

The Secret Lives of Molecular Clouds

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Back to the Galaxy II

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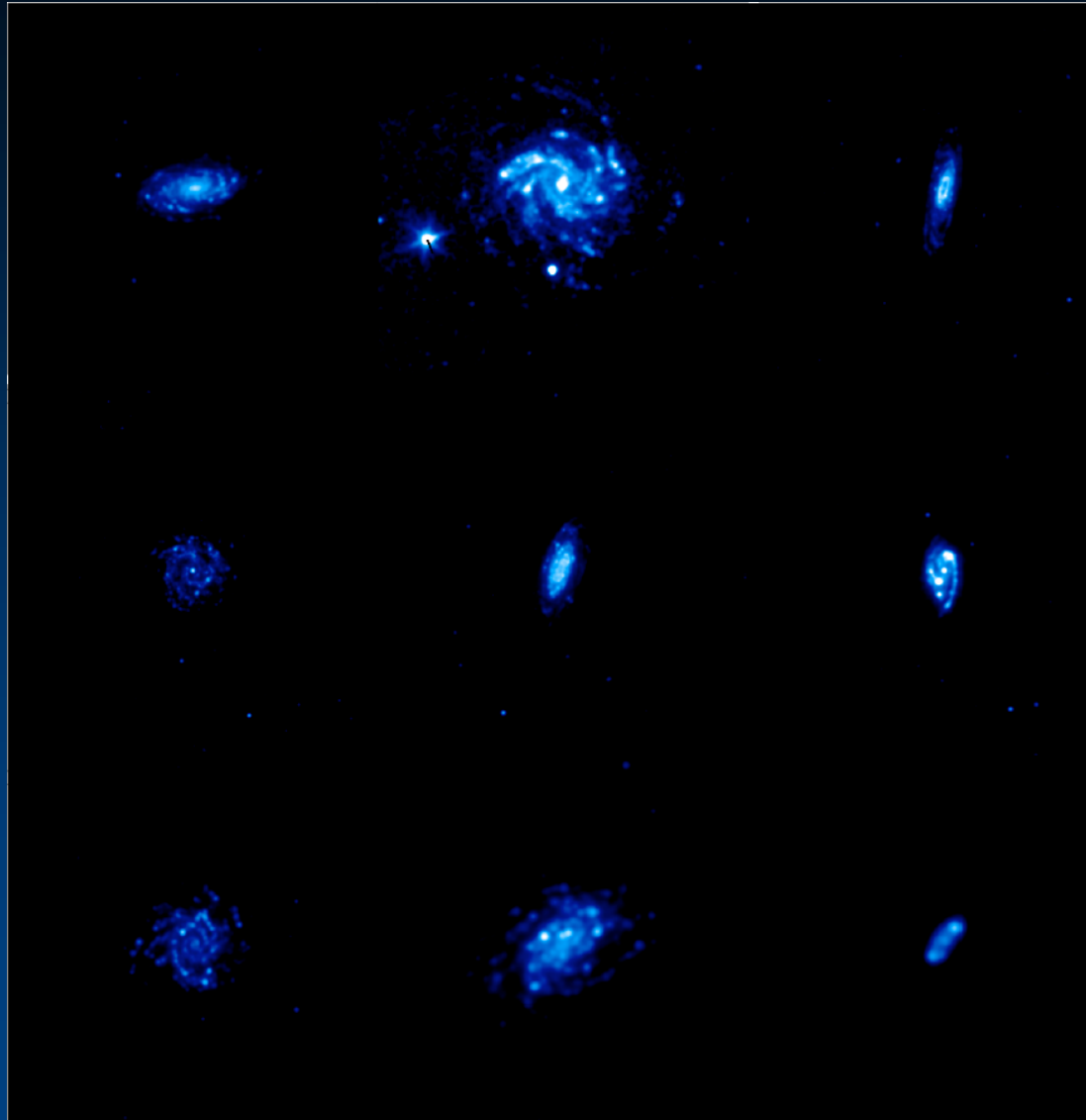
Outline

- Embarrassing observational facts about star formation
- Turning gas into stars
 - Making molecular clouds
 - Turning molecular clouds into stars
 - Feedback in molecular clouds

Observations



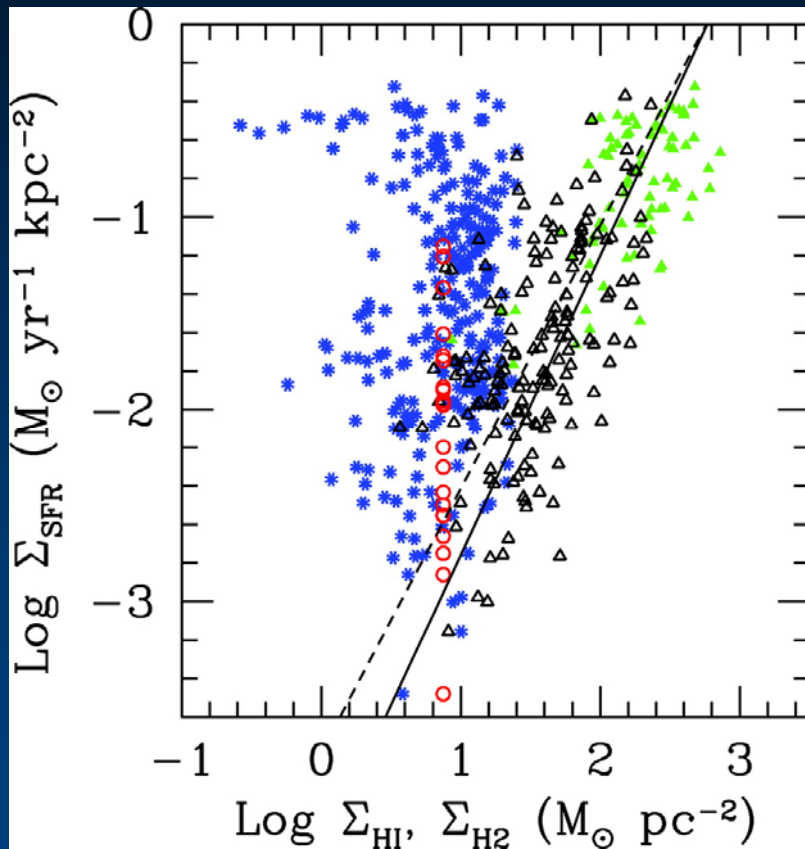
Stars Do Not Form in Gas



SINGS + GALEX
+ THINGS +
SONG (animation
borrowed from N.
Gnedin)

SFR distributions from 24 μm SINGS + GALEX

Stars Form in Molecular Gas



The SFR in a galaxy correlates well with the molecular gas surface density, and only poorly with the HI.

SFR vs. surface densities of HI (blue asterisks) and H₂ (black and green triangles) in M51a (Kennicutt et al. 2007)

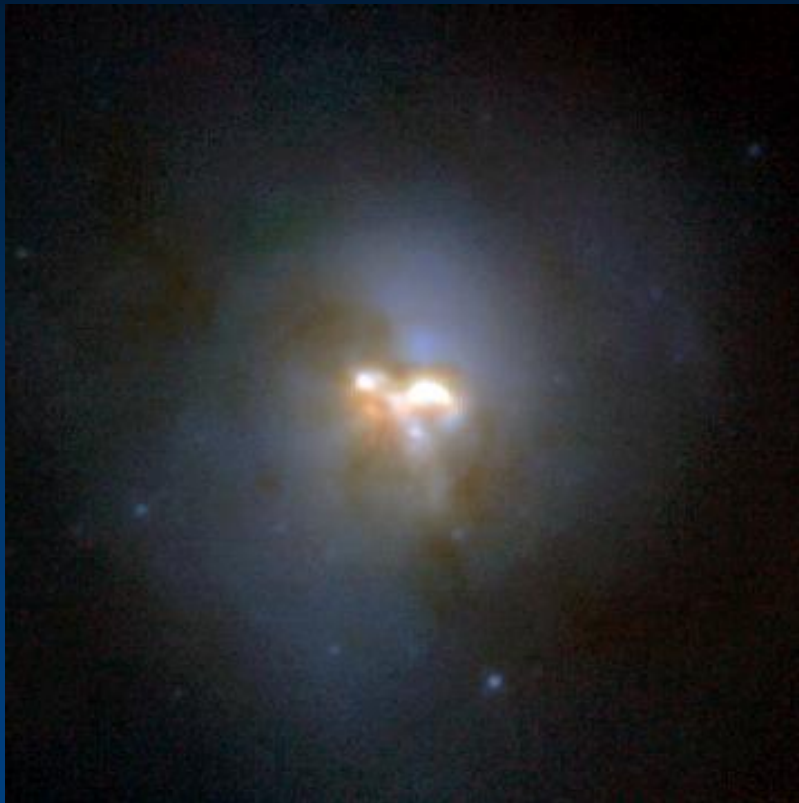
Even once H_2 forms, SF is slow...

(Zuckerman & Evans 1974; Rownd & Young 1999; Wong & Blitz 2002)

- The MW disk contains $\sim 10^9 M_\odot$ of gas in giant molecular clouds
- GMCs have $n_{\text{H}} \sim 100 \text{ cm}^{-3}$, $t_{\text{ff}} \sim 4 \text{ Myr}$
- If GMCs were collapsing, the SFR would be $\sim 10^9 M_\odot / 4 \text{ Myr} = 250 M_\odot / \text{yr}$
- Observed SFR in MW is $\sim 3 M_\odot / \text{yr}$, lower by a factor of ~ 100
- Numbers similar in nearby galaxies

...even in starbursts...

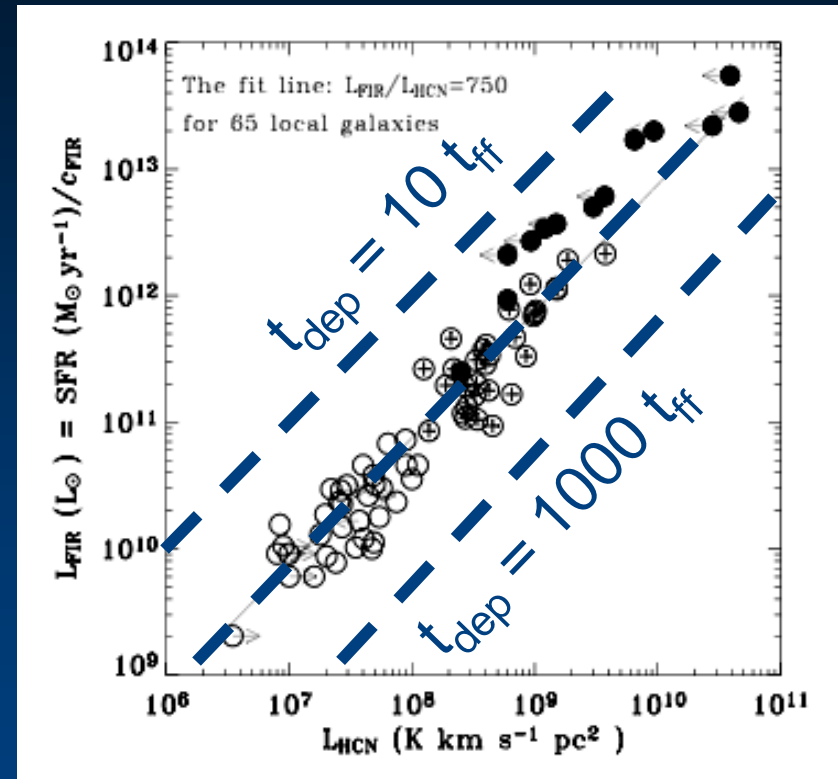
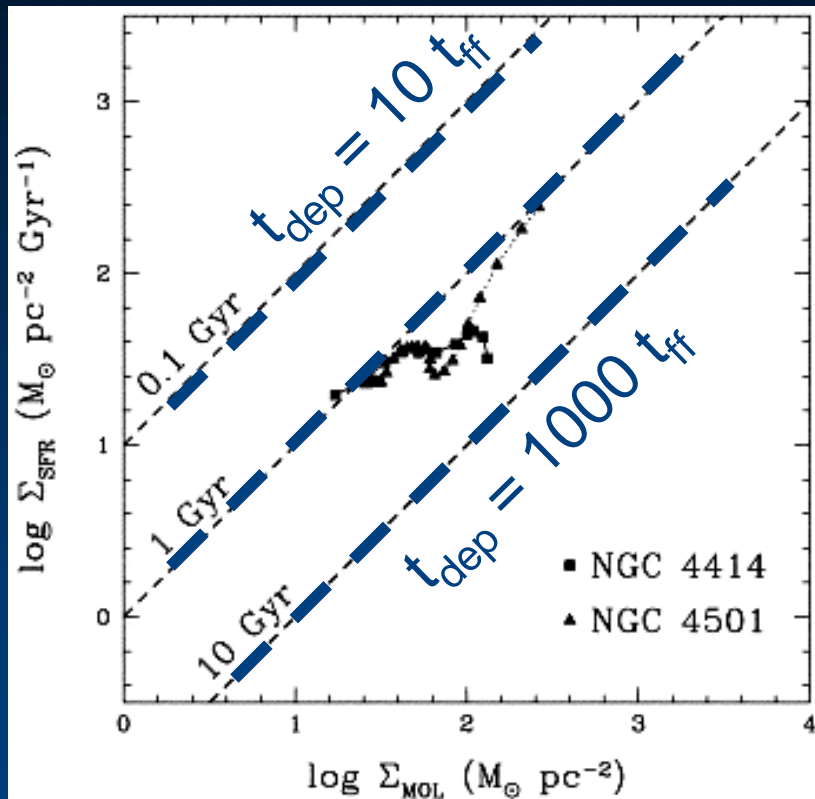
(Downes & Solomon 1998)



Arp 220 imaged by HST/NICMOS,
Thompson et al. 1997

- Example: Arp 220
- ISM mass $2 \times 10^9 M_{\odot}$ in molecular gas
- ISM density 10^4 cm^{-3} , $t_{\text{ff}} \sim 0.4 \text{ Myr}$
- Suggested SFR $\sim 5000 M_{\odot} / \text{yr}$
- Actual SFR $\sim 50 M_{\odot} / \text{yr}$: too small by factor of 100

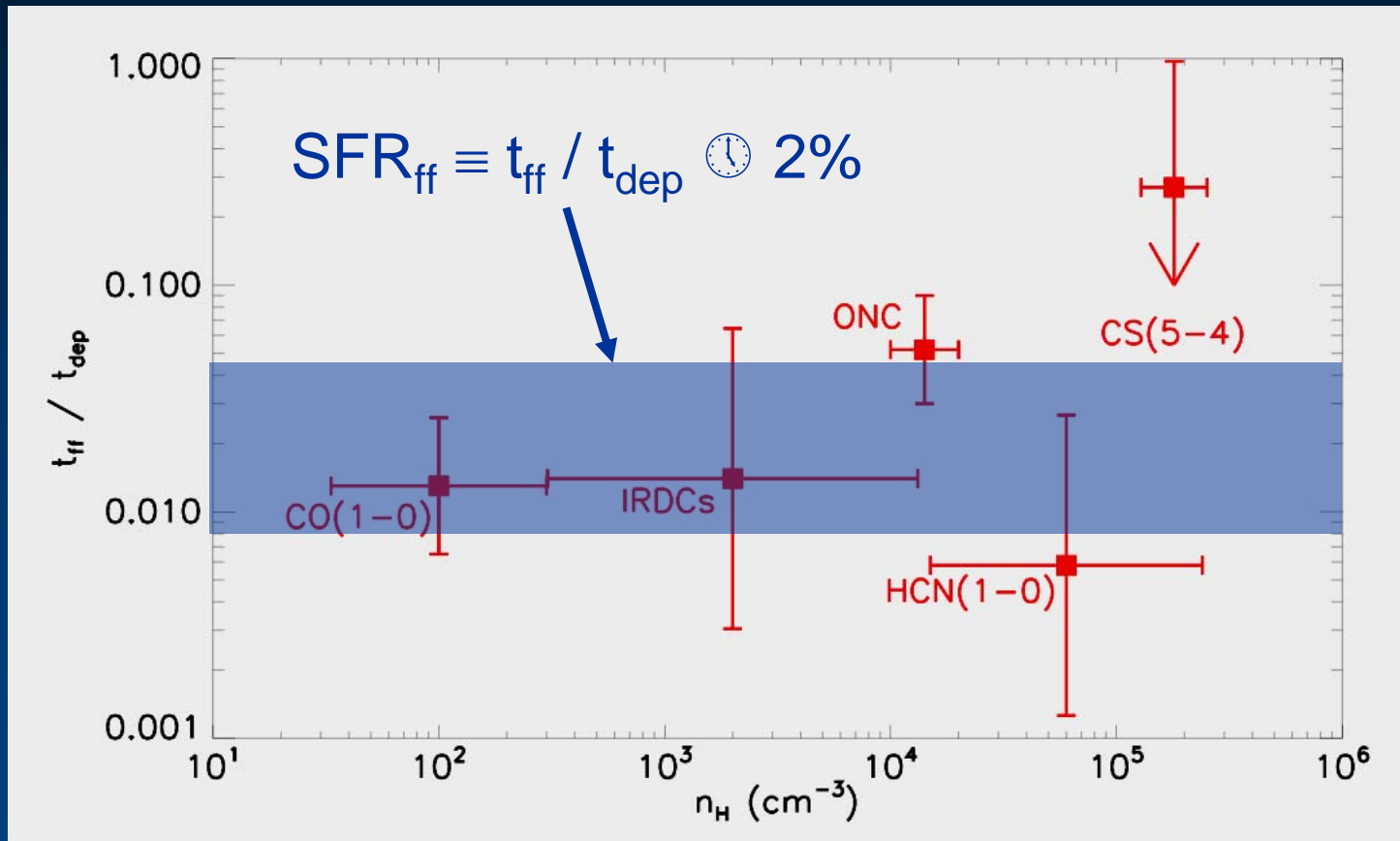
...even in dense gas...

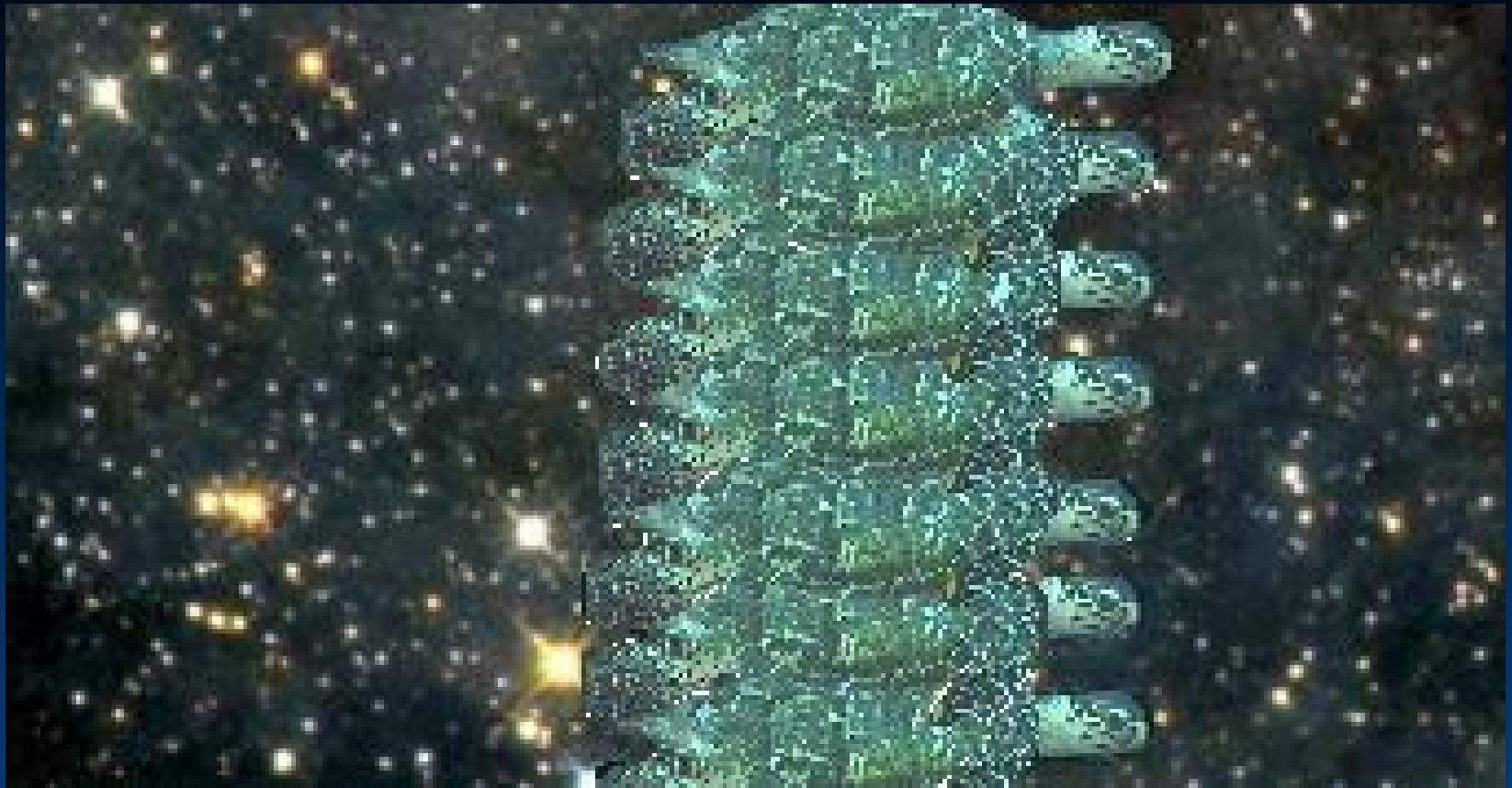


Depletion time as a function of Σ_{H_2} for 2 local galaxies (left, Wong & Blitz 2002) and as a function of L_{HCN} for a sample of local and $z \sim 2$ galaxies (below, Gao et al. 2007)

Now the Good News: There is a Universal SFR!

(Tan, Krumholz, & McKee 2006; Krumholz & Tan 2007)





In other words:
so far it's turtles all the way down...

Implications of Slow Star Formation

- For people who care about galaxies:
 - **Bad news:** you can simulate formation of GMCs with an approximate treatment of H_2 formation, but the SFR in GMCs is set at very small scales. Galaxy-scale simulations are stuck with subgrid models for that.
 - **Good news:** once molecules form, the SFR seems to follow a universal law that $\sim 1 - 5\%$ of the gas goes into stars per t_{ff} , independent of density.

Implications of Slow Star Formation, Part II

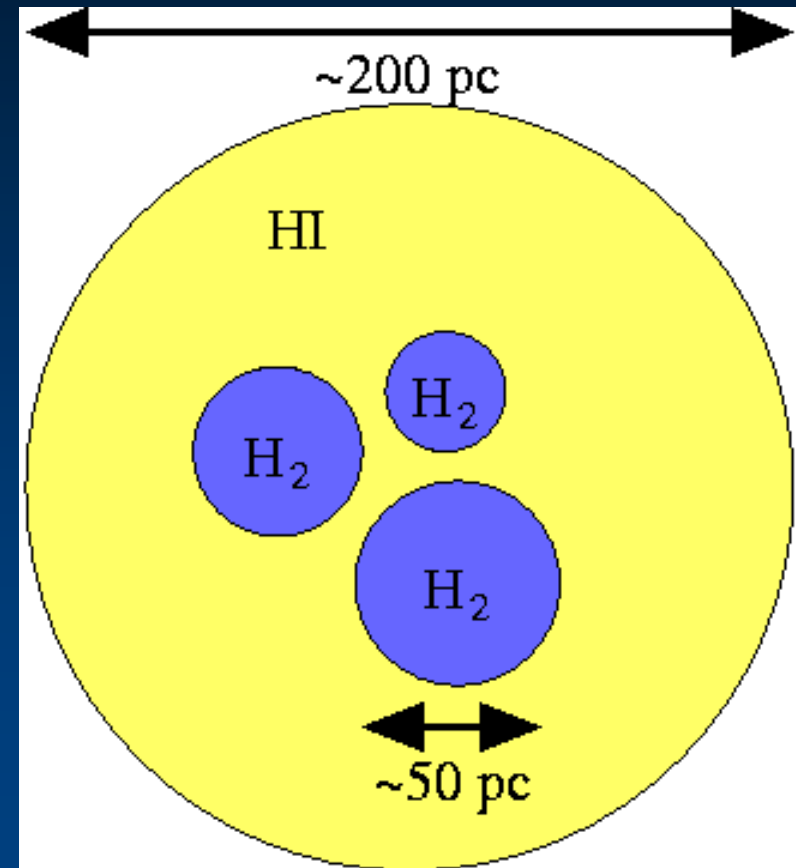
- For star formation theorists:
 - **Task 1:** figure out what determines what fraction of the gas will form molecules, since this controls which gas is “eligible” for form stars
 - **Task 2:** understand what sets the universal few percent per t_{ff} in the molecular gas

From HI to Stars



Step 1: Making Molecules

- Molecules reside in giant molecular clouds (GMCs) that are part of atomic-molecular complexes
- The outer parts are dissociated by interstellar Lyman-Werner photons
- Inner parts are shielded by dust and H_2 self-shielding
- Goal: compute HI and H_2 mass fractions



Dissociation Balance in Atomic-Molecular Complexes

(Krumholz, McKee, & Tumlinson, 2008a, ApJ, in press)

The basic equations for this system are *chemical equilibrium* and *radiative transfer*.

$$n_{\text{HI}} n_{\mathcal{R}} = n_{\text{H}_2} \int d\Omega \int d\nu \sigma_{\text{H}_2} f_{\text{diss}} I_{\nu} / (h\nu)$$

$$\hat{e} \cdot \nabla I_{\nu} = -(n_{\text{H}_2} \sigma_{\text{H}_2} + n \sigma_{\text{d}}) I_{\nu}$$

Idealized problem: spherical cloud of radius R , density n , dust opacity σ_{d} , H_2 formation rate coefficient \mathcal{R} , immersed in radiation field with UV photon number density E_0^* , find fraction of mass in HI and H_2 .

Decrease in radiation intensity = Absorptions by H_2 molecules + dust grains

Calculating Molecular Fractions

To good approximation, solution only depends on two dimensionless numbers:

$$\tau_R = n\sigma_d R$$

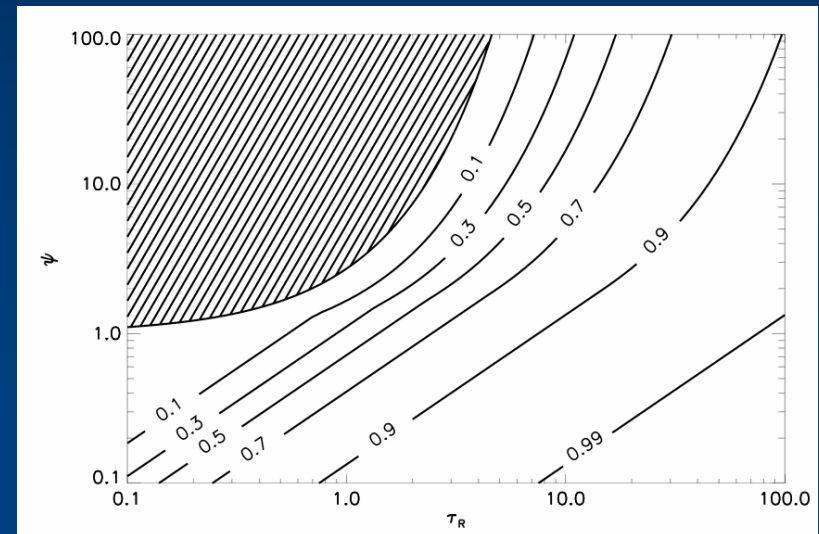
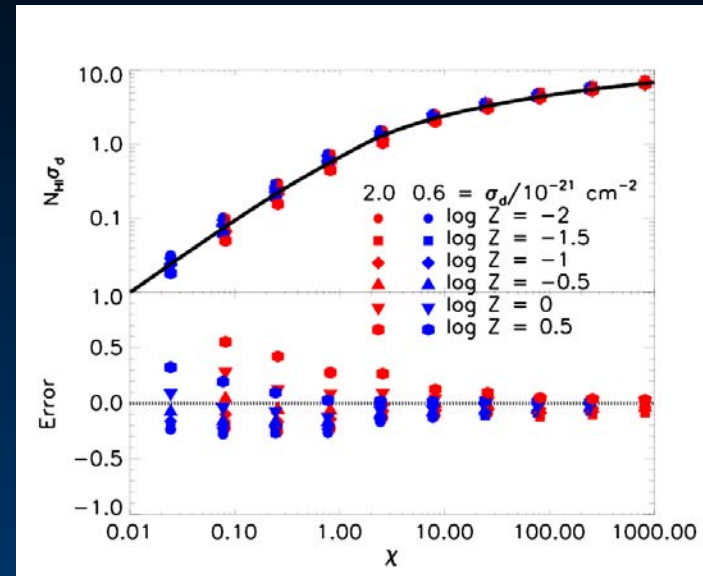
$$\psi = \phi \frac{f_{\text{diss}} \sigma_d E_0^*}{n\mathcal{R}}$$

Approximate solution:

$$f_{\text{H}_2, \text{vol}} \approx 1 - \frac{3\psi}{4(\tau_R + 0.2\psi)}$$

Top: analytic solution for location of HI / H₂ transition vs. exact numerical result

Bottom: H₂ volume fraction vs. ψ , τ_R



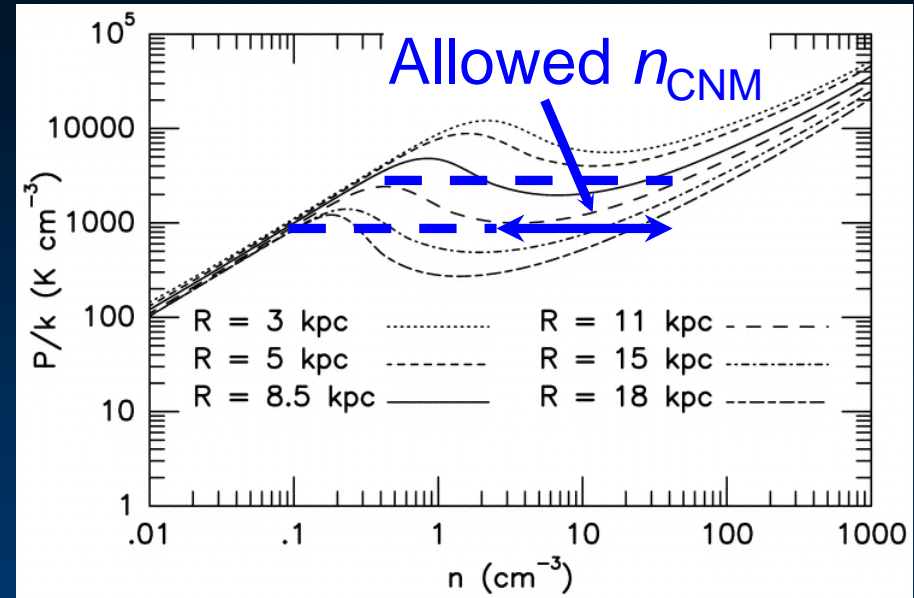
Shielding Layers in Galaxies

(Krumholz, McKee, & Tumlinson, 2008b, submitted)

What is $\psi \propto \sigma_d E_0^* / nR$?

- Dust opacity σ_d and H_2 formation rate R both $\propto Z$, so $\sigma_d / R \sim \text{const}$
- CNM dominates shielding, so n is the CNM density
- CNM density set by pressure balance with WNM, and $n_{\text{CNM}} \propto E_0^*$, with weak Z dependence.

$\Rightarrow \psi \propto \sigma_d E_0^* / nR \sim 1$ in all galaxies!

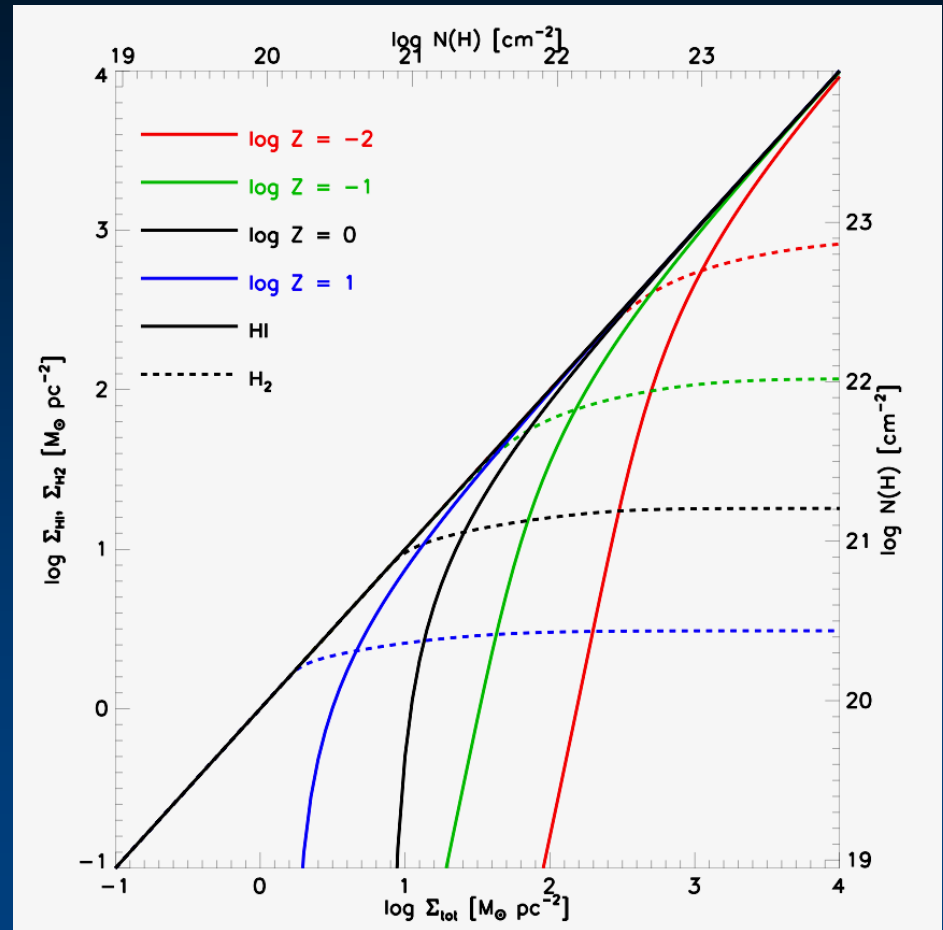


FGH curves for MW (Wolfire et al. 2003)

Predictions for H₂ Content

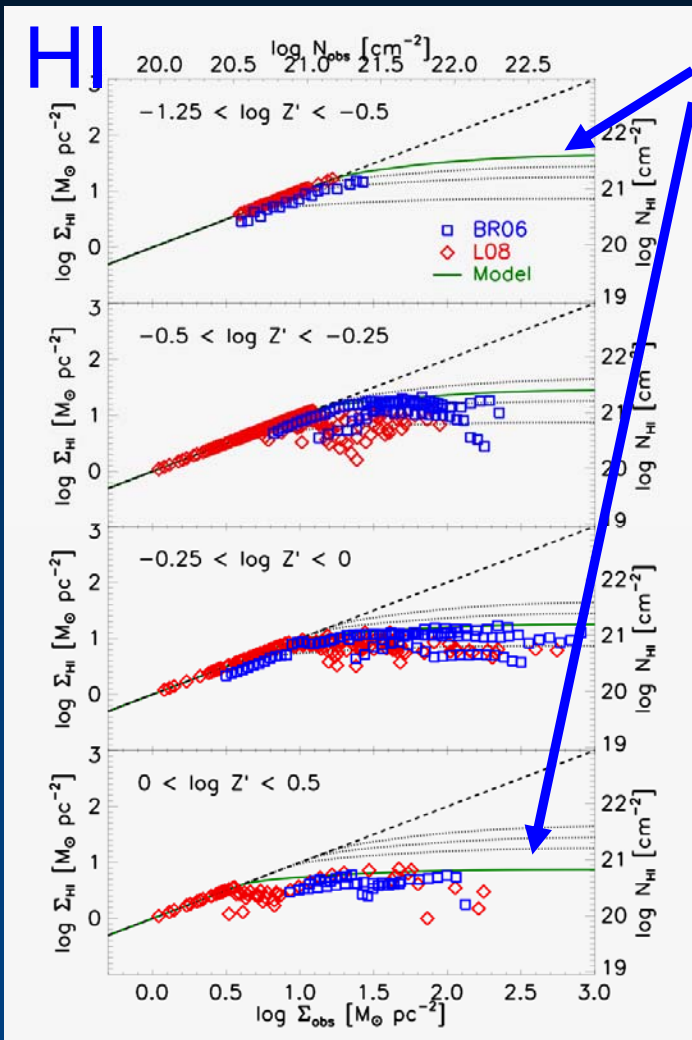
Compute τ_R from column density Σ , metallicity Z , and pressure balance between molecules and CNM.

Then use solution for H₂ fraction vs. ψ , τ_R to compute molecular content as a function of Σ , Z



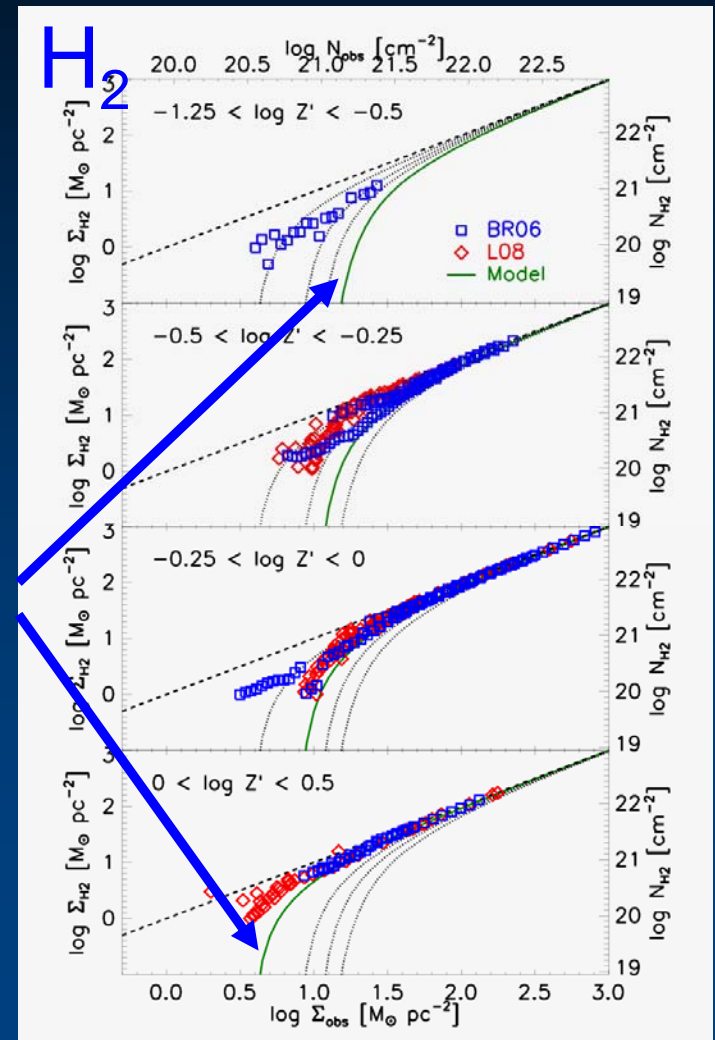
Reality Check...

Compare model to BIMA SONG (Blitz & Rosolowsky 2006) and THINGS (Leroy et al. 2008) surveys, with galaxies binned by metallicity



Matches
observed
saturation of HI
at $\sim 10 \text{ M}_{\odot} \text{pc}^{-2}$,
with higher Σ_{HI} at
low metallicity!

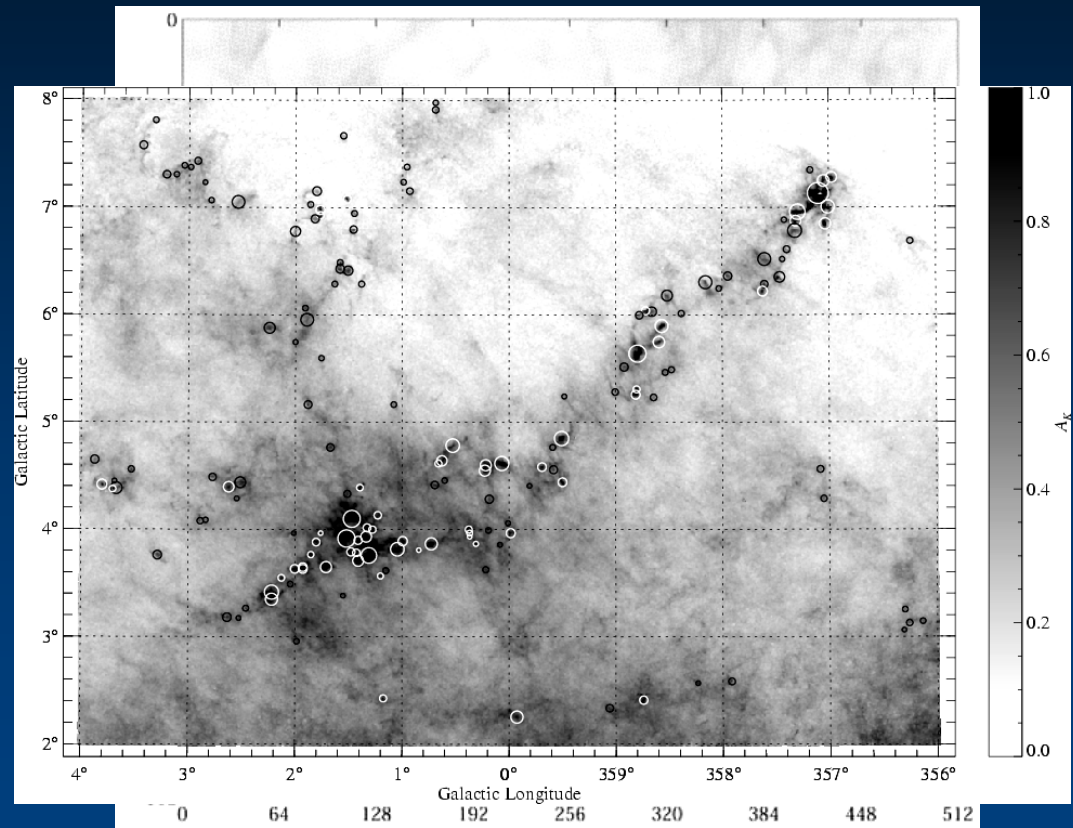
Matches column
needed for
molecules to
appear, with
higher Σ at lower
metallicity!



Step 2: Turning Molecules into Stars (Slowly)

(Krumholz & McKee 2005)

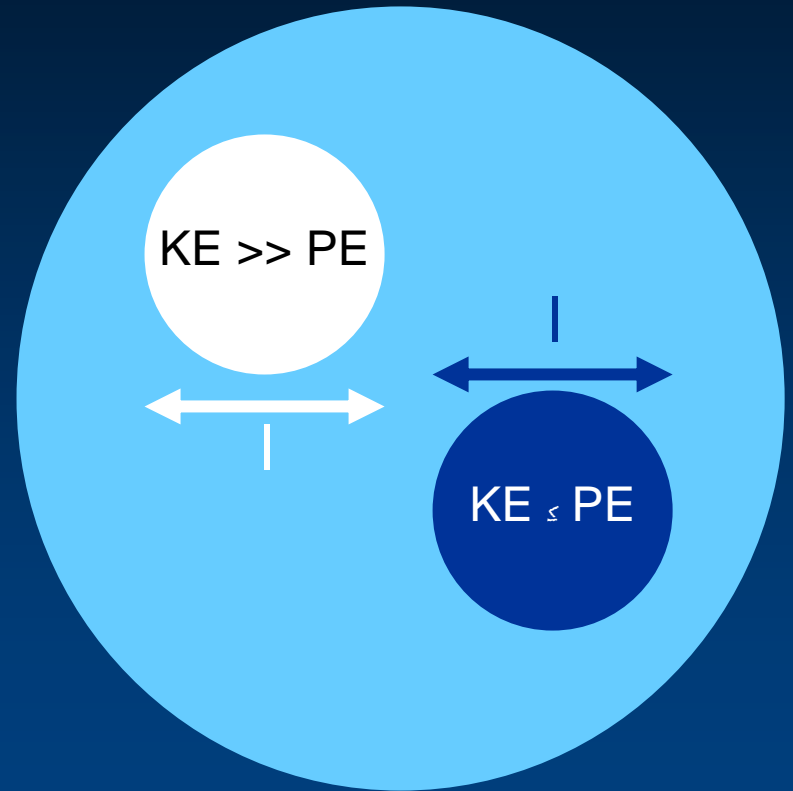
- Most GMC gas is in low density “envelopes”, not dense “cores”
- GMCs are very turbulent, $\mathcal{M} \sim 30$
- Simulations of turbulence give core-envelope structure



Column density map of the Pipe Nebula, with cores circled, Alves, Lombardi, & Lada 2007
Column density map of a simulation of MHD turbulence with $\mathcal{M} = 10$, Li et al. 2004

How Turbulence Sets the SFR

- On large scales, GMCs have $\alpha \approx 1$ (i.e. $PE \approx KE$)
- Linewidth-size relation: $\sigma_v \approx c_s (l / \lambda_s)^{1/2}$
- In average region, $M \propto l^3$
 $\Rightarrow KE \propto l^4$, $PE \propto l^5$
 $\Rightarrow KE \gg PE$
- **Hypothesis:** SF only occurs in regions where $PE \gtrsim KE$ and $P_{th} \gtrsim P_{ram}$
- Only overdense regions meet these conditions
- Required overdensity is given by $\lambda_J \leq \lambda_s$, where $\lambda_J = c_s [\pi / (G\rho)]^{1/2}$



Calculating the SFR

QuickTime™ and a
YUV420 codec decompressor
are needed to see this picture.

$$\lambda_s = \lambda_j$$

- Density PDF in turbulent clouds is lognormal; width set by \mathcal{M}
- Integrate over region where $\lambda_j \leq \lambda_s$, to get mass in “cores”, then divide by t_{ff} to get SFR
- Result:

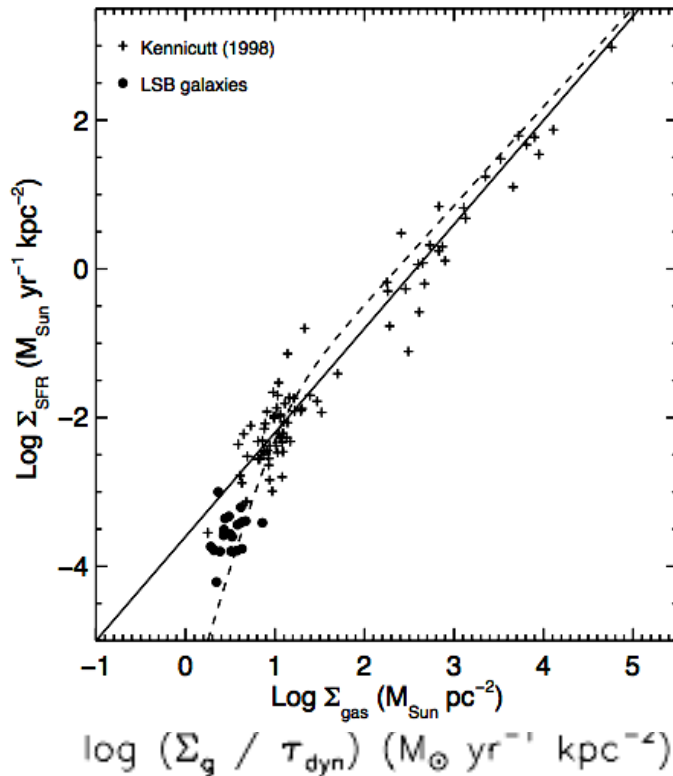
$$\text{SFR}_{\text{ff}} \approx 0.073 \alpha^{-0.68} \mathcal{M}^{-0.32}$$

$\text{SFR}_{\text{ff}} \sim 1\text{-}5\%$ for any turbulent, virialized object

Predicting the Kennicutt Law

Plus signs: solid line: Kennicutt 1998
(1998) solid line: LSB galaxies: KM05 model 2008

Dashed line: KM05 model



- We don't know t_{ff} , \mathcal{M} in GMCs in other galaxies
- Estimate by assuming
 - GMC mass \sim Jeans mass
 - GMCs are virialized
 - GMC internal pressure balances weight of ISM
- Result: a Kennicutt law in terms of observables

$$\frac{\dot{\Sigma}_*}{M_\odot/\text{yr}/\text{kpc}^2} = 130 f_{\text{mol}} \left(\frac{\Omega}{\text{Myr}^{-1}} \right)^{1.3} \left(\frac{\Sigma_g}{\text{g}/\text{cm}^2} \right)^{0.7}$$

Step 3: Stellar Feedback

(Krumholz, Matzer, & McKee 2007)



30 Doradus HII region, MCELS

- All observed GMCs turbulent, $\alpha \approx 1$
- GMC lifetime is ~ 30 Myr (Blitz et al. 2007), $t_{\text{cr}} \sim 7$ Myr \Rightarrow need energy injection
- GMCs in have very similar properties in all nearby galaxies
- **Hypothesis:** SF feedback regulates GMC density and virial balance

A Semi-Analytic GMC Model

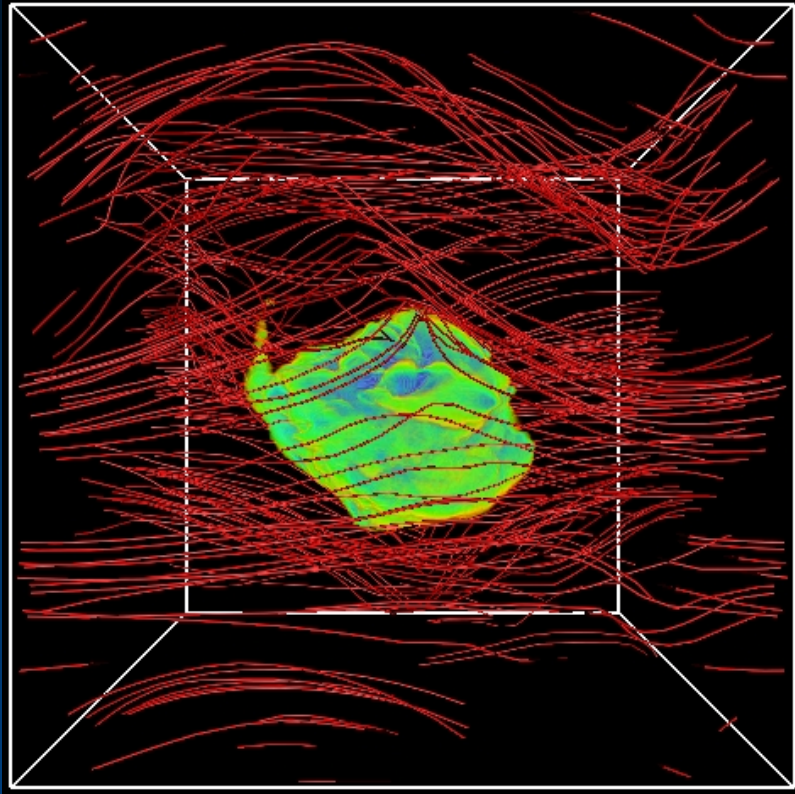
Follow evolution of:

$M_{\text{gas}}, M_*,$
 $R, dR/dt, \sigma$

- Model GMC mass, energy, momentum budgets
- Include feedback driving turbulence, mass loss
- Use 1D model
 - Bad: real GMCs not spheres
 - Good: can solve exact equations: non-equilibrium virial, energy equations

$$\ddot{I}/2 = 2(\mathcal{T} - \mathcal{T}_0) + \mathcal{W} + \mathcal{B} - (1/2)(d/dt) \int (\rho \mathbf{v} r^2) \cdot d\mathbf{S}$$
$$\dot{E} + \int \rho(v^2/2 + e + \phi + P_s/\rho) \mathbf{v} \cdot d\mathbf{S} = \Gamma - \Lambda$$

HII Region Feedback



Simulation of HII region in a magnetized cloud, Krumholz & Stone 2008, in preparation

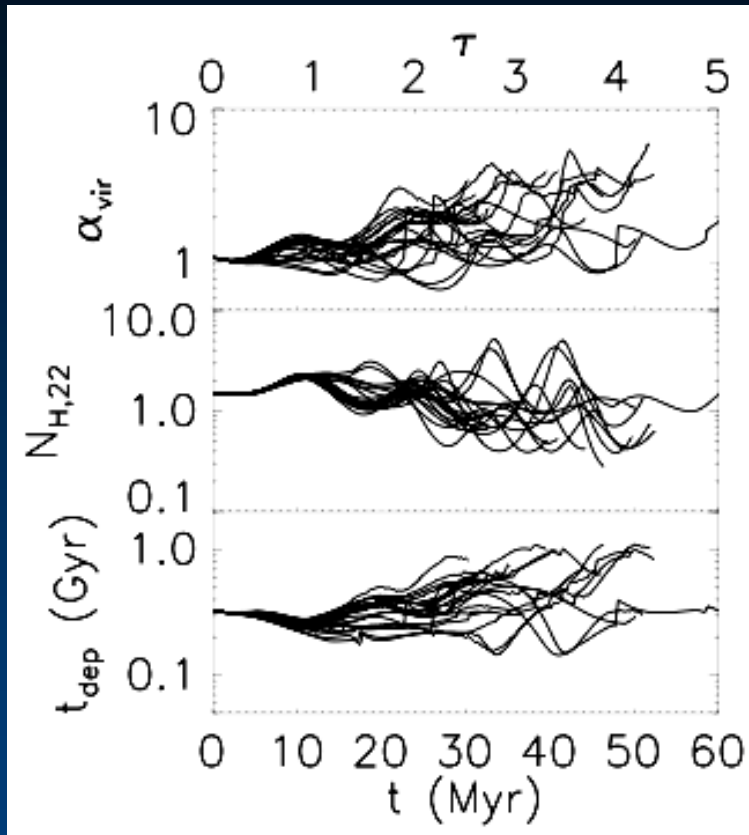
- HII regions dominate feedback (even beat SNe)
- Use modified Spitzer solution to get HII region expansion
- Assume all HII regions blister, lead to mass loss
- Compute energy injection assuming shells break up, merge with turbulence

Simulation of HII Region Driving

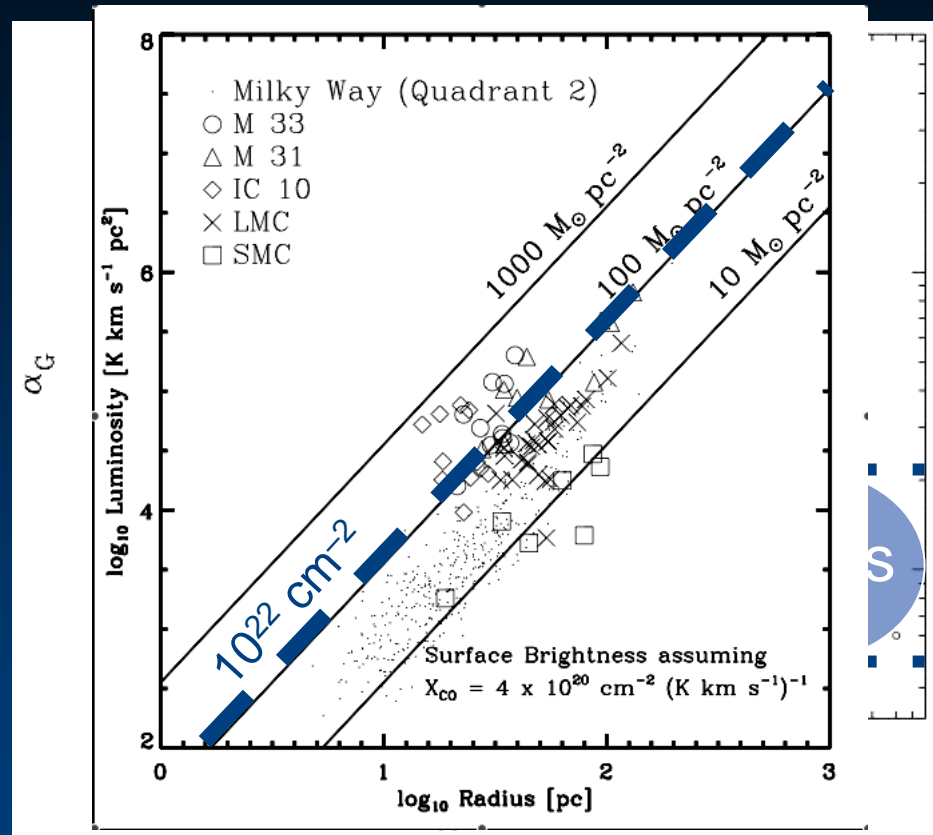
(Krumholz, Stone, & Gardiner 2007; Krumholz & Stone 2008, in preparation)

QuickTime™ and a
YUV420 codec decompressor
are needed to see this picture.

Quasi-Equilibrium Clouds



Sample of runs for $M_{\text{cl}} = 5 \times 10^6 M_{\odot}$



GMC mass vs. radius in the Local Group, Blitz et al. 2007
 α vs. mass for GMCs, Heyer et al. 2001

Feedback keeps GMCs close to equilibrium ($\alpha \sim 1$) at a preferred (column) density

Conclusions

- Star formation is slow because (1) **only molecular gas** makes stars; (2) **even this gas** forms stars at only $\sim 2\%$ per t_{ff}
- The molecular fraction is determined by column density and metallicity; **low Z galaxies require very high Σ to make H_2**
- The SFR in the H_2 is determined by **turbulence** driven by SF feedback
- Feedback energy balance imposes a **column density $\sim 10^{22} \text{ cm}^{-2}$ in GMCs**