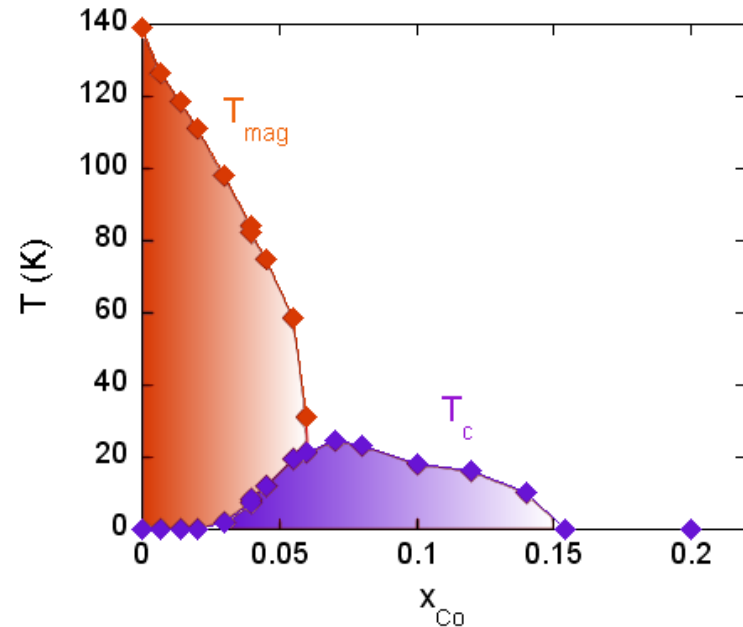
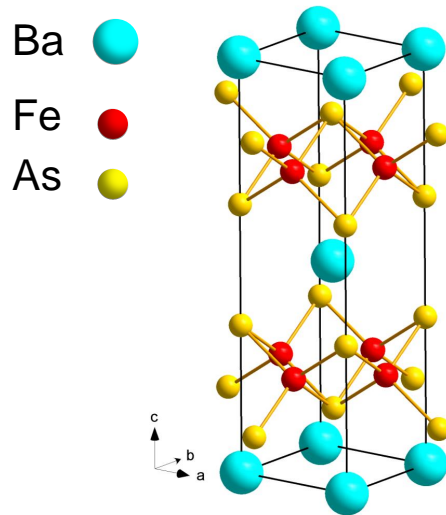


# Respective roles of electrons and holes in the transport properties of pnictides

F. Rullier-Albenque<sup>1</sup>, D. Colson<sup>1</sup>, A. Forget<sup>1</sup> and H. Alloul<sup>2</sup>

<sup>1</sup>SPEC, Orme des Merisiers, CE Saclay

<sup>2</sup>Laboratoire de Physique des Solides, Université d'Orsay



F. Rullier-Albenque et al. PRL 103 (2009)

V. Brouet et al., PRB 80 (2009)

Y. Laplace et al., PRB 80 (2009)

EPJB (2009)

B. Mansart et al. PRB 80 (2009)

L. Chauvière et al., PRB 80 (2009)

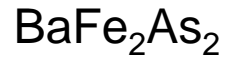
P. Bonville et al., arxiv 1002.0931

# Outline

- The “122” phase:  $BaFe_2As_2$  : *electron doped*  
 $Ba(Fe_{1-x}Co_x)_2As_2$
- Analysis of the transport properties of  
 $Ba(Fe_{1-x}Co_x)_2As_2$   
*Resistivity and Hall effect*  
*Carrier content, scattering processes and*  
*magnetic transitions*
- Isovalent substitution :  $Ba(Fe_{1-x}Ru_x)_2As_2$ :  
*Contribution of holes and electrons to the transport*  
*properties*
- Summary

# 122 Phase : $\text{BaFe}_2\text{As}_2$

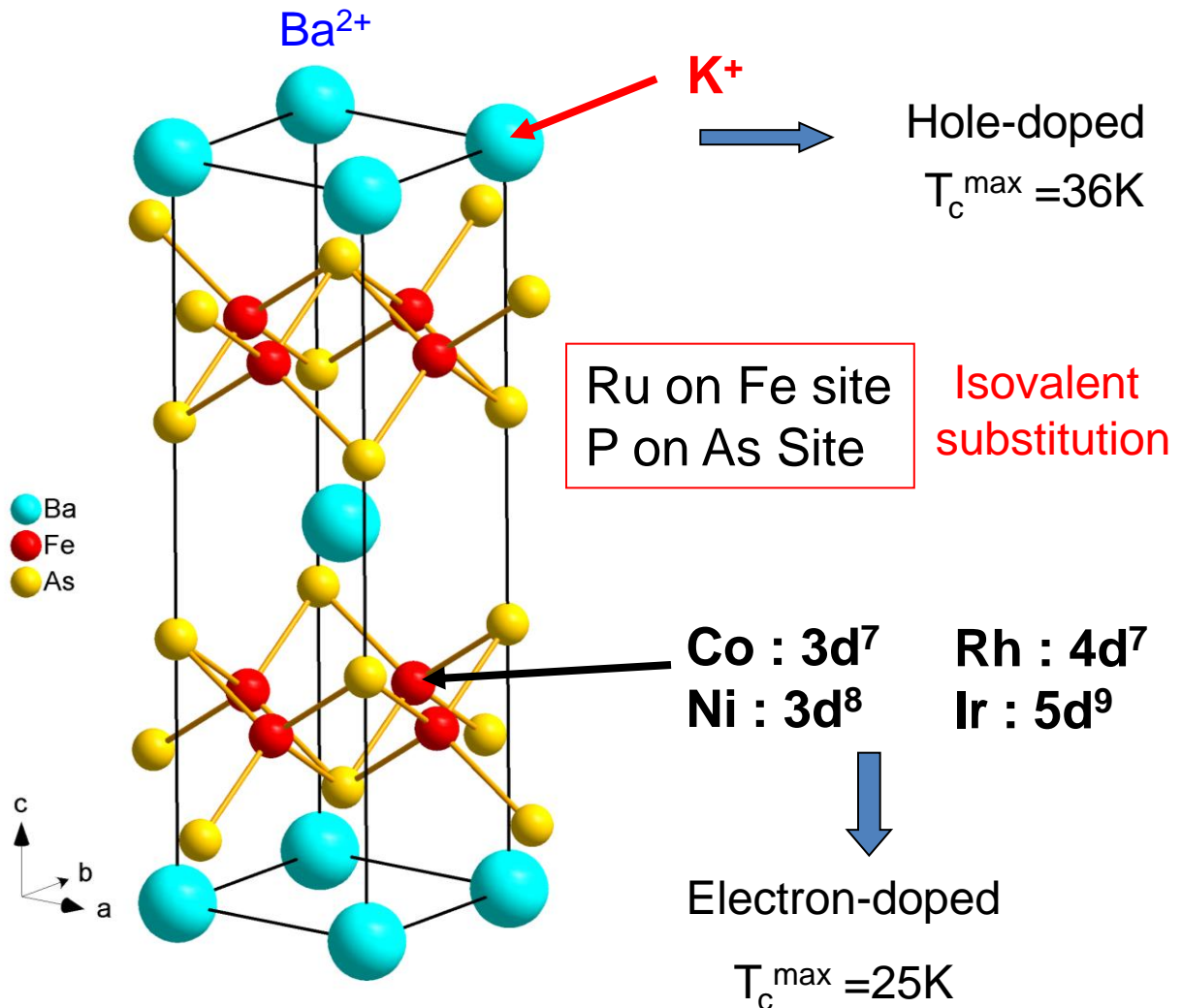
Possibility to get large single crystals



Tetragonal ( $I4/m\ m\ m$ )  
 $a = 3,9625(1) \text{ \AA}$   
 $c = 13,0168(3) \text{ \AA}$

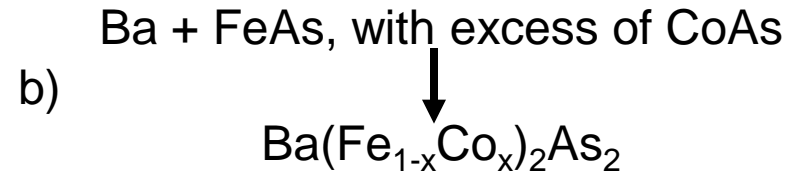
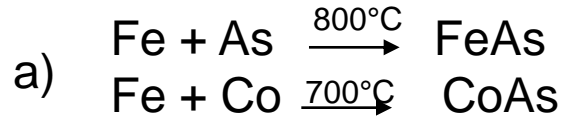
Ba-As:  $3,38 \text{ \AA}$   
Fe-As:  $2,40 \text{ \AA}$   
Fe-Fe:  $2,80 \text{ \AA}$   
As-Fe-As:  $111,1$  et  $108,7^\circ$

*Rotter et al., PRL 2008*



# 122 Phase $\text{BaFe}_2\text{As}_2$ : crystal growth

## Self flux method

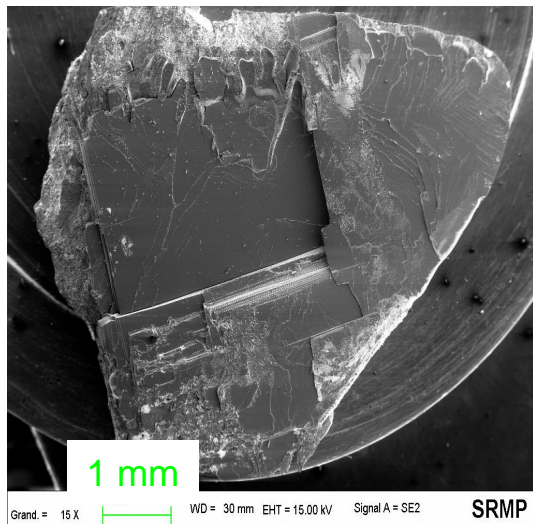


Weighing and mixing in a glove box under Ar atmosphere : ~2 g powder,  $\text{Al}_2\text{O}_3$  crucibles

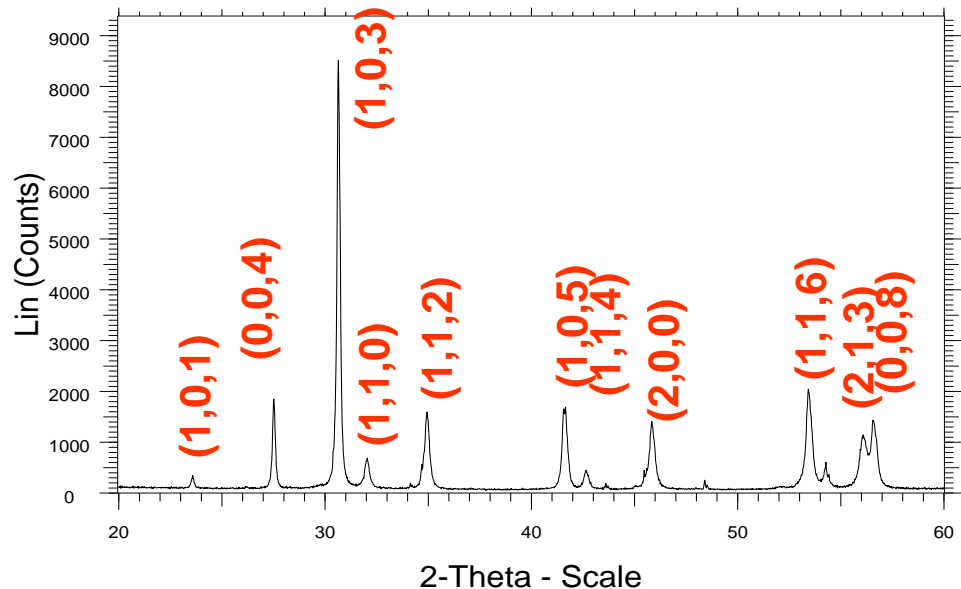
Synthesis in quartz tubes, sealed under vacuum

Kept at  $1180^\circ\text{C}$  for 4h and cooling at  $5^\circ\text{C/h}$   $\longrightarrow$   $1000^\circ\text{C}$  for 6h  
Cooled down to room T

Mechanical extraction of  
platelets single crystals,



RX diagram



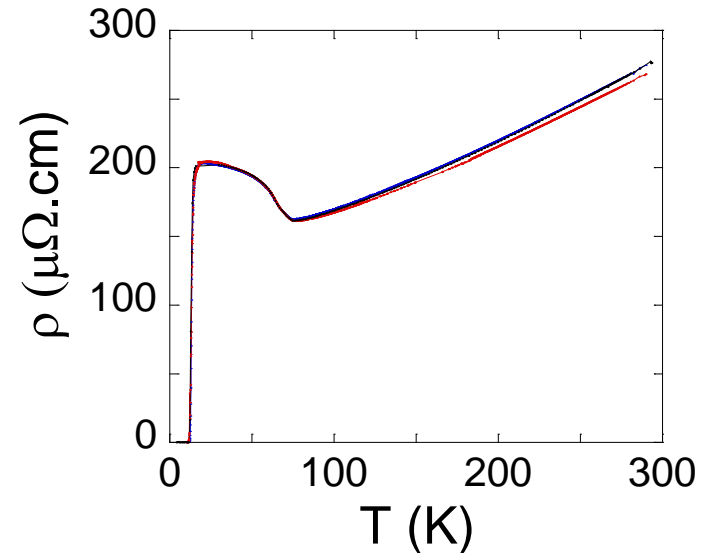
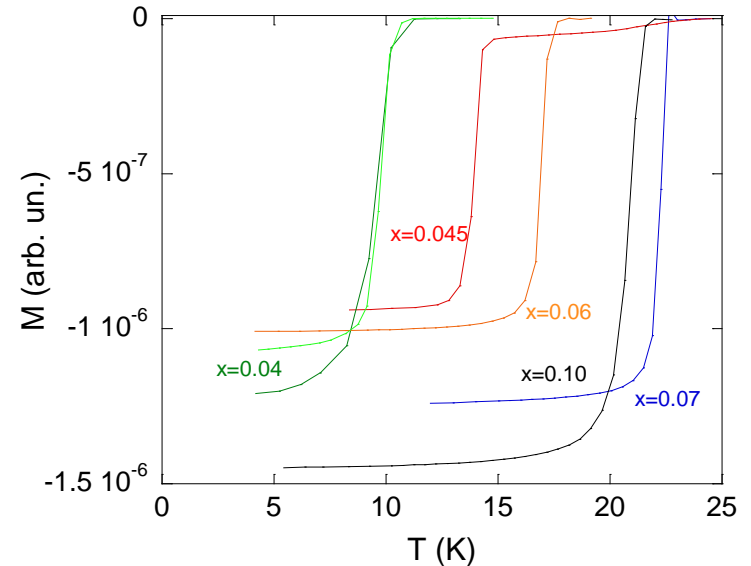
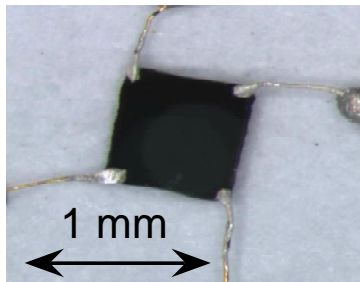
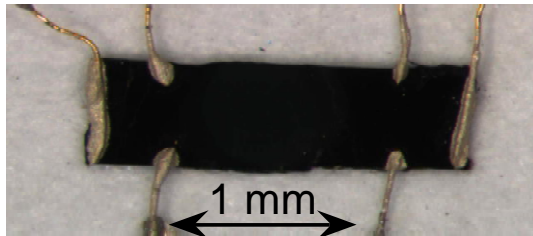
# Co-doped: $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$

Co content determined by wave length dispersive X-ray spectroscopy

$0 \leq x \leq 0.3$  With a precision of 10%

Resistivity and Hall effects measurements

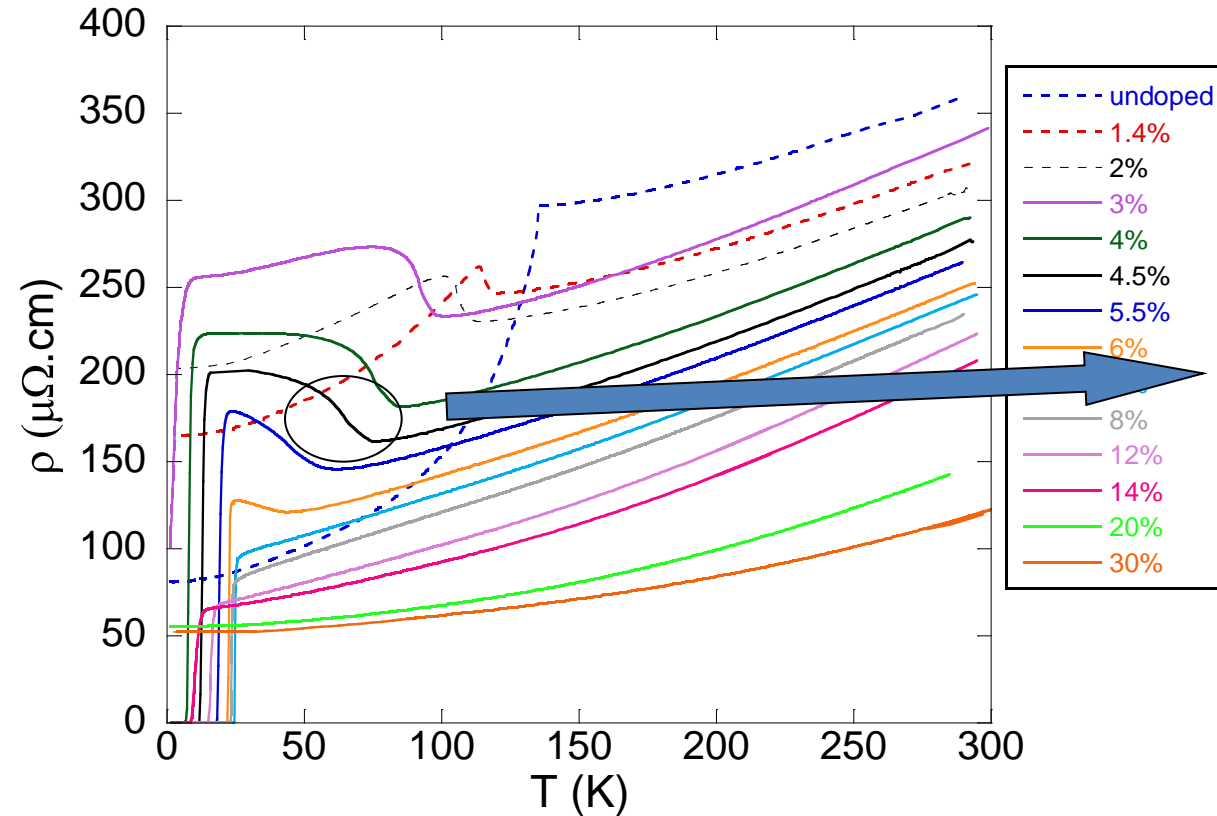
Thickness between 10 to 30  $\mu\text{m}$



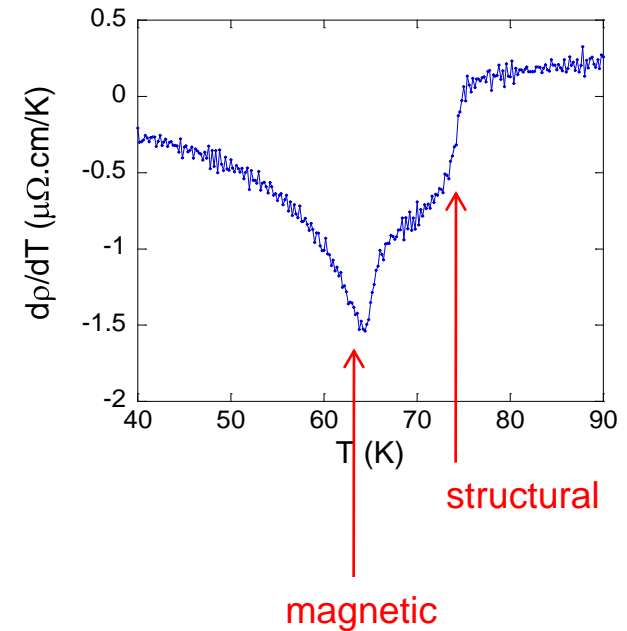
Good reproducibility of the measurements  
Absolute values of the resistivity

# Resistivity curves all over the phase diagram

Co-doped:  $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$



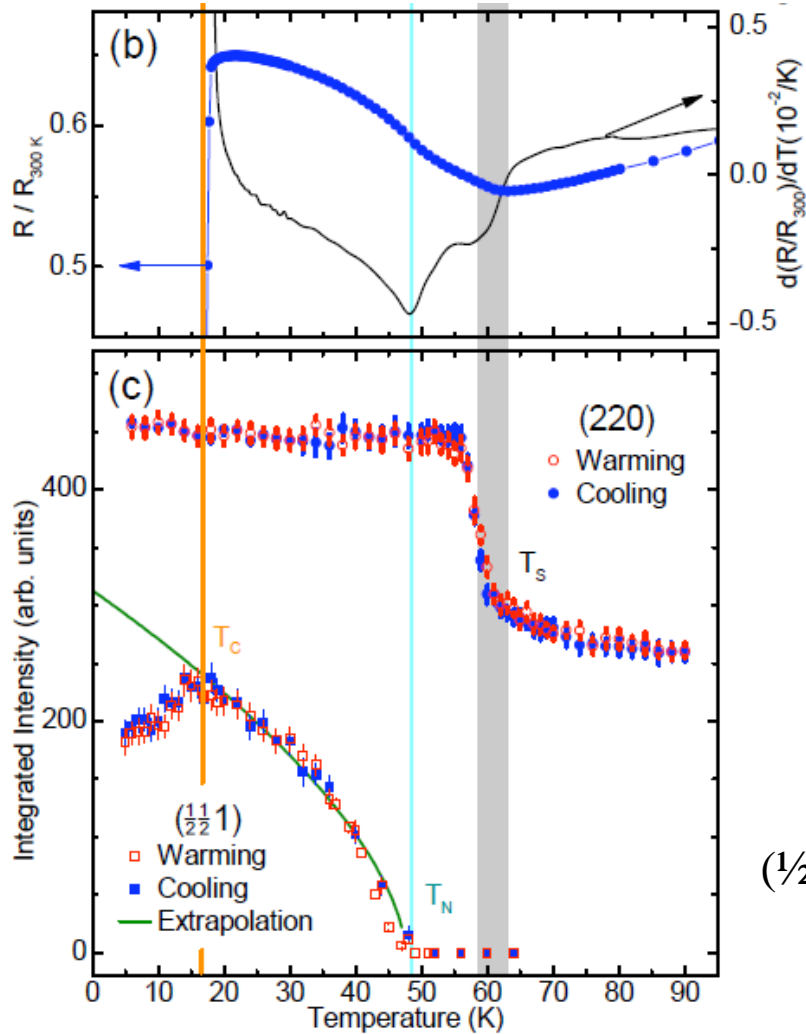
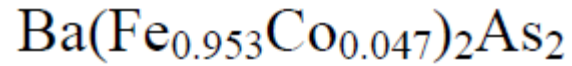
## Structural and magnetic transitions



Very homogeneous samples

See also Y. Laplace et al., EJPB (2009)

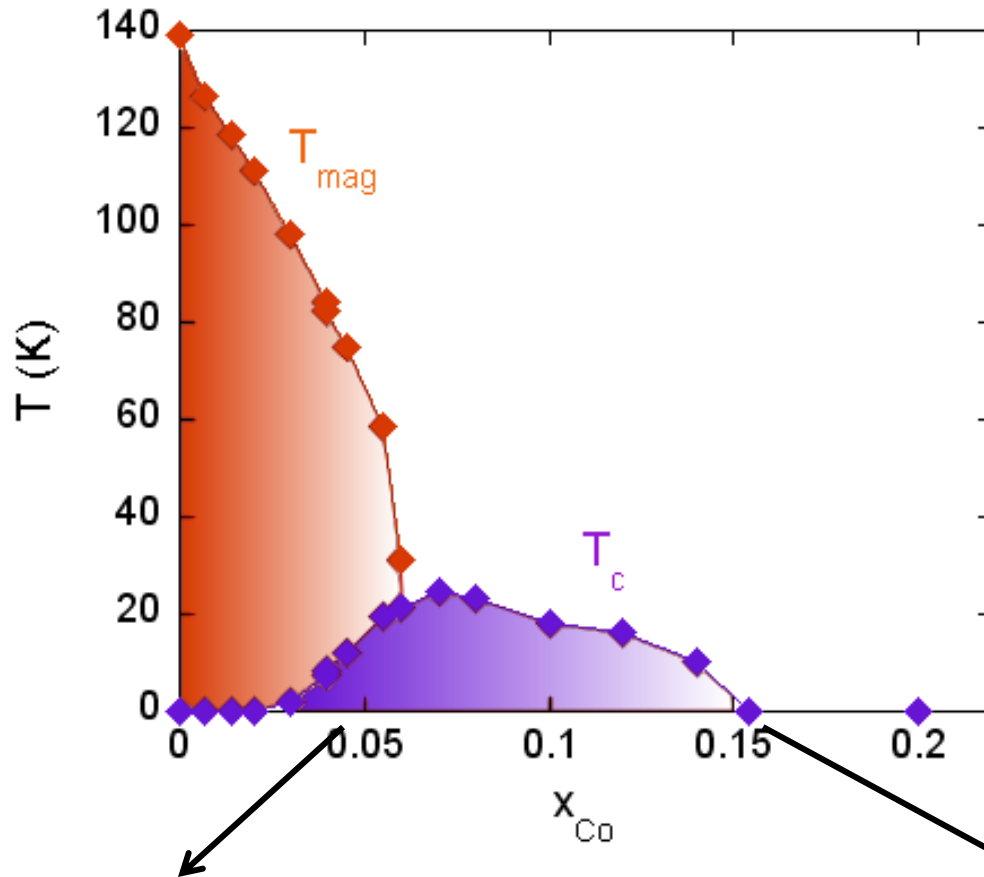
# Structural and magnetic transitions



(220) nuclear reflection

$(\frac{1}{2}, \frac{1}{2}, 1)$  magnetic reflection

# Phase diagram of $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$



$3\% < x < 7\%$  coexistence magnetism and superconductivity

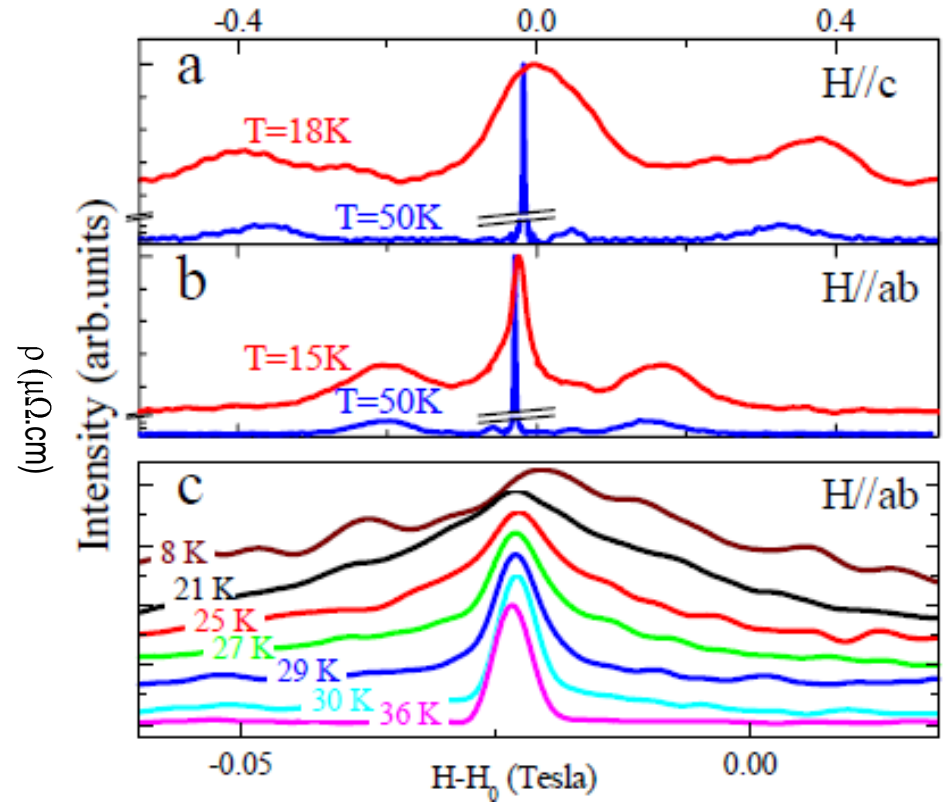
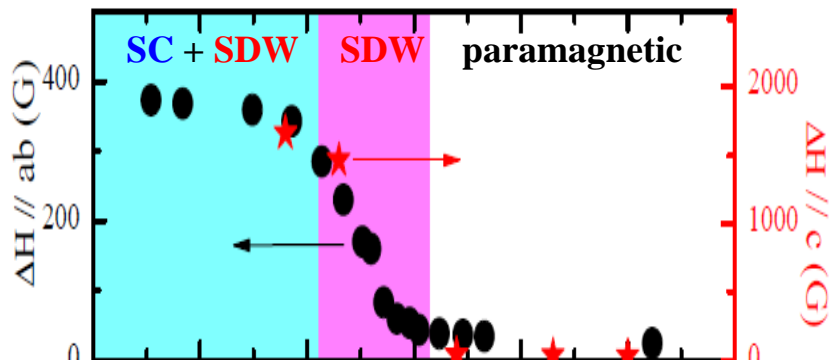
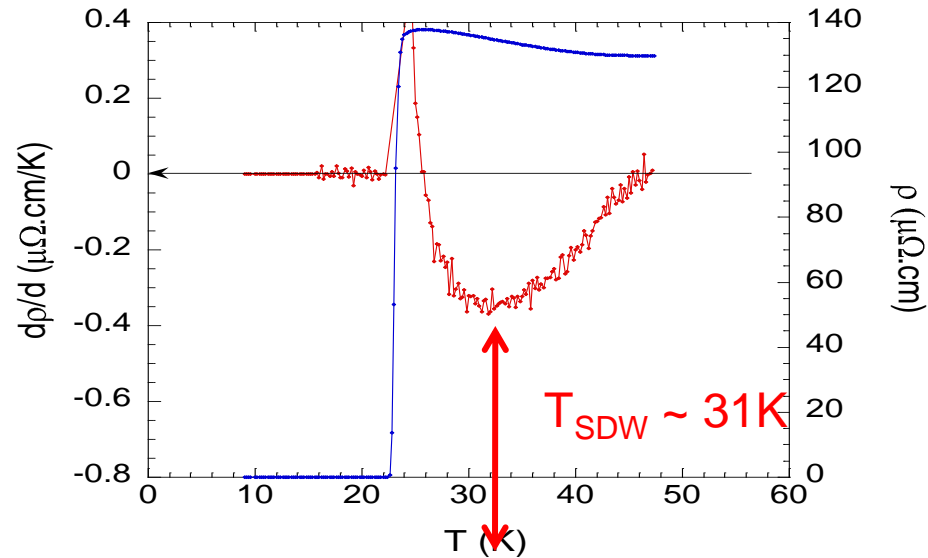
$T_c = 0$  for  $x > 15.5\%$



# Coexistence of superconductivity and magnetism at the atomic scale

## NMR on single crystal with 6%Co

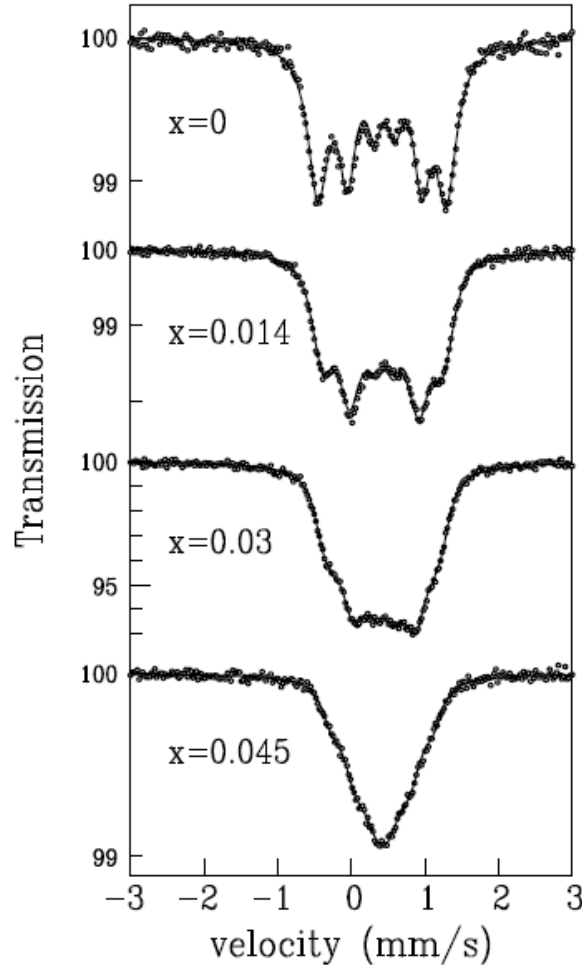
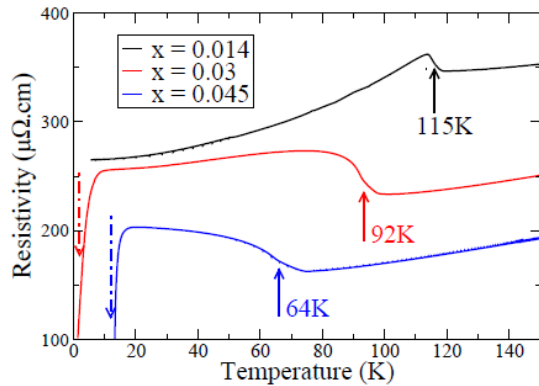
Y. Laplace et al. PRB 80 (2009)



Incommensurate magnetic order  
Very small deviation  $\varepsilon \sim 4\%$

# Mössbauer spectroscopy - Incommensurate magnetic order

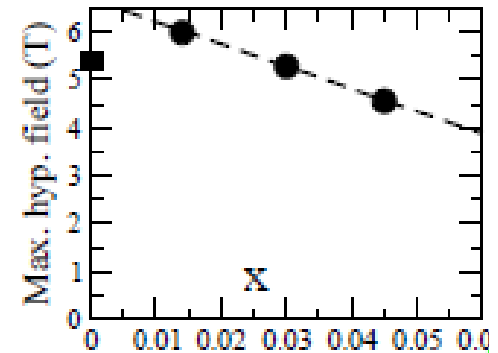
B. Bonville et al., arXiv 1002.0931



Single magnetic hyperfine pattern

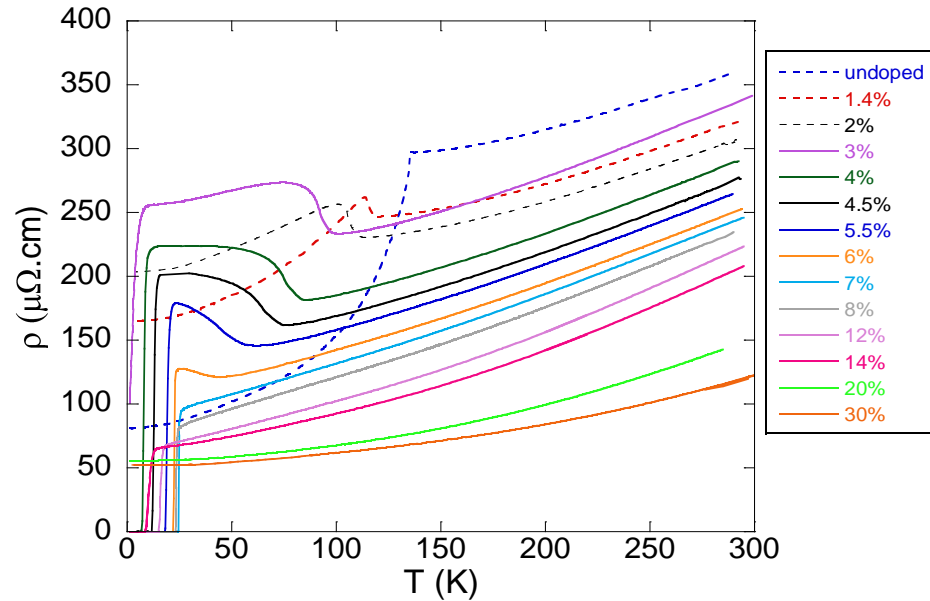
6 lines  
 $H_{\text{hf}} \sim 5.4\text{T}$

Distribution of hyperfine fields  
Incommensurate modulated  
magnetic structure



This incommensurability has not been seen by neutron diffraction up to now.

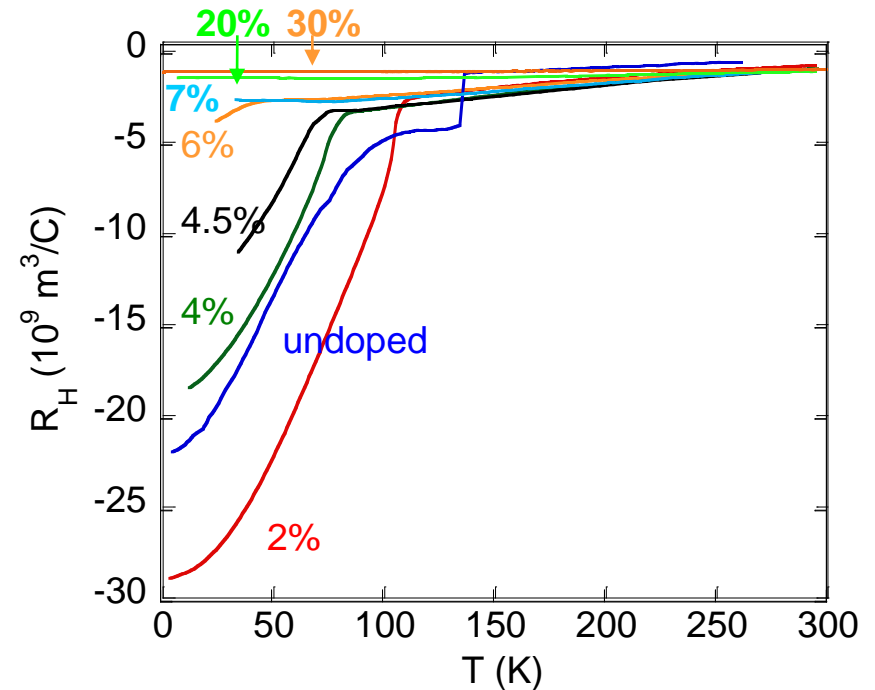
# Resistivity and Hall effect all over the phase diagram



$$\rho = \frac{m^*}{ne^2\tau}$$

Multiband system:

$$\sigma = \sigma_e + \sigma_h$$



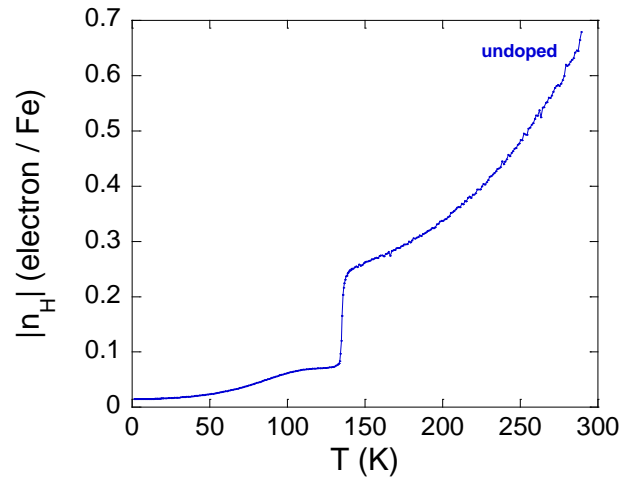
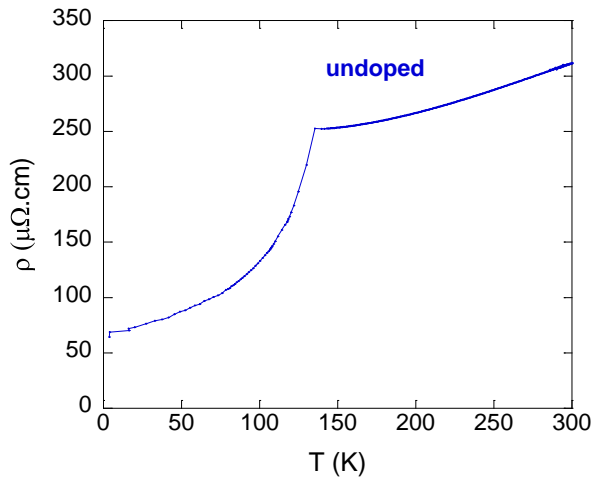
Drop in the absolute value of  $R_H$  at the magnetic transition.

$R_H$  always negative



Transport dominated by electrons all over the phase diagram

# Undoped compound $\text{BaFe}_2\text{As}_2$

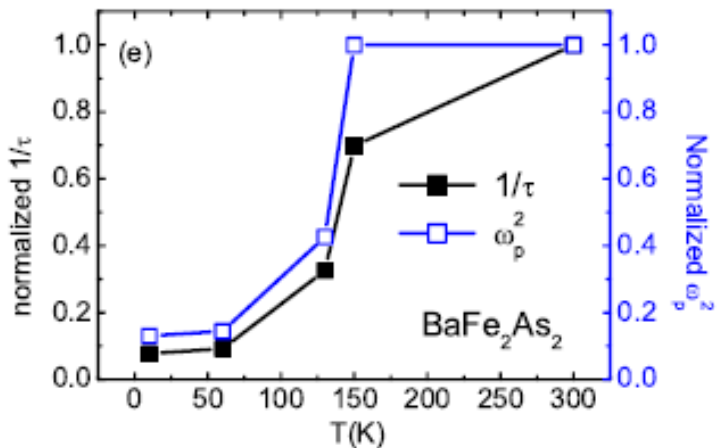


Double step  
in the SDW transition

Drop of  $\rho$  attributed to a decrease  
of the scattering rate at  $T_{\text{SDW}}$

Overcomes the reduction of carrier contents due  
to a partial gapping of the Fermi surface

Confirmed by optical measurements

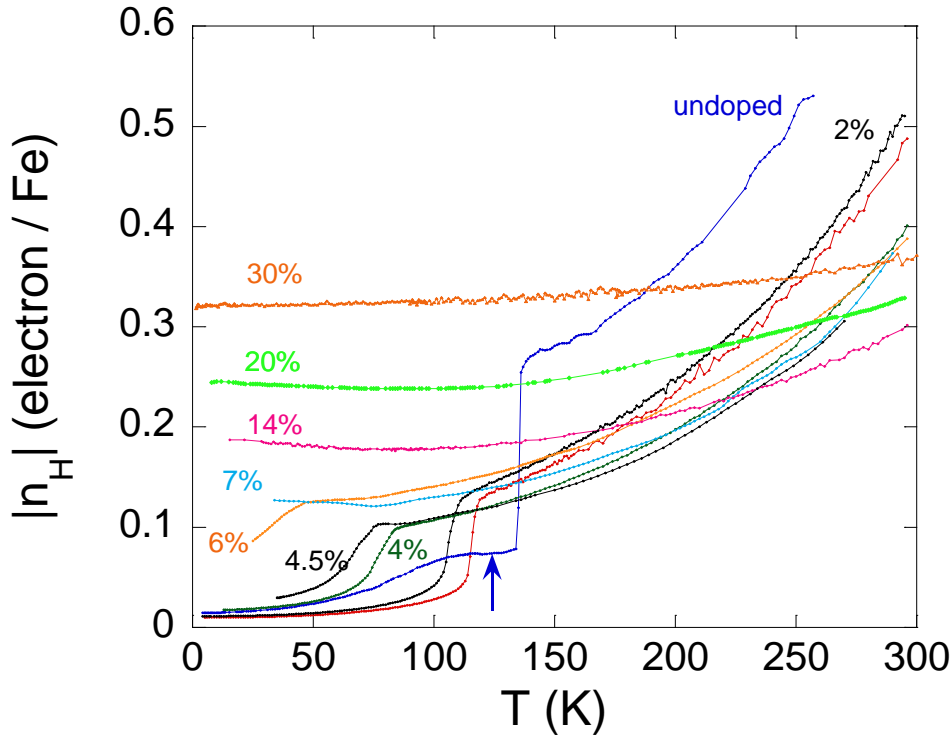


Reduction of the carrier content by a factor 8  
Reduction of the scattering rate by a factor 13

This can account for a factor  $\sim 2$  reduction  
of the resistivity

# Hall effect measurements

$$|n_H| = \frac{1}{e|R_H|}$$



In a two-band model:

$$\frac{1}{n_H e} = R_H = \frac{-R_e \sigma_e^2 + R_h \sigma_h^2}{(\sigma_e + \sigma_h)^2}$$

electrons

holes

$$R_e = \frac{1}{n_e e}$$

$$R_h = \frac{1}{n_h e}$$



$|n_H|$  upper bound for  $n_e$

4 unknown quantities:

$n_e, n_h, \sigma_e, \sigma_h$

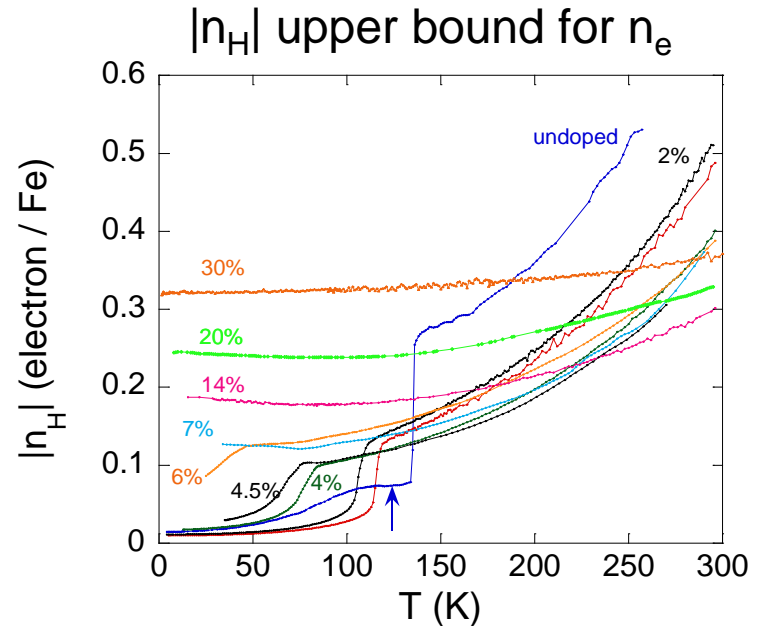
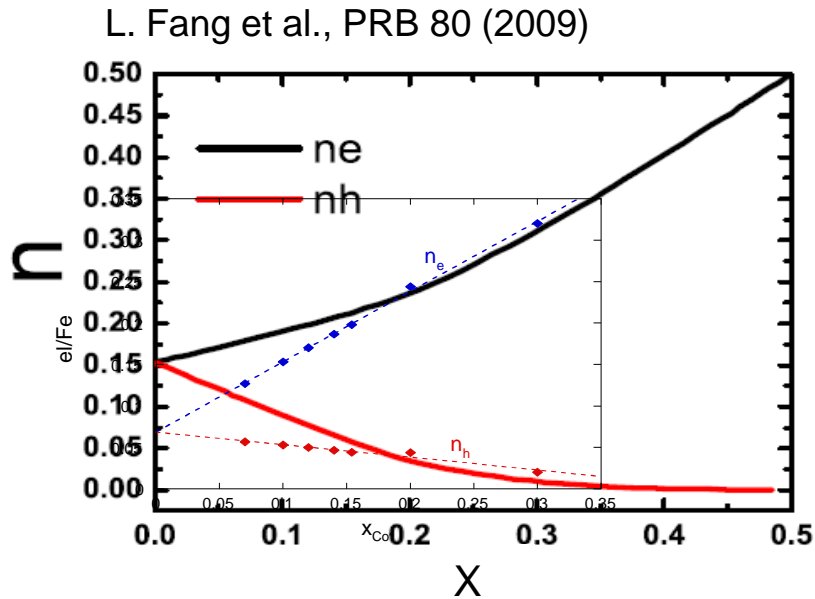
But only 3 relations

Conductivity:  $\sigma = \sigma_e + \sigma_h$

Hall effect

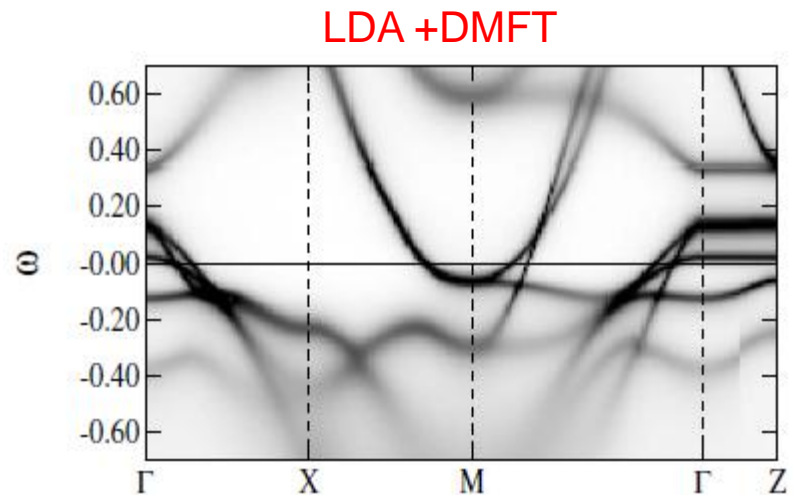
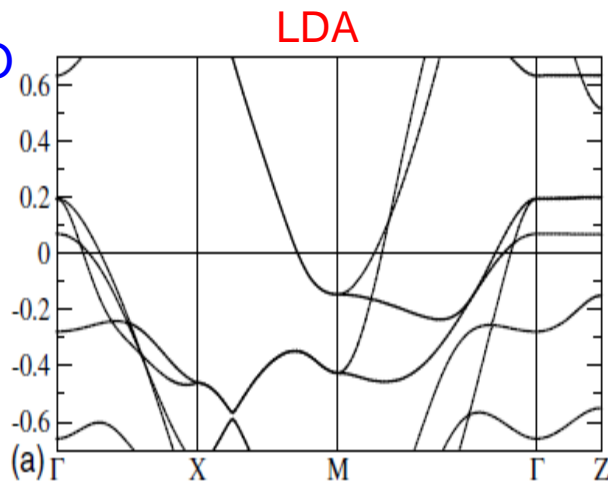
$$n_e = n_h + x_{Co}$$

# Using LDA calculations to determine $n_e$ and $n_h$



LDA calculations appear to overestimate the number of electrons and holes up to  $x \sim 0.2$

LaFeAsO

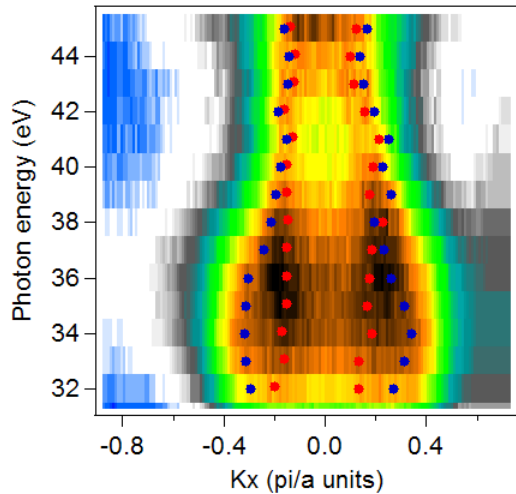


M. Aichhorn et al. PRB 80 (2009)

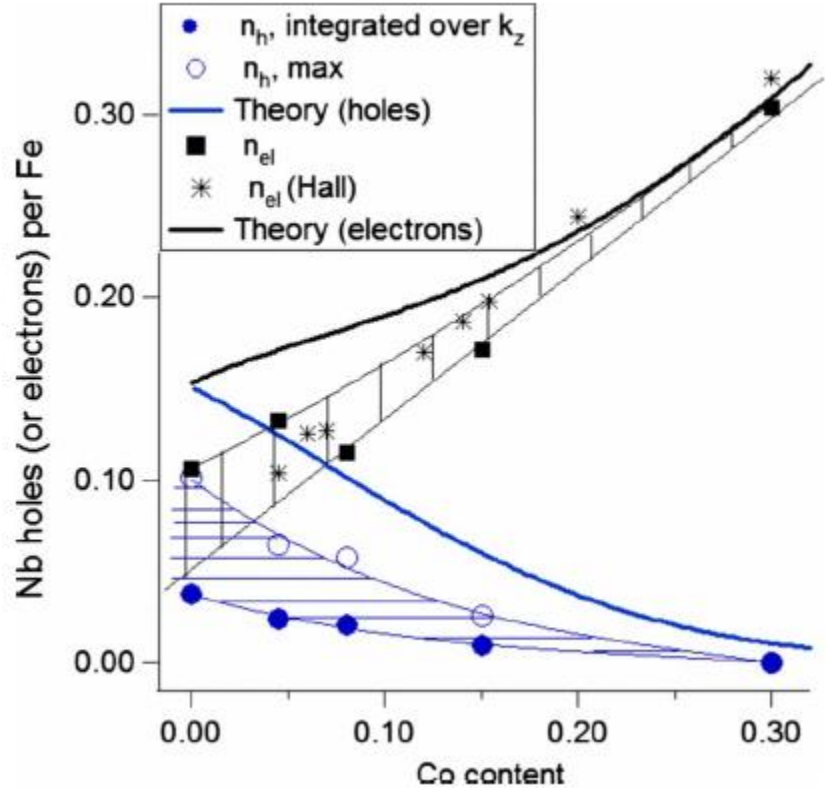
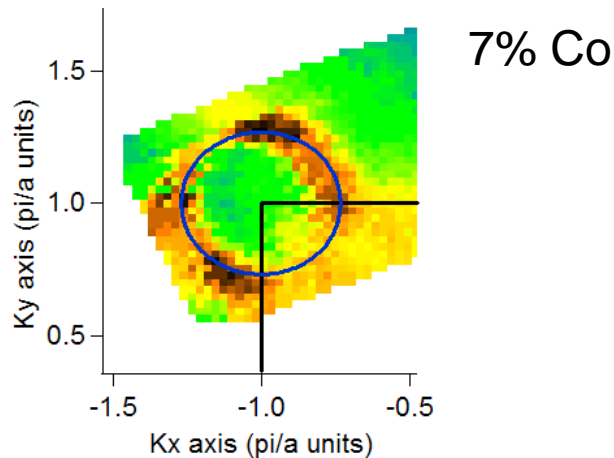
# Using ARPES measurements to determine $n_e$ and $n_h$

V. Brouet et al., PRB (2009)

Hole bands: two 2D, one 3D



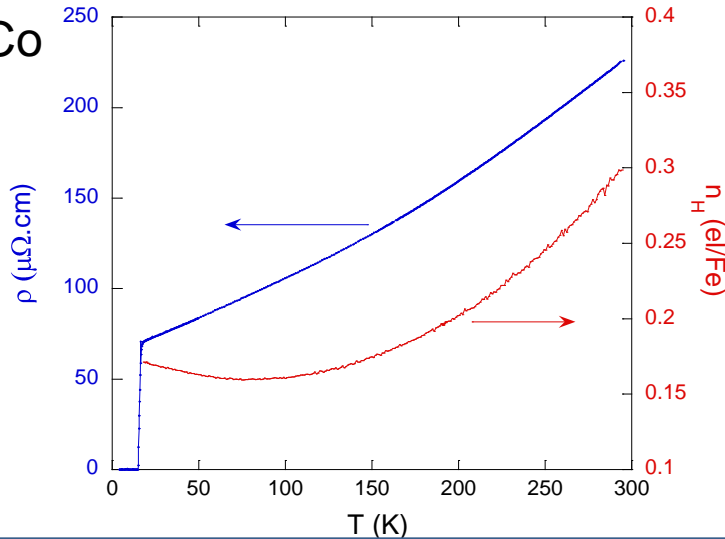
2 degenerate electron bands: 2D



Good agreement between ARPES and transport data

# Analysis of resistivity and Hall effect in a two-band model with T independent $n_e$ and $n_h$

10% Co

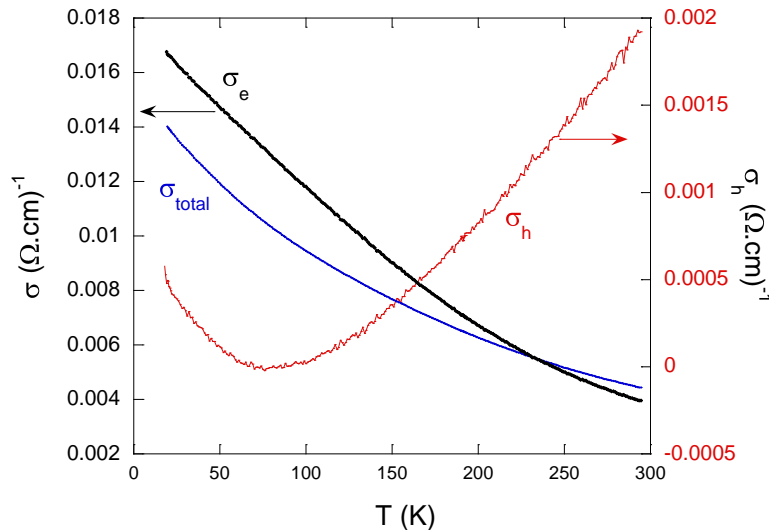


Minimal value of  $n_H \approx n_e$

For large Co doping :  $n_h \ll n_e$

$$n_H \approx n_e \left(1 + 2 \frac{\sigma_h}{\sigma}\right)$$

$$\sigma = \sigma_e + \sigma_h$$



Minimal value for  $\sigma_h$

Mobility of the holes increases with T!

Opposite to what expected  
for a metallic system

Assumption of constant  $n_e$  and  $n_h$   
with T is not valid



Holes not directly visible in the transport properties

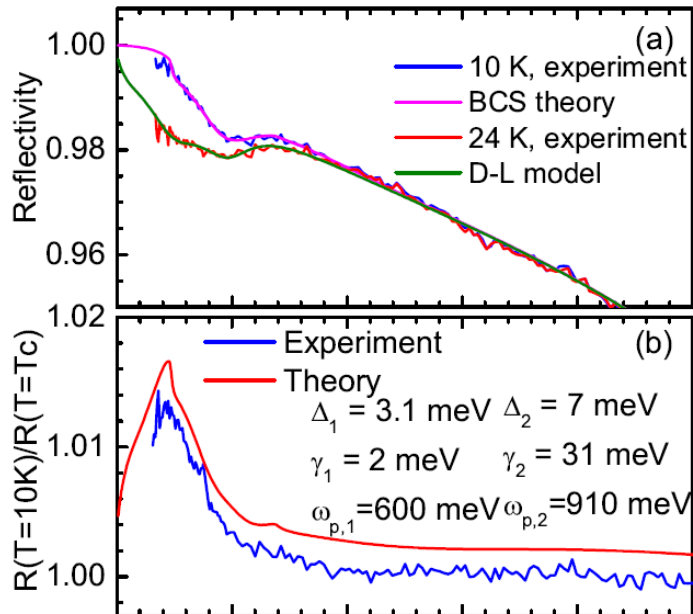


# Optical measurements in $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)\text{As}_2$

➔ One type of carriers more strongly scattered than the other one

## Superconducting state for $x=7\%$

*E. Van Heumen et al., arXiv 0912.0636*

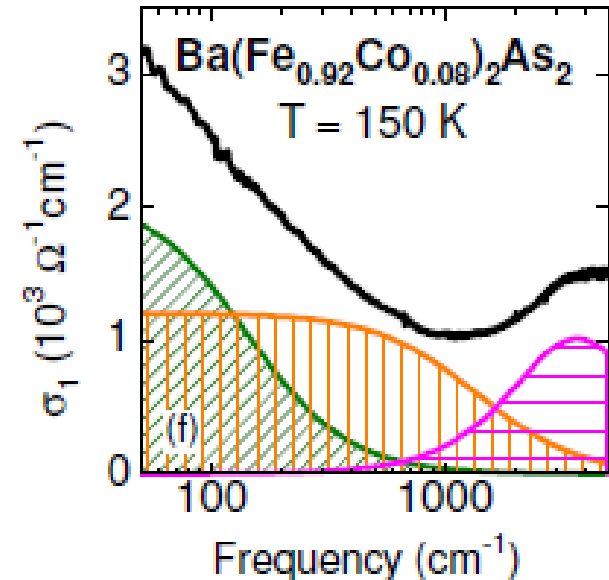


The holes are 15 times more scattered than the electrons

Interband scattering more efficient for the holes than for the electrons

## Normal state for $x=7\%$

*D. Wu et al., arXiv 0912.3334*

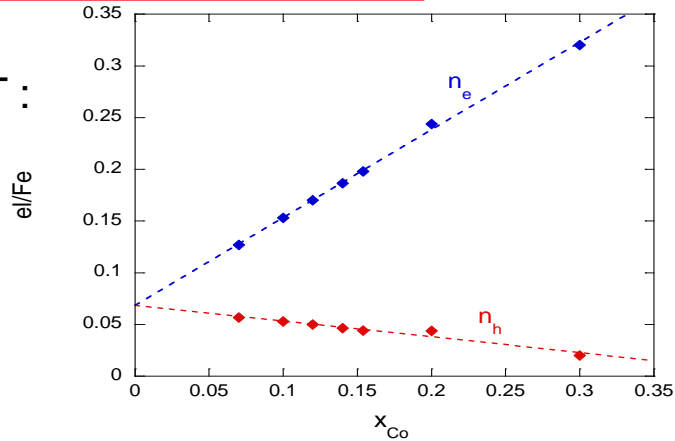


Two Drude components with very different scattering rates

# Carrier content vs Co doping and temperature

## Variation with Co content

At low T :



Each Co adds  $\sim 0.9e/Fe$

$$n_e = n_h + x_{Co} \longrightarrow n_h$$

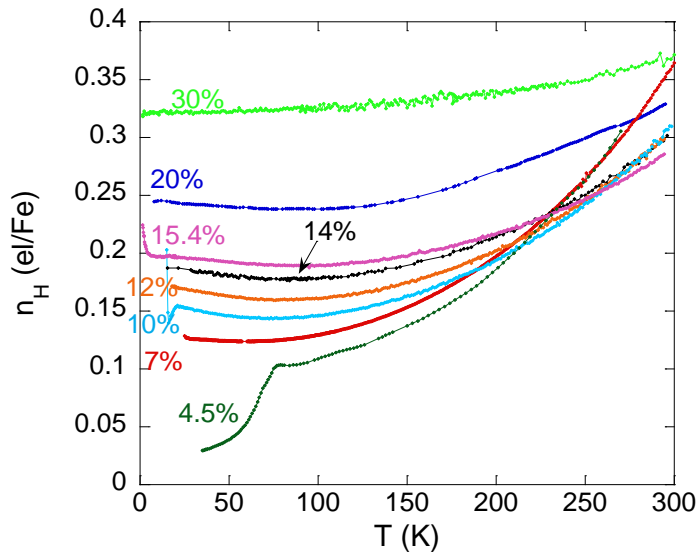
This result can be explained if the shapes of electron and hole pockets are different

Shift of the Fermi level



Different changes of  $n_e$  and  $n_h$ .

## Variation with T



For  $T \leq 150$  K and  $x \geq 4.5\%$  :  
rather weak T dependence

BaFe<sub>2</sub>As<sub>2</sub> : semimetal with the bottom of the electron bands and the top of the hole bands very near  $E_F$  (20-40meV)

Increasing T Shift of the chemical potential

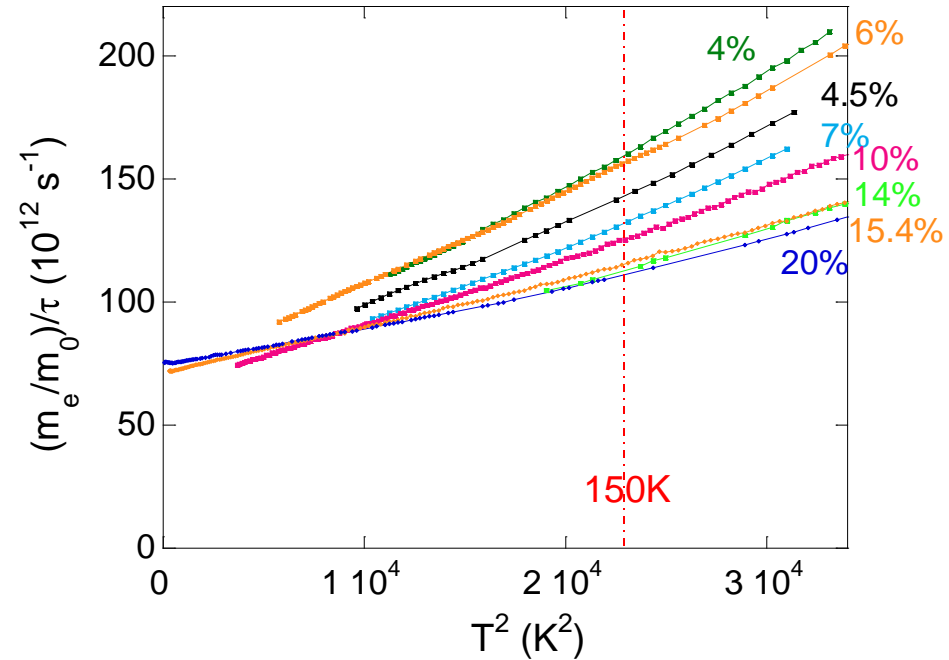
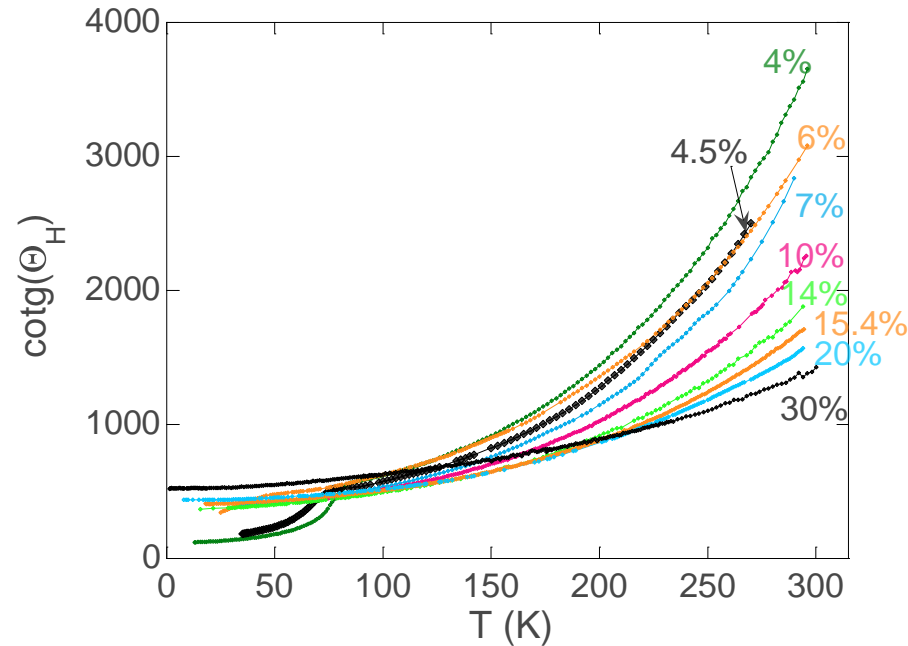
Equal modification of  $n_e$  and  $n_h$

# Electronic scattering rates

$$n_e(T) = n_H(T)$$

$$\rho(T) = \frac{m_e}{n_e(T)e^2\tau(T)}$$

$$\cot(\Theta_H) = \rho / |R_H| = m_e / e\tau$$

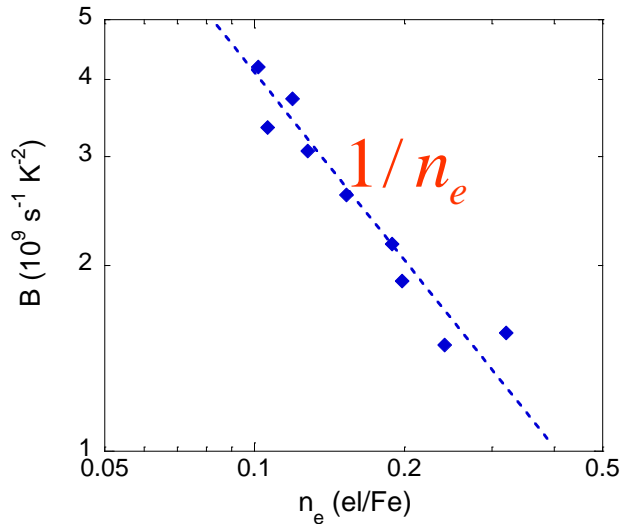


Up to 150K:  $1/\tau = 1/\tau_0 + BT^2$  ➔ Electron-electron interaction

$1/\tau_0$  independent on Co doping for  $x \leq 20\%$

No contribution of phonons nor spin fluctuations to the electronic scattering rates

# Fermi liquid behavior



Electron-electron scattering process:

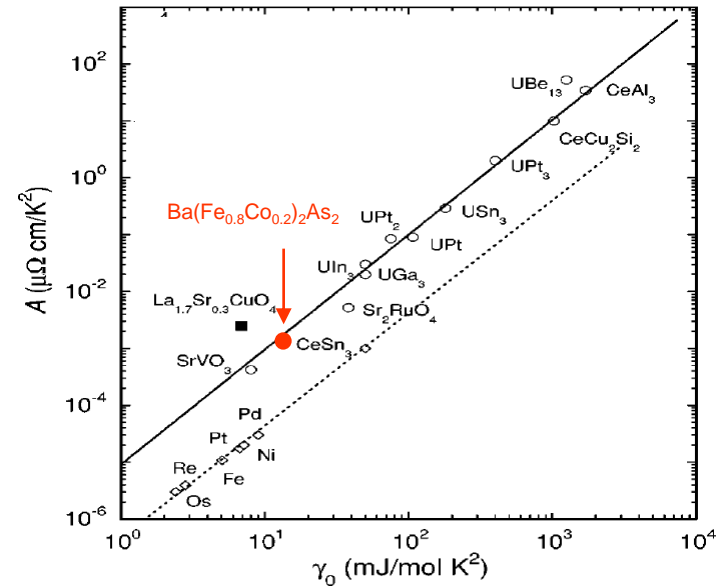
$$\frac{\hbar}{\tau_e(T)} \propto \frac{(k_B T)^2}{E_F}$$

In agreement with  $E_F \propto n_e$  in 2D

Comparison with Kadowaki woods ratio :

relation between  
the  $T^2$  coefficient of the resistivity  $A$   
and  
the electronic specific heat  $\gamma_0$

$x = 20\%$        $A = 1.4 \text{ n}\Omega \cdot \text{cm/K}^2$   
 $\gamma_0 \sim 10\text{-}15 \text{ mJ/mol K}^2$



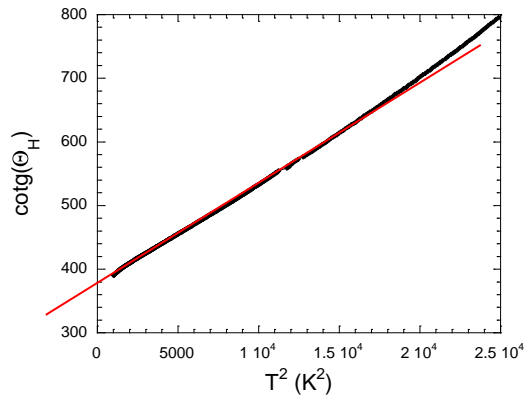
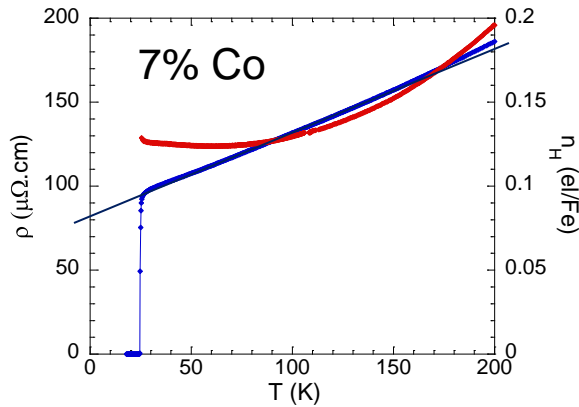
This corroborates our analysis in terms of transport dominated by electrons  
with a  $T$  variation of carriers

# Fermi liquid vs non Fermi liquid behavior

Observation of a T linear resistivity near optimal doping



Non Fermi liquid behavior

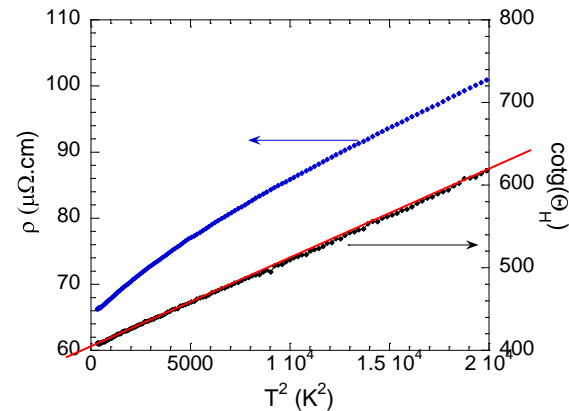
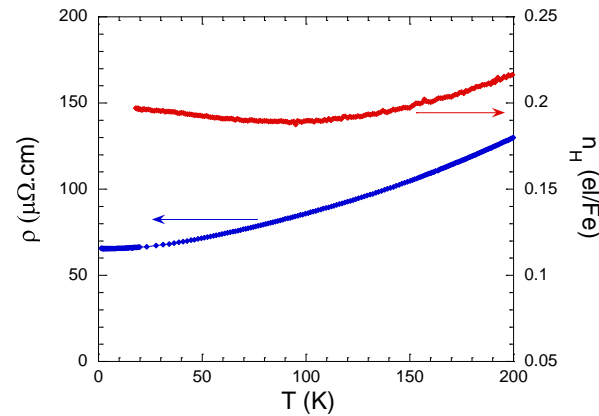


More precisely in the overdoped regime

$$\rho(T) = \rho_0 + AT + BT^2 \quad \text{with } A = 0 \text{ when } T_c = 0$$

*N. Doiron-Leyraud et al., PRB (2009)*

15.5% Co –  $T_c = 0\text{K}$

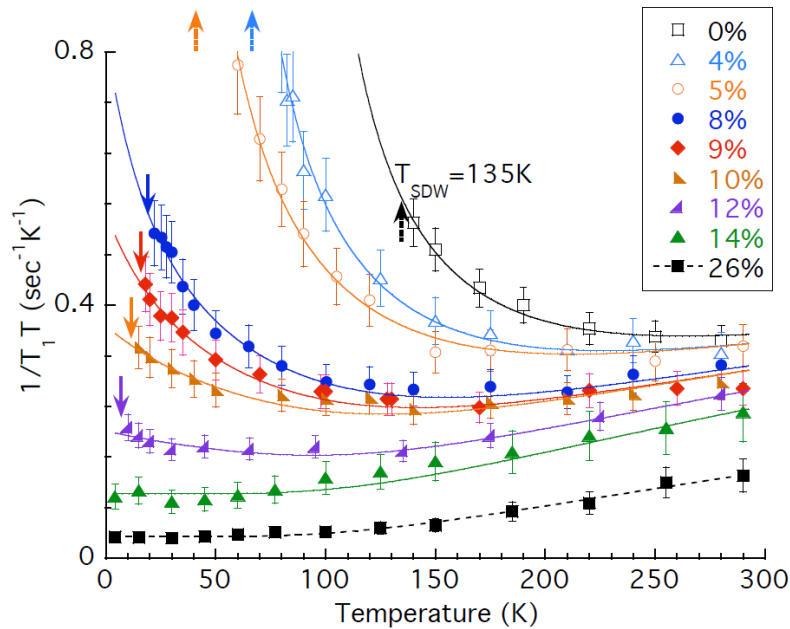


A very small variation of the number of carriers with temperature can be responsible for an apparent T linear resistivity

# Why the spin fluctuations are not directly apparent in the transport?

As NMR in underdoped and overdoped  $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)\text{As}_2$

Ning et al. Arxiv 0907.3875



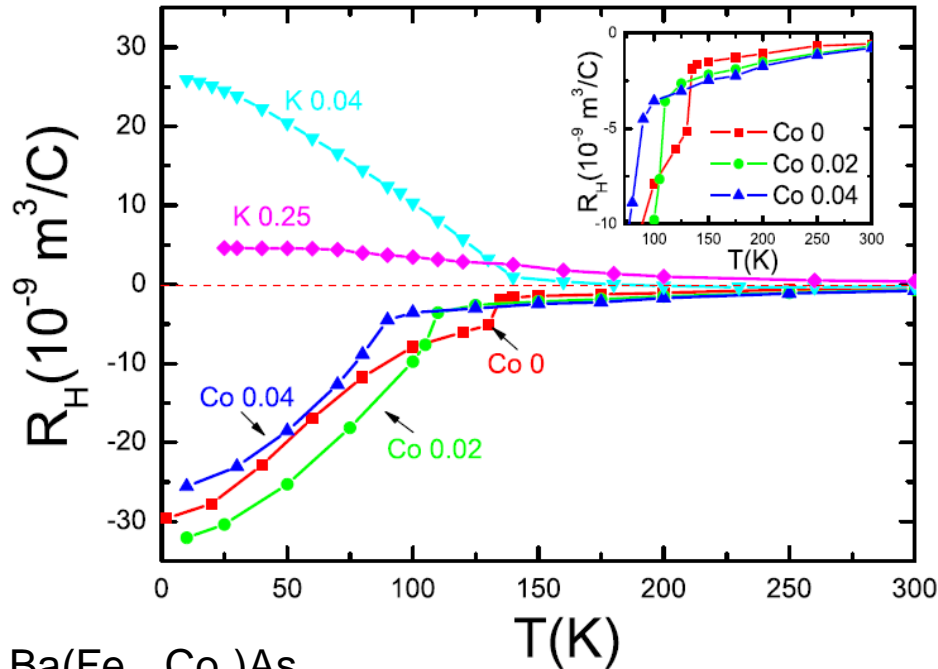
Huge increase of  $1/T_1T$  below 10%

Are the holes localized by the spin fluctuations?

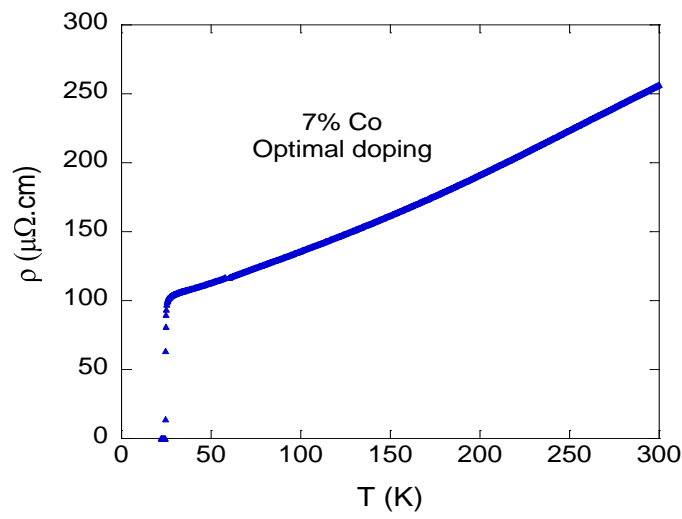
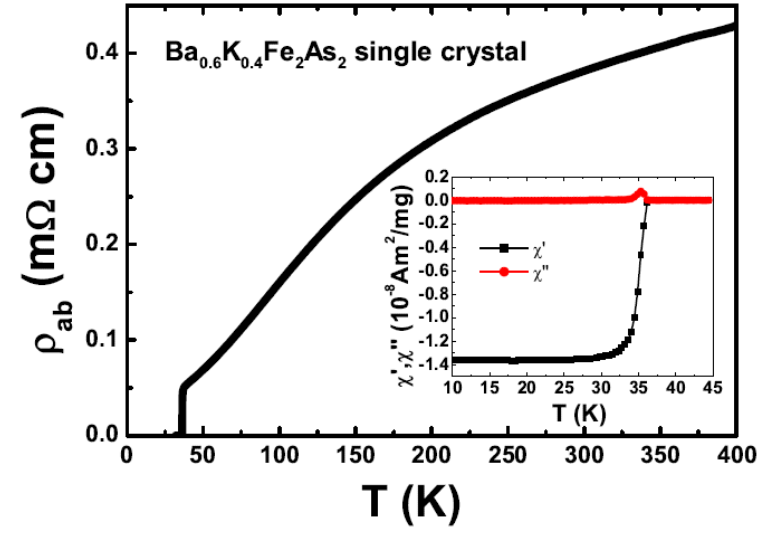
Why the electron transport is not affected?

# Electron-Hole asymmetry in the 122 family

L. Fang et al., PRB (2009)



$\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$



Very different  $\rho(T)$  curves



# Ru-doped: $\text{Ba}(\text{Fe}_{1-x}\text{Ru}_x)_2\text{As}_2$

Isovalent substitution

$$n_e = n_h = n$$

LDA calculations :

*L. Zhang and D.J. Singh, PRB (2009)*

Electronic structure near the Fermi level  
quite similar in  $\text{BaFe}_2\text{As}_2$  and  $\text{Ru-BaFe}_2\text{As}_2$

Results on polycrystalline samples

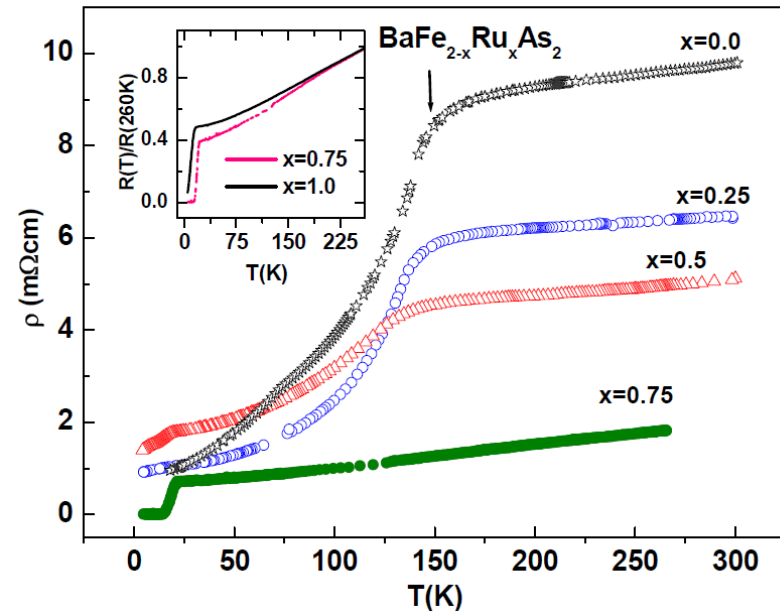
*S. Paulraj et al., arXiv 0902.2728*

Weak suppression of magnetism

Superconductivity for  $x \sim 0.35$

:  
But **negative Hall coefficient**

Electron doped?

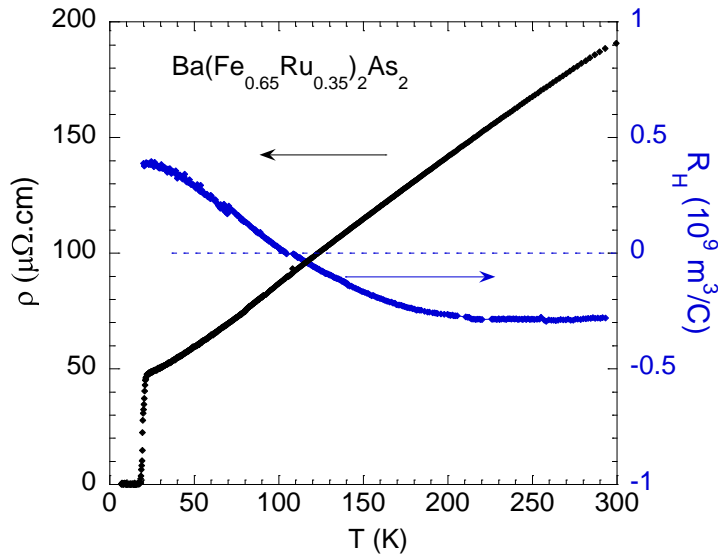


Negative Hall coefficient also in single crystals of  $\text{BaFe}_2(\text{As}_{1-x}\text{P}_x)_2$

*S. Kasahara et al., arXiv0905.4427*



# Transport and ARPES in $\text{Ba}(\text{Fe}_{1.65}\text{Ru}_{0.35})_2\text{As}_2$

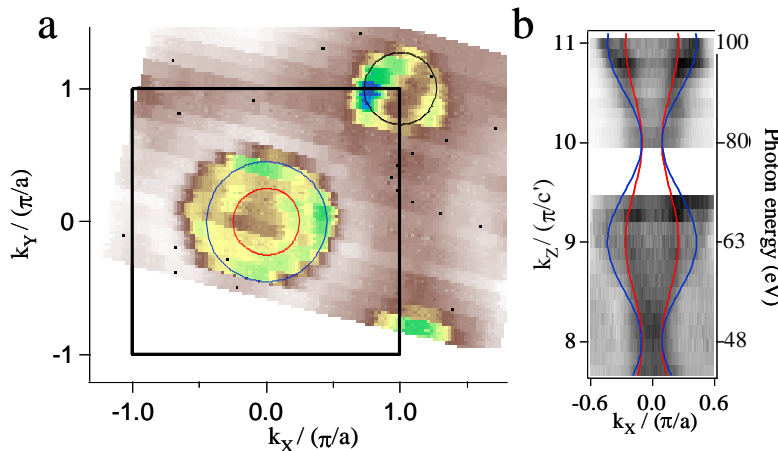


The Hall coefficient changes sign at low temperature



Mobility of holes larger than the one of electrons

ARPES measurements : See Poster by V. Brouet

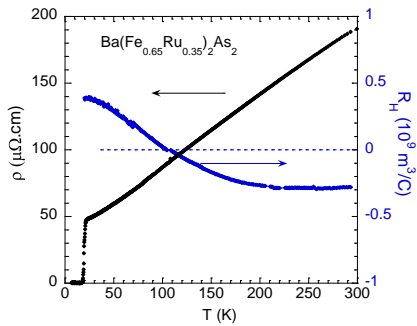


For the hole pockets:  
FS significantly warped along  $k_z$

$$n_e = n_h = n = 0.11 / Fe$$

Number of carriers nearly twice the one found by ARPES in  $\text{BaFe}_2\text{As}_2$

# Analysis of the transport data for $\text{Ba}(\text{FeRu}_{0.35})_2\text{As}_2$

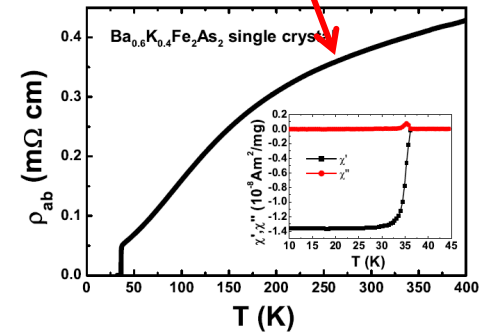
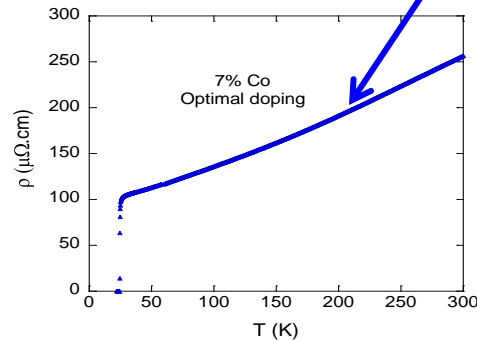
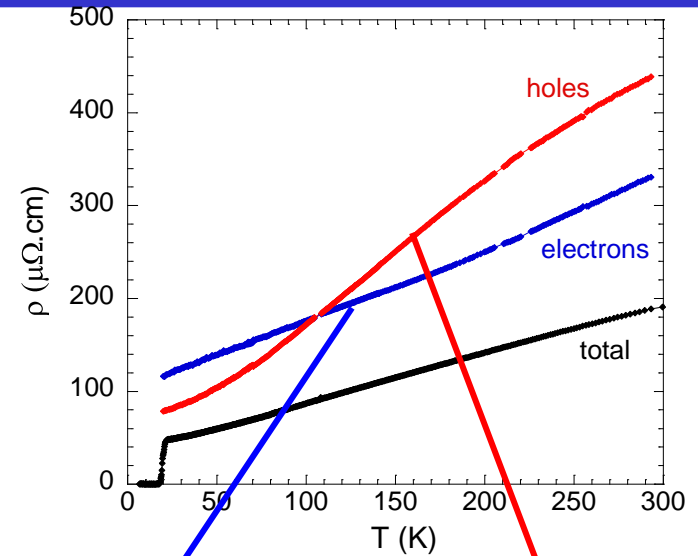


Compensated semi-metal  
In a two-band model:

$$R_H = \frac{1}{ne} \frac{(\sigma_h - \sigma_e)}{(\sigma_h + \sigma_e)}$$

$$\sigma = \sigma_e + \sigma_h$$

$$n = 0.11 e / \text{Fe}$$



Why are the holes less scattered in Ru doped compounds than in the undoped parent?

Spin fluctuations?

Modification of the electronic structure : weaker correlations – See poster by V. Brouet

# Summary

## □ Coherent picture of the charge transport in $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ at low T

Transport dominated by the electrons

The T variation of the Hall number explained by a T variation of the number of carriers

## □ Scattering processes : Fermi liquid behavior

*Spin fluctuations not directly visible in the transport properties*

Residual scattering rates independent on Co doping

## □ Ru-substituted $\text{BaFe}_2\text{As}_2$

Phase diagram very similar to the one found for Co or K doped compounds

Electrons and holes contribute to the transport.

Modification of the electronic structure

