

Plasticity and memory effects in glasses

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Memory Formation in Matter
KITP, UC Santa Barbara, 05 february 2018

Credits

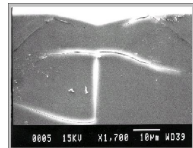
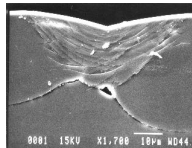
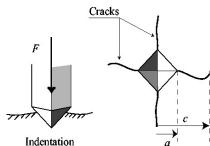
- Sylvain Patinet (PMMH, ESPCI), Stéphane Roux (ENS Cachan)
- Mehdi Talamali, Botond Tyukodi, Armand Barbot, Matthias Lerbinger (PMMH, ESPCI)
- Antoine Perriot, Etienne Barthel (CNRS/Saint-Gobain)
- Cindy L. Rountree, Elisabeth Bouchaud (CEA Saclay/ESPCI)
- Craig Maloney (NorthEastern University, Boston)
- Michael L. Falk (Johns Hopkins University, Baltimore)

Introduction: Plasticity and memory

Mechanical memory effects in bulk amorphous silica

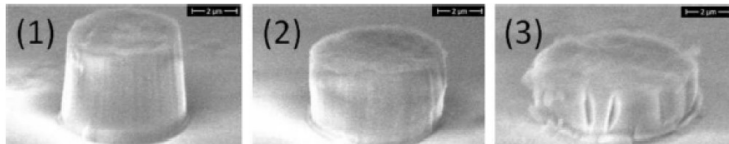
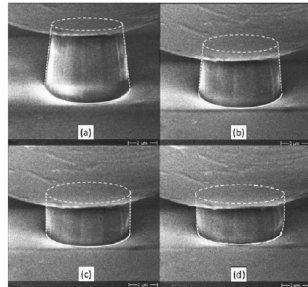
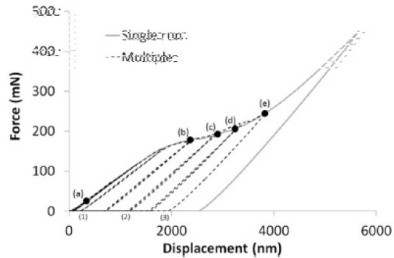
Plasticity of oxide glasses

- **Indentation induces plastic imprints** No crack under low enough loading. Plastic deformation confined at micro-scale
- **Permanent densification** under indentation or hydrostatic test, amplitude depends on glass nature, up to 20 % in the case of silica glass ; strong coupling shear/pressure ;

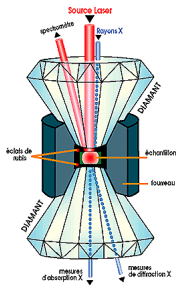


G. Kermouche et al. *Acta Mat.* **56**, 3222 (2008)

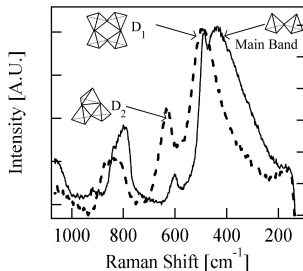
Compression of micro-pillars of silica glass



Raman spectroscopy as a probe of glass density



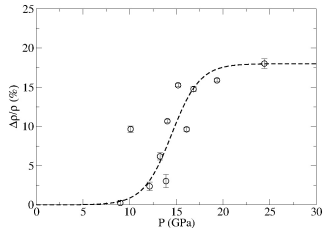
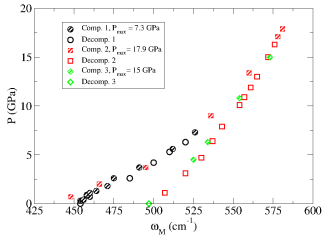
Diamond Anvil Cell – Pressure
up to 18-20 GPa



Raman spectrum of normal
and densified silica

A. Perriot et al. *J. Am. Ceram Soc* **89**, 596 (2006)

Amorphous silica under cycles of increasing pressure



Beyond $P_Y \approx 10$ GPa, the Raman shift is no longer reversible: densification is permanent.

Amorphous silica keeps memory of the highest Pressure level experienced in the past.

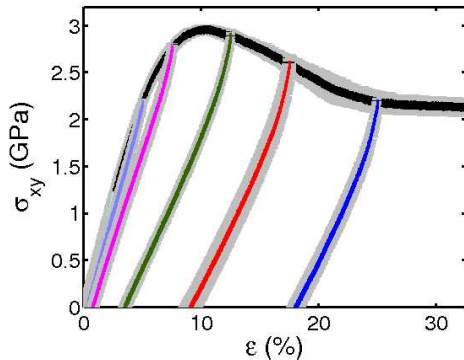
The higher P_{Max} the higher the final density. Saturation at $\Delta\rho/\rho \approx 20\%$ for $P_{Max} > 20$ GPa

D. Vandembroucq et al. *J. Phys. Cond. Mat.* **20**, 485221 (2008)

Pressure induced permanent densification of amorphous silica

- Amorphous silica keeps memory of the highest Pressure level experienced in the past $\Delta\rho/\rho = f(P_{Max})$
- The higher P_{Max} the higher the final density. Saturation at $\Delta\rho/\rho \approx 20\%$ for $P_{Max} \approx 20$ GPa
- Below 20 GPa, amorphous silica = network of tetrahedra SiO_4 connected by their vertices (4-fold coordination)
- Above 20 GPa, *reversible* transition to 6-fold coordination
- **Question:** Can we compute the maximum density of a packing of tetrahedra connected by their vertices only ?

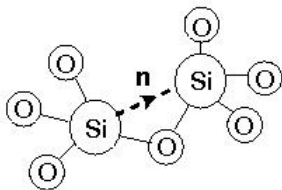
Amorphous silica under cycles of increasing shear strain



A molecular dynamic study of a model silica glass (3-body potential, Vashishta et al PRB 90)

C.L. Rountree, DV, M. Talamali, E. Bouchaud, S. Roux, PRL 09

A probe of structural anisotropy at medium range order

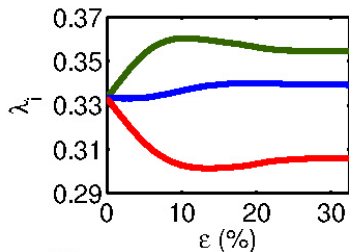
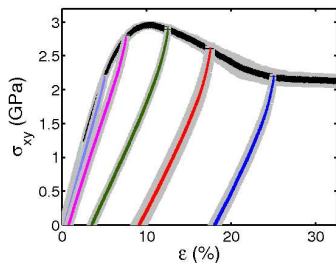


Anisotropy parameter:

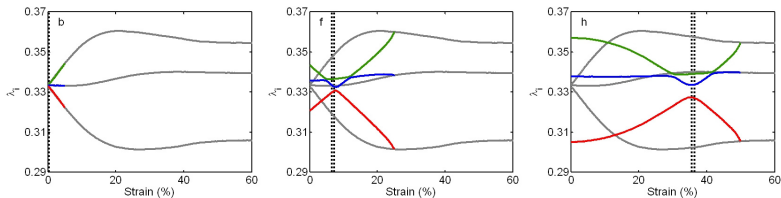
$$\alpha = \frac{3}{2} \sqrt{\sum_i \left(\lambda_i - \frac{1}{3} \right)^2}$$

- Analogy with the fabric tensor in granular materials
- $F = \langle \vec{n} \otimes \vec{n} \rangle$ is the tensor of “contacts” between neighboring Si atoms, Eigenvalues λ_i
- $\lambda_i = \left\{ \frac{1}{3}, \frac{1}{3}, \frac{1}{3} \right\}$: Isotropy
- $\lambda_i = \{1, 0, 0\}$: Full Anisotropy

Shear strain induces structural anisotropy

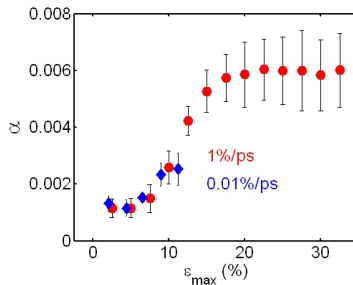
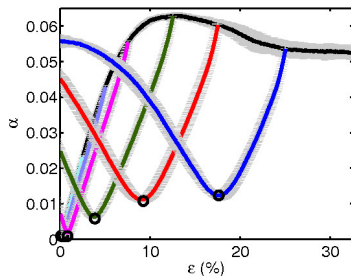


Plasticity induces *persistent* structural anisotropy



Isotropy is not recovered after unloading at zero shear stress

An orientational memory of amorphous silica



Experimental confirmation: Sato et al. JAP 13, PRB 15

Ultrafast laser photo inscription in amorphous silica

JOURNAL OF APPLIED PHYSICS

VOLUME 94, NUMBER 3

1 AUGUST 2003

Demonstration of high-density three-dimensional storage in fused silica by femtosecond laser pulses

Guanghua Cheng,^{a)} Yishan Wang, J. D. White, Qing Liu, Wei Zhao, and Guofu Chen
State Key Laboratory of Transient Optics and Technology, Xi'an Institute of Optics and Precision Mechanics, Chinese Academy of Sciences, Xi'an 710068, China

(Received 16 December 2002; accepted 13 May 2003)

Three-dimensional optical recording of high contrast spherical bits (diameter < 300 nm) at a density of 500 G/cm³ in fused silica using a Ti:sapphire femtosecond laser is demonstrated. Bits are optically read out using both a confocal and a phase-contrast scheme. The recording density for different materials and recording mechanisms are discussed. © 2003 American Institute of Physics. [DOI: 10.1063/1.1589596]

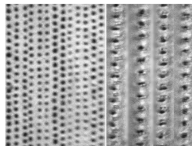
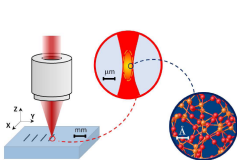


FIG. 2. Optical image of bits written inside fused silica as viewed parallel to the 400 nm 200 fs excitation pulse. Light is focused by a NA=0.85 objective. In-plane bit separation is 1 μm and the layer separation is 2 μm . (Left) bits as viewed parallel to the excitation. (Right) bits as viewed orthogonal to the excitation pulse.

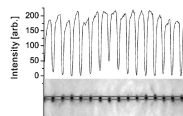
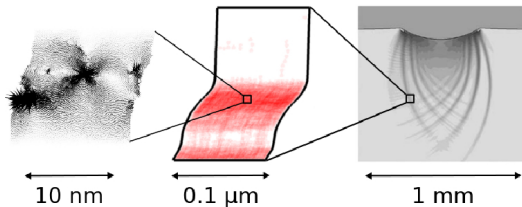


FIG. 4. Read out of data through a phase-contrast microscope. The upper part of the figure is the signal recorded by a single row of the CCD camera. The pits being imaged are shown in the bottom portion, with the two lines denoting the edges of the row. High contrast is evident even for those bits not perfectly centered.

N. Shcheblanov et al, PRB 2018

Mesoscopic models of amorphous plasticity

Lattice models of amorphous plasticity



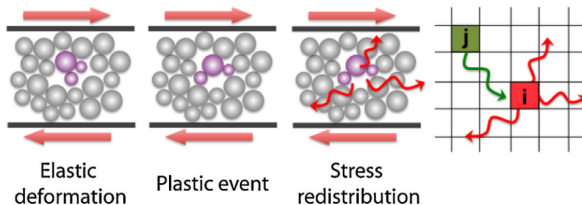
Can we build at mesoscopic scale a *minimal model* that reproduces at large scale the important features of amorphous plasticity ?

Two main ingredients:

- Local threshold dynamics
- Elastic interactions

$$\mu \frac{\partial \varepsilon_p}{\partial t} = \Sigma^{ext} + \sigma^{el} [\mathbf{x}, \{\varepsilon_p(\mathbf{x})\}] - \sigma^Y [\mathbf{x}, \varepsilon_p(\mathbf{x})]$$

Lattice models of amorphous plasticity

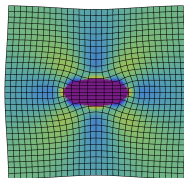
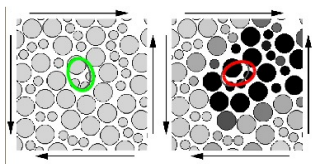


Bocquet et al. [PRL 2009](#)

Various implementations: Boston, Erlangen, Grenoble, Helsinki, Lausanne, Milano, Paris...

Recent review: A. Nicolas et al. [arxiv:1708.09194](#)

Local reorganizations and anisotropic elastic interactions



Local rearrangement induced by shear stress, Falk and Langer PRE 98
The (far field) internal stress induced by a plastic reorganization obeys a quadrupolar symmetry (Eshelby):

$$\sigma_{xy}(0,0) = -\mu^* \gamma_p; \quad \sigma_{xy}(r,\theta) = A \frac{\cos 4\theta}{r^2}; \quad A = \frac{2\mu^*}{\pi} S \gamma_p$$

Potential Energy Landscape: distribution of random barriers

Internal stress induces avalanches and cascades of plastic events

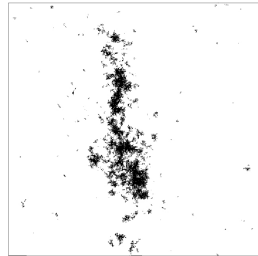
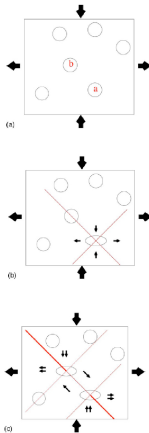


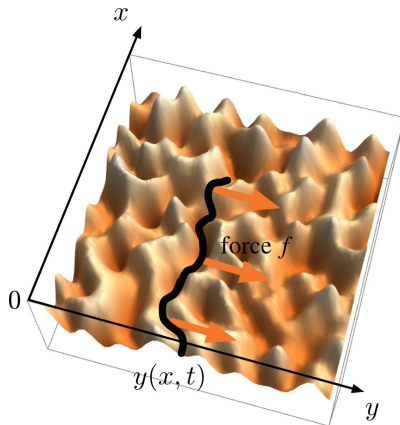
FIG. 16. Local slip (as defined in the text) which occurs during the entire plastic event. Arrow lengths are equal to the magnitude of the particular slip scaled by a factor of 10. Slips of amplitude less than 10^{-3} are not shown for clarity.

A cellular automaton for amorphous plasticity

- Discretization on a lattice at mesoscopic scale
- A scalar plastic criterion $\sigma > \sigma_Y$
- Structural disorder Local criterion : $\sigma(i,j) > \overline{\sigma_Y} + \delta\sigma_Y(i,j)$
- Local reorganization local slip increment $\delta\gamma$ and **update** of local plastic threshold: $\delta\sigma_Y$ is **renewed** after slip
- Anisotropic elastic response A local slip induces a stress redistribution σ^{el} all over the system, $\sigma^{el} \propto \mu\delta\gamma \cos 4\theta/r^2$
- Evolution rule: extremal dynamics, stress or strain control...
- Parameters of the model: $\mu = 1$, $\delta\sigma_Y \in \text{rand}[0, 1]$,
 $\delta\gamma \in \text{rand}[0, \mathbf{d}]$; \mathbf{d} = coupling parameter

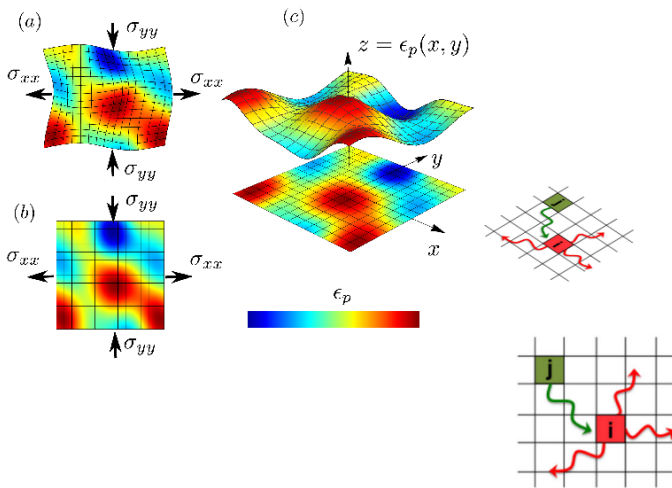
Baret et al. [PRL 2002](#), Talamali et al. [PRE 2011](#), [C.R. Mecanique 2012](#)

Depinning of an elastic line/manifold



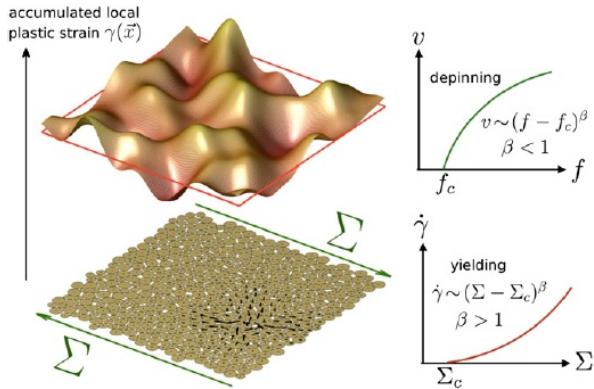
$$\mu \frac{\partial h}{\partial t} = f^{ext} + f^{el} [\mathbf{x}, \{h(\mathbf{x})\}] - \frac{\partial V}{\partial h} [\mathbf{x}, h(\mathbf{x})]$$

Amorphous plasticity as an elastic manifold: Depinning



$$\partial \epsilon_p = \sum \text{ext} \dots \sigma^{el}[\mathbf{y}, \mathbf{c}(\mathbf{y})] - \sigma^Y[\mathbf{y}, \mathbf{c}(\mathbf{y})]$$

Depinning vs Plastic yielding: scaling properties

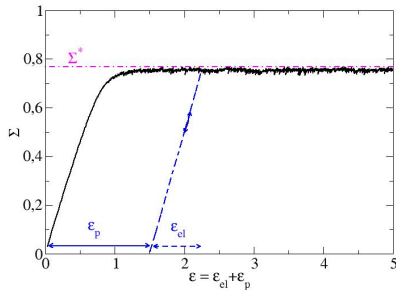


J. Lin, E. Lerner, A. Rosso and M. Wyart, PNAS **111**, 40 (2014)

a mechanically induced aging process
Transient hardening:

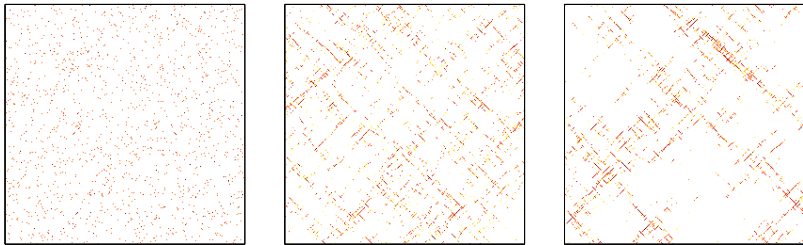
A transient “hardening” stage precedes a stress plateau

The macroscopic flow stress Σ^* can be interpreted as a **depinning threshold** :



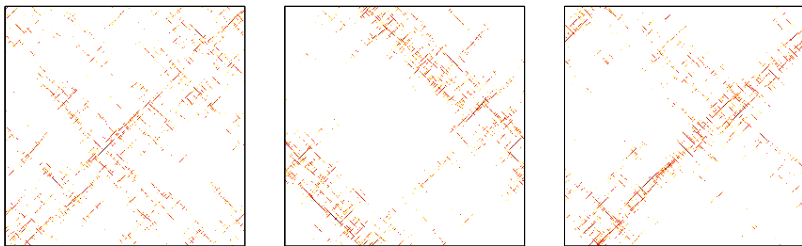
Below Σ^* the material starts deforming and eventually stops ;
Above Σ^* it flows indefinitely.

Elastic interactions matter: strain localization



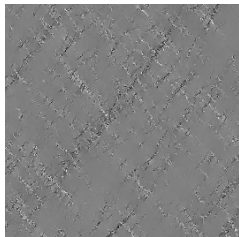
During the transient/hardening stage, plastic deformation gets gradually more correlated with the same quadrupolar symmetry as the elastic interaction.

Localization vs diffusion



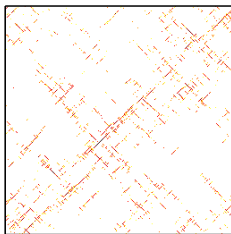
In the stationary regime, plastic deformation remains localized but, localization patterns are not persistent, rather they diffuse throughout the system.

Localization vs diffusion



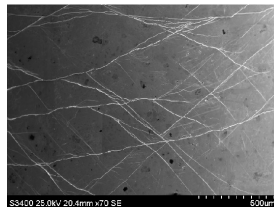
Atomistic scale

Maloney & Robbins, PRL 09



Mesoscopic Scale

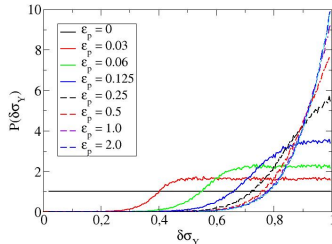
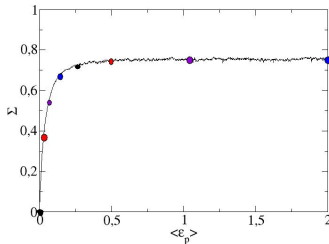
Talamali et al, PRE 11



Experimental Results

Zhang et al, Scripta Mat. 09

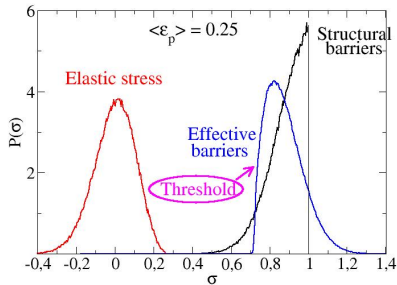
Exhaustion of low barriers induces transient hardening



Hardening: the plastic yield stress gets higher with deformation.

Interpretation: “Darwinian” or SOC-like dynamics; progressive exhaustion of the weakest sites during the transient stage induces a systematic bias in the distribution of local plastic thresholds.

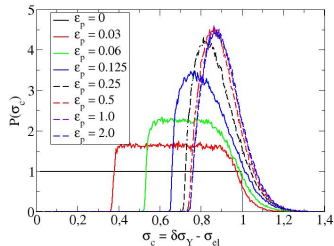
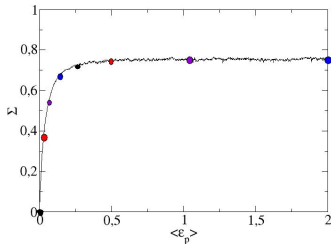
Elastic interactions matter: effective barriers



Local effective barriers: $\sigma_c(i,j) = \delta\sigma_Y(i,j) - \sigma^{el}(i,j)$

A well defined stress threshold bounds the distribution of effective barriers: yield stress

Exhaustion of effective barriers



Exhaustion induces a progressive increase of yield stress: transient hardening

Again a “mechanical memory”: dependence on the past maximum stress

Talamali et al, C.R. Mecanique 2012

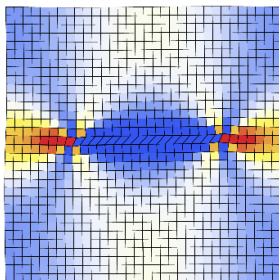
Plasticity vs Depinning

- **Depinning**: Any fluctuation of the line/manifold induces a restoring force
- **Plasticity**: Any unit shear-band in a maximum shear stress direction induces no elastic stress

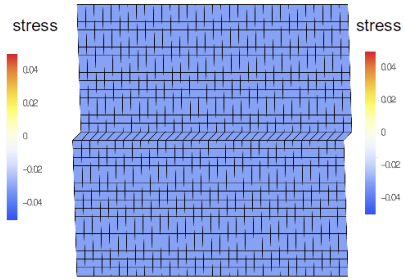
Specific features of plastic yielding models: anisotropic strain fluctuations, localization, shear-banding

Tyukodi et al, PRE **93**, 063005 (2016)

Spanning shear-bands induce no stress

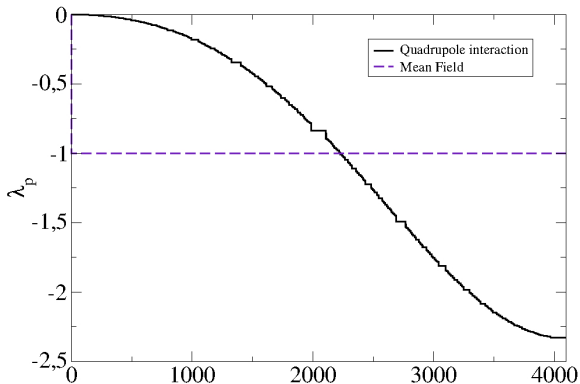


partial shear band

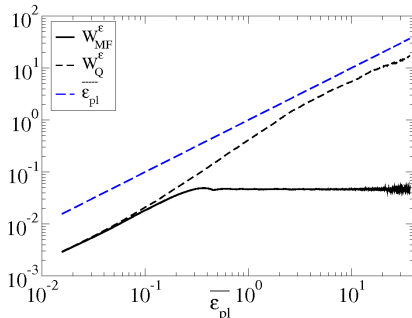


full shear band

Spectrum of eigenvalues: quadrupolar vs Mean Field interaction



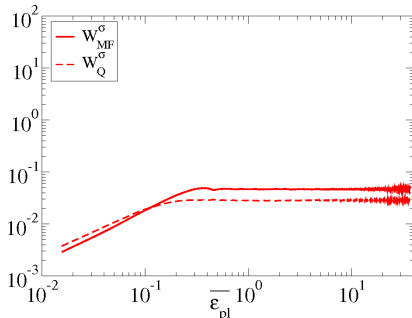
Family-Vicsek scaling vs diffusive regime



Depinning: Variance of interface fluctuation saturates when correlation length reaches system size

Plasticity: Variance of strain fluctuations does not saturate, diffusive regime

Family-Vicsek scaling vs diffusive regime

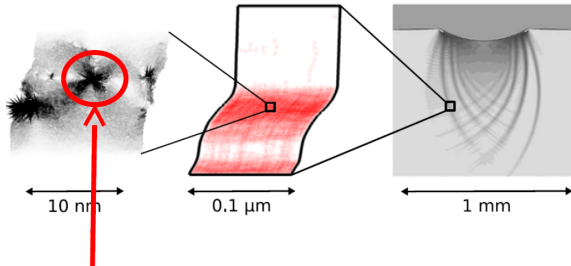


Depinning: Variance of elastic force fluctuation saturates when correlation length reaches system size

Plasticity: Variance of stress fluctuations saturates, soft modes prevents divergence of stress fluctuations

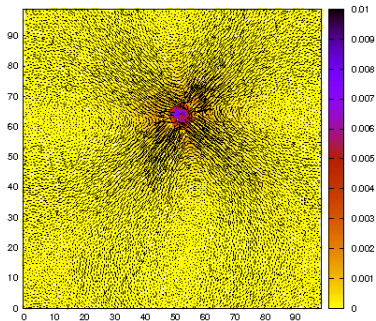
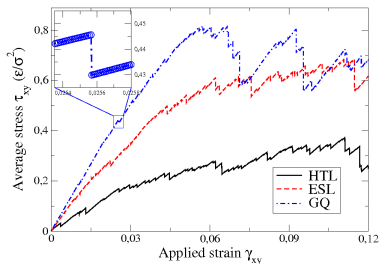
Yield stress disorder under the mesoscope

Characterization of local plastic thresholds



Question: how to localize (*before plastic deformation*) and characterize the population of STZ at the atomistic scale ?

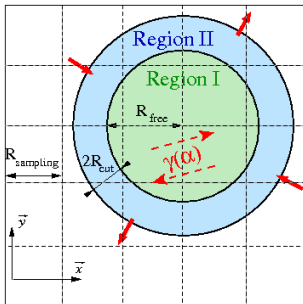
AQS simulations of a 2D LJ model glass



Three protocols of preparation:

- HTL “liquid-like”
- ESL: “fast quench”
- GQ: “slow quench”

A mesoscopic probe of local yield stress



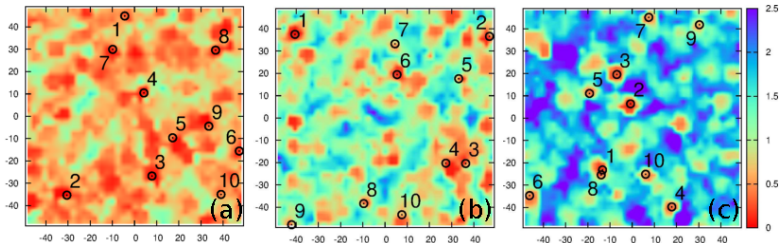
- Shear along direction α
- Initial shear stress along α : $\tau^0(\alpha)$
- Local shear stress at the onset of the instability: $\tau^{inst}(\alpha)$
- Shear stress threshold along α : $\Delta\tau^C(\alpha) = \tau^{inst}(\alpha) - \tau^0(\alpha)$

Projected yield stress along α_ℓ :

$$\Delta\tau^Y(\alpha_\ell) = \min_{\alpha} \frac{\Delta\tau^C(\alpha)}{2 \cos(\alpha_\ell - \alpha)}$$

S. Patinet, DV, M.L. Falk PRL16, Barbot et al submitted 2017

Dependence on thermal history



HTL-“liquid-like”

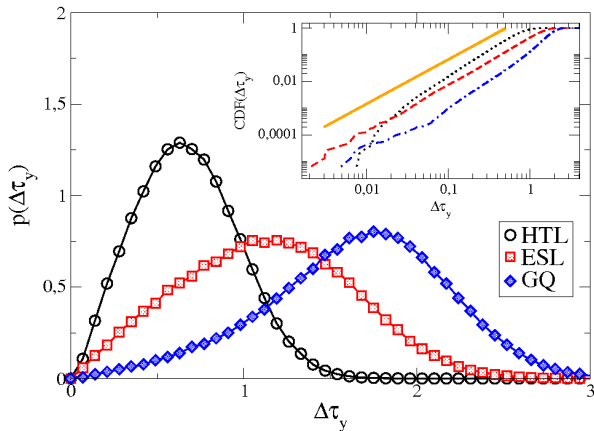
ESL-“fast quench”

GQ-“slow quench”

The more equilibrated the glass, the larger the local yield stresses

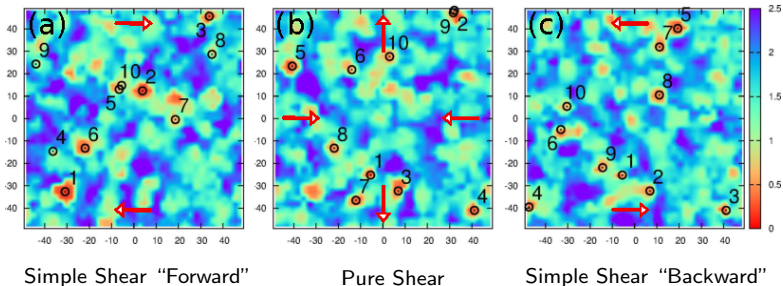
S. Patinet, DV, M.L. Falk PRL16, Barbot et al submitted 2017

Distributions of local yield stress

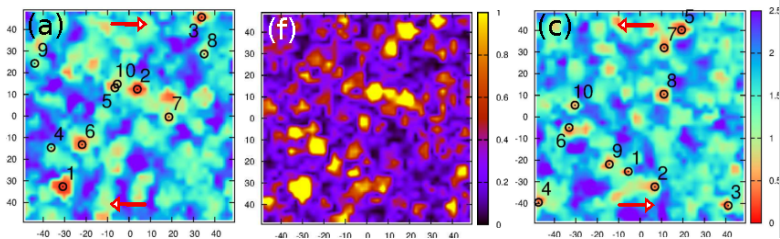


S. Patinet, DV, M.L. Falk PRL16, Barbot et al submitted 2017

Anisotropy of yield stress disorder



The location of the weakest sites is highly dependent on the orientation of the remote applied stress.



Forward thresholds $\Delta\sigma_+^Y$ $\Delta\sigma_{\pm}^Y = \Delta\sigma_+^Y - \Delta\sigma_-^Y$ Backward thresholds $\Delta\sigma_-^Y$

A tool to study memory effects in glasses ?