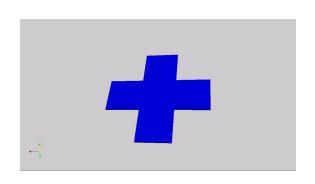
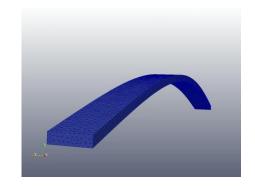
Modeling liquid crystal elastomers: from auto-origami to light-driven autonomous soft robotics



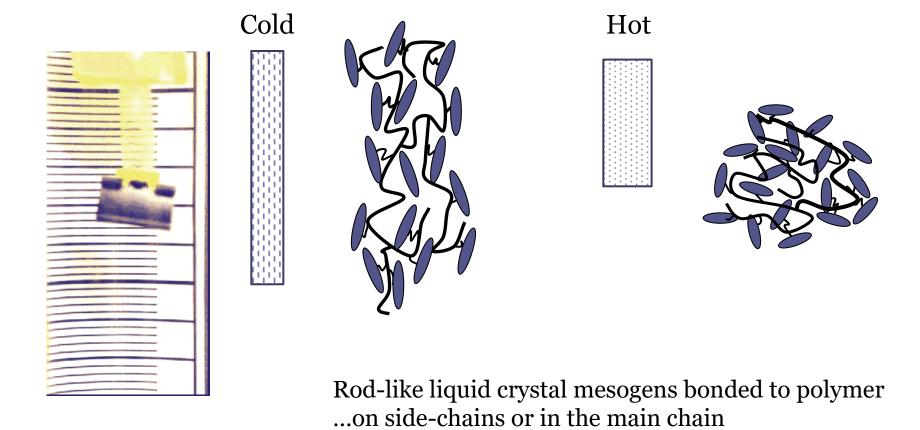


Robin L. B. Selinger





Liquid crystal elastomers: programmable shape-morphing materials





E. M. Terentjev

Local average molecular orientation of liquid crystal defines the **nematic director**

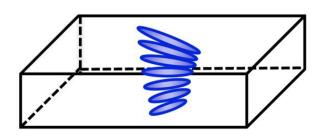
Heating reduces nematic order→ induces strain

Non-uniform nematic director encodes complex shape change

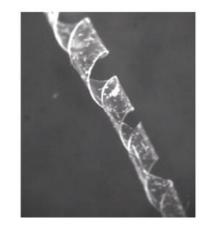


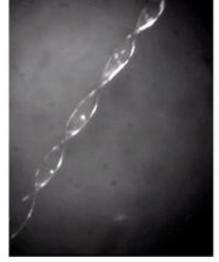


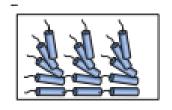
De Haan et al Angewandte Chemie 2012



Twisted or splayed director → spirals, helicoids, curves









Y. Sawa, F. Ye, K. Urayama, T. Takigawa, V. Gimenez-Pinto, RLB Selinger and JV Selinger PNAS 2011

Ways to encode memory: "Blueprinting"



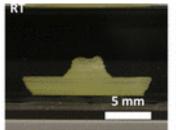
OLD: Form polymer between patterned substrates that impose surface anchoring

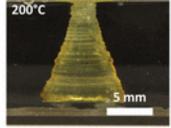
- Director uniform through thickness; or with twist or splay
- Shape transformations: 2-D to 3D

NEW: "4-D printing" recently developed by...

- Taylor Ware ACS Applied Materials & Interfaces 2017
- Dirk Broer Macromolecular Rapid Communications 2017
- TJ White and Jennifer Lewis—Advanced Materials 2018

Ware 2017





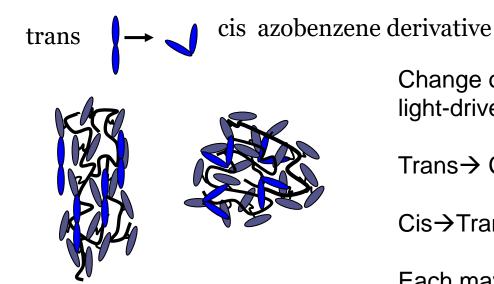
Director set by direction of extrusion

Shape transforms from one 3D shape to another

Inducing motion with light



Nathalie Katsonis Univ of Twente



Change of nematic order produces light-driven mechanical strain

Trans→ Cis photoactuation

Cis→Trans relaxation

Each may be slow or fast, depending on which azo-compound is used, light wavelength and intensity

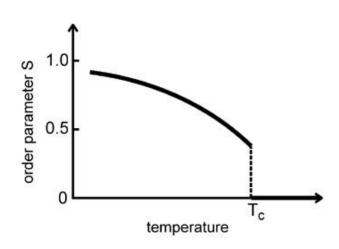
How to model shape change? Often too complex for analytical solution. Finite element elastodynamics with a tetrahedral mesh

$$H = \frac{1}{2} \sum_{t} V^{t} C_{ijkl} \varepsilon_{ij}^{t} \varepsilon_{kl}^{t} - \alpha \sum_{t} V^{t} \varepsilon_{ij}^{t} (S^{t} - S^{0}) \frac{1}{2} \left(3n_{i}^{t} n_{j}^{t} - \delta_{ij} \right) + \frac{1}{2} \sum_{p} m_{p} v_{p}^{2}$$

Elastic potential energy (function of Green-Lagrange strain) Summed over elements Coupling between strain and nematic order Summed over elements

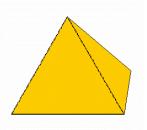
Kinetic energy Summed over nodes

Change in scalar order parameter S(T) describes loss of orientational order

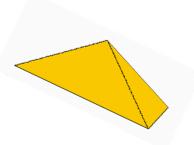


Each tetrahedral element has a REFERENCE state (positions of nodes if free of external stress→ fully relaxed)

and a CURRENT state (defined by current node positions)



Reference state: Relaxed, in mechanical equilibrium



Current state Stretched, sheared, translated, rotated...

This approach can be adapted to model materials that undergo "training" via evolution of the reference state as a function of mechanical stress \rightarrow e.g. via plastic deformation

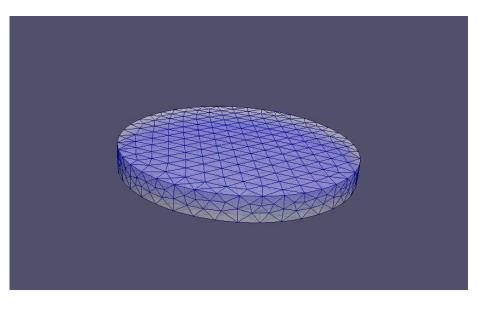
Finite Element Elastodynamics Simulation (FEM)

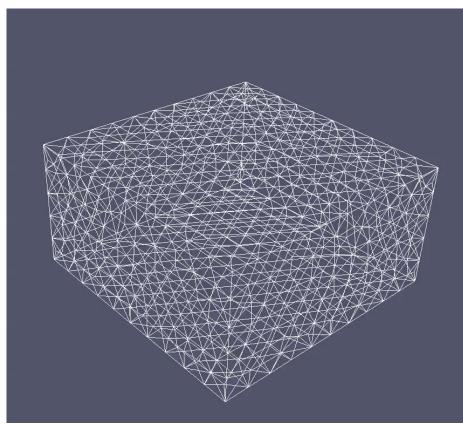
Elastodynamics: Mesh nodes move via f = ma

$$m_n = -\frac{\partial F_n}{\partial r_n}$$
 ——— Force on each node

Home-made code implemented in CUDA for GPU-enabled computer

Fast nonlinear 3-d elastodynamics



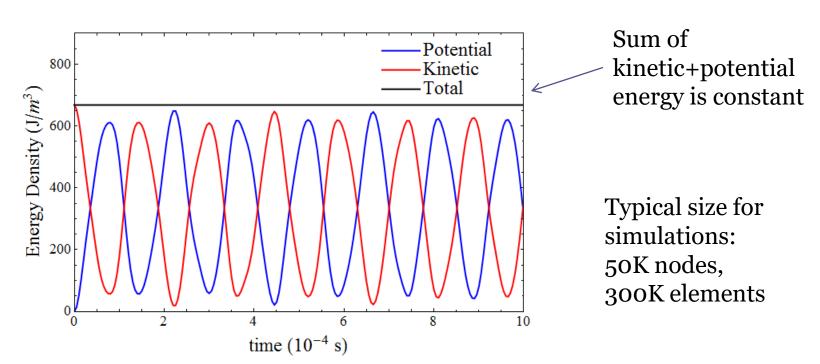




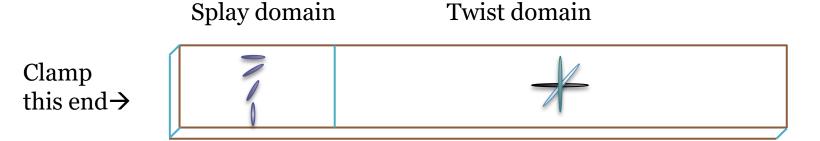
Code runs at 700 frames/second

Without dissipative forces added, kinetic+potential energy is well conserved

Must add friction/dissipation to relax to mechanical equilibrium



Nematic director is "programmed" when the polymer is cross-linked, encodes not just initial and final states, but a whole choreography of motion. Some simple cases...



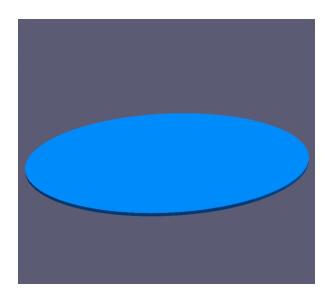






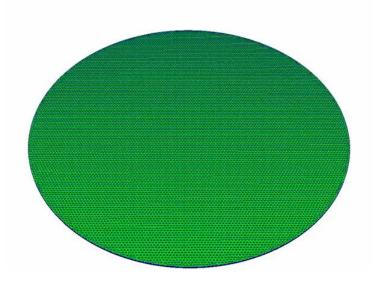
Degenerate final state... disk can pop up or down

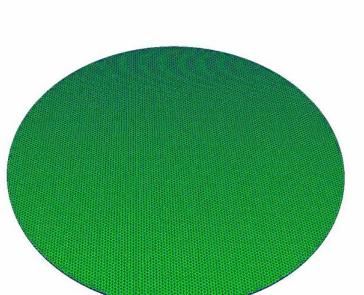
Pop-through → hysteresis





Degenerate final states can include both stable and metastable configurations







-4 defect

Heat uniformly...
Up-down symmetry→
asymmetic, metastable state

Heat from one side→ breaks up/down symmetry

Reaches symmetric ground state



Nematic director can be programmed to enable photoactuated wave motion





UV light from left side at oblique angle

Experiment by Anne Helene Gelebart, Dirk Broer, and coworkers Thin ribbon with splay director: planar on one side, homeotropic on the other



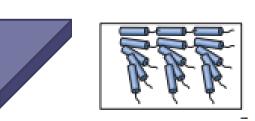
Anne Hélène Gélébart



Dirk Broer

Blueprinted nematic polymer film with azobenzene derivative

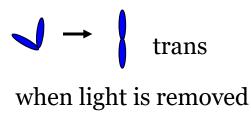
Director: Planar on top



Homeotropic on bottom

photoactuation
trans cis
under UV light

thermal relaxation



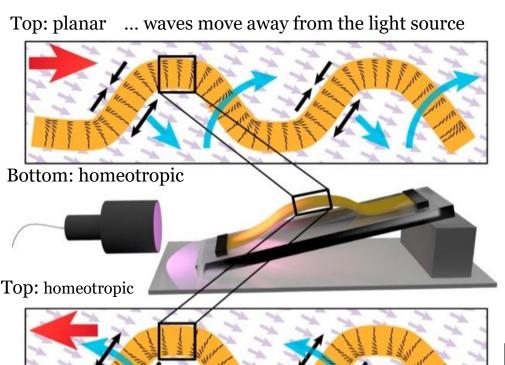
Both transitions take less than 1 sec



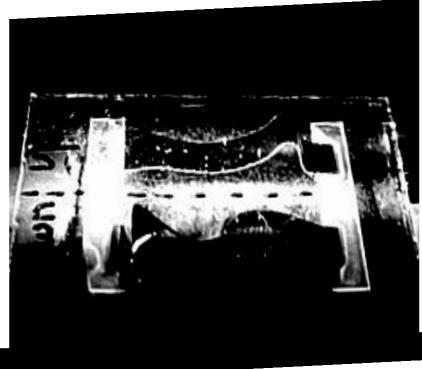
Anne Hélène Gélébart



Dirk Broer



Bottom: planar... waves move toward the light source



UV light from left side

Questions:

What produces continuous photoactuated oscillation?

Why does the direction of wave motion change with orientation?

Clues:

- → Oscillation occurs ONLY with splay director geometry ... Not uniform planar, not uniform homeotropic
- → Oscillation only occurs with oblique light angle (e.g. 10° to the horizontal) ... self-shadowing plays an important role

Dynamics matter, motion is not overdamped...

Changes in *S* are not uniform in the sample...depend on local light exposure and vary with time as the material self-shadows

Assumptions in our model:

S = Nematic scalar order parameter in each tetrahedral volume element

S varies between upper and lower limits according to light exposure

Upper limit= highest degree of nematic order, all azo-dye in trans state

Lower limit = lowest degree of nematic order, maximum of azo-dye in cis state

Perform **ray tracing** to figure out which surface volume elements are illuminated

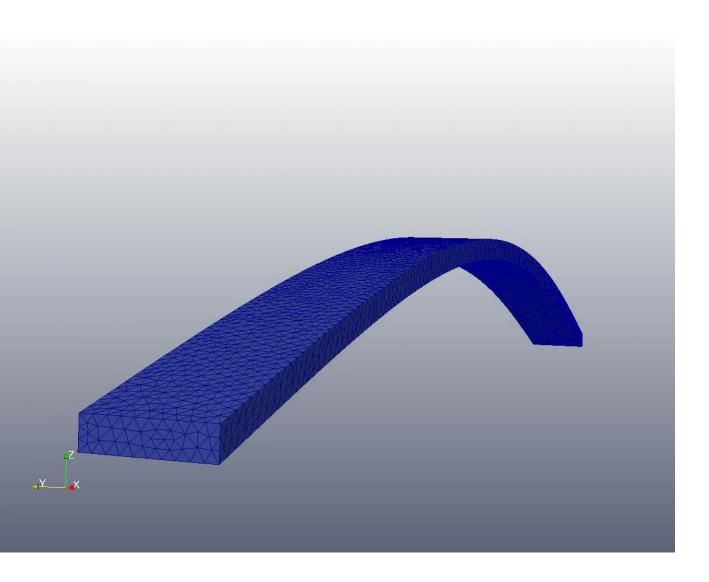
Photoactuation:

S drops linearly in time when material illuminated, reaches equilibrium in \sim 1 sec

Thermal Relaxation:

S increases linearly in time when material shadowed, reaches equilibrium in ~1 sec

Rather simplified model ...

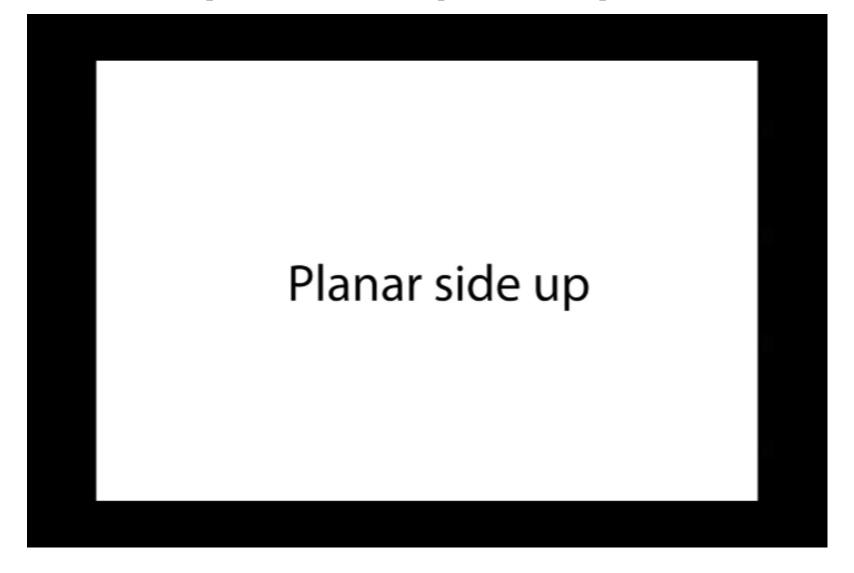




Michael Varga Kent State University

Planar side up: when illuminated, top shrinks along the long axis and expands weakly along the other two axes... light-exposed area curves downwards.

Comparison: Planar side up model and experiment



Time step: 2×10^{-5} s, 2×10^{6} t ime steps \rightarrow 40 seconds of motion

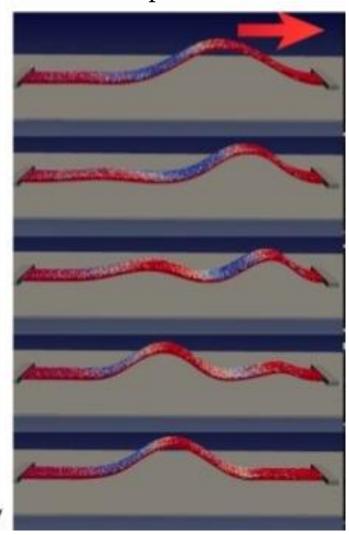
Execution: 25 min on a single CPU with GPU acceleration

Tetrahedral mesh: 5,717 nodes and 19,478 volume elements

Comparison: Homeotropic side up model and experiment



Planar side up



Pop-through instability drives wave regeneration

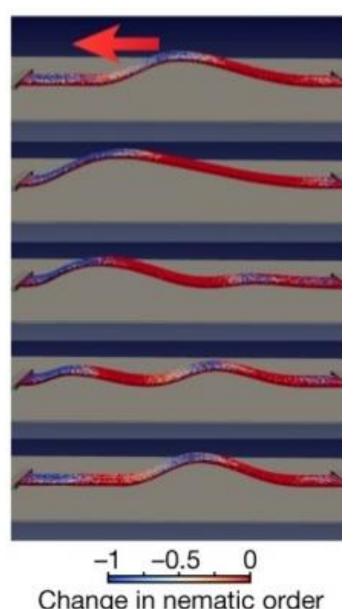
Light shines from the left and strikes the film. Local order parameter decreases in area shown in blue.

There, the top surface of the film contracts along its long axis and the film curves downwards, pushing crest toward the right.

Crest reaches end of the sample; its motion halts

New crest begins to form near the other end, grows to a critical size

Pop-through transition: elastic energy in previous crest is released into the new crest. Process repeats...



Illuminated area in blue has reduced nematic order; homeotropic surface shrinks along surface normal, inducing upward curvature.

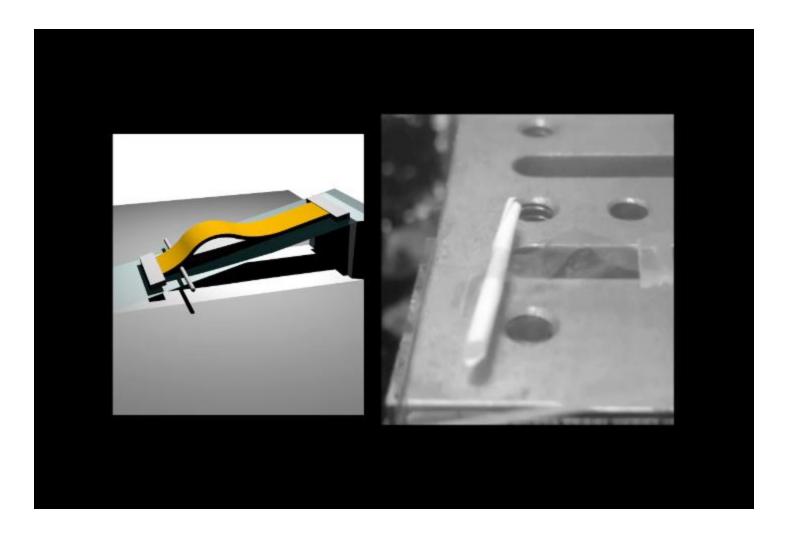
Resulting stress pulls the crest of the bump towards the left, towards the light source

It arrests when it reaches the clamped end and motion stops.

Shadowed region relaxes towards equilibrium; crest shrinks until light passes over it, nucleating a new crest to the right.

Time delay while the new crest forms, then a pop-through transition

What can this device potentially accomplish?



Photoactuated autonomous locomotion:

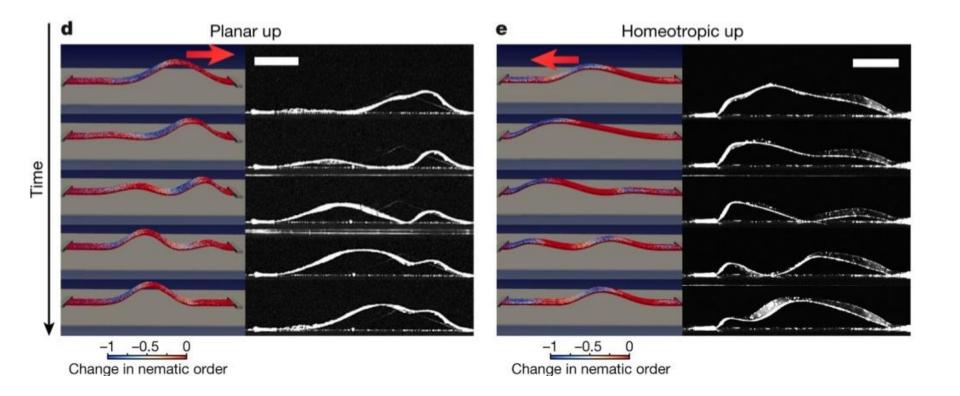
Planar side up ... Moves away from light source



Photoactuated autonomous locomotion:

Homeotropic side up ... Moves toward light source





Anne Helene Gelebart, Dirk Jan Mulder, Michael Varga, Andrew Konya, Ghislaine Vantomme, E. W. Meijer, Robin L. B. Selinger, and Dirk J. Broer, "Making waves in a photoactive polymer film," *Nature* v. 546, p. 632, 29 JUNE 2017

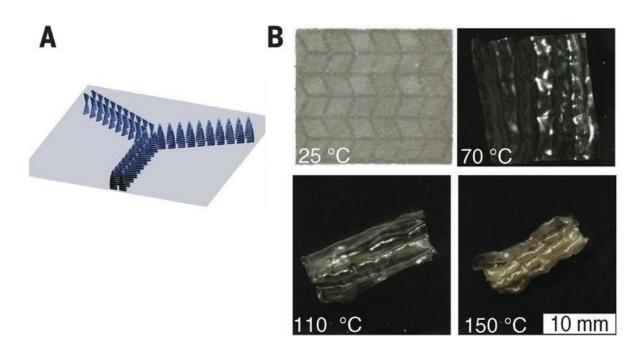
Link to paper and videos at tiny.cc/rselinger2018



Shape change includes origami-like folding, slits and irises that open/close

Taylor Ware et al Science 2015

Origami-inspired actuators.(A) Schematic of an edge portion of the Miura-ori pattern with a localized twisted nematic region bounded by unordered regions.

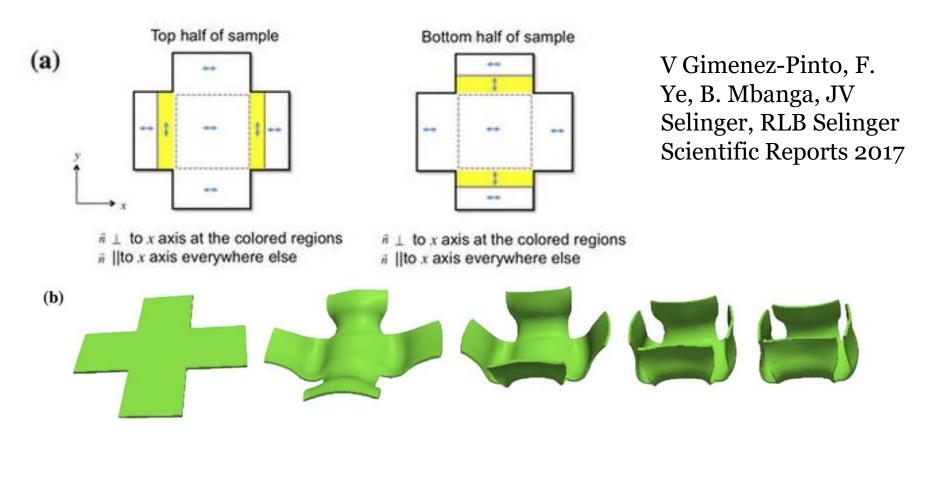


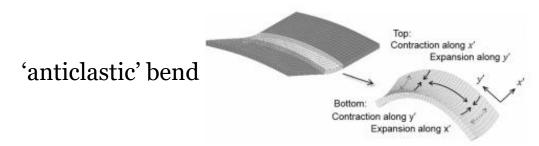
Taylor H. Ware et al. Science 2015;347:982-984



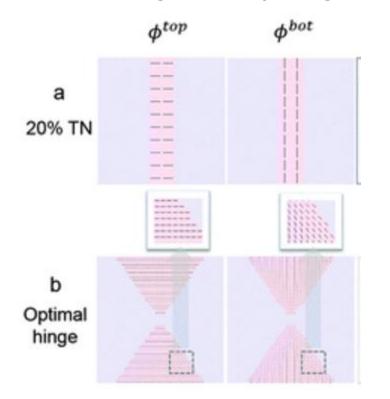
Hinges are narrow regions with director twist...

Can narrow twisted domains create sharp folds? Not very well!

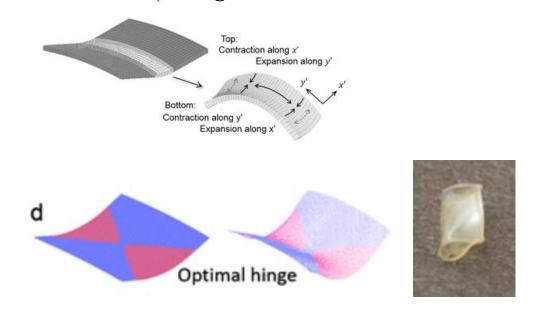




K. Fuchi et al Soft Matter 2015 How to design a sharp hinge for origami? Fuchi et al propose a structure



Twisted nematic hinge creates saddle curvature



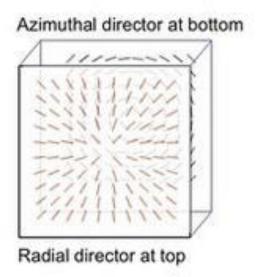
Topology optimization for the design of folding liquid crystal elastomer actuators†

Kazuko Fuchi,‡^{ab} Taylor H. Ware,§‡^{cd} Philip R. Buskohl,^c Gregory W. Reich,^a Richard A. Vaia,^c Timothy J. White^c and James J. Joo*^a

Soft Matter 2015

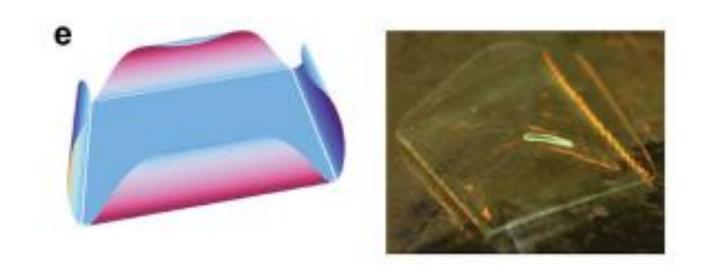
PROBLEM: This optimized solution isn't workable if spacing between folds is small compared to their length

Twisted director –square sample with a +1 defect



Twist through the thickness ...add chiral dopant to keep twist direction uniform

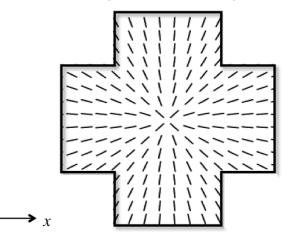
Azimuthal + radial = "Radimuthal".... Term coined by Laurens de Haan



"Radimuthal" disk: makes sharp folds

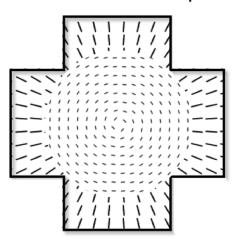
Predicted by theory of Carl Modes and Mark Warner (C Modes, PRE 2012)

Top half of sample

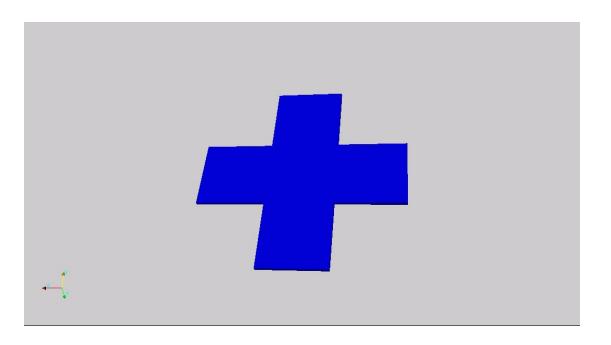


Radial director on the top

Bottom half of sample



Azimuthal director in the middle circle Radial director elsewhere



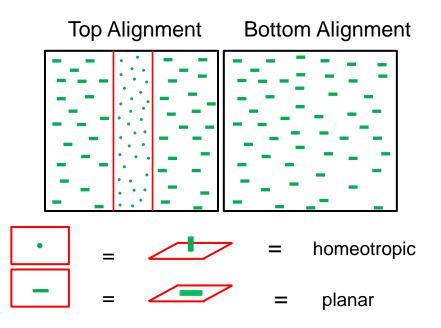
"Radimuthal" box

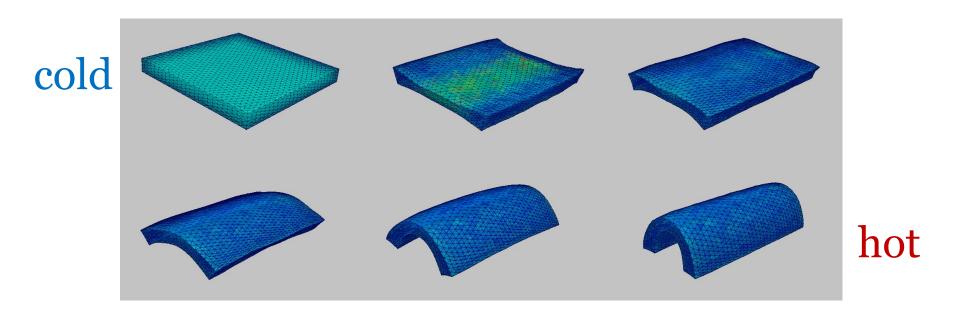
V Gimenez-Pinto et al Scientific Reports 2017

Splay hinge design:

Homeotropic anchoring on top

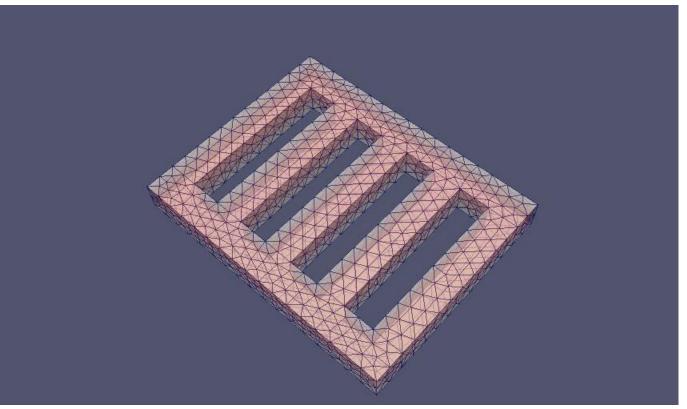






"Kirigami" approach: cut-outs accommodate strain in splay hinge

Cut-outs accommodate transverse strain



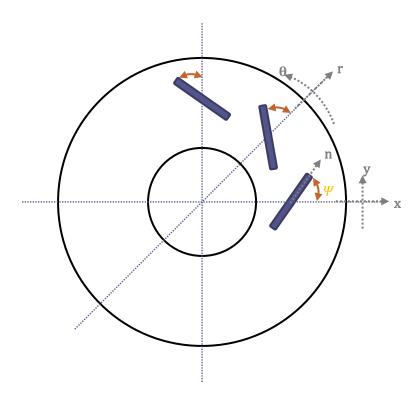
Cross beams: homeotropic on top, planar on bottom



Long edges: isotropic



Modeling an iris...annulus with director at an angle to radial; clamp outer edge



Initial (nematic) director field: $n(\Phi) = (\cos \Phi, \sin \Phi, 0)$ at (x,y,z) $\theta = \arctan(y/x)$ $\Phi = \theta + \Psi$

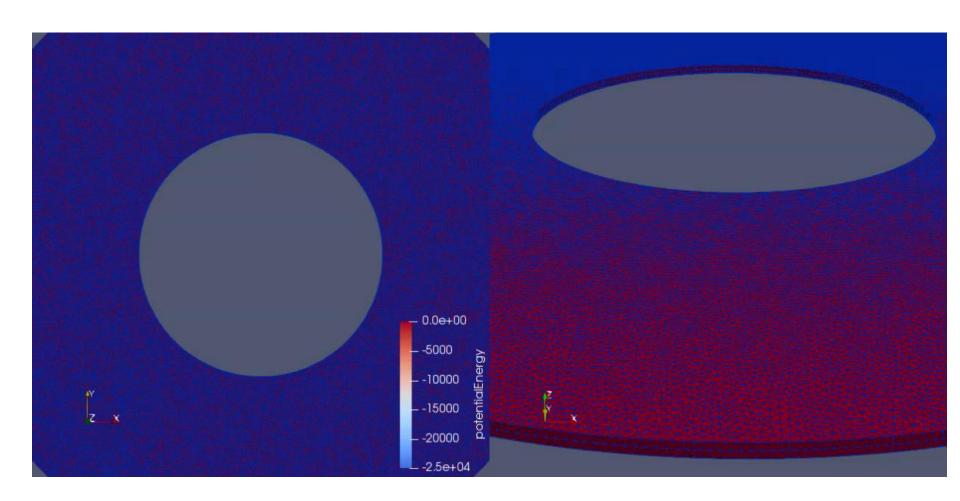
Opening the iris on heating: if director at an angle to radial direction, actuate with both dilation and rotation

Director is 30deg from radial 30,000 time steps

Mesh:

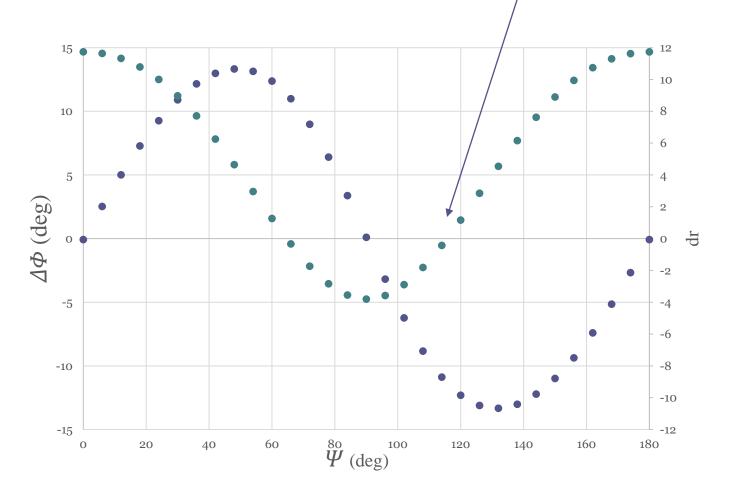
nodes:88,816

tetrahedra: 315,090

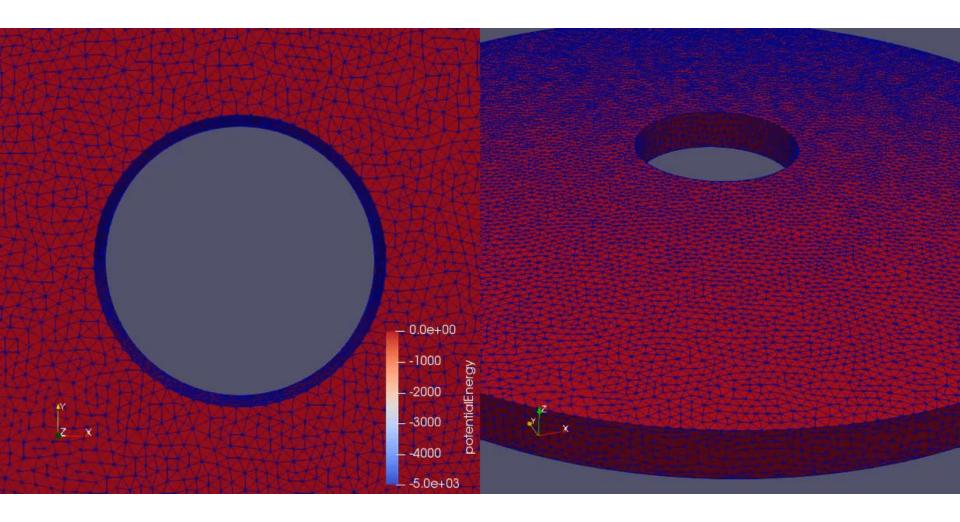




At special angle, no dilation... pure rotation!



Pure rotation > Torsional actuator

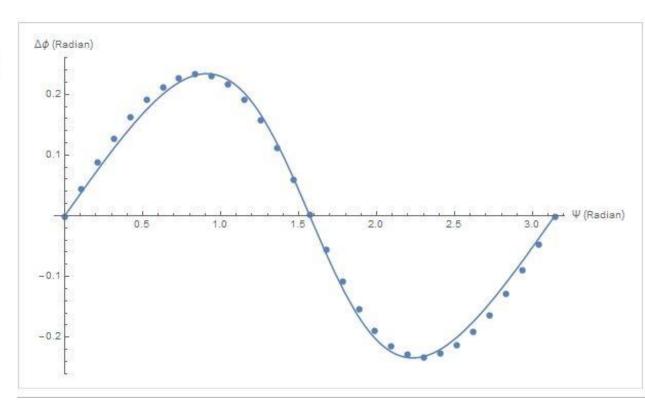


Analytical solution by Hillel Aharoni

$$\Delta \phi = \frac{\epsilon \sin(2\theta)}{1 + \epsilon \cos(2\theta)} \log \left(\frac{r_{out}}{r_{in}}\right)$$

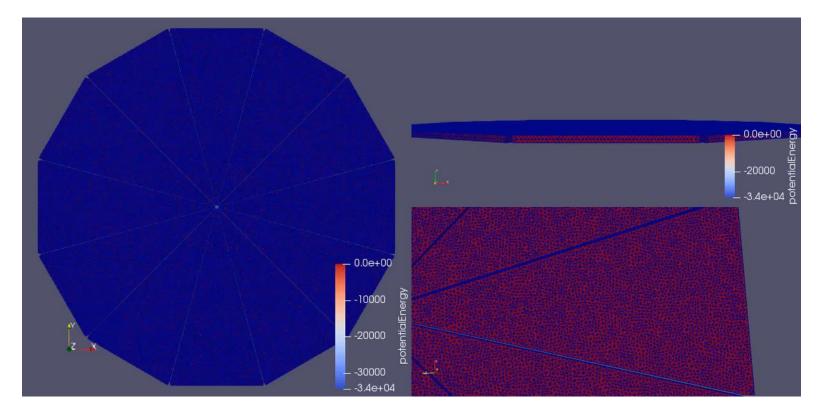
$$\epsilon = \frac{\lambda^{-2\nu} - \lambda^2}{\lambda^{-2\nu} + \lambda^2}$$

 $\begin{array}{l} \underline{\text{Input}} \\ R_{\text{out}}(r_{\text{out}}) = 100 & \underline{\text{Matched}} \\ R_{\text{in}}(r_{\text{in}}) = 40 & \lambda = 0.844 \\ v = 0.49 & \end{array}$



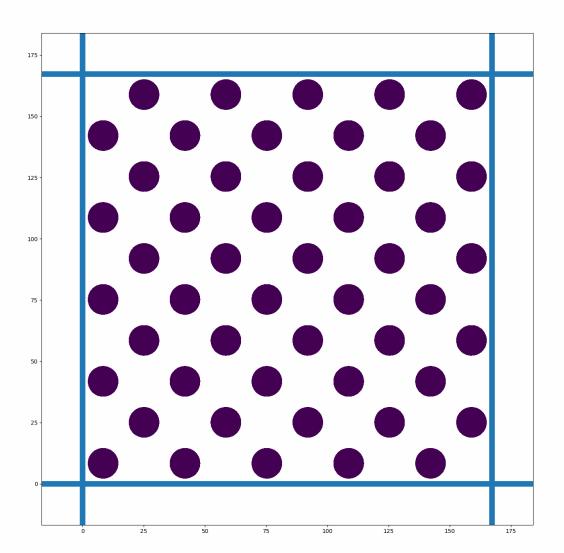
 λ measured at Ψ =0, where there is no rotation, is $(R_{out}-dr)/(R_{out}) = 0.805$. Plot above was obtained after matching λ =0.844

Another iris design ...



Design studied in experiments by Priimagi, Advanced Materials 2017





Using FEM to model deformable colloids

Preliminary 2D results...

RLBS Collaboration with:

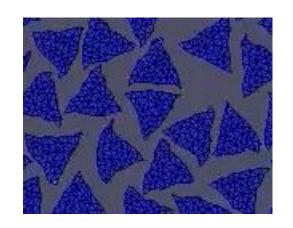
Manuel Valera-Slippery Rock Craig Maloney –Northeastern

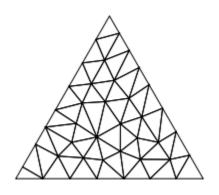
Can add...

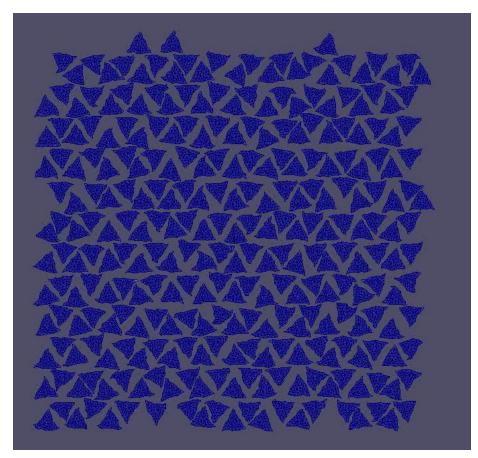
- swelling effects
- internal dissipation
- friction
- hydrodynamic interactions
- → Implement in 3D with CUDA

Possible application: visualize force chains in flexible granular media, different grain shapes Particles don't have to be disks...

Could model any particle shape in 2D or 3D

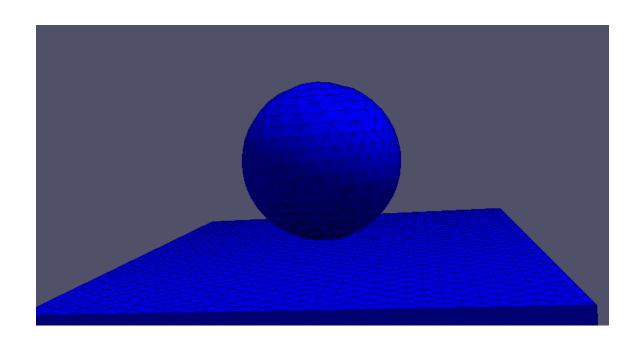






Repulsive forces on surface nodes prevent objects from passing through each other

Surface node interactions → bumpy equipotential surfaces



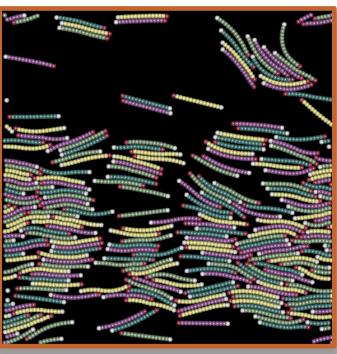


Other fun things we've been working on....

(if time allows)

Active matter: flexible self-propelled fibers

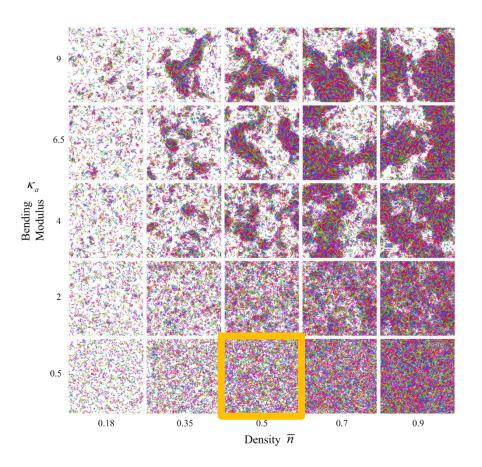


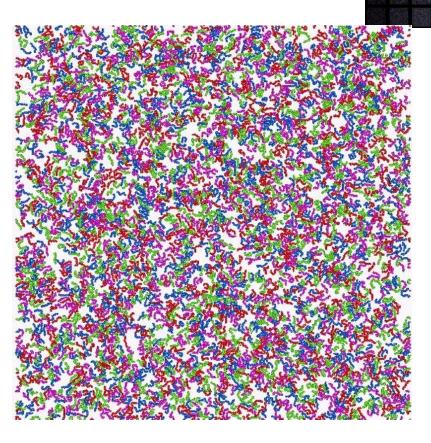


Lower bending modulus

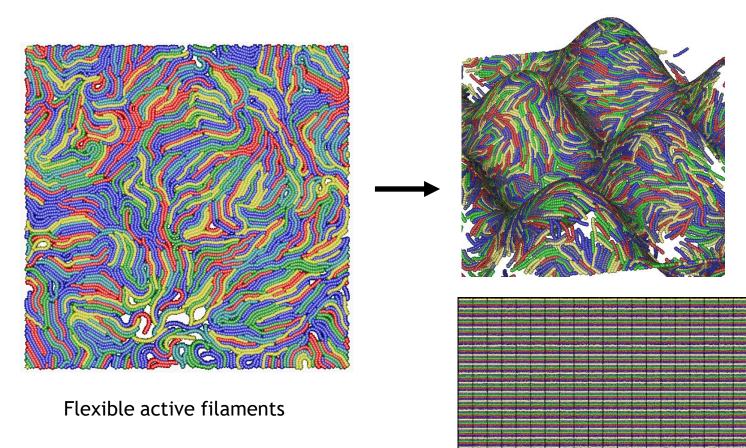
Higher bending modulus

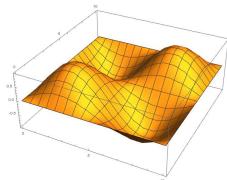
Phase diagram: bending modulus vs density

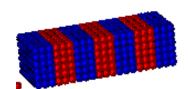




Flexible Active Filaments on Curved Surfaces



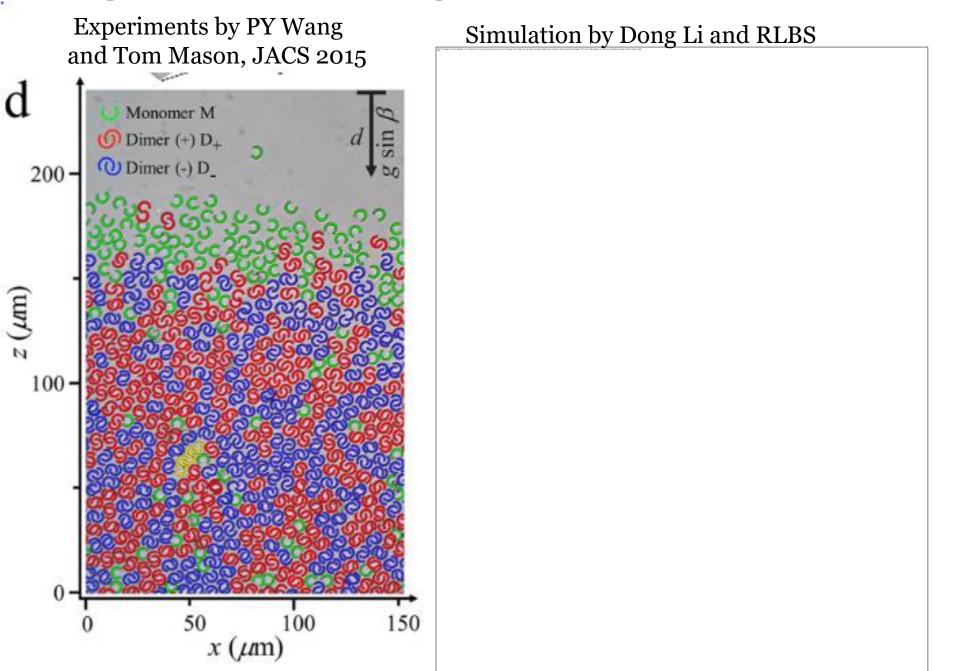


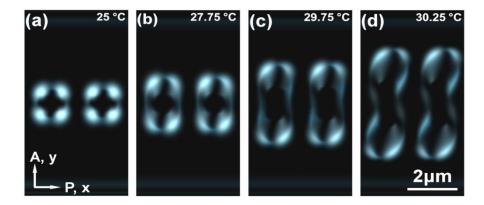


More detailed models of soft active particles using FEM elastodynamics

... animal locomotion?

C-shaped colloids → chiral dimers: Spontaneous formation of homochiral domains



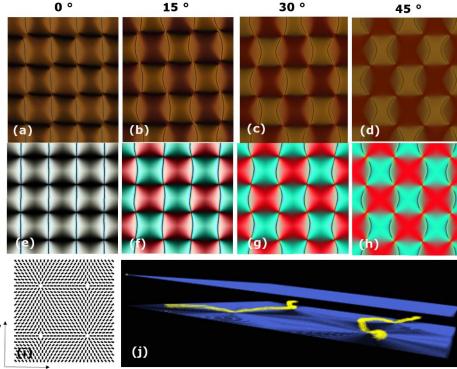


"Cooking **Skyrmions** until they pop"

Bubble defects in cholesteric LC's confined in microfluidic channels expand when heated

Simulations by Sajedeh Afghah





Cholesteric liquid crystals in rectangular microchannels: skyrmions and stripes

Yubing Guo,† Sajedeh Afghah,† Jie Xiang,† Oleg D. Lavrentovich, Robin L. B. Selinger and Qi-Huo Wei*

Soft Matter, 2016

Nematic between substrates with surfacepatterned point defects Collaboration with Qi-Huo Wei

Collaborators...



Anne Hélène Gélébart and Dirk Broer TU Eindhoven



Jonathan Selinger Kent State



Manuel Valera Slippery Rock University



Andrew Konya



Michael Varga



Vianney Gimenez-Pinto



Youssef Golestani

Current and former students...

Conclusions

Liquid crystal elastomers encode complex shape actuation trajectories

- Simulation studies of the forward problem: for a given sample shape and director field, calculate trajectory
 - → Degeneracy, metastable states, snap-through, hysteresis
 - → Design motifs for actuators and simple machines
- Director field can encode photoactuated wave motion
- Origami folding: designing the best fold
- Simple machines: iris, torsion actuator
- Preliminary studies of flexible colloids



NSF-DMR 1409658, and NSF CMMI-1436565, NSF-CMMI -1663041

References: tiny.cc/selinger2018