



Why are magnetism and ferroelectricity contra-indicated?

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Outline

Multiferroics: definition

Conventional mechanism for ferroelectricity (" do-ness")

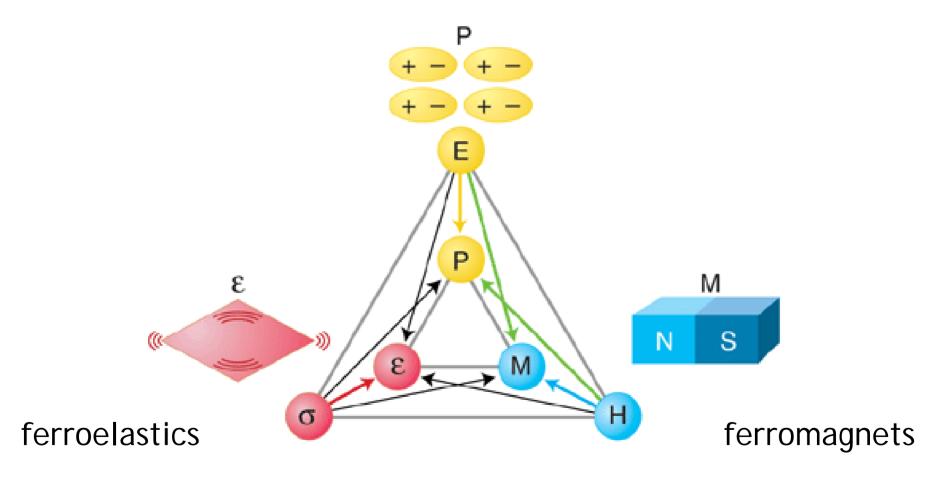
Combining magnetism and ferroelectricity alternative mechanisms for ferroelectricity non-d-electron magnets



Multiferroics



ferroelectrics

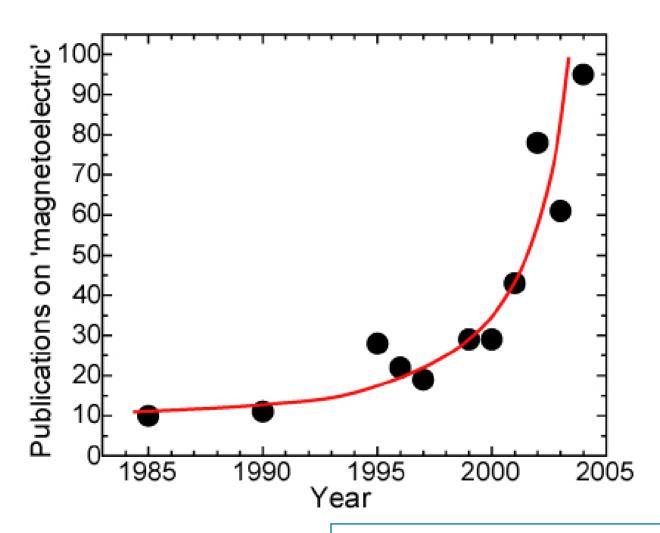


The renaissance of magnetoelectric multiferroics, N. A. Spaldin and M. Fiebig, Science 15, 5733 (2005)



Increasing interest in magnetoelectric multiferroics





The revival of the magnetoelectric effect, M. Fiebig, J. Phys. D 38, R123 (2005)



Problem

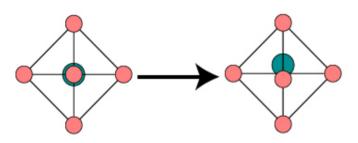


There are (almost) no magnetic ferroelectrics

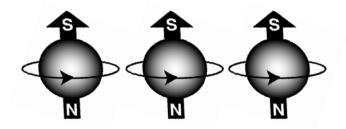
Magnetism and ferroelectricity are *chemically contra-indicated*:

Magnetism requires localized (transition metal d) electrons

Atoms with localized *d* electrons don't off-center to form polar ferroelectric states



requires empty *d* orbitals
Second-order Jahn-Teller effect



Ferromagnetism requires filled *d* orbitals Stoner instability

N.A. Hill, Why are there so few magnetic ferroelectrics? J. Phys. Chem. B 104, 6694-6709 (2000)



Conventional ferroelectricity mechanism I. Hand-waving explanation of "do-ness"



+

+

paraelectric

ferroelectric

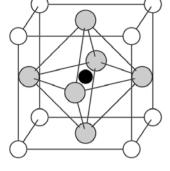


Conventional ferroelectricity mechanism I. Hand-waving explanation of "do-ness"





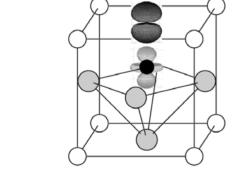
+



paraelectric



+



ferroelectric

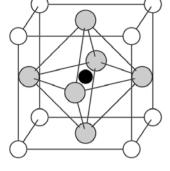
Unfavorable Coulomb repulsion between oxygen and transition metal valence electrons



Conventional ferroelectricity mechanism I. Hand-waving explanation of "do-ness"

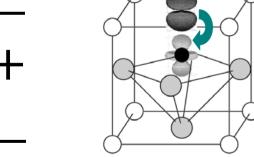






paraelectric





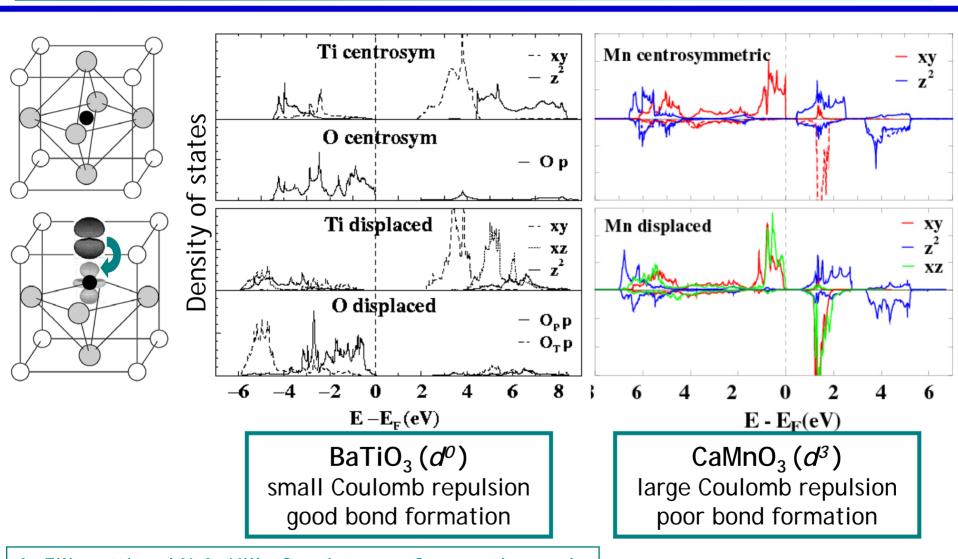
ferroelectric

Favorable "ligand field stabilization" of *empty* cation *d* orbitals by oxygen *p* electrons



Illustration within LDA





A. Filippetti and N.A. Hill, *Coexistence of magnetism and ferroelectricity in perovskites*, PRB **65**, 195120 (2002).



Perturbation theory



Expand Hamiltonian as function of atomic distortion (normal coordinate), Q:

$$H=H^{(0)}+H^{(1)}Q+rac{1}{2}H^{(2)}Q^2 \quad ext{ where } \qquad egin{aligned} H^{(1)}Q&=(\delta H/\delta Q)_0Q\ H^{(2)}Q^2&=(\delta^2 H/\delta Q^2)_0Q^2 \end{aligned}$$

then

$$E(Q) = E(0) + <0 |(\delta H/\delta Q)_0|0 > Q + \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0 > -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 + ...$$

1st-order JT Non-zero for orbitally degenerate states

> always positive (moving nuclei with fixed electrons); want this to be small

always negative (relaxation of electron distribution); want this to be large 1) need a non-zero matrix element; 2) need En close to E(0)

Second-order Jahn-Teller effect



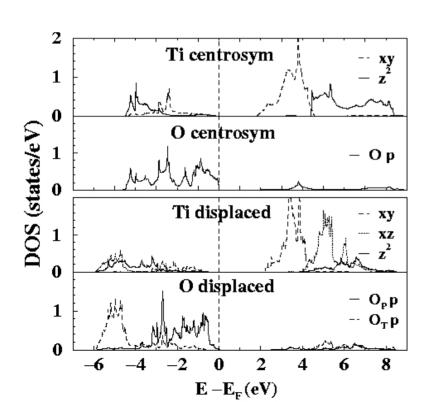
Perturbation theory (slightly hand-waving)



$$+ \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q^2)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} \right) Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/\delta Q)_0|n>|^2}{E_n - E(0)} Q^2 - \frac{1}{2} \left(<0 |(\delta^2 H/\delta Q)_0|0> -2\Sigma_n' \frac{|<0|(\delta^2 H/$$

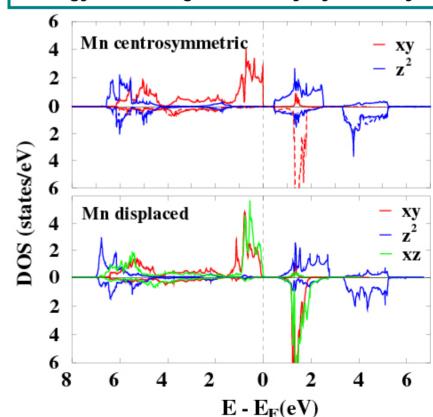
$BaTiO_3(d^0)$

Repulsive term small Energy-lowering term non-zero



$CaMnO_3(d^3)$

Repulsive term large Energy-lowering term 0 by symmetry





BUT magnetism requires localized electrons!



In perovskite structure oxides the source of magnetic, localized electrons is usually the transition metal *d* electrons e.g. LaMnO₃, SrRuO₃, etc.

Photo: Anna Karin-Axelsson, Imperial College, London





How to combine M and P?



either

- 1) use an alternative mechanism for P or
- 2) use an alternative mechanism for M



How to combine M and P?



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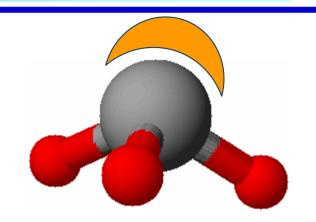


"Lone-pair active" multiferroics



Ferroelectricity from the "stereochemically active lone pair" on Bi³⁺ (cf ammonia, NH₃)

Magnetism from a 3d transition metal (Mn³⁺ or Fe³⁺)



<u>BiMnO₃:</u>

Ferromagnetic
Polar instability from Bi lone pairs
Anti-polar? (C2/c) Bi Bi

P. Baettig, R. Seshadri and N. A. Spaldin, *Anti-polarity in ideal BiMnO₃*, JACS **129**, 9854-9855 (2007).

BiFeO₃:

Ferroelectric, $P = 90 \mu C/cm^2$ Polar instability from Bi lone pairs *Anti*-ferromagnetic (weak FM)

Epitaxial BiFeO₃ multiferroic thin film heterostructures, Wang, Spaldin, Ramesh et al., Science 299, 1719 (2003)

<u>Another idea:</u> Combining ferroelectricity with ferromagnetism is HARD! ferr *i* magnetism might be easier...

P. Baettig, C. Ederer and N.A. Spaldin, PRB 72, 214105 (2005)

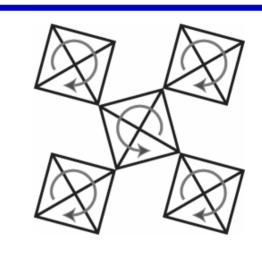


Geometric ferroelectrics

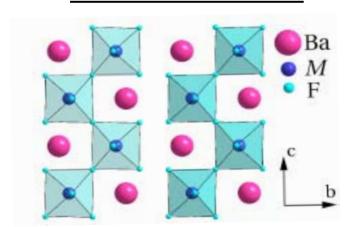


rotations driven by non-ideal ion packing doesn't yield a net P with 3D connectivity

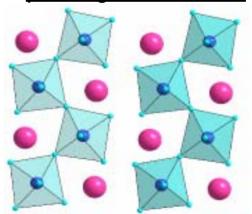
in 2D, inversion center can be lifted e.g. BaNiF₄ (FE $T_C = 1200K$; AFM, $T_N = 60K$)



reference structure



polar ground state



C. Ederer and N.A. Spaldin, Electric-field switchable magnets: The case of BaNiF₄, PRB **74**, 020401(R) (2006)



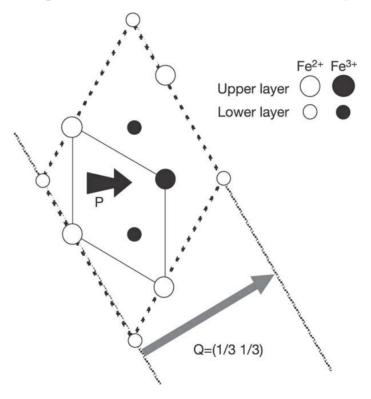
Charge ordered



e.g. LuFe₂O₄ (FE, T_C = 330K; frustrated magnet)

Asymmetric charge ordering of Fe²⁺ and Fe³⁺ causes polarization

Proposed structure:



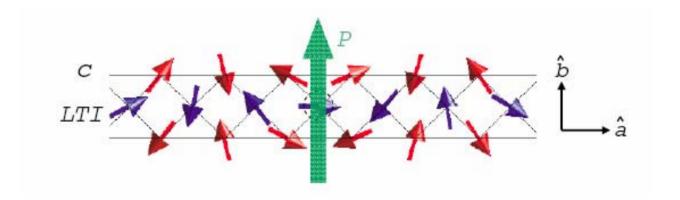
N. Ikeda et al., Nature 436, 1136 (2005)



Magnetically-driven (spiral) ferroelectricity



e.g. TbMnO₃ P very small BUT coupled to magnetic axis



T. Kimura et al., *Magnetic control of ferroelectric polarization*, Nature **426**, 55 (2004).



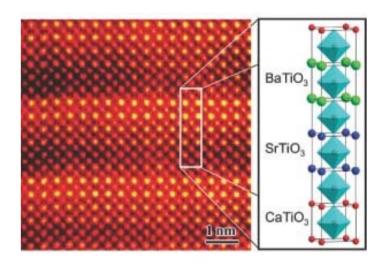
Another idea



Inversion symmetry breaking by three-component layering

Strong polarization enhancement in asymmetric three-component ferroelectric superlattices, H. N. Lee et al., Nature 433, 395 (2005).

A.J. Hatt and N.A. Spaldin, Tri-layer superlattices: A route to magnetoelectric multiferroics? APL **90**, 242916 (2007)





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either

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- 2) use an alternative mechanism for M

f electron magnetism

EuTiO₃ (Eu²⁺; Ti⁴⁺ is d^0) prediction (Fennie and Rabe): FM under strain



Summary



do cations can provide ferroelectricity but not magnetism

Alternative mechanisms for ferroelectricity are compatible with magnetism

f electron magnetism is compatible with d^0 ferroelectricity

Multiferroic with large magnetization and large polarization at room temperature not yet achieved

Magnetoelectric coupling an additional challenge