

The Superconductor-Insulator Transition Far From Equilibrium

Dan Shahar
The Weizmann Institute

KITP, 2015



Maoz Ovadia (Harvard)

Benjamin Sacépé (CNRS Grenoble)

Idan Tamir

Adam Doron

Tal Levinson

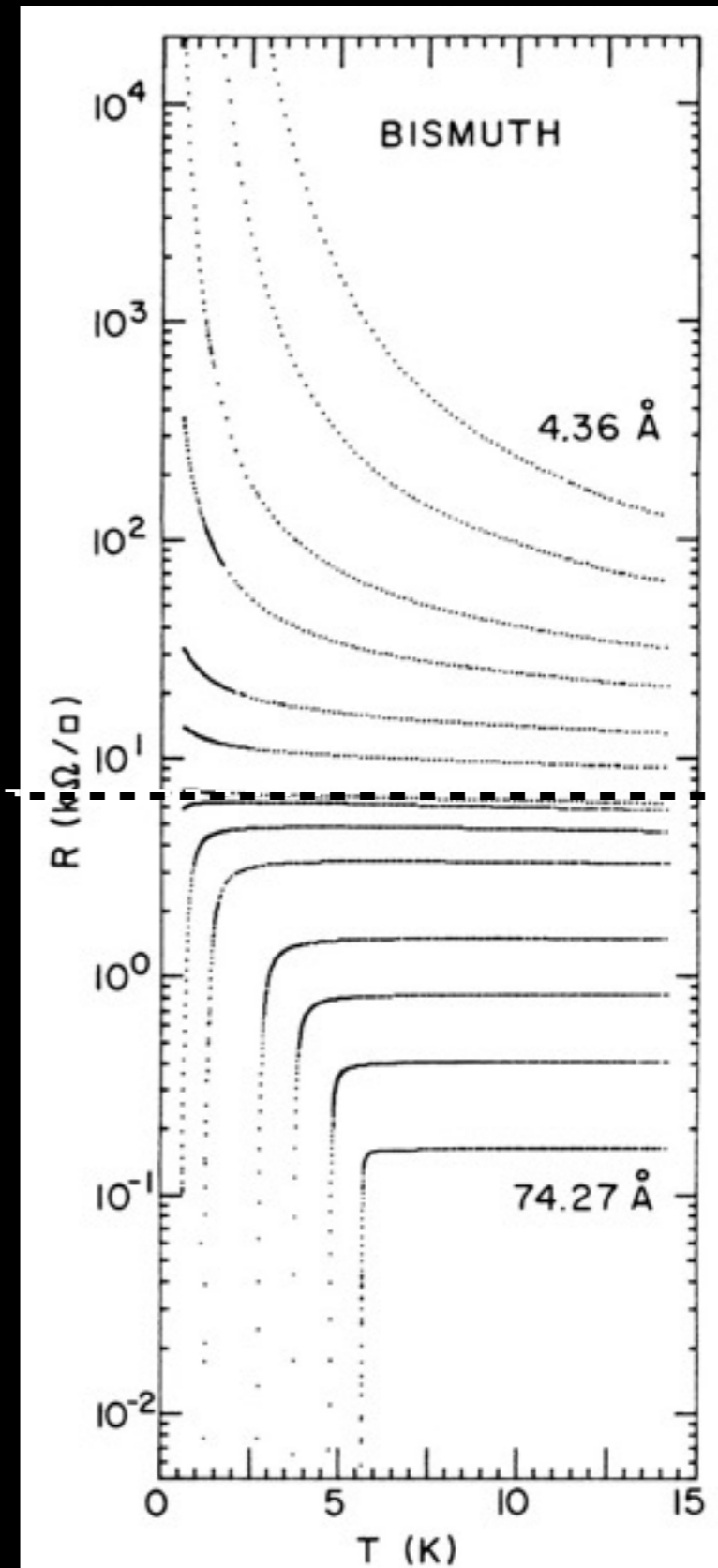
The Superconductor-Insulator Transition Far From Equilibrium

The Superconductor-Insulator Transition

M. Strongin, et. al., Phys.
Rev. B1, 1078 (1970).

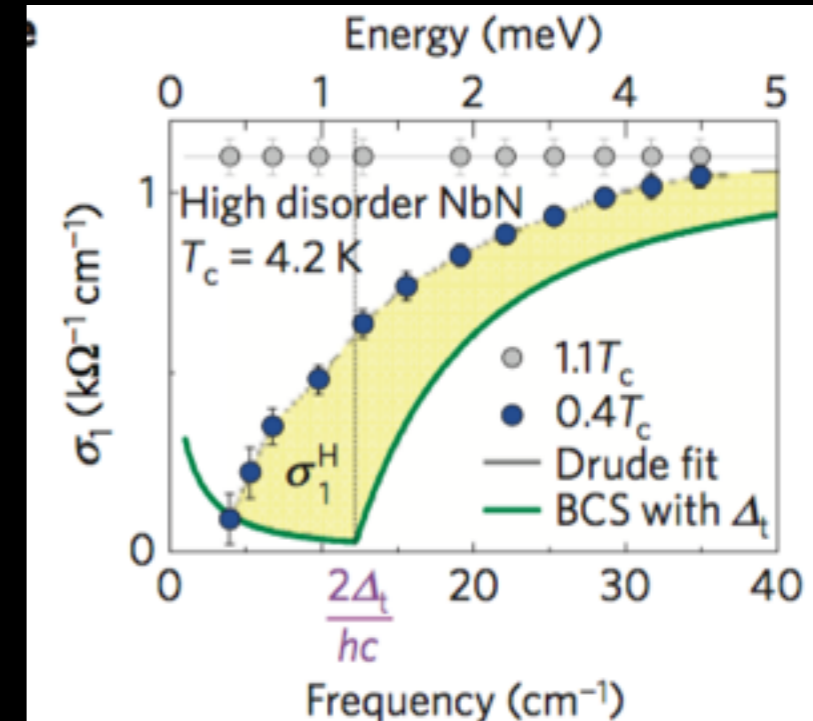
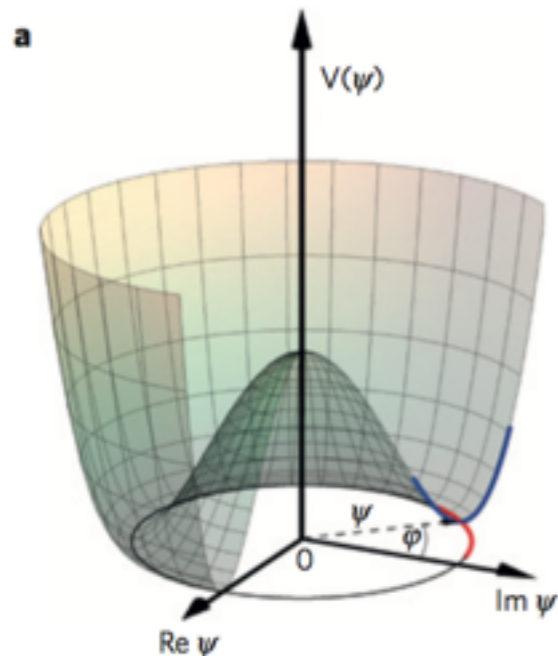
Haviland, Liu and Goldman,
Phys. Rev. Lett. 62, 2180 (1989)...

Reviews: Finkl'stein ('94)
Markovic and Goldman ('98)
Gantmakher and Dolgopolov ('10)



The Higgs mode in disordered superconductors close to a quantum phase transition

Daniel Sherman^{1,2†}, Uwe S. Pracht², Boris Gorshunov^{2,3,4}, Shachaf Poran¹, John Jesudasan⁵, Madhavi Chand⁵, Pratap Raychaudhuri⁵, Mason Swanson⁶, Nandini Trivedi⁶, Assa Auerbach⁷, Marc Scheffler², Aviad Frydman^{1*} and Martin Dressel²



A quantum phase transition

Strictly $T=0$.

driven by:

Thickness

Magnetic field

Disorder

Carrier density

Pressure

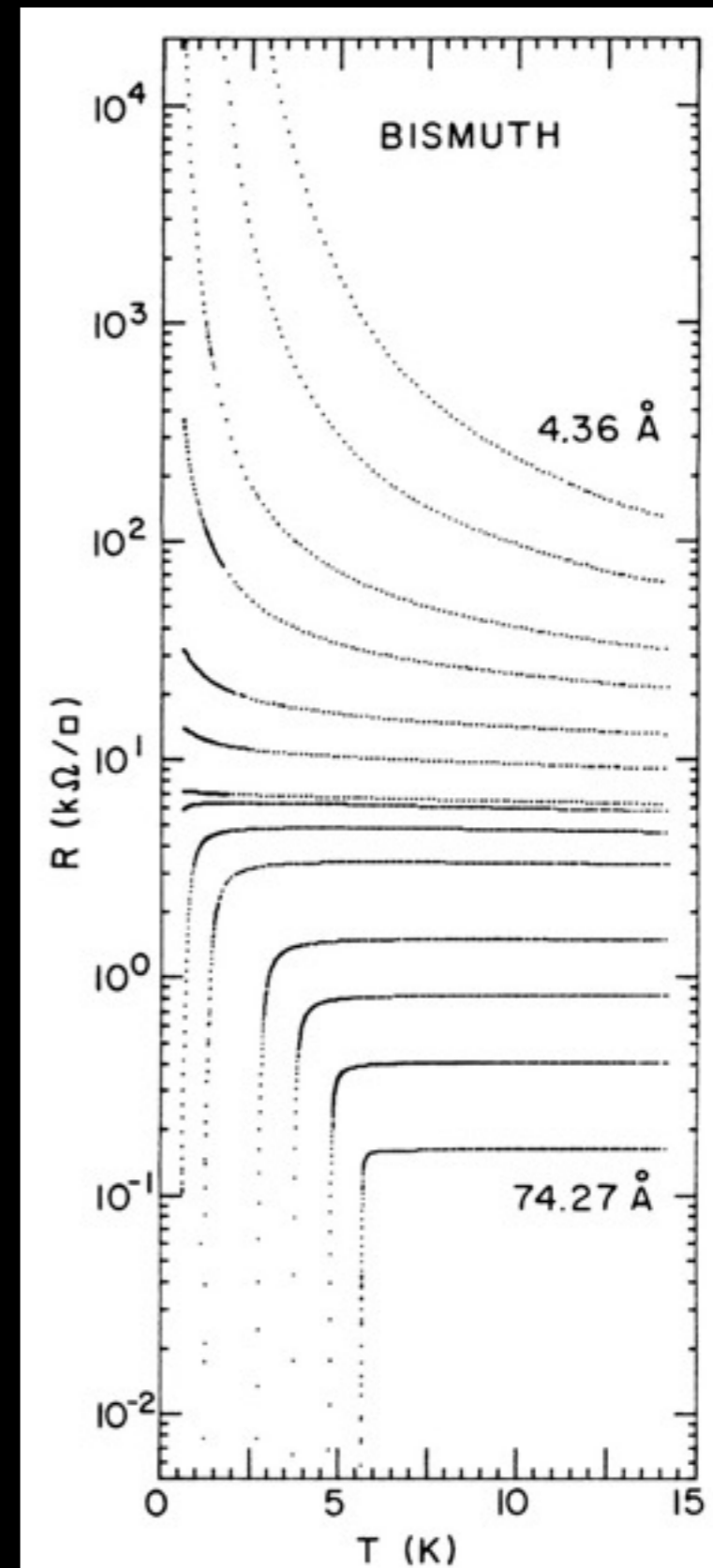
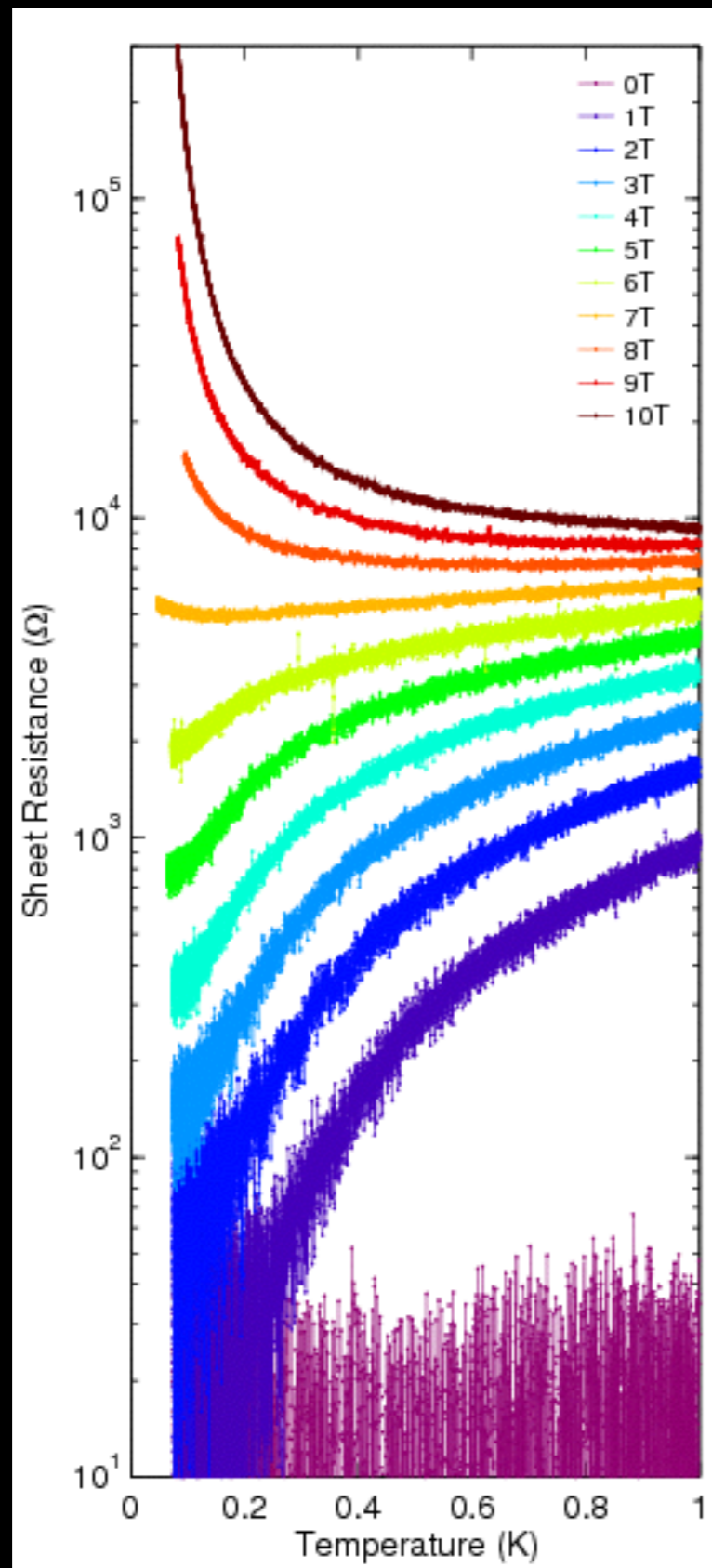
Chemical composition

Structural composition

...

Magnetic-field dependence

Magnetic-field dependence

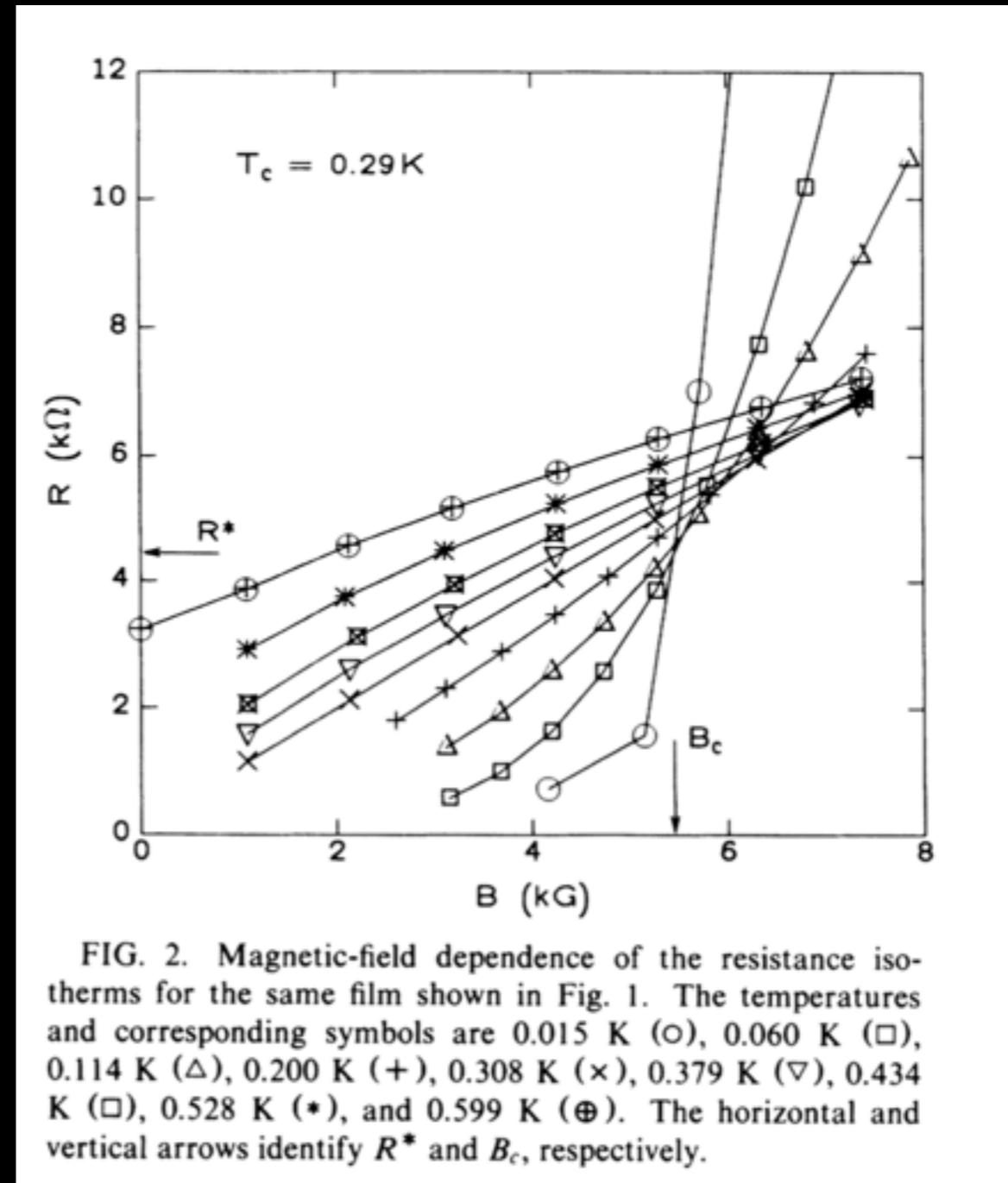


A quantum phase transition

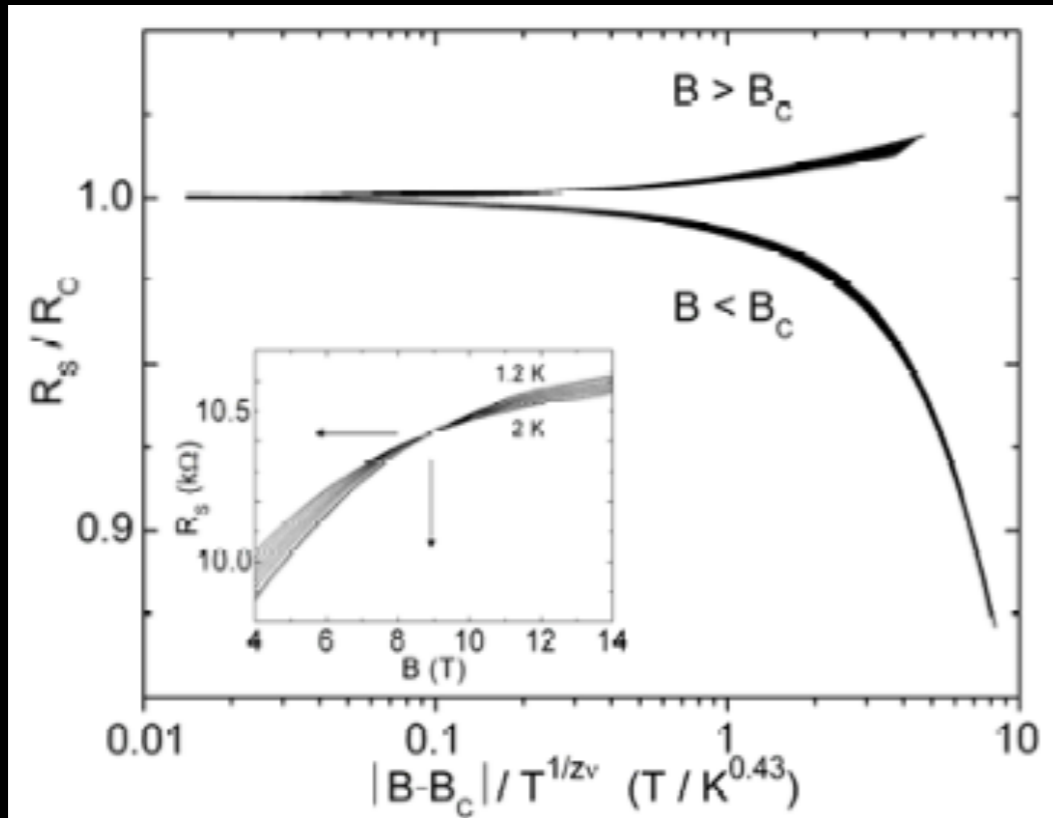
How do we know?

$$z\nu = 1.26-1.31$$

Hebard and Paalanen 1990



Scaling analysis



Hebard and Paalanen 1990

Yazdany and Kapitulnik 1995

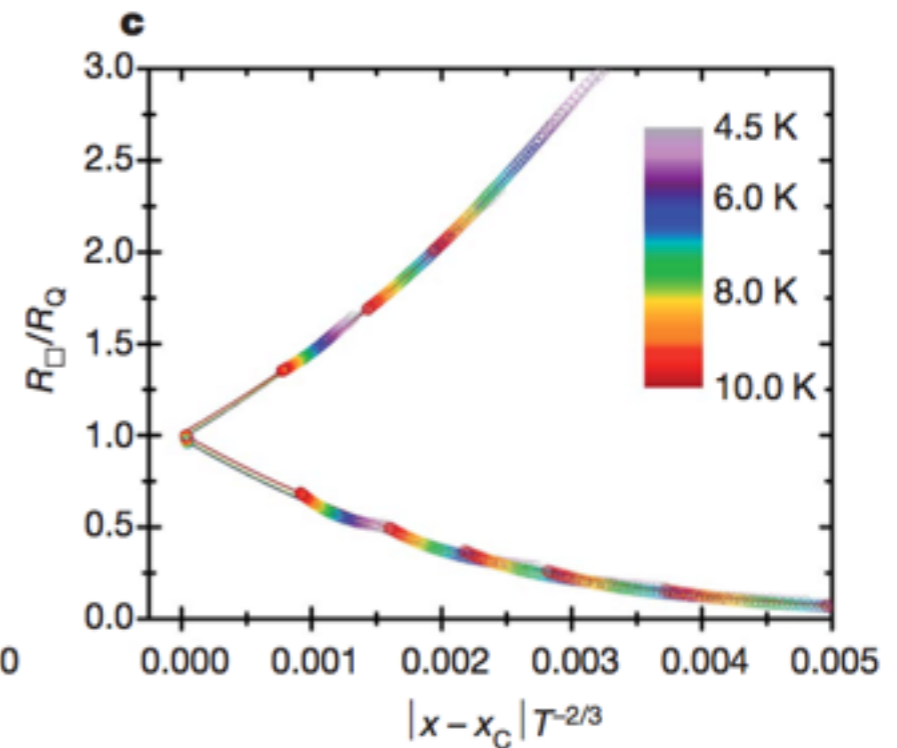
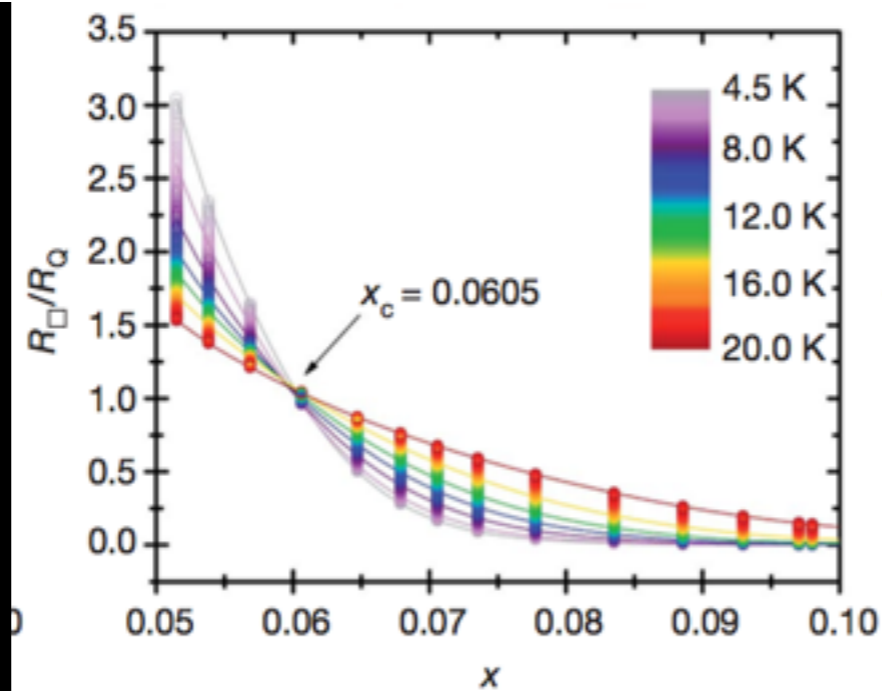
...

Bollinger et al, Nature 2011

Schneider et al,

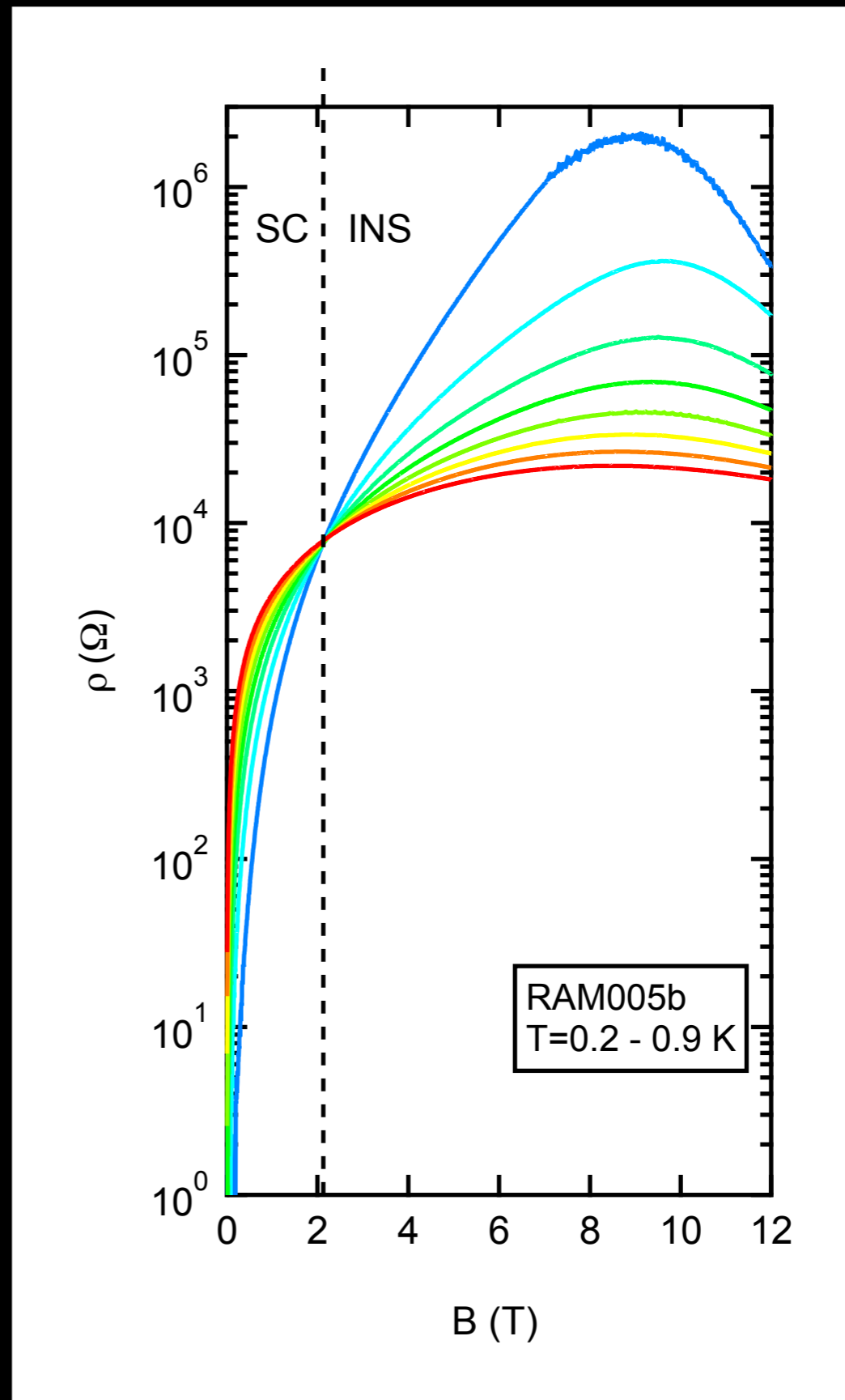
Phys. Rev. Lett. 108, 257003 (2012)

...



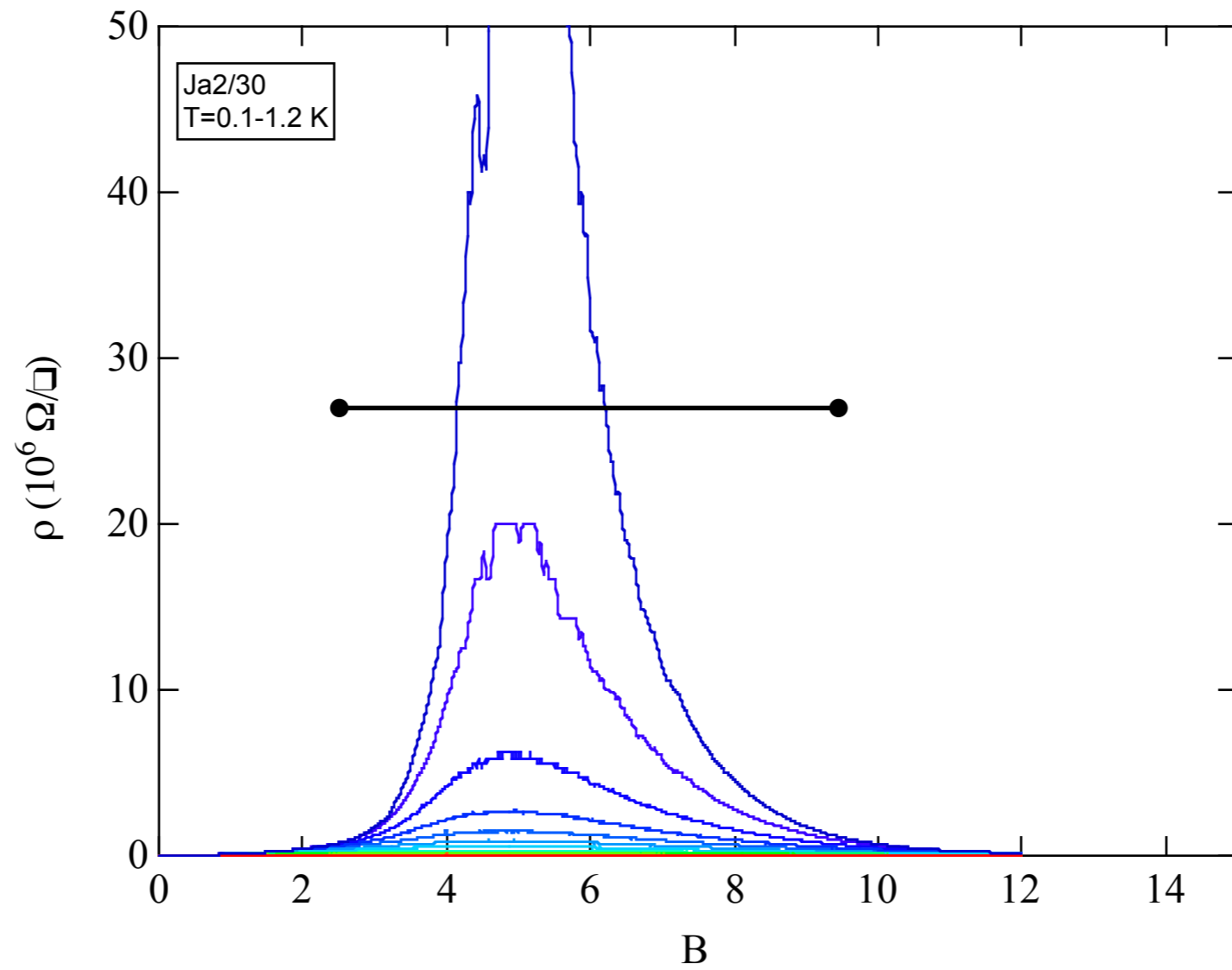
Magnetic-field dependence

a:InO

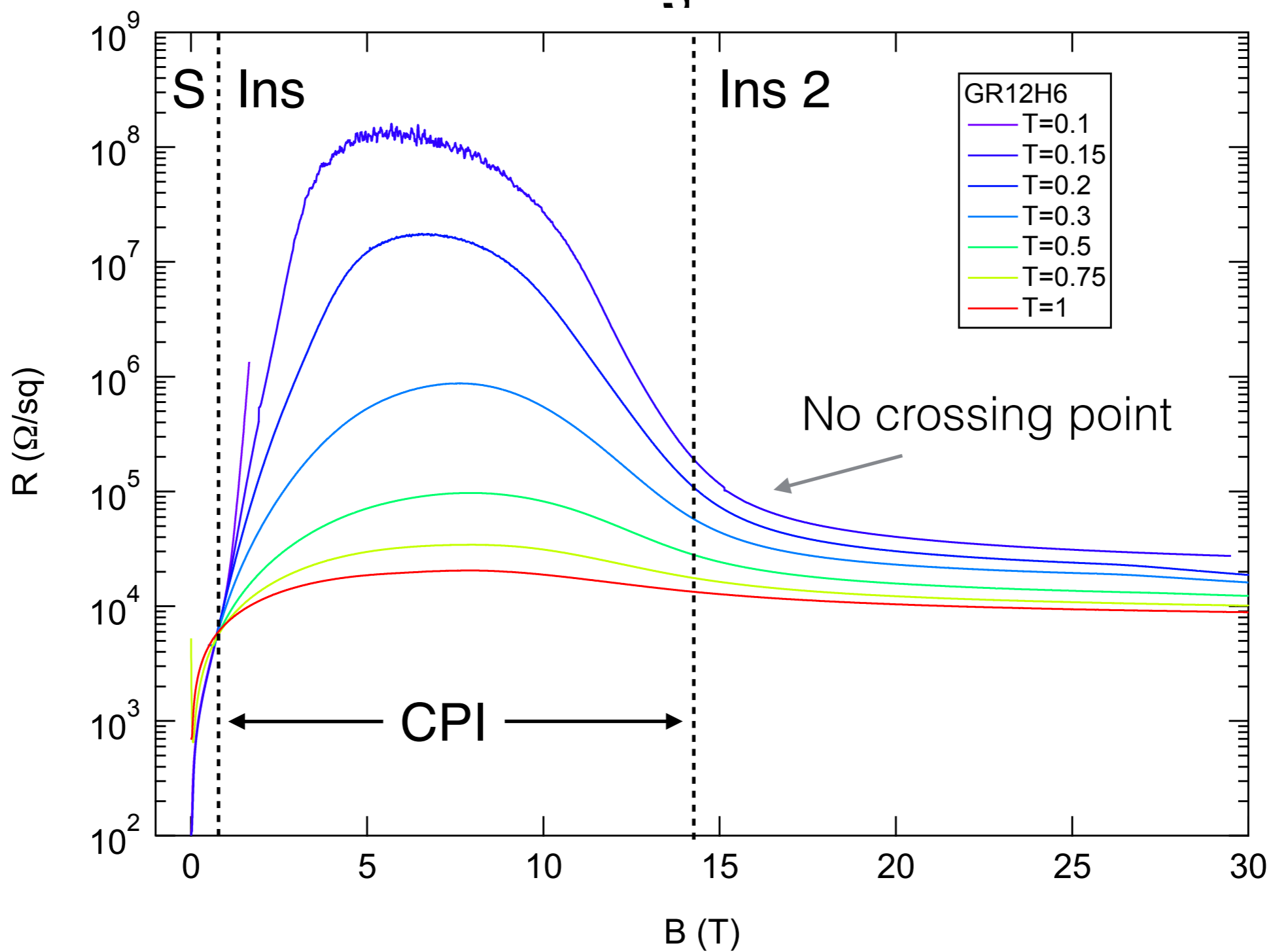


Narrow range in B

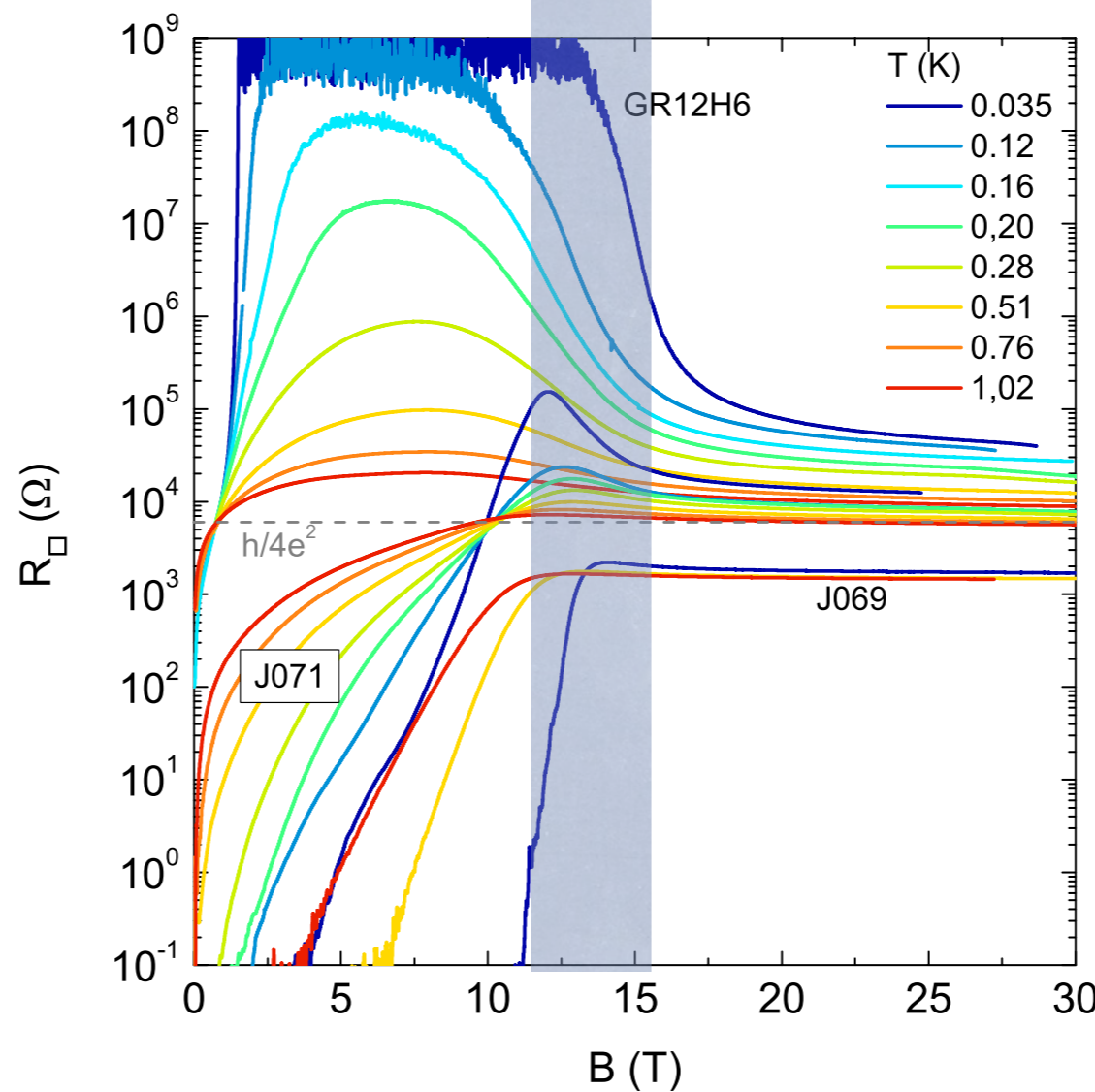
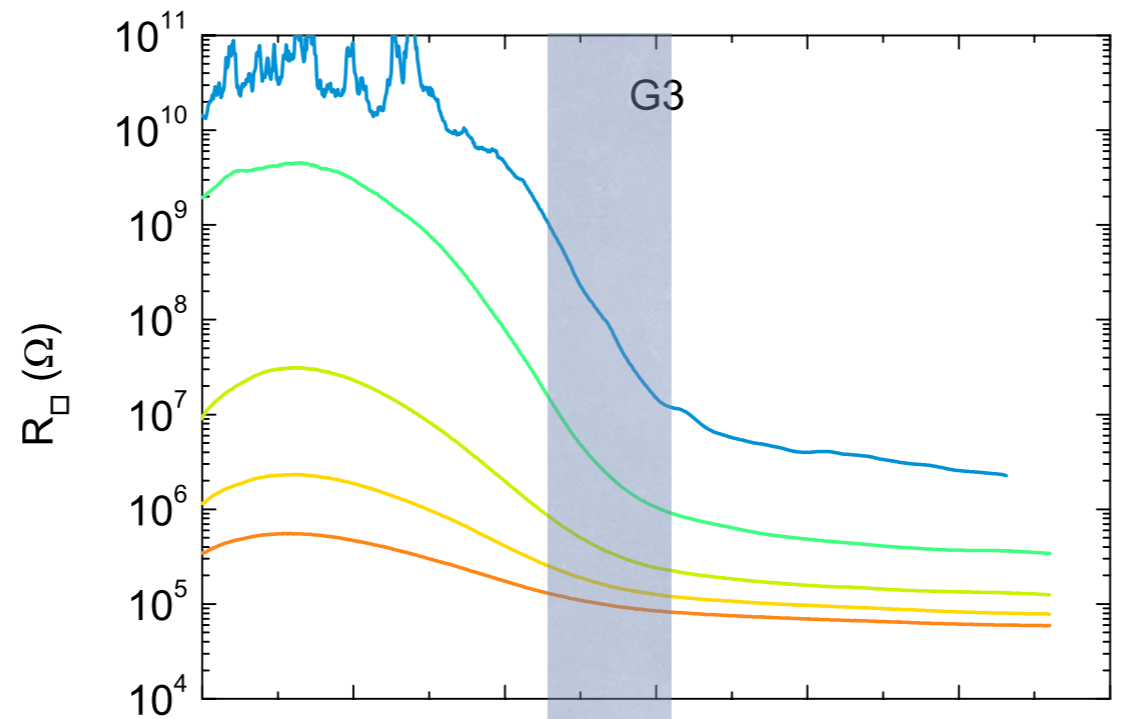
The insulating peak



The insulating peak



a:InO disorder range

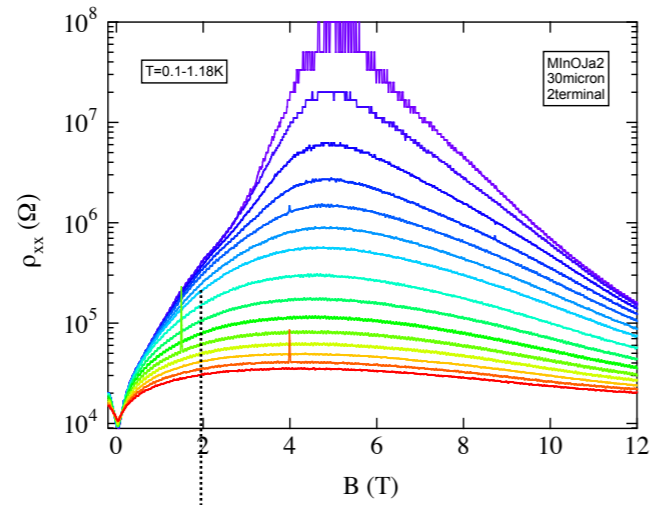


Disorder ↑

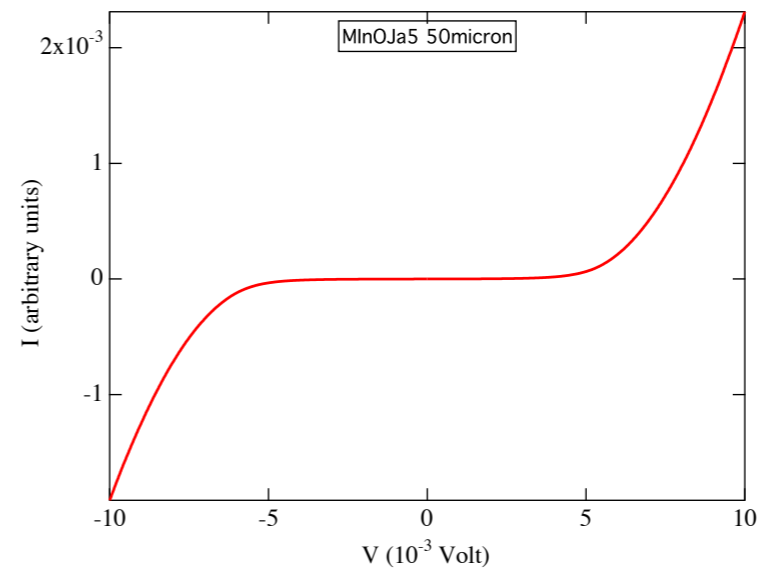
The Insulator at $T < 0.2$ K:

Current-voltage characteristics

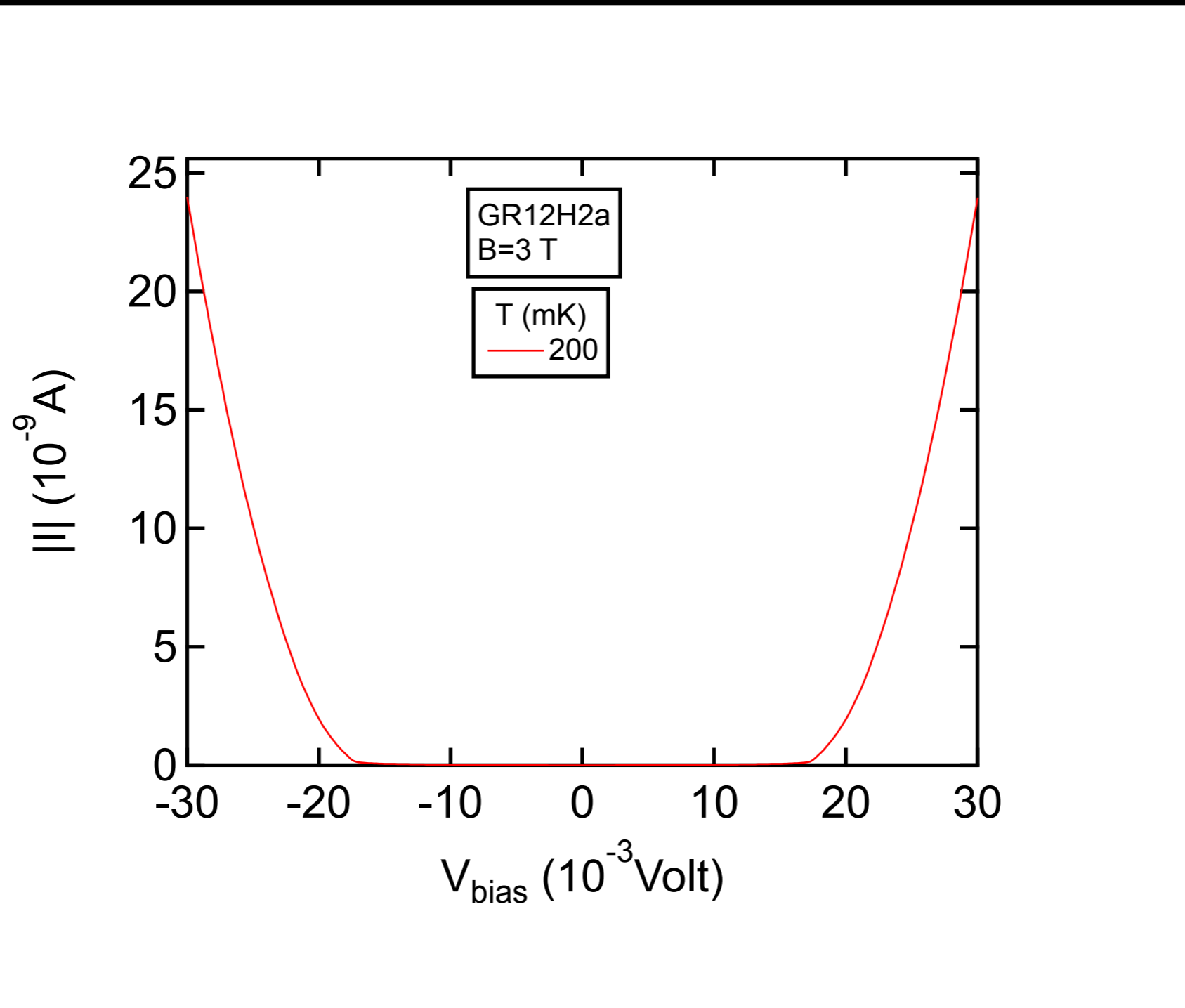
Current-voltage characteristics



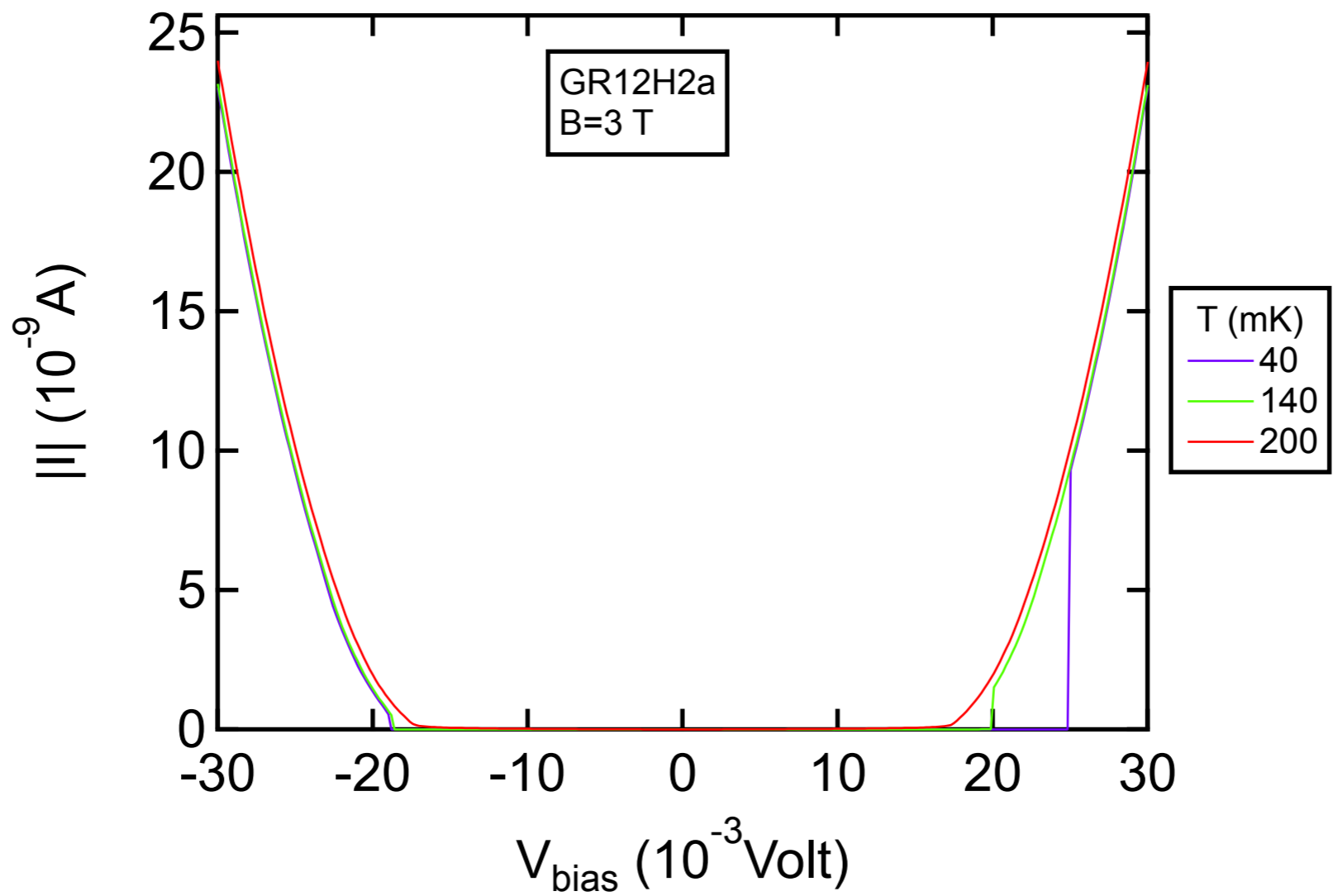
Sambandamurthy et. al,
PRL **94**, 017003 (2005).



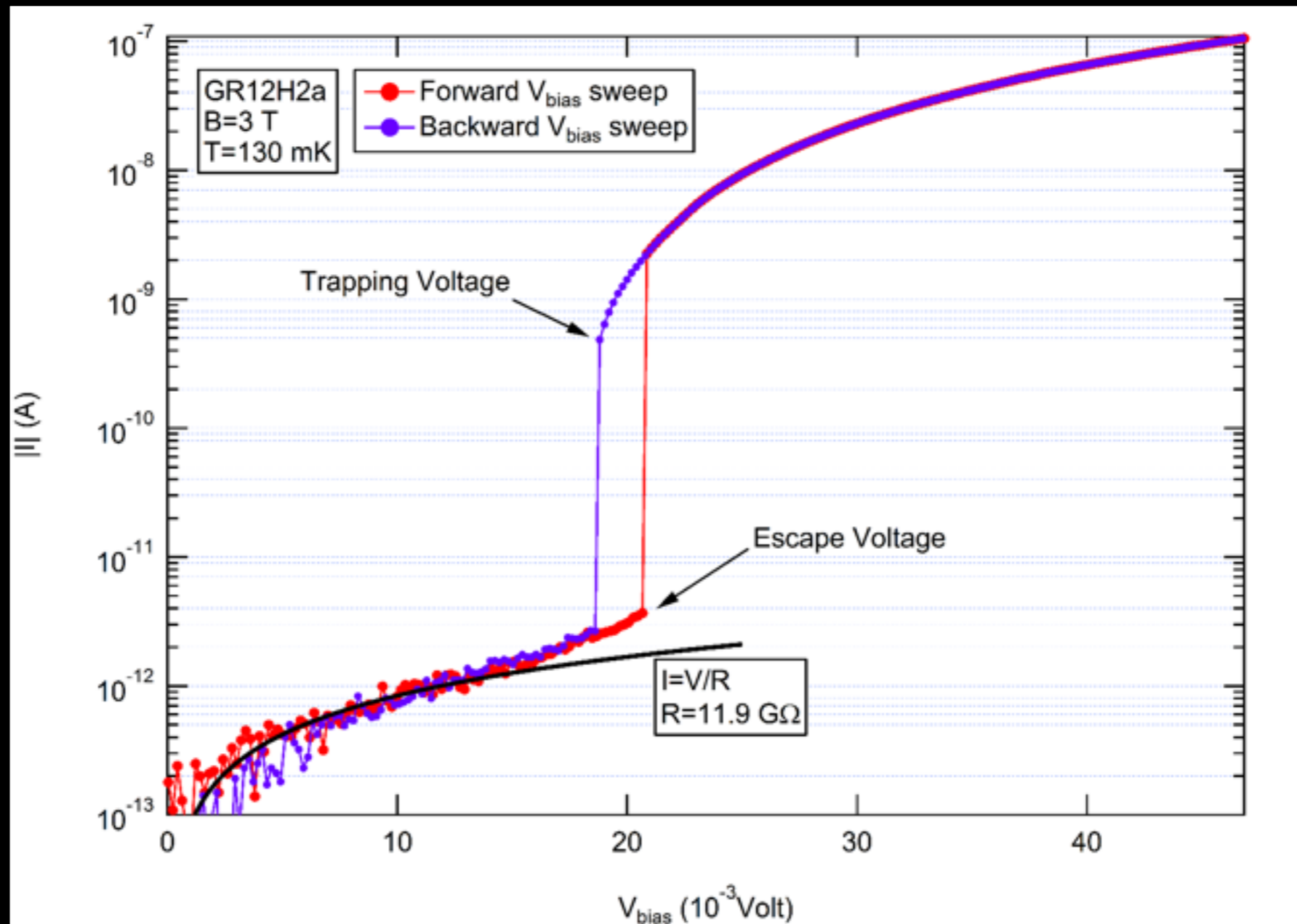
Highly non-linear



Transition to abrupt I-V



Hysteresis

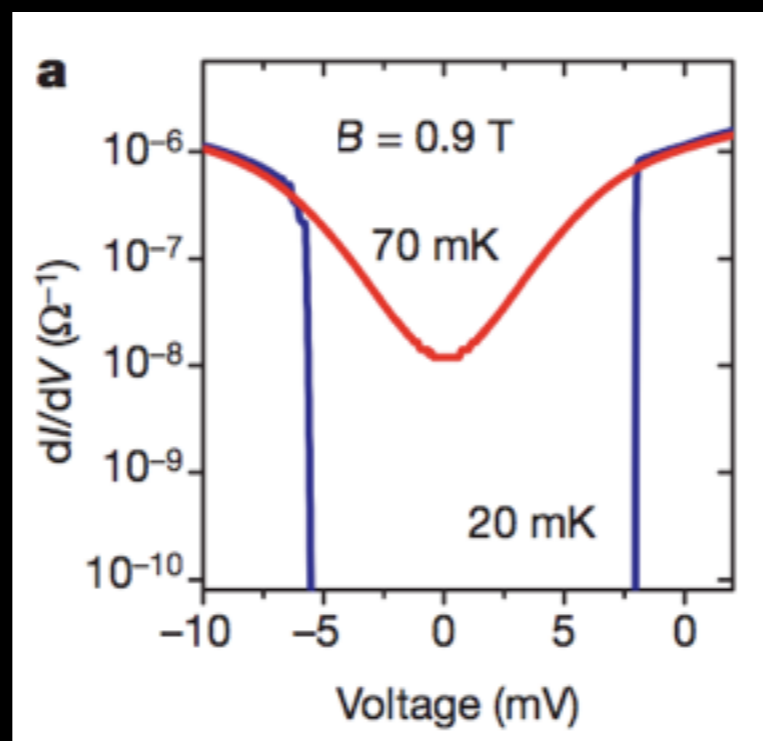


Experimental Evidence for a Collective Insulating State in Two-Dimensional SuperconductorsG. Sambandamurthy,¹ L. W. Engel,² A. Johansson,¹ E. Peled,¹ and D. Shahar¹¹*Department of Condensed Matter Physics, Weizmann Institute of Science, Rehovot 76100, Israel,*²*National High Magnetic Field Laboratory, Florida State University, Tallahassee, Florida 32306, USA*

(Received 18 March 2004; published 12 January 2005)

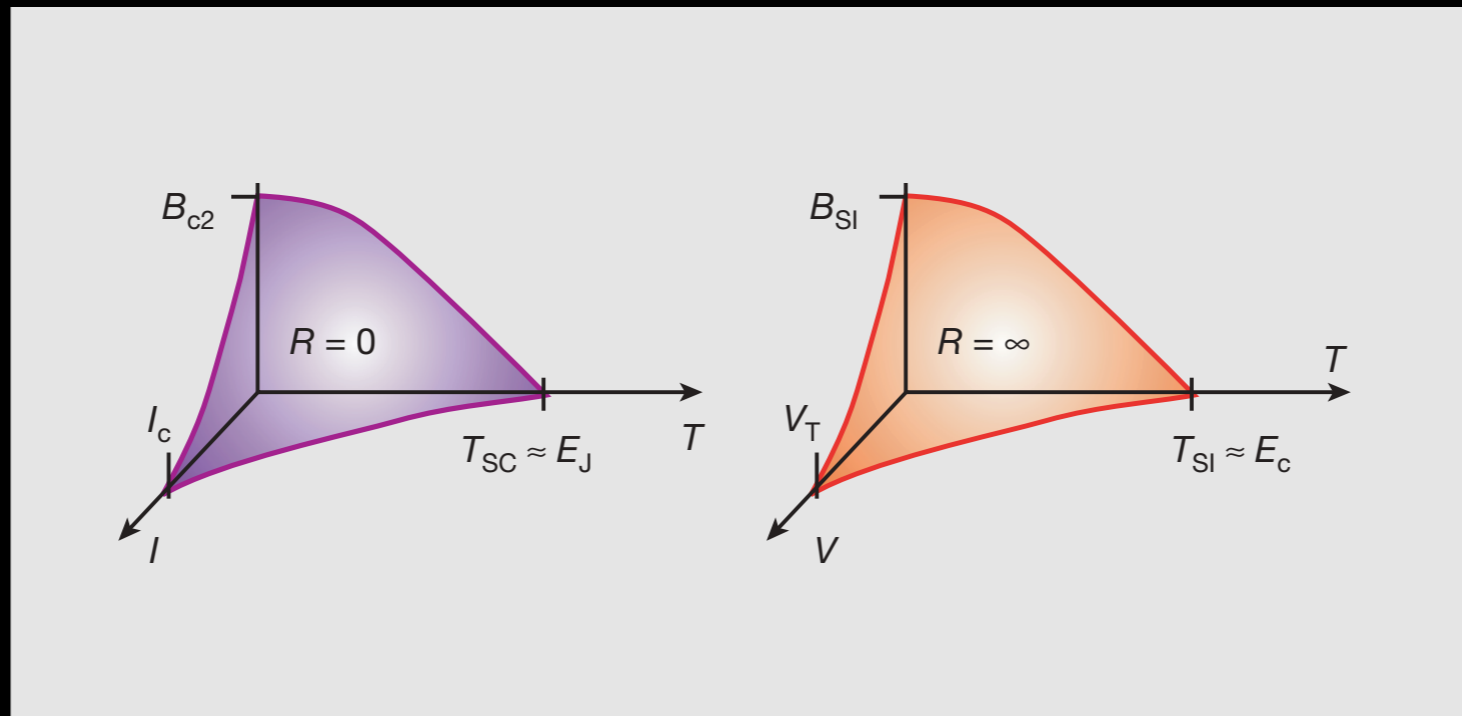
TiN

Baturina et al, 2007



Superinsulator and quantum synchronization

Vinokur et al., NATURE | Vol 452 | 3 April 2008





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Metal–insulator transition in a weakly interacting many-electron system with localized single-particle states

D.M. Basko^{a,b,*}, I.L. Aleiner^b, B.L. Altshuler^{a,b,c}

^a *Department of Physics, Princeton University, Princeton, NJ 08544, USA*

^b *Physics Department, Columbia University, New York, NY 10027, USA*

^c *NEC-Laboratories America, 4 Independence Way, Princeton, NJ 085540, USA*

Received 14 August 2005; accepted 30 November 2005

Available online 23 January 2006

$$\hat{\mathcal{G}} = \begin{bmatrix} \mathcal{G}_l^R(\boldsymbol{\rho}) & \mathcal{G}_l^K(\boldsymbol{\rho}) \\ 0 & \mathcal{G}_l^A(\boldsymbol{\rho}) \end{bmatrix}_K; \quad \hat{\Sigma} = \begin{bmatrix} \Sigma_l^R(\boldsymbol{\rho}) & \Sigma_l^K(\boldsymbol{\rho}) \\ 0 & \Sigma_l^A(\boldsymbol{\rho}) \end{bmatrix}_K; \quad (54)$$

$$\hat{\tau}^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}_K; \quad \hat{\tau}^2 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}_K.$$

A

$$\overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \overleftarrow{\mu_2} = i [\hat{\mathcal{G}}_l(\boldsymbol{\rho})]_{\mu_1 \mu_2} = i \begin{bmatrix} G_l^R(\boldsymbol{\rho}) & G_l^K(\boldsymbol{\rho}) \\ 0 & G_l^A(\boldsymbol{\rho}) \end{bmatrix}_{\mu_1 \mu_2}$$

$$\overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \neq \overleftarrow{\mu_1} \xrightarrow{(l', \boldsymbol{\rho}') \neq (l, \boldsymbol{\rho})}$$

$$G_l^R(\boldsymbol{\rho}) = [G_l^A(\boldsymbol{\rho})]^* = \frac{1}{\epsilon - \xi_l(\boldsymbol{\rho}) + i0^+}$$

$$G_l^K(\boldsymbol{\rho}) = -2\pi i n_l(\boldsymbol{\rho}, \epsilon) \delta[\epsilon - \xi_l(\boldsymbol{\rho})]$$

B

$$\overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \overleftarrow{\mu_2} = i [\hat{\mathcal{G}}_l(\boldsymbol{\rho})]_{\mu_1 \mu_2};$$

C

$$\overleftarrow{\mu_1} \xrightarrow{(l_1, \boldsymbol{\rho}_1)} \blacksquare \overleftarrow{\mu_2} \xrightarrow{(l_2, \boldsymbol{\rho}_2)} = -i I \delta_\xi [\hat{\tau}^0]_{\mu_1 \mu_2} (\boldsymbol{\rho}_1, \boldsymbol{\rho}_2)_{nn};$$

D

$$\begin{array}{c} (j_1, \boldsymbol{\rho}_1) \\ \mu_2 \\ \downarrow \\ (l_1, \boldsymbol{\rho}) \\ \mu_1 \leftarrow \text{diamond} \rightarrow (l_2, \boldsymbol{\rho}) \\ \mu_4 \\ \uparrow \\ (j_2, \boldsymbol{\rho}_2) \\ \mu_3 \end{array} = -\frac{i}{2} V_{l_1 l_2}^{j_1 j_2}(\boldsymbol{\rho}) \left([\hat{\tau}^2]_{\mu_1 \mu_2} [\hat{\tau}^0]_{\mu_3 \mu_4} + [\hat{\tau}^2]_{\mu_1 \mu_2} [\hat{\tau}^0]_{\mu_3 \mu_4} \right)$$

E

$$\overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \overleftarrow{\mu_2} = \overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \overleftarrow{\mu_2} + \overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \text{hexagon} \overleftarrow{\mu_2} + \overleftarrow{\mu_1} \xrightarrow{(l, \boldsymbol{\rho})} \text{circle} \overleftarrow{\mu_2}$$

Possible experimental manifestations of the many-body localization

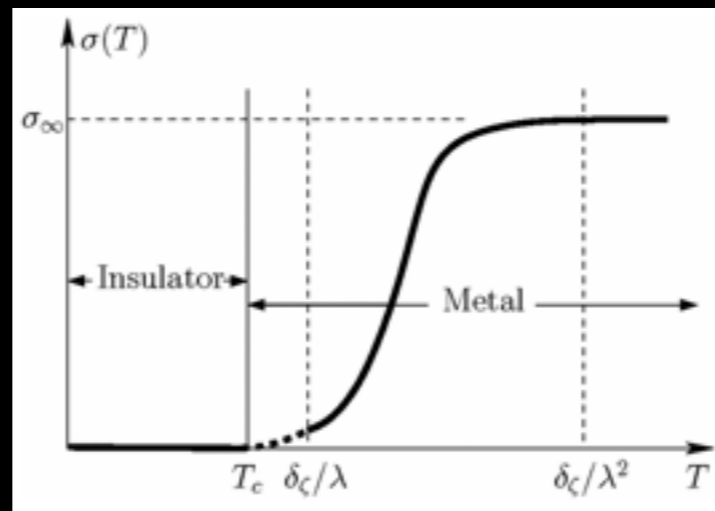
D. M. Basko,^{1,*} I. L. Aleiner,¹ and B. L. Altshuler^{1,2}

¹*Physics Department, Columbia University, New York, New York 10027, USA*

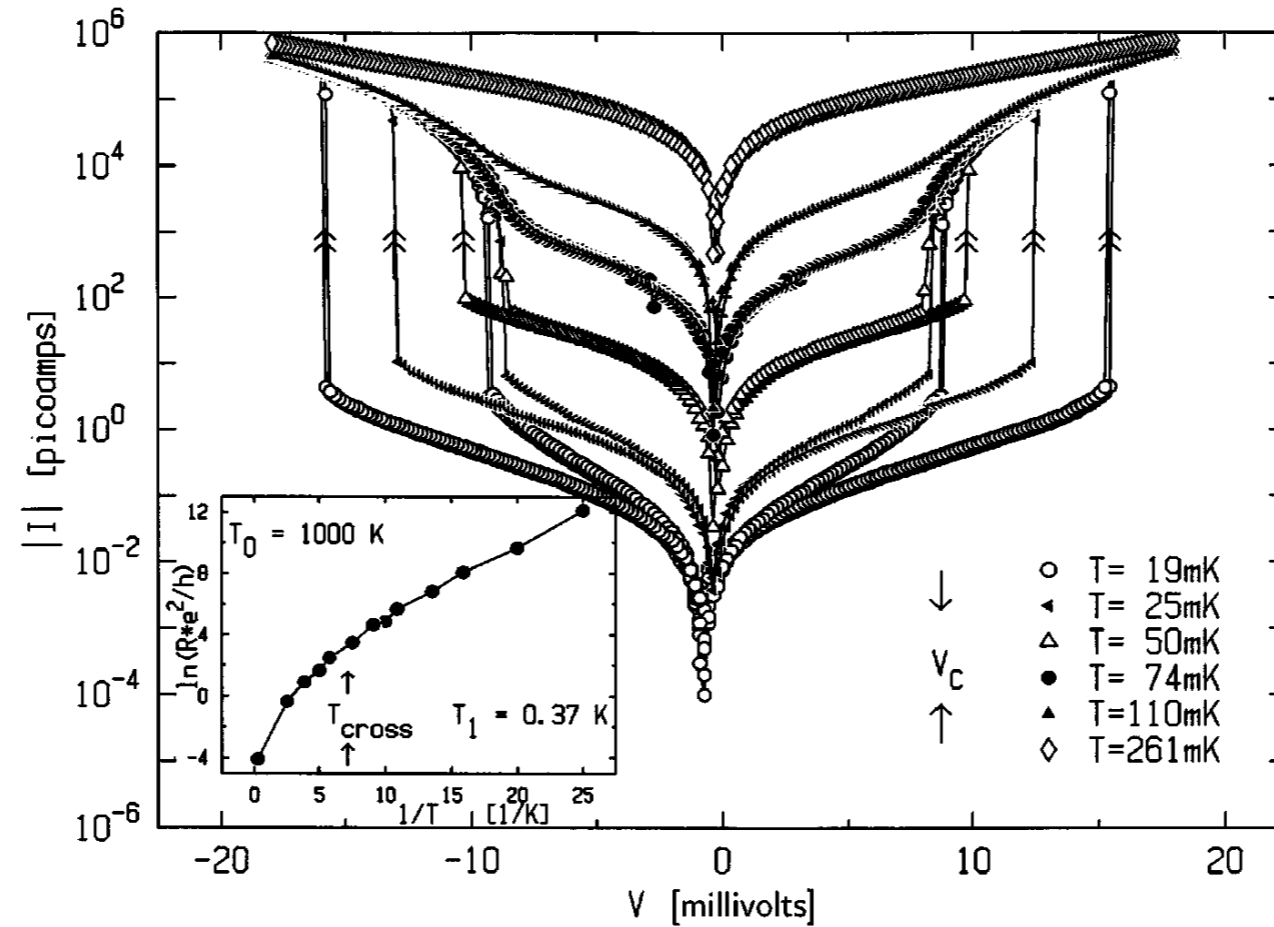
²*NEC-Laboratories America, Inc., 4 Independence Way, Princeton, New Jersey 08540, USA*

(Received 24 July 2007; published 23 August 2007; publisher error corrected 14 September 2007)

Recently, it was predicted that if all one-electron states in a noninteracting disordered system are localized, the interaction between electrons in the absence of coupling to phonons leads to a finite-temperature metal-insulator transition. Here, we show that even in the presence of a weak coupling to phonons the transition manifests itself (i) in the nonlinear conduction, leading to a bistable I–V curve, and (ii) by a dramatic enhancement of the nonequilibrium current noise near the transition.



No Superconductivity.



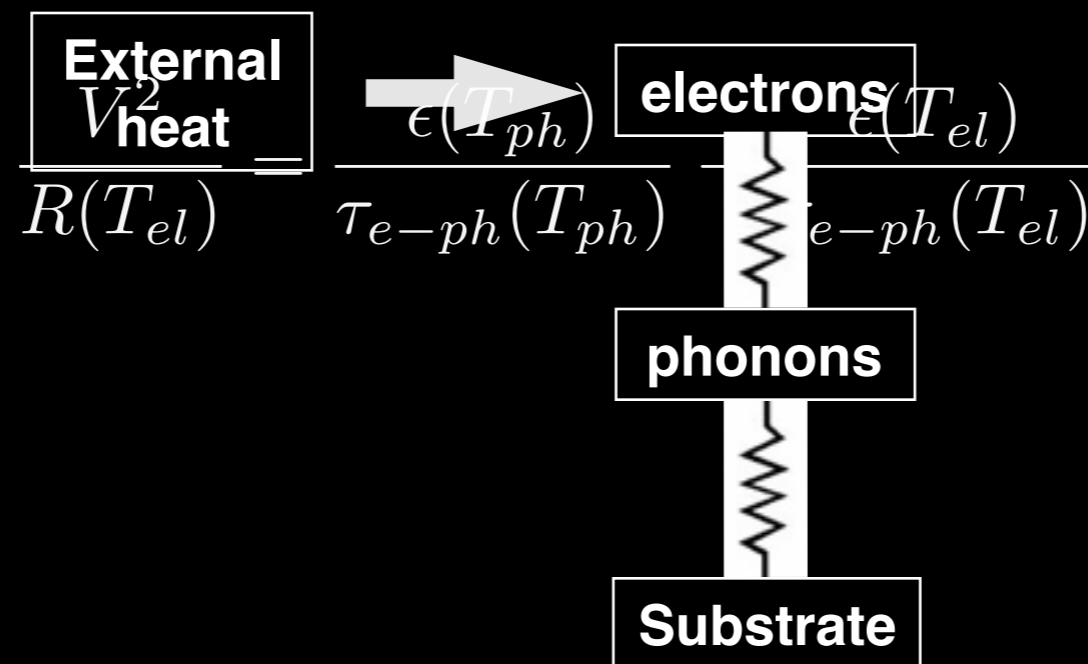
Ladieu, Sanquer, and Bouchaud, 1996

a: $\text{Y}_x\text{Si}_{1-x}$

Jumps in current-voltage characteristics in disordered films

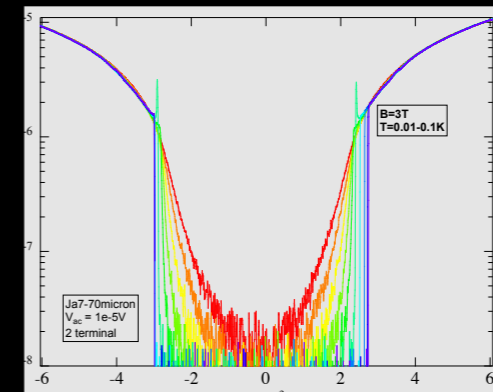
B. L. Altshuler, V. E. Kravtsov, I. V. Lerner, and I. L. Aleiner,
Phys. Rev. Lett. 102, 176803 (2009)

Heat balance equation:



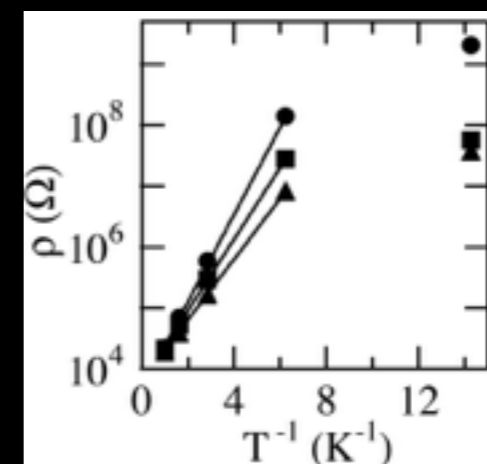
Jumps in current-voltage characteristics in disordered films

1. Electrons and phonons are decoupled
2. Electrons are strongly interacting
(electrons can have $T_{el} \neq T_{ph}$)
3. I-V is linear
(Apparent non-linearity from electron heating)



4. $R(T)$ is a fast function

$$\rho(T) = \rho_0 e^{T_1/T}$$

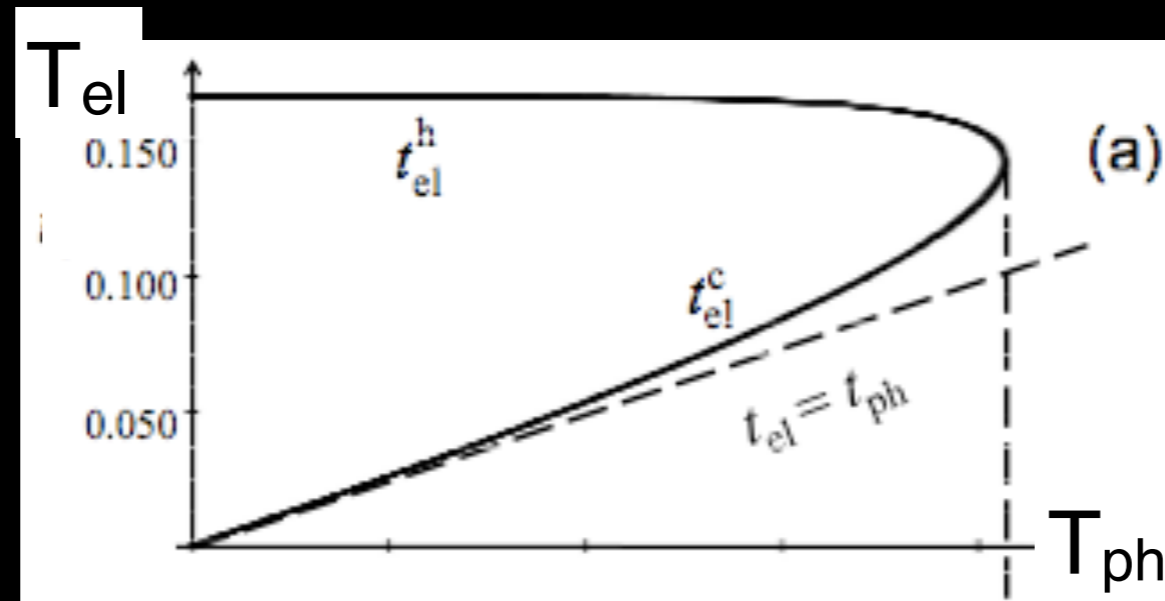


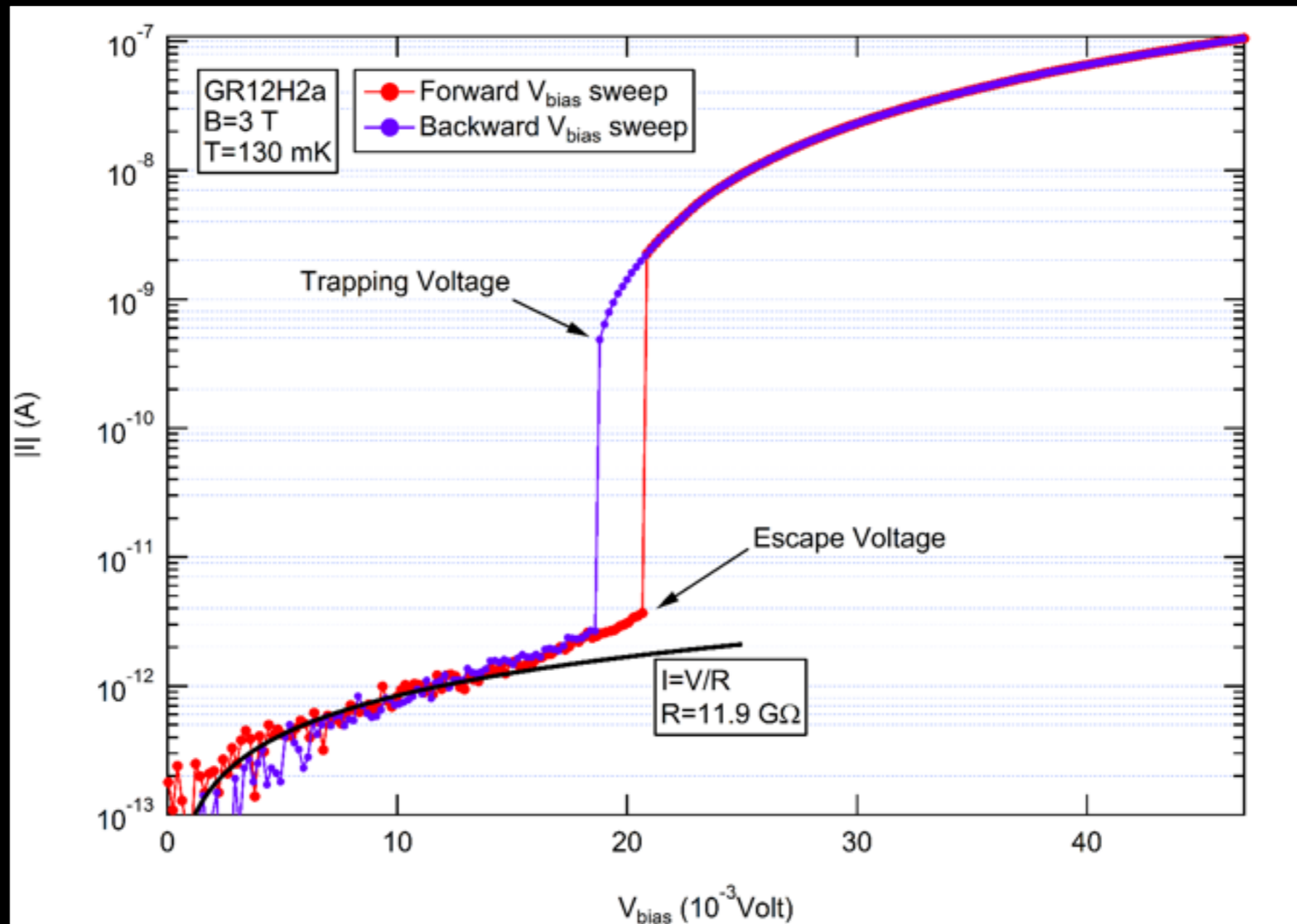
Jumps in current-voltage characteristics in disordered films

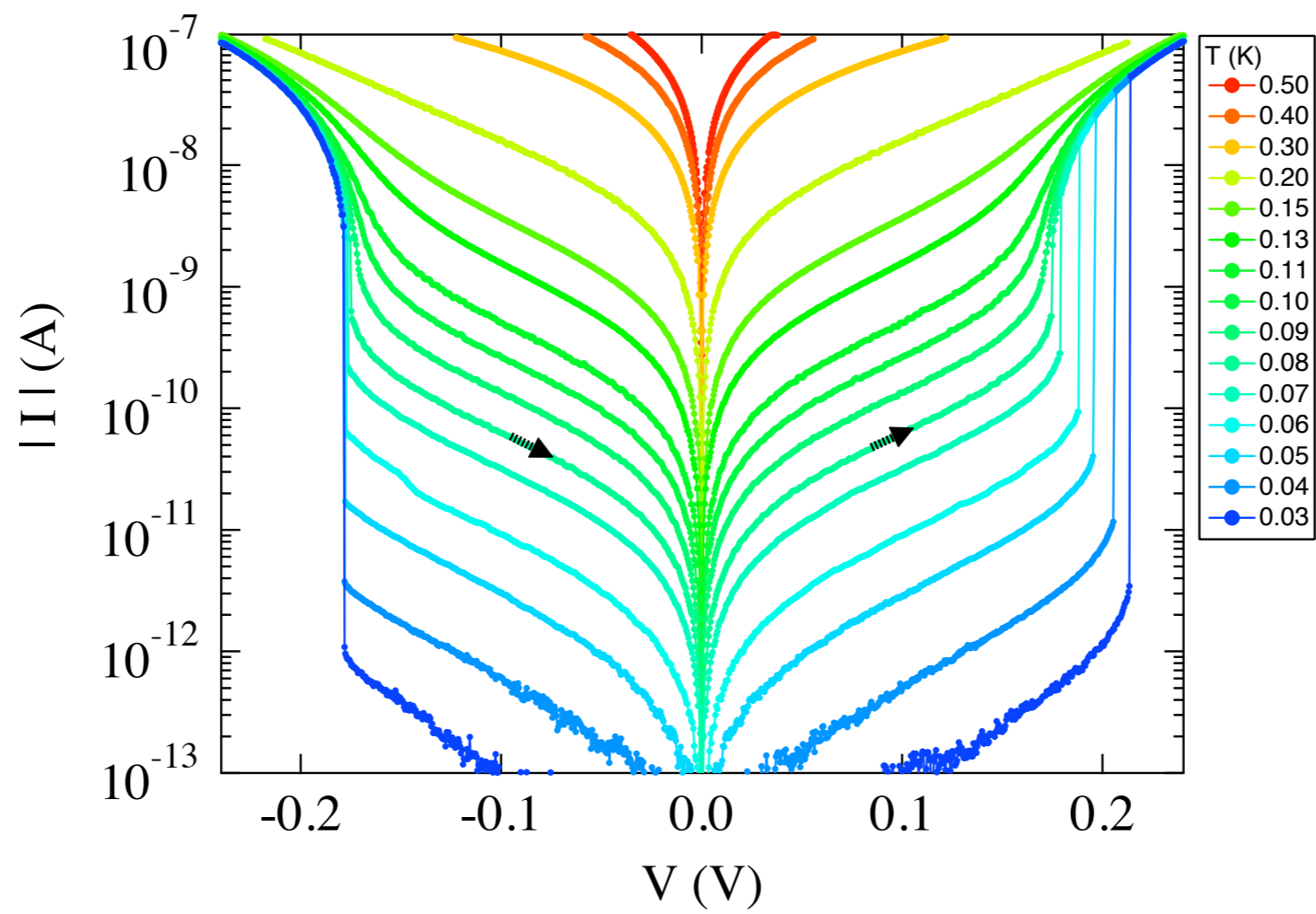
B. L. Altshuler, V. E. Kravtsov, I. V. Lerner, and I. L. Aleiner,
Phys. Rev. Lett. 102, 176803 (2009)

Heat balance equation:

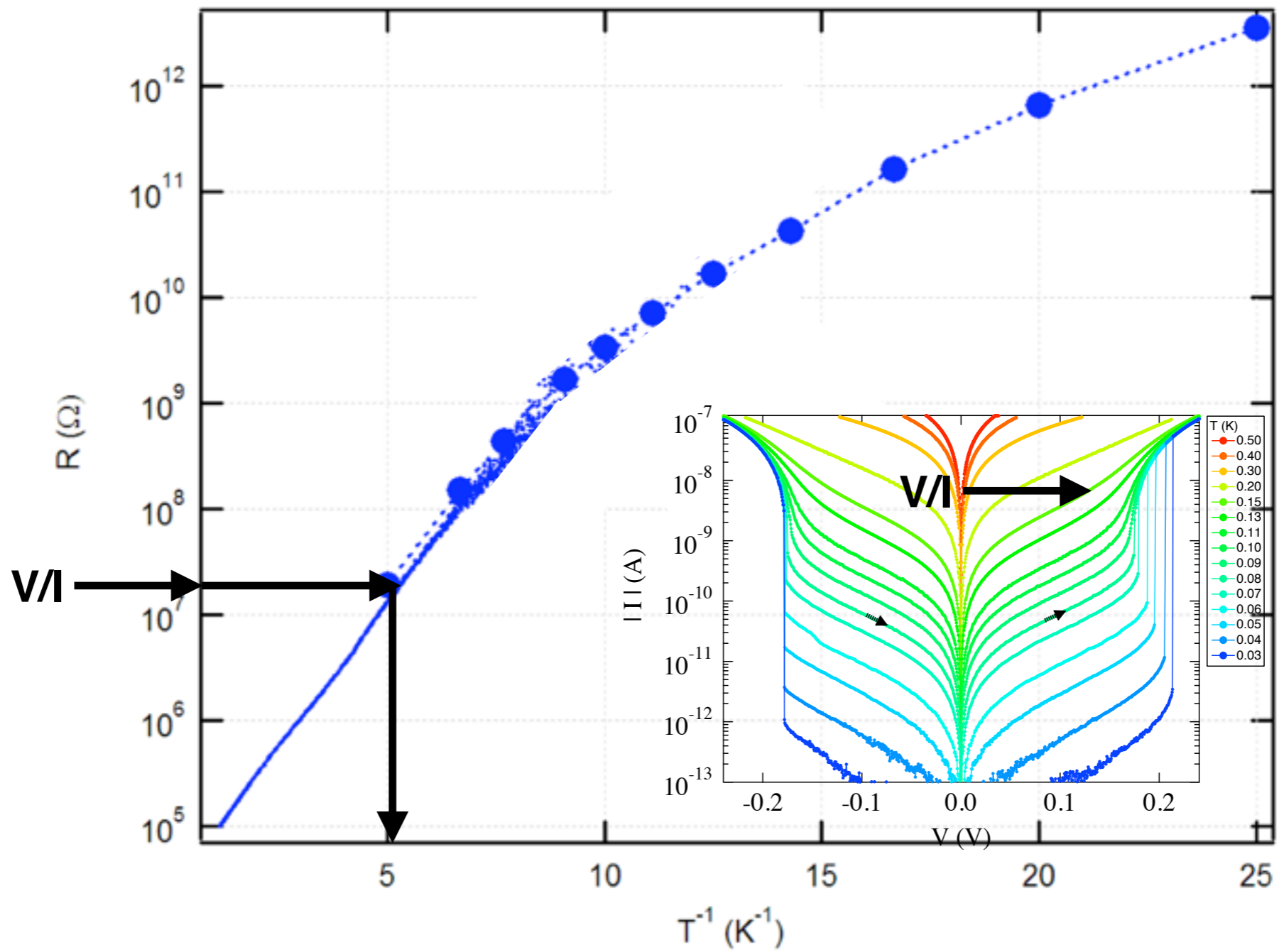
$$\frac{V^2}{R(T_{el})} = AT_{el}^6 - BT_{ph}^6$$



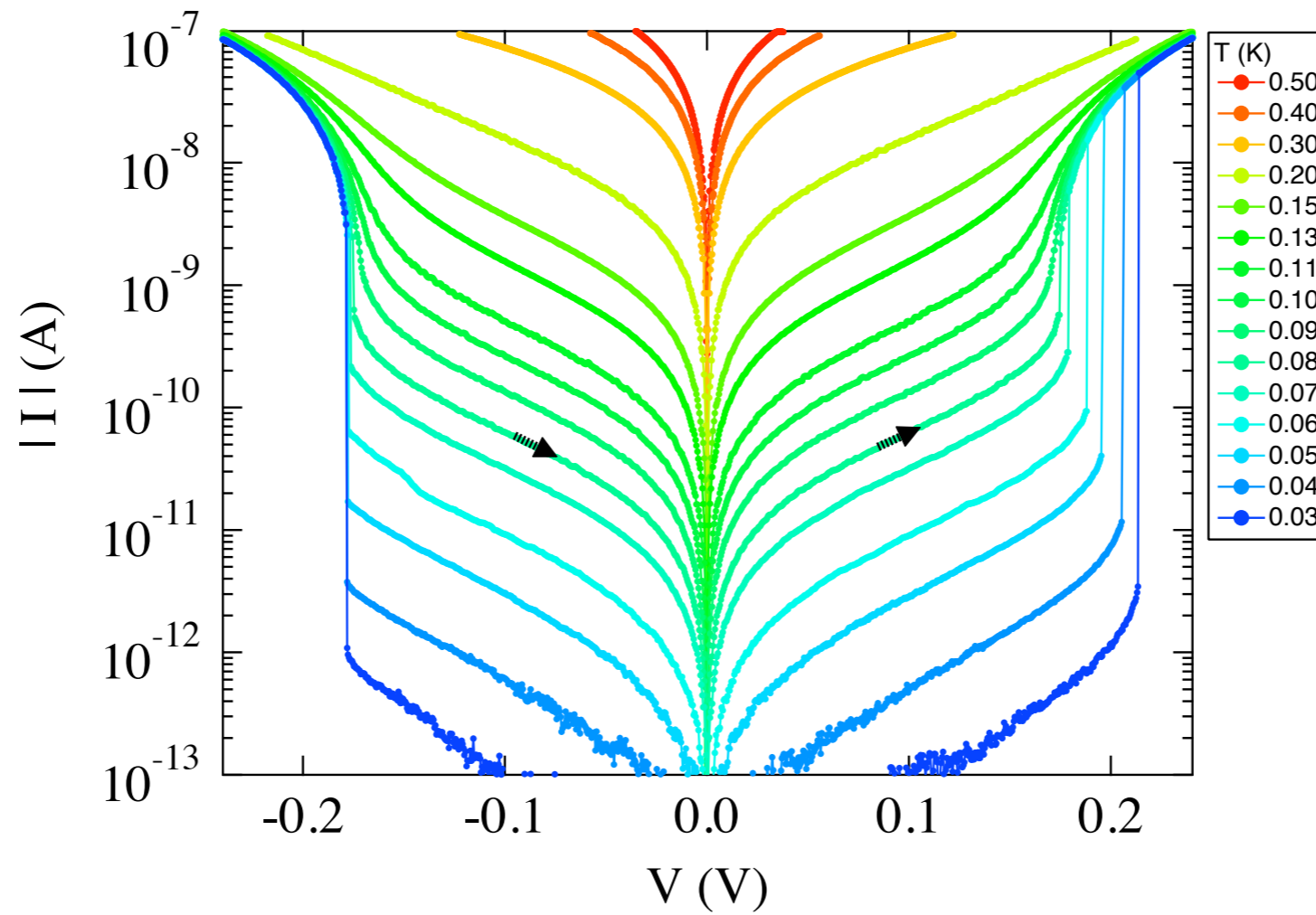


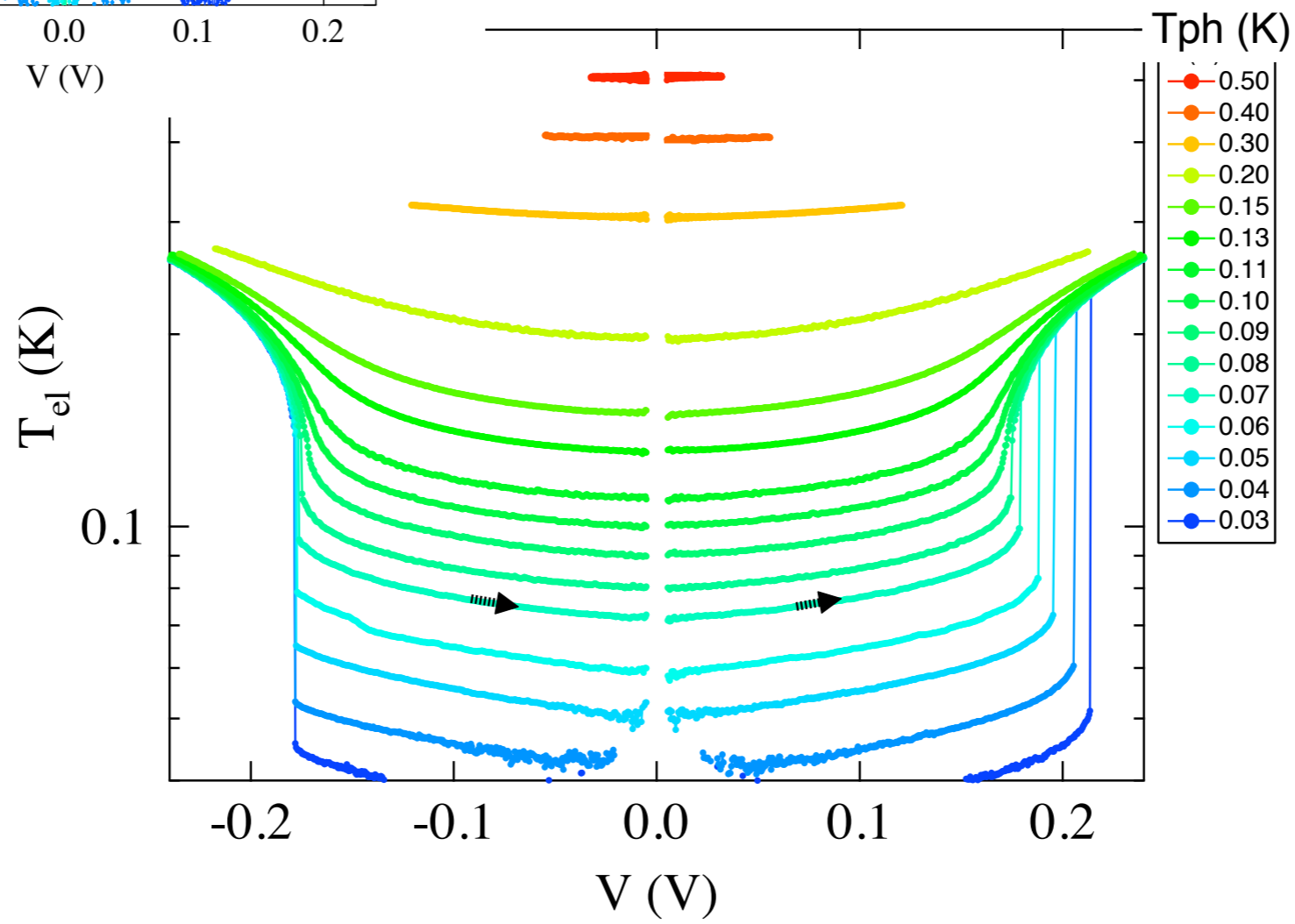
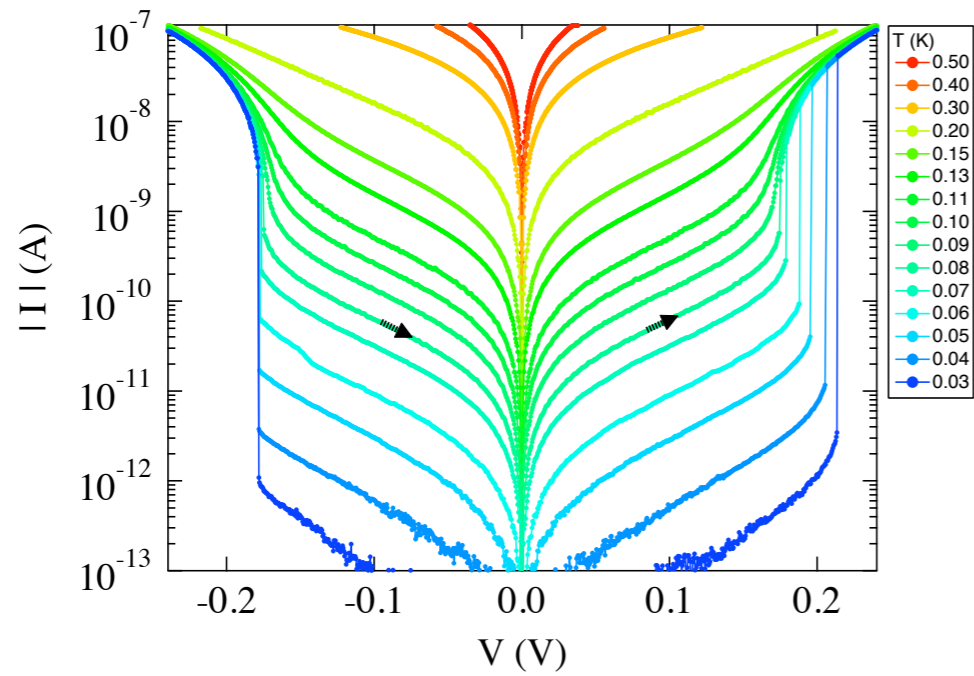


Determine T_{el}

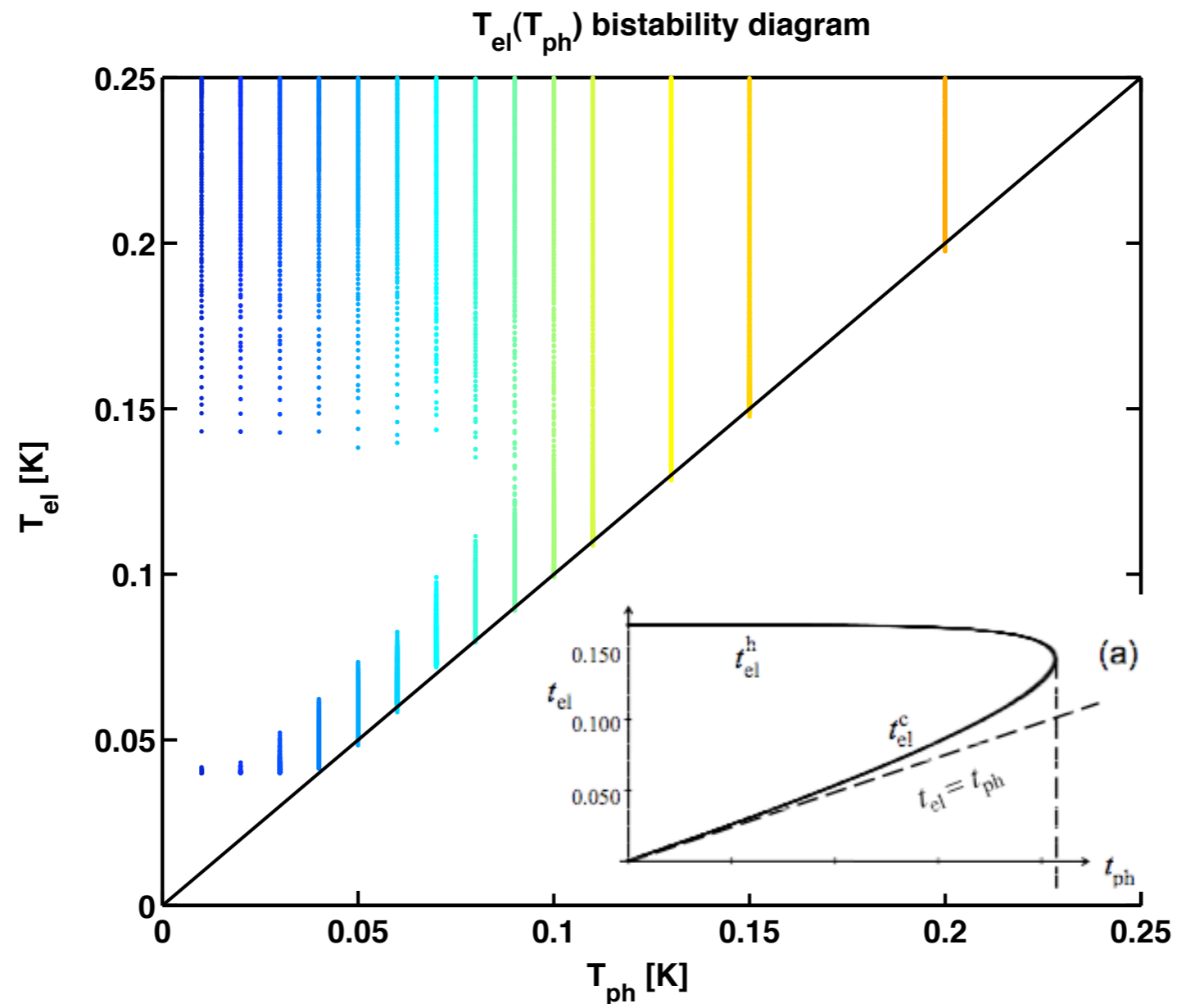


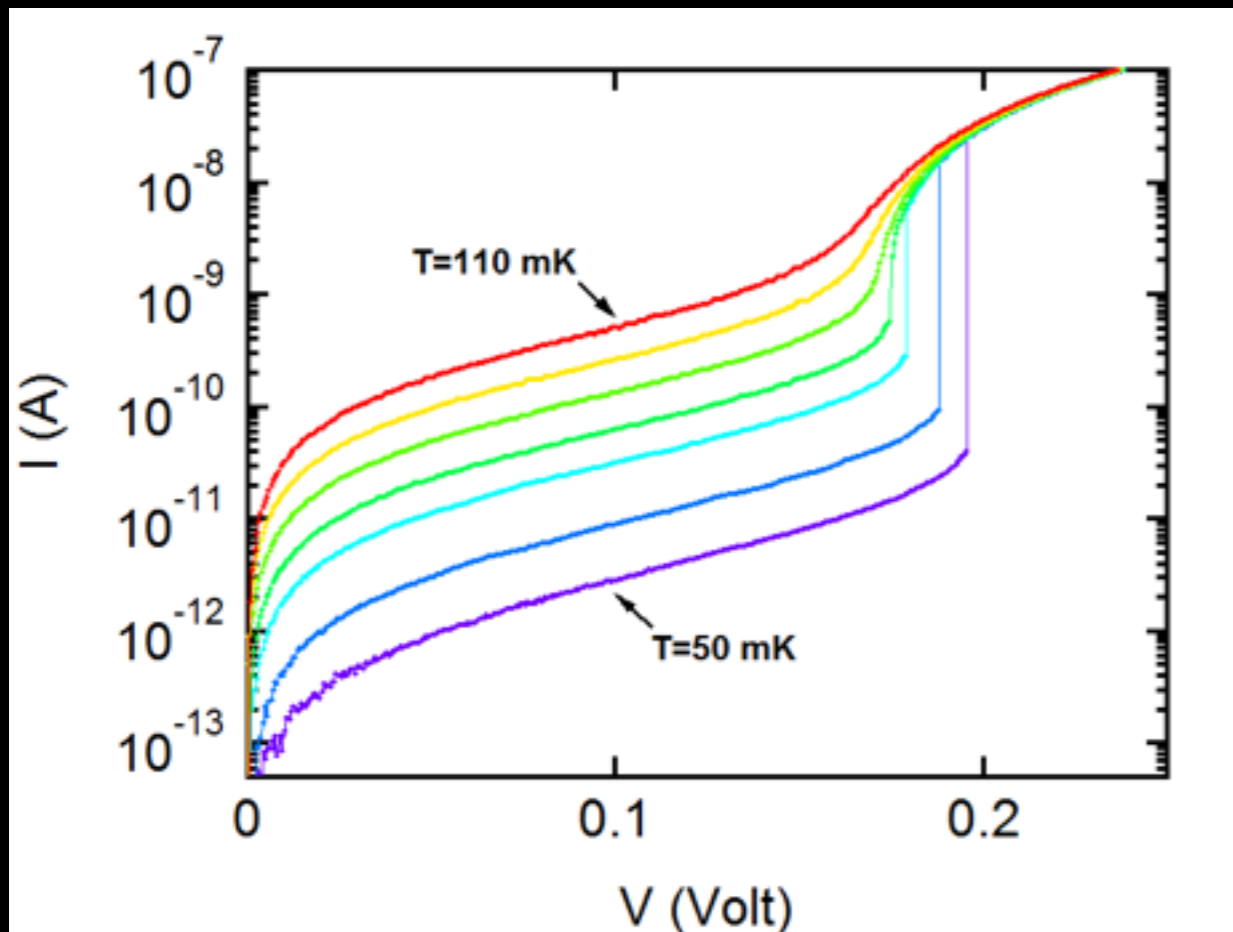
Applying to I-V



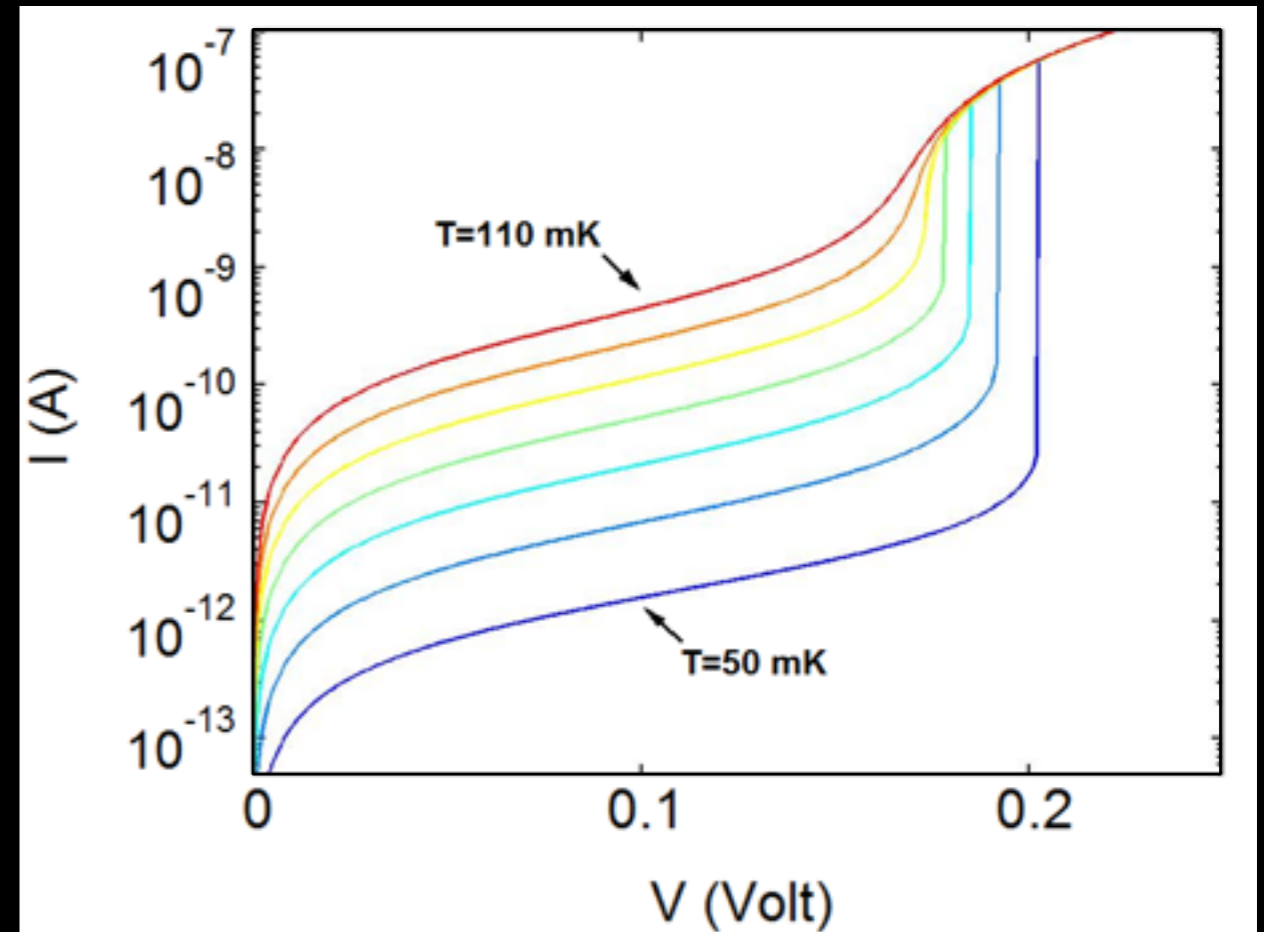


Excluded Region

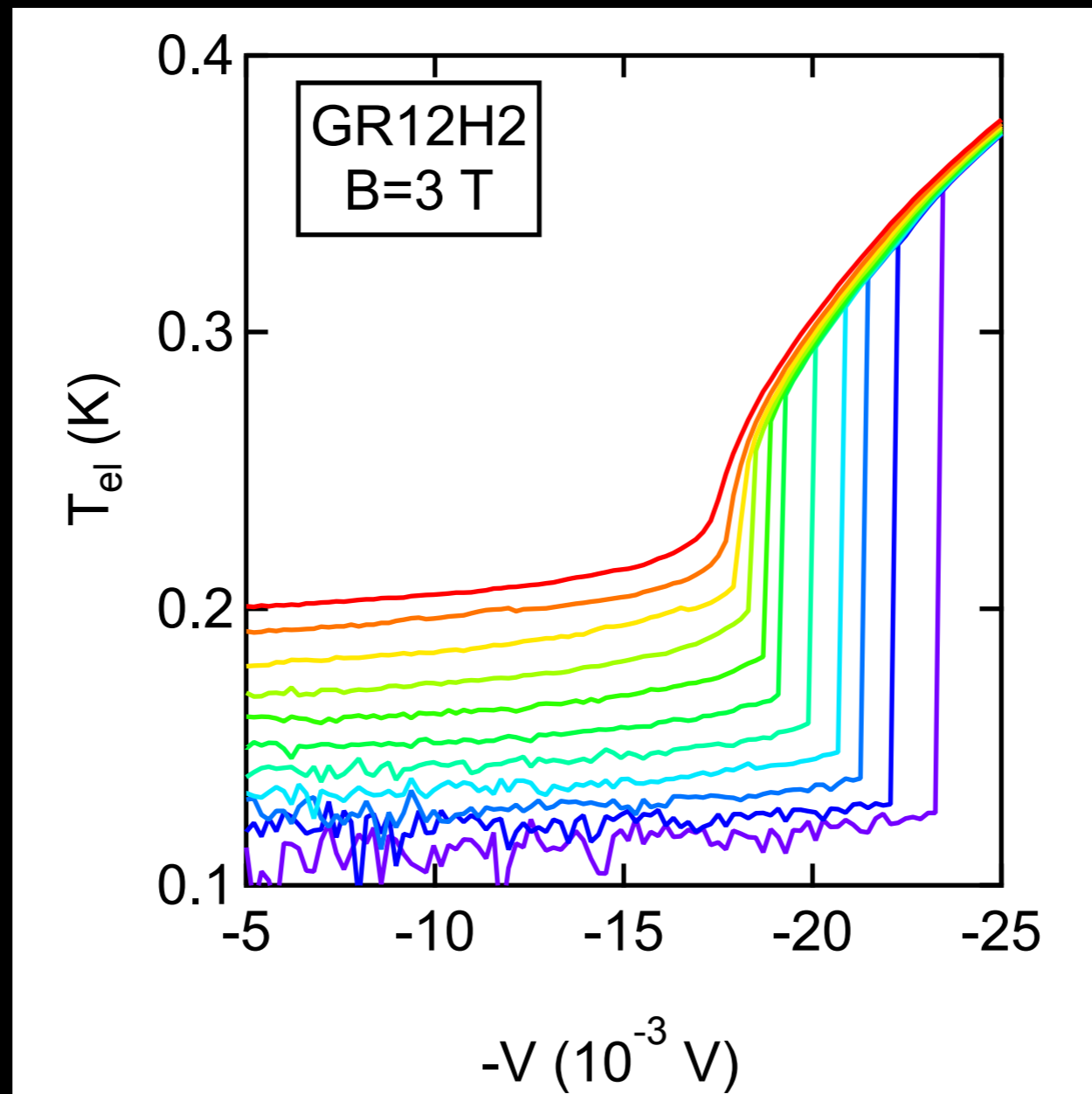


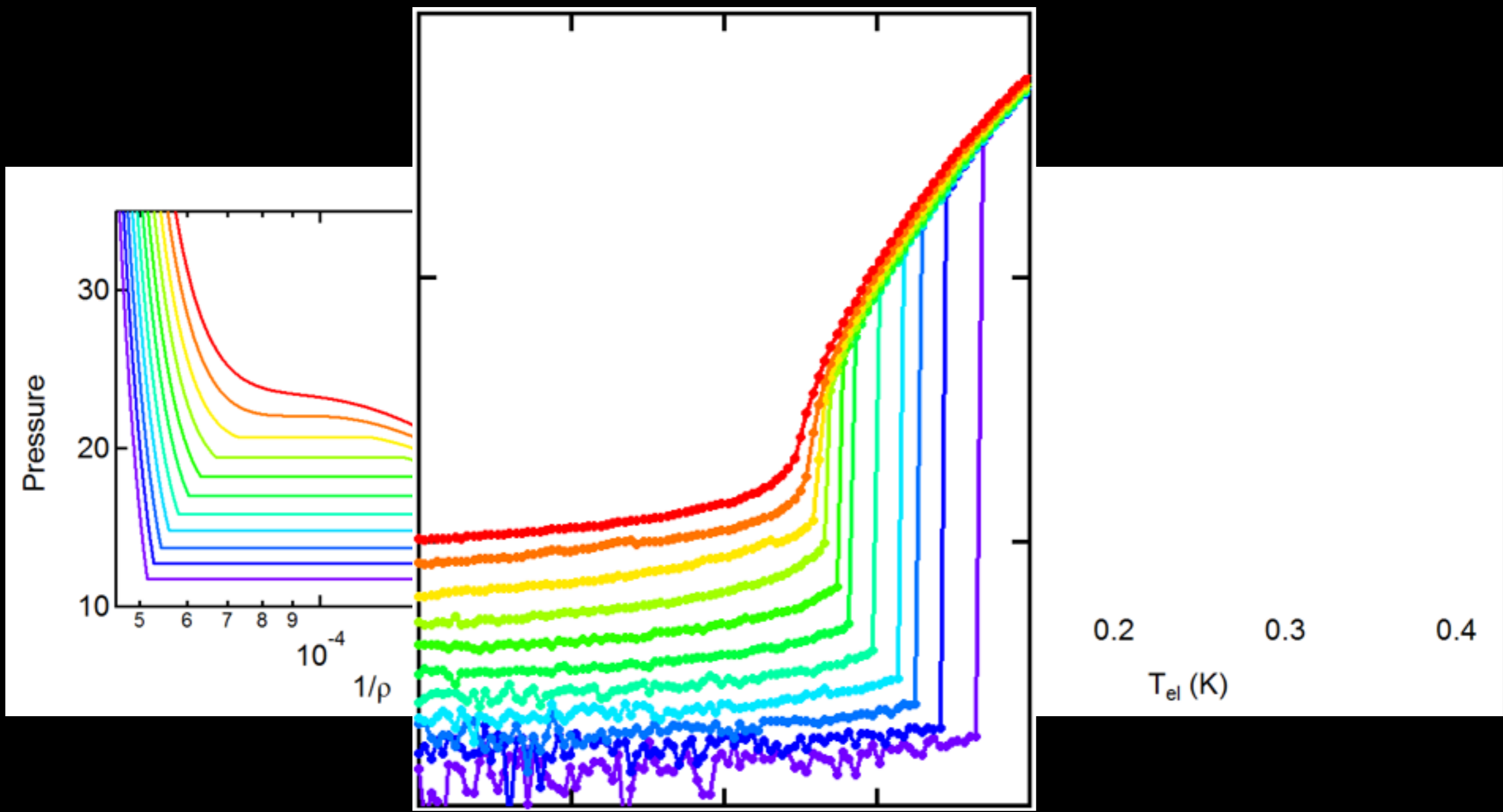


M. Ovadia, B. Sacepe, and D. Shahar, Phys. Rev. Lett. 102, 176802 (2009)



B. L. Altshuler, V. E. Kravtsov, I. V. Lerner, and I. L. Aleiner, Phys. Rev. Lett. 102, 176803 (2009)



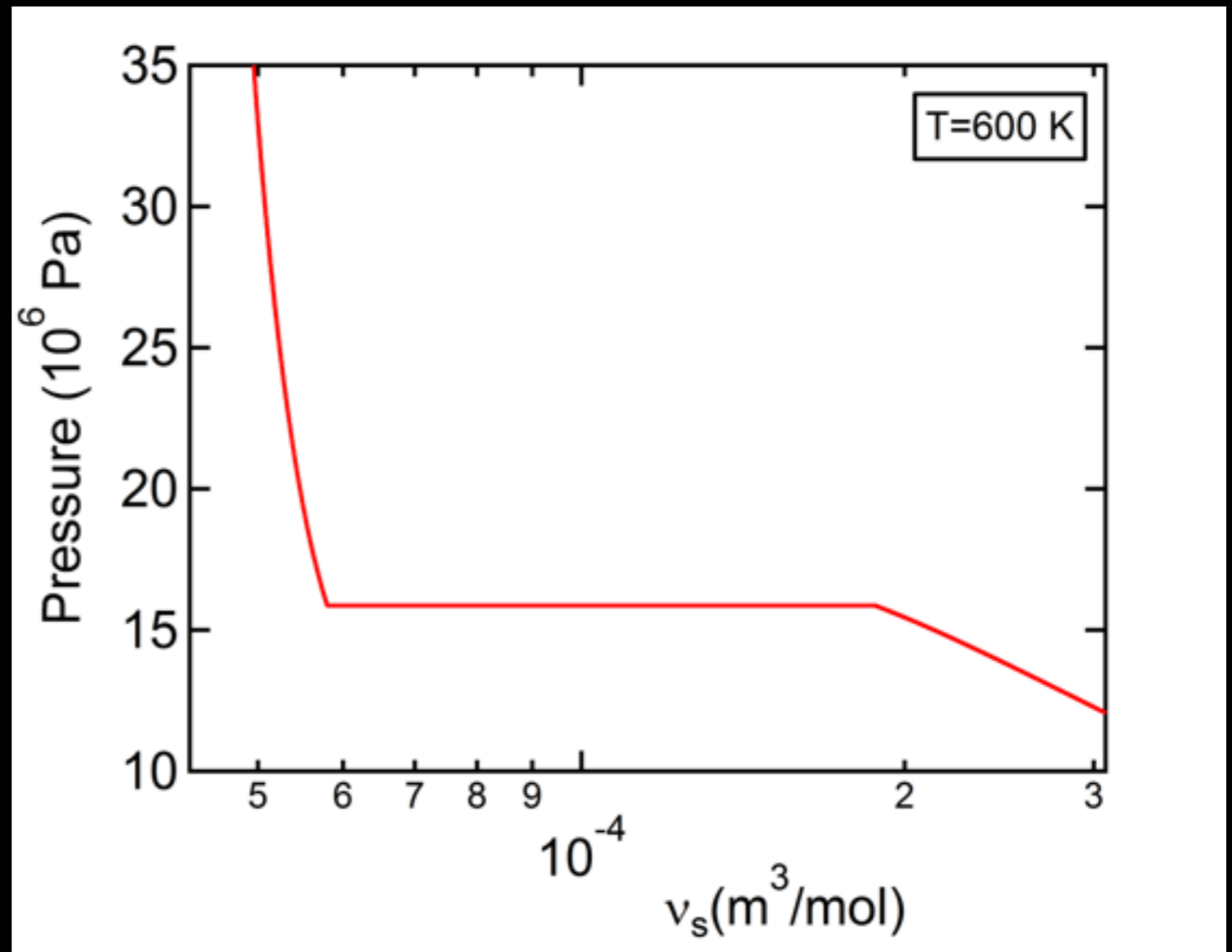


The Van der Waals equation

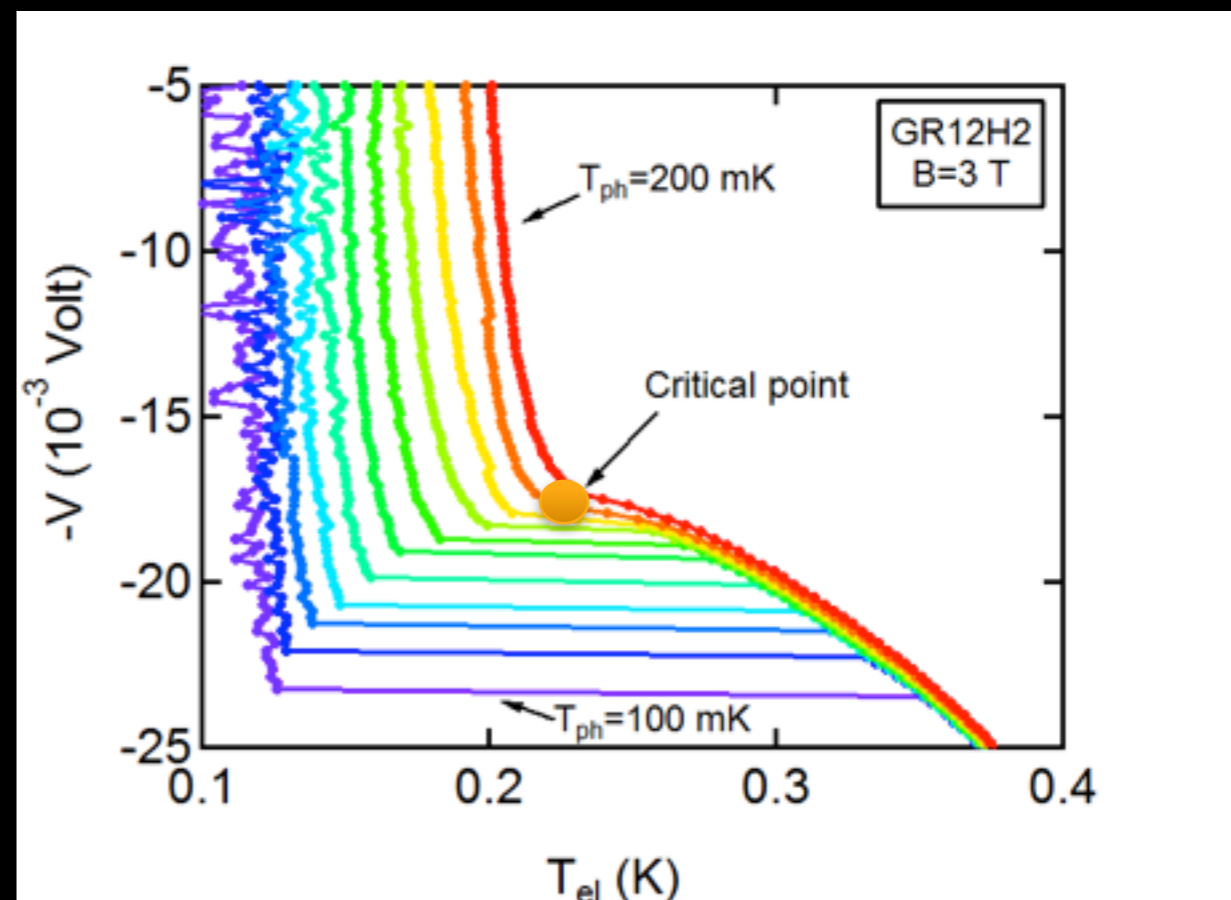
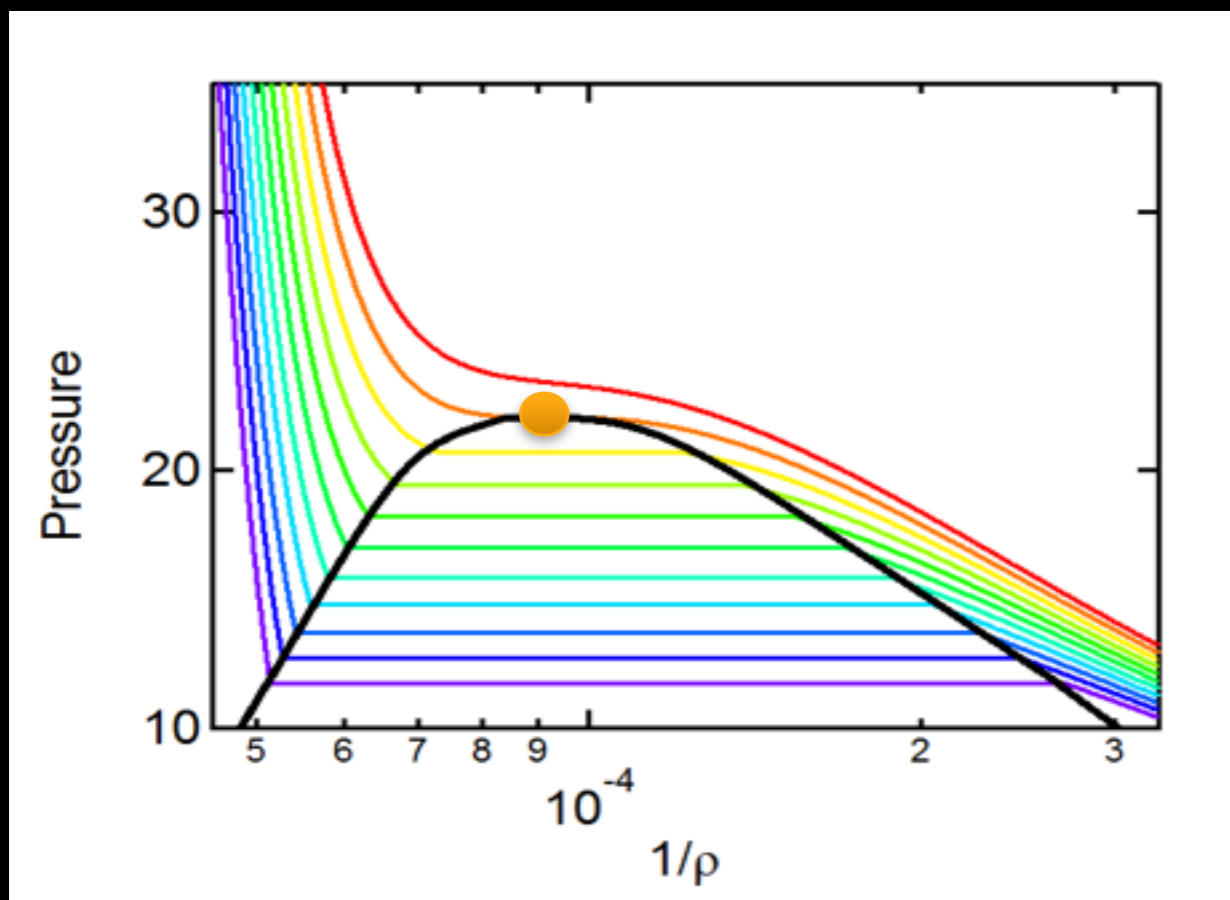
- First correction to the pressure equation of state

$$P = \frac{Nk_B T}{V - b} - \frac{N^2 a}{V^2}$$

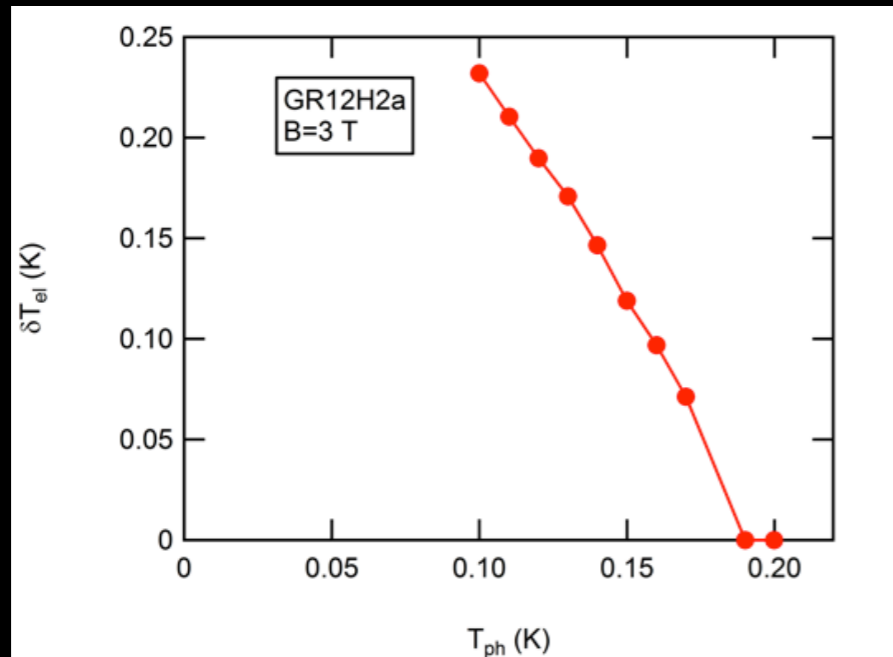
- Maxwell construction



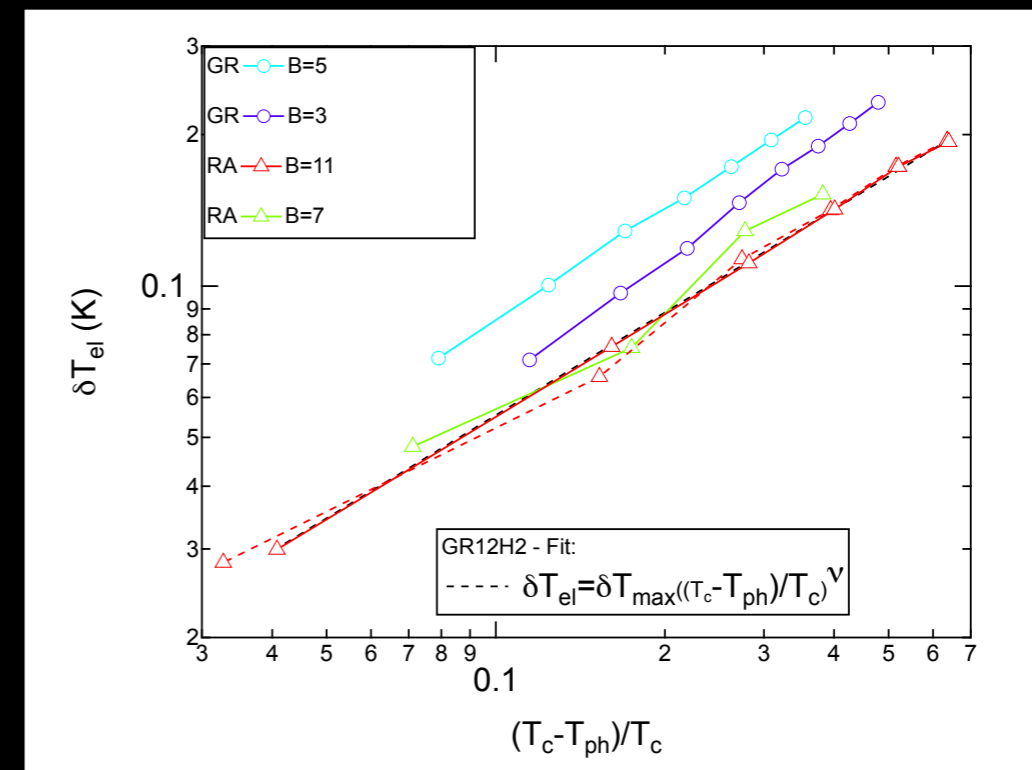
An Analogy



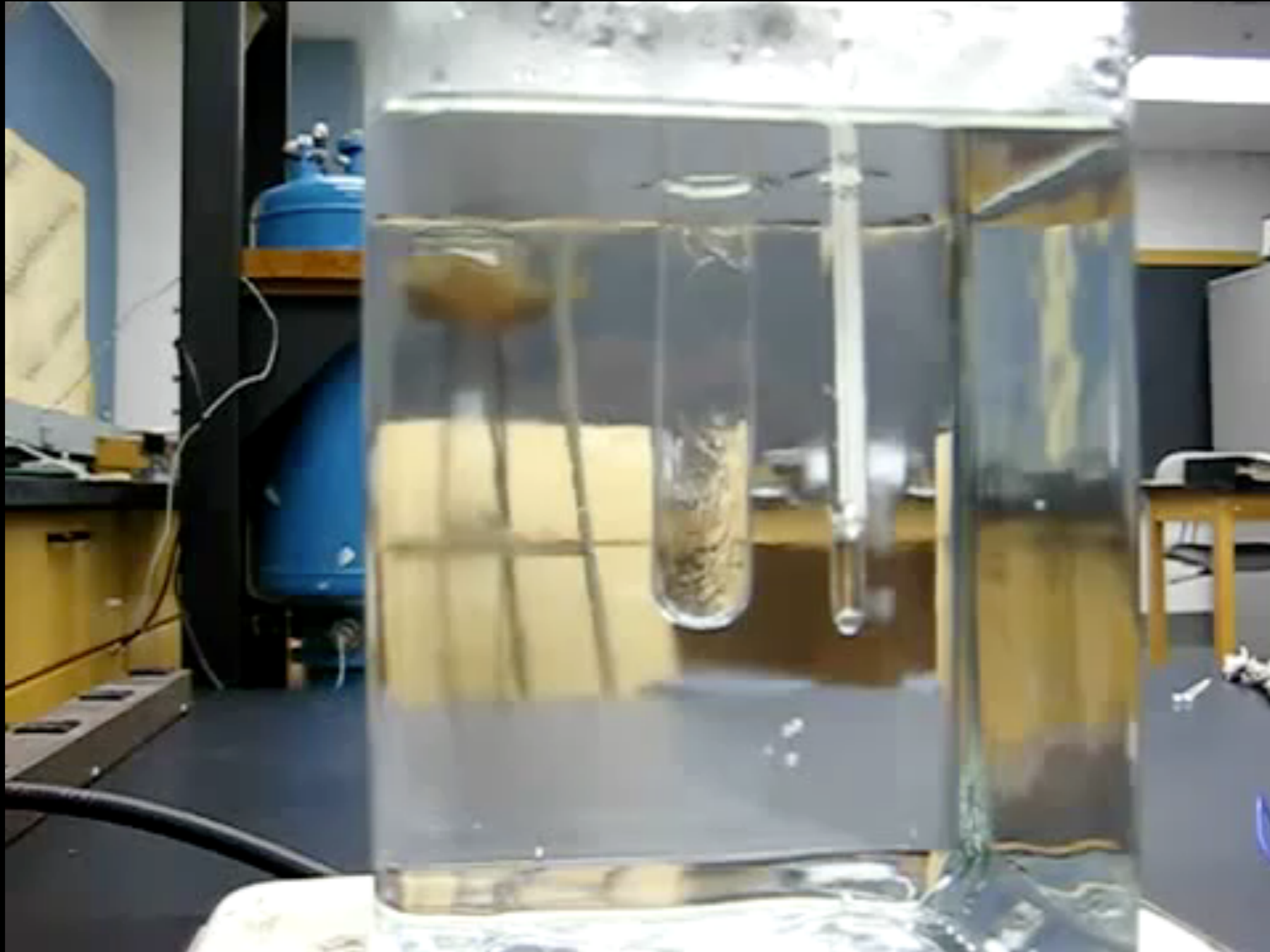
Order parameter



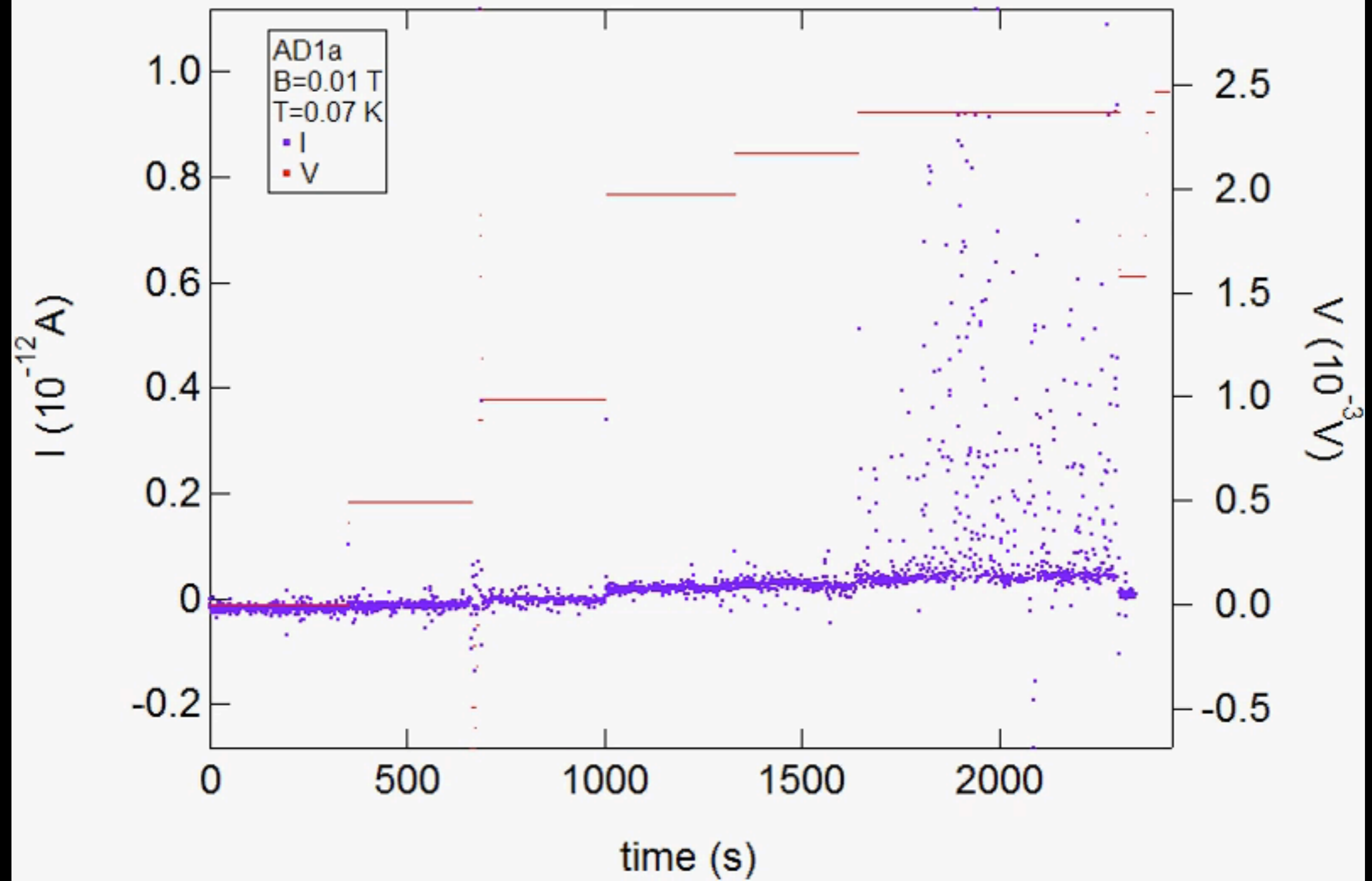
Critical exponents



Critical Opalescence



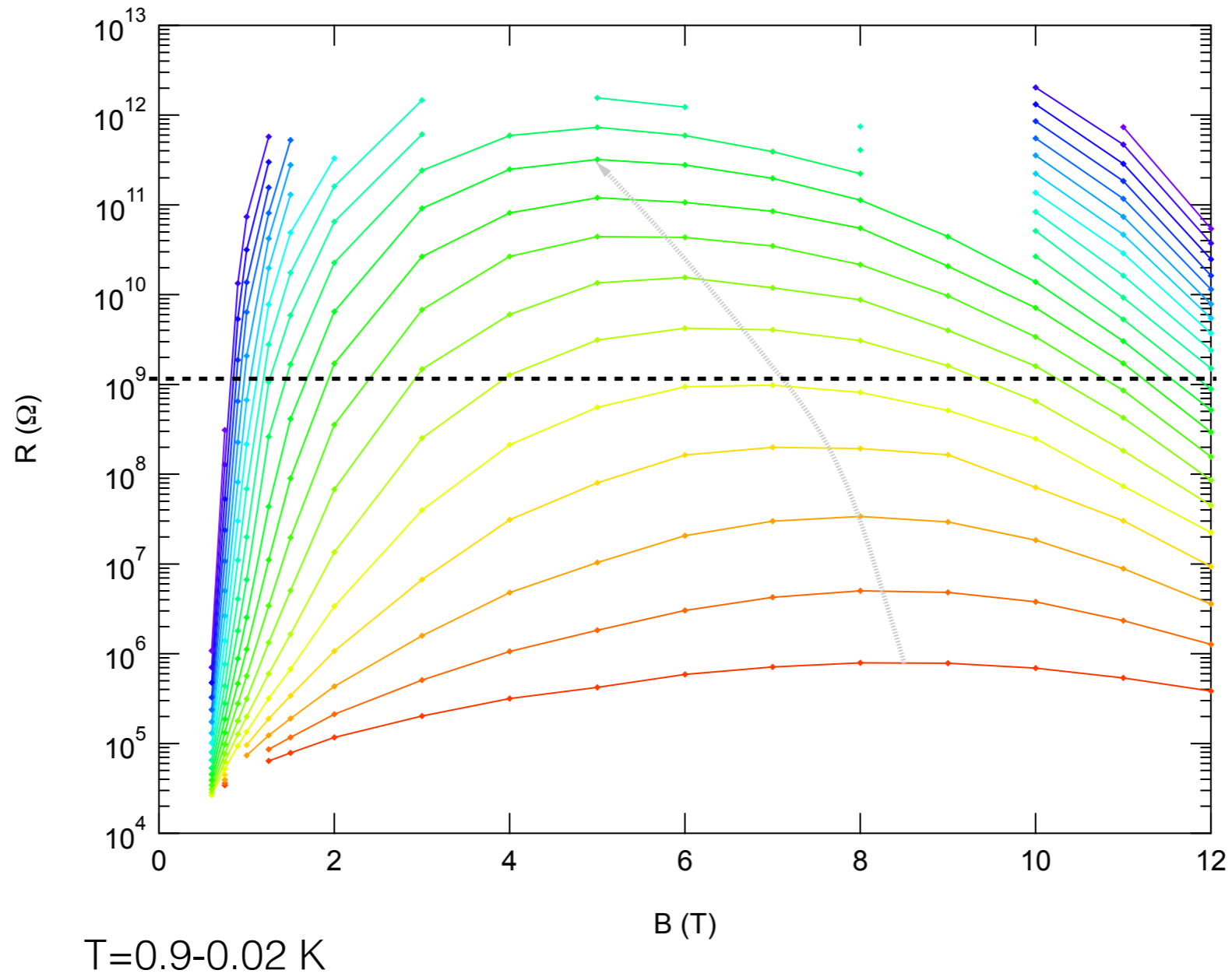
Time dependence



The Insulator at $T < 0.2$ K:

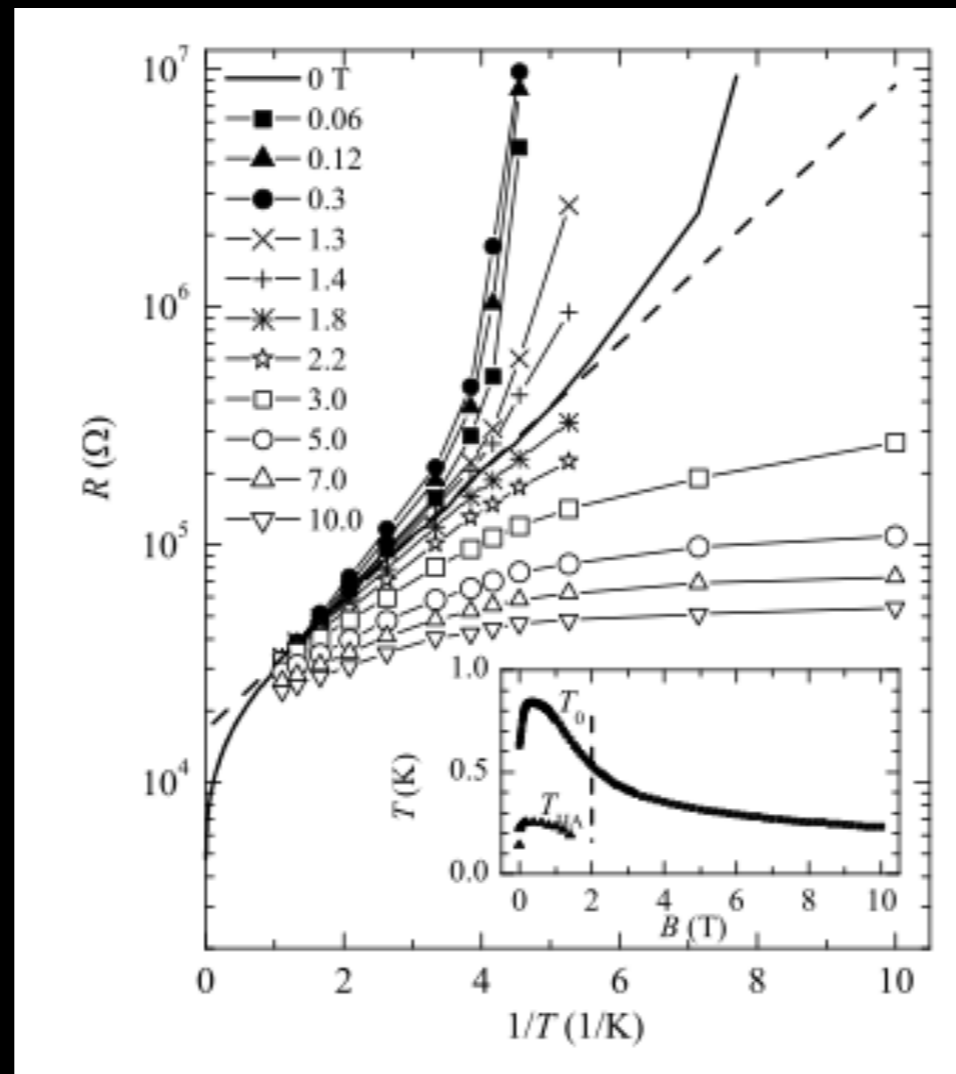
Ohmic transport

Ohmic transport



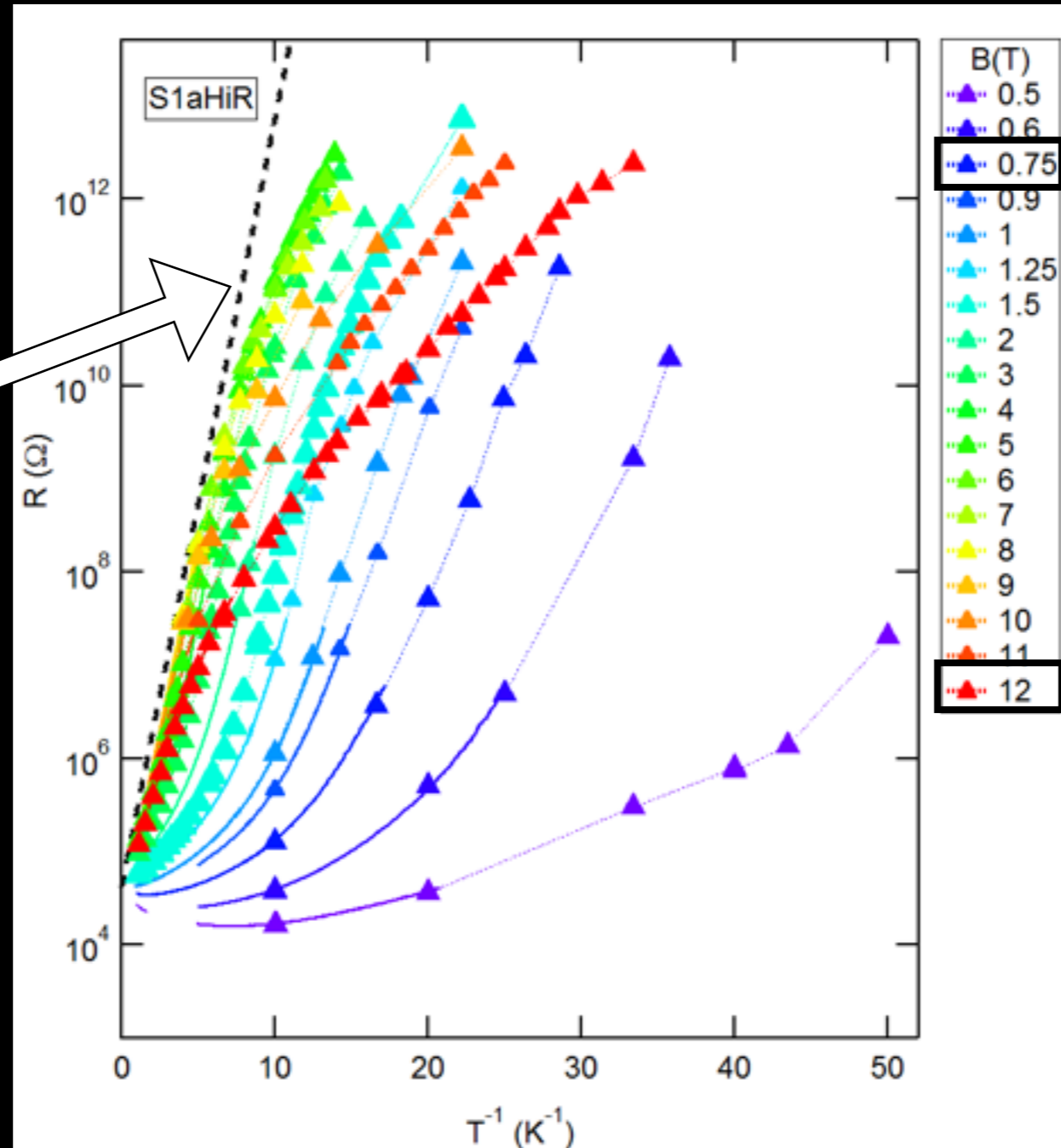
Hyperactivated Resistance in TiN Films on the Insulating Side of the Disorder-Driven Superconductor-Insulator Transition[¶]

T. I. Baturina^a, A. Yu. Mironov^a, V. M. Vinokur^b, M. R. Baklanov^c, and C. Strunk^d

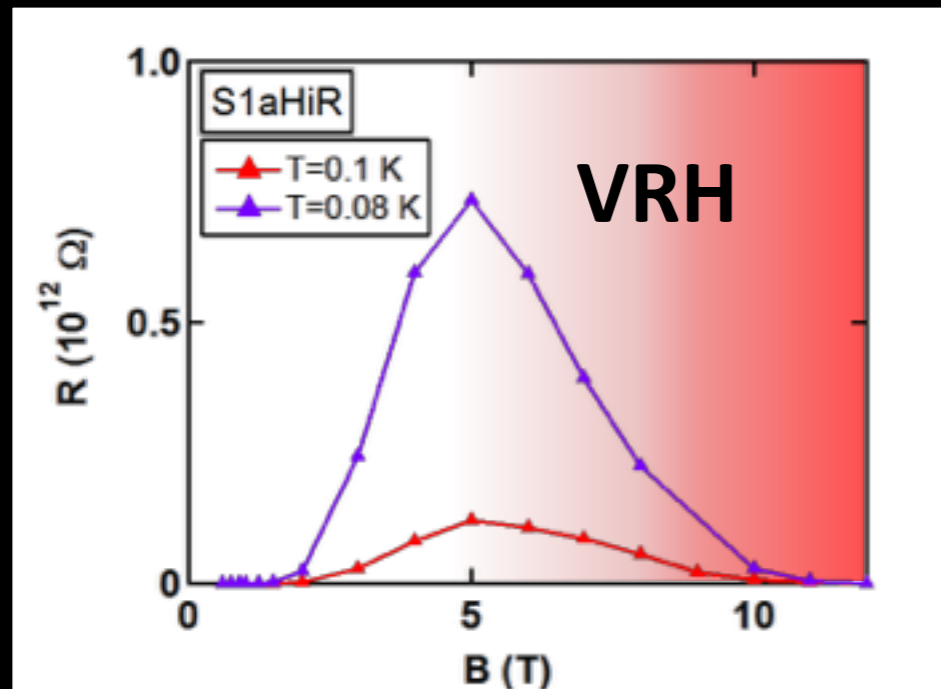


Arrhenius plot

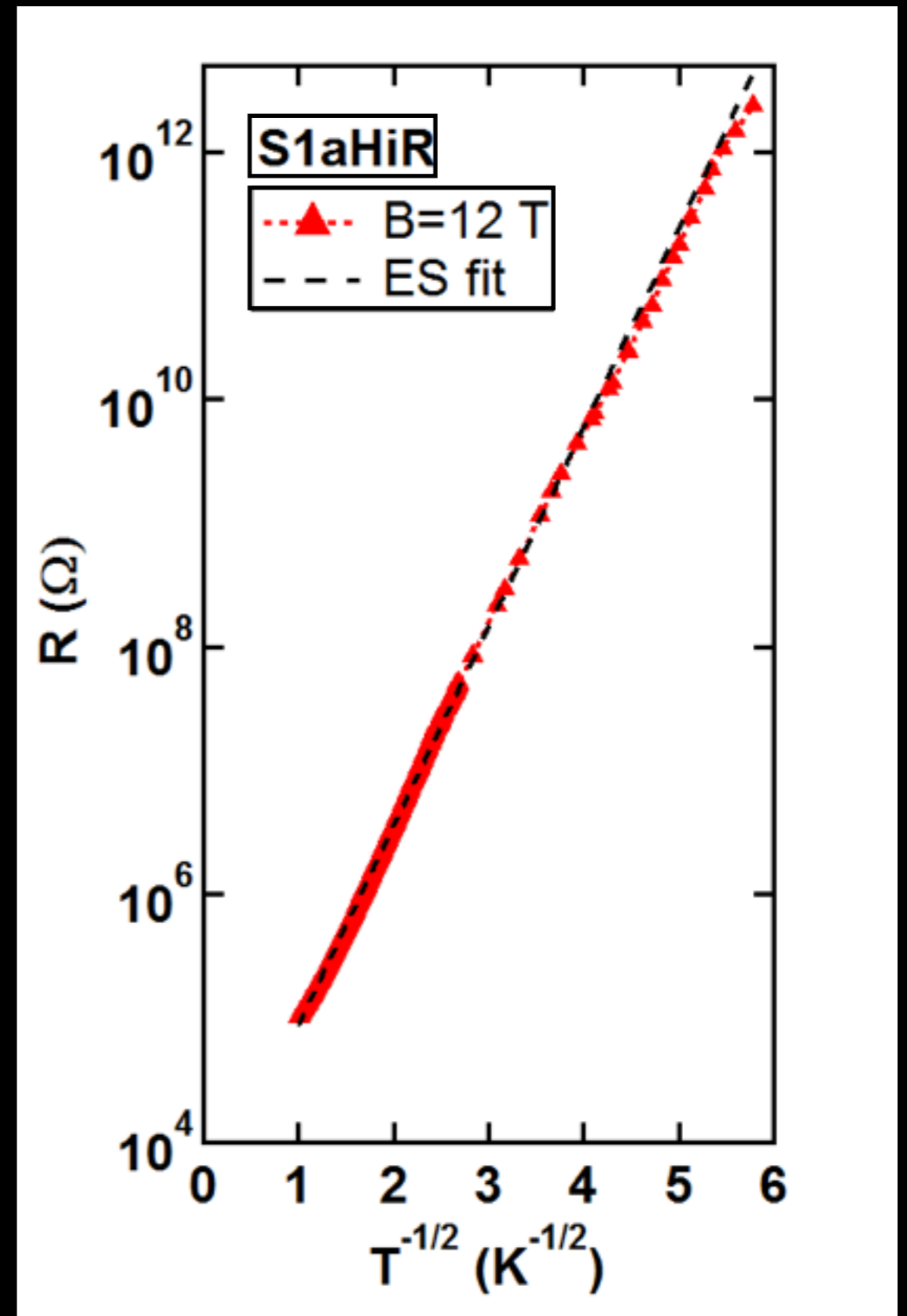
Activation



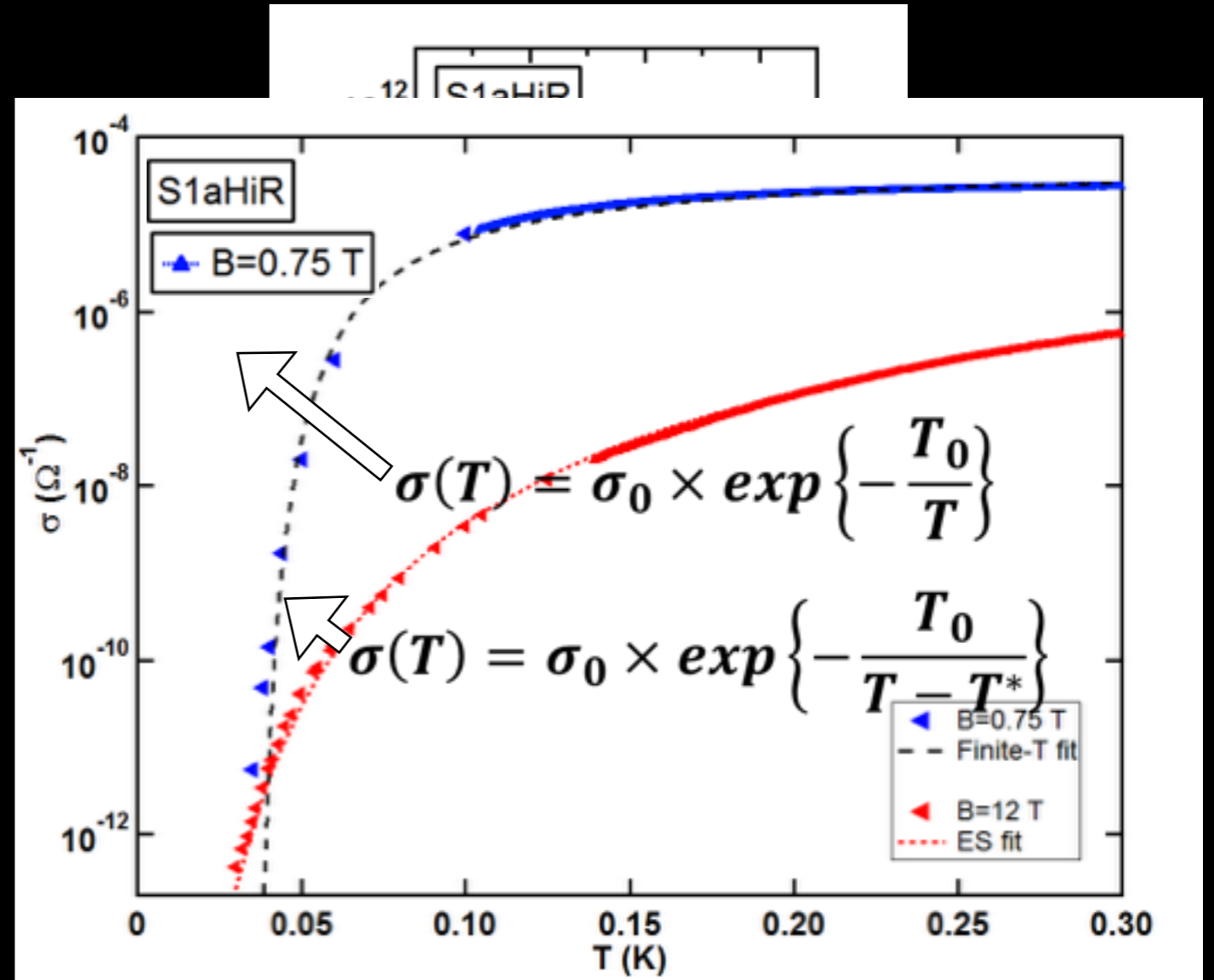
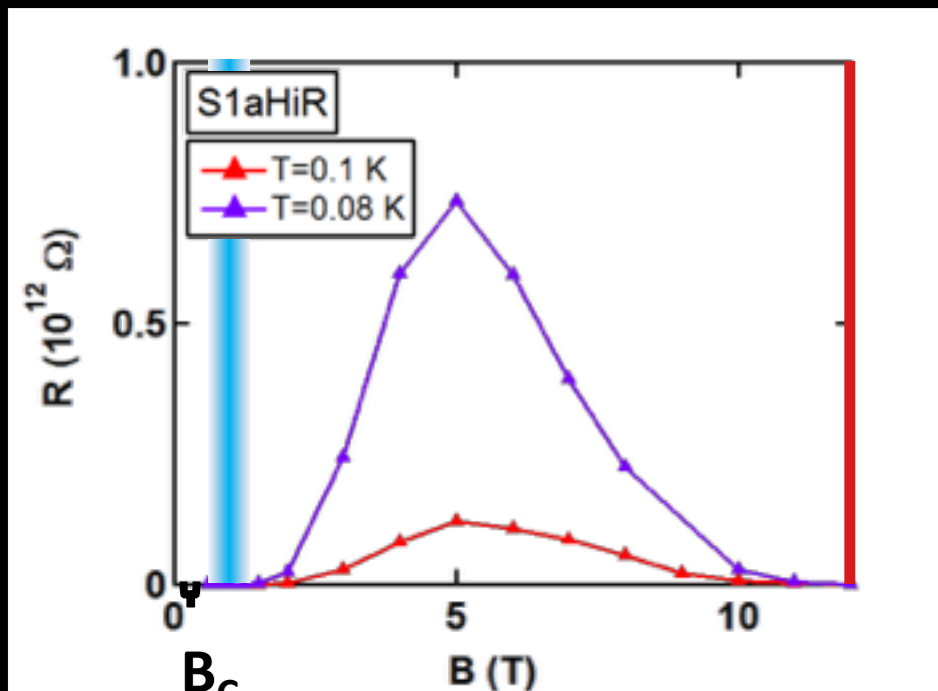
Efros-Shklovskii Hopping at High B



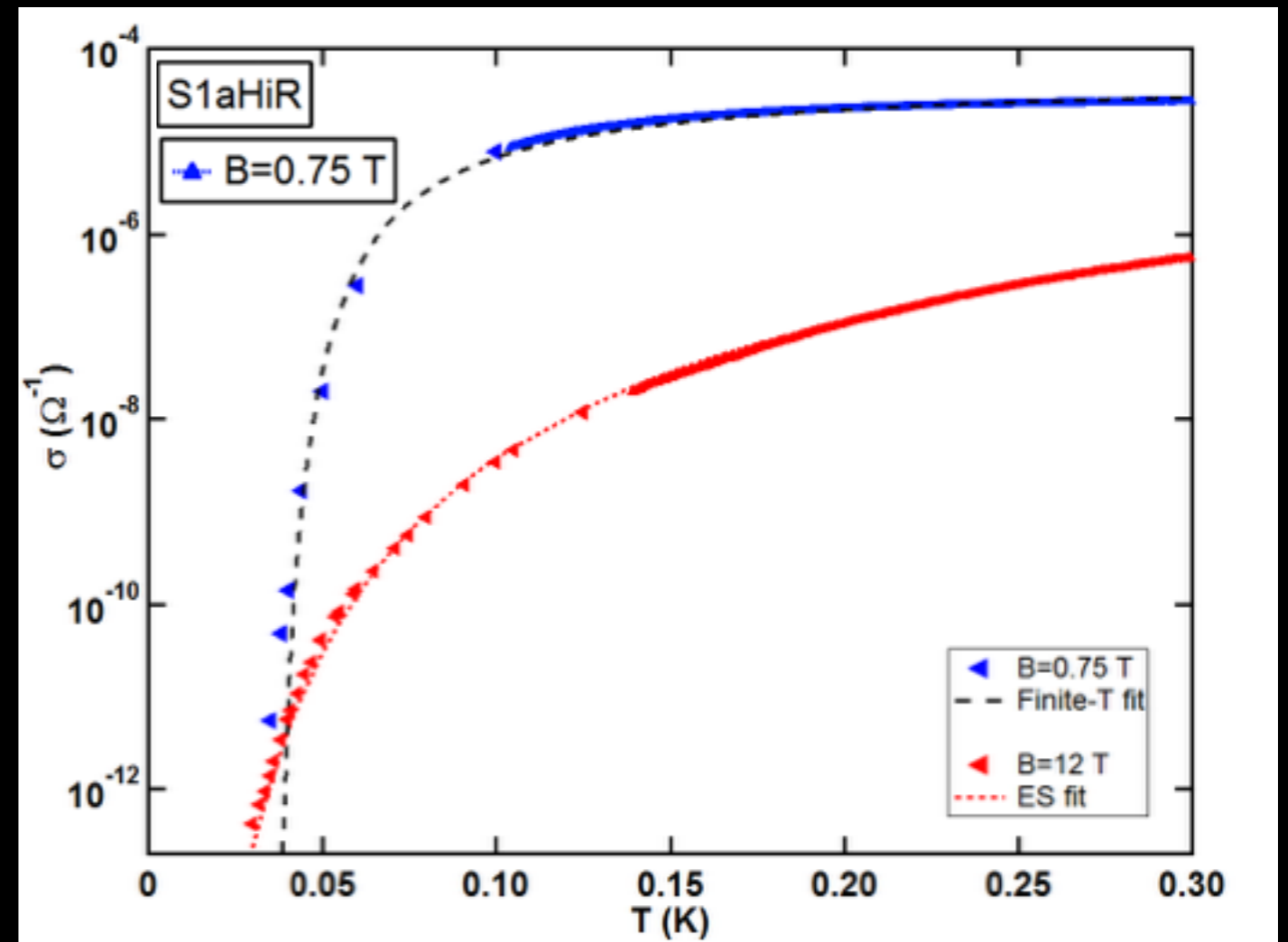
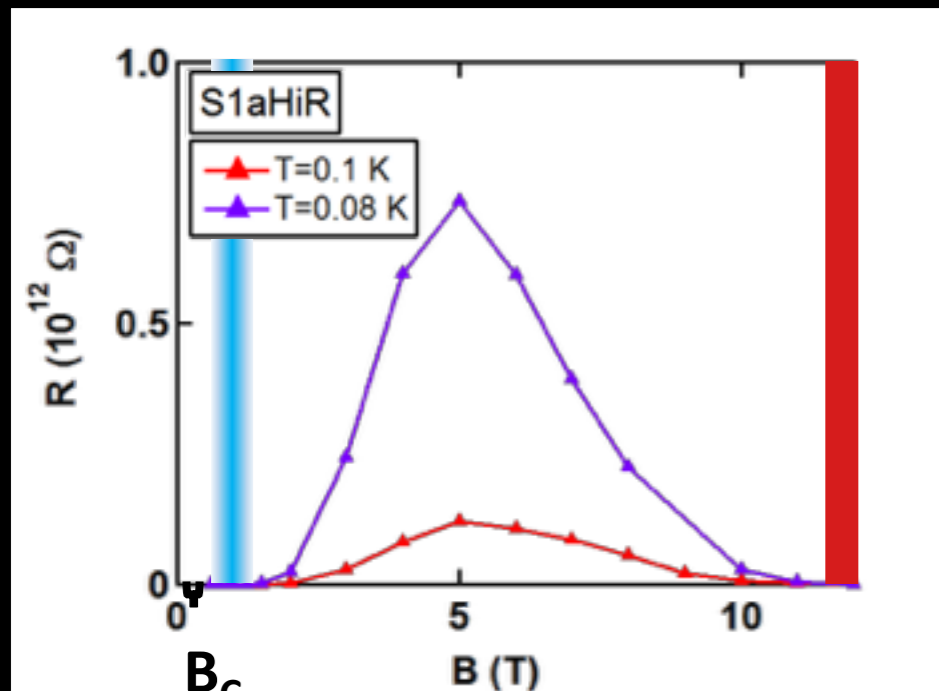
$$R(T) = R_{ES} \times \exp \left\{ \left(\frac{T_{ES}}{T} \right)^{1/2} \right\}$$



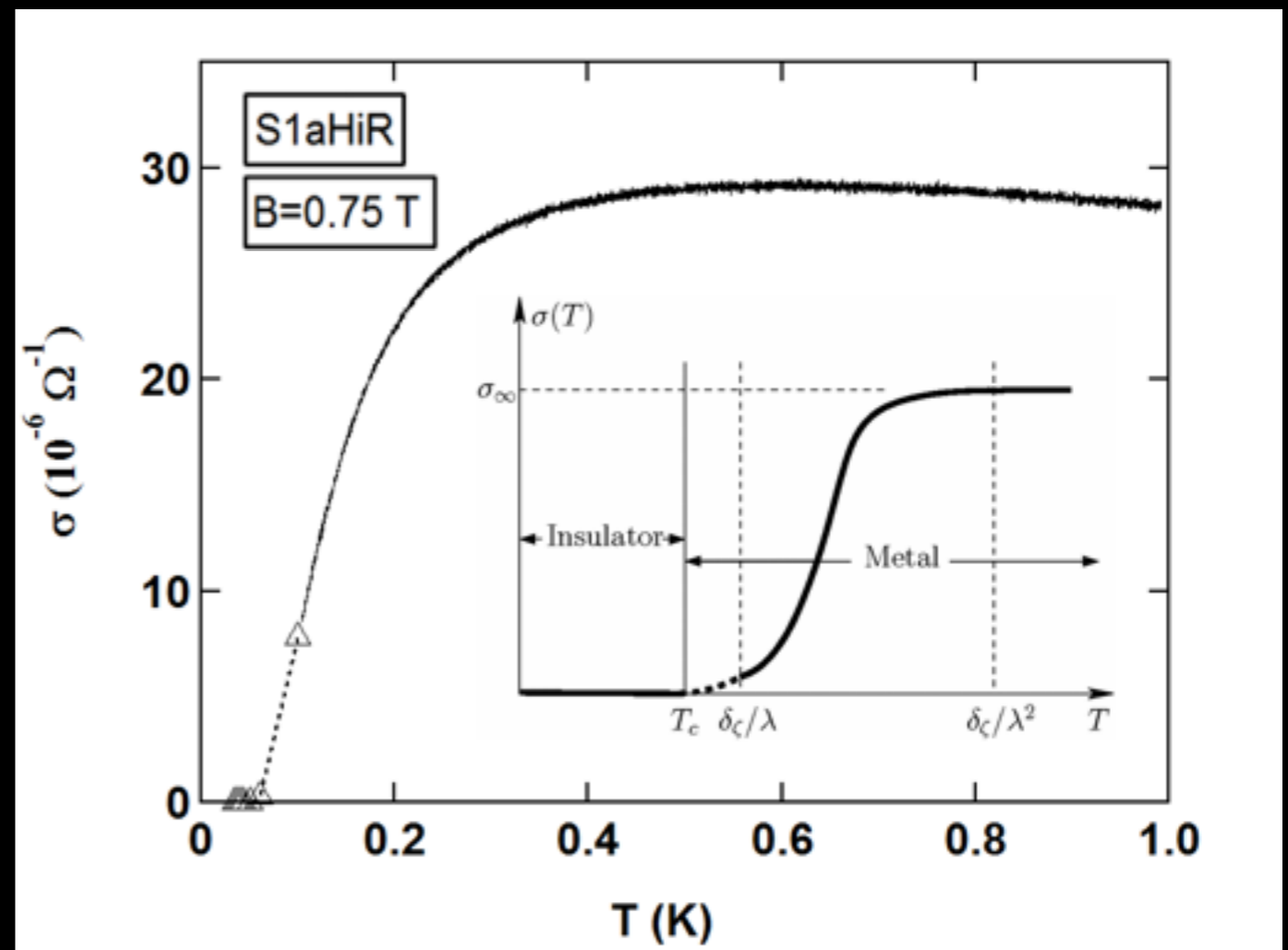
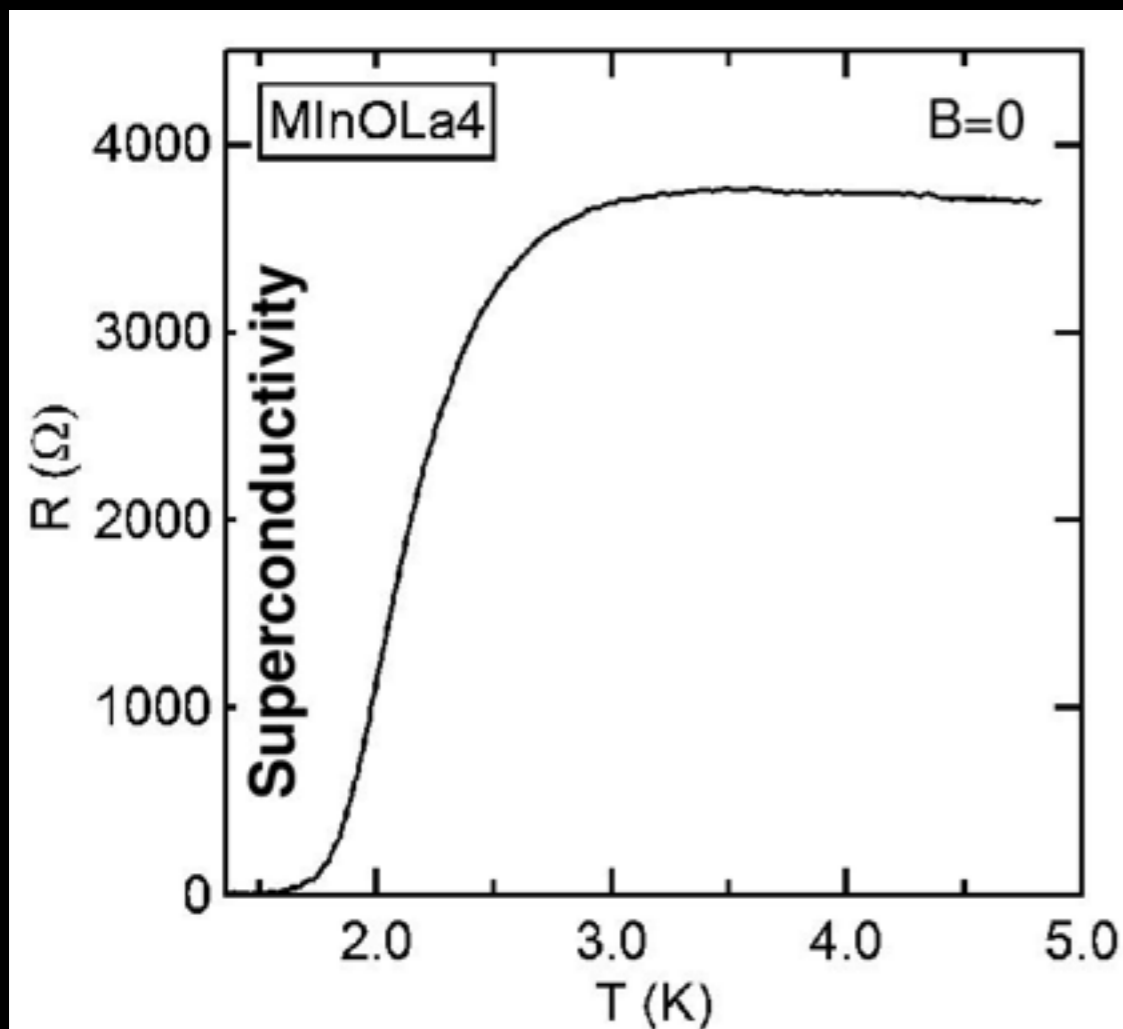
Resistivity Near $B_c B_c$



Conductivity

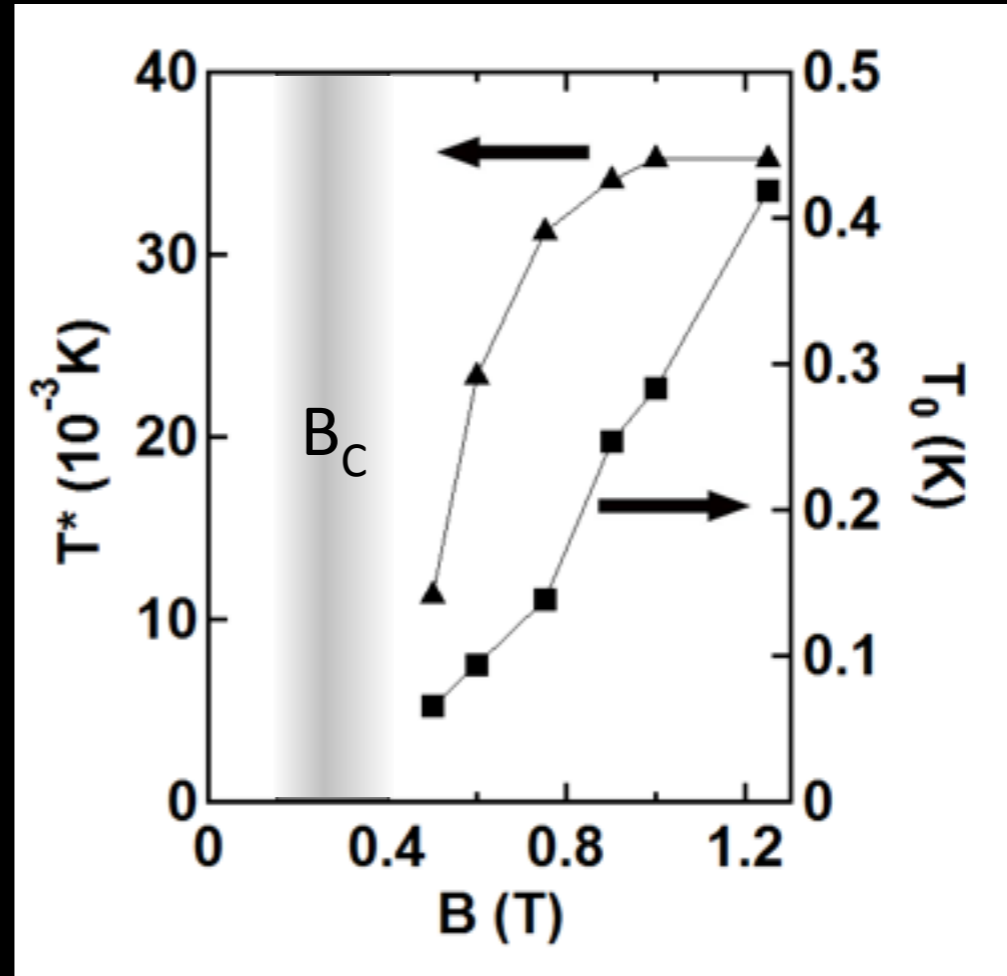


Many-body localization?



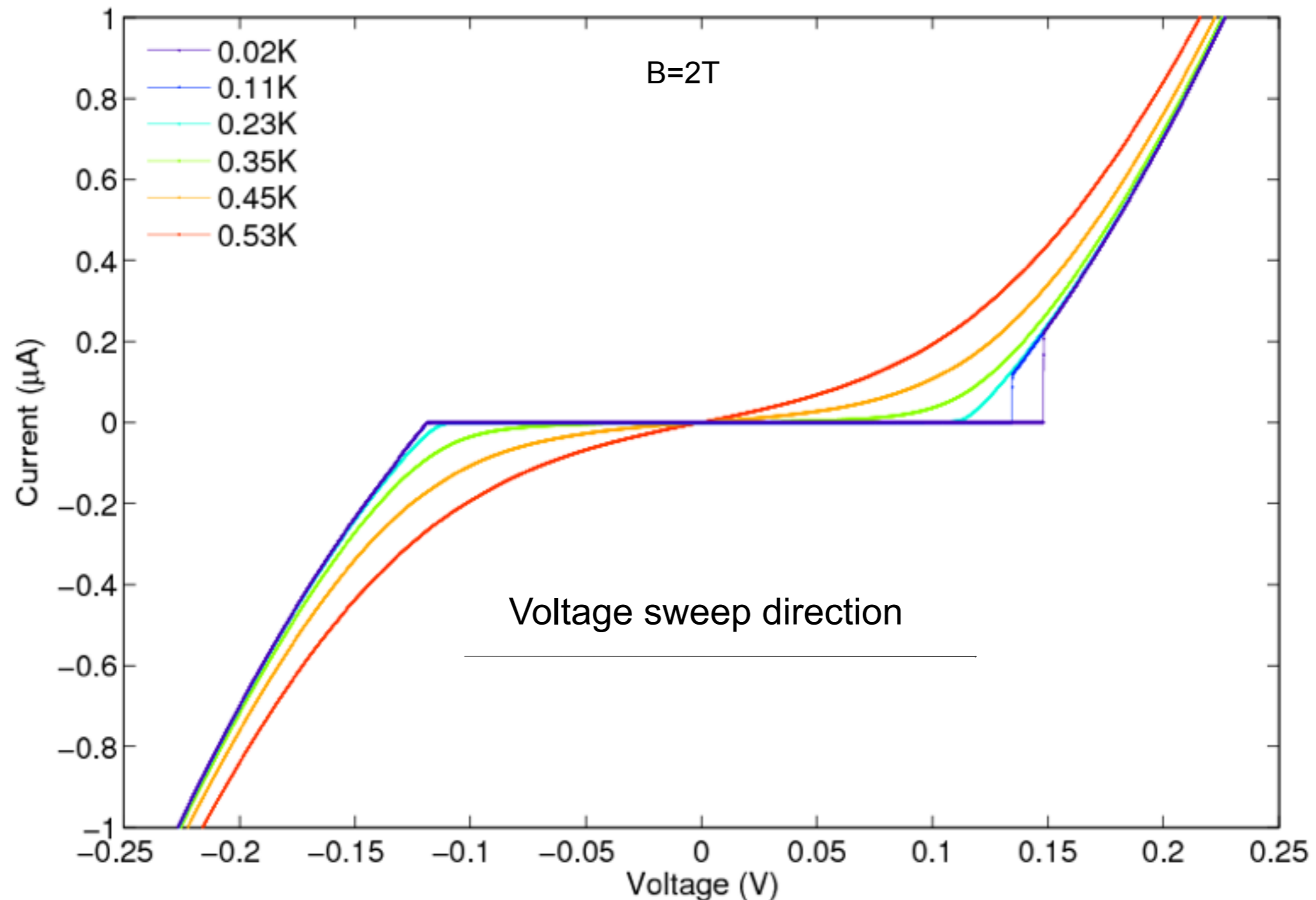
Thank you.

Fit Parameters

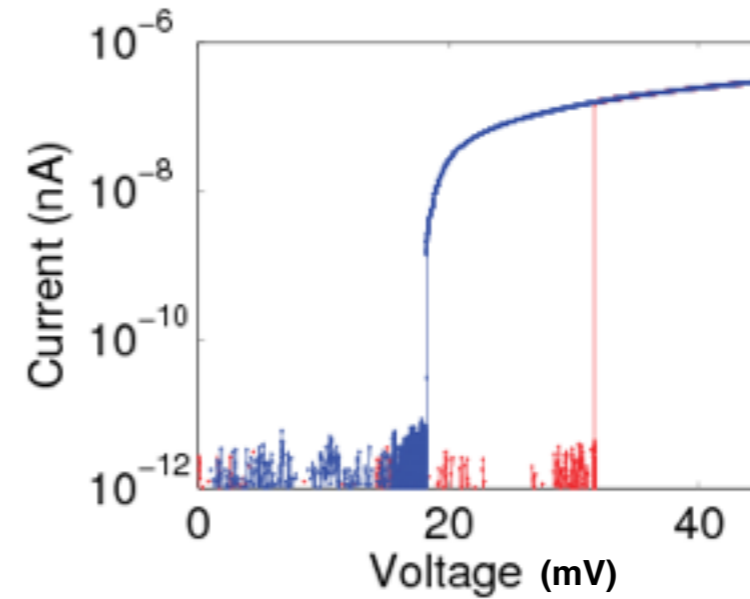
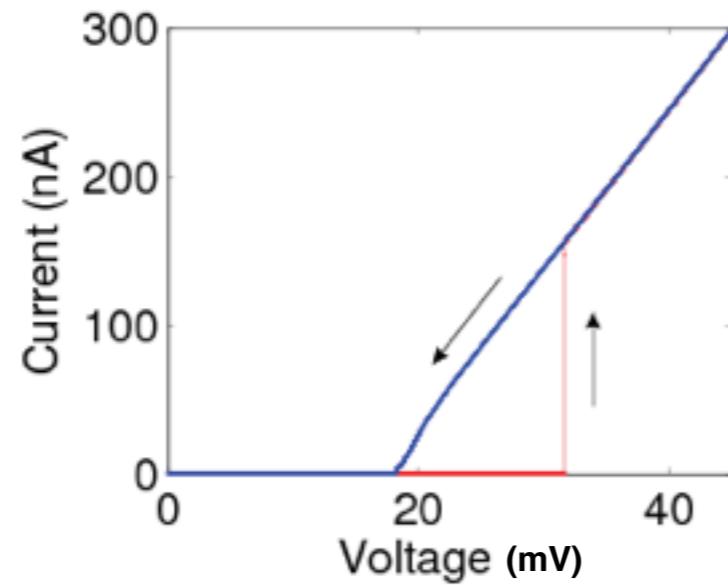


$$\sigma(T) = \sigma_0 \times \exp\left\{-\frac{T_0}{T - T^*}\right\}$$

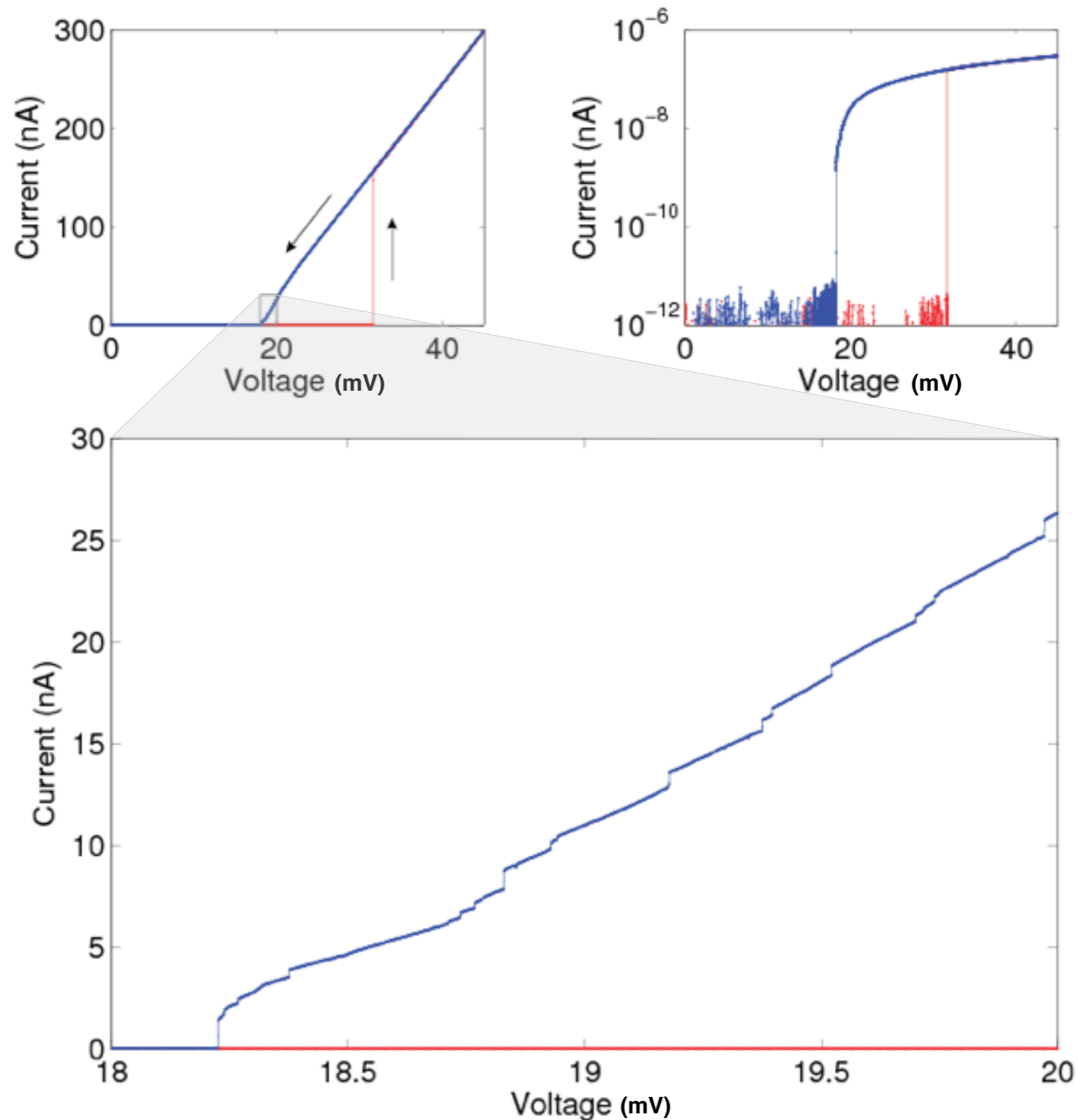
In the strongly insulating regime the I-V curves exhibit a voltage threshold



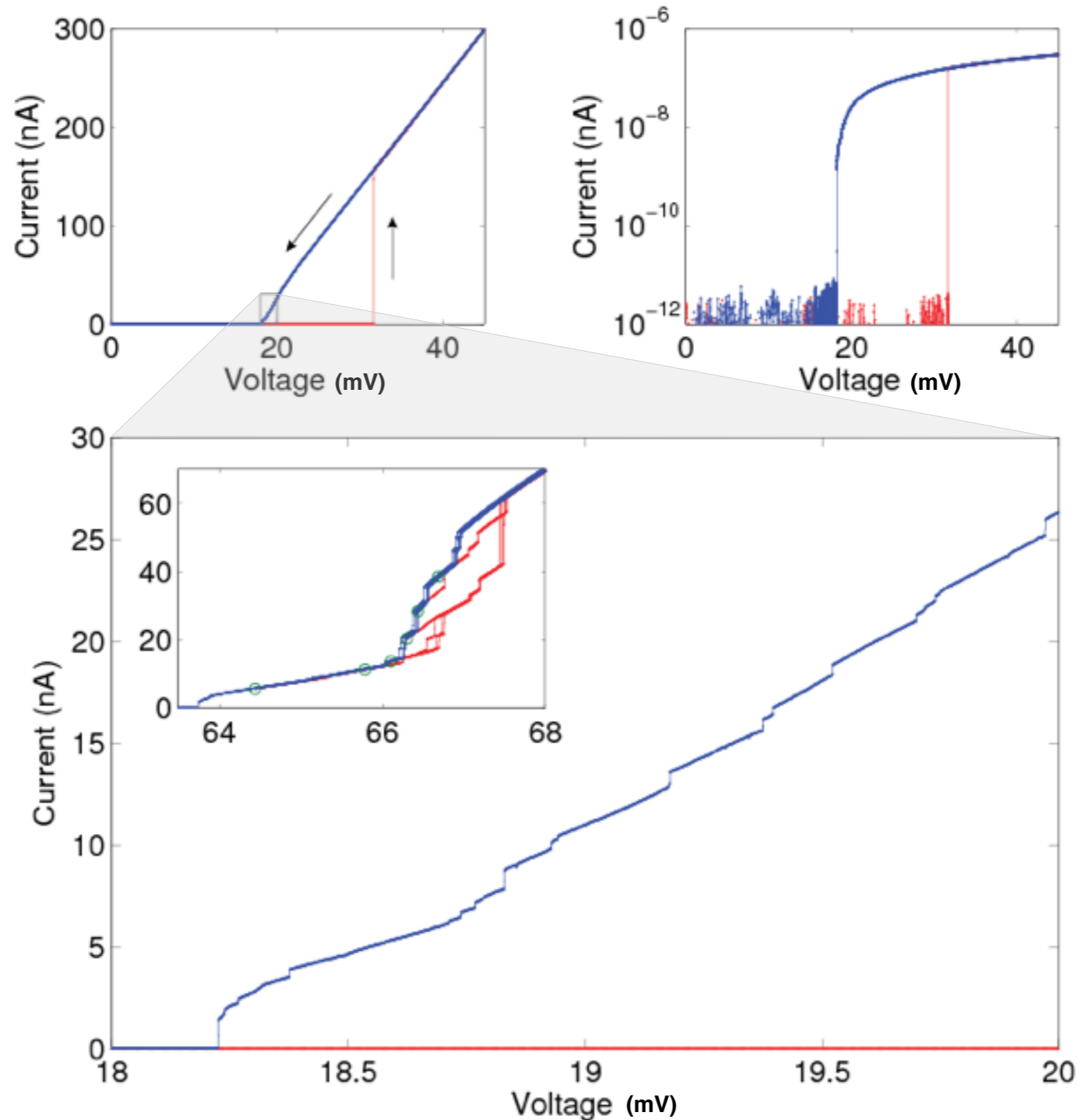
At low temperatures ($T < 0.2$ K), the threshold is accompanied by switching and hysteresis



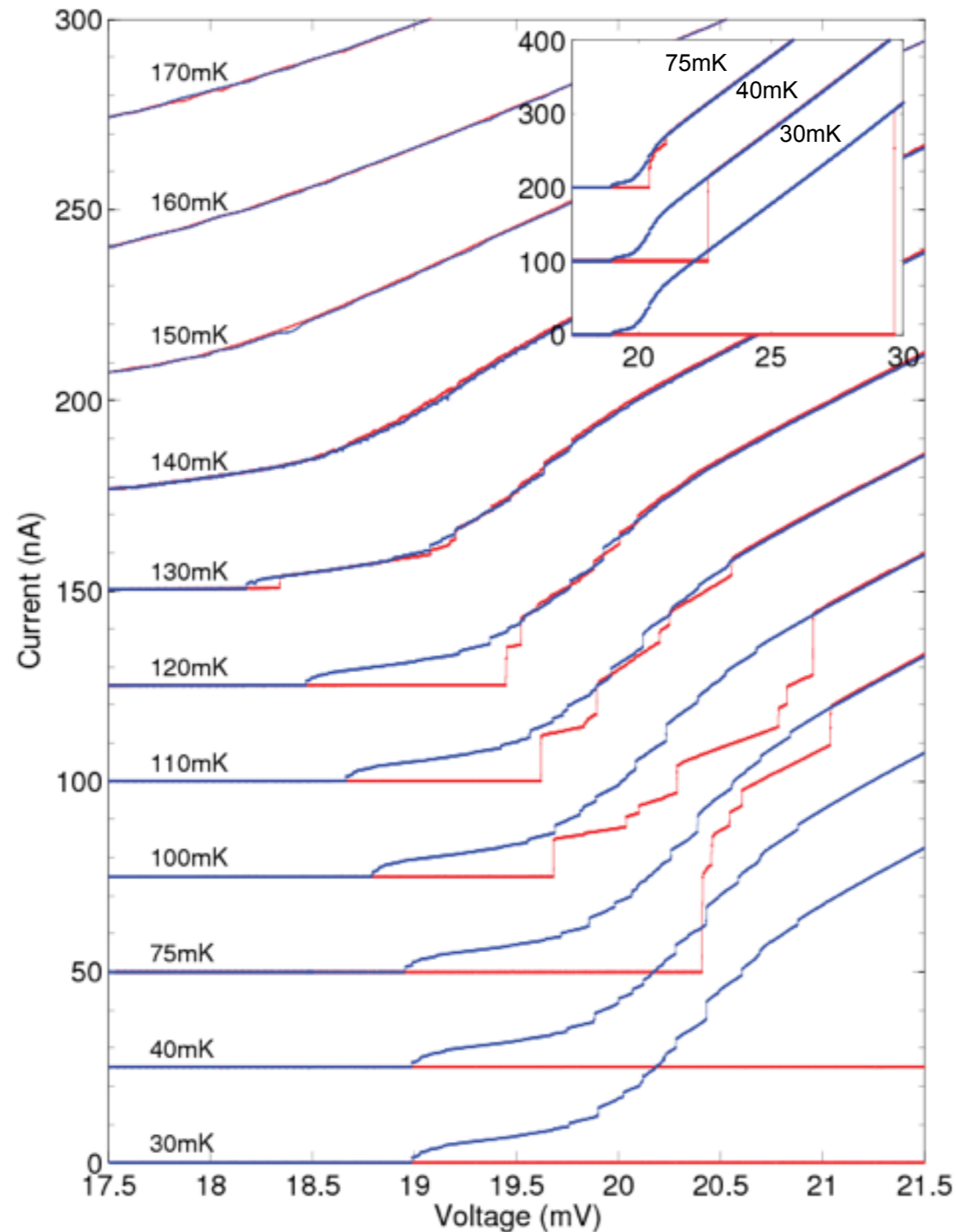
A series of discontinuous steps in current is revealed in the negative sweep direction



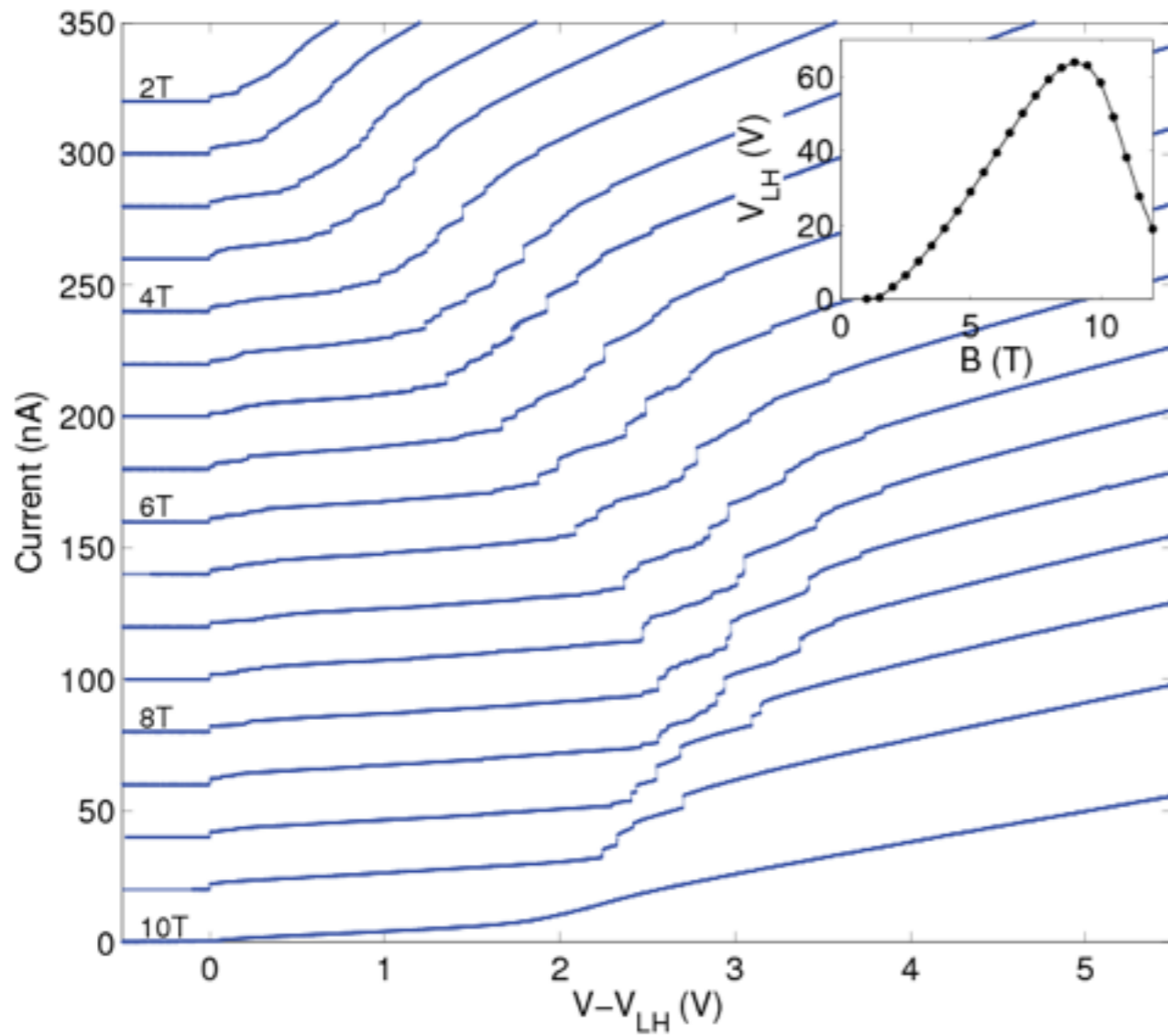
The current steps are accompanied by a complex hysteresis phenomenon



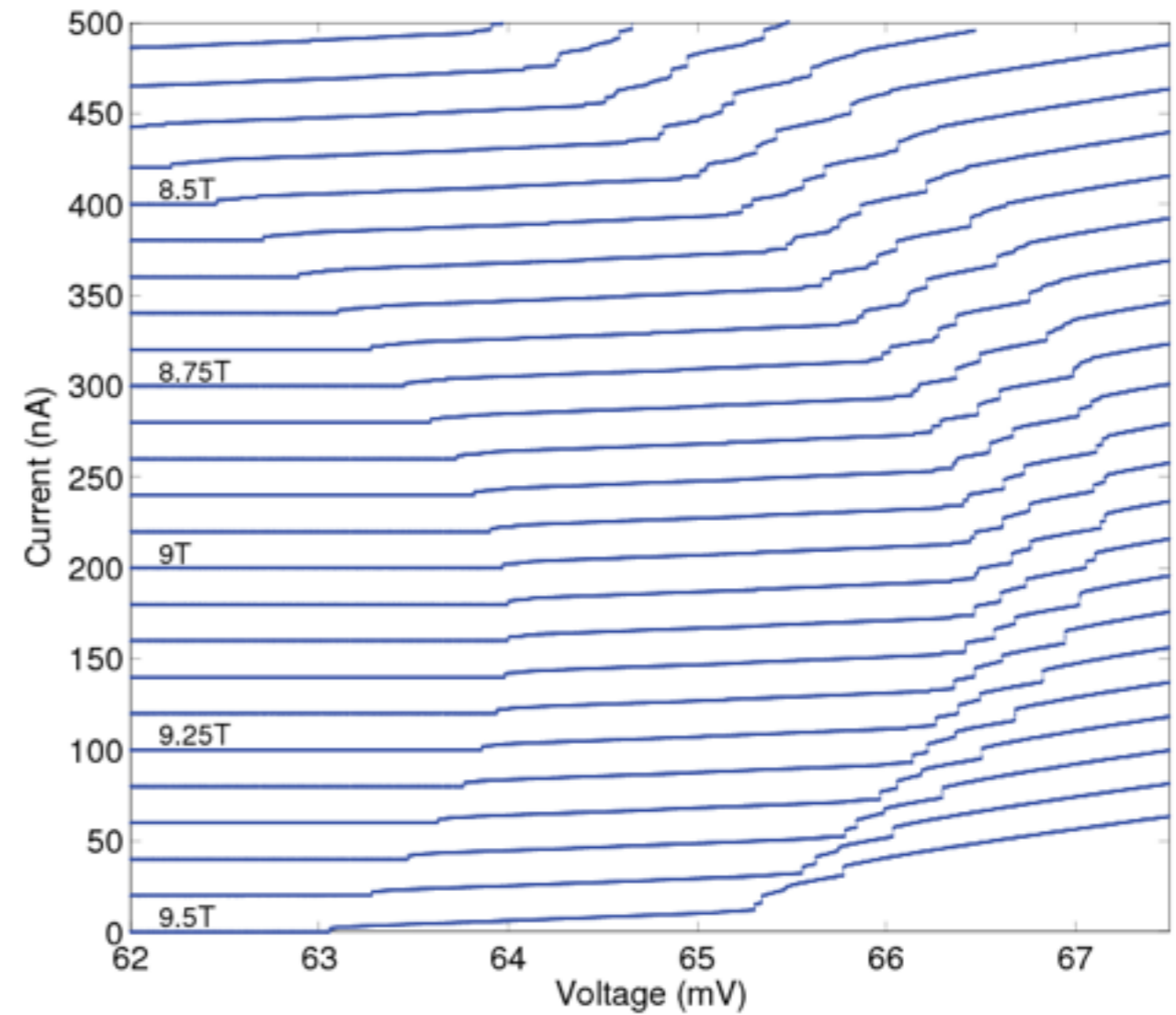
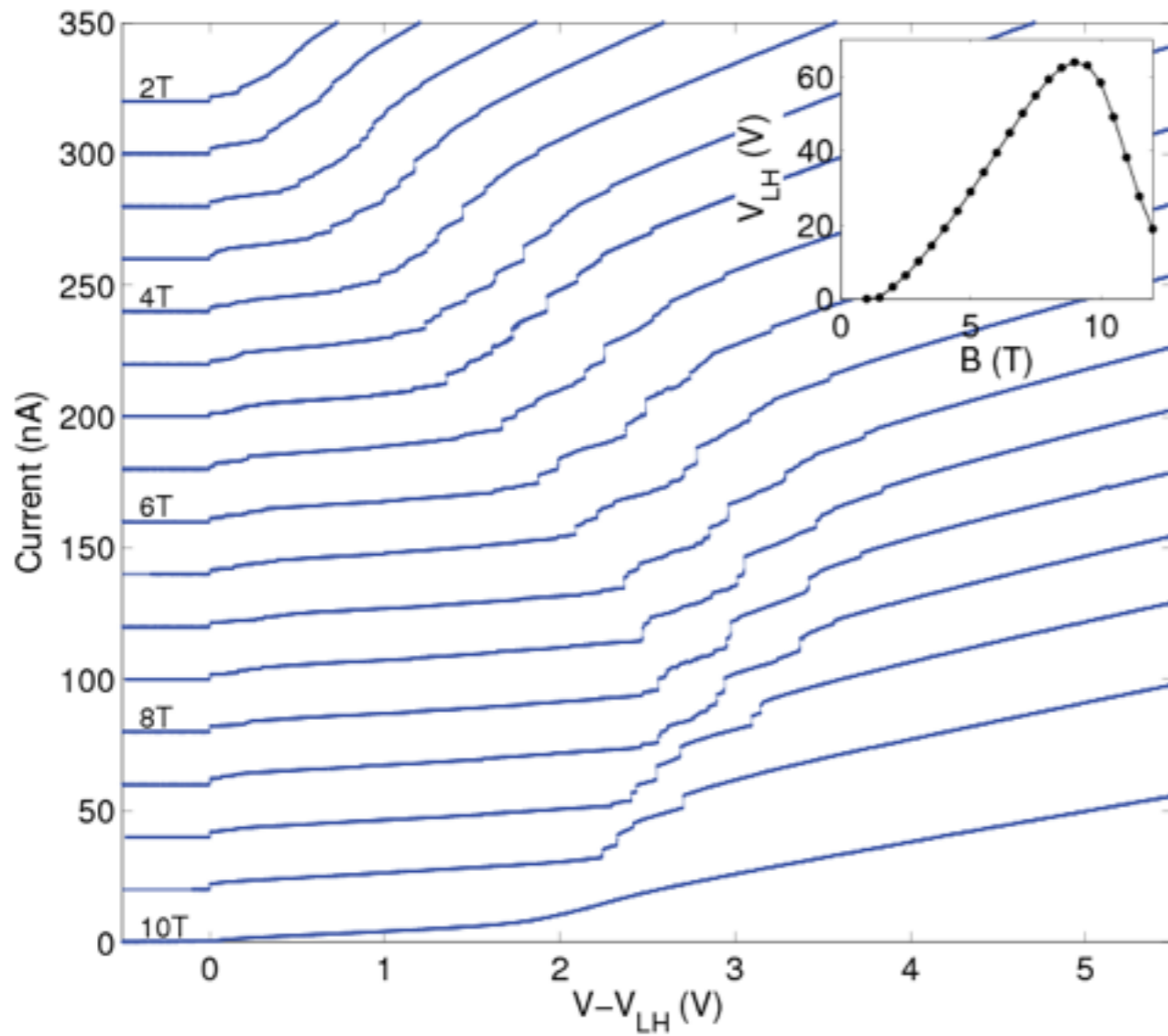
As T increases, the steps pattern deforms gradually



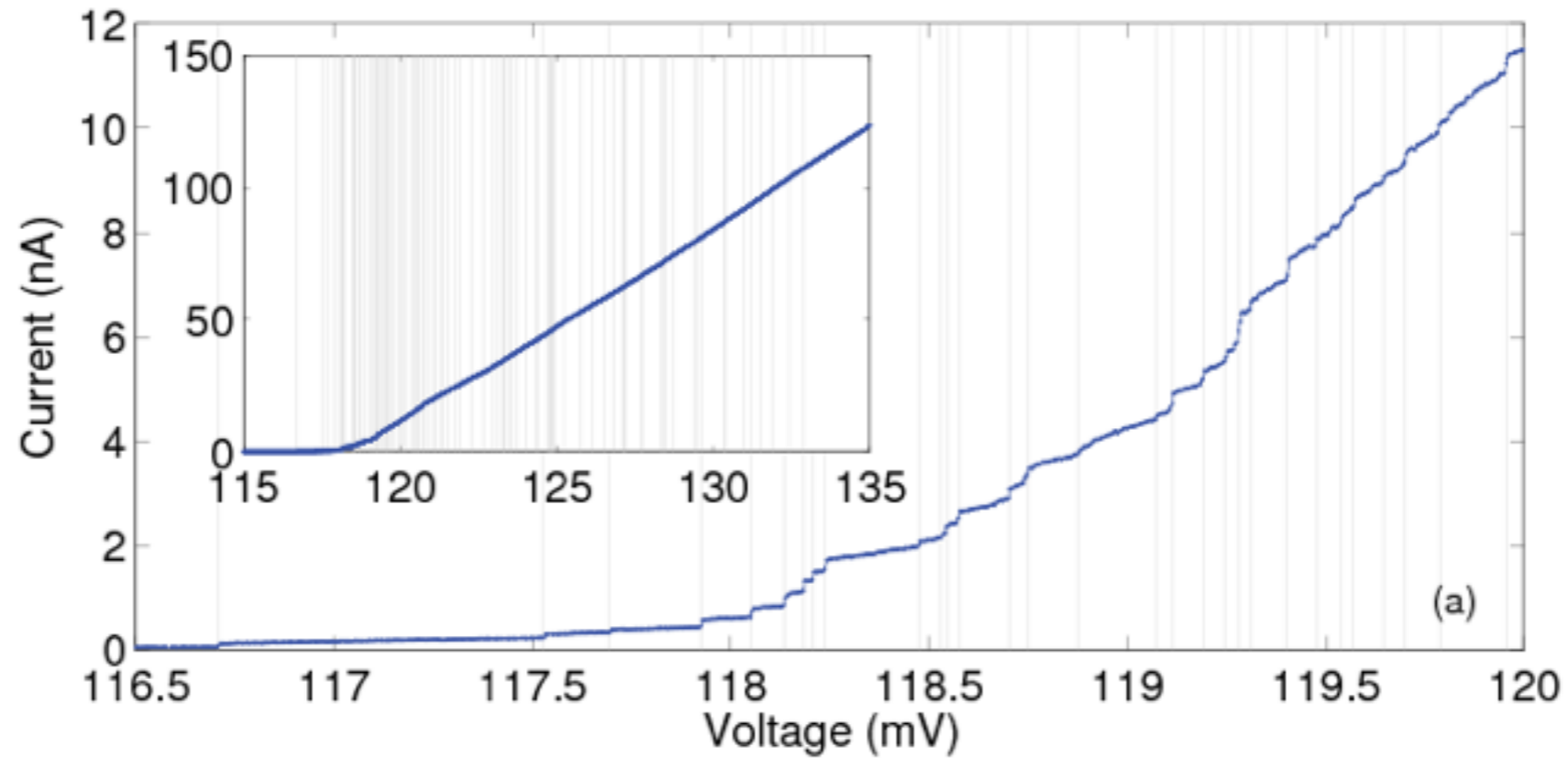
The magnetic field is seen to have only a moderate influence on the hierarchy of the steps pattern



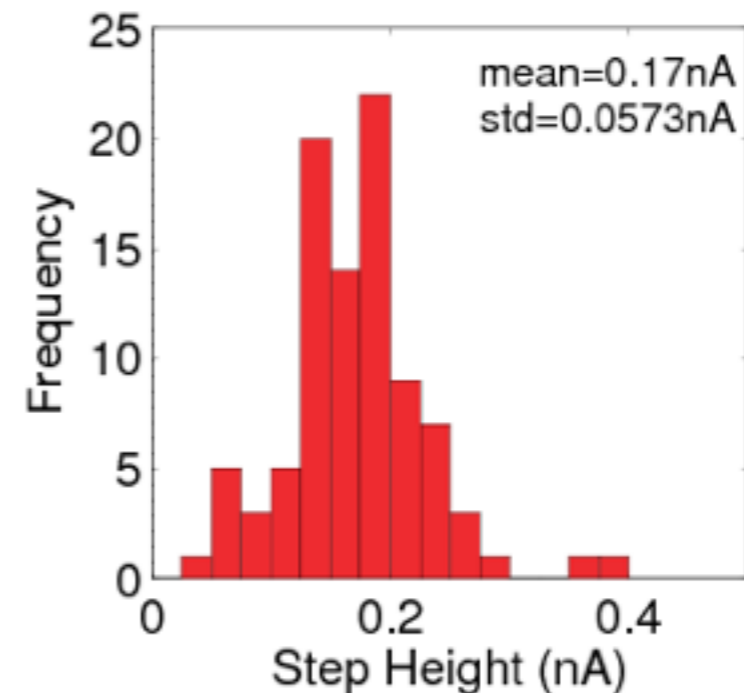
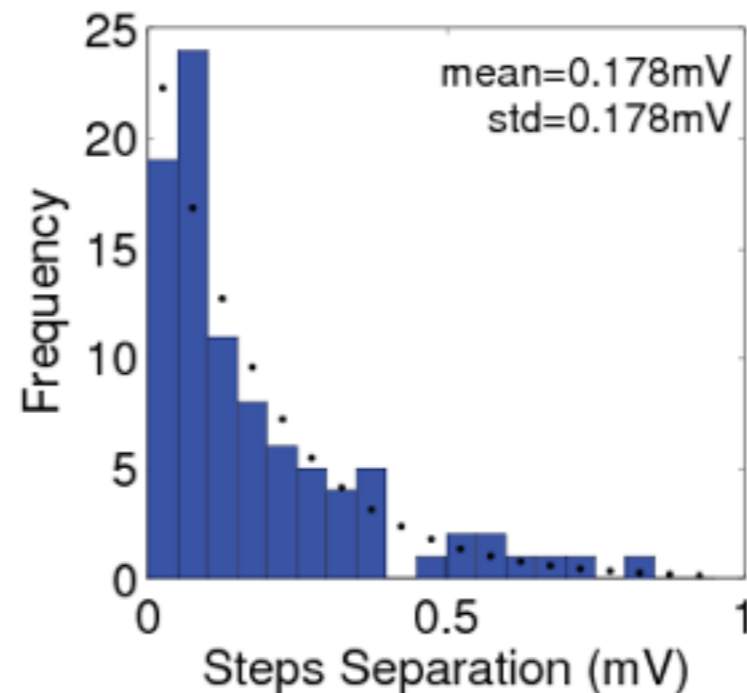
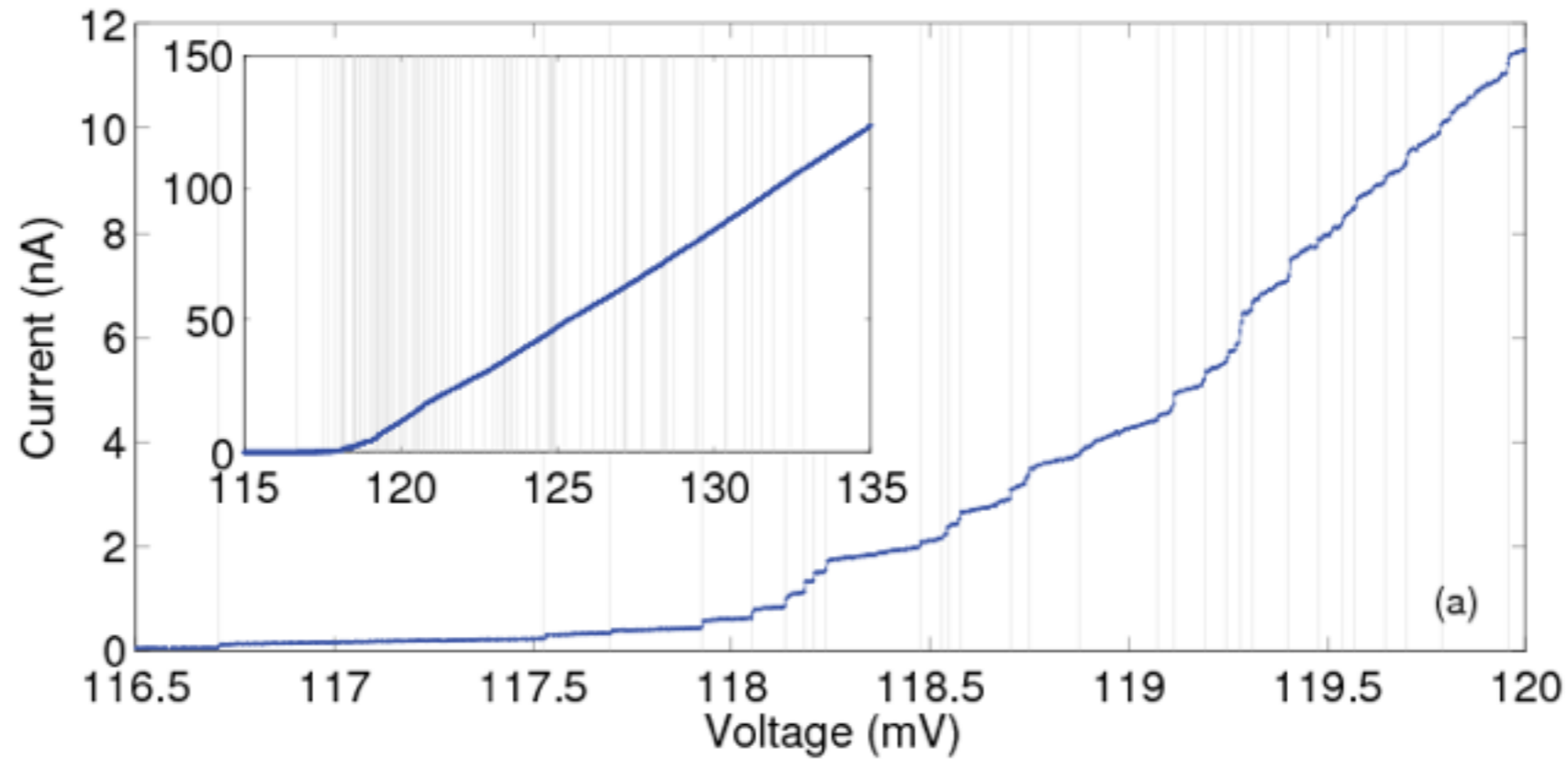
The magnetic field is seen to have only a moderate influence on the hierarchy of the steps pattern



Films with higher degree of disorder show larger number of steps

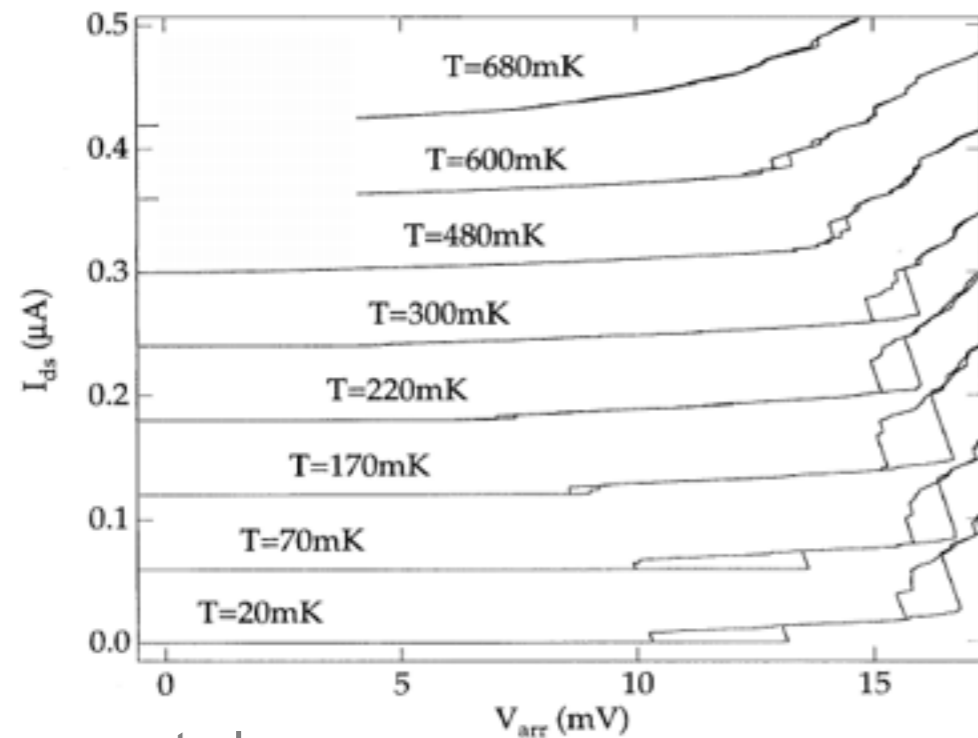
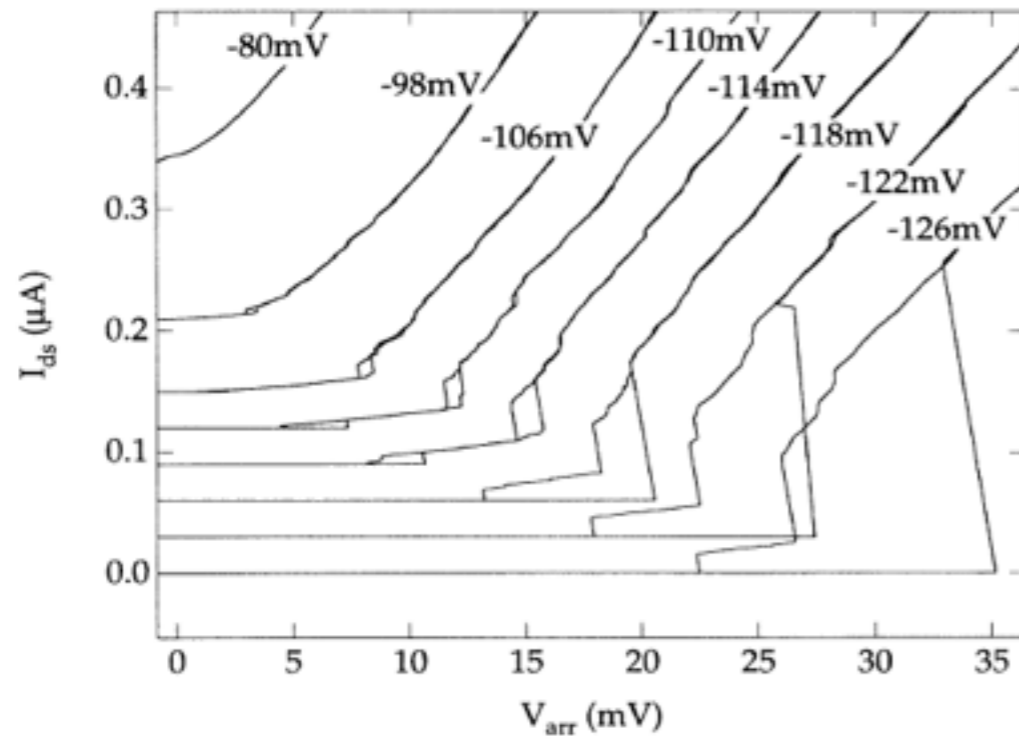


Films with higher degree of disorder show larger number of steps

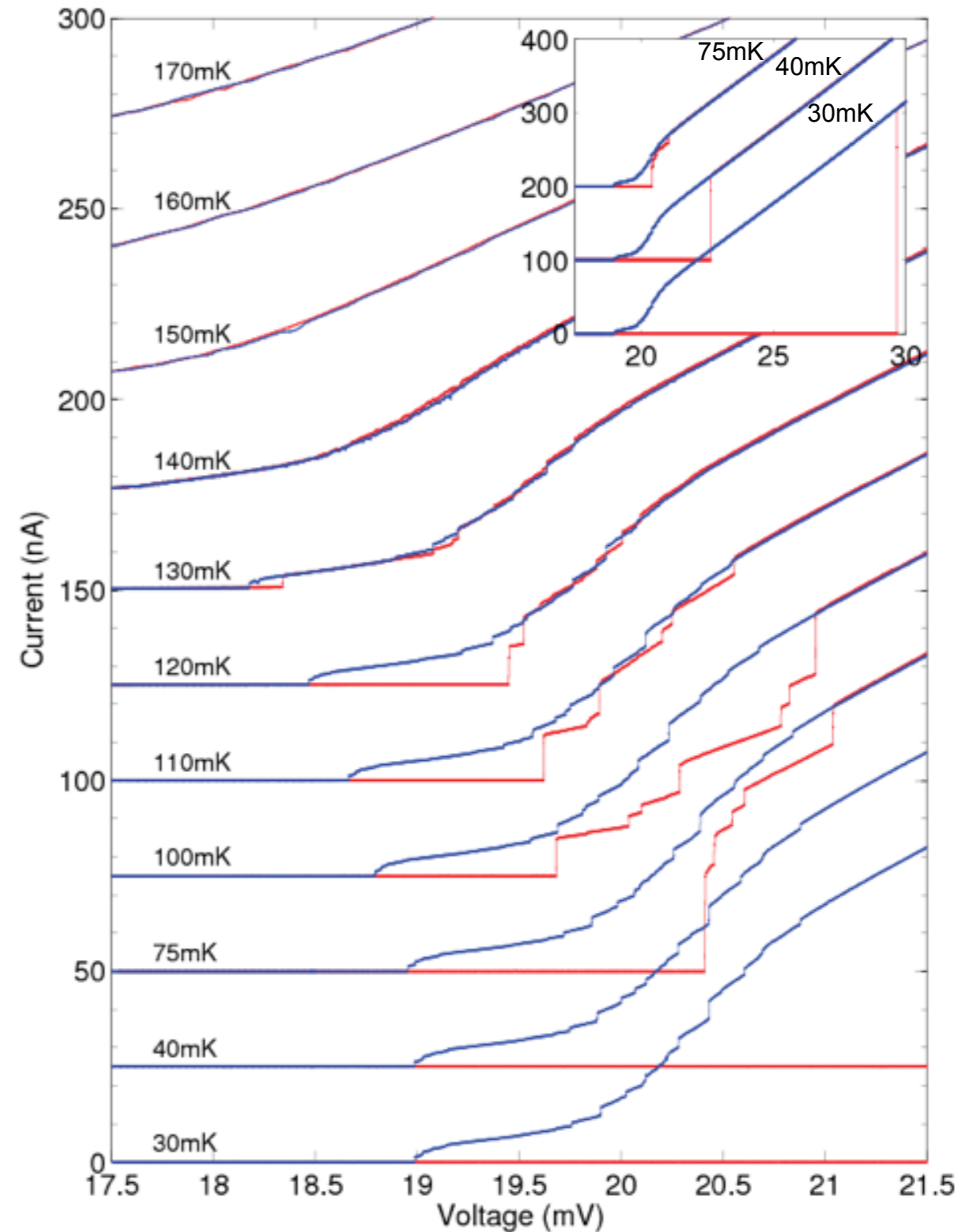


Experiment: QD array vs Thin Film

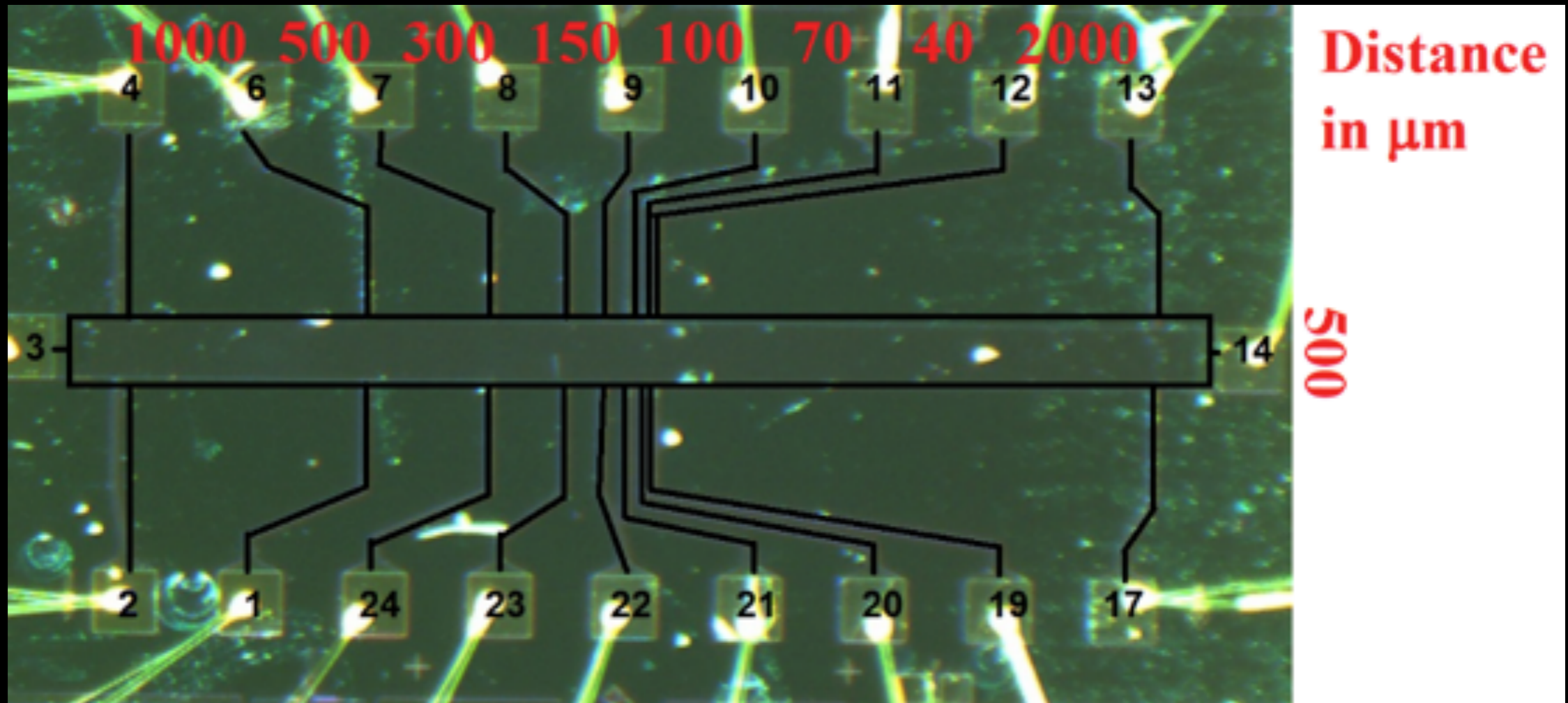
2D QD array

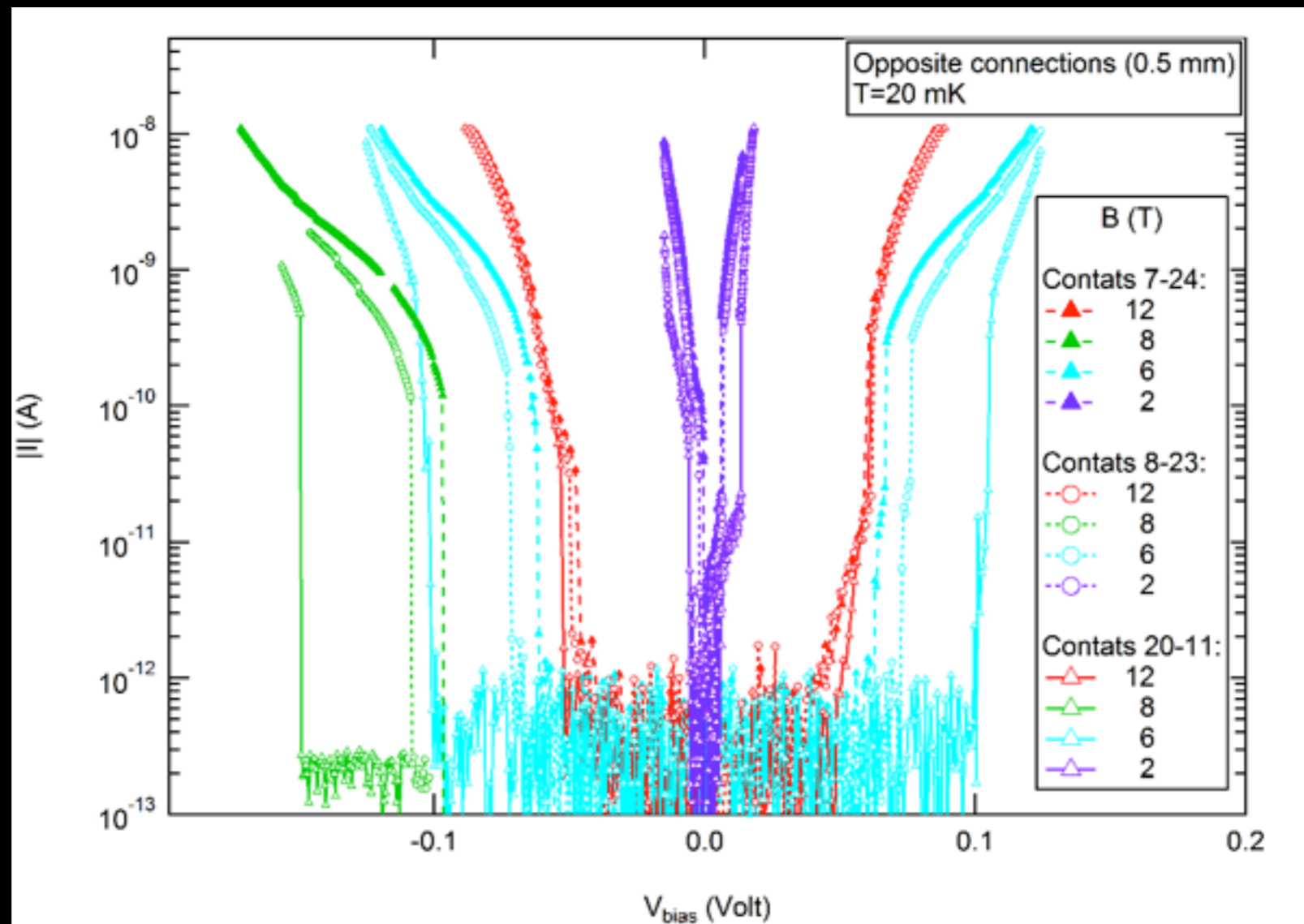


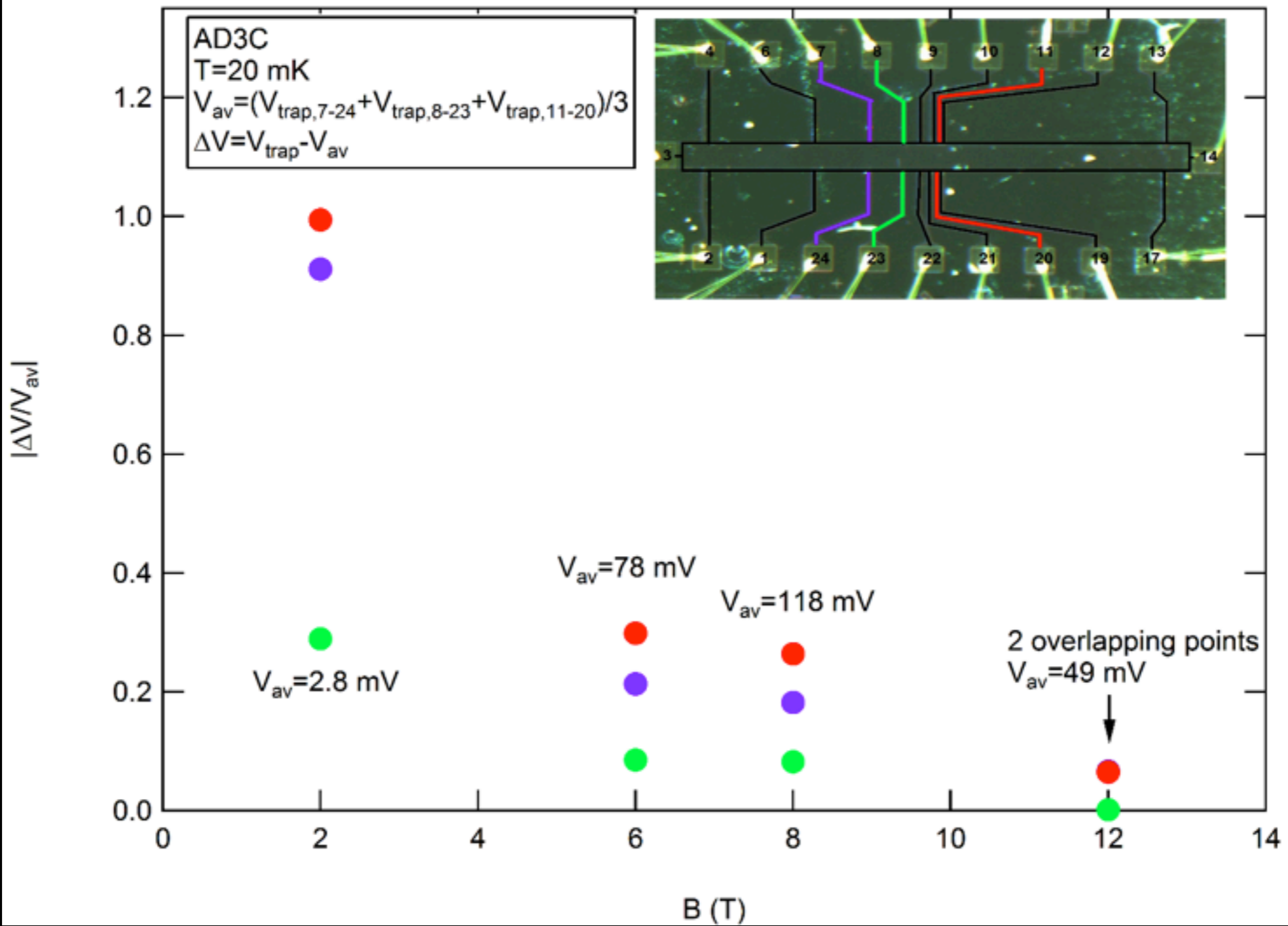
InO Thin Film



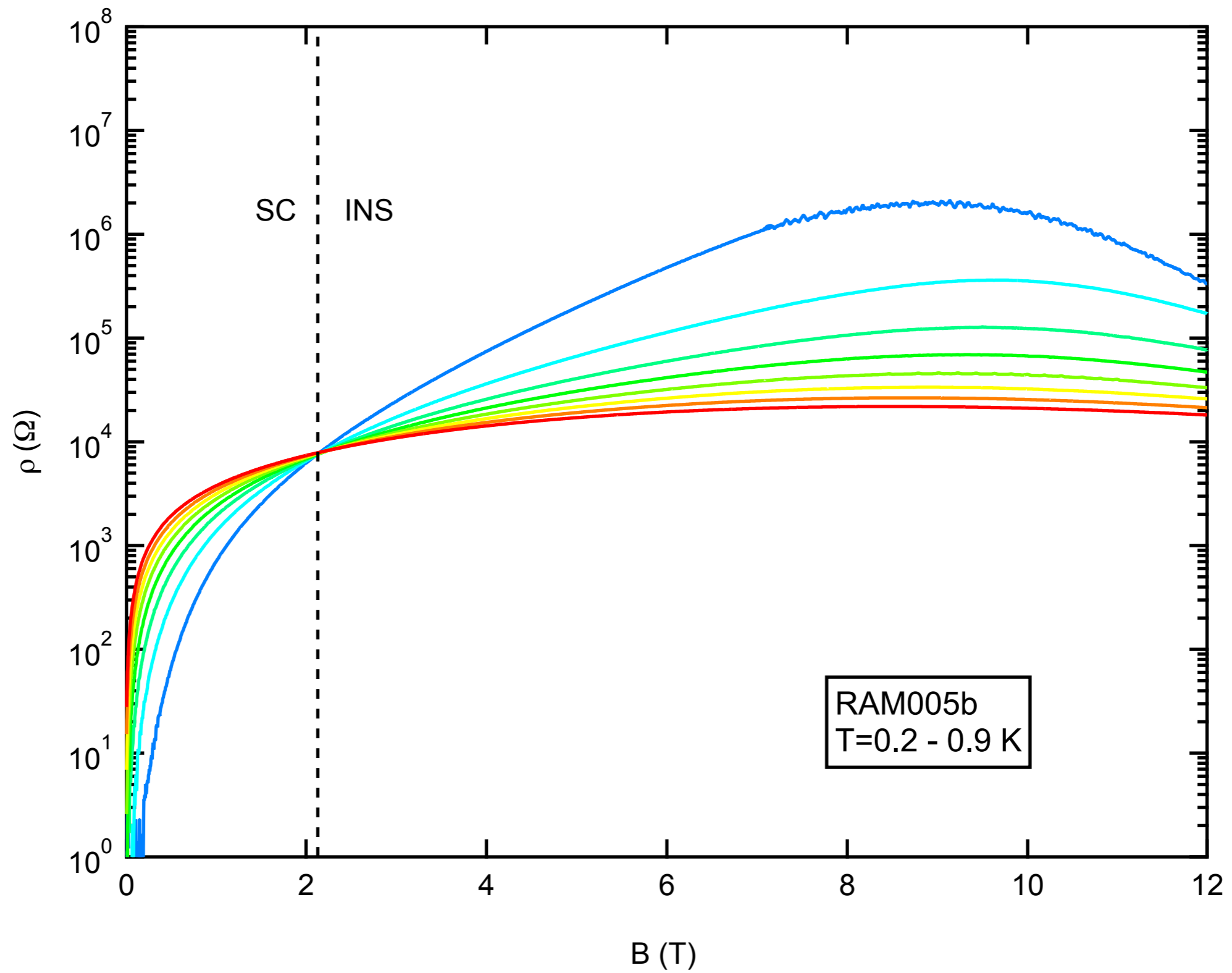
Inhomogeneity

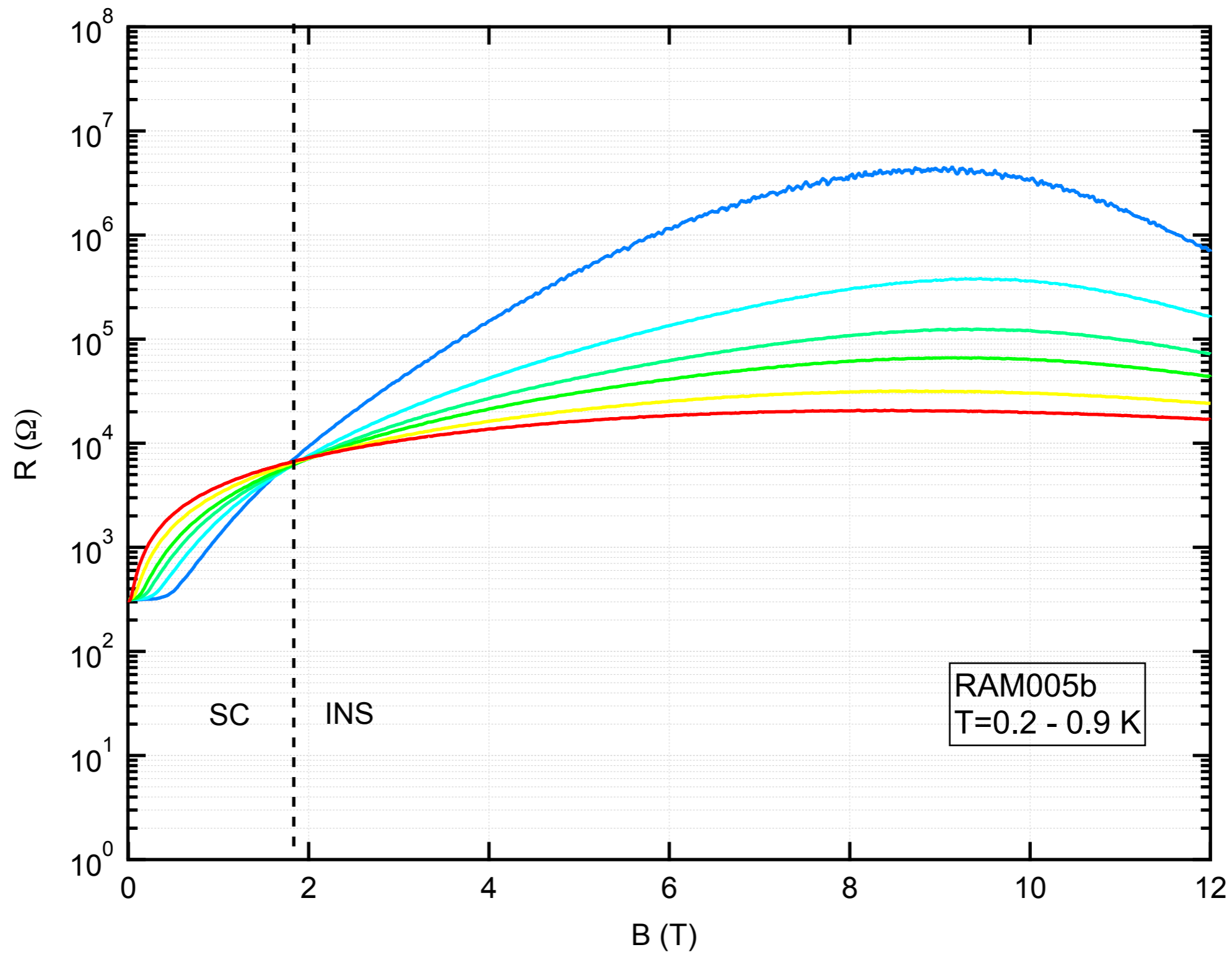


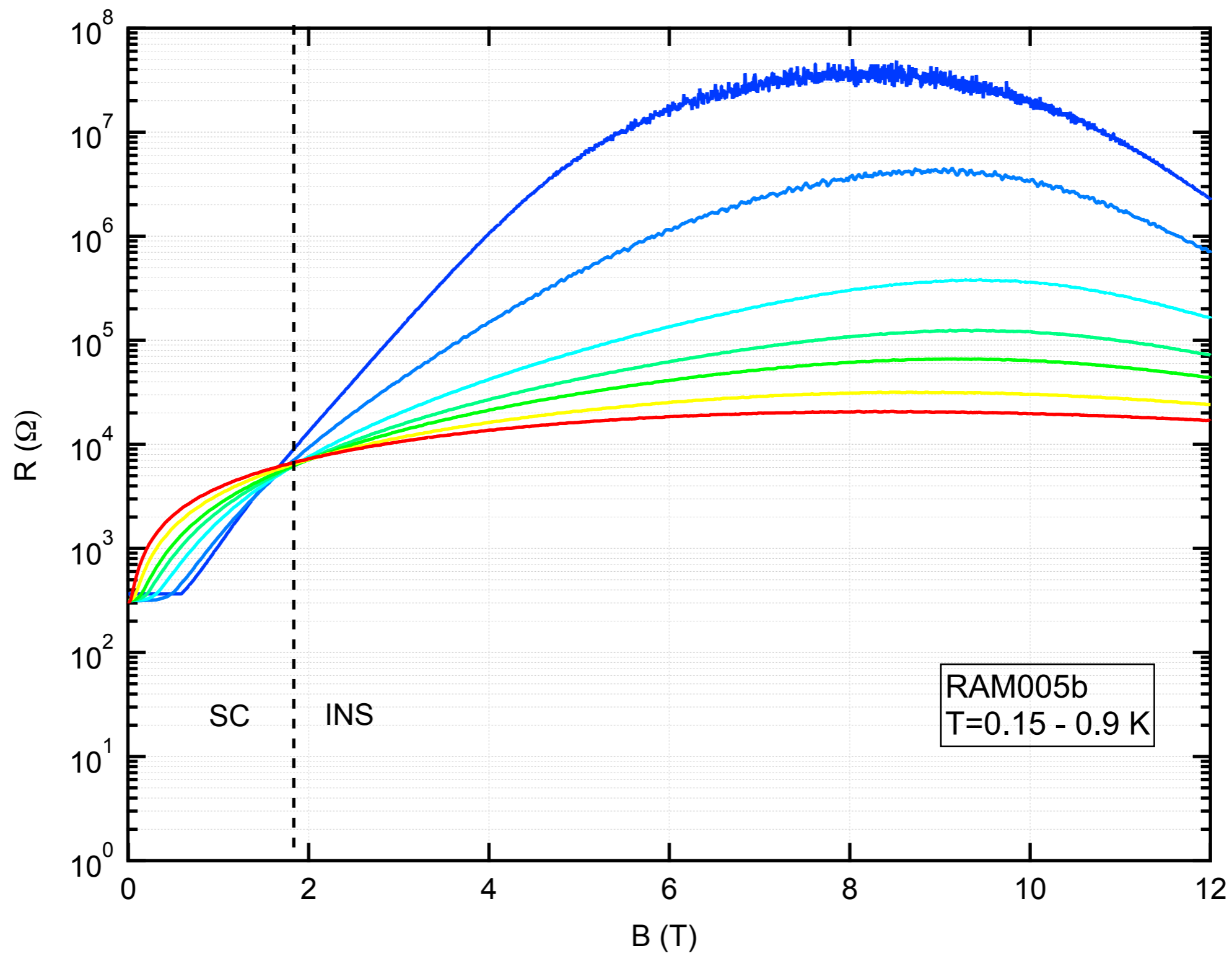


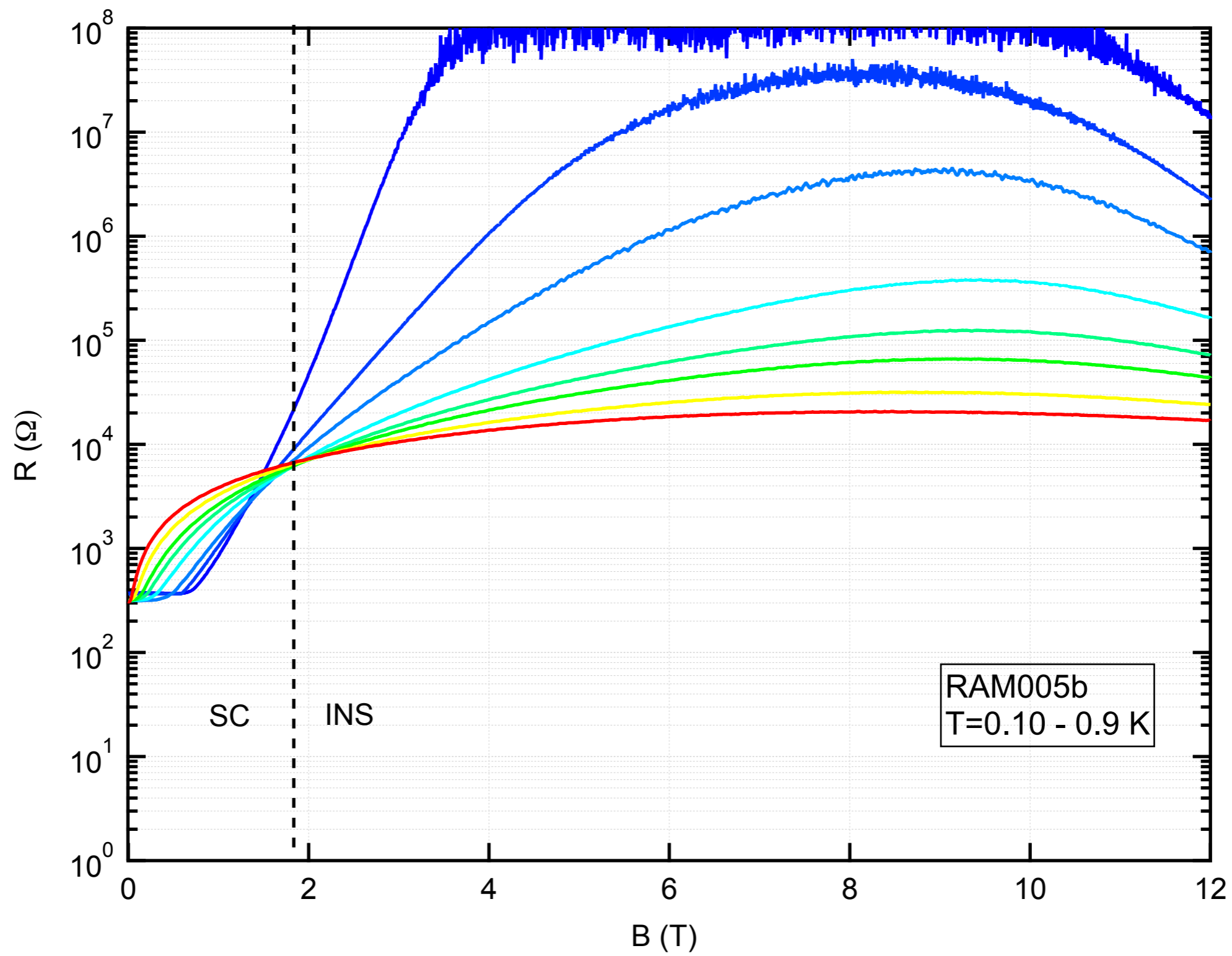


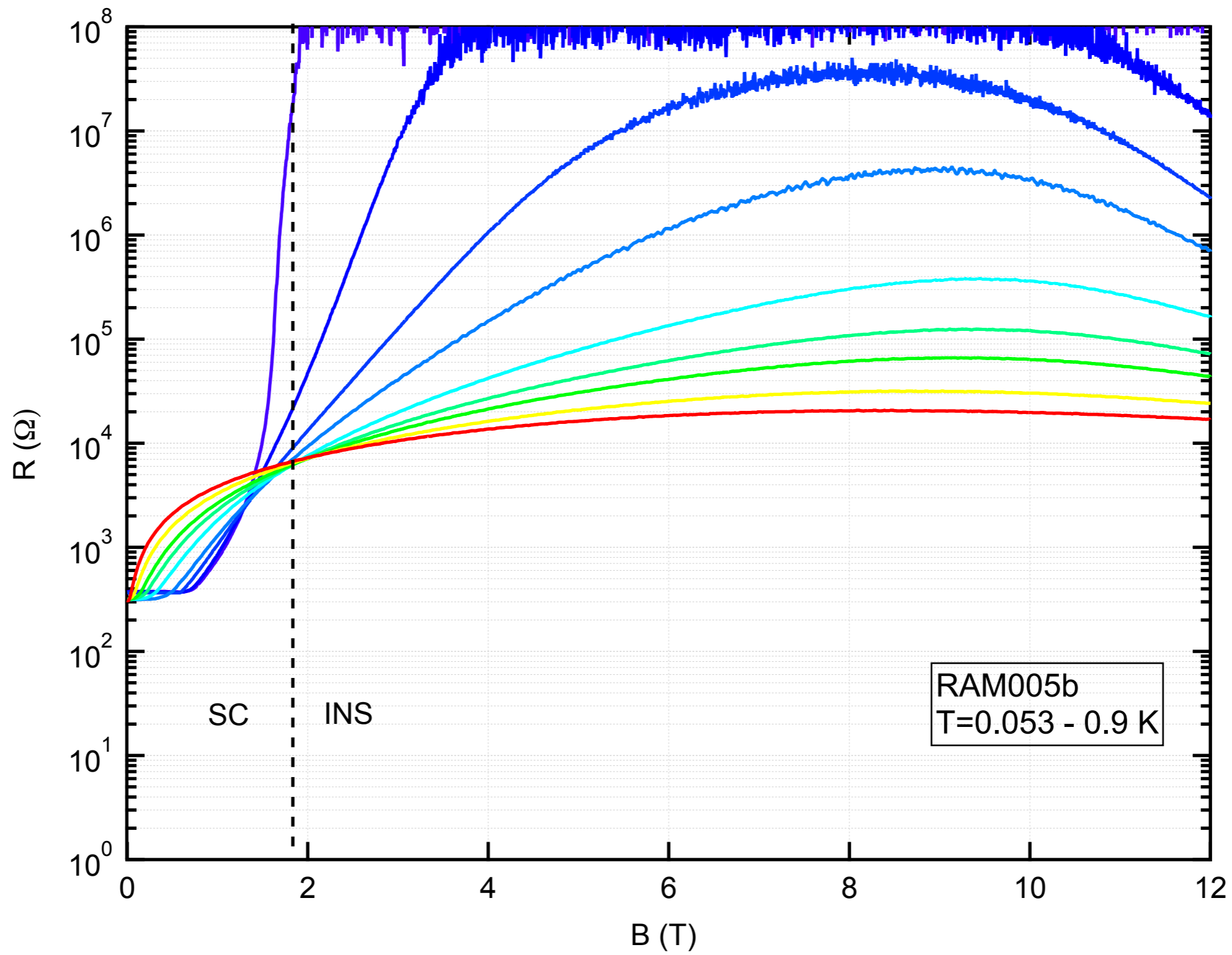
The SIT below $T=0.2$ K:

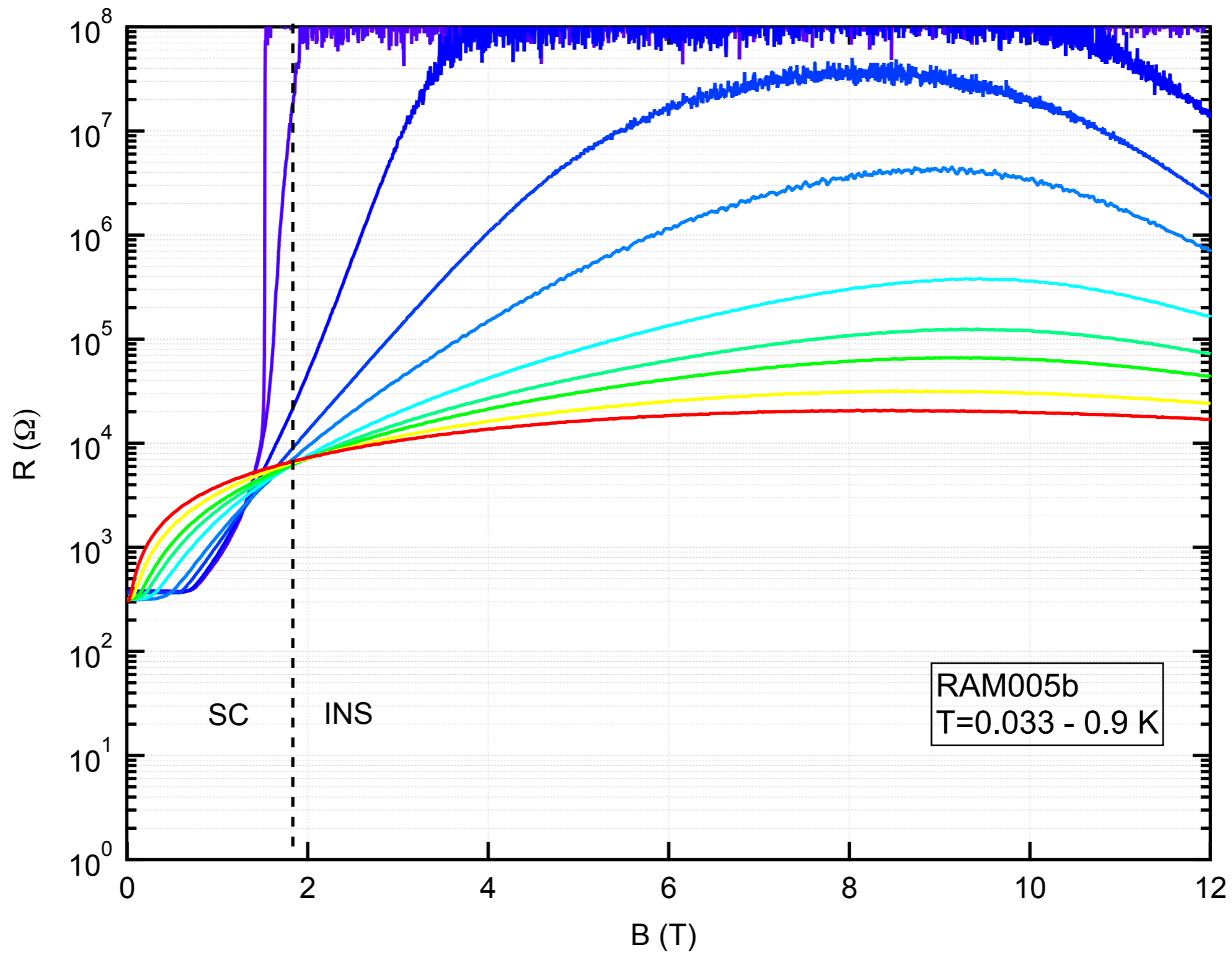


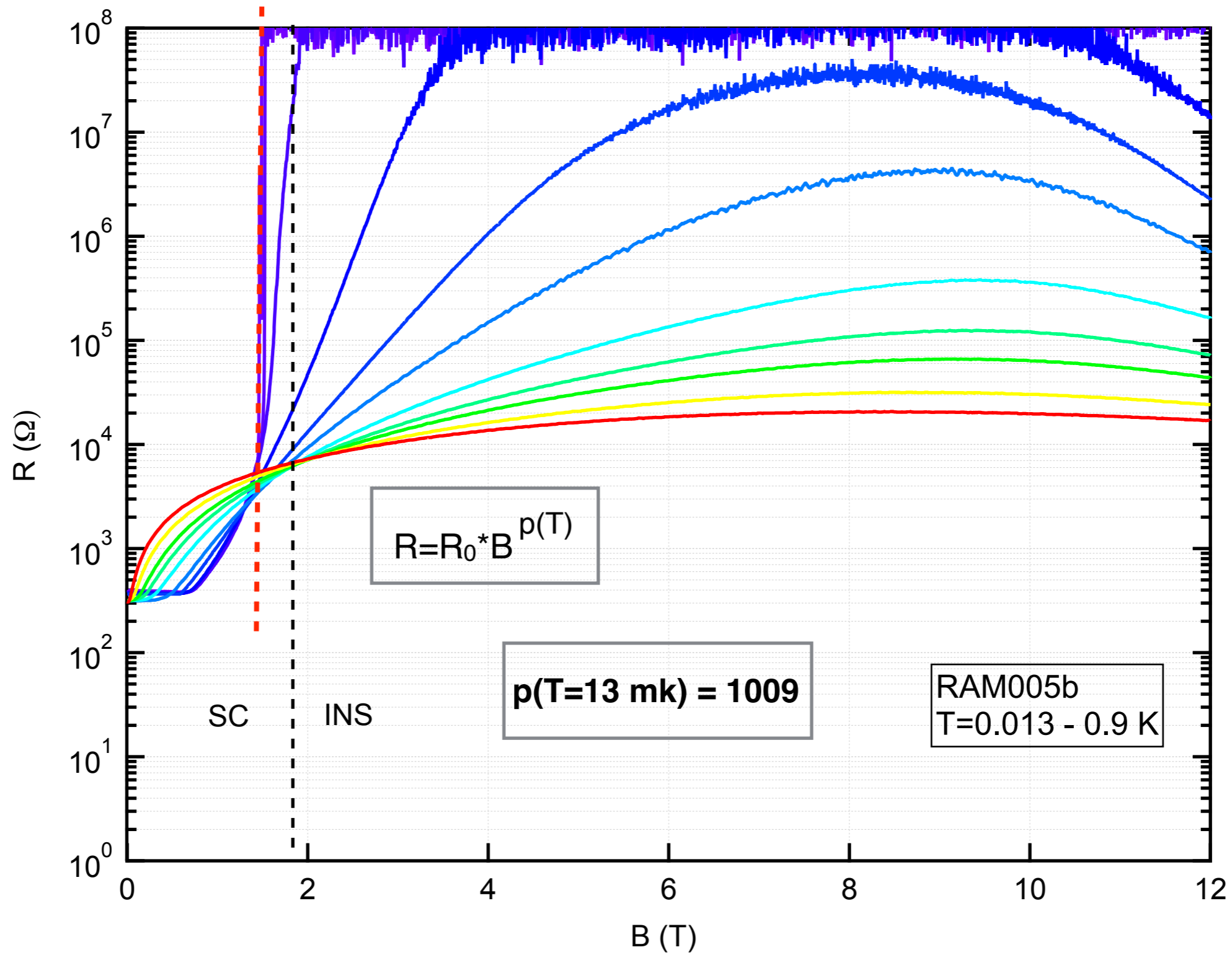












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