



# A Novel Coarse-Grained Force Field for Unbiased Protein-Membrane Simulations



Zun-Jing Wang,  
Tristan Bereau,  
& Markus Deserno



Department of Physics  
Carnegie Mellon University

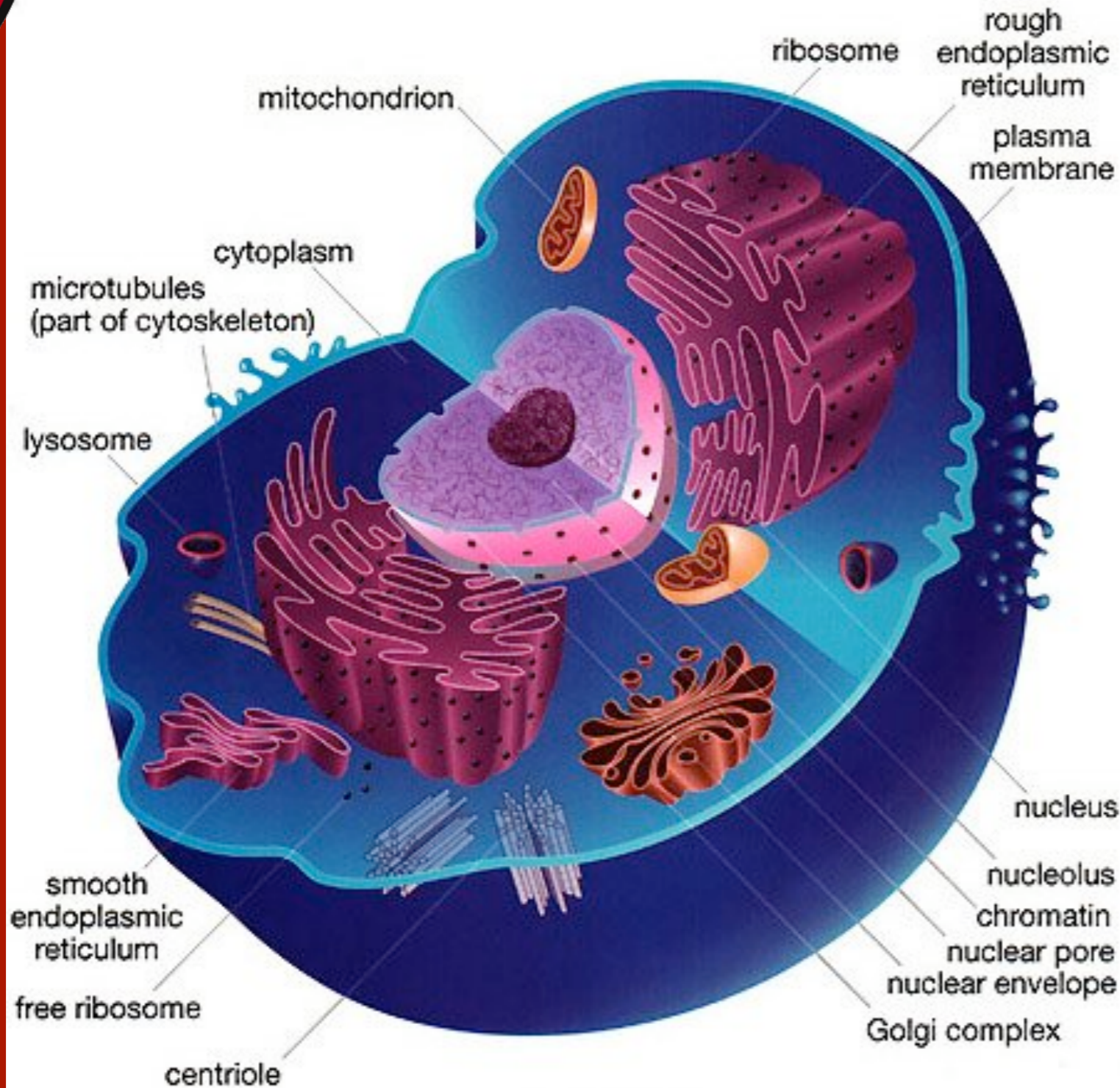
KITP program on  
“Modeling Soft Matter: Linking Multiple Length and Time Scales”

June 6, 2012, Santa Barbara

Carnegie Mellon



# Introduction



Membranes are crucial for the functioning of cells.

Proteins do the work. They also help shaping membranes.

Things often happen on large length- and time-scales.

Protein structure not always known and not always fixed.

Are there any generic principles at work?

<http://www.animalport.com/img/Animal-Cell.jpg>



# Introduction

We would like to look into this, using tools of the following ilk:

chemically realistic but coarse-grained  
modeling of both membranes and proteins

no bias on secondary structure

a modicum of transferability

implicit solvent



# Introduction

We think we have accomplished a little bit of that tall order. Let me thus talk about:

1. Our membrane model
2. Our peptide model
3. Our method of cross-parametrizing them
4. Some examples of what we can do



# 1. Our membrane model

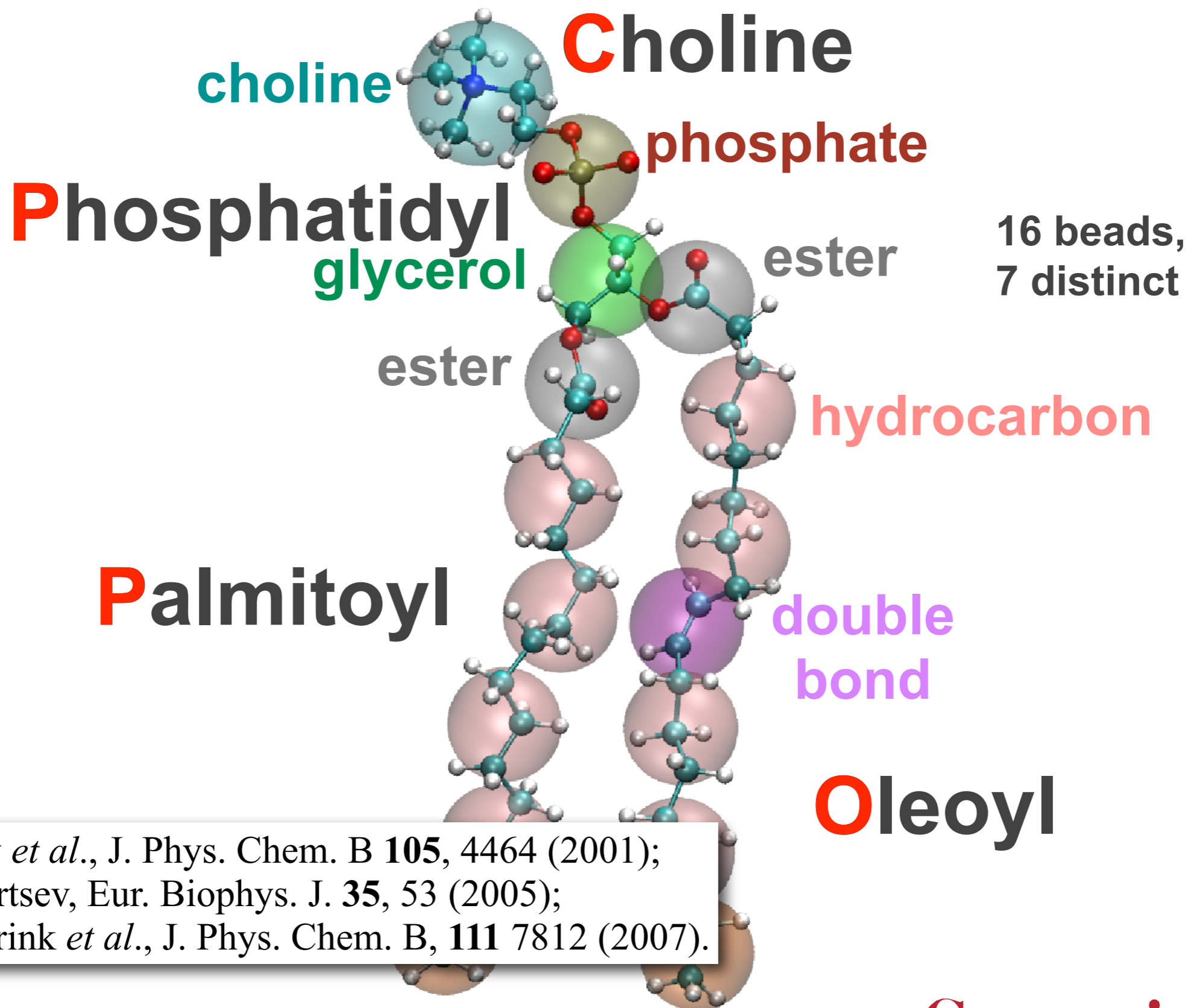


Z.-J. Wang and M. Deserno, *J. Phys. Chem. B* **114**, 11207 (2010).

Z.-J. Wang and M. Deserno, *New J. Phys.* **12**, 095004 (2010).



# Mapping



J. Shelley *et al.*, J. Phys. Chem. B **105**, 4464 (2001);  
A. Lyubartsev, Eur. Biophys. J. **35**, 53 (2005);  
S.-J. Marrink *et al.*, J. Phys. Chem. B, **111** 7812 (2007).



# Construction principles:

Structure-based coarse-graining  
(for most non-bonded interactions)

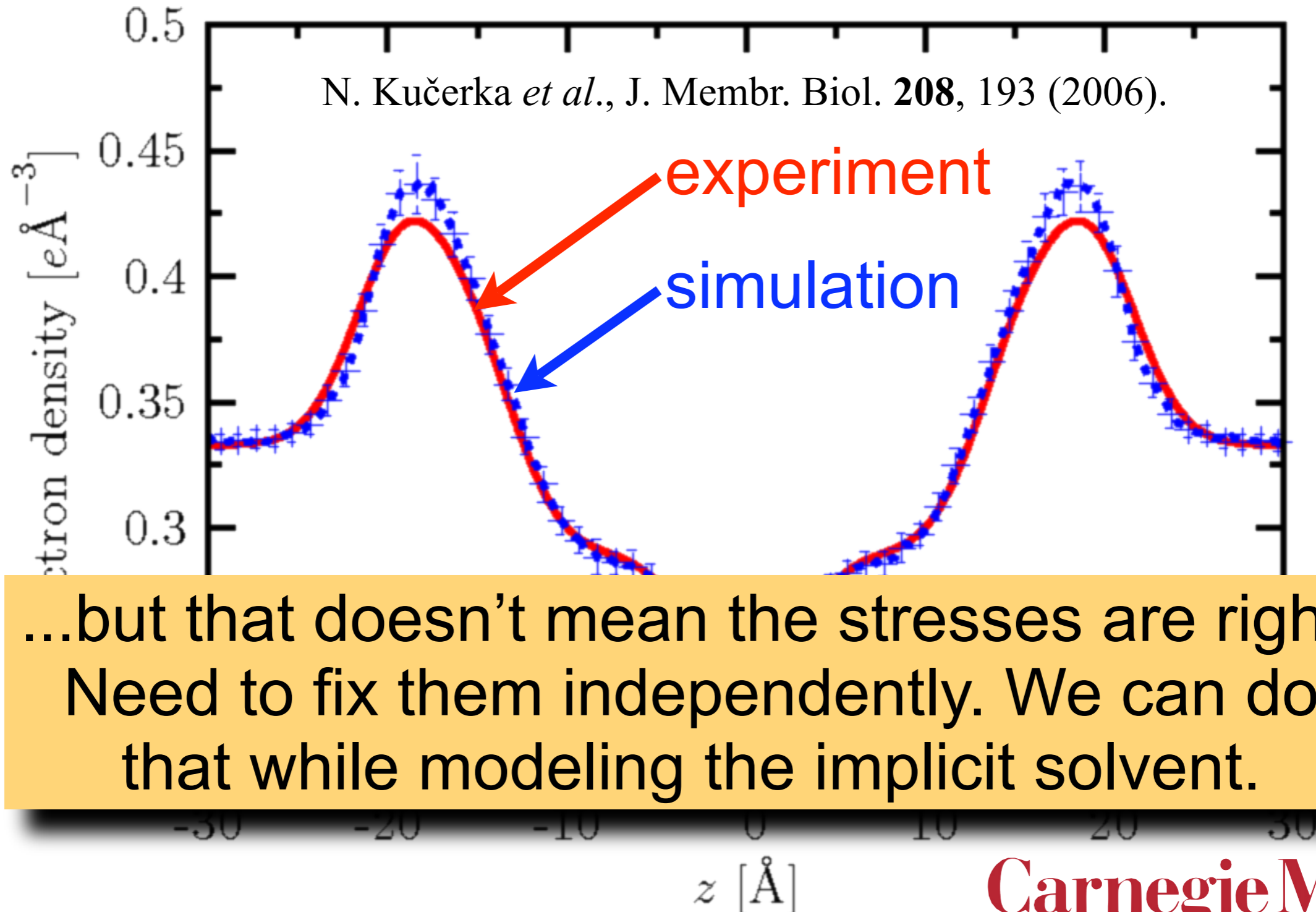
density and pressure profiles  
(from experiment and simulation)

area per lipid  
(from experiment)



# Area per lipid

The electron density at an area-per-lipid fixed to the correct value is reproduced very well!



...but that doesn't mean the stresses are right. Need to fix them independently. We can do that while modeling the implicit solvent.





# Strategy:

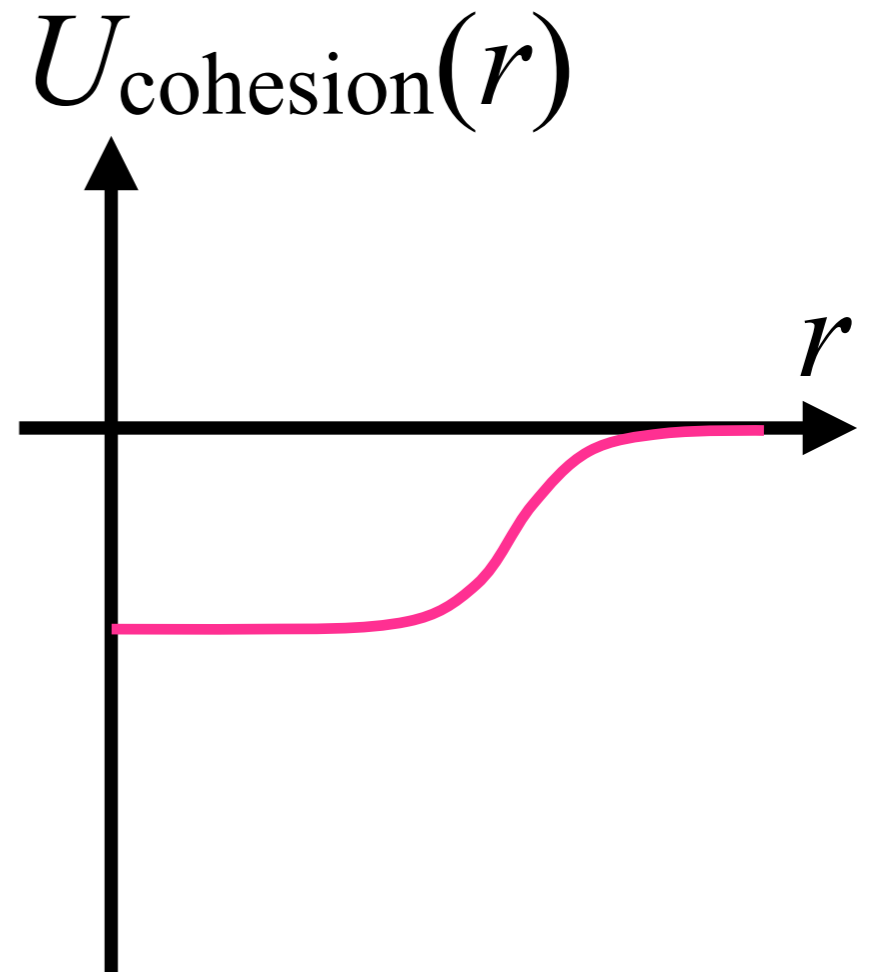
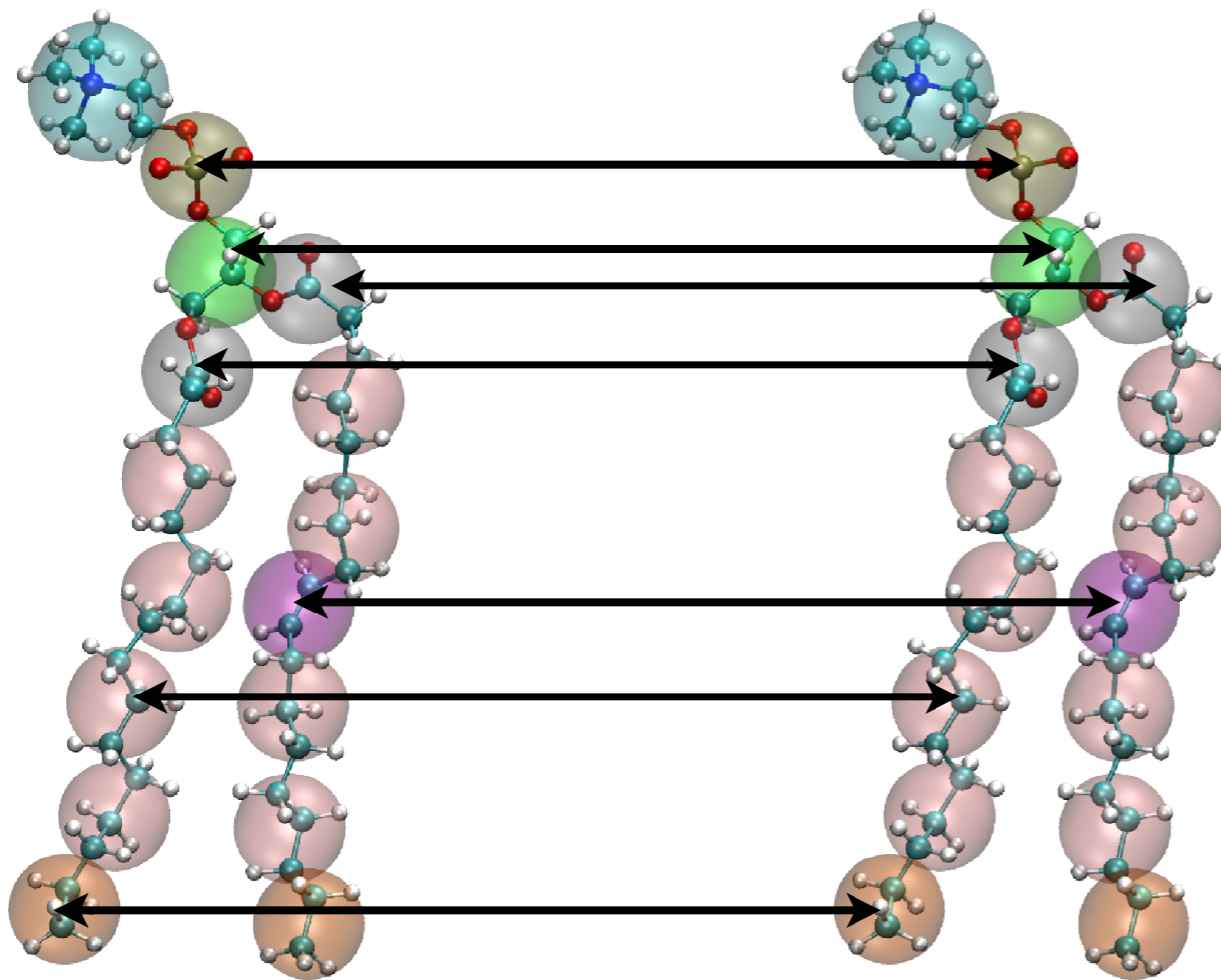
Use aa simulations at the correct area to infer structure

Introduce phenomenological cohesion in CG model to drive aggregation

Choose value of that cohesion to fix the zero-tension area per lipid.



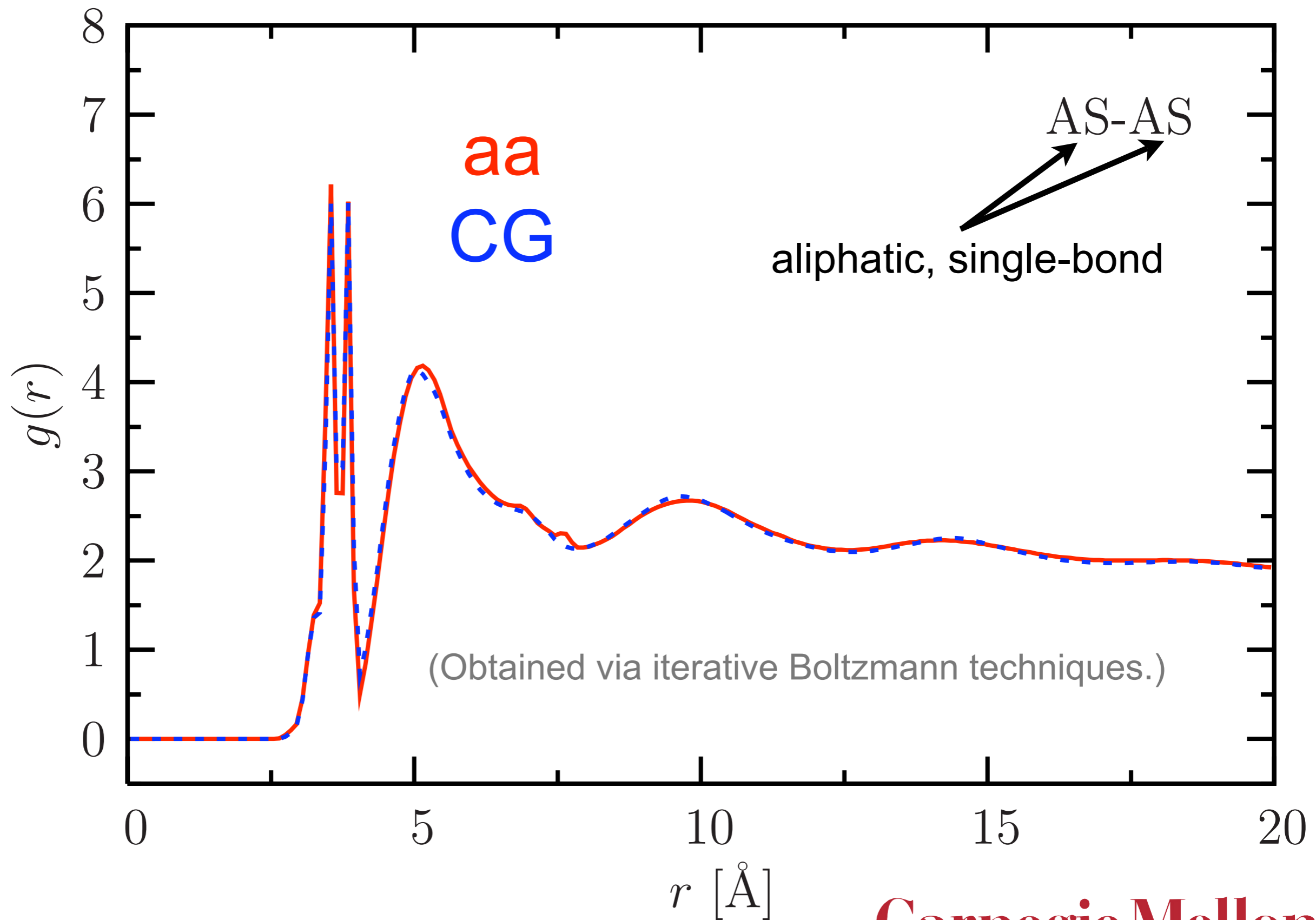
# Strategy:



Range and amplitude can vary between different bead types.  
This is optimized to reproduce the pressure profile.

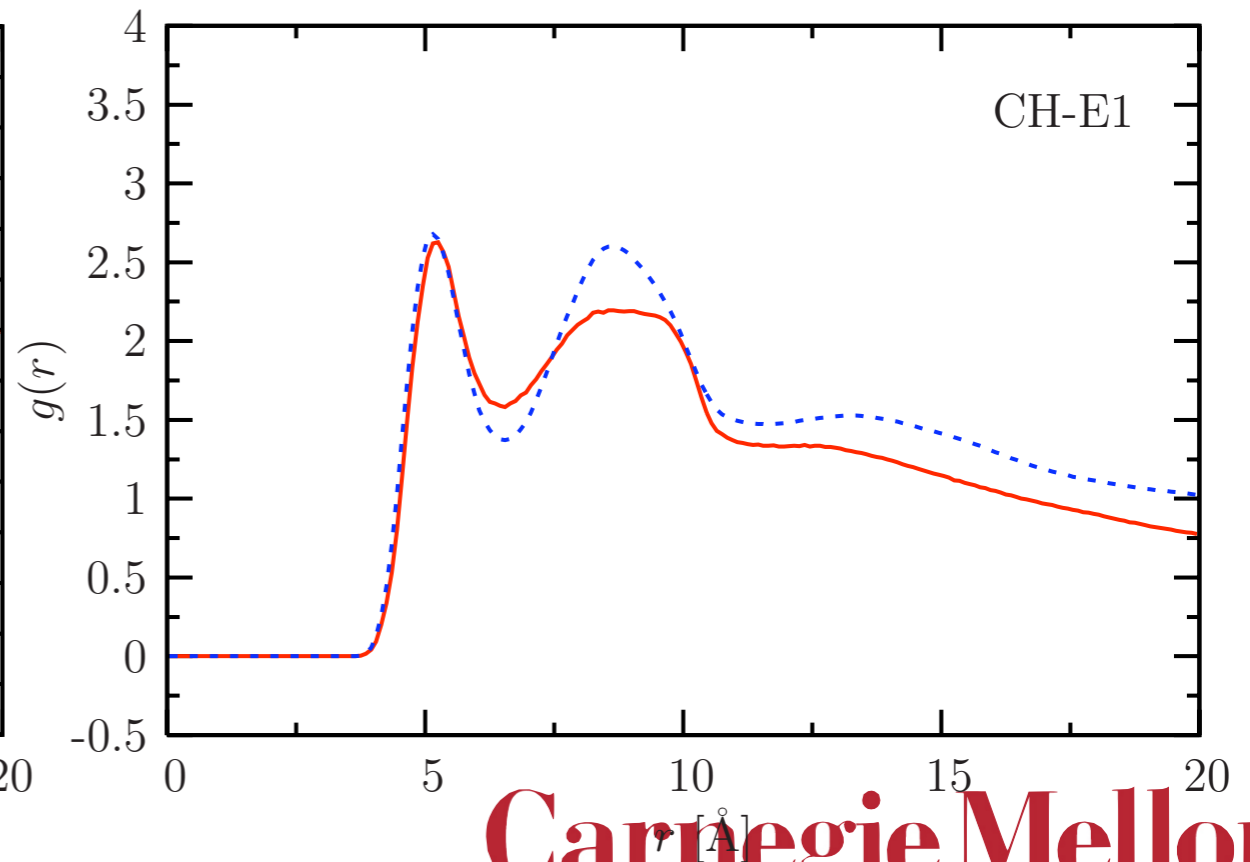
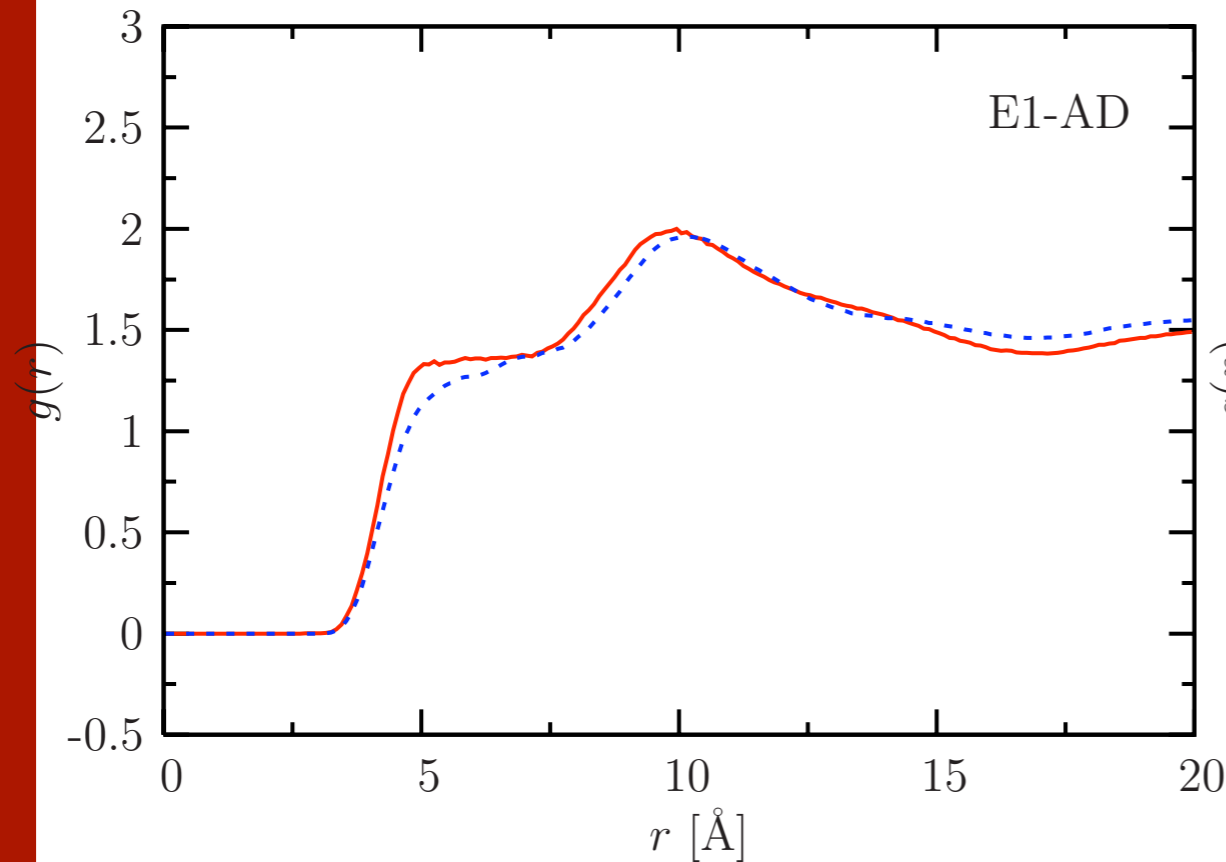
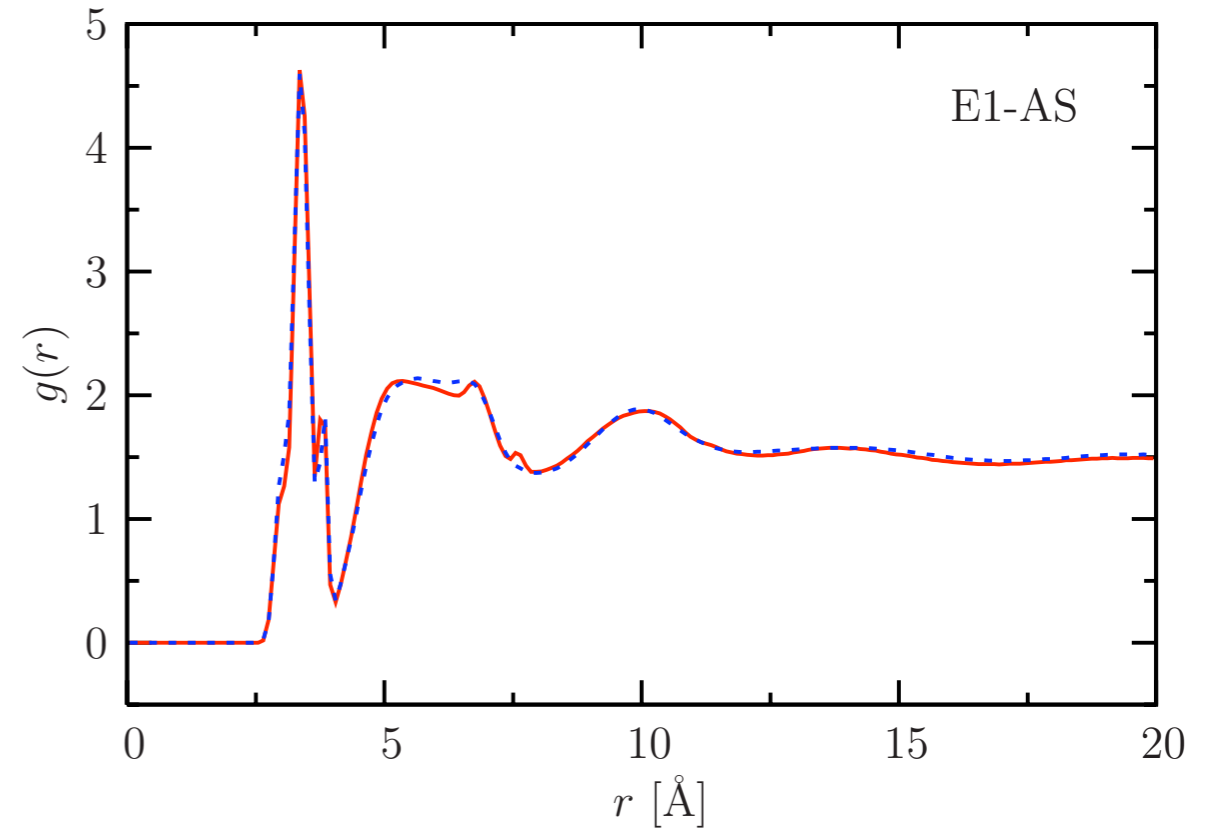
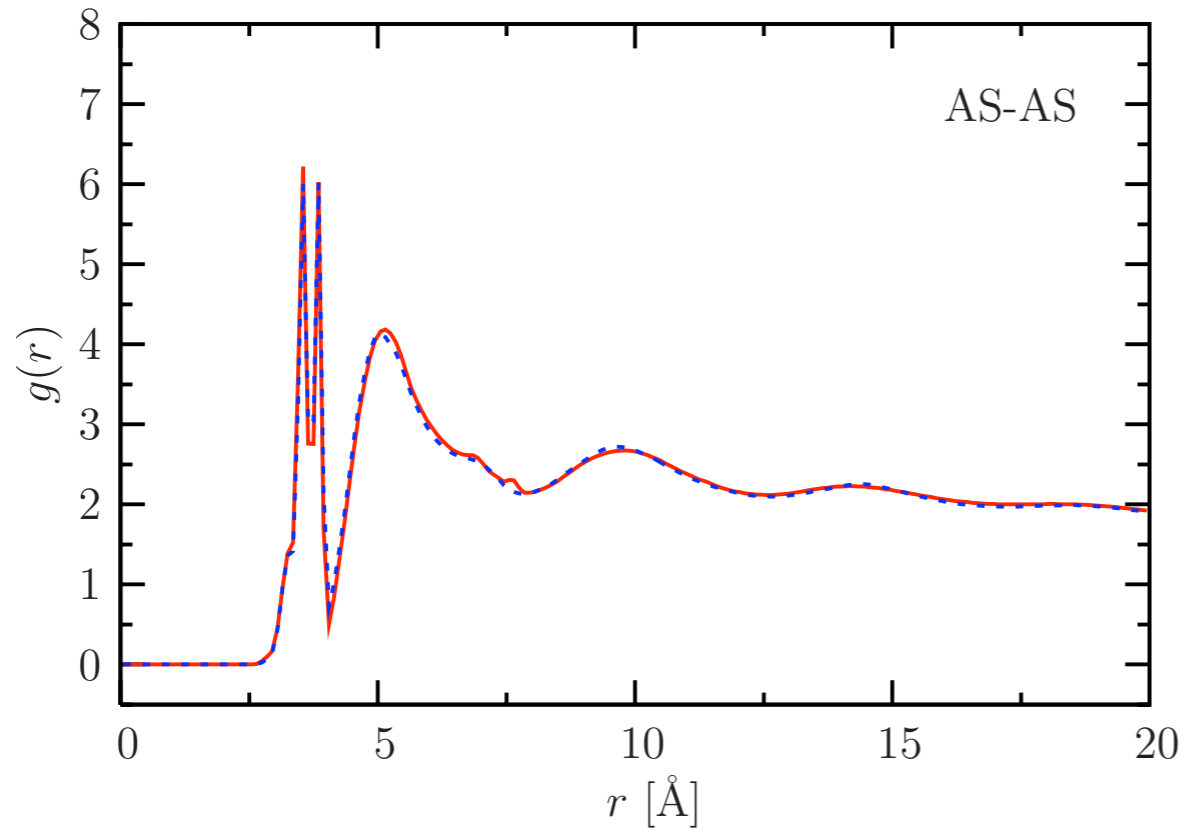


# Example for structure



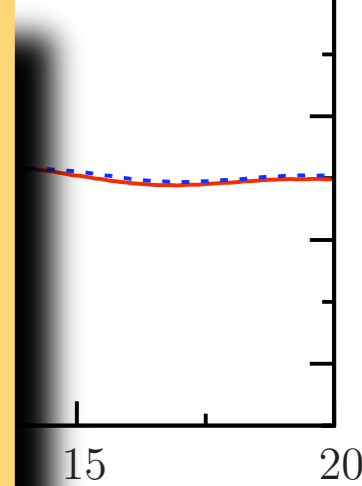
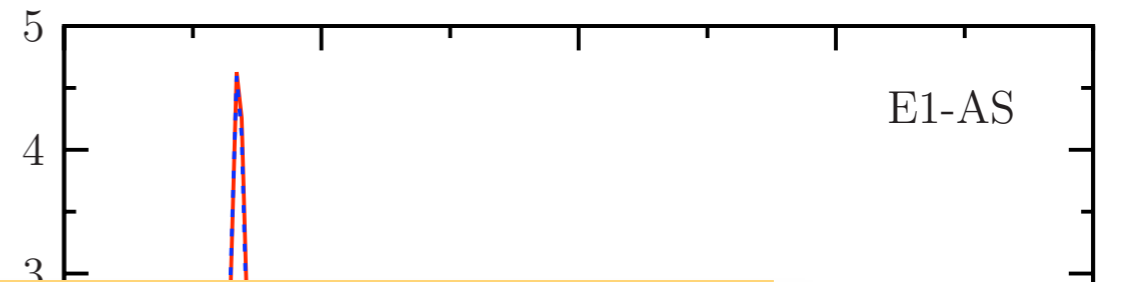
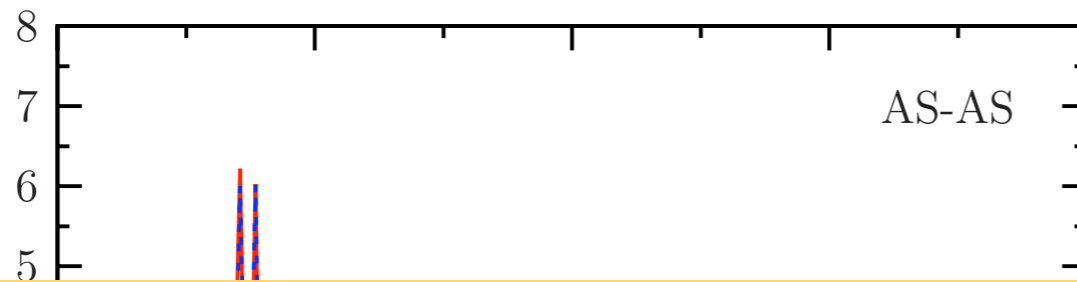


# Example for structure aa CG

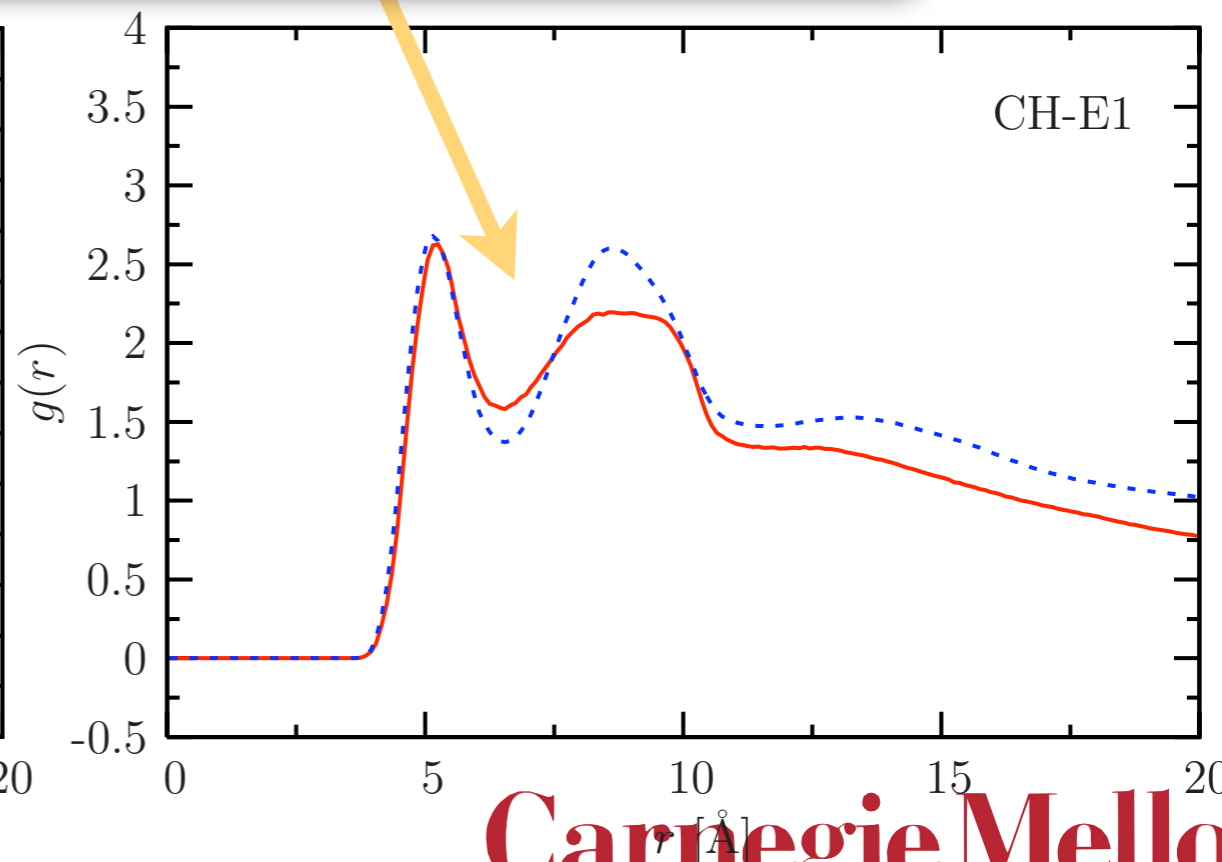
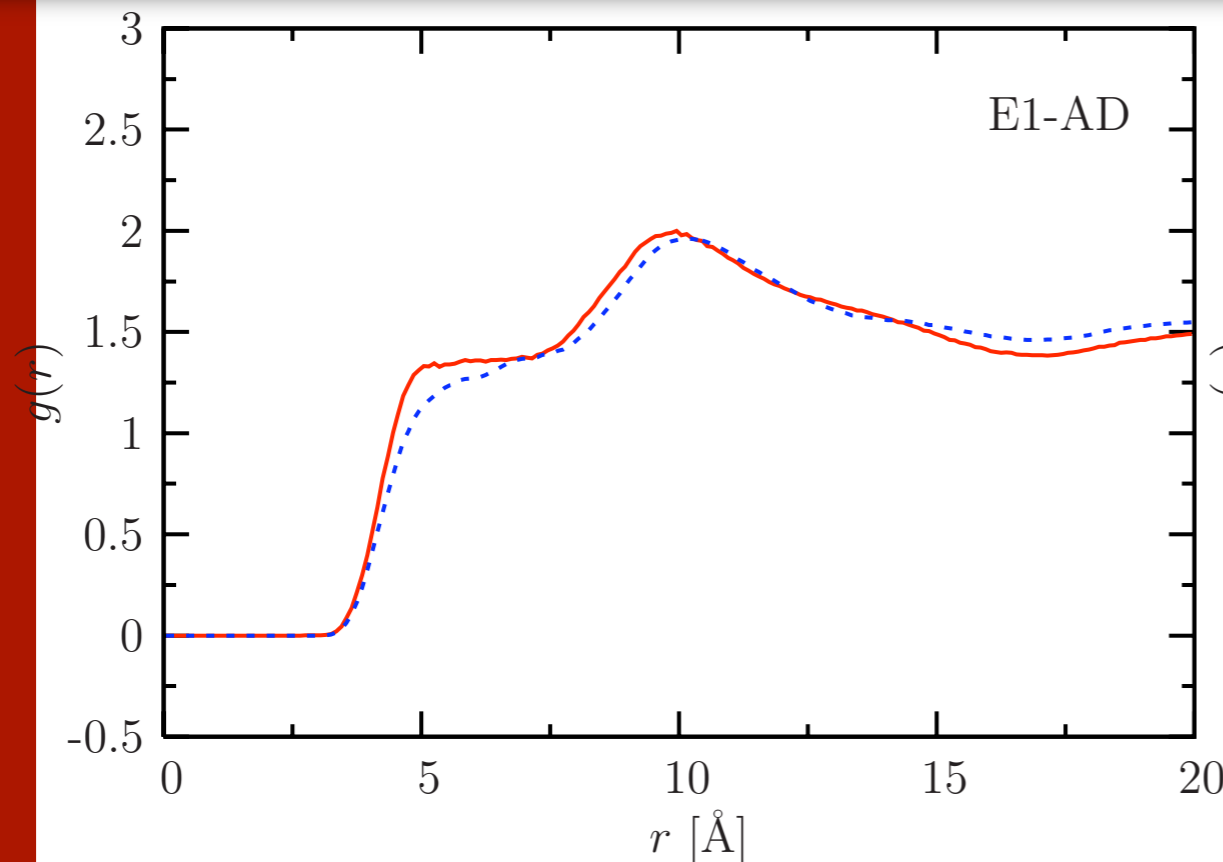




# Example for structure aa CG



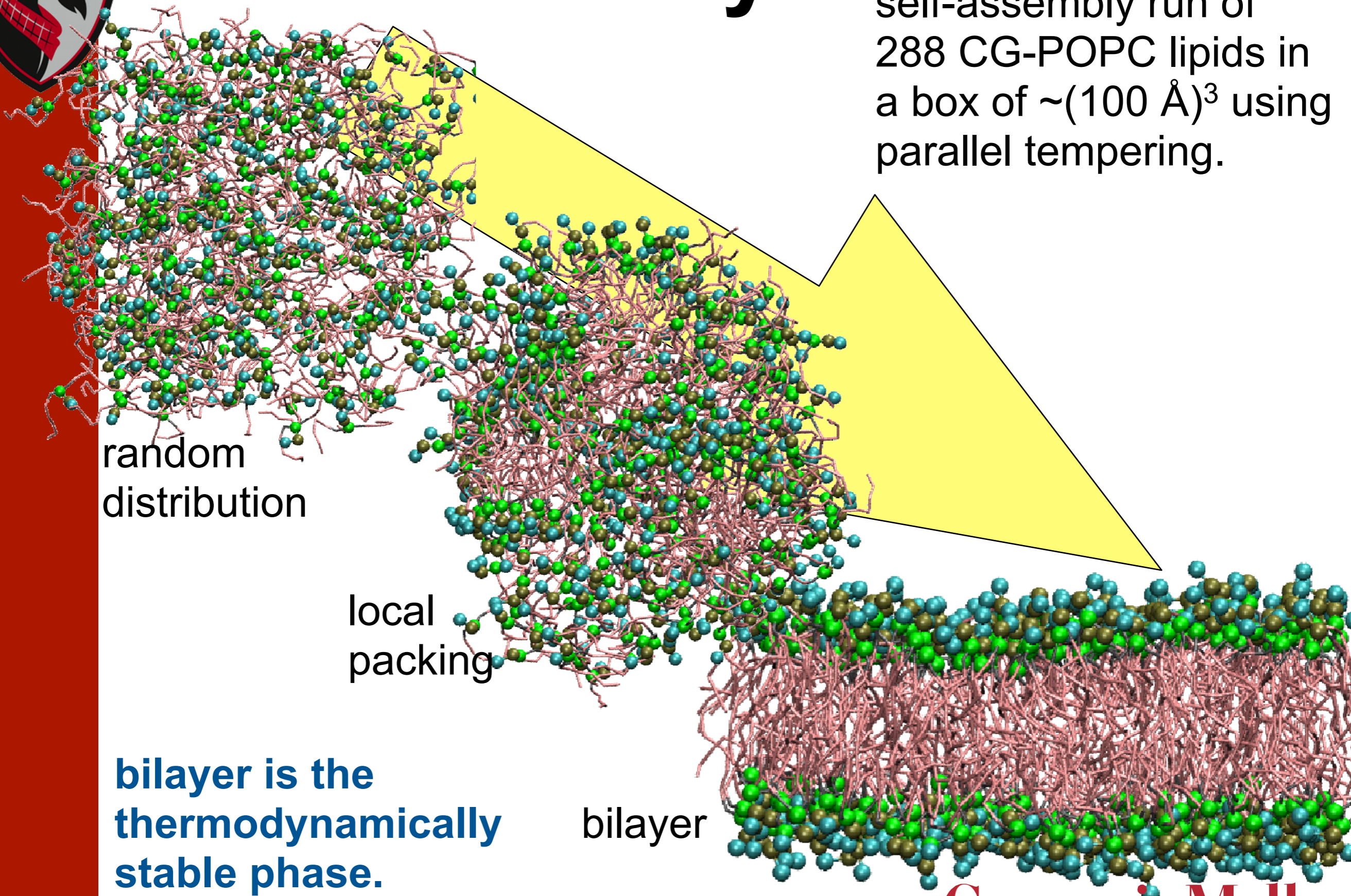
Some quantitative discrepancies remain.  
**However:** Not clear whether “poor convergence” or genuine inability to get the right structure.  
(Note: Henderson’s theorem does *not* apply to this complicated case!)





# Self-assembly

self-assembly run of 288 CG-POPC lipids in a box of  $\sim(100 \text{ \AA})^3$  using parallel tempering.



random distribution

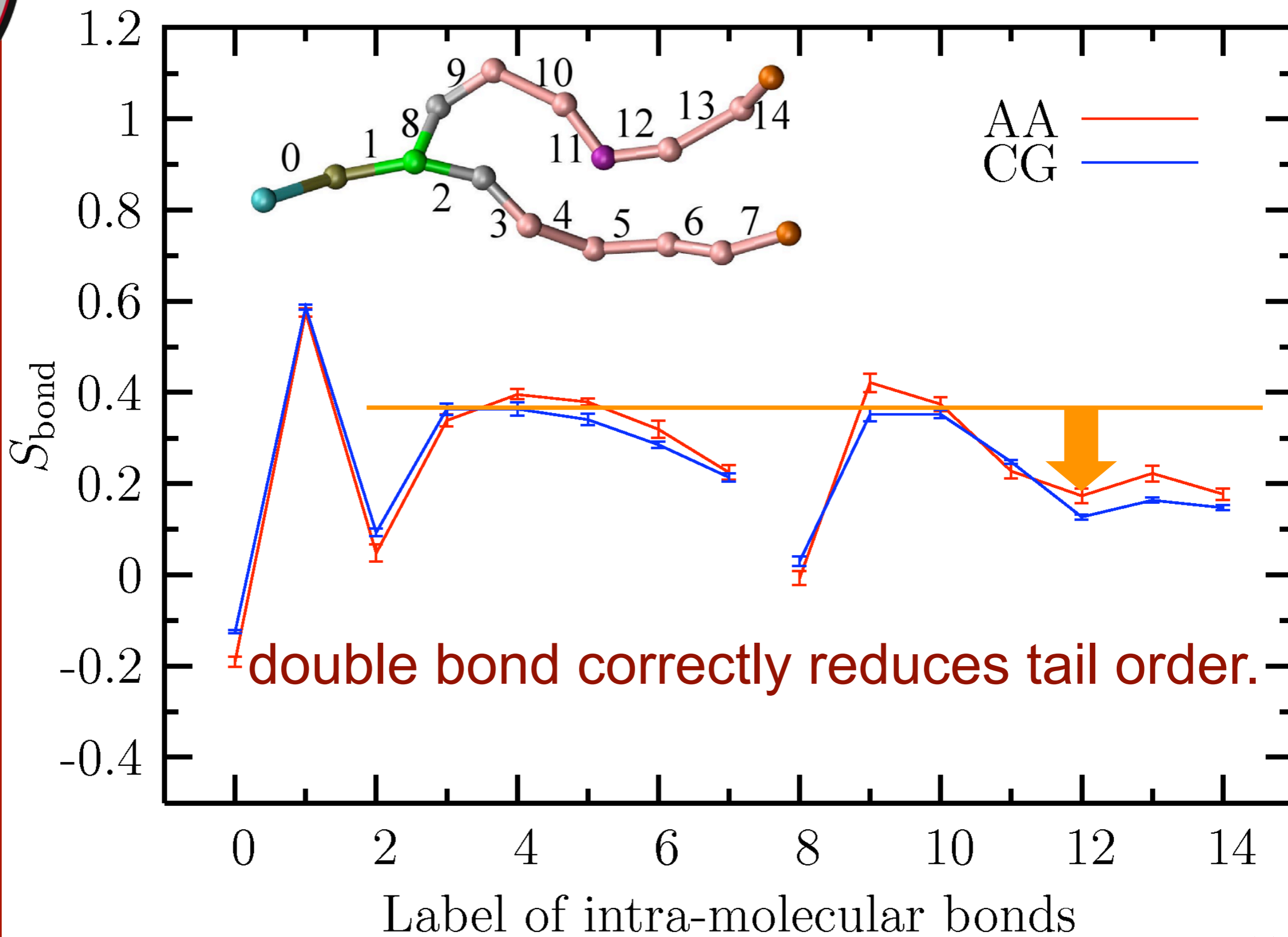
local packing

bilayer

**bilayer is the thermodynamically stable phase.**



# Tail order

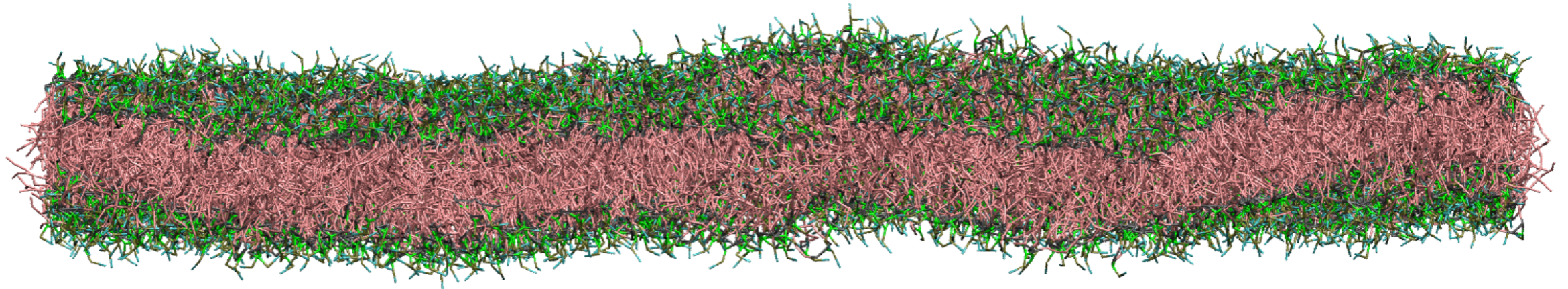




# Elastic properties

## Bending modulus $\kappa$

4608 lipids,  $(\sim 40\text{nm})^2$  box, several  $\mu\text{s}$



from fluctuations:  $\kappa = \frac{k_B T}{L^2 q^4 \langle |h_q|^2 \rangle}$   $(8...10) \times 10^{-20} \text{ J}$

experimentally, it's about:  $8.5 \times 10^{-20} \text{ mN/m}$

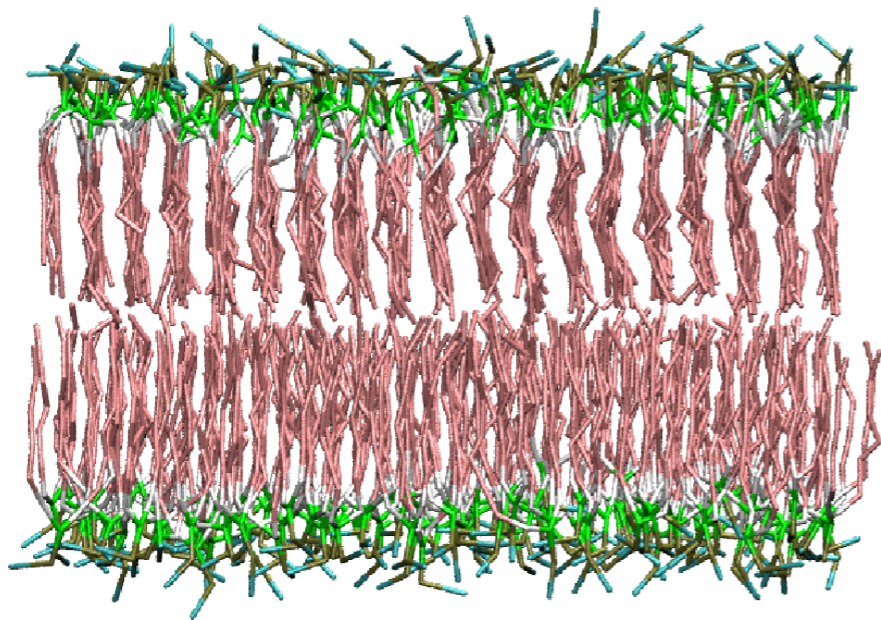
N. Kučerka, S. Tristram-Nagle & J.F. Nagle, J. Mem. Biol. **208**, 193 (2005).



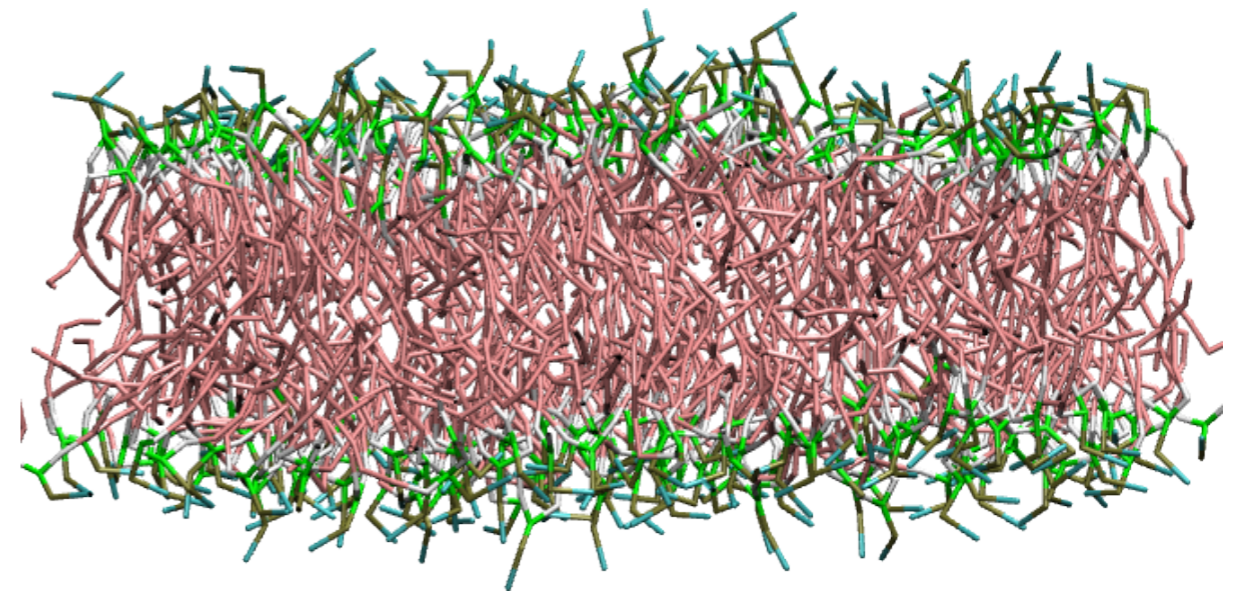


# Phase behavior

Our bilayer shows a liquid-gel transition.



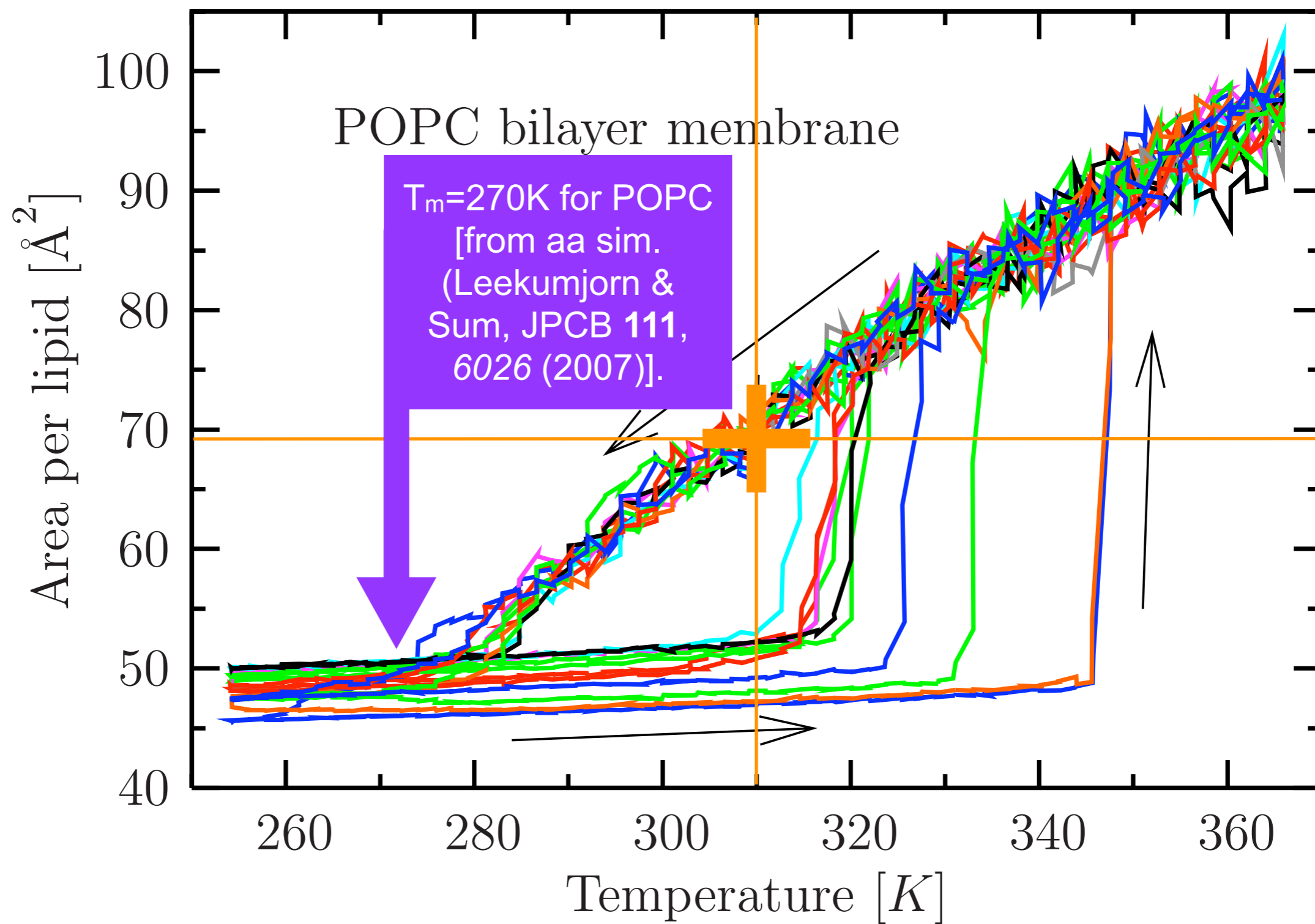
$T=279$  K



$T=310$  K



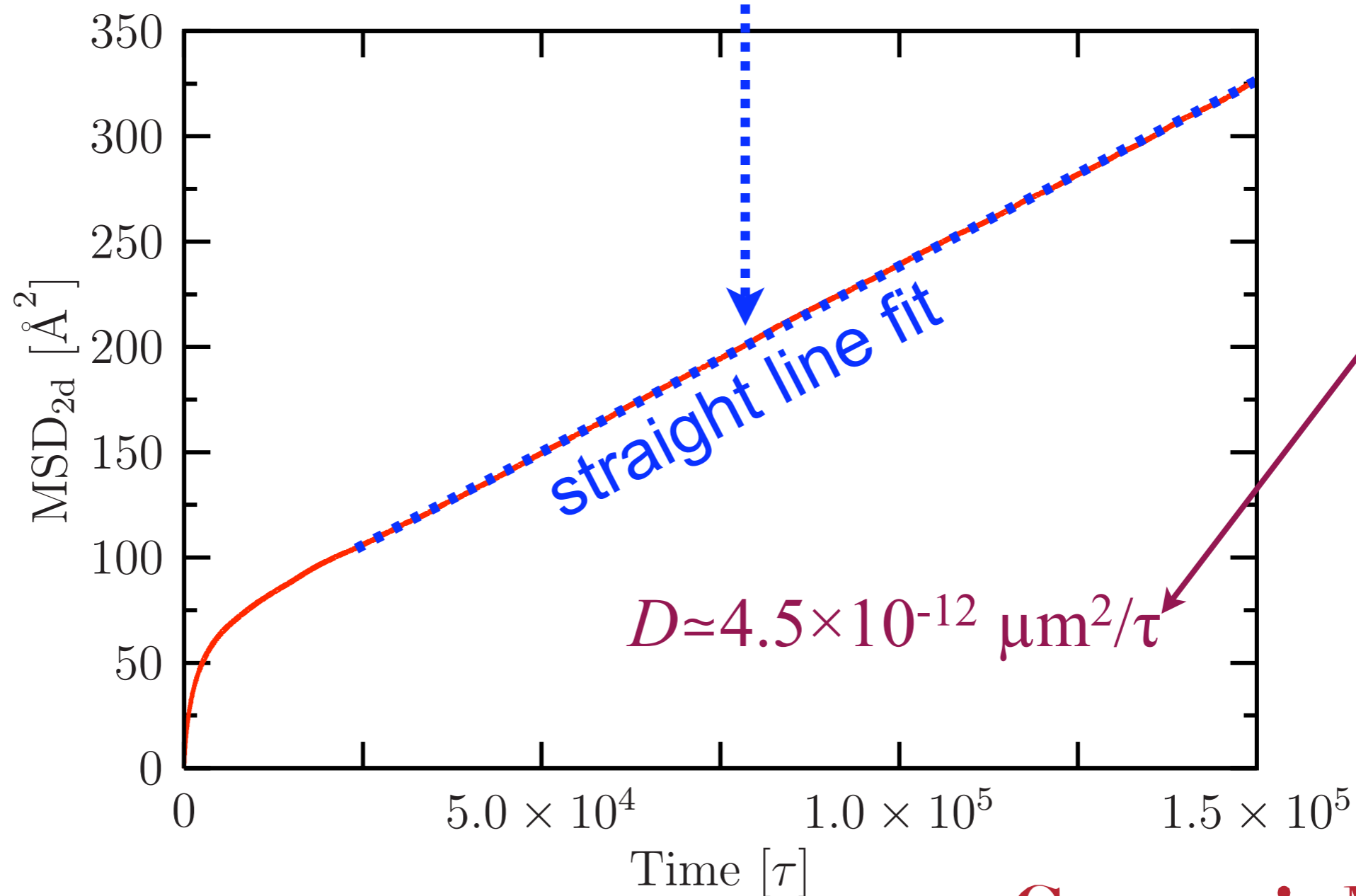
# Phase behavior





# Diffusion

Diffusion constant from:  $D = \lim_{t \rightarrow \infty} \left\{ \frac{1}{4N} \frac{d}{dt} \left\langle \sum_{i=1}^N \Delta r_i^2 \right\rangle \right\}$





# Time mapping

intrinsic computer time:  $\tau = \sigma\sqrt{m/\varepsilon} = 0.062\text{ps}$

Using this to translate times gives

$$D \approx 4.5 \times 10^{-12} \mu\text{m}^2 / (0.062 \text{ ps}) \approx 73 \mu\text{m}^2/\text{s}$$

Experimentally, more like  $D \approx (1 \dots 10) \mu\text{m}^2/\text{s}$

**However:**

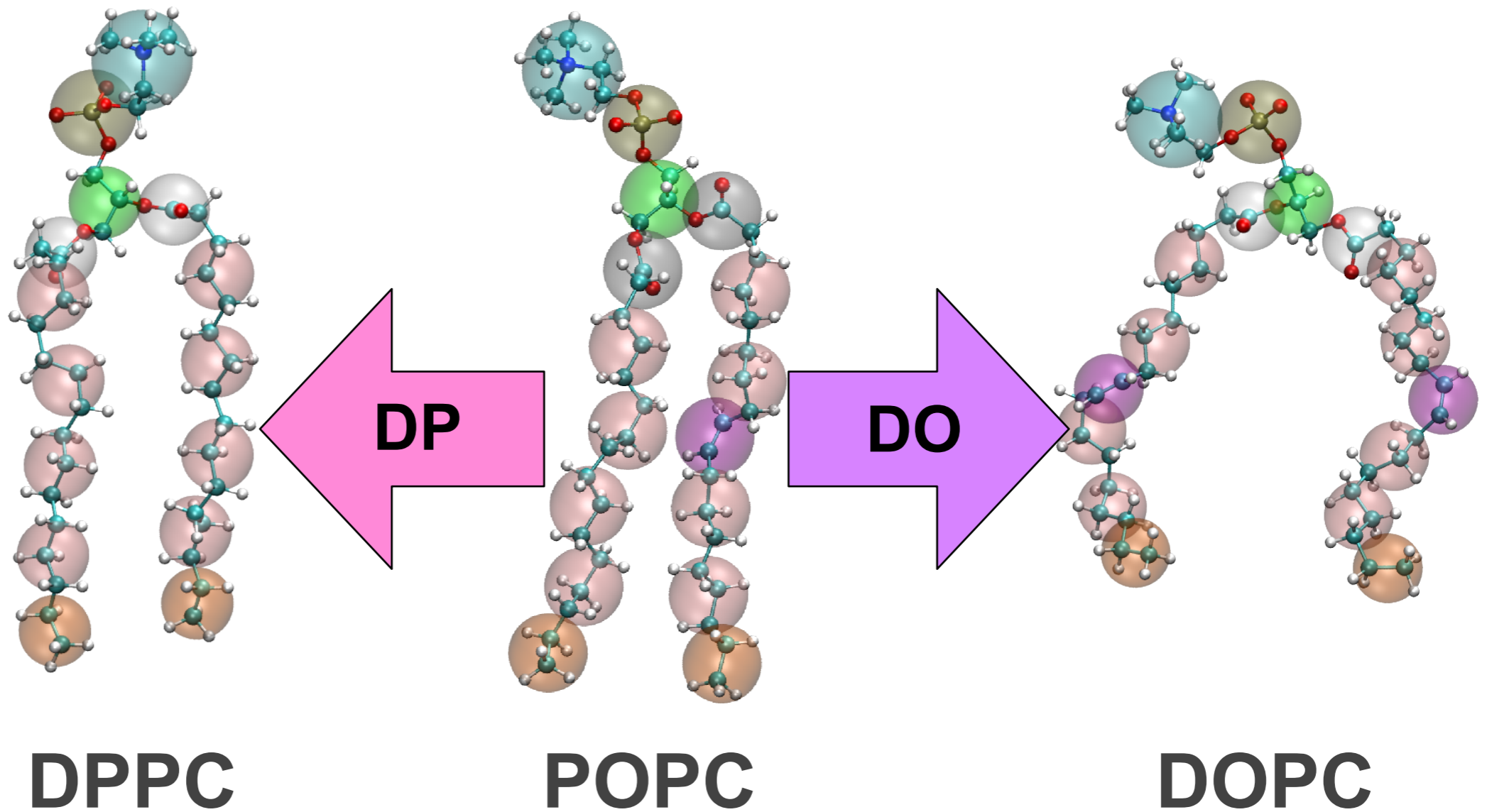
Reduced CG resolution implies reduced friction  
and thus an *expected* speed-up!

(Here, between one and two orders of magnitude)



# Transferability

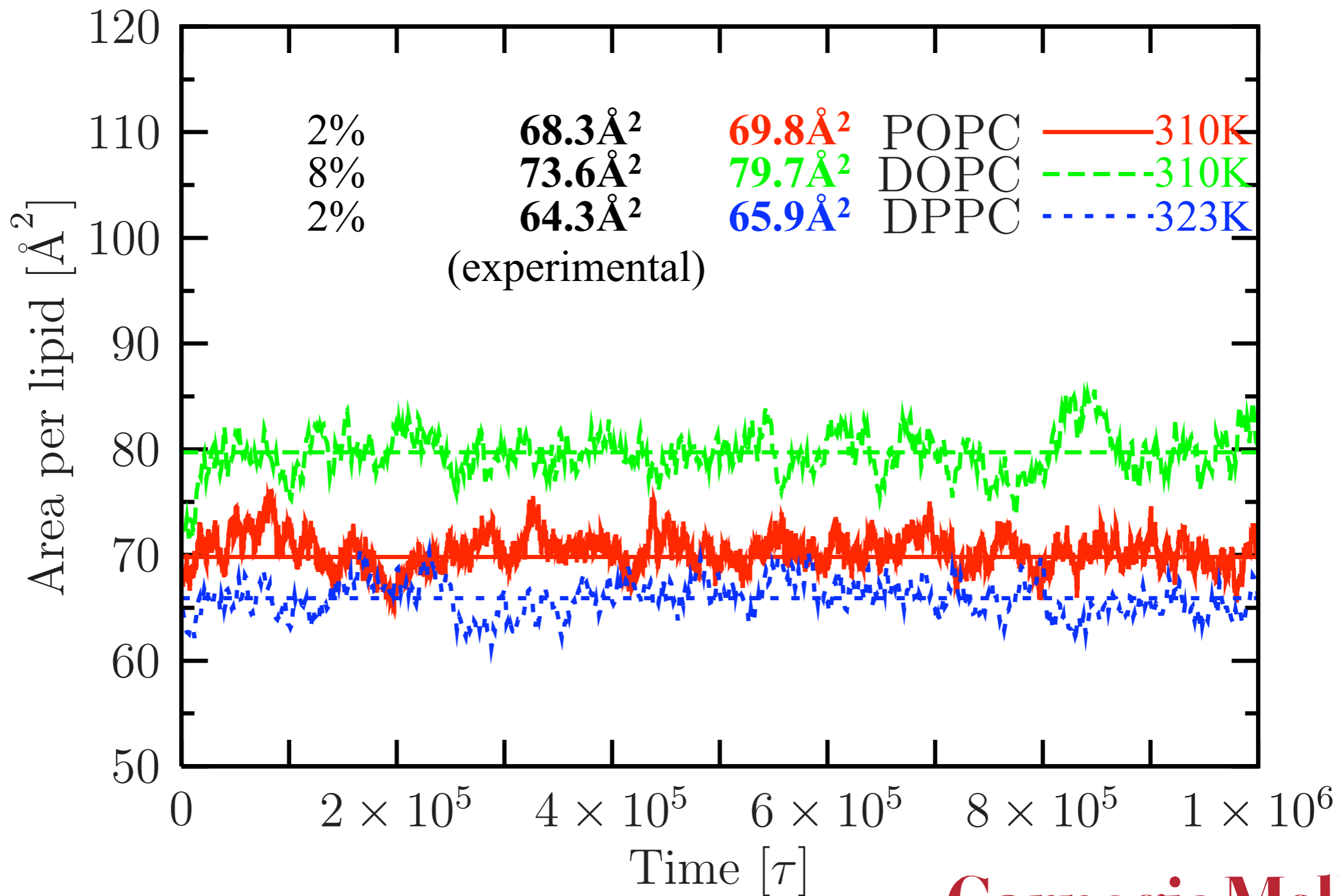
Example: “tail swapping”





# Transferability

## Area per lipid





# 2. Our peptide model

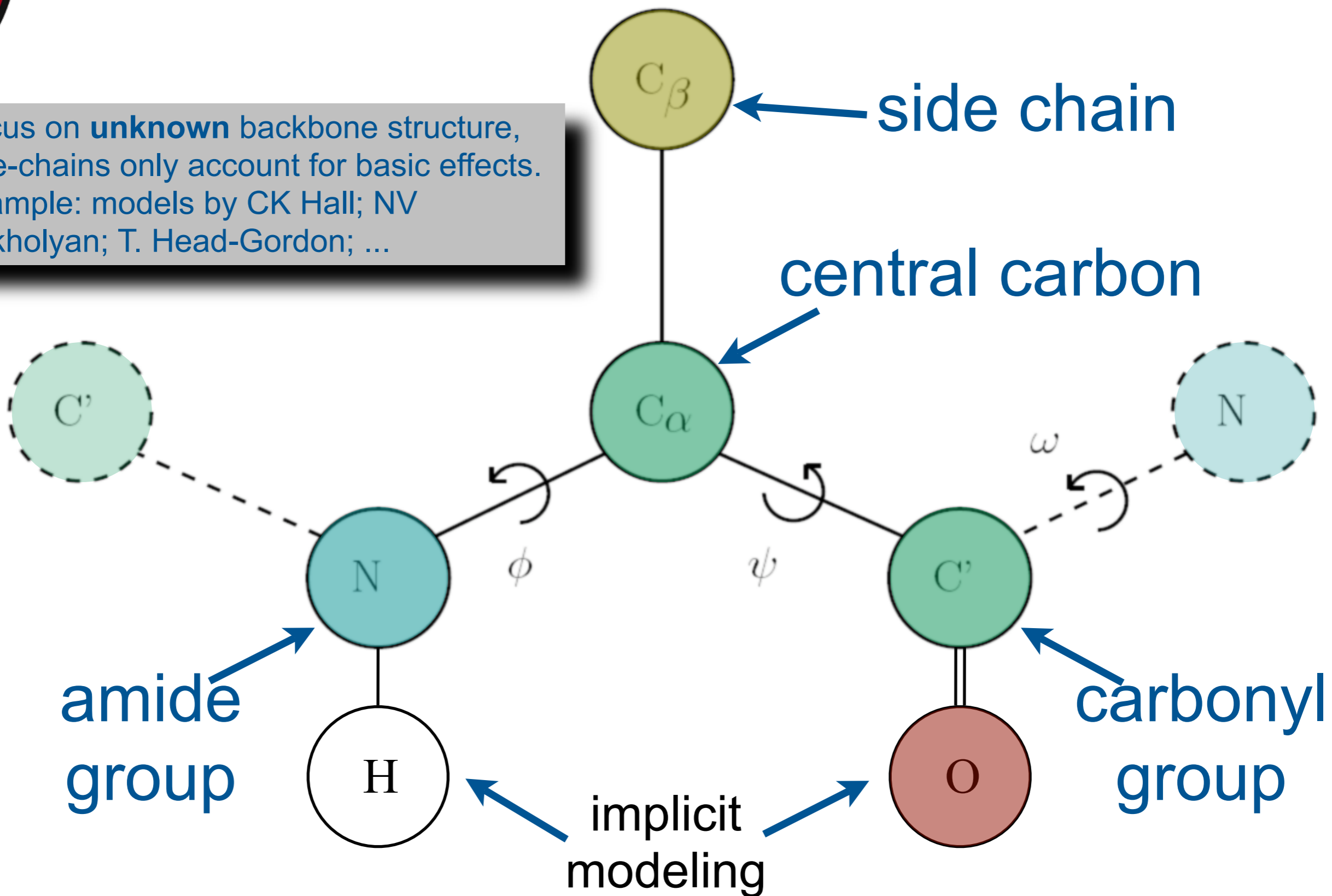


- T. Bereau and M. Deserno, J. Chem. Phys. **130**, 235106 (2009);  
T. Bereau, M. Bachmann, and M. Deserno, J. Am. Chem. Soc. **132**, 13129 (2010);  
T. Bereau, M. Deserno, and M. Bachmann, Biophys. J. **100**, 2764 (2011);  
T. Bereau, C. Globisch, M. Deserno, and C. Peter, J. Chem. Theo. Comput. ASAP (2012).



# Our model

Focus on **unknown** backbone structure, side-chains only account for basic effects. Example: models by CK Hall; NV Dokholyan; T. Head-Gordon; ...







# Interactions

bonded

- harmonic bonds
- harmonic angles
- dihedrals: simple cosine

improper dihedrals  
fix chirality at  $C_\alpha$

non-bonded

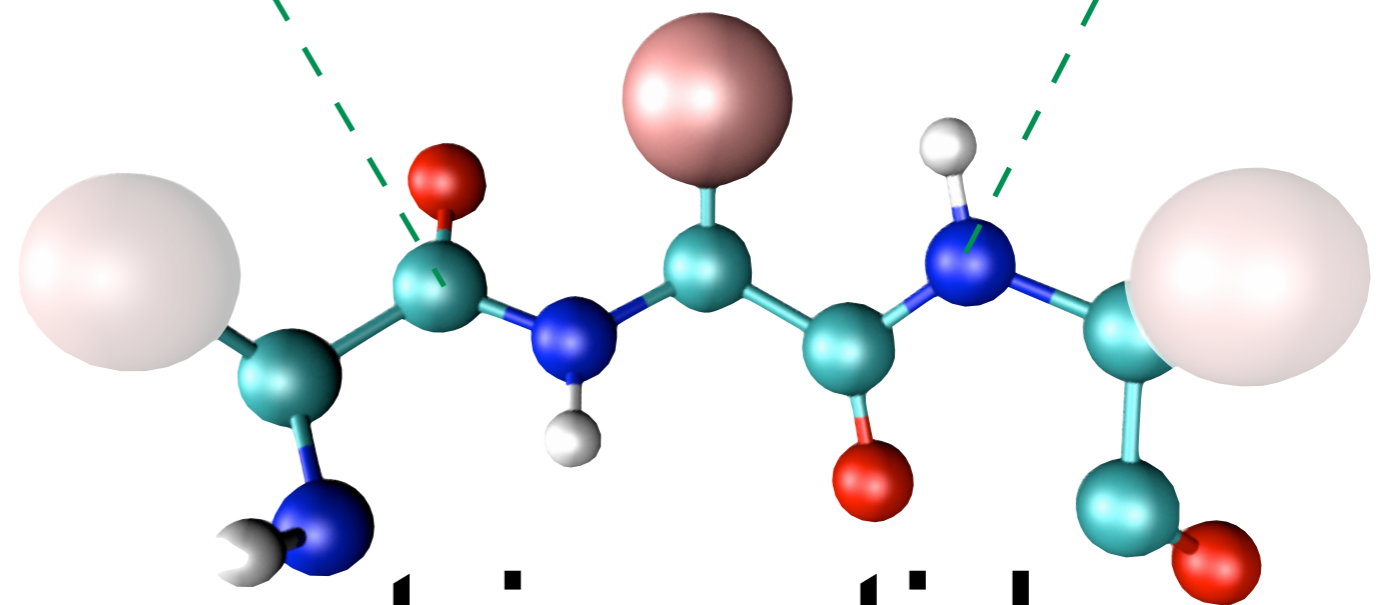
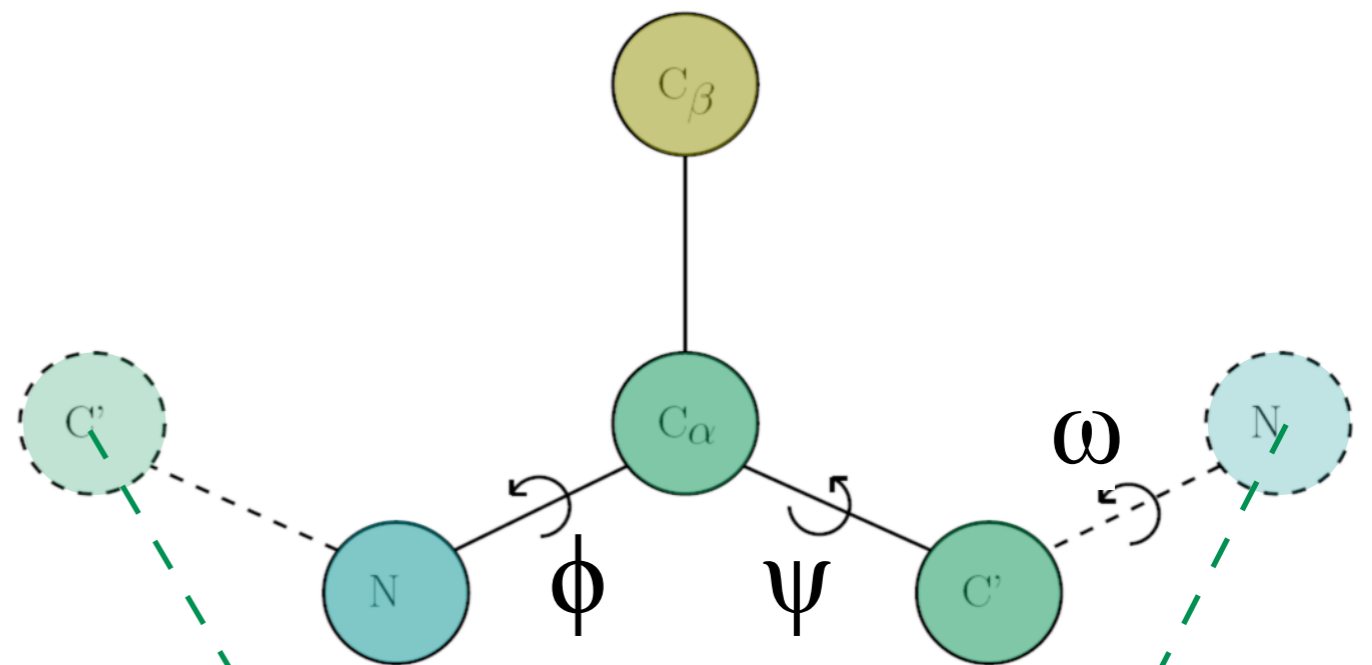
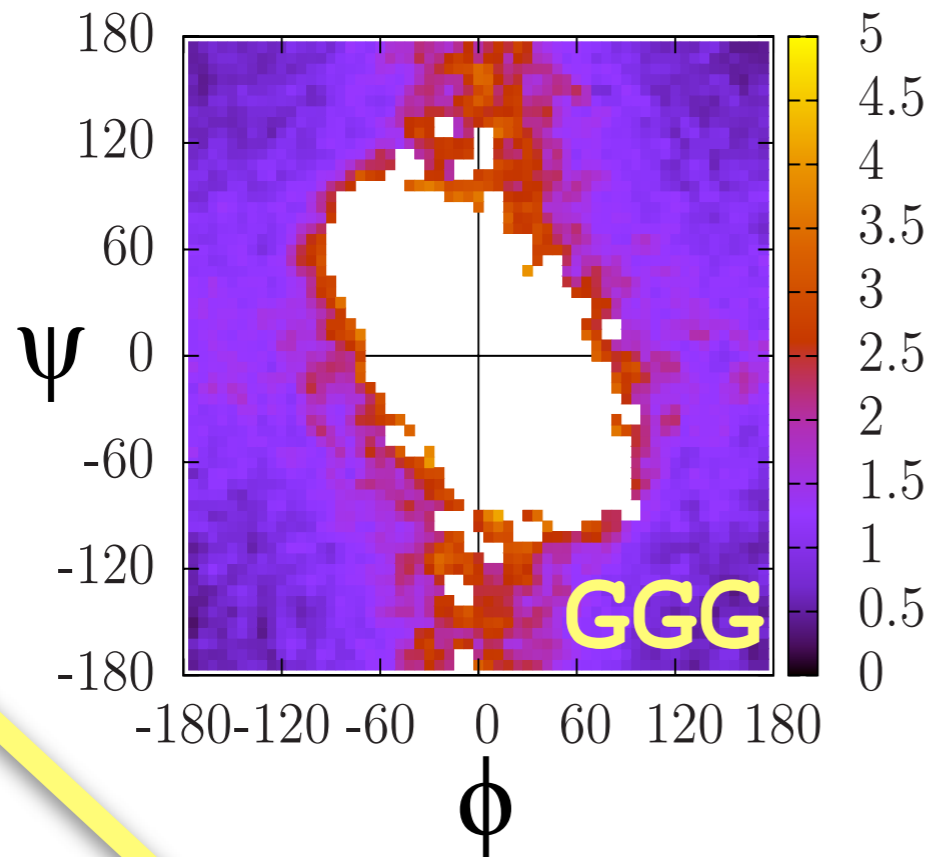
- excluded volume
- hydrophobic interactions
- hydrogen bonds
- peptide dipoles

Ramachandran plot

folding  
characteristics



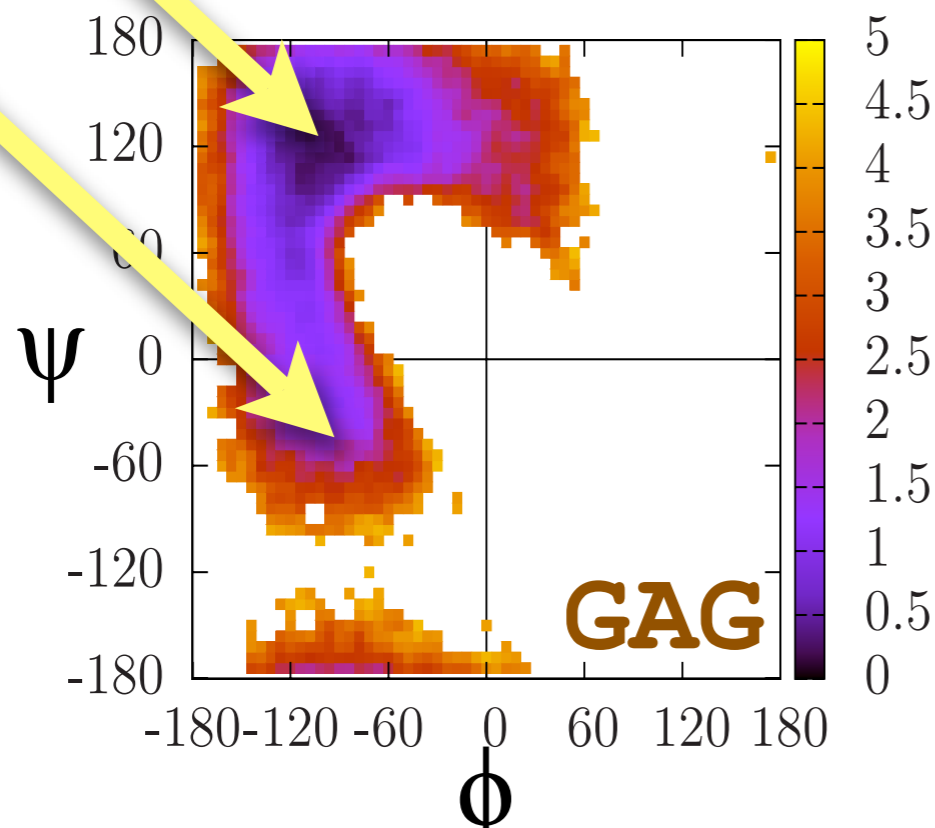
# Ramachandran plots



tripeptide

$\beta$

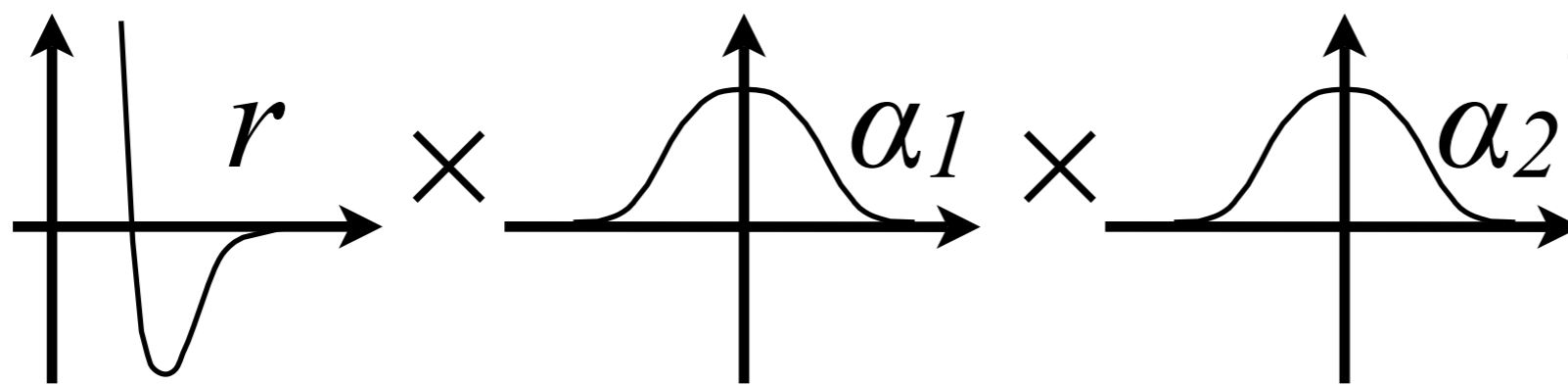
$\alpha$





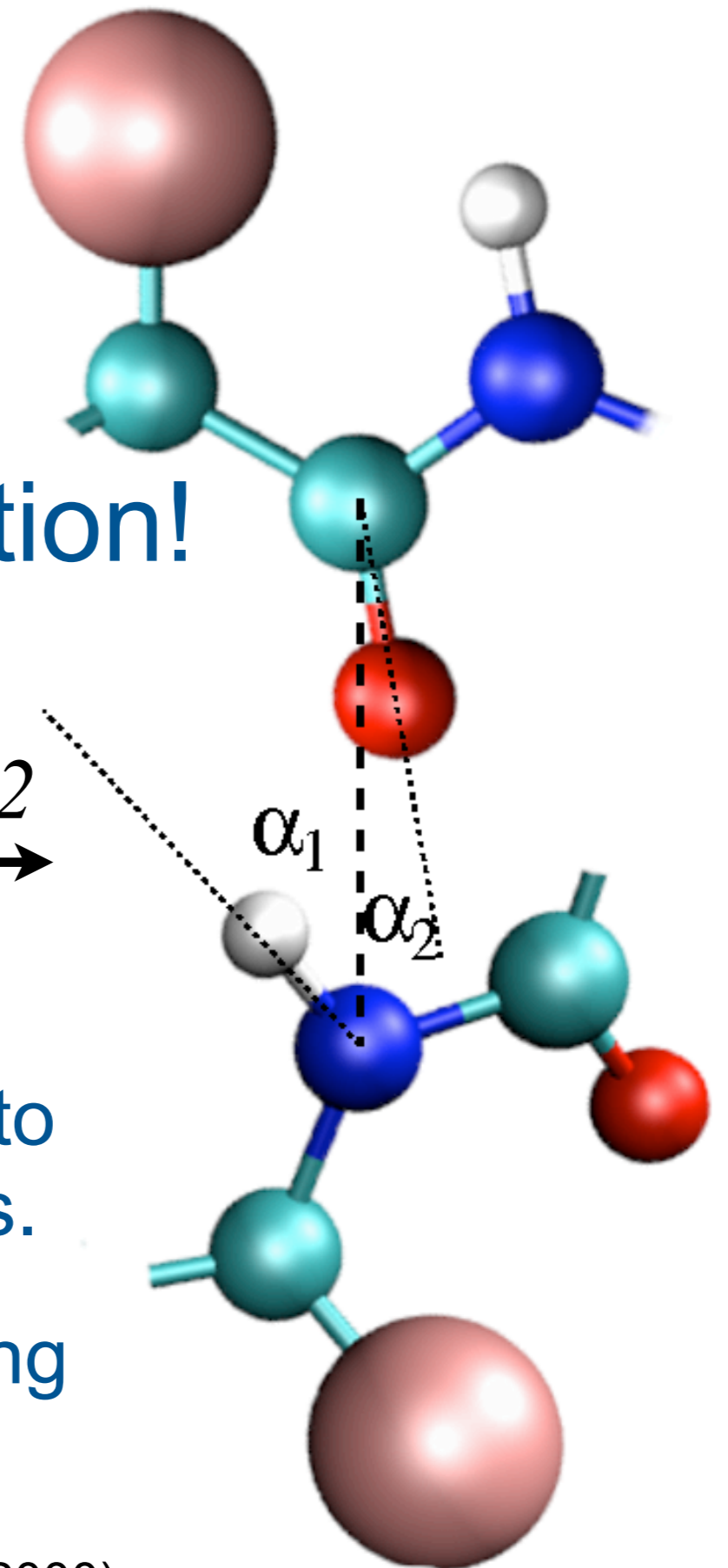
# Hydrogen bonds

extremely important for secondary structure formation!



Potential *range* adjusted according to known properties of hydrogen bonds.

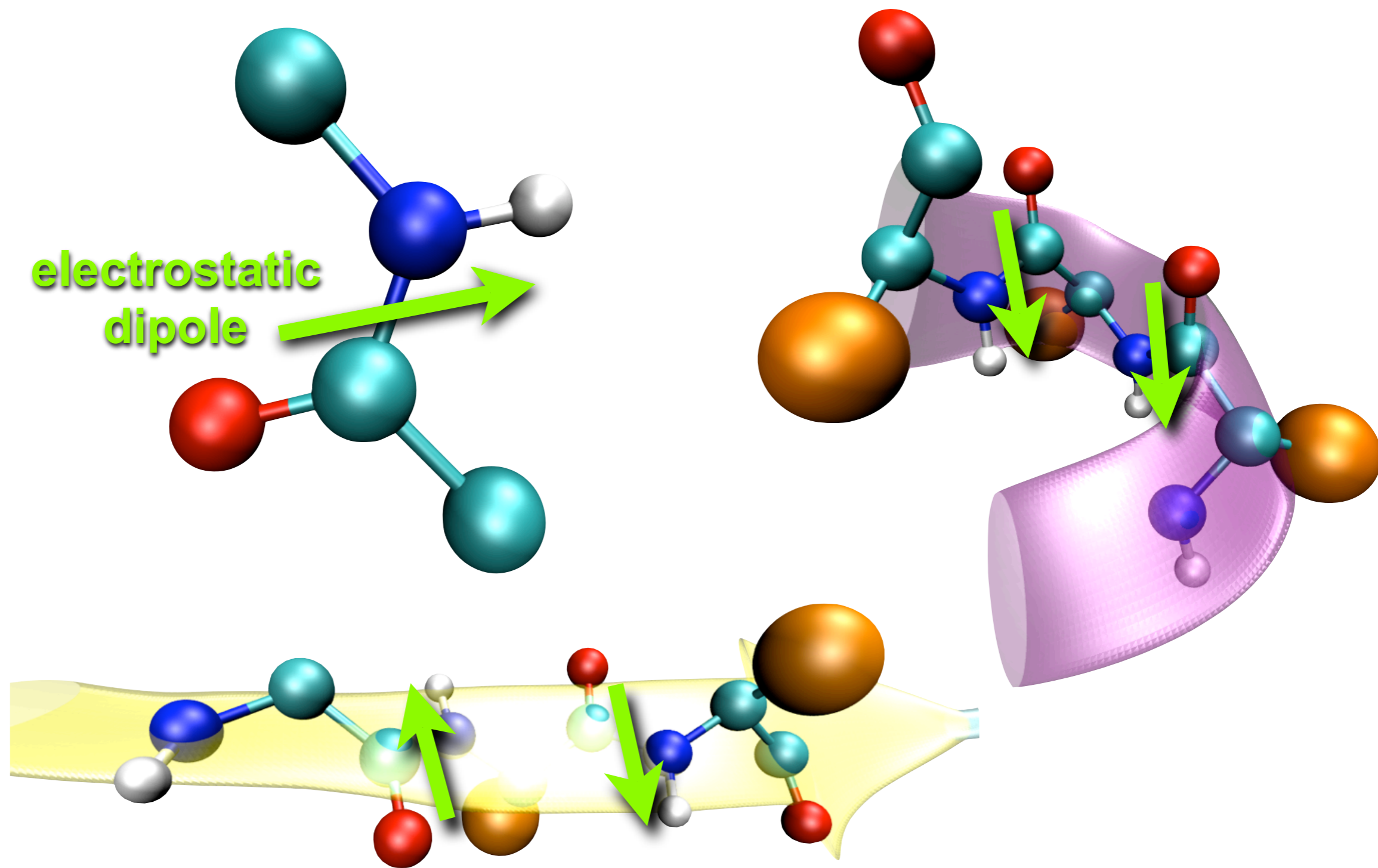
Potential *strength* adjusted according to folding properties of the model.



A. Irbäck, F. Sjunneson, and S. Wallin, PNAS 97, 13614 (2000)



# Peptide bond dipoles



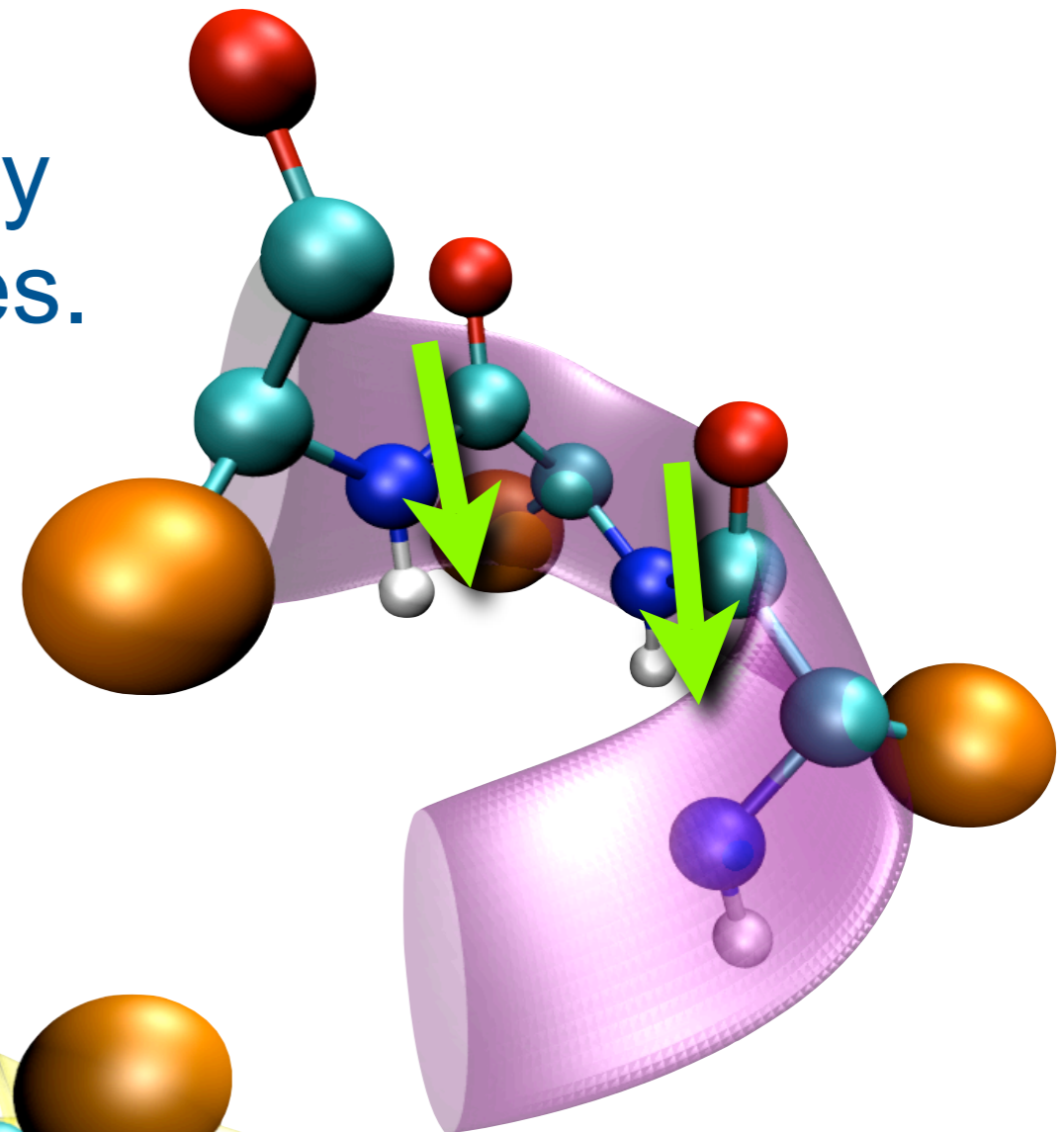
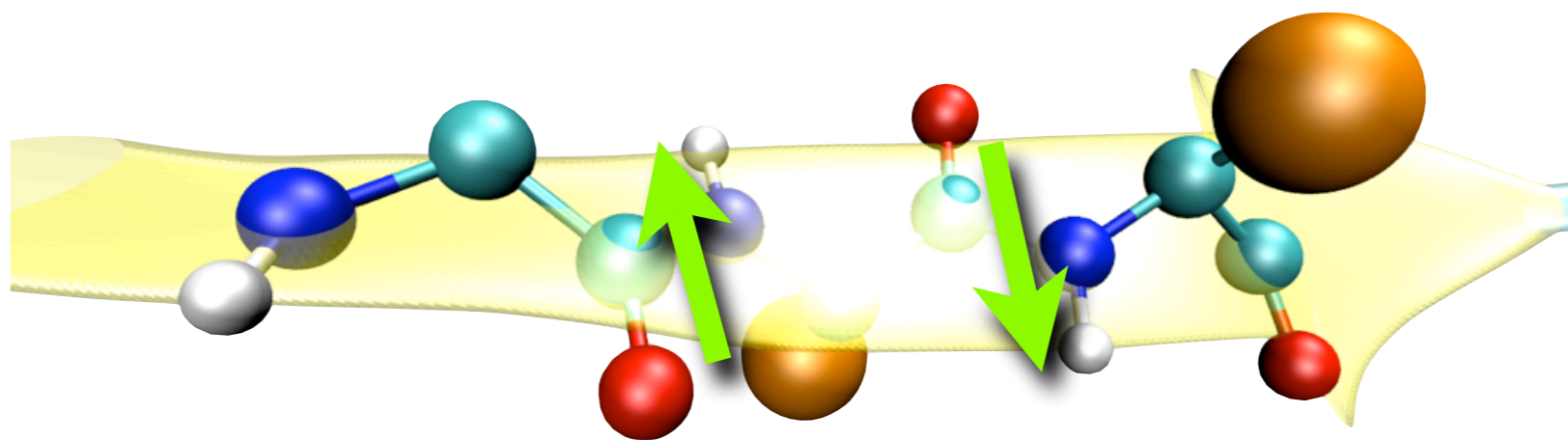
N.Y. Chen, Z.Y. Su, and C.Y. Mou, Phys. Rev. Lett. **96**, 078103 (2006)



# Peptide bond dipoles

$\beta$ -sheets are electrostatically more favorable than  $\alpha$ -helices.

Incorporate this effect in the  $(\phi, \psi)$ -dihedrals!

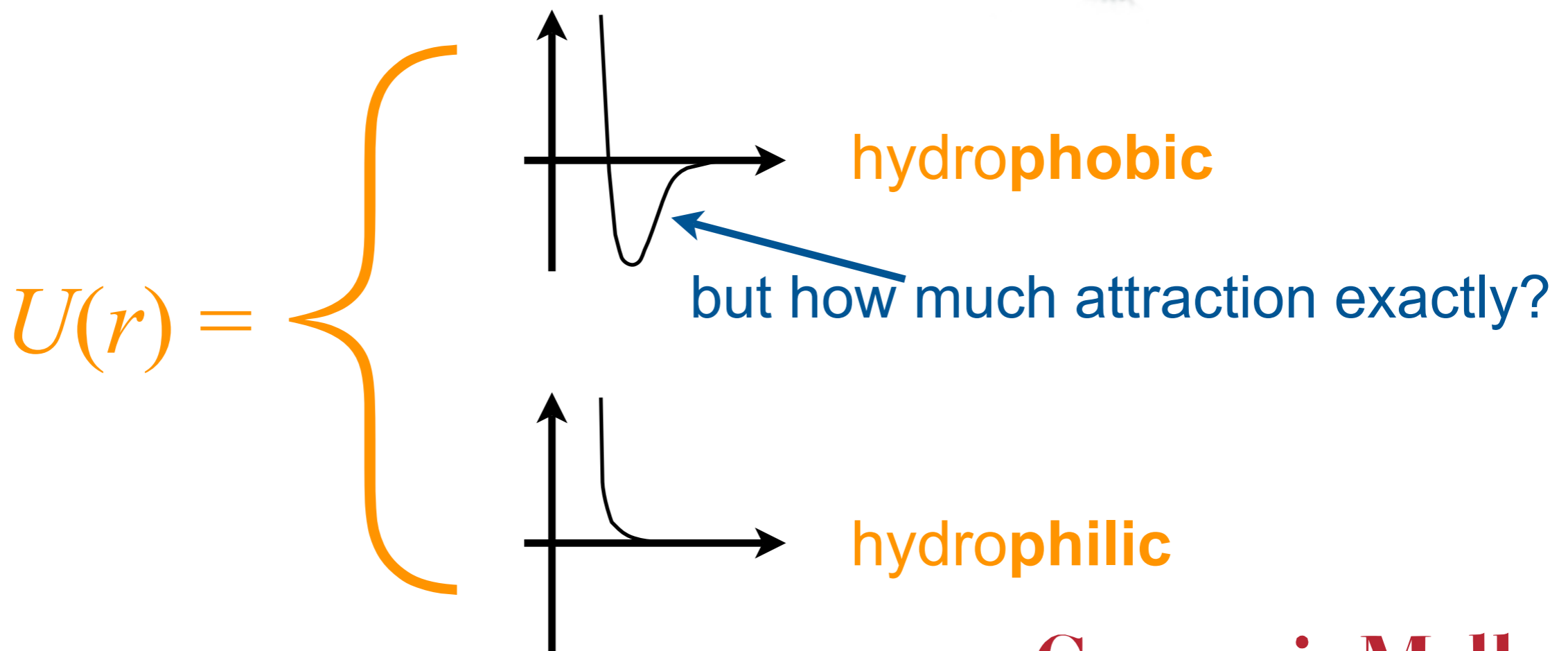
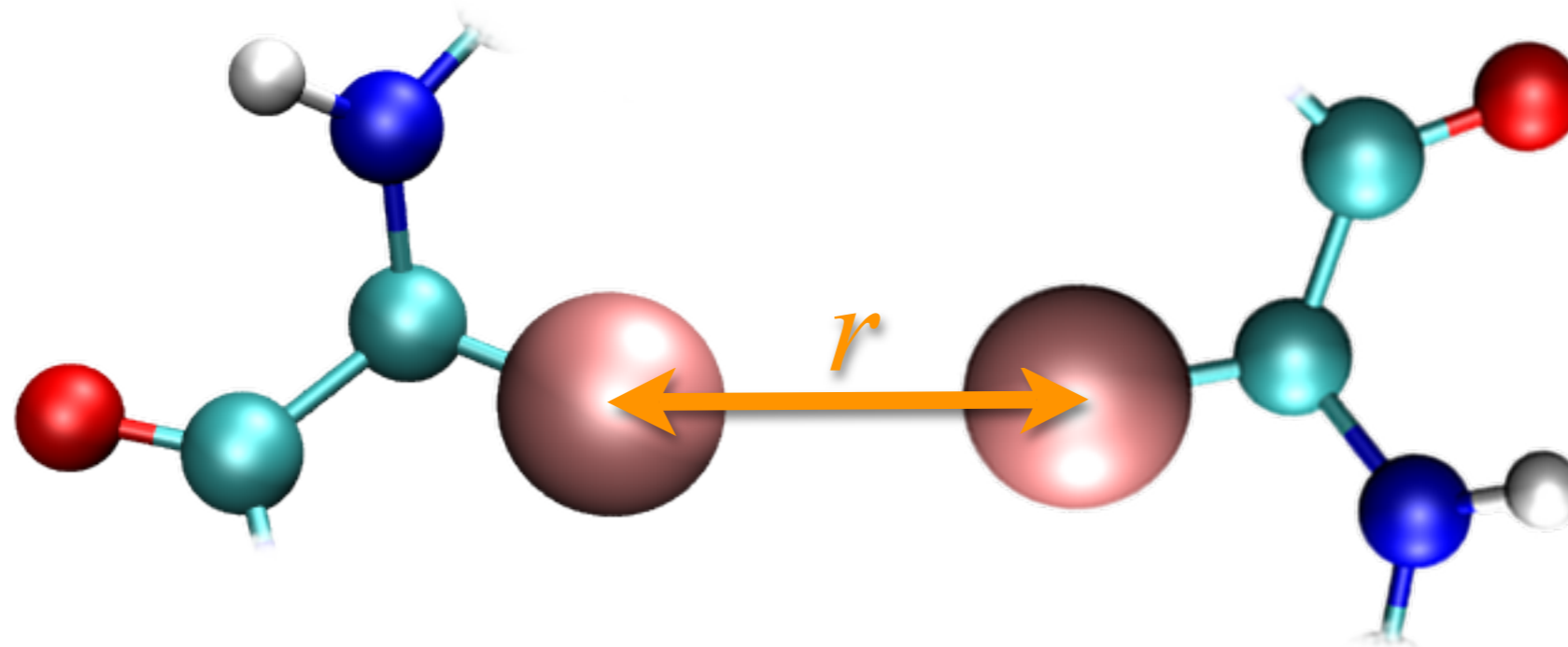


N.Y. Chen, Z.Y. Su, and C.Y. Mou, Phys. Rev. Lett. **96**, 078103 (2006)

Y. Mu and Y.Q. Gao, J. Chem. Phys. **127**, 105102 (2007)



# Hydrophobic interactions





# Miyazawa-Jernigan matrix

$$\epsilon_{ij}^{MJ}$$

	Ser	Gly	Phe	
Ser	-1.67	-1.82	-4.02	...
Gly	-1.82	-2.24	-4.13	...
Phe	-4.02	-4.13	-7.26	...

S. Miyazawa and R.L. Jernigan, J. Mol. Biol. **256**, 623 (1996)



# Miyazawa-Jernigan matrix

$$\epsilon_{ij}^{\text{MJ}} \approx \epsilon_{ij} = \sqrt{\epsilon_i \epsilon_j} \quad (\text{Lorentz-Berthelot})$$

normalize:  $0 \longleftrightarrow 1$

$$\epsilon_i' = \frac{\epsilon_i - \min \epsilon_k}{\max \epsilon_k - \min \epsilon_k}$$

correlation between  $\epsilon_{ij}^{\text{MJ}}$  and  $\epsilon_{ij}$  : 98%

correlation between  $\epsilon_i'$  and  $\epsilon_i$  : 87%

hydrophobicity scale by Fauchere and Pliska

J.L. Fauchere and V. Pliska, Eur. J. Med. Chem. **18**, 369 (1983).

S. Miyazawa and R.L. Jernigan, J. Mol. Biol. **256**, 623 (1996)





# Miyazawa-Jernigan matrix

$$\epsilon_{ij}^{\text{MJ}} \approx \epsilon_{ij} = \sqrt{\epsilon_i \epsilon_j} \quad (\text{Lorentz-Berthelot})$$

normalize:  $0 \longleftrightarrow 1$

$$\epsilon_i' = \frac{\epsilon_i - \min \epsilon_k}{\max \epsilon_k - \min \epsilon_k}$$

Use  $\epsilon_{\text{overall}} \times \epsilon_i'$  as the strength in a Lennard-Jones potential.

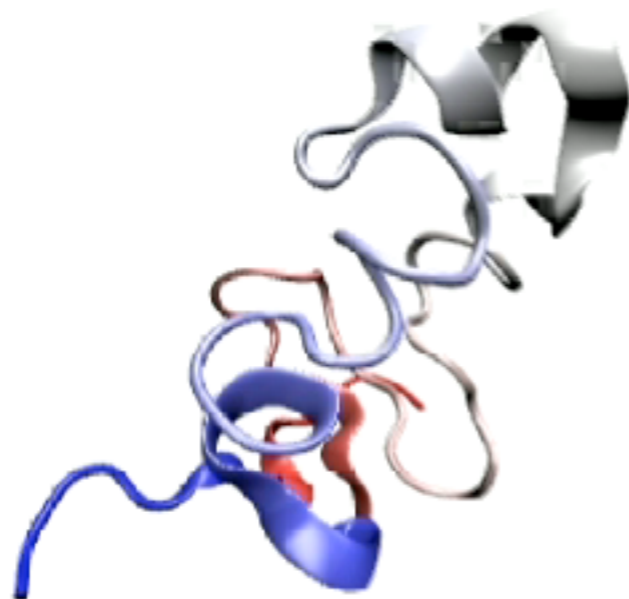


73 aa

# Folding of a 3-helix bundle

2A3D

MGSWAEEFKQRLAAIKTRLQALGGSEAELAAFEKEIAA...  
...FESELQAYKKGKGNPEVEALRKEAAAIRDELQAYRHN



Parallel tempering at 8 temperatures;  
shown: configuration in lowest replica.  
Total simulation time: ~10 hours.



73 aa

# Folding of a 3-helix bundle

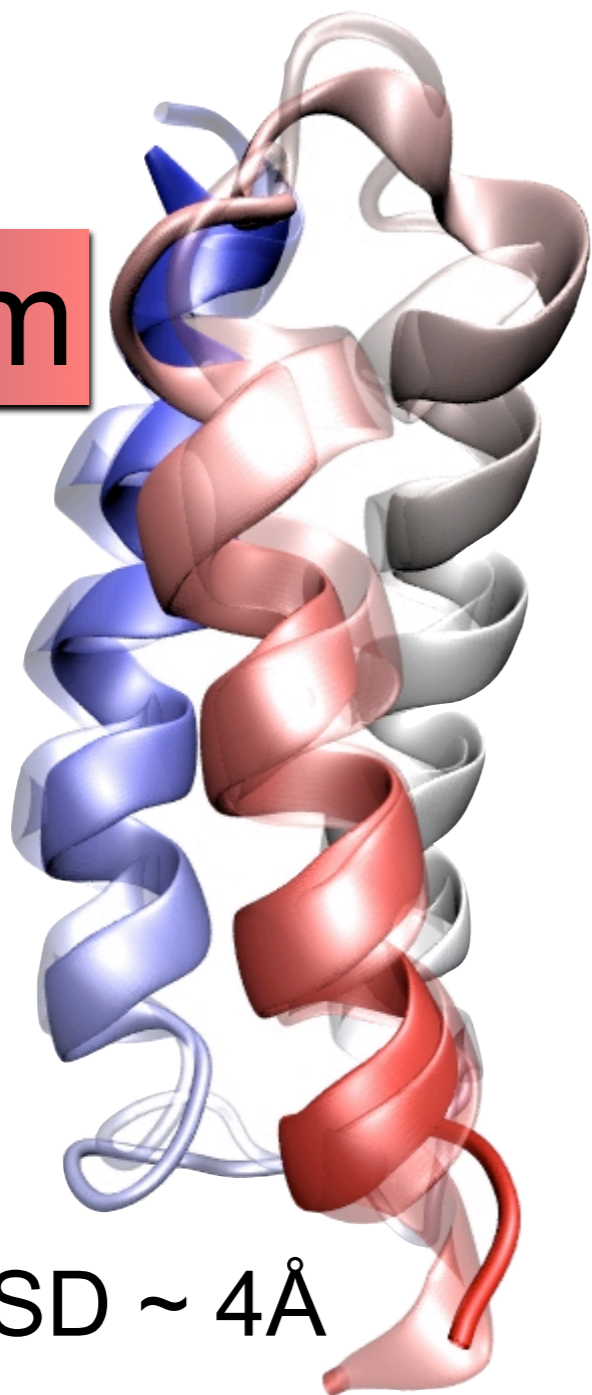
2A3D

MGSWAEFKQRLAAIKTRLQALGGSEAELAAFEKEIAA...  
...FESELQAYKKGKGNPEVEALRKEAAAIRDELQAYRHN

CG sim

NMR

RMSD ~ 4Å



hydrophobicity:

proper helix-alignment,  
but not a collapsed globule

hydrogen bond:

secondary structure formation,  
but not just one long  $\alpha$ -helix.

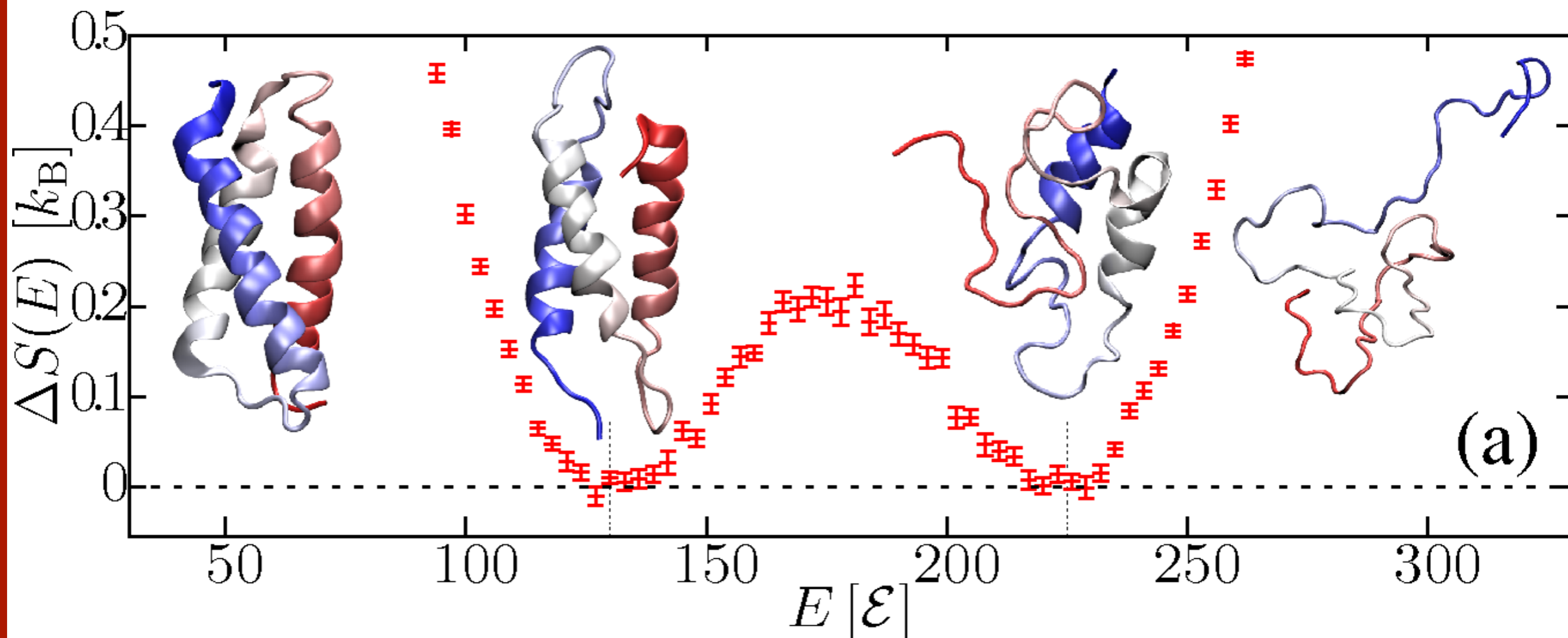
dipole term:

$\beta$ -sheets maximally stable  
without disabling  $\alpha$ -helices.



# Example application

Microcanonical folding analysis of this 3-helix bundle.



T. Bereau, M. Bachmann, and M. Deserno, *J. Am. Chem. Soc.* **132**, 13129 (2010);  
T. Bereau, M. Deserno, and M. Bachmann, *Biophys. J.* **100**, 2764 (2011);

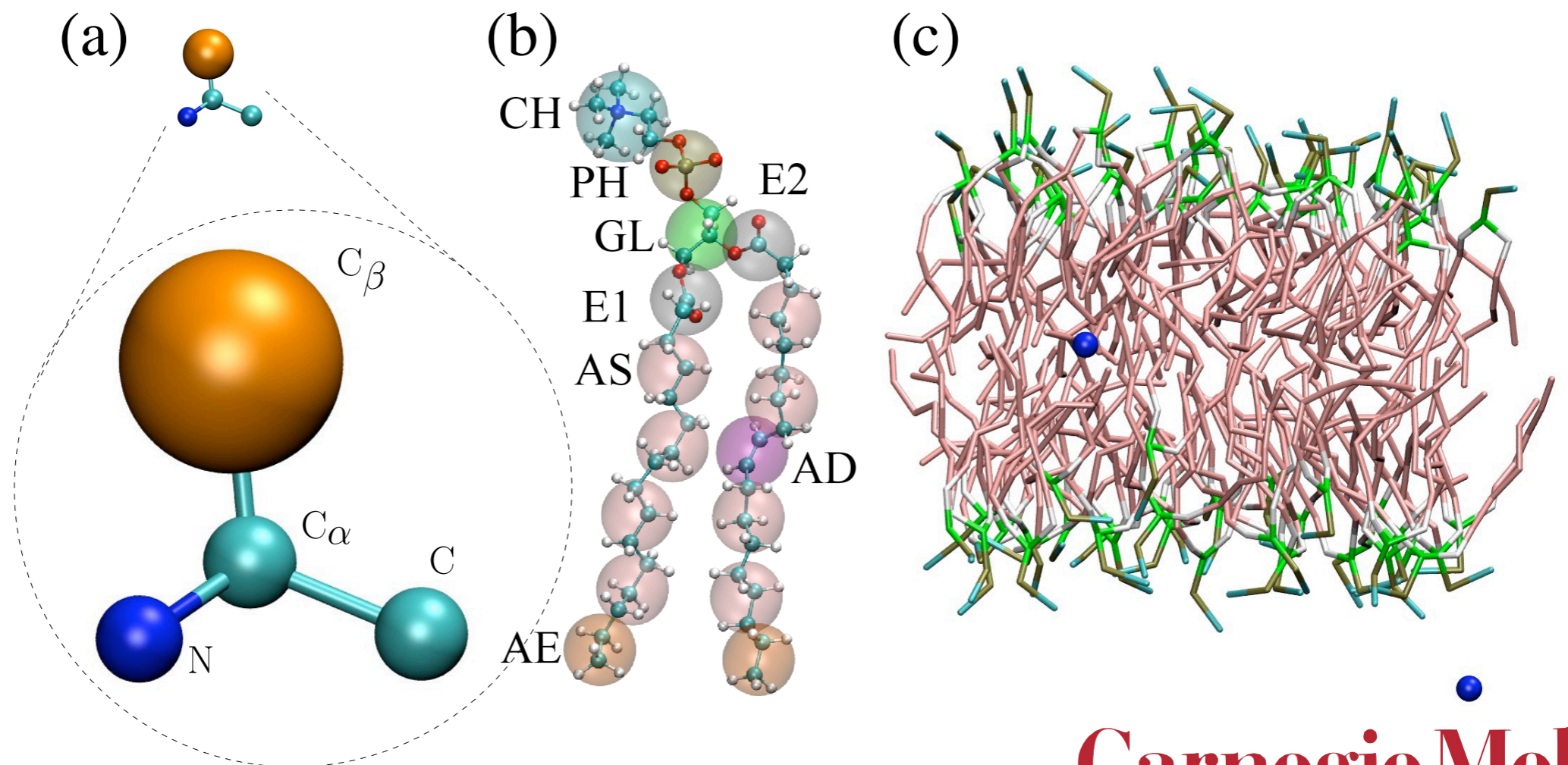


# 3. Cross-parametrization

**Aim:** Reproduce *free energy of insertion* of individual amino acids into the bilayer.

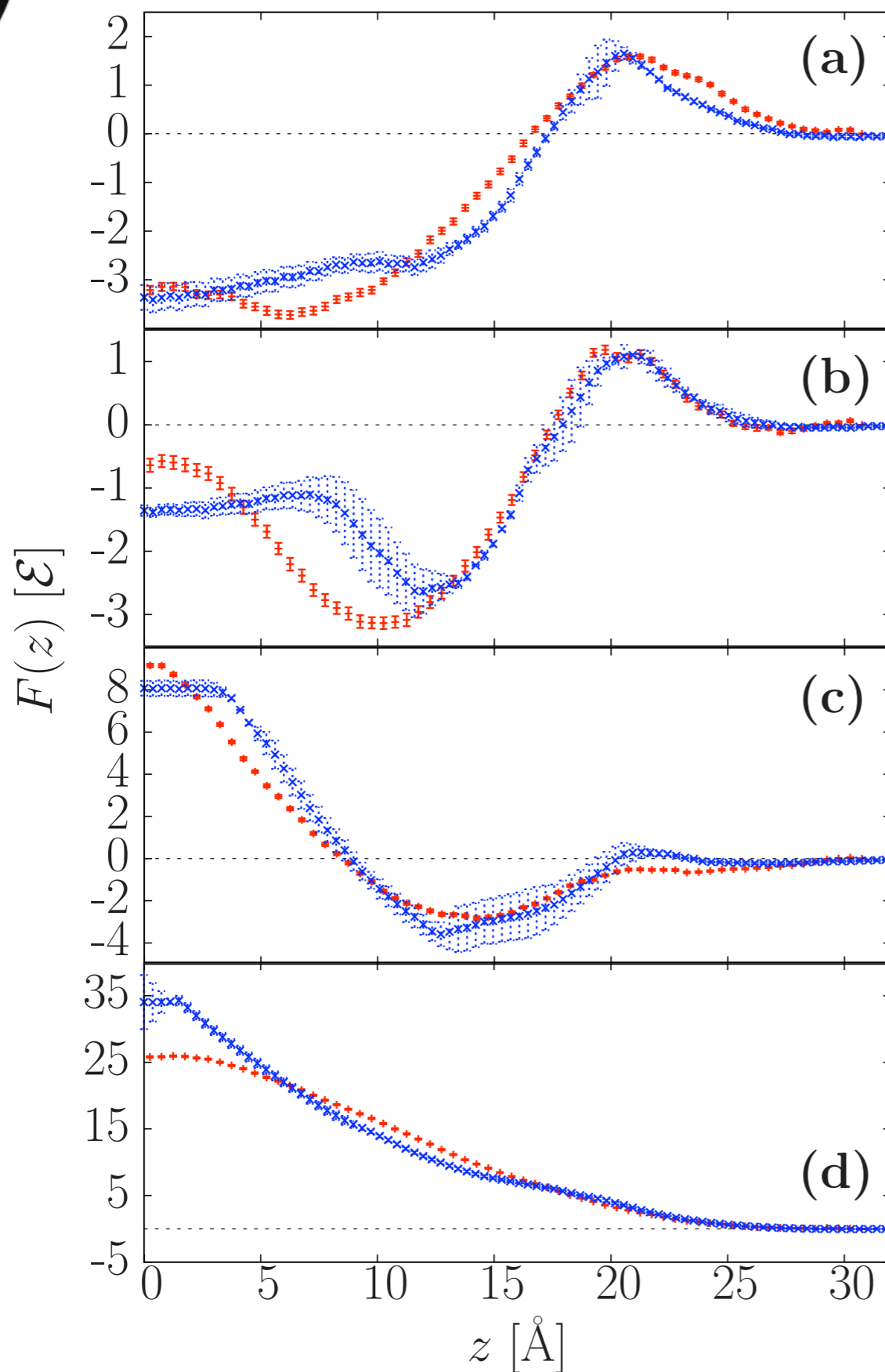
*Target data* for this:

J. L. MacCallum, W. F. D. Bennett, and D. P. Tieleman, *Biophys. J.* **94**, 3393 (2008).





# Free energy of insertion



Ala      atomistic  
            CG

Cys      Generic features  
            captured semi-  
            quantitatively.

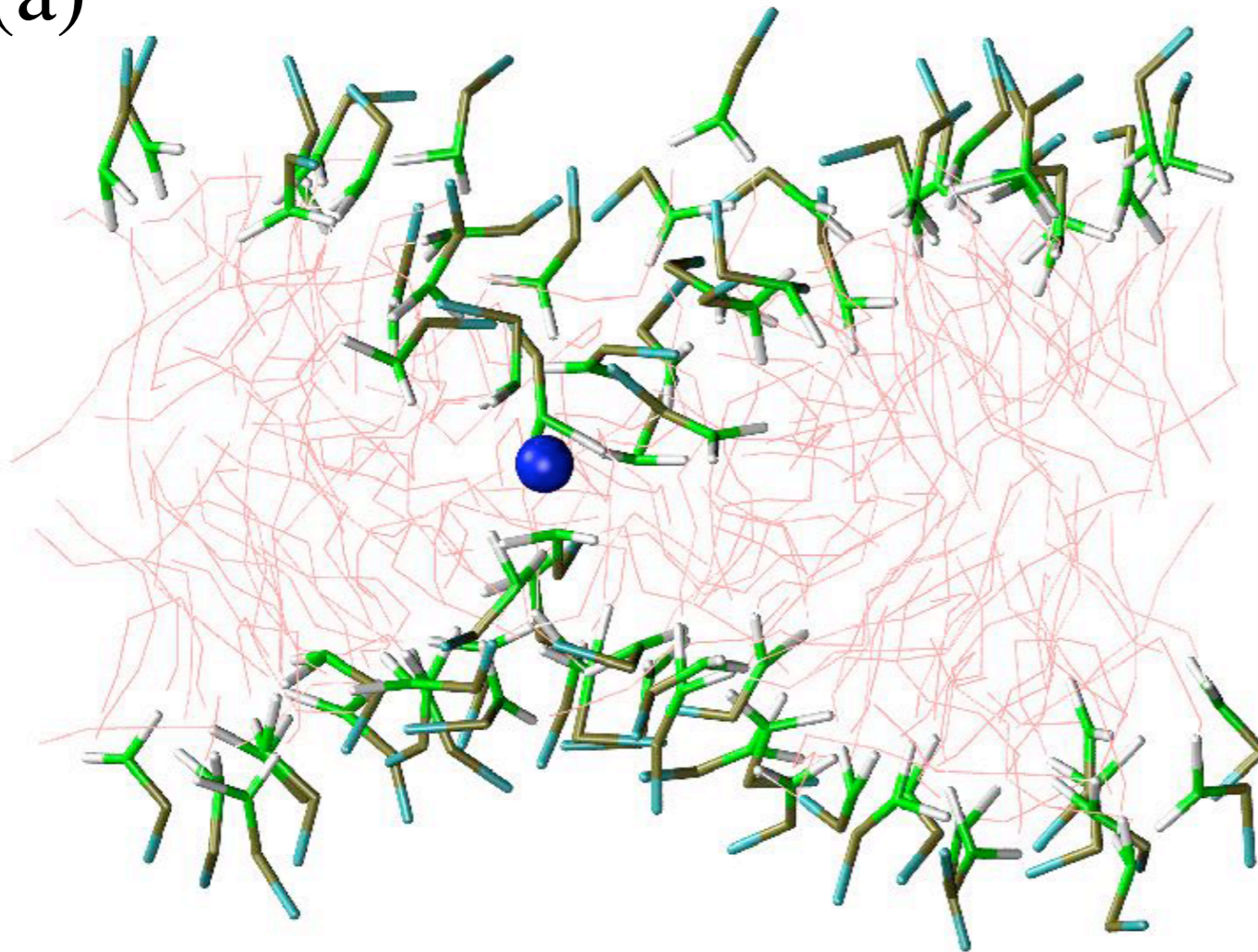
Gln      “Sharp” features  
            difficult to match with  
            low-resolution model.

$\text{Glu}^-$       Final interaction  
            potentials likely not  
            unique, but this is not a  
            problem.



# Arg<sup>+</sup> in a DOPC bilayer

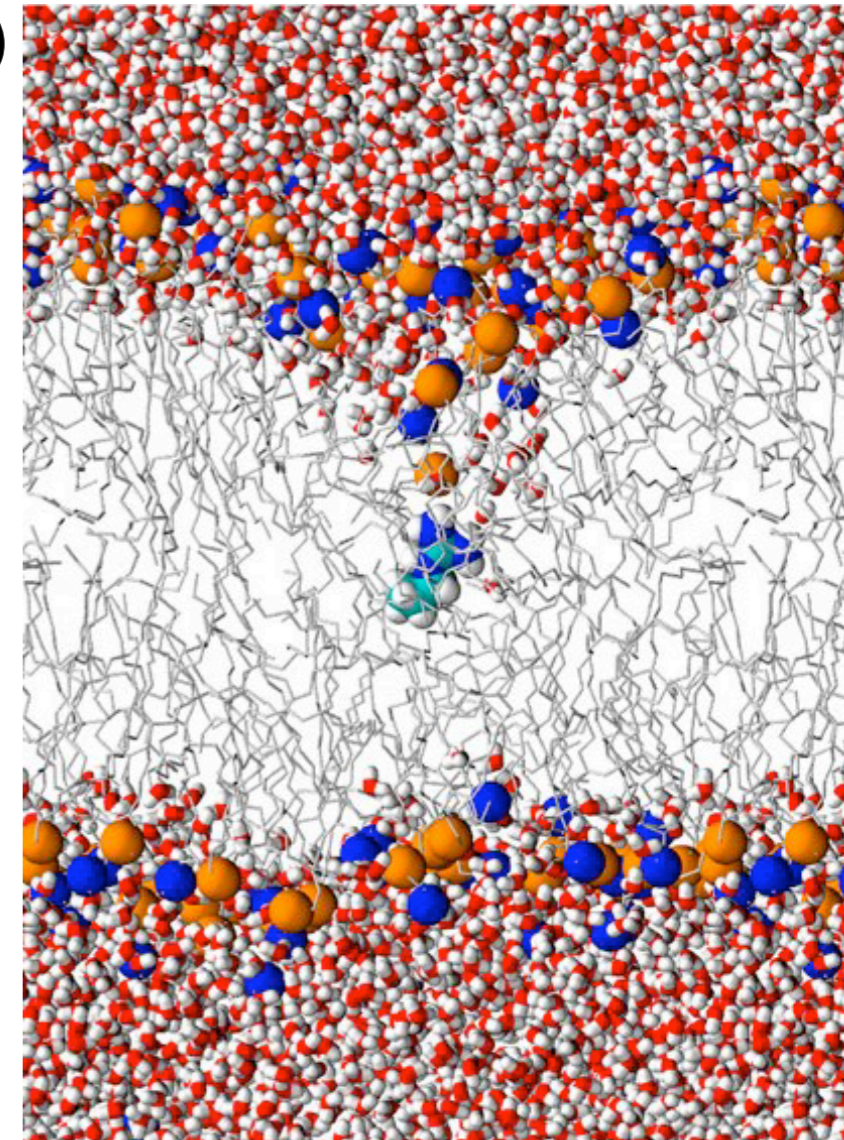
(a)



coarse grained

can observe “water defects” (without water)

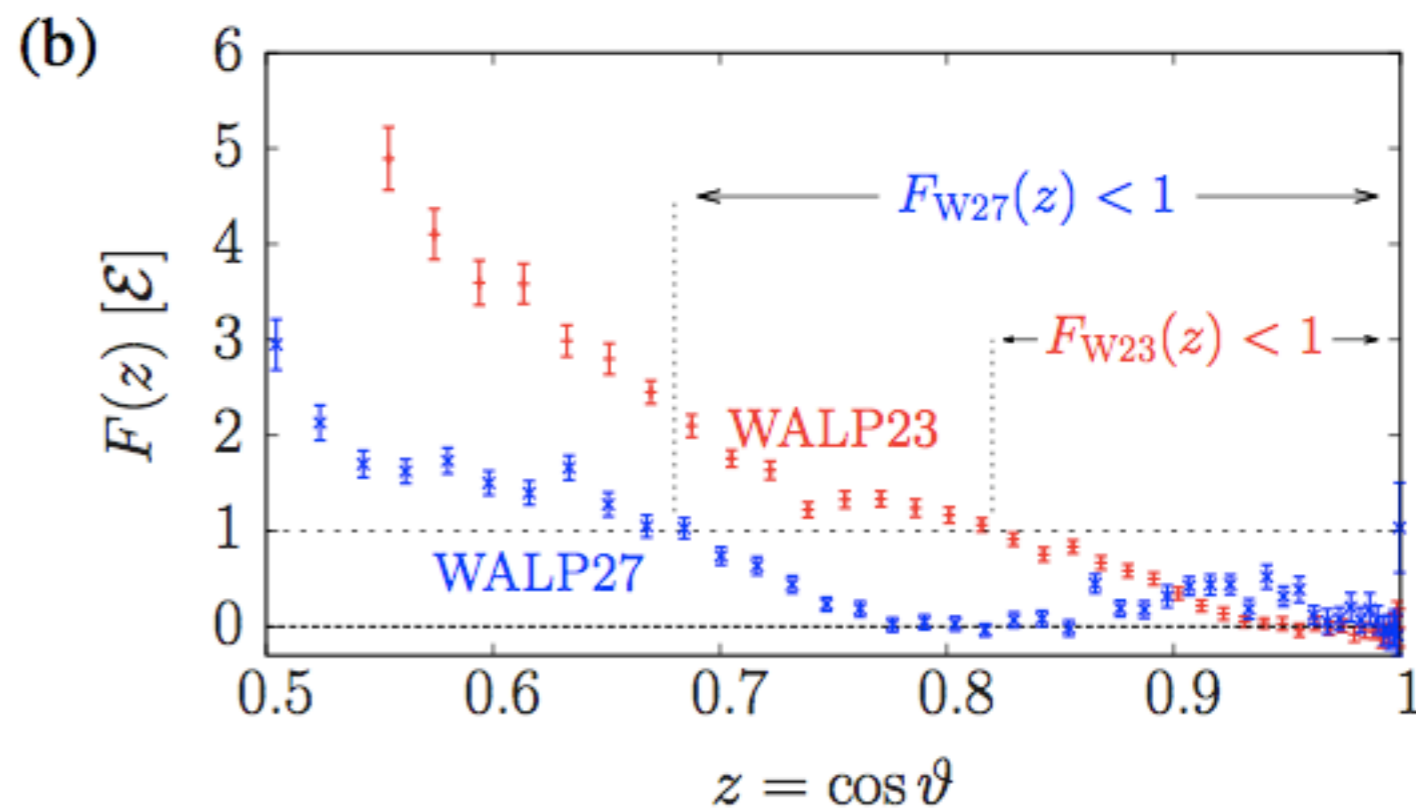
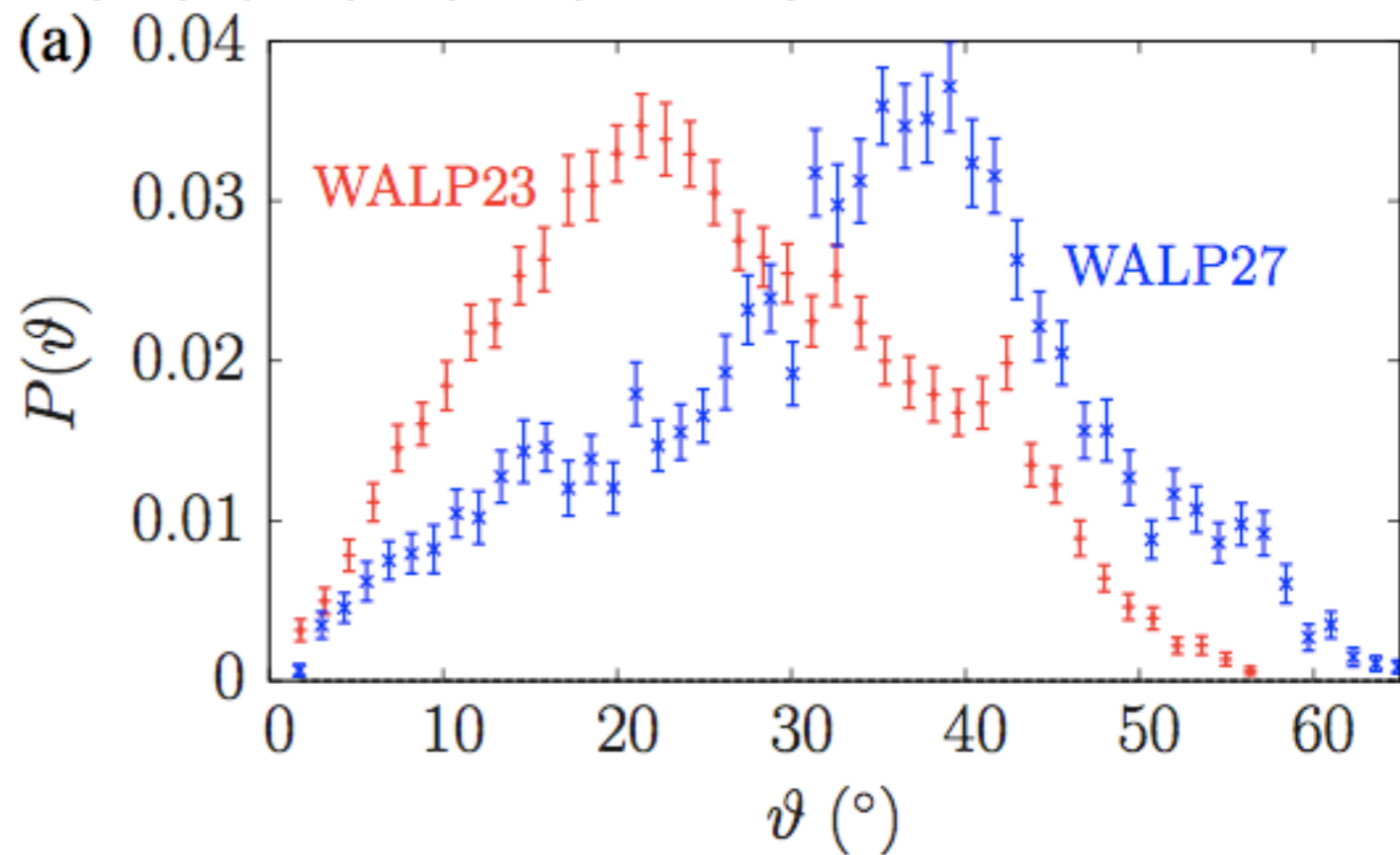
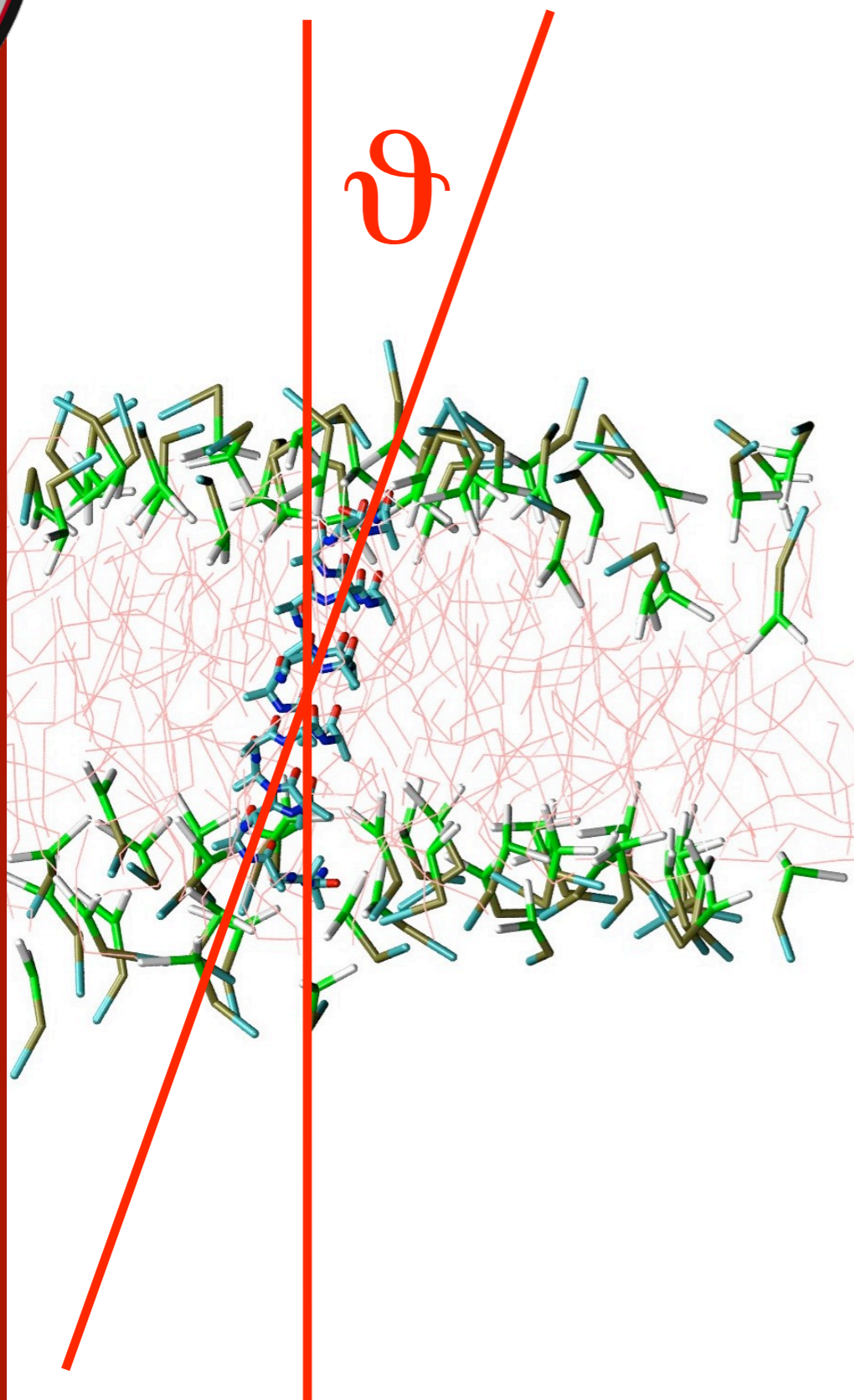
(b)



atomistic



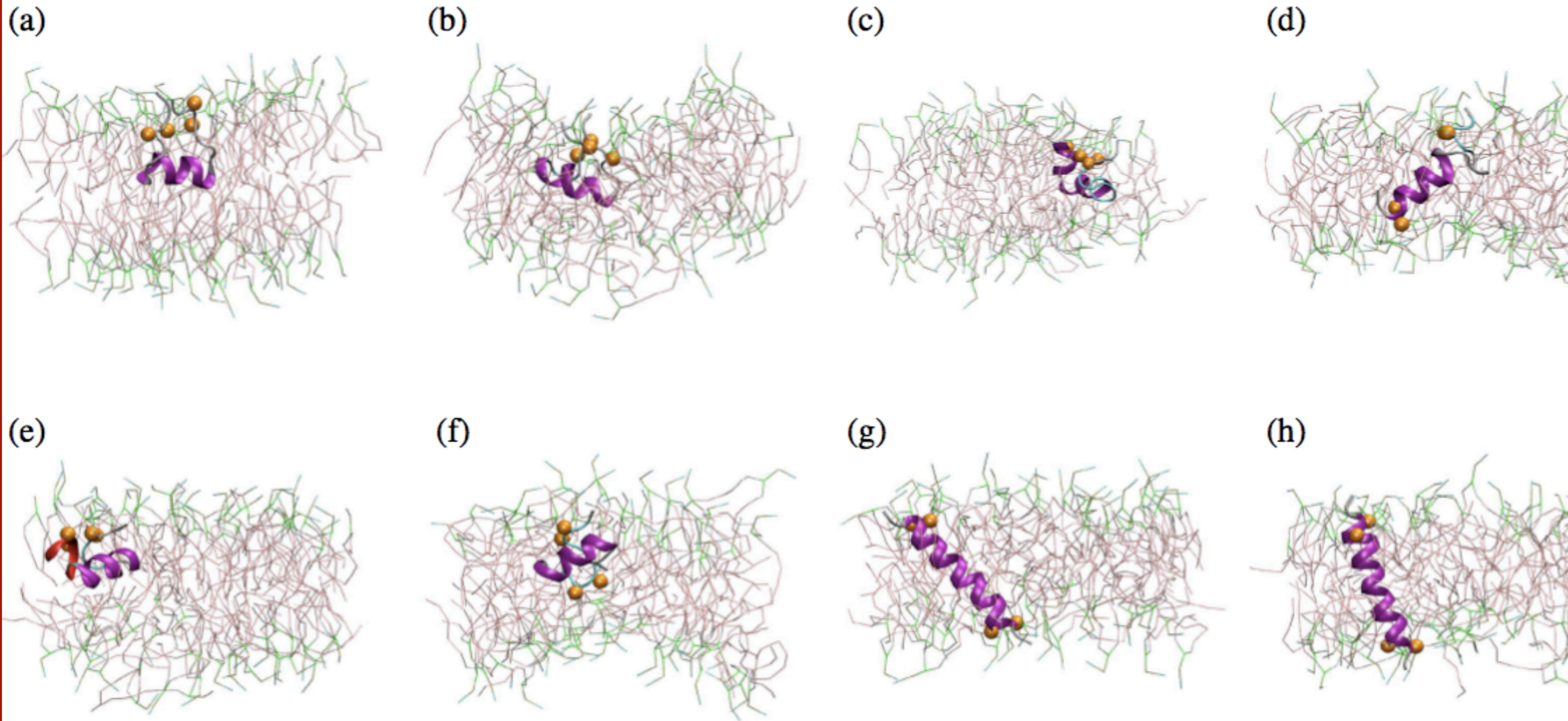
# Helix tilt fluctuations







# Folding of WALP23 in membrane environment



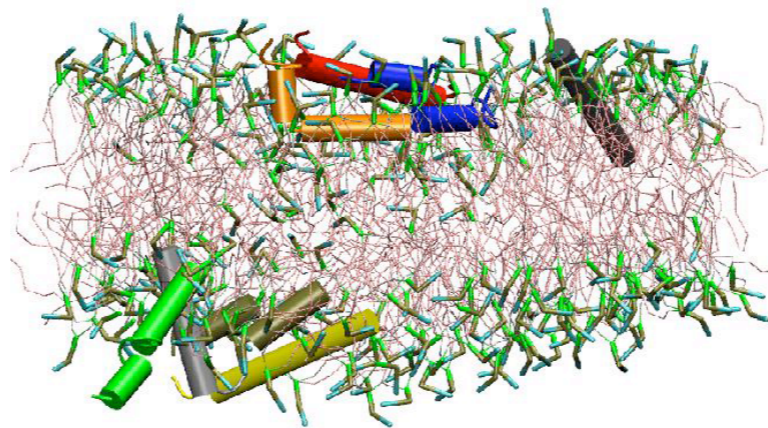
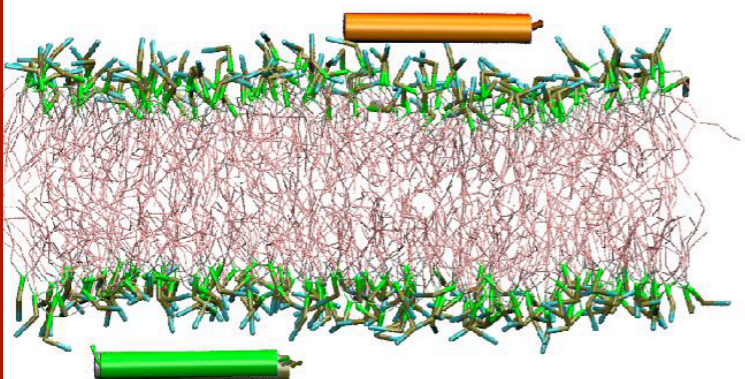
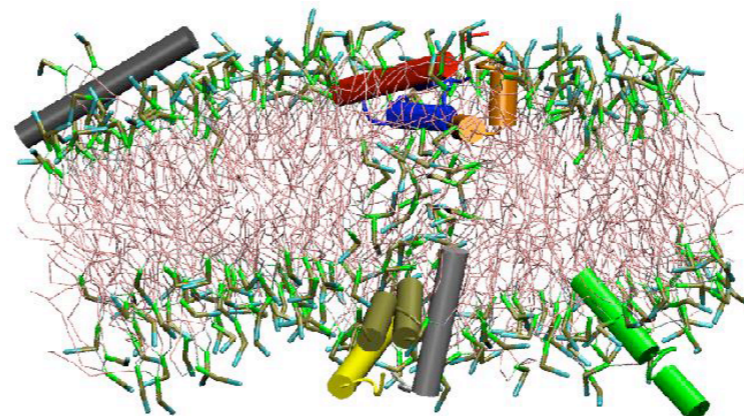
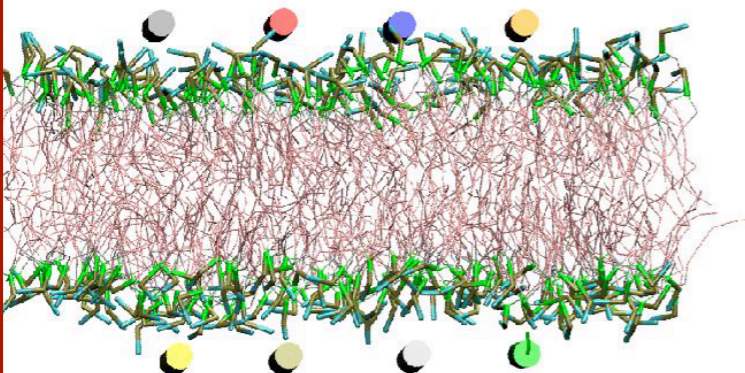
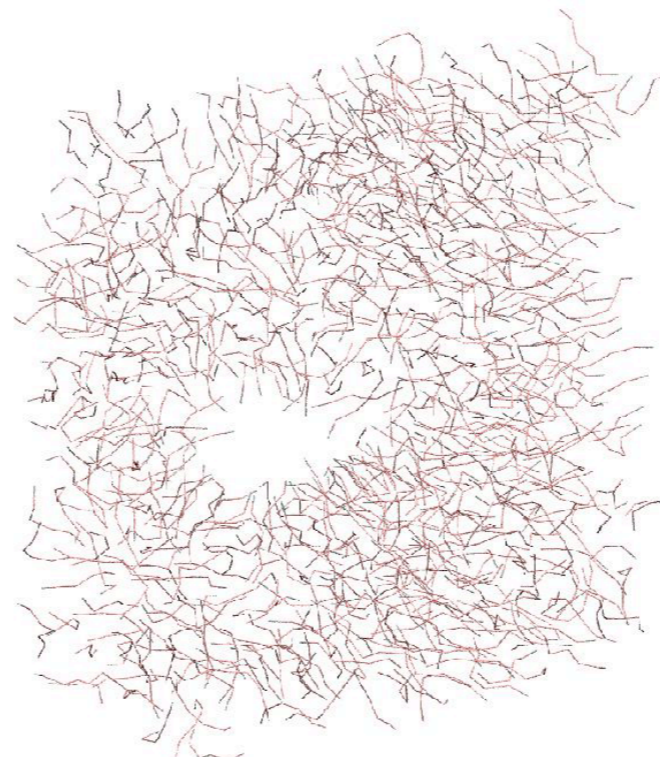
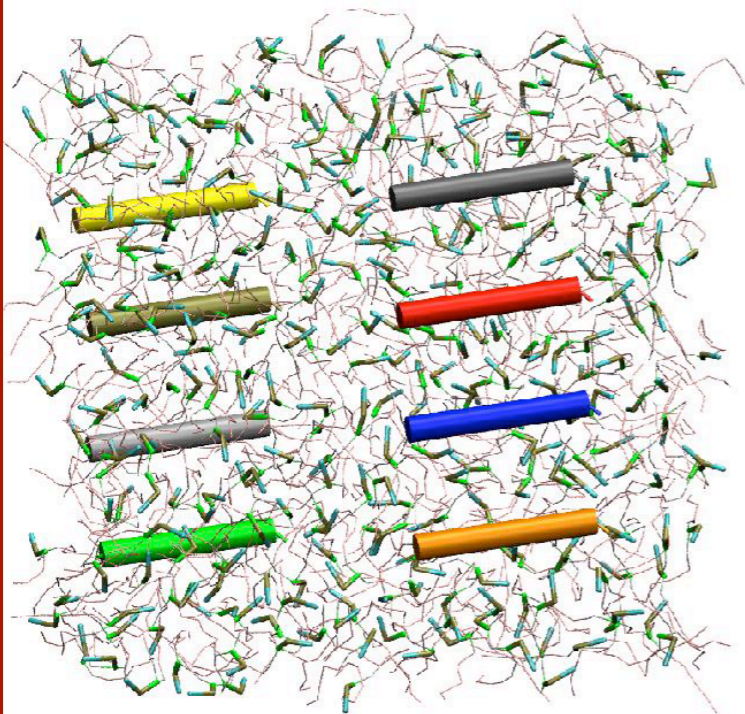
WALP23 in 62 lipid POPC bilayer at  $T=356\text{K}$ .  
CG REMD, total time  $\sim 10^6 \tau$  ( $\sim\sim 8\mu\text{s}$ )



# Pore formation by APs

$t = 0 \tau$

$t = 291\,000 \tau$



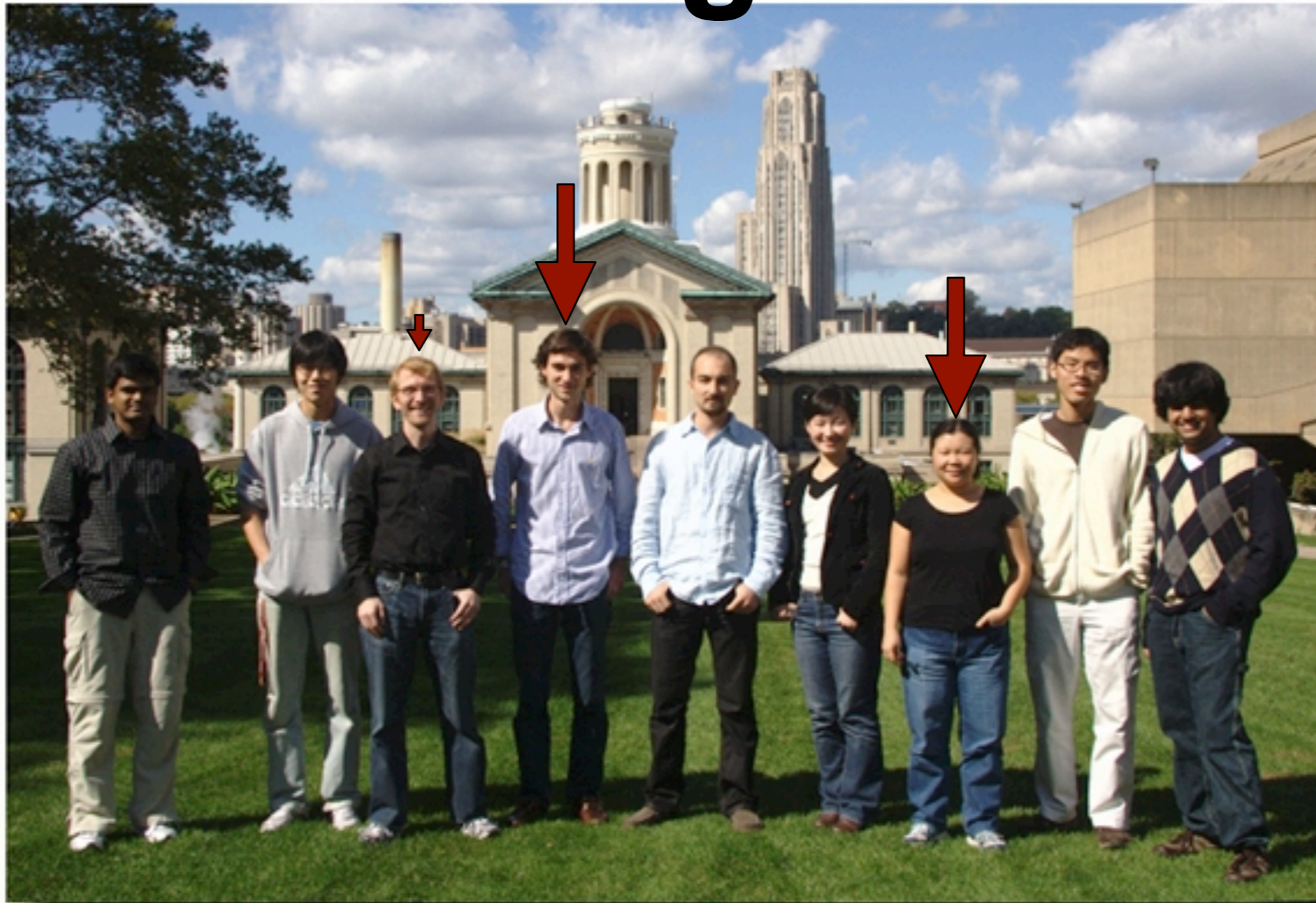
8 magainin, 288  
POPC,  $T=308\text{K}$ .  
CG REMD.  
Total time  $\sim\sim 2\mu\text{s}$

Pore is created without  
even one AP actually  
spanning it. Helices  
break.

Results secondary at  
this point. But notice the  
kinds of problems you  
could treat in this way.



# Acknowledgements



Senthil Muthiah, Mingyang Hu, **Tristan Bereau**, Cem Yolcu,  
Xiaofei Li, **Zun-Jing Wang**, Eddie Chua, Karpur Shukla

\$\$\$: DFG, NSF, NIH, Volkswagenstiftung, MPG, CMU

**Carnegie Mellon**