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07-29-03

Lecture 2

Yesterday: considered the geometry of  $T \oplus T^*$ .

We defined the notion of a generalized complex manifold: manifold  $M$  of dimension  $n$  together with  $E \subset (T \oplus T^*) \otimes \mathbb{C}$  so that

- $E \subset (T \oplus T^*) \otimes \mathbb{C}$  is maximally isotropic
- $E \cap \bar{E} = 0$
- sections of  $E$  are closed under the Courant bracket.

Basic examples:

- symplectic manifold:  $E$  is the transform of  $T \subset T \oplus T^*$  by  $i\omega$
- complex manifold  

$$E = T^{1,0} \oplus T^{*,0,1}$$
- transform either of the above by a closed B-field

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We would like to have an easy way to check the closeness of  $E$  w.r.t. Courant bracket. This is possible if we use some facts about pure spinors.

Pure spinors: We interpreted  $\Lambda^n T^*$  as spinors:  
 $(X + \zeta) \cdot \psi = i_X \psi + \zeta \wedge \psi$   
 Clifford multiplication

Suppose

$$(X+\zeta) \cdot \psi = 0 \Rightarrow (X+\zeta)^2 \cdot \psi = 0$$

$$\Rightarrow (X+\zeta, X+\zeta) = 0$$

Thus the annihilator of any spinor  $\psi \in \Lambda^* T^*$  is an isotropic subspace in  $T \oplus T^*$  (or  $(T \oplus T^*) \otimes \mathbb{C}$ ).

Def: A spinor  $\psi$  is called a pure spinor if its annihilator is maximally isotropic

Note: The property of being pure is invariant under the action of the Spin group.

Examples: •  $1 \in \Lambda^0 T^* \subset \Lambda^* T^*$

$$\Rightarrow i_X 1 + \zeta \wedge 1 = 0 \Leftrightarrow \zeta = 0$$

hence annihilator(1) =  $T \subset T \oplus T^*$   
i.e. 1 is a pure spinor

•  $e^B \cdot 1$  is pure

• If  $\dim M = m$ , then decomposable  $m$ -forms are pure spinors:

$$i_X (dx_1 \wedge \dots \wedge dx_m) + \zeta \wedge (dx_1 \wedge \dots \wedge dx_m) = 0$$

$$X = \sum_{i=m+1}^n a_i \frac{\partial}{\partial x_i} \quad \zeta = \sum_{i=1}^m b_i dx_i$$

•  $e^B \cdot (dx_1 \wedge \dots \wedge dx_m)$  - pure

Proposition There is a 1-to-1 correspondence between pure spinors (taken up to scalar multiplication) and maximally isotropic subspaces

$$\varrho \longrightarrow E_{\varrho} \subset T \oplus T^*$$

Proof: E. Cartan, C. Chevalley  $\square$

Fact:  $\dim(E_{\varrho} \cap E_{\psi}) > 0$  iff  $\langle \varrho, \psi \rangle = 0$   
Here  $\langle, \rangle$  is the 'Mukai pairing' on spinors.

If now  $M$  - generalized complex manifold then  $E \subset (T \oplus T^*) \otimes \mathbb{C}$  corresponds to some pure spinor  $\varrho$  and the Fact above implies

$$E \cap \bar{E} = 0 \iff \langle \varrho, \bar{\varrho} \rangle \neq 0$$

Note that if  $M$  is a generalized complex manifold, then  $\forall m \in M$   $\exists$  a 4-dim space of pure spinors  $\varrho_m \in \Lambda^0 T_m^*$  s.t.  
 $\text{Ann}(\varrho_m) = E_m$ .

These spaces  $L_m$  fit together into a line bundle  $L \rightarrow M$  - the canonical bundle of the generalized complex structure.

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Example: If  $E = T^{0,1} \oplus T^{1,0}$   
 (gives a usual complex structure)  
 then  $L = K_M$ .

Proposition: If  $\psi$  is a pure spinor, then  
 $E_\psi$  is closed under the Courant bracket  
 iff  $d\psi = (x + \zeta)\psi$  for some  $x, \zeta$ .  
 In particular if  $\psi$  is a  
 pure spinor and  $d\psi = 0$ , then  $E_\psi$  is  
 closed under the Courant bracket.

Def: A generalized Calabi-Yau manifold  
 is a manifold  $M^{2n}$  together with  
 a closed form  $\psi \in \Lambda^n T^* M$  s.t.  
 •  $\psi$  is a pure form  
 •  $\langle \psi, \bar{\psi} \rangle \neq 0$

Examples: (1) If  $M$  is an ordinary CY manifold  
 (here this means a complex manifold  
 $M$  together with a section  $\Omega \in H^0(K_M)$   
 trivializing  $K_M$ ),  
 $\Omega$  - top degree holomorphic form  $\Rightarrow$  decomposable  
 $\Rightarrow \Omega$  - pure spinor  
 $\langle \Omega, \bar{\Omega} \rangle = \pm \Omega \wedge \bar{\Omega} \neq 0$ .

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2. If  $(M, \omega)$  symplectic manifold then

$e = e^{i\omega}$  is a pure spinor

and

$$\langle e^{i\omega}, e^{-i\omega} \rangle = \langle e^{2i\omega}, 1 \rangle = \pm (2i\omega)^m \neq 0$$

3. A B-field transform of a symplectic manifold is also a generalized CY:

$$e^{B+i\omega} = e^B e \quad - \text{pure spinor.}$$

Comment: At a first glance it sounds strange that one can treat symplectic and CY manifolds on an equal footing.

Indeed:

- If  $(M, \omega)$  - compact symplectic  
 $\Rightarrow$  Lie algebra automorphism group  $= C^\infty(M)/\mathbb{R}$   
 is infinite dimensional.

- If  $(M, \Omega)$  - compact CY with  $b_1 = 0$   
 $\Rightarrow$   $M$  has no vector fields (holomorphic).

However if we add the action of B-fields we get similar pictures:

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If  $(M, \omega)$  symplectic and we consider the action of  $\text{Diff}(M) \times \Omega^2_{\text{closed}}$   $\Rightarrow$  on the level of Lie algebras we get

$$\begin{aligned} (\mathcal{L}_X + d\zeta) e^{i\omega} = 0 &\Rightarrow \mathcal{L}_X i\omega + d\zeta = 0 \\ &\Rightarrow \mathcal{L}_X \omega = 0 \\ &\quad d\zeta = 0 \end{aligned}$$

i.e. we get again  $C^\infty(M)/\mathbb{R}$ .

If  $(M, \Omega)$  - compact complex CY (assume also Kähler) . We get

$$\begin{aligned} (\mathcal{L}_X + d\zeta) \Omega = 0 &\Rightarrow \mathcal{L}_X \Omega = 0 \quad d\zeta \wedge \Omega = 0 \\ &\Rightarrow X = 0 \\ &\quad d\zeta \text{ is of type } (1,1) \\ &\quad \Rightarrow \text{by } \partial\bar{\partial} \text{ lemma} \\ &\quad d\zeta = \partial\bar{\partial} f \end{aligned}$$

Hence we get again  $C^\infty(M)/\mathbb{R}$ .

Caution: This gives Lie algebras which are isomorphic as vector spaces but have different Lie brackets: Poisson bracket in the symplectic case and trivial bracket in the CY case

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Question (Greg Moore): Can one incorporate quantum corrections on the symplectic side in order to make the Lie brackets match?

If  $M^{4k}$  is hyperkähler we have  
 $\rightarrow$  symplectic forms

$$\omega_1, \omega_2, \omega_3$$

$\omega_1$  - Kähler

$\omega_2 + i\omega_3$  - holo symplectic

Now  $B$ -field  $\rightarrow$   $e^{\omega_2 + i\omega_3}$   $\leftarrow$  symplectic  
 $e^{\omega_2 + i\omega_3}$  - generalized complex structure

The spinor

$$e = t^k e^{\frac{\omega_2 + i\omega_3}{t}}$$

defines a  $B$ -field transform of a symplectic manifold

Now  $e \xrightarrow{t \rightarrow 0} (\omega_2 + i\omega_3)^k$   
 $\uparrow$   
 pure spinor defining an ordinary CY structure.

§.

Thus: Generalized complex structures allow us in some cases to interpolate between complex and symplectic manifolds.

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As we mentioned before: If  $\psi = \psi_0 + \psi_2 + \dots$  is a pure spinor, then  $\psi_0 \neq 0$  and

$$\psi = \psi_0 e^{B+i\omega}$$

for some  $B, \omega$ .

The GCY condition is simply that  $d\psi = 0 \Rightarrow \psi_0 = \text{const.}$

However for a general generalized complex manifold this is not the case and we can see a type change for our pure spinor from point to point on the manifold.

Type change for pure spinors - even case

$$\begin{aligned} L &= \Lambda^k T^* \otimes \mathbb{C} &\longrightarrow & \Lambda^0 T^* \otimes \mathbb{C} \\ \psi & &\longrightarrow & \psi_0 \end{aligned}$$



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$\varphi_0$  defines a section of  $L^*$  which vanishes where the type changes.

In particular if  $c_1(L) \neq 0 \Rightarrow$  type change must occur

Examples: (1)  $M = \mathbb{R}^4 = \mathbb{C}^2$

$$\varphi = z_1 + dz_1 \wedge dz_2 = z_1 \exp\left(\frac{dz_1}{z_1} \wedge dz_2\right)$$

When  $z_1 \neq 0 \Rightarrow \varphi =$  B-field transform of a symplectic manifold

When  $z_1 = 0 \Rightarrow \varphi = dz_1 \wedge dz_2 =$  complex manifold

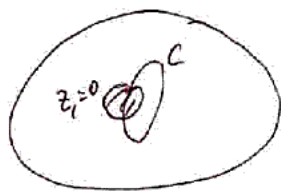
Note that this is a generalized complex manifold. Indeed

- $\varphi$  - pure  $\Rightarrow E_\varphi$  closed under the Courant bracket if  $d(\sharp\varphi) = 0$  for some function.

- $\langle \varphi, \bar{\varphi} \rangle = \langle z_1 + dz_1 \wedge dz_2, \bar{z}_1 + d\bar{z}_1 \wedge d\bar{z}_2 \rangle$   
 $= dz_1 dz_2 d\bar{z}_1 d\bar{z}_2 \neq 0$

- $d(z_1 \varphi) = 0 \Rightarrow E_\varphi$  - integrable

(2) If  $M$  is a complex surface with  $\beta$  a holomorphic section of  $K^{-1} \Rightarrow \beta^{-1}$  is a holomorphic 2-form outside  $C = \text{zero}(\beta)$ .



$$z_1 + z_1 \beta^{-1}$$

Topological conditions: Assume that we have  
 a generalized almost complex manifold:  
 $M$  together with

- $E \subset (T \oplus T^*) \otimes \mathbb{C}$
- $E$  - isotropic
- $E \cap \bar{E} = 0$ .

In other words we have an almost  
 complex structure

$$J: T \oplus T^* \rightarrow T \oplus T^*$$

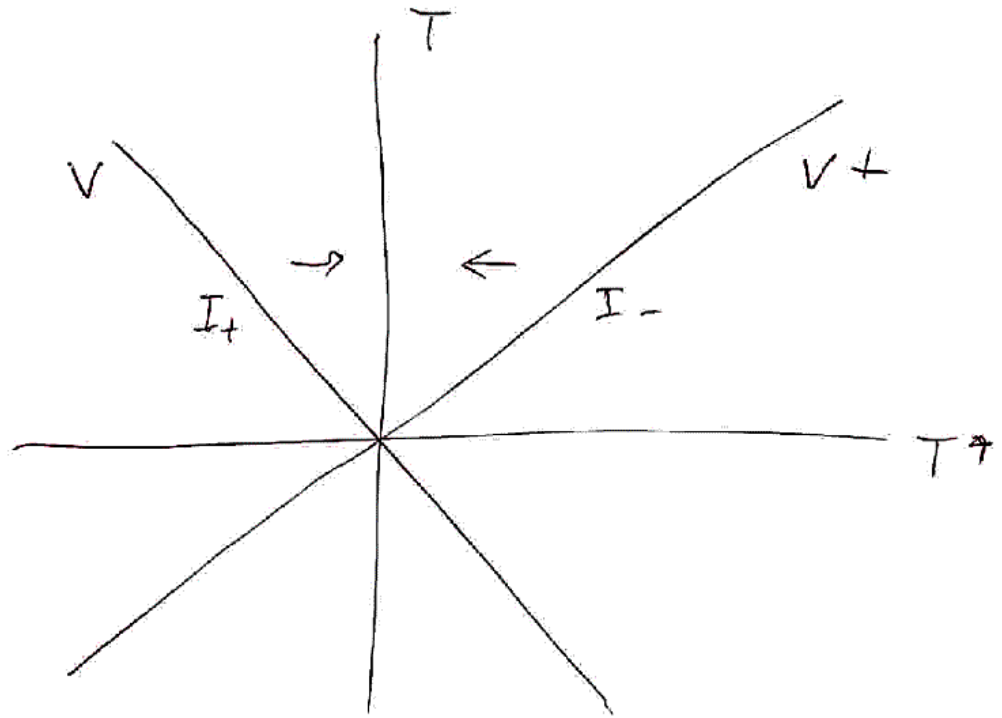
and  $J$  is compatible with the  
 indefinite metric on  $T \oplus T^*$ . In other  
 words  $M +$  (generalized almost  
 complex structure)

is the same as

$$M + \left( \begin{array}{l} \text{reduction of the structure} \\ \text{group } O(m, m) \text{ of } M \\ \text{to } U(m, m) \end{array} \right)$$

Now  $U(m, m) \sim U(m) \times U(m)$   
 homotopy  
 equivalent

This gives the following picture:



where  
with

$$V \oplus V^+ = T \oplus T^*$$

metric  $|_V =$  positive definite

metric  $|_{V^+} =$  negative definite

Also projections  $V \rightarrow T$   
 $V^+ \rightarrow T$

are isomorphisms and we get

two almost complex structures  $(V, I_+)$   
 $(V^+, I_-)$

In particular: if  $M$  - generalized almost  
complex  $\Rightarrow T$  has two natural  
almost complex structures.

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Thus: Spheres can not be generalized  
almost complex if  $\dim > 2$ .

Now going back to  $I_+, I_-$   
note that if  $L$  - canonical bundle on  $M$ ,  
then

$$-2c_1(L) = c_1^+ + c_1^-$$

Note: Symplectic case:  $c_1^+, c_1^-$  - opposite  
Complex case:  $c_1^+, c_1^-$  - same

### Generalized complex submanifolds

If  $Y \subset M$  - submanifold

$$TY \subset TM|_Y$$

$$N^*Y = (TY)^{\text{ann}} \subset T^*M|_Y$$

$$\Rightarrow TY \oplus N^*Y \subset T \oplus T^*|_Y$$

Def: A generalized complex manifold is  
a submanifold  $Y \subset M$  in a generalized  
complex manifold  $p$  s.t.  $TY \oplus N^*Y$   
is preserved by  $J$ .

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Examples: (1)  $M$ -complex

$$J = \begin{pmatrix} I_T & 0 \\ 0 & -I_{T^*} \end{pmatrix}$$

$\Rightarrow Y \subset M$  is a generalized complex submanifold  $\Leftrightarrow Y$  is a complex manifold

(2) B-field transform of a complex manifold.

$$\begin{pmatrix} 1 & 0 \\ -B & 1 \end{pmatrix} \begin{pmatrix} I_T & 0 \\ 0 & -I_{T^*} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ B & 0 \end{pmatrix} =$$

$$= \begin{pmatrix} I_T & 0 \\ -(BI + IB) & -I_{T^*} \end{pmatrix}$$

$\Rightarrow Y \subset M$  is a generalized complex manifold iff:

- $Y \subset M$  - complex submanifold
- $B|_Y$  is of type  $(1,1)$

(because

$$B(Ix_1, x_2) + B(x_1, Ix_2) = 0)$$

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(3.) If  $M$  - symplectic,  $Y \subset M$  - submanifold

$$J = \begin{pmatrix} 0 & -\omega^{-1} \\ \omega & 0 \end{pmatrix}$$

$\Rightarrow \omega^{-1}\zeta \in TY \Rightarrow Y$  - coisotropic

$\omega \chi \in (TY)^{\text{ann}} \Rightarrow Y$  - isotropic

$\Rightarrow Y$  - generalized complex submanifold

$\Leftrightarrow Y$  - Lagrangian

(4) B-field deformation of symplectic manifold

$$J \begin{pmatrix} x \\ \zeta \end{pmatrix} = \begin{pmatrix} -\omega^{-1}(\zeta + Bx) \\ (\omega + B\omega^{-1}B)x + B\omega^{-1}\zeta \end{pmatrix}$$

$\Rightarrow Y \subset M$  - generalized complex  $\Leftrightarrow$

•  $Y$  - coisotropic

• B-field is basic i.e. is a pull back from the leaf space of the null foliation on  $Y$  corresponding to  $\omega$ .

•  $(B + i\omega)$  is the  $(2,0)$  form of a complex structure on the leaf space

In other words:  $Y$  is equipped with a transverse holomorphic symplectic structure for the same foliation of  $W$ .

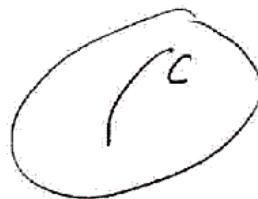
Remark: • This is almost the same as Kapustin-Oslov generalized A-branes only there one requires

$$B = F \nabla$$

for some connection  $\nabla$  on a line bundle on  $Y$ ,

- Points are not always generalized complex submanifolds.

In the example with the Poisson structure on a surface



- => points on  $C$  are generalized complex submanifolds but points not on  $C$  are not generalized complex submanifolds.