

KITP Program: Moments and Multiplets in Mott Materials

Supersolids in Lattice Boson Models



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Supersolid

Is it possible to have BEC in a Solid? (Quantum solid: ^4He)

ODLRO (Superfluid)

DLRO (Crystal Order)

$$\langle \widehat{\Psi}(\vec{r}) \rangle \neq 0, S_\rho(Q) \neq 0$$

In an ideal commensurate boson crystal at $T=0, N_{\text{atom}} = Z^* N_{\text{lattice}}$

$\langle \widehat{\Psi}(\vec{r}) \rangle \neq 0, S_\rho(Q) \neq 0$ is impossible!
Argue $T > 0$ case is also impossible.

End of Story?

Penrose and Onsager, P.R. 104, 576 (1956).

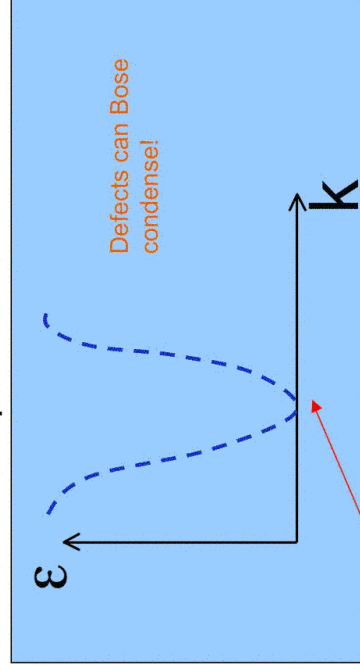
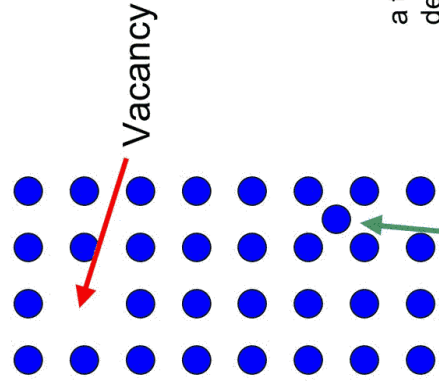


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Defect Condensation

- Andreev and Lifshitz, proposed “zero-point vibration” \rightarrow “defectons” and “impuritons”



a tiny density of quantum defects can give superfluidity while preserving crystalline order

Background crystal
+ Defect superflow
= Supersolid



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Andreev and Lifshitz, Sov. Phys. JETP **29**, 1107
Chester, PRA, **2**, 250

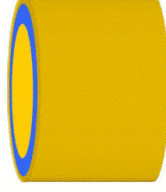


PRL 25, 1543 (1970)

Can a Solid Be “Superfluid”?

A. J. Leggett
 School of Mathematical and Physical Sciences, University of Sussex, Falmer, Brighton, Sussex, England
 (Received 15 September 1970)

It is suggested that the property of nonclassical rotational inertia possessed by superfluid liquid helium may be shared by some solids. In particular, nonclassical rotational inertia very probably occurs if the solid is Bose-condensed as recently proposed by Chester. Anomalous macroscopic effects are then predicted. However, the associated superfluid fraction is shown to be very small (probably $\leq 10^{-4}$) even at $T = 0$, so that these effects could well have been missed. Direct tests are proposed.



Leggett proposed to detect supersolid by measuring NCRI (Nonclassical rotational inertia)

$$I(T) = I_{\text{classical}}[1 - f_s(T)], f_s(T) = (\rho_s / \rho) \text{ is the “superfluid fraction”}$$

Estimate $\rho_s / \rho \approx 10^{-4}$, too small to observe (at that time).

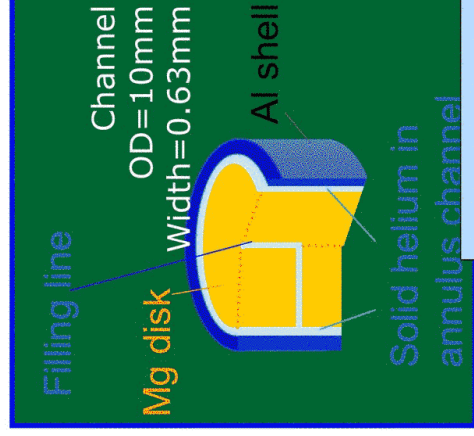
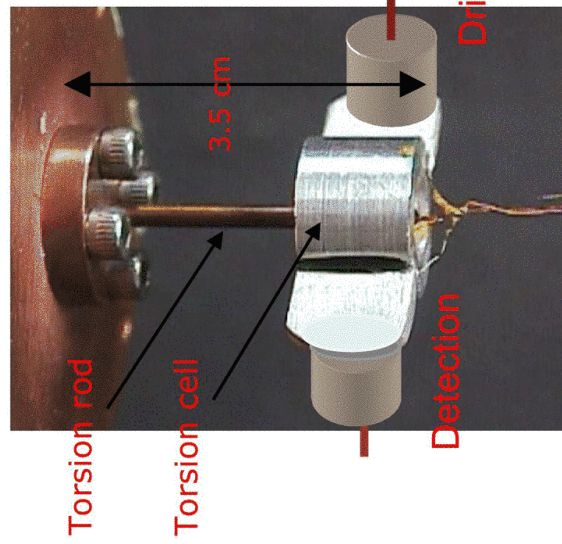


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Supersolid in ^4He

Torsion cell with helium in annulus



$$\tau_o = 2\pi \sqrt{\frac{I}{K}}$$

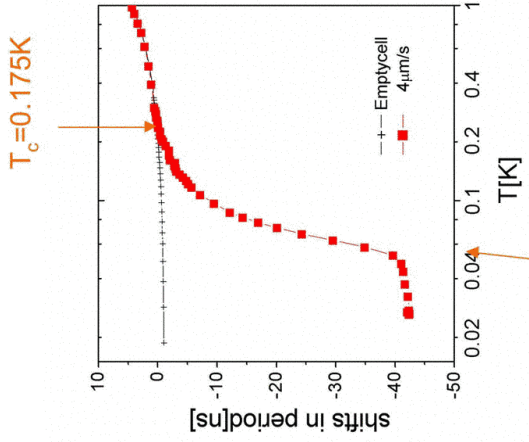


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Kim and Chan, Science 305, 1941 (2004);
 Nature 427, 225 (2004).

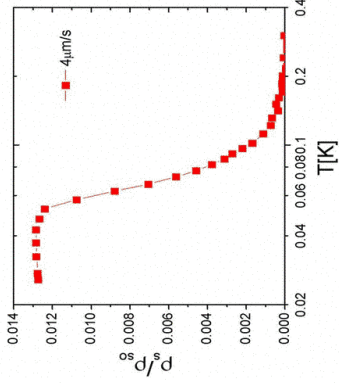


Supersolid in ^4He



$T_c = 0.175\text{K}$

The supersolid fraction is on the order of 1.3%



$$\tau_o = 2\pi\sqrt{\frac{I}{K}}$$

$$I(T) = I_{\text{classical}}[1 - f_s(T)]$$



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Supersolid in ^4He

Existence : Phase separation? Surface phenomenon?

Mechanism : Defect/Impurity condensate? (Annealing, doping with He^3) Grain boundary superfluidity

Does supersolid phase exist in a **Perfect Crystal (lattice)**?



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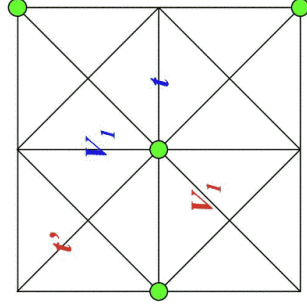


Hard-core Boson Hubbard model

$$\mathcal{H} = \sum_{ij} -t_{ij}(b_i^\dagger b_j + c. c.) + \sum_{ij} V_{ij} n_i n_j - \mu \sum_i n_i$$

XXZ model

$$\mathcal{H} = J \sum_{i,j} \left[\Delta S_i^z S_j^z + \frac{1}{2}(S_i^+ S_j^- + S_i^- S_j^+) \right] - h \sum_i S_i^z$$



$$S^+ = b^\dagger, S^- = b, S_z = (b^\dagger b - \frac{1}{2})$$

$$J_x = -2t, J'_x = -2t'$$

$$J_z = V_1, J'_z = V_2$$

$$h = \mu - \frac{1}{2}(zV_1 + z'V_2)$$



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Quantum Lattice Gas

	$\langle S_i^x S_j^x + S_i^y S_j^y \rangle$	ODLRO	Breaks rotational (gauge)U(1) symmetry
η DLRO Breaks translation symmetry	NO	NO	YES
	YES	NO	YES

NO	<p>Paraferro Normal Fluid (Gas, Liquid)</p>	<p>Spin Flop Superfluid</p>
YES	<p>Antiferro Normal Solid</p>	<p>Intermediate state Supersolid</p>

$\eta = 1$: perfect crystal ($n_B = 0$);
 $\eta = 0$: liquid state ($n_A = n_B$).

$$\eta \sim S_z^A - S_z^B$$



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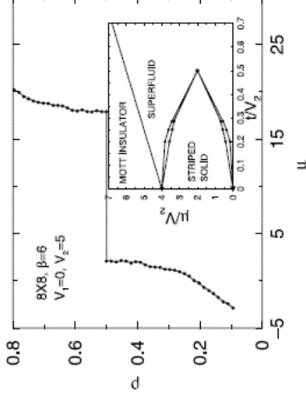
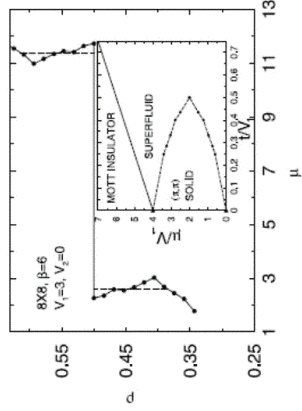
Matsuda and Tsueto, Supp. Prog. Theo. Phys. 46, 411 (1970)

Results of numerical simulations

Hard-core Boson model

$$H = -t \sum_{\langle ij \rangle} (a_i^\dagger a_j + a_j^\dagger a_i) + V_1 \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j + V_2 \sum_{\langle\langle ik \rangle\rangle} \hat{n}_i \hat{n}_k.$$

With N.N. interaction, the SS phase is **unstable** against phase separation. With N.N.N. interaction, a stable **striped SS** is found.



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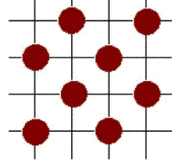
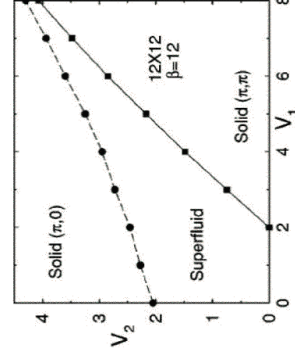
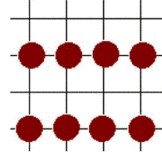
Batrouni *et al.*, PRL 84, 1599 (2000)



Half-filling

F. Hebert *et al.*, Phys. Rev. B 65, 014513 (2002)

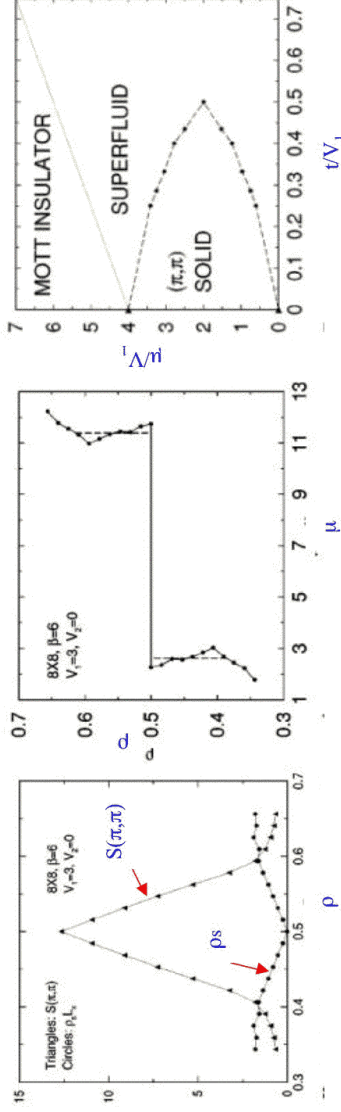
- V_1 dominant, **checkerboard solid** (π, π) .
- V_2 dominant, **striped solid** $(\pi, 0)$.
- In between, **superfluid**.



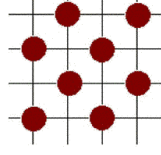
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Doped System ($V_2=0$)



- Coexistence of $S(\pi,\pi)$ and superfluidity.
- Negative compressibility $d\rho/d\mu$ indicates the phase is unstable. **Phase separation.**



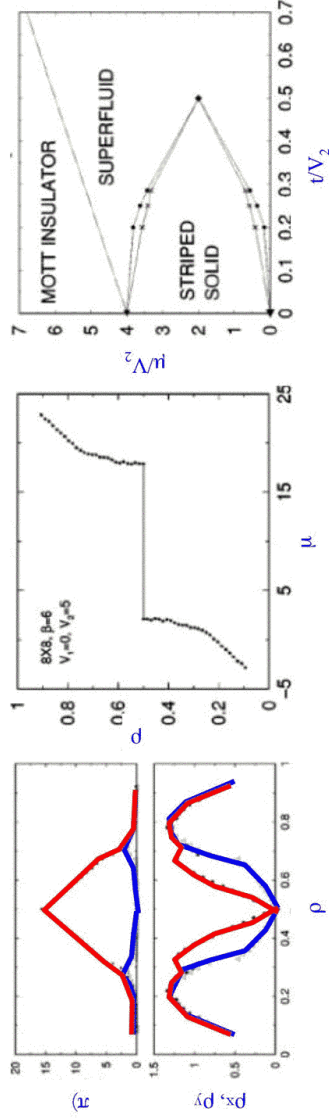
F. Hebert *et al.*, Phys. Rev. B **65**, 014513 (2002)



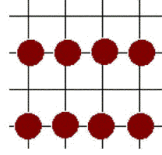
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Doped System ($V_1=0$)



- Superfluidity in x and y direction are both finite, not through one-dimensional channels. However, strong anisotropy.
- Positive compressibility, and the phase is stable.
- A stable supersolid is observed!



F. Hebert *et al.*, Phys. Rev. B **65**, 014513 (2002)



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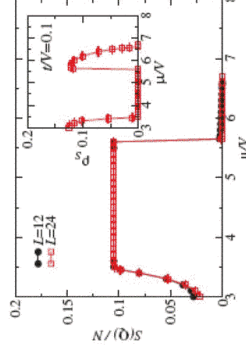
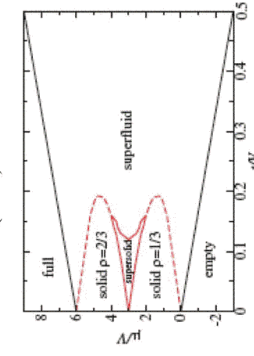
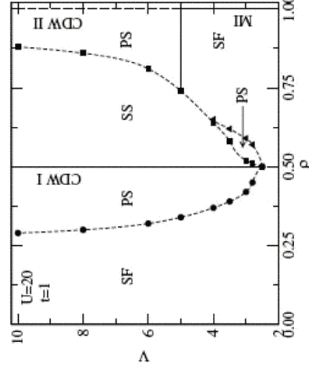


Soft-core Boson model

$$H = -t \sum_{\langle i,j \rangle} (a_i^\dagger a_j + a_j^\dagger a_i) - \mu \sum_i n_i + V \sum_{\langle i,j \rangle} n_i n_j + \frac{U}{2} \sum_i n_i (n_i - 1),$$

P. Sengupta *et al.*, PRL 94, 207202 (2005)

Supersolid phase of hard-core boson is found to be stabilized in **frustrated triangular lattice** by mechanism of **order by disorder**
 S. Wessel *et al.*, PRL 95, 127205; D. Heidarian *et al.*, PRL 95, 127206; R.G. Melko *et al.*, PRL 95 127207 (2005)



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VMFT, QMC

- Gutzwiller variational MF $|\phi_g\rangle = \prod_i f(n_i)|n_i\rangle$
- Quantum Monte Carlo (SSE)

Superfluid density: $\rho_s = \frac{\langle W^2 \rangle}{2t\beta L^{d-2}}$

Structure factor: $S(k) = \sum_r e^{ikr} \langle n(r_0)n(r_0 + r) \rangle$

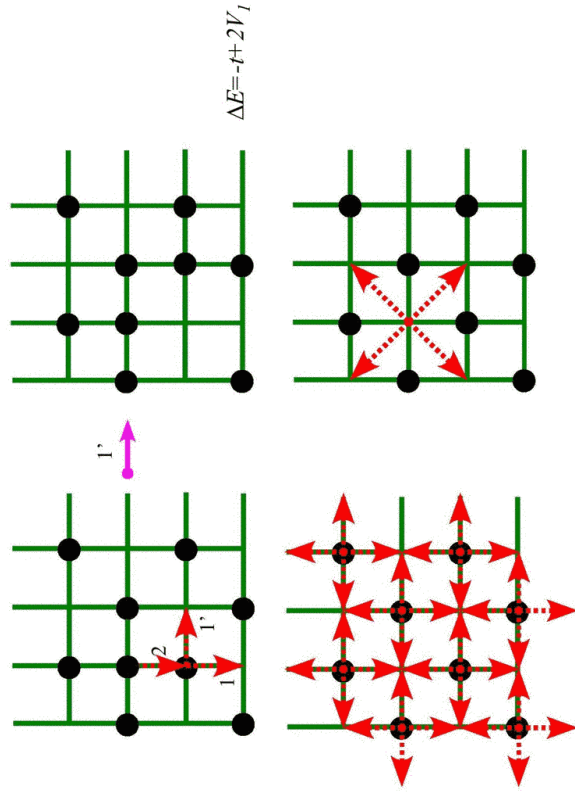


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M. Gutzwiller, Phys. Rev. Lett. 10, 159 (1963)
 A.W. Sandvik and J. Kurkijärvi, PRB 43, 5950 (1991).

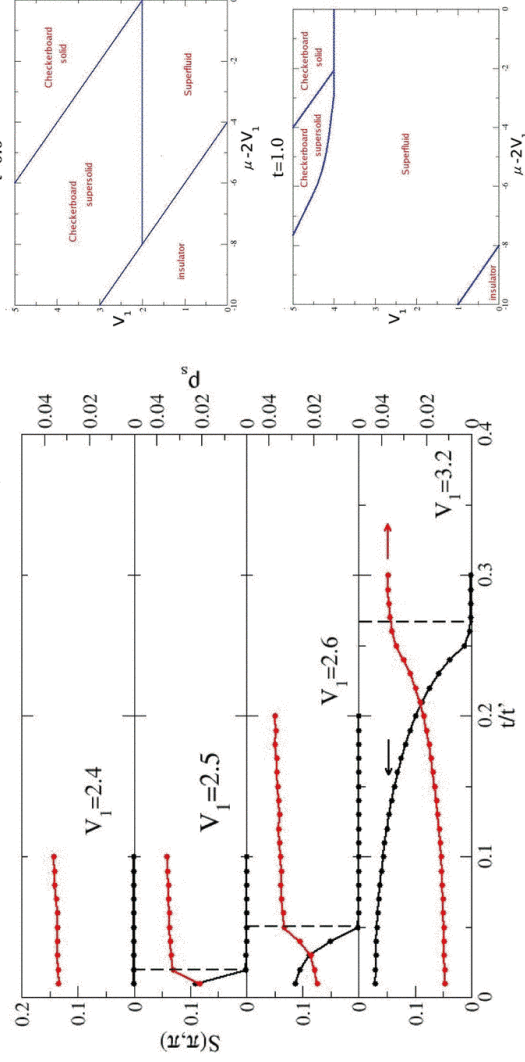
Effects of nnn hopping



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Finite nn hopping t ($V_2=0$)



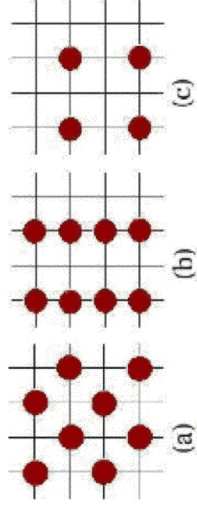
- Supersolid state emerges when V_1 is larger than $2.5t'$.
- When t turns on, supersolid melts into a superfluid.
- What is the effect of nnn repulsion V_2 ?



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Solid States



Structure factor	checkerboard	striped	quarter-filled
$S(\pi, \pi)$	$\neq 0$	0	$\neq 0$
$O_x + O_y$	0	$\neq 0$	$\neq 0$
$ O_x - O_y $	0	$\simeq O_x + O_y$	0

$$S(\mathbf{q}) = \frac{1}{N} \sum_j e^{iq \cdot \mathbf{r}_j} \langle n_j n_{j+1} \rangle$$

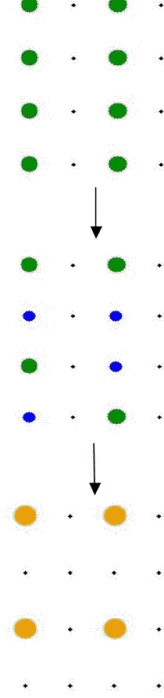
$$S(\pi, \pi), O_x = S(\pi, 0), O_y = S(0, \pi)$$



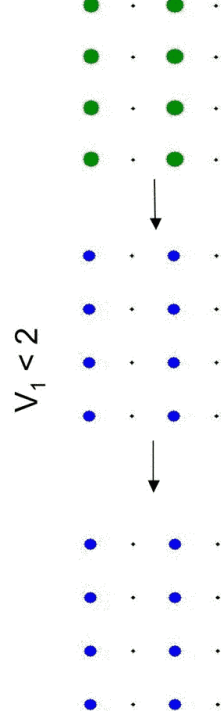
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Doping toward $\rho=0.25$



$$V_1 > 2$$



$$V_1 < 2$$

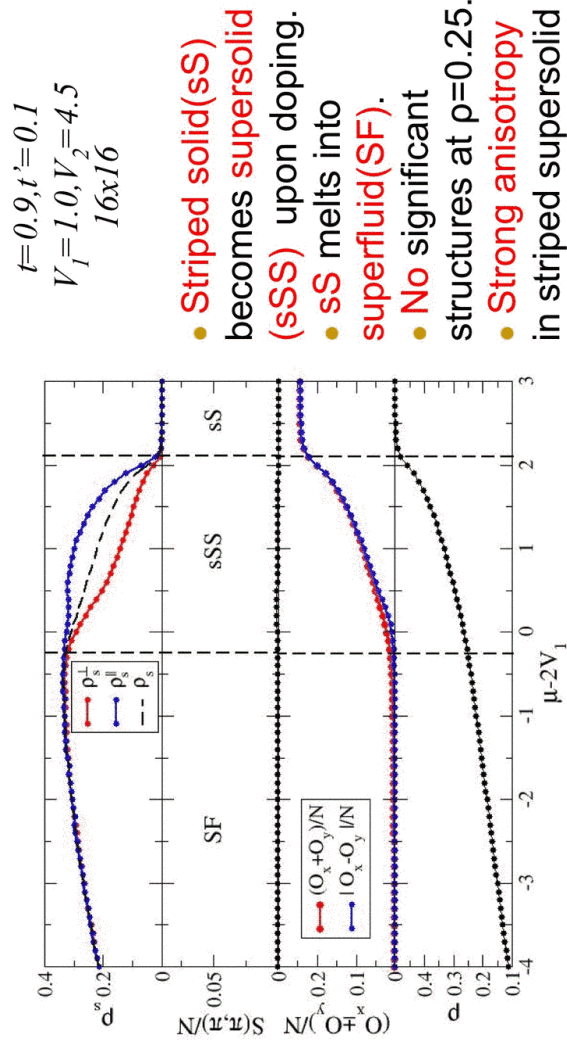
$$V_2 > 2V_1$$



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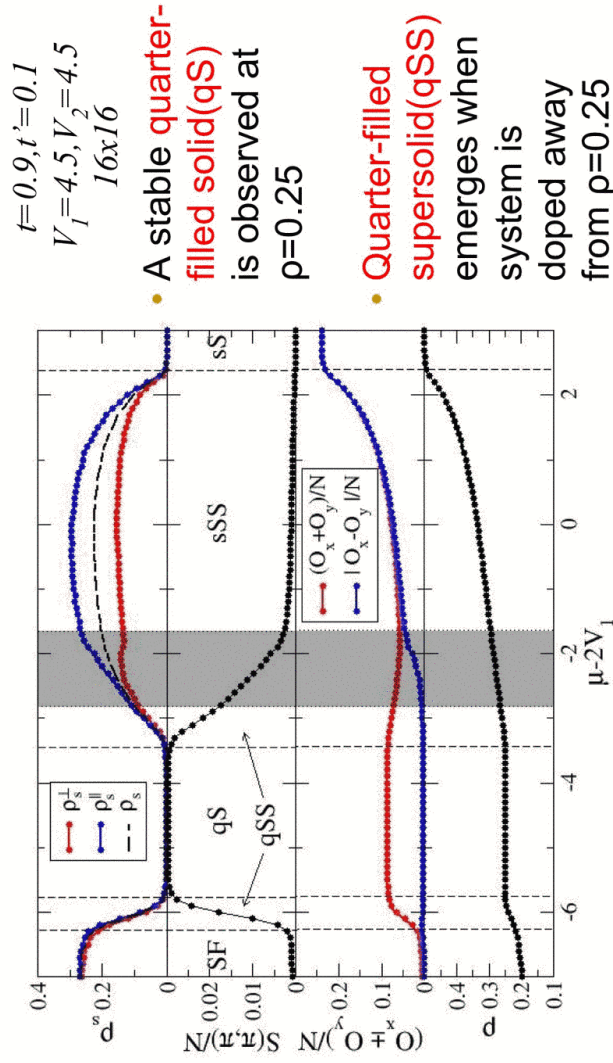
Phases at small V_1



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Phases at large V_1



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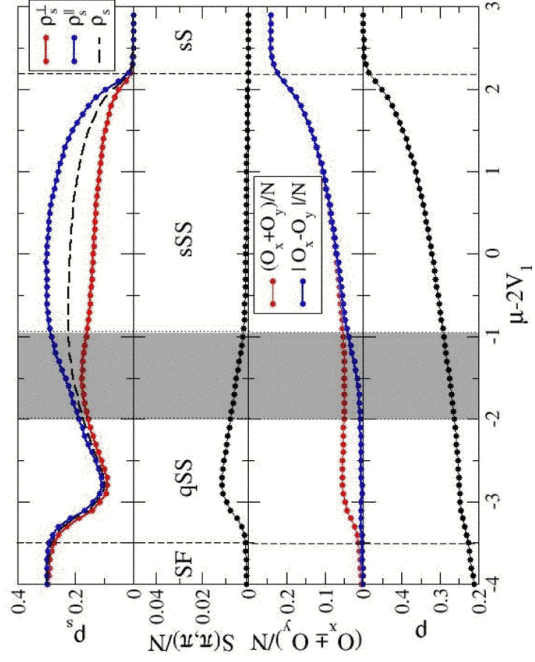


Phases at intermediate V_1

$$t=0.9, t'=0.1$$

$$V_1=3.0, V_2=4.5$$

$$16 \times 16$$



- Finite size scaling shows superfluidity vanishes near $\rho=0.25 \rightarrow$ **solid state**
- Grand canonical ensemble simulation, ρ is **not fixed**
- Canonical ensemble results

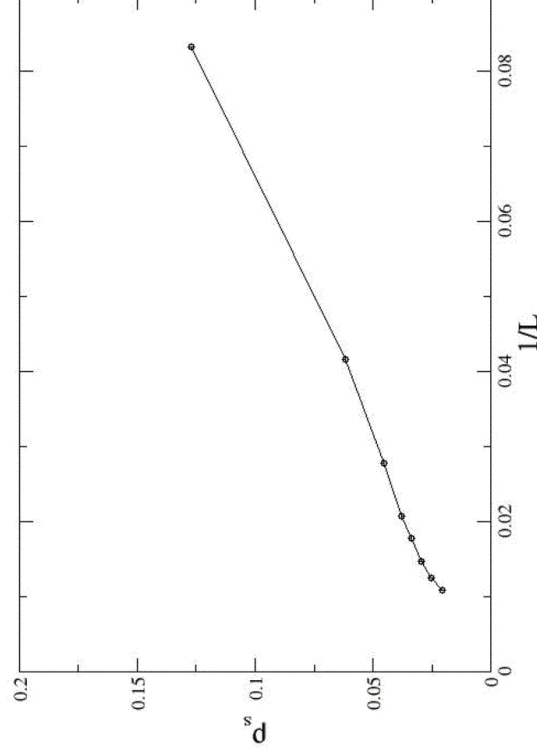
- A dip in superfluid density near $\rho=0.25$, but still finite.
- Is it possible to have supersolid at exactly $\rho=0.25$?



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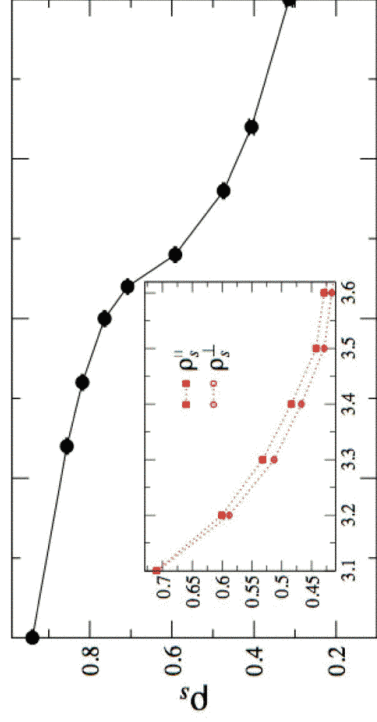
Finite Size Scaling at fixed $\rho=1/4$



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Triangular Lattice t - V_1 - V_2 model



Melko et al, PRB 74, 214517 (2006)



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Summary

- At the limit where V_1 is dominant, we observe a **stable checkerboard supersolid** when t' is introduced.
- When V_2 is dominant, but V_1 finite, we observe a **quarter-filled solid** and **corresponding supersolid** phase near $p=0.25$.
- We observe a transition/crossover from quarter-filled solid and striped solid via intermediate supersolid state(s).
- Suggests a **Defect Condensation** mechanism, but not the simple microscopic phase separation picture.
- Details of the transitions/crossovers? Finite temperature? Long-range dipolar interaction?



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