

Anton KapustinLecture 4Last time we looked at:

$X = \mathbb{C}$

$W = \mathbb{Z}^n$

Consider the brane

$$\mathcal{O}_X \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{G} \end{array} \mathcal{O}_X$$

$F = z$

$G = z^{n-1}$

We saw that

$$\text{End} \left(\mathcal{O}_X \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{G} \end{array} \mathcal{O}_X \right) = \langle 1, \theta \mid \theta^2 = 0 \rangle$$

for all $n \geq 2$.

We got the algebra $\text{End} \left(\mathcal{O}_X \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{G} \end{array} \mathcal{O}_X \right)$
 as the cohomology of a dg-algebra
 with differential

$$\mathcal{D} = [\mathcal{Q}_{\text{boundary}}, 0]$$

Here ~~$\mathcal{Q}_{\text{boundary}}$~~ $\mathcal{Q}_{\text{boundary}} = \begin{pmatrix} 0 & z^{n-1} \\ z & 0 \end{pmatrix}$

In particular we have an A_∞-structure
 on $\text{End} \left(\mathcal{O}_X \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{G} \end{array} \mathcal{O}_X \right)$ which carries
 additional physical information.

Consider $n=3$. In this case $m_2(\theta, \theta) = \theta^2 = 0$
 What is $m_3(\theta, \theta, \theta) = ?$

In general if we have a dg-algebra \mathcal{A} and the A_∞ -structure on the cohomology is such that

$$m_2 = m_3 = \dots = m_k = 0$$

$\Rightarrow m_{k+1}$ is ~~unambiguously~~ unambiguously defined.

In particular if $m_2 = 0 \Rightarrow$

$$m_3([a], [b], [c]) = [dc - (-1)^{|a|} a\beta]$$

where

$$\alpha : ab = \partial a \quad ([a] \cdot [b] = 0)$$

$$\beta : bc = \partial b \quad ([b] \cdot [c] = 0)$$

In our case:

$$\theta = \begin{pmatrix} 0 & 1 \\ -2 & 0 \end{pmatrix}$$

$$\theta^2 = \begin{pmatrix} -2 & 0 \\ 0 & -2 \end{pmatrix} = -2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Hence $m_3(\theta, \theta, \theta) = \left[\begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} \right] = -1.$

In general (for arbitrary n) we expect that

$$m_k(\theta, \dots, \theta) = 0$$

$$m_n(\theta, \dots, \theta) = -1$$

for all $k < n$
(This observation, as well as the preceding computation of m_3 , is due to Kontsevich)

This suggests that the theory has a degree n superpotential in the open string moduli space (= moduli space of deformations of $\mathcal{O}_X \rightarrow \mathcal{O}_X$ in our original dg category)

Recall that a brane in the \mathcal{B} -model was

$$E_0 \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{G} \end{array} E_1$$

with $F \circ G = w \cdot \text{id}_{E_1} - w_0 \cdot \text{id}_{E_1}$

$$G \circ F = w \cdot \text{id}_{E_0} - w_0 \cdot \text{id}_{E_0}$$

where w_0 is a constant.

Now our notes:

- branes with different w_0 's do not interact (no ^{non-zero} morphisms between them)
- the only non-trivial branes correspond to w_0 which is a critical value of w .

This was analyzed by Hori in some cases: linear models of D-branes ---

In particular: Category of branes breaks into pieces = branes corresponding to the critical values w_0 of w .

Orlov made this mathematically precise:

If $w_0 \in \mathbb{C}$ - any value then the fiber has junctions

$$\Gamma(\mathcal{O}(X_{w_0})) = \mathbb{C}[x_1, \dots, x_n] / (w(x) - w_0)$$

Now consider

$$D_{w_0} := D^b(X_{w_0}) / \text{Perf}(X_{w_0})$$

\Rightarrow if $w_0 \neq$ critical value $\Rightarrow X_{w_0}$ - smooth
and hence $D_{w_0} = 0$

~~if $w_0 =$ critical value \Rightarrow Orlov showed~~

$D_{w_0} \cong$ category of branes for value w_0

Hence $D(\text{B-branes}) = \coprod_{\substack{w_0\text{-critical} \\ \text{for } w}} D_{w_0}$

Note: One can show that for $X = \mathbb{C}$
 $w = z^n$

then all branes are of the form

$$\oplus \left(\text{of } \begin{array}{ccc} \mathcal{O}_X & \xrightarrow{F} & \mathcal{O}_X \\ & \xleftarrow{G} & \end{array} \right. \begin{array}{l} F = z^k \\ G = z^{n-k} \\ k = 1, \dots, n-1 \end{array} \left. \right)$$

This follows from Orlov's theorem and the fact that the singular fibers of w_0 are 0-dimensional and have simple 'singular sheaves' e.g.

$$\mathbb{C}[x]/x^i \quad i \in \mathbb{N}$$

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Example: Let $X = \mathbb{C}^2$

$$W = x^n + xy^2 \quad (\mathcal{D}\text{-type minimal model})$$

what are the branes here?

Note: If we have a general X and $W = W_1 + W_2$ and a brane with differential

$$D = \begin{pmatrix} 0 & G \\ F & 0 \end{pmatrix}$$

we can try to break D into D_1, D_2 s.t.

$$D_1^2 = W_1$$

$$D_2^2 = W_2$$

for this write

$$D = \zeta_1 D_1 \otimes 1 + 1 \otimes D_2 \zeta_2$$

$$\zeta_1^2 = 1 \quad \zeta_2^2 = 1 \quad \zeta_1 \zeta_2 + \zeta_2 \zeta_1 = 0$$

$$\Rightarrow D^2 = W.$$

\Rightarrow can apply this to $W_1 = x^n$ on \mathbb{C}
 $W_2 = xy^2$ on \mathbb{C}^2

in order to compute the branes



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Recall: Vafa's computation of the one point function in closed string theory:



$$f \in \mathbb{C}[\vec{x}] / \partial_i W \quad \mapsto \quad \langle f \rangle = \text{Res} \frac{f}{\partial_1 W \cdots \partial_n W}$$

where

$$\begin{aligned} \text{Res} \frac{f}{\partial_1 W \cdots \partial_n W} &= \\ &= \int_{\Gamma} \frac{f dx_1 \cdots dx_n}{\partial_1 W \cdots \partial_n W} \end{aligned}$$

Γ = Lagrangian cycle

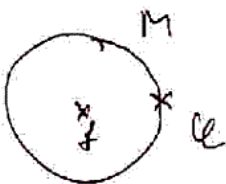
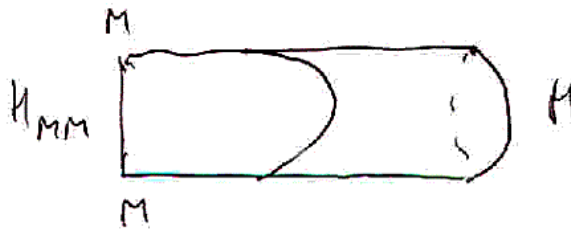
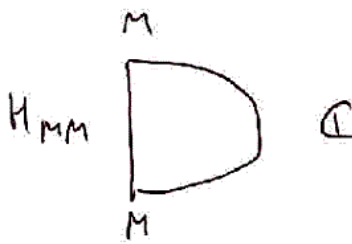
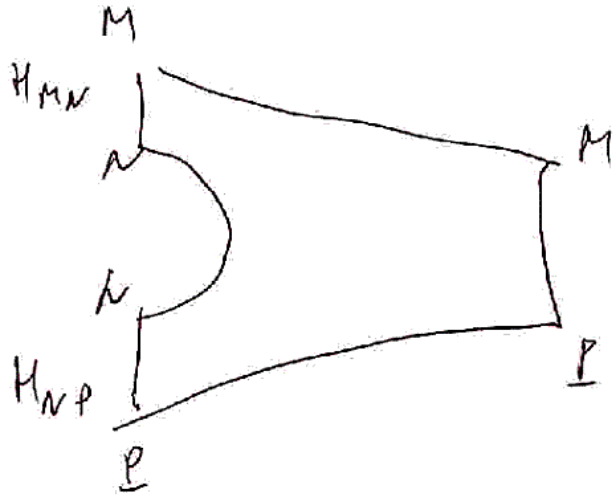
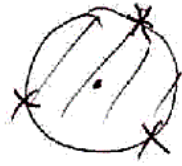
$$\begin{aligned} |\partial_i W| &= \varepsilon_i \\ \varepsilon_i &= \text{small} \end{aligned}$$

We get a bilinear form

$$\langle f, g \rangle := \langle fg \rangle$$

Note: It is non-trivial that this form is non-degenerate.

We would like now to do the same for ~~closed~~ open strings:



$$D = \begin{pmatrix} 0 & G \\ F & 0 \end{pmatrix}$$

$$D = z_i \Gamma_i + \dots$$

$$\{\Gamma_i, \Gamma_j\} = Q_{ij} \quad - \text{Clifford algebra relations}$$

$$\Rightarrow \langle f | \mathcal{L} \rangle = \text{res} \frac{f \text{str} ((\partial \bar{\partial})^n \mathcal{L})}{z_1 w z_2 w \dots z_n w}$$

Here $(\not{D})^n \sim dx_1 \dots dx_n$

Since this expectation value was defined in terms of the path integral it will be good if we can independently check that the formula makes sense and has the right properties:

$$\text{Res}_{\partial W \sim \partial nW} \frac{\int \text{str}((\not{D})^n \psi)}{\partial W \sim \partial nW} \text{ is well defined}$$

Indeed if $\psi = [D, \alpha] \Rightarrow \text{OK}$
 Since

$$D^2 = W$$

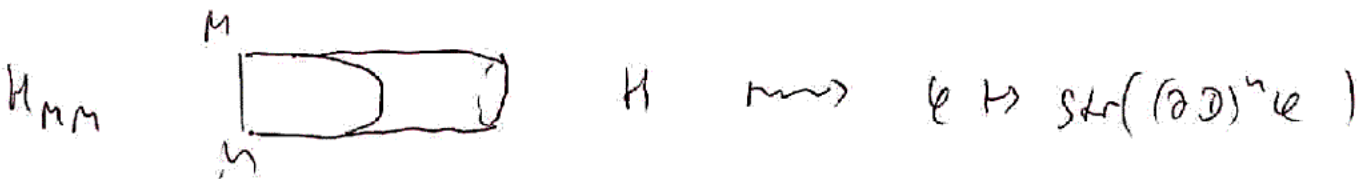
$$dD \cdot D + D \cdot dD = \not{d}W$$

$$\langle \int \psi_1 \psi_2 \rangle \stackrel{?}{=} \pm \langle \int \psi_2 \psi_1 \rangle$$

$$[D, \psi_1] = 0$$

$$[dD, \psi_1] + [D, d\psi_1] = 0$$

We can use the residue formula to read off the boundary bulk maps:

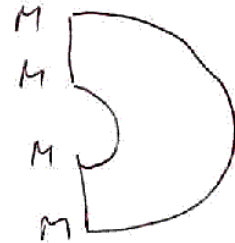


Again we will have the consistency conditions⁹ for all this data that need to be checked mathematically:

$$\langle \psi_1, \psi_2 \rangle = \text{res} \frac{\text{str}((\partial D)^n \psi_1 \psi_2)}{\partial_1 W \dots \partial_n W}$$

Should be non-degenerate

(this is just the pairing in the endomorphism algebra)



Cardy conditions

~~XXXXXXXXXX~~

If $\alpha \in H_{MM}$ $\beta \in H_{NN}$ \Rightarrow should get $T_{\alpha\beta} : H_{MN} \rightarrow H_{MN}$ and the Cardy condition expresses $T_{\alpha\beta}$ in terms of inner product of bulk states corresponding to α, β .
 For example, In the case $\alpha=1$, $\beta=1$ then should get

$$\chi(H_{MN}) = \text{res} \frac{\text{str}((\partial D_M)^m) \text{str}((\partial D_N)^n)}{\partial_1 W \dots \partial_n W}$$

This can be checked in examples.



We would like to understand how the closed strings can be identified with the Hochschild cohomology of the category of branes in this case.

Start with a brane $E_0 \xrightleftharpoons[G]{F} E_1$
 $F \circ G = G \circ F = W$

Recall: a dga (A, d) : $d(a_1 a_2) = da_1 a_2 \pm a_1 da_2$
 $d^2 = 0$
 a module over (A, d) : $M, \mathcal{D}_M: M \rightarrow M$
 $s.t. \mathcal{D}_M^2 = 0$
 $\mathcal{D}_M(a m) = da \cdot m \pm a \mathcal{D}_M(m)$

Similarly we have: cdg algebras
 $(A, d: A \rightarrow A, B) \times \text{Leibnitz}$
 $d^2 = [B, \cdot]$
 cdg module
 $(M, \mathcal{D}_M) \times \text{Leibnitz}$
 $\mathcal{D}_M^2 = B \cdot$

Note: A cdg algebra is not ~~a module~~
 a CDG module over itself in general.

Now consider $(A = \mathbb{C}[x_1, \dots, x_n], d=0, W)$
 as a cdg algebra

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→ the branes are just the derived category of cdg modules over this cdg algebra.

The HH of $\mathcal{D}^b(\text{cdg mod } (A, 0, w))$ is simply the HH of $(A, 0, w)$

$$\begin{aligned} \text{Explicitly } (A, 0, w) \otimes (A, 0, w)^{\text{op}} \\ \parallel \\ (A, 0, w) \otimes (A, 0, -w) \\ \parallel \\ (A \otimes A, 0, w \otimes (-w)) \end{aligned}$$

$$\Rightarrow \text{HH}^*(A, 0, w) := \text{Ext}_{(A, 0, w) \otimes (A, 0, w)^{\text{op}}}^0((A, 0, w), (A, 0, w))$$

Note: $(A, 0, w)$ is now a cdg module over $(A, 0, w) \otimes (A, 0, w)^{\text{op}}$ since $d^2a = wa - aw = 0$.

To compute this HH look e.g. at

$$\begin{aligned} X &= \mathbb{C} \\ w &= w(x) \end{aligned}$$

Then can look at $\mathbb{C}^2(x, y)$
 $w(x, y) = w(x) - w(y)$

$$\mathcal{D} = \begin{pmatrix} 0 & \frac{w(x)-w(y)}{x-y} \\ x-y & \end{pmatrix}$$

← this deformation is resolution of the diagonal in \mathbb{C}^2 has been ~~explained~~ found by Lev Rozansky.

If now we look at a module

⇒ have

$$\begin{pmatrix} d(x,y) & 0 \\ 0 & d(x,y) \end{pmatrix}$$

$d \in \mathbb{C}[x,y]$ and

$$d(x,y) \sim d(x,y) + \star (x-y) + \star \frac{w(x)-w(y)}{x-y}$$

⇒ get the Jacobi ring of $w(x)$.

~~(this computation was explained to us by is)~~

Now we can recover Vafa's formula

$$\text{res} \frac{\text{str}((\partial \mathcal{D})^2 \epsilon)}{\partial_x w \partial_y w} = \text{res} \frac{d(x,x)}{\partial_x w}$$