

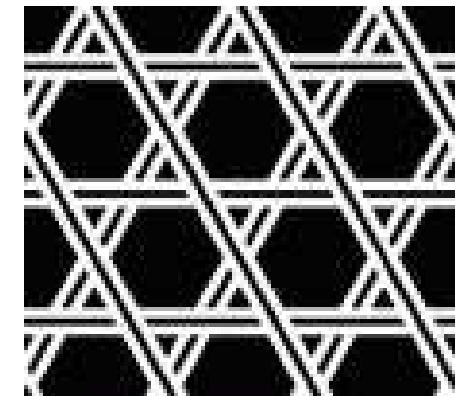
Algebraic vortex liquid theory of a quantum antiferromagnet on the kagome lattice

Shinsei Ryu (KITP, UCSB)

Olexei I. Motrunich (Caltech)

Jason Alicea (UCSB -> Caltech)

Matthew P. A. Fisher (Microsoft Station Q, UCSB)



Kagome kagome
Kago no naka no tori wa
Itsuitsu deyaru
Yoake no ban ni
Tsuru to kame ga subetta
Ushiro no shoumen - Daare

Surround, surround the bird.
You are a trapped bird inside the cage.
When oh when will you come out?
At the dawn's twilight,
The crane and the turtle both slid.
Right behind you - Who __ s it?

-- a Japanese nursery rhyme

Kagome materials (-> Young Lee and Philippe Sindzingre talks)

volborthite Cu₃V₂O₇(OH)₂ 2H₂O

herbertsmithite ZnCu₃(OH)₆Cl₂

$$\left. \right\} S = 1/2$$

SCGO SrCrGaO $S = 3/2$

jarosites KM₃(OH)₆(SO₄)₂

K=Cr ($S = 3/2$) K=Fe ($S = 5/2$)

Classical kagome magnets

$$\mathcal{H} = J \sum_{\mathbf{r}, \mathbf{r}'}^{\text{n.n.}} \mathbf{S}_{\mathbf{r}} \cdot \mathbf{S}_{\mathbf{r}'}$$

continuous local GS degeneracy

$$\mathbf{S}^A + \mathbf{S}^B + \mathbf{S}^C = 0 \forall \Delta$$

$$\mathcal{H} = J \sum_{\mathbf{r}, \mathbf{r}'}^{\text{n.n.}} \cos(\varphi_{\mathbf{r}} - \varphi_{\mathbf{r}'})$$

infinitely many degenerate GSs

Huse and Rutenberg (92)

Numerics on quantum kagome magnets (-> Philippe Sindzingre talk)

$S=1/2$ quantum kagome Heisenberg model exact diagonalization upto N=36

disordered ground state

Lecheminant et al (1997) Waldtmann et al (1998)

small (or zero) spin gap

$$\Delta \sim J/20 \quad \text{or} \quad 0$$

gapless spin liquid ?

many low-lying singlets :

$$\sim 1.15^{N_{\text{site}}}$$

c.f quantum XY kagome AF Sindzingre (unpublished)

Approach in terms of vortices

easy-plane quantum AF: $J \gg J^z$

$$\mathcal{H} = \frac{1}{2} \sum_{\mathbf{r}, \mathbf{r}'} [J_{\mathbf{r}, \mathbf{r}'} S_{\mathbf{r}}^+ S_{\mathbf{r}'}^- + \text{h.c.}] + \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'}^z S_{\mathbf{r}}^z S_{\mathbf{r}'}^z$$

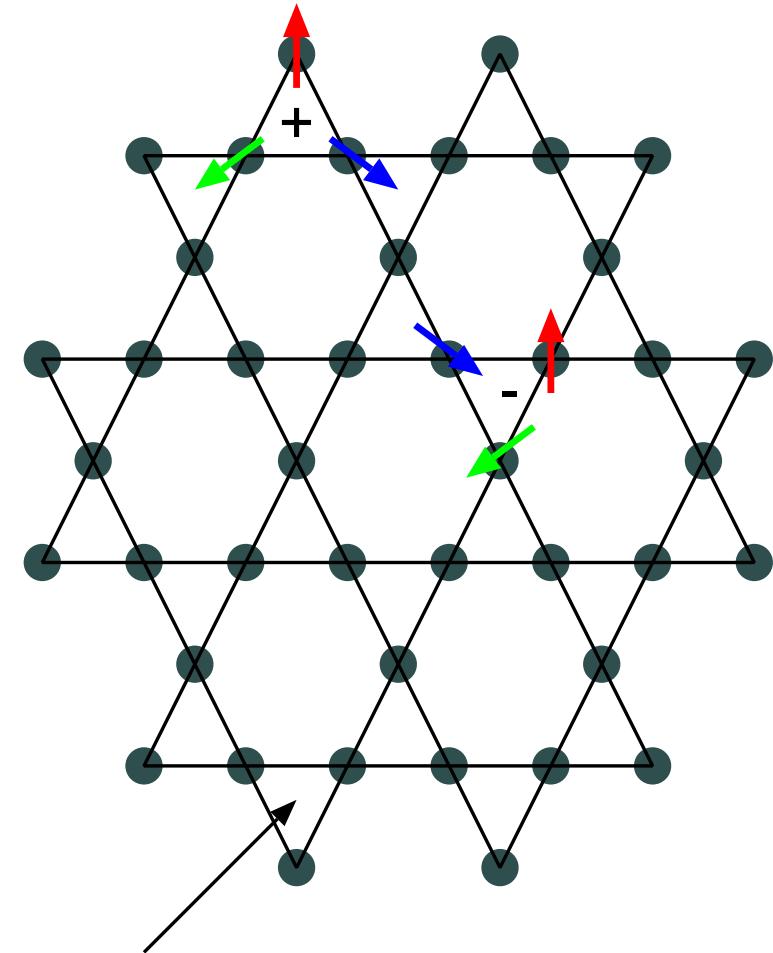
rotor Hamiltonian:

$$S_{\mathbf{r}}^+ \sim e^{i\varphi_{\mathbf{r}}}$$

$$S_{\mathbf{r}}^z \sim n_{\mathbf{r}} - n_{\mathbf{r}}^0$$

$$[\varphi_{\mathbf{r}}, n_{\mathbf{r}}] = i\delta_{\mathbf{r}, \mathbf{r}'} \quad n^0 = 1/2 \quad \mathbf{S}=1/2$$

$$\mathcal{H} = - \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'} \cos(\varphi_{\mathbf{r}} - \varphi_{\mathbf{r}'} - b_{\mathbf{r}, \mathbf{r}'}^{(0)}) + \sum_{\mathbf{r}} U_{\mathbf{r}} (n_{\mathbf{r}} - n_{\mathbf{r}}^0)^2$$



geometrical frustration :

$$\sum_{\text{triangle}} b_{\mathbf{r}, \mathbf{r}'}^{(0)} = \pi$$

Approach in terms of vortices

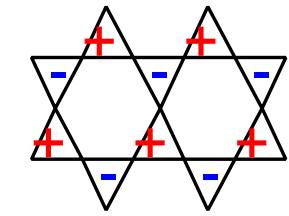
easy-plane quantum AF: $J \gg J^z$

$$\mathcal{H} = \frac{1}{2} \sum_{\mathbf{r}, \mathbf{r}'} [J_{\mathbf{r}, \mathbf{r}'} S_{\mathbf{r}}^+ S_{\mathbf{r}'}^- + \text{h.c.}] + \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'}^z S_{\mathbf{r}}^z S_{\mathbf{r}'}^z$$

rotor Hamiltonian:

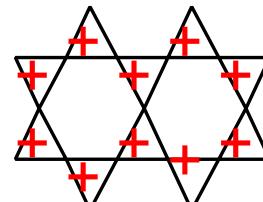
$$\mathcal{H} = - \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'} \cos(\varphi_{\mathbf{r}} - \varphi_{\mathbf{r}'} - b_{\mathbf{r}, \mathbf{r}'}^{(0)}) + \sum_{\mathbf{r}} U_{\mathbf{r}} (n_{\mathbf{r}} - n_{\mathbf{r}}^0)^2$$

$$n^0 = \begin{cases} 1/2 & @ \text{kagome sites} \\ 0 & @ \text{extra sites} \end{cases}$$



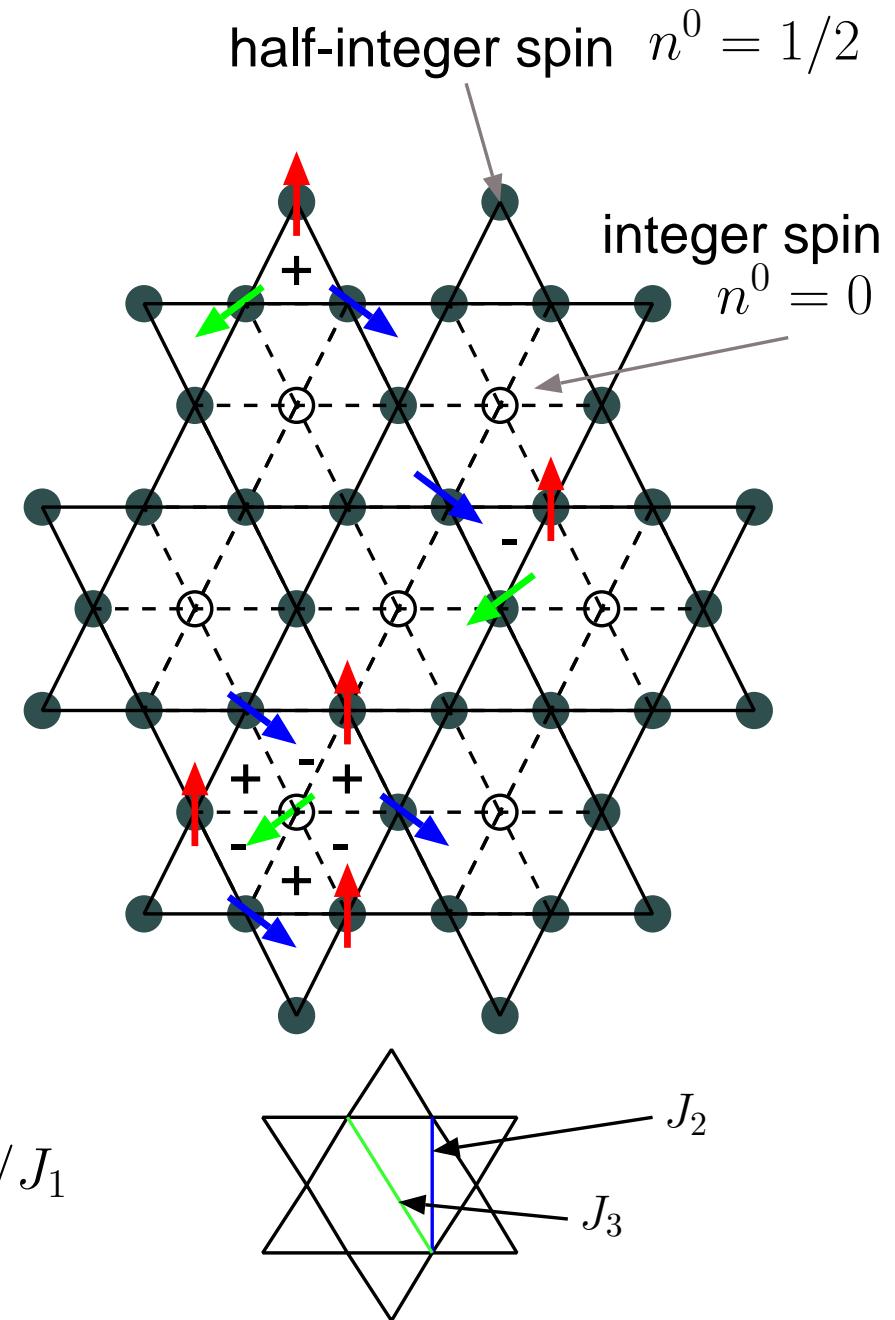
$\sqrt{3} \times \sqrt{3}$ state

???



$q = 0$ state

J_2/J_1



XY <-> scalar QED3 duality

Dasgupta and Halperin (81)

Nelson (88) D.-H. Lee and M.P.A Fisher (89)

rotor Hamiltonian: $\varphi_{\mathbf{r}} \ n_{\mathbf{r}}$

$$\mathcal{H} = - \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'} \cos(\varphi_{\mathbf{r}} - \varphi_{\mathbf{r}'} - b_{\mathbf{r}, \mathbf{r}'}^{(0)}) + \sum_{\mathbf{r}} U_{\mathbf{r}} (n_{\mathbf{r}} - n_{\mathbf{r}}^0)^2$$

spin

$$n_{\mathbf{r}}^{(0)}$$

geometrical
frustration

$$\sum_{\text{triangle}} b_{\mathbf{r}, \mathbf{r}'}^{(0)} = \pi$$



dual flux

$$\frac{1}{2\pi} (\nabla \times \mathbf{a})_{\mathbf{r}} = n_{\mathbf{r}} - n_{\mathbf{r}}^{(0)}$$



vortex density
frustration

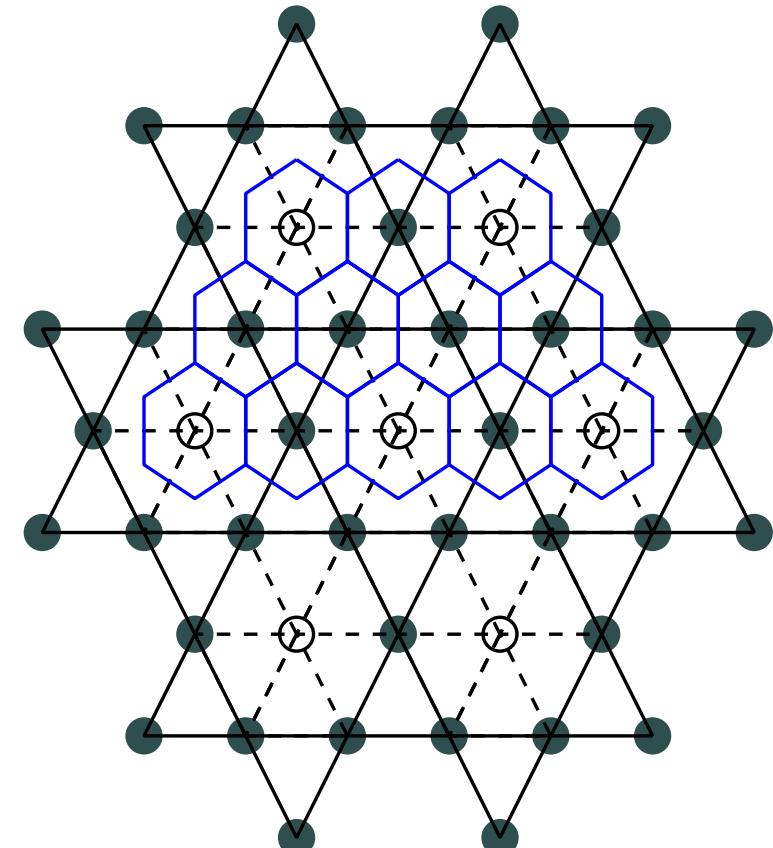
$$(\nabla \cdot \mathbf{e})_{\mathbf{x}} = N_{\mathbf{x}} - 1/2$$

scalar QED3: $a_{\mathbf{xx}'} \ e_{\mathbf{xx}'} \ \theta_{\mathbf{x}}$

$$\mathcal{H} = 2\pi^2 \sum_{\mathbf{xx}'} J_{\mathbf{xx}'} e_{\mathbf{xx}'}^2 + \frac{U}{(2\pi)^2} \sum_{\mathbf{r}} (\nabla \times \mathbf{a})_{\mathbf{r}}^2 - \sum_{\mathbf{xx}'} t_{\mathbf{xx}'} \cos(\theta_{\mathbf{x}} - \theta_{\mathbf{x}'} - a_{\mathbf{xx}'})$$

$e^{\pm i\theta_{\mathbf{x}}}$ vortex creation/annihilation operator

$a_{\mathbf{xx}'}$ U(1) gauge field



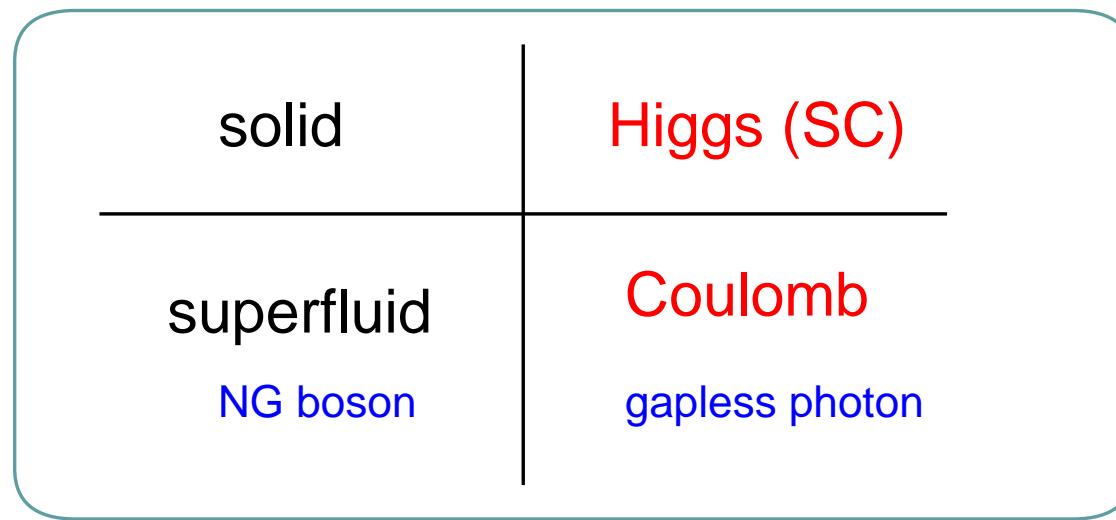
XY \leftrightarrow scalar QED3 duality

Dasgupta and Halperin (81)

Nelson (88) M.P.A Fisher, D.-H. Lee (89)

rotor Hamiltonian: $\varphi_{\mathbf{r}} n_{\mathbf{r}}$

$$\mathcal{H} = - \sum_{\mathbf{r}, \mathbf{r}'} J_{\mathbf{r}, \mathbf{r}'} \cos(\varphi_{\mathbf{r}} - \varphi_{\mathbf{r}'} - b_{\mathbf{r}, \mathbf{r}'}^{(0)}) + \sum_{\mathbf{r}} U_{\mathbf{r}} (n_{\mathbf{r}} - n_{\mathbf{r}}^0)^2$$

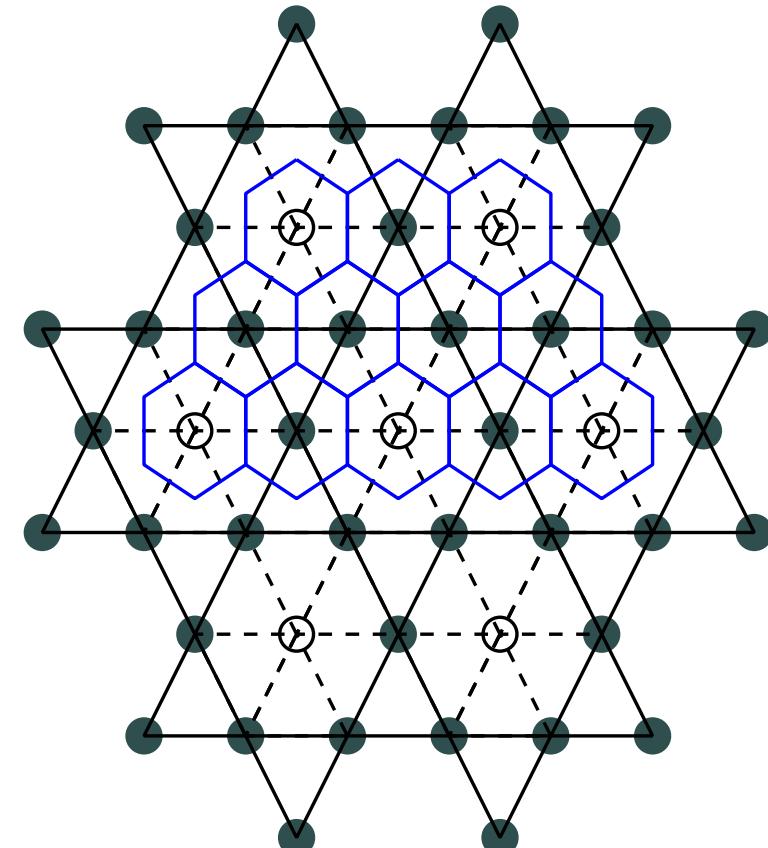


scalar QED3: $a_{\mathbf{xx}'} e_{\mathbf{xx}'} \theta_{\mathbf{x}}$

$$\mathcal{H} = 2\pi^2 \sum_{\mathbf{xx}'} J_{\mathbf{xx}'} e_{\mathbf{xx}'}^2 + \frac{U}{(2\pi)^2} \sum_{\mathbf{r}} (\nabla \times \mathbf{a})_{\mathbf{r}}^2 - \sum_{\mathbf{xx}'} t_{\mathbf{xx}'} \cos(\theta_{\mathbf{x}} - \theta_{\mathbf{x}'} - a_{\mathbf{xx}'})$$

$e^{\pm i\theta_{\mathbf{x}}}$ vortex creation/annihilation operator

$a_{\mathbf{xx}'}$ U(1) gauge field



looking for a saddle point : extracting effective theory by fermionization

Fradkin (89)

$$e^{i\theta_x} = d_x^\dagger \exp \left[i \sum_{x \neq x'} \arg(x, x') N_{x'} \right]$$

vortex(fermion)

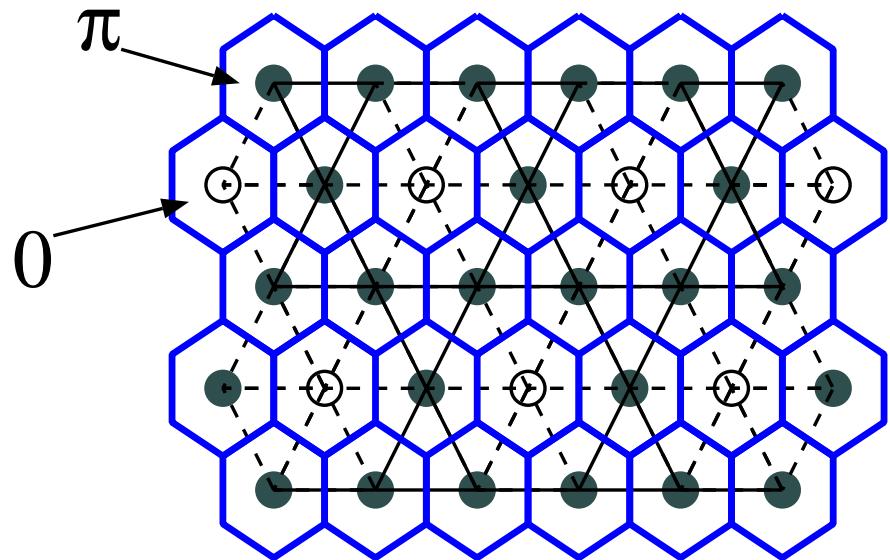
vortex(boson)

$$\mathcal{H}_{\text{ferm}} = - \sum_{xx'} \left[t_{xx'} d_x^\dagger d_{x'} e^{-i(a_{xx'} + A_{xx'})} + \text{h.c.} \right]$$

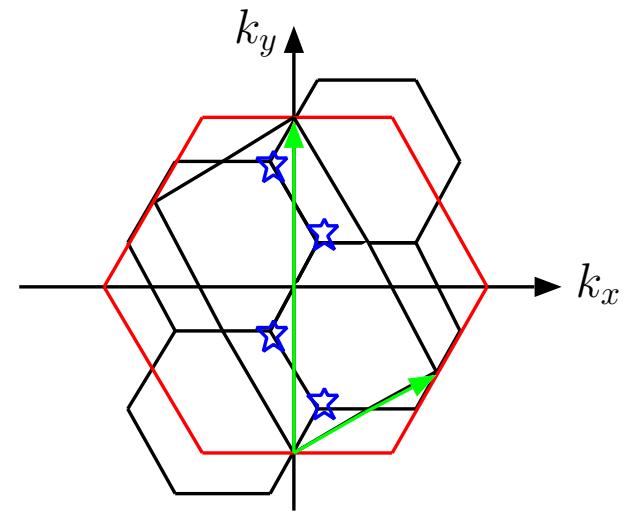
Chern-Simons gauge field

$$(\nabla \times A)_x = 2\pi N_x$$

'flux smearing' mean field



vortex band structure



8 Dirac cones

critical spin liquid: algebraic vortex liquid phase (AVL)

N=8 non-compact QED3

$$S_{\text{eff}} = \int d^3x \sum_{(\alpha)=1}^8 \bar{\psi}_{(\alpha)} \gamma_\mu (\partial_\mu - ia_\mu) \psi_{(\alpha)} + \frac{1}{2e^2} (\epsilon_{\mu\nu\lambda} \partial_\nu a_\lambda)^2$$

- CS gauge field is irrelevant
- $N=8 > N_c$ (?) gauge field is strongly screened
- all fermion bilinears are prohibited by PSG
- there is no monopole <- microscopic S^z conservation law (non-compact = no-monopole event)

Algebraic Spin Liquid (ASL)	AVL
<p>SU(2) symmetric</p> <p>fermionic spinons</p> <p>compact U(1) gauge field</p> <p>always confining</p> <p>bare gauge coupling is large (infinity)</p> <p>many mean field candidates -> energetics "landscape problem"</p>	<p>U(1) symmetric</p> <p>fermionic vortices</p> <p>non-compact U(1) gauge field</p> <p>"natural" mean field, but no energetics subtlities in time-reversal symmetry</p>

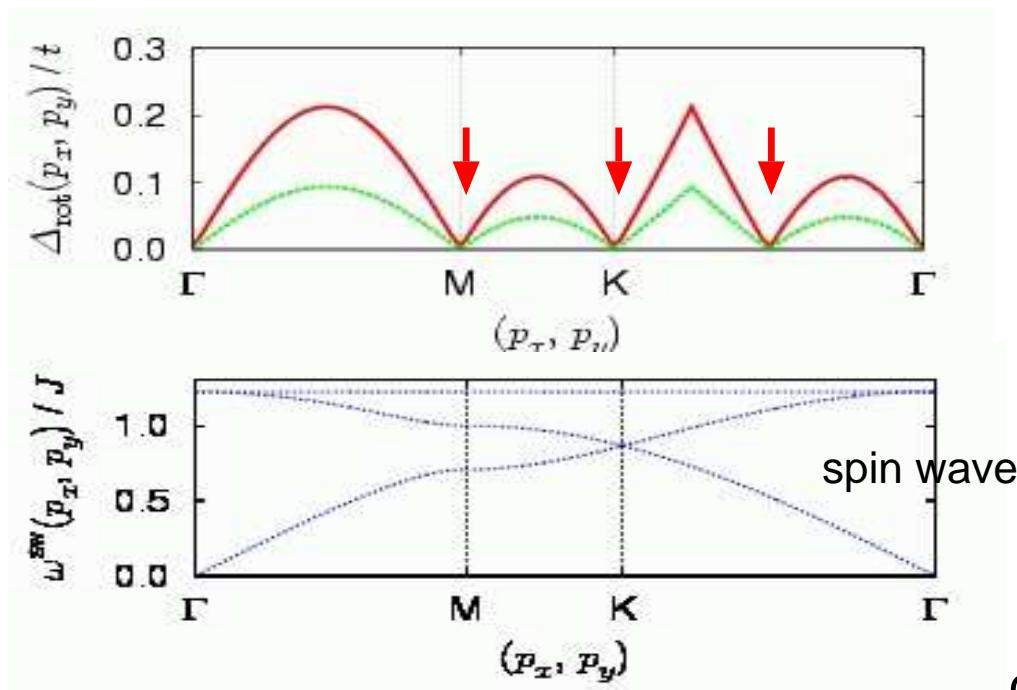
roton minima from vortex-antivortex excitations

N=8 non-compact QED3

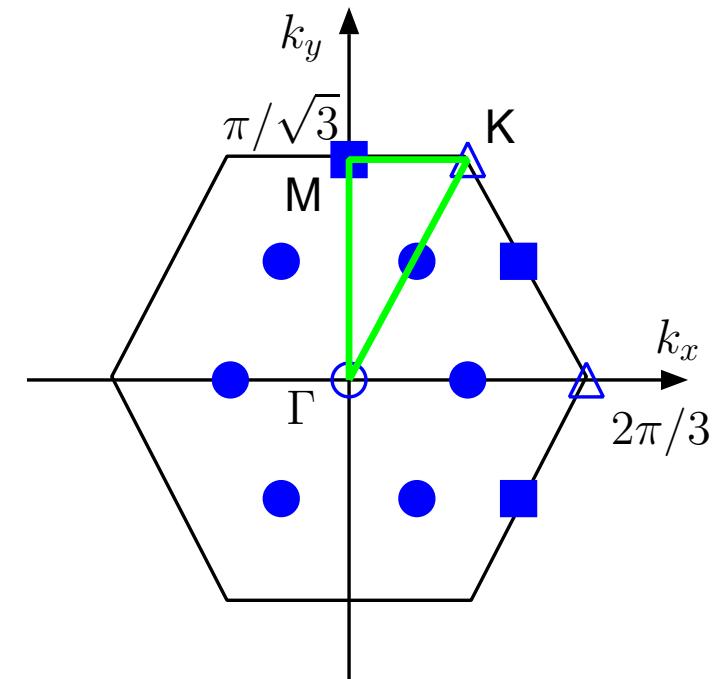
$$S_{\text{eff}} = \int d^3x \sum_{(\alpha)=1}^8 \bar{\psi}_{(\alpha)} \gamma_\mu (\partial_\mu - ia_\mu) \psi_{(\alpha)} + \frac{1}{2e^2} (\epsilon_{\mu\nu\lambda} \partial_\nu a_\lambda)^2$$

vortex(fermion)

$$\Delta_{\text{rot}}(\mathbf{p}) = \min_{\mathbf{k}} \{ E^+(\mathbf{k} + \mathbf{p}) - E^-(\mathbf{k}) \}$$



c.f. triangular lattice



dynamical structure factors from PSG

N=8 QED3 'remembers' the original lattice model through its **PSG**

symmetry transformation + gauge transformation

$$S^{zz}(\mathbf{k} = \mathbf{K} + \mathbf{q}, \omega) \propto \frac{\Theta(\omega^2 - \mathbf{q}^2)}{(\omega^2 - \mathbf{q}^2)^{(2-\eta_z)/2}}$$

ferminon bilinears

$$S_{\mathbf{r}}^z = \frac{\nabla \times \mathbf{a}}{2\pi} + \sum_i C_i \psi^\dagger M_i \psi + \dots$$

$$\eta_z \sim 3 - \frac{128}{3\pi^2 N} = 2.46 \quad @ \pm \mathbf{Q} \quad \pm \mathbf{P}_{1,2,3}$$

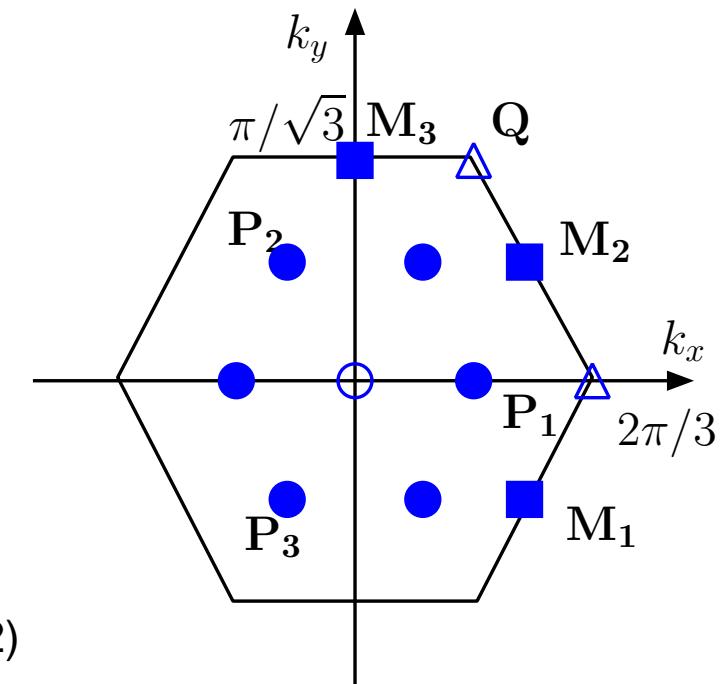
Rantner and Wen (02)

$$S^{+-}(\mathbf{k} = \mathbf{K} + \mathbf{q}, \omega) \propto \frac{\Theta(\omega^2 - \mathbf{q}^2)}{(\omega^2 - \mathbf{q}^2)^{(2-\eta_\pm)/2}}$$

monopole insertion

$$\eta_\pm \sim 0.53N - 1 = 3.24 \quad @ \text{all 12 momenta}$$

Borokhov, Kapustin, Wu (02)



Summary and future issues

effective field theory: N=8 QED3 + PSG

large N \longleftrightarrow many competing ordered states nearby

stability of spin liquid state (no chiral symmetry breaking)

roton minima, dynamical suseptibilities

specific heat $C \sim T^2$ thermal conductivity $\kappa \sim T$

- low-T specific data of herbertsmithite $C \sim T^\alpha$ $\alpha < 1$
 - including distortion -- application to volborthite
 - connection to other SL states ?
 - U(1) compact QED3 with N=4, Ran et al.
 - gapped Z2 SL, Sachdev, Wang and Vishwanath

discussions : dynamical structure factors

$$S^{\alpha\beta}(\mathbf{q}, \omega) = \langle S^\alpha(-\mathbf{q}, -\omega) S^\beta(\mathbf{q}, \omega) \rangle$$

gapless spin liquid scenarios

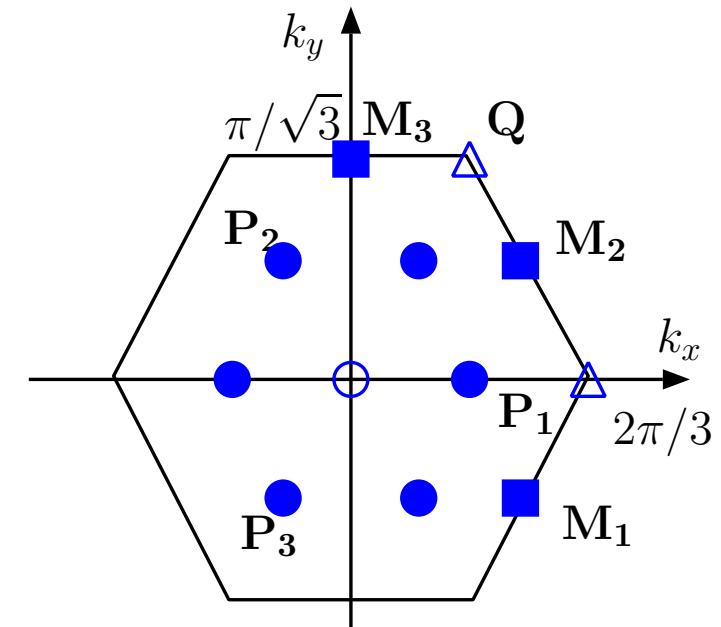
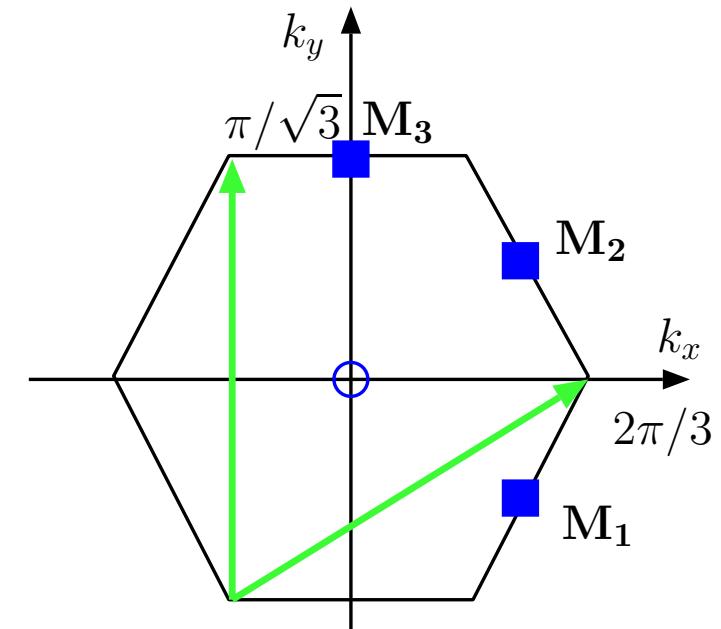
$$S^{\alpha\beta}(\mathbf{k} = \mathbf{K} + \mathbf{q}, \omega) \propto \frac{\Theta(\omega^2 - \mathbf{q}^2)}{(\omega^2 - \mathbf{q}^2)^{(2-\eta)/2}} \quad @ \mathbf{K}$$

slave particle approach: (Hastings, Ran et al.)
(fermionic spinons, 4 Dirac + U(1) gauge field)

algebraic votex approach (easy-plane):
(fermionic vortices, 8 Dirac+ U(1) gauge field)

$$\eta_z \sim 3 - \frac{128}{3\pi^2 N} = 2.46 \quad @ \pm\mathbf{Q} \quad \pm\mathbf{P}_{1,2,3}$$

$$\eta_{\pm} \sim 0.53N - 1 = 3.24 \quad @ \text{all 12 momenta}$$

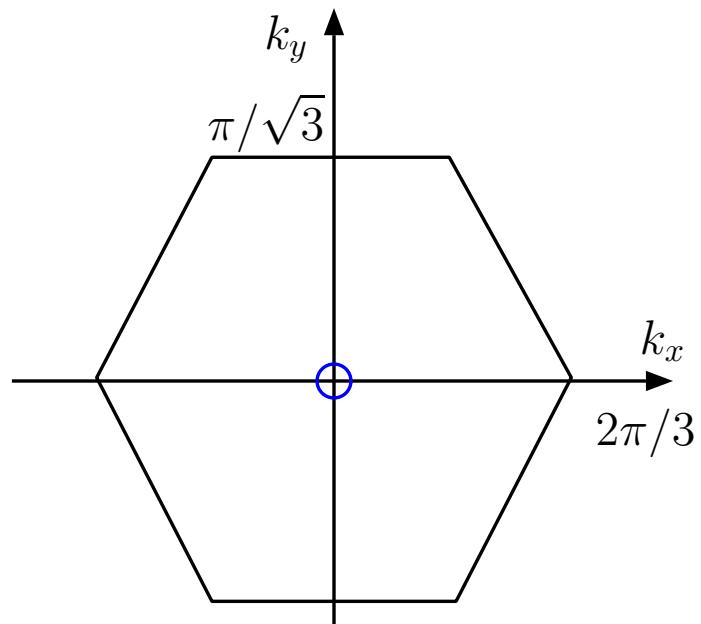


discussions : dynamical structure factors

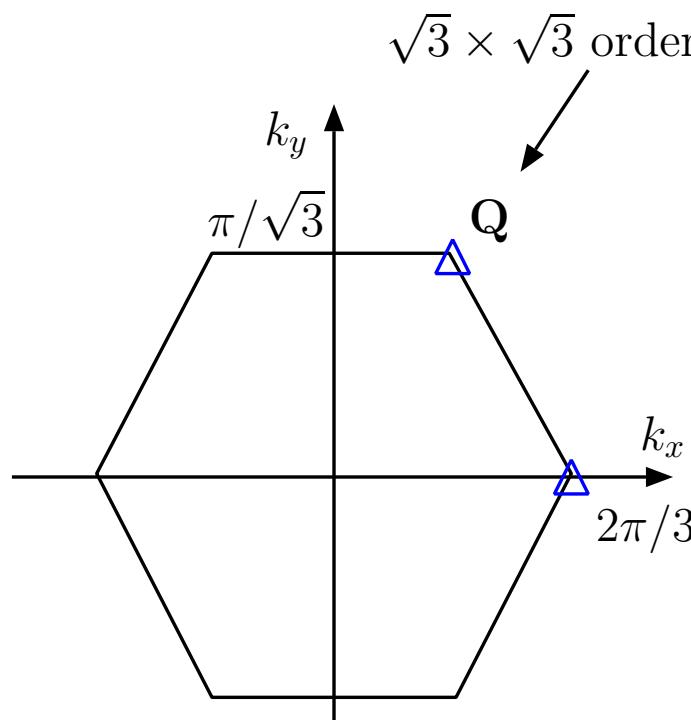
gapped spin liquid scenarios -> look for minima

slave particle approach: (Sachdev, Wang and Vishwanath)
(gapped Z2 spin liquid, bosonic spinon)

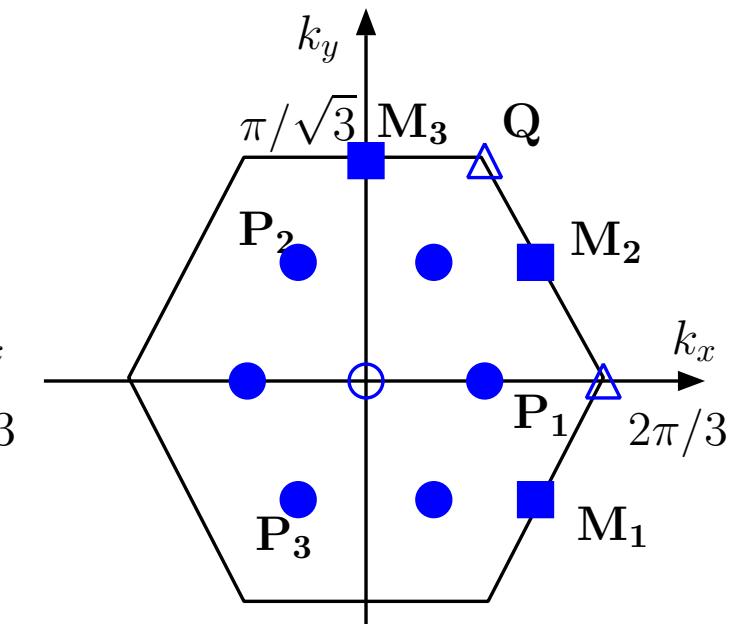
four different saddle points :



zero flux state $Q_1 = -Q_2$ state



zero flux state $Q_1 = Q_2$ state

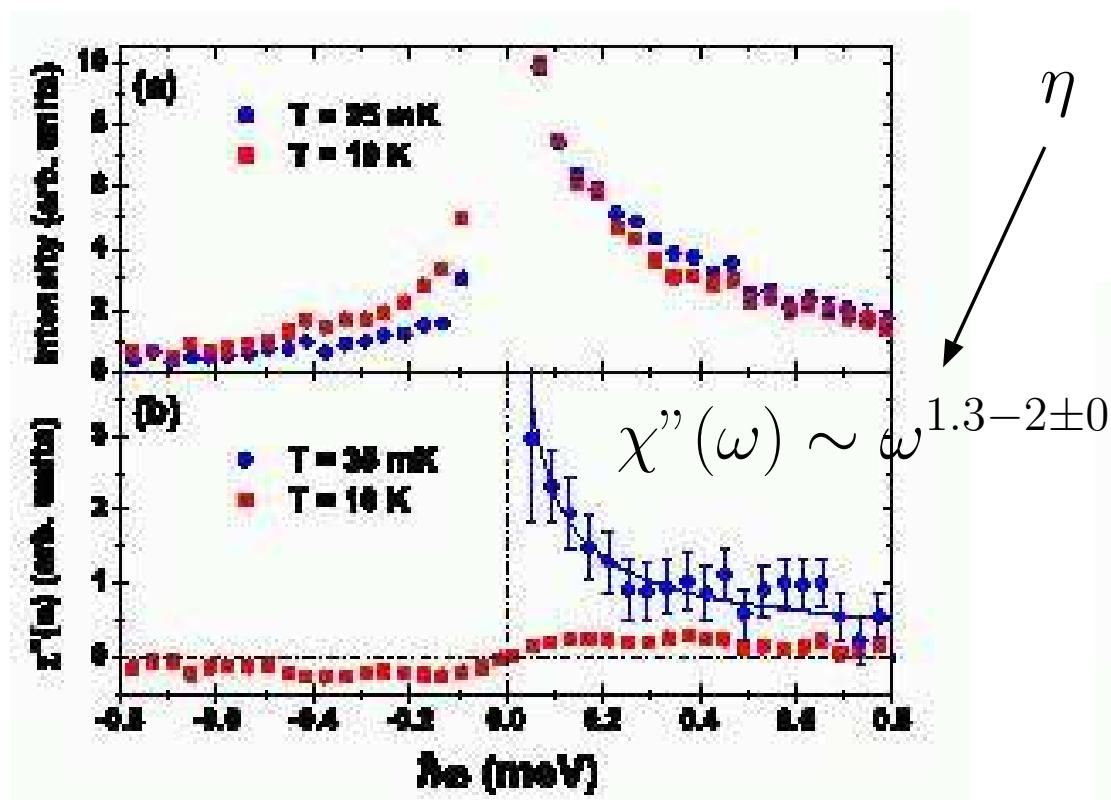


pi flux states 0Hex, π Rhom state
 π Hex, π Rhom state

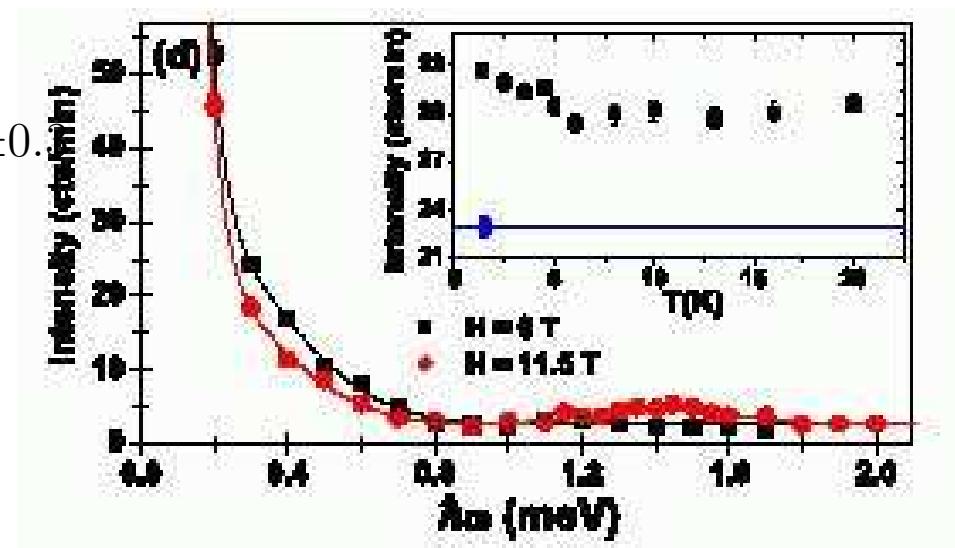
discussions : dynamical structure factors

Numerical dynamical structure factor by Andreas Laeuchli (Philippe Sindzingre talk)

Helton et al. PRL (2007)



c.f. $1/T_1 \sim T^\eta$
T. Imai et al. cond-mat/0703141
O. Ofer et al., cond-mat/0610540



integrated over $0.25 < |Q| < 1.5 \text{ \AA}^{-1}$