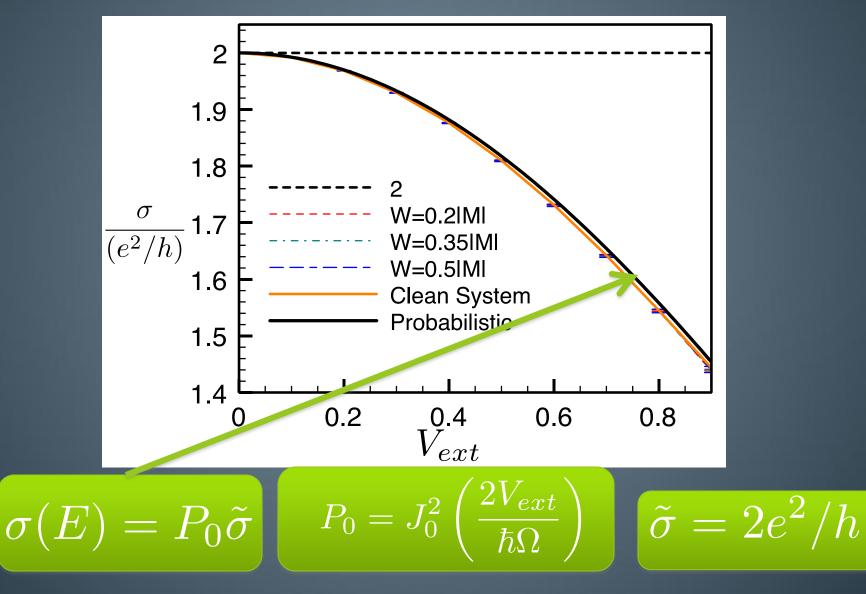


Bands of $H_{\mathbf{k}}^{eff}$



Conductivity of $\underline{H}_{\mathbf{k}}^{eff}$

Bands of $H_{\mathbf{k}}^{eff}$

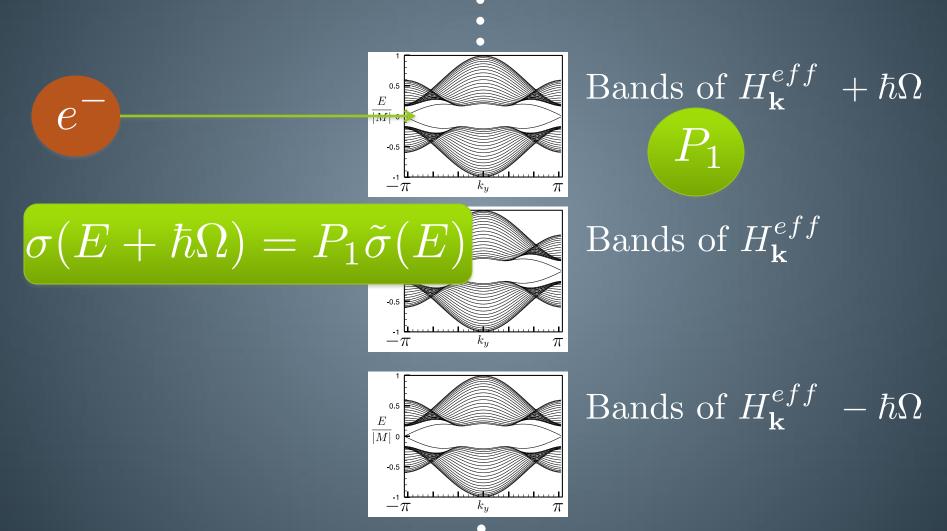


 $\sigma_F(E) = \sum \sigma(E + n\hbar\Omega)$ n

Kundu and Seradjeh, PRL 111 (2013)

Bands of $H^{eff}_{\mathbf{k}} + \hbar \Omega$ 0.5 EeM -0.5 k_{n} Bands of $H^{eff}_{\mathbf{k}}$ 0.5 E $\overline{|M|}$ o -0.5 $-\pi$ k_u π Bands of $H_{\mathbf{k}}^{eff} - \hbar \Omega$ 0.5 E $\overline{|M|}$ o -0.5 -1 $-\pi$ k_u π

Bands of $H^{eff}_{\mathbf{k}} + \hbar\Omega$ 0.5 EeM P_1 -0.5 Bands of $H^{eff}_{\mathbf{k}}$ 0.5 E $\overline{|M|}$ o -0.5 $-\pi$ k_u π Bands of $H^{eff}_{\mathbf{k}} - \hbar \Omega$ 0.5 E $\overline{|M|}$ o -0.5 -1 $-\pi$ k_u π



 $\sigma(E + n\hbar\Omega) = P_n \tilde{\sigma}(E)$

 $\sigma(E + n\hbar\Omega) = P_n\tilde{\sigma}(E)$

 $\sigma_F(E) = \sum \sigma(E + n\hbar\Omega)$ n

 $\sigma(E + n\hbar\Omega) = P_n\tilde{\sigma}(E)$



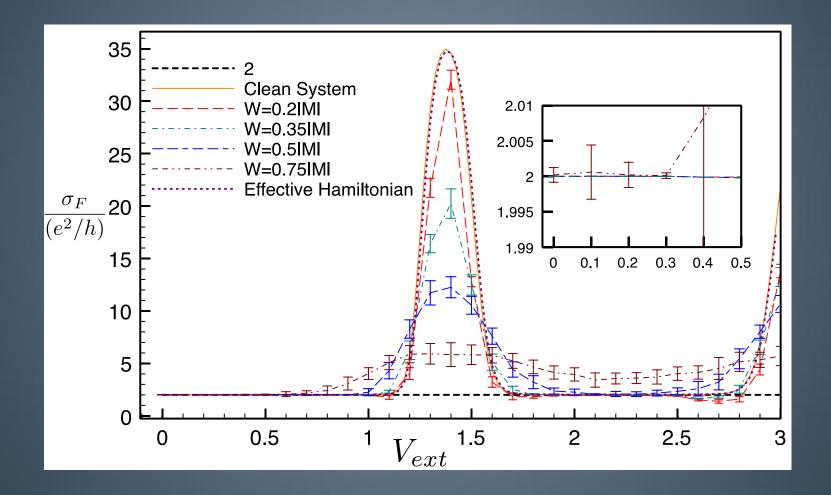
 $\sigma_F(E) = \tilde{\sigma}(E) \sum P_n = \tilde{\sigma}(E)$ n

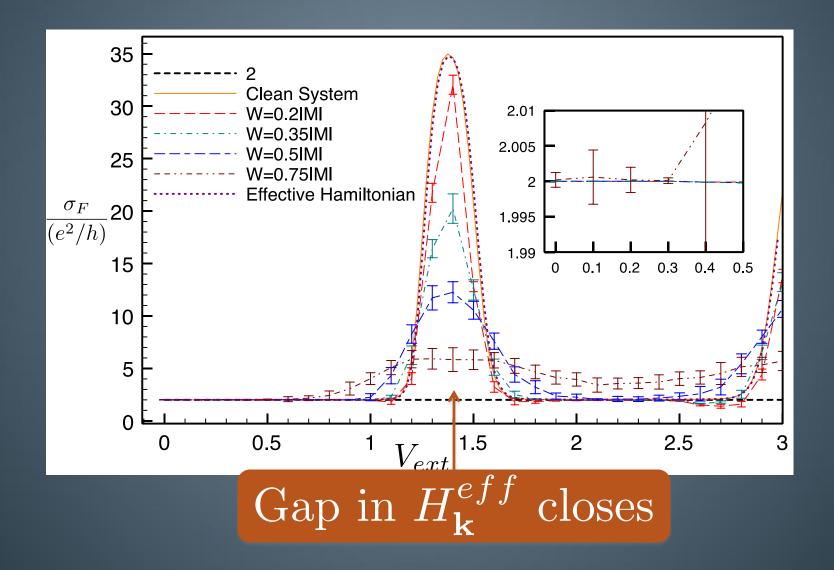
 $\sigma(\overline{E+n}\overline{h}\Omega) = P_n\tilde{\sigma}(E)$

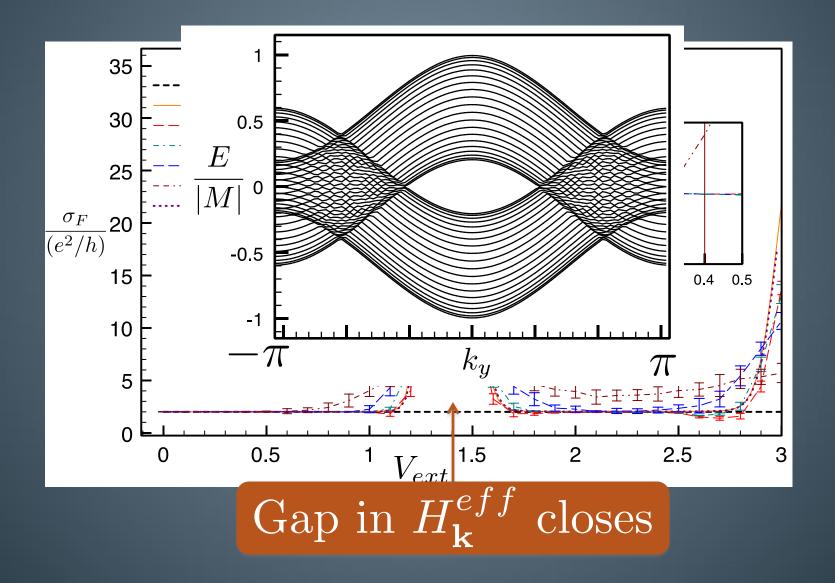
 $\sigma_F(E) = \sum \sigma(E + n\hbar\Omega)$ \mathcal{N}

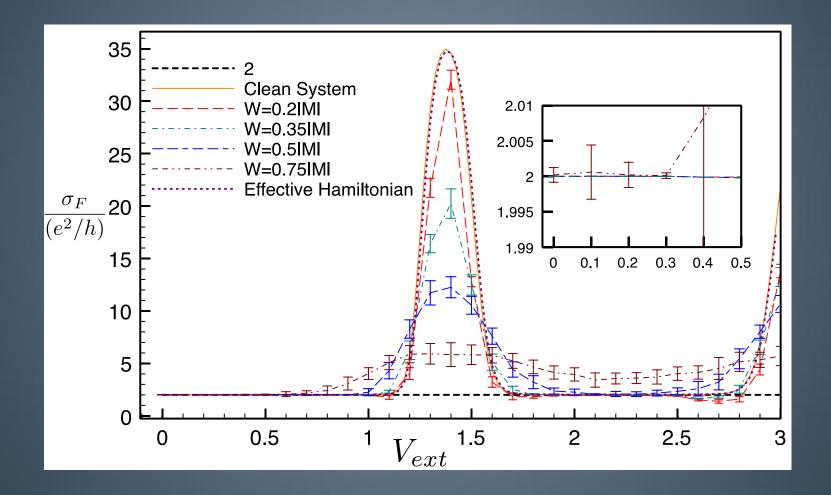
 $\sigma_F(E) = 2e^2/h$

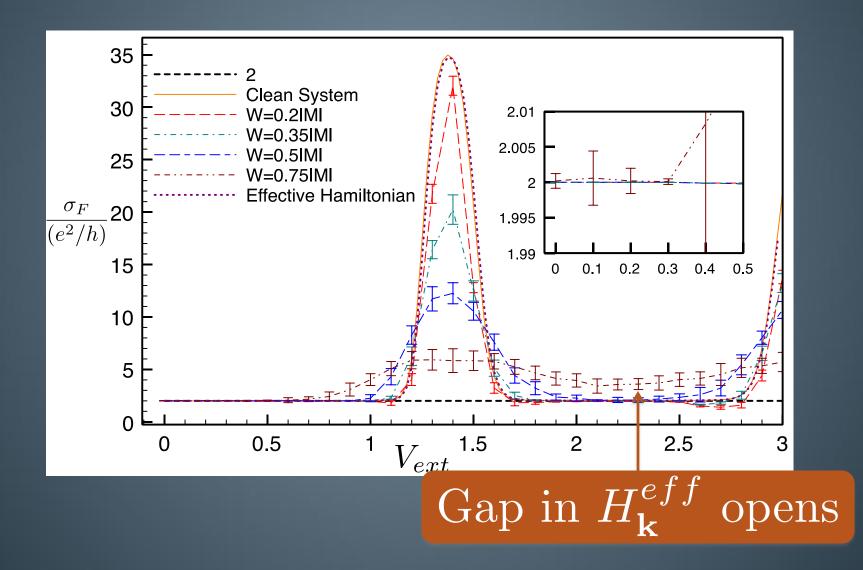
When edge states are present at energy *E* in effective Hamiltonian

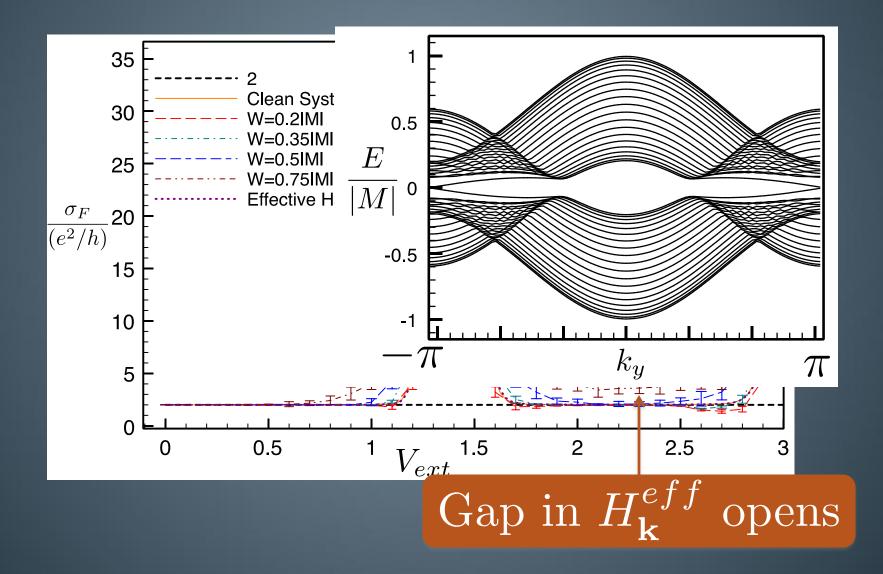






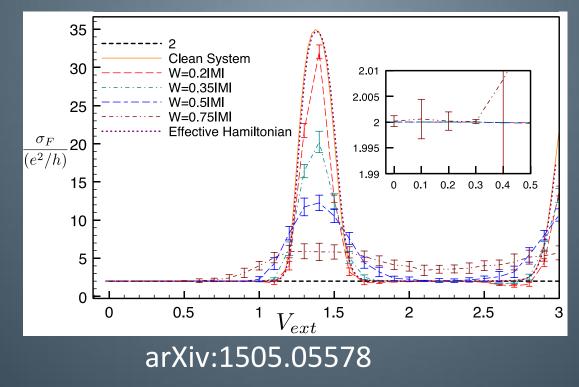


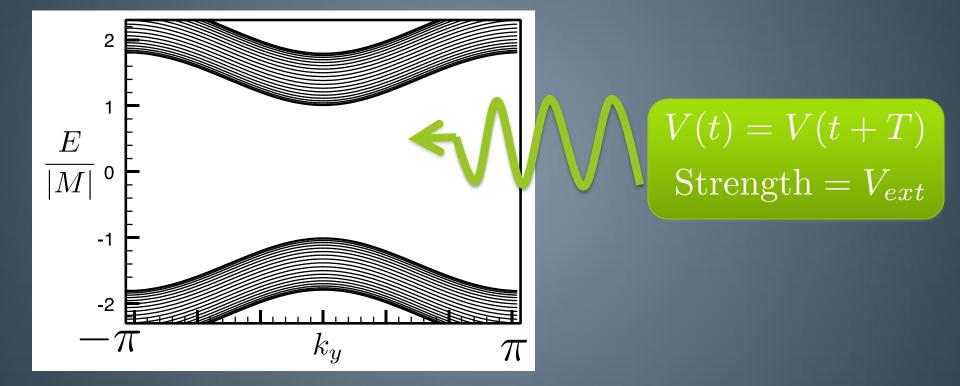


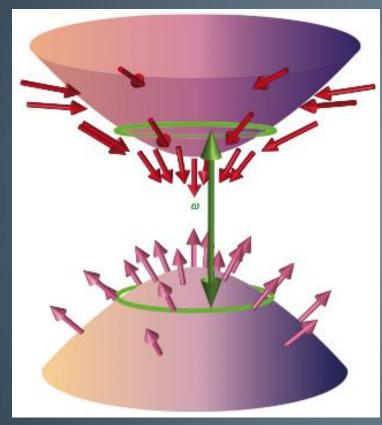


Periodic Topological Insulator

- Conductivity lost due to photon scattering
- Conductivity is still robust
- Lost conductivity reclaimed by sum rule

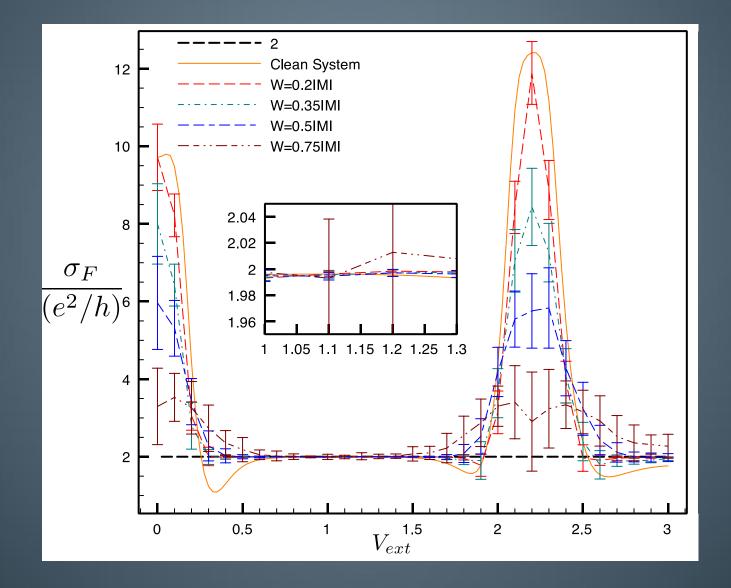


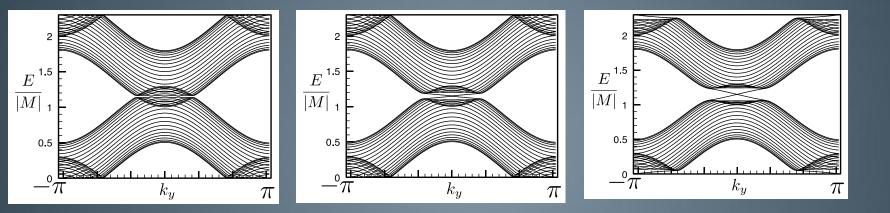




Lindner et al, Nature Physics 7 (2011)

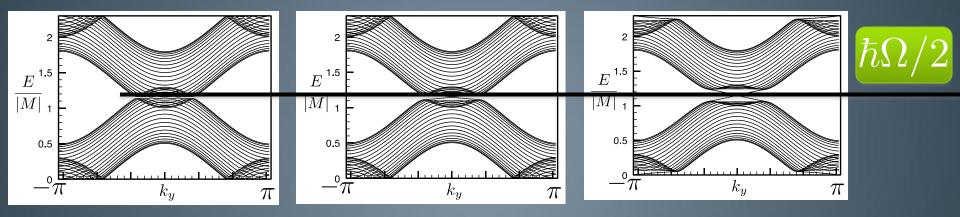
- "On-resonant" light induces edge states
- Effective Hamiltonian description now not possible
- Edge states still occupied probabilistically
- Begin with the sum rule





Increasing V_{ext}

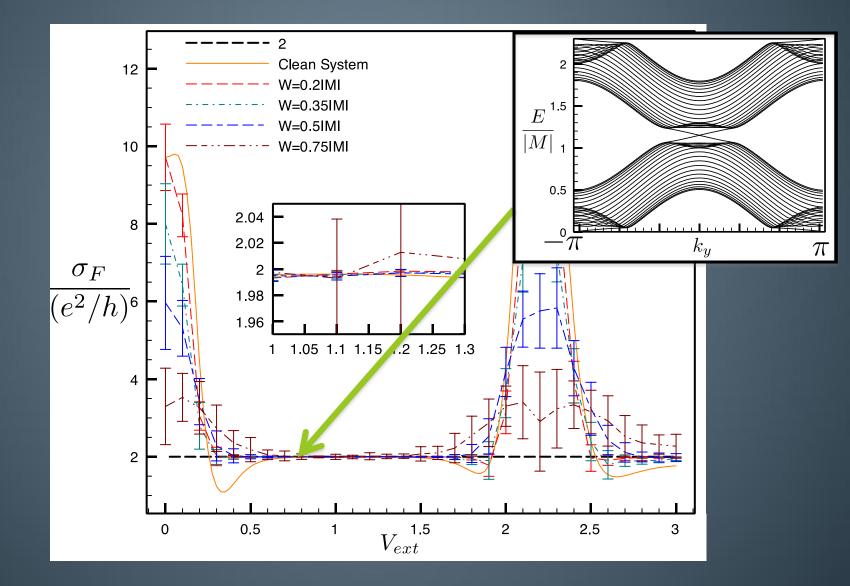
How do these states effect transport?



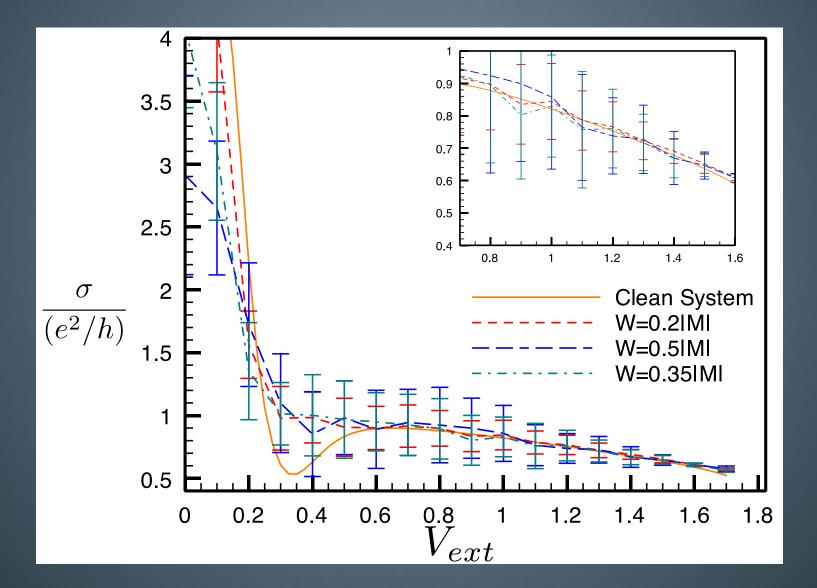
Increasing V_{ext}

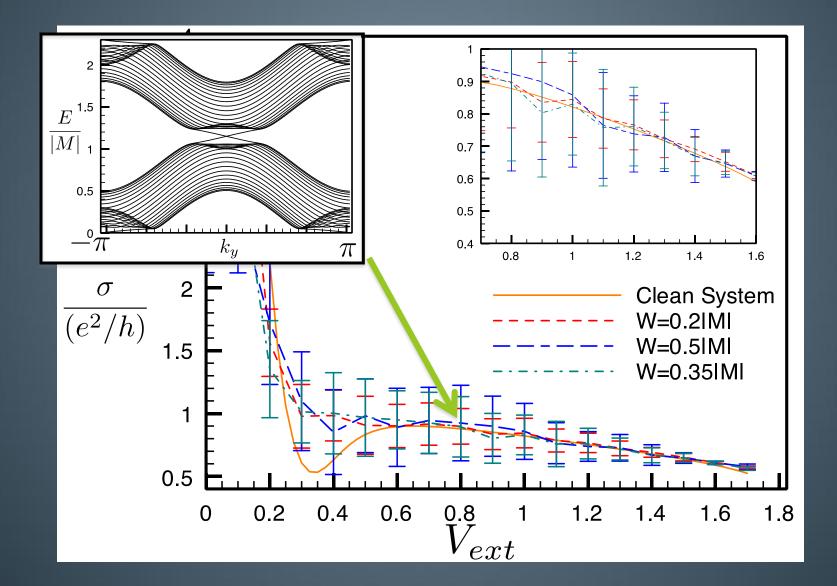
How do these states effect transport?

Summed Conductivity



Conductivity (not summed)



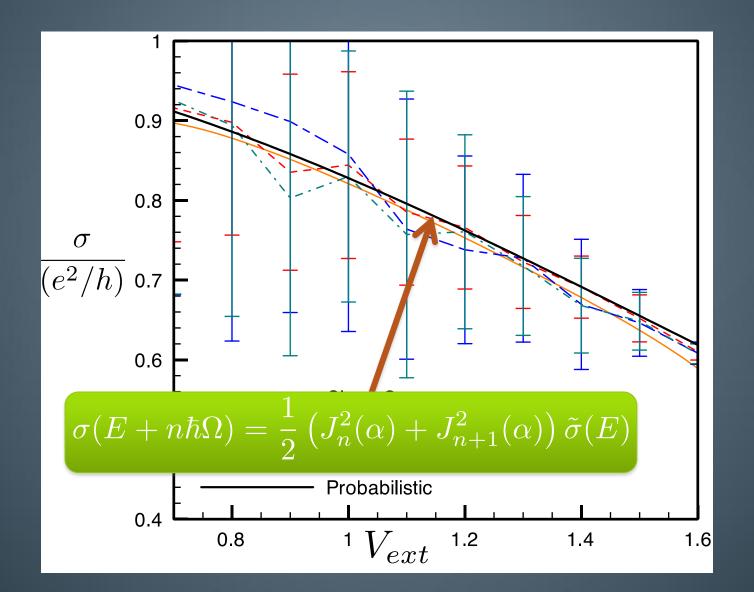


$\sigma(E + n\hbar\Omega) = \frac{1}{2} \left(J_n^2(\alpha) + J_{n+1}^2(\alpha) \right) \tilde{\sigma}(E)$

Probability of accessing edge states

$\sigma(E + n\hbar\Omega) = \frac{1}{2} \left(J_n^2(\alpha) + J_{n+1}^2(\alpha) \right) \tilde{\sigma}(E)$

Conductivity of edge states (complicated)

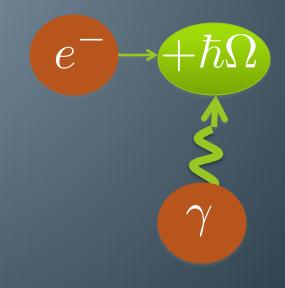


- Edge states induced using "on-resonant" light
- Quantization of Floquet sum rule, presence of edge states and robustness coexist
- Edge states are accessed with a certain probability

arXiv:1505.05584

Summary

- Periodic driving reduces transport signatures of topological edge states, maintains robustness
- Can recollect these lost signatures using a "Floquet sum rule"



arXiv:1505.05578 arXiv:1505.05584

Outlook

- Effects of additional bands
- Transport signatures at small frequency
- Signatures in the AC conductivity

