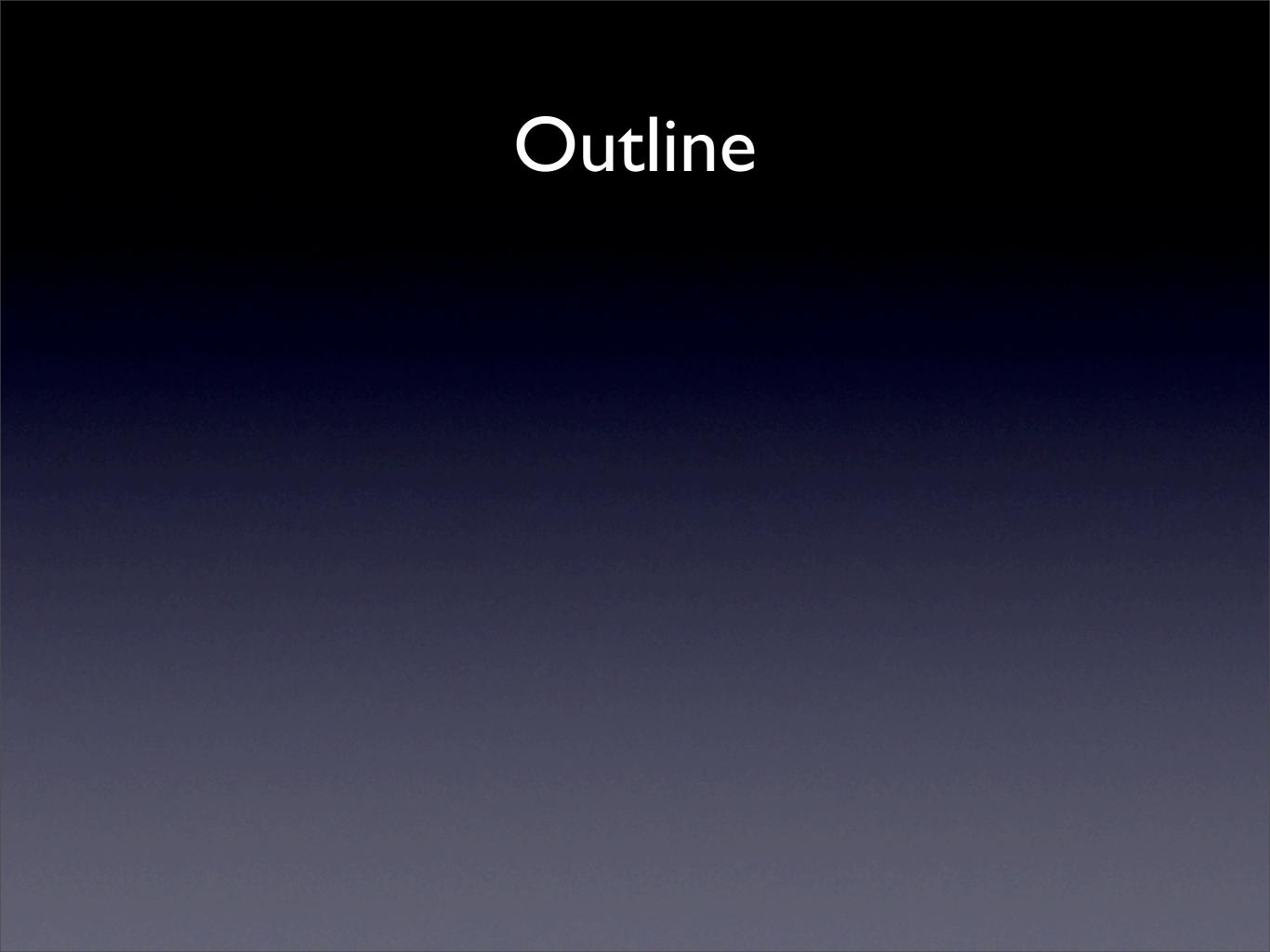
## Orbital and Magnetic Physics in Vanadium Spinels

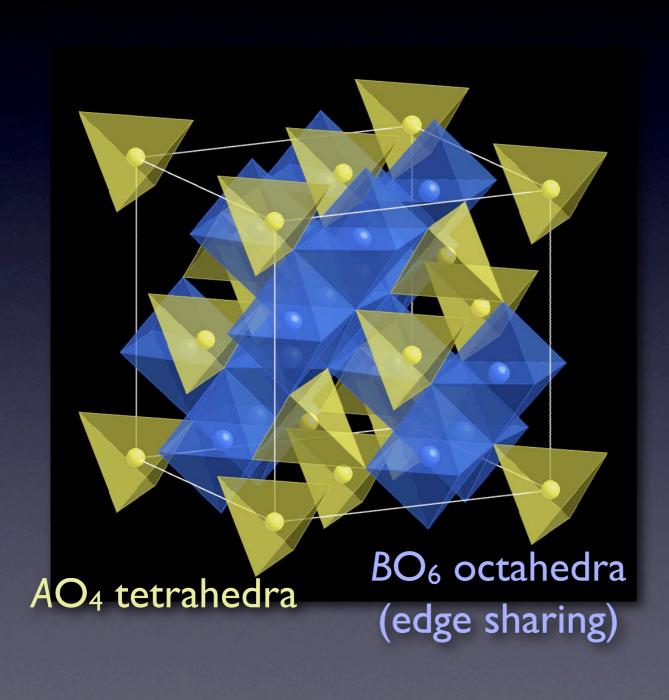
Yukitoshi Motome (Univ. of Tokyo)

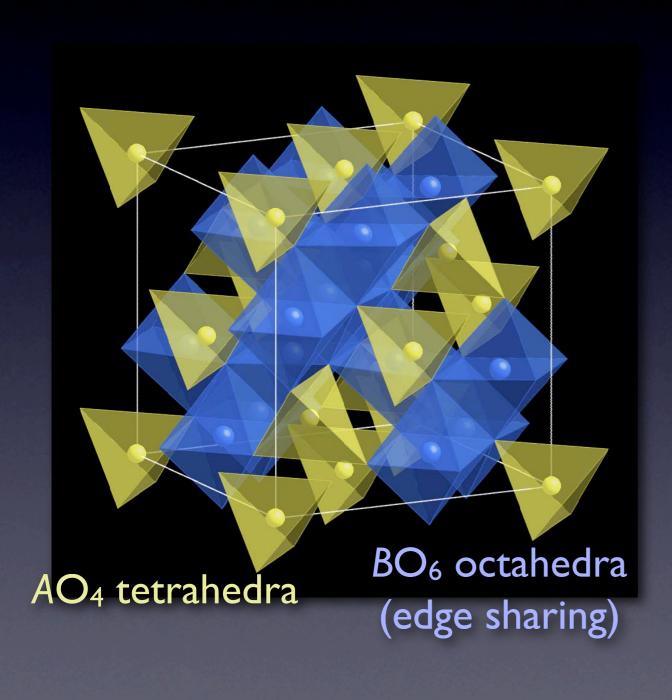
Moments and Multiplets in Mott Materials
Sep. 25, 2007



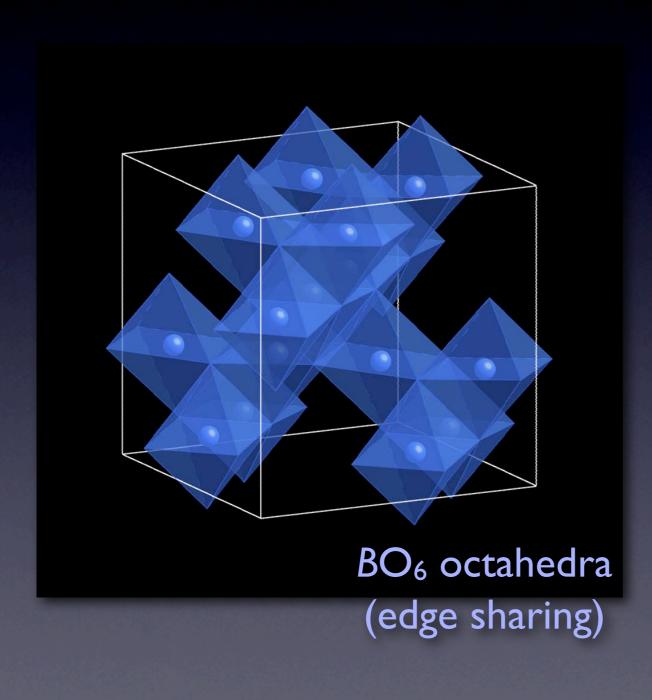
#### Outline

- introduction to spinels and t<sub>2g</sub> orbital physics
- controversy on orbital ordering in ZnV2O4
  - different models for spin/orbital order in ZnV<sub>2</sub>O<sub>4</sub>: relative importance of Kugel-Khomskii superexchange, Jahn-Teller and relativistic spin-orbit couplings
  - symmetry analysis: lesson from experiments in MnV<sub>2</sub>O<sub>4</sub>
- self-organized 7-site cluster (heptamer) in AIV<sub>2</sub>O<sub>4</sub>
  - heptamer scenario: 'molecule' of bonding states with anisotropic t<sub>2g</sub> orbitals
  - implication to heavy-fermion compound LiV<sub>2</sub>O<sub>4</sub>

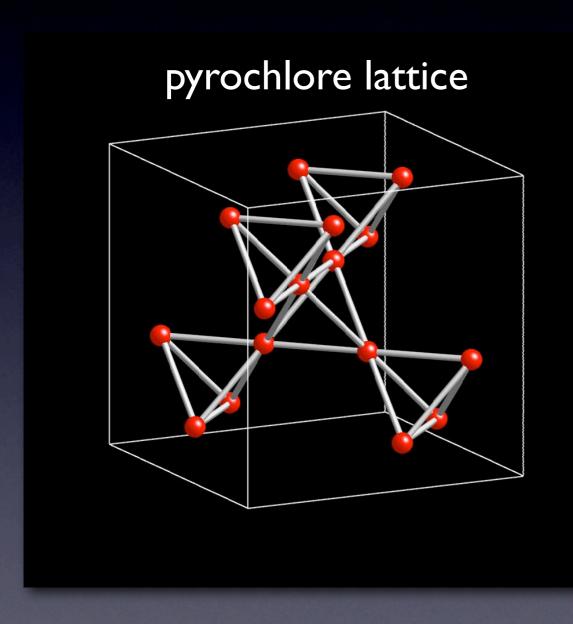




B spinels: A-site cations are nonmagnetic



- B spinels: A-site cations are nonmagnetic
- 3D network of edgesharing BO<sub>6</sub> octahedra



- B spinels: A-site cations are nonmagnetic
- 3D network of edgesharing BO<sub>6</sub> octahedra
- 3D network of cornersharing B<sub>4</sub> tetrahedra → pyrochlore lattice: strong geometrical frustration

### B Spinels with t2g Electrons

d <sup>1</sup> MgTi <sub>2</sub> O <sub>4</sub>	d <sup>2</sup> AV <sub>2</sub> O <sub>4</sub> (A=Zn,Mg)	d <sup>3</sup> ACr <sub>2</sub> O <sub>4</sub> (A=Cd,Hg,Zn)
<ul> <li>metal-insulator transition</li> <li>spin-singlet ground state</li> <li>helical dimerization</li> <li>orbital-Peierls scenario</li> </ul>	<ul> <li>two successive transitions</li> <li>complicated AF ordering</li> <li>dimensionality reduction</li> <li>competition between spin and orbital degrees of freedom</li> </ul>	<ul> <li>single transition</li> <li>half-magnetization plateau</li> <li>spin-lattice coupling (spin Jahn-Teller mechanism)</li> <li>self-organized 'hexamer' in high-T para phase</li> </ul>
d <sup>0.5</sup> LiTi <sub>2</sub> O <sub>4</sub>	d <sup>1.5</sup> LiV <sub>2</sub> O <sub>4</sub>	$d^{2.5}$ AIV $_2$ O $_4$
• superconductivity below 12.4 K (BCS mechanism)	<ul> <li>metallic down to 300 mK</li> <li>absence of any transition</li> <li>heavy-fermion behavior</li> <li>metal-insulator transition by applying pressure</li> </ul>	<ul> <li>structural transition with spin-singlet formation</li> <li>self-organized 7-site cluster 'heptamer'?</li> </ul>

### B Spinels with t2g Electrons

d <sup>1</sup> MgTi <sub>2</sub> O <sub>4</sub>	d <sup>2</sup> AV <sub>2</sub> O <sub>4</sub> (A=Zn,Mg)	d <sup>3</sup> ACr <sub>2</sub> O <sub>4</sub> (A=Cd,Hg,Zn)
<ul> <li>metal-insulator transition</li> <li>spin-singlet ground state</li> <li>helical dimerization</li> <li>orbital-Peierls scenario</li> </ul>	<ul> <li>two successive transitions</li> <li>complicated AF ordering</li> <li>dimensionality reduction</li> <li>competition between spin and orbital degrees of freedom</li> </ul>	<ul> <li>single transition</li> <li>half-magnetization plateau</li> <li>spin-lattice coupling (spin Jahn-Teller mechanism)</li> <li>self-organized 'hexamer' in high-T para phase</li> </ul>
d <sup>0.5</sup> LiTi <sub>2</sub> O <sub>4</sub>	d <sup>1.5</sup> LiV <sub>2</sub> O <sub>4</sub>	$d^{2.5}$ AIV $_2$ O $_4$
• superconductivity below 12.4 K (BCS mechanism)	<ul> <li>metallic down to 300 mK</li> <li>absence of any transition</li> <li>heavy-fermion behavior</li> <li>metal-insulator transition by applying pressure</li> </ul>	<ul> <li>structural transition with spin-singlet formation</li> <li>self-organized 7-site cluster 'heptamer'?</li> </ul>



#### Actions



Join this Space



**Recent Changes** 



Manage Space

Search



#### **Navigation**

#### Wiki Home

Seminar Schedule
Experimenter of the week
Participants' Interests
Social activities

#### **Links to KITP sites:**

KITP program page
Conference page
Participant Directory
Participant's photos
Recorded Talks
Transport info
edit navigation

#### guest · Join · Help · Sign In ·





page 🌞

discussion

history

notify me

Three main questions related to the physics of orbital degrees of freedom came to the fore in the discussion of Wednesday (biased view --JvdB).

We know that orbitals can order and that they couple to the lattice, but the questions are:

- 1. is there any material in which the *quantum* character of orbital degrees of freedom become relevant?
- 2. are there any cases where orbital fluctuations, either quantum or classical are relevent?
- 3. does orbital ordering have interesting textures, symmetries and/or excitations?

Also 15 more or less detailed discussion topics came up:

- 1. What is the role of vibronic coupling in cooperative Jahn-Teller systems
- 2. The importance of relativistic spin orbit coupling in eg and t2g systems
- 3. Orbital and frustration: frustration due to orbital degrees of freedom --- orbitals in frustrated lattices
- 4. Relative importance of electron-lattice effects (Jahn Teller) versus electronic effects (superexchange).
- 5. Role of geometry: differences for the situation of 180 degree O-TM-O bonds, 90 degree O-TM-O bonds and edge sharing octahedra
- 6. Reduced dimensionality due to orbitals
- 7. Importance of direct d-d electronic hopping versus d-oxygen-d hopping, especially in t2g spinels
- 8. Orbitals in charge transfer insulators
- 9. Role of orbital degrees of freedom at metal-insulator transitions
- 10. Orbital liquids -- quantum effects
- 11. Orbital waves -- orbitons
- 12. Importance of long-range interactions in short-range orbital (cooperative Jahn Teller) models
- 13. Multiplets en Mottiplets
- 14. Orbital textures, orbital domains and their effect on electronic degrees of freedom
- 15. What happens to orbital order when going to metallic states -- orbital melting



#### **Actions**



Join this Space



**Recent Changes** 



Manage Space

Search



#### **Navigation**

#### Wiki Home

Seminar Schedule Experimenter of the week Participants' Interests Social activities

#### **Links to KITP sites:**

KITP program page Conference page **Participant Directory** Participant's photos Recorded Talks Transport info edit navigation guest · Join · Help · Sign In ·





page 🗼

discussion

history

notify me

Three main questions related to the physics of orbital degrees of freedom came to the fore in the discussion of Wednesday (biased view --JvdB).

We know that orbitals can order and that they couple to the lattice, but the questions are:

- 1. is there any material in which the *quantum* character of orbital degrees of freedom become relevant?
- 2. are there any cases where orbital *fluctuations*, either quantum or classical are relevent?
- 3. does orbital ordering have interesting textures, symmetries and/or excitations?

Also 15 more or less detailed discussion topics came up:

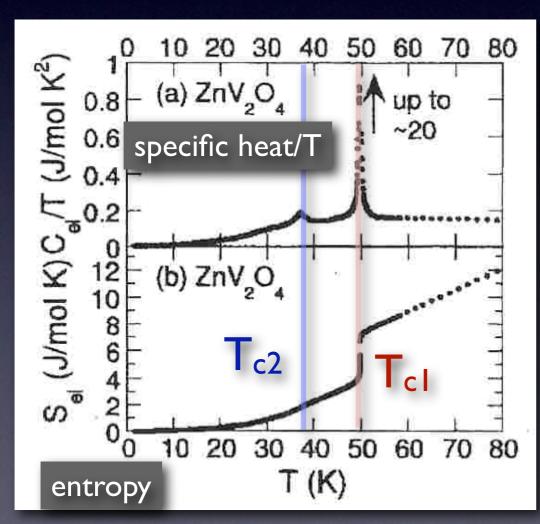
- 1. What is the role of vibronic coupling in cooperative Jahn-Teller systems
- 2. The importance of relativistic spin orbit coupling in eg and t2g systems
- 3. Orbital and frustration: frustration due to orbital degrees of freedom --- orbitals in frustrated lattices
- 4. Relative importance of electron-lattice effects (Jahn Teller) versus electronic effects (superexchange).
- 5. Role of geometry: differences for the situation of 180 degree O-TM-O bonds, 90 degree O-TM-O bonds and edge sharing octahedra
- 6. Reduced dimensionality due to orbitals
- 7. Importance of direct d-d electronic hopping versus d-oxygen-d hopping, especially in t2g spinels
- 8. Orbitals in charge transfer insulators
- 9. Role of orbital degrees of freedom at metal-insulator transitions
- 10. Orbital liquids -- quantum effects
- 11. Orbital waves -- orbitons
- 12. Importance of long-range interactions in short-range orbital (cooperative Jahn Teller) models
- 13. Multiplets en Mottiplets
- 14. Orbital textures, orbital domains and their effect on electronic degrees of freedom
- 15. What happens to orbital order when going to metallic states -- orbital melting

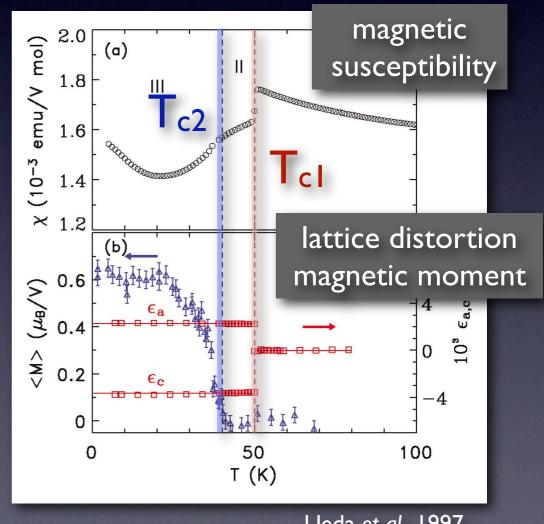
# Two Transitions and Controversy on Orbital Ordering in ZnV<sub>2</sub>O<sub>4</sub>

in collaboration with Hirokazu Tsunetsugu

#### Two Transitions in ZnV<sub>2</sub>O<sub>4</sub>

- cubic to tetragonal transition at  $T_{cl}$  ~50K (1st order)
- antiferromagnetic transition at  $T_{c2}$ ~40K (2nd order)



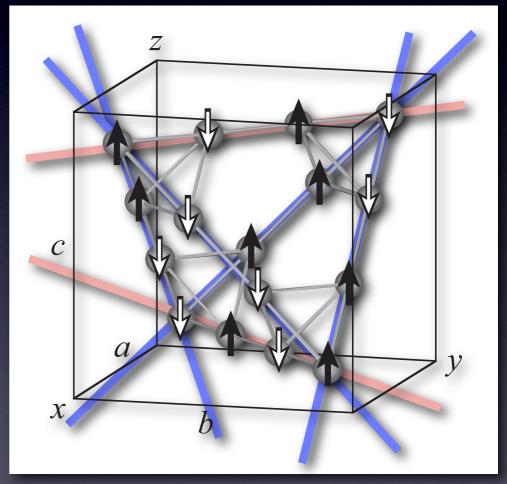


Kondo et al., 2000

Ueda et al., 1997 Lee et al., 2004

## Lattice symmetry and Magnetic Order

- lattice symmetry: I4<sub>1</sub>/amd (powder sample)
- orbital order: undetermined
- spin order: antiferromagnetic
   †-↓-†-↓-... in the xy chains
   †-†-↓-↓-... in the yz/zx chains
- $\bullet$  moment at T=0  $\sim$  0.6 $\mu$ B

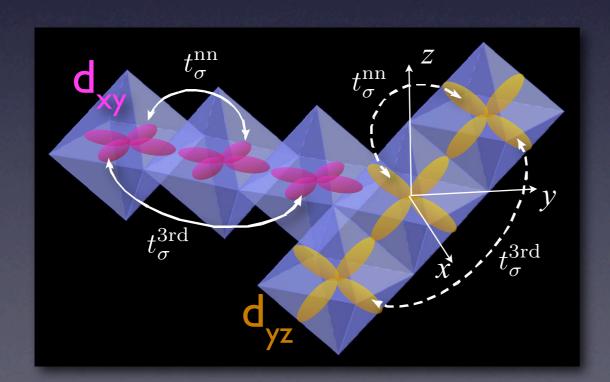


#### Questions

- What is the microscopic mechanism of two transitions? Who is the main player? Kugel-Khomskii superexchanges, Jahn-Teller or relativistic spin-orbit coupling?
- How is the complex AF ordering stabilized? Why is the moment at T=0 reduced so largely?
- What is the role of orbital degree of freedom? Is there orbital ordering? If yes, what type of ordering sets in?

Tsunetsugu and Motome (2003, 2004, 2005)

- Kugel-Khomskii type model derived from 3-fold multi-orbital
   Hubbard model + tetragonal Jahn-Teller coupling
  - assumptions: σ-type transfer integrals only, classical phonon, neglecting spin-orbit coupling and trigonal distortion



$$t_{\sigma}^{\text{nn}} = \sim -0.32 \text{eV}$$
 $t_{\sigma}^{\text{3rd}} = \sim -0.045 \text{eV}$ 
(Matsuno et al., 1999: for LiV<sub>2</sub>O<sub>4</sub>)

Tsunetsugu and Motome (2003, 2004, 2005)

$$H_{\mathrm{SO}}^{\mathrm{nn}} = -J \sum_{\langle ij \rangle} \left[ h_{\mathrm{o-AF}}^{(ij)} + h_{\mathrm{o-F}}^{(ij)} \right] \text{ : nearest neighbor term } J = (t_{\sigma}^{\mathrm{nn}})^2/U \quad A = (1-\eta)/(1-3\eta) \\ J_3 = (t_{\sigma}^{\mathrm{3rd}})^2/U \quad B = \eta/(1-3\eta) \\ J_3 = (t_{\sigma}^{\mathrm{3rd}})^2/U \quad B = \eta/(1-3\eta) \\ \bar{n}_{i\alpha} = 1 - n_{i\alpha} \quad C = (1+\eta)/(1+2\eta) \\ \bar{\eta} = J_{\mathrm{H}}/U \\ h_{\mathrm{o-AF}}^{(ij)} = (A+B\vec{S}_i \cdot \vec{S}_j) (n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)}) \\ h_{\mathrm{o-AF}}^{(ij)} = (A+B\vec{S}_i \cdot \vec{S}_j) (n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)})$$

 $h_{\alpha}^{(ij)} = C(1 - \vec{S}_i \cdot \vec{S}_j) n_{i\alpha(ij)} n_{j\alpha(ij)}$ 

Tsunetsugu and Motome (2003, 2004, 2005)

$$H_{
m SO}^{
m nn} = -J\sum_{\langle ij
angle} \left[h_{
m o-AF}^{(ij)} + h_{
m o-F}^{(ij)}
ight] \; :$$
 nearest neighbor term  $J = (t_{\sigma}^{
m nn})^2/U \quad A = (1-\eta)/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nn})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m nrd})^2/U \quad J_3 = (t_{\sigma}^{$ 

$$h_{\mathrm{o-AF}}^{(ij)} = (A + B\vec{S}_i \cdot \vec{S}_j)(n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)}) \quad \text{spin F / orbital AF}$$

$$h_{\text{o-F}}^{(ij)} = C(1 - \vec{S}_i \cdot \vec{S}_j) n_{i\alpha(ij)} n_{j\alpha(ij)}$$

Tsunetsugu and Motome (2003, 2004, 2005)

$$H_{
m SO}^{
m nn} = -J\sum_{\langle ij
angle} \left[h_{
m o-AF}^{(ij)} + h_{
m o-F}^{(ij)}
ight] \; :$$
 nearest neighbor term  $J = (t_{\sigma}^{
m nn})^2/U \quad A = (1-\eta)/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad B = \eta/(1-3\eta) \quad J_3 = (t_{\sigma}^{
m 3rd})^2/U \quad J_3 = (t_{\sigma}^$ 

$$h_{\mathrm{o-AF}}^{(ij)} = (A + B\vec{S}_i \cdot \vec{S}_j)(n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)}) \quad \text{spin F / orbital AF}$$

$$h_{\text{o-F}}^{(ij)} = C(1 - \vec{S}_i \cdot \vec{S}_j) n_{i\alpha(ij)} n_{j\alpha(ij)}$$

spin AF / orbital F

Tsunetsugu and Motome (2003, 2004, 2005)

$$H_{\rm SO}^{\rm nn} = -J \sum_{\langle ij \rangle} \left[ h_{\rm o-AF}^{(ij)} + h_{\rm o-F}^{(ij)} \right] \; : \text{nearest neighbor term} \qquad \begin{array}{c} J = (t_{\sigma}^{\rm nn})^2/U \quad A = (1-\eta)/(1-3\eta) \\ J_3 = (t_{\sigma}^{\rm 3rd})^2/U \quad B = \eta/(1-3\eta) \\ J_3 = (t_{\sigma}^{\rm 3rd})^2/U \quad B = \eta/(1-3\eta) \\ \hline n_{i\alpha} = 1 - n_{i\alpha} \quad C = (1+\eta)/(1+2\eta) \\ \hline \eta = J_{\rm H}/U \end{array}$$

$$h_{\mathrm{o-AF}}^{(ij)} = (A + B\vec{S}_i \cdot \vec{S}_j)(n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)}) \quad \text{spin F / orbital AF}$$

$$h_{\text{o-F}}^{(ij)} = C(1 - \vec{S}_i \cdot \vec{S}_j) n_{i\alpha(ij)} n_{j\alpha(ij)}$$

spin AF / orbital F

spin part: Heisenberg / orbital part: 3-state Potts

Tsunetsugu and Motome (2003, 2004, 2005)

$$H_{\rm SO}^{\rm nn} = -J\sum_{\langle ij\rangle} \left[h_{\rm o-AF}^{(ij)} + h_{\rm o-F}^{(ij)}\right] \text{ : nearest neighbor term} \qquad J = (t_{\sigma}^{\rm nn})^2/U \quad A = (1-\eta)/(1-3\eta) \\ J_3 = (t_{\sigma}^{\rm 3rd})^2/U \quad B = \eta/(1-3\eta) \\ H_{\rm SO}^{\rm 3rd} = -J_3\sum_{\langle\langle ij\rangle\rangle} \left[h_{\rm o-AF}^{(ij)} + h_{\rm o-F}^{(ij)}\right] \quad \text{: 3rd neighbor term} \qquad \begin{matrix} J = (t_{\sigma}^{\rm nn})^2/U \quad A = (1-\eta)/(1-3\eta) \\ J_3 = (t_{\sigma}^{\rm 3rd})^2/U \quad B = \eta/(1-3\eta) \\ \bar{n}_{i\alpha} = 1 - n_{i\alpha} \quad C = (1+\eta)/(1+2\eta) \\ \eta = J_{\rm H}/U \end{matrix}$$

$$h_{\mathrm{o-AF}}^{(ij)} = (A + B\vec{S}_i \cdot \vec{S}_j)(n_{i\alpha(ij)}\bar{n}_{j\alpha(ij)} + \bar{n}_{i\alpha(ij)}n_{j\alpha(ij)}) \quad \text{spin F / orbital AF}$$

$$h_{\text{o-F}}^{(ij)} = C(1 - \vec{S}_i \cdot \vec{S}_j) n_{i\alpha(ij)} n_{j\alpha(ij)}$$

spin AF / orbital F

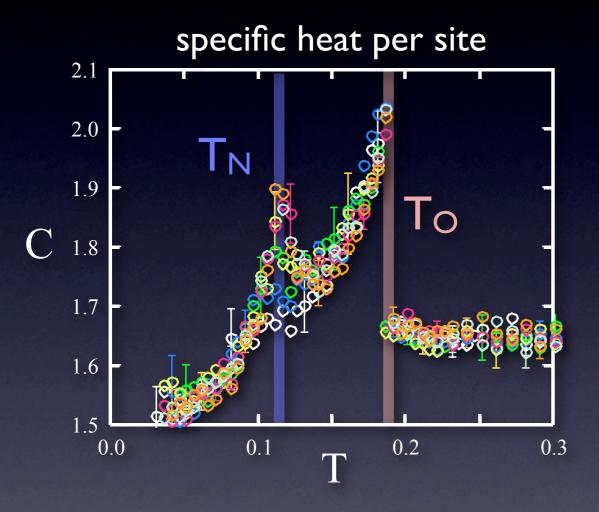
spin part: Heisenberg / orbital part: 3-state Potts

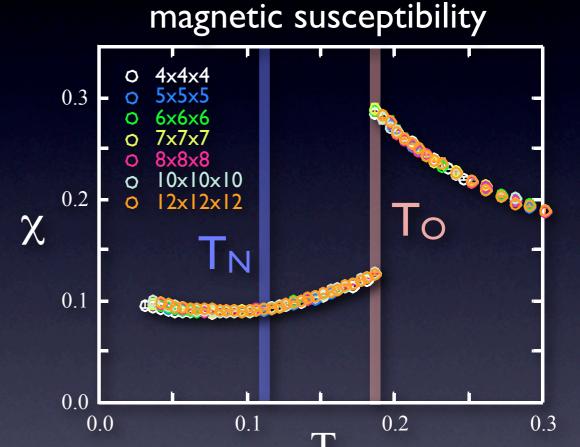
#### + tetragonal Jahn-Teller coupling

+ tetragonal Jahn- Ieller coupling
$$H_{\text{JT}} = g \sum_{i} Q_{i}(n_{i,yz} + n_{i,zx} - 2n_{i,xy}) + \sum_{i} Q_{i}^{2}/2 - \lambda \sum_{\langle ij \rangle} Q_{i}Q_{j} \qquad xy \longrightarrow \bullet$$

$$\begin{array}{c} yz,zx = \\ xy = \\ \end{array}$$

#### Monte Carlo Results

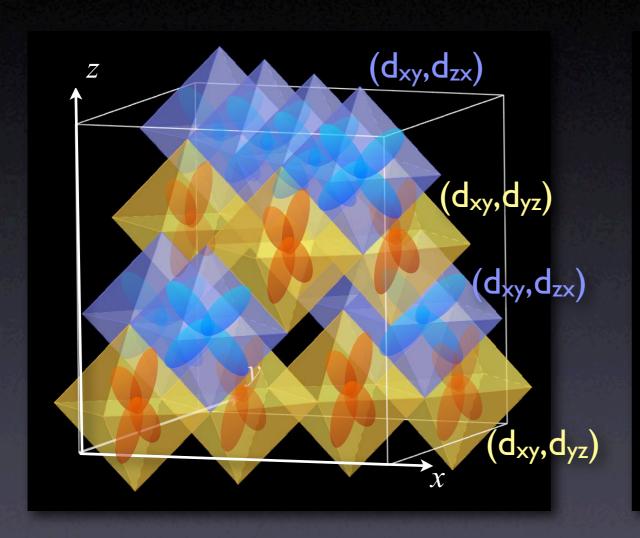


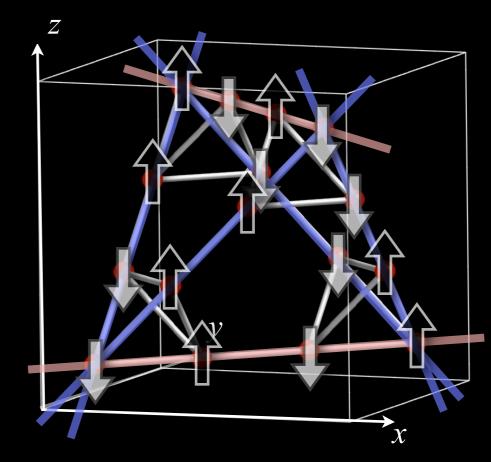


- Ist order at T=T<sub>O</sub>,
   2nd order at T=T<sub>N</sub>
- consistent estimates
   of entropy changes

- sudden drop at T=To
- tiny change at T=T<sub>N</sub>

### Orbital and Spin Structure





- orbital: alternative stacking of  $(d_{xy}, d_{zx})$  and  $(d_{xy}, d_{yz})$  states
- spin:  $\uparrow \downarrow \uparrow \downarrow -$  in the xy chains and  $\uparrow \uparrow \downarrow \downarrow -$  in the yz/zx chains

instability in the high-T (para) phase

instability in the high-T (para) phase

assuming orbital para:  $n_{i\alpha} \rightarrow \langle n_{i\alpha} \rangle = 2/3$ 

instability in the high-T (para) phase

assuming orbital para:  $n_{i\alpha} \rightarrow \langle n_{i\alpha} \rangle = 2/3$ 

$$H_{\mathrm{SO}} \to H_{\mathrm{spin}}^{\mathrm{eff}} = \tilde{J}_{\mathrm{S}} \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + \tilde{J}_{\mathrm{S}}^{(3)} \sum_{\langle \langle ij \rangle \rangle} \vec{S}_i \cdot \vec{S}_j$$

isotropic Heisenberg model with AF nearest- and third-neighbor exchanges

no long-range order at T=0 (Reimers et al., 1991)

instability in the high-T (para) phase

assuming orbital para:  $n_{i\alpha} \rightarrow \langle n_{i\alpha} \rangle = 2/3$ 

$$H_{\mathrm{SO}} \to H_{\mathrm{spin}}^{\mathrm{eff}} = \tilde{J}_{\mathrm{S}} \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + \tilde{J}_{\mathrm{S}}^{(3)} \sum_{\langle \langle ij \rangle \rangle} \vec{S}_i \cdot \vec{S}_j$$

isotropic Heisenberg model with AF nearest- and third-neighbor exchanges

no long-range order at T=0 (Reimers et al., 1991)

spin correlations hardly develop by themselves alone

instability in the high-T (para) phase assuming spin para:

instability in the high-T (para) phase

assuming spin para:  $ec{S}_i \cdot ec{S}_j 
ightarrow \langle ec{S}_i \cdot ec{S}_j 
angle = 0$ 

instability in the high-T (para) phase

assuming spin para:  $\vec{S}_i \cdot \vec{S}_j o \langle \vec{S}_i \cdot \vec{S}_j \rangle = 0$ 

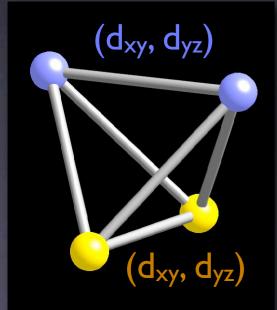
$$H_{\rm SO} \to H_{\rm orbital}^{\rm eff} = \tilde{J}_{\rm O} \sum_{\langle ij \rangle} n_{i\alpha(ij)} n_{j\alpha(ij)} + \tilde{J}_{\rm O}^{(3)} \sum_{\langle\langle ij \rangle\rangle} n_{i\alpha(ij)} n_{j\alpha(ij)}$$

3-state Potts model with AF interactions which depend on both the bond direction and the orbital states

instability in the high-T (para) phase

assuming spin para:  $ec{S}_i \cdot ec{S}_j 
ightarrow \langle ec{S}_i \cdot ec{S}_j 
angle = 0$ 

$$H_{\rm SO} \to H_{\rm orbital}^{\rm eff} = \tilde{J}_{\rm O} \sum_{\langle ij \rangle} n_{i\alpha(ij)} n_{j\alpha(ij)} + \tilde{J}_{\rm O}^{(3)} \sum_{\langle\langle ij \rangle\rangle} n_{i\alpha(ij)} n_{j\alpha(ij)}$$



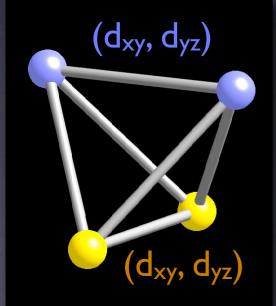
3-state Potts model with AF interactions which depend on both the bond direction and the orbital states

→ (partial) lifting of degeneracy

instability in the high-T (para) phase

assuming spin para:  $ec{S}_i \cdot ec{S}_j 
ightarrow \langle ec{S}_i \cdot ec{S}_j 
angle = 0$ 

$$H_{\rm SO} \to H_{\rm orbital}^{\rm eff} = \tilde{J}_{\rm O} \sum_{\langle ij \rangle} n_{i\alpha(ij)} n_{j\alpha(ij)} + \tilde{J}_{\rm O}^{(3)} \sum_{\langle\langle ij \rangle\rangle} n_{i\alpha(ij)} n_{j\alpha(ij)}$$

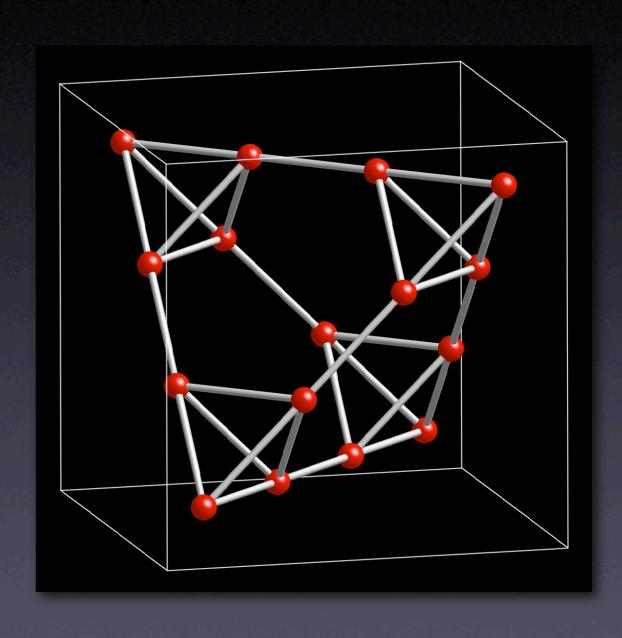


3-state Potts model with AF interactions which depend on both the bond direction and the orbital states

→ (partial) lifting of degeneracy

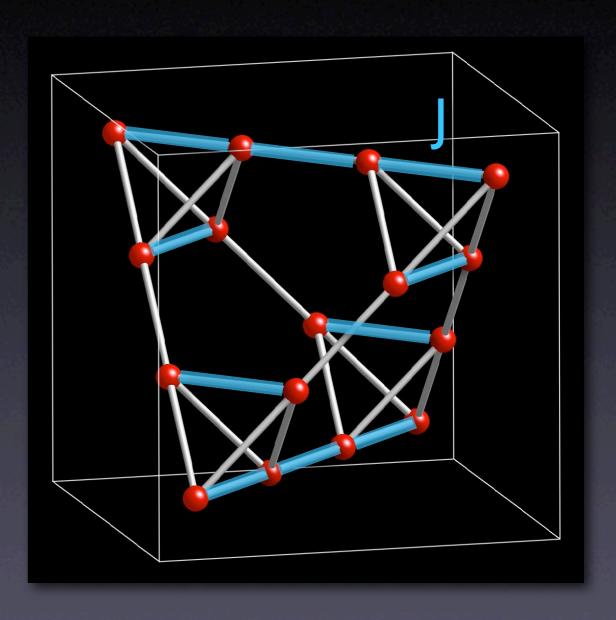
tetragonal Jahn-Teller distortion assists to stabilize this orbital configuration

## Effective Spin Exchanges under the Orbital Order



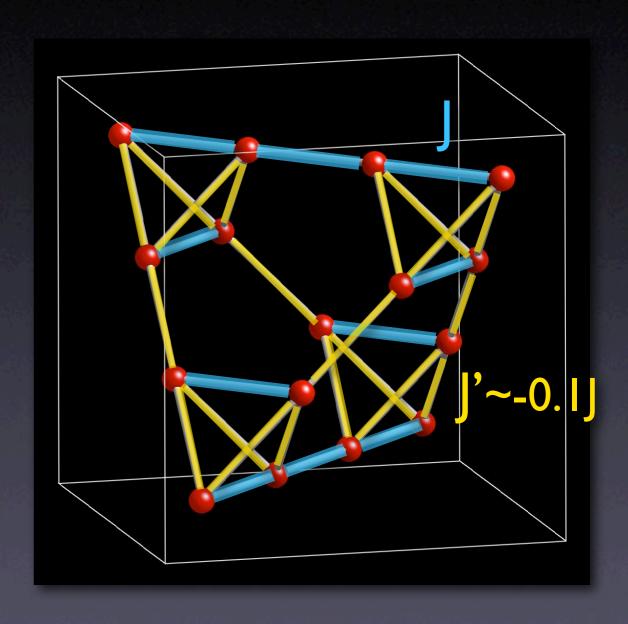
### Effective Spin Exchanges under the Orbital Order

 d<sub>xy</sub> is singly occupied at all the sites → strong AF exchange in the xy chains J



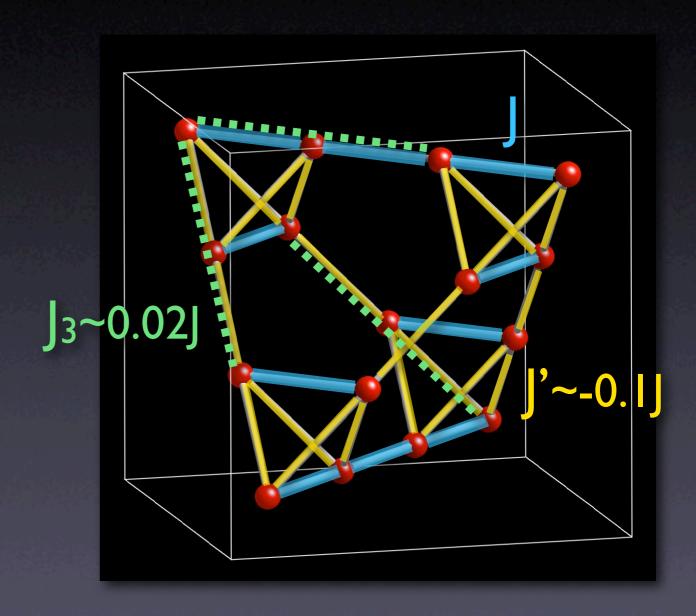
### Effective Spin Exchanges under the Orbital Order

- d<sub>xy</sub> is singly occupied at all the sites → strong AF exchange in the xy chains J
- n.n. exchange couplings in the yz/zx chains J' are ferromagnetic and about IO times weaker than the AF exchange in the xy chains J



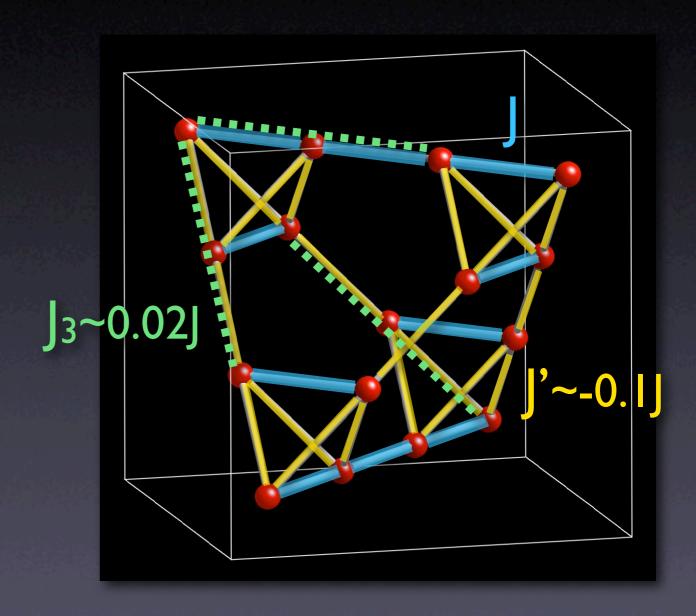
## Effective Spin Exchanges under the Orbital Order

- d<sub>xy</sub> is singly occupied at all the sites → strong AF exchange in the xy chains J
- n.n. exchange couplings in the yz/zx chains J' are ferromagnetic and about 10 times weaker than the AF exchange in the xy chains J
- 3rd-neighbor exchange  $\frac{1}{3}$  is  $\sim 0.02 \longrightarrow AF$  order at  $T_N$



## Effective Spin Exchanges under the Orbital Order

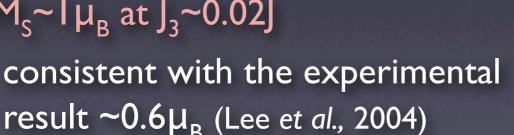
- d<sub>xy</sub> is singly occupied at all the sites → strong AF exchange in the xy chains J
- n.n. exchange couplings in the yz/zx chains J' are ferromagnetic and about 10 times weaker than the AF exchange in the xy chains J
- 3rd-neighbor exchange  $\int_3$  is  $\sim 0.02 \rightarrow AF$  order at  $T_N$

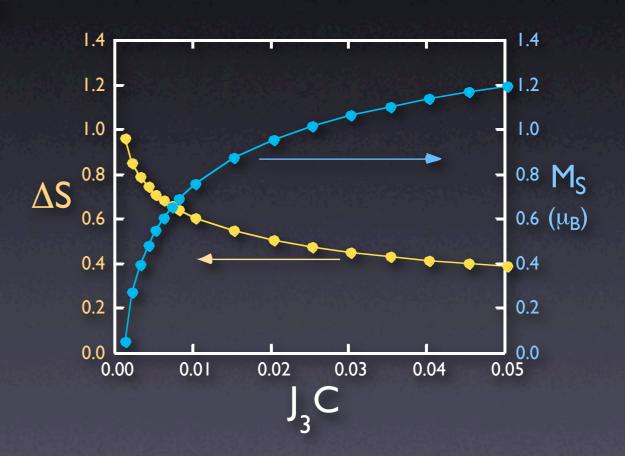


weakly-coupled ID spin chains (dimensionality reduction)

## Quasi-ID Quantum Fluctuation: Large Reduction of AF Moment

- linear spin-wave analysis for the spin and orbital ordered ground state
- moment reduction  $\Delta S$  diverges logarithmically at  $J_3=0$  due to the zero modes
- $\Delta S$  is large in the small  $J_3$  region:  $M_{s} \sim 1 \mu_{B}$  at  $J_{s} \sim 0.02J$ consistent with the experimental





## Short Summary...

### Short Summary...

- Kugel-Khomskii spin-orbital exchange + tetragonal Jahn-Teller
- classical Monte Carlo simulation and mean-field type analysis
- linear spin-wave analysis of effective spin model

### Short Summary...

- Kugel-Khomskii spin-orbital exchange + tetragonal Jahn-Teller
- classical Monte Carlo simulation and mean-field type analysis
- linear spin-wave analysis of effective spin model
- two transitions with reasonable estimates of trasition temperatures as well as entropy changes
- T-dep of magnetic susceptibility consistent with experiment
- magnetic order consistent with the neutron scattering result
- reduced magnetic moment at T=0
- A-type antiferro orbital order with tetragonal distortion

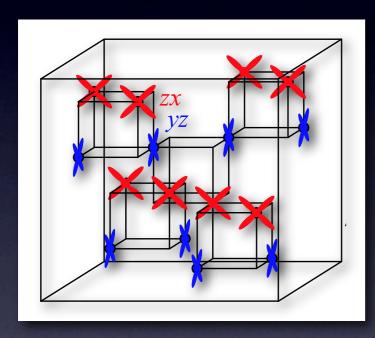
#### Three Different Models

In all models, xy orbital is singly occupied at all the sites (not shown in the figures)

#### Three Different Models

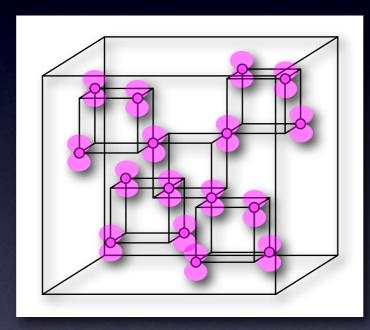
In all models, xy orbital is singly occupied at all the sites (not shown in the figures)

Tsunetsugu-Motome, 2003



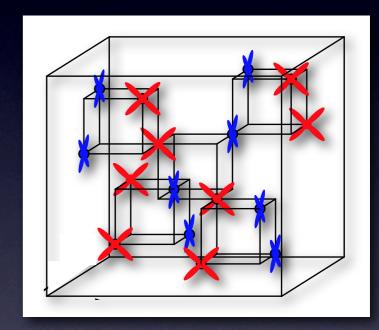
- A-type orbital order
- *I*4<sub>1</sub>/*a*
- spin-orbital superexchanges

Tchernyshyov, 2004



- uniform orbital order
- *I*4<sub>1</sub>/*amd*
- relativistic spin-orbit coupling

Khomskii-Mizokawa, 2005

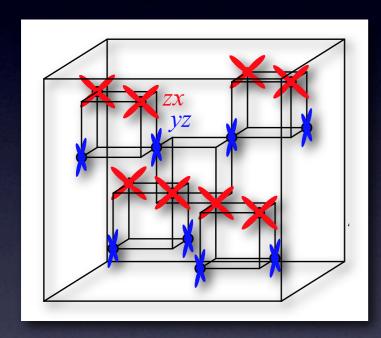


- orbitally-drivenPeierls order
- *P*4<sub>1</sub>2<sub>1</sub>2
- approach from itinerant picture (band Jahn-Teller)

#### Three Different Models

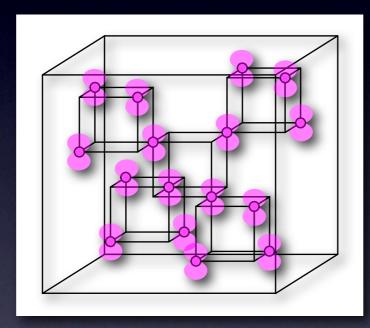
In all models, xy orbital is singly occupied at all the sites (not shown in the figures)

Tsunetsugu-Motome, 2003



- A-type orbital order
- *I*4<sub>1</sub>/*a*
- spin-orbital superexchanges

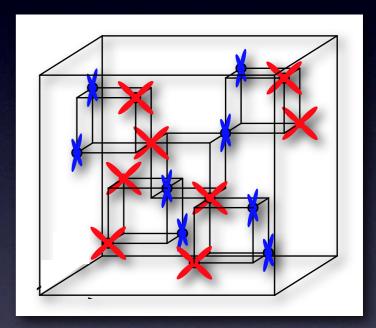
Tchernyshyov, 2004



- uniform orbital order
- *I*4<sub>1</sub>/*amd*
- relativistic spin-orbit coupling

mean-field (Di Matteo et al.)
LSDA+U+SO (T. Maitra and R. Valenti)

Khomskii-Mizokawa, 2005



- orbitally-drivenPeierls order
- *P*4<sub>1</sub>2<sub>1</sub>2
- approach from itinerant picture (band Jahn-Teller)



#### Issue...



- role of relativistic spin-orbit interaction
- orbital ordering at T=0: mean-field analysis and firstprinciple calculation suggest the relevant role
- thermodynamics: single or two transitions? In general, systems with dominant spin-orbit coupling shows a single transition with concomitant ordering of spin and orbital.
- reduced AF moment: due to dimensionality reduction and/or L-S coupling?

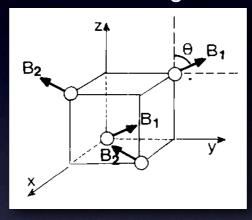
#### Issue...

- ole of relativistic spin-orbit interaction
  - orbital ordering at T=0: mean-field analysis and firstprinciple calculation suggest the relevant role
  - thermodynamics: single or two transitions? In general, systems with dominant spin-orbit coupling shows a single transition with concomitant ordering of spin and orbital.
  - reduced AF moment: due to dimensionality reduction and/or L-S coupling?
- Remark: X-ray diffraction has been done only for powder samples...

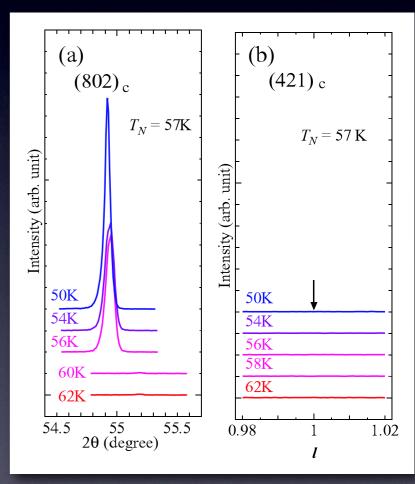
## Lesson from Related Spinel MnV<sub>2</sub>O<sub>4</sub>

- $Mn^{2+} = (3d)^5, V^{3+} = (3d)^2$
- single transition at 57K
  - cubic → tetragonal
  - non-collinear ferri

Plumier and Sougi, 1987



- low-T phase:  $I4_1/a$  (large single crystal)
  - diamond-glide symmetry is broken, but face-center symmetry is hold
  - peak intensity is ~10<sup>-4</sup> times smaller compared to the fundamental peaks, difficult to observe in powder samples



Suzuki et al., 2007

#### Other Issues...

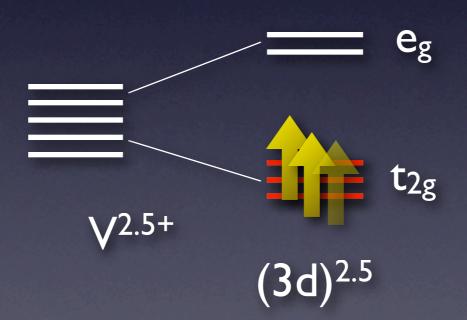
- role of trigonal distortion
  - quantitative difference in Cd compound
- d-d direct vs d-p-d (d-p-p-d) indirect transfers
- orbital and spin ordering in MnV<sub>2</sub>O<sub>4</sub>
- single crystal of ZnV<sub>2</sub>O<sub>4</sub>!

# Self-organized 7-site Cluster (heptamer) in AIV<sub>2</sub>O<sub>4</sub>

in collaboration with Keisuke Matsuda and Nobuo Furukawa

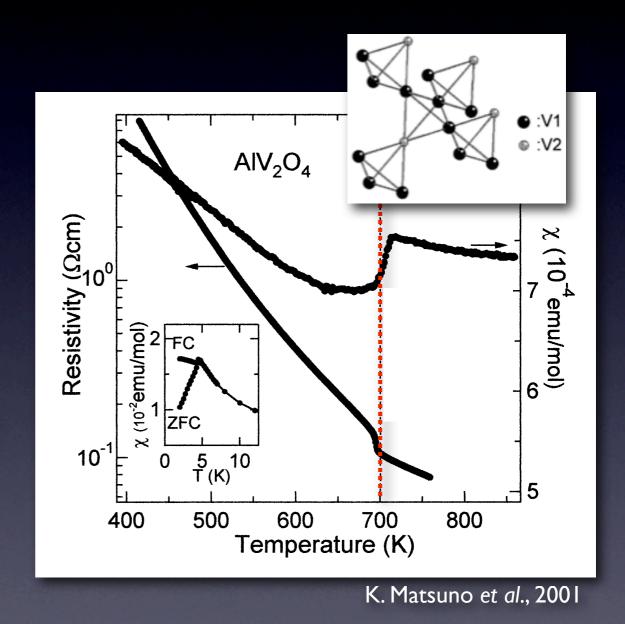
## (Atomic) Electronic Structure in AIV<sub>2</sub>O<sub>4</sub>

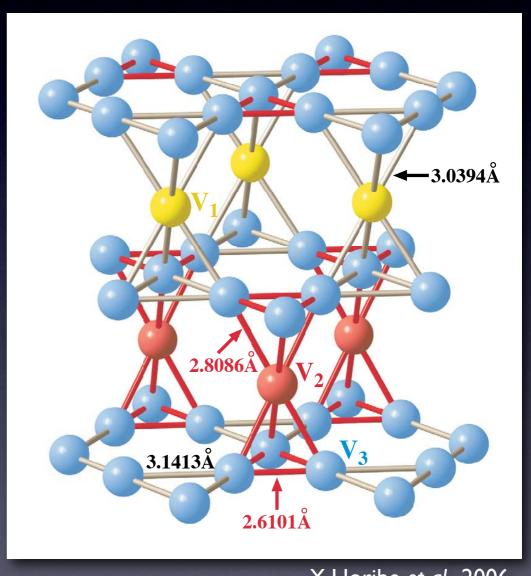
- mixed valence:  $V^{2.5+} = (3d)^{2.5}$
- charge, spin and orbital degrees of freedom are all active



#### Phase Transition at T~700K

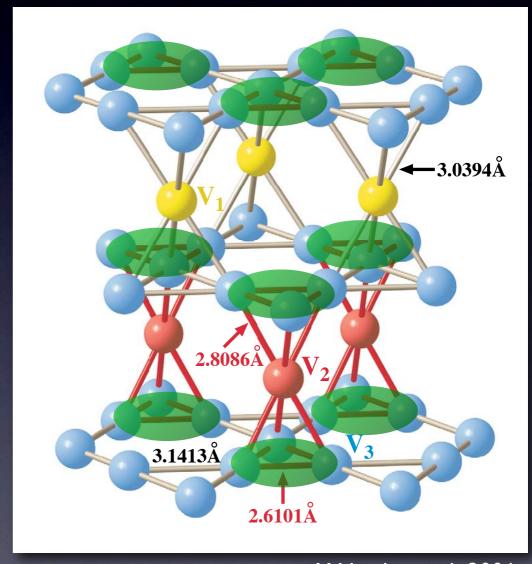
- structural change: doubling of the unit cell along the [111] direction
- shoulder in the resistivity
- sudden drop in the magnetic susceptibility followed by Curie behavior at lower temperatures
- valence-skipping-type charge ordering?





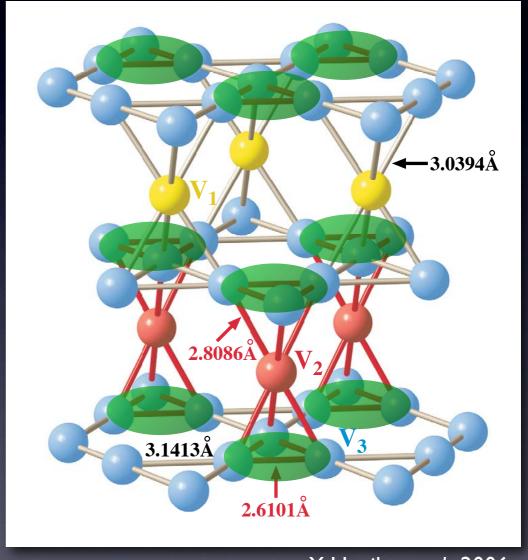
Y. Horibe et al., 2006

new experimental finding: trimer formation in Kagome layers below T<sub>c</sub>



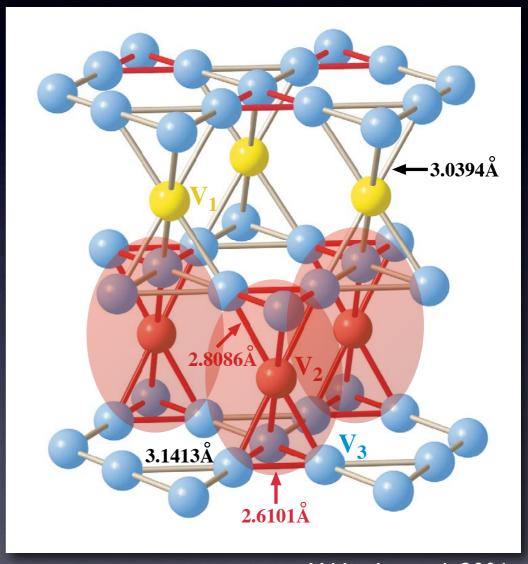
Y. Horibe et al., 2006

- new experimental finding: trimer formation in Kagome layers below T<sub>c</sub>
- spin-singlet formation in trimers? → sharp drop of the magnetic susceptibility?



Y. Horibe et al., 2006

- new experimental finding: trimer formation in Kagome layers below T<sub>c</sub>
- spin-singlet formation in trimers? → sharp drop of the magnetic susceptibility?
- We propose a singlet state emerging from the 7-site clusters (heptamers)

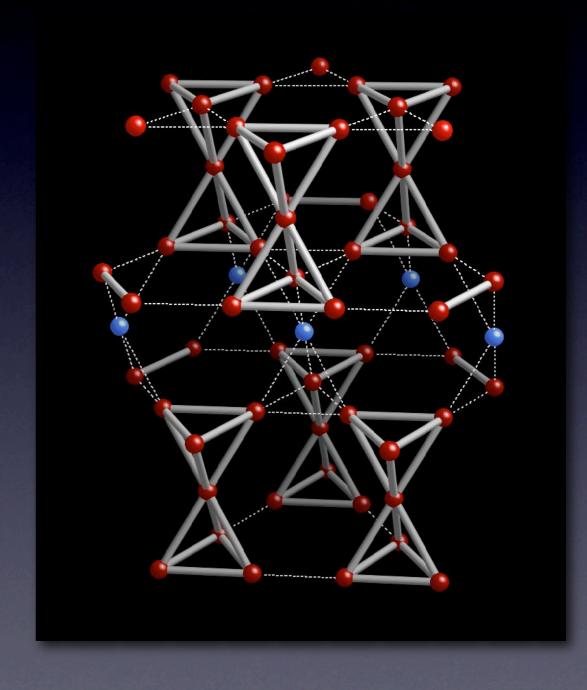


Y. Horibe et al., 2006

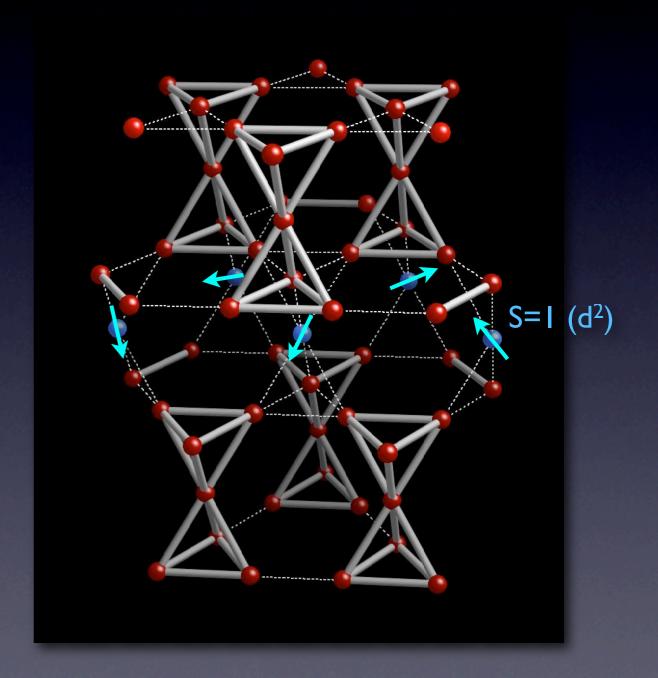
#### Questions

What is the mechanism of the heptamer formation? How is the degeneracy in the frustrated pyrochlore system lifted?

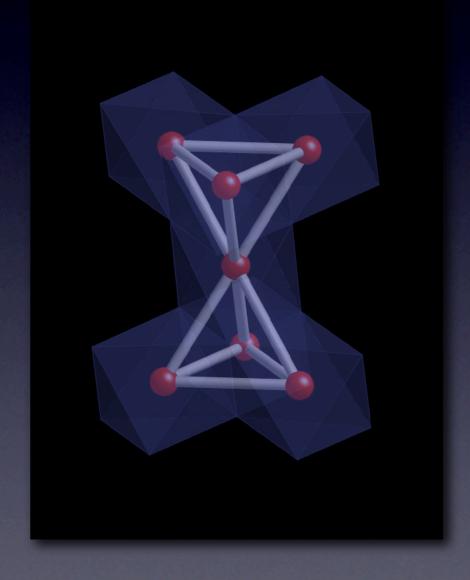
Is the heptamer in a spin-singlet state? How does the singlet state emerge in each heptamer?



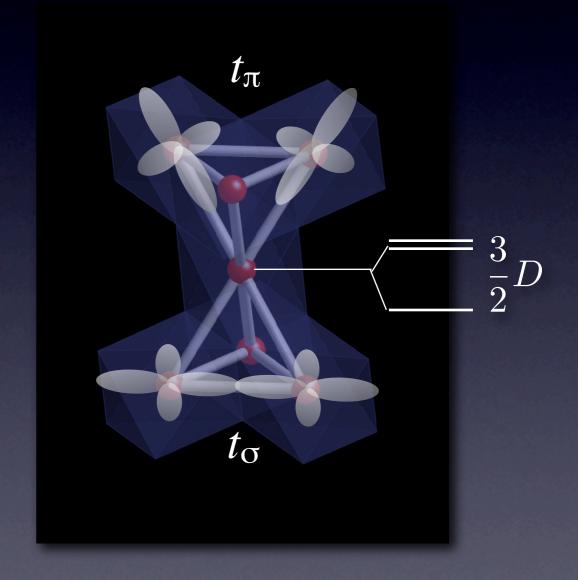
- assumption: S=1 localized moments at isolated V sites (leading Curie behavior at low T)
  - → 18 electrons per heptamer



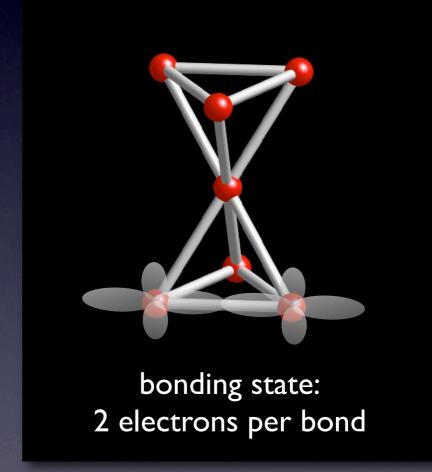
- assumption: S=1 localized moments at isolated V sites (leading Curie behavior at low T)
  - → 18 electrons per heptamer
- t<sub>2g</sub> multi-orbital Hubbard model for each heptamer
  - $\sigma$  and  $\pi$  transfer integrals
  - trigonal lattice distortion at the central site
  - Coulomb interactions



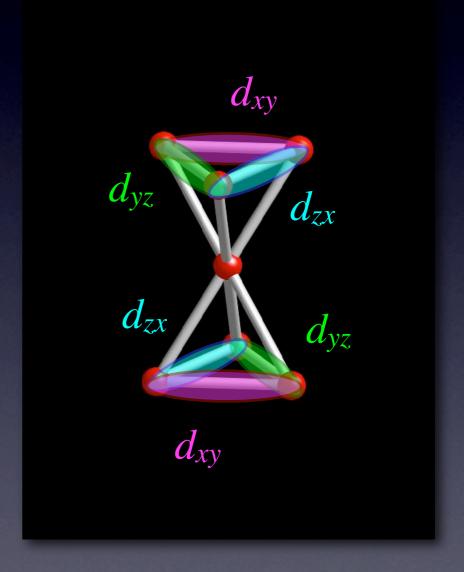
- assumption: S=I localized moments at isolated V sites (leading Curie behavior at low T)
  - → 18 electrons per heptamer
- t<sub>2g</sub> multi-orbital Hubbard model for each heptamer
  - $\sigma$  and  $\pi$  transfer integrals
  - trigonal lattice distortion at the central site
  - Coulomb interactions
- **assumption**: σ-type bonding states for shortest V-V bonds
  - → 6 electrons remaining



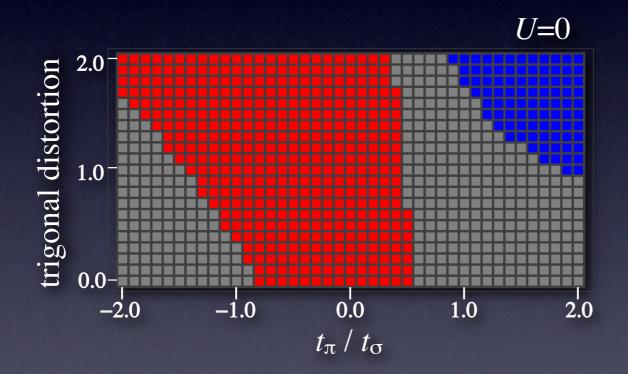
- assumption: S=1 localized moments at isolated V sites (leading Curie behavior at low T)
  - → 18 electrons per heptamer
- t<sub>2g</sub> multi-orbital Hubbard model for each heptamer
  - $\sigma$  and  $\pi$  transfer integrals
  - trigonal lattice distortion at the central site
  - Coulomb interactions
- **assumption**: σ-type bonding states for shortest V-V bonds
  - → 6 electrons remaining



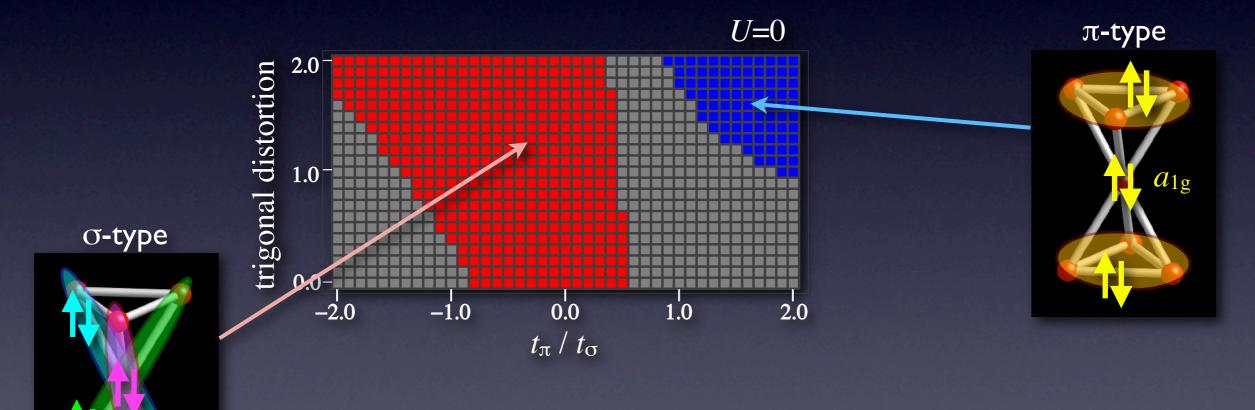
- assumption: S=1 localized moments at isolated V sites (leading Curie behavior at low T)
  - → 18 electrons per heptamer
- t<sub>2g</sub> multi-orbital Hubbard model for each heptamer
  - $\sigma$  and  $\pi$  transfer integrals
  - trigonal lattice distortion at the central site
  - Coulomb interactions
- assumption: σ-type bonding states for shortest V-V bonds
  - → 6 electrons remaining



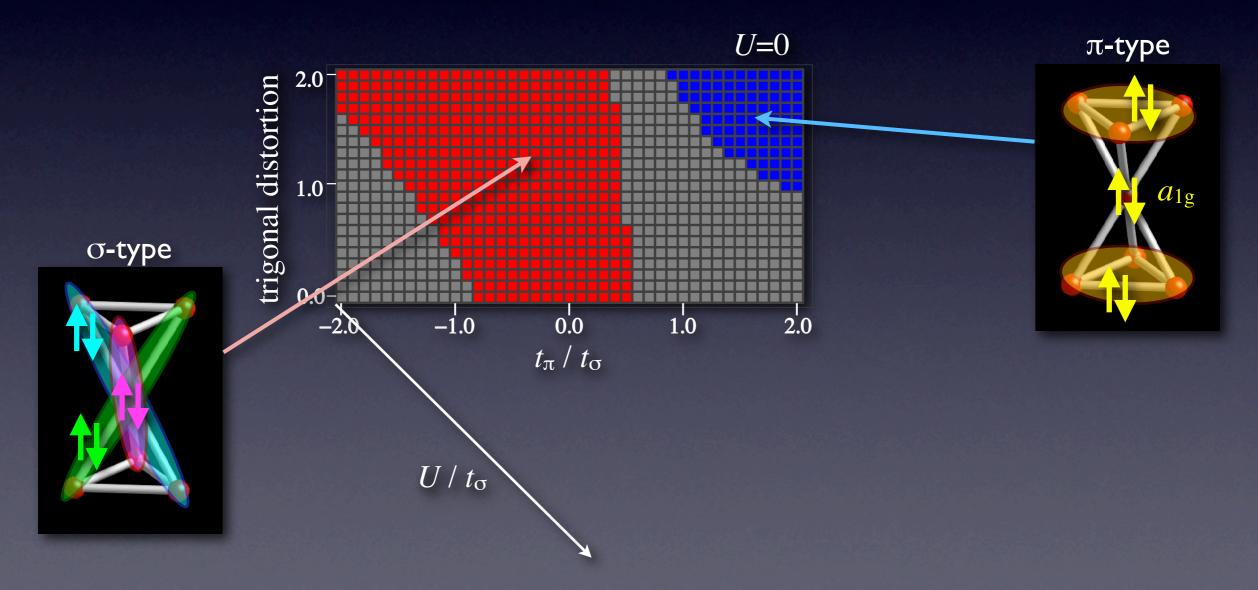
- exact diagonalization of the effective heptamer model
- two different singlet regimes:  $\sigma$ -type and  $\pi$ -type



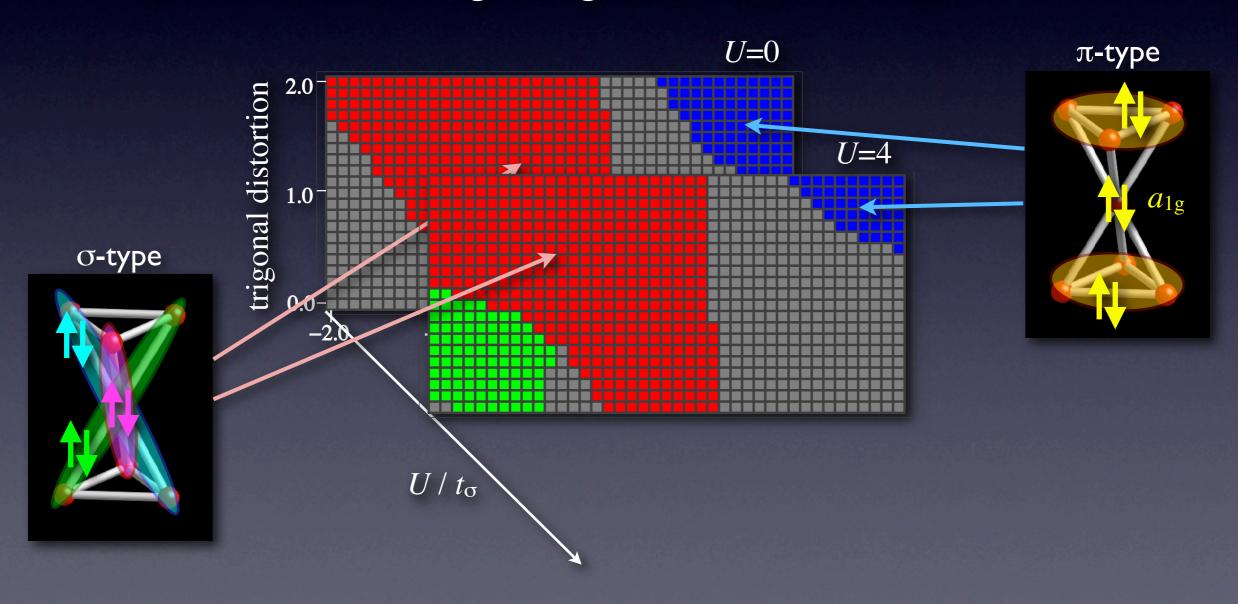
- exact diagonalization of the effective heptamer model
- **e** two different singlet regimes:  $\sigma$ -type and  $\pi$ -type



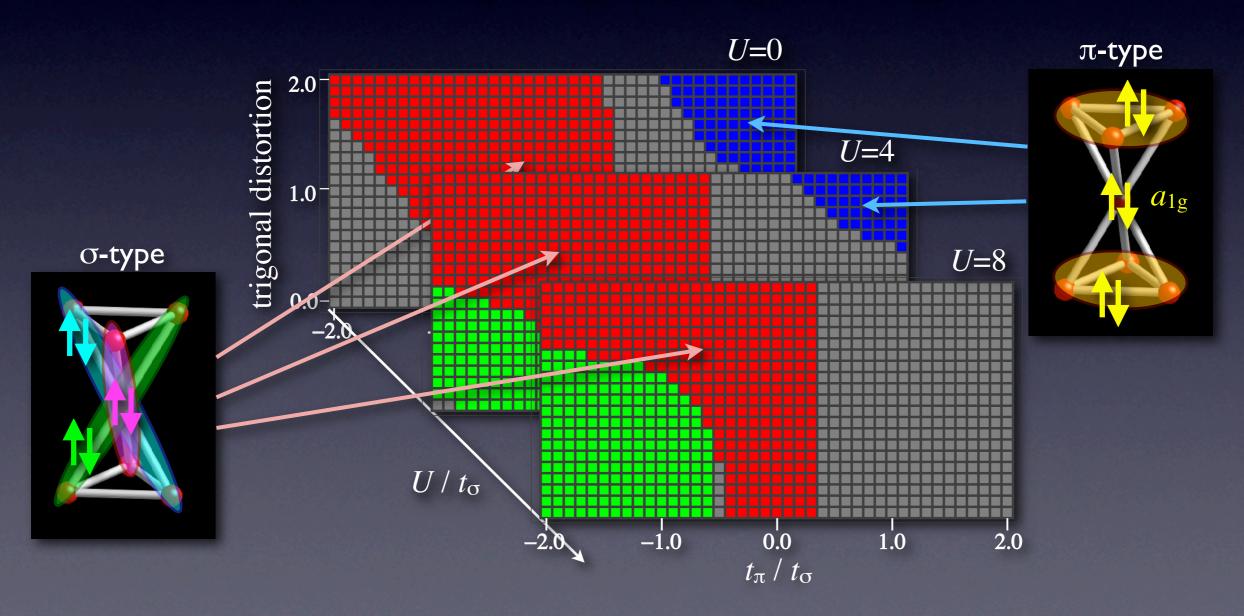
- exact diagonalization of the effective heptamer model
- **e** two different singlet regimes:  $\sigma$ -type and  $\pi$ -type



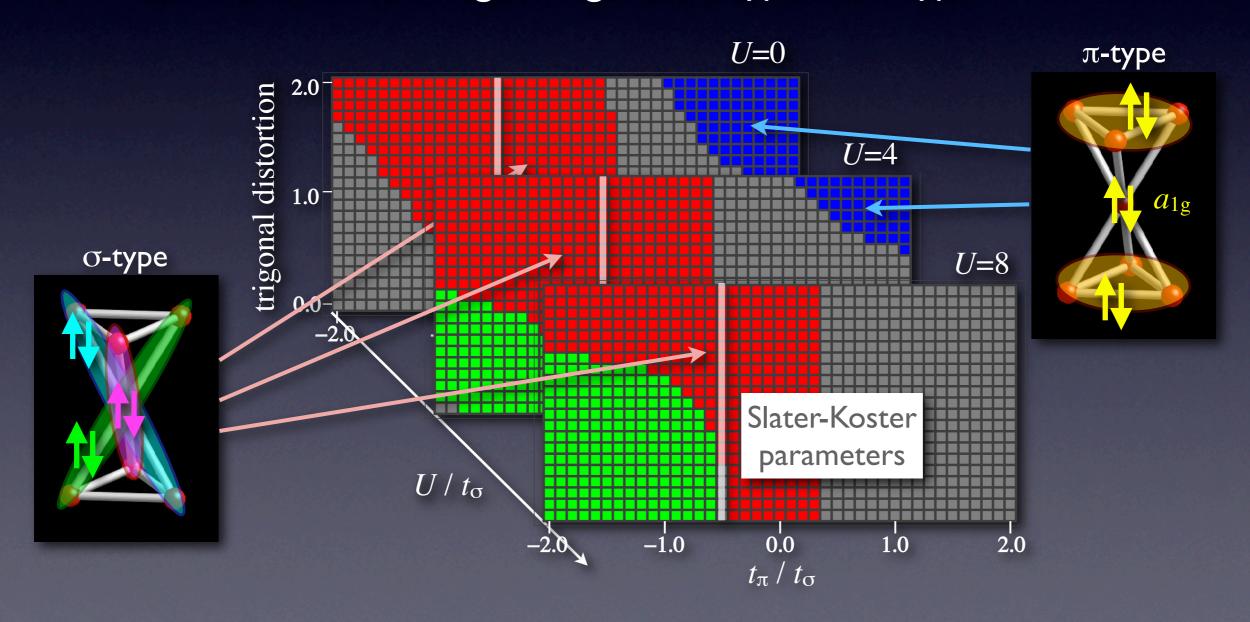
- exact diagonalization of the effective heptamer model
- **e** two different singlet regimes:  $\sigma$ -type and  $\pi$ -type



- exact diagonalization of the effective heptamer model
- **e** two different singlet regimes:  $\sigma$ -type and  $\pi$ -type

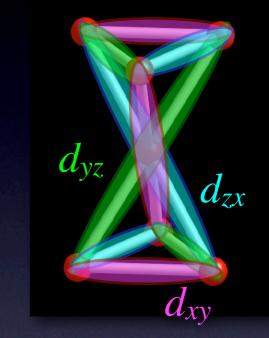


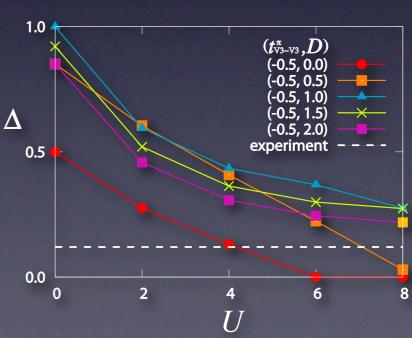
- exact diagonalization of the effective heptamer model
- two different singlet regimes:  $\sigma$ -type and  $\pi$ -type



## Singlet State in Heptamer

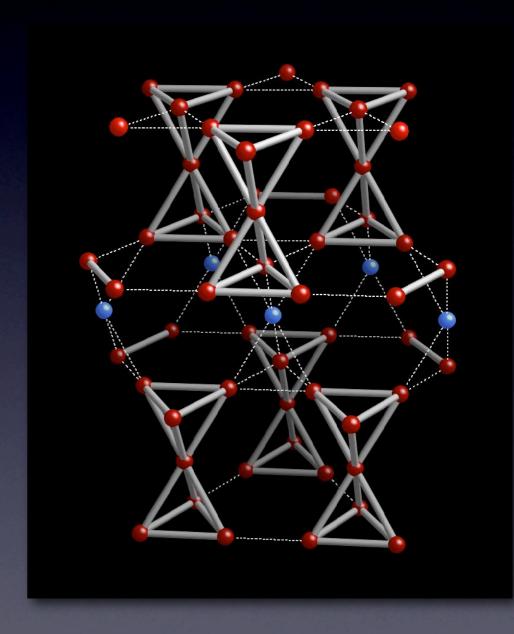
- singlet ground state for realistic parameters =  $\sigma$ -type
  - 'molecule' of the bonding states with three t<sub>2g</sub> orbitals
- estimate of the spin gap is larger than the experimental one: heptamer-heptamer coupling?
- comprehensive understanding of the T-dependence of magnetic susceptibility



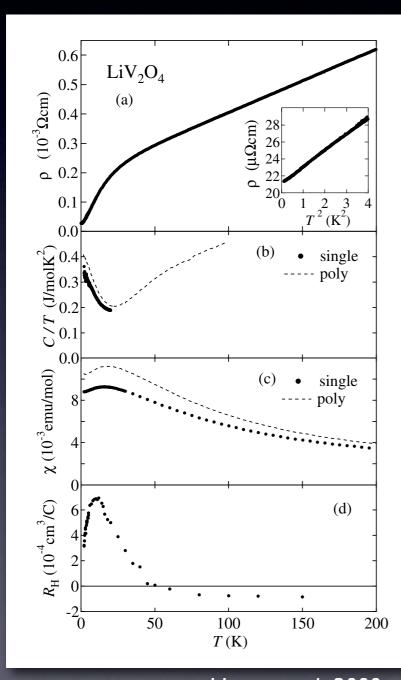


## Open Issues...

- Once the heptamers are assumed to be stable, experimental results at low-T phase are explained comprehensively.
- What is the mechanism of the heptamer formation? How is the degeneracy in the frustrated pyrochlore system lifted?
- Is similar phenomenon seen in other mixed-valence compounds?



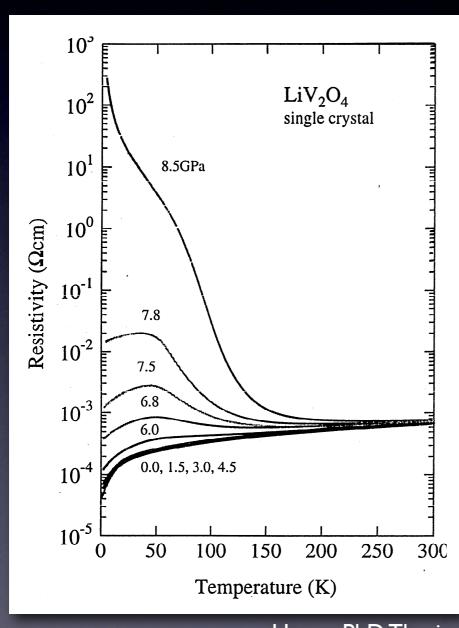
## Another Mixed-Valence Compound LiV<sub>2</sub>O<sub>4</sub>: Heavy-Fermion Behavior



Urano et al., 2000

- mixed valence: $V^{3.5+} = (3d)^{1.5}$
- heavy mass (Kondo et al., 1997)
- cubic, metallic, no magnetic ordering (Rogers et al., 1967; Chmaissem et al., 1997; Mahajan et al., 1997; Merrin et al., 1998)
- only t<sub>2g</sub> electrons: new mechanism for heavy fermion behavior?
  - Kondo scenario
  - geometrical frustration + correlation

### Implication of AIV<sub>2</sub>O<sub>4</sub>?



Urano, PhD Thesis

- metal-to-insulator transition by applying pressure: opposite to usual pressure effect
- short and long V-V bonds in the insulating state (EXAFS by Niitaka et al., unpublished)
- possibility: some cluster formation similar to AIV<sub>2</sub>O<sub>4</sub>

### Summary

- introduction to spinels and t<sub>2g</sub> orbital physics
- $^{\$}$  controversy on orbital ordering in  $ZnV_2O_4$ 
  - different models for spin/orbital order in ZnV<sub>2</sub>O<sub>4</sub>: relative importance of Kugel-Khomskii superexchange, Jahn-Teller and relativistic spin-orbit couplings
  - symmetry analysis: lesson from experiments in MnV<sub>2</sub>O<sub>4</sub>
- self-organized 7-site cluster (heptamer) in AIV<sub>2</sub>O<sub>4</sub>
  - heptamer scenario: 'molecule' of bonding states with anisotropic t<sub>2g</sub> orbitals
  - implication to heavy-fermion compound LiV2O4