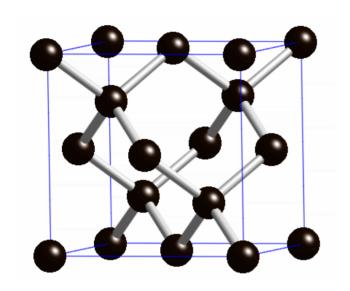
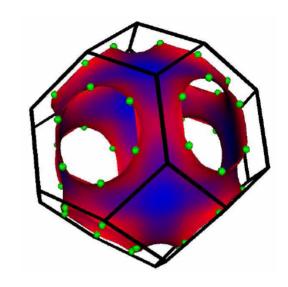
Frustrated diamond lattice antiferromagnets





Jason Alicea (Caltech)

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Simon Trebst (Station Q)

Bergman et al., Nature Physics 3, 487 (2007).

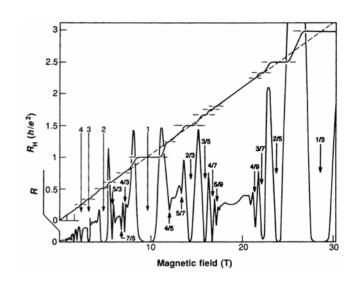
Outline

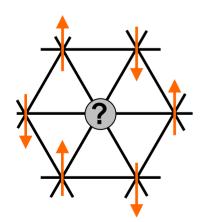
Introduction

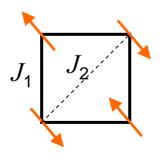
- Frustration, degeneracy, & emergent phenomena
- Diamond lattice antiferromagnets
- Overview of experiments
- Theory of frustrated diamond lattice antiferromagnets
 - Ground states (highly degenerate spirals)
 - Stability & order-by-disorder
 - Monte Carlo simulations
 - Spiral spin liquid
- Comparison to experiment
- Summary & future directions

Frustration, degeneracy, & emergent phenomena

- Frustration = all interactions not satisfiable simultaneously
- General principle: The presence of many competing states often leads to interesting physics
 - Quantum Hall effect
 - High- T_c superconductors
 - Frustrated magnets (Mott insulators)

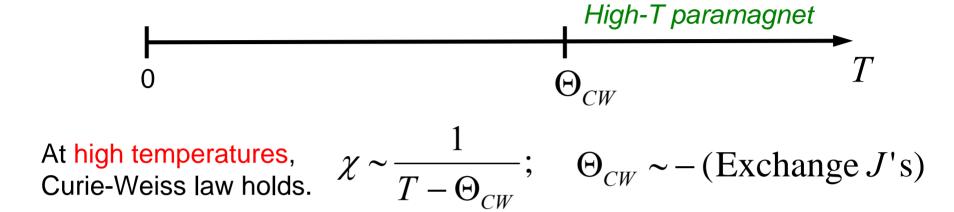




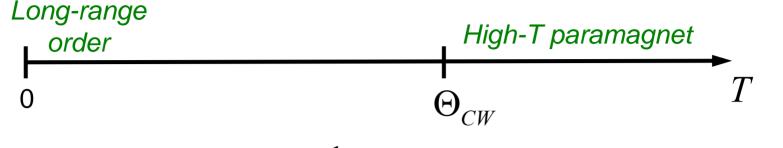


- ⇒ Highly degenerate ground states (pyrochlore, kagome, FCC, etc.)
 - High sensitivity to perturbations
 - Spin-glass behavior
 - Spin-liquid physics
 - Order-by-disorder

Experimental signatures of frustration



Experimental signatures of frustration

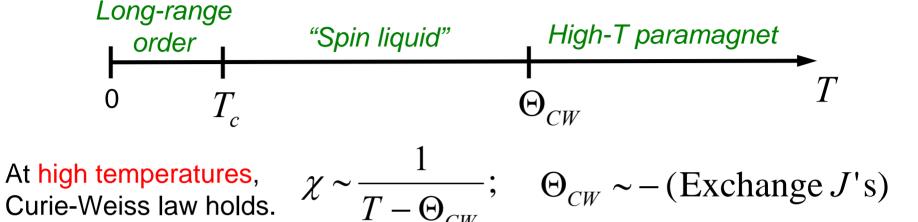


At high temperatures, Curie-Weiss law holds.
$$\chi \sim \frac{1}{T - \Theta_{CW}}$$
; $\Theta_{CW} \sim - (\text{Exchange } J's)$

At low temperatures, systems typically order.

Useful diagnostic: "frustration parameter" $f = |\Theta_{CW}|/T_c$

Experimental signatures of frustration



At low temperatures, systems typically order.

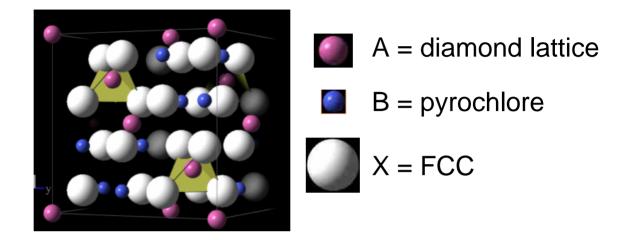
Useful diagnostic: "frustration parameter" $f = |\Theta_{\rm CW}|/T_c$

Highly frustrated systems $\Rightarrow f > 5 - 10$

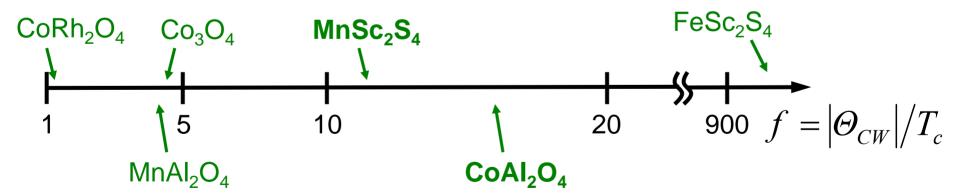
- Key challenges: Low-*T* ordering mechanisms Characterizing "spin liquid" correlations

Frustrated diamond lattice antiferromagnets: Materials

Many materials take on the normal spinel structure: AB₂X₄



Focus: spinels with magnetic A-sites (only)

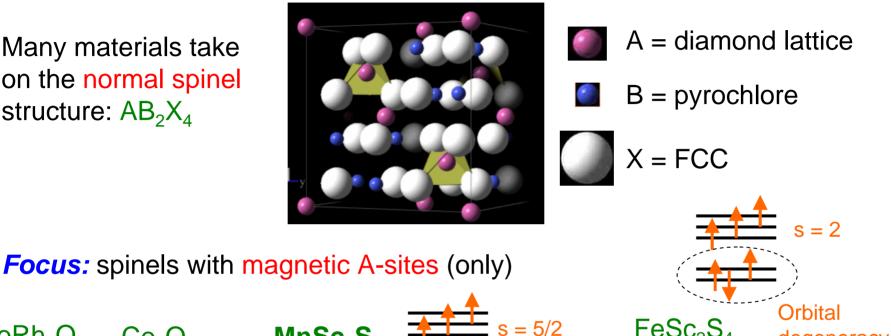


Very limited theoretical understanding...

V. Fritsch et al. PRL **92**, 116401 (2004); N. Tristan et al. PRB **72**, 174404 (2005); T. Suzuki *et al.* (unpublished)

Frustrated diamond lattice antiferromagnets: Materials

Many materials take on the normal spinel structure: AB₂X₄



FeSc₂S₄ MnSc₂S₄ CoRh₂O₄ Co_3O_4 degeneracy 900 $f = |\Theta_{CW}|/T_{c}$ 10

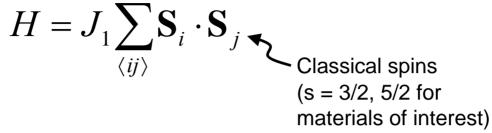
CoAl₂O₄ = MnAl₂O₄

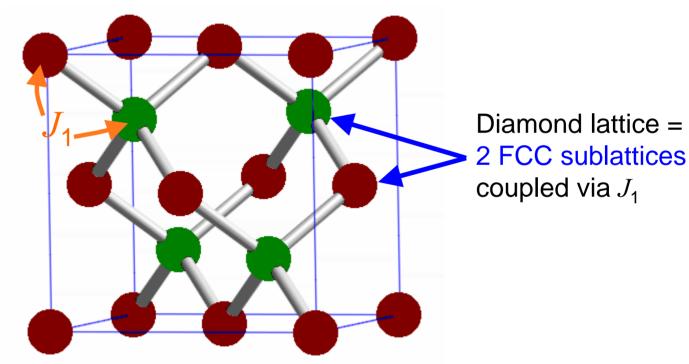
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Frustration on the bipartite diamond lattice??

Naïve Hamiltonian:

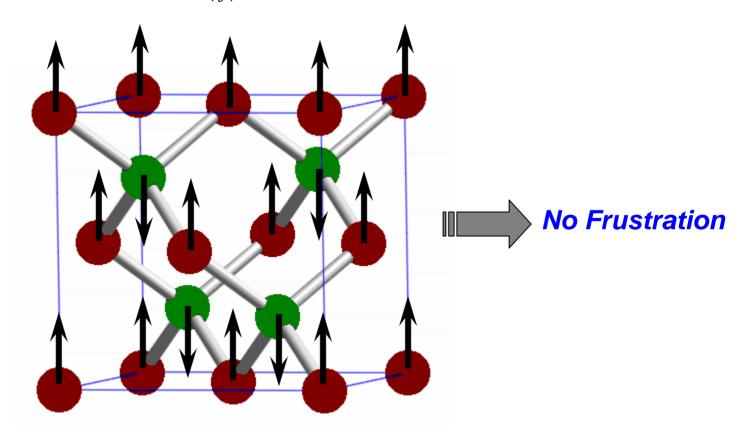




Frustration on the bipartite diamond lattice??

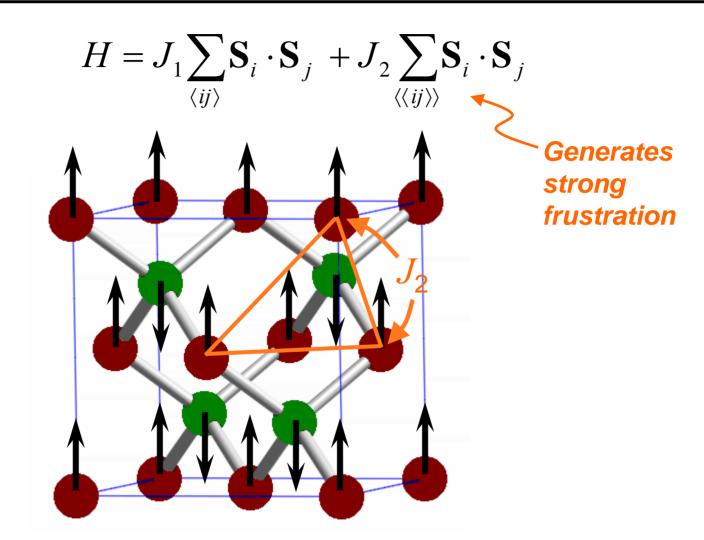
Naïve Hamiltonian:

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$



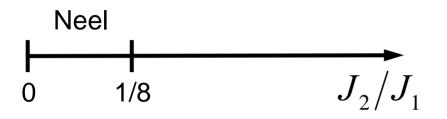
Frustration on the diamond lattice

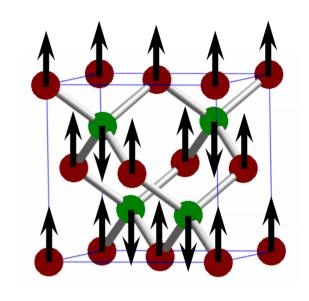
Remedy: 2nd neighbor exchange



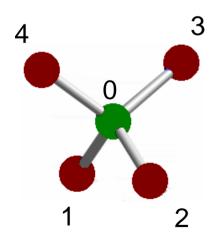
 J_1 and J_2 expected to be comparable due to similarity in exchange paths

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle \langle ij \rangle \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$



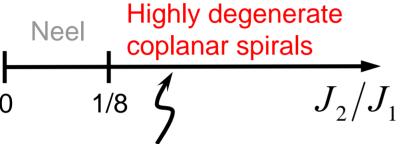


Useful rewrite of Hamiltonian:

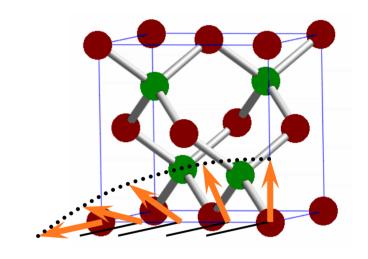


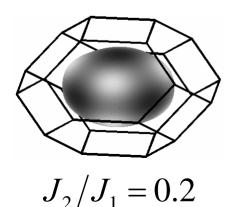
$$H = J_1 \sum_{t} [\mathbf{S}_0 + (\mathbf{S}_1 + \mathbf{S}_2 + \mathbf{S}_3 + \mathbf{S}_4)/4]^2$$
$$+ (J_2 - J_1/8) \sum_{\langle\langle ij \rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle \langle ij \rangle \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

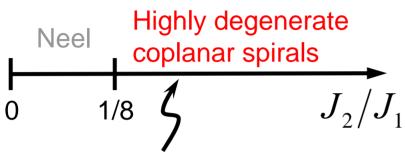


Direction & pitch of spirals characterized by a wavevector residing on a surface in momentum space!

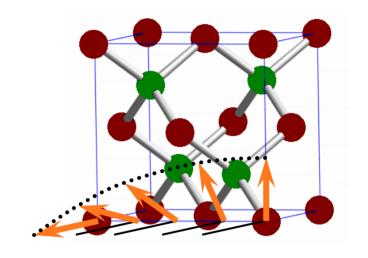


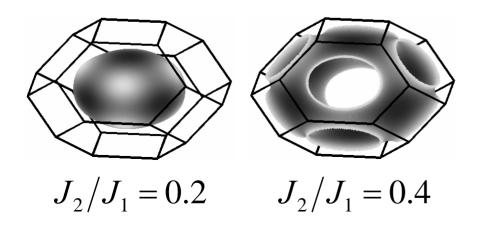


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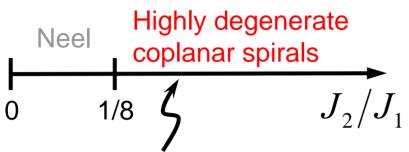


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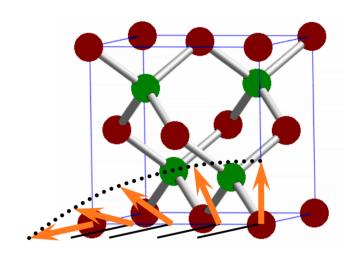


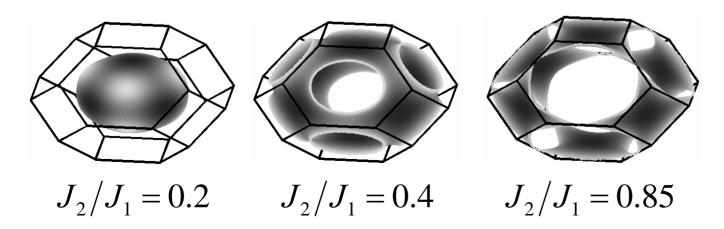


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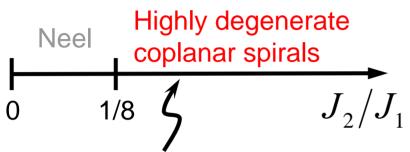


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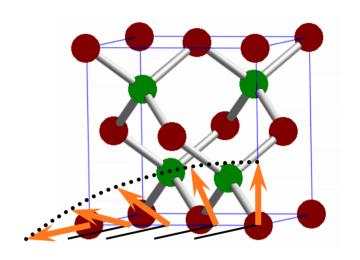


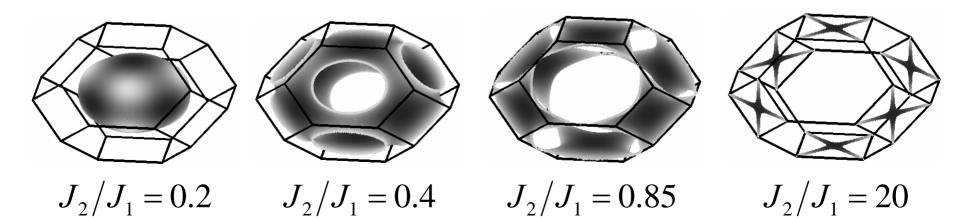


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Direction & pitch of spirals characterized by a wavevector residing on a surface in momentum space!





Low-T physics: Can long-range order occur?

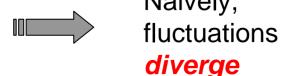
Stability nontrivial due to massive spiral degeneracy.

- Expand Hamiltonian in **fluctuations**:

$$\delta \mathbf{S}_i = \mathbf{S}_i - \langle \mathbf{S}_i \rangle$$
 Arbitrary ground state order

- At *T* = 0, branch of normal modes has infinite # of zeros!

$$\omega_0({f q})=0$$
 For all ${f q}$ on surface

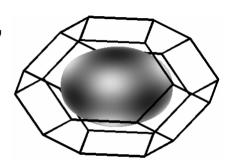


Naively, fluctuations
$$\langle \delta \mathbf{S}_{i}^{2} \rangle \sim T \int \frac{d^{3}\mathbf{q}}{\omega_{0}^{2}(\mathbf{q})} \rightarrow \infty$$

"Order-by-disorder" stabilization

Key ideas:

- Only symmetry-required zeros in $\omega_0(\mathbf{q})$ are the "Goldstone modes"
- Thermal fluctuations lifts the remaining "accidental" zeros ⇒ entropy stabilizes long-range order!
- **Needed:** finite-T corrections to $\omega_0(\mathbf{q})$ on "spiral surface"
 - Perturbation theory insufficient
 - Use self-consistent approach instead



• Answer: $\omega_T^2(\mathbf{q}) = \omega_0^2(\mathbf{q}) + T^{2/3}\Sigma(\mathbf{q})$

$$\langle \delta \mathbf{S}_{i}^{2} \rangle \sim T \int \frac{d^{3}\mathbf{q}}{\omega_{T}^{2}(\mathbf{q})} \sim T^{1/3}$$

(Fluctuations small at low T)

• Non-analytic *T*-dependence ⇒ unconventional thermodynamic behavior, e.g.,

$$C_{v} = A + BT^{1/3}$$

Aside on self-consistent approach

Expand Hamiltonian in fluctuations:

"Interaction" terms

$$\delta \mathbf{S}_{i} = \mathbf{S}_{i} - \langle \mathbf{S}_{i} \rangle; \qquad H = H_{2} + H_{3} + H_{4} + \dots$$

- Get self-energy self-consistently for divergent mode

$$\overline{\Sigma}(\mathbf{q}) = + + + \dots$$

$$\overline{\Sigma}(\mathbf{q}) = T \int_{\mathbf{k}} \Gamma(\mathbf{q}, \mathbf{k}) G(\mathbf{k}); \quad G(\mathbf{k}) = \frac{1}{\omega_0^2(\mathbf{k}) + \overline{\Sigma}(\mathbf{k})}$$

- For \mathbf{q} on surface, assume $\overline{\Sigma}(\mathbf{q}) \sim T^{\alpha} \Sigma(\mathbf{q})$

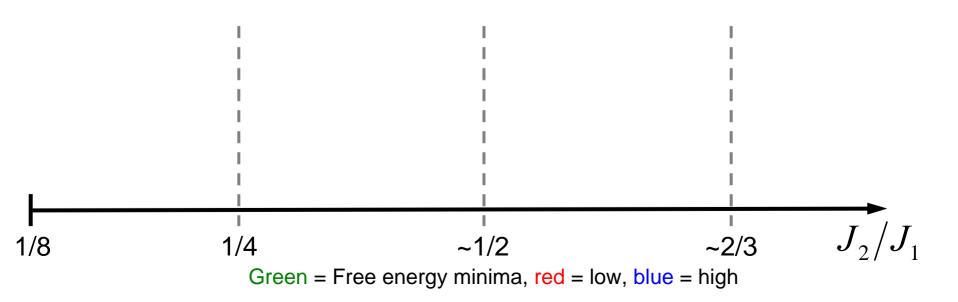
$$\omega_T^2(\mathbf{q}) = \omega_0^2(\mathbf{q}) + T^{2/3}\Sigma(\mathbf{q})$$

Long-range order occurs—but which state does entropy select?

 Need Free Energy for all Q on spiral surface

$$F(\mathbf{Q}) = E - TS(\mathbf{Q})$$

- Entropy favors states with highest density of nearby low-energy states
- Complex phase structure emerges:

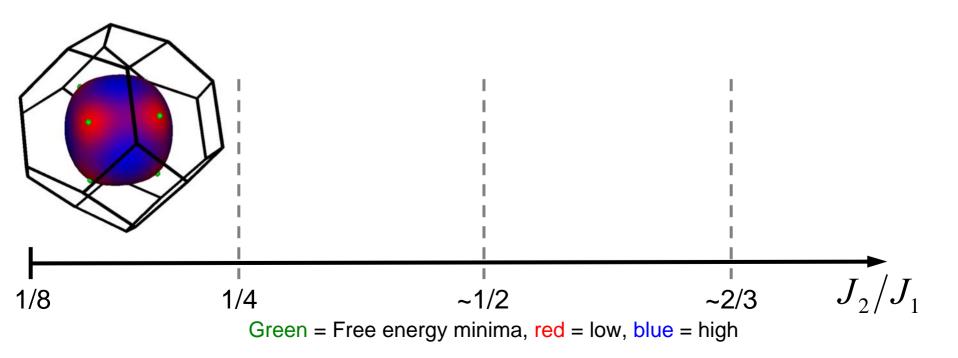


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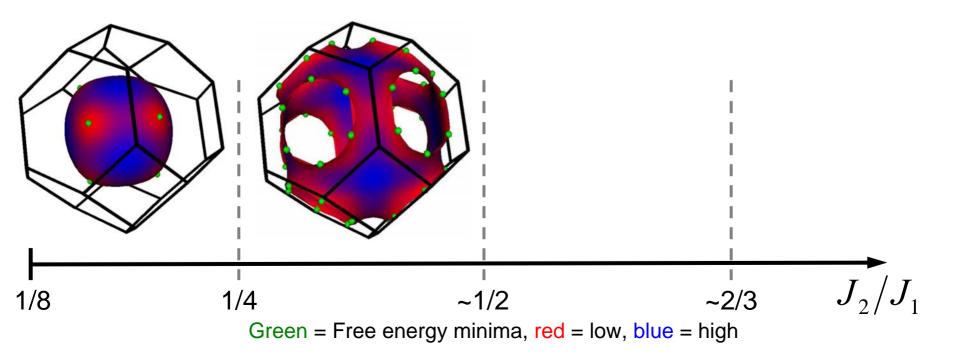


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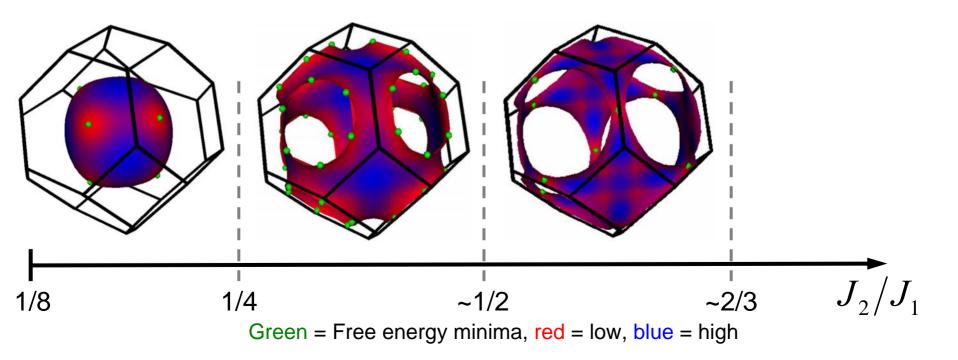


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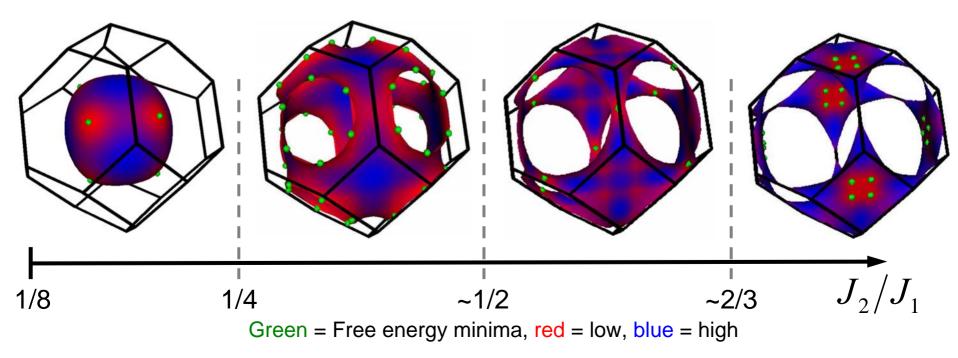


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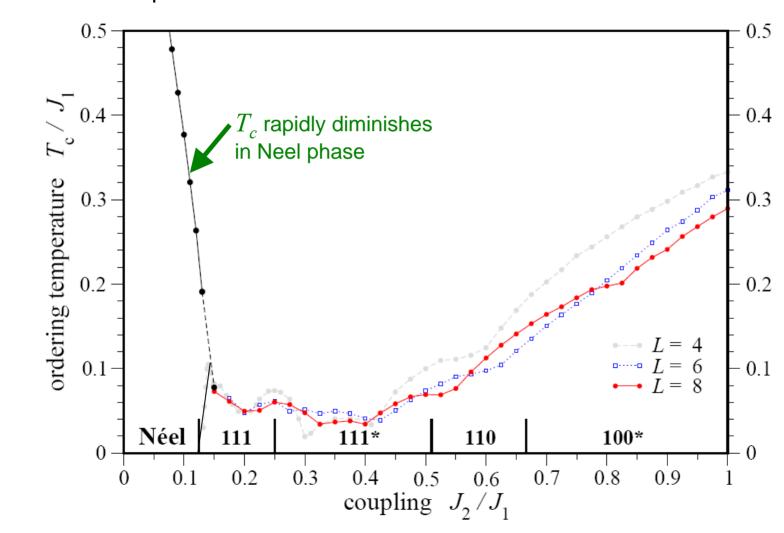
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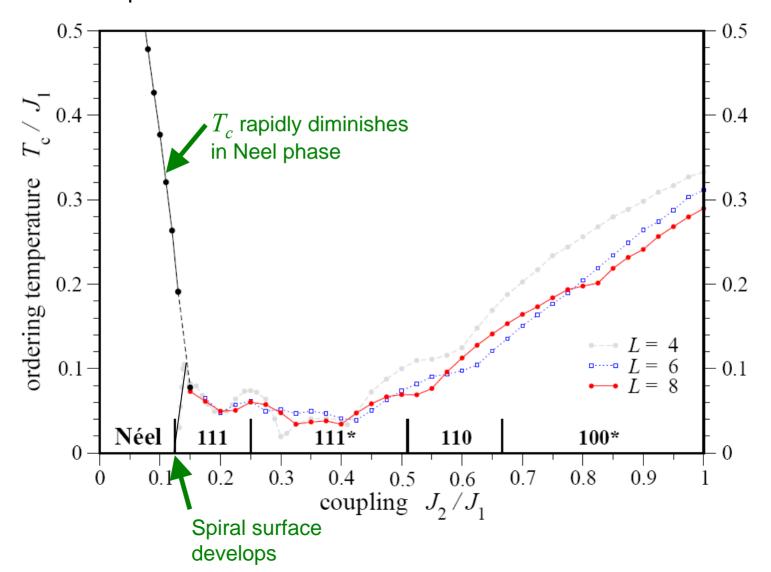
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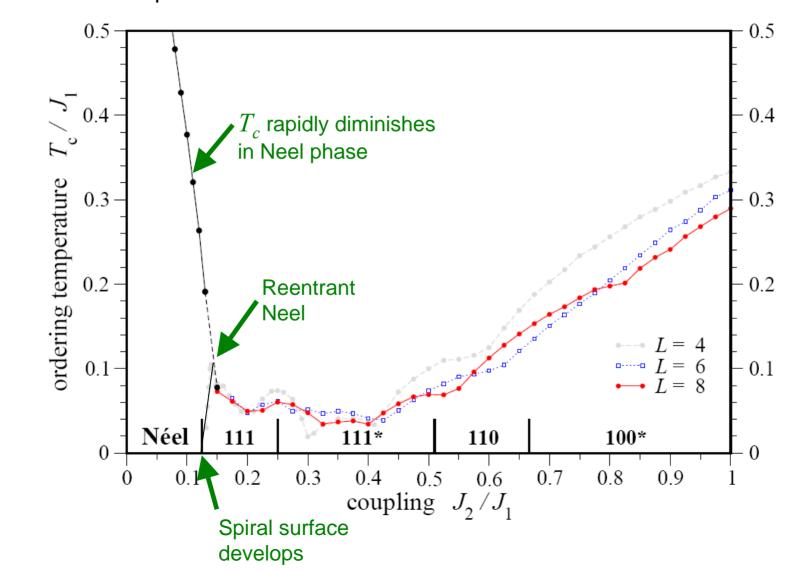
Parallel tempering algorithm employed to dramatically improve thermal equilibration



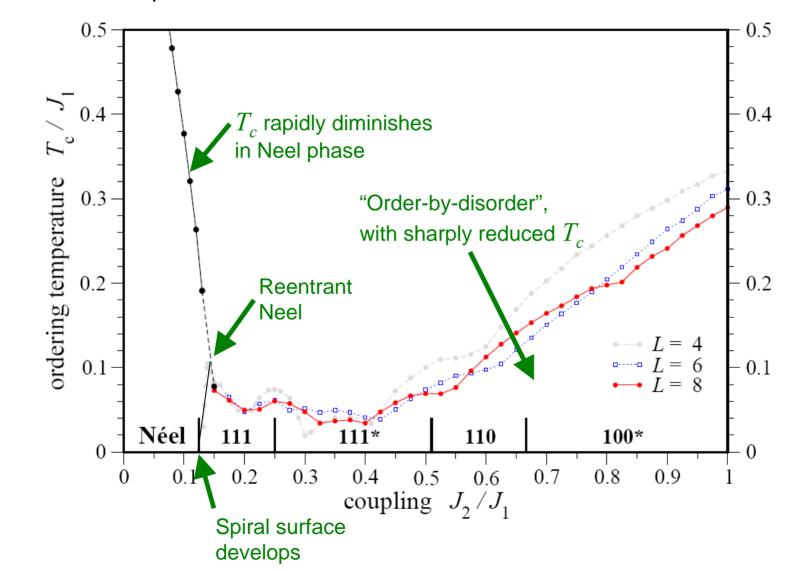
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Parallel tempering algorithm employed to dramatically improve thermal equilibration

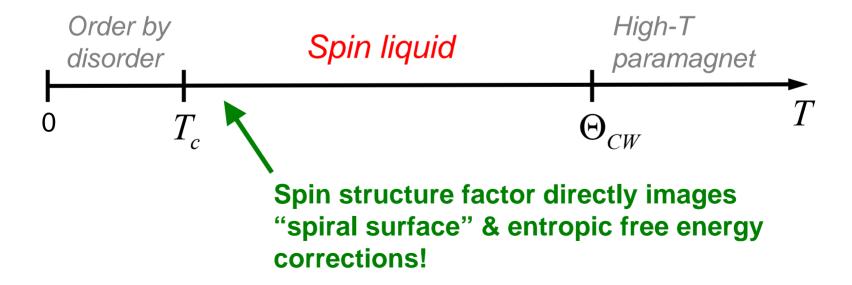


Parallel tempering algorithm employed to dramatically improve thermal equilibration



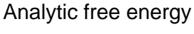
Spin liquid physics

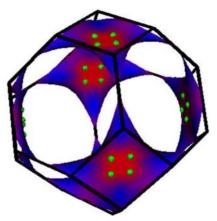
- Order-by-disorder occurs at low temperatures
- Broad spin liquid regime emerges due to low T_c
- Can probe this physics experimentally via neutron scattering



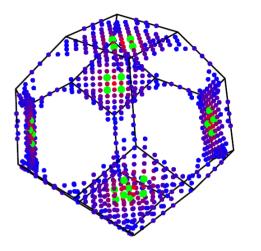
Structure factor in spiral spin liquid regime

$J_2/J_1 = 0.85$ MnSc₂S₄

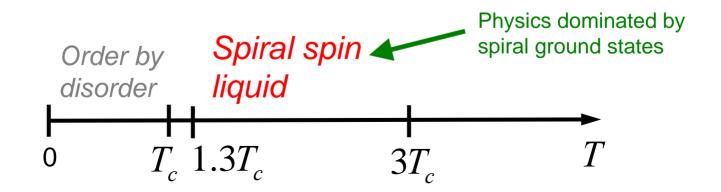




Numerical structure factor



- Free energy corrections visible for $T_c < T < 1.3 T_c$
- "Spiral surface" more robust: persists for $T_c < T < 3 T_c$

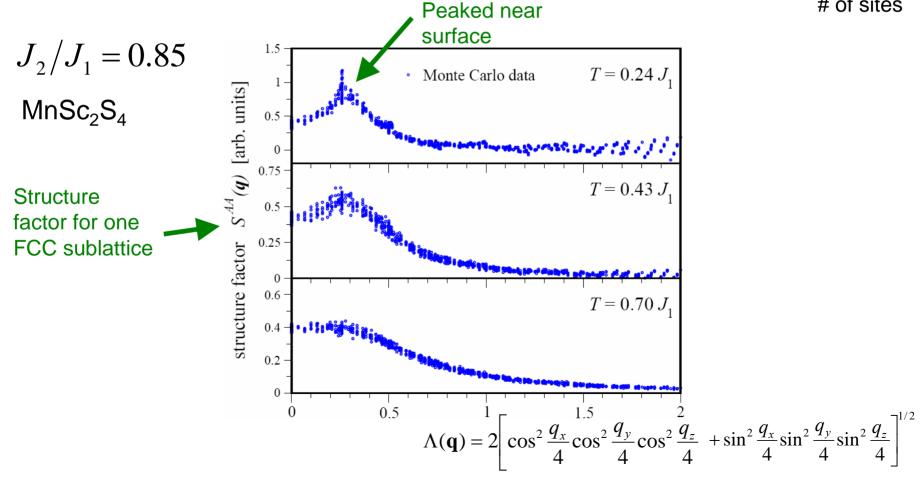


Spin liquid correlations analytically

"Spherical model"

- Describes spin liquids in kagome, pyrochlore antiferromagnets
- Predicts structure factor data collapse

$$\mathbf{S}_{j}^{2} = 1 \rightarrow \sum_{j} \mathbf{S}_{j}^{2} = N$$
of sites

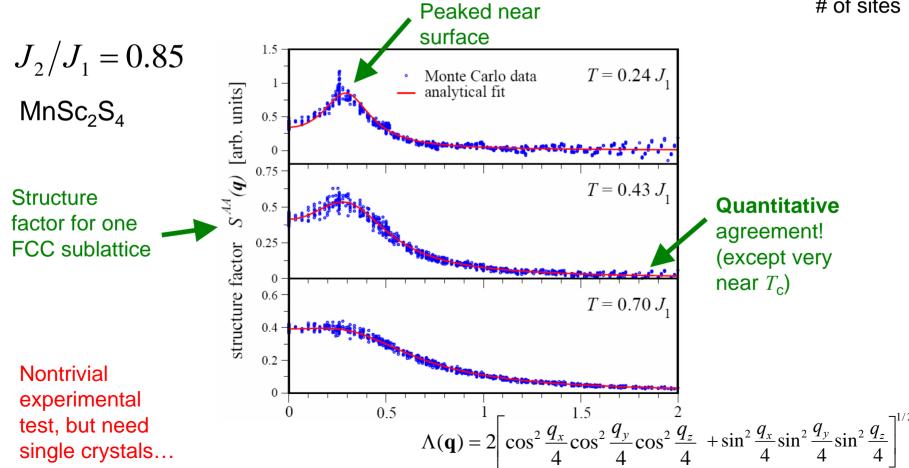


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What can we expect for experiments?

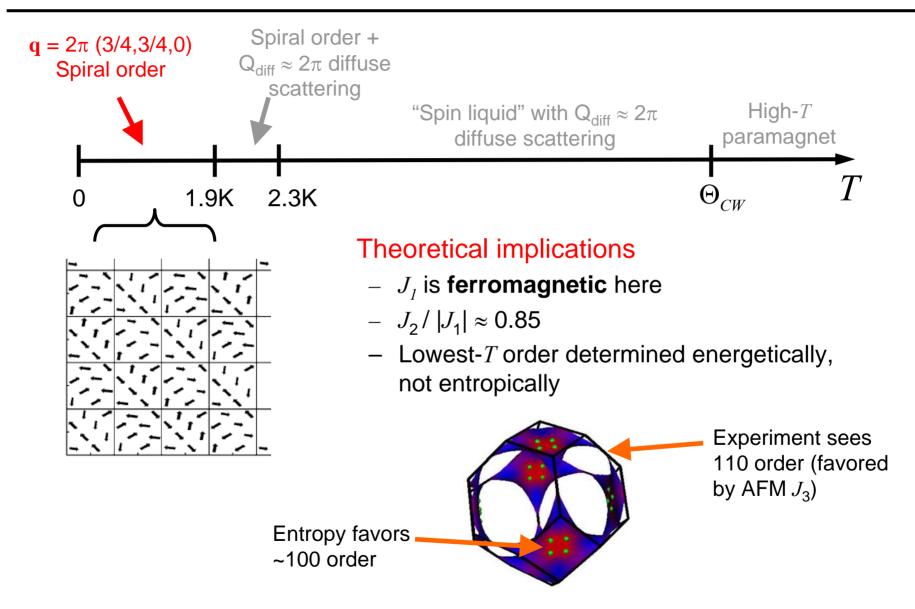
Realistic Hamiltonian:

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle \langle ij \rangle \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + \delta H$$

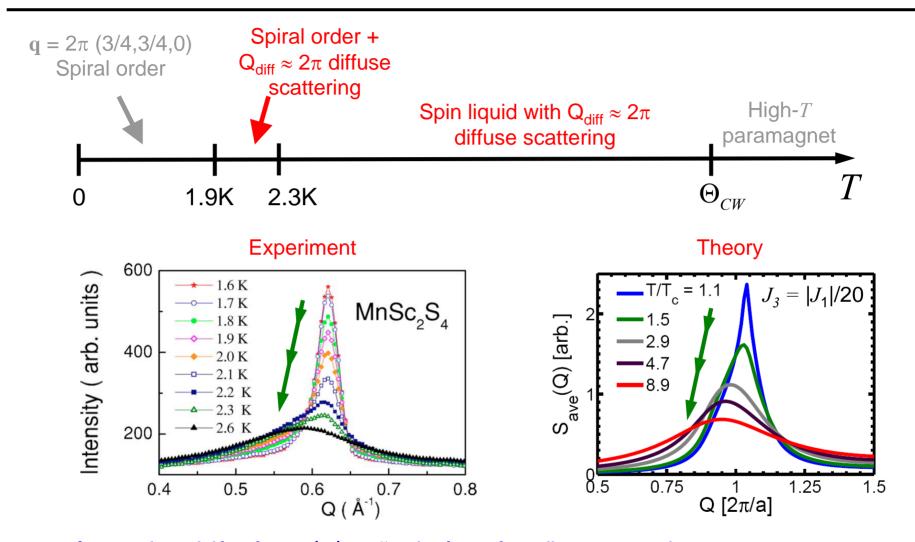
Degeneracybreaking perturbations

- Entropic free energy corrections vanish as $T \rightarrow 0$
- Energetic corrections from δH inevitably dominate at lowest T
- If δH small enough, expect order-by-disorder phase to appear at higher T

Comparison with experiment: MnSc₂S₄



Comparison with experiment: MnSc₂S₄ (cont'd)

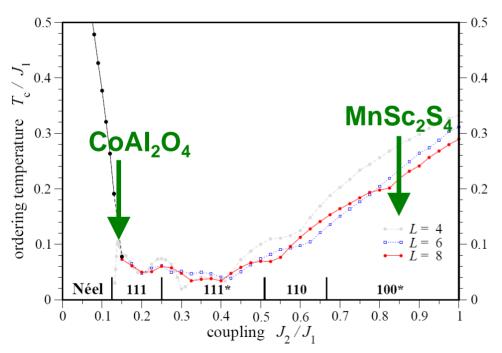


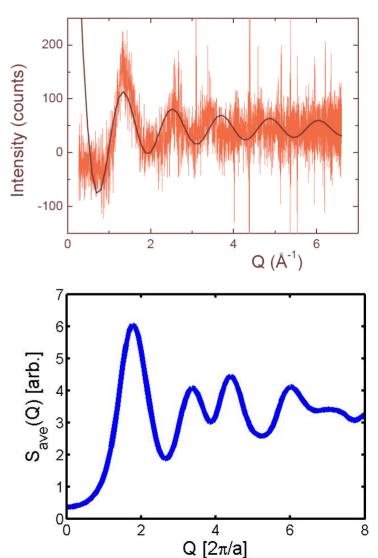
- Intensity shifts from $|\mathbf{q}|$ to "spiral surface" as T washes out J_3
- Consistent with "spiral spin liquid"

A. Krimmel et al. PRB **73**, 014413 (2006); M. Mucksch et al. (unpublished)

Comparison with experiment: CoAl₂O₄

- Much less known here
 - Strong frustration, sample dependent
 - No sharp transition observed yet
- Powder neutron data + frustration suggest $J_2/J_1 \approx 1/8$ for this material





Summary

- Many spinels constitute frustrated diamond lattice antiferromagnets
 - MnSc₂S₄, CoAl₂O₄, etc.
- Simple J_1 - J_2 model captures essential physics
 - Continuous spiral ground state degeneracy
 - Important ordering mechanism is order-by-disorder
 - Spin correlations in "spiral spin liquid" reveals surface + entropic effects
- Theoretical predictions consistent with existing experiments

Future Directions

- Single crystals wanted
 - Allow for more direct comparison
 - Concrete experimental realization of order-by-disorder??
- Explore spin dynamics for inelastic neutron scattering?
- Effects of disorder?
- Details of low-T order in MnSc₂S₄? Commensurate lock-in?
- Physics of spin + orbitally frustrated FeSc₂S₄? Exotic quantum ground state?

Acknowledgments

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