



The Abdus Salam
**International Centre
for Theoretical Physics**

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RANDOM MATRIX MODEL WITH LOCALIZATION AND ERGODIC TRANSITIONS

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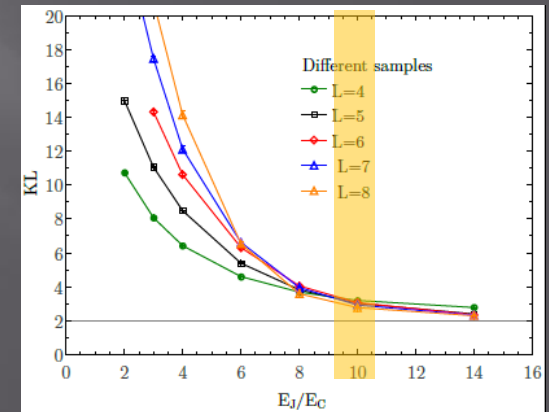
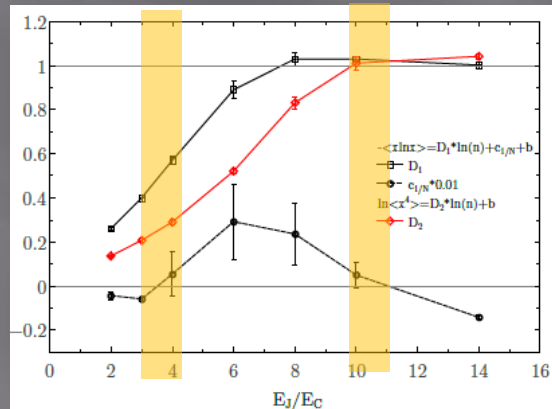
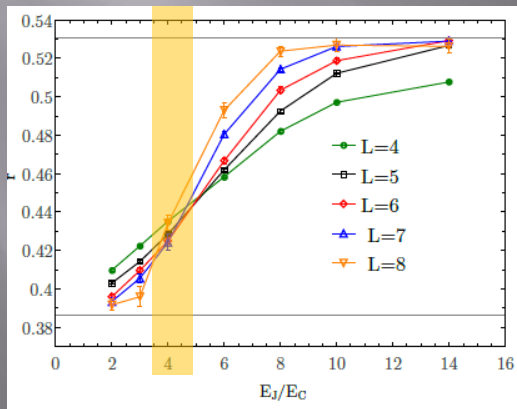


Two localization transitions in JJA?

many-body
insulator

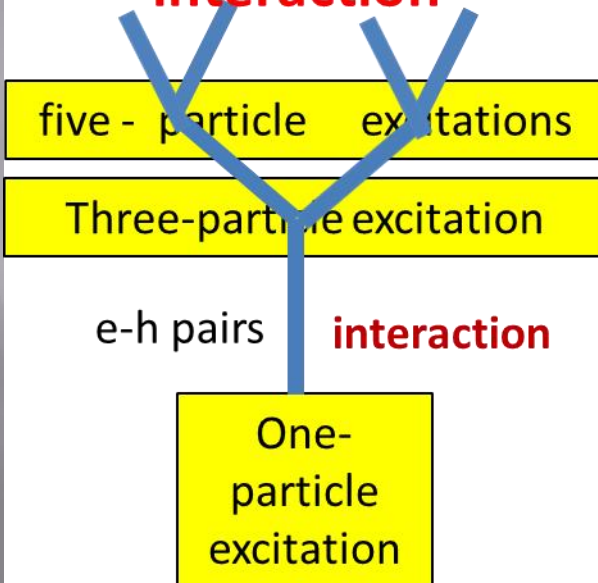
bad metal

good
metal

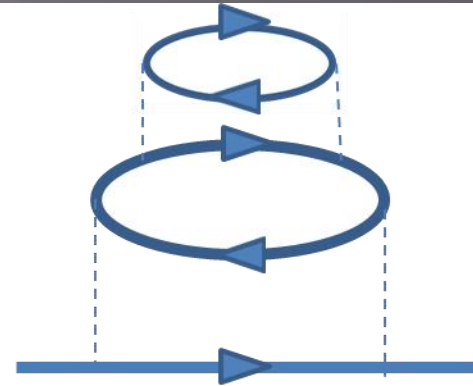


Hierarchical structure of many-body Fock space

**Tree-like structure
of many-body
interaction**



Altshuler, Gefen, Kamenev,
Levitov, 1997



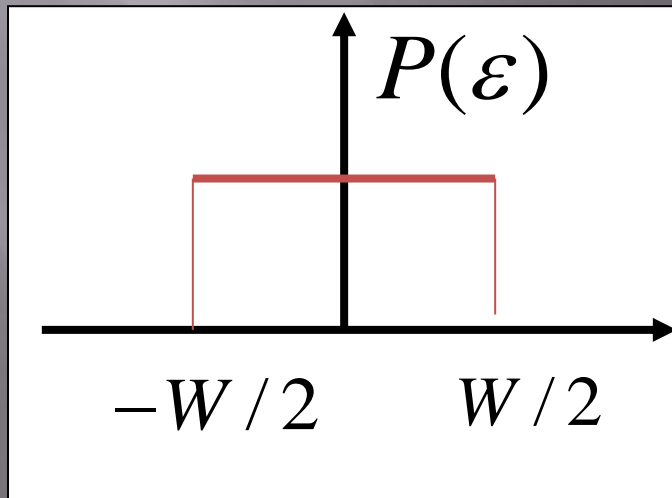
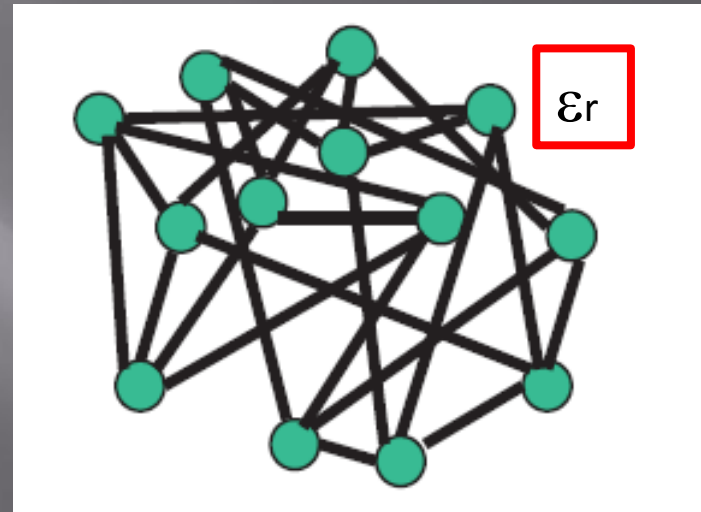
Basko, Aleiner, Altshuler,
2005



Anderson model on random regular graphs (RRG)

$$H = -I \sum_{\langle r, r' \rangle} c_r^+ c_{r'} + \sum_r \varepsilon_r n_r$$

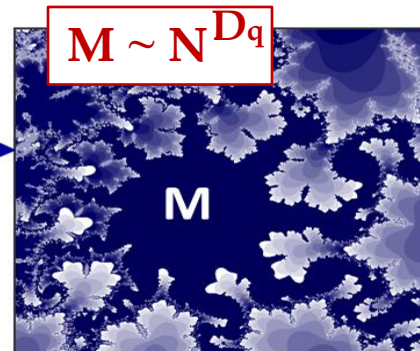
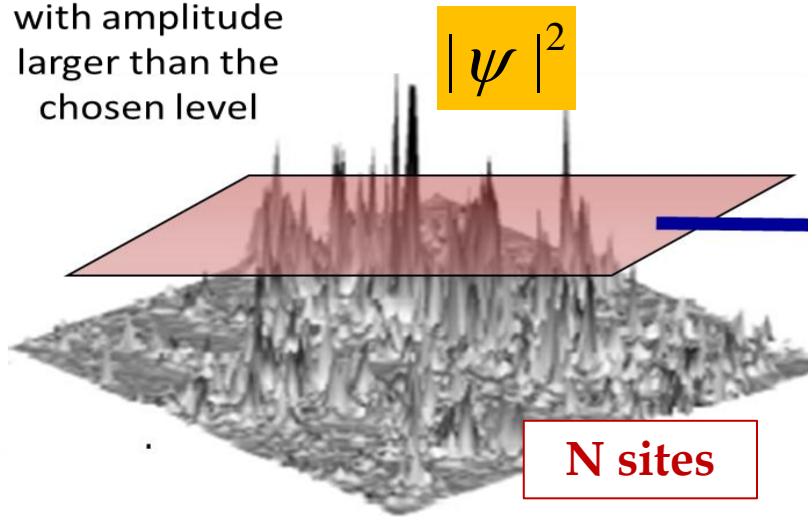
branching number K



Disorder strength W

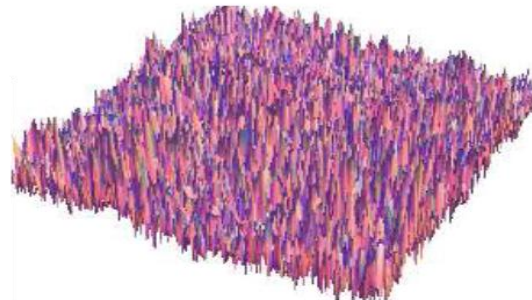
Ergodic and non-ergodic extended states

Map of the regions with amplitude larger than the chosen level



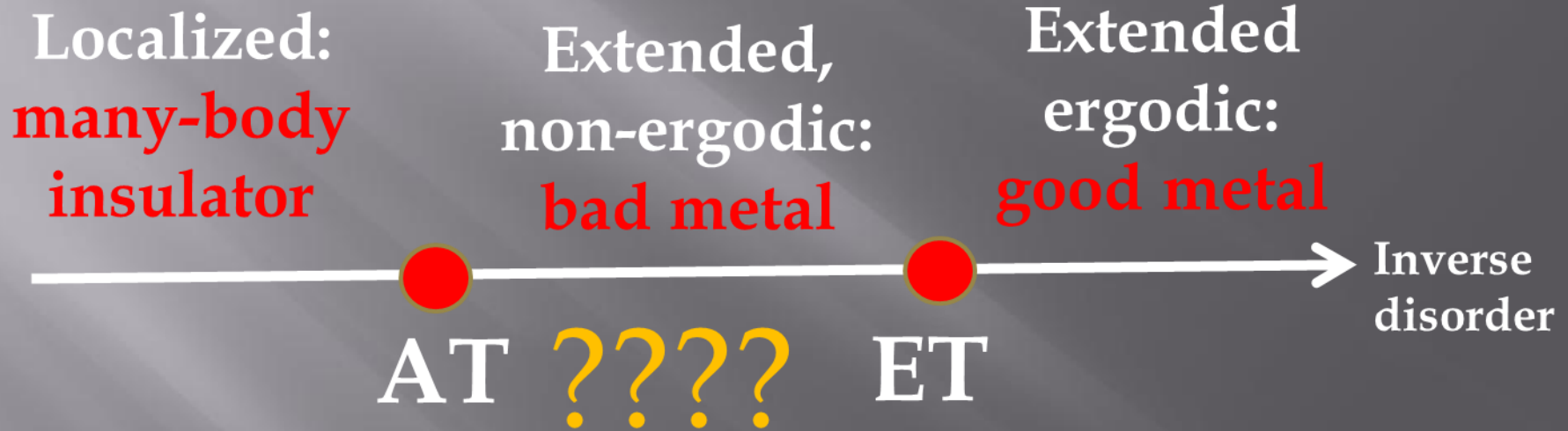
Non-ergodic, fractal states:
 $M/N \rightarrow 0$
 $M \rightarrow \infty$

$$I_q = \left\langle \sum_r |\Psi(r)|^{2q} \right\rangle \propto \frac{1}{N^{(q-1)D_q}}$$



Ergodic states:
 $M/N \rightarrow \text{cst}$
 $M \rightarrow \infty$

Search for a model with both localization and ergodic transitions

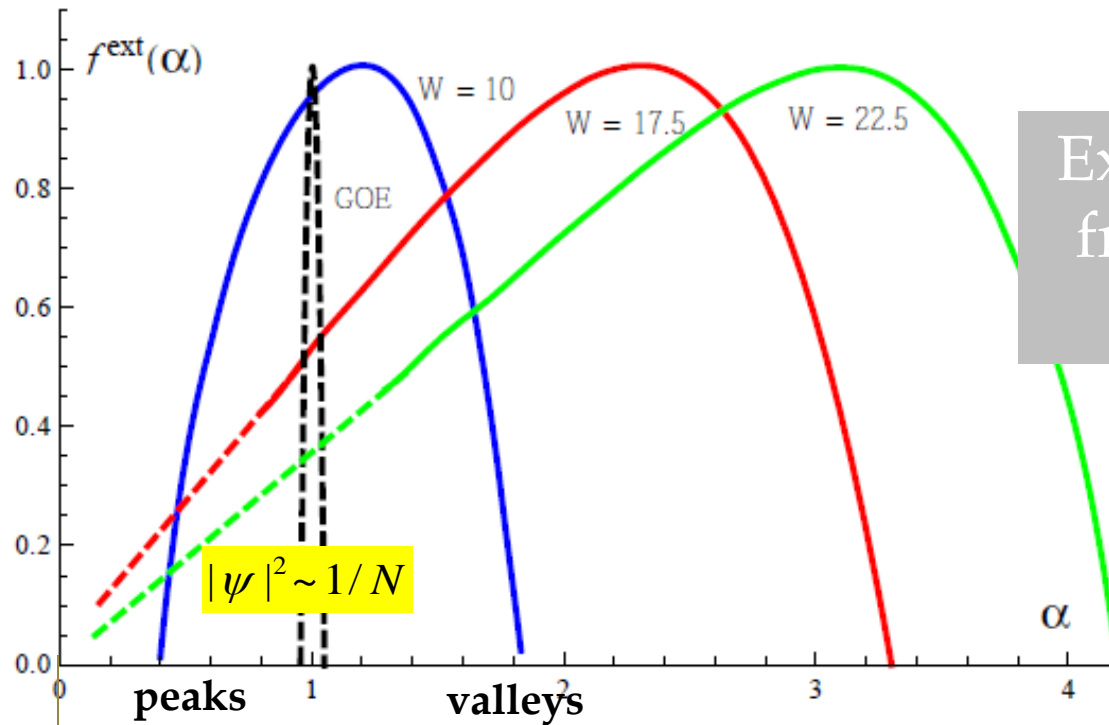
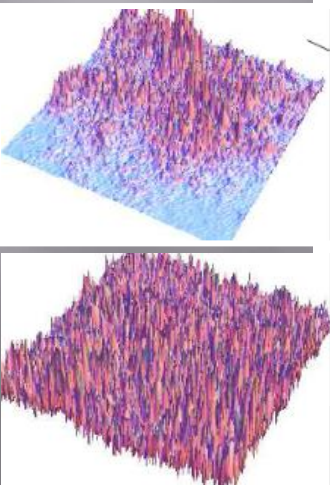


Search for ergodic transition (ET) on Bethe lattice (RRG)

Mirlin & Fyodorov: 1991-1997:

ET: **No, only ergodic and localized phases**

$$M = N^{f(\alpha)}$$



Extrapolation from $N=2K-32K$

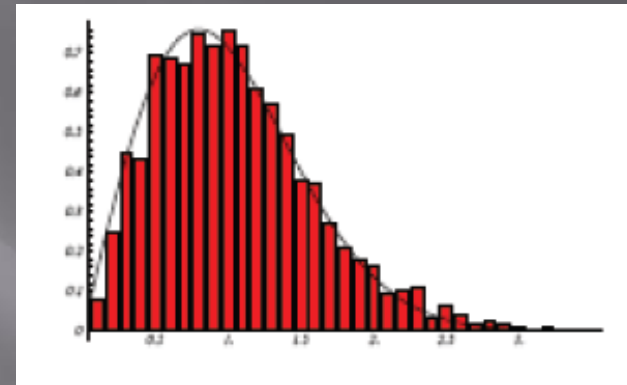
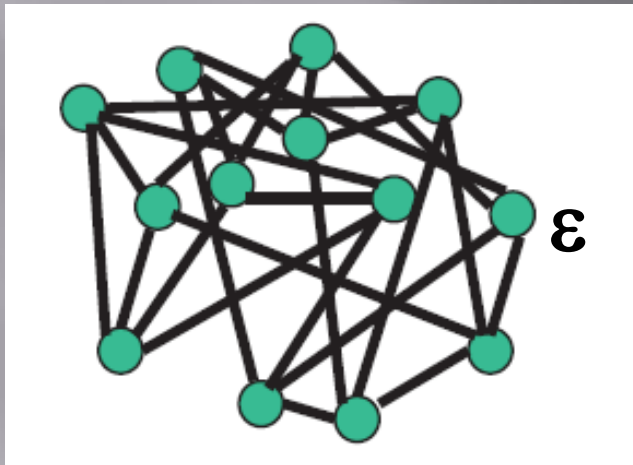
$$|\psi|^2 \sim N^{-\alpha}$$

Giulio Biroli, 2011: ET
- **YES, but...**

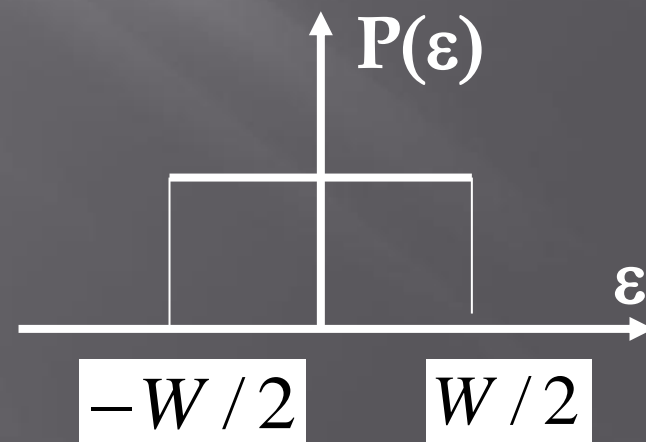
De Luca et al. 2014: ET -
No, only non-ergodic and localized phases

RRG: two disorder ensembles

Structural disorder: WD (Uzy Smiliansky)



On-site energy disorder:



A simpler model-relative of RRG?

WD RMT:

$$\sigma = \frac{\lambda^2}{N^\gamma} \lll 1$$

Special diagonal:
Rosenzweig-
Porter (1960)
ensemble

Mimics pristine RRG

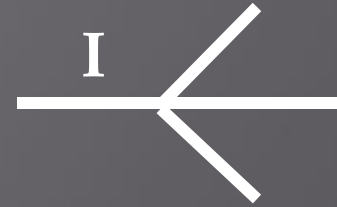
$$\langle H_{nm}^2 \rangle = \begin{pmatrix} \sigma & \sigma & \sigma & \sigma \\ \sigma & \sigma & \sigma & \sigma \\ \sigma & \sigma & \sigma & \sigma \\ \sigma & \sigma & \sigma & \sigma \end{pmatrix}$$

Adding diagonal disorder

$$\langle H_{nm}^2 \rangle = \begin{pmatrix} 1 & \sigma & \sigma & \sigma \\ \sigma & 1 & \sigma & \sigma \\ \sigma & \sigma & 1 & \sigma \\ \sigma & \sigma & \sigma & 1 \end{pmatrix}$$

Critical values of γ

$$W \sim I d$$



$$\langle H_{nm}^2 \rangle = \begin{pmatrix} 1 & \sigma & \sigma & \sigma \\ \sigma & 1 & \sigma & \sigma \\ \sigma & \sigma & 1 & \sigma \\ \sigma & \sigma & \sigma & 1 \end{pmatrix}$$

disorder

Hopping
integral

Coordination
number

$$\sigma = \frac{\lambda^2}{N^\gamma} \ll 1$$

$$1 \sim \frac{\lambda}{N^{\gamma/2}} N$$



$$\gamma = 2$$

I is of random sign!

$$d_{\text{eff}} \sim \sqrt{N}$$



$$\gamma = 1$$

FACTOR

$$K(t, t') = \langle \exp[itE_n - it' E_m] \rangle$$

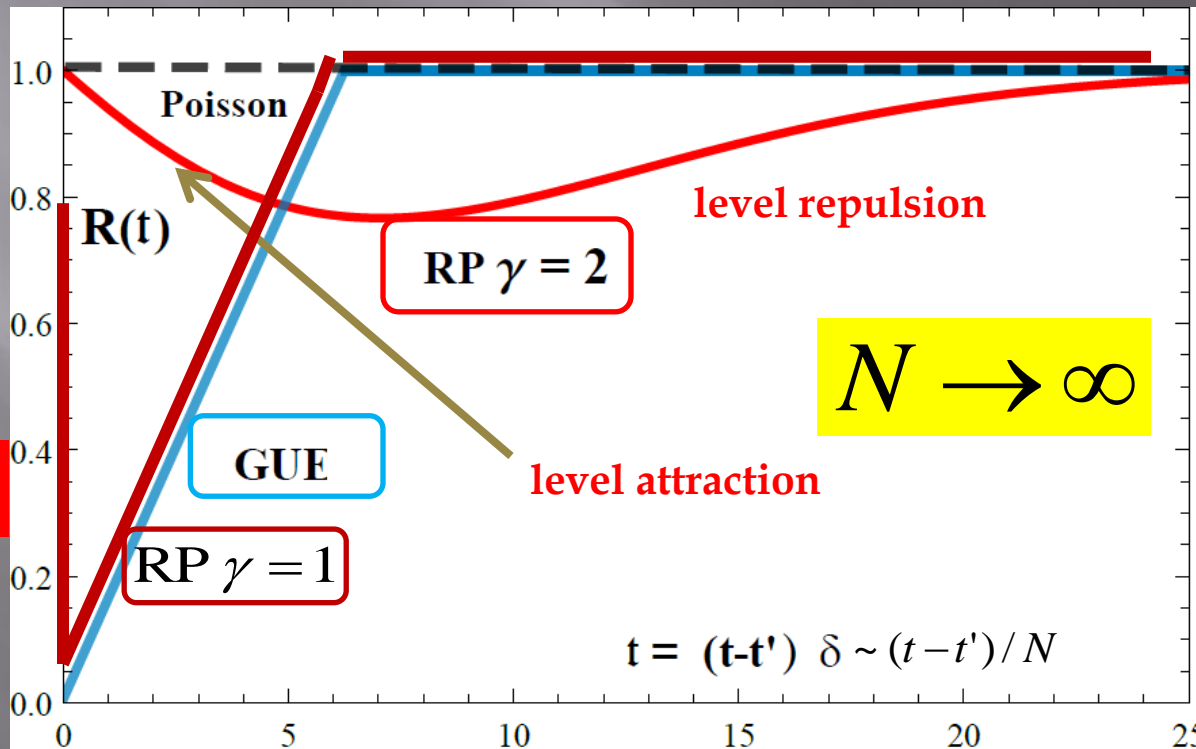
Pandey 1995,
Kunz & Shapiro, 1998

$$\sigma = \frac{\lambda^2}{N^\gamma} \lll 1$$

Brezin &
Hikami, 1995

$\gamma=2$

$\gamma=1$



Jump!

Generalization of Kunz & Shapiro

$$H = H_0 + A$$

$$H_0 = H - A, P(H_0) = \exp\left[-\frac{\text{tr}H_0^2}{2\sigma}\right]$$

GUE

Diagonal RM

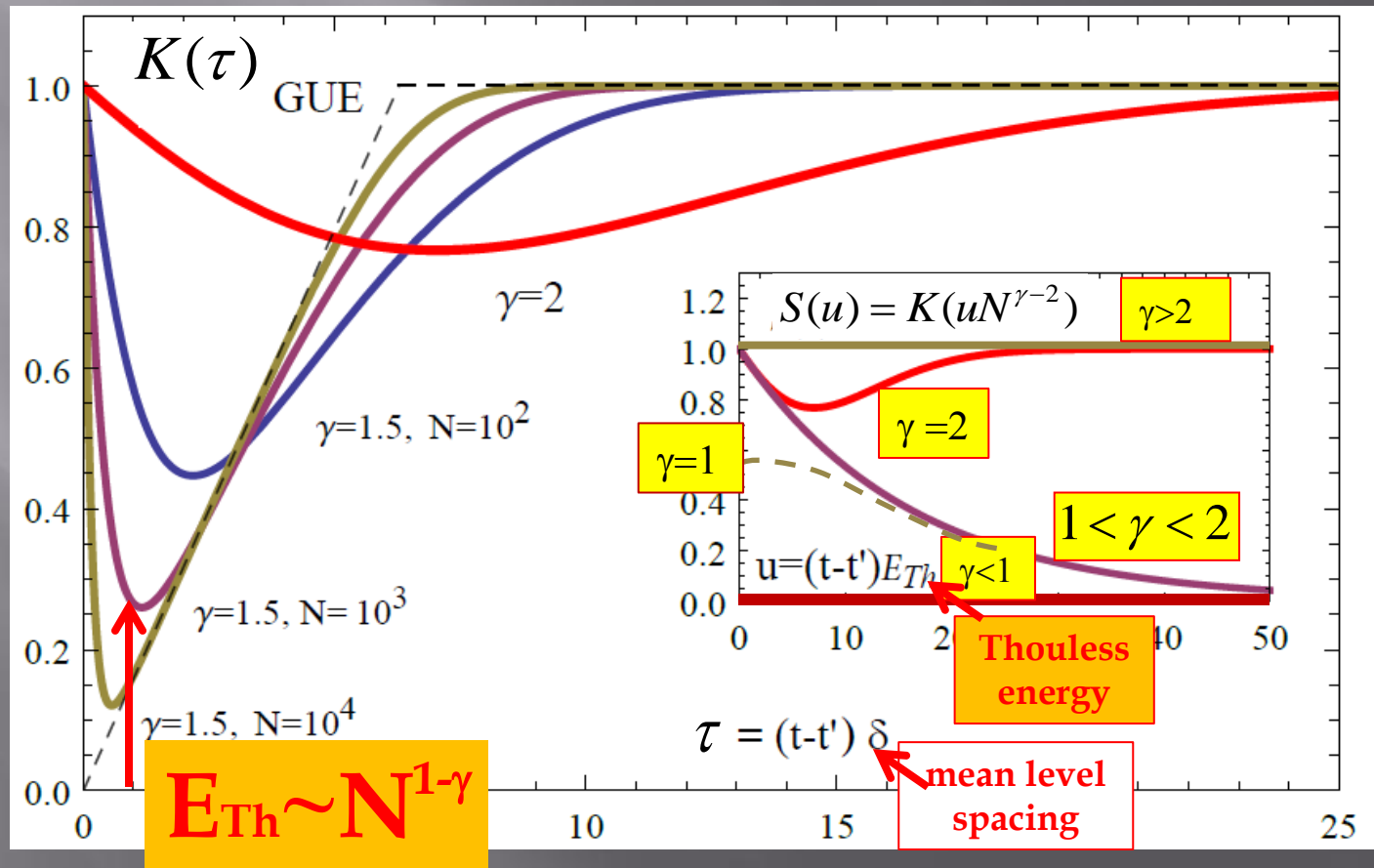
Izikson-Zuber formula of integrating over unitary matrices diagonalizing H

$$\int dU \exp\left[\frac{1}{\sigma} \text{tr} AU\lambda U^\dagger\right] = \frac{c}{\Delta(\lambda)\Delta(a)} \det \exp\left(\frac{a_i \lambda_j}{\sigma}\right)$$

ANALYTICAL SOLUTION FOR SPECTRAL FORM-FACTOR

$$S(u) = 1 + e^{-2\pi\Lambda^2 u} e^{-\Lambda^2 u^2 N^{\gamma-2}} \left[\frac{2I_1(\kappa u^{3/2})}{\kappa u^{3/2}} - \frac{1}{4\pi} \kappa u^{5/2} N^{\gamma-2} \int_0^\infty \frac{x dx}{\sqrt{x+1}} I_1(\kappa u^{3/2} \sqrt{x+1}) e^{-x u^2 \Lambda^2 N^{\gamma-2}} \right]$$

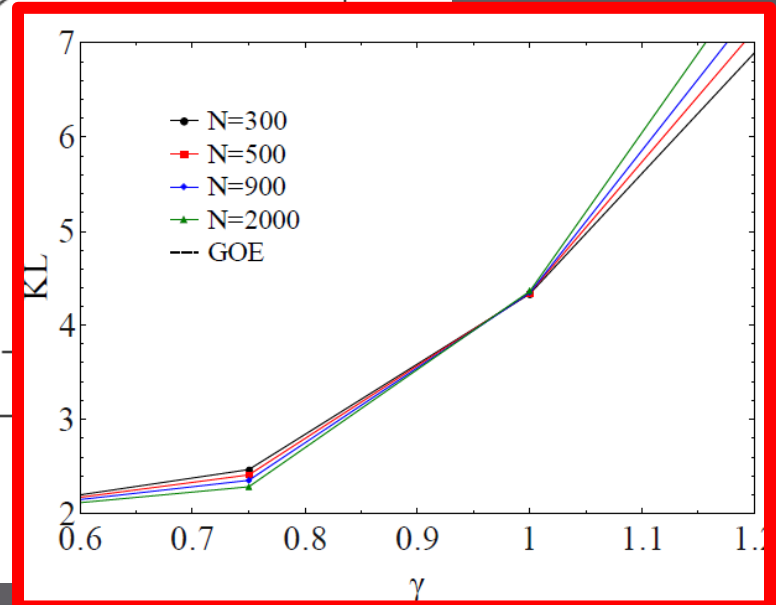
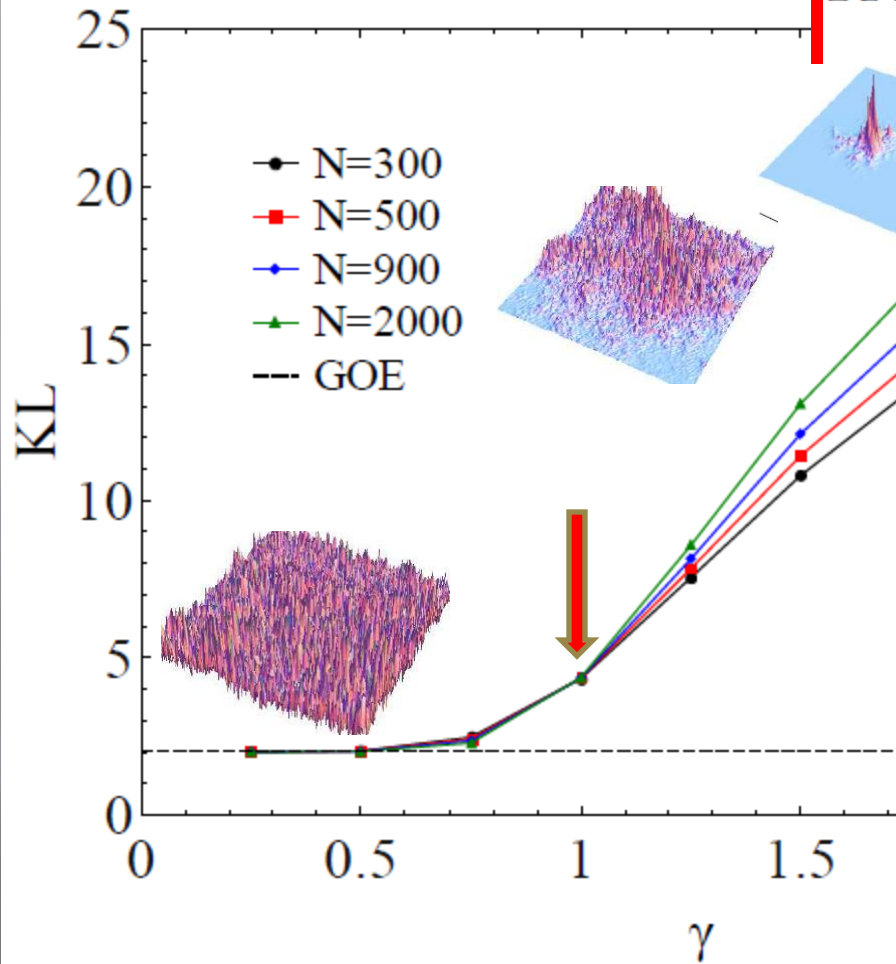
$$\kappa = \sqrt{8\pi} N^{\gamma-2} \Lambda^2 \text{ and } \Lambda = \lambda p(0)$$



KL-statistics

Kullback-Lieber

$$KL(\psi, \varphi) = \sum_{i=1}^N \psi_i^2 \log\left(\frac{\psi_i^2}{\varphi_i^2}\right)$$

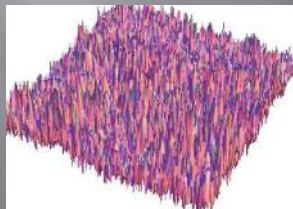


Two transition points

$$\sigma = \frac{\lambda^2}{N^\gamma} \ll 1$$

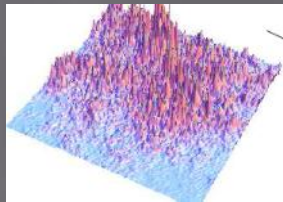
$$\langle H_{nm}^2 \rangle = \begin{pmatrix} 1 & \sigma & \sigma & \sigma \\ \sigma & 1 & \sigma & \sigma \\ \sigma & \sigma & 1 & \sigma \\ \sigma & \sigma & \sigma & 1 \end{pmatrix}$$

Extended
ergodic



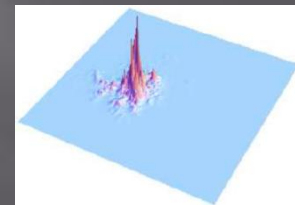
1

Extended,
non-ergodic



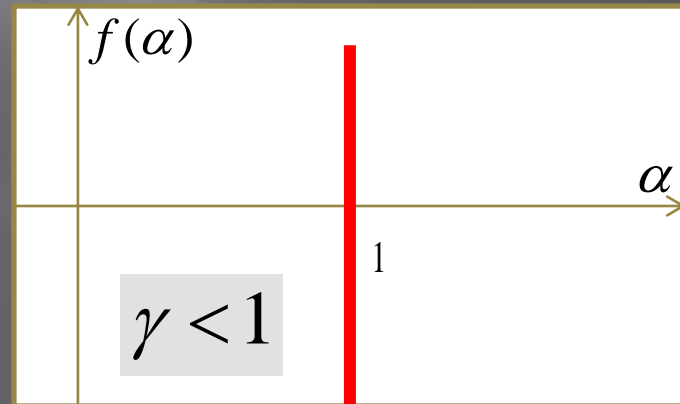
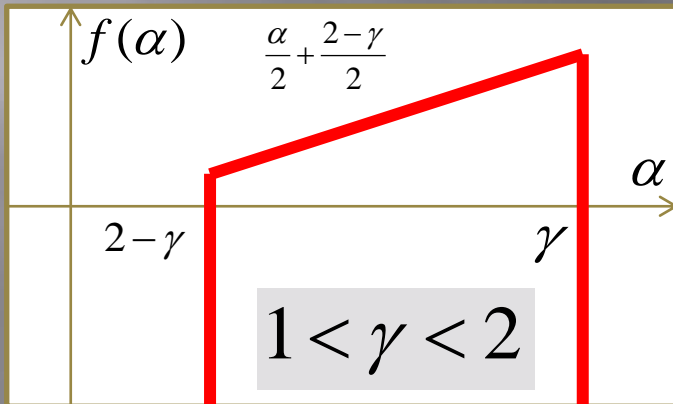
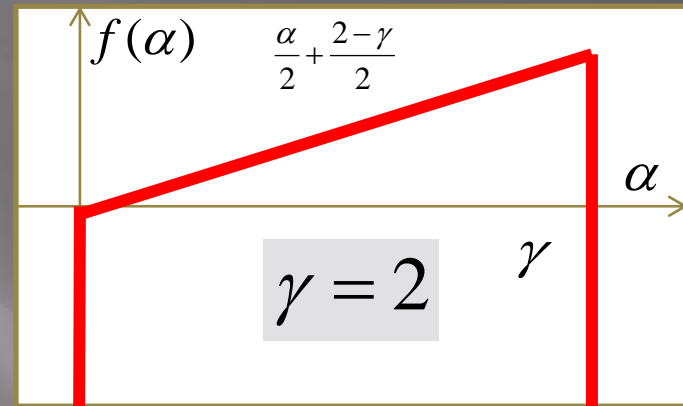
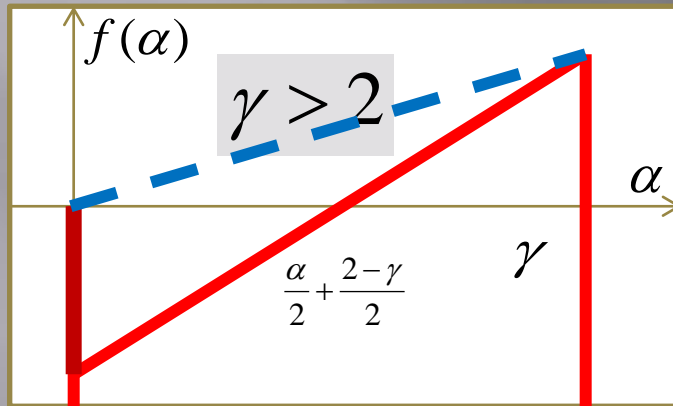
2

localized



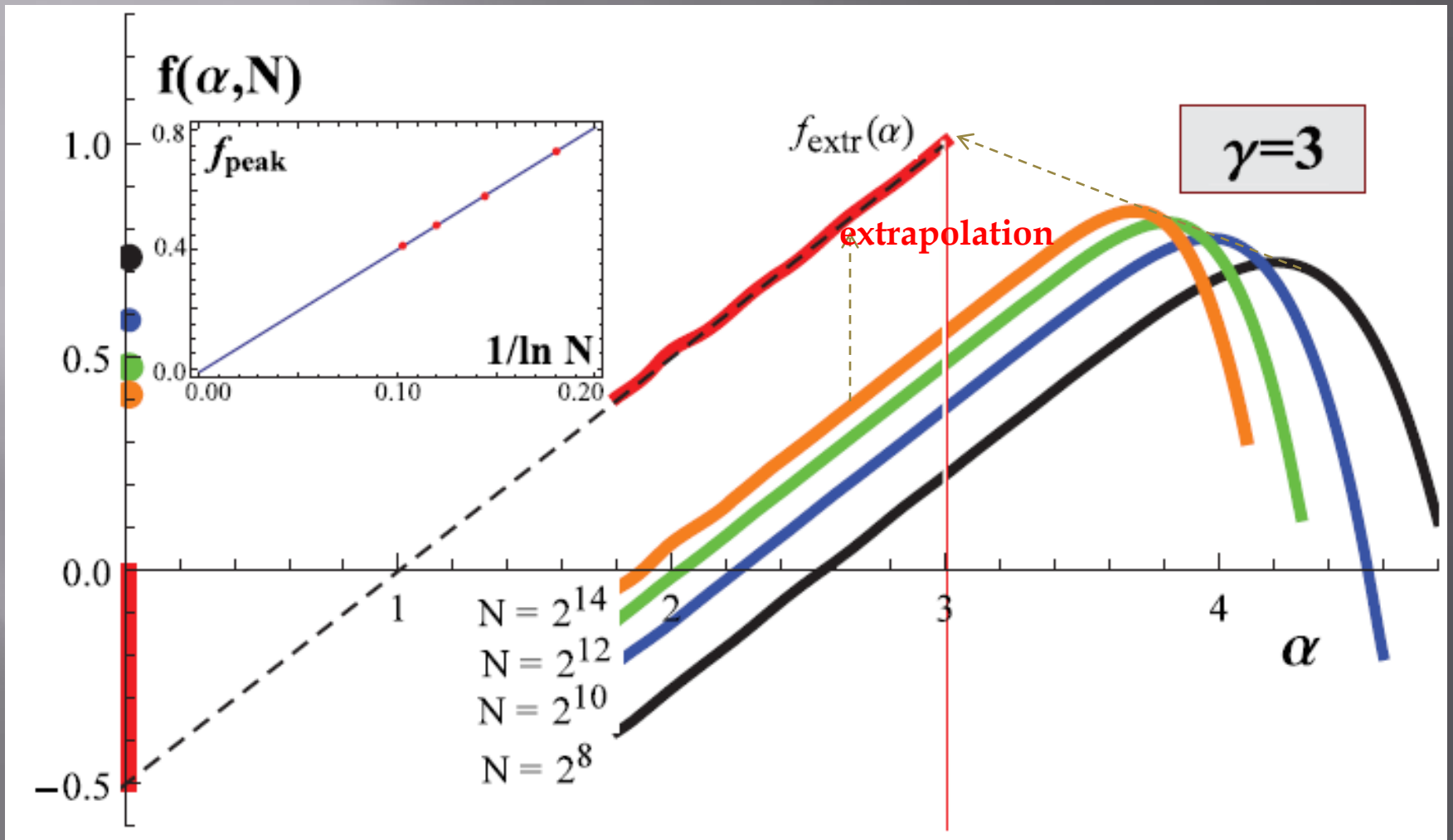
γ

Multifractality spectrum $f(\alpha)$



$$D_{q>1/2} = 2 - \gamma$$

Numerics for $f(\alpha, N)$

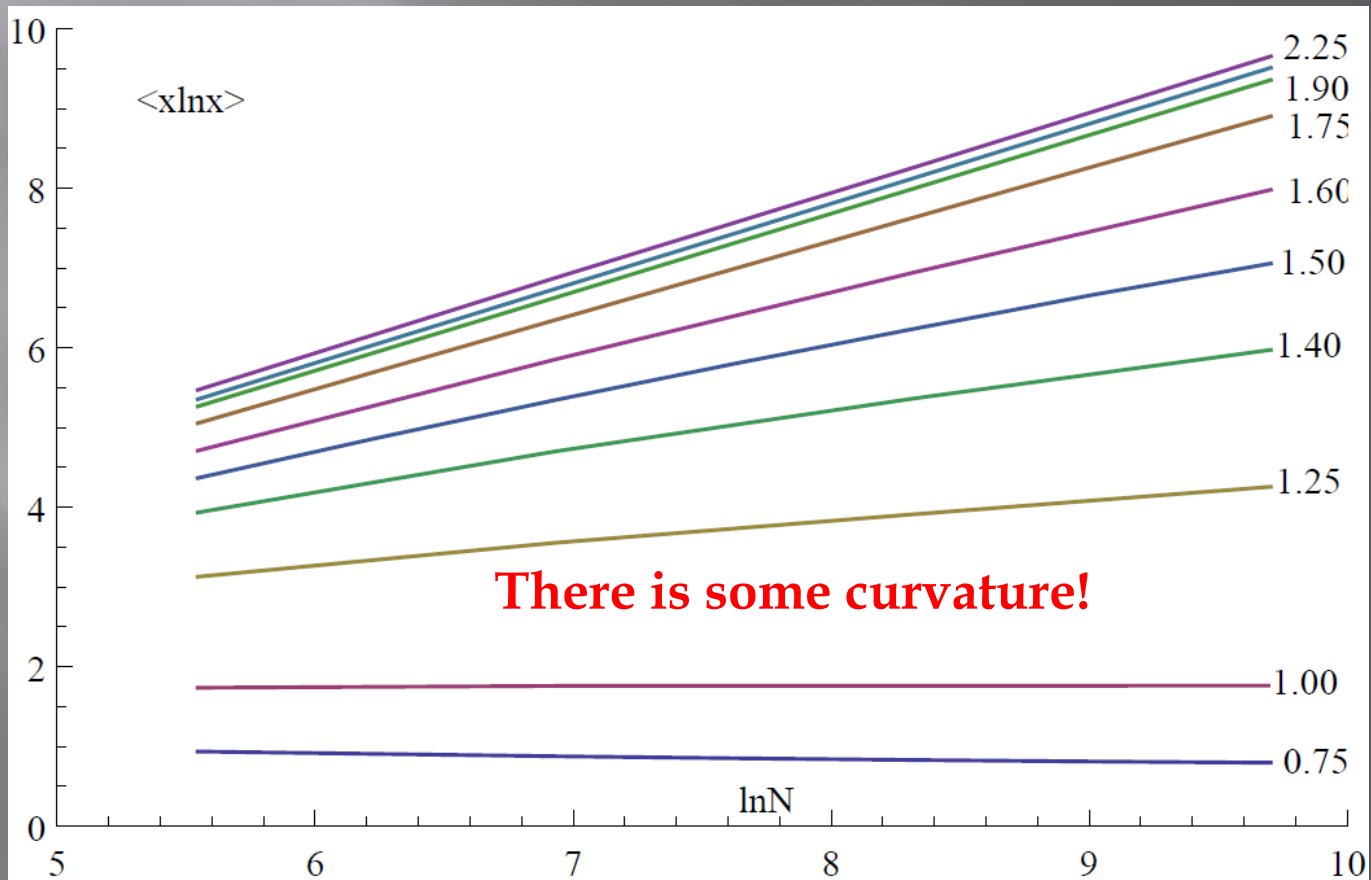


Wave function support set dimension D_1

Shannon entropy:
 $\ln N - \langle x \ln x \rangle$

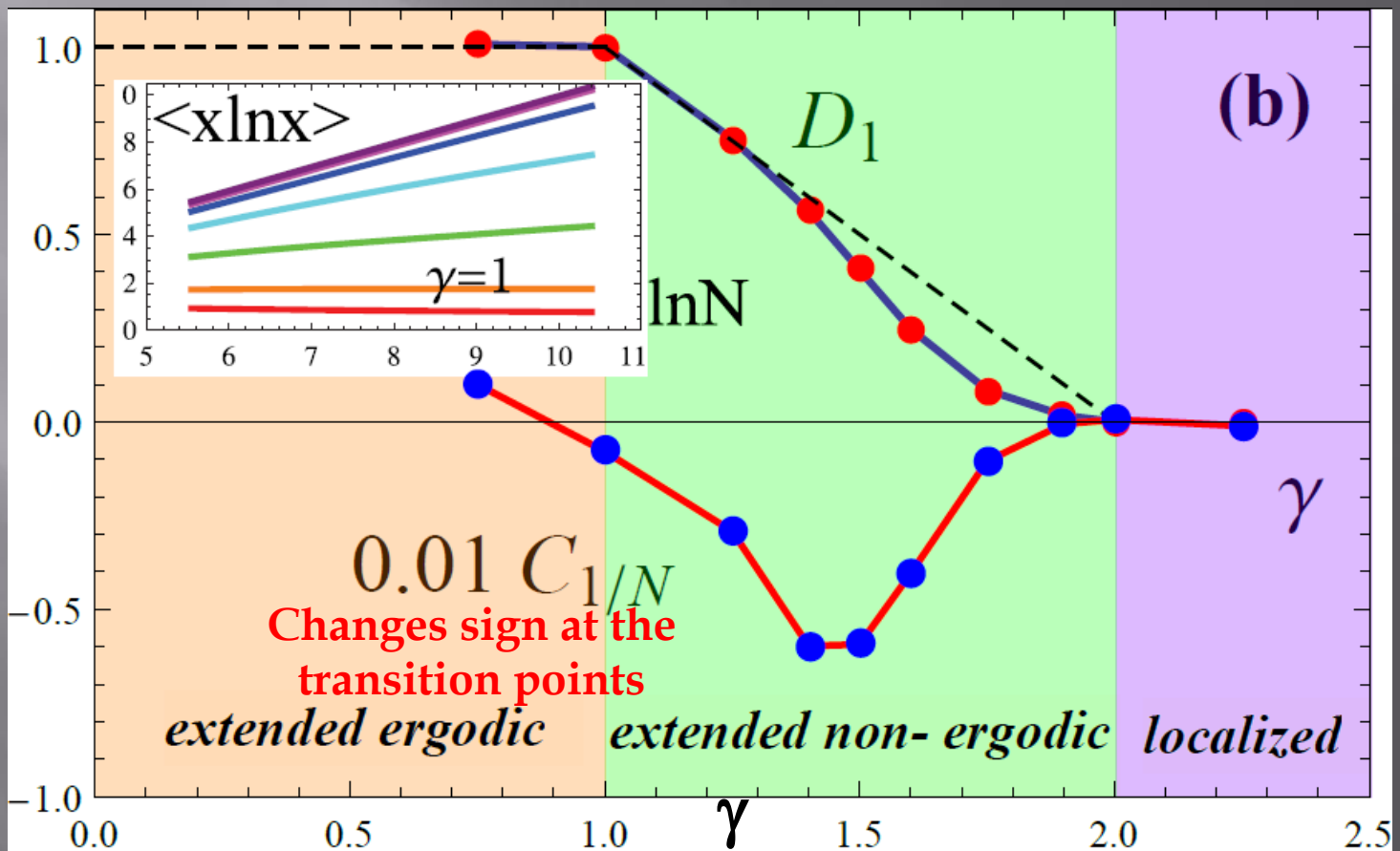
$$\langle x \ln x \rangle = (1 - D_1) \ln N + \text{const}$$

$$x = N |\Psi|^2$$

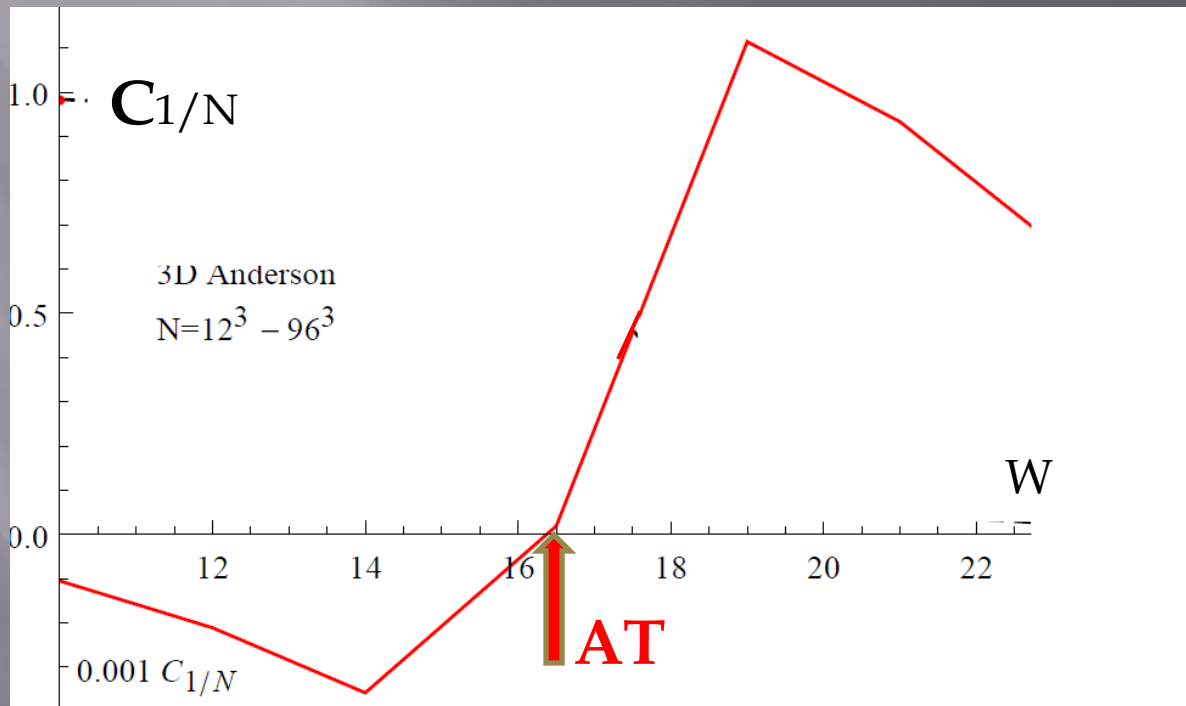


Curvature as signature of transitions

$$\langle x \ln x \rangle = (1 - D_1) \ln N + b + \frac{C_{1/N}}{N}$$



3D Anderson transition



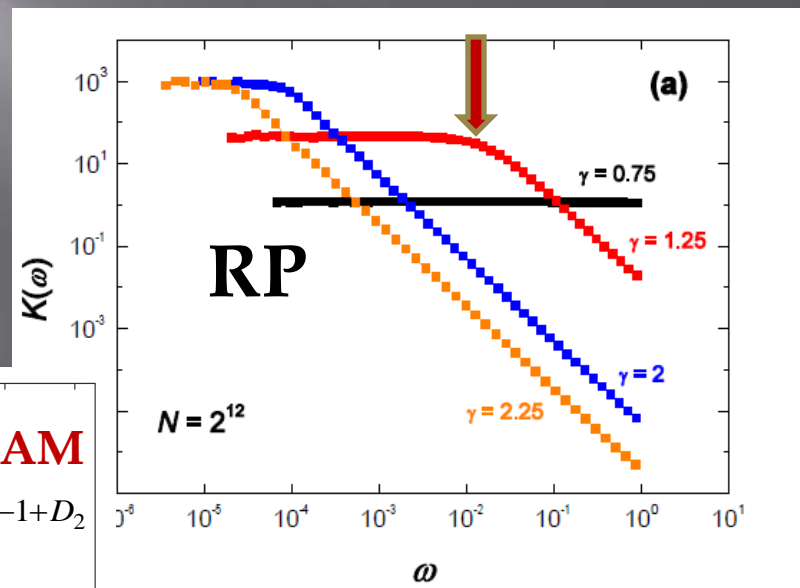
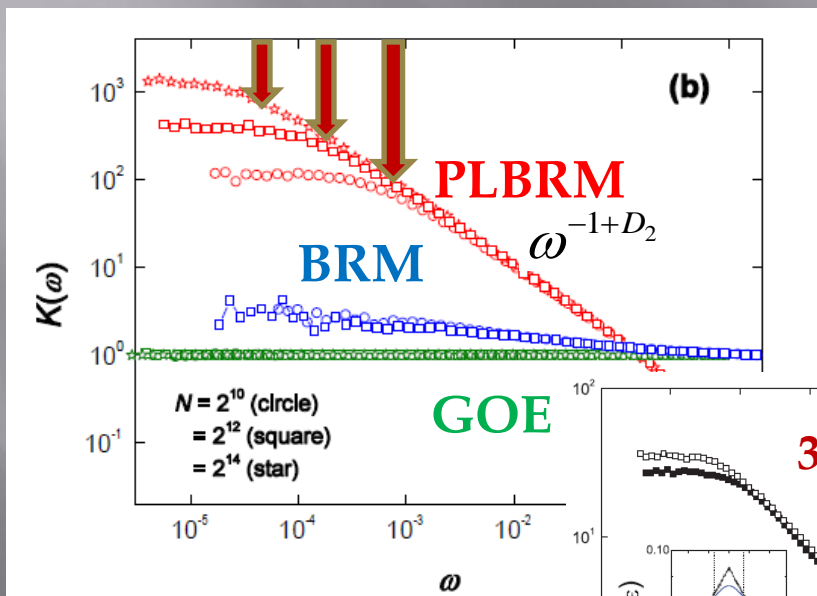
OVERLAP CORRELATION FUNCTION AND THE THOULESS ENERGY

Cuevas & VEK,
PRB, 2007

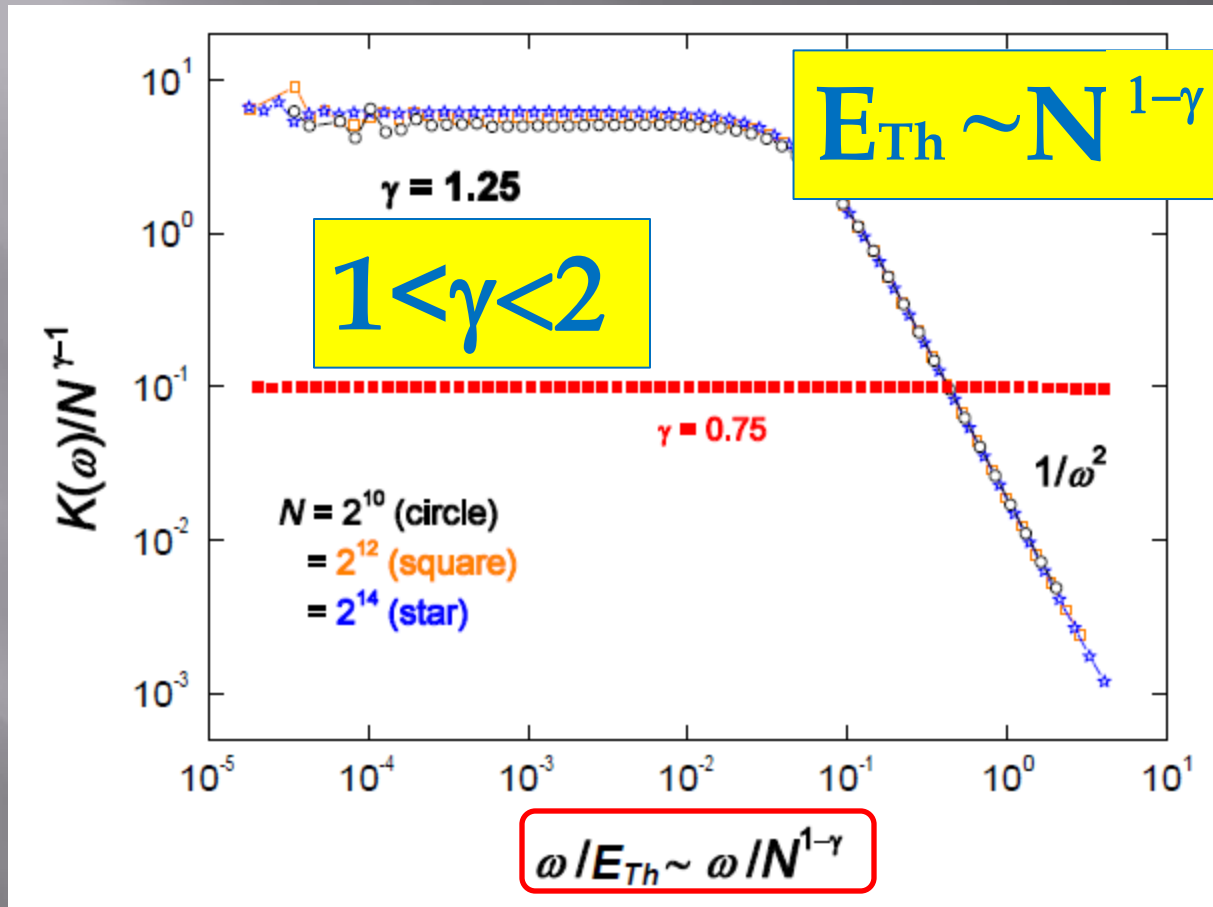
$$K(\omega) = N \sum_r |\Psi_E(r)|^2 |\Psi_{E+\omega}(r)|^2$$

$E_{Th} \sim \delta \sim 1/N$

$E_{th} \gg \delta$



HOW DOES ETH SCALE?



The same scaling of Thouless energy
as for spectral form-factor

THOULESS ENERGY AND FRACTAL DIMENSION

$$D_{q>1/2} = 2 - \gamma$$

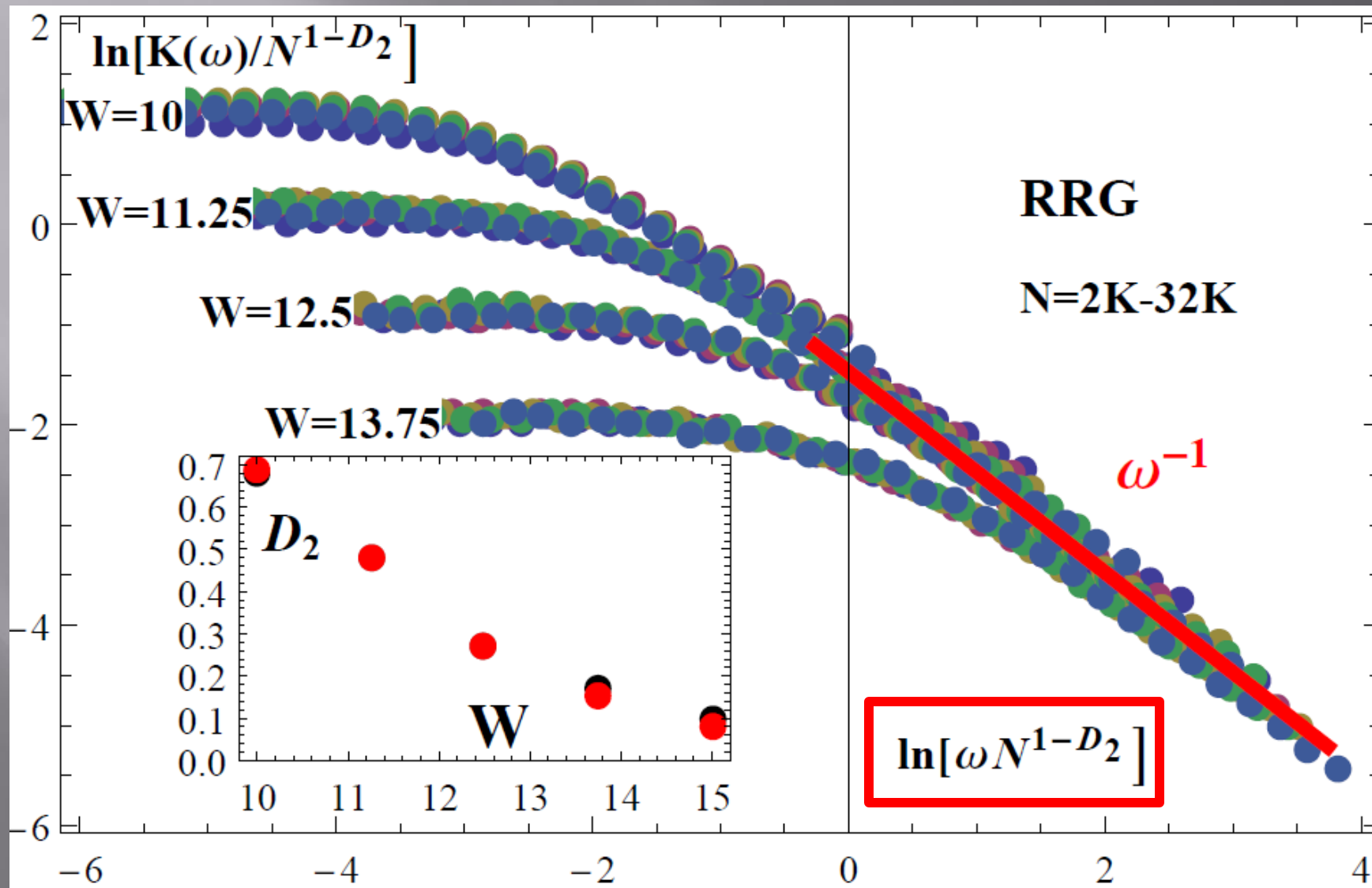
$$E_{Th} = \delta N^D$$

N sites in a sample
 N^{1-D} fractal domains
 N^D sites in each

Each fractal domain is a support of N^D wave functions leaving on it



THOULESS ENERGY ON RRG

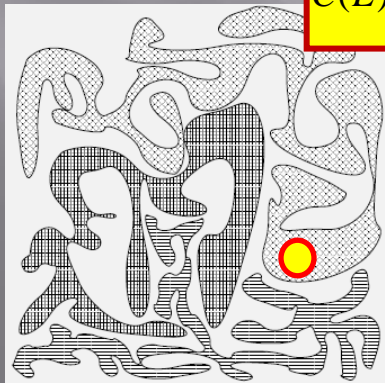
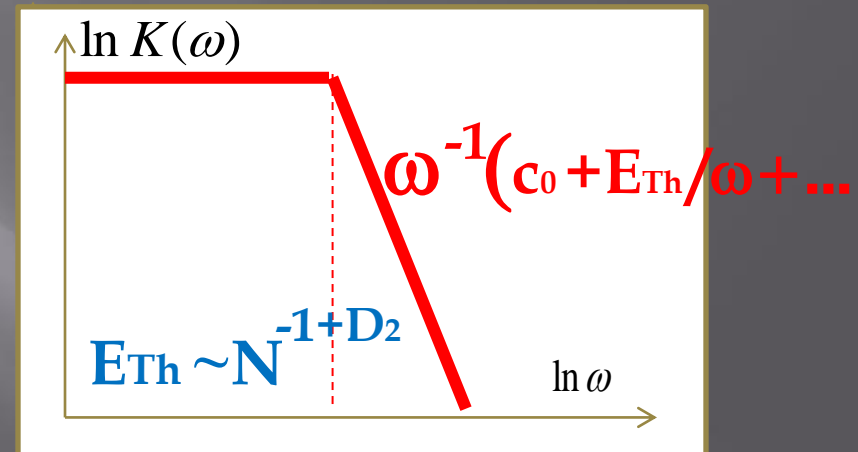
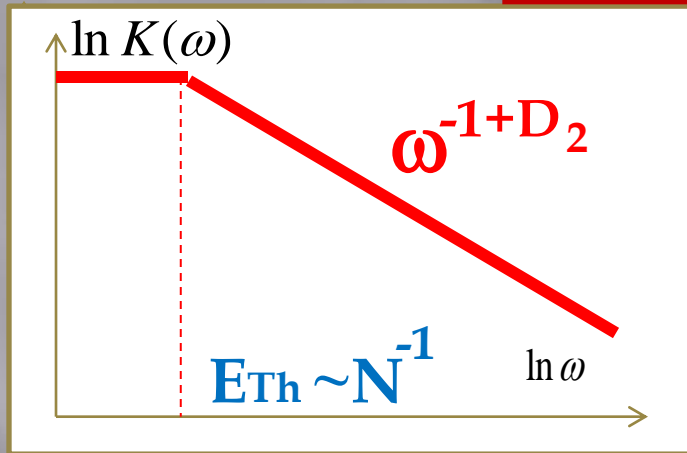


TWO TYPES OF SCALING FOR NON-ERGODIC EXTENDED STATES

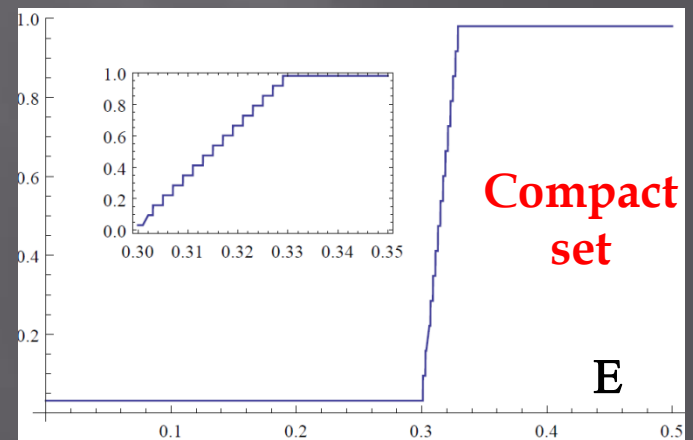
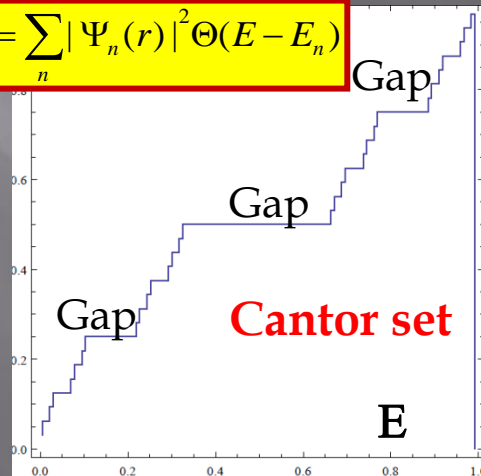
Standard Chalker's scaling:
3D AT point, PRBRM

"Hierarchical" lattices: RRG,

$$K(\omega) = -\partial_{\omega}^2 \langle C(E - \omega/2) C(E + \omega/2) \rangle_E$$



$$C(E) = \sum_n |\Psi_n(r)|^2 \Theta(E - E_n)$$



Conclusions

- ▣ Rosenzweig-Porter RM model shares many of the properties of RRG
- ▣ RP model contains both Anderson and Ergodic transitions
- ▣ They are seen in the rigorous theory of the two-level correlation function
- ▣ Perturbative treatment of the eigenfunction statistics gives the three phases: localized, non-ergodic extended, ergodic extended
- ▣ Numerics confirm existence of three phases
- ▣ $\langle \ln x \rangle$ moments (Shannon entropy) vs $\ln N$ has a curvature which changes sign at the transitions.
- ▣ Two different scalings of Thouless energy

Statistics of $x = N\psi^2$: perturbative arguments

$$(\psi_n(m))^2 = \frac{H_{nm}^2}{(E_n - E_m)^2}.$$

$$\int dt \langle \delta(x - N\psi^2) \rangle e^{itx} = \langle e^{itN\psi^2} \rangle = \frac{1}{N} e^{itN} + \langle e^{itNV^2/\Delta^2} \rangle.$$

$$P(V = H_{nm}, \Delta = H_{nn} - H_{mm}) = \frac{1}{2\pi\sqrt{2\sigma}} e^{-\frac{\Delta^2}{4}} e^{-\frac{V^2}{2\sigma}}$$

$$P_{reg}(x) = \frac{1}{\sqrt{2\pi}} \frac{\sqrt{N\sigma}}{x^{3/2}}$$

$$+ c\delta(x - N)$$

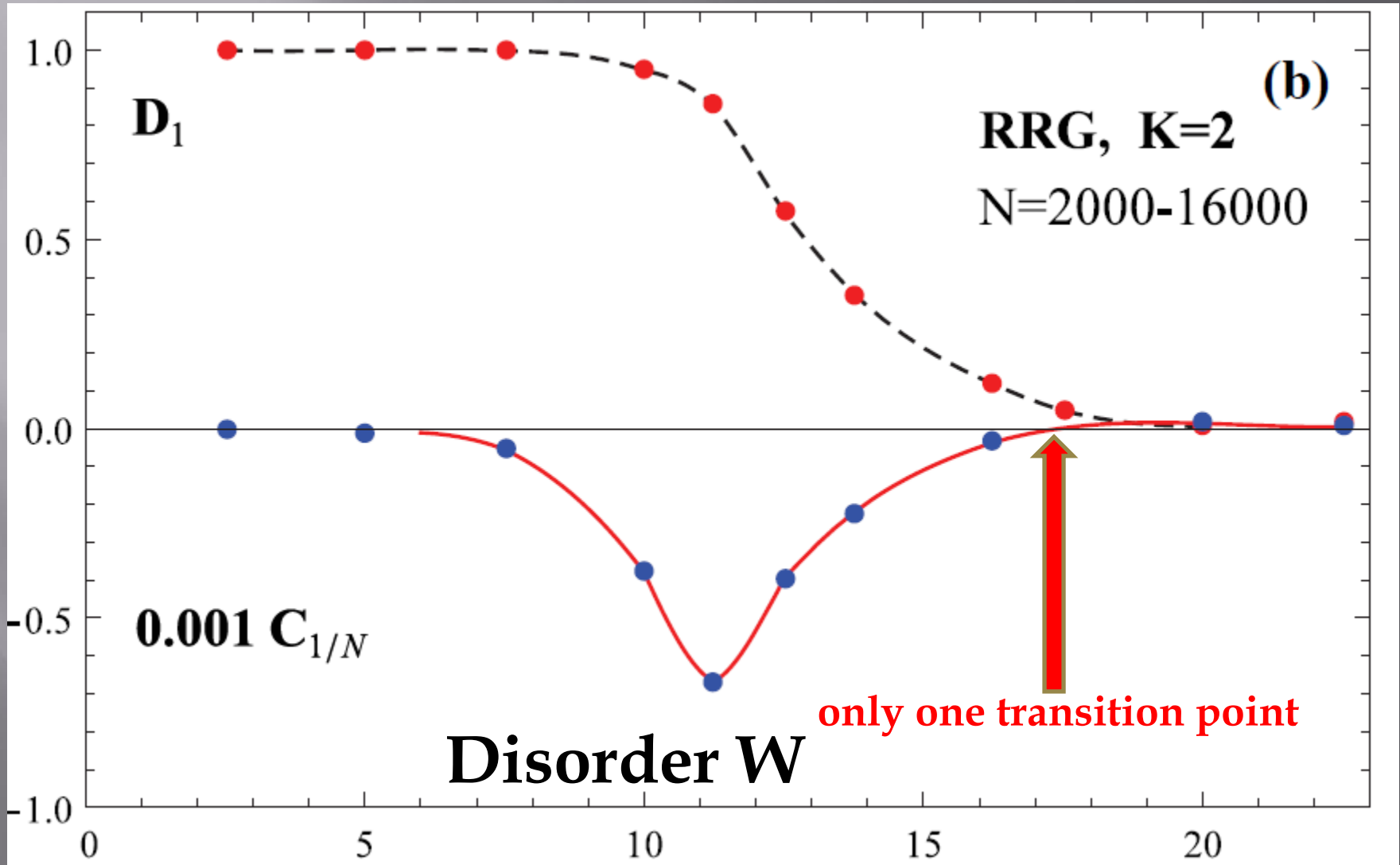
$$\int_0^\infty P(x) dx = 1$$

$$\int_0^\infty x P(x) dx = 1.$$

$$x_{\min} = N^{-(\gamma-1)}$$

$$x_{\max} = \begin{cases} N^{(\gamma-1)}, & \gamma < 2 \\ N, & \gamma > 2. \end{cases}$$

D1 and curvature on the RRG



Generalization of Kunz & Shapiro

$$K_1(t, t') = \frac{1}{N\tau\tau'} \oint_{\Gamma_R} \frac{dz}{2\pi i} \oint_{\Gamma_R} \frac{dz'}{2\pi i} e^{i(tz+t'z')} (g(z, z')^N - \rho(z/\tau)^N \rho(z'/\tau')^N)$$

$$K_2(t, t') = \frac{1}{N} \oint_{\Gamma_R} \frac{dz}{2\pi i} \oint_{\Gamma_R} \frac{dz'}{2\pi i} e^{i(tz+t'z')} \frac{g(z, z')^N}{(z' - z - \tau)(z' - z + \tau')} .$$

$$g(z, z') = 1 + \tau\alpha(z) + \tau'\alpha(z') + \frac{\tau\tau'}{z-z'} [\alpha(z') - \alpha(z)]$$

$$\rho = 1 + \tau\alpha(z)$$

$$\alpha(z) = \left\langle \frac{1}{z-a} \right\rangle$$

Depends on the distribution function of the diagonal matrix
 $A = \text{diag}\{a\}$

$$(t, t') = i\sigma(\tau, \tau') \rightarrow \frac{T}{2} \boxed{\pm N^{\gamma-1} s}$$

$$(z, z') \rightarrow x \pm \boxed{\frac{y}{2N}}$$

Rescaling and $N \rightarrow \infty$

Generalization of Kunz & Shapiro

$$K_1(t, t') = \frac{1}{N\tau\tau'} \oint_{\Gamma_R} \frac{dz}{2\pi i} \oint_{\Gamma_R} \frac{dz'}{2\pi i} e^{i(tz+t'z')} \left(\gamma(z, z')^N - \rho(z/\tau)^N \rho(z'/\tau')^N \right)$$

$$K_2(t, t') = \frac{1}{N} \oint_{\Gamma_R} \frac{dz}{2\pi i} \oint_{\Gamma_R} \frac{dz'}{2\pi i} e^{i(tz+t'z')} \frac{q(z, z')^N}{(z' - z - \tau)(z' - z + \tau')}$$

cancels out

$\propto N^{\gamma-2}$

tend to a finite limit
independent of γ

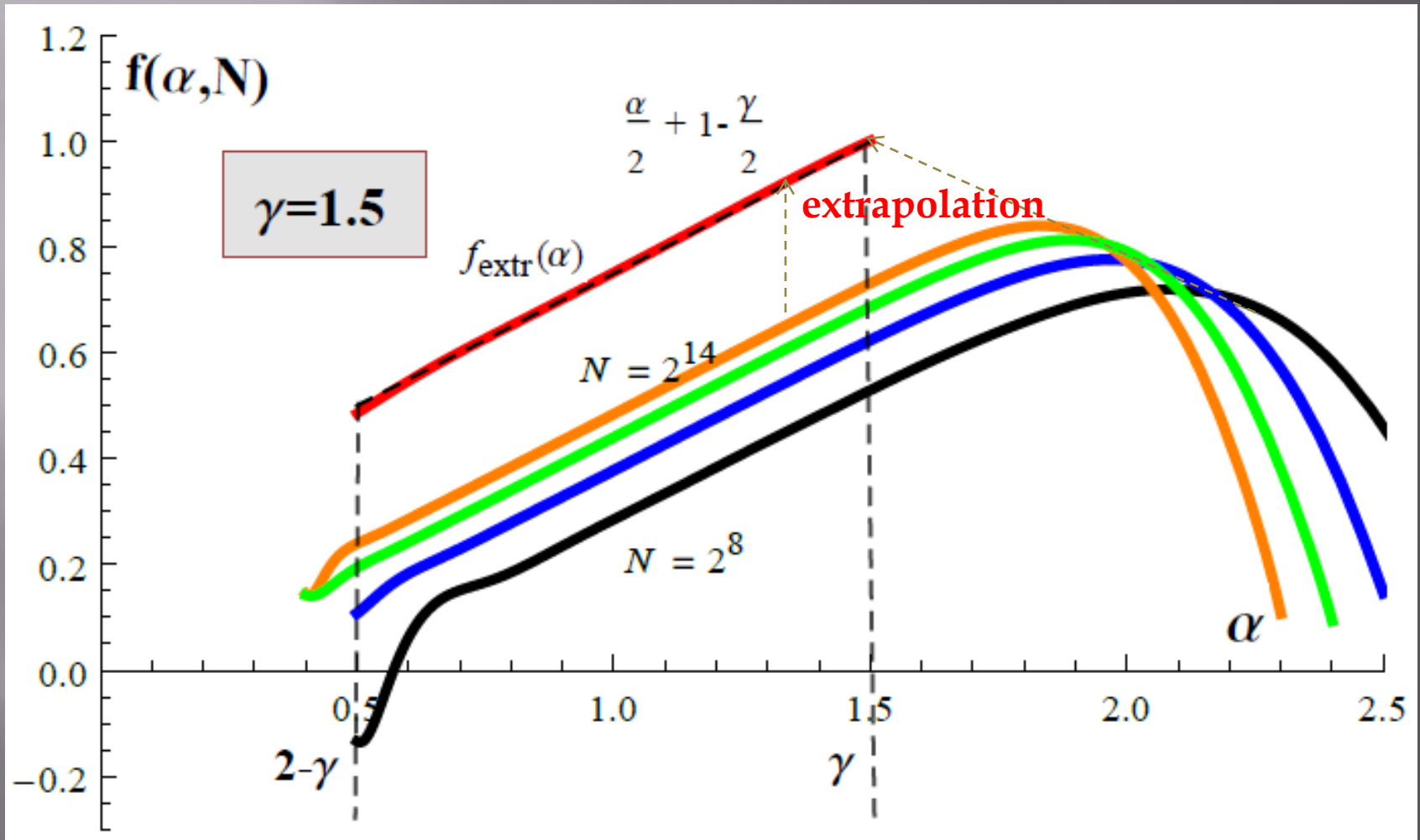
$\gamma > 2$: Poisson due to infinitely fast oscillations

$\gamma = 2$: regular oscillations

$1 < \gamma < 2$: no oscillations

$0 < \gamma < 1$: rescaling does not work

Numerics for $f(\alpha, N)$



Numerics for $f(\alpha, N)$

